

# *Low-energy $K^-$ -nucleus/nuclei interactions with light nuclei with AMADEUS*

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*On the behalf of the AMADEUS collaboration*

*ROCKSTAR: Towards a ROadmap of the Crucial measurements of Key observables in Strangeness reactions for neutron sTARs equation of state  
9-13 October 2023, ECT\* - Trento (Italy)*

# AMADEUS scientific case

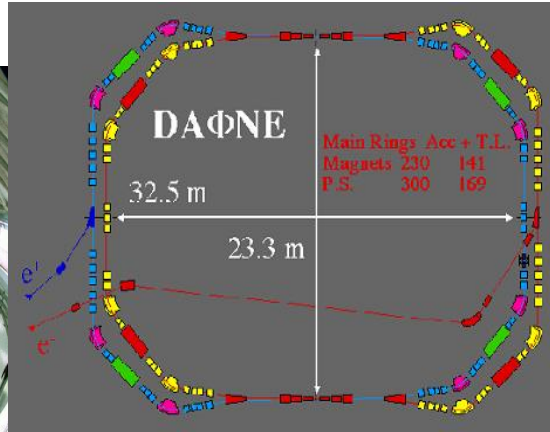
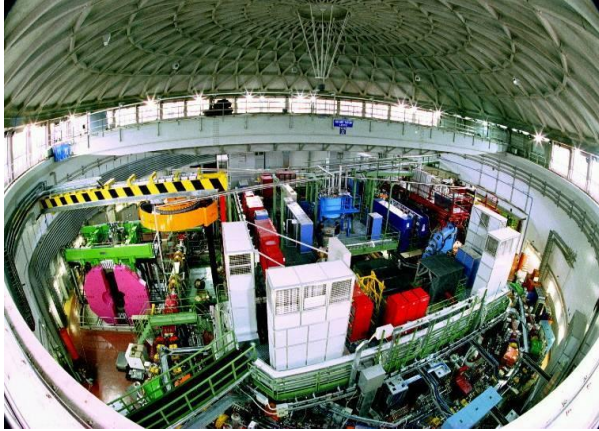
AMADEUS (Antikaonic Matter At DAΦNE: an Experiment with Unravelling Spectroscopy ) investigates **low-energy  $K^-$  absorption in nuclei** with the aim to extract information on:

- $K^-N$  interaction above and below threshold
  - $\Lambda(1405)$  nature
  - $K^-N$  scattering amplitudes and cross sections
- $K^-NN$ ,  $K^-NNN$ ,  $K^-NNNN$  (multi-nucleon) interactions
  - $K^-$ -multi nucleon cross sections
  - essential for the determination of  $K^-$ -nuclei optical potential
  - kaonic bound states
- Hyperon-nucleon/(multi-nucleons) interaction cross sections

# DAΦNE the $\Phi$ factory

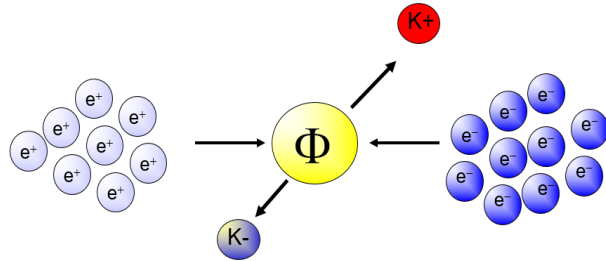


Istituto Nazionale di Fisica Nucleare  
LABORATORI NAZIONALI DI FRASCATI



- $e^+ e^-$  at 510 MeV
- $\phi$  resonance decays at 49 % in  $K^+ K^-$  back to back pair
- Very low momentum ( $\approx 127$  MeV/c)  $K^-$  beam

Unique low momentum  $K^-$  factory



Suitable for low-energy kaon physics:

→ **Kaonic atoms** (**SIDDHARTA-2**)

→ **Kaon-nucleons/nuclei interaction** studies  
(**AMADEUS**)

# LNf - $e^+e^-$ Accelerator Complex



DAΦNE

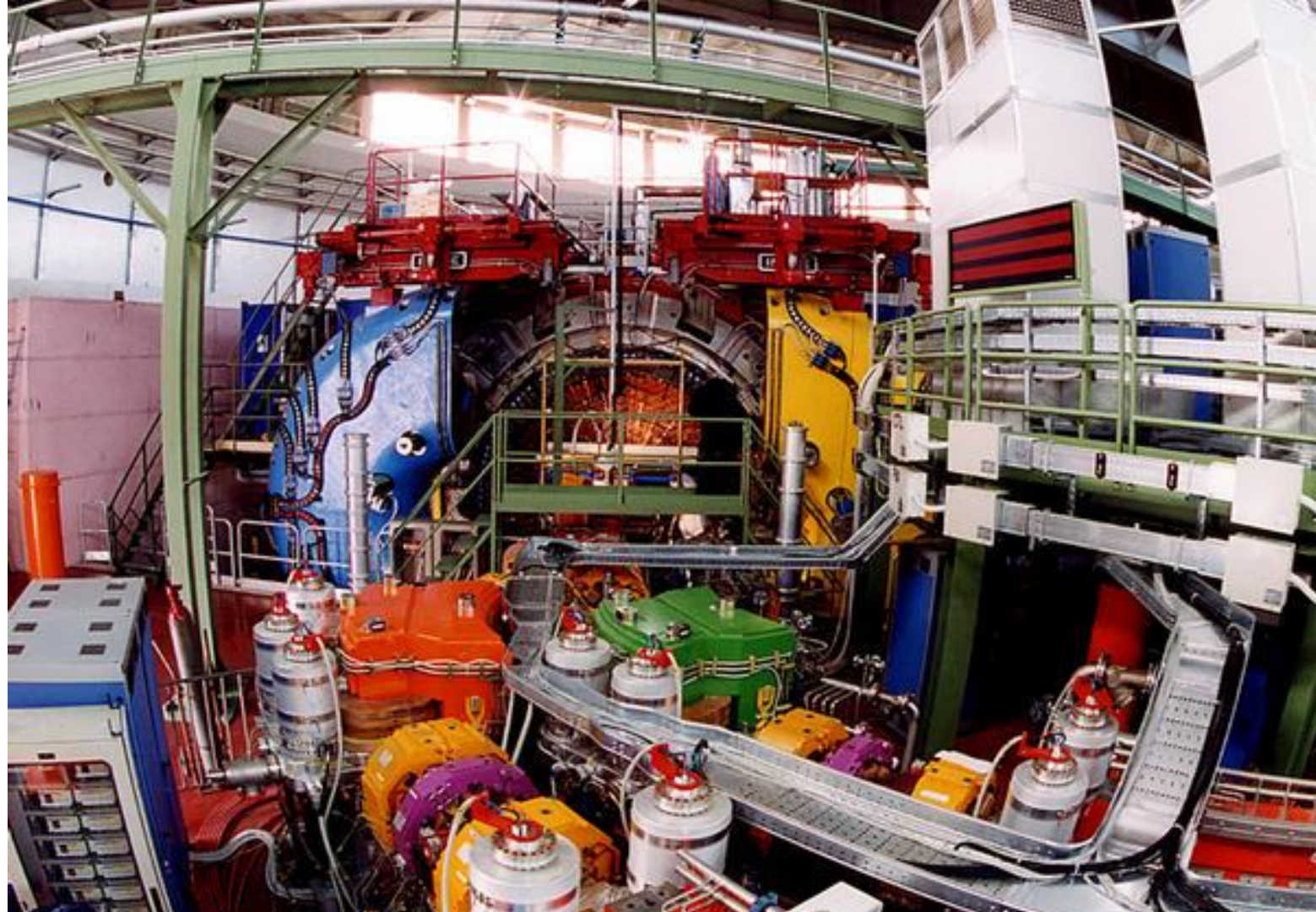
LINAC

DAMPING RING

$e^+$

$e^-$

$\Phi$



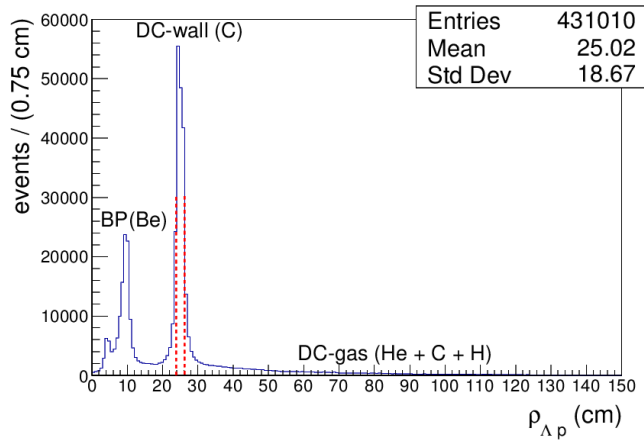
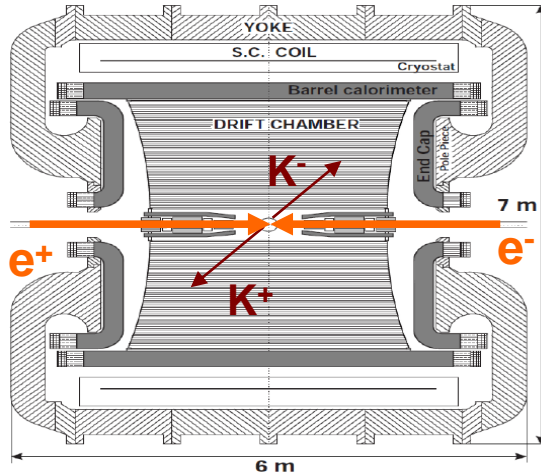
# AMADEUS

## The KLOE detector

Cylindrical drift chamber with a  $4\pi$  geometry and electromagnetic calorimeter, **96% acceptance**

- optimized in the energy range of all **charged particles** involved
- **good performance** in detecting **photons and neutrons** checked by kloNe group

[M. Anelli et al., Nucl Inst. Meth. A 581, 368 (2007)]



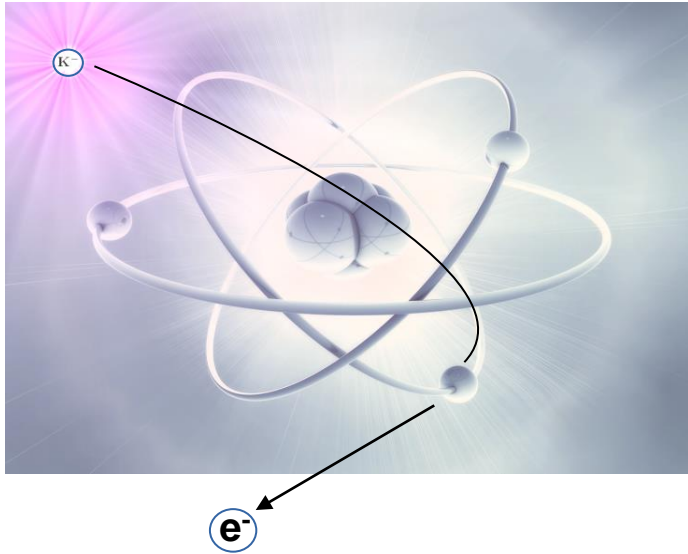
### KLOE used as an active target

- DC wall (750  $\mu\text{m}$  C foil , 150  $\mu\text{m}$  Al foil);
  - DC gas (90% He, 10%  $\text{C}_4\text{H}_{10}$ ).
- +  
pure sample of  $\text{K}^-$   $^{12}\text{C}$  absorptions at-rest

# K<sup>-</sup> absorptions at-rest and in-flight

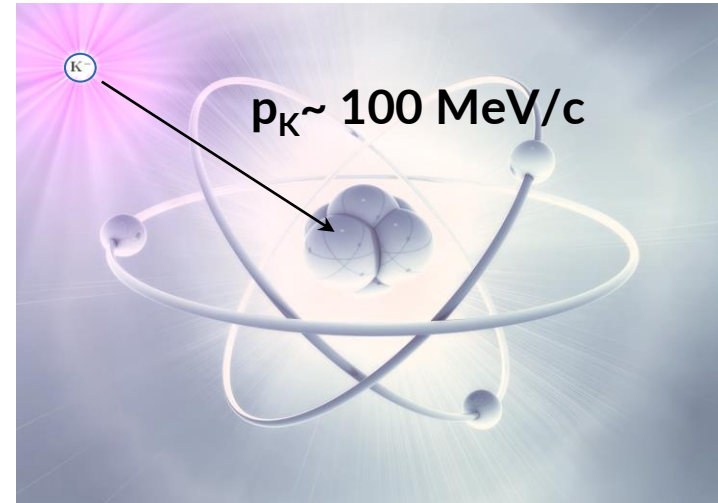
## AT-REST

K<sup>-</sup> absorbed from atomic orbitals  
( $p_K \sim 0 \text{ MeV}/c$ )



## IN-FLIGHT

( $p_K \sim 100 \text{ MeV}/c$ )



# $K^- n \rightarrow \Lambda \pi^-$ events selection and interpretation

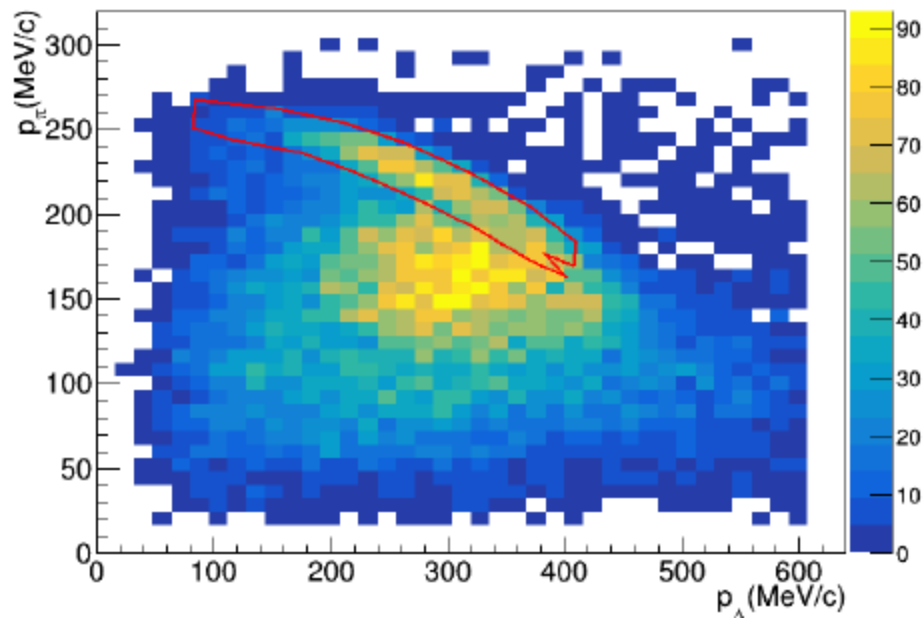


FIG. 2. (Color online) Experimental distribution of the  $\pi^-$  vs  $\Lambda$  momenta. The red line represents the selection of the direct  $\Lambda \pi^-$  production events. See the text for details.

