

Results from the JLab CLAS EG4 experiment

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The EG4 experiment Group

Main goal: measurement of the generalized **Gerasimov-Drell-Hearn** (GDH) sum for the **proton** and **deuteron** at low Q^2 .

E03-006 (NH₃):

Spokespeople: M. R., M. Battaglieri, A. Deur, R. de Vita

Students: H. Kang (Seoul U.), K. Kovacs[◆] (UVa)

E06-017 (ND₃)

Spokespeople: A. Deur, G. Dodge, M. R., K. Slifer

Students: K. Adhikari[◆] (ODU)

EG4 ran from Feb. to May 2006.

Main goal: inclusive analyses. Also, exclusive analysis by X. Zheng

X. Zheng et al. (CLAS Collaboration), PRC 94, 045206 (2016)

◆ Graduated.

The GDH and Generalized GDH Sum Rules

Sum rule: **relation** between an **integral** of a dynamical quantity (cross section, structure function,...) and a global property of the target (mass, spin,...).

Can be used to:

- Test theory (e.g. QCD) and hypotheses with which they are derived. Ex: GDH, Ellis-Jaffe, Bjorken sum rules.
- Measure the global property (e.g. spin polarizability sum rules)

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GDH sum rule: derived for real photons ($Q^2=0$):

$$\int_{\nu_{\text{thr}}}^{\infty} \frac{\sigma_A(\nu) - \sigma_P(\nu)}{\nu} d\nu = \frac{-4\pi^2 S \alpha \kappa^2}{M^2}$$

The equation is annotated with the following labels:

- fine structure constant**: points to α
- target anomalous magnetic moment**: points to S
- target spin**: points to S
- target mass**: points to M^2
- photon spin parallel to S**: points to $\sigma_P(\nu)$
- photon spin anti-parallel to S**: points to $\sigma_A(\nu)$
- photoprod. cross section with photon spin anti-parallel to S**: points to $\sigma_A(\nu)$

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fine structure constant
 target anomalous magnetic moment
 target mass
 target spin
 photon spin anti-parallel to S
 photoprod. cross section with photon spin anti-parallel to S

Generalized GDH sum rule: valid for any Q^2 . Recover the original GDH sum rule at $Q^2 = 0$

$$\Gamma_1(Q^2) = \int_0^{x_{\text{th}}} g_1(x, Q^2) dx = \frac{Q^2}{2M^2} I_1(0, Q^2)$$

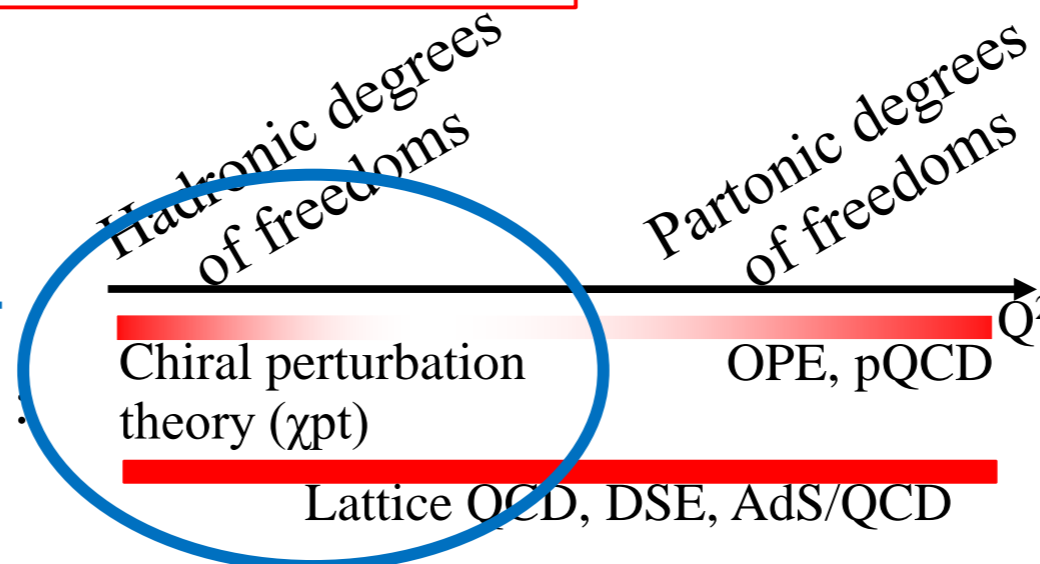
$g_1(\nu, Q^2)$: first spin structure function (mostly a longit. target pol. observable)

$I_1(\nu, Q^2)$: first covariant polarized VVCS amplitude

⇒ Study QCD at any scale

EG4

$I_1(0, Q^2)$



Spin Polarizabilities Sum Rules

Sum rule: relation between an **integral** of a dynamical quantity (cross section, structure function,...) and a global property of the target (mass, spin,...).

Can be used to:

- Test theory (e.g. QCD) and hypotheses with which they are derived. Ex: GDH, Ellis-Jaffe, Bjorken sum rules.
- Measure the global property (e.g. spin polarizability sum rules)

Spin polarizability sum rules involve higher moments:

Generalized forward spin polarizability:

$$\gamma_0 = \frac{4e^2 M^2}{\pi Q^6} \int x^2 \left(g_1 - \frac{4M^2}{Q^2} x^2 g_2 \right) dx = \int x^2 A_1 F_1 dx$$

$g_2(\nu, Q^2)$: second spin structure function (mostly a perp. target pol. observable)

Contribution suppressed in γ_0

Longitudinal-transverse spin polarizability:

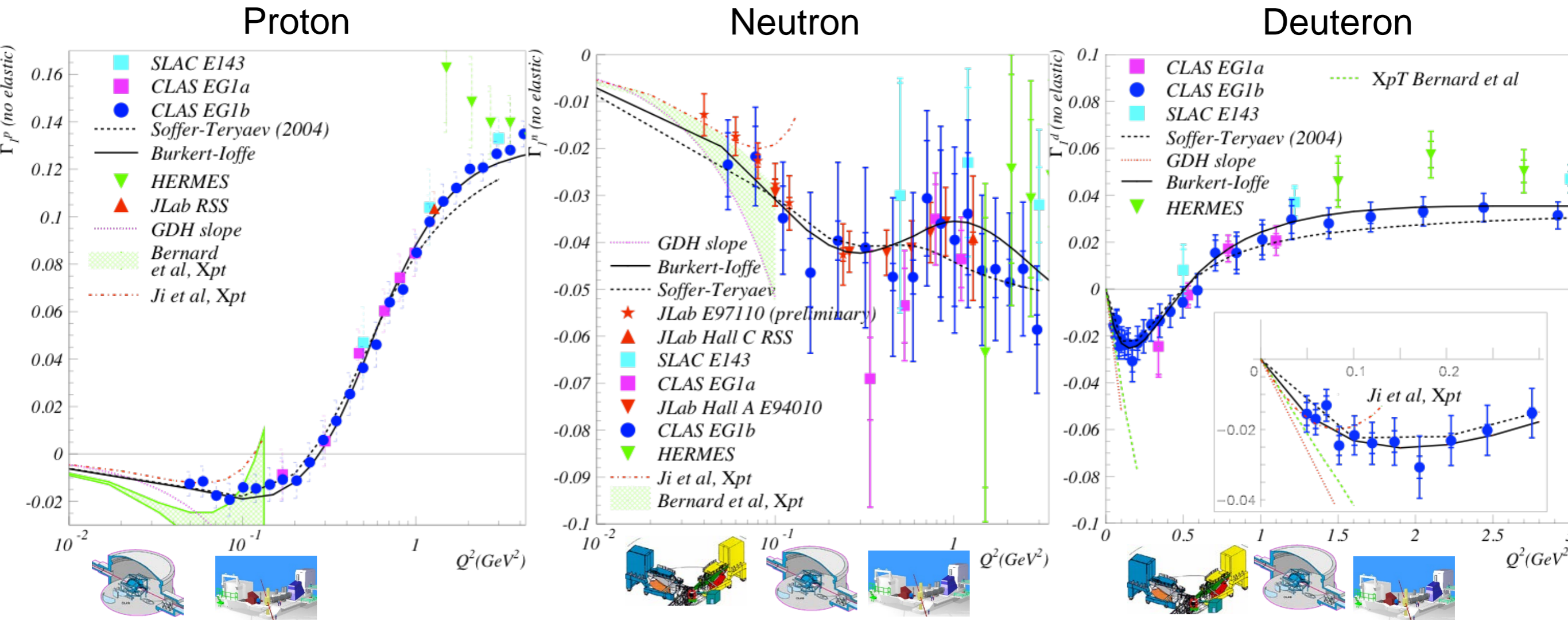
$$\delta_{LT} = \frac{4e^2 M^2}{\pi Q^6} \int x^2 (g_1 + g_2) dx$$

Waiting for g_2 data

Not further discussed in this presentation

Previous data: high to intermediate Q^2

Before EG4 run (2006):



Precise mapping of intermediate Q^2 region for p, n and d.
 pQCD, models and data agree.
 Not so clear for χ pT.

Test of χ pt

A = agree

X = disagree

- = no calculation available

No significant low-x contribution
(More robust observables)



Ref.	Γ_1^p	Γ_1^n	Γ_1^{p-n}	Γ_1^{p+n}	γ_0^p	γ_0^n	γ_0^{p-n}	γ_0^{p+n}	δ_{LT}^n
Ji 1999	X	X	A	X	-	-	-	-	-
Bernard 2002	X	X	A	X	X	A	X	X	X
Kao 2002	-	-	-	-	X	A	X	X	X
Bernard 2012	X	X	A	X	X	A	X	X	X
Lensky 2014	X	A	A	A	A	X	X	X	\sim A

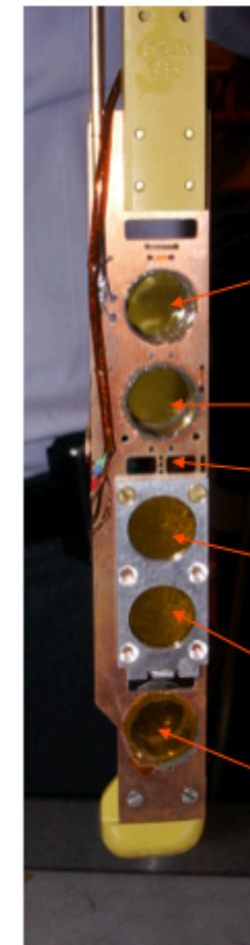
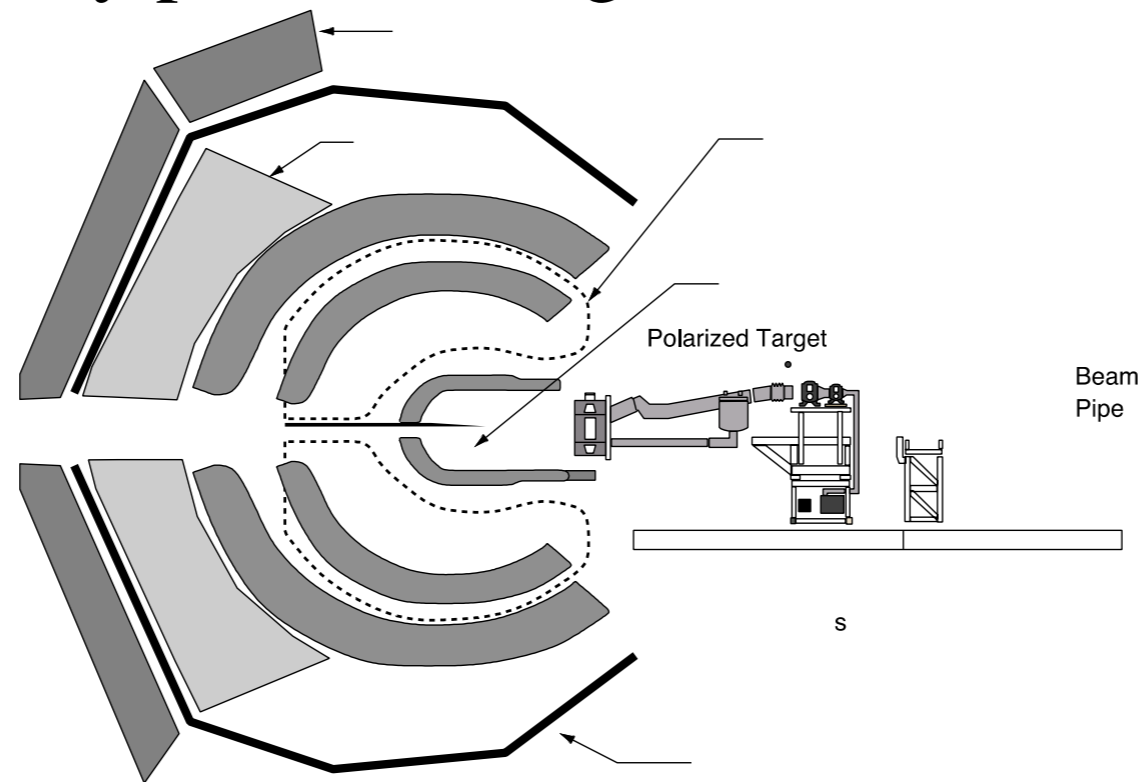
No Δ_{1232} contributions
(More robust χ pt calculations)



