Meson Structure Program at EicC (some updates)

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Workshop "Parton Distribution Functions at A Crossroad" Trento, 18- 22 September 2023

Contents



Motivations

- Sullivan process and event generators
- Simulation of pion form factor experiment at EicC
- Simulation of pion structure function experiment at EicC
- Simulation of kaon form factor experiment at EicC

Motivations



• The pion and kaon are Goldstone bosons, of quite light masses. Measuring their structures are helpful in checking the Emergent Hadron Mass (EHM) mechanism in Continuum Schwinger Function Methods (CSM), and the interplay between EHM and Higgs Boson mechanism. [J. Arrington et al, J. Phys. G: Nucl. Part. Phys. 48 (2021) 075106]

Understand the strong interaction

- The pion and kaon play the important roles in nuclear physics as the key nuclear fore carriers.
- Experimentally, there are TOO few data on the pion and kaon structures (form factor and structure function). Some Drell-Yan data from CERN & Fermilab decades ago, and the pion form factor data at low Q² from JLab and CERN pion-electron scattering. Very scarce data for the kaon structure.

Much fewer data compared to proton

Motivations

Understanding the internal meson structure and the dynamics in Lattice QCD



Fast advancement

[X. Gao, et al, Phys. Rev. D 104 (2021) 114515]

From Qi SHI's (**CCNU & BNL**) presentation at Lattice2023.

Also **DSE** predictions: [Lei Chang et al, Phys. Rev. Lett. 111 (2013) 141802; Fei Gao et al, Phys. Rev. D 96 (2017) 034024]

Instanton prediction: [E. Shuryak et al, PRD 103 (2021) 054028] Many other predictions

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Motivations



Understanding the internal meson structure and the dynamics in DSE/CSM



Fast advancement

[M. Ding et al, Phys. Rev. D 101
(2020) 054014; Zh.-F. Cui et al, Eur.
Phys. J. A 57 (2021) 5; Zh.-F. Cui et al,
Eur. Phys. J. C 80 (2020) 1064;T.
Nguyen et al, Phys. Rev. C 83 (2011)
062201; Chen Chen et al, Phys.Rev.D
93 (2016) 074021; Chao Shi et al,
Phys.Rev.D 98 (2018) 5, 054029]
Reviews: [C. D. Roberts et al, Prog.
Part. Nucl. Phys. 120 (2021) 103883;
J. Arrington et al, J. Phys. G: Nucl.
Part. Phys. 48 (2021) 075106]

Also **LQCD** predictions: [X. Gao et al, Phys. Rev. Lett. 128 (2022) 142003; H.-W. Lin et al, Phys. Rev. D 103 (2021) 014516; Z. Fan, H.-W. Lin, Phys. Lett. B 823 (2021) 136778; A. Salas-Chavira et al, Phys. Rev. D 106 (2022) 094510].....



Sullivan processes at small t (<0.6/0.9 GeV2) is sensitive to pion and kaon structures.





Exclusive processes for meson form factor measurements.





Leading baryon semi-inclusive deep inelastic scattering processes for meson structure measurements.

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To write an event generator and to estimate the statistics, we adapt the π -pole model for the differential cross-section:



Pion form factor measurement

$$\frac{d^{3}\sigma}{dQ^{2}dx_{B}dt} = \Gamma(Q^{2}, x_{B}, s) \left[\frac{d\sigma_{T}}{dt} + \epsilon \frac{d\sigma_{L}}{dt}\right]$$

$$\Gamma(Q^{2}, x_{B}, s) = \frac{\alpha y^{2}(1 - x_{B})}{2\pi x_{B}(1 - \epsilon)Q^{2}}$$

$$\epsilon = \frac{1 - y - \frac{Q^{2}}{4E^{2}}}{1 - y + \frac{y^{2}}{2} + \frac{Q^{2}}{4E^{4}}}$$
Pion pole and pion form factor
$$N\frac{d\sigma_{L}}{dt} = 4\hbar c(eg_{\pi NN}(t))^{2} \frac{-t}{(t - m_{\pi}^{2})^{2}} Q^{2}F_{\pi}^{2}(Q^{2})$$