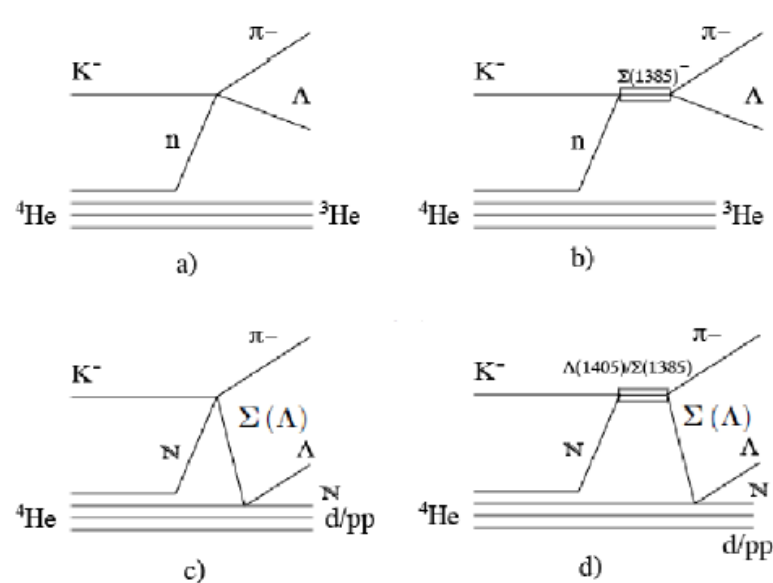
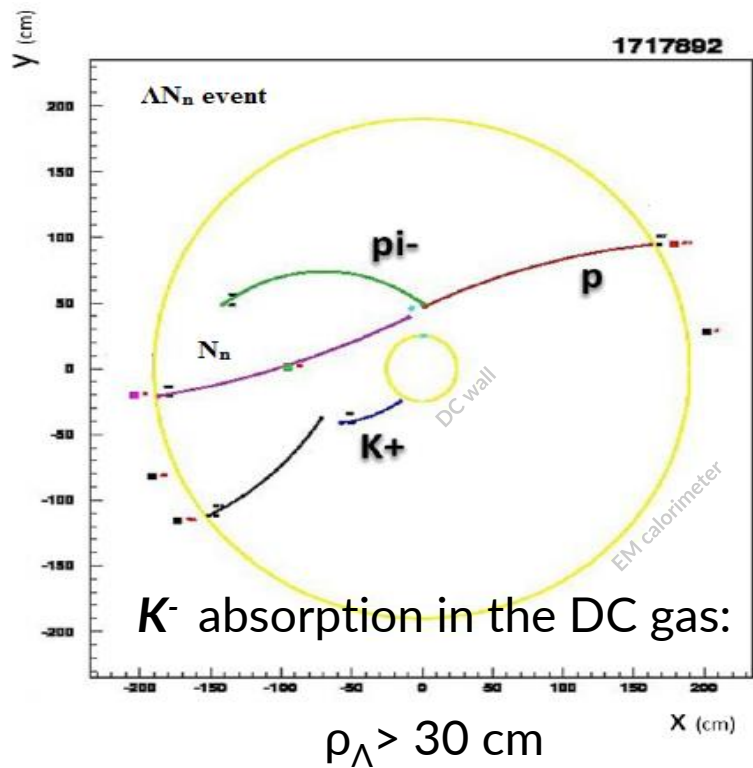


FIG. 1. Panels a) and b) show the non-resonant and resonant  $\Lambda \pi^-$  direct productions, respectively. Panels c) and d) show the primary hyperon-pion formation, followed by the inelastic/elastic scattering of the  $\Sigma/\Lambda$  hyperon on a single nucleon, for the resonant and non-resonant cases, respectively.

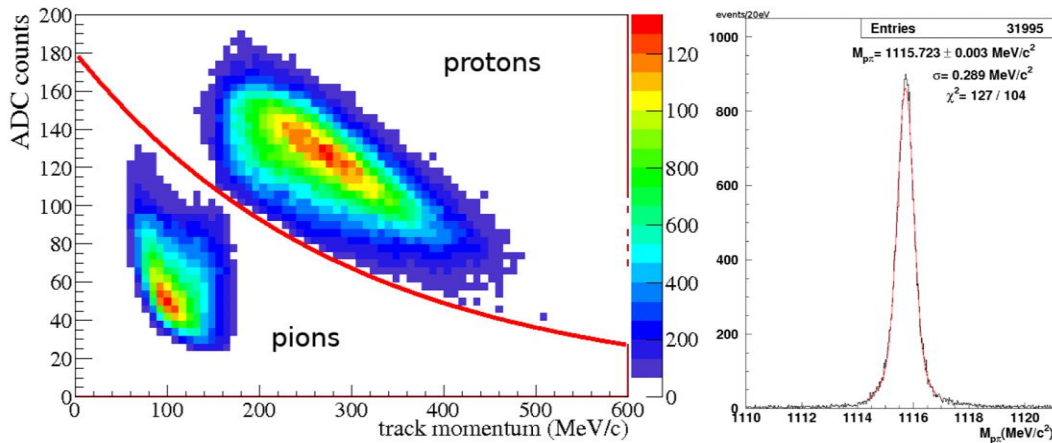
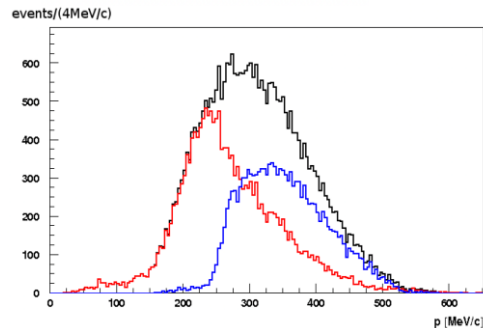


# Events selection - $\Lambda \rightarrow p \pi^-$ (BR = $63.9 \pm 0.5\%$ )

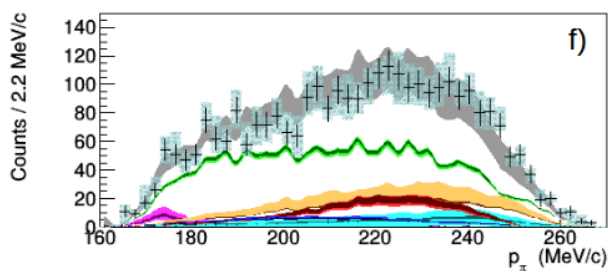
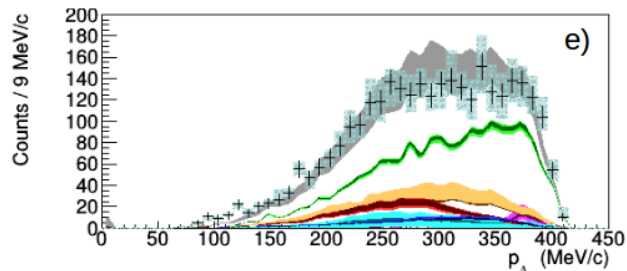
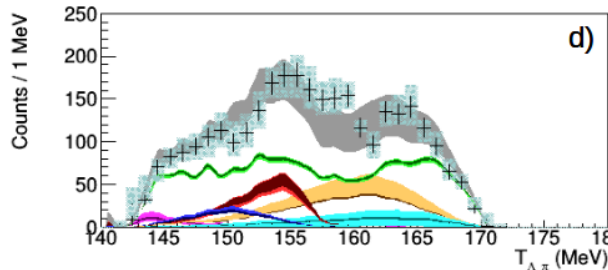
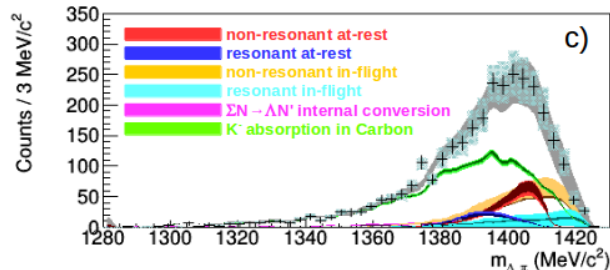
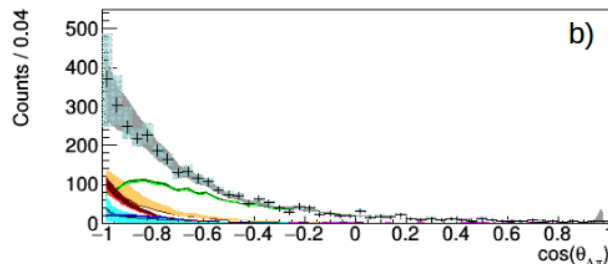
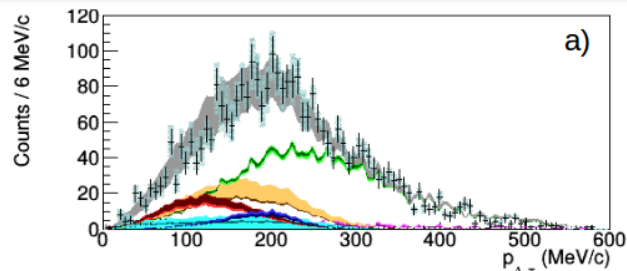


contamination with  $K^-$  absorption  
in the DC wall < 3%

Opposite charged tracks with common vertex,  
main background: ( $K^\pm \rightarrow \pi^\pm \pi^\pm \pi^\mp$ )



# Simultaneous fit : $p_{\Lambda\pi^-} - m_{\Lambda\pi^-} - \cos\theta_{\Lambda\pi^-}$



Investigated using:  
 $K^- - n - {}^3\text{He} \rightarrow \Lambda\pi^- - {}^3\text{He}$

$$E_{K_n} \sim -B_n - \left\langle \frac{p_{\Lambda\pi}^2}{2\mu_{\pi,\Lambda,3\text{He}}} \right\rangle$$

$33 \pm 6$  MeV below threshold  
 see also

A. Cieply et al., Phys. Lett. B 702 (2011) 402  
 T. Hoshino et al., Phys. Rev. C 96 (2017) 045204  
 N. Barnea, E. Friedman,  
 A. Gal, Nucl. Phys. A968 (2017)

[K. P., S. Wycech, L. Fabbietti et al. Phys.Lett. B782 (2018) 339-345]

