An hyperfine view to the TPE



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From the 2S-2P to HFS measurements



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Hyperfine splitting theory and goals

Measure

the 1S-HFS in μ p with 1-2 ppm accuracy

Goals

- TPE contribution with 3x10⁻⁴ rel. accuracy
- Zemach radius and polarisability contributions

$\Delta E_{\rm HFS}^{\rm th} = 183.788(7) + 1.0040 \Delta E_{\rm TPE} \,[{\rm meV}]$

Pineda & Peset (2017)

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The principle of the μp HFS experiment



- Laser pulse: $\mu p(F=0) \longrightarrow \mu p(F=1)$
- Collision: $\mu p(F=1) + H_2 \longrightarrow H_2 + \mu p(F=0) + E_{kin}$
- Diffusion: the faster μp reach the target walls
- Resonance: plot number of X-rays vs. frequency



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• muon beam

thermalisation

- laser excitation
- diffusion

X-ray detection



- Stopping probability & Target • Length : 1-3 mm • Pressure : 0.5-2 bar
 - Temperature : 30-50 K

Stopping probability: 10-20 %





- Stopped μ⁻ form μp in highly excited state
- During the de-excitation to the ground state, the µp win kinetic energy
- μp thermalise through collision with H₂ gas
- A considerable fraction of µp reach the target walls prior the laser pulse

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• muon beam

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laser excitation

diffusion

Optical Bloch equations $\frac{d\rho_{11}}{dt} = -\frac{d\rho_{22}}{dt},$ $\frac{d\rho_{22}}{dt} = -\frac{i}{2} \left(\varpi \rho_{12} e^{i(\omega_r - \omega)t} - \varpi^* \rho_{12}^* e^{-i(\omega_r - \omega)t} \right) - \Gamma \rho_{22},$ $\frac{d\rho_{12}}{dt} = \frac{i\varpi^*}{2} (1 - 2\rho_{22}) e^{-i(\omega_r - \omega)t} - \frac{\Gamma'}{2} \rho_{12},$





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X-ray detection

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• muon beam

- target
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X-ray detection



Thermalised µp close to the target wall may diffuse to to the target walls in the signal time window,

⇒ intrinsic background

Signal and background simulations for optimistic laser fluence



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X-ray emission in the μ Au de-excitation

Transition (n→n')	Energy	Probability		
<u>2</u> →1	5.6 MeV	90%		
3→2	2.4 MeV	84%		
4→3	0.9 MeV	76%		

• target



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• muon beam

Requirements:

- ▶ X-ray detection eff. > 50%
- False identification of $\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$ as X-ray < 1x10⁻³

Beam line background suppression (to be investigated)



target

thermalisation

laser excitation

diffusion

X-ray detection

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The challenge: the laser system



Simulations of cavity started...



Challenges

- Wavelength: 6.7 μm
- Cryogenic temperatures
- Large laser fluence
- Toroidal geometry

(sparse laser technology)

- (coating stability?)
- (damage threshold, laser energy, reflectivity)

Summary of systematic





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HFS contributions and uncertainties



$$\Delta E_{\rm TPE} = \Delta E_{\rm Z} + \Delta E_{\rm Recoil} + \Delta E_{\rm pol}$$

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Two main ways to the TPE



Dispersion relation+ data: $g_1(x,Q^2), g_2(x,Q^2), F_1, G_E...$

Chiral EFT Chiral + dispersion Carlson, Vanderhaeghen, Martynenko, Tomalak, Pascalutsa

Pascalutsa, Pineda, Peset Hagelstein



pion-production cross section:

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J. M. Alarcon, V. Lensky, V. Pascalutsa, Eur. Phys. J. C 74, 2852 (2014).

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TPE: dispersion based approach

$$\begin{split} & \mathsf{Elastic part (Zemach)} \\ & \Delta_{Z} = \frac{8Z\alpha m_{r}}{\pi} \int_{0}^{\infty} \frac{\mathrm{d}Q}{Q^{2}} \left[\frac{G_{E}(Q^{2})G_{M}(Q^{2})}{1+\kappa} - 1 \right] \equiv -2Z\alpha m_{r}R_{Z}, \\ & \mathsf{Distler, Bernauer} \\ \\ & \mathsf{Recoil finite-size} \\ & \Delta_{\mathrm{recoil}} = \frac{Z\alpha}{\pi(1+\kappa)} \int_{0}^{\infty} \frac{\mathrm{d}Q}{Q} \left\{ \frac{8mM}{Q} \frac{G_{M}(Q^{2})}{Q^{2}} \left(2F_{1}(Q^{2}) + \frac{F_{1}(Q^{2}) + 3F_{2}(Q^{2})}{(v_{l}+1)(v+1)} \right) \\ & - \frac{8m_{r}G_{M}(Q^{2})G_{E}(Q^{2})}{Q} - \frac{m}{M} \frac{5+4v_{l}}{(1+v_{l})^{2}} F_{2}^{2}(Q^{2}) \right\}. \end{split}$$

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Alternative approach

$$\Delta_{Z} = \frac{4\alpha m_{r}Q_{0}}{3\pi} \left(-r_{E}^{2} - r_{M}^{2} + \frac{r_{E}^{2}r_{M}^{2}}{18}Q_{0}^{2} \right) + \frac{8\alpha m_{r}}{\pi} \int_{Q_{0}}^{\infty} \frac{\mathrm{d}Q}{Q^{2}} \left(\frac{G_{M}\left(Q^{2}\right)G_{E}\left(Q^{2}\right)}{\mu_{P}} - 1 \right)$$