$$N = 32\pi (W^{2} - m_{p}^{2})\sqrt{(W^{2} - m_{p}^{2})^{2} + Q^{4} + 2Q^{2}(W^{2} + m_{p}^{2})}$$

$$g_{\pi NN}(t) = g_{\pi NN}(m_{\pi}^{2}) \left(\frac{\Lambda_{\pi}^{2} - m_{\pi}^{2}}{\Lambda_{\pi}^{2} - t}\right)$$

$$F_{\pi}(Q^{2}) = \frac{1}{1 + Q^{2}/\Lambda_{\pi}^{2}}$$
Comparisons to difference while encoded and the encoded an

Comparisons to EIC generator is made by Weizhi. The event distributions show a little difference while the phasespace shapes are the same.





G. Xie et al., Chin. Phys. C 45, 053002 (2021)

Comparisons to EIC generator (RAPGAP) is made by Jixie. The cross sections and event distributions are very similar.

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Distributions of the invariant kinematics Q^2 , x_B , t, W^2 in the region of interests for various kinds of Sullivan processes.





Cross-section weighted distributions of the invariant kinematics

50 50 10⁶ 10⁵ 40 40 10⁵ 10⁴ 0² (GeV²) 10⁴ (GeV^2) 30 10³ 10³ \mathbf{Q}^2 20 10² 10² 10 10 10 10 0¹ 0.2 0.6 0.8 200 300 0.4 100 W^2 (GeV²) X_B

PionExculsiveElectroproduction demp_pion; demp_pion.SetQ2max(50); demp_pion.SetQ2min(1); demp_pion.SetxBmax(0.8); demp_pion.SetxBmin(0.001); demp_pion.SetTmax(0.6); demp_pion.SetTmin(0.01);



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Cross-section weighted kinematical distributions of the final electron and pion

No beam crossing



Beam crossing angle 50 mrad

IM

10⁴

 10^{3}

10²

10







In pion form factor experiment, the forward neutrons have higher energies and smaller scattering angles, comparing to the pion structure function experiments.

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Cross-section weighted kinematical distributions of the final neutron









The final-state electron and pion go to the central rapidity region due to the elastic scattering between the electron and the "pion cloud". Neutron as a spectator goes to the small angle very close to the proton beam.

Forward neutron detector is needed!

Cross-section weighted kinematical distributions of the final neutron



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Without ZDC cut

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 $N_{\rm i} = \epsilon L \overline{\sigma}_{\rm i} \Delta Q^2 \Delta t \Delta W^2 \frac{\partial x_{\rm B}}{\partial W^2}$ **Binning and stat. error estimation** \in [12, 16] GeV², with ZDC cut \in [16, 20] GeV², with ZDC cut $= 10^2$ \in [8, 12] GeV², with ZDC cut 10³ 0.6 0.6 0.6 10² -t (GeV²) (GeV^2) GeV 10² 10 - 10 0.2 0.2 10 0.2 100 200 W² (GeV²) 100 200 W² (GeV²) 100 200 W² (GeV²) 300 300 100 300 10 [20, 25] GeV², with ZDC cut \in [25, 30] GeV², with ZDC cut [30, 40] GeV², with ZDC cut 0.6 0.6 0.6 10 -t (GeV²) -t (GeV²) -t (GeV²) 0.2 10 0.2 0.2 300¹ 100 200 W² (GeV²) 100 200 W² (GeV²) 300 300 200 100 200 W² (GeV²) 200 ኻ

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Pion Form Factor is extracted from the model-dependent fits of longitudinal cross section.