EG4 setup

- $Q^2 > 0$: electron beam (polarized). Energies: 3.0, 2.3, 2.0, 1.3 & 1.0 GeV
- $g_1^{p,n}$: ~longitudinally polarized target

DNP NH_3 and
 ND_3 target:



Long NH_3 (1.0 cm)

Short NH_3 (0.5 cm)

Raster

Long Carbon (2.30 mm)

Short Carbon (1.15 mm)

Empty

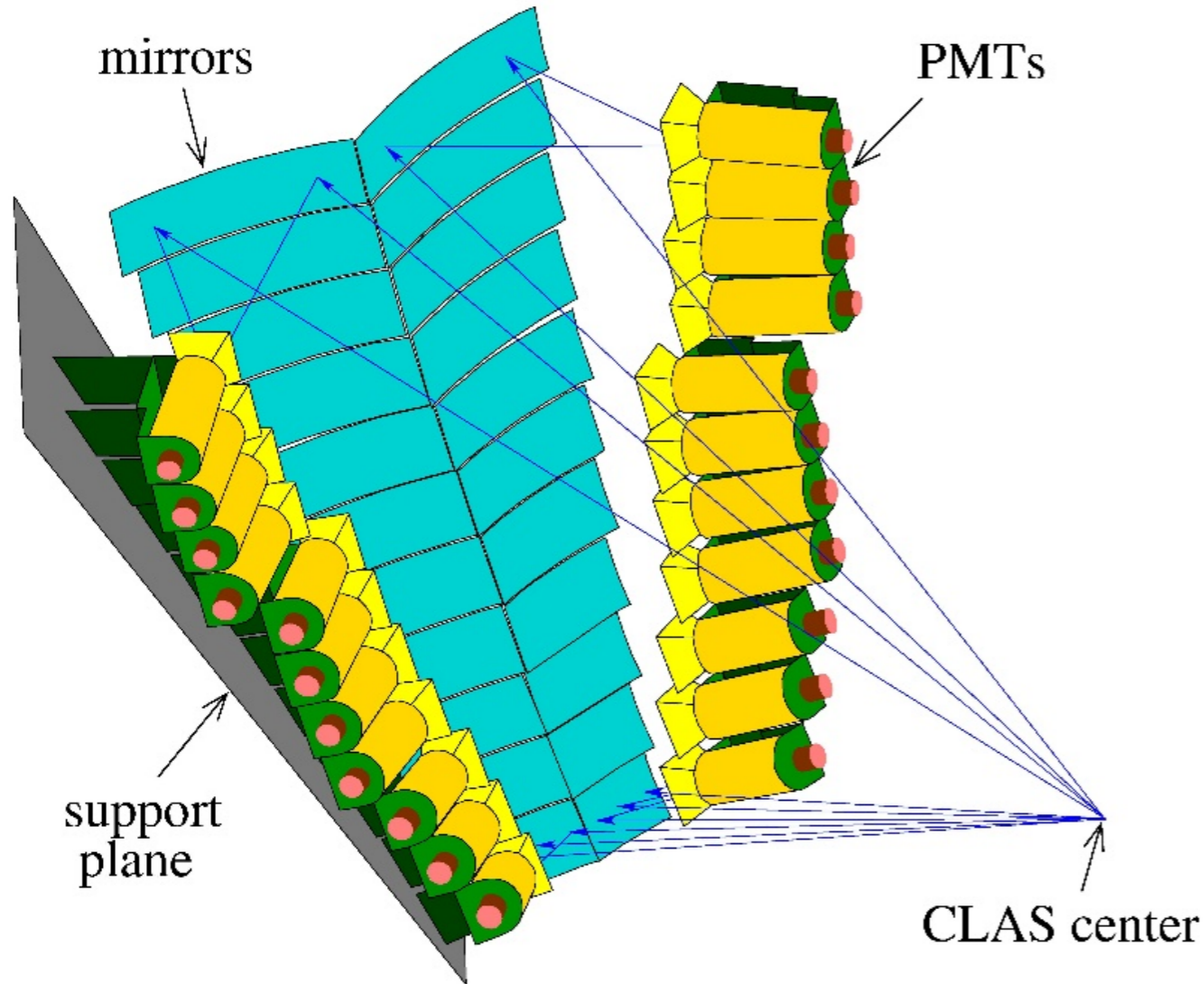
- g_1 from pol. cross-section differences (not asymmetries, as in EG1, EG1dvcs)
Advantage: dilution from unpol. target material cancels out

- Small angles: outbending torus field, new Møller shield; target at -1m

- Cross-section \Rightarrow controlled (i.e high) efficiency at small angles.

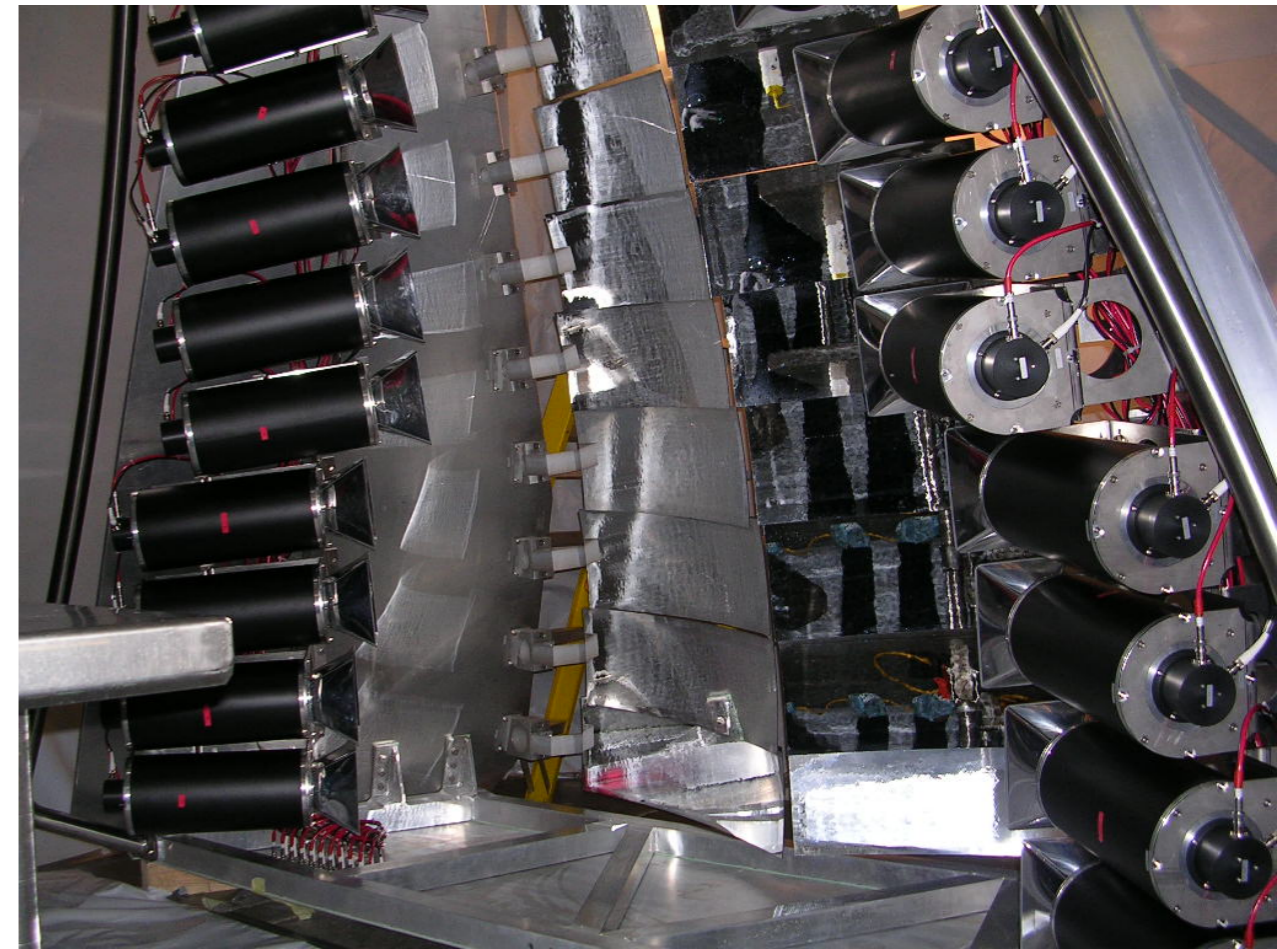
\rightarrow New Cherenkov detector (INFN). Installed in sector 6. Cover down to 6°

EG1: Largest possible kinematic coverage -> inbending and outbending configuration, $E = 1.6 \dots 5.8$ GeV



