

Fast and Collisional Neutrino Flavor Instabilities in Core-collapse Supernovae with Boltzmann Neutrino Transport

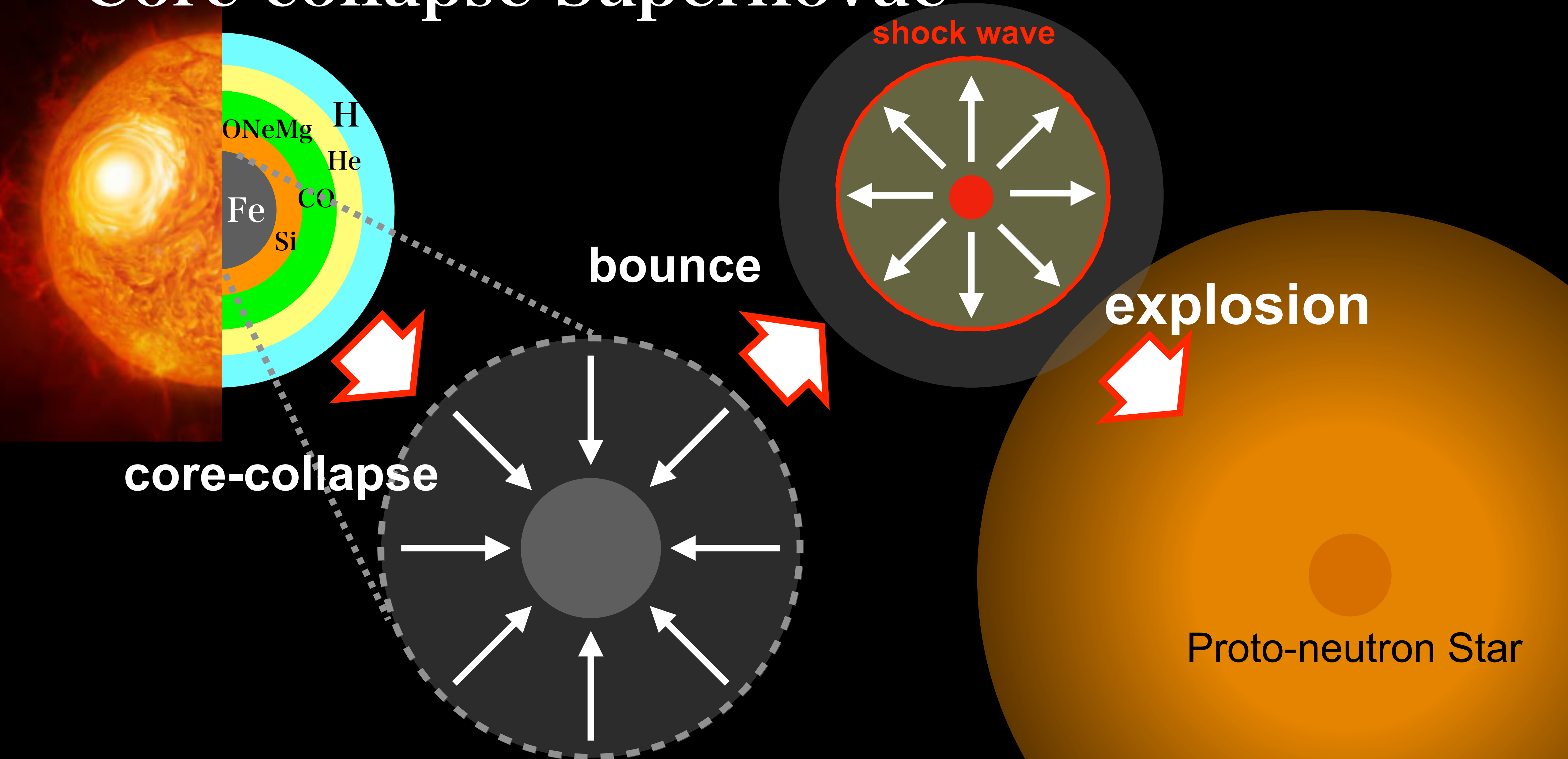
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The simulation data is available thanks to

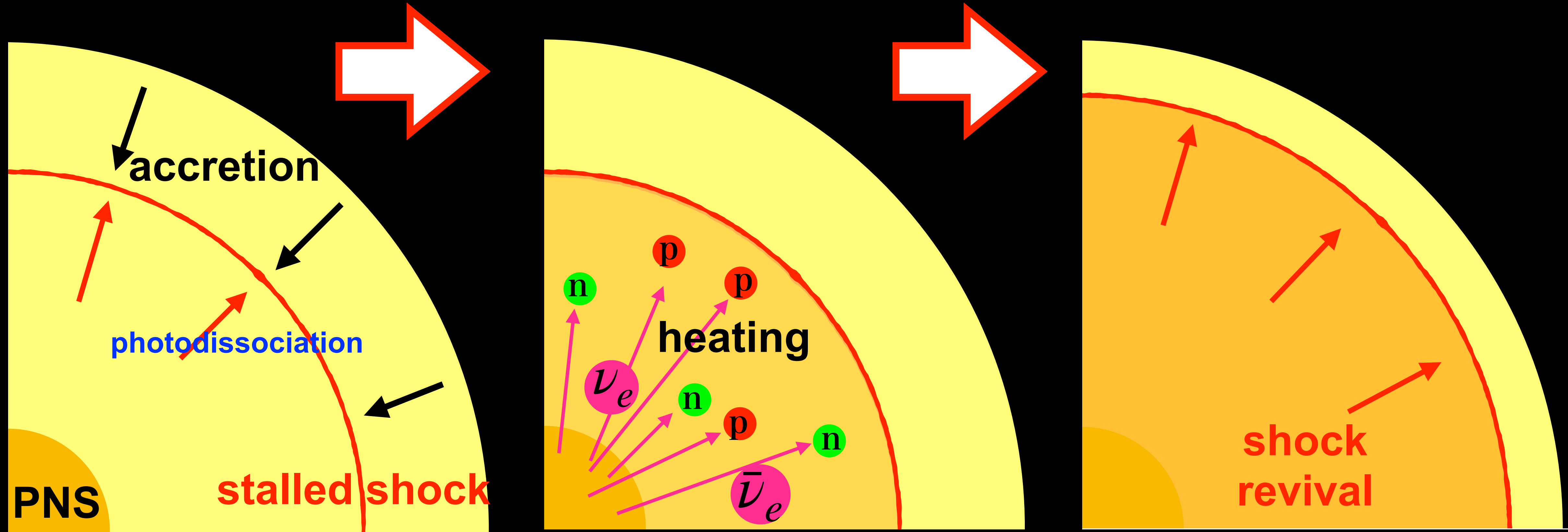
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Kohsuke Sumiyoshi

Core-collapse Supernovae



Neutrino Heating Mechanism

Neutrinos carry $\gtrsim 99\%$ of supernova energy. Most of them escape, but some of them reluctantly give energy to surrounding matter.



shock wave stalls due to energy loss and accretion ram pressure

neutrinos from the center heats the matter behind the shock

shock gradually propagates outward

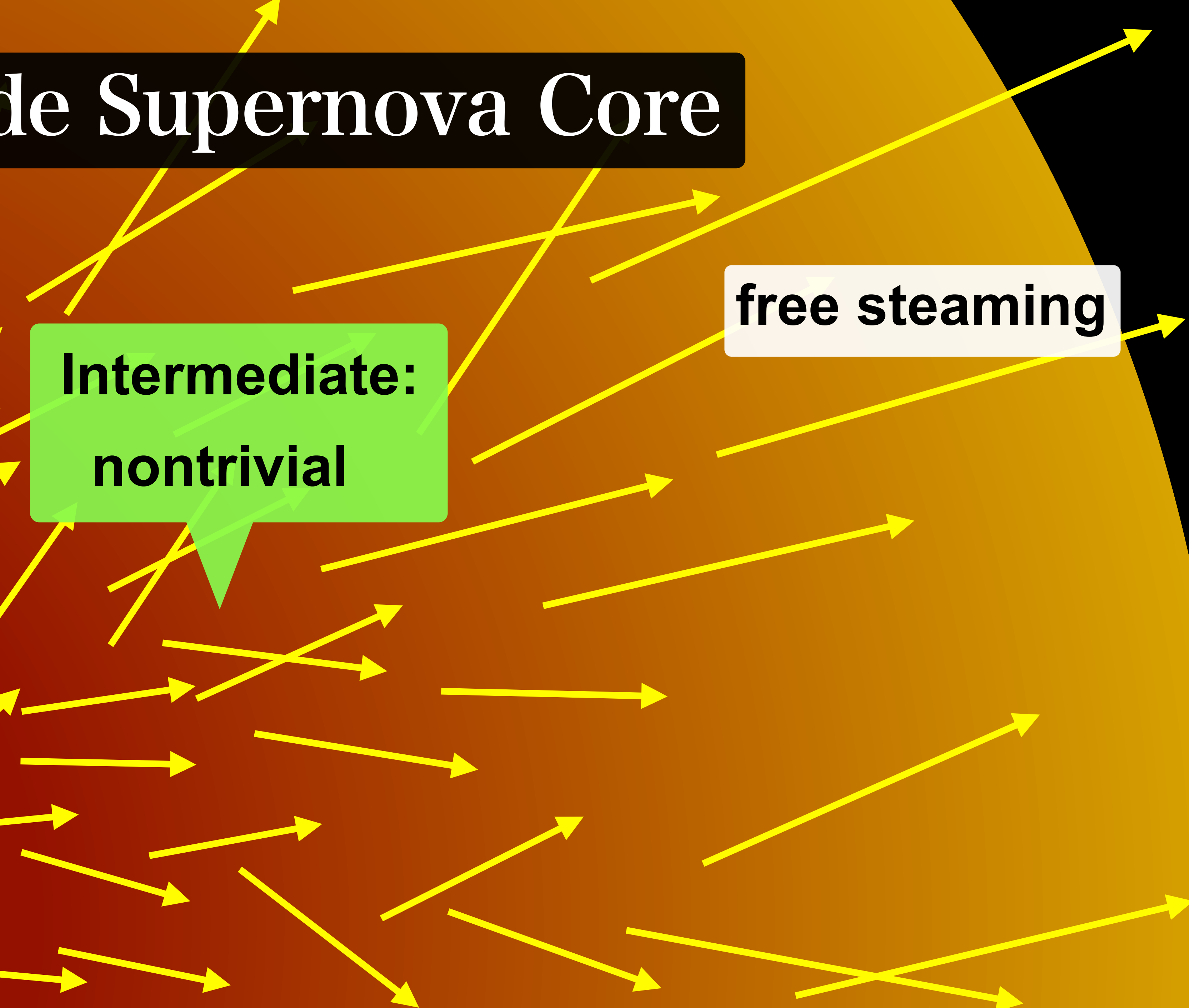
Neutrinos inside Supernova Core



**Thermal-equilibrium
Isotropic**

**Intermediate:
nontrivial**

free steaming



Boltzmann Equation

$$p^\alpha \frac{\partial f}{\partial x^\alpha} - \Gamma_{\alpha\beta}^i p^\alpha p^\beta \frac{\partial f}{\partial p^i} = \left[\frac{\delta f}{\delta t} \right]_{\text{collision}}$$

