

PAUL SCHERRER INSTITUT



Diego Alejandro Sanz-Becerra - on behalf of the muonEDM collaboration :: ECT\* workshop on EDMs

# Toward an improved measurement of the muon EDM

## Project funded by



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**State Secretariat for Education,  
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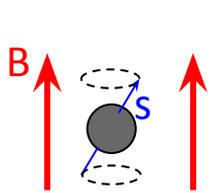
**FONDS NATIONAL SUISSE  
SCHWEIZERISCHER NATIONALFONDS  
FONDO NAZIONALE SVIZZERO  
SWISS NATIONAL SCIENCE FOUNDATION**

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# THE BASICS



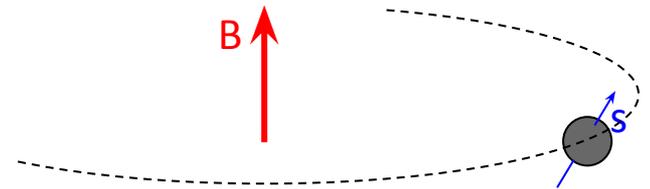
$$\vec{\omega}_0 = -\frac{q a}{m} \left( \left(1 + \frac{1}{\gamma a}\right) \vec{B} - \frac{\gamma}{\gamma + 1} (\vec{B} \cdot \vec{\beta}) \vec{\beta} - \left(1 + \frac{1}{a(1 + \gamma)}\right) \frac{\vec{\beta} \times \vec{E}}{c} \right)$$

Larmor precession

$$\hat{H} = -\vec{\mu} \cdot \vec{B} - \vec{d} \cdot \vec{E}$$

$$\vec{\mu} = g \frac{q}{2 m_\mu} \vec{s} \quad \vec{d} = \eta \frac{q}{2 m_\mu c} \vec{s}$$

$$\vec{\omega}_c = -\frac{q a}{m} \left( \frac{\vec{B}}{\gamma a} - \frac{\gamma}{a(\gamma^2 - 1)} \frac{\vec{\beta} \times \vec{E}}{c} \right)$$



Cyclotron frequency

$$\vec{\omega}_a = \vec{\omega}_0 - \vec{\omega}_c = \frac{q a}{m} \left( \vec{B} - \frac{\gamma}{\gamma + 1} (\vec{B} \cdot \vec{\beta}) \vec{\beta} - \left( 1 + \frac{1}{a (1 - \gamma^2)} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right)$$

Larmor
Cyclotron

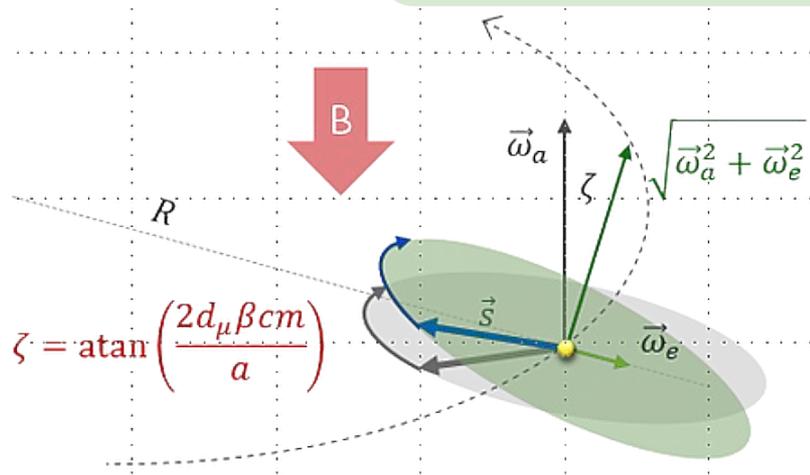
- This precession is only due to the magnetic dipole moment
- What if it has an EDM?

# MDM And EDM



$$\vec{\omega} = \vec{\omega}_a + \vec{\omega}_e = \frac{q a}{m} \left( \vec{B} - \frac{\gamma}{\gamma+1} (\vec{B} \cdot \vec{\beta}) \vec{\beta} - \left( 1 + \frac{1}{a(1-\gamma^2)} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right) + \frac{q \eta}{2 m} \left( \frac{\vec{E}}{c} - \frac{\gamma c}{\gamma+1} (\vec{E} \cdot \vec{\beta}) \vec{\beta} + \vec{\beta} \times \vec{B} \right)$$

$\vec{B} \perp \vec{\beta} \perp \vec{E}$

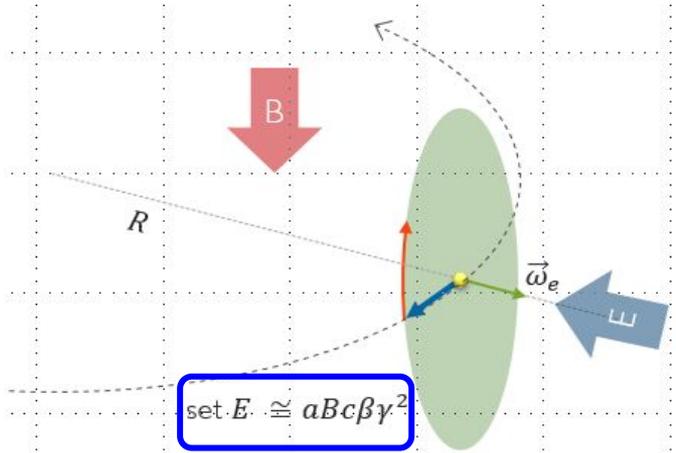


# The Frozen-Spin Technique

$$\vec{\omega} = \vec{\omega}_a + \vec{\omega}_e = \frac{q a}{m} \left( \vec{B} - \frac{\gamma}{\gamma + 1} (\vec{B} \cdot \vec{\beta}) \vec{\beta} - \left( 1 + \frac{1}{a (1 - \gamma^2)} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right) + \frac{q \eta}{2 m} \left( \frac{\vec{E}}{c} - \frac{\gamma c}{\gamma + 1} (\vec{E} \cdot \vec{\beta}) \vec{\beta} + \vec{\beta} \times \vec{B} \right)$$

$\vec{B} \perp \vec{\beta} \perp \vec{E}$

$E \approx a B c \beta \gamma^2$



# The Frozen-Spin Technique



$$\vec{\omega}_{FS} = \frac{q \eta}{2 m} \left( \frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} \right)$$

Muon's rest frame:

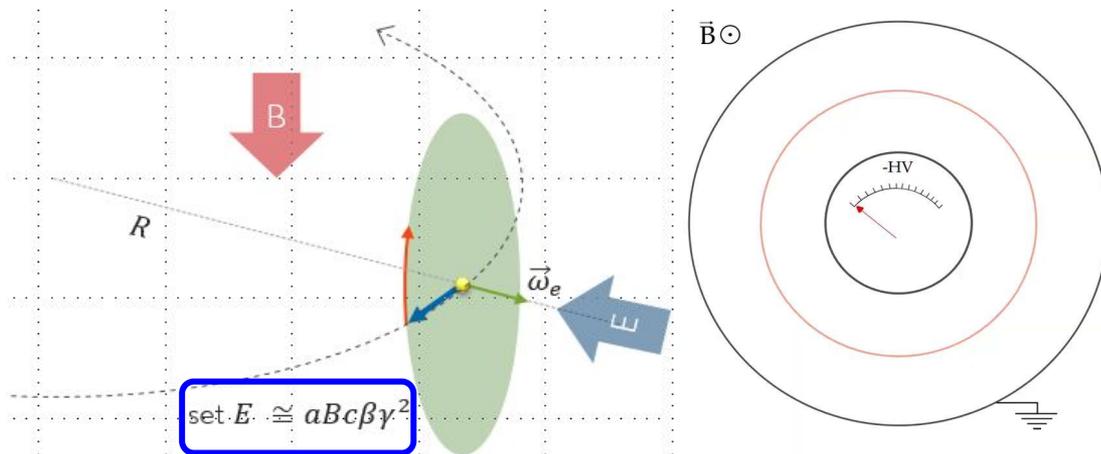
Lorentz boosted B field

$$\vec{E}'_B = \gamma c \vec{\beta} \times \vec{B}$$

$$\vec{B} \perp \vec{\beta} \perp \vec{E} \quad E \approx a B c \beta \gamma^2$$

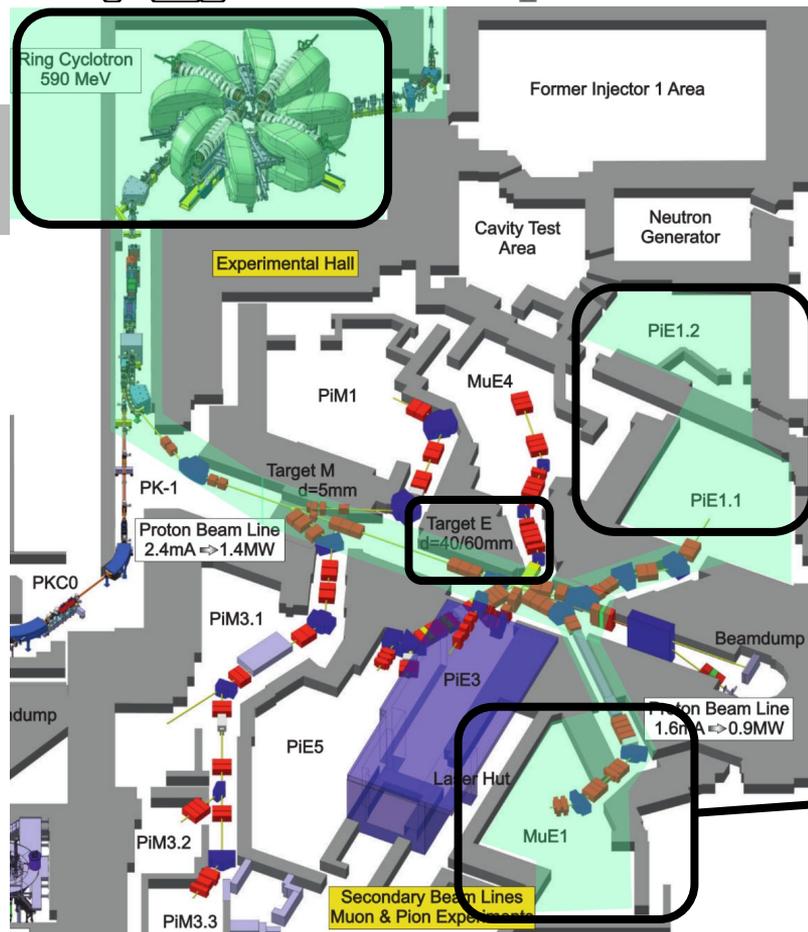
For the experiment:

- B = 3 T
- p: O(100 MeV/c)
- E: O(1 MV/m)
- E'\_B: O(1 GV/m)



# THE EXPERIMENT

# The Experiment - The Muons



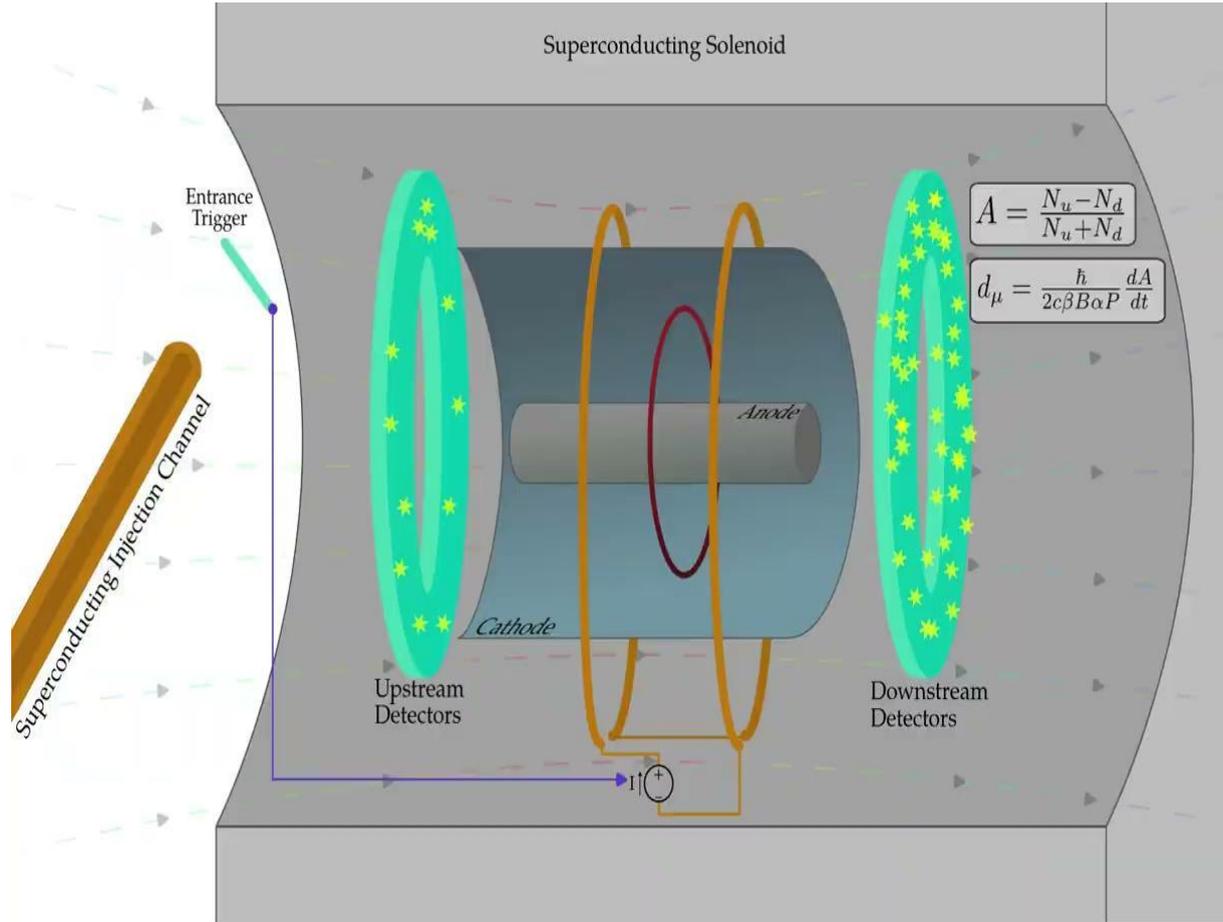
## Phase 1:

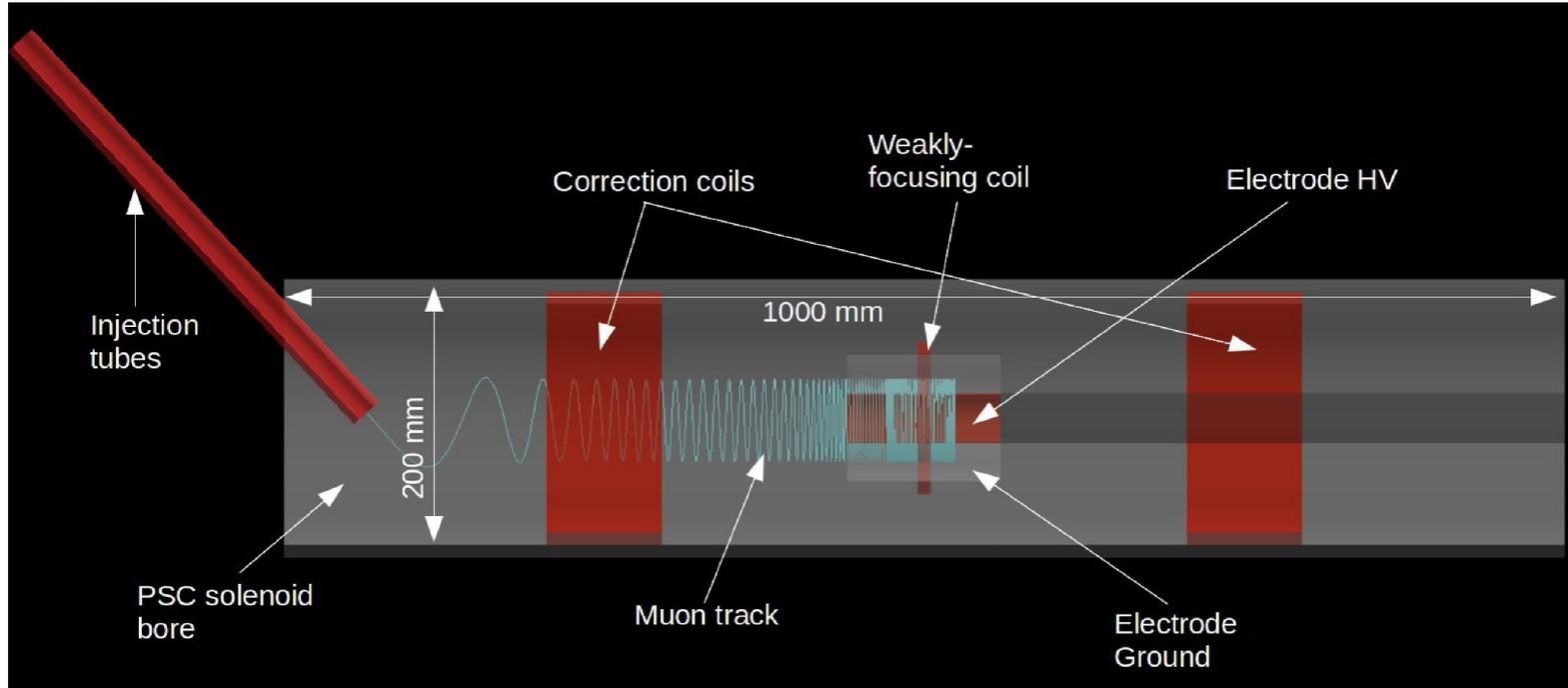
- Negative helicity (95%)
- $p=28 \text{ MeV}/c$
- Flux:  $O(10^6 \mu^+/\text{s})$
- Sensitivity/year  $< 3 \times 10^{-21} \text{ e} \cdot \text{cm}$

## Phase 2:

- Positive helicity (95%)
- $p=125 \text{ MeV}/c$
- Flux:  $O(10^8 \mu^+/\text{s})$
- Sensitivity/year  $< 6 \times 10^{-23} \text{ e} \cdot \text{cm}$

