



2p2h implementation in NuWro

Kajetan Niewczas

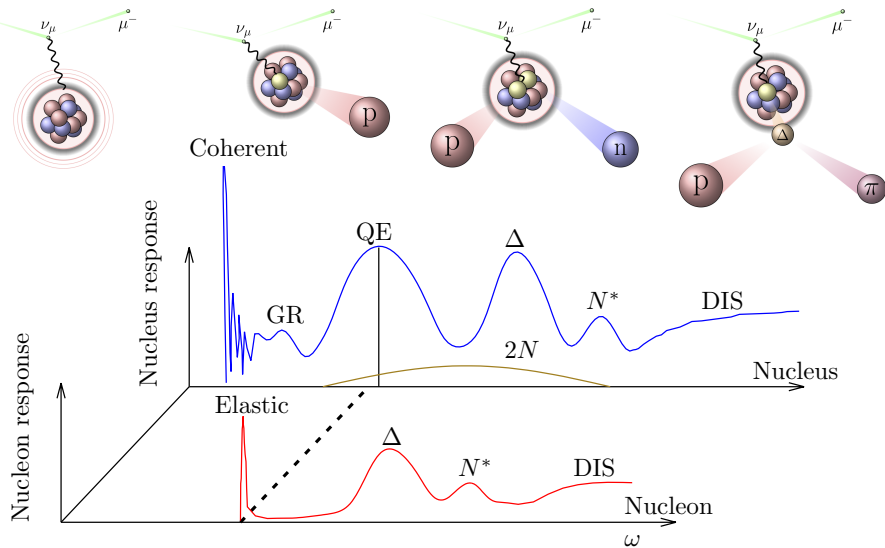


Uniwersytet
Wrocławski



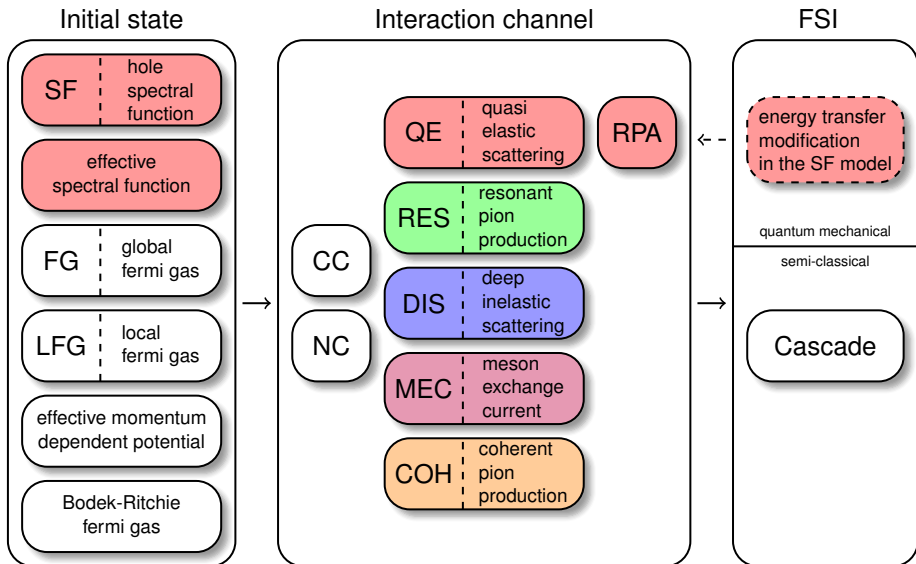
UNIVERSITEIT
GENT

Nuclear response

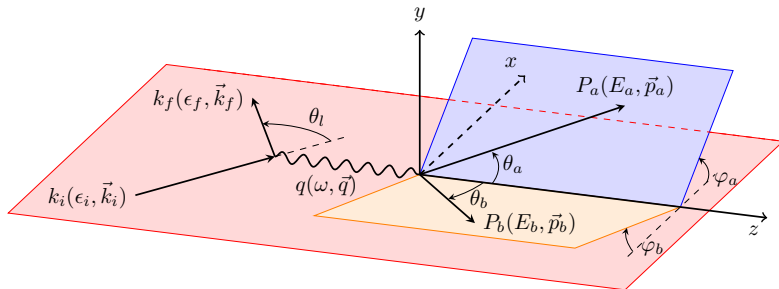


T. Van Cuyck

NuWro blueprint



2p2h cross section



T. Van Cuyck

Exclusive

$$\frac{d^8\sigma}{d\epsilon_f d\Omega_f dT_a d\Omega_a d\Omega_b}$$

→

Semi-inclusive

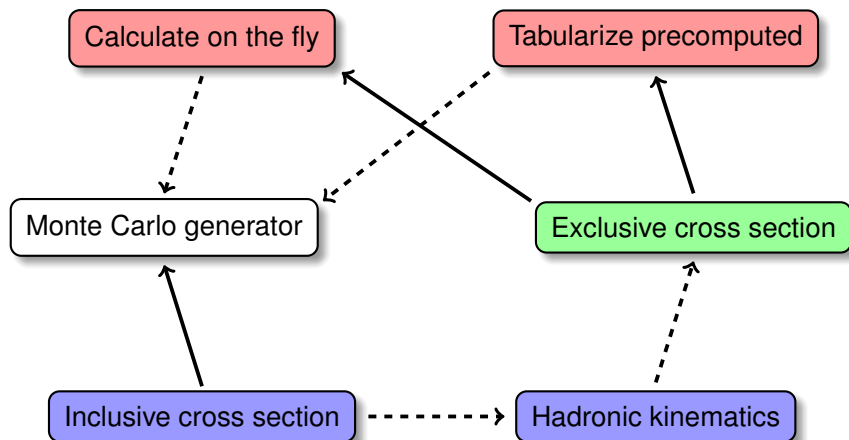
$$\frac{d^6\sigma}{d\epsilon_f d\Omega_f dT_a d\Omega_a}$$

→

Inclusive

$$\frac{d^3\sigma}{d\epsilon_f d\Omega_f}$$

Strategy of implementation



MEC: Valencia model

Implemented in terms of 5 hadronic **structure functions** $W_i(\omega, |\vec{q}|)$:

$$\begin{aligned} \frac{d^2\sigma}{d\epsilon_f d\Omega_f} = & \frac{2G_F^2 |\vec{k}_f| \epsilon_f}{\pi} \left\{ 2W_1 \sin^2 \frac{\theta_l}{2} + W_2 \cos^2 \frac{\theta_l}{2} \mp W_3 (\epsilon_i + \epsilon_f) \sin^2 \frac{\theta_l}{2} \right. \\ & + \frac{m_f^2}{\epsilon_f (\epsilon_f + |\vec{k}_f|)} \left[W_1 \cos \theta_l - \frac{W_2}{2} \cos \theta_l \pm \frac{W_3}{2} (\epsilon_f (1 - \cos \theta_l) + |\vec{k}_f| - \epsilon_i) \right. \\ & \left. \left. + \frac{W_4}{2} (m_f^2 \cos \theta_l + 2\epsilon_f (\epsilon_f + |\vec{k}_f|) \sin^2 \theta_l) - \frac{W_5}{2} (\epsilon_f + |\vec{k}_f|) \right] \right\} \end{aligned}$$

Using 5 hadronic **tensor elements**:

$$\begin{aligned} W_1 &= \frac{1}{2} W^{11}; \quad W_2 = \frac{1}{2} \left(W^{00} + W^{11} + \frac{\omega}{|\vec{q}|^2} (W^{33} - W^{11}) - 2 \frac{\omega}{|\vec{q}|} \Re W^{03} \right) \\ W_3 &= -i \frac{1}{|\vec{q}|} W^{12}; \quad W_4 = \frac{1}{2|\vec{q}|^2} (W^{33} - W^{11}); \quad W_5 = \frac{1}{|\vec{q}|} \left(\Re W^{03} - \frac{\omega}{|\vec{q}|} (W^{33} - W^{11}) \right) \end{aligned}$$

