

# Neutrino kinetics in core-collapse supernova

Hiroki Nagakura

(National Astronomical Observatory of Japan)

# A Chronological table: progress of SN (and NS) research

Future

2015-: Multi-dimensional SN models with high-fidelity of input physics  
Successful SN explosions on big iron → Connecting observations

2019-: Diversity  
(SL-SNe, FBOT etc..)

2001-: Establishing 1D-Boltzmann SN models  
(Liebendörfer et al., Sumiyoshi et al.....)

2015: Dawn of GW-astronomy

1990-: Recognizing importance of fluid instabilities  
on SN (Mezacappa, Janka, and Burrows.....)

2010: Discovery of 2 Msun NS

1985-: Bruenn documented “Core” of SN theory

1998: GRB-SN connection

1985: Neutrino-heating explosion  
was proposed by Bethe and Wilson

1987: IMB, Kamiokande-II made the first direct  
detections of SN neutrinos

1966: Colgate and White  
Neutrino emission from  
stellar implosion

1974: Observation of Hulse-Taylor binary

1967-: Discovery of the first radio pulsar (Hewish et al. and Gold)

1938-: Observations of extragalactic supernova and their remnants  
(See e.g., Baade 1938)

1933-: Baade and Zwicky  
Hypothesized Connection between neutron star and “super-nova”

Past

# Neutrino-heating mechanism of core-collapse supernova (CCSN)

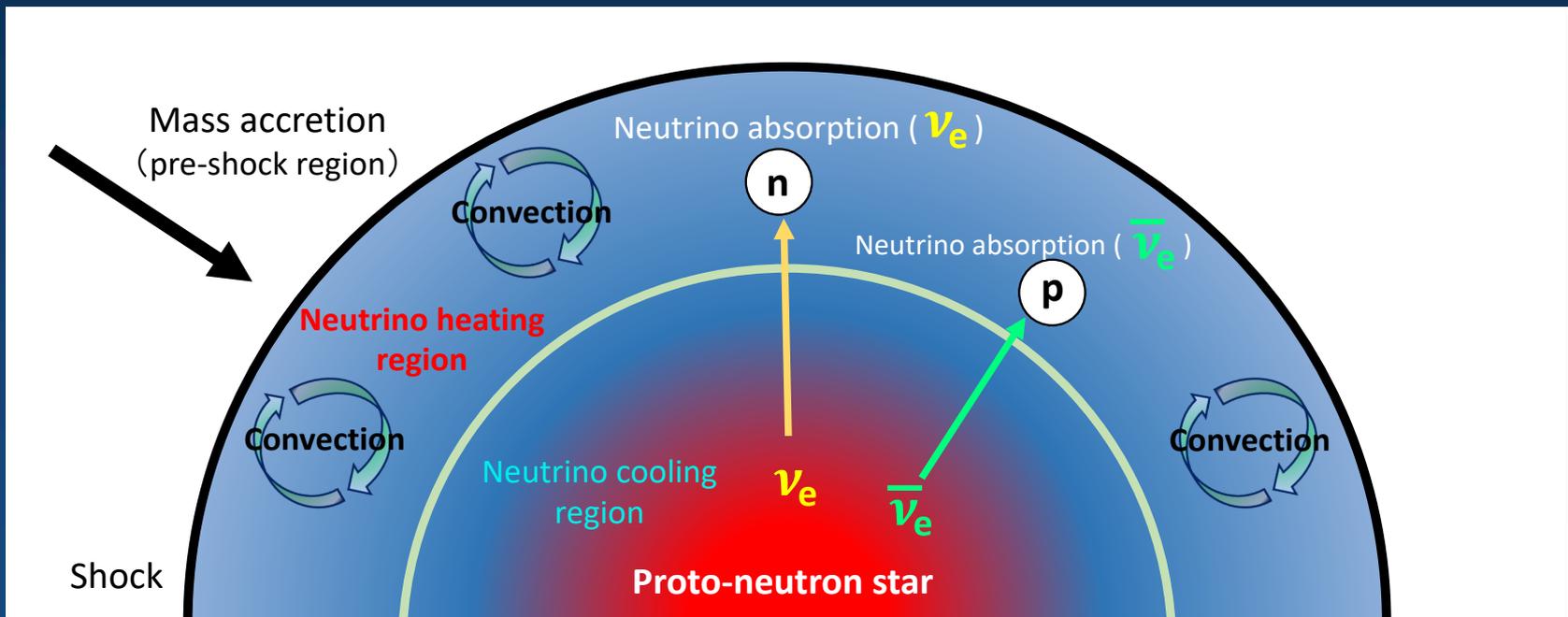
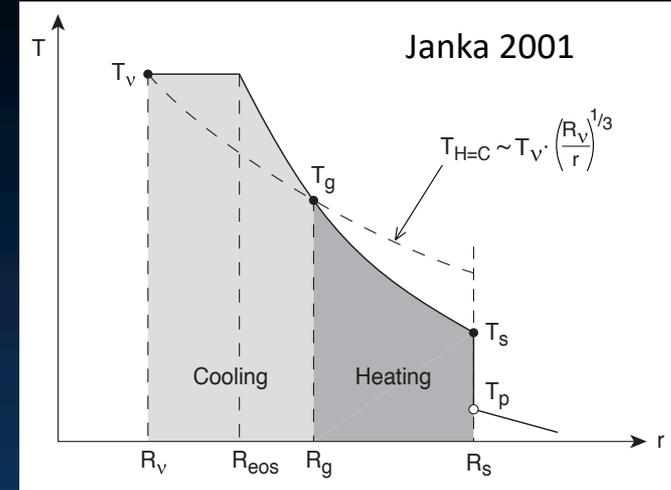
See also talk by Haakon Andresen