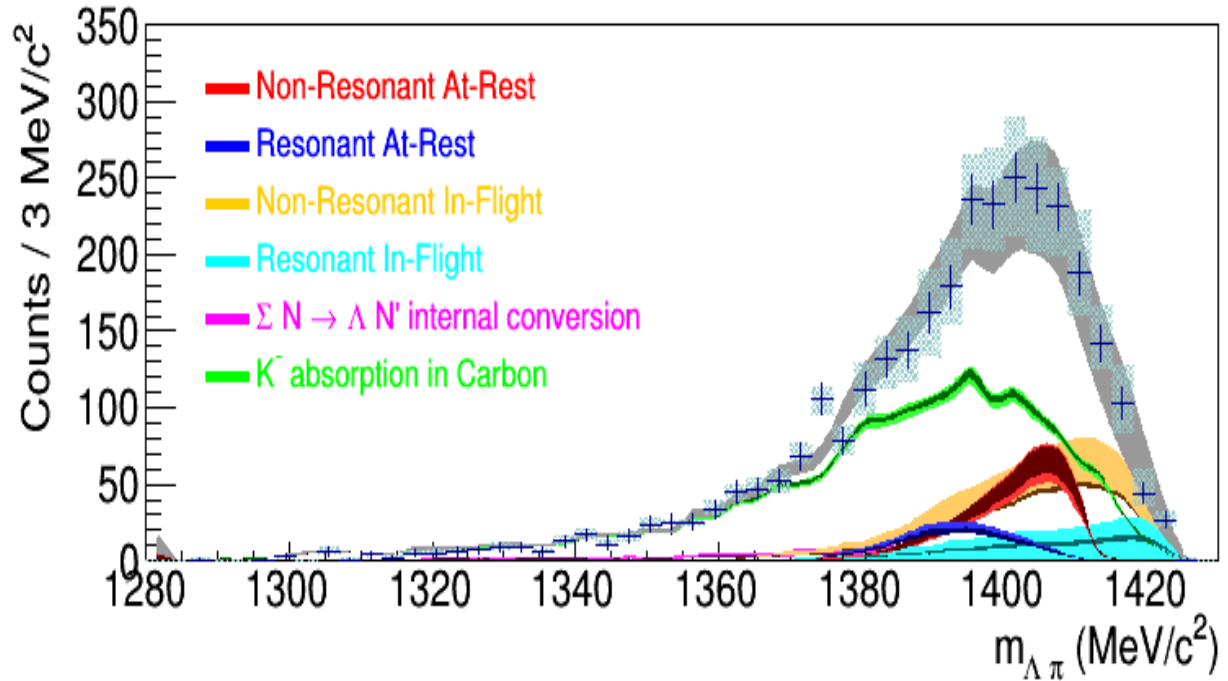
[K. P., S. Wycech, C. Curceanu, Nucl. Phys. A 954 (2016) 75-93]

# Outcome of the measurement

Investigated using:  
 $K^- - n \rightarrow \Lambda \pi^- \text{ } ^3\text{He}$

Energy of the  $K^-n$  system:

$$E_{K_n} \sim -B_n - \left\langle \frac{p_{\Lambda\pi}^2}{2\mu_{\pi,\Lambda,^3\text{He}}} \right\rangle$$

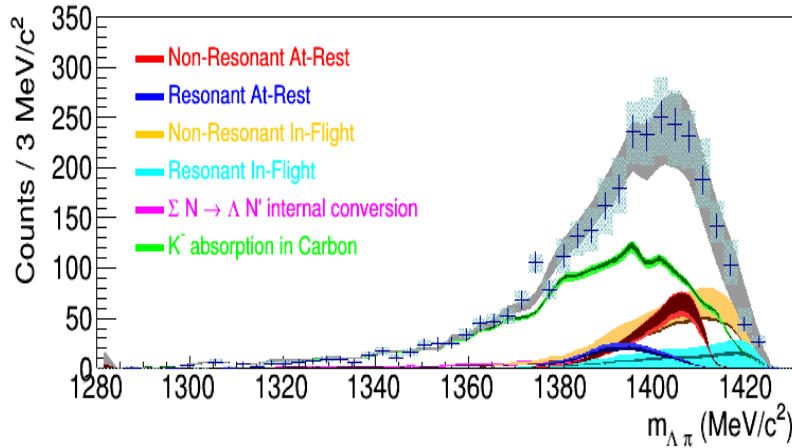


[K. P., S. Wycech, L. Fabbietti et al. Phys.Lett. B782 (2018) 339-345]

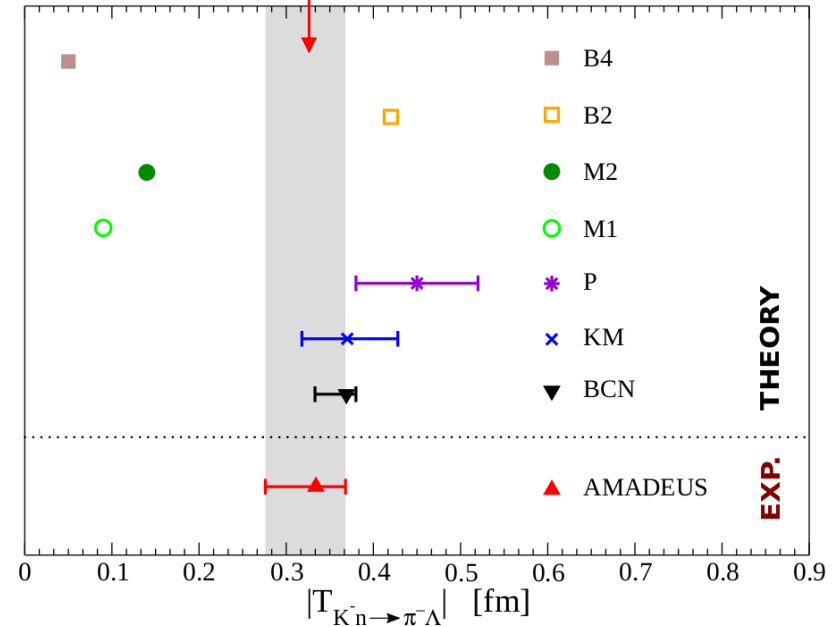
[K. P., S. Wycech, C. Curceanu, Nucl. Phys. A 954 (2016) 75-93]

# Outcome of the measurement

Investigated using:  $K^- - n \text{ } ^3\text{He} \rightarrow \Lambda \pi^- \text{ } ^3\text{He}$



$$|f_{ar}^{nr}| = (0.334 \pm 0.018 \text{ stat}^{+0.034}_{-0.058} \text{ syst}) \text{ fm}$$



[K. P., S. Wycech, L. Fabbietti et al. Phys.Lett. B782 (2018) 339-345]

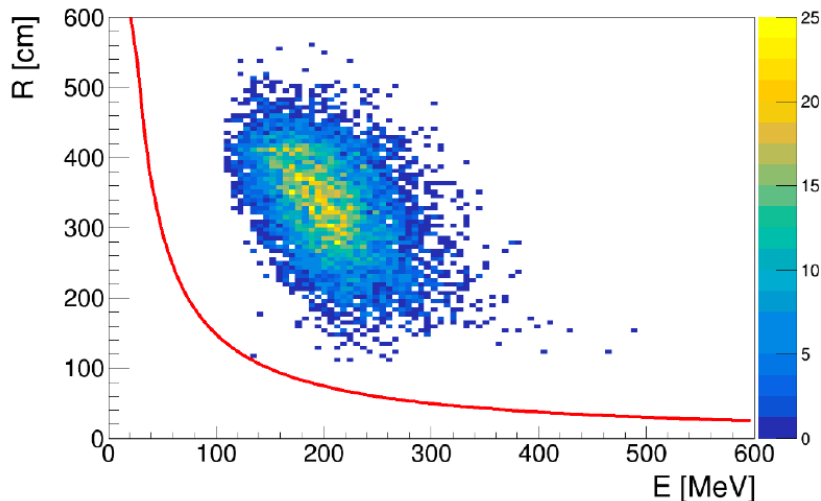
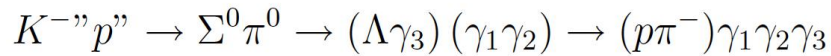
[K. P., S. Wycech, C. Curceanu, Nucl. Phys. A 954 (2016) 75-93]

**Simultaneous/independent measurement of the**



**cross sections at  $p_{K^-} = 98 \pm 10 \text{ MeV}/c$**

# Events selection - $\gamma_1, \gamma_2$ & $\gamma_3$



- $K^+ \rightarrow \pi^+ \pi^0$  background is rejected
- three photon clusters are selected by TOF

$$\chi_t^2 = t^T V_t^{-1} t$$

with  $t = t_i - t_j$  ;  $t_i = t_{cli} - r_i/c$

- disentangling  $\gamma_1, \gamma_2$  &  $\gamma_3$  :

$$\chi_{\pi\Sigma}^2 = \frac{(m_{\pi^0} - m_{ij})^2}{\sigma_{ij}^2} + \frac{(m_{\Sigma^0} - m_{k\Lambda})^2}{\sigma_{k\Lambda}^2}$$

- MC based rejection criteria:

$$\chi_t^2 \leq 20, \chi_{m_{\gamma_1 \gamma_2}}^2 \leq 5 \text{ and } \chi_{m_{\Lambda \gamma_3}}^2 \leq 4$$

- Cluster splitting background free! Algorithm overall efficiency for  $\gamma$  detection:  $0.98 \pm 0.01$ .

# Events selection - $\Sigma^0$ & $\pi^0$

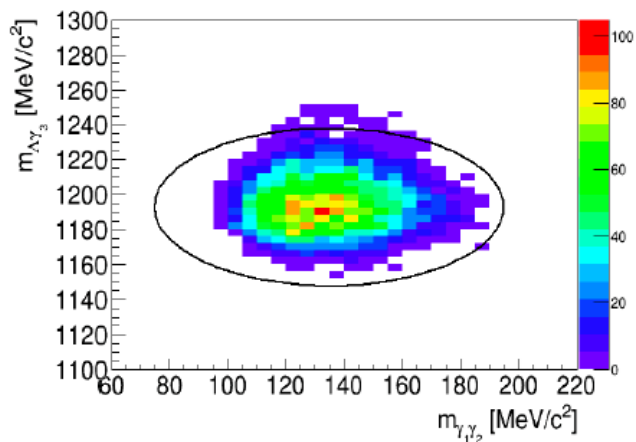
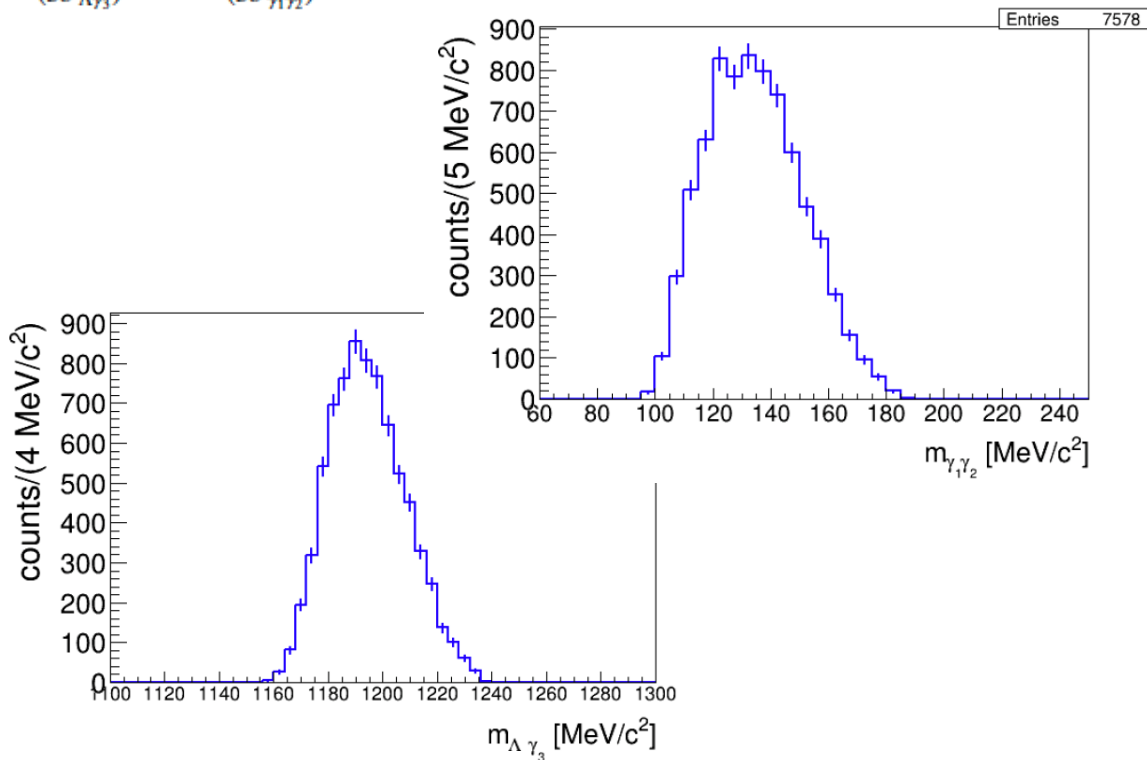


FIG. 1. The plot shows the  $m_{\Lambda\gamma_3}$  distribution as a function of  $m_{\gamma_1\gamma_2}$ . A black ellipse represents the applied selection in the two invariant masses space.

$$\frac{(m_{\Lambda\gamma_3} - m_{\Sigma^0})^2}{(3\sigma_{\Lambda\gamma_3})^2} + \frac{(m_{\gamma_1\gamma_2} - m_{\pi^0})^2}{(3\sigma_{\gamma_1\gamma_2})^2} < 1$$