Spatial advection

Momentum advection

Collision terms

Distribution function in
phase space $f(x^\mu, p^i)$

Where

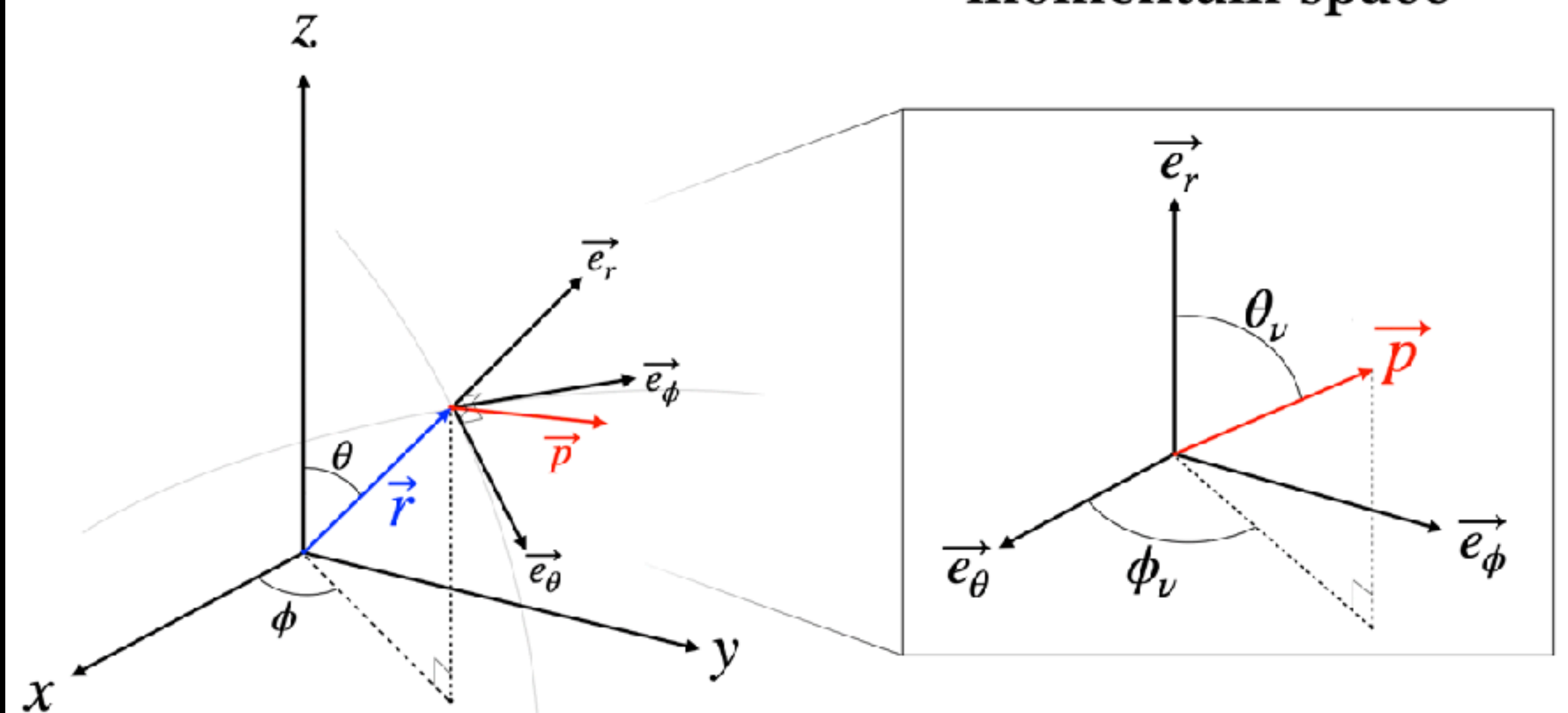
7 (1+3+3) dimensional equation

Time

Momentum

configuration space

momentum space



Example of Alternative Method

Two-moment transport

Save numerical cost by

$$f(r, \theta, \phi, \epsilon, \theta_\nu, \phi_\nu)$$

Boltzmann equation

$$\frac{\partial f}{\partial t} + p^i \frac{\partial f}{\partial x^i} + \dot{p}^i \frac{\partial f}{\partial p^i} = C$$



Multiply

directional cosines
and angle-integrate

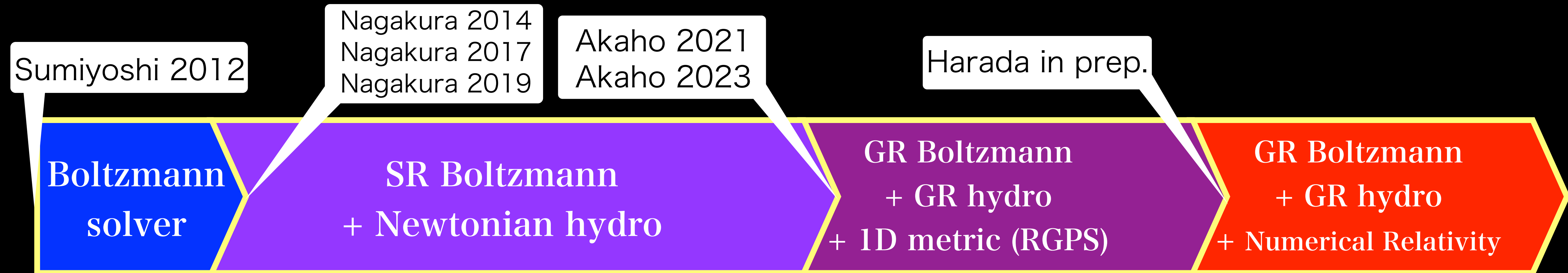
$$\begin{aligned} E(r, \theta, \phi, \epsilon) \\ F^i(r, \theta, \phi, \epsilon) \\ P^{ij}(r, \theta, \phi, \epsilon) \end{aligned}$$

Moment equations

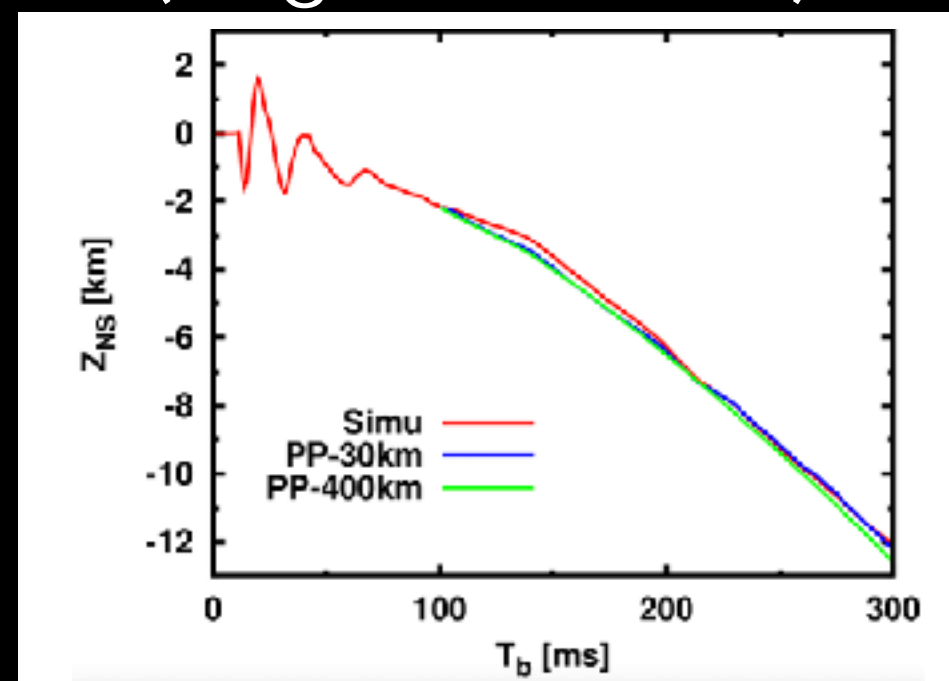
$$\frac{\partial E}{\partial t} = L_1(E, F^i, P^{ij}) \quad \frac{\partial F^i}{\partial t} = L_2(E, F^i, P^{ij})$$

$$\text{Closure: } P_{M1}^{ij} = \frac{3\chi - 1}{2} P_{thin}^{ij} + \frac{3(1 - \chi)}{2} P_{thick}^{ij}$$

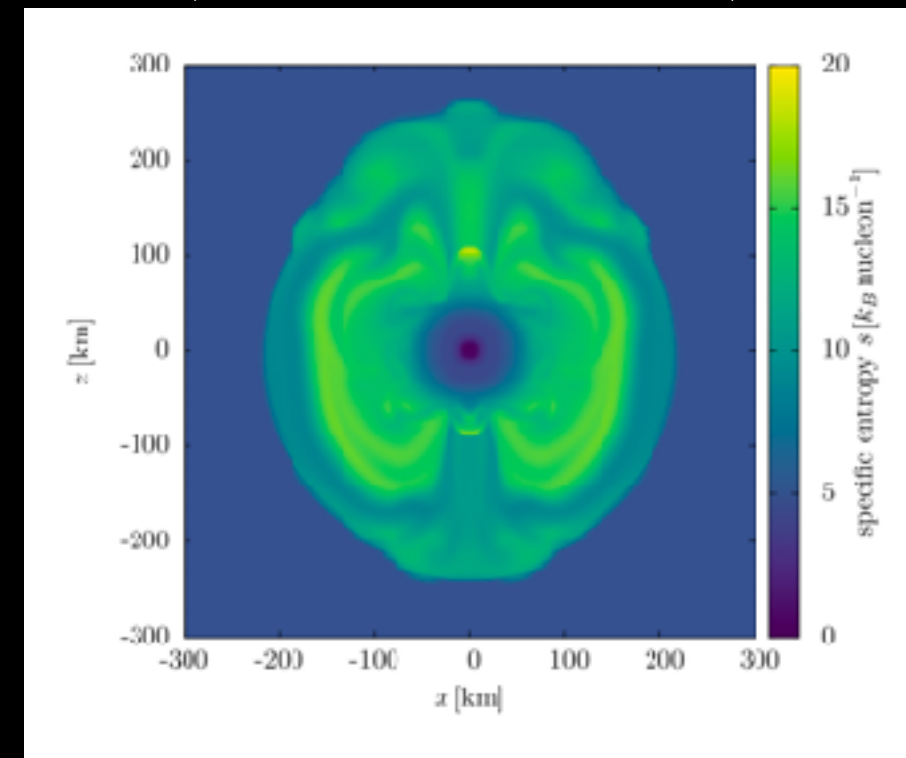
Boltzmann Radiation-hydro Simulation



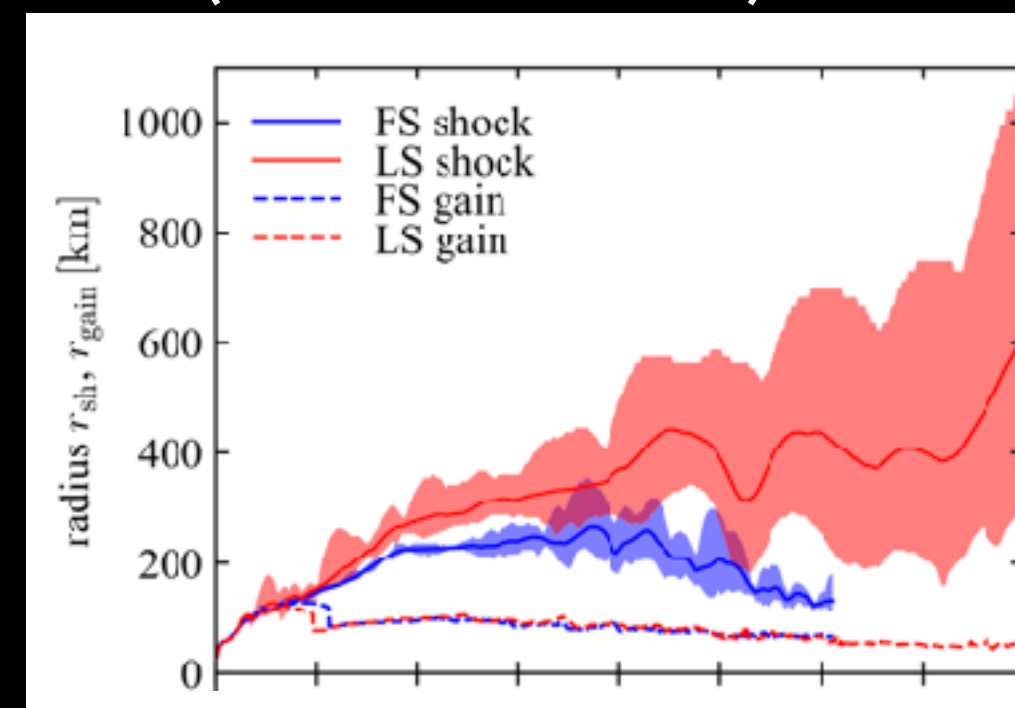
PNS kick
(Nagakura 2019)