✓ Neutrino Heating Rate

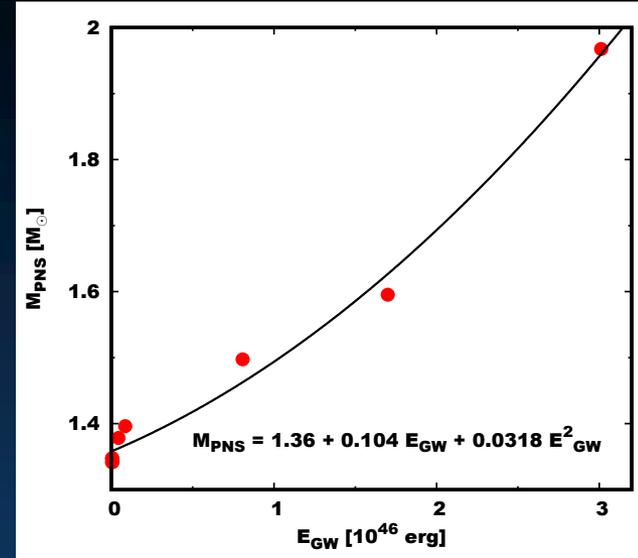
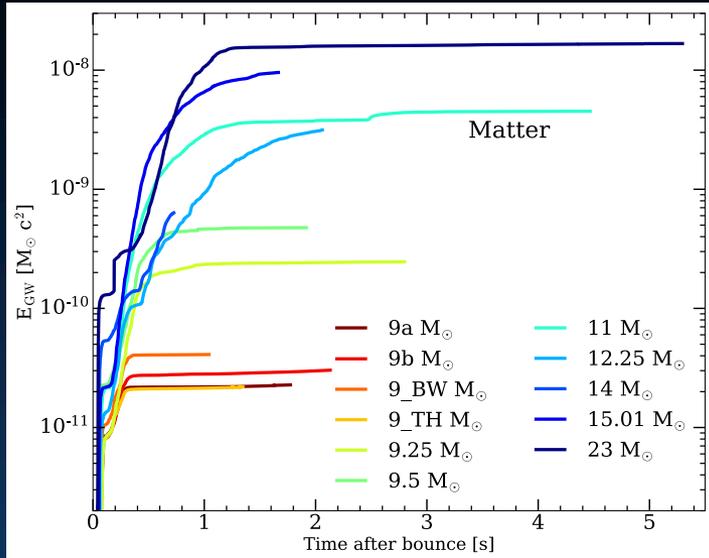
$$Q_{\nu}^{+} \approx 160 \text{Mev/s} \frac{\rho}{m_a} \frac{L_{\nu_e,52}}{r_7^2 \langle \mu_{\nu} \rangle} \left( \frac{T_{\nu_e}}{4 \text{MeV}} \right)^2$$

✓ Neutrino Cooling Rate

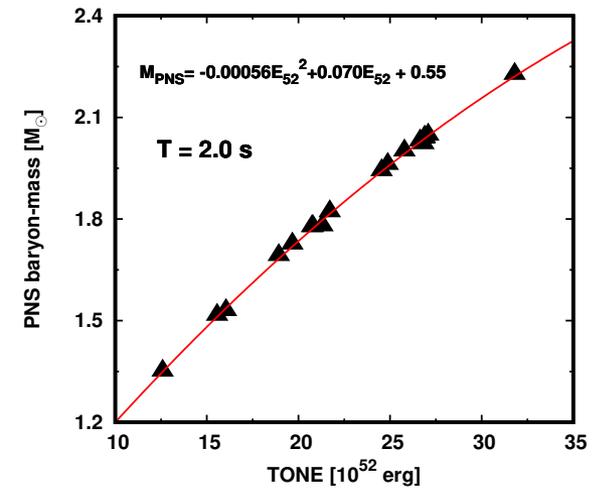
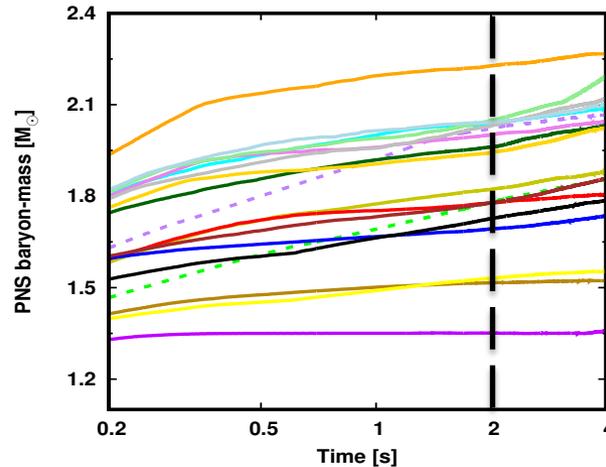
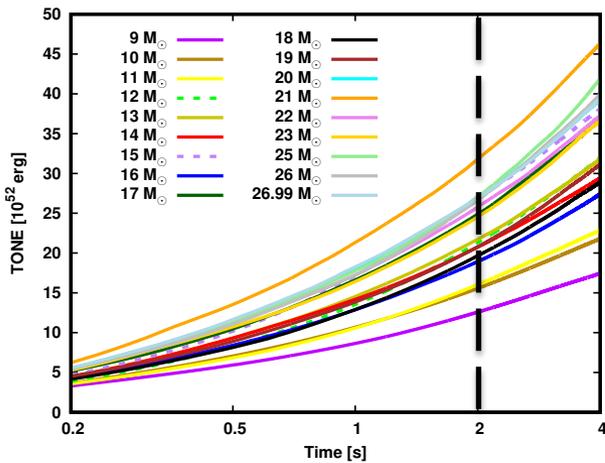
$$Q_{\nu}^{-} \approx 145 \text{Mev/s} \frac{\rho}{m_a} \left( \frac{T}{2 \text{MeV}} \right)^6$$