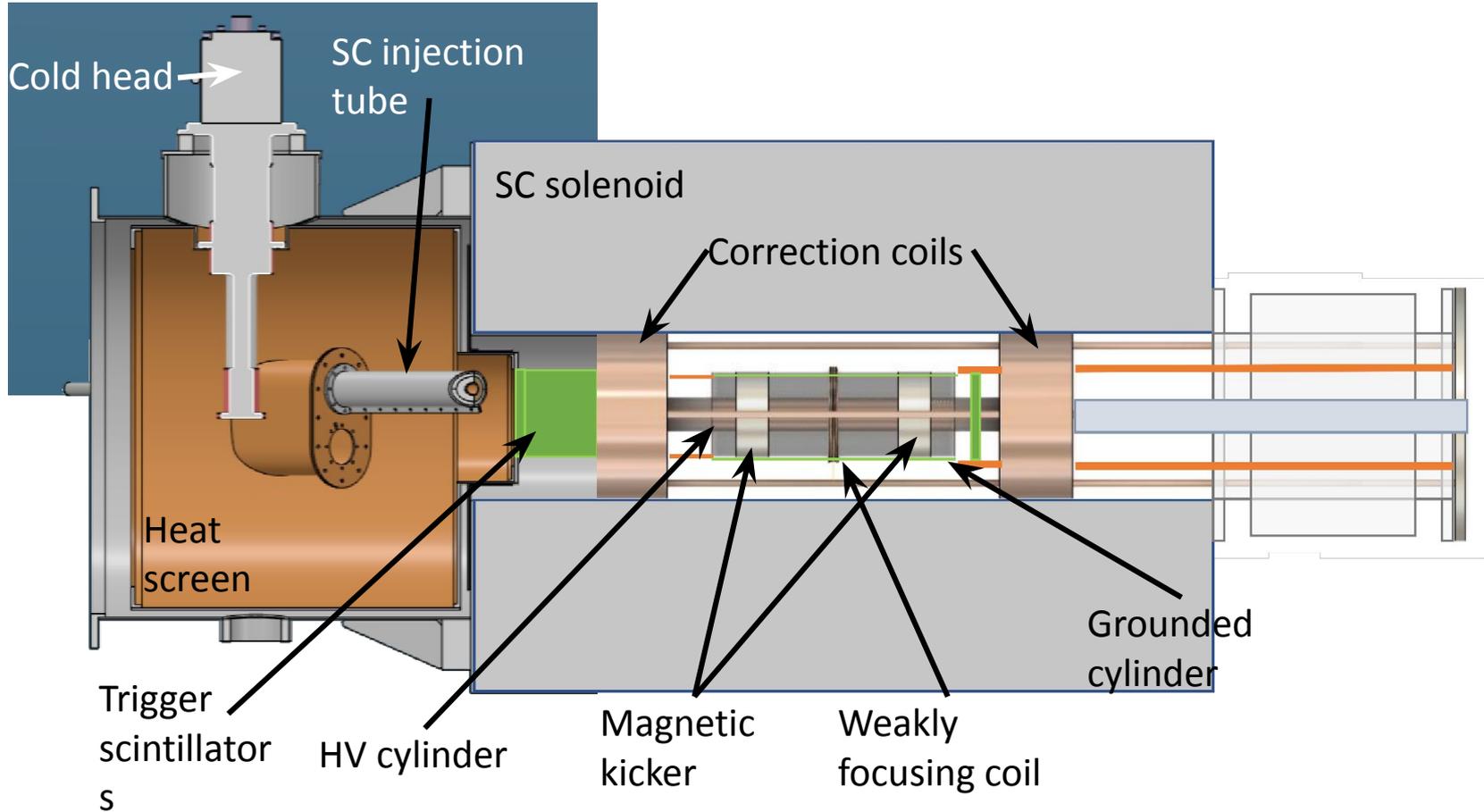
# The Experiment



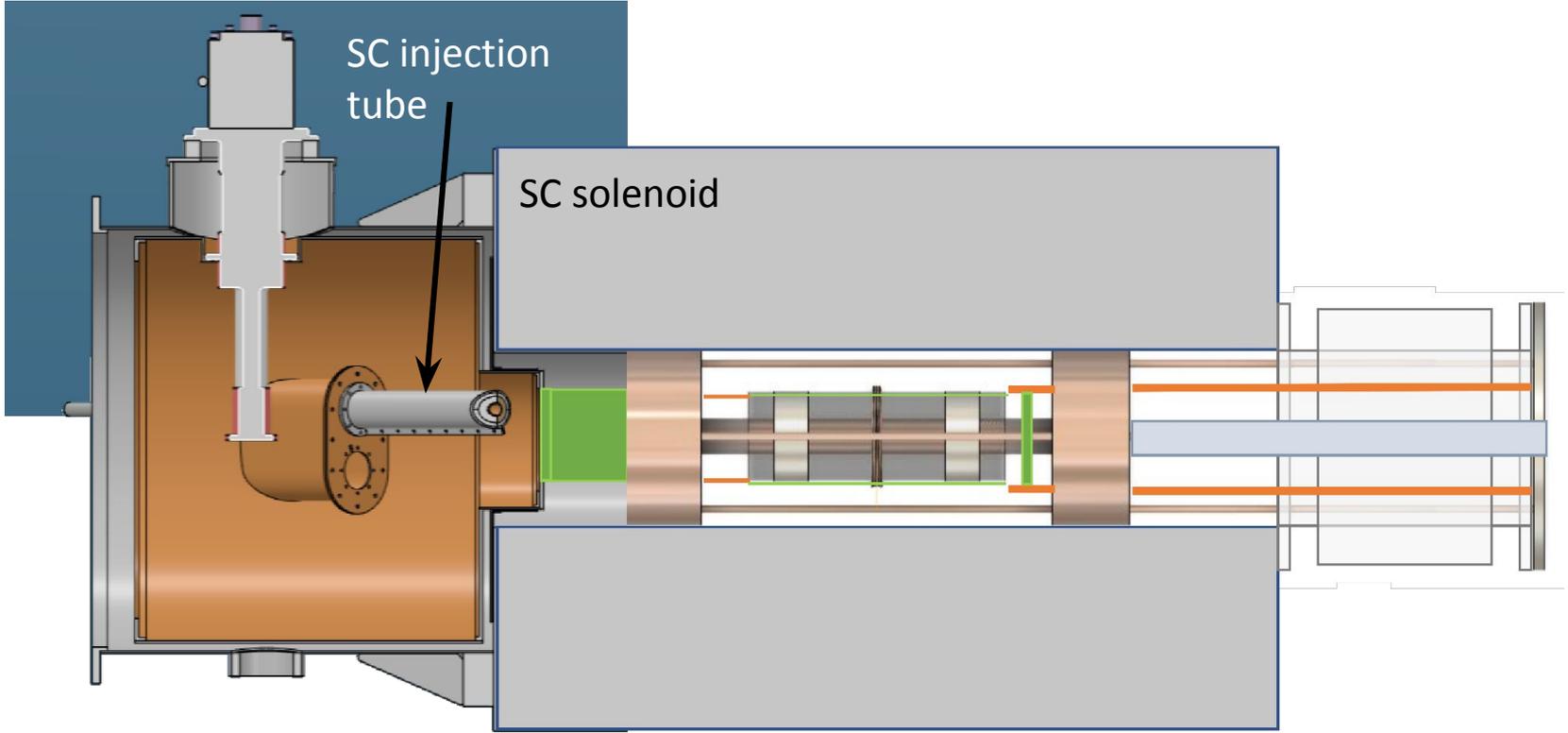


- Using G4beamline and surrogate models, we optimized the storage efficiency of muons up to  $\sim 0.4\%$ .

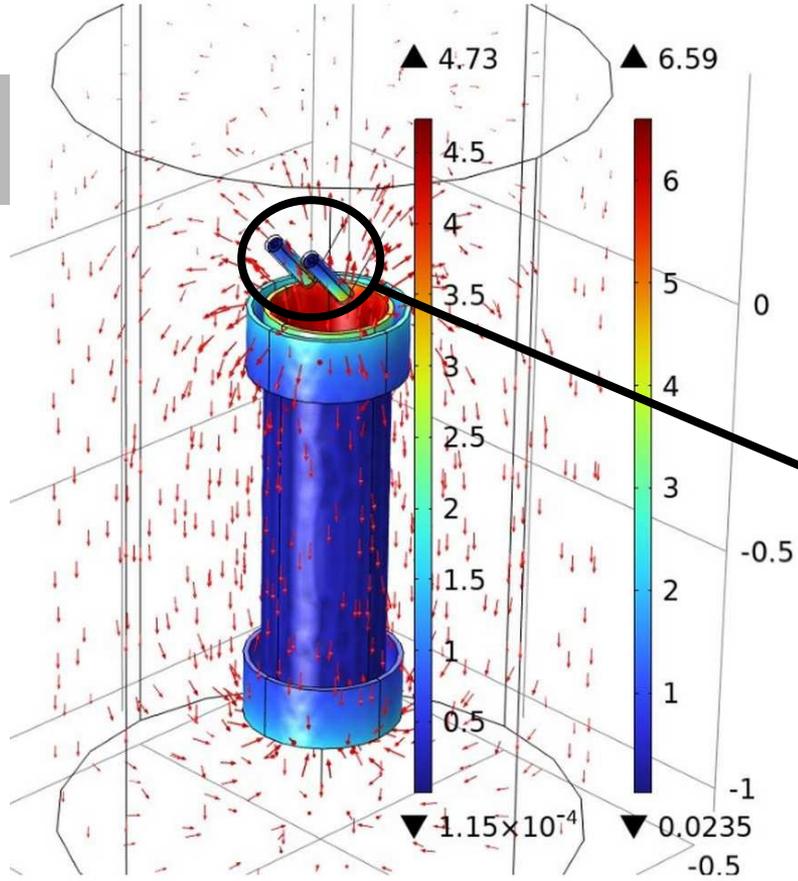
# The Experiment - Parts



# Injection Channels

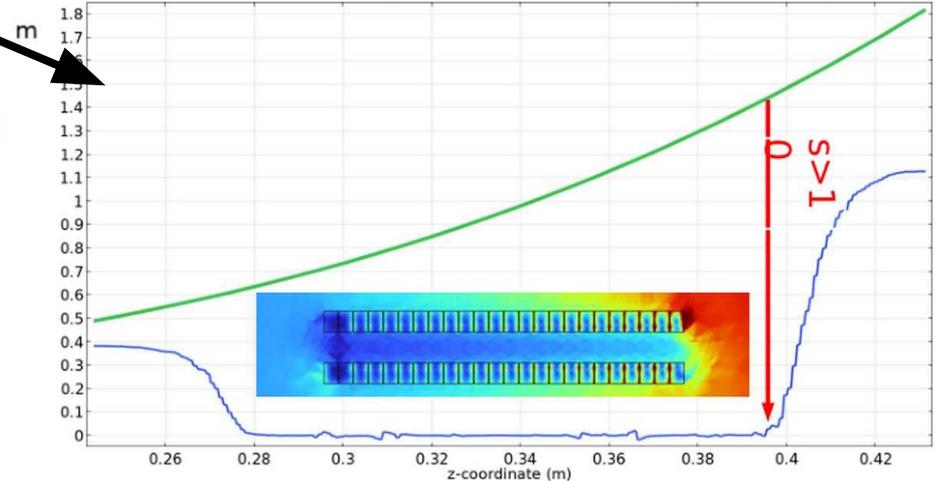


# Injection Channels - High Fields

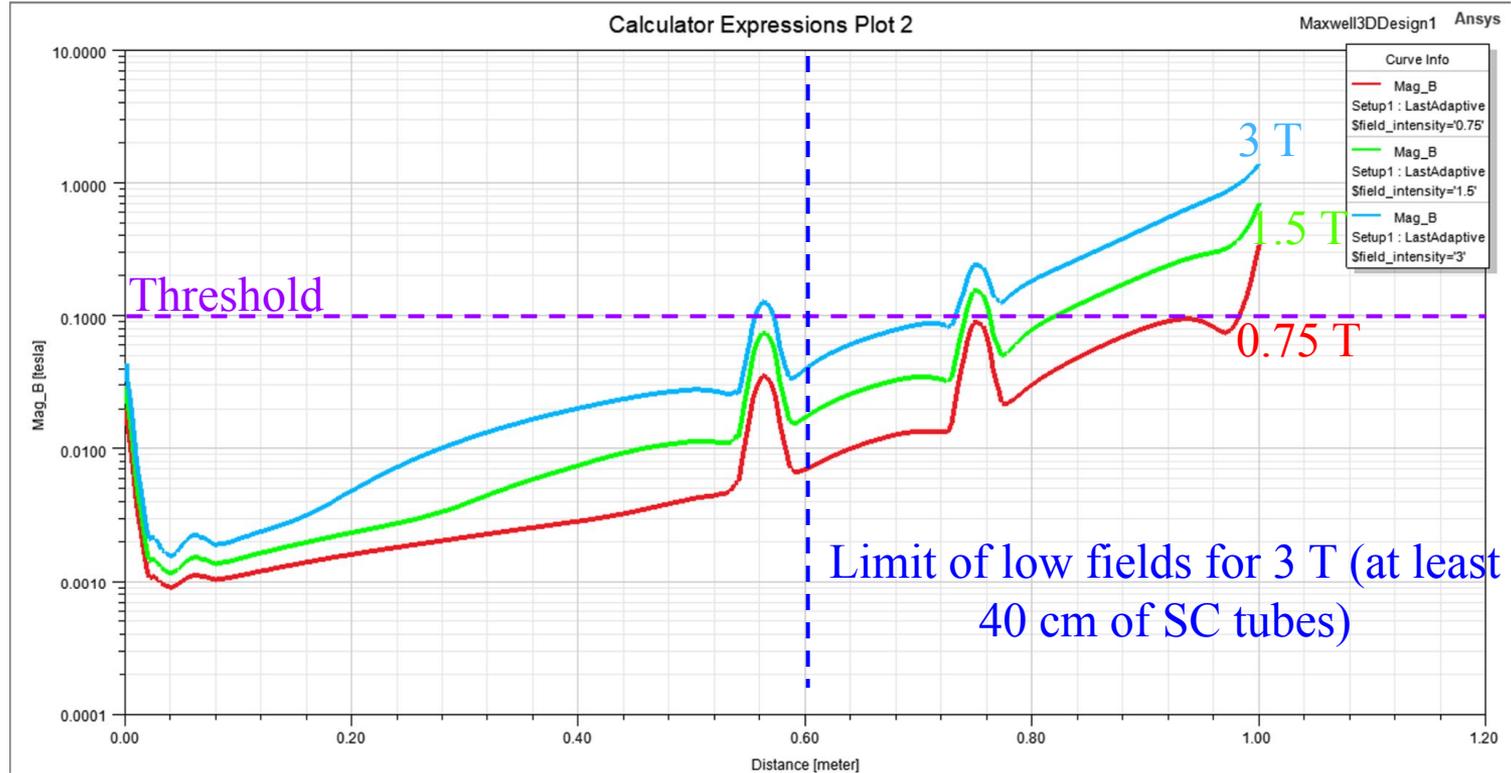


- Shield muons during the transport to the magnet's bore ( $B \gtrsim 0.8$  T) with SC materials.
- SC channels inside cryostat.
- Simulating different concepts and SC materials.
- Expected prototype tests this year.

Unshielded and shielded magnetic field intensity (T)