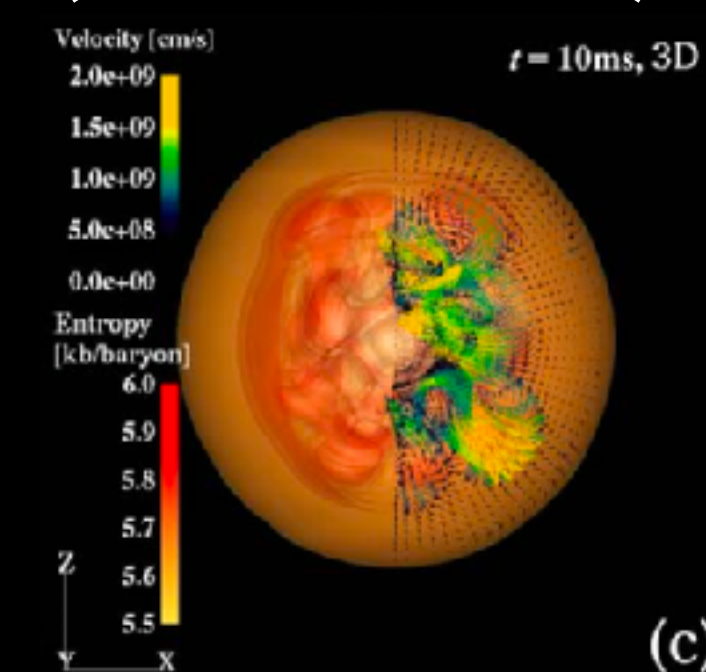
Rotating progenitor
(Harada 2019)



EOS dependence
(Harada 2020)



3D simulation
(Iwakami 2020)



We are open for collaboration by using our data (e.g. calibrating approximate methods)

Collective Neutrino Oscillation

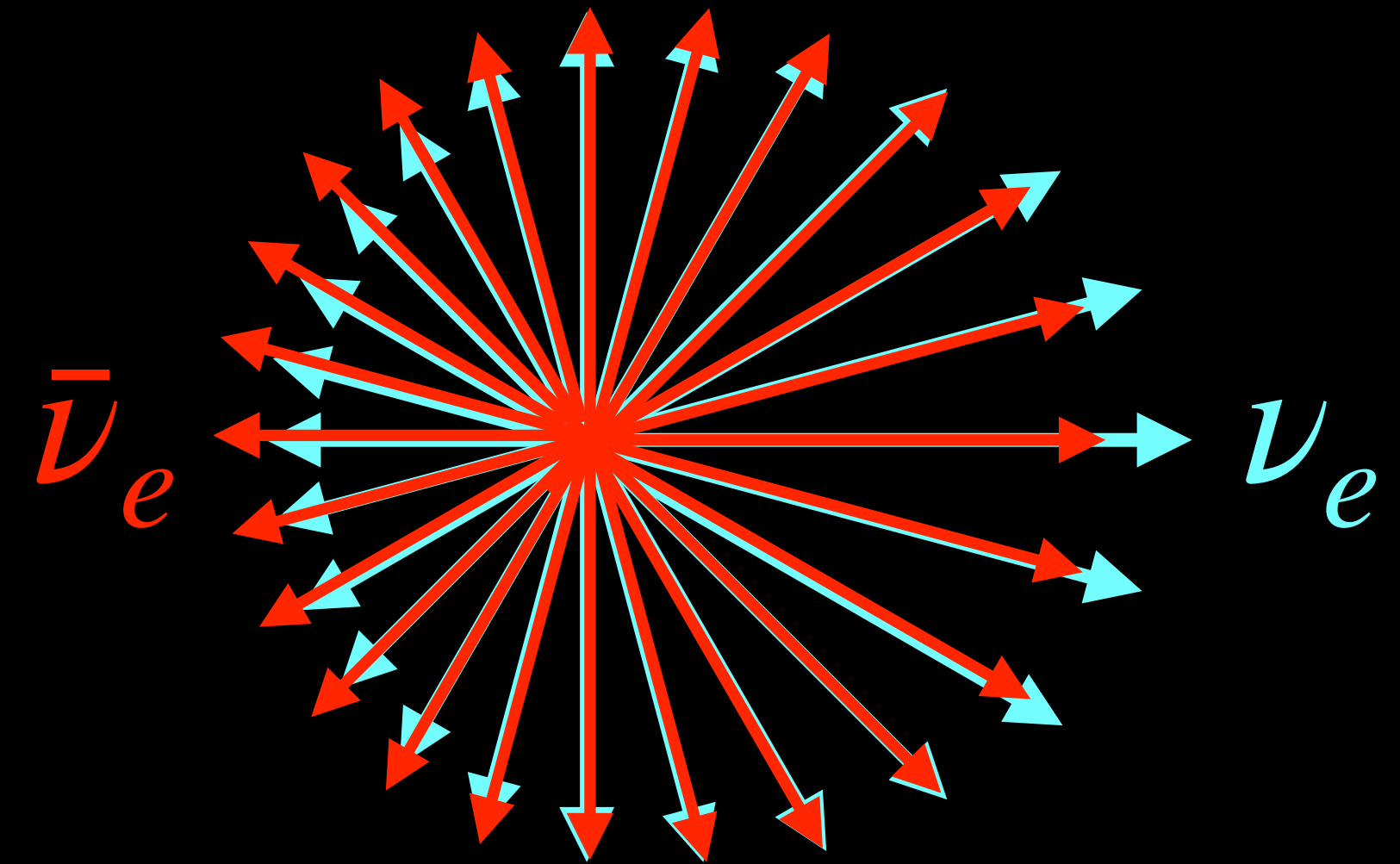
$$i(\partial_t + \vec{v} \cdot \nabla)\rho = [H_{\text{vac}} + H_{\text{mat}} + H_{\nu\nu}, \rho] + iC$$

Vacuum potential Matter potential Neutrino self-interaction potential Collision terms

- **Slow instability:** induced by neutrino energy crossing
- **Fast instability:** induced by neutrino angular crossing
- **Collisional instability:** induced by difference of collision rates

Fast Flavor Instability

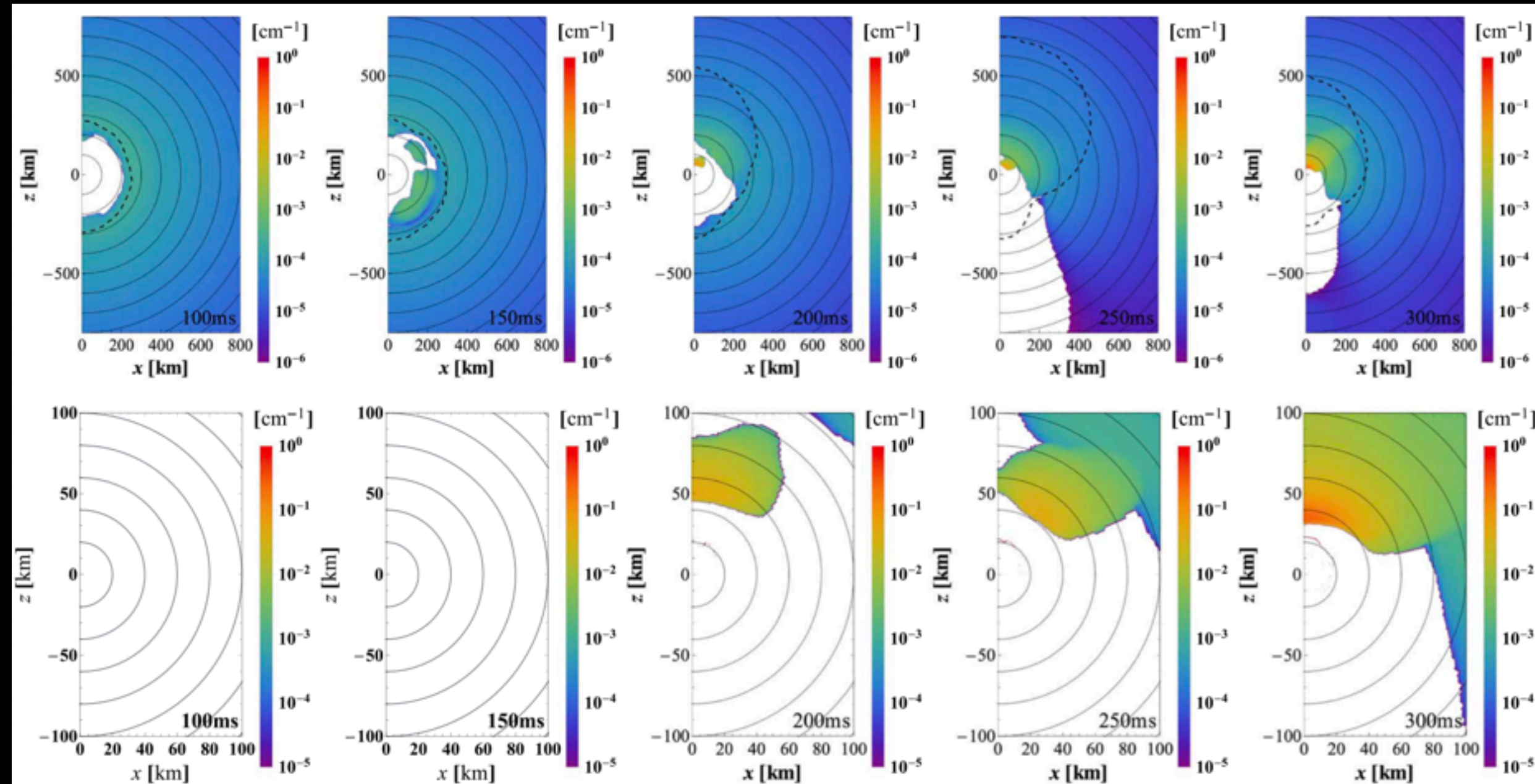
- Fast flavor instability (FFI) is expected to be important ingredient for CCSN theory.
- Existence of FFI is known to be equivalent to the existence of neutrino flavor lepton number (NFLN) crossing (Morinaga 2022, Dasgupta 2022)