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Stat. error plus some sys. error







Cross-section weighted distributions of the invariant kinematics

TaggedN_DIS dis; dis.SetQ2max(50); dis.SetQ2min(1); dis.SetxBmax(0.8); dis.SetxBmin(0.001); dis.SetTmax(0.6); dis.SetTmin(0.01); dis.SetxLmax(0.995); dis.SetxLmin(0.5);

Cross-section weighted kinematical distributions of the final electron and neutron

100 100F 10⁵ 10⁴ 10⁴ 10⁴ ¢∣of electron (°) φ₁of electron (°) 001 001 of neutron (°) −0 00 00 (∘) 10⁴ 50 50F 10³ of neutron 10³ 10³ 10³ 10² 10² 10² 10² -50 Ð -0-10 10 10 10 -100F -100 172 174 176 178 180 182 50 100 150 172 174 176 178 180 182 50 100 150 0 0 θ of neutron (°) θ of electron (°) θ of electron (°) θ of neutron (°) ZDC acceptance:±15 mrad ZDC acceptance:±15 mrad 10⁶ 10 10⁶ 10⁶ 10⁵ 10⁵ 10⁴ 10³ 10⁵ Counts Counts ⁵00 Counts 10³ 10³ 10⁴ 10⁴ 10² 10² 0 1 2 3 4 5 Angle between beam and neut. (°) 0 1 2 3 4 5 Angle between beam and neut. (°) -2 2 -2 0 2 0 η of electron η of electron

Beam crossing angle 50 mrad

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No beam crossing

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 $\begin{array}{l} \text{ZDC cut: } \theta_n < 15 \text{ mrad} \\ \text{Leading-neutron Cut (picking up} \\ \text{high energy neutron): } x_L > 0.75 \end{array}$

The combination of the above strict cuts reduces t-range significantly.

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Binning and stat. error estimation $N_{\rm i} = \epsilon L \overline{\sigma}_{\rm i} \Delta x_{\pi} \Delta Q^2 \Delta x_{\rm L} \Delta t (1 - x_{\rm L})$ $\epsilon = 50\%$





Binning and stat. error estimation $N_{\rm i} = \epsilon L \overline{\sigma}_{\rm i} \Delta x_{\pi} \Delta Q^2 \Delta x_{\rm L} \Delta t (1 - x_{\rm L})$ $\epsilon = 50\%$

















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Fitting $F_2^{\pi} \times f_{\pi}$ in different x_{π} bins, Q²: [5, 10] GeV²

The pion flux \mathbf{f}_{π} is t-dependent:



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Fitting $F_2^{\pi} \times f_{\pi}$ in different x_{π} bins, Q²: [10, 20] GeV²

The pion flux \mathbf{f}_{π} is t-dependent:

$$f_{\pi^+/p}(x_{\rm L},t) = \frac{1}{2\pi} \frac{g_{pn\pi}^2}{4\pi} (1-x_{\rm L}) \frac{-t}{(m_{\pi}^2 - t)^2} \exp\left(R_{n\pi}^2 \frac{t-m_{\pi}^2}{1-x_{\rm L}}\right)$$

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x_π = 0.275, Q2 = 25 GeV²

 $F_2^{\pi}(x_{\pi}, Q^2) f_{\pi}(t) / f_{\pi}(t=0.6)$

 $F_2^{\pi}(x_{\pi}, Q^2) \ f_{\pi}(t) \ / \ f_{\pi}(t=0.6)$

 $x_{\pi} = 0.325, Q2 = 25 \text{ GeV}^2$

Pion SF Experiment at EicC

 $x_{\pi} = 0.425, Q2 = 25 \text{ GeV}^2$

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x_π = 0.475, Q2 = 25 GeV²



 $x_{\pi} = 0.375, Q2 = 25 \text{ GeV}^2$

6

Fitting $F_2^{\pi} \times f_{\pi}$ in different x_{π} bins, Q²: [20, 30] GeV²

The pion flux \mathbf{f}_{π} is t-dependent:

$$f_{\pi^+/p}(x_{\rm L},t) = \frac{1}{2\pi} \frac{g_{pn\pi}^2}{4\pi} (1-x_{\rm L}) \frac{-t}{(m_{\pi}^2-t)^2} \exp\left(R_{n\pi}^2 \frac{t-m_{\pi}^2}{1-x_{\rm L}}\right)$$