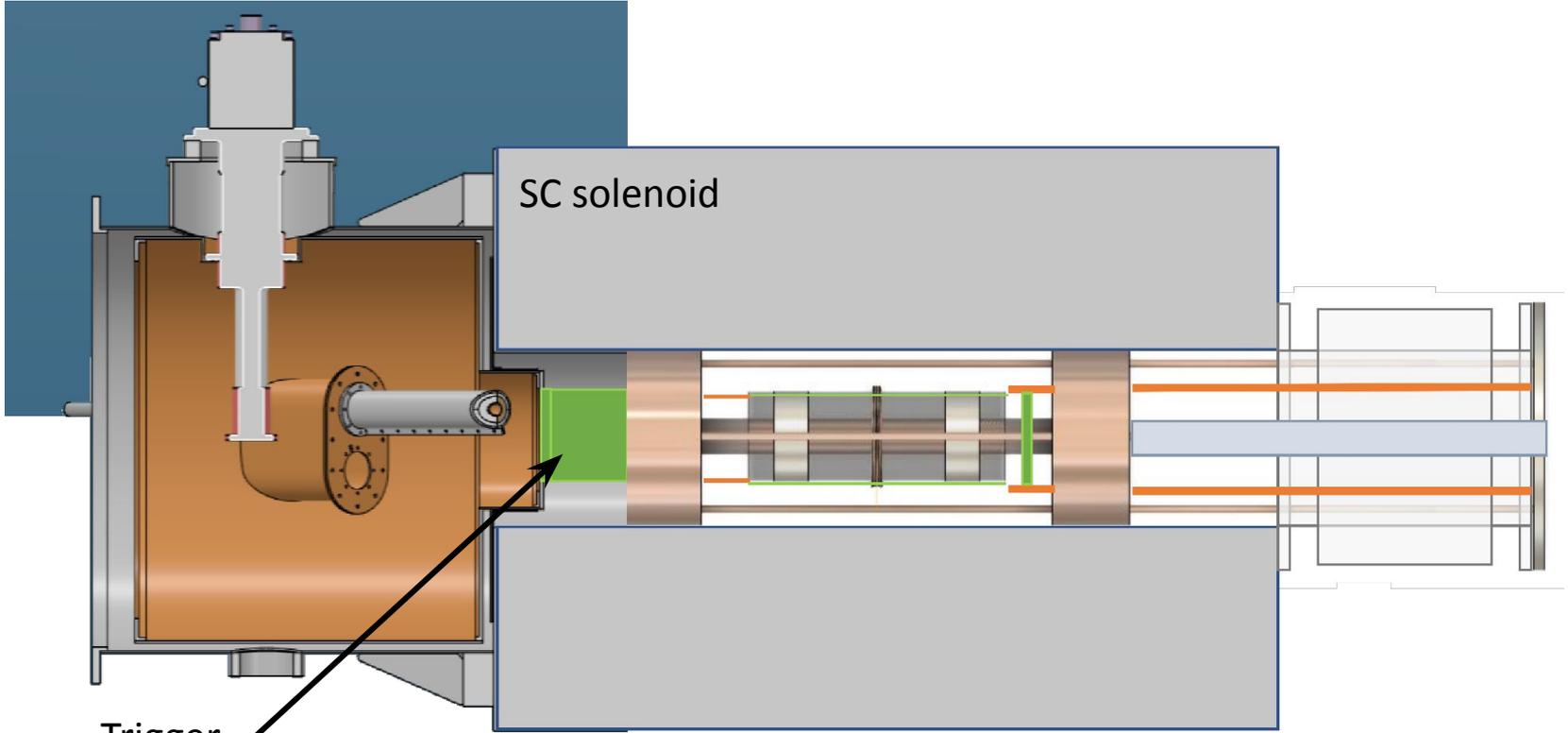


# Injection Channels - Low Fields



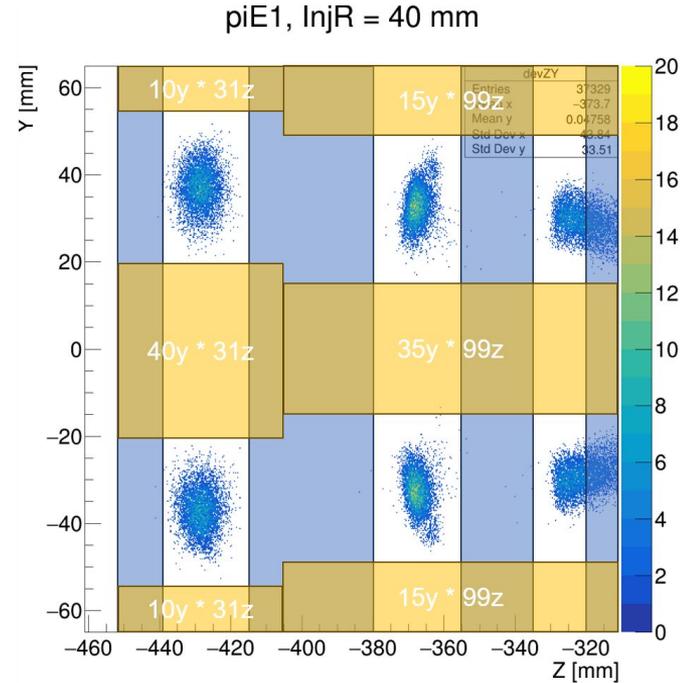
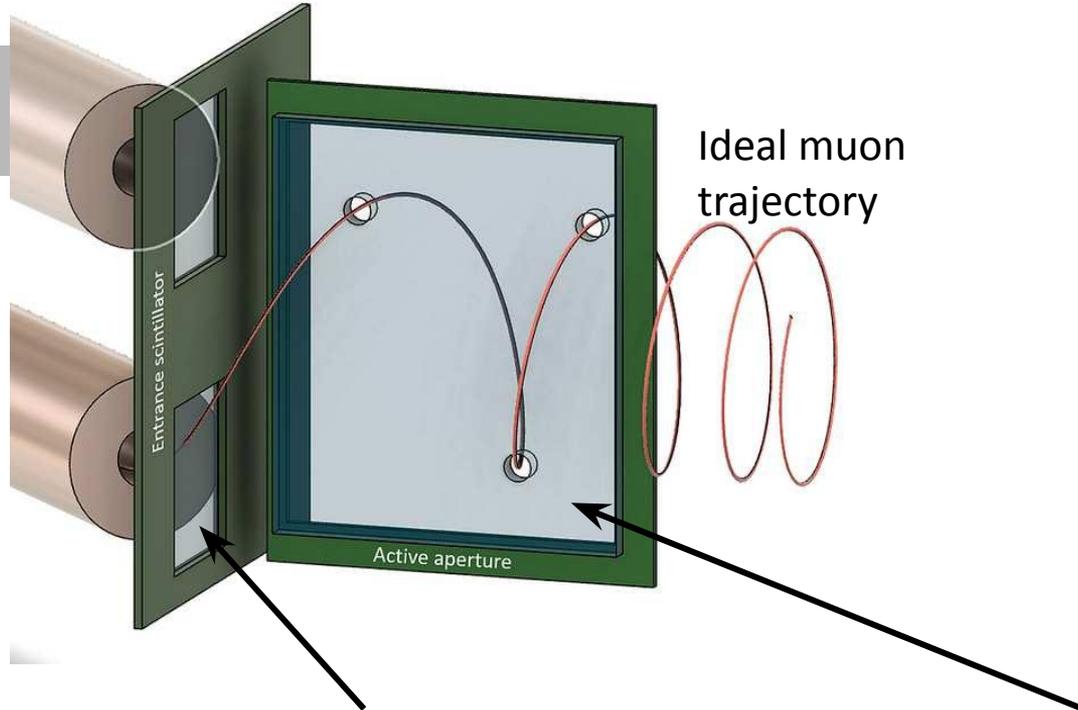
- Use magnetic steel to prolong the SC tubes outside cryostat to the beamline window.
- Shield the injected muons to below  $\sim 100$  mT.

# Trigger Scintillators



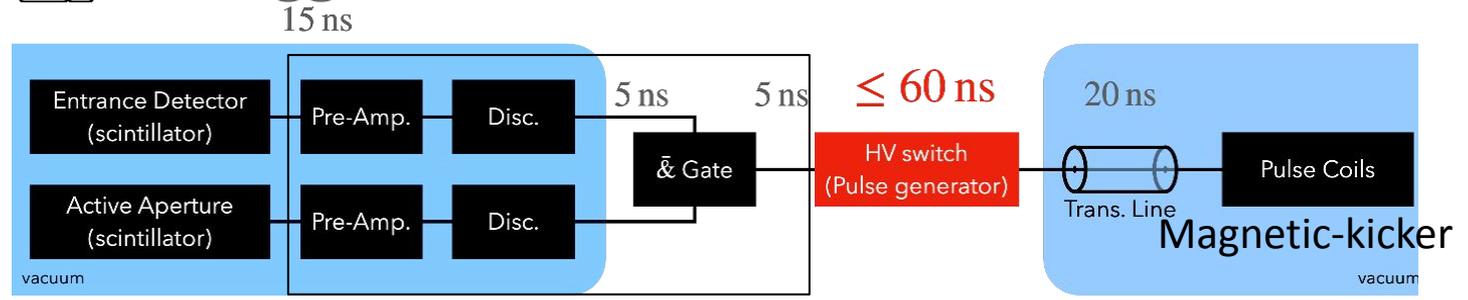
Trigger  
scintillator

S

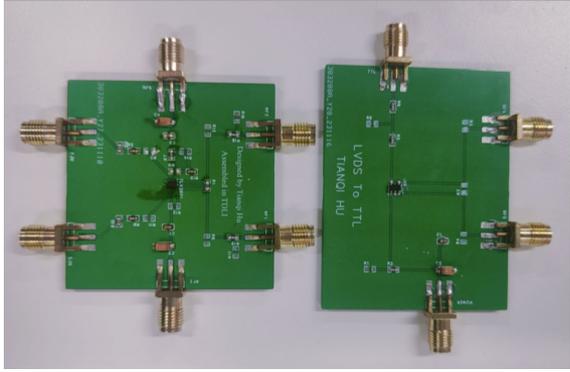


- Thin entrance detector in anti-coincidence with veto-detectors
- Meshed array of scintillator strips determine the trajectory of the storable-muons

# Trigger Scintillators - Fast Electronics R&D

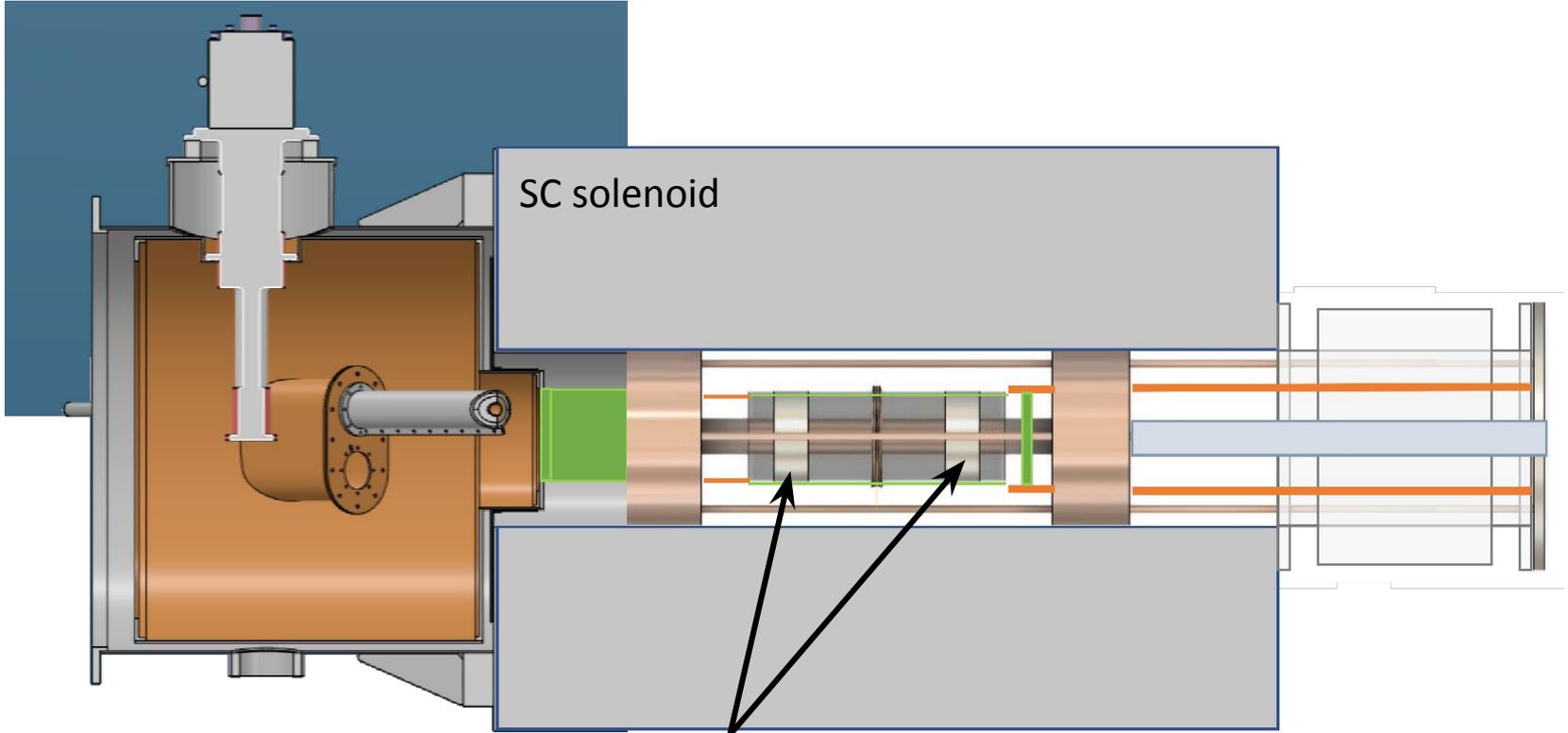


- The trigger signal has to reach the Magnetic kicker within  $O(100$  ns)



- Custom electronics with 3 ns discriminator and 2 ns pre-amplifiers and splitter.

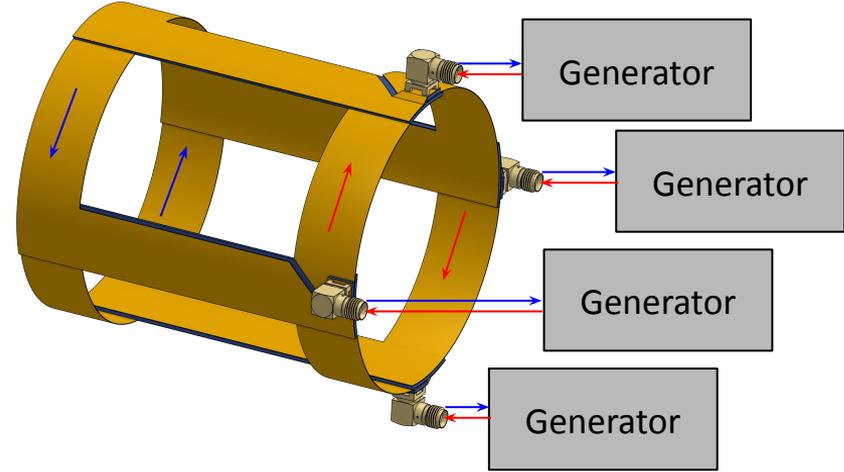
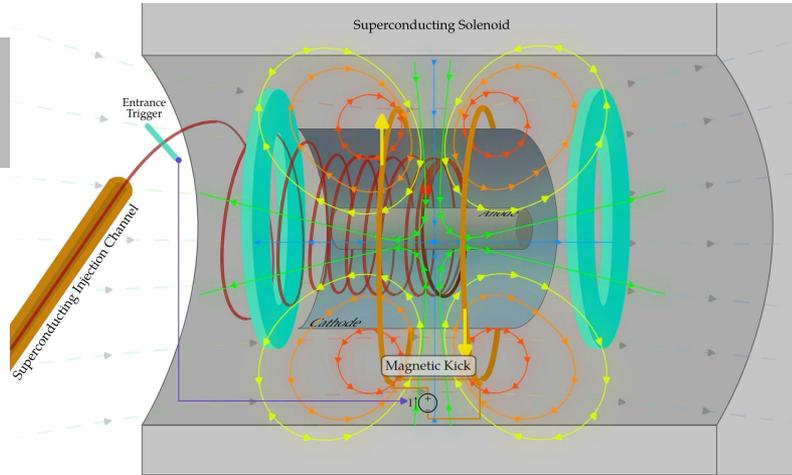
# Trigger Scintillators



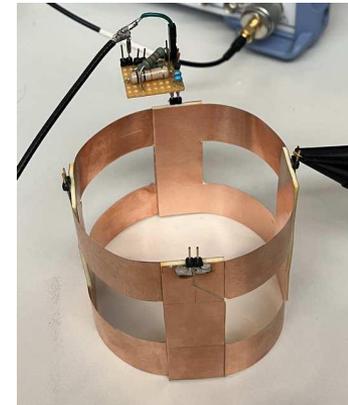
SC solenoid

Magnetic  
kicker

# Magnetic Kicker R&D



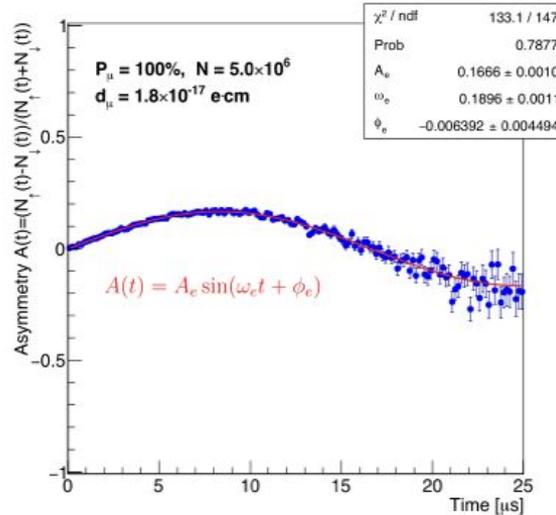
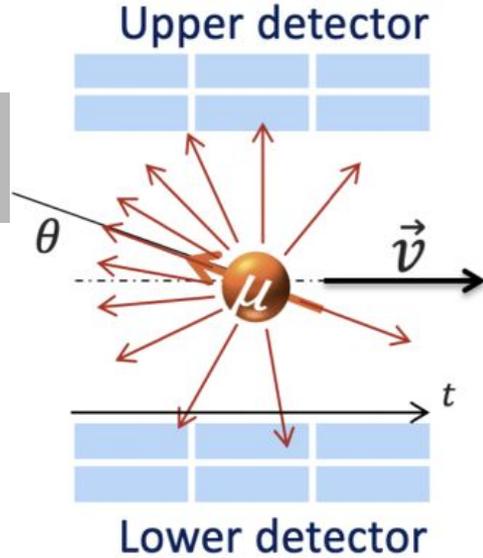
- Magnetic kick generated by Anti-Helmholtz coils.
- Segmented into quadrants to reduce inductance:
  - Lower driving voltage.
  - Faster pulse.



