

Bayesian inference of the dense matter equation of state built within mean field models

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Overview

- Neutron Star Equation of State & Bayesian Inferences
- Setup
 - covariant density functional vs. non-relativistic mean field
 - constraints
- Results & Conclusions



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Equation of State (EOS)

Neutron Star (NS) properties depend on a one parameter EOS, $P(e)$

- any astrophys. measurement, e.g., radius, tidal deformability, moment of inertia, oscillation modes frequencies, maximum mass probes the EOS
- inference of NS EOS is difficult
- current observational measurements: $0.8 \lesssim M/M_{\odot} \lesssim 2.1$; combined tidal deformability of NS in GW170817; joint mass and radius measurements
- depending on NS mass and the *unknown* EOS, different n_B and Y_p domains are explored
- global parameters depend in a convoluted way on the EOS
- information from nuclear physics is limited $n_B \approx n_{\text{sat}}$, $Y_p \lesssim 0.5$
- over limited n_B domains, different particle compositions may provide similar EOS

[Doroshenko+, Nature (2022)] [Antoniadis+, Science (2013)] [Arzoumanian+, ApJSS (2018)] [Cromartie+, Nature (2020)] [Fonseca+, ApJL (2021)] [Abbott+, PRL (2017)] [Abbott+, PRX (2019)] [Miller+, ApJL (2019)] [Riley+, ApJL (2019)] [Miller+, ApJL (2021)] [Riley+, ApJL (2021)] [Vinciguerra+, arXiv:2308.09469]

NS EOS - the "traditional" approach

- ▷ Schematic parametrizations: piecewise polytropes; param. of the speed of sound
 - ✓ computation cost, flexibility
 - ✗ no composition info, no physical pinning, no composition
- ▷ From an (effective) interaction
 - ▷ phenomenological: non-relativistic mean field (Skyrme, Gogny), covariant density functionals,
 - ▷ ab initio: variational, quantum Monte Carlo, coupled cluster expansion, diagrammatic, Brueckner Hartree-Fock, lattice, χ EFT models,
 - ✓ physical underpinning, inclusion of heavy baryons and mesons
 - ? validity away from initial conditions; compliance with causality and astrophys. measurements

[Dutra+, PhysRevC (2012); PhysRevC (2014)] [Oertel+, RMP (2017)] [Burgio+, ProgrPartNuclPhys (2021)]

Large model spaces and inhomogeneous sets of constraints require **statistical tools**

Bayes theorem: $P(\Theta|D) = \mathcal{L}(D|\Theta)P(\Theta)/\mathcal{I}$,

Θ : the set of model params, i.e., model (schematic, phenomenological) + priors on parameters

D : the fit data, i.e., constraints (NM param., E/A or P in neutron matter as computed by ab initio, astro. meas.)

\mathcal{I} = the evidence

Results: Posterior PDF for Θ , $\xi_i(\Theta)$

Conclusions:

- dependence on the EOS model;
- sensitivity of posterior distrib. to prior distrib.;
- narrowing down of the parameter space upon progressive incorporation of constr.;
- tension between constraints

[Lim & Holt, EPJA (2019)] [Raaijmakers+ (2019)] [Güven+, PhysRevC (2020)] [Raaijmakers+, ApJ (2021)]

[Miller+, ApJL (2021)] [Malik+, ApJ (2022)] [Malik+, PhysRevC (2023)] [Beznogov & Raduta, PhysRevC (2023)]

Derive and confront minimally constrained EOS built within two mean field models

Models:

- M1. non-relativistic mean-field model of nuclear matter with Skyrme effective interactions,
- M2. covariant density functional model with (simplified) density dependent couplings,

Minimal set of constraints:

- C1. four best known nuclear empirical parameters (NEP): n_{sat} , E_{sat} , K_{sat} , J_{sym} ,
- C2. $(E/A)(n)$ in pure neutron matter (PNM), as computed by χ EFT,
- C3. $M_{\text{max}}/M_{\odot} \geq 2$,
- C4. for non-relativistic mean-field model: causality for $n \leq n_{\text{c}}^*$

Results: posterior PDF of eff. int. parameters, NEP, global NS properties

- R1. Model dependence,
- R2. The role of correlations among E/A in PNM at various densities

[Beznogov & Raduta, PhysRevC (2023)] [Beznogov & Raduta, arXiv:2308.1535]

Non-relativistic mean field model with Skyrme eff. int.