MEC: Valencia model

- Thanks to the courtesy of **Juan Nieves** and **Manuel Vicente Vacas**: **code** producing hadronic **tensor elements** on a **grid** in ω and $|\vec{q}|$
- **Grid size**: 10 MeV steps up to 1.2 GeV in $|\vec{q}|$ ($|\vec{q}_{\text{cut}}|$), same spacing in ω with $\omega \leq |\vec{q}|$ (**triangular physical region**)
- Using $|\vec{q}_{\text{cut}}|$, the original model **extended** up to $E_\nu \simeq 10$ GeV
R. Gran at al., Phys. Rev. D88 (2013) 113007
- **CC** (anti)neutrinos for two targets: **carbon** and **oxygen**
(cross section per nucleon for $Z \leq 6$: carbon grid; for $Z > 6$: oxygen grid)
- For **asymmetric** targets ($A \geq 40$): additional **combinatoric factors**
R. Gran at al., Phys. Rev. D88 (2013) 113007

J. Nieves, I. Ruiz Simo, and M. Vicente Vacas, Phys. Rev. C83 (2011) 045501

MEC: Transverse Enhancement model

- An **effective approach** based on the analysis of **electron** scattering
- np - nh contribution to the **QEL** cross section is introduced by the following **modification** of the **vector magnetic form factor**:

$$G_M^{p/n} \rightarrow \tilde{G}_M^{p/n} = \sqrt{1 + \textcolor{red}{A} Q^2 \exp(-Q^2/\textcolor{blue}{B})} G_M^{p/n}$$

where $\textcolor{red}{A} = 6 \text{ GeV}^{-2}$, $\textcolor{blue}{B} = 0.34 \text{ GeV}^2$ and $M_A^{np-nh} = 1013 \text{ MeV}$

- np - nh **cross section** is obtained by

$$\left(\frac{d\sigma^{\nu/\bar{\nu}}}{dQ^2} \right)_{CC}^{np-nh} \equiv \frac{d\sigma^{\text{QEL}}}{dQ^2}(\tilde{G}_M^{p/n}) - \frac{d\sigma^{\text{QEL}}}{dQ^2}(G_M^{p/n})$$

$$\left(\frac{d\sigma^{\nu/\bar{\nu}}}{dQ^2} \right)_{NC}^{np-nh} \equiv \frac{1}{2} \left[\left(\frac{d\sigma^{\text{EL}}}{dQ^2}(\tilde{G}_M^p) - \frac{d\sigma^{\text{EL}}}{dQ^2}(G_M^p) \right) + \left(\frac{d\sigma^{\text{EL}}}{dQ^2}(\tilde{G}_M^n) - \frac{d\sigma^{\text{EL}}}{dQ^2}(G_M^n) \right) \right]$$

- **Lepton kinematics** based on **QEL**, assuming target **nucleon at rest**

A. Bodek, H.S. Budd, and M.E. Christy, Eur. Phys. J. C71 (2011) 1726

MEC: Marteau-like model

$$\begin{aligned} \frac{d^2\sigma}{d\epsilon_f d\Omega_f} = & \frac{G_F^2 \cos\theta_c |\vec{k}_f|^2}{\pi} \cos^2 \frac{\theta_l}{2} \left[G_E^2 \frac{q^4}{|\vec{q}|^4} R_{\tau}^{NN} \right. \\ & + G_A^2 \frac{(M_\Delta - M_N)^2}{2|\vec{q}|^2} R_{\sigma\tau(L)}^{N\Delta} + G_A^2 \frac{(M_\Delta - M_N)^2}{|\vec{q}|^2} R_{\sigma\tau(L)}^{\Delta\Delta} \\ & + (G_M^2 \frac{\omega^2}{|\vec{q}|^2} + G_A^2) \left(-\frac{q^2}{|\vec{q}|^2} + 2 \tan^2 \frac{\theta_l}{2} \right) (R_{\sigma\tau(T)}^{NN} + 2R_{\sigma\tau(T)}^{N\Delta} + R_{\sigma\tau(T)}^{\Delta\Delta}) \\ & \left. \pm G_A G_M \frac{|\vec{k}_i| + |\vec{k}_f|}{M_N} \tan^2 \frac{\theta_l}{2} (R_{\sigma\tau(T)}^{NN} + 2R_{\sigma\tau(T)}^{N\Delta} + R_{\sigma\tau(T)}^{\Delta\Delta}) \right] \end{aligned}$$