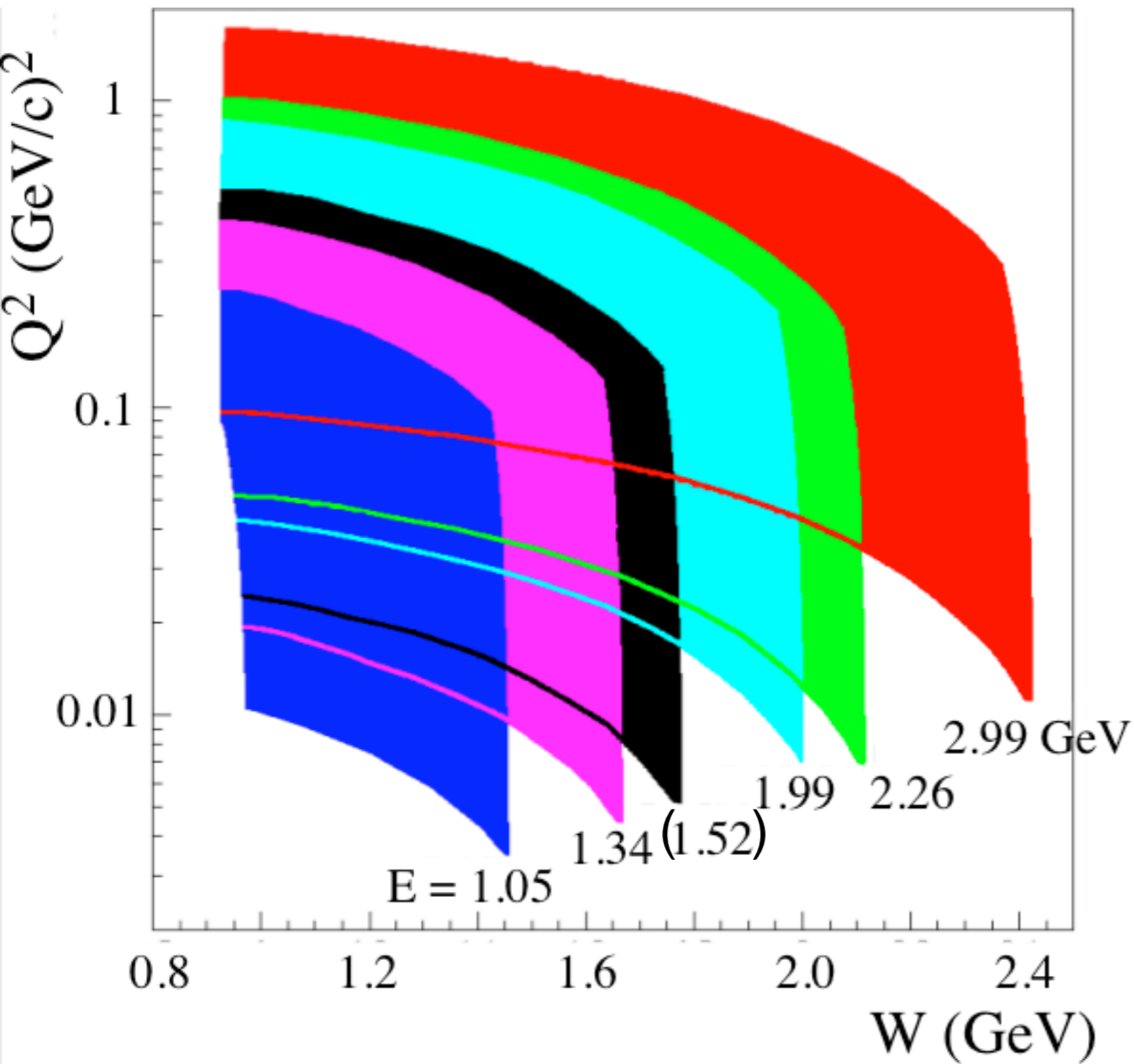
EG4: Focus on low Q^2 => lower beam energies, new Cherenkov for optimal acceptance in outbending configuration, θ_e as small as 6 degrees

... and EG1-DVCS: Highest statistics at large x , Q^2

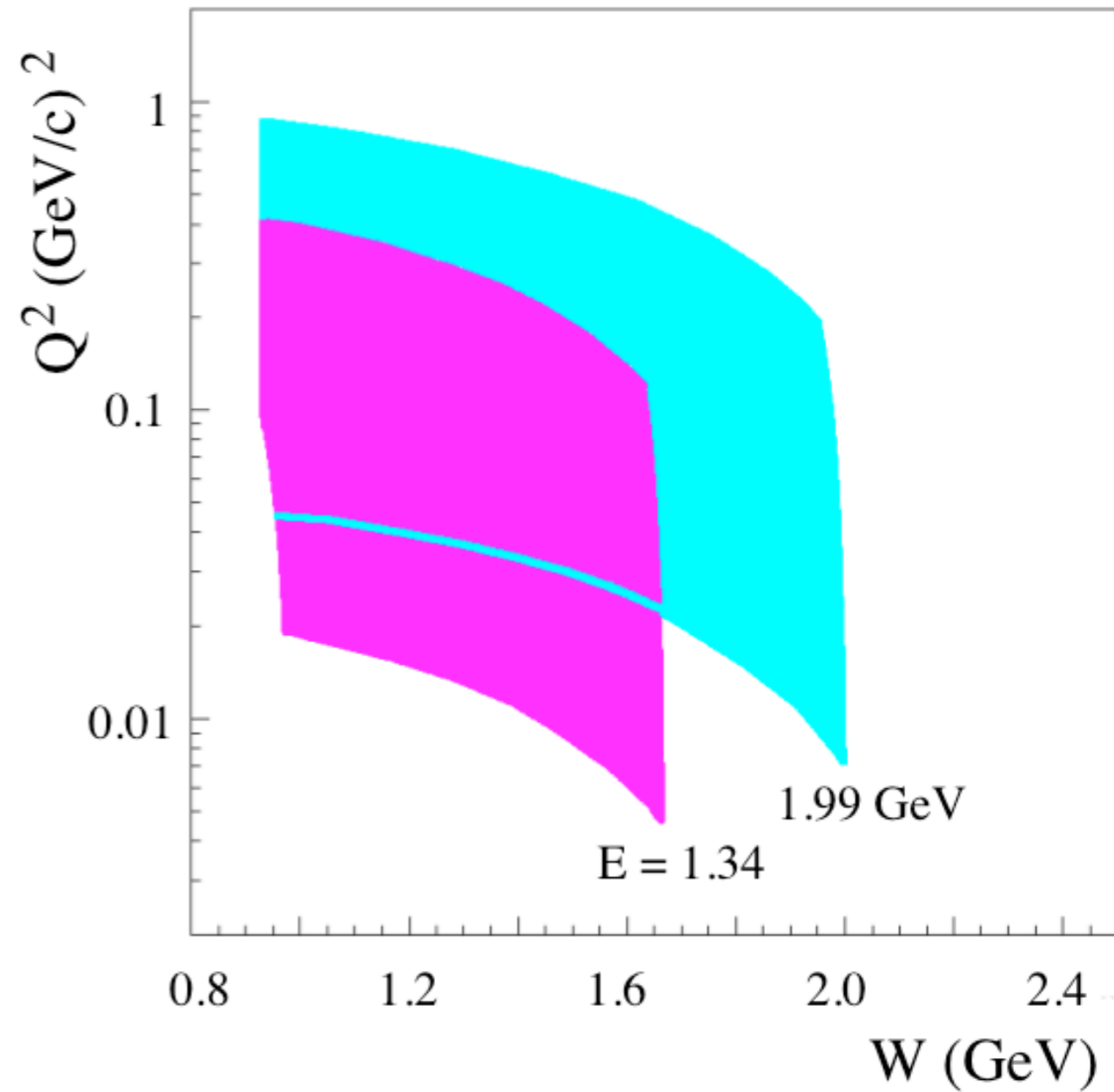


EG4 kinematic coverage

Proton



Deuteron



1.52 GeV proton only for commissioning purposes

Experimental procedure

- g_1 from pol. cross-section differences (not asymmetries, as in EG1, EG1dvcs)

Advantage: dilution from unpol. target material cancels out

$$\frac{\Delta d\sigma^{theor}}{d\Omega dE'} \equiv \frac{d\sigma^{\rightarrow\Rightarrow}}{d\Omega dE'} - \frac{d\sigma^{\leftarrow\Rightarrow}}{d\Omega dE'} = \frac{4\alpha^2 E'^2}{ME\nu Q^2} [(E - E' \cos\vartheta)g_1(x, Q^2) - 2Mxg_2(x, Q^2)]$$

- Small $Q^2 \rightarrow$ small $x \rightarrow g_2$ contribution suppressed

- Actual experimental quantity measured:

$$\Delta N^{exp}(\Delta E', \Delta\Omega) \equiv N^{\rightarrow\Rightarrow}(\Delta E', \Delta\Omega) - N^{\leftarrow\Rightarrow}(\Delta E', \Delta\Omega) = P_f \mathcal{L} P_b P_T \frac{\Delta d\sigma^{theor}}{d\Omega dE'} \Delta E' \Delta\Omega \epsilon_{det} \quad (*)$$

where

- P_f = packing fraction (how much target cell filling with ammonia beads)
- \mathcal{L} = integrated luminosity (how many electrons on target times nominal target surface thickness) (*)
- P_b = beam polarization ($85 \pm 2\%$)
- P_b = target polarization (59 to 71 % for H, 30 to 45 % for D)
- ϵ_{det} = detector acceptance/efficiency

(*) there is a small correction due to the beam charge asymmetry, not shown here

Experimental procedure

First step

$$\Delta N^{exp,(quasi)elastic}(\Delta E', \Delta \Omega)$$

What CLAS actually measures

$$\equiv N^{\rightarrow\Rightarrow,(quasi)elastic}(\Delta E', \Delta \Omega) - N^{\leftarrow\Rightarrow,(quasi)elastic}(\Delta E', \Delta \Omega)$$

$$= P_f \mathcal{L} P_b P_T \frac{\Delta d\sigma^{theor,(quasi)elastic}}{d\Omega dE'} \Delta E' \Delta \Omega \epsilon_{det}$$

Monte Carlo simulation

Extracted from data

Well known in terms of elastic FF or quasielastic parameterization

Experimental procedure

Second step

1 $\Delta N^{exp,inelastic}(\Delta E', \Delta \Omega)$ What CLAS actually measures

$$\equiv N^{\rightarrow\Rightarrow,inelastic}(\Delta E', \Delta \Omega) - N^{\leftarrow\Rightarrow,inelastic}(\Delta E', \Delta \Omega)$$

2 $= P_f \mathcal{L} P_b P_T \frac{\Delta d\sigma^{theor,inelastic}}{d\Omega dE'} \Delta E' \Delta \Omega \epsilon_{det}$

Input from elastic Monte Carlo simulation

Calculated from parameterization, g_1 or $A_1 F_1$ varied until 1 and 2 agree at best

$g_1, A_1 F_1$

Experimental procedure

Sum rule integrals

Deuteron

$$\int_0^{x_{th}} \dots \rightarrow \int_{0.001}^{x_{min}} Model + \int_{x_{min}}^{x(W=1.15 GeV)} data + \int_{x(W=1.15 GeV)}^{x(W=1.07 GeV)} Model$$

For the lowest Q^2 bin, 0.020 GeV^2 , $x_{min} = 0.0073$

For the largest Q^2 bin considered for integration, 0.592 GeV^2 , $x_{min} = 0.280$