Prototype

# THE MEASUREMENTS

# The Measurements - The Signal



$$A = \frac{N_u - N_d}{N_u + N_d}$$

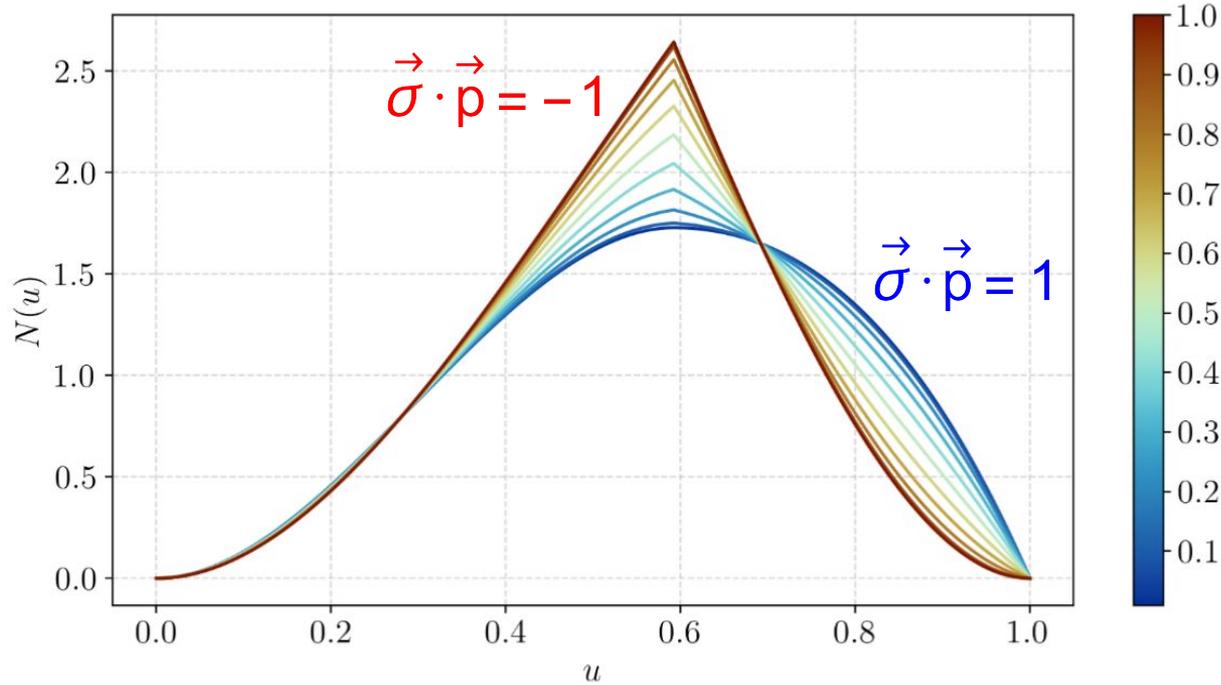
$$\dot{A}(\Psi_0, t) = \frac{\partial A}{\partial \Psi} \frac{\partial \Psi}{\partial t} \Big|_{\Psi_0} = \alpha(\Psi_0) \dot{\Psi}(t)$$

$$\sigma_{d\mu} = \frac{\hbar}{2 c \beta B \alpha P} \frac{1}{\gamma \tau \sqrt{N}}$$

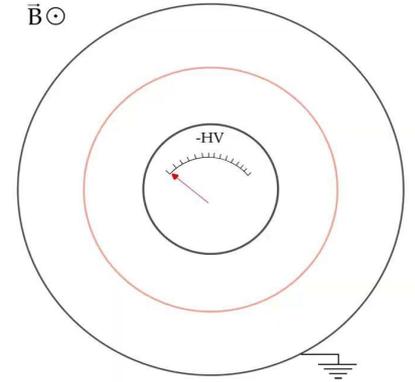
- Asymmetry precession due to EDM is too slow.
- The change of asymmetry with respect to time is relevant.
- The sensitivity is optimized by maximizing  $\alpha\sqrt{N}$ : (Figure-of-merit)

- P: Initial polarization
- $E_B'$ :  $c\beta B\gamma$  (Boosted B)
- N: Detected positrons
- T: Muon life-time
- $\alpha$ : Analysis parameter

# The Measurements - Freezing The Spin



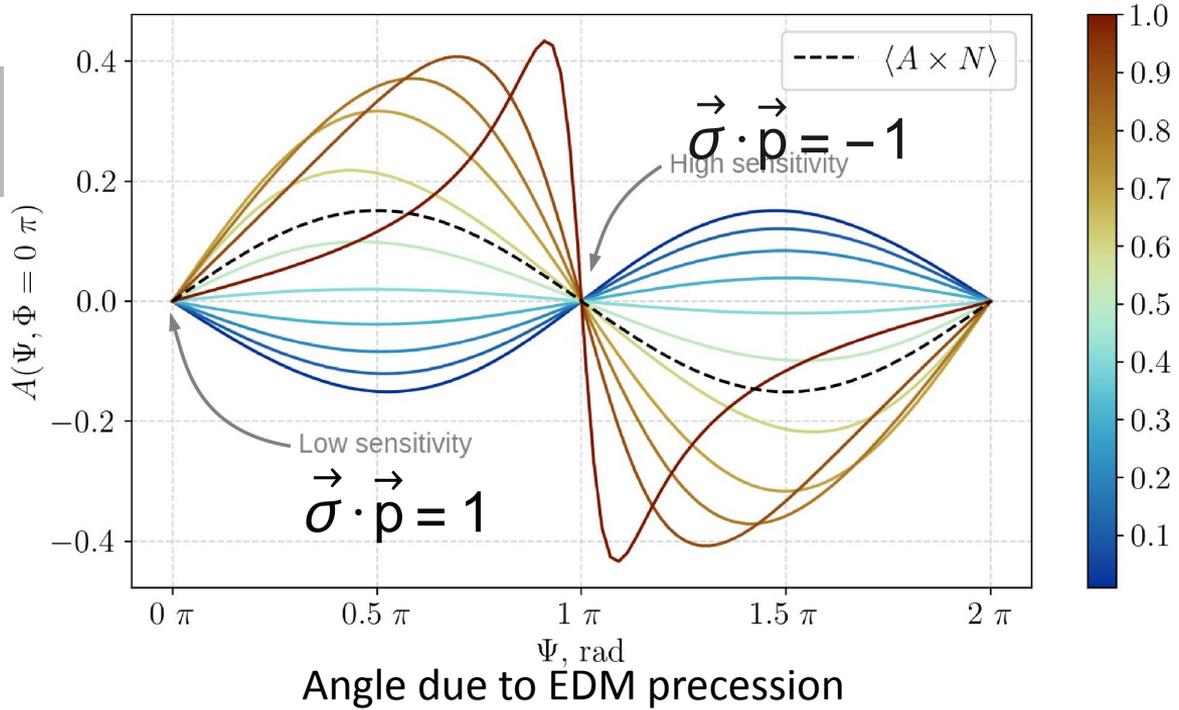
Angle due to g-2 precession  $\Phi, \pi \text{ rad}$



Fractional positron energy  $E/E_{\text{max}}$

- Counts oscillate between the red curve and the blue curve.
- Counts above a threshold would give the g-2 oscillation signal.
- Increase the electric field until the g-2 oscillations “freeze”

# The Measurements - Measuring The EDM

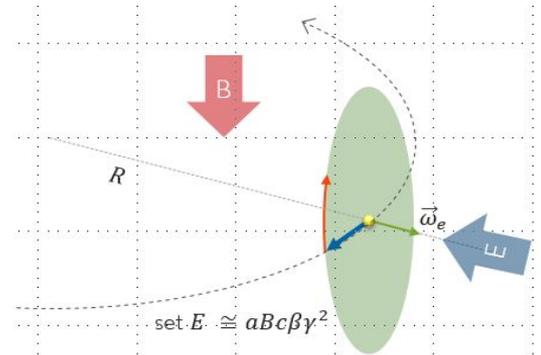


- Low sensitivity working point: Negative helicity.
- High sensitivity working point: Positive helicity.
  - The configuration for phase 1.

Fractional positron energy  $E/E_{\text{max}}$   

$$\dot{A}(\Psi_0, t) = \frac{\partial A}{\partial \Psi} \frac{\partial \Psi}{\partial t} \Big|_{\Psi_0} = \alpha(\Psi_0) \dot{\Psi}(t)$$
 Working point  
 EDM precession  

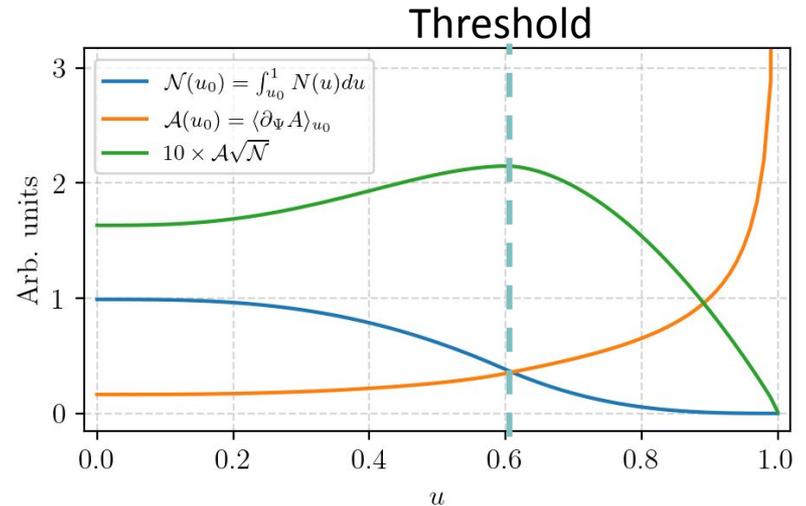
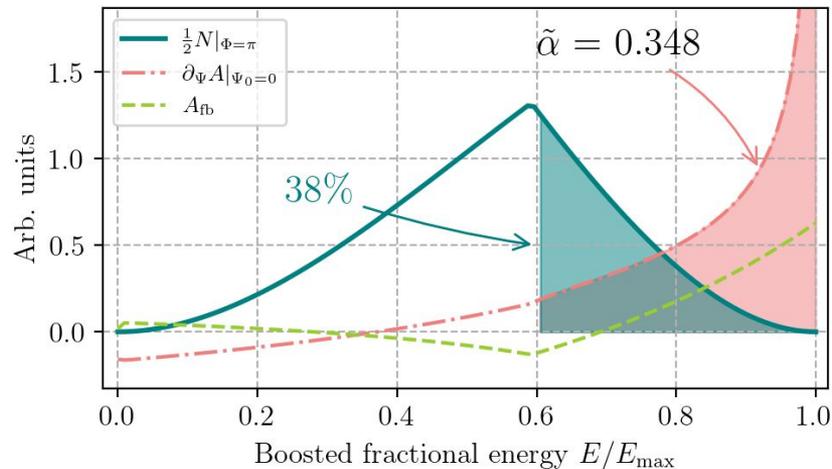
$$d_\mu = \frac{\hbar}{2 c \beta \alpha P} \dot{A}$$



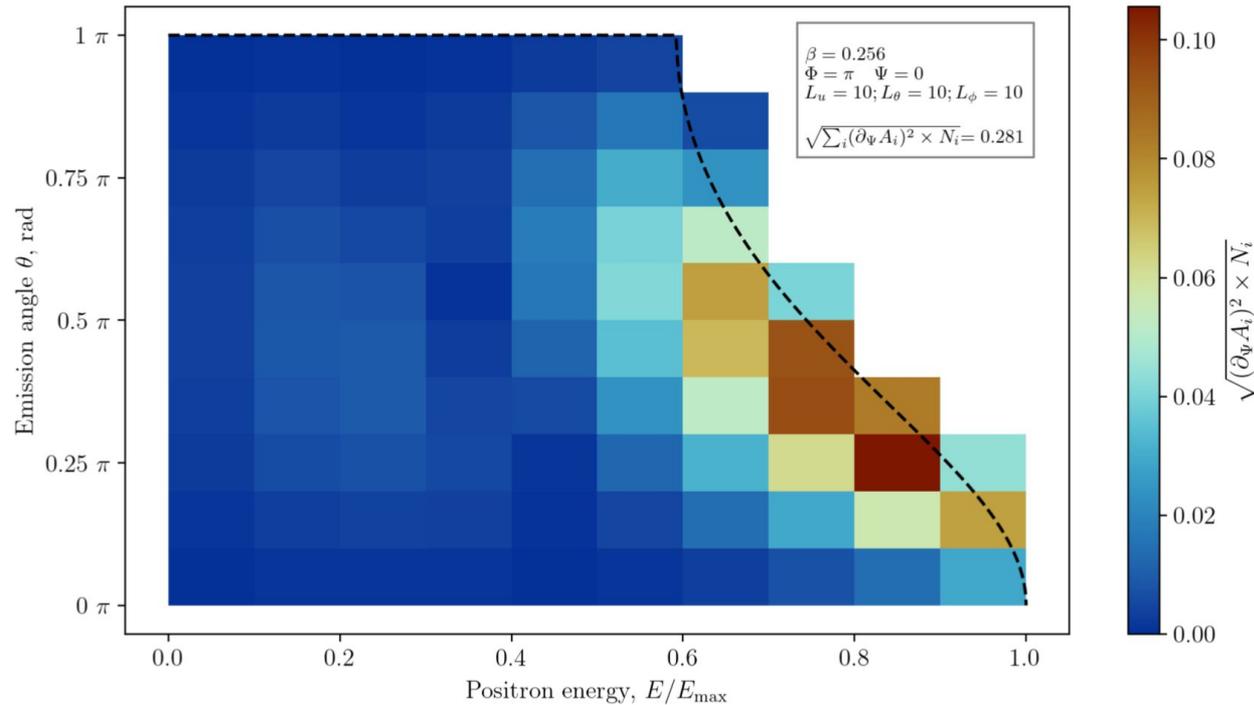
We studied three methods with different complexity:

1. Simple: Use all measured events. 
$$d_\mu = \frac{\hbar}{2 c \beta \alpha P} \dot{A}$$

2. T-Method: Count events only over a threshold that maximizes the FoM. 
$$d_\mu = \frac{\hbar}{2 c \beta \tilde{\alpha} P} \dot{A}_{Th}$$



### 3. W-Method: Weight the measured events for each voxel in energy and direction.



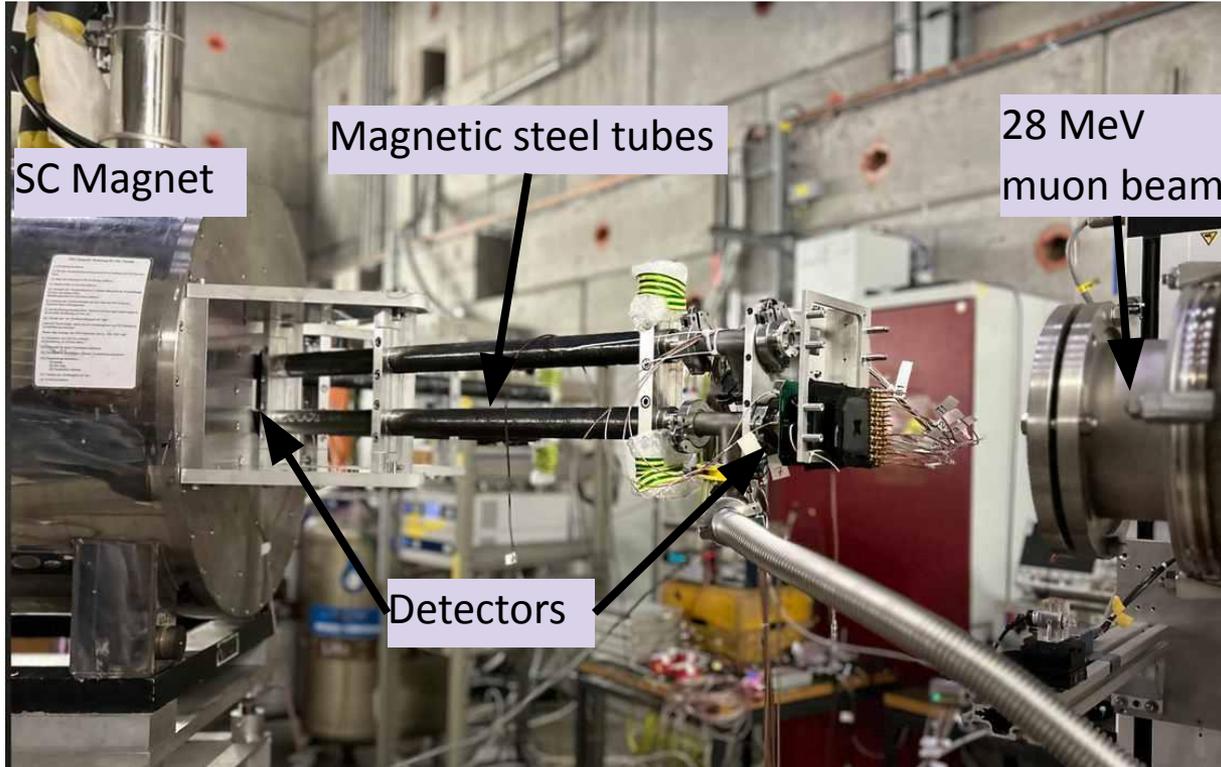
$$W_i = \alpha_i \sqrt{N_i}$$

$$W = \sqrt{\sum_i W_i^2} = \sqrt{\sum_i \alpha_i^2 N_i}$$

$$d_\mu = \frac{\hbar}{2 c \beta B P} \frac{\sum_i \frac{\dot{A}_i}{\alpha_i} W_i^2}{\sum_i W_i^2}$$

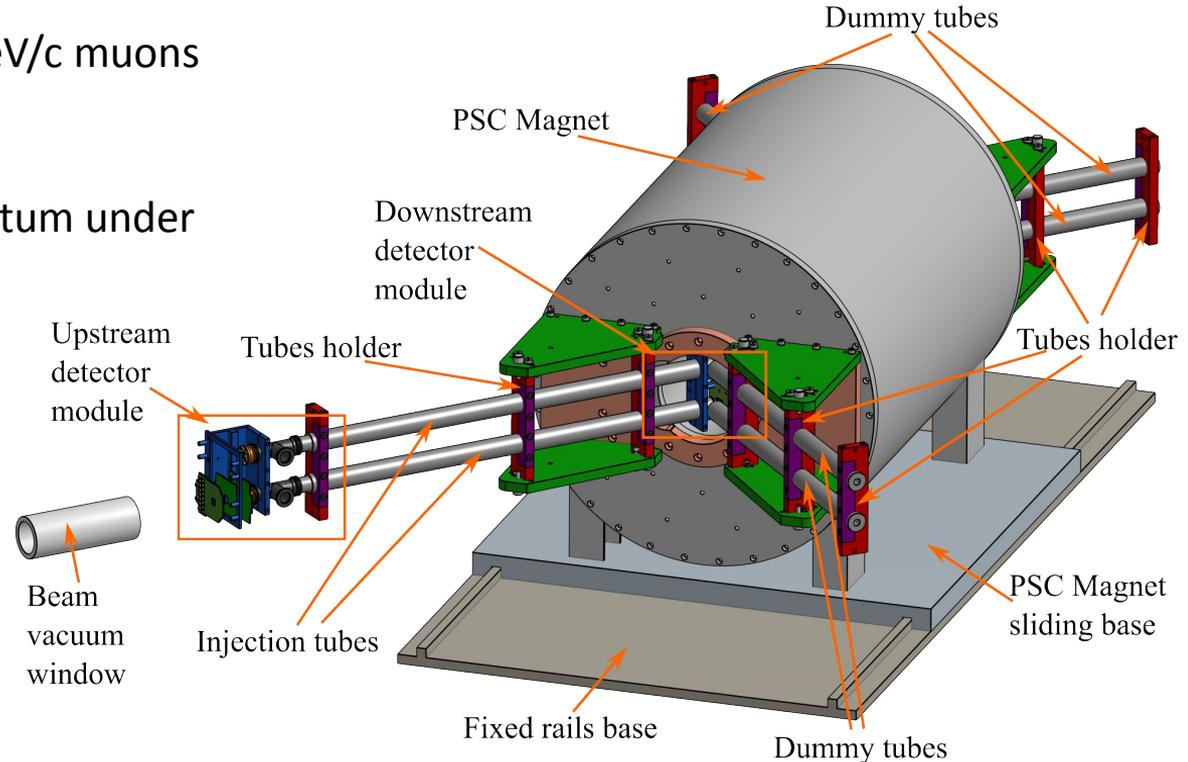
Method	Phase I	Phase II
	FoM	FoM
Simple	0.17	0.17
T-method	0.22	0.18
W-method (20x20x20bins)	0.29	0.28