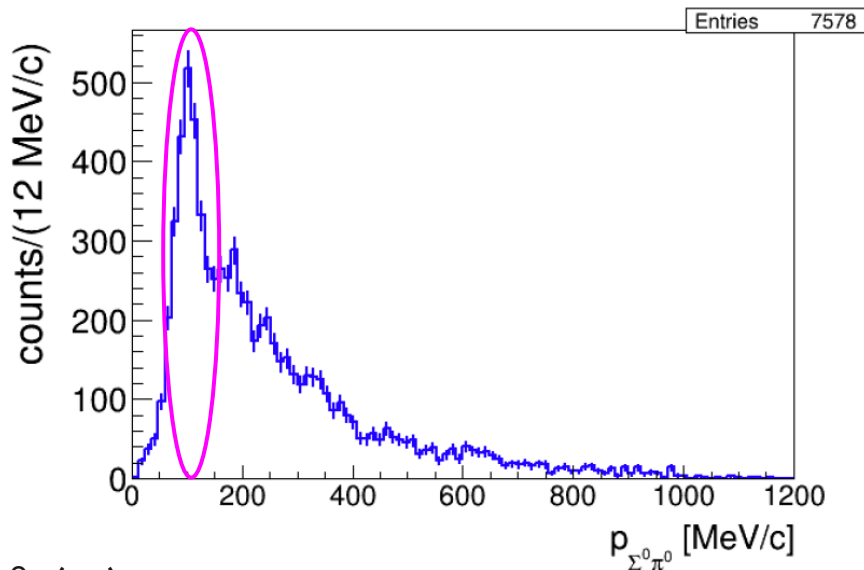
- RESOLUTIONS:

$$\sigma_{\Lambda\gamma_3} \sim 15 \text{ MeV/c} \text{ and } \sigma_{\gamma_1\gamma_2} \sim 20 \text{ MeV/c}$$



# Events selection - $K^- H \rightarrow (\Sigma^0 / \Lambda) \pi^0$ in flight

- **SIGNAL:**  $K^- H \rightarrow (\Sigma^0 / \Lambda) \pi^0$  (*if*)



- **BACKGROUND:**  $K^- H \rightarrow (\Sigma^0 / \Lambda) \pi^0$  (*ar*)

$K^- (^4\text{He}/^{12}\text{C}) \rightarrow (\Sigma^0 / \Lambda) \pi^0 + \text{Residual}$  (*ar + if*)

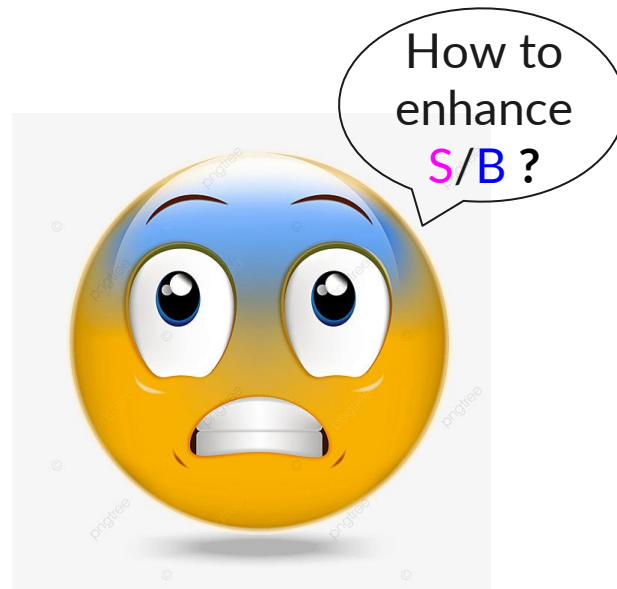
elastic or inelastic FSI of the hyperon ( $Y$ ) and/or the  $\pi^0$

- **FURTHERMORE** in  $\Lambda \pi^0$  the direct production is affected by the background:  $\Sigma^0 \pi^0$  primary production followed by  $\Sigma^0 \rightarrow \Lambda \gamma$  for ALL the channels



# Events selection - $K^- H$ abs. in flight

- **SIGNAL:**  $K^- H \rightarrow (\Sigma^0 / \Lambda) \pi^0$  (if)



- **BACKGROUND:**  $K^- H \rightarrow (\Sigma^0 / \Lambda) \pi^0$  (ar)

$K^- (^4\text{He}/^{12}\text{C}) \rightarrow (\Sigma^0 / \Lambda) \pi^0 + \text{Residual}$  (ar + if)

elastic or inelastic FSI of the hyperon ( $\Upsilon$ ) and/or the  $\pi^0$

- **FURTHERMORE** in  $\Lambda \pi^0$  the direct production is affected by the background:  $\Sigma^0$   
 $\pi^0$  primary production followed by  $\Sigma^0 \rightarrow \Lambda \gamma$  for ALL the channels

# Events selection - $K^- H$ abs. in flight

**SIGNAL:**  $K^- H \rightarrow (\Sigma^0 / \Lambda) \pi^0$  (*if*)

characteristic features:

a) the kinematics (for both *ar* & *if*) is completely determined by E-p cons.  
 signal is almost back to back,

b)  $K^- H \rightarrow \Lambda \pi^0$  (*ar* & *if*) events can be sampled exploiting the resolution on  $p_\Lambda$

$$\sigma_{p\Lambda} = 1.9 \pm 0.2 \text{ MeV}/c$$

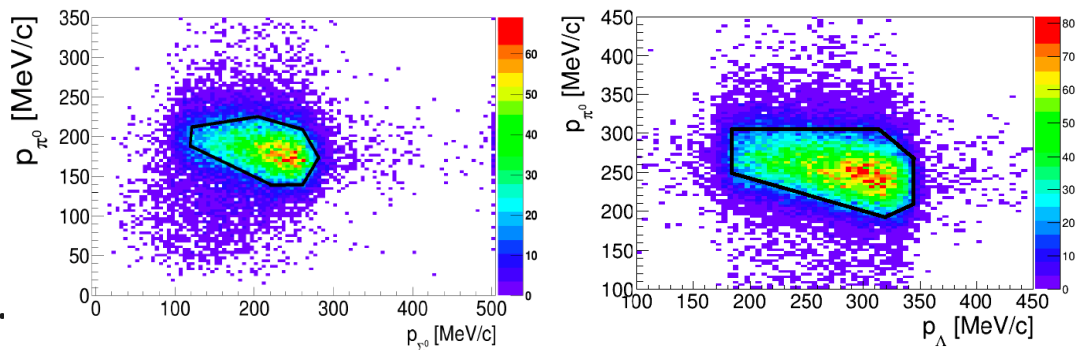
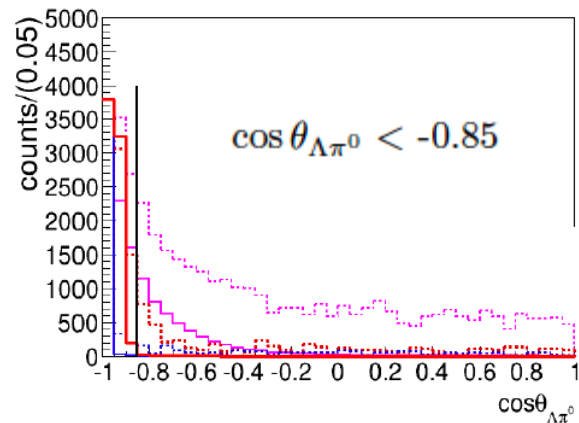


FIG. 2. The plot shows reconstructed MC  $p_{\pi^0}$  vs.  $p_Y$  distributions for the  $K^- H \rightarrow \Sigma^0 \pi^0$  *if* reaction (top) and  $K^- H \rightarrow \Lambda \pi^0$  *if* reaction (bottom). The phase space selections are represented as black contours.

- $K^- H \rightarrow \Lambda \pi^0$  (*if*)
- $K^- H \rightarrow \Lambda \pi^0$  (*ar*)
- He/C *ar/if* background  $\Lambda \pi^0$
- $K^- H \rightarrow \Sigma^0 \pi^0$  (*if*)
- $K^- H \rightarrow \Sigma^0 \pi^0$  (*ar*)
- He/C *ar/if* background  $\Sigma^0 \pi^0$



# Events selection - $K^-$ H abs. in flight

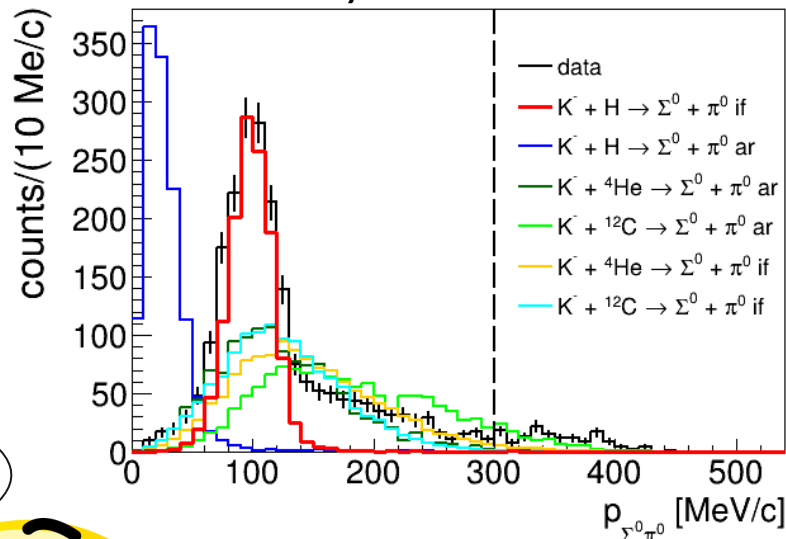
c) due to E-p cons. the total  $Y\pi^0$  momentum distribution for  $K^-$  H  $\rightarrow$   $Y \pi^0$  (if) events reflects the **original  $K^-$  momentum spectrum**

FSI is negligible after phase-space selections, is considered in systematics.

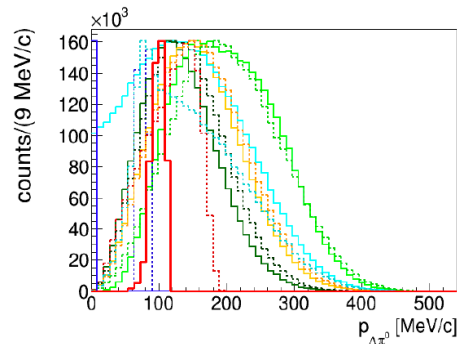
residual bkg:  $K^-$  H (ar) and  $K^-$  ( $^4\text{He}/^{12}\text{C}$ ) can be efficiently disentangled by means of a simultaneous fit



arbitrary normalization



- $K^-$  H  $\rightarrow$   $\Lambda\pi^0$  (if)
- $K^-$  H  $\rightarrow$   $\Lambda\pi^0$  (ar)
- $K^-$   $^4\text{He}$   $\rightarrow$   $\Lambda\pi^0 + ^3\text{H}$  (ar)
- $K^-$   $^{12}\text{C}$   $\rightarrow$   $\Lambda\pi^0 + ^{11}\text{B}$  (ar)
- $K^-$   $^4\text{He}$   $\rightarrow$   $\Lambda\pi^0 + ^3\text{H}$  (if)
- $K^-$   $^{12}\text{C}$   $\rightarrow$   $\Lambda\pi^0 + ^{11}\text{B}$  (if)
- ⋯  $K^-$  H  $\rightarrow$   $\Sigma^0\pi^0$  (if)
- ⋯  $K^-$  H  $\rightarrow$   $\Sigma^0\pi^0$  (ar)
- ⋯  $K^-$   $^4\text{He}$   $\rightarrow$   $\Sigma^0\pi^0 + ^3\text{H}$  (ar)
- ⋯  $K^-$   $^{12}\text{C}$   $\rightarrow$   $\Sigma^0\pi^0 + ^{11}\text{B}$  (ar)
- ⋯  $K^-$   $^4\text{He}$   $\rightarrow$   $\Sigma^0\pi^0 + ^3\text{H}$  (if)
- ⋯  $K^-$   $^{12}\text{C}$   $\rightarrow$   $\Sigma^0\pi^0 + ^{11}\text{B}$  (if)



# Simultaneous fit and cross sections ( $\Sigma^0 \pi^0$ )

$$\chi^2 = \sum_q \sum_{n=1}^{N_{bins}^q} \frac{(N_n^q - \mathcal{F}^q(q_n))^2}{\sigma_n^{q2}} \quad \text{with} \quad \mathcal{F}^q(q_n) = \sum_{i=1}^{N_{par}} \alpha_i \cdot h_i^q(q_n)$$

physical processes:

1.  $K^- H \rightarrow (\Sigma^0/\Lambda) \pi^0$  if (red),
2.  $K^- H \rightarrow (\Sigma^0/\Lambda) \pi^0$  ar (blue),
3.  $K^- + {}^4\text{He}/{}^{12}\text{C} \rightarrow \Sigma^0/\Lambda + \pi^0 + {}^3\text{H}/{}^{11}\text{B}$  (magenta).

