On the theory interpretation of upcoming **EECs measurements**





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Scaling argument for EECs: PYTHIA simulation



Minimum angle accessed perturbatively $\sim \lambda_{\rm QCD}/p_{t,\rm jet}$







Medv JHEP 04



Connection between EEC's peak and infrared scales in MC



Comment on Barbara's talk: peak-position at parton level \neq hadron one





Connection between EEC's peak and infrared scales in MC

[M. Leitão, G. Milhano and ASO, in preparation]



EEC peak independent of shower cutoff (as it should). It shifts when decreasing string tension \Rightarrow later hadronization





Flavour-dependence of hadronization time



On the theory interpretation of upcoming **EECs measurements**



Alba Soto-Ontoso

New jet quenching tools to explore equilibrium and non equilibrium dynamics in heavy-ion collisions

ECT*, 13th February, 2024





Some (not serious) crystal-balling on EEC (AA) / EEC (pp)





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Some (not serious) crystal-balling on EEC (AA) / EEC (pp)



Some (more serious) considerations

- unfortunately, out-of-the-scope.

- #GoldenChannel, #PinDownCriticalAngle, #WakeDiscovery,...

[Barata, Caucal, ASO and Szafron arXiv:2312.12527] [Barata, Caucal, Monni, ASO, Szafron work in progress]

Precise (à la pp) theoretical calculations of jet observables in heavy-ions are,

 Recent proposals for new observables have been based on over-simplistic (mainly leading-order) predictions \Rightarrow disappointment once measurement is out.

• We (as a community) have a new opportunity to do things better with EECs.

• Keywords to avoid without a systematic study: #SmokingGun, #Unravel,

Rest of this talk: point out a few effects beyond the leading-order picture

s despite energy weight in the EEC

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Ratio

Leading-logarithmic calculation of the EEC in vacuum

The cumulative distribution for the abelian EEC in the collinear limit at $O(\alpha_s)$ is

$$\sum_{q=1}^{\infty} \Sigma^{R}(\chi) \Big|_{\mathcal{O}(\alpha_{s})} = \bar{\alpha} \int_{0}^{1} \frac{d\theta}{\theta} \int_{0}^{1} dz P_{gq}(z) \Big[z^{2} \Theta(\chi) + (1-z)^{2} \Theta(\chi) + 2z(1-z) \Theta(\chi-\theta) \Big]_{\mathcal{O}(\alpha_{s})}$$

$$\sum_{q=1}^{\infty} \Sigma^{V}(\chi) \Big|_{\mathcal{O}(\alpha_{s})} = -\bar{\alpha} \int_{0}^{1} \frac{d\theta}{\theta} \int_{0}^{1} dz P_{gq}(z) \Theta(\chi)$$

Iterating this procedure one finds the all-orders result at leading-log accuracy

$$\Sigma(\chi) = \sum_{k=0}^{\infty} (-1)^{k+1} \frac{\bar{\alpha}^{k+1}}{(k+1)!} \ln^{k+1} \frac{1}{\chi} [\gamma_{qq}^{k+1}(3) + \gamma_{qq}^{k}(3)\gamma_{gq}(3)] = \left(-1 + \chi^{\bar{\alpha}\gamma_{qq}(3)}\right) \frac{\gamma_{gq}(3) + \gamma_{qq}(3)}{\gamma_{qq}(3)}$$
where $\gamma_{ik}(j) = -\int_{0}^{1} dz \, z^{j-1} \hat{P}_{ik}(z)$ are the so-called anomalous dimensions
 ζ_{j} regularized

Effect #1: impact of phase-space constraints on VLEs

The phase-space for vacuum-like emissions affects anomalous dimensions

[Caucal et al. PRL 120 (2018) 232001]

$$\gamma^{\text{med}}(j,\theta) = -\int_0^1 dz \, z^{j-1} \hat{P}(z) [1 - \Theta_{\text{veto}}(z,\theta)]$$

Effect #1: impact of phase-space constraints on VLEs

The phase-space for vacuum-like emissions affects anomalous dimensions

[Caucal et al. PRL 120 (2018) 232001]

Effect #2: in-medium matrix element

The in-medium EEC at leading order ($\mathcal{O}(\alpha_s)$) can be expressed as:

$$\frac{d\Sigma}{d\theta dp_t} = \sum_{\{i,j\}} \int_0^1 dz \left[z(1-z) \frac{dP_{ij}^{\text{vac}}}{d\theta dz} \left(1 + F_{\text{med}}^{ij}(\theta, z) \right) \right] \frac{d\sigma_j}{\sigma_j dp_t}$$

Available options for the $1 \rightarrow 2$ splitting function (F_{med}):

- Semi-classical approx: see Jack and Fabio's talks. Validity: $z \sim 1/2$ [Dominguez, Milhano, Salgado, Vila EPJC 80 (2020) 1, 11]
- Soft approximation: BDMPS-Z spectrum. Validity: $z \rightarrow 0$ [BDMPS, NPB 483 (1997) 291-320]

• Exact (with all caveats that Fabio explained): only available for $\gamma
ightarrow q ar{q}$ [Isaksen, Tywoniuk JHEP 09 (2023) 049]

Simple question: Does the $z \rightarrow 0$ matter for the evaluation of the EEC?

Effect #2: in-medium matrix element ($q \rightarrow qg$)

Simple answer: The $z \rightarrow 0$ does matter for the evaluation of the EEC

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Effect #2: in-medium matrix element ($\gamma \rightarrow q\bar{q}$)

Wide-angle enhancement substantially reduced when using exact ME. Another indication of the importance of describing full-z dependence.

Effect #3: better description of balanced splittings in JetMed

We modified the description of the **first** emission inside the medium

 $\frac{\alpha_s(k_{\perp})}{\pi} P_{ij}(z) dz \frac{dt}{\theta}$

with $F_{med}(z, \theta)$ an overestimate of the semi-classical matrix-element

- + Angular-ordered vacuum shower

+ Gaussian transverse momentum broadening [See Paul's talk]

$$z \frac{d\theta}{\theta} \left[1 + F_{\text{med}}(z,\theta) \right]$$

+ Phase-space constraints for vacuum-like emissions

+ BDMPS-Z cascade of medium induced emissions

Effect #3: better description of balanced splittings in JetMed

Energy-loss dominates over medium-induced emissions, i.e. narrowing

Instead of conclusions, a new observable! Lund-based EEC

Several reasons to allow for different energy-weights on the EEC, e.g.

- Access higher moments of the splitting function
- Mitigate background contamination
- Possible in the canonical EEC after introducing (non-perturbative) track functions [Barata and Szafron arXiv:2401.04164
- Our approach: minimal modification without sacrificing calculability

$$\frac{1}{\sigma} \sum_{\{i,j\}\in \text{declust.}} \int_{0}^{1} dz \frac{d\sigma}{d\theta_{ij}dz} \frac{E_{i}^{n}E_{j}^{n}}{Q^{2n}} \delta\left(\chi - \frac{\theta_{ij}}{R}\right)$$

