

Fundação para a Ciência e a Tecnologia

QGP tomography Helena Santos LIP, FCUL



CMS Heavy Ion Workshop 2023, 29-30 May, Trento

Experimental observables reviewed in this talk

How do jets allow us to see the QGP and vice-versa?

• Do jets disturb the QGP measurably?

Explore:

- the jet substructure
- the path length dependence of quenching
- the recoil region







Searching for the resolution scale

Substructure techniques like soft-drop or reclustering using subjets allow to experimentally access scales associated with individual branchings in the parton shower.









Jet substructure - rg



Jets are narrower with increasing p_{T} , independently on centrality





Jet substructure - r_g



- Jets are narrower with increasing p_{T} , independently on centrality
- Jets with wider hard splittings are significantly more suppressed in central collisions Result points to decoherent energy loss

JETSCAPE compatible with exception for very low r_q (underestimates the suppression)

Caucal et al. overestimates the suppression for intermediate r_g







Jet substructure - *z*_g







Jet substructure - θ_g



- Enhancement of narrow splittings and suppression of wider ones. No difference for R=0.4 jets
- Hybrid model assuming fully incoherent energy loss reproduces the data for $\theta_{q} > 0.2$
- Yuan (quark only) describes the data as well





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0.3 $\theta_{g}^{0.8}$



Sensitivity to medium response

JHEP 05 (2021) 116



Jet radial momentum distributions

$$P(\Delta r) = \frac{1}{\delta r} \frac{1}{N_{\text{jet}}} \Sigma_{\text{jets}} \Sigma_{\text{trk} \in (\Delta r_{a}, \Delta r_{b})} p_{\text{T}}^{\text{trk}}$$

Jet shapes
$$\rho(\Delta r) = \frac{P(\Delta r)}{\sum_{jets} \sum_{tracks \in \Delta r < 1} p_T^{ch}}$$



Sensitivity to medium response





Jet radial momentum distributions

$$-\Sigma_{\text{jets}}\Sigma_{\text{trk}\in(\Delta r_a,\Delta r_b)}p_{\text{T}}^{\text{trk}}$$
et

et shapes
$$\rho(\Delta r) = \frac{P(\Delta r)}{\sum_{jets} \sum_{tracks \in \Delta r < 1} p_T^{ch}}$$

Angular scan of jet-QGP interactions



Jet radial momentum flow







One common explanation is that the energy lost at high p_T resulting from interactions between partons of both jet and QGP reappears in the form of low-*p*T particles far away from the jet axis.





















Difference of p_T profiles in pp and PbPb

For inclusive and *b*-jets



- $(\Delta r \sim 0.3 - 0.4 \text{ for } R = 0.4 \text{ jets}).$
- strongly enhanced for *b*-jets

In PbPb more p_T from low momentum particles is accumulated at the edge of the jet

• Core of the jet is unmodified, independently on the flavour, but then the redistribution is





Jet shape in dijets as a function of x_i



- Enhancement of low- p_T charged particles in central PbPb relative to pp at large Δr .
- For leading jets the enhancement is larger for balanced dijets (no surface bias).

anti-k	5.02 TeV	pp 320 pb ⁻¹	PbPb 1.7 i	nb ⁻¹
	_T R = 0.4, η _{jet}	< 1.6, p _{T,1} > 120 0	GeV, p _{T,2} > 50 0	GeV, Δφ _{1,2} > <u>5π</u>
eV	4 < p	o ^{ch} < 8 GeV	12 <	p ^{ch} < 300 GeV
eV		o ^{ch} < 12 GeV	0.7 <	p ^{ch} < 300 GeV





Jet shape in dijets as a function of x_i



- Enhancement of low- p_T charged particles in central PbPb relative to pp at large Δr .
- For leading jets the enhancement is larger for balanced dijets (no surface bias).
- Higher modification for sub-leading jets in unbalanced dijets.





Jet shape ratios as a function of x_i



For subleading jets a clear effect at $\Delta r \simeq 0.3$ is observed, explained by the presence of a 3rd recoiled jet in *pp* that eventually disappeared in PbPb.

Enhancement of the PbPb leading (and subleading) jet shape modification at large Δr in balanced dijets.

Significant effects from different path length crossed.

Jet shape ratios as a function of x_i

Besides, the jet core remains unmodified. Then depletion of particles at low radial region. This depletion is there independently on centrality (backup).

For subleading jets a clear effect at $\Delta r \simeq 0.3$ is observed, explained by the presence of a 3rd recoiled jet in *pp* that eventually disappeared in PbPb.

Enhancement of the PbPb leading (and subleading) jet shape modification at large Δr in balanced dijets.

Significant effects from different path length crossed.

Z-recoil particles - *I*_{AA} as a function of *p*_T^{ch}

The p_T spectrum of charged particles recoinling against Z is strongly modified in PbPb.

SCET_G (no medium response) describes CMS data but not available at low- p_T^{trk} . Hybrid w/o wake disagrees. Global improvement if including wake. Medium response is needed to describe the excess at low- p_T^{trk} .

- CollBT (medium feeds back the quenched energy deposited) describes well CMS and ATLAS.

