Novel calculation of the nucleon form factors with Dispersively Improved Chiral EFT

Jose Manuel Alarcón



Works done in collaboration with Christian Weiss (JLab)

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 - Important to understand and solve the "Proton Radius Puzzle".

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Unified approach

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$$Im t$$

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Space-like region (t < 0)
$$e^{-} \underbrace{e^{-}}_{\gamma*} \underbrace{t = -Q^{2}}_{K} \underbrace{e^{-}}_{K} \underbrace{e^{-}}_$$

$$\langle N(p',s')|J_{\mu}(0)|N(p,s)\rangle = u(p',s') \Big[\gamma_{\mu}F_{1}(t) + \frac{i\sigma_{\mu\nu}q^{\nu}}{2m_{N}}F_{2}(t) \Big] u(p,s) \qquad J_{\mu}(x) \equiv \sum_{q=u,d,...} c_{q}q(x)\gamma_{\mu}q(x)$$

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Nucleon Form Factors in DIXEFT

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Nucleon Form Factors in DIXEFT

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Nucleon Form Factors in DIXEFT

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• We write this representation in a more convenient form

 $\operatorname{Im} G^{V}_{\{E,M\}}(t) = \frac{k_{cm}^{3}}{\{m_{N},\sqrt{2}\}\sqrt{t}} F^{*}_{\pi}(t)f^{1}_{\pm}(t)$

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• $DI\chi EFT \longrightarrow$ The spectral function is factorized into two parts:

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$$J_{\pm}^{1} = \overbrace{F_{\pi}}^{f_{\pm}^{1}} \overbrace{F_{\pi}}^{\eta} \overbrace{\Lambda}^{\eta} \overbrace{I}^{\eta} \overbrace{I}^{\eta}$$

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Nucleon Form Factors in $DI\chi EFT$

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Nucleon Form Factors in DIXEFT

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Nucleon Form Factors in DIXEFT

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• Electromagnetic FF: Since t >1 GeV² is far away from the space-like region, we parametrize the contribution from this region by an effective pole P_V [Höhler and Pietarinen, Phys. Lett. 53B (1975)] : $\mathrm{Im}G_{E,M}^V = -\pi a_{E,M}^{P_V} \delta(t - M_{P_V}^2)$

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• We fix the free parameters by imposing:

$$G_{E,M}^{V}(0) = \frac{1}{\pi} \int_{4M_{\pi}^{2}}^{\infty} dt' \frac{\mathrm{Im}G_{E,M}^{V}(t')}{t'} \qquad \langle r_{E,M}^{2} \rangle^{V} = \frac{6}{\pi} \int_{4M_{\pi}^{2}}^{\infty} dt' \frac{\mathrm{Im}G_{E,M}^{V}(t')}{t'^{2}}$$

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Nucleon Form Factors in DIXEFT

$$\mathrm{Im}G_{E}^{V}(t) = \frac{k_{cm}^{3}}{m_{N}\sqrt{t}}|F_{\pi}(t)|^{2}J_{+}^{1}(t)$$

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[]. M. Alarcón, C. Weiss, arXiv: 1803.09748]

• Higher order corrections are important for $t > 0.2 \text{ GeV}^2$.

• Error bands shown correspond to the uncertainties in the LECs.

• Systematic errors are inferred from the difference between NLO and NLO+pN2LO.

[1] Obtained from: Hoferichter, Kubis, Ruiz de Elvira, Hammer, Meißner EPJA 52 (2016)

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Nucleon Form Factors in DIXEFT

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Nucleon Form Factors in $\mathsf{DI}\chi\mathsf{EFT}$

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- Pion EM FF \rightarrow related to $e^+e^- \rightarrow \pi^+\pi^-$ cross sections
 - Related to measured quantities.
 - Dispersion Theory $\rightarrow \pi\pi$ phase shifts.
 - LQCD
 - We use the GS parametrization of [Lorenz, Hammer, Meißner, EPJ A 48 (2012)]

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Nucleon Form Factors in $DI\chi EFT$

Spectral Functions

$$Im G_E^V(t) = \frac{k_{cm}^3}{m_N \sqrt{t}} (F_{\pi}(t))^2 J_{+}^1(t)$$

ChEFT

$$Im G_M^V(t) = \frac{k_{cm}^3}{\sqrt{2t}} F_{\pi}(t) J_{-}^1(t)$$
ChEFT



Nucleon Form Factors in $\mathsf{DI}\chi\mathsf{EFT}$



[1] Belushkin, Hammer and Meißner, PRC 75 (2007) [2] Hoferichter, Kubis, Ruiz de Elvira, Hammer, Meißner EPJA 52 (2016)

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Nucleon Form Factors in $DI\chi EFT$

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• Comparison with respect to the old results



Comparison with respect to the old results



Conclusions:

- Brute force calculations are hopeless.
- Non-perturbative effects are visible in the near-threshold region.
- Based on unitarity one achieves a factorization suitable for perturbative calculations.