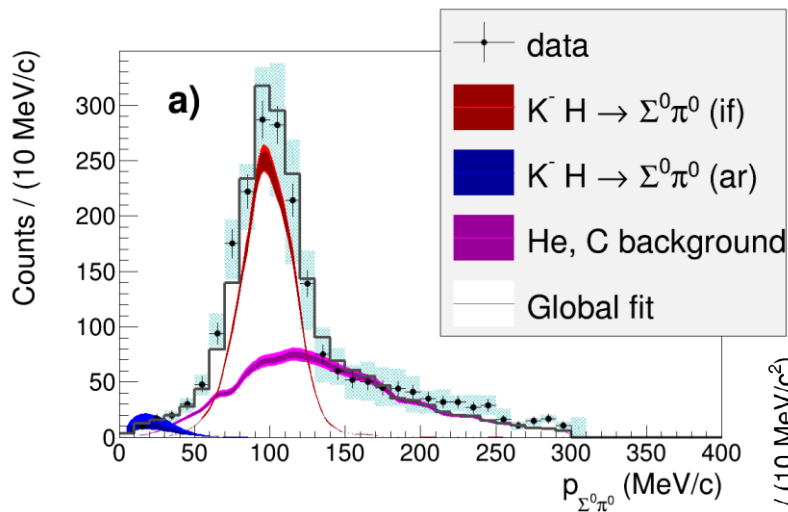
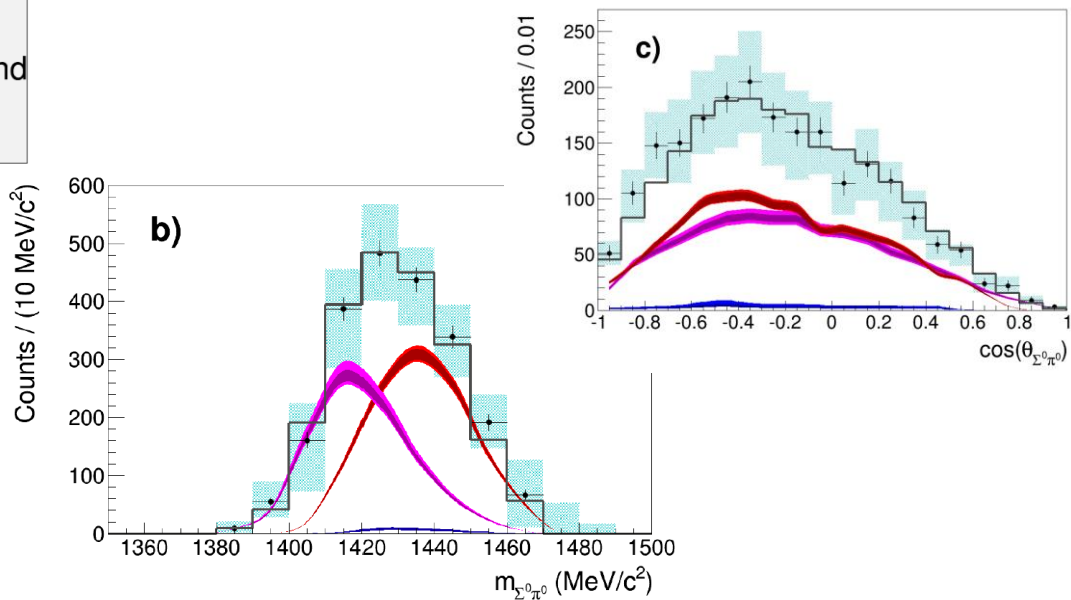


FIG. 3. From top to bottom the figure shows the result of the simultaneous fit of  $p_{\Sigma^0 \pi^0}$ ,  $m_{\Sigma^0 \pi^0}$  and  $\cos \theta_{\Sigma^0 \pi^0}$ . The experimental data and the corresponding statistical errors are represented by black crosses, the systematic errors are light blue boxes. The contributions of the various physical processes are shown as colored histograms, according to the color code shown in the caption. The light and dark bands correspond to systematic and statistical errors, respectively. The gray distribution reproduces the global fit function.



# Simultaneous fit and cross sections ( $\Lambda \pi^0$ )

+ the same processes initiated by primary  $\Sigma^0 \pi^0$   
and followed by  $\Sigma^0 \rightarrow \Lambda \gamma$

physical processes:

1.  $K^- H \rightarrow (\Sigma^0/\Lambda) \pi^0$  if (red),
2.  $K^- H \rightarrow (\Sigma^0/\Lambda) \pi^0$  ar (blue),
3.  $K^- + {}^4\text{He}/{}^{12}\text{C} \rightarrow \Sigma^0/\Lambda + \pi^0 + {}^3\text{H}/{}^{11}\text{B}$  (magenta).

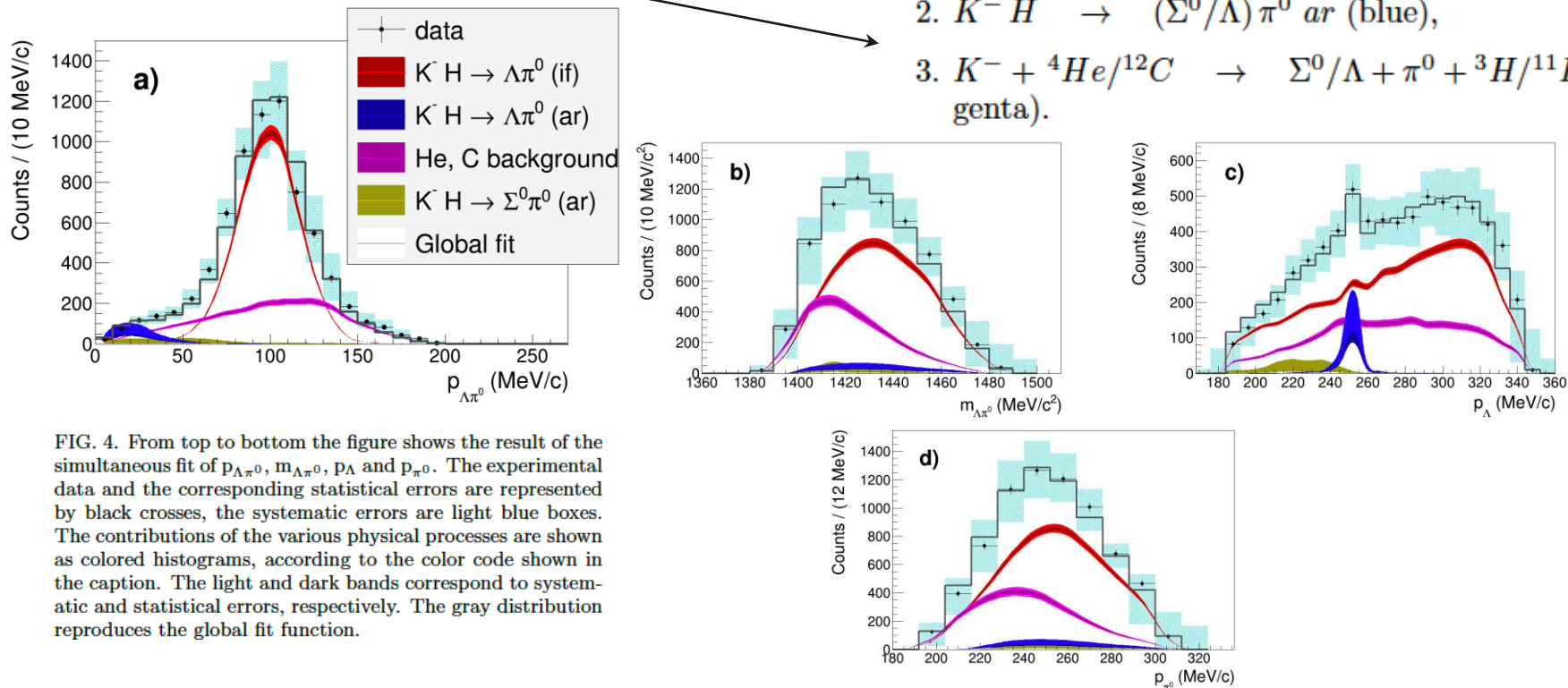


FIG. 4. From top to bottom the figure shows the result of the simultaneous fit of  $p_{\Lambda \pi^0}$ ,  $m_{\Lambda \pi^0}$ ,  $p_{\Lambda}$  and  $p_{\pi^0}$ . The experimental data and the corresponding statistical errors are represented by black crosses, the systematic errors are light blue boxes. The contributions of the various physical processes are shown as colored histograms, according to the color code shown in the caption. The light and dark bands correspond to systematic and statistical errors, respectively. The gray distribution reproduces the global fit function.

# Cross sections results

1 o.m. improvement on  
the relative error

cross section at  $p_{K^-} = 98 \pm 10$  MeV/c

:

$$\bullet \sigma_{K^-p \rightarrow \Sigma^0 \pi^0} = 42.8 \pm 1.5(\text{stat.})_{-2.0}^{+2.4}(\text{syst.}) \text{ mb}$$

$$\bullet \sigma_{K^-p \rightarrow \Lambda \pi^0} = 31.0 \pm 0.5(\text{stat.})_{-1.2}^{+1.2}(\text{syst.}) \text{ mb},$$

arXiv:2210.10342 [nucl-ex]

Accepted PRC

$\Sigma^0 - \pi^0$ CHANNEL $\frac{\chi^2}{(dof-np)} = \frac{92}{54} = 1.71$			
process	fit par.	value	$\sigma_{\text{stat.}}$
$K^-H \rightarrow \Sigma^0 \pi^0$ (if)	0.511		$\pm 0.018$
$K^-H \rightarrow \Sigma^0 \pi^0$ (ar)	0.017		$\pm 0.005$
$K^- + ^4\text{He}/^{12}\text{C} \rightarrow \Sigma^0 \pi^0$			
+ residual (ar/if)	0.463		$\pm 0.018$
$\Lambda - \pi^0$ CHANNEL $\frac{\chi^2}{(dof-np)} = \frac{165}{37} = 2.95$			
process	fit par.	value	$\sigma_{\text{stat.}}$
$K^-H \rightarrow \Lambda \pi^0$ (if)	0.659		$\pm 0.011$
$K^-H \rightarrow \Lambda \pi^0$ (ar)	0.021		$\pm 0.003$
$K^- + ^4\text{He}/^{12}\text{C} \rightarrow \Lambda \pi^0$			
+ residual (ar/if)	0.298		$\pm 0.012$
$K^-H \rightarrow \Sigma^0 \pi^0$			
$\rightarrow \Lambda \gamma \pi^0$ (ar)	0.018		$\pm 0.006$

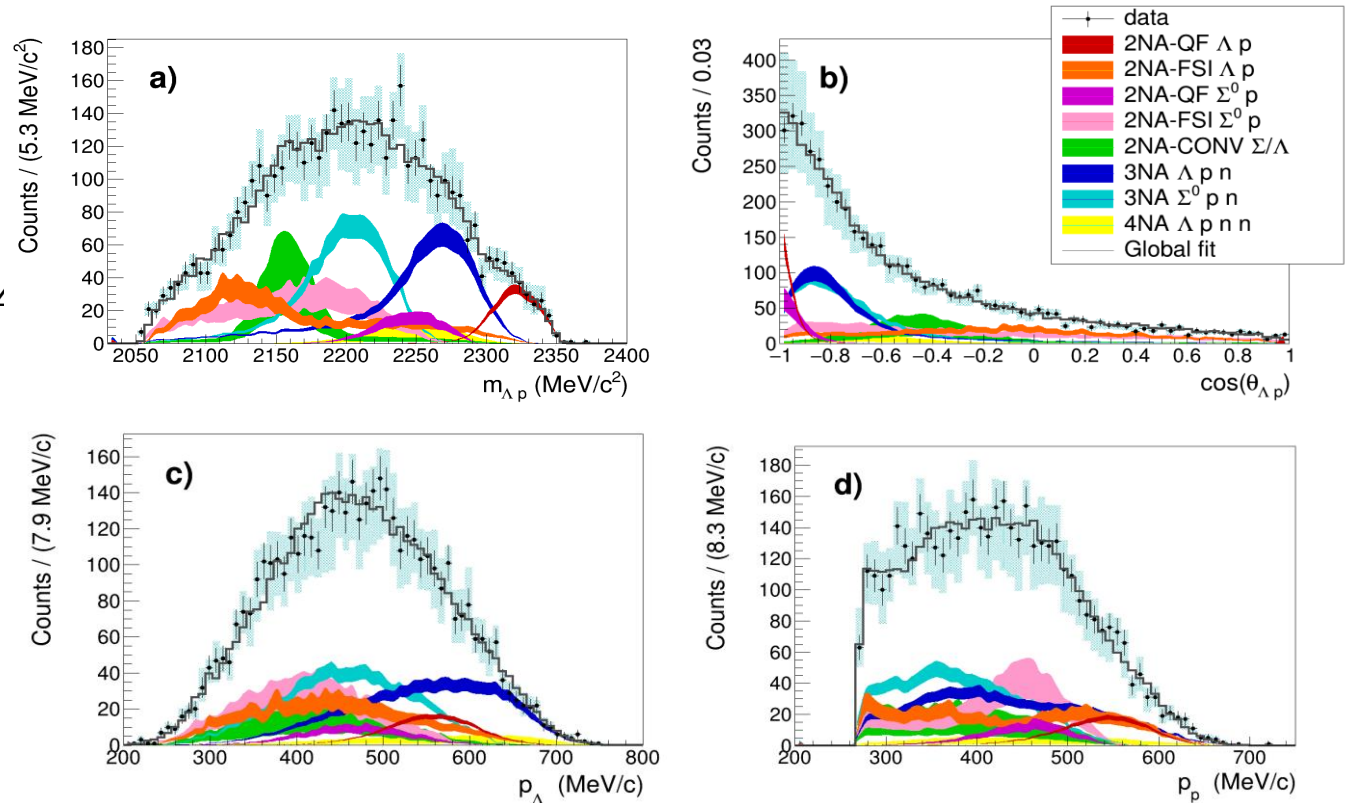
TABLE I. The table summarizes the results obtained from the fits of the  $\Sigma^0 \pi^0$  and  $\Lambda \pi^0$  samples. The values of the reduced chi-squares and of the fit parameters are summarized.