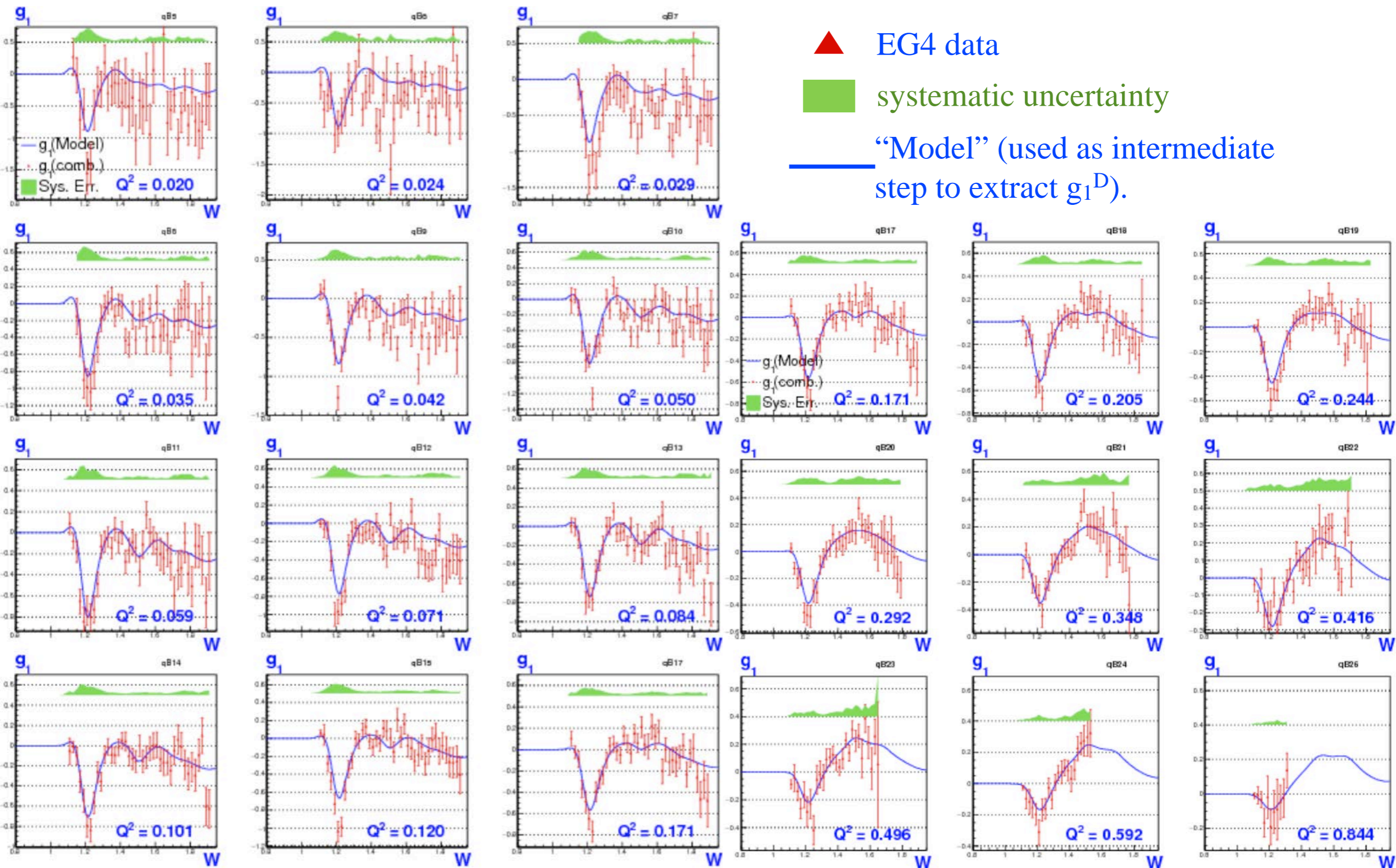
In the 3rd integral, the model is used rather than data to avoid quasielastic scattering and radiative tail contaminations

Proton

$$\int_0^{x_{th}} \dots \rightarrow \int_{0.001}^{x_{min}} Model + \int_{x_{min}}^{x(W=1.08 GeV)} data$$

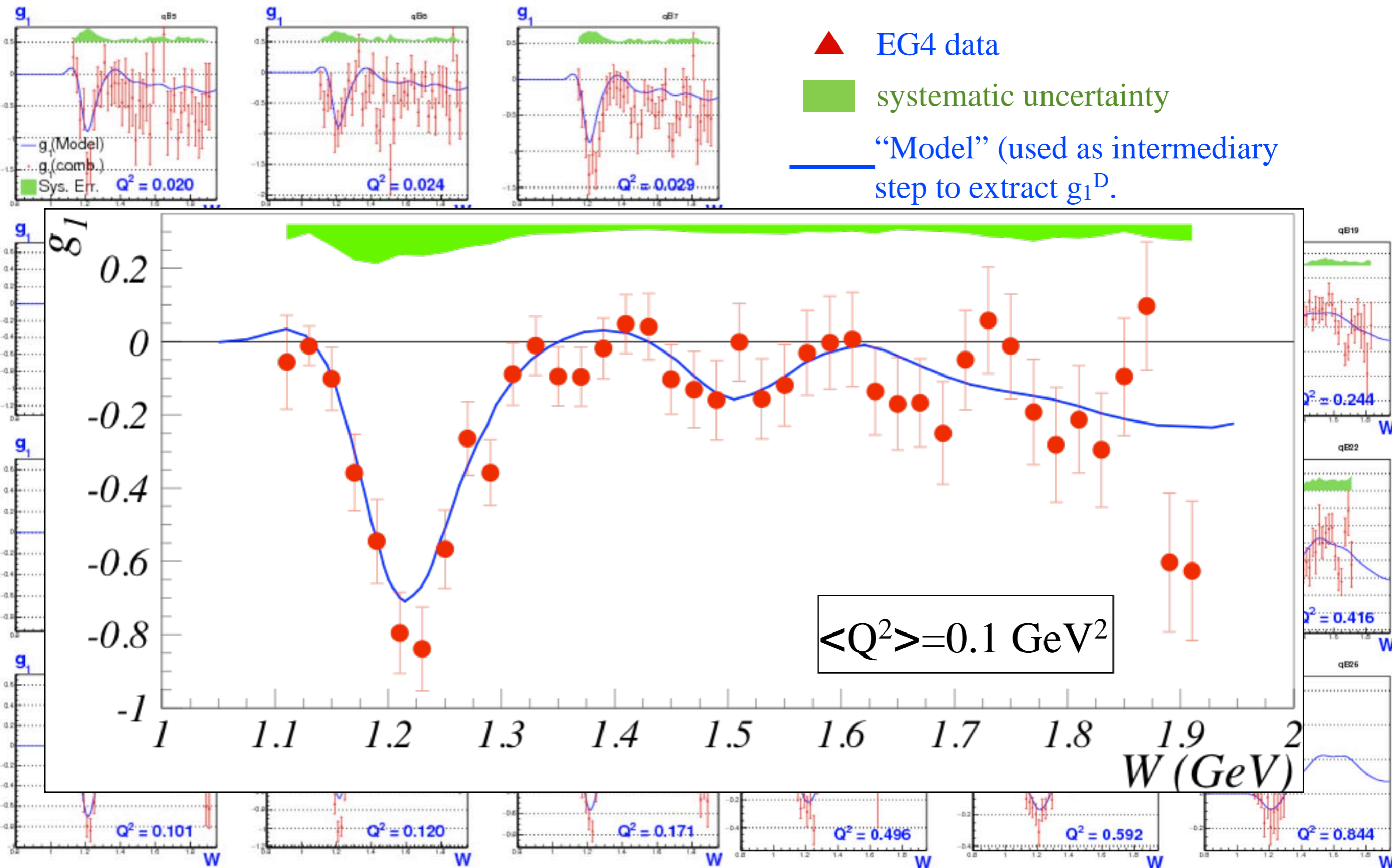
g_1^d from EG4 polarized cross-section difference

K. Adhikari, S. Kuhn



g_1^d from EG4 polarized cross-section difference

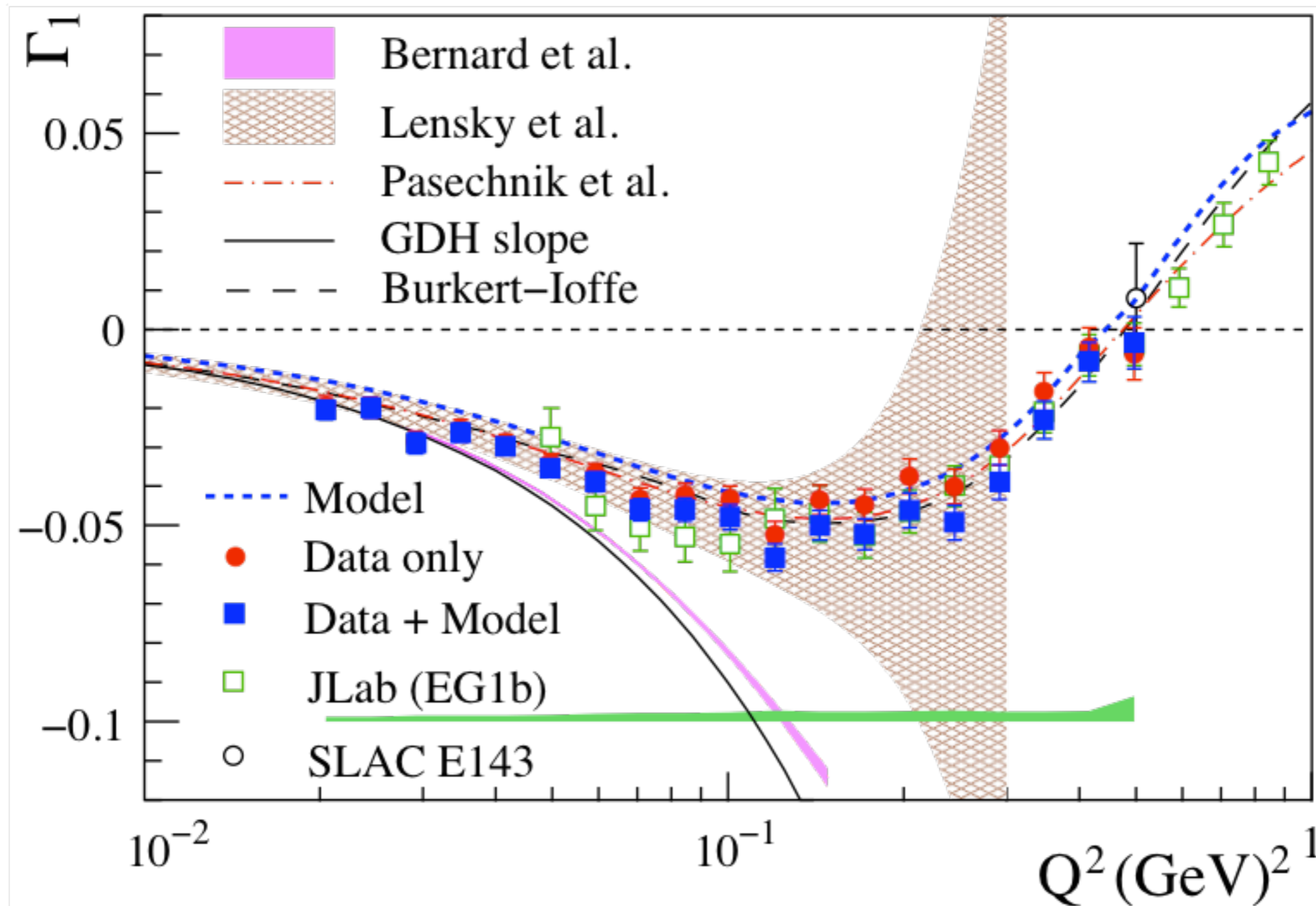
K. Adhikari, S. Kuhn



$$\Gamma_1^d = \int g_1^d(x, Q^2) dx$$

K. Adhikari, S. Kuhn

K.P. Adhikari et al. (CLAS Collaboration). PRL 120, 062501 (2018)