Electromagnetic Form Factor

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$$\operatorname{Im} G_{E,M}^{S} = -\pi \sum_{V=\omega, P_{S}} a_{i}^{E,M} \delta(t - M_{i}^{2})$$

• We fix the couplings by imposing the charge and radii sum rules:

$$G_{E,M}^{S}(0) = \frac{1}{\pi} \int_{4M_{\pi}^{2}}^{\infty} dt' \frac{\text{Im}G_{i}^{S}(t')}{t'} \qquad \langle r_{E,M}^{2} \rangle^{S} = \frac{6}{\pi} \int_{4M_{\pi}^{2}}^{\infty} dt' \frac{\text{Im}G_{E,M}^{S}(t')}{t'^{2}}$$

Nucleon Form Factors in DIXEFT

• Reconstructing the form factors with $G_{E,M}^{p,n}(t) = \frac{1}{\pi} \int_{4M_{\pi}^2}^{\infty} dt' \frac{\text{Im} G_{E,M}^{p,n}(t')}{t'-t-i0^+}$
• Reconstructing the form factors with $G_{E,M}^{p,n}(t) = \frac{1}{\pi} \int_{4M_{\pi}^2}^{\infty} dt' \frac{\text{Im}G_{E,M}^{p,n}(t')}{t'-t-i0^+}$



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Nucleon Form Factors in DIXEFT

• Reconstructing the form factors with $G_{E,M}^{p,n}(t) = \frac{1}{\pi} \int_{4M_{\pi}^2}^{\infty} dt' \frac{\text{Im}G_{E,M}^{p,n}(t')}{t'-t-i0^+}$



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Nucleon Form Factors in DIXEFT

• Proton radius puzzle: Discrepancy between e^{-p} scattering [0.8751(61) fm] and μ H [0.84087(39) fm]. [PDG 2016]

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Parametrization of the FF and Q² range used affects the extracted radius.

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Nucleon Form Factors in DIXEFT

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$$\rho_{1}(b) = \int_{0}^{\infty} d\Delta_{T} \frac{\Delta_{T} J_{0}(\Delta_{T} b)}{2\pi} F_{1}^{B}(t = -\Delta_{T}^{2}) = \int_{t_{thr}}^{\infty} dt \frac{K_{0}(\sqrt{tb})}{2\pi} \frac{\mathrm{Im} F_{1}^{B}(t)}{\pi}$$
$$\tilde{\rho}_{2}(b) = \int_{0}^{\infty} d\Delta_{T} \frac{-\Delta_{T}^{2} J_{1}(\Delta_{T} b)}{4\pi m_{B}} F_{2}^{B}(t = -\Delta_{T}^{2}) = \int_{t_{thr}}^{\infty} dt \frac{-\sqrt{t} K_{1}(\sqrt{tb})}{4\pi m_{B}} \frac{\mathrm{Im} F_{2}^{B}(t)}{\pi}$$

Nucleon Form Factors in DIXEFT

 $(\overline{D}) = \overline{D}$

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$$\langle B'|J^+(b)|B\rangle = [\ldots] \Big[\rho_1(b) + (2S^y)\cos\phi\tilde{\rho}_2(b)\Big]$$

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• The input necessary to compute the densities can be taken from experimental data (parametrizations) or theory.

• Relation between spectral functions and transverse densities.



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Nucleon Form Factors in DIXEFT

Charge Densities





Magnetization Densities





[J. M. Alarcón, C. Weiss, in preparation]

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Nucleon Form Factors in DIXEFT

Summary and Conclusions

Summary and Conclusions

- Through unitarity, it is possible to find a representation suited for
 ChEFT --> Predictions of the Nucleon Form factors.
- The results improve previous ChEFT calculations and are competitive with dispersion theory calculations.
- EM FFs have a much complex structure that what it seems.
- DIXEFT implements the constrains that allow to reconstruct the FFs with its full complexity:
 - Analyses of FF data.
 - Two photon exchange corrections to e⁻p scattering.
- Results used to understand "Proton Radius Puzzle" (PRad).
- New promising method to compute nucleon matrix elements from first principles (EM tensor, D-term, extension to G-parity odd, ...).

FIN



ullet Checking the parametrization of the spectral function at high t.

$$\Delta_E(t) \equiv G_E^V(t)[\exp] - \frac{1}{\pi} \int_{4M_\pi^2}^{t_{\max}} dt' \frac{\text{Im}G_E^V(t')}{t' - t - i0^+}$$



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Comparison with respect to the old results

[T. Bauer, J. Bernauer, S. Scherer, PRC 86 (2012)]



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Talk by Marko Horbatsch (JLab, 12/8/2017)

$(Q_{max}^2 = 0.2 \text{ GeV}^2) \chi_{red}^2$: green<1.08, blue<1.10, red<1.14



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Talk by Marko Horbatsch (JLab, 12/8/2017)

Is it consistent for the higher moments ?



(Courtesy of Marko Horbatsch)

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• We study the naturalness of the isovector moments by defining:

$$a_n = \frac{\langle r^{2n} \rangle^V}{(2n+1)!} = \frac{1}{\pi} \int_{4M_\pi^2}^\infty dt' \frac{\mathrm{Im}G^V(t')}{t'^{n+1}}$$

• If the integral were dominated by a certain region t', the ratio $\frac{a_{n+1}}{a_n}$ would be given by the average of 1/t' over this region.



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