# $\Lambda p$ analysis: $K^- + {}^{12}\text{C} \rightarrow \Lambda + p + \text{R}$

Simultaneous fit of:

- $\Lambda p$  invariant mass;
- angular correlation;
- proton momentum;
- $\Lambda$  momentum.

Total reduced  $\chi^2$ :  $\chi^2/dof = 0.94$



[R. Del Grande, K. P., O. Vazquez Doce et al., Eur.Phys.J. C79 (2019) no.3, 190]

[R. Del Grande, K. P., S. Wycech, Acta Phys. Pol. B 48 (2017) 1881]

[O. Vazquez Doce, L. Fabbietti et al., Phys.Lett. B 758, 134-139 (2016)]

# $\Lambda p$ analysis: $K^-$ multi-nucleon absorption BRs and $\sigma$

[R. Del Grande, K. P., O. Vazquez Doce et al., Eur.Phys.J. C79 (2019) no.3, 190]

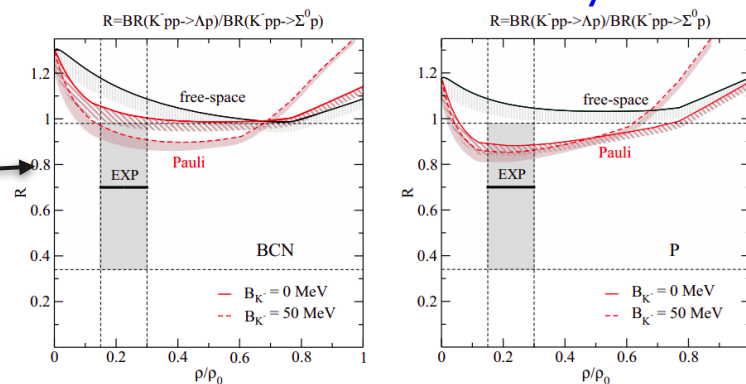
Process	Branching Ratio (%)	$\sigma$ (mb)	@	$p_K$ (MeV/c)
2NA-QF $\Lambda p$	$0.25 \pm 0.02$ (stat.) $^{+0.01}_{-0.02}$ (syst.)	$2.8 \pm 0.3$ (stat.) $^{+0.1}_{-0.2}$ (syst.)	@	$128 \pm 29$
2NA-FSI $\Lambda p$	$6.2 \pm 1.4$ (stat.) $^{+0.5}_{-0.6}$ (syst.)	$69 \pm 15$ (stat.) $\pm 6$ (syst.)	@	$128 \pm 29$
2NA-QF $\Sigma^0 p$	$0.35 \pm 0.09$ (stat.) $^{+0.13}_{-0.06}$ (syst.)	$3.9 \pm 1.0$ (stat.) $^{+1.4}_{-0.7}$ (syst.)	@	$128 \pm 29$
2NA-FSI $\Sigma^0 p$	$7.2 \pm 2.2$ (stat.) $^{+4.2}_{-5.4}$ (syst.)	$80 \pm 25$ (stat.) $^{+46}_{-60}$ (syst.)	@	$128 \pm 29$
2NA-CONV $\Sigma/\Lambda$	$2.1 \pm 1.2$ (stat.) $^{+0.9}_{-0.5}$ (syst.)	-		
3NA $\Lambda pn$	$1.4 \pm 0.2$ (stat.) $^{+0.1}_{-0.2}$ (syst.)	$15 \pm 2$ (stat.) $\pm 2$ (syst.)	@	$117 \pm 23$
3NA $\Sigma^0 pn$	$3.7 \pm 0.4$ (stat.) $^{+0.2}_{-0.4}$ (syst.)	$41 \pm 4$ (stat.) $^{+2}_{-5}$ (syst.)	@	$117 \pm 23$
4NA $\Lambda pnn$	$0.13 \pm 0.09$ (stat.) $^{+0.08}_{-0.07}$ (syst.)	-		
Global $\Lambda(\Sigma^0)p$	$21 \pm 3$ (stat.) $^{+5}_{-6}$ (syst.)	-		

The ratio between the branching ratios of the 2NA-QF in the  $\Lambda p$  channel and in the  $\Sigma^0 p$  is measured to be:

$$\mathcal{R} = \frac{BR(K^- pp \rightarrow \Lambda p)}{BR(K^- pp \rightarrow \Sigma^0 p)} = 0.7 \pm 0.2(\text{stat.})^{+0.2}_{-0.3}(\text{syst.})$$

and the ratio between the corresponding phase spaces is  $\mathcal{R}' \simeq 1.22$ .

## Information on the in-medium dynamics



[J. Hrtánková and A. Ramos. Phys. Rev. C, 101(3):035204, 2020]



# Total BR of the $K^-$ 2NA process in $^{12}\text{C}$

the only missing components are:

- $\text{BR}(\Sigma^- n) = (0.12 \pm 0.01(\text{syst.}))\%$
- $\text{BR}(\text{QF-}\Lambda n + \text{QF-}\Sigma^0 n) = (0.76 \pm 0.09(\text{stat.})^{+0.13}_{-0.06}(\text{syst.}))\%$
- $\text{BR}(\text{FSI-}\Lambda n + \text{FSI-}\Sigma^0 n) = (1.62 \pm 0.04(\text{stat.})^{+0.22}_{-0.21}(\text{syst.}))\%$
- $\text{BR}(\text{no conv } \Sigma^+ \text{ and } \Sigma^-) = (3.04 \pm 0.03(\text{stat.}) \pm 0.92(\text{syst.}))\%$

→  $(5.5 \pm 0.1(\text{stat.})^{+1.0}_{-0.9}(\text{syst.}))\%$

[R. Del Grande, K. P., et al., 2020 Phys. Scr.95 084012]

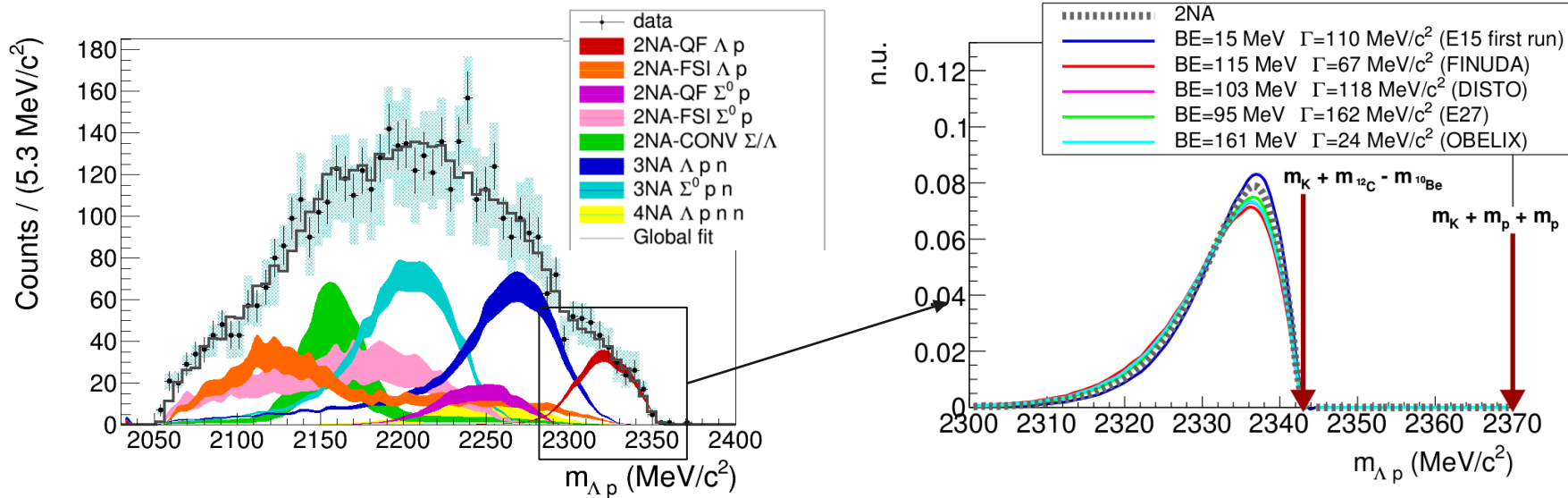
[R. Del Grande, K. P., et al., *Few Body Syst.* 62 (2021) 1, 7]

**Including the missing components the total BR of the  $K^-$ 2NA is:**

$$\text{BR}(K^- 2\text{NA} \rightarrow \text{YN}) = (21.6 \pm 2.9(\text{stat.})^{+4.4}_{-5.6}(\text{syst.}))\%$$

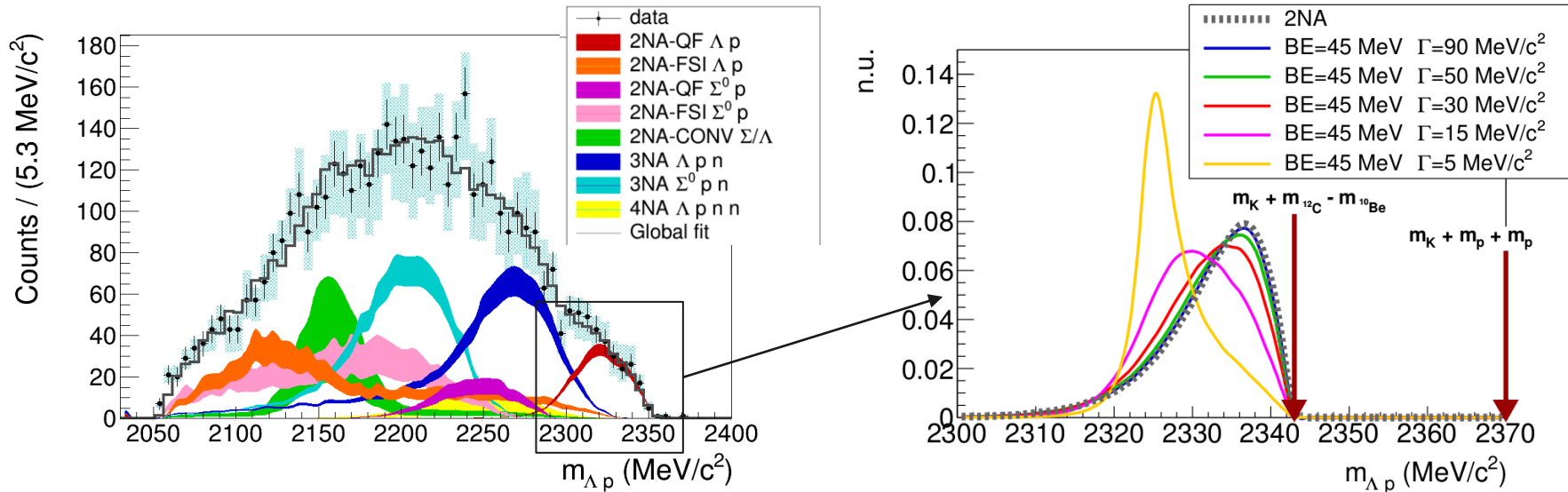
to be compared with [J. Hrtánková and A. Ramos. Phys. Rev. C, 101(3):035204, 2020]

# $\Lambda$ p analysis: $K^-$ pp bound state



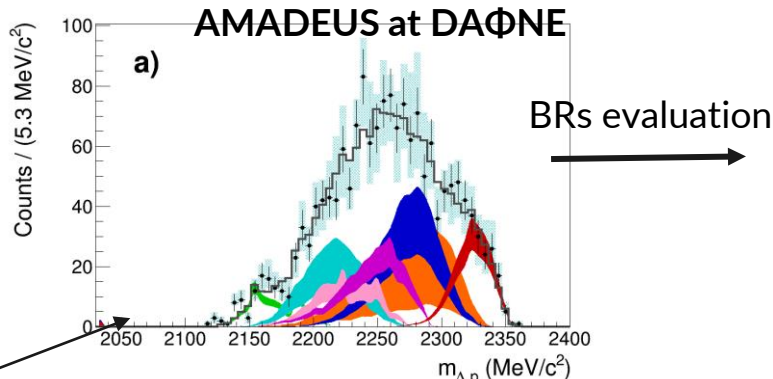
$K^-$ pp bound state contribution **completely overlaps** with the  $K^-$ 2NA

# $\Lambda p$ analysis: $K^- pp$ bound state



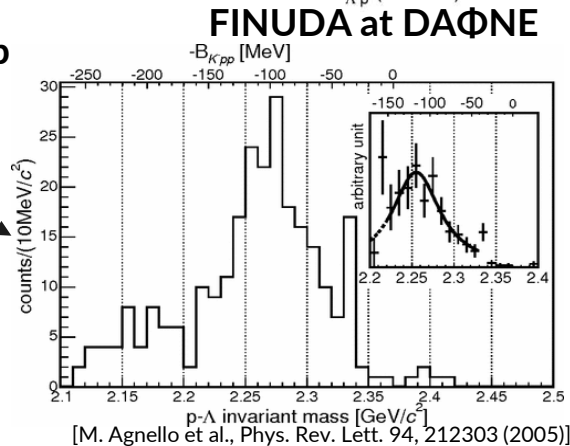
$K^- pp$  bound state contribution **completely overlaps** with the  $K^- 2NA$