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Polarisability

$$\Delta_{\text{pol.}} = \frac{Z\alpha m}{2\pi(1+\kappa)M} \left[\delta_1 + \delta_2\right] =$$

with:

$$\begin{split} \delta_{1} &= 2 \int_{0}^{\infty} \frac{\mathrm{d}Q}{Q} \left(\frac{5 + 4v_{l}}{(v_{l}+1)^{2}} \left[4I_{1}(Q^{2})/Z^{2} + F_{2}^{2}(Q^{2}) \right] + \frac{8M^{2}}{Q^{2}} \int_{0}^{x_{0}} \mathrm{d}x \, g_{1}(x, Q^{2}) \\ &\left\{ \frac{4}{v_{l} + \sqrt{1 + x^{2}\tau^{-1}}} \left[1 + \frac{1}{2(v_{l}+1)(1 + \sqrt{1 + x^{2}\tau^{-1}})} \right] - \frac{5 + 4v_{l}}{(v_{l}+1)^{2}} \right\} \right), \\ &= 2 \int_{0}^{\infty} \frac{\mathrm{d}Q}{Q} \left(\frac{5 + 4v_{l}}{(v_{l}+1)^{2}} \left[4I_{1}(Q^{2})/Z^{2} + F_{2}^{2}(Q^{2}) \right] - \frac{32M^{4}}{Q^{4}} \int_{0}^{x_{0}} \mathrm{d}x \, x^{2}g_{1}(x, Q^{2}) \\ &\left\{ \frac{1}{(v_{l} + \sqrt{1 + x^{2}\tau^{-1}})(1 + \sqrt{1 + x^{2}\tau^{-1}})(1 + v_{l})} \left[4 + \frac{1}{1 + \sqrt{1 + x^{2}\tau^{-1}}} + \frac{1}{v_{l} + 1} \right] \right\} \right) \end{split}$$
 Need also g_{1}, g_{2}
$$\delta_{2} &= 96M^{2} \int_{0}^{\infty} \frac{\mathrm{d}Q}{Q^{3}} \int_{0}^{x_{0}} \mathrm{d}x \, g_{2}(x, Q^{2}) \left\{ \frac{1}{v_{l} + \sqrt{1 + x^{2}\tau^{-1}}} - \frac{1}{v_{l} + 1} \right\} \cdot \quad \text{Hagelstein, Pascalutsa, Carlson, Martynenko, Tomalak Faustov, Vanderhaegen....}$$

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TPE contribution on the market

	Δ , (ppm)	Δ_{Z}	$\Delta^{\mathrm{p}}_{\mathrm{R}}$	$\Delta_Z + \Delta_R^p$	Δ_0^{pol}	$\Delta_{ m HFS}$
	this work, $\mu H r_E, r_M^W$	-7415(84)	844(7)	-6571(87)	364(89)	-6207(127)
alak	this work, electron r_E , r_M^W	-7487(95)	844(7)	-6643(98)	364(89)	-6279(135)
Tom	this work, $\mu H r_E, r_M^e$	-7333(48)	846(6)	-6486(49)	364(89)	-6122(105)
	this work, electron r_E , r_M^e	-7406(56)	847(6)	-6559(57)	364(89)	-6195(109)
	Hagelstein et al. $[59]$				-61^{+70}_{-52}	
	Peset et al. $[29]$					-6247(109)
	Carlson et al. $[28, 39]$	-7587	835	-6752(180)	351(114)	-6401(213)
	Martynenko et al. $[38]$	-7180		-6656	410(80)	-6246(342)
	Pachucki [7]	-8024		-6358	0(658)	-6358(658)

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All dispersive approaches needs to be re-evaluated with the new g₁ and g₂ data

TPE contribution on the market

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Interesting tension between dispersion-based and ChPT predictions

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TPE: from H to µp

Recently predictions of the TPE contribution has been achieved scaling the results from the TPE measured in H with m_r and correcting for small deviations from this scaling.

Extract hydrogen TPE

$$\Delta E_{\rm HFS}^{\rm th}(H) = \Delta E_{\rm QED}^{\rm th}(H) + \Delta E_{\rm TPE}(H)$$
$$\Delta E_{\rm HFS}^{\rm th}(H) = \Delta E_{\rm HFS}^{\rm exp}(H)$$

Scale TPE from H to μp

Pineda & Peset Tomalak

$$\implies \Delta E_{\rm TPE}(H)$$

- HFS in H measured with 7x10⁻¹³ rel. acc.
- TPE contribution in H: 50 ppm of HFS

 $\Delta E_{\text{TPE}}(H) \implies \Delta E_{\text{TPE}}^{\text{th}}(\mu p) = \text{scaling}(\Delta E_{\text{TPE}}(H)) + \varepsilon$ $\bullet \text{ Model independent}$ $\bullet \text{ Smaller uncertainties than from "direct "calculations"}$

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Recent values of the TPE

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Recent values of the TPE

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Meson exchanges

Hagelstein & Pascalutsa, arXiv 1511.0430

Cancellation: contribution is very small

Unanticipated large contributions. Needs to be verified by independent group. Already accounted in TPE?

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Dorokhov et al., PLB 776, 105 (2018)

Uncertainties and scanning range

Large BG/Signal ratio

Uncertainties and scanning range

Uncertainties and scanning range

Conclusions: wish list to find the line

QED

- check QED contributions in H to improve the TPE(H)
- higher-order QED corrections in µp
- Summary of all contributions would be very helpful (at 1 ppm level).

Is the meson exchange already included in the TPE computed with dispersion relations?

Zemach radius

- improve determination of Zemach radius, mainly through magnetic FF
- Study correlations R_z vs R_p

Polarisability contribution

- re-evaluate the pol contribution given the new g₁ and g₂ data
- improve chPT prediction also in view of interpretation of HFS measurement
- subtraction term really absent?

A TPE contribution with an accuracy of 25 ppm of HFS is needed to find the line