Responses due to *nucleon-hole* or *delta-hole* excitations:

$$R_{\alpha}^{PP'} = \sum_n \langle n | \sum_{j=1}^A O_{\alpha}^P(j) e^{i\vec{q} \cdot \vec{x}_j} | 0 \rangle \langle n | \sum_{k=1}^A O_{\alpha}^{P'}(k) e^{i\vec{q} \cdot \vec{x}_k} | 0 \rangle^* \delta(\omega - E_n + E_0)$$

where

$$\begin{aligned} O_{\tau, \sigma\tau(L), \sigma\tau(T)}^N(j) &= \tau_j^{\pm}, (\sigma_j \cdot \hat{q}) \tau_j^{\pm}, ((\sigma_j \times \hat{q}) \times \hat{q}) \tau_j^{\pm} \\ O_{\sigma\tau(L), \sigma\tau(T)}^{\Delta}(j) &= (\mathbf{S}_j \cdot \hat{q}) T_j^{\pm}, ((\mathbf{S}_j \times \hat{q}) \times \hat{q}) T_j^{\pm} \end{aligned}$$

J. Marteau, Eur. Jour. Phys. A5 (1999) 183

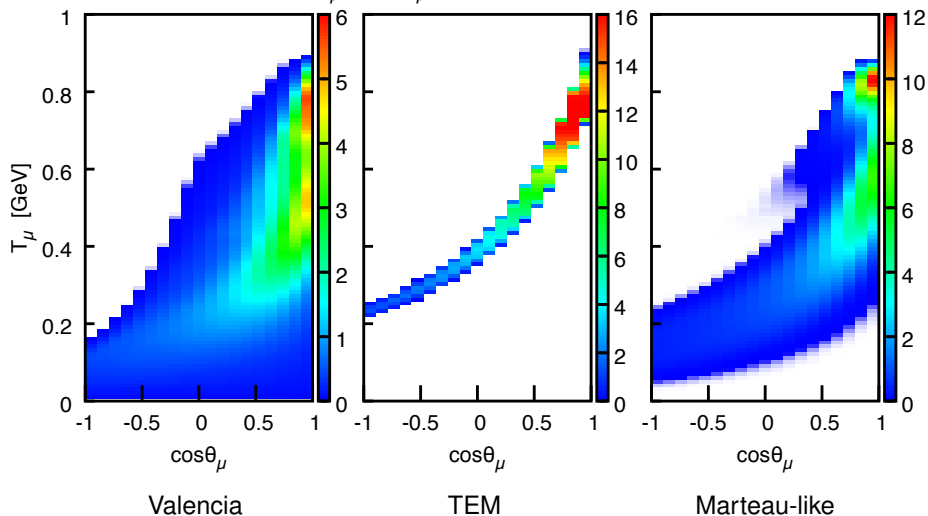
MEC: Marteau-like model

- **Custom implementation** with **tabularized responses** up to $\omega < 3 \text{ GeV}$, $|\vec{q}| < 3.6 \text{ GeV}$, then actually **solving integrals**
- The **relativistic kinematics** of nucleons is used
- **Constant** value of **Fermi momentum** is kept
- **Local density effects** are included in the **approximate** way

J. Sobczyk, arXiv:nucl-th/0307047 (2003)

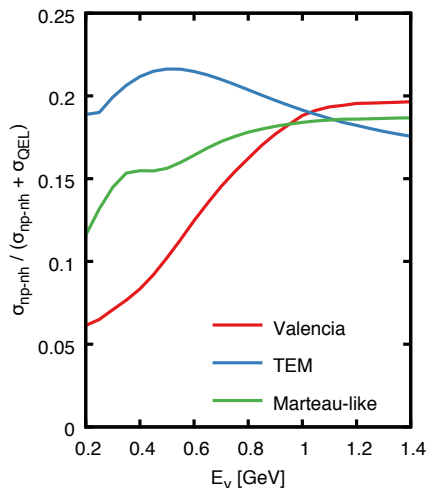
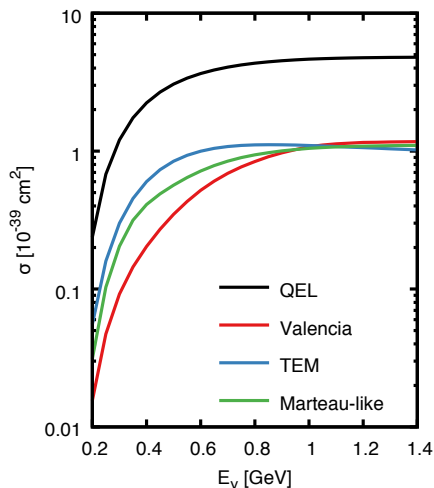
MEC: Inclusive cross section

$$d\sigma / dT_\mu / d\cos\theta_\mu [10^{-39} \text{ cm}^2 / \text{GeV} / \text{nucleon}]$$