Effective interaction:

$$V(r_1, r_2) = t_0 (1 + x_0 P_\sigma) \delta(r) + \frac{t_1}{2} (1 + x_1 P_\sigma) [k'^2 \delta(r) + \delta(r) k^2] \\ + t_2 (1 + x_2 P_\sigma) k' \cdot \delta(r) k + \frac{t_3}{6} (1 + x_3 P_\sigma) [n(R)]^\sigma \delta(r),$$

Energy density: $\mathcal{H} = k + h_0 + h_3 + h_{\text{eff}}$, with $k = \frac{\hbar^2}{2m} \tau$, $h_0 = C_0 n^2 + D_0 n_3^2$,

$$h_3 = C_3 n^{\sigma+2} + D_3 n^\sigma n_3^2, \quad h_{\text{eff}} = C_{\text{eff}} n \tau + D_{\text{eff}} n_3 \tau_3,$$

Analytic expressions for all thermo quantities at $T = 0$, including the pressure and

NEPs: $X_{\text{sat}}^{(i)} = \left(\partial^i E_0(n_B, 0) / \partial \chi^i \right) \Big|_{n=n_{\text{sat}}}$ with $\chi = (n_B - n_{\text{sat}}) / 3n_{\text{sat}}$

$$X_{\text{sym}; k}^{(j)} = \left(\partial^j E_{\text{sym}; k}(n_B, 0) / \partial \chi^j \right) \Big|_{n=n_{\text{sat}}}$$

7 parameters: σ , C_0 , D_0 , C_3 , D_3 , C_{eff} , D_{eff} , or, alternatively,
 n_{sat} , E_{sat} , J_{sym} , D_3 , C_{eff} , D_{eff} , σ .

Widely used of nuclei, nuclear matter, NS

Covariant density functional with density dep. couplings

The Lagrangian:

$$\begin{aligned}\mathcal{L} = & \Psi \left[\gamma^\mu \left(i \partial_\mu - \Gamma_\omega A_\mu^{(\omega)} - \Gamma_\rho \boldsymbol{\tau} \cdot \mathbf{A}_\mu^{(\rho)} \right) - (m - \Gamma_\sigma \phi) \right] \Psi + \frac{1}{2} \left\{ \partial_\mu \phi \partial^\mu \phi - m_\sigma^2 \phi^2 \right\} \\ & - \frac{1}{4} F_{\mu\nu}^{(\omega)} F^{(\omega)\mu\nu} + \frac{1}{2} m_\omega^2 A_\mu^{(\omega)} A^{(\omega)\mu} - \frac{1}{4} \mathbf{F}_{\mu\nu}^{(\rho)} \cdot \mathbf{F}^{(\rho)\mu\nu} + \frac{1}{2} m_\rho^2 \mathbf{A}_\mu^{(\rho)} \cdot \mathbf{A}^{(\rho)\mu},\end{aligned}$$

Interactions via exchange of σ , ω , ρ mesons.

Density-dependent couplings: $\Gamma_M(n) = \Gamma_{M,0} h_M(x)$, $x = n/n_{\text{sat}}$, with

$$h_M(x) = \exp[-(x^{a_M} - 1)], \quad M = \sigma, \omega; \quad h_\rho(x) = \exp[-a_\rho(x - 1)]$$

Energy density:

$$e = \frac{1}{\pi^2} \sum_{B=n,p} e_{kin;B} + \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{1}{2} m_\omega^2 \omega^2 + \frac{1}{2} m_\rho^2 (\rho_3)^2,$$

6 parameters: Γ_σ , Γ_ω , Γ_ρ , a_σ , a_ω , a_ρ

Widely used of nuclei, nuclear matter, NS

See M. Beznogov's talk

Constraints

Quantity	Units	Skyrme			CDF		
		Value	Std. deviation	Ref.	Value	Std. deviation	Ref.
n_{sat}	fm^{-3}	0.16	0.004	[1]	0.153	0.005	[4]
E_{sat}	MeV	-15.9	0.2	[1]	-16.1	0.2	[5]
K_{sat}	MeV	240	30	[1]	230	40	[6,7]
J_{sym}	MeV	30.8	1.6	[1]	32.5	1.8	[8]
$(E/A)_1$	MeV	9.212	0.226	[2]	9.50	0.52	[9]
$(E/A)_2$	MeV	12.356	0.512	[2]	12.68	1.20	[9]
$(E/A)_3$	MeV	15.877	0.872	[2]	16.31	2.13	[9]
M_G^*	M_\odot	> 2.0	—	[3]	> 2.0	—	[3]
c_s^{*2}	c^2	< 1	—	—	—	—	—