## Correlation: GWs - PNS mass Nagakura et al. 2020, Vartanyan et al. 2023, Nagakura and Vartanyan 2023

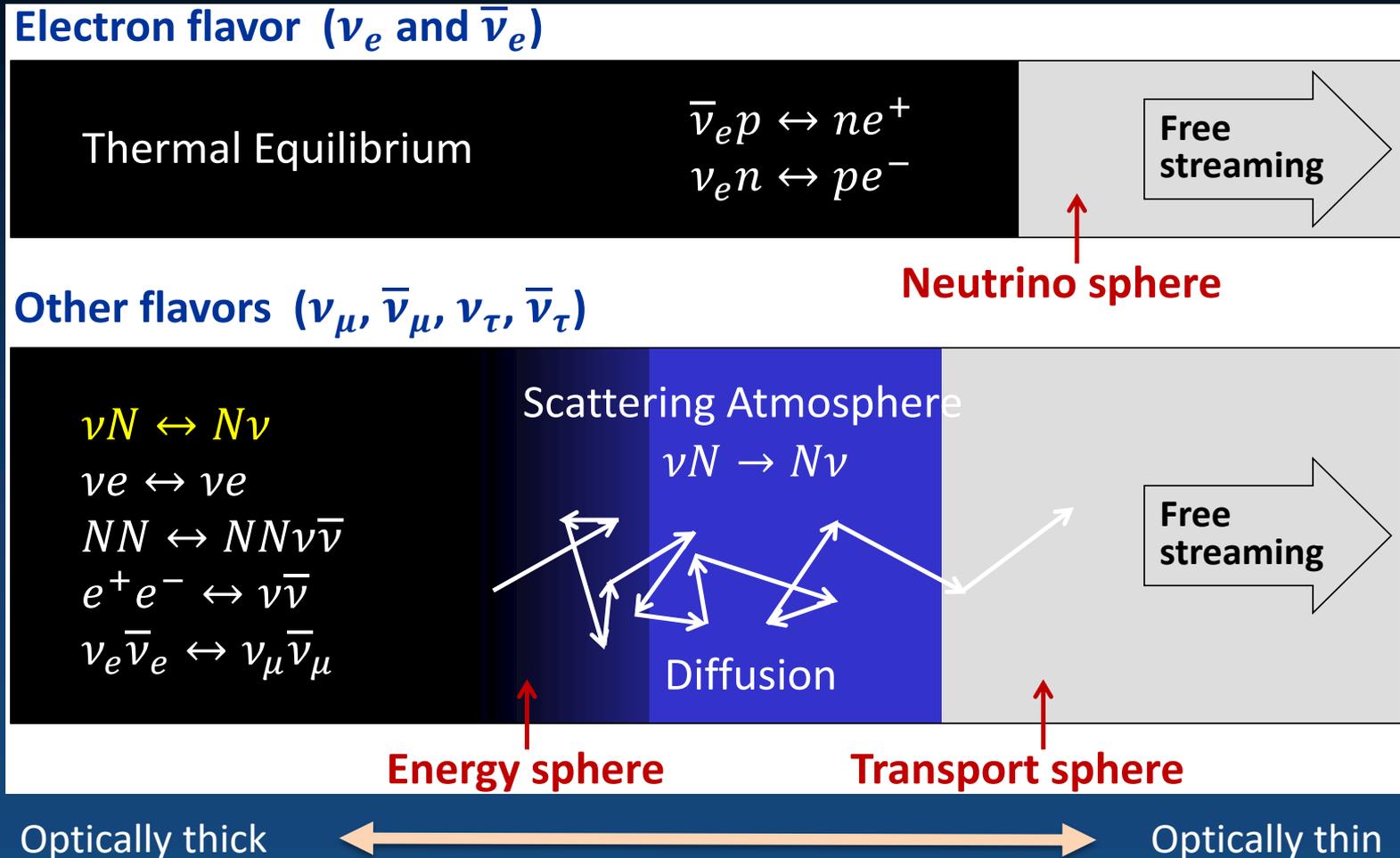


## Correlation: Neutrinos - PNS mass Nagakura and Vartanyan 2021, 2023



# Modeling of neutrino radiation field requires **kinetic theory**

Figure by Janka 2017



# General relativistic full Boltzmann neutrino transport

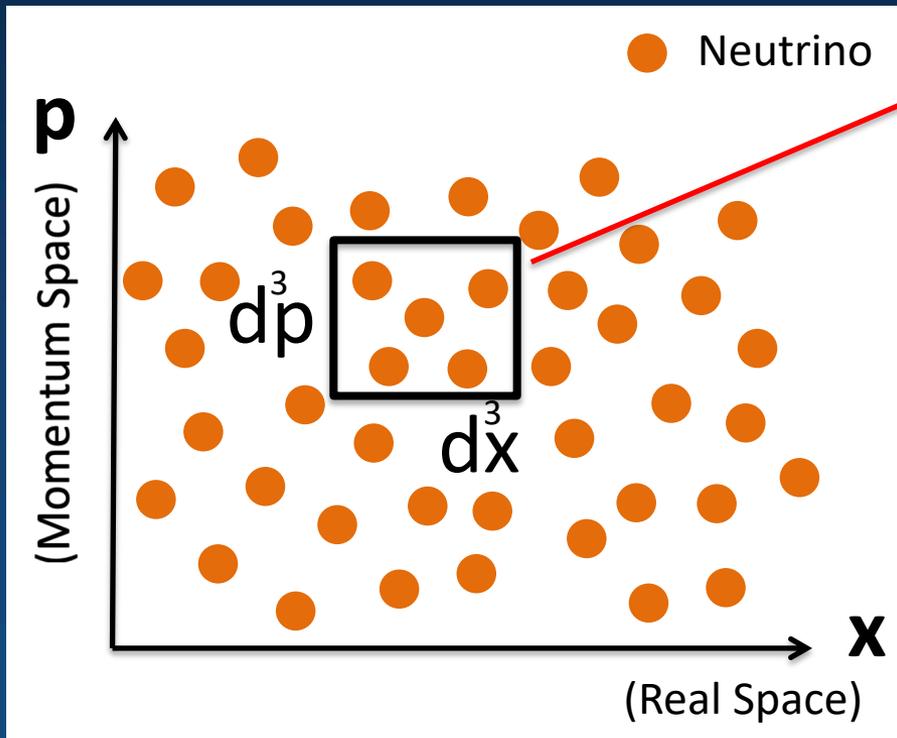