# $\Lambda p$ analysis: $K^- pp$ bound state search



Process	Branching Ratio (%)
2NA-QF $\Lambda p$	$0.20 \pm 0.04(\text{stat.}) \pm 0.02(\text{syst.})$
2NA-FSI $\Lambda p$	$3.8 \pm 2.3(\text{stat.}) \pm 1.1(\text{syst.})$
2NA-QF $\Sigma^0 p$	$0.54 \pm 0.20(\text{stat.})^{+0.20}_{-0.16}(\text{syst.})$
2NA-FSI $\Sigma^0 p$	$5.4 \pm 1.5(\text{stat.})^{+1.0}_{-2.7}(\text{syst.})$
2NA-CONV $\Sigma/\Lambda$	$22 \pm 4(\text{stat.})^{+1}_{-12}(\text{syst.})$
3NA $\Lambda pn$	$1.1 \pm 0.3(\text{stat.}) \pm 0.2(\text{syst.})$
3NA $\Sigma^0 pn$	$1.9 \pm 0.7(\text{stat.})^{+0.8}_{-0.4}(\text{syst.})$

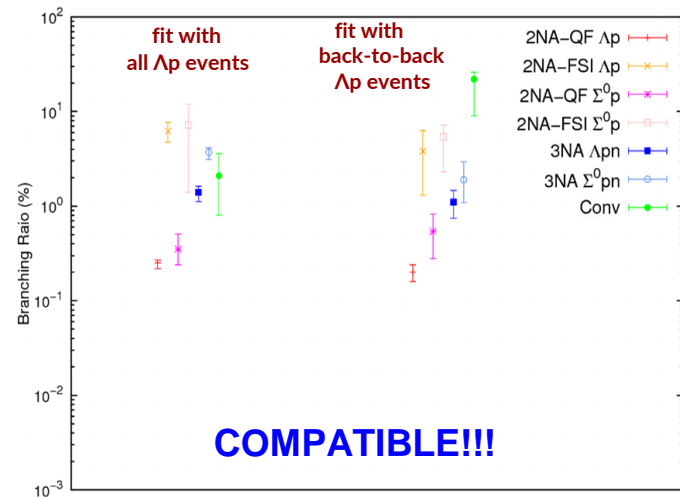
only **back-to-back**  $\Lambda p$  pairs ( $\cos\theta_{\Lambda p} < -0.8$ )



In agreement with

V.K. Magas, E. Oset, A. Ramos, H. Toki, Phys. Rev. C 74, 025206 (2006)

V.K. Magas, E. Oset, A. Ramos, Phys. Rev. C 77, 065210 (2008)



# $\Lambda$ t analysis: Cross section and BR for 4NA

GOLDEN CHANNEL to extrapolate the  $K^-$  4NA



## Previous data:

- in  $^4\text{He}$ : bubble chamber experiment

/M. Roosen, J. H. Wickens, II Nuovo Cimento 66, 101 (1981)/

only 3 events compatible with  $\Lambda$ t kinematics found

$$\text{BR}(K^- ^4\text{He} \rightarrow \Lambda t) = (3 \pm 2) \times 10^{-4} / K_{\text{stop}}^- \rightarrow \text{global, no 4NA}$$

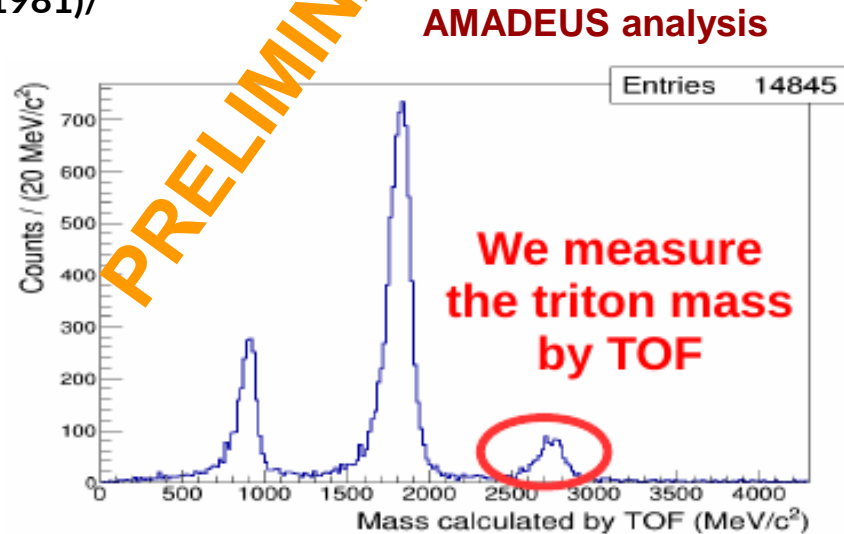
- in solid targets:  $^6,^7\text{Li}$ ,  $^9\text{Be}$  (FINUDA)

/Phys. Lett. B, 229 (2008)/

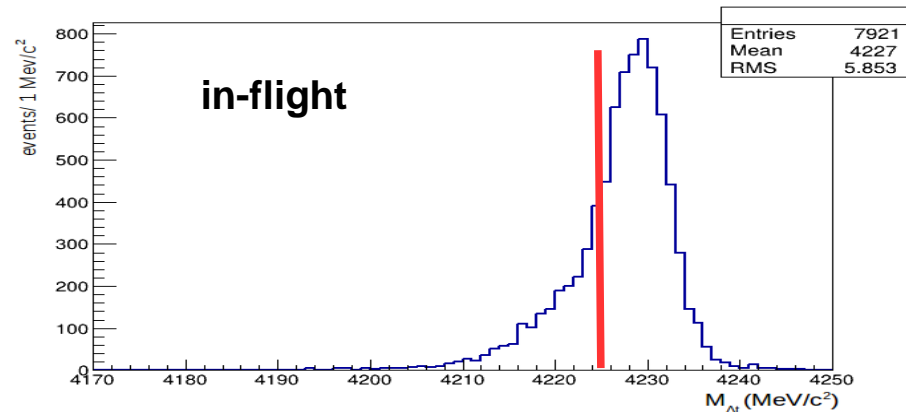
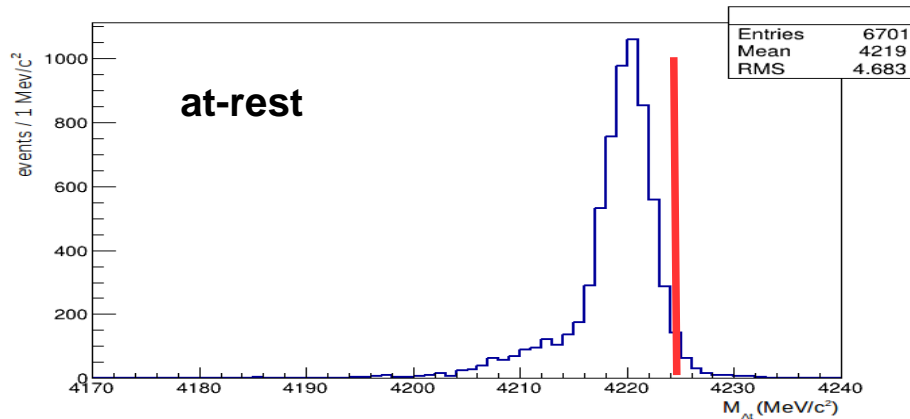
40 events, only back-to-back data

$$\Lambda t \text{ emission yield} \rightarrow 10^{-3} - 10^{-4} / K_{\text{stop}}^-$$

$\rightarrow$  global, no 4NA



# MC simulations: efficiency & resolution



**mass threshold at-rest**

**M<sub>Λt</sub> invariant mass resolution = 2.2 MeV/c<sup>2</sup>**

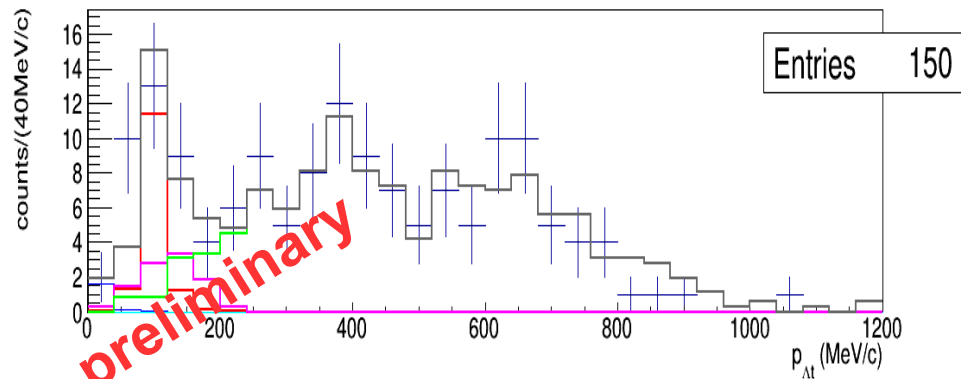
**overall detection + reconstruction efficiency for 4NA direct Λt production :**

$$\epsilon_{4NA,ar,\Lambda t} = 0.0493 \pm 0.0006 \quad ; \quad \epsilon_{4NA,if,\Lambda t} = 0.0578 \pm 0.0006,$$

**at-rest**

**in-flight**

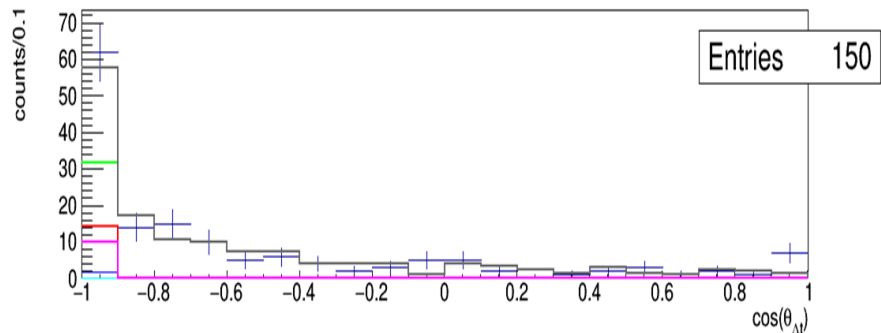
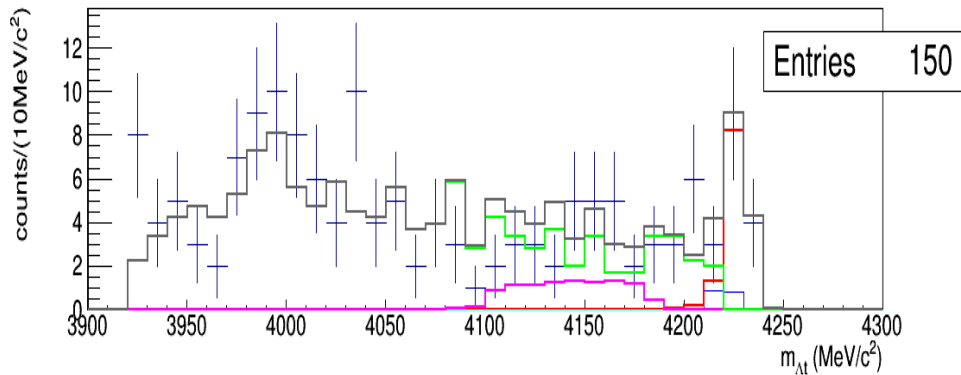
# Cross section and BR for 4NA in $K^- \text{}^4\text{He} \rightarrow \Lambda t$ process



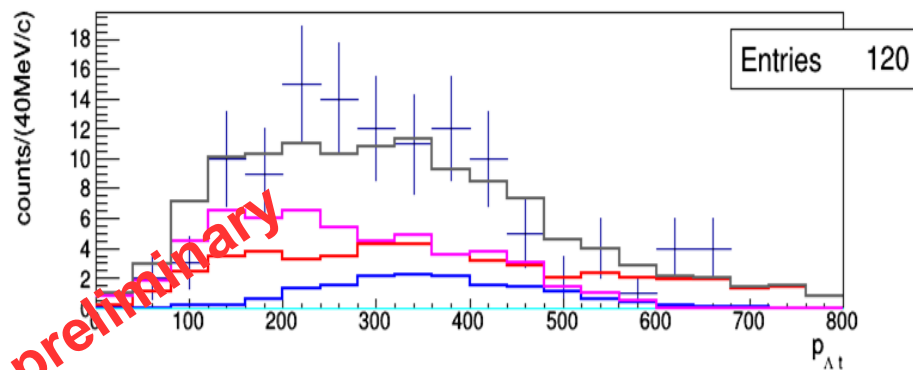
BR( $K^- \text{}^4\text{He} \rightarrow \Lambda t$ )  $< 2.0 \times 10^{-4} / K_{\text{stop}}$  (95% c.)

I.)