- The T-method is considered good enough for phase 1, but not for phase 2.
- The W-method is considered for phase 2.
  - Strong requirements for the detectors (momentum resolution, and good tracking)

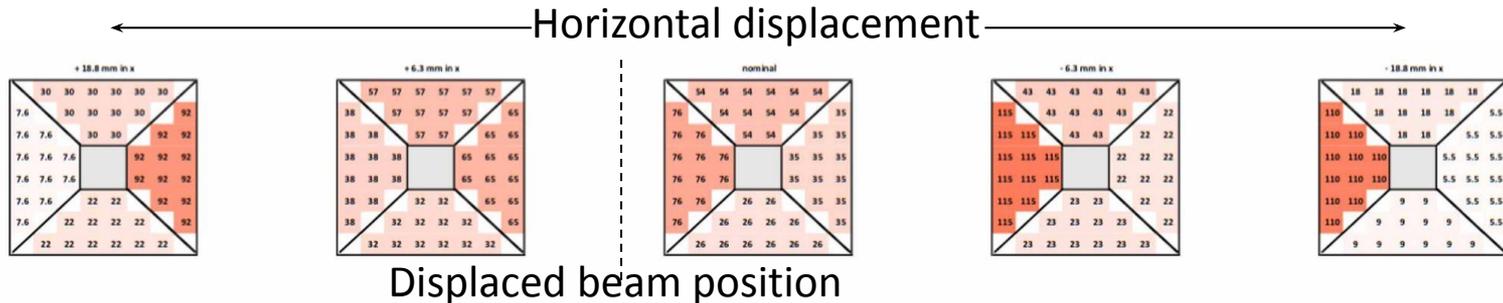
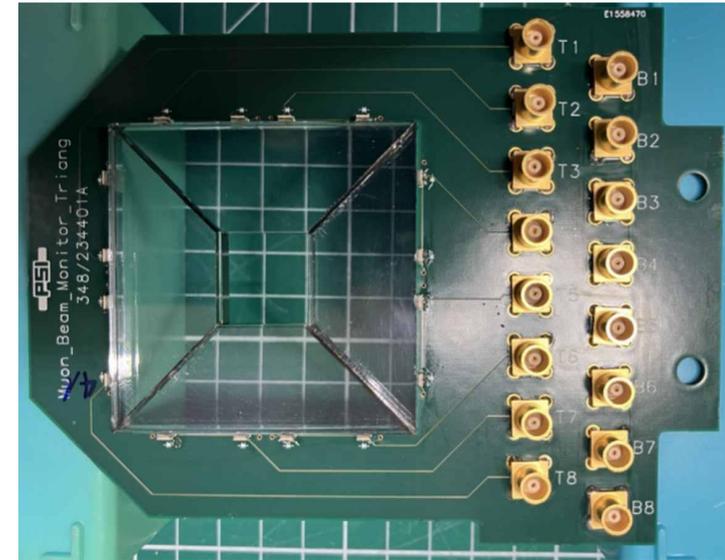


# LATEST PROGRESS

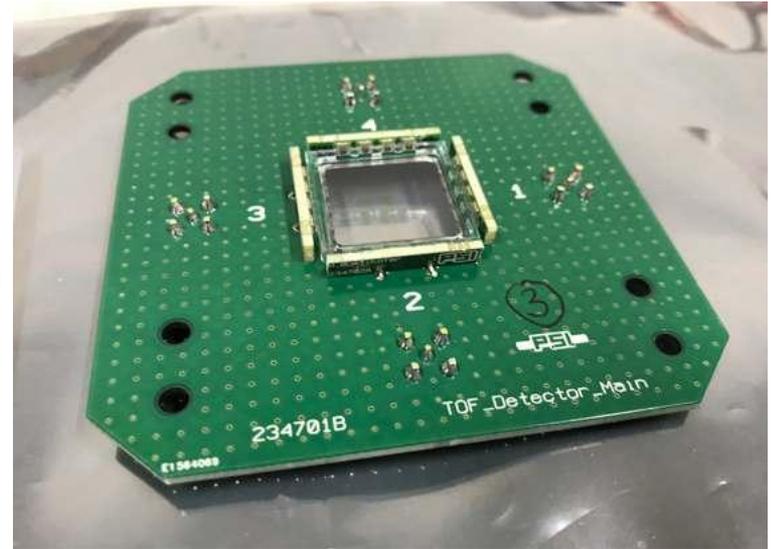
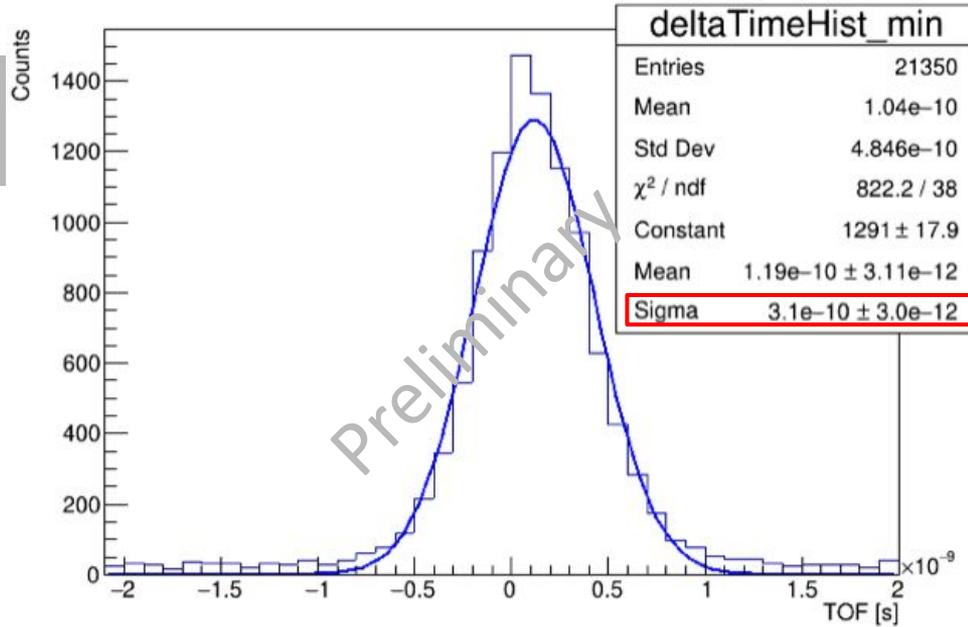
- Test Beam-Monitor detector.
- Test Entrance detectors.
- Test the injection of 28 MeV/c muons at low B fields.
- Test control of the momentum under external changes.
  - Systematic effect



- Objective:
  - Align the beam with the injection channel.
  - Monitor the stability of the beam (e.g. during change of B-field direction).
- Tests:
  - Performance of different detector geometries.
  - Stability under external changes (e.g. position)
  - Sensitivity to beam displacement.



# Latest Progress - Entrance Detector

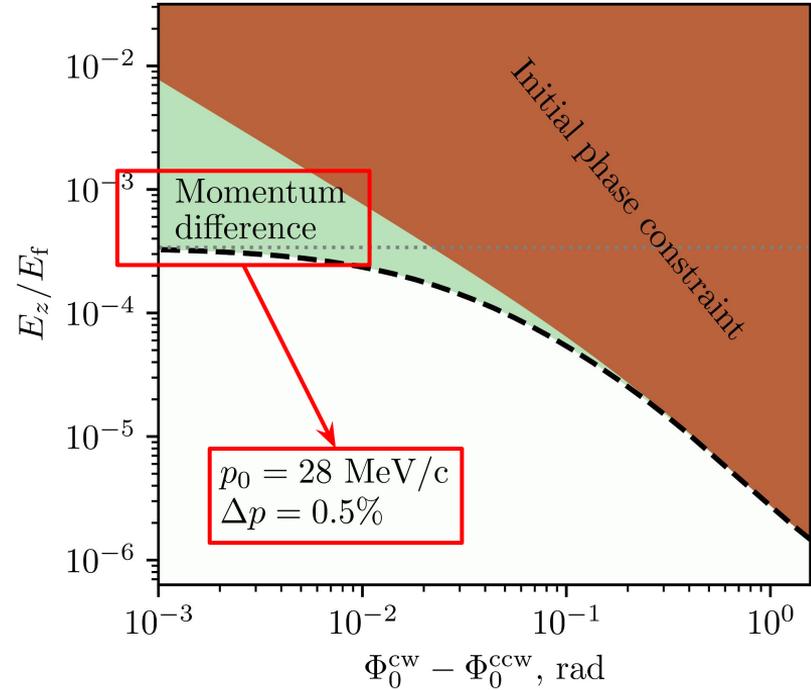


- Benchmarked different scintillators for future use as trigger for the magnetic kick.
- Obtained timing resolution on individual muons of  $\sim 300$  ps.

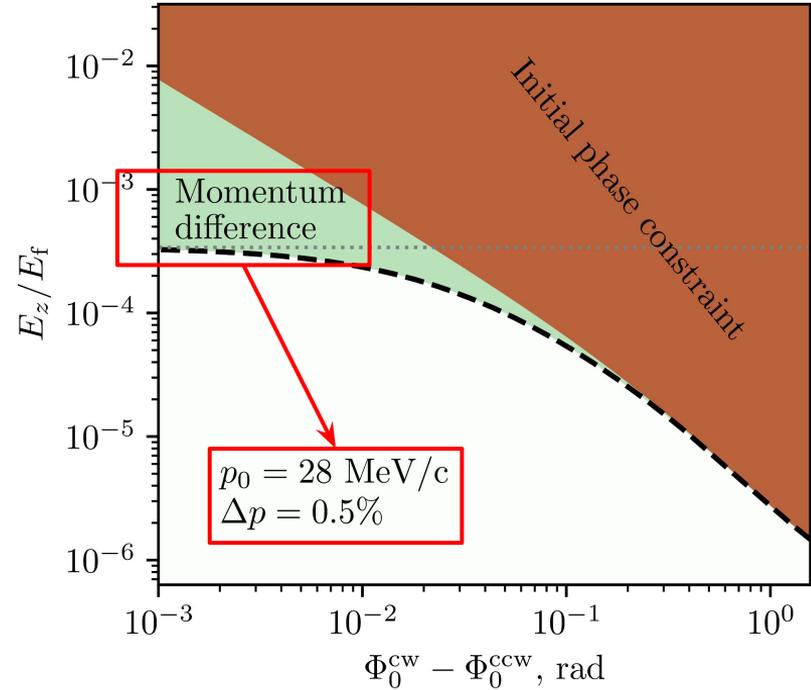
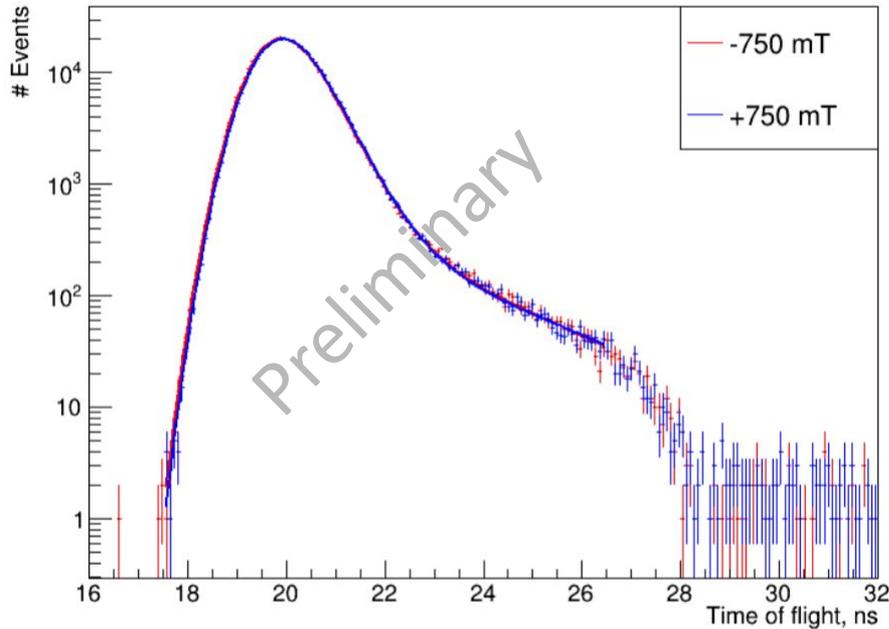
- False EDM-like precession ( $\Omega \propto \beta \times E$ ) due to non-zero longitudinal electric field ( $E_z$ ):

$$\Omega_{\rho}^{E_z} = -\frac{e}{mc} \left( a - \frac{1}{\gamma^2 - 1} \right) \beta_{\theta} E_z,$$

- One of the criteria to cancel the false signal:
  - Momentum distribution for clockwise (CW) and anticlockwise (CCW) injection should be “equal”.
  - CW and CCW injections achieved by inverting the B field.



ToF histogram DS: 100  $\mu$ m US: 50  $\mu$ m



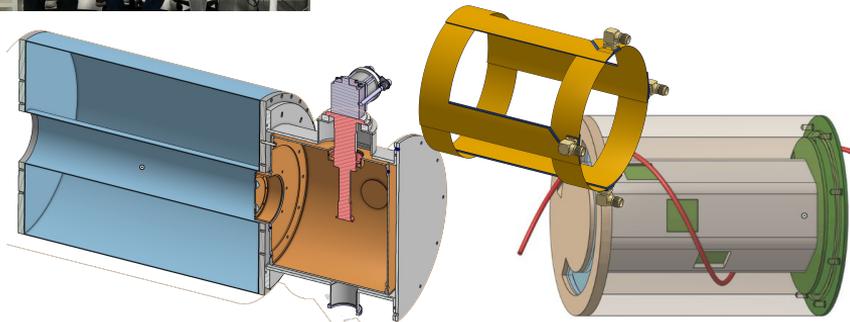
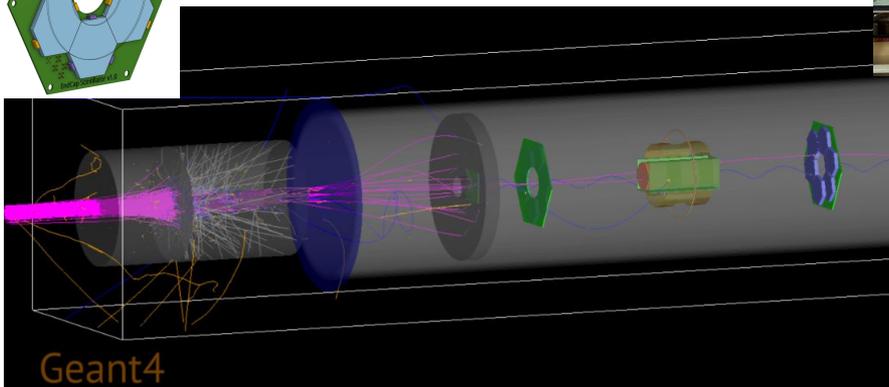
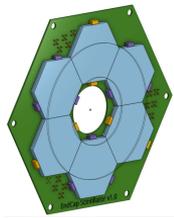
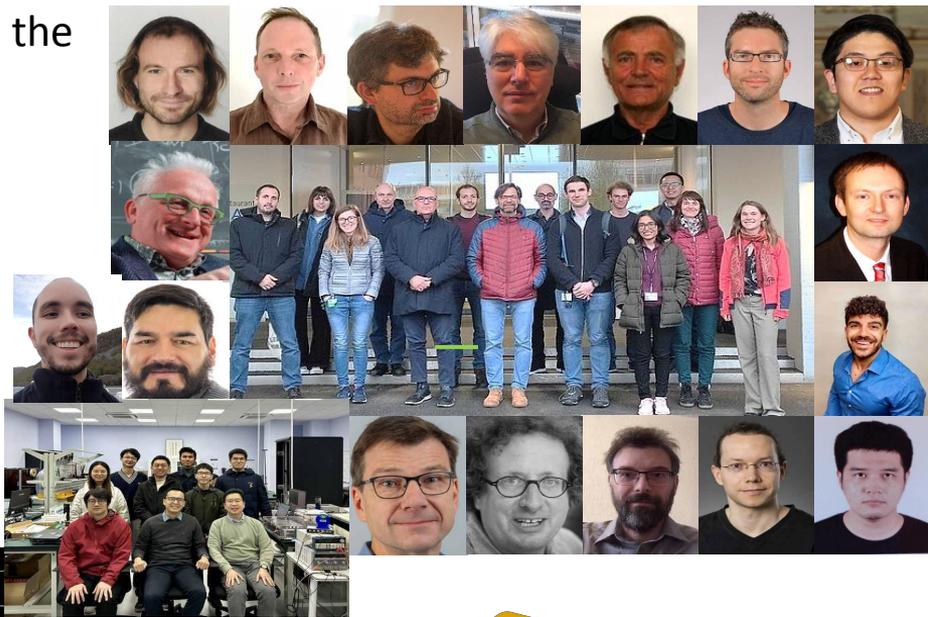
- Agreement of the ToF spectra for positive/negative B field configurations within 0.2%.
- Momentum control better than 0.5%: Necessary condition to avoid false signal from  $E_z \neq 0$