Trent

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Pion SF Experiment at EicC

Fitting $F_2^{\pi} \times f_{\pi}$ in different x_{π} bins, Q²: [30, 50] GeV²

The pion flux \mathbf{f}_{π} is t-dependent:

$$f_{\pi^+/p}(x_{\rm L},t) = \frac{1}{2\pi} \frac{g_{pn\pi}^2}{4\pi} (1-x_{\rm L}) \frac{-t}{(m_{\pi}^2-t)^2} \exp\left(R_{n\pi}^2 \frac{t-m_{\pi}^2}{1-x_{\rm L}}\right)$$

IMP

 F_2^{π} are extracted from the model-dependent fits of the differential cross sections above.

QCD fit of the pseudo F_2^{π} data

Parametrization of initial PDF: same as JAM's input (π^+ meson)

JAM collaboration, Phys. Rev. Lett. 121 (2018) 15, 152001

. .

 $Q_0^2 = 1.0 \text{ GeV}^2$

$$q_v^{\pi} \equiv u_v^{\pi^+} = u^{\pi^+} - \bar{u}^{\pi^+} = \bar{d}_v^{\pi^+} = \bar{u}_v^{\pi^-} = d_v^{\pi^-} \qquad \qquad f(x_{\pi}, Q_0^2; \boldsymbol{a}) = \frac{N}{B(2+\alpha, \beta)} x_{\pi}^{\alpha} (1-x_{\pi})^{\beta}$$
$$q_s^{\pi} \equiv \bar{u}^{\pi^+} = d^{\pi^+} = s^{\pi^+} = \bar{s}^{\pi^+} \qquad \qquad \int_0^1 v(x) dx = 2, \int_0^1 x(v(x) + S(x) + g(x)) dx = 1.$$

Evolution with Q²: DGLAP evolution equation at NNLO

Utilize the package qcdnum18.00

M. Botje, Comput. Phys. Commun. 182 (2011) 490, arXiv:1005.1481, Erratum arXiv:1602.08383 (2016)
Pion SF Experiment at EicC



QCD fit of the pseudo F_2^{π} data

Fitting: least-square method, Utilize the "minuit" package

$$\chi^2 = \sum_i^N \frac{(D_i - T_i)^2}{\sigma_i^2}$$

 $C_{ii}(a) = (H^{-1})_{ii}$

TMinuit in ROOT

Error analysis method: Hessian matrix

$$(\Delta X)^2 = \Delta \chi^2 \sum_{i=1}^n \sum_{j=1}^n \frac{\partial X}{\partial a_i} C_{ij}(a) \frac{\partial X}{\partial a_j}$$

C. Han et al., Eur. Phys. J. C 81, 302 (2021); J. Pumplin et al., Phys. Rev. D 65, 014013 (2001); A.D. Martin et al., Eur. Phys. J. C 28, 455–473 (2003)

$$\Delta \chi^2 = \chi^2 - \chi_0^2 = \sum_{i=1}^n \sum_{j=1}^n H_{ij} (a_i - a_i^0) (a_j - a_j^0)$$

Pion SF Experiment at EicC



QCD fit of the pseudo F_2^{π} data



Systematic errors would be dominant!

One sees that the quark distributions are precisely determined with only the EicC pseudo-data, in the x range above 0.05. One also nds that the pion gluon distribution is given by the QCD evolution equations.