$\sigma(100 \pm 19 \text{ MeV/c}) (K^- \text{}^4\text{He} \rightarrow \Lambda t) =$   
 $= (0.81 \pm 0.21 \text{ (stat)}^{+0.03}_{-0.04} \text{ (syst)}) \text{ mb}$



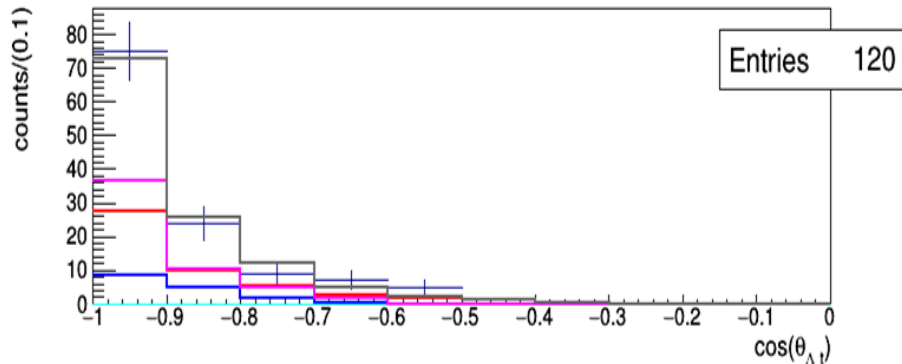
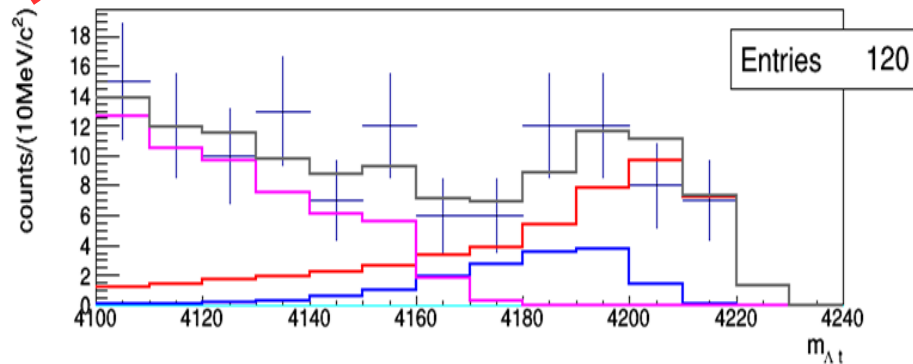
# Cross section and BR for 4NA in $K^- ^{12}\text{C} \rightarrow \Lambda/\Sigma^0 t$ processes



$$\text{BR}(K^- ^{12}\text{C}(4\text{NA}) \rightarrow \Lambda t \text{ } ^8\text{Be}) = 1.5 \pm 0.5 \times 10^{-4} \text{ (stat)}$$

$$\sigma(K^- ^{12}\text{C}(4\text{NA}) \rightarrow \Lambda t \text{ } ^8\text{Be}) = 0.58 \pm 0.11 \text{ (stat) mb}$$

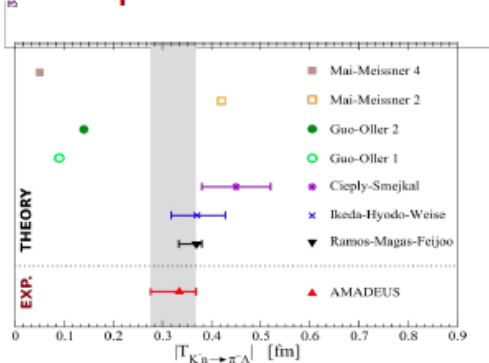
$$\sigma(K^- ^{12}\text{C}(4\text{NA}) \rightarrow \Sigma^0 t \text{ } ^8\text{Be}) = 1.88 \pm 0.35 \text{ (stat) mb}$$





# Highlights of AMADEUS results

## K<sup>-</sup>n amplitude below threshold



## Λp channel: 2NA, 3NA and 4NA BRs and σ

Process	Branching Ratio (%)	σ (mb)	@	p <sub>K</sub> (MeV/c)
2NA-QF Λp	0.25 ± 0.02 (stat.) <sup>+0.01</sup> <sub>-0.02</sub> (syst.)	2.8 ± 0.3 (stat.) <sup>+0.1</sup> <sub>-0.2</sub> (syst.)	@	128 ± 29
2NA-FSI Λp	6.2 ± 1.4 (stat.) <sup>+0.5</sup> <sub>-0.6</sub> (syst.)	69 ± 15 (stat.) ± 6 (syst.)	@	128 ± 29
2NA-QF Σ <sup>0</sup> p	0.35 ± 0.09 (stat.) <sup>+0.13</sup> <sub>-0.06</sub> (syst.)	3.9 ± 1.0 (stat.) <sup>+1.4</sup> <sub>-0.7</sub> (syst.)	@	128 ± 29
2NA-FSI Σ <sup>0</sup> p	7.2 ± 2.2 (stat.) <sup>+4.2</sup> <sub>-5.4</sub> (syst.)	80 ± 25 (stat.) <sup>+46</sup> <sub>-60</sub> (syst.)	@	128 ± 29
2NA-CONV Σ/Λ	2.1 ± 1.2 (stat.) <sup>+0.9</sup> <sub>-0.5</sub> (syst.)	-	-	-
3NA Λpn	1.4 ± 0.2 (stat.) <sup>+0.1</sup> <sub>-0.2</sub> (syst.)	15 ± 2 (stat.) ± 2 (syst.)	@	117 ± 23
3NA Σ <sup>0</sup> pn	3.7 ± 0.4 (stat.) <sup>+0.2</sup> <sub>-0.4</sub> (syst.)	41 ± 4 (stat.) <sup>+2</sup> <sub>-5</sub> (syst.)	@	117 ± 23
4NA Λpnn	0.13 ± 0.09 (stat.) <sup>+0.08</sup> <sub>-0.07</sub> (syst.)	-	-	-
Global Λ(Σ <sup>0</sup> )p	21 ± 3 (stat.) <sup>+5</sup> <sub>-6</sub> (syst.)	-	-	-

The ratio between the branching ratios of the 2NA-QF in the Λp channel and in the Σ<sup>0</sup>p is measured to be:

$$R = \frac{BR(K^-pp \rightarrow \Lambda p)}{BR(K^-pp \rightarrow \Sigma^0 p)} = 0.7 \pm 0.2 (stat.)^{+0.2}_{-0.3} (syst.)$$

$$BR(K^-2NA \rightarrow YN) = (21.6 \pm 2.9 (stat.)^{+4.4}_{-5.6} (syst.))\%$$

## K<sup>-</sup>p → (Σ<sup>0</sup>/Λ) π<sup>0</sup>

cross section at p<sub>K<sup>-</sup></sub> = 98 ± 10 MeV/c :

- $\sigma_{K^-p \rightarrow \Sigma^0 \pi^0} = 42.8 \pm 1.5 (stat.)^{+2.4}_{-2.0} (syst.)$  mb
- $\sigma_{K^-p \rightarrow \Lambda \pi^0} = 31.0 \pm 0.5 (stat.)^{+1.2}_{-1.2} (syst.)$  mb,

## Λt channel: 4NA BRs and σ

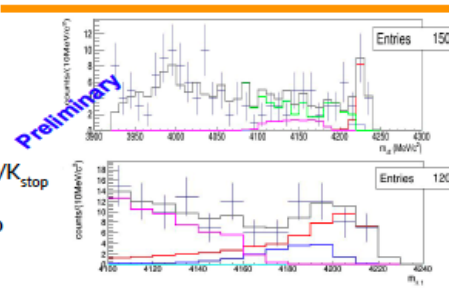
$$BR(K^4\text{He}(4NA) \rightarrow \Lambda t) < 2.0 \times 10^{-4} / K_{\text{stop}} \quad (95\% \text{ c. l.})$$

$$\sigma(100 \pm 19 \text{ MeV/c } (K^4\text{He}(4NA) \rightarrow \Lambda t) = (0.81 \pm 0.21 (stat.)^{+0.03}_{-0.04} (syst.) \text{ mb}$$

$$BR(K^{12}\text{C}(4NA) \rightarrow \Lambda t \text{ } ^8\text{Be}) = 1.5 \pm 0.5 \times 10^{-4} (stat) / K_{\text{stop}}$$

$$\sigma(K^{12}\text{C}(4NA) \rightarrow \Lambda t \text{ } ^8\text{Be}) = 0.58 \pm 0.11 (stat) \text{ mb}$$

$$\sigma(K^{12}\text{C}(4NA) \rightarrow \Sigma^0 t \text{ } ^8\text{Be}) = 1.88 \pm 0.35 (stat) \text{ mb}$$



# Future perspectives

- The present knowledge of total and differential cross sections of low energy kaon-nucleon reactions is **very limited**: below 150 MeV/c there is a “desert” - the experimental data are very scarce and with large errors and practically no data exist below 100 MeV/c.
- **Kaon-nucleon scattering/interaction data are fundamental to validate theories**: chiral symmetries; lattice calculations; potential models etc.

## New $\bar{K}N$ potentials, $K^-p$ scattering

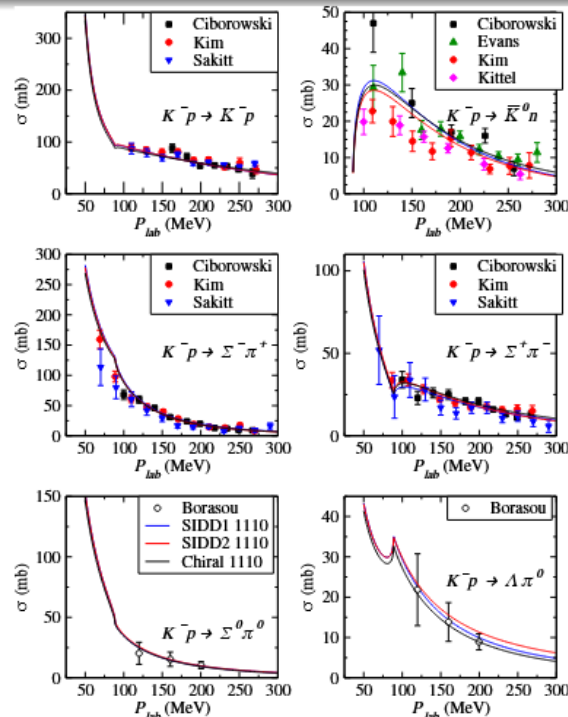
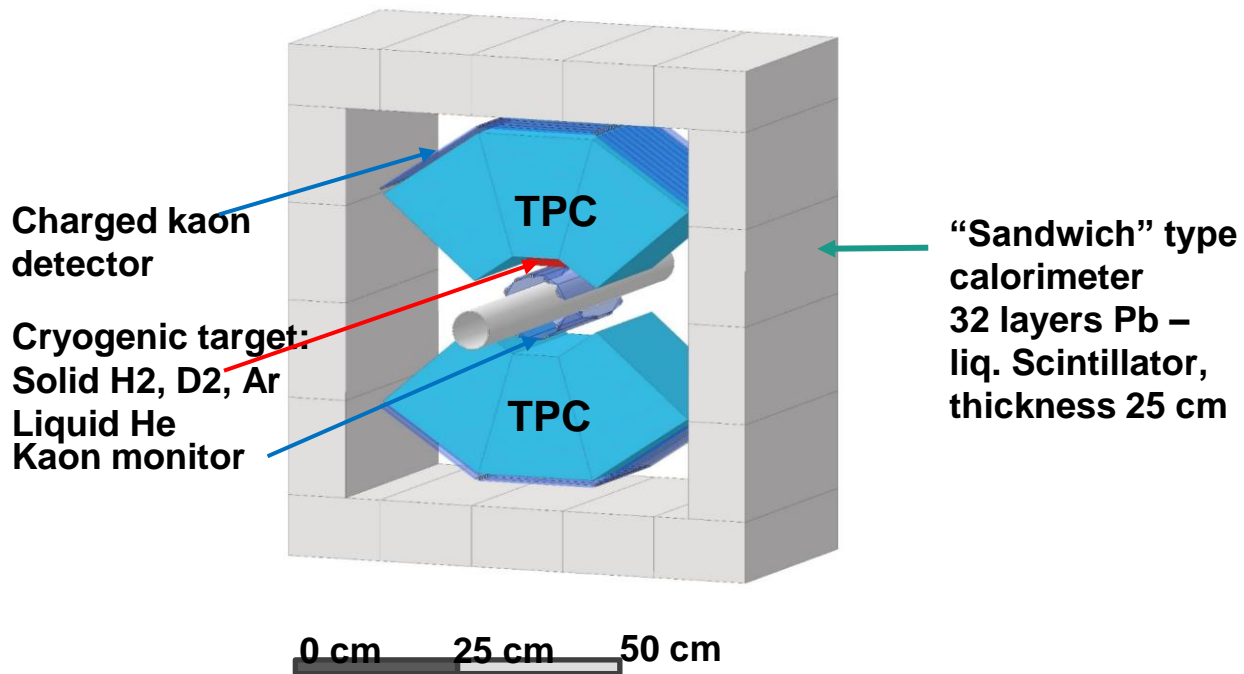


Figure: New  $V_{\bar{K}N}$  potentials: one-pole, two-pole phenomenological and chirally motivated

## Experiment for kaon-nuclei interaction studies



**Thank You**

***Edgar Allan Poe – 'There is no exquisite beauty...  
without some strangeness in the proportion.'***