# STATUS AND OUTLOOK

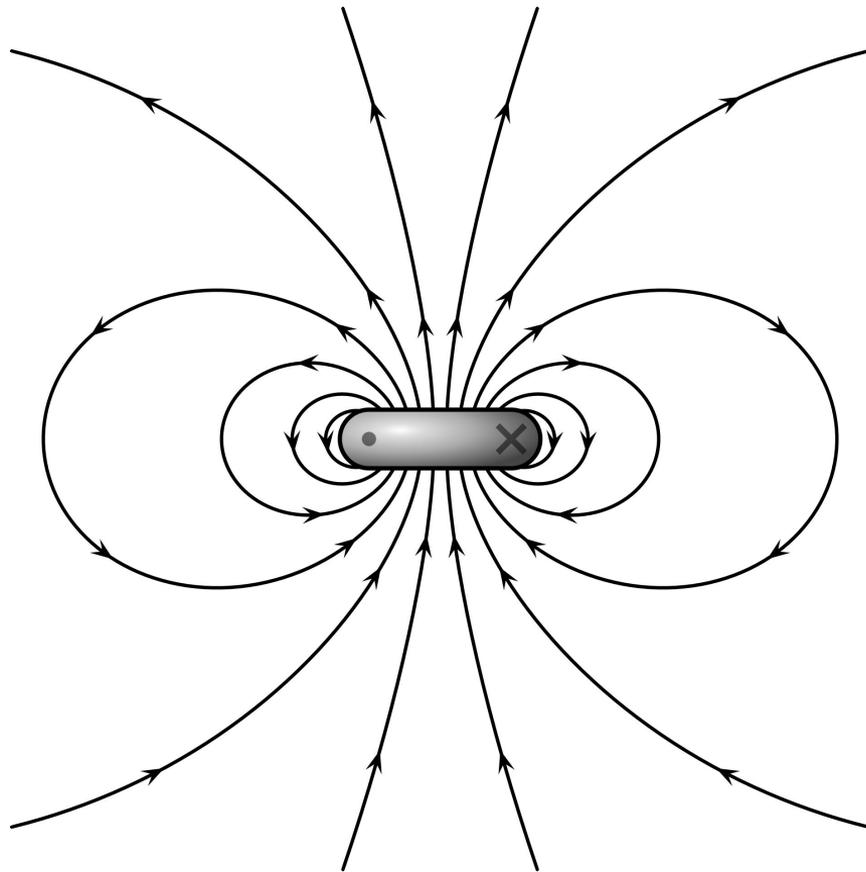
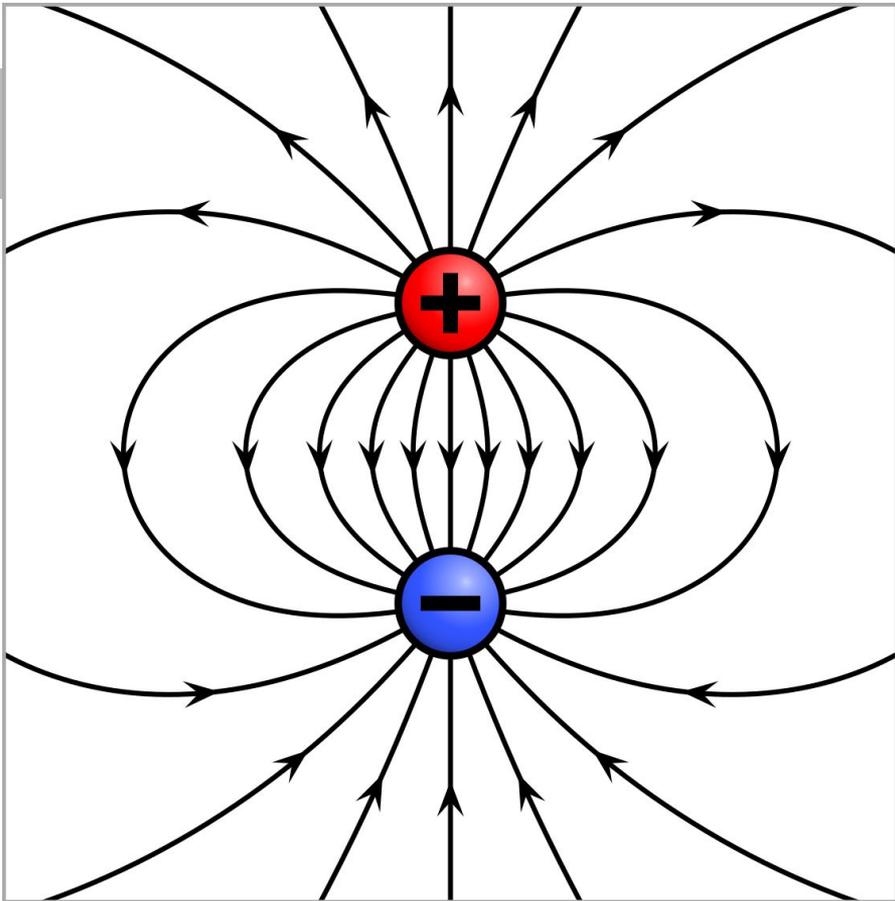


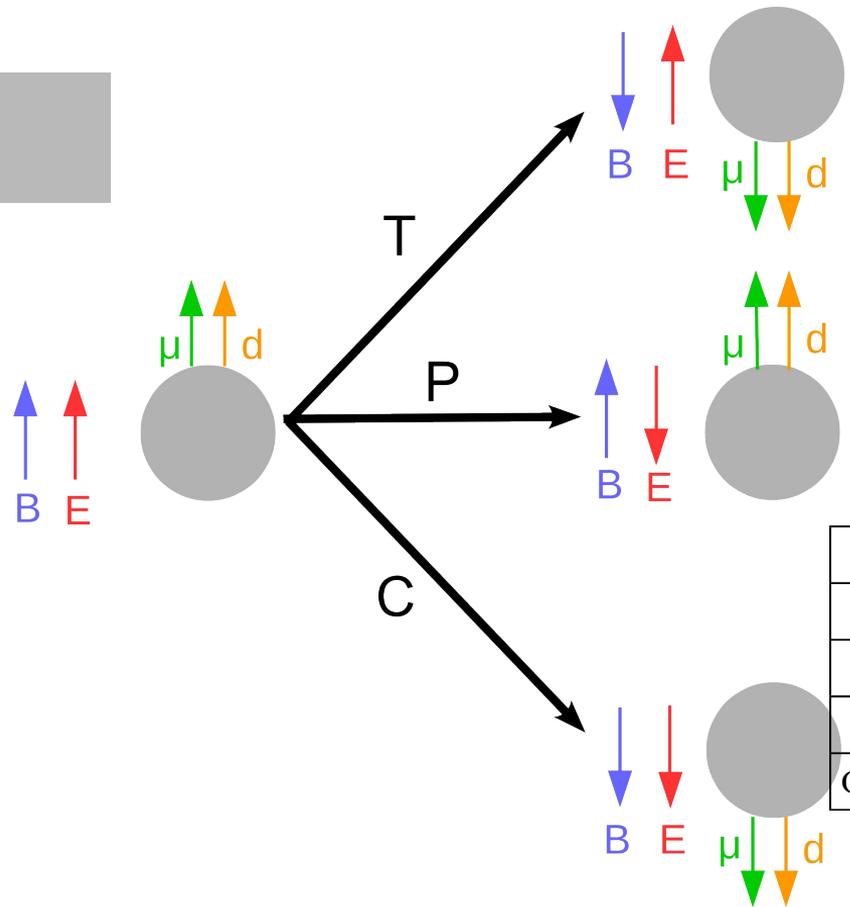
# Outlook

- Study possible effects of magnetic kick on the positron detection.
- Characterize muons trajectory.
- Test cryogenics and SC tubes.
- Precise mapping of the magnet's B field.
- Stop muons in orbit with magnetic kicker.



# Backup

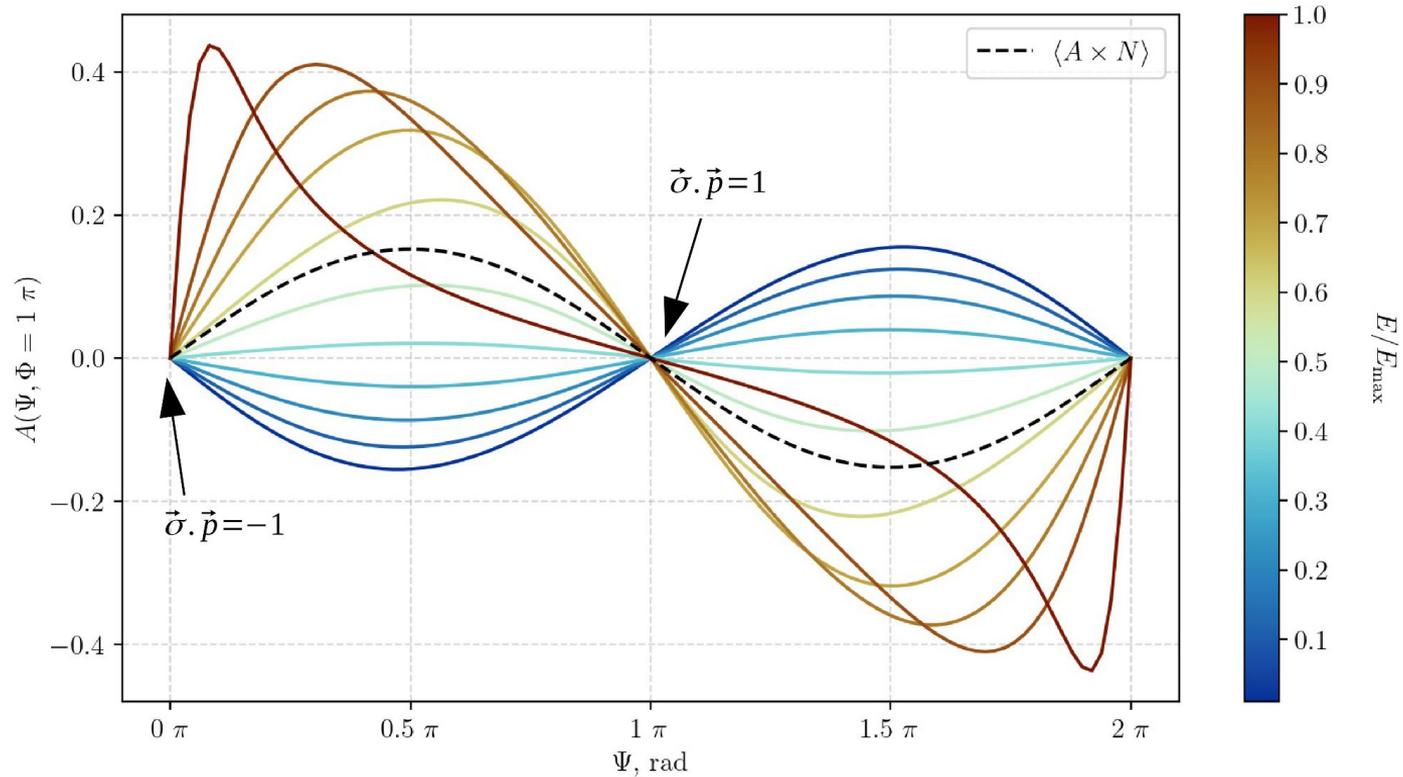


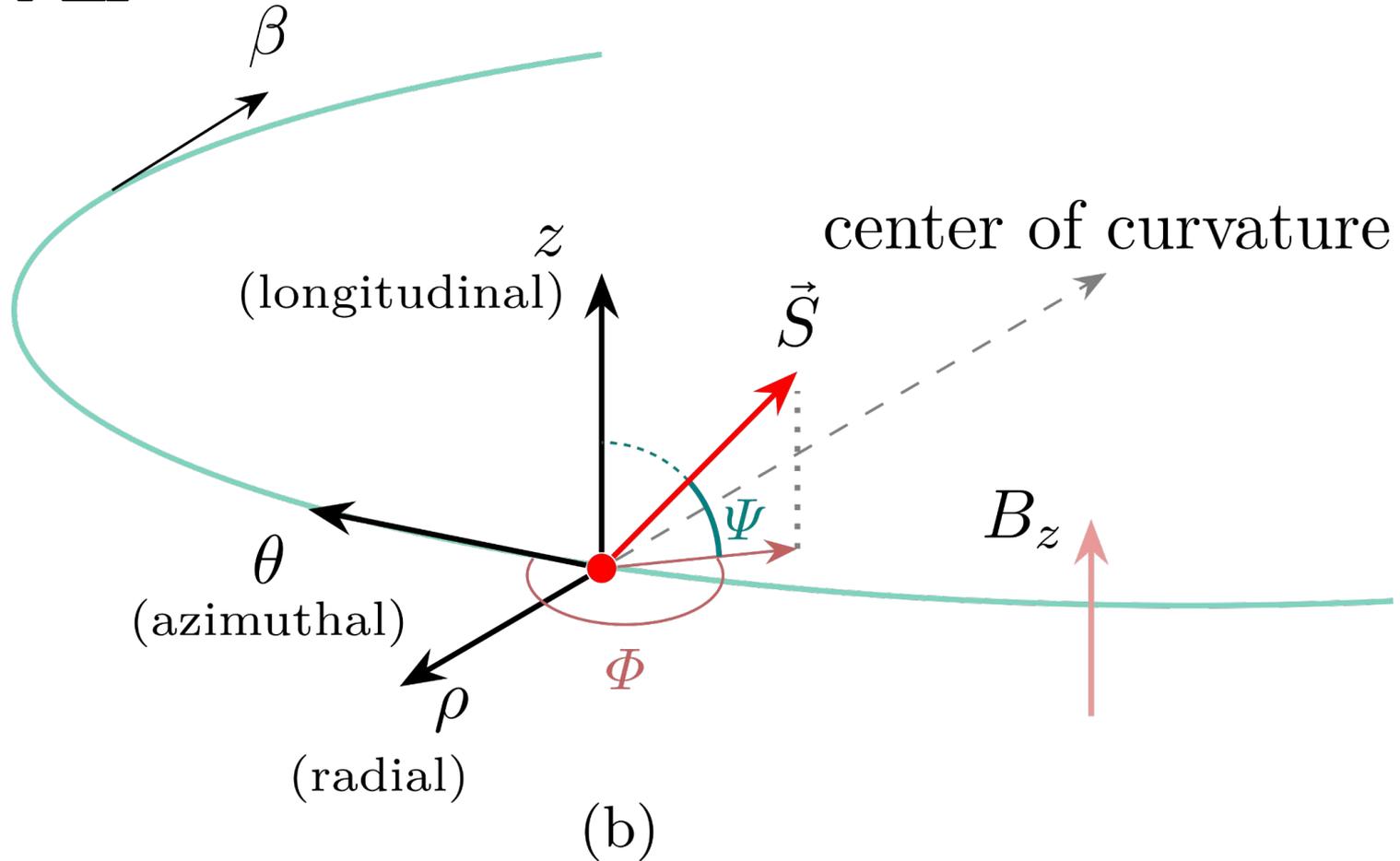


$$\hat{H} = -\vec{\mu} \cdot \vec{B} - \vec{d} \cdot \vec{E}$$

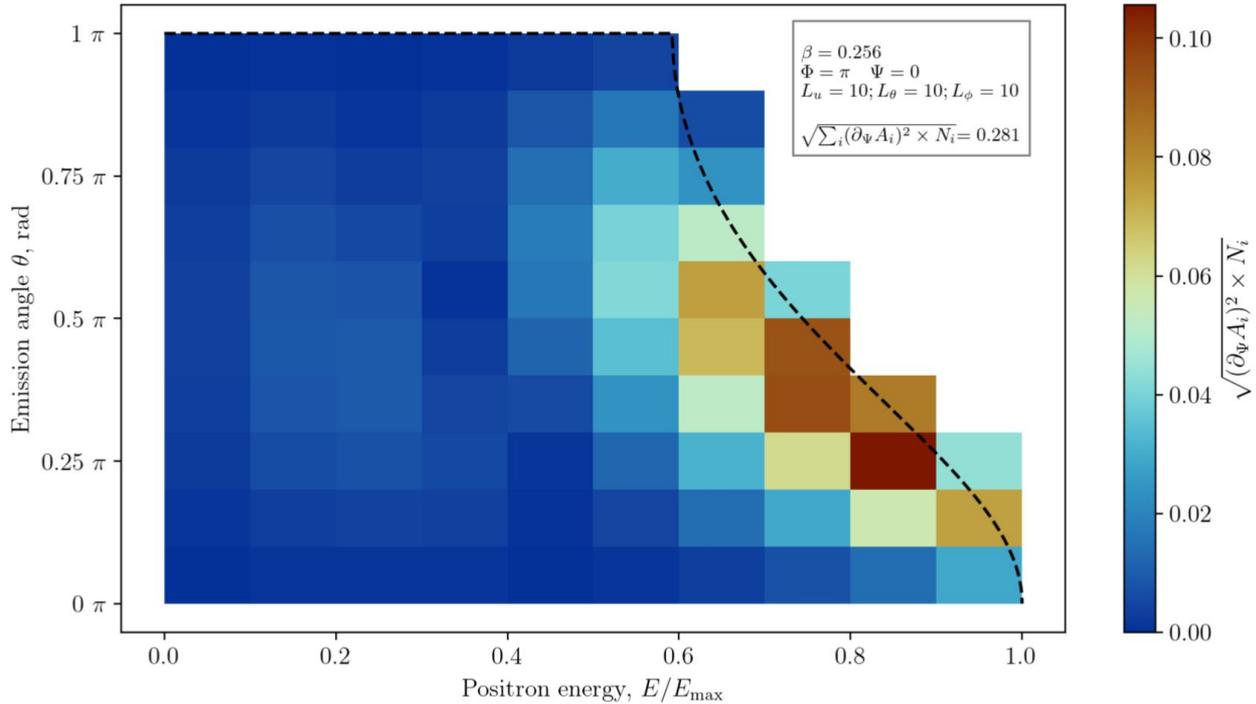
$$\vec{\mu} = g \frac{q}{2 m_{\mu}} \vec{s} \quad \vec{d} = \eta \frac{q}{2 m_{\mu} c} \vec{s}$$