NM parameters; astro. contr. and causality.
density behavior of PNM as predicted by χ EFT;

1, 2, 3: $n_B = 0.08, 0.12, 0.16 \text{ fm}^{-3}$

[1] Margueron+, PhysRevC (2018); [2] Somasundaram+, PhysRevC (2021); [3] Fonseca+, ApJL (2021); [4] Typel & Wolter, NuclPhysA (1999); [5] Dutra+, PhysRevC (2014); [6] Todd-Rutel & Piekarewicz, PhysRevLett (2005); [7] Shlomo+, EPJA (2006); [8] Essik+, PhysRevC (2021); [9] Hebeler+, ApJ (2013);

Priors: uniform (uninformative) distributions

Skyrme: $n_{\text{sat}}, E_{\text{sat}}, J_{\text{sym}}, D_3, C_{\text{eff}}, D_{\text{eff}}, \sigma$ ← mixed eff. int. and NM parameters
CDF: $\Gamma_\sigma, \Gamma_\omega, \Gamma_\rho, a_\sigma, a_\omega, a_\rho$ ← eff. int. parameters

For domains, see [Beznogov & Raduta, PhysRevC (2023); arXiv:2308.1535],
M. Beznogov's talk

$$\log \mathcal{L}_q \propto -\chi_q^2 = -\sum_{n=1}^N \chi_n^2 - \sum_{m=1}^M \chi_m'^2 - \sum_{p=1}^P \chi_p''^2,$$

1) Uncorrelated obs., e.g., n_{sat} , E_{sat} , K_{sat} , L_{sym} :

$$-\chi_n^2 = -\frac{1}{2} \left(\frac{d_i - \xi_i(\Theta)}{\mathcal{X}_i} \right)^2$$

2) Correlated obs., e.g., the values that E/A in PNM takes at various n_i ,

$$-\chi_m'^2 = -\chi_m^2 - \frac{1}{2} \sum_{i=1}^3 \sum_{j=1}^3 (\text{cov}^{-1})_{ij} \delta \mathcal{E}_i \delta \mathcal{E}_j, \quad [\text{Somasundaram+}, \text{PhysRevC (2021)}]$$

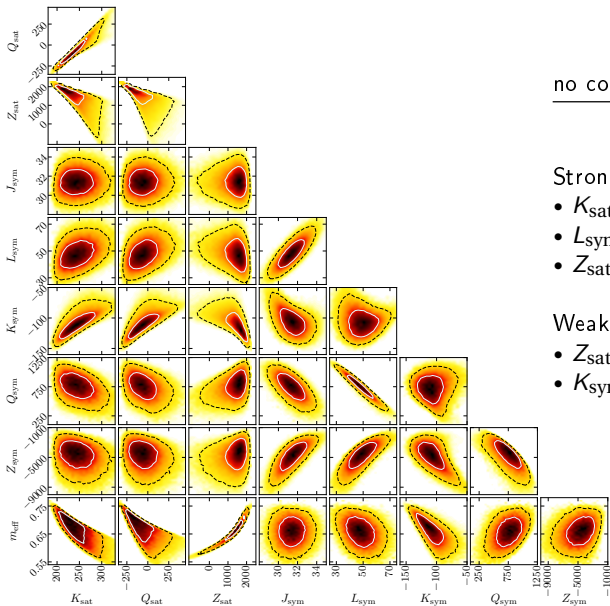
3) "Hard-wall", e.g., M_G^* , c_s^{*2} ,

$$-\chi_p''^2 = -10^{10}, \text{ if } M_G^*/M_\odot < 2 \text{ or } c_s^{*2}/c^2 \geq 1$$

see M. Beznogov's talk

- EOS of nuclear matter & NS
 - posterior PDF for nuclear matter ($n_{\text{sat}}, \chi_{\text{sat}}^{(i)}, \chi_{\text{sym}}^{(j)}, m_{\text{eff}}$)
 - correlations among NM params.
 - posterior PDF for NS properties ($M_{\text{max}}^*, R_{1.4}, R_{2.0}, \Lambda_{1.4}, \Lambda_{2.0}, M_{\text{DU}}$)
 - correlations among params of MN and properties of NS
- model dependence: CDF versus non-rel. mean field model with Skyrme
- constraints dependence: correlations between E/A in PNM