Scattering of 1 GeV neutrinos on carbon target

MEC: Total cross sections



Scattering of 1 GeV neutrinos on carbon target

A little more exclusive...

- We need to guess the **isospin configuration** of **target nucleons**

$$\nu_l + p + n \rightarrow l^- + p + p$$

$$\bar{\nu}_l + p + n \rightarrow l^+ + n + n$$

$$\nu_l + n + n \rightarrow l^- + p + n$$

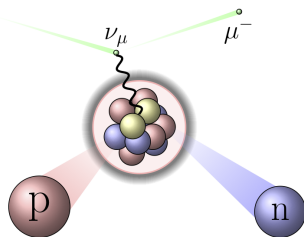
$$\bar{\nu}_l + p + p \rightarrow l^+ + p + n$$

$$\nu_l + 2N \rightarrow \nu_l + 2N$$

$$\bar{\nu}_l + 2N \rightarrow \bar{\nu}_l + 2N$$

- Fraction of **mixed isospin** pairs for CC is a **free parameter**: p_{CC}
- For NC it is obtained by: $p_{NC} = (2/p_{CC} - 1)^{-1}$

	p_{CC}
Nieves et al.	$\simeq 0.70$
GFMC (SRC)	$\simeq 0.95$
our guess	≈ 0.85



Hadronic phase-space model

1.



Choose p_1, p_2 back-to-back

2.



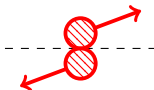
$$(p_1 + p_2 + q)^2 > 4M_N^2$$

3.



Boost to CMS frame

4.



Decay isotropically

5.



Boost back to LAB frame

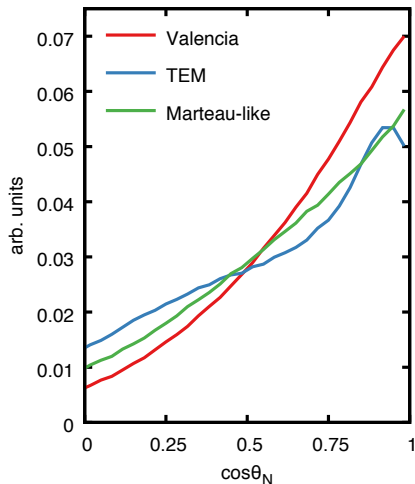
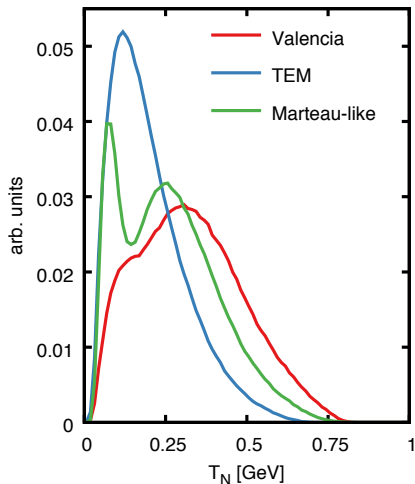
6.

$$\Theta(k_F - |\vec{p}'_1|)\Theta(k_F - |\vec{p}'_2|)$$

Check Pauli blocking

Hadronic phase-space model

Distributions for the **most energetic nucleon** before FSI



Scattering of 1 GeV neutrinos on carbon target

Enhancing initial hadronic kinematics

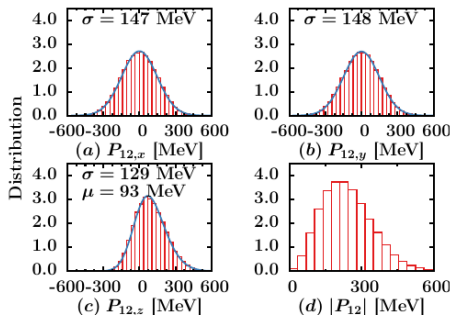
- Add nucleon pair
CM momentum

`mec_central_motion`
parameterizes the gaussian
distribution

- Add **smearing** of the
back-to-back configuration
(when no CM motion)