• Lowest Q^2 decreased by factor of ~ 2.5

• Much improved precision

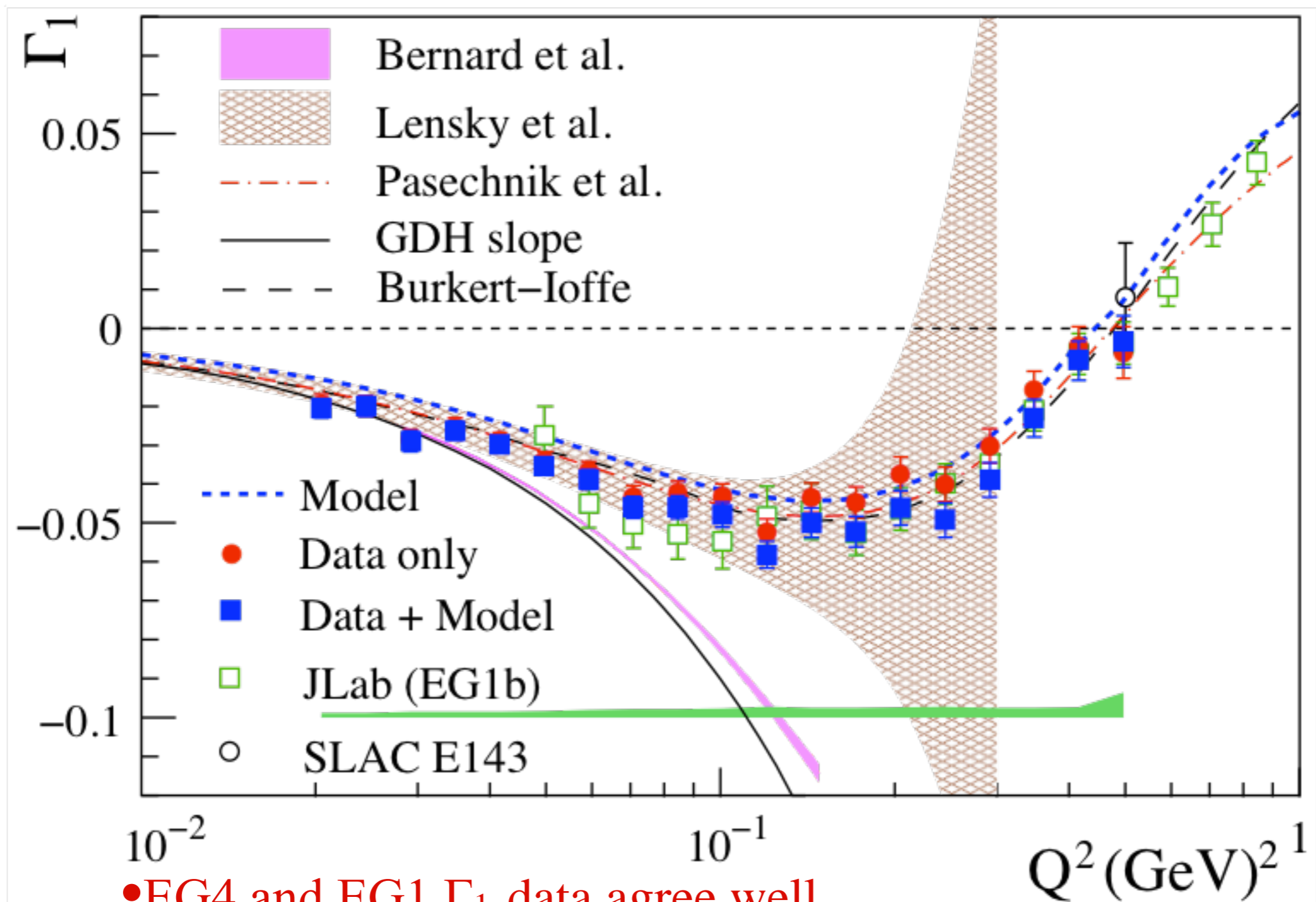
\Rightarrow Clean test of χ pt

• Small unmeasured low- x and large- x contributions

$$\Gamma_1^d = \int g_1^d(x, Q^2) dx$$

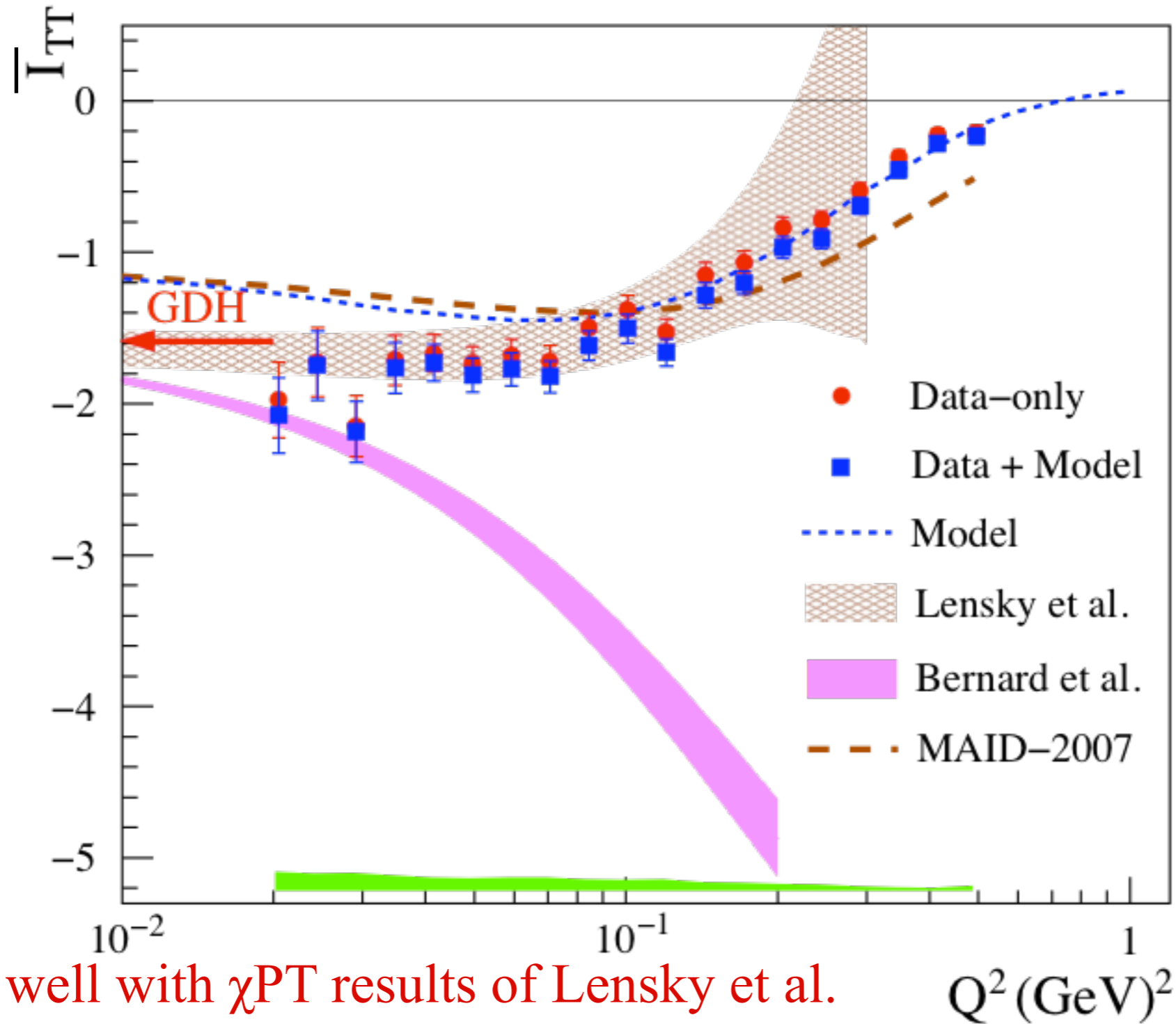
K. Adhikari, S. Kuhn

K.P. Adhikari et al. (CLAS Collaboration). PRL 120, 062501 (2018)



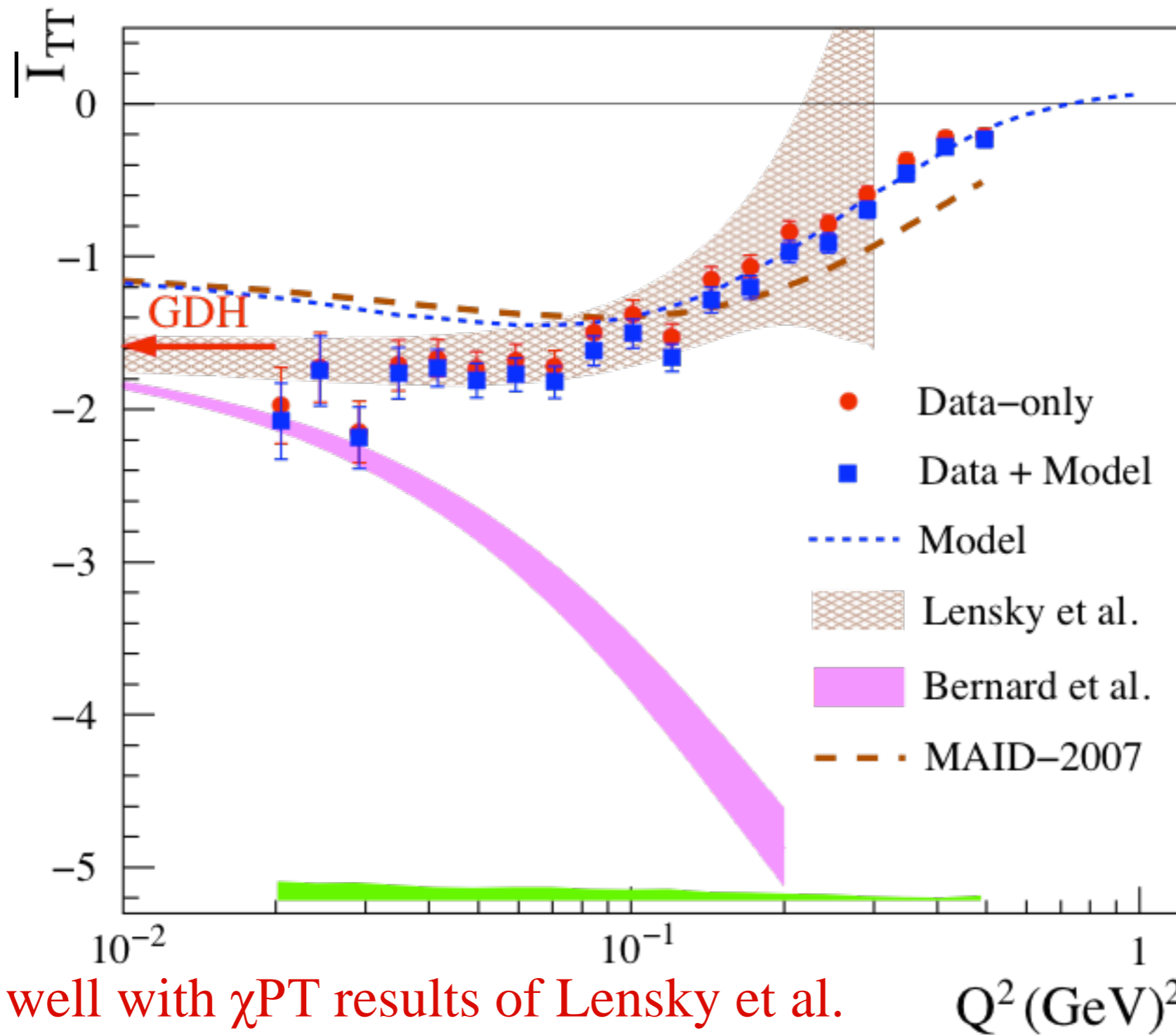
- EG4 and EG1 Γ_1 data agree well.
- EG4 data agree well with χ PT results of Lensky et al.
- Bernard et al. χ PT calculation agrees only for the lowest Q^2 points.
- Phenomenological models (Pasechnik et al, Burkert-Ioffe) agree well.

Generalized GDH sum $\bar{I}_{TT} = \int \frac{\sigma_A(\nu) - \sigma_P(\nu)}{\nu} d\nu$



- Data agree well with χ PT results of Lensky et al.
- Bernard et al. χ PT calculation does not agree as well.
- Maid model disagrees at low $Q^2=0$.