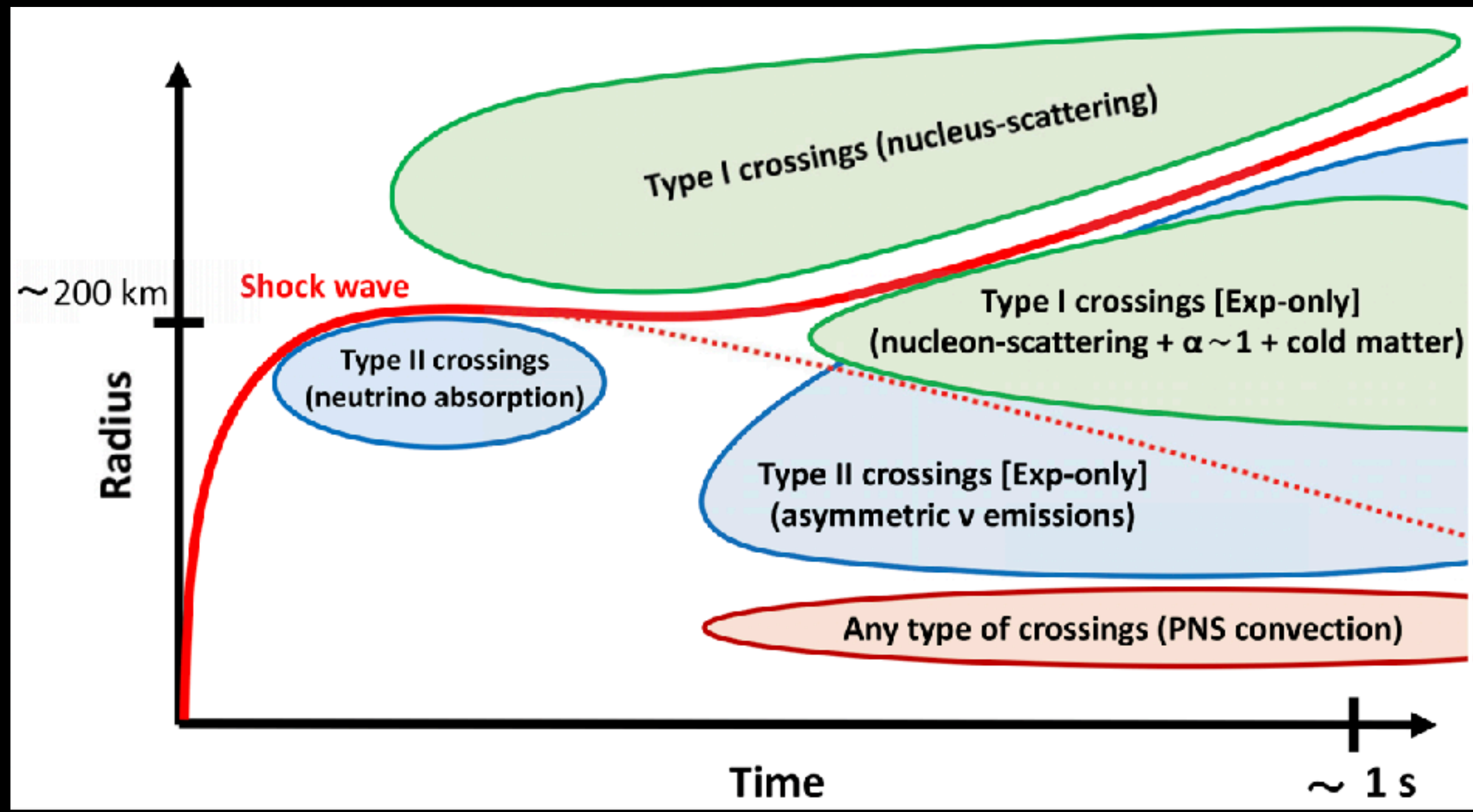
Several methods have been proposed to detect ELN crossing only from angular moments

- Zero-mode search (Dasgupta 2018)
- Polynomial method (Abbar 2020)
- Fitting formulae (Nagakura 2021a,b, Johns 2021)
- Machine learning (Abbar 2023)