Nonperturbative input:

$$xq_v^{\pi} = 2.35x^{1.0}(1-x)^{1.34}$$

$$xq_s^{\pi} = 1.32x^{0.48}(1-x)^{9.36}$$

 $xg^{\pi} = 3.25x^{0.03}(1-x)^{15}$

Current ZDC performance and the spatial resolution requirement

ZDC Performance and Sys. Errors

How the energy and angular resolutions of ZDC affect the resolution of the t variable:

$$t = (P_{\rm p} - P_{\rm n})^2 \approx 2E_{\rm p}E_{\rm n}(\cos(\theta) - 1)$$

(with, $E_{\rm p} \approx p_{\rm p}$ and $E_{\rm n} \approx p_{\rm n}$)

$$\delta t = \frac{\partial t}{\partial E_{n}} \delta E_{n} + \frac{\partial t}{\partial \theta} \delta \theta$$
$$\frac{\delta t}{t} = \frac{\delta E_{n}}{E_{n}} + \frac{\sin(\theta)}{1 - \cos(\theta)} \delta \theta$$

$$\frac{\delta E_{\rm n}}{E_{\rm n}} = \frac{47\%}{\sqrt{E_{\rm n}}} + 2.5\%$$

$$\frac{\delta E_{\rm n}}{E_{\rm n}} \approx 14.5\%$$
 for 15 GeV neutron

 $\frac{\sin(\theta)}{1 - \cos(\theta)} \delta\theta \text{ is suggested to be around } 14.5\%$

At $\theta \sim 7$ mrad, the angular resolution $\delta\theta$ should be around 0.00051. And the position resolution should be around 12 m $\times \delta\theta$ = 0.6 cm.

This kind of position resolution is very challenging!

ZDC Performance and Sys. Errors



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ZDC Performance and Sys. Errors



With the central detector, the neutron information also can be given with the missing-particle method.

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Systematic Uncertainties from ZDC Resolutions





For the form factor experiment, the momentum transfer t also can be evaluated with the virtual photon and final pion. Need more studies!

For the structure function experiment, the t can only be measured with ZDC. Thus ZDC introduces more uncertainties to the pion structure functions.



The typical resolution of t: $\sqrt{2} \times 14.5\% = 21\%$

Preliminary MC study shows that the systematic uncertainty from the resolution of t is 6-19% for the pion structure function.

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The invariant kinematical distributions of $ep \rightarrow eK^+\Lambda$ from the event generator



KaonExclusiveElectroproduction demp_kaon;

demp_kaon.SetTmax(0.6); demp_kaon.SetTmin(0.01); demp_kaon.SetQ2max(50); demp_kaon.SetQ2min(1); demp_kaon.SetxBmax(0.8); demp_kaon.SetxBmin(0.001); demp_kaon.SetElecBeamEnergy(3.5); demp_kaon.SetProtBeamEnergy(20); demp_kaon.SetBeamCrossAngle(0.0);

demp_kaon.SetOutputFileName("DEMP-kaon-pole-at-EicC.root");
demp_kaon.SetEvtFileOutput(1);
demp_kaon.SetQuiet(1);
demp_kaon.SetSamplingMode(1);
demp_kaon.Generate(1000000);

The evt file can be easily used in the EicCRoot and geant4 simulations.

0	7						
Ν	Id	Ist	M1	M2	DF	DL	px py pz E t x y z
Θ	11	1	- 1	- 1	- 1	- 1	-0.275297 1.14401 -3.31999 3.52234 0 0 0 0
1	321	1	- 1	- 1	- 1	- 1	0.628167 -0.733735 4.47504 4.60464 0 0 0 0
2	2212	1	- 1	- 1	- 1	- 1	-0.304412 -0.354582 11.7698 11.8163 0 -0.17107 -0.1989 7.42851
3	-211	1	- 1	- 1	- 1	- 1	-0.0484579 -0.0556922 3.55318 3.55668 0 -0.17107 -0.1989 7.42851
4	2112	1	- 1	- 1	- 1	- 1	-0.376788 -0.374849 14.2433 14.2841 0 -0.17107 -0.1989 7.42851
5	22	1	- 1	- 1	- 1	- 1	0.0713131 -0.0310818 0.888843 0.892241 0 -0.17107 -0.1989 7.42851
6	22	1	- 1	- 1	- 1	- 1	-0.0473955 -0.00434271 0.190813 0.19666 0 -0.17107 -0.1989 7.4285
1	7						
Ν	Id	Ist	M1	M2	DF	DL	px py pz E t x y z
Θ	11	1	- 1	- 1	- 1	- 1	-0.465765 -1.19213 -3.29303 3.53301 0 0 0 0
1	321	1	- 1	- 1	- 1	- 1	0.450413 0.950829 4.31165 4.46553 0 0 0 0