$$p^\mu \frac{\partial f}{\partial x^\mu} + \frac{dp^i}{d\tau} \frac{\partial f}{\partial p^i} = \left( \frac{\delta f}{\delta \tau} \right)_{\text{col}}$$

(Time evolution + Advection Term)

(Collision Term)

See also Lindquist 1966  
 Ehlers 1971

6 dimensional Phase Space



$$dN = f(t, \mathbf{p}, \mathbf{x}) d^3 p d^3 x$$

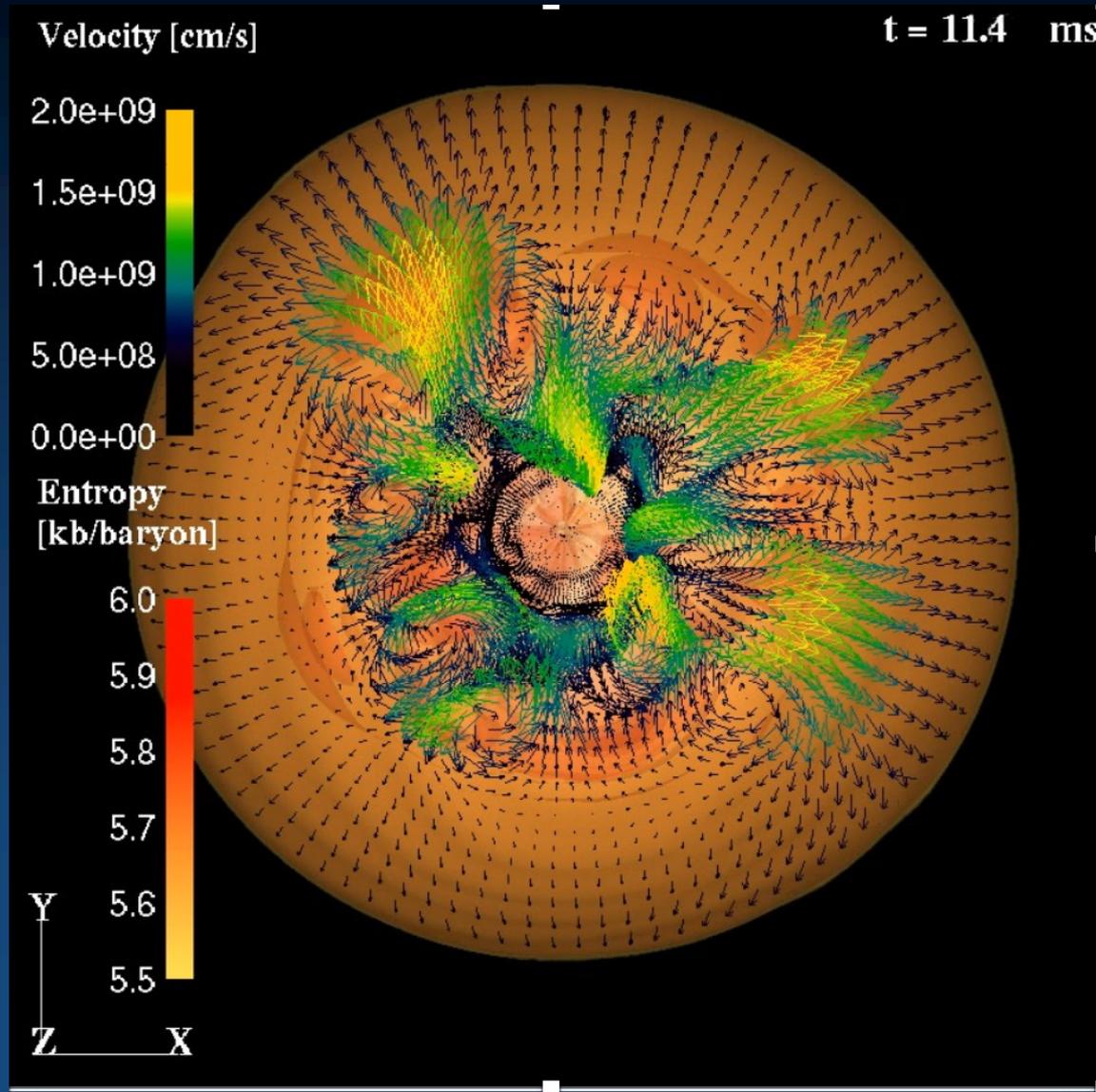
Conservative form of GR Boltzmann eq.