`mec_back_to_back_smearing`
parameterizes the gaussian
distribution

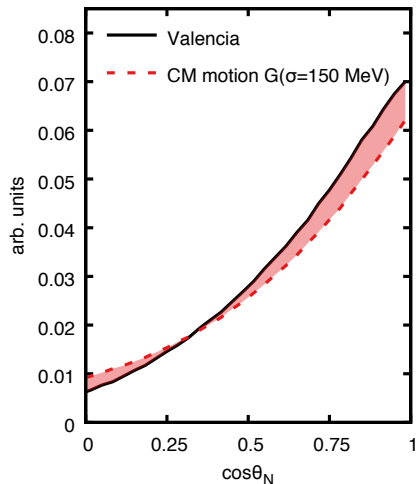
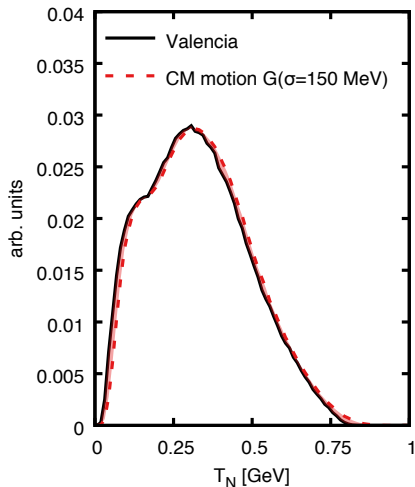
Nucleon pair CM distributions



C. Colle et al., Phys.Rev. C89 (2014) 024603

Enhancing initial hadronic kinematics

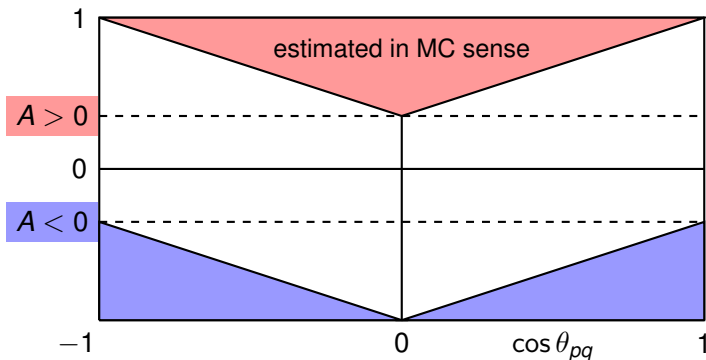
Distributions for the **most energetic nucleon** before FSI



Scattering of 1 GeV neutrinos on carbon target

Enhancing final hadronic kinematics

- Nucleon pair **CM distribution** can be **different** than **uniform**
- **Direction** of **momentum transfer** $\vec{q}$ can be distinguished

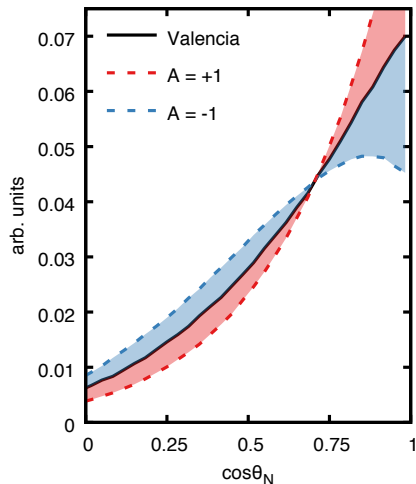
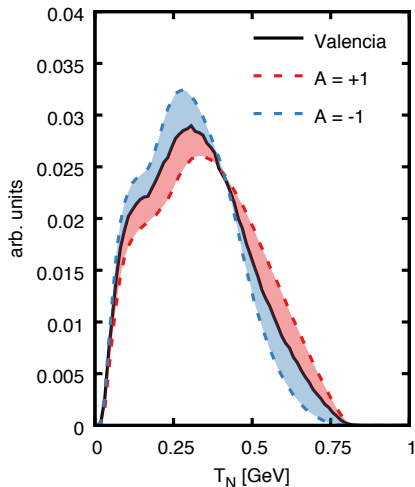