Cross-section weighted kinematical distributions of the final electron and kaon

No beam crossing



Beam crossing angle 50 mrad



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For exclusive kaon production in the kaon pole model, the final baryons (Lambda and its decay) have larger scattering angle and wider distributions.

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Cross-section weighted distributions of energy, angle, decay vertex of Lambda





10

 10^{2}

10

0

-2

-2

Kinematical distributions of Lamba decays ($p\pi^-$, $n\gamma\gamma$)

Blue: center detector system Green: EDT





Long Λ decay length reduces the EDT acceptance.

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0.2

Acceptance (%)

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0.8

Kaon FF Experiment at EicC

The left plot shows the acceptances of the channel $ep \to eK^+\Lambda$ by measuring the decay protons and decay neutrons. The proton channel and neutron channel are complementary.

0.4

-t (GeV²)

nd neutron channel are complementary.

 $\Lambda \rightarrow n\gamma\gamma$ channel, Q2 \in [10, 20] GeV²

0.6









The differential cross sections of $ep \rightarrow eK^+\Lambda$ in different Q² bins, from measurements of both the proton decay channel and the neutron decay channel of Lambda.

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这里仅仅展示了统计误差预测的初步结果。统计误差会和模型有关,需要更多结果横向比较。 对于Pion的形状因子实验,系统误差可能是主要的误差来源。

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https://github.com/rong-wang-impcas/tagged-neutron-DIS https://github.com/rong-wang-impcas/DEMP-generator-pion-pole-model https://github.com/rong-wang-impcas/tagged-Lambda-DIS https://github.com/rong-wang-impcas/DEMP-generator-kaon-pole-model

PionExclusiveElectroproduction demp_pion; demp_pion.SetOutputFileName("DEMP-pion-pole-at-EicC.root"); demp_pion.SetEvtFileOutput(1);

Options for output ROOT and evt files

demp_pion.SetElecBeamEnergy(3.5); demp_pion.SetProtBeamEnergy(20); //demp_pion.SetBeamCrossAngle(0.05); demp_pion.SetSamplingMode(1); demp_pion.Generate(20000);

double dsigmaT();Just modifying the following functions,
the users can apply any models they like.double dsigmaT();the users can apply any models they like.double dsigmaLT();double d4sigma_dQ2dxBdtdPhi(double Q2, double xB, double t, double Phi);double d3sigma_dQ2dxBdt(double Q2, double xB, double t);

Summary





Assumptions:

- 5(e⁻) x 100(*p*)
- Integrated L=20 fb⁻¹/yr
- Clean identification of exclusive p(e,e'π⁺n) events
- t reconstruction resolution based on ECCE detector design
- Syst. Unc: 2.5% pt–pt and 12% scale
- $R = \sigma_L / \sigma_T = 0.013 0.14$ at lowest -t from VR model, and $\delta R = R$ syst. unc. in model subtraction to isolate σ_I .
- π pole dominance at small -t confirmed in ²H π⁻/π⁺ ratios.

Garth Huber's slide for US-EIC (2022-04-29)



Up-to-date simulation result With realistic ZDC acceptance

Systematic uncertainty dominates for the measurement!

Summary





Acceptance for the deep exclusive kaon production process $ep \rightarrow eK^+\Lambda$ at EicC. This result should have small model dependence, as it is pure geometric.

The statistic error projection for the kaon form factor at EicC, updated with the current designs of ZDC and EDT. This result is model-dependent.