	$\mu$	$d$	$s$	$\mathbf{E}$	$\mathbf{B}$	$-d \mathbf{s} \cdot \mathbf{E}$	$-\mu \mathbf{s} \cdot \mathbf{B}$
parity	+	+	+	-	+	(+ + -) = -	(+ + +) = +
time	+	+	-	+	-	(+ - +) = -	(+ - -) = +
charge	-	-	+	-	-	(- + -) = +	(- + -) = +
charge & parity	-	-	+	+	-	(- + +) = -	(- + -) = +





# Minimize The Uncertainty - Sensitivity

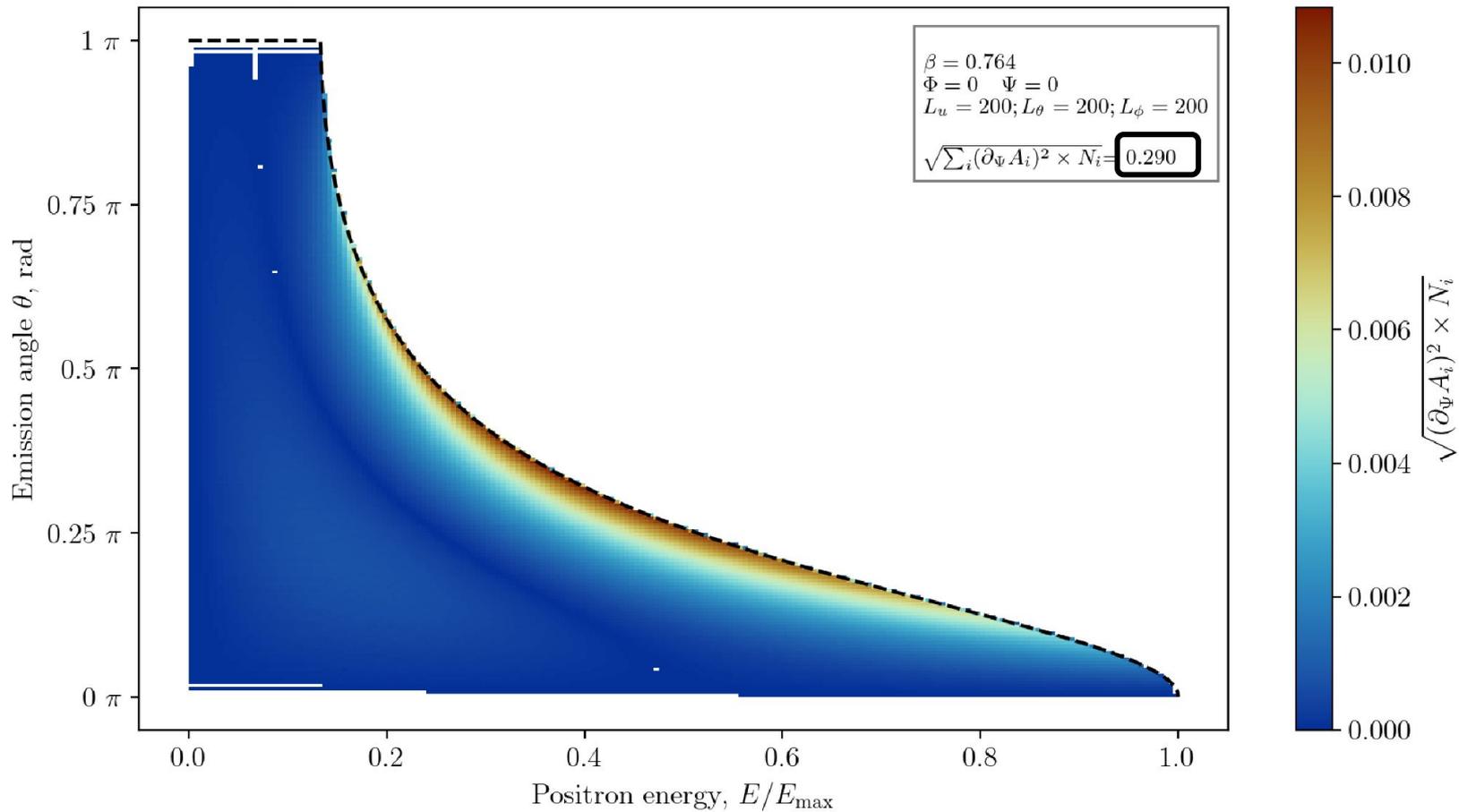


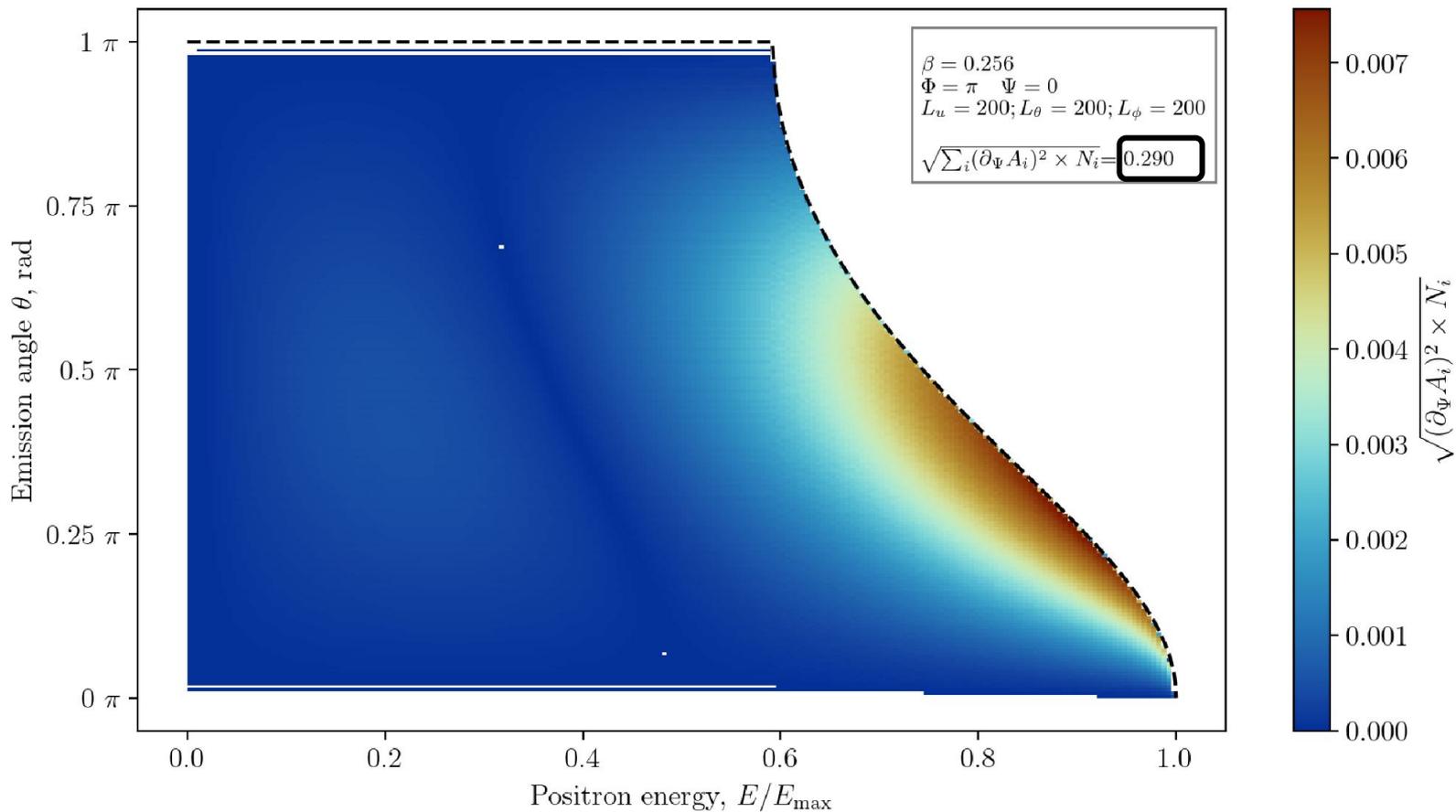
$$W_i = \alpha_i \sqrt{N_i}$$

$$W = \sqrt{\sum_i W_i^2} = \sqrt{\sum_i \alpha_i^2 N_i}$$

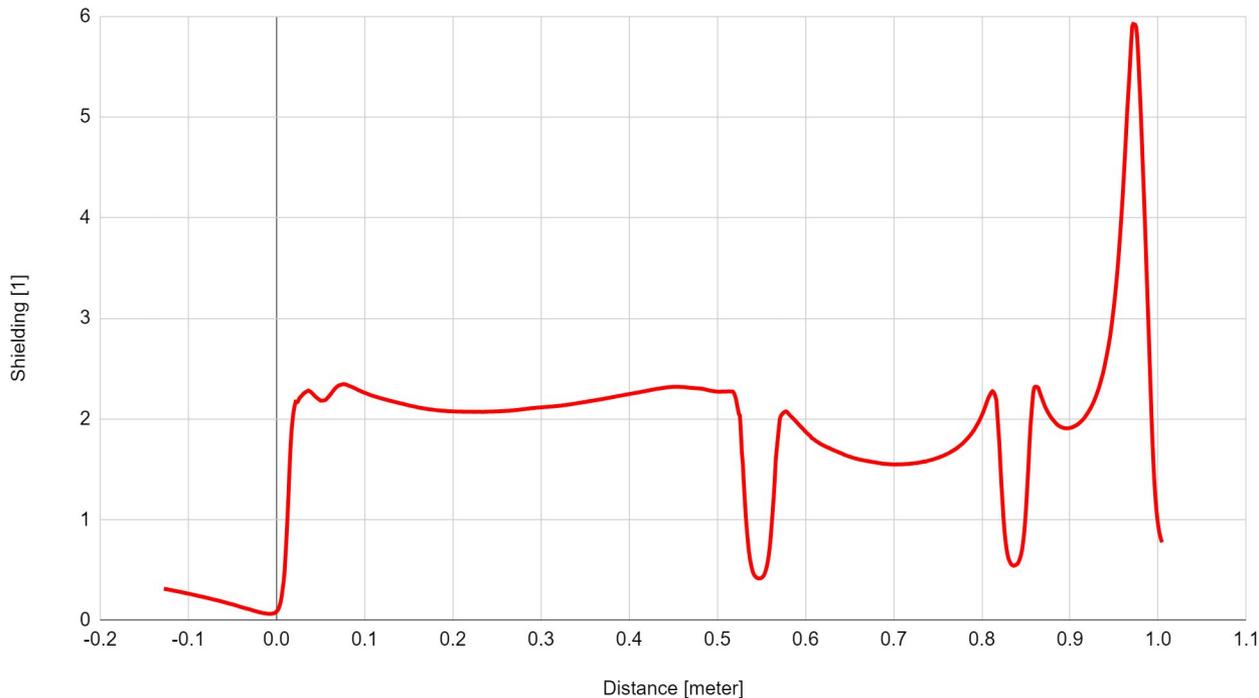
$$d_\mu = \frac{\hbar}{2 c \beta B P} \frac{\sum_i \dot{A}_i W_i^2}{\sum_i W_i^2}$$

- Binning in energy and direction.
- The W-method is preferred, which imposes strong requirements for tracking and momentum discrimination of the detectors.

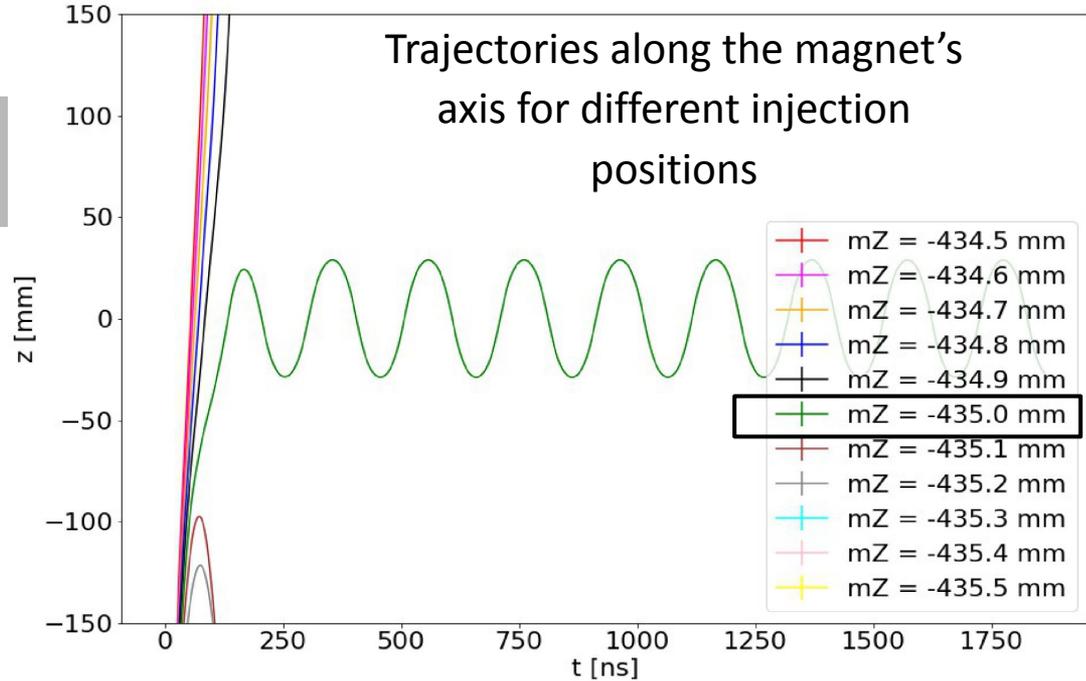




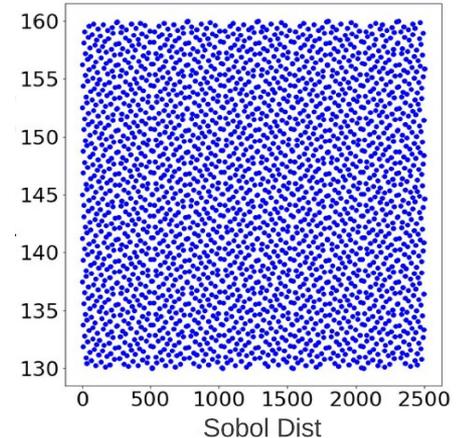
Shielding [1] vs. Distance [meter]



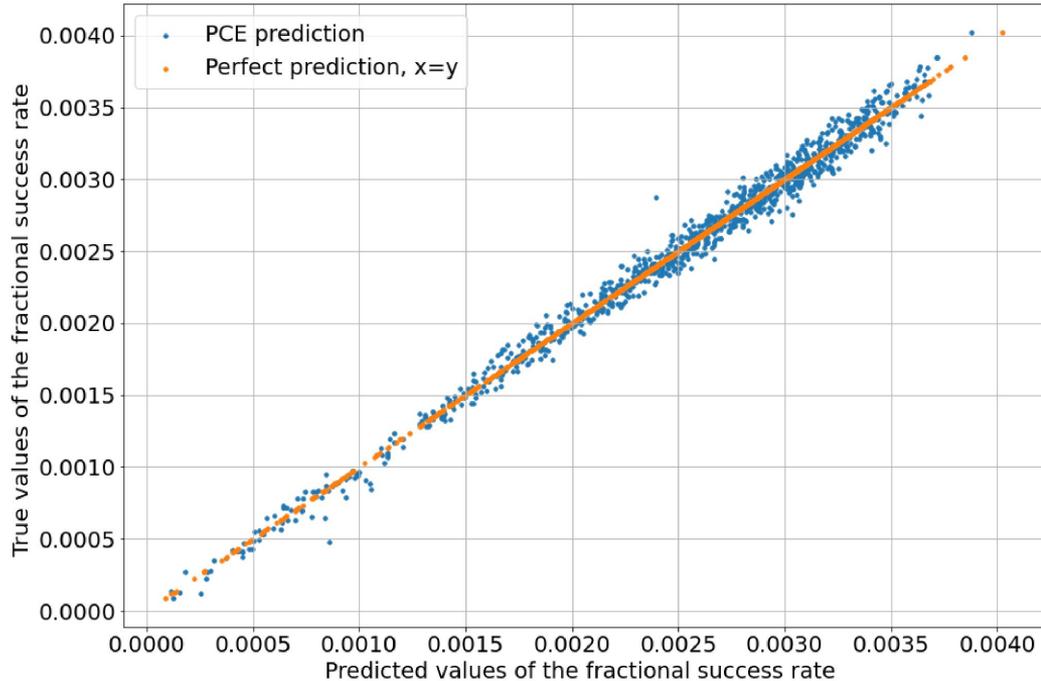
- Operate magnet at a field high enough without saturating the injection tubes
- Shield the injected muons to below  $\sim 100$  mT



- Highly restrictive parameter space for successful injection of muons.
- Need for efficient methods to sample the parameter space and generate estimations.