Generalized GDH sum $\bar{I}_{TT} = \int \frac{\sigma_A(\nu) - \sigma_P(\nu)}{\nu} d\nu$



- Data agree well with χ PT results of Lensky et al.
- Bernard et al. χ PT calculation does not agree as well.
- Maid model disagrees at low $Q^2=0$.
- Extrapolation to $Q^2=0$ tests original GDH sum rule:

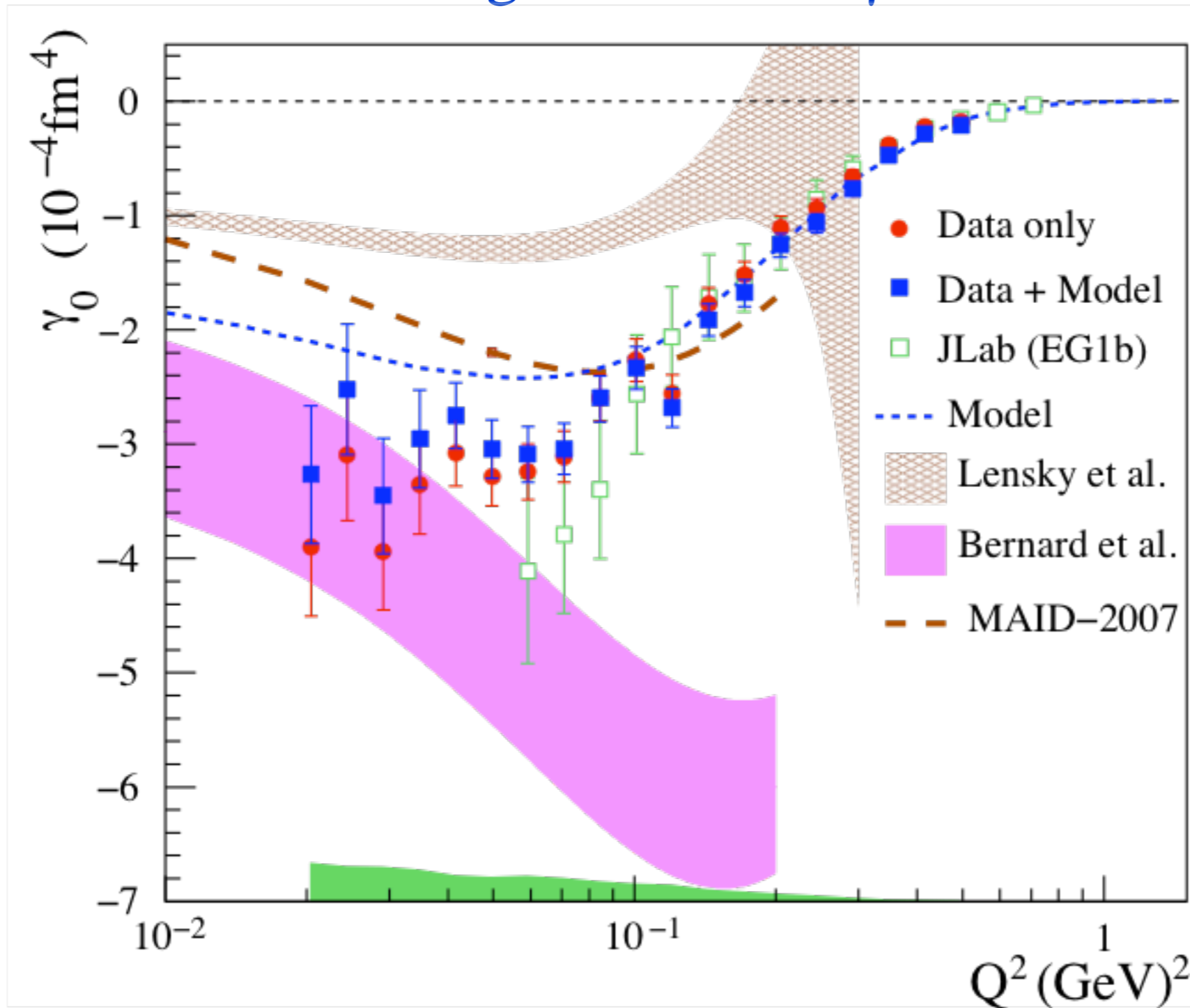
$$\bar{I}_{TT}^d = -1.724 \pm 0.027(\text{stat}) \pm 0.050(\text{syst})$$

Sum rule expectation: -1.574 ± 0.026

$$I_{TT}^n = -0.955 \pm 0.040(\text{stat}) \pm 0.113(\text{syst})$$

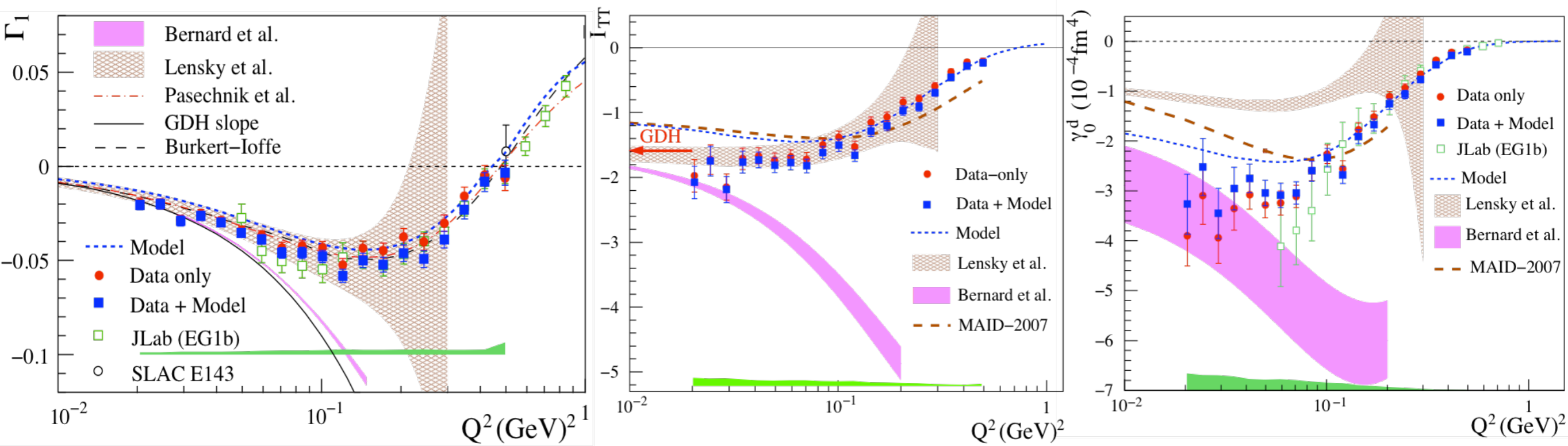
Sum rule expectation: -0.803

Higher moment γ_0



- Incoherent sum of p and n $\rightarrow \chi$ PT results of Lensky et al. disagree with data.
- Bernard et al. χ PT calculation agree for lowest Q^2 points only.
- Maid model disagrees at low Q^2 .

Conclusion from deuteron data

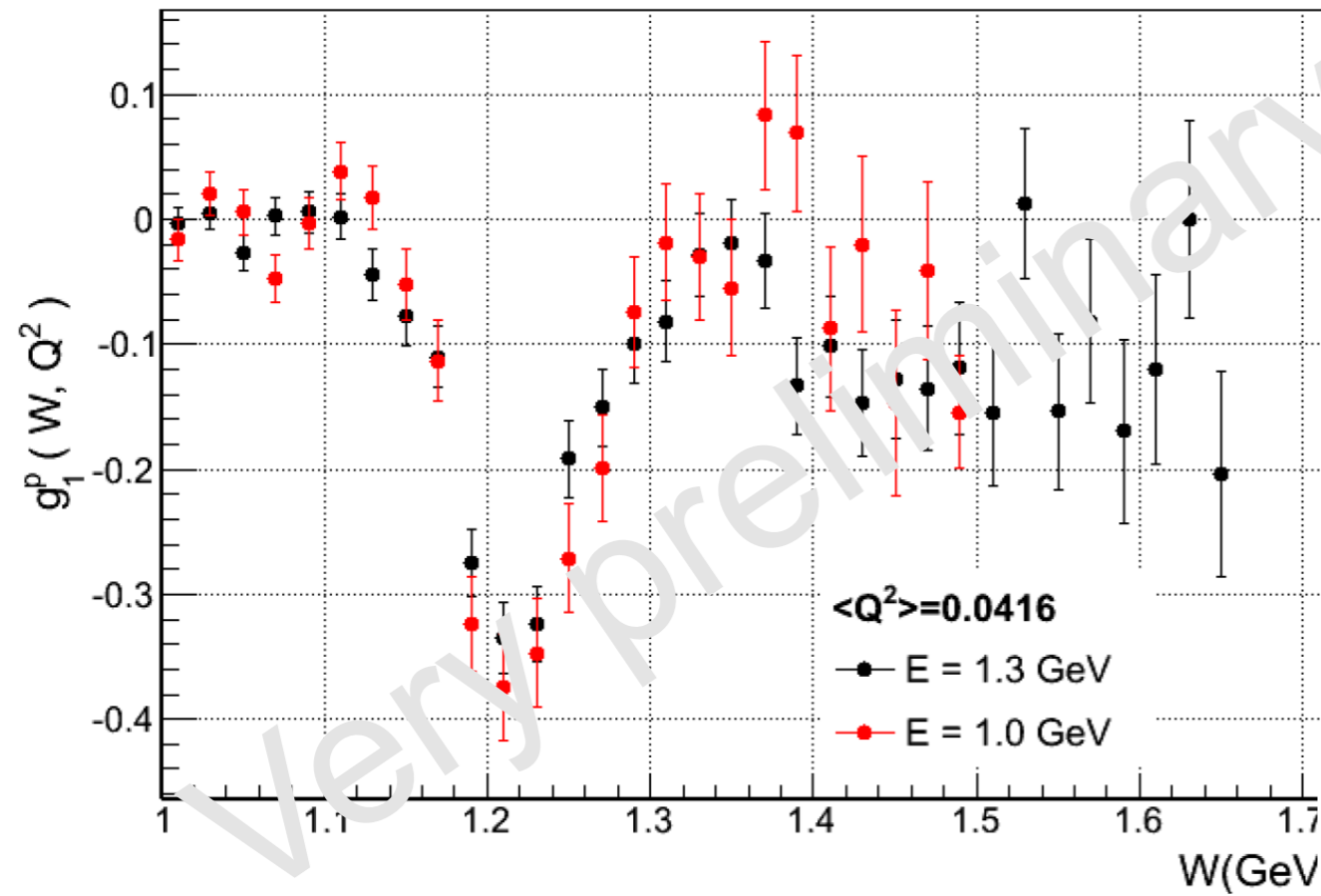


No χ PT single method describes well both Γ_1 , $\overline{I_{TT}}$, and γ_0 , except at the lowest Q^2 .