$$\begin{aligned} & \frac{1}{\sqrt{-g}} \frac{\partial(\sqrt{-g} \nu^{-1} p^\alpha f)}{\partial x^\alpha} \Big|_{q(i)} + \frac{1}{\nu^2} \frac{\partial}{\partial \nu} (-\nu f p^\alpha p_\beta \nabla_\alpha e^\beta_{(0)}) \\ & + \frac{1}{\sin \bar{\theta}} \frac{\partial}{\partial \bar{\theta}} \left( \nu^{-2} \sin \bar{\theta} f \sum_{j=1}^3 p^\alpha p_\beta \nabla_\alpha e^\beta_{(j)} \frac{\partial \ell_{(j)}}{\partial \bar{\theta}} \right) \\ & + \frac{1}{\sin^2 \bar{\theta}} \frac{\partial}{\partial \bar{\varphi}} \left( \nu^{-2} f \sum_{j=2}^3 p^\alpha p_\beta \nabla_\alpha e^\beta_{(j)} \frac{\partial \ell_{(j)}}{\partial \bar{\varphi}} \right) = S_{\text{rad}}, \end{aligned}$$

Shibata and Nagakura et al. 2014, Cardall et al. 2013

# - 3D CCSN simulations with full Boltzmann neutrino transport

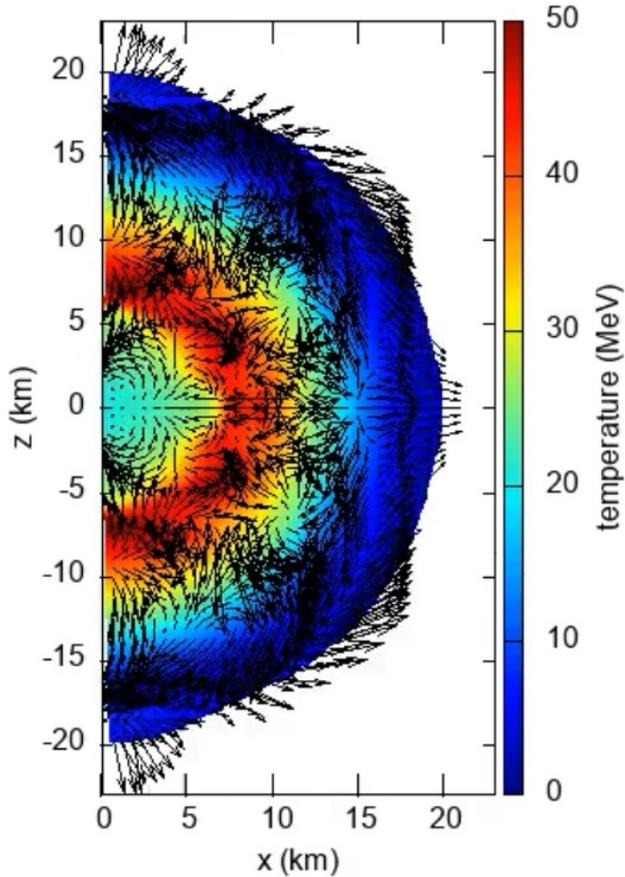
Iwakami, Nagakura et al. 2020, 2021



# ✓ GR simulations with full Boltzmann neutrino transport

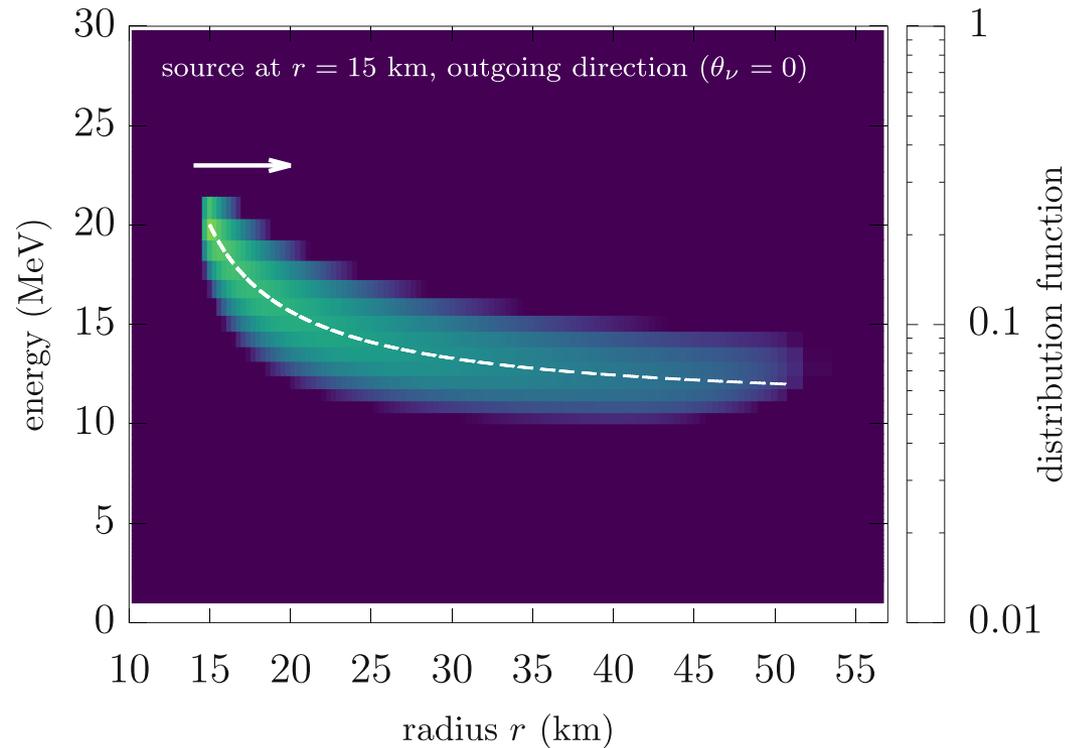
## PNS convection

0.0044 s



Akaho, Nagakura et al. 2023

## Gravitational redshift in Black hole spacetime



Akaho, Nagakura et al. 2020