# Optimize The Experiment - Preliminary Result

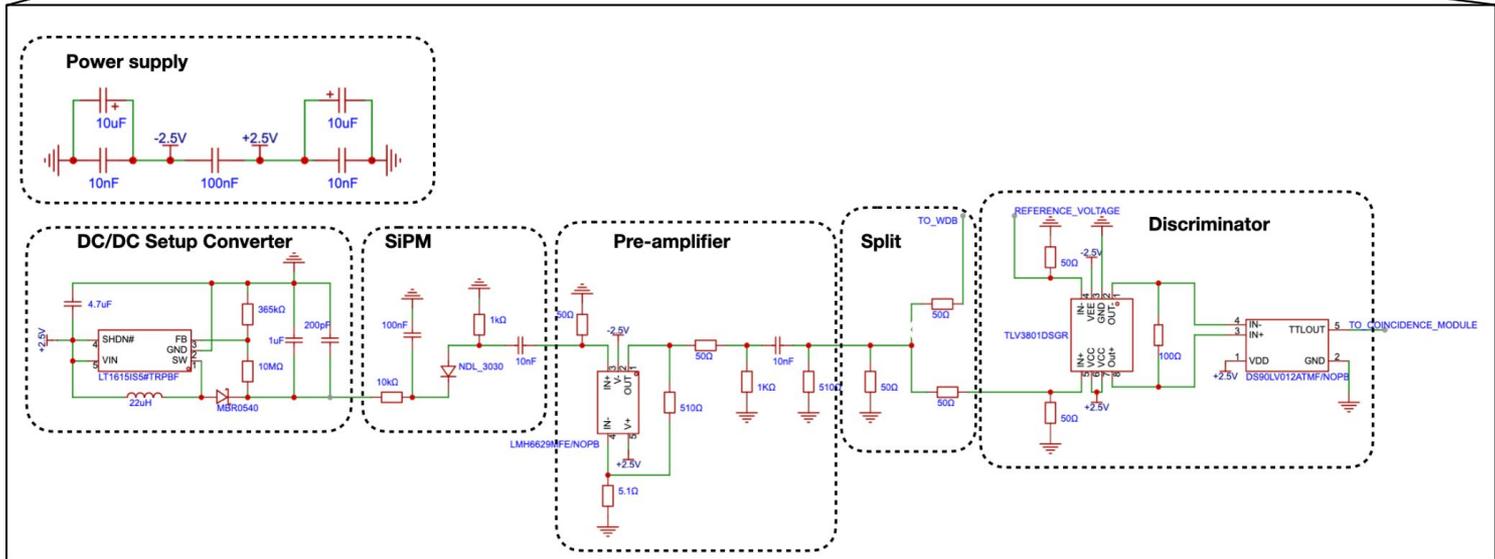
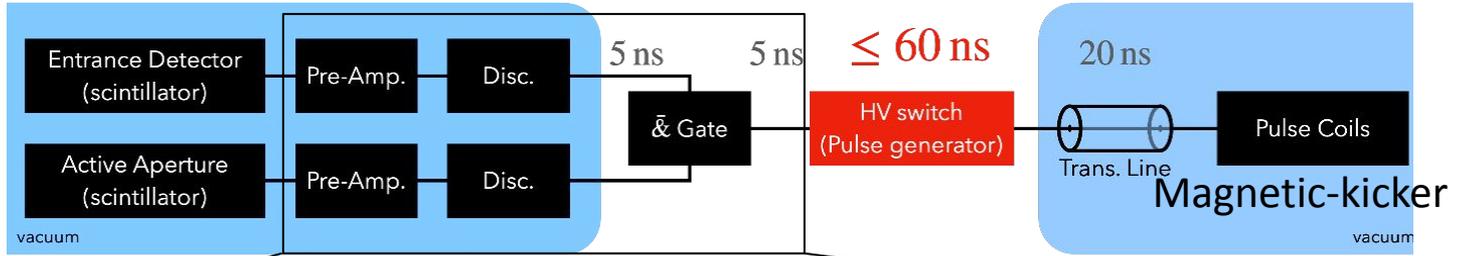


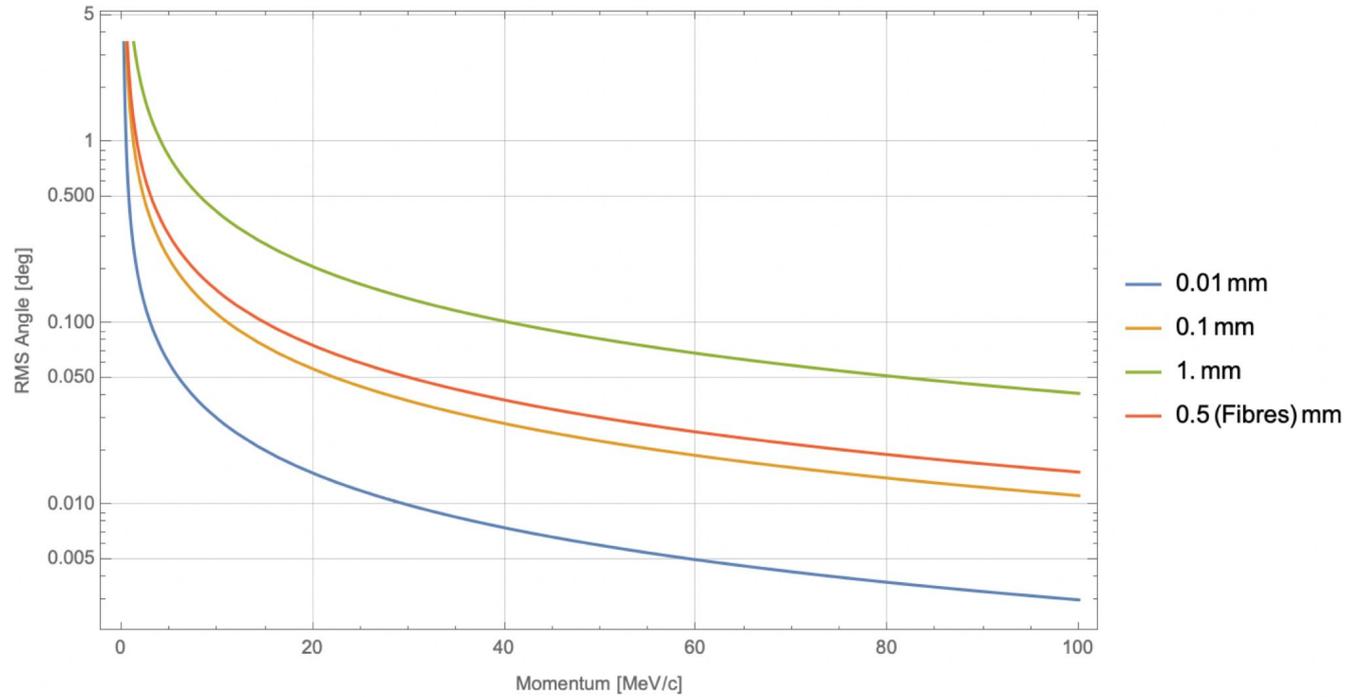
- 6<sup>th</sup> degree polynomial-chaos-expansion surrogate model performs well for predicting the efficiency.
- Using optimization based in Genetic-Algorithm, found experiment parameters for storing 0.4 % of the injected muons (0.47% with G4beamline)

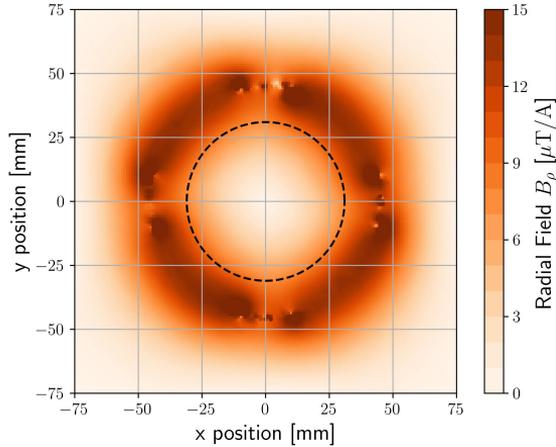
# Fast Electronics Development



15 ns



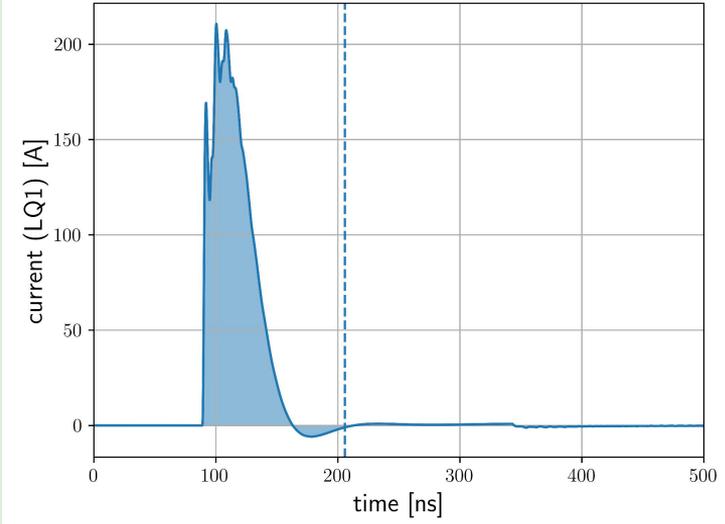
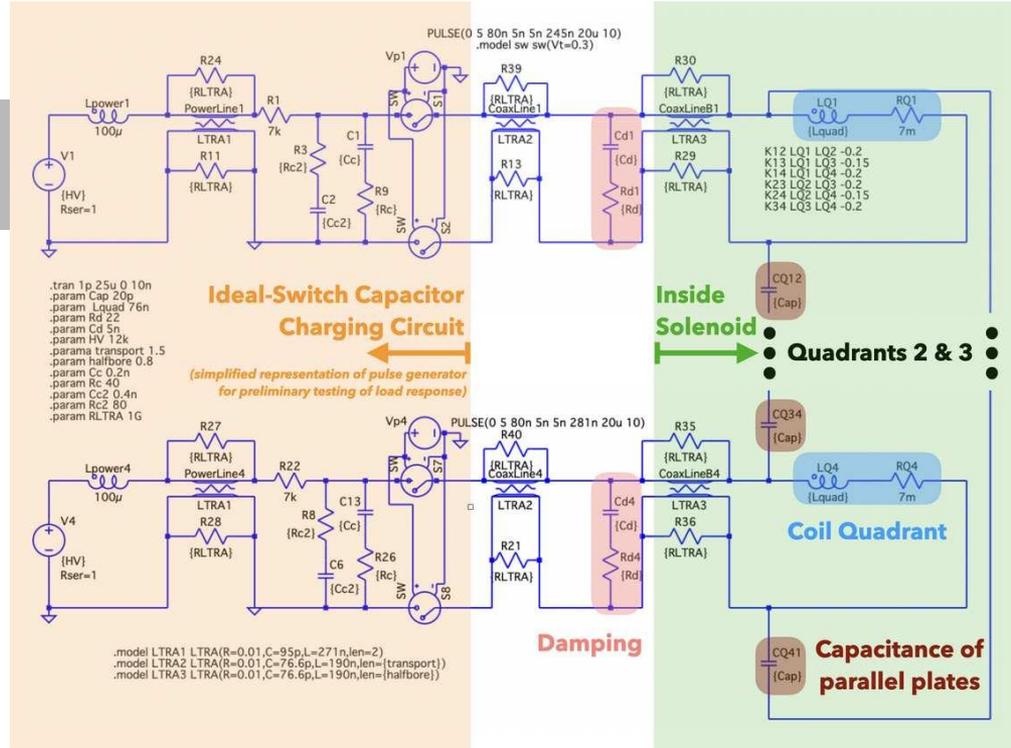




$$\frac{dp_\phi}{dz} = -\frac{e\rho}{2} \frac{dB}{dz}$$

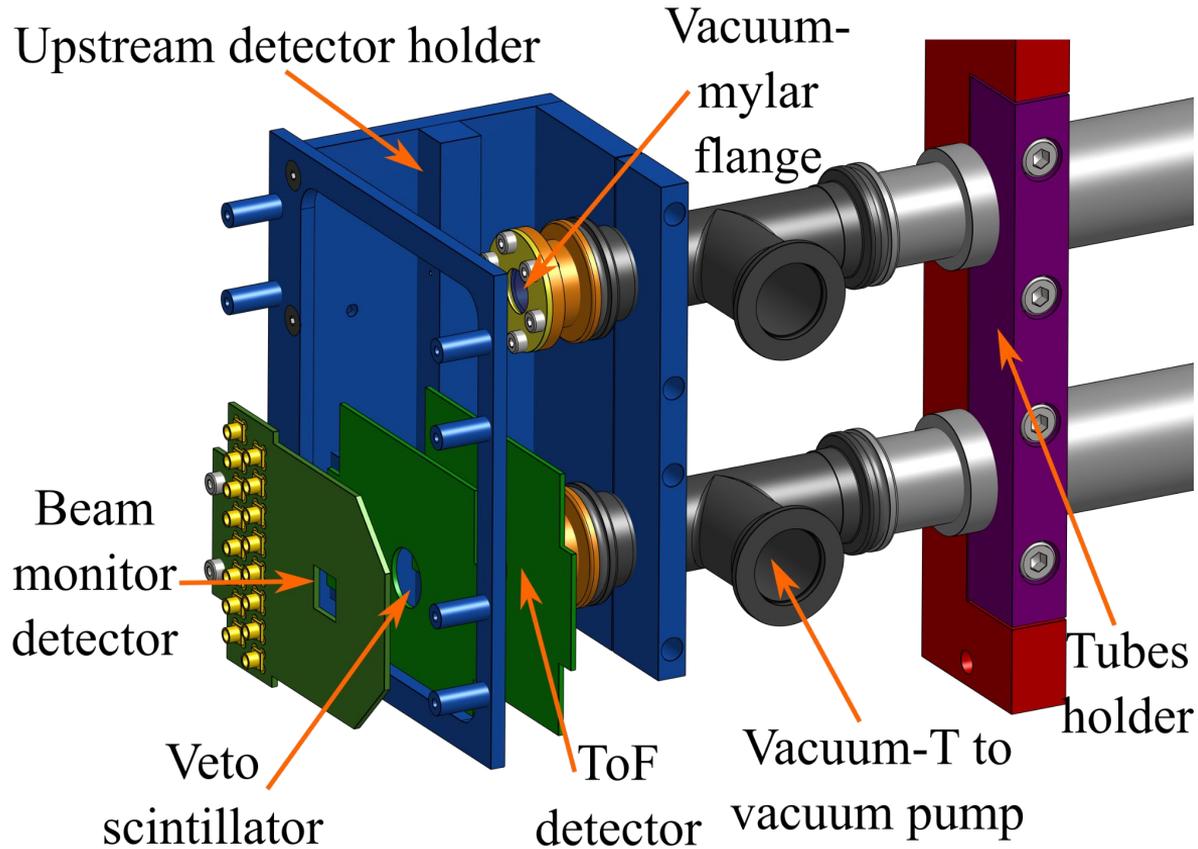
$$\dot{\vec{p}}(t) = \frac{e}{\gamma m} \vec{p} \times \vec{B}_{\text{pulse}}(t)$$

$$\implies \dot{p}_z(t) = \frac{e}{\gamma m} p_\phi \vec{B}_{\text{pulse}}(t) \cdot \hat{\rho}$$

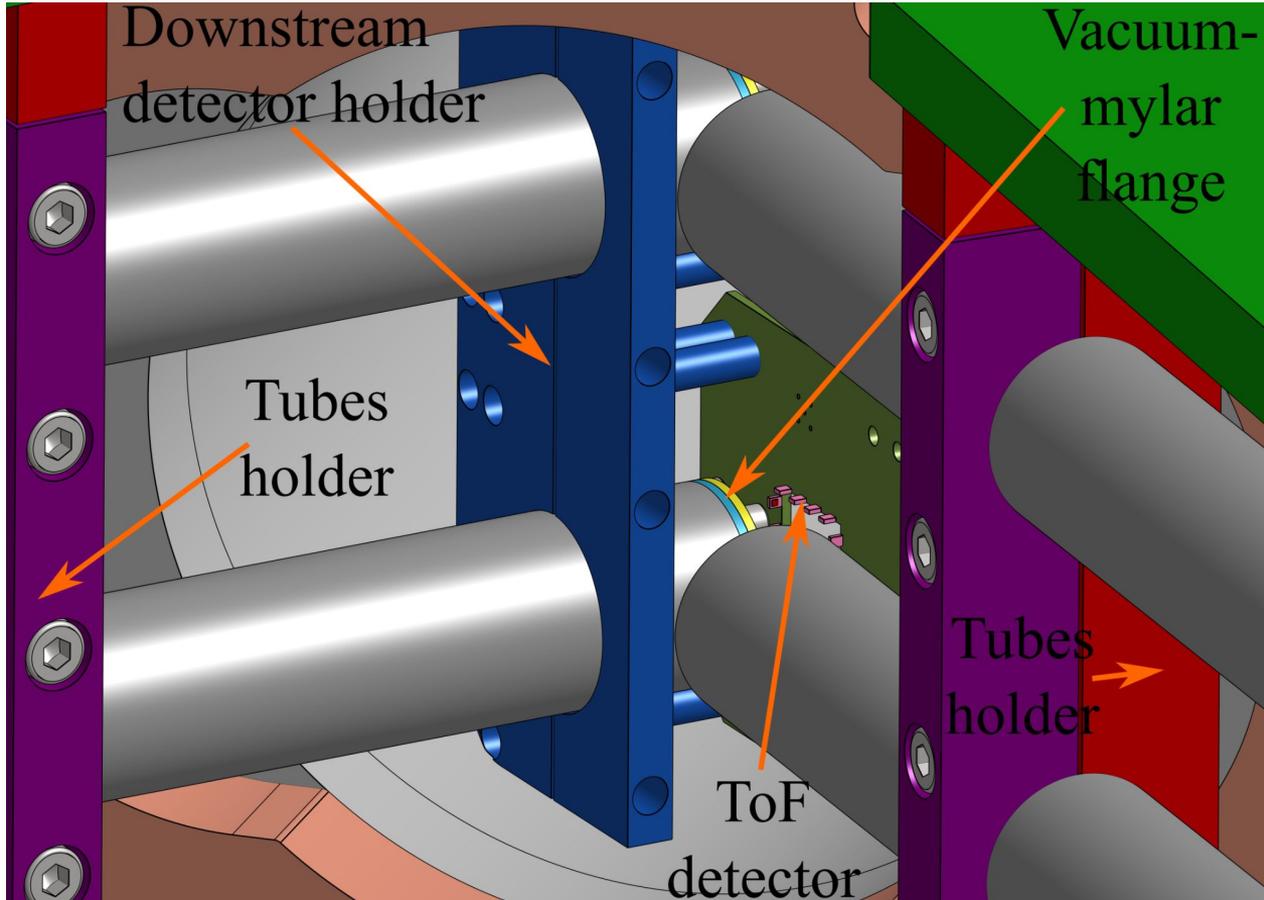


- 200 A - 35 ns FWHM pulse.
- Prototype under development for testing in summer.

# Latest Progress - Test Beam 2023



# Latest Progress - Test Beam 2023



- Study possible effects of magnetic kick on the positron detection.
- Stop 200 MeV/c pions inside 3 T field.
  - Uniform distribution of positrons from muons at rest.
  - Hexagonal-tiling scintillating detectors before and after the stopping target.