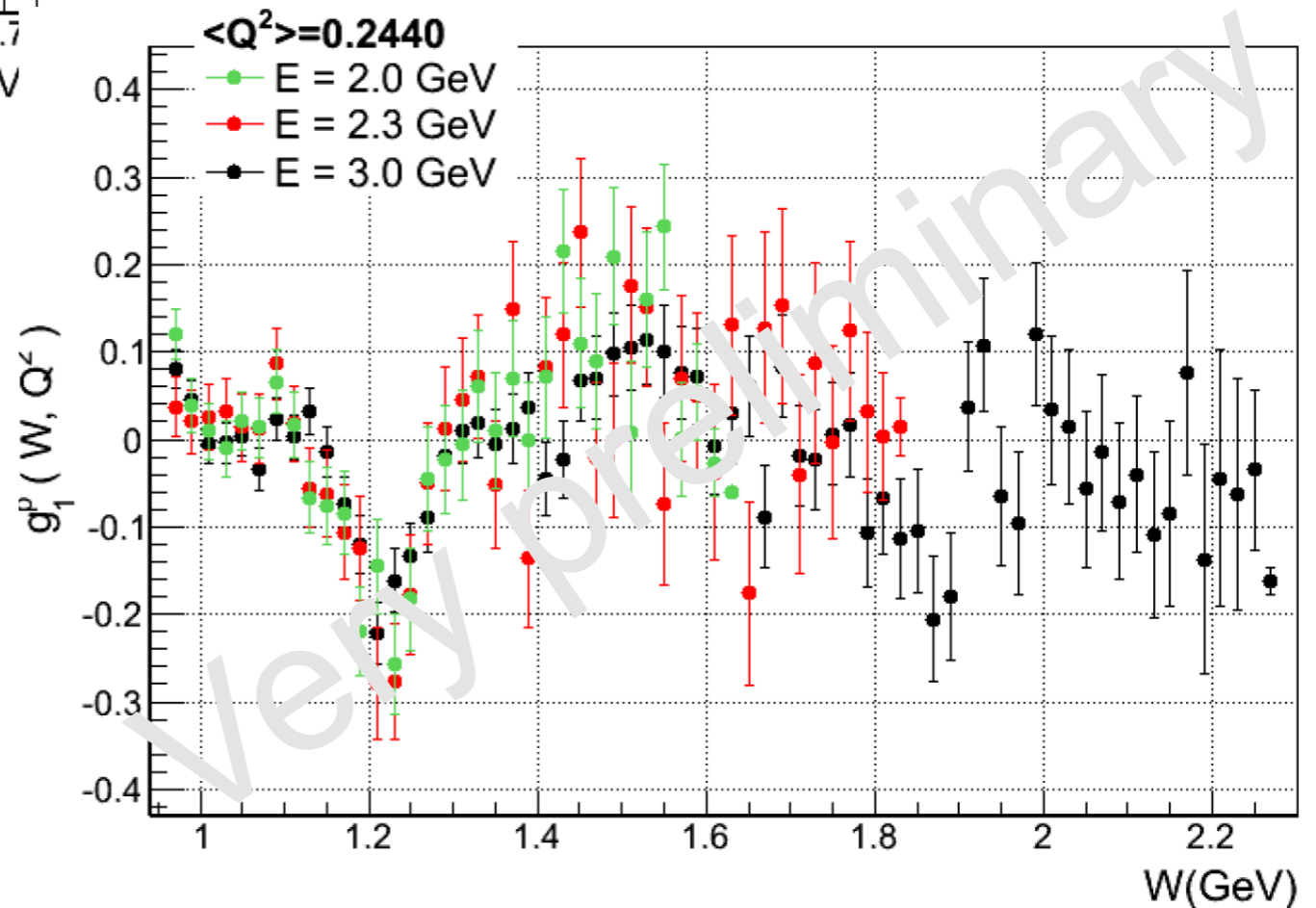
A satisfactory theoretical description of spin observables at low Q^2 remains challenging.

g_1^p from EG4 polarized cross-section difference

H. Kang

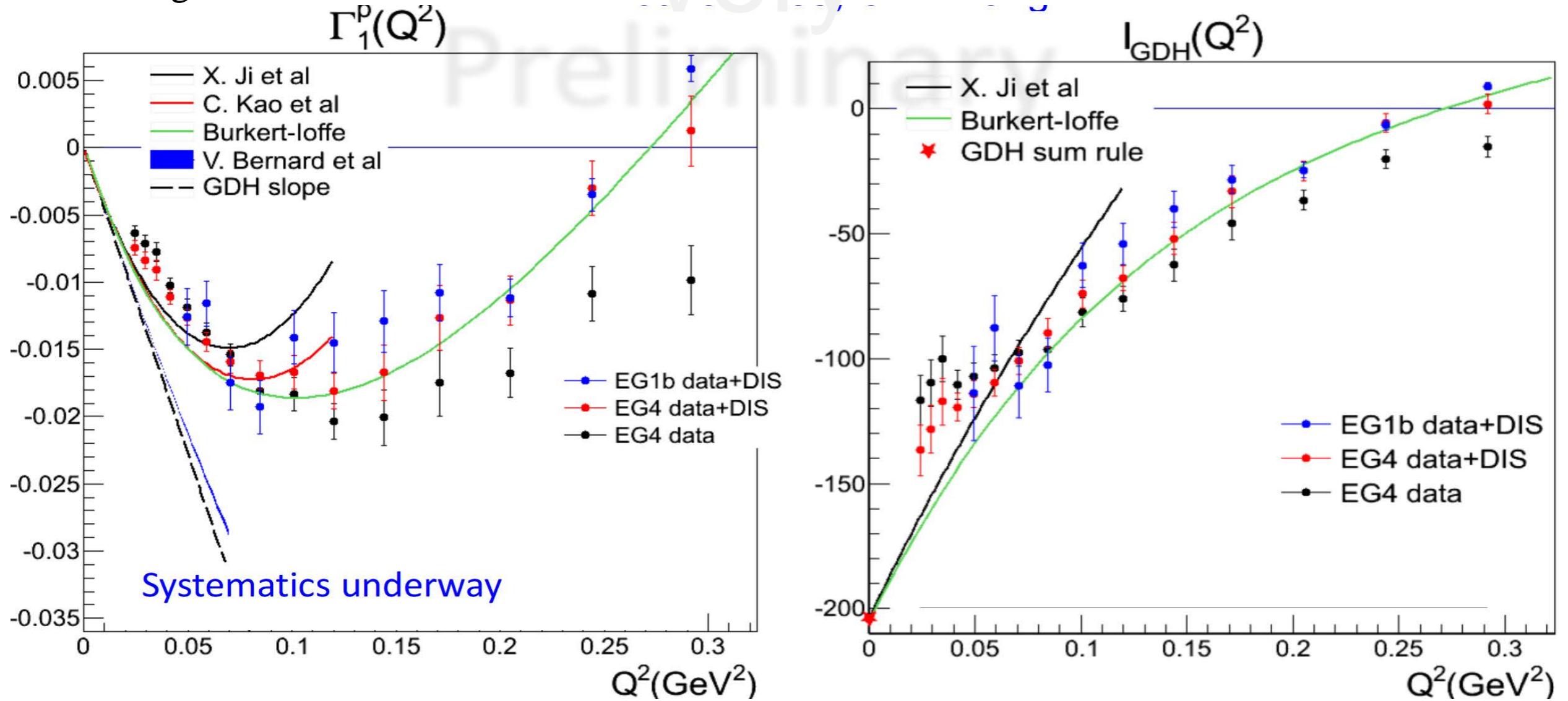


Initial approach was to correct expl data with combination of factors derived from elastic analysis, then invert equations to get g_1



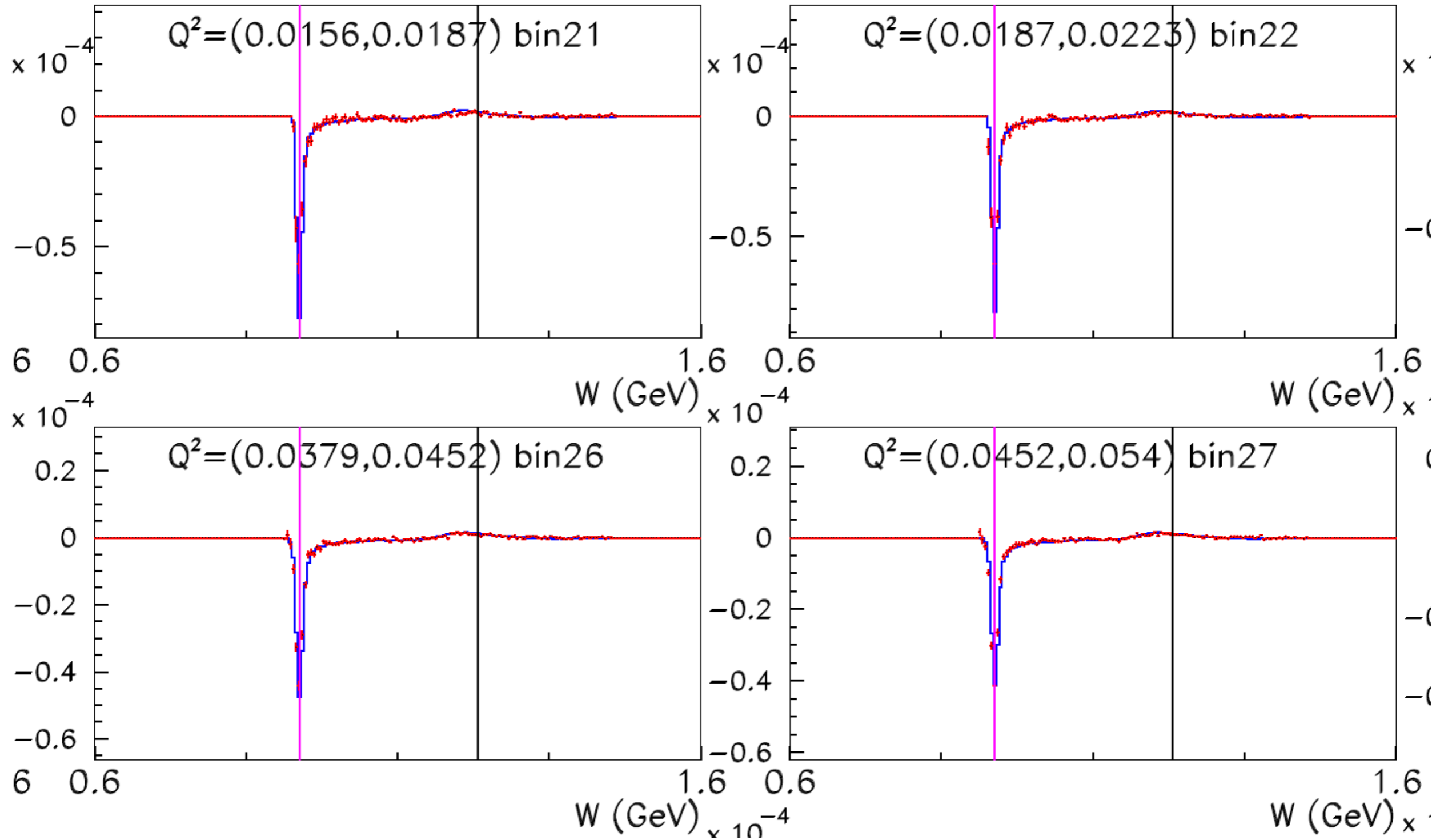
g_1^p from EG4 polarized cross-section difference

H. Kang



Currently revisiting analysis with same technique as deuteron (X. Zheng with essential support from L. El Fassi and J. Zhang)