Summary





Pion structure function extracted from differential cross sections (stat. errors only)

[Please find more information from the talk in the last CDR workshop and R. Wang, W. Xiong, Y. Liang, X. Chen, Few Body Syst. 64 (2023) 28]



Pion PDFs determined by a QCD analysis of the pseudo-data (for π^+ meson, at $Q^2 = 10$ GeV²)

Summary and Outlook



- The feasibilities of the meson structure experiments at EicC are studied, and the statistical error projections are given.
- Controlling the systematical errors are very challenging. Without any corrections, the systematic errors from ZDC would be around 20%. We need more studies on the ZDC performance in the future.
- The systematic uncertainty from the background contamination also should be studied in the future.

Backup: Ion Forward Detector Complex



Without detector acceptance



With detector acceptance

Without detector acceptance



With detector acceptance





The differential cross sections of $ep \rightarrow eK^+\Lambda$ in different Q² bins, from measurements of only the proton decay channel of Lambda.

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 $Q^2 = 4 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有6-19%的 系统误差!

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 $Q^2 = 7.5 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有8-27%的 系统误差!



 $Q^2 = 15 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有7-22%的 系统误差!



 $Q^2 = 25 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有10-60%的 系统误差!

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 $Q^2 = 40 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有10-50%的 系统误差!



 $Q^2 = 4 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有6-19%的 系统误差!



 $Q^2 = 7.5 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有8-27%的 系统误差!



 $Q^2 = 15 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有7-22%的 系统误差!



 $Q^2 = 25 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有10-60%的 系统误差!



 $Q^2 = 40 \text{ GeV}^2$

典型的动量转移t的相对误差 为√2×14.5% = 21%

初步MC研究表明,由于我 们ZDC测量到的t的范围较小, t变量21%的相对误差导致提 取到的结构函数有10-50%的 系统误差!

Backup: Justifications of the MC Models



Rong WANG (王 荼)

Trento, 18-22 September 2023
LN DIS产生器



IMP



Deep Exclusive Pion Production产生器

我们的π-pole模型与 纵向截面实验数据比较

数据来自 JLab Fpi collaboration, Phys. Rev. C 78 (2008) 045202

Rong WANG (王 茶)

IMP



Deep Exclusive Pion Production产生器

参数化的横向截面模型 与横向截面实验数据比

模型A:

参数化的模型和实验数据均来自 Phys. Rev. C 78 (2008) 045202

$$\frac{d\sigma_T}{dt} = \left(\frac{0.74}{Q^2} + \frac{1.25}{Q^4} + \frac{0.57|t|}{(|t| + m_\pi^2)^2}\right) \frac{8.54}{\left(W^2 - m_N^2\right)^2}$$

IMP

In the hard scattering regime, QCD scaling predicts $\sigma_L \propto 1/Q^6$ and $\sigma_T \propto 1/Q^8.$



Deep Exclusive Pion Production产生器







Backup

$\sqrt{2}$
(IMP)

	Confidence level (probability contents desired inside				
Number of	hypercontour of $\chi^2 = \chi^2_{ m min} + UP$)				
Parameters	50%	70%	90%	95%	99%
1	0.46	1.07	2.70	3.84	6.63
2	1.39	2.41	4.61	5.99	9.21
3	2.37	3.67	6.25	7.82	11.36
4	3.36	4.88	7.78	9.49	13.28
5	4.35	6.06	9.24	11.07	15.09
6	5.35	7.23	10.65	12.59	16.81
7	6.35	8.38	12.02	14.07	18.49
8	7.34	9.52	13.36	15.51	20.09
9	8.34	10.66	14.68	16.92	21.67
10	9.34	11.78	15.99	18.31	23.21
11	10.34	12.88	17.29	19.68	24.71
	If FCN is $-\log(\text{likelihood})$ instead of χ^2 , all values of UP				
	should be divided by 2.				

From minuit manual

Table 7.1: Table of UP for multi-parameter confidence regions