Sensitivity to the wake and diffusion wake — $\Delta \varphi_{hZ}$

Looking forward to jet-hadron correlations in γ/Z +jet events

Enhancement of recoiling particles for $p_T^{ch} > 1$ GeV in CMS and $p_T^{ch} > 2 \text{ GeV in ATLAS (for } p_T^Z > 60$ GeV only).

CMS data supports medium response predicted by ColLBT. Hybrid with wake describes the data in the away-side region but not in the near side, for which predicts a depletion of particles in PbPb.

Conclusions and outlook

- branching width.
- Enhancement of low- p_T particles at the edge of the jet supporting medium response is observed. Jet core remains unmodified.

the QGP.

But

. . . .

In a steep falling p_T spectrum ($\sigma \sim p_T^6$) the population in each bin is composed essentially by jets that suffered few or no energy loss (e.g. 1812.05111 and references therein). Comparing observables as a function of reconstructed jet p_T ranges imply: - lower p_T ranges have larger mixing of jets that born differently. More than a physics feature, it is a mess and makes data interpretation very model dependent.

-> enriched samples of jets that were modified are needed for complex observables.

• LHC data show evidence for decoherent energy loss. Jet quenching depends strongly on the

We are starting to relate the jet modifications observed in the experiments with the properties of

Backup

Jet reconstruction in Pb+Pb collisions

Figure 1: Fraction of photon-tagged jets (filled markers) and inclusive jets (open markers) initiated by a quark, as a function of p_T^{jet} , in the PYTHIA (red), HERWIG (black), and SHERPA (blue) event generators.

Quark jet fraction

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ATLAS-CONF-2022-019

Jet substructure uncertainties

ATLAS-CONF-2022-026

(second and third panels). The legend applies to all the panels.

Figure 3: The relative systematic uncertainties on inclusive r_g cross-section and per-event jet yield measurements in pp collisions at $\sqrt{s} = 5.02$ TeV (left) and for different event centralities in Pb+Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV (second and third panels) shown for soft-drop parameters $z_{cut} = 0.2$ and $\beta = 0$. The legend applies to all the panels.

Figure 2: The relative systematic uncertainties on inclusive p_{T}^{jet} cross-section and per-event jet yield measurements in pp collisions at $\sqrt{s} = 5.02$ TeV (left) and for different event centralities in Pb+Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV

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Jet substructure

Inclusive jets

 $R_{AA}/R_{AA}^{0.2}$ for R=0.3, 0.4, 0.6, 0.8, 1.0

on jet radii.

2102.120

- The suppression for jet $p_T > 300$ GeV does not depend
- -> Look for two-prong structures as they might be sensitive to coherence effects in the QGP.

J. Casalderrey-Solana, Y. Mehtar-Tani, C. A. Salgado, K. Tywoniuk, PLB725 (2013) 357 Helena Santos, HI CMS Workshop, Trento 2023

30 [hep-ex]

24 Large R jets reconstructed from R = 0.2 jets

M. Rybar QM19

conventional R=1.0.

Trimming & 35 GeV threshold remove soft

Jet substructure - reclustered large jets (R = 1.0)

R = 0.2 jets with $p_T > 35$ GeV reclustered into anti-k_t R = 1.0

study of k_t splitting scale

Recluster jets and remove soft contributions

- Significant change of the R_{AA} magnitude between jets with SSJ and those with more complex substructure
- Then R_{AA} is not dependent on $\sqrt{d_{12}}$
- Result points to decoherent energy loss

How do particles redistribute within the jet and beyond? 26

Study *FF* as a function of the angular distance between the charged particle and the jet axis.

In central collisions $R_{D(pT,r)}$ is above unity at all r for all $p_T < 4 \text{ GeV} \longrightarrow \text{Energy lost by jets is being transferred to}$ particles with $p_T < 4$ GeV with larger radial distance. Expected from jet FF.

Depletion of high- p_{T} particles over all range.

Jet core remains unmodified.

Sensitivity to angle between jet axes

SD: axis defined by the particles

left after grooming.

WTA: merged branch takes the p_T sum of the two sub-branches and the direction of the hardest one. Insensitive to soft radiation.

ΔR distributions are narrower in PbPb Grooming (zcut = 0.2, β =0) doesn't change the jet axis visibly. Not sensitive to grooming. MATTER+LBT and Hybrid ($L_{res}=0$) consistent with the data. Intra-jet p_T broadening ruled out.

JHEP 05 Leading and subleading jet shape ratios $\rho(\Delta r)PbPb/\rho(\Delta r)pp$, in different centrality and x_j ranges

Significant effects from different path length crossed.

Jet shape ratios in dijets

Clear effect at $\Delta r \approx 0.3$, and for imbalance dijets (x_i < 0.8). Enhancement at large Δr less pronounced in imbalanced dijets. Opposite effect in balanced.

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Z-recoil particles - I_{AA} as a function of p_T^{ch}

yield in PbPb / yield in pp

Sensitivity to the wake and diffusion wake — $\Delta \varphi_{hz}$

Looking forward to jet-hadron correlations in γ/Z +jet events

CoLBT, Xin-Nian et al, PRL 130, 052301 (2023)

at the photon side

The DF-wake valley on top of the MPI ridge gives rise to a double peak feature in the rapidity distribution of the jet-hadron correlation.