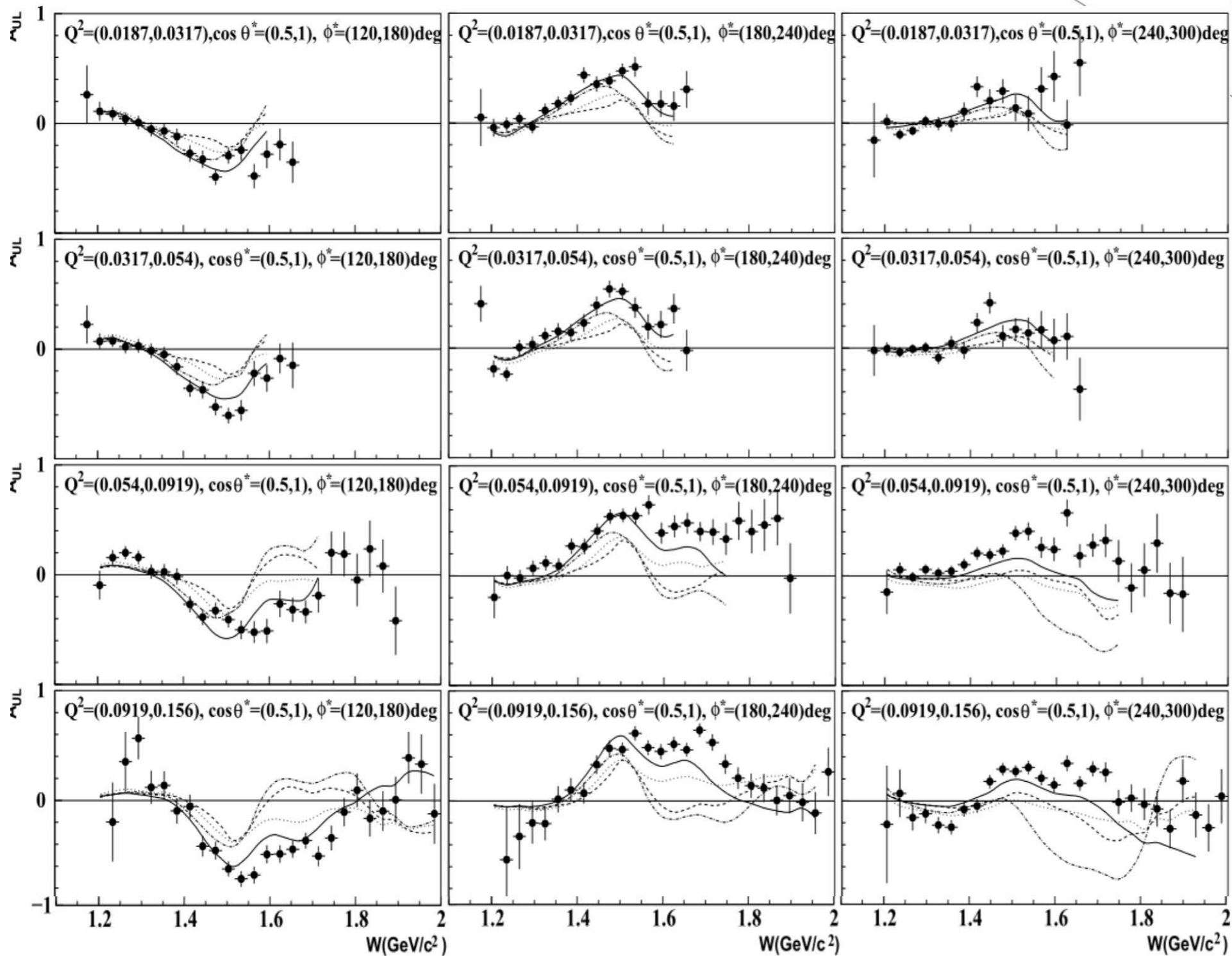
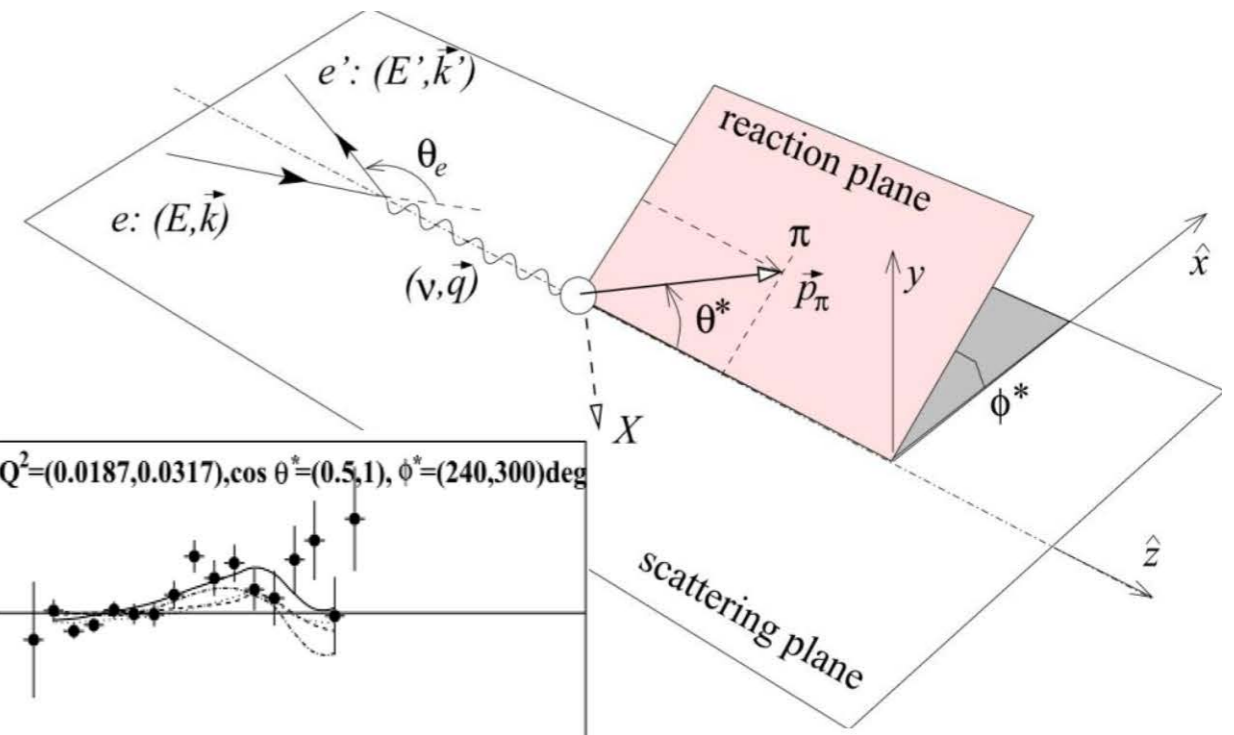
Comparison between data and simulations in new approach



Data
Simulations

EG4: Spin Asymmetry A_{UL} Results on $p(e, e'\pi^+)n$

Plots courtesy of Xiaochao Zheng



The target spin asymmetries A_{UL} for the π^+n channel

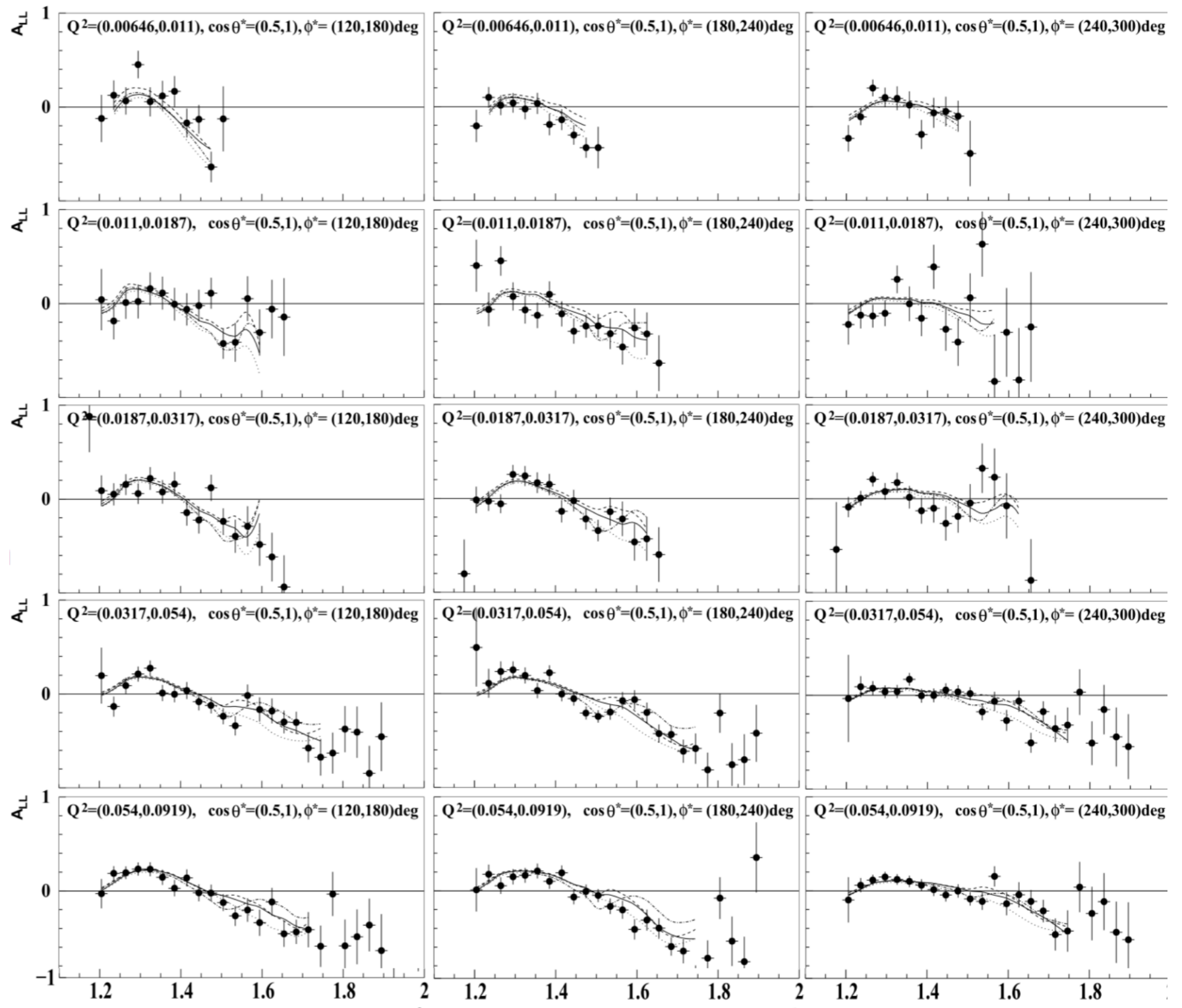
EG4: Spin Asymmetry A_{LL} Results on $p(e,e'\pi^+)n$



Plots courtesy of
Xiaochao Zheng

Results on the double-
spin symmetries A_{LL} for
 π^+n channel in different
 Q^2 bins

X. Zheng et al.
(CLAS Collaboration),
PRC 94, 045206 (2016)



Summary and perspectives

- EG4: **Low Q^2** measurement using polarized e^- on polarized p and d, over a large x-range in order to study **spin sum rules**.
- **New detector** necessary to reach these kinematics.
- Main goal: **unambiguous test of χ PT**.
- Doubly polarized inclusive cross-section analysis.
- Exclusive data for π^+ and π^- spin-dep. electroprod. on p published in 2016 (asym. analysis).
X. Zheng et al. (CLAS Collaboration), PRC 94, 045206 (2016)
- **Inclusive analysis on d recently published** K.P. Adhikari et al. (CLAS Collaboration), PRL 120, 062501 (2018)
- Data on Γ_1 , $\overline{I_{TT}}$, and γ_0 for the deuteron shows that **χ PT has mixed success**, depending on the χ PT method and observable.
- Original GDH sum rule ($Q^2=0$) checked on d and n.
- First result of larger **JLab program to measure benchmark spin observables for χ PT**
 \Rightarrow More low Q^2 data to come:
 - g_1 , Γ_1 , I_{TT} , and γ_0 for the **proton** (CLAS EG4).
 - g_1 , g_2 , Γ_1 , Γ_2 , I_{TT} , γ_0 and δ_{LT} for the **neutron** and **^3He** (Hall A E97110).
 - g_2 , g_1 , Γ_2 , Γ_1 , I_{TT} , δ_{LT} and γ_0 for the **proton** (Hall A E08027).