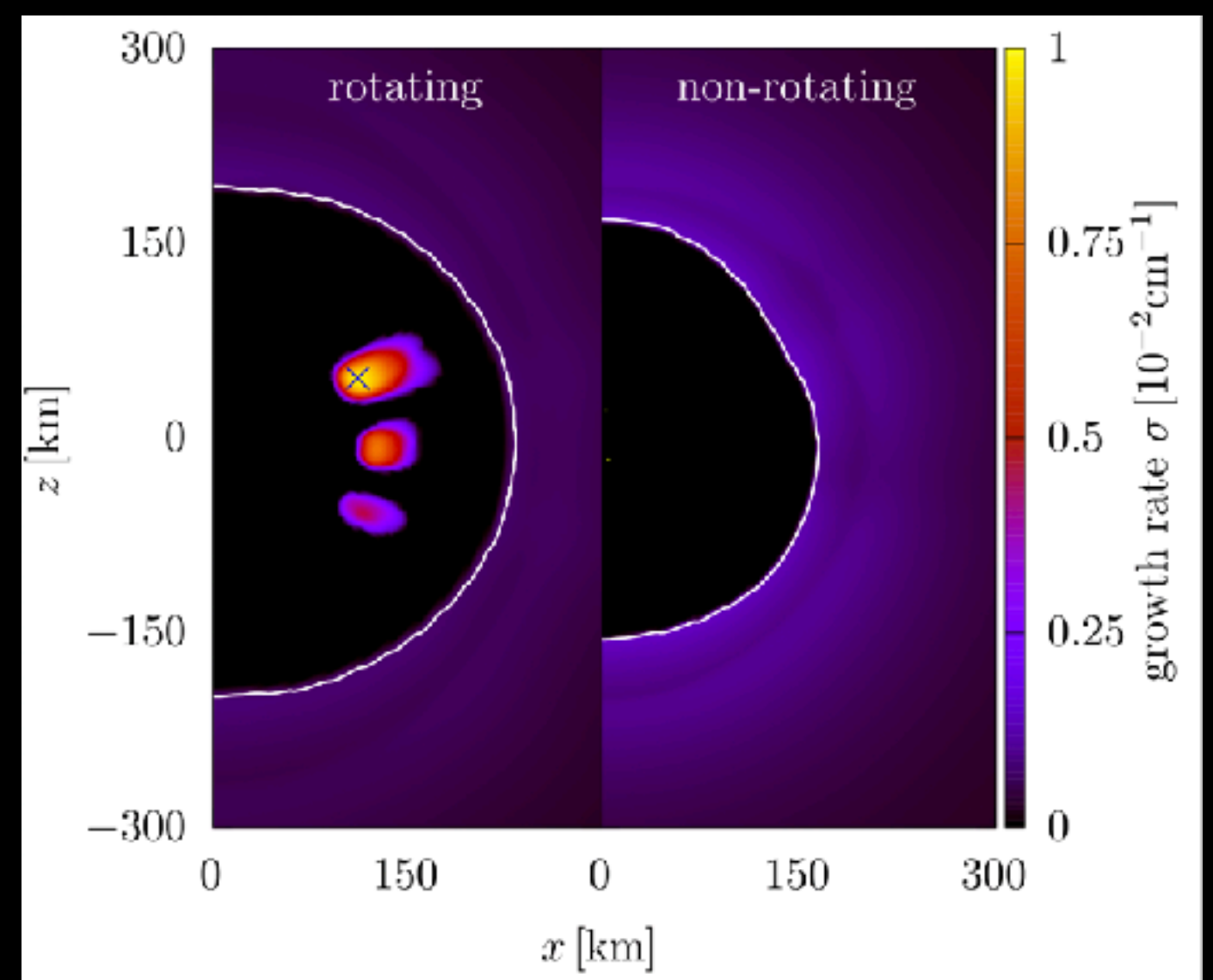
FFI in CCSN Models

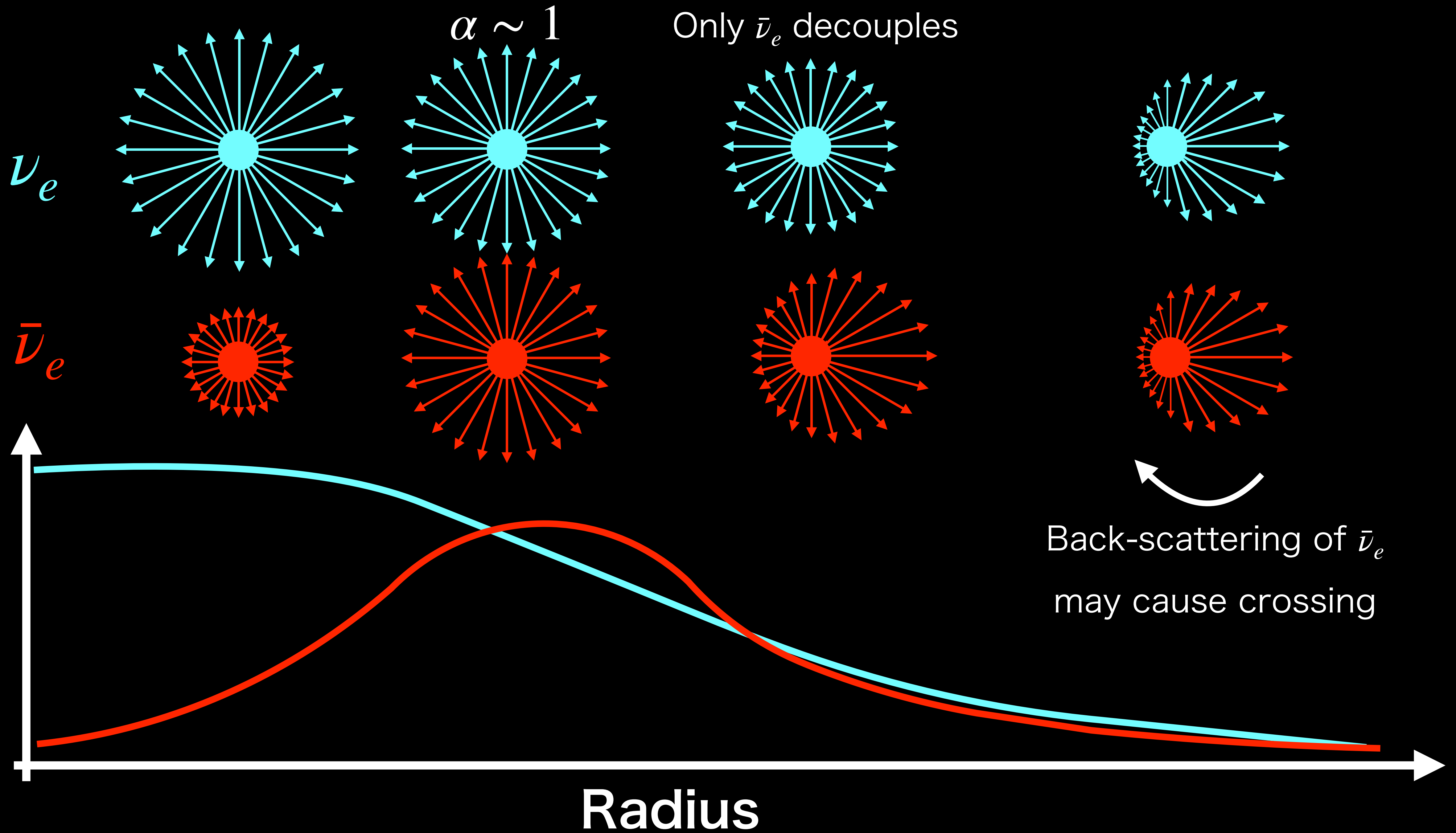


Nagakura 2021



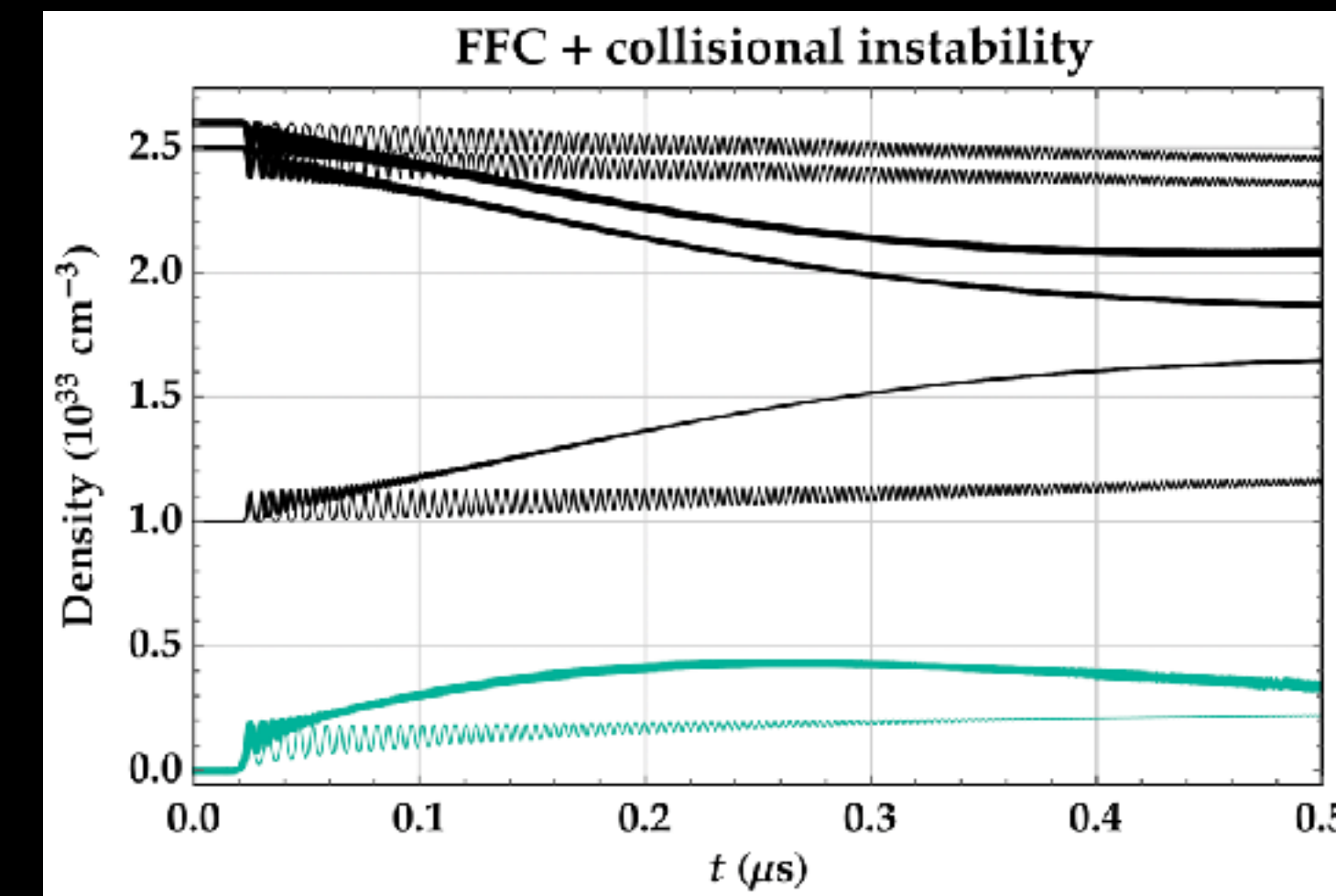
Harada 2022



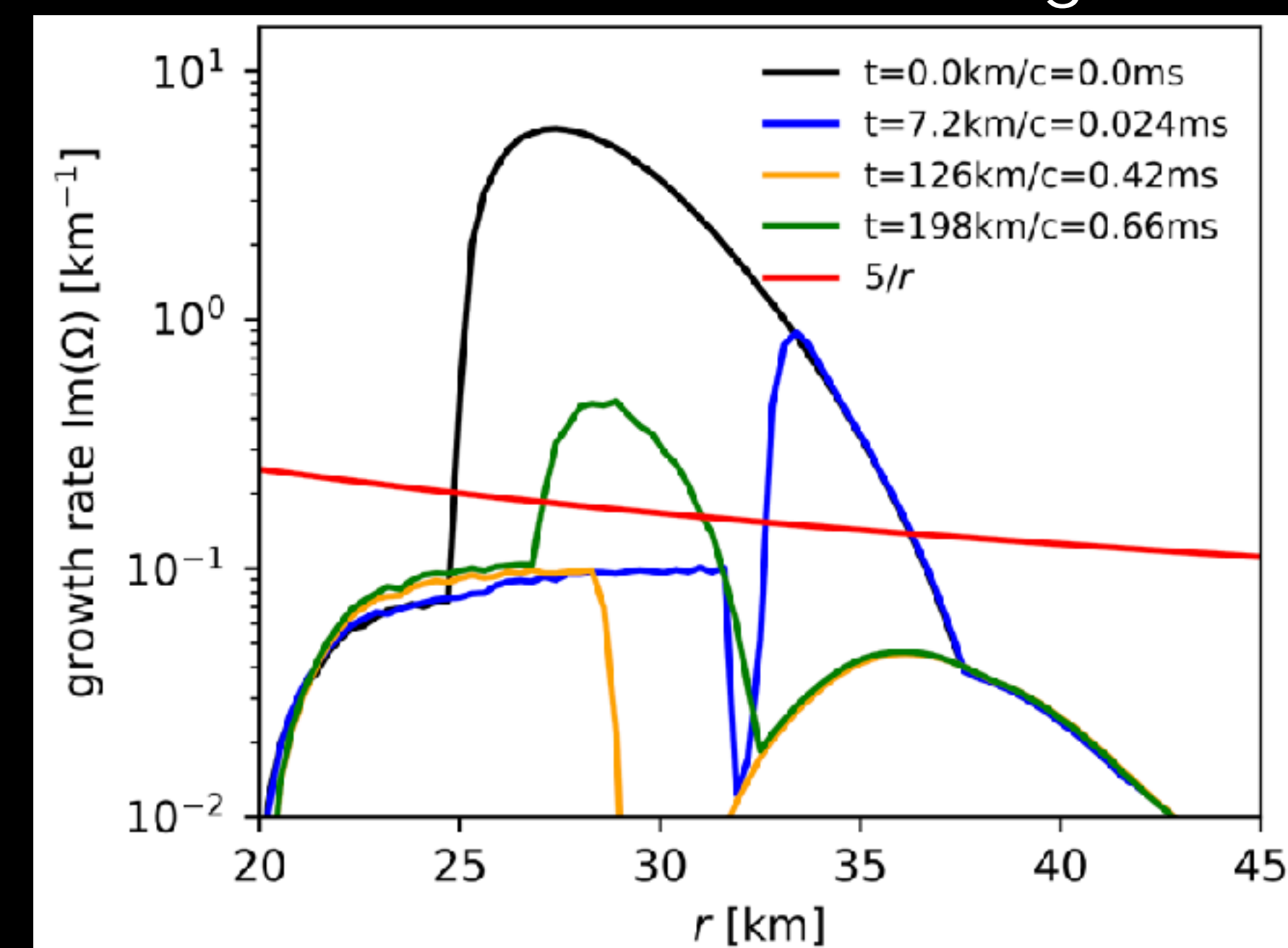


Collisional Flavor Instability

- Recently, existence of the “collisional flavor instability” (CFI) has been pointed out (Johns 2023).
- CFI is induced by asymmetry of collision terms, which can exist without FFI.
- If CFI and FFI coexists, they may enhance each other.

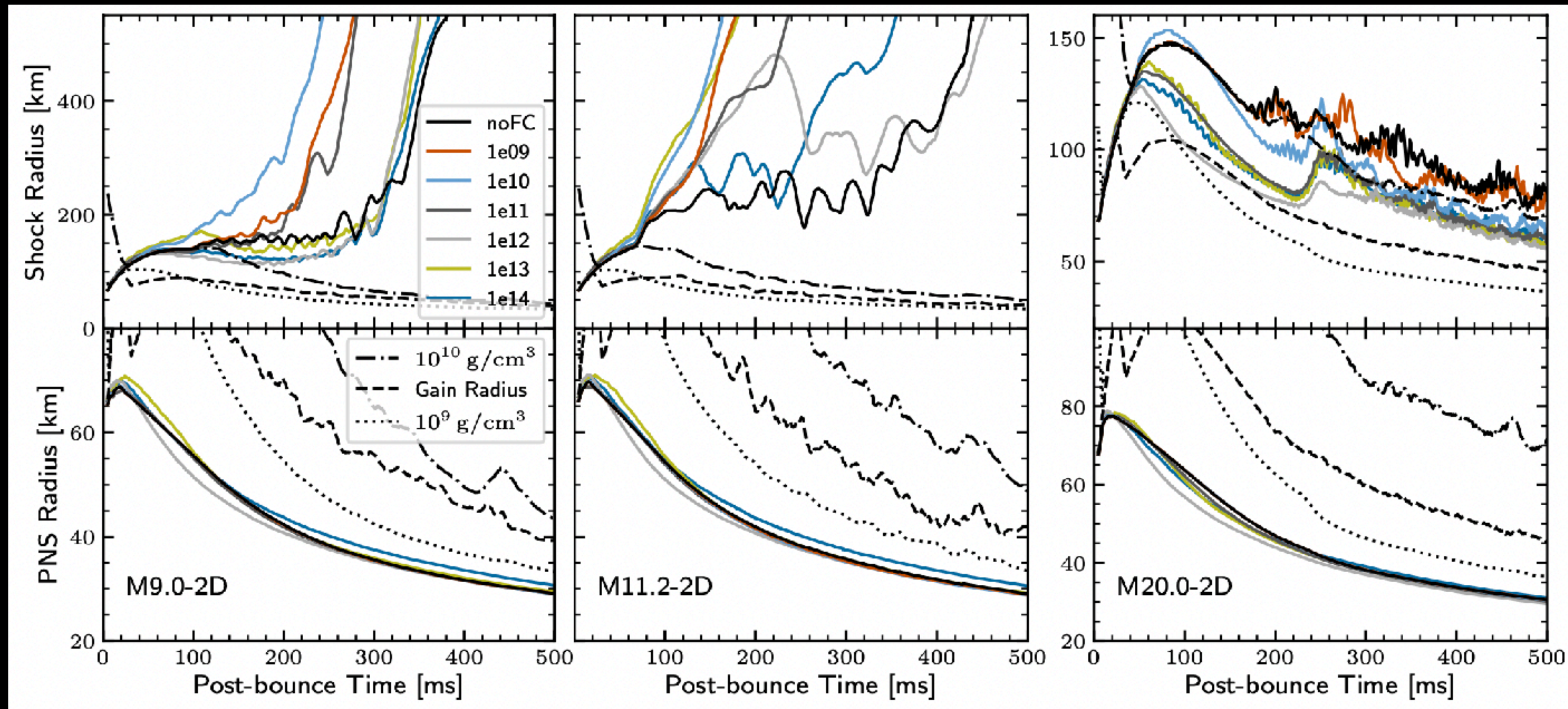


Xiong 2023



CCSN Simulations with Flavor Equilibration

- CCSN simulations are performed by forcing flavor equilibrium for the threshold matter density, which is parametrically changed.