# General relativistic full Boltzmann neutrino transport

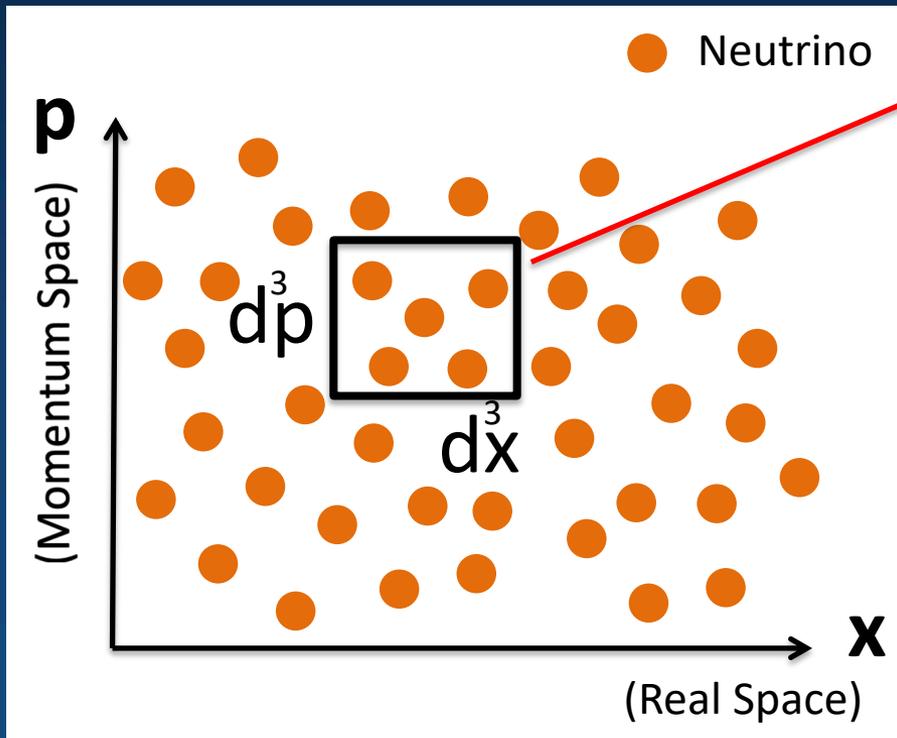
$$p^\mu \frac{\partial f}{\partial x^\mu} + \frac{dp^i}{d\tau} \frac{\partial f}{\partial p^i} = \left( \frac{\delta f}{\delta \tau} \right)_{\text{col}}$$

(Time evolution + Advection Term)

(Collision Term)

See also Lindquist 1966  
 Ehlers 1971

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$$dN = f(t, \mathbf{p}, \mathbf{x}) d^3 p d^3 x$$

Conservative form of GR Boltzmann eq.

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Shibata and Nagakura et al. 2014, Cardall et al. 2013

# Weak Interactions

See talks on Tuesday

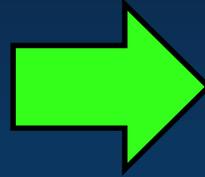
## Basic Sets:

$\nu_e n \rightleftharpoons e^- p$	Bruenn (1985)
$\bar{\nu}_e p \rightleftharpoons e^+ n$	Bruenn (1985)
$\nu_e A' \rightleftharpoons e^- A$	Bruenn (1985)
$\nu N \rightleftharpoons \nu N$	Bruenn (1985)
$\nu A \rightleftharpoons \nu A$	Bruenn (1985), Horowitz (1997)
$\nu e^\pm \rightleftharpoons \nu e^\pm$	Bruenn (1985)
$e^- e^+ \rightleftharpoons \nu \bar{\nu}$	Bruenn (1985)
$NN \rightleftharpoons \nu \bar{\nu} NN$	Hannestad & Raffelt (1998)