MEC_cm_direction (denoted as A)

allows for selecting directions on average more parallel or perpendicular

Enhancing final hadronic kinematics

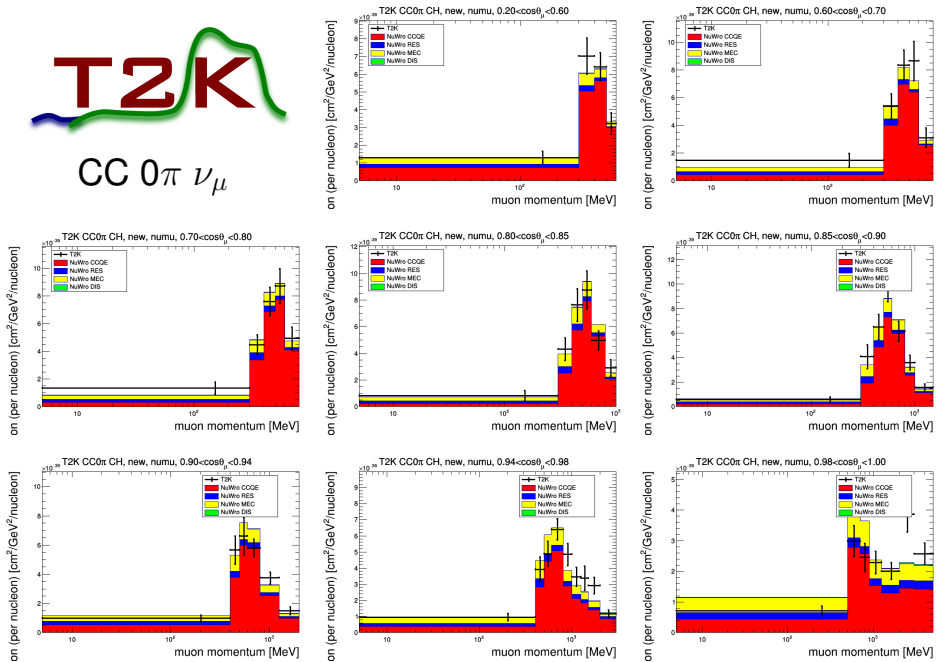
Distributions for the **most energetic nucleon** before FSI