NM parameters: Covariant Density Functional



no correl between $(E/A)_i$ in PNM

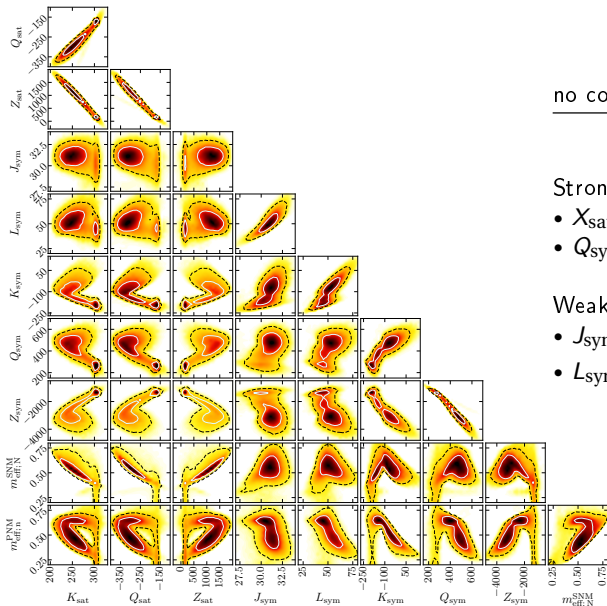
Strong correlations:

- K_{sat} and Q_{sat}
- L_{sym} and Q_{sym}
- Z_{sat} and m_{eff}

Weak correlations:

- Z_{sat} with K_{sym} , Q_{sym}
- K_{sym} and X_{sat} , with $X = K, Q, Z$

NM parameters: Skyrme



no correl between $(E/A)_i$ in PNM

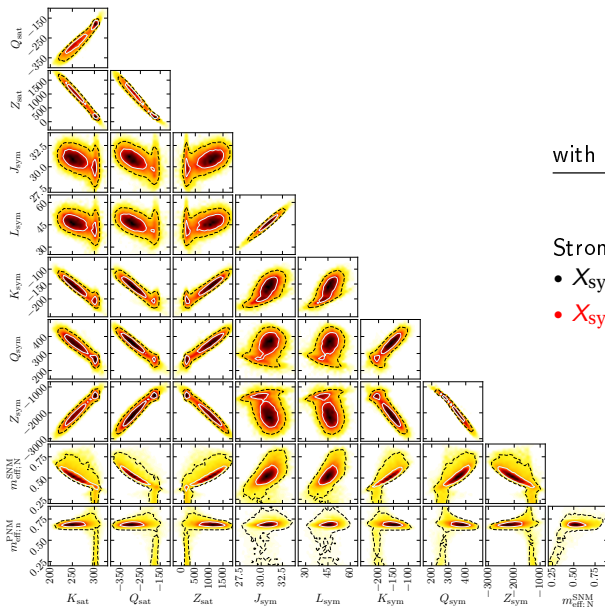
Strong correlations:

- X_{sat} and Y_{sat} , $X, Y = K, Q, Z$
- Q_{sym} and Z_{sym}

Weak correlations:

- J_{sym} and L_{sym}
- L_{sym} and K_{sym} ; K_{sym} and Q_{sym}

NM parameters: Skyrme

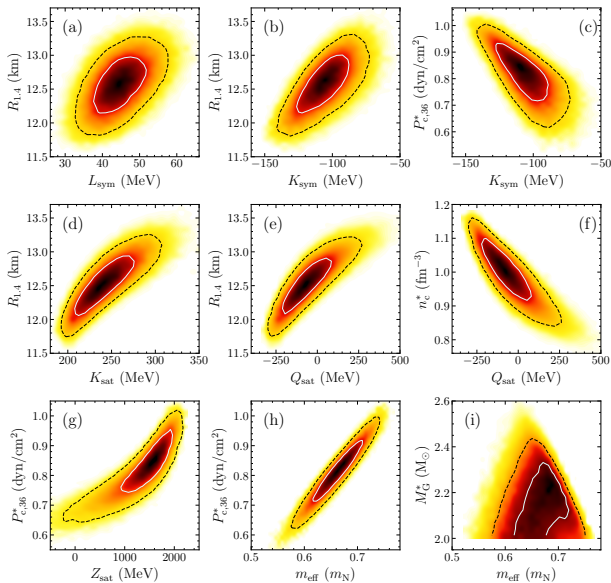


with correl between $(E/A)_i$ in PNM

Strong correlations:

- X_{sym} with Y_{sym} , $X, Y = J, L, K, Q, Z$
- X_{sym} with Y_{sat} , $X, Y = K, Q, Z$

NM parameters vs. NS properties: CDF



- L_{sym} , K_{sym} are weakly correl. with $R_{1.4}$

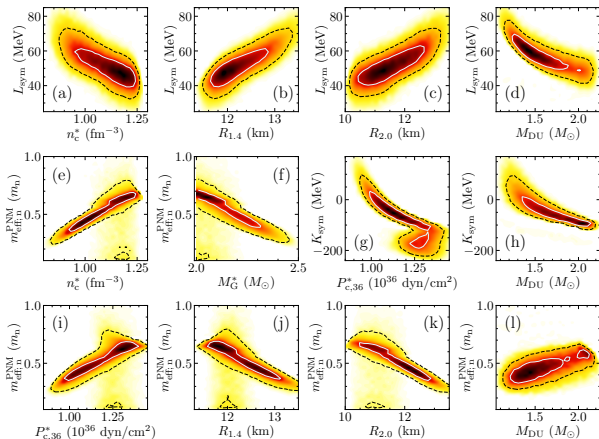
- K_{sat} , Q_{sat} are weakly correl. with $R_{1.4}$

- K_{sym} , Z_{sat} are weakly correl. with P_c^*

- m_{eff} is strongly corr. with P_c^*

no correl (E/A)_i in PNM

NM parameters vs. NS properties: Skyrme

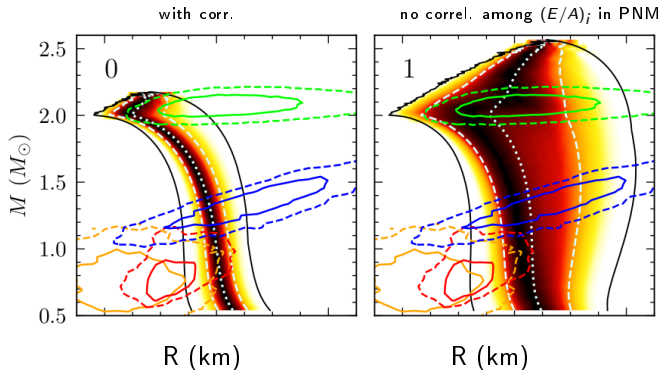


- $L_{\text{sym}}, K_{\text{sym}}$ are correl. with $R_{1.4}, R_{2.0}, M_{DU}$

- L_{sym} is correl. with n_c^*

- m_{eff} is correl. with $n_c^*, P_{c,36}^*, M_G^*, R_{1.4}, R_{2.0}, M_{DU}$

no correl (E/A); in PNM



green: PSR J0740+6620 [Miller+, ApJL (2021)]

blue: PSR J0030+0451 [Miller+, ApJL (2019)]

red, orange: HESS J1731-34 [Doroshenko+, Nature (2022)]

- Bayesian inference of dense matter EOS derived within
 - Non-rel. mean field approach with Skyrme eff. int.
 - Covariant density functional model with density dep. couplings
- minimally constrained
 - well constrained NEP
 - $E/A(n)$ in PNM as predicted by χ EFT
 - $M_{\max}/M_{\odot} \geq 2$, $c_s^2/c^2 \leq 1$
- posterior PDF of NEP and their correlations depend on the model; on the constraints
- NS prop. show correlations with Landau/Dirac effective masses

[Beznogov & Raduta, PhysRevC (2023)] [Beznogov & Raduta, arXiv:2308.1535]