Lentz et al. 2011, Kotake et al. 2018

See also Grang et al. 2020, Fisher et al. 2020,  
Sugiura et al. 2022

Extensions



## Lepton Sectors (including muons):

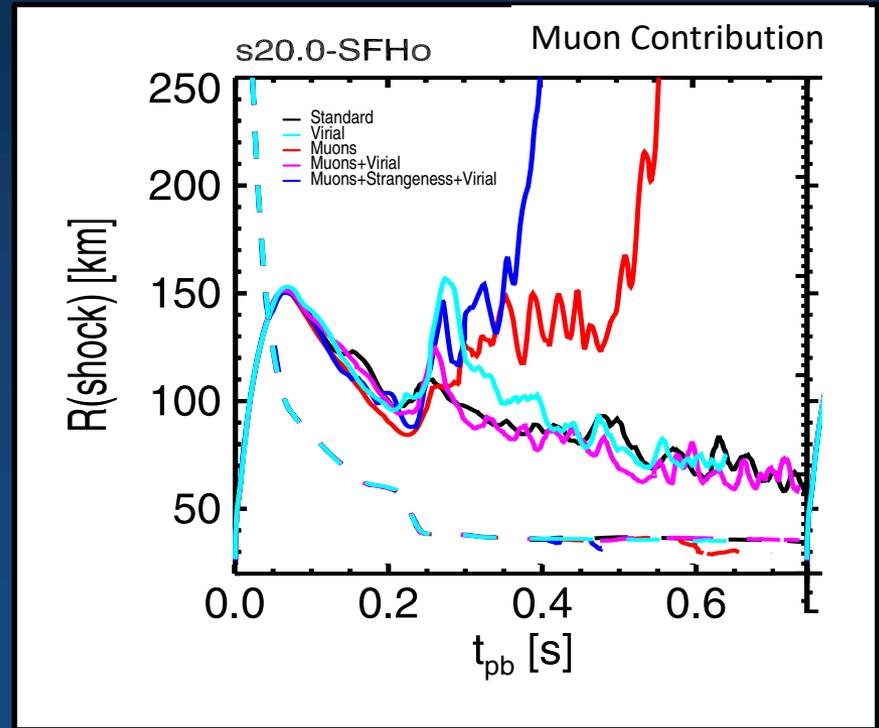
$$\nu_e + \bar{\nu}_e \rightleftharpoons \nu_x + \bar{\nu}_x \quad \text{Buras et al. (2003),}$$

$$\nu_x + \nu_e(\bar{\nu}_e) \rightleftharpoons \nu'_x + \nu'_e(\bar{\nu}'_e) \quad \text{Fischer et al. (2009)}$$

$\nu + \mu^- \rightleftharpoons \nu' + \mu^{-'}$	$\nu + \mu^+ \rightleftharpoons \nu' + \mu^{+'}$
$\nu_\mu + e^- \rightleftharpoons \nu_e + \mu^-$	$\bar{\nu}_\mu + e^+ \rightleftharpoons \bar{\nu}_e + \mu^+$
$\nu_\mu + \bar{\nu}_e + e^- \rightleftharpoons \mu^-$	$\bar{\nu}_\mu + \nu_e + e^+ \rightleftharpoons \mu^+$
$\bar{\nu}_e + e^- \rightleftharpoons \bar{\nu}_\mu + \mu^-$	$\nu_e + e^+ \rightleftharpoons \nu_\mu + \mu^+$

$$\nu_\mu + n \rightleftharpoons p + \mu^- \quad \bar{\nu}_\mu + p \rightleftharpoons n + \mu^+$$

Bollig et al. '18



# Weak Interactions

See talks on Tuesday

## Basic Sets:

$\nu_e n \rightleftharpoons e^- p$	Bruenn (1985)
$\bar{\nu}_e p \rightleftharpoons e^+ n$	Bruenn (1985)
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$\nu N \rightleftharpoons \nu N$	Bruenn (1985)
$\nu A \rightleftharpoons \nu A$	Bruenn (1985), Horowitz (1997)
$\nu e^\pm \rightleftharpoons \nu e^\pm$	Bruenn (1985)
$e^- e^+ \rightleftharpoons \nu \bar{\nu}$	Bruenn (1985)
$NN \rightleftharpoons \nu \bar{\nu} NN$	Hannestad & Raffelt (1998)

Lentz et al. 2011, Kotake et al. 2018

## Hadron Sectors (**Nucleon scattering**):

### Nucleon Neutral Weak Current

$$J_\mu = \langle N(p') | F_1(Q^2) \gamma_\mu + \underline{F_2(Q^2) \sigma_{\mu\nu} q^\nu} + G_A(Q^2) \gamma_\mu \gamma_5 | N(p) \rangle$$

Weak magnetism

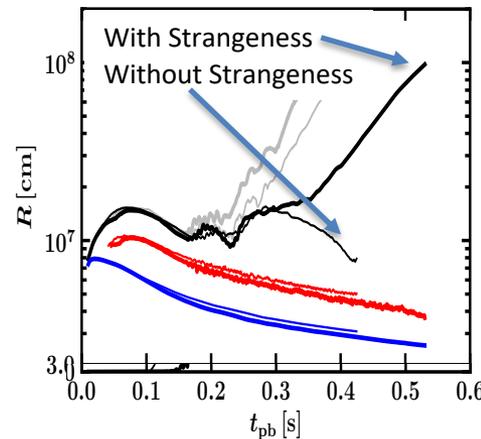
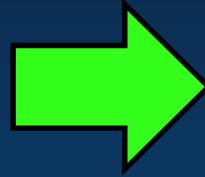
$$G_A(Q^2) = \frac{1}{2} \frac{G_A(0)}{(1 + Q^2/M_A^2)} \tau_3 + \underline{G_A^s(Q^2)}$$

Strangeness contribution

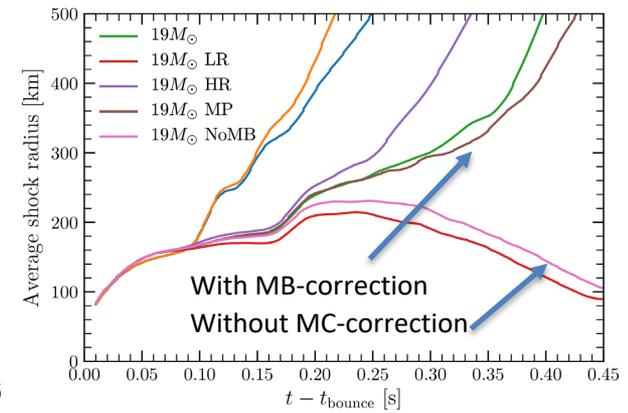
$$\frac{1}{V} \frac{d\sigma}{d\Omega} = \frac{G_F^2 E_\nu^2}{16\pi^2} \left[ g_a^2 (3 - \cos\theta) (n_n + n_p) \underline{S_A} + (1 + \cos\theta) n_n \underline{S_V} \right].$$