Scattering of 1 GeV neutrinos on carbon target

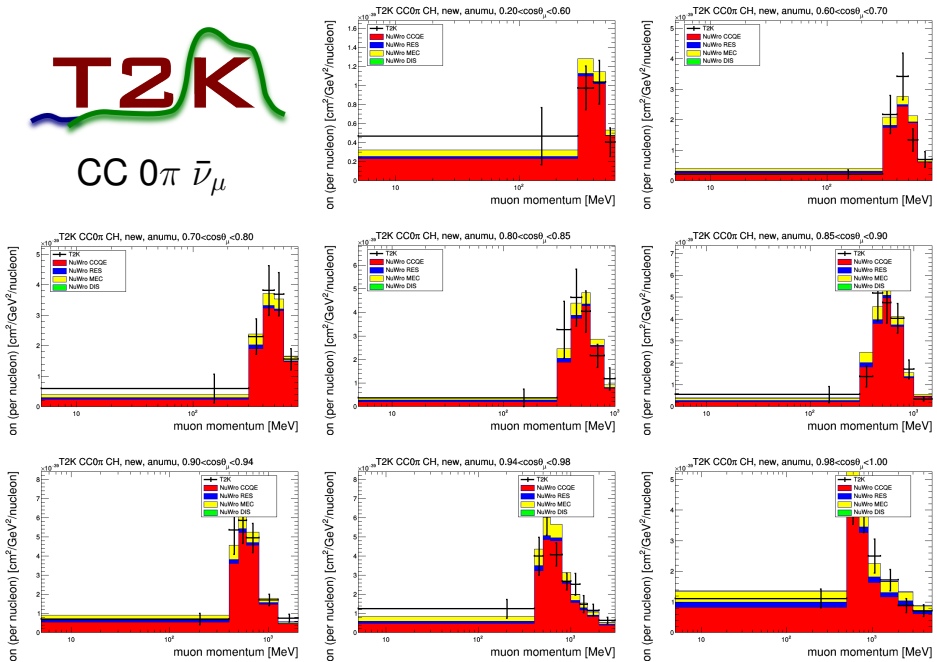
T2K

CC $0\pi \nu_\mu$



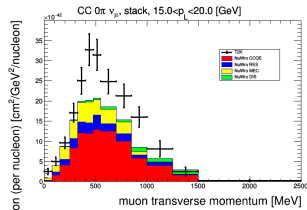
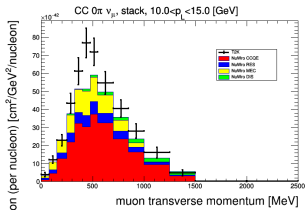
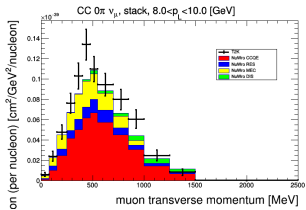
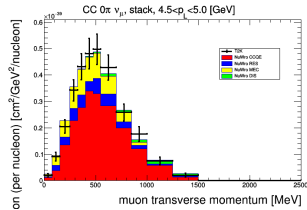
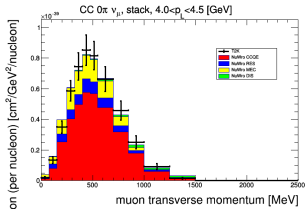
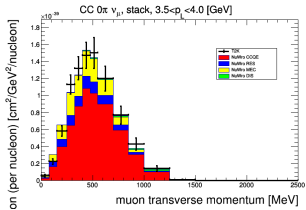
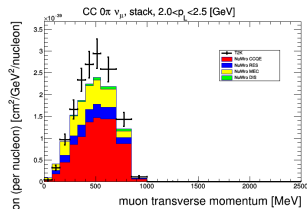
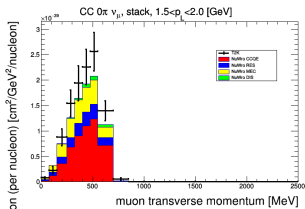
T2K

CC $0\pi \bar{\nu}_\mu$



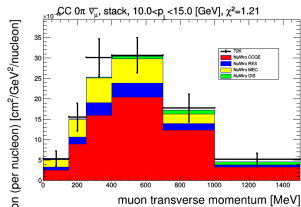
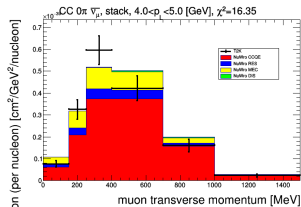
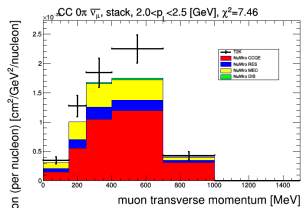
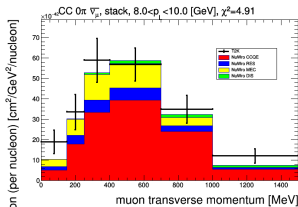
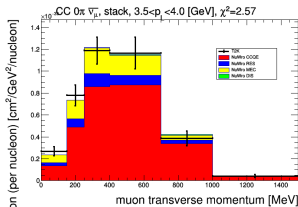
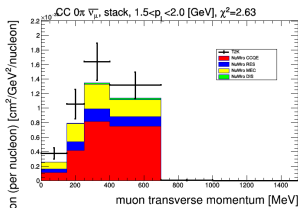
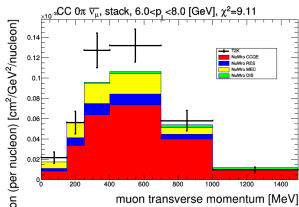
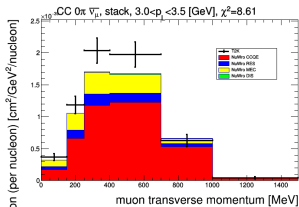


CC $0\pi \nu_\mu$





CC $0\pi \bar{\nu}_\mu$



Conclusions

We are working on it!