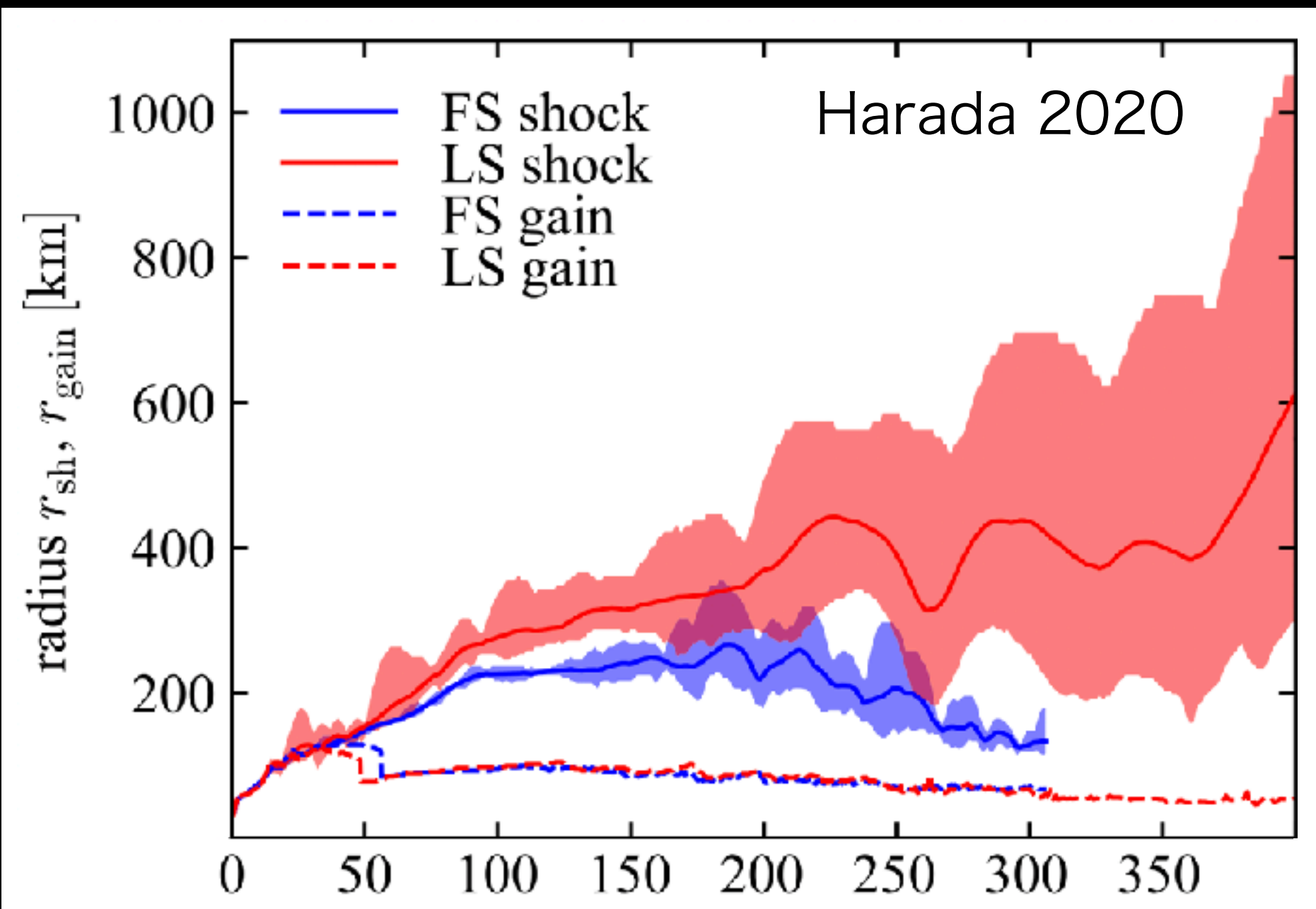


Ehring 2023

FFI and CFI on Boltzmann CCSN Simulation

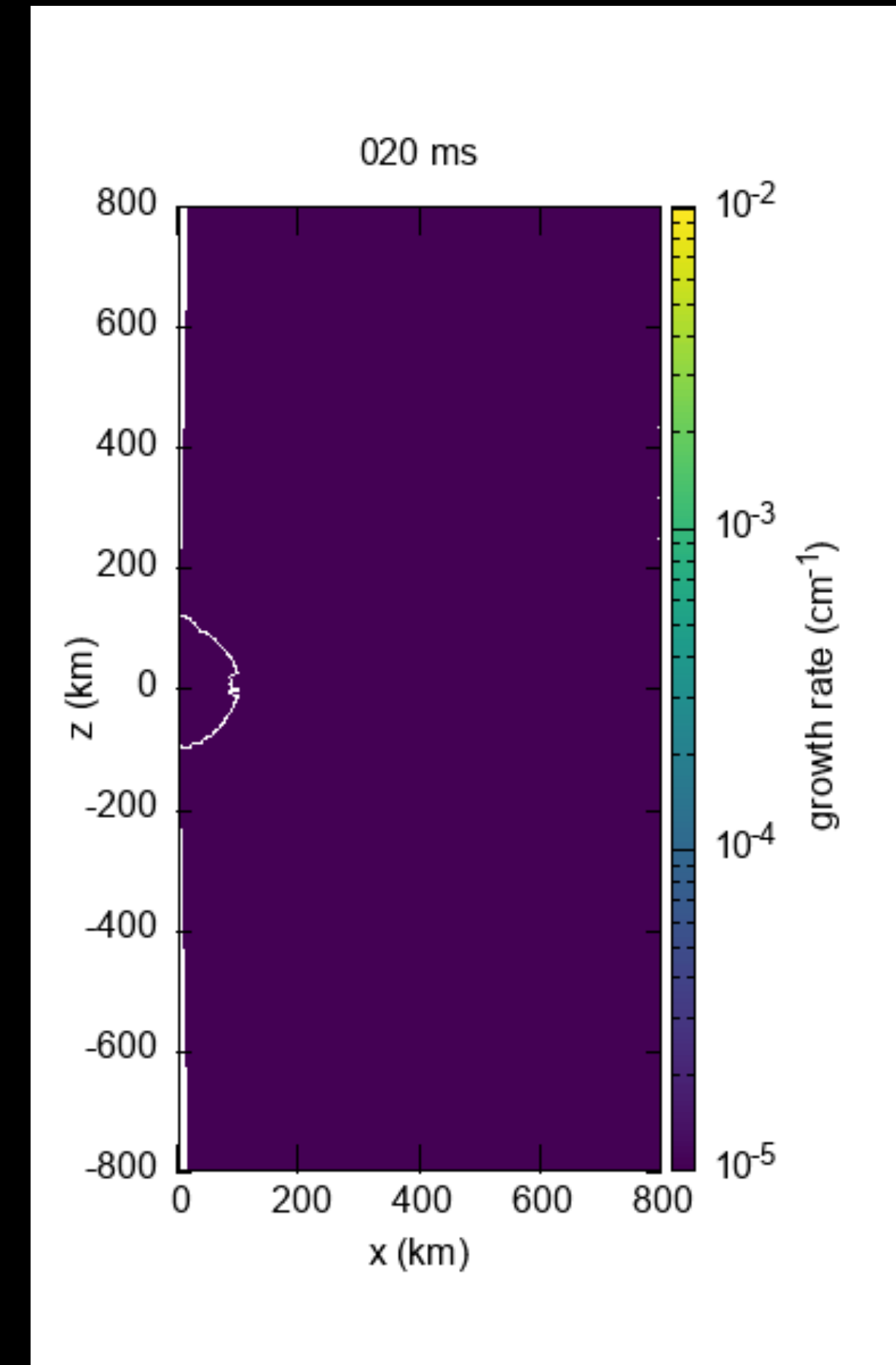
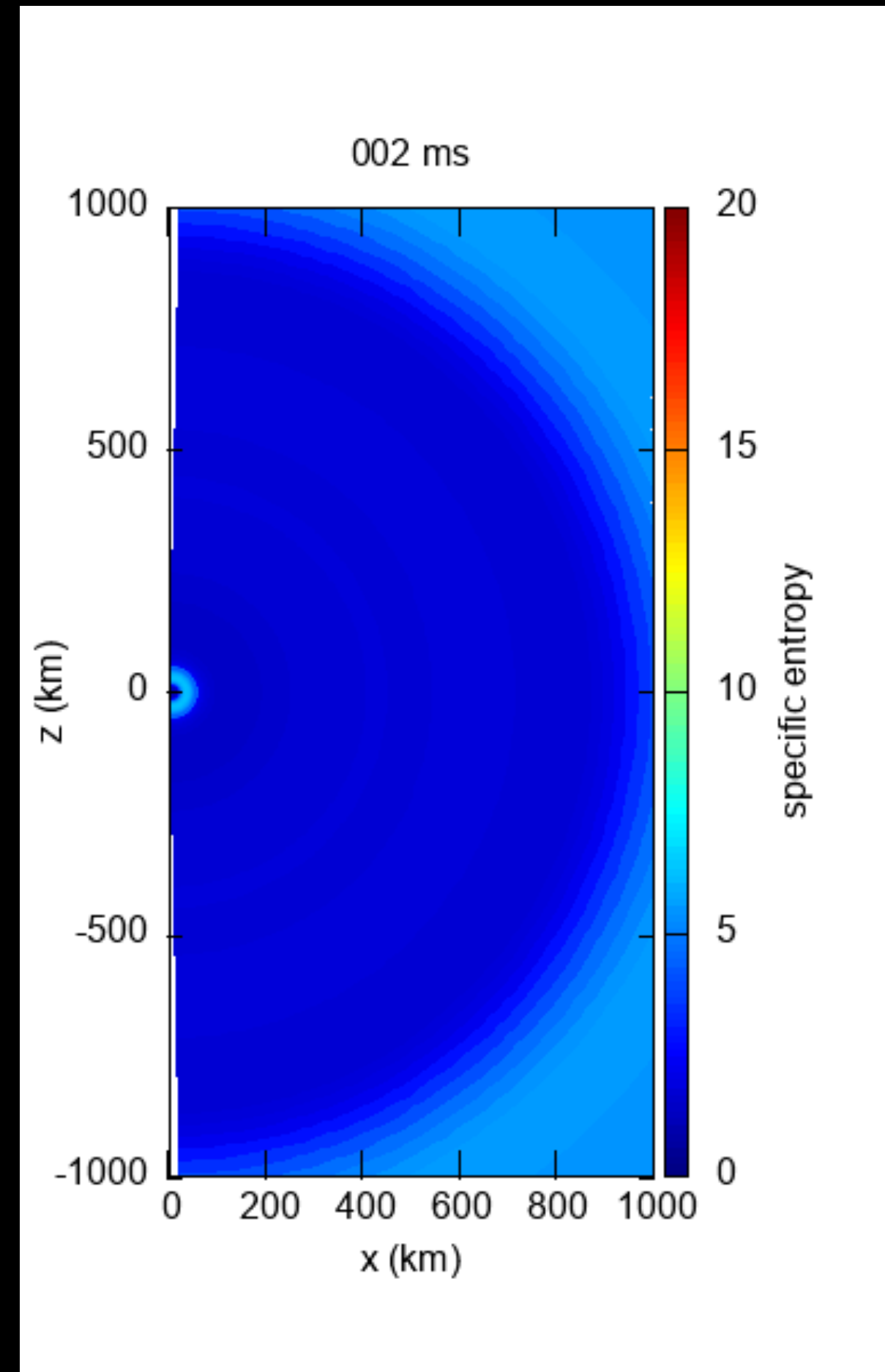
Entropy

FFI growth rate



$M = 11.2M_{\odot}$ progenitor

EOS: LS220



CFI

$$\omega_{\pm} = -A - i\gamma \pm \sqrt{A^2 - \alpha^2 + 2iG\alpha},$$

$$G \equiv \frac{\mathbf{g} + \bar{\mathbf{g}}}{2} \quad A \equiv \frac{\mathbf{g} - \bar{\mathbf{g}}}{2}$$

$$\gamma \equiv \frac{\Gamma + \bar{\Gamma}}{2} \quad \alpha \equiv \frac{\Gamma - \bar{\Gamma}}{2}$$

$$\mathbf{g} \equiv n_{\nu_e} - n_{\bar{\nu}_e} \quad \bar{\mathbf{g}} \equiv n_{\bar{\nu}_e} - n_{\nu_e}$$

$$\Gamma \equiv \frac{\Gamma_e + \Gamma_x}{2} \quad \bar{\Gamma} \equiv \frac{\bar{\Gamma}_e + \bar{\Gamma}_x}{2}$$

(Liu 2023)