Many-body corrections

Extensions



Melson et al. 2015

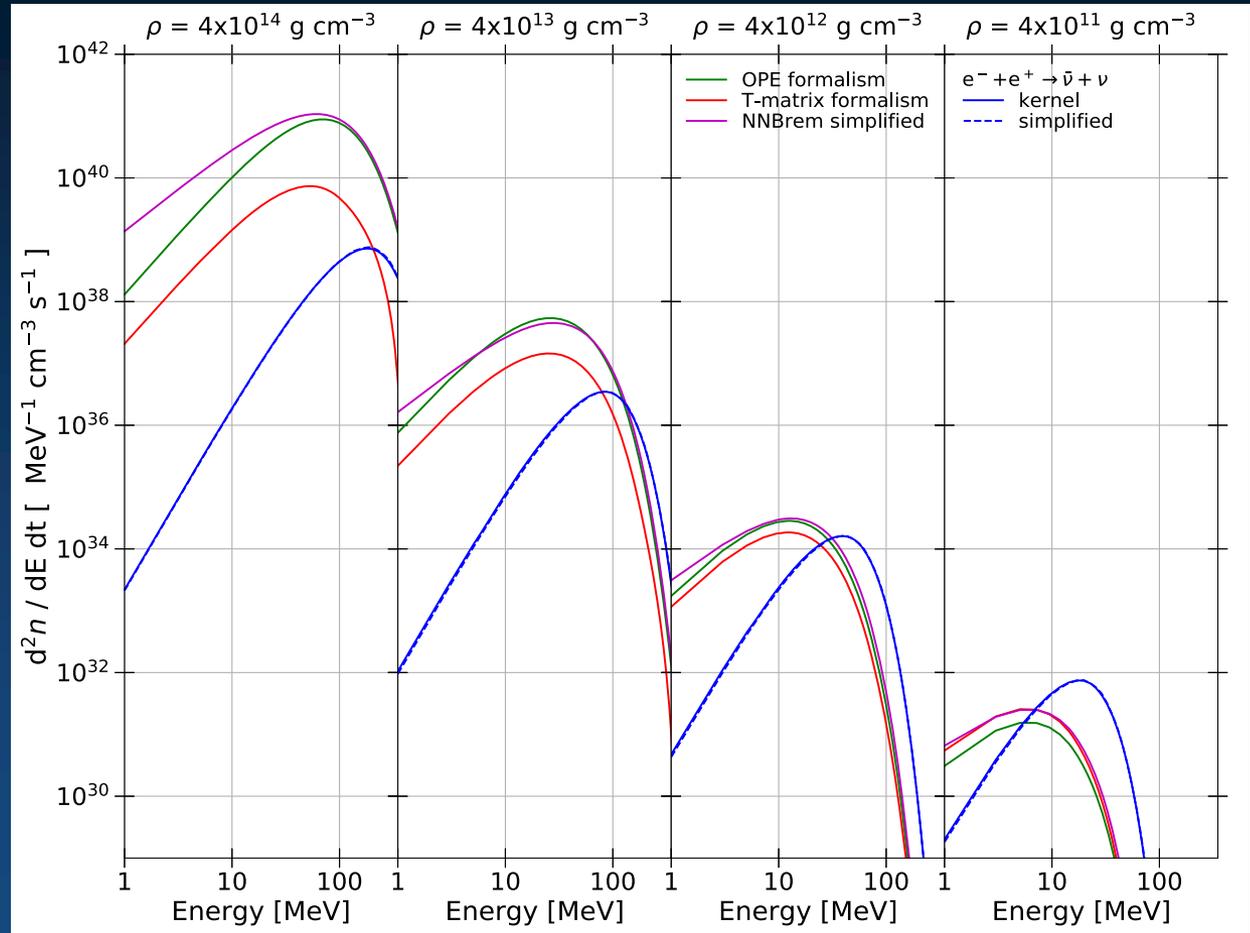
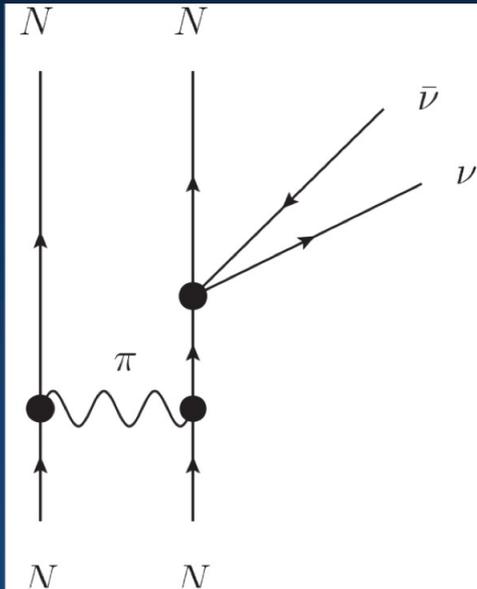


Burrows et al. 2020

# - Nucleon bremsstrahlung of neutrino pairs

See talk by Aurore Betranhandy on Thursday