FFI

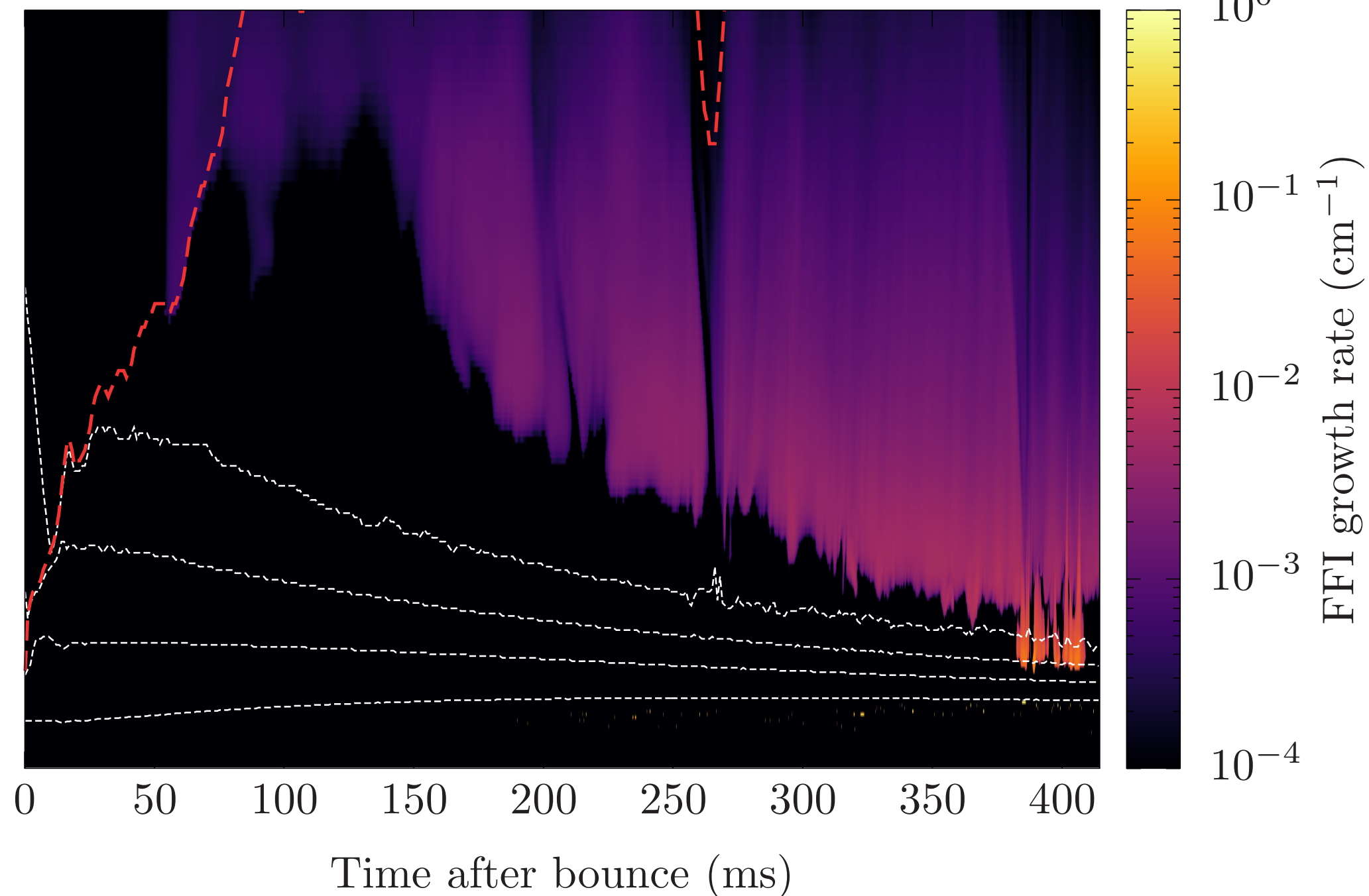
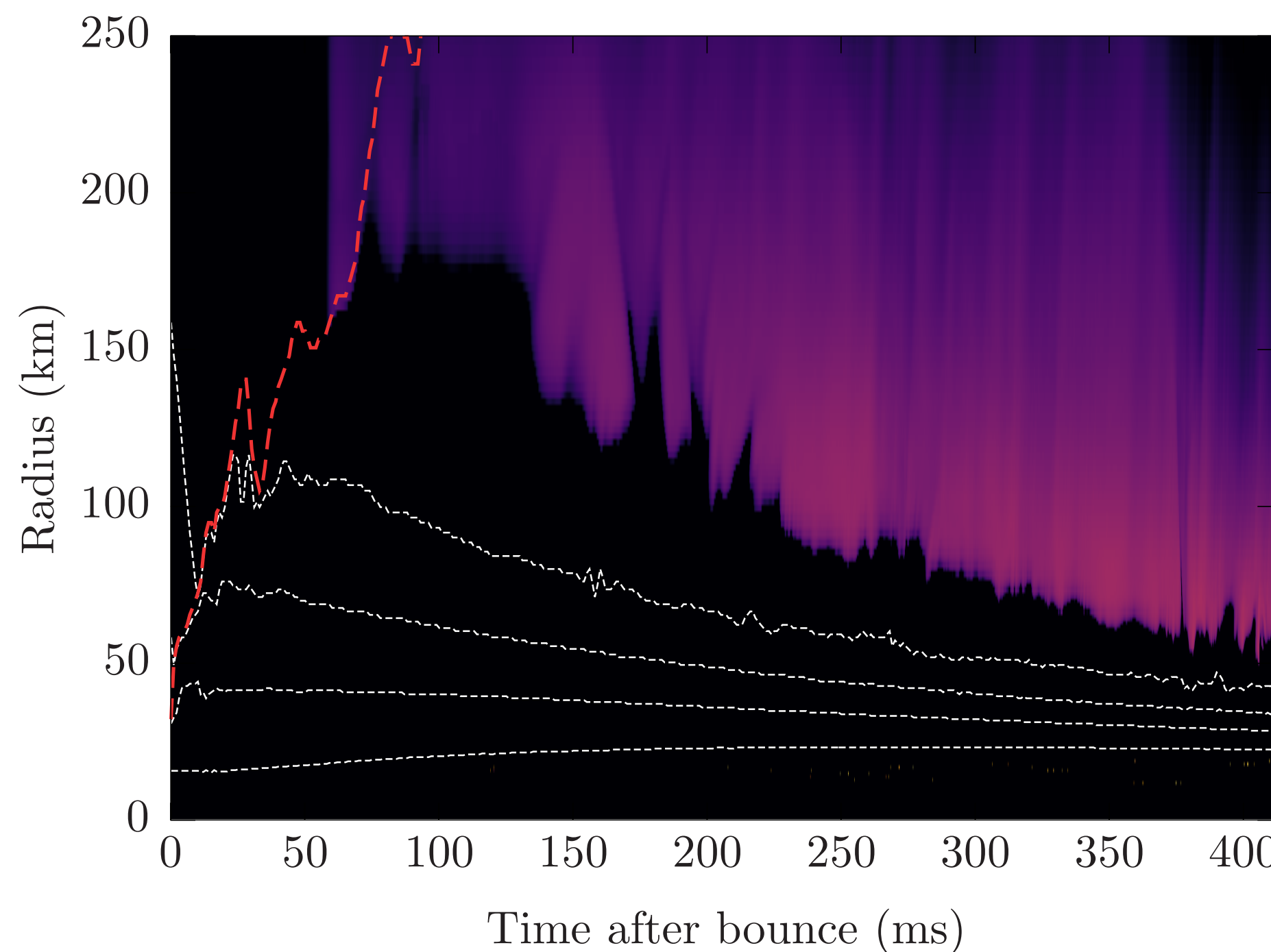
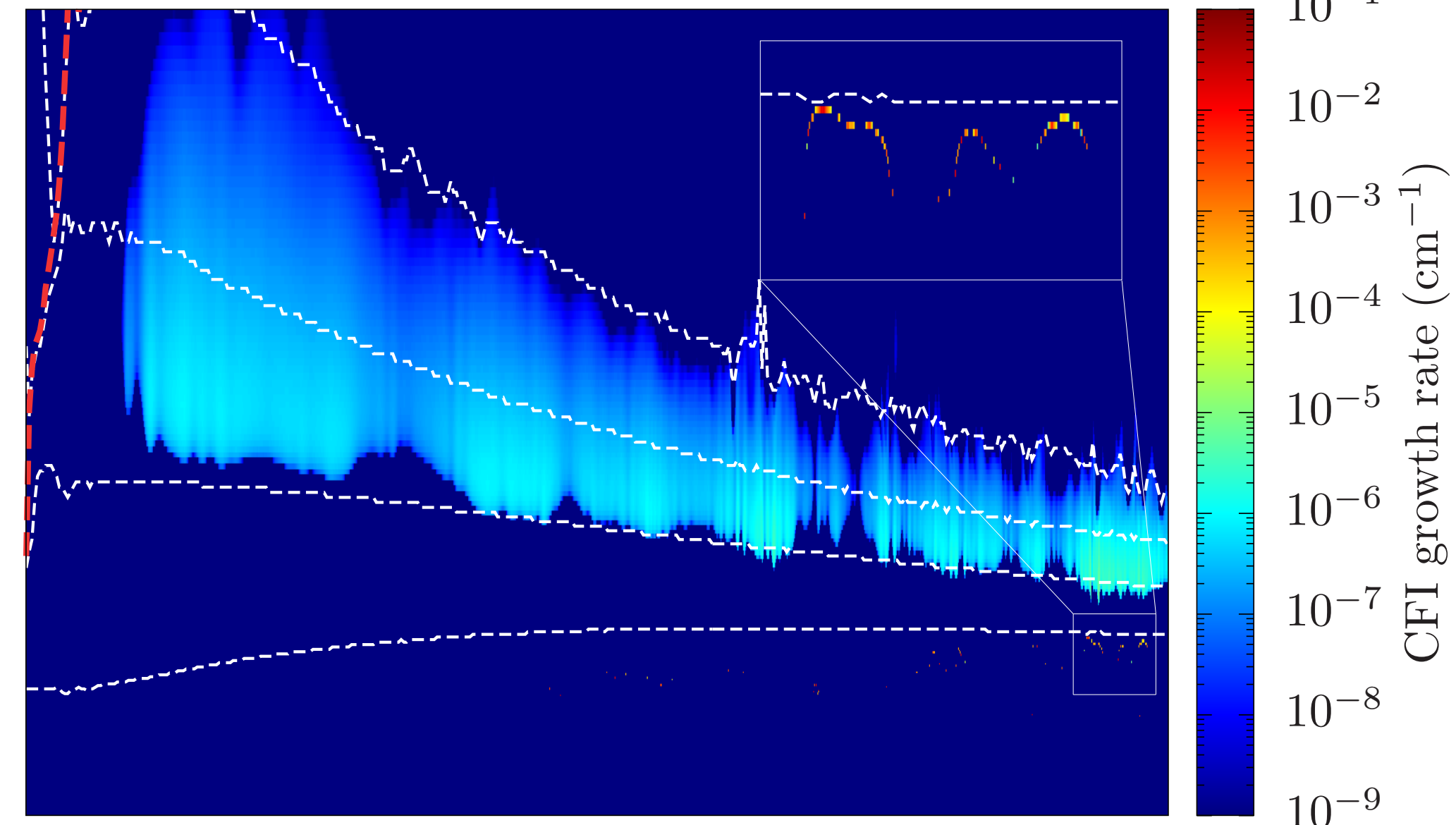
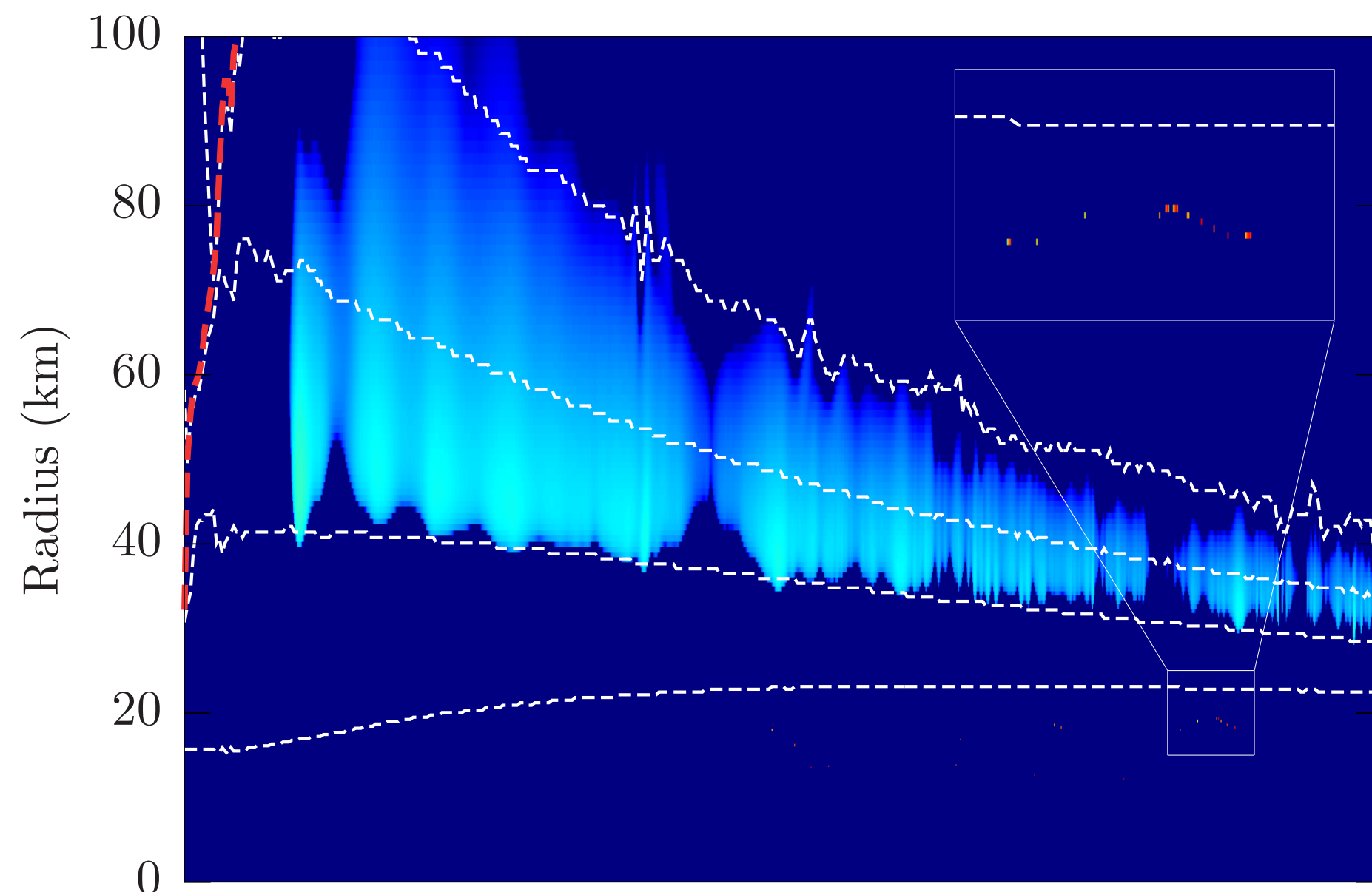
$$\sigma = \sqrt{-\left(\int_{\Delta G > 0} \frac{d\Omega}{4\pi} \Delta G\right) \left(\int_{\Delta G < 0} \frac{d\Omega}{4\pi} \Delta G\right)}$$

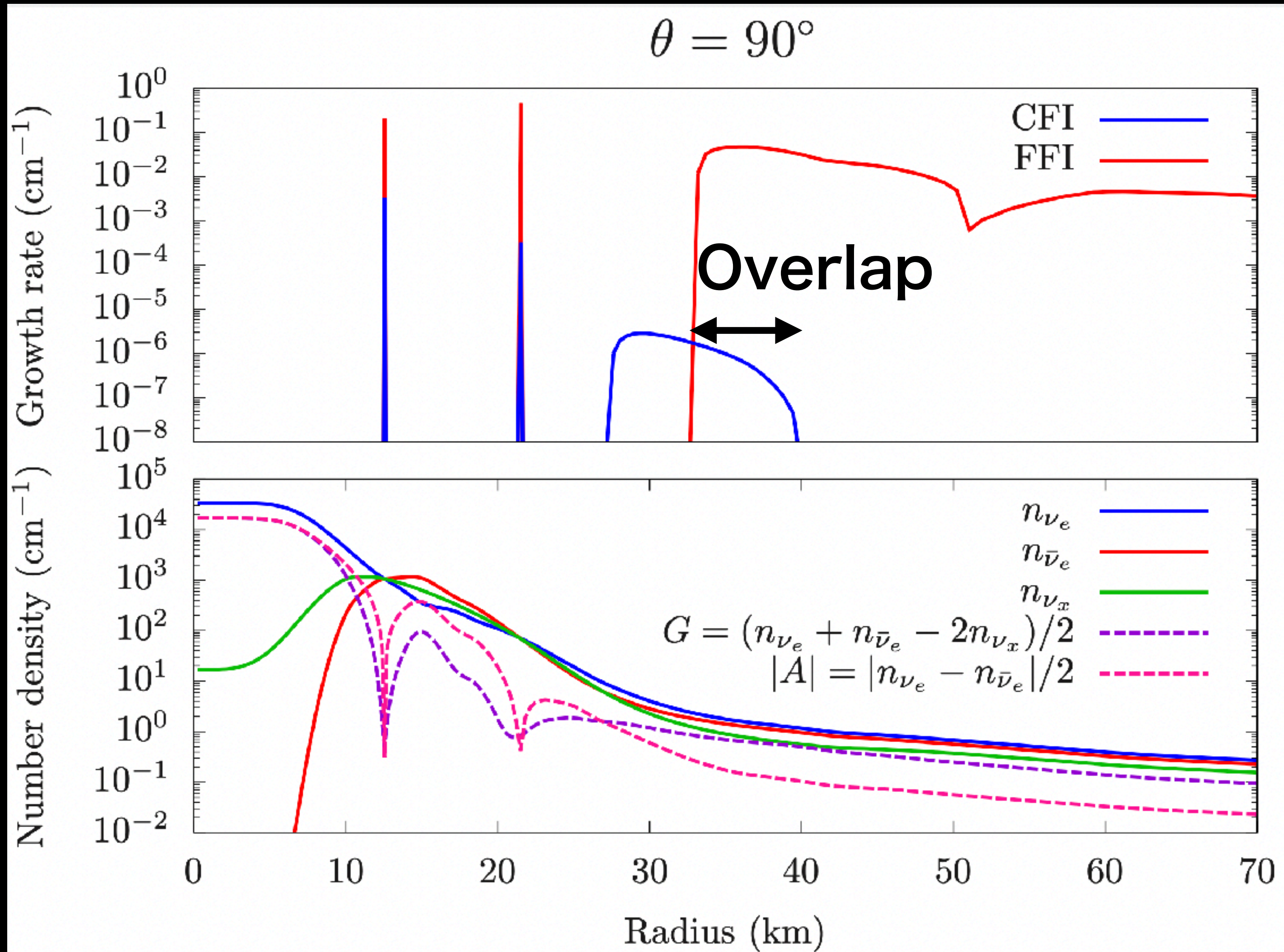
$$\Delta G = \frac{\sqrt{2}G_F}{2\pi^2} \int (f_{\nu_e} - f_{\bar{\nu}_e}) \nu^2 d\nu$$

(Morinaga 2020)

$\theta = 45^\circ$

$\theta = 90^\circ$





Growth rate (Liu 2023)

$$\sigma = \text{Im}\omega$$

$$\omega_{\pm} = -A - i\gamma \pm \sqrt{A^2 - \alpha^2 + 2iG\alpha},$$

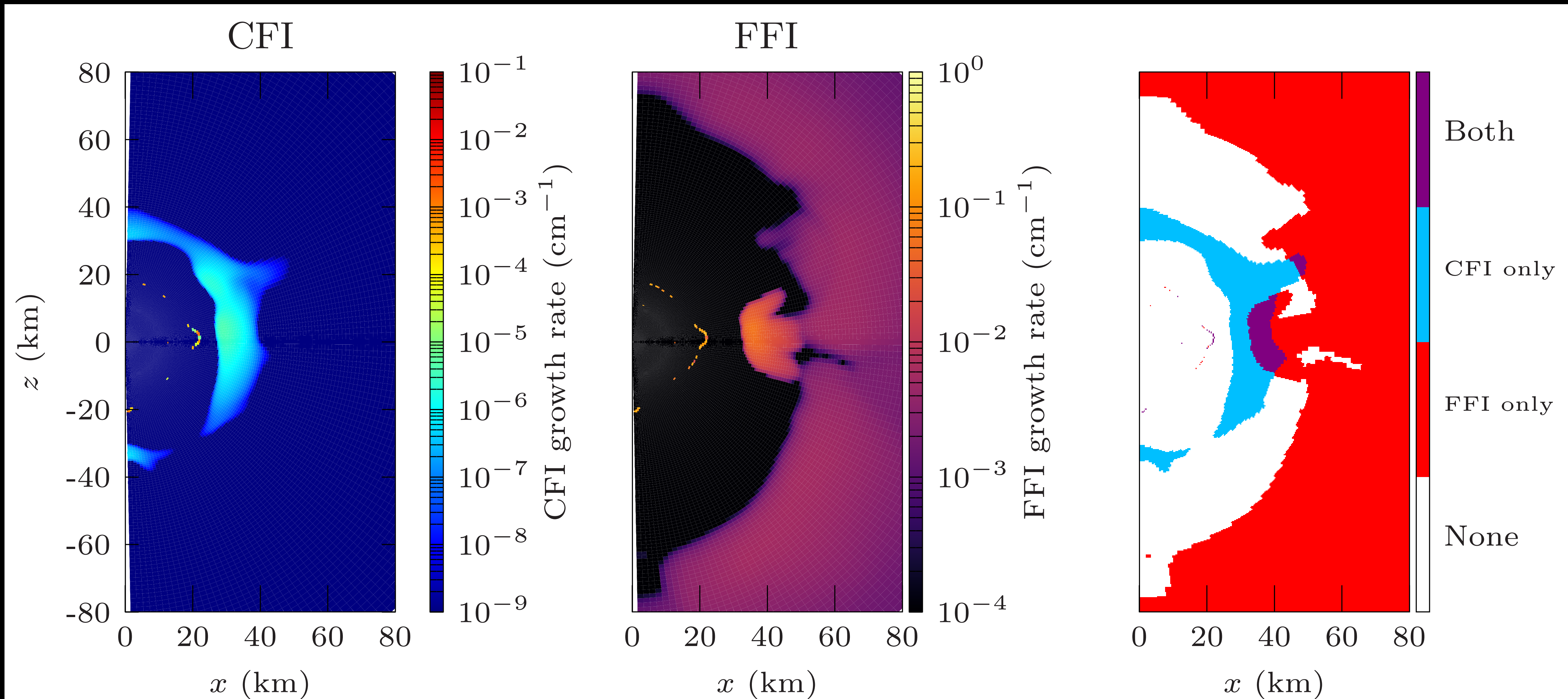
$$G \equiv \frac{\mathbf{g} + \bar{\mathbf{g}}}{2} \quad A \equiv \frac{\mathbf{g} - \bar{\mathbf{g}}}{2}$$

$$\gamma \equiv \frac{\Gamma + \bar{\Gamma}}{2} \quad \alpha \equiv \frac{\Gamma - \bar{\Gamma}}{2}$$

$$\mathbf{g} \equiv n_{\nu_e} - n_{\nu_x} \quad \bar{\mathbf{g}} \equiv n_{\bar{\nu}_e} - n_{\bar{\nu}_x}$$

$$\Gamma \equiv \frac{\Gamma_e + \Gamma_x}{2} \quad \bar{\Gamma} \equiv \frac{\bar{\Gamma}_e + \bar{\Gamma}_x}{2}$$

$$\Gamma_i \equiv \sqrt{2}G_F \int \frac{E^2 dE}{2\pi^2} \Gamma(E) f(E)$$



Summary

- We analyzed the occurrence of fast and collisional flavor instabilities in 2D CCSN model performed with Boltzmann neutrino transport code.
- We both found region with FFI and CFI, and overlap region is also observed.

Future Prospects

- Our study is a post-process analysis, hence the neutrino distribution would be different if oscillation is taken into account in the time evolution.
- We will run CCSN simulation with fast and collisional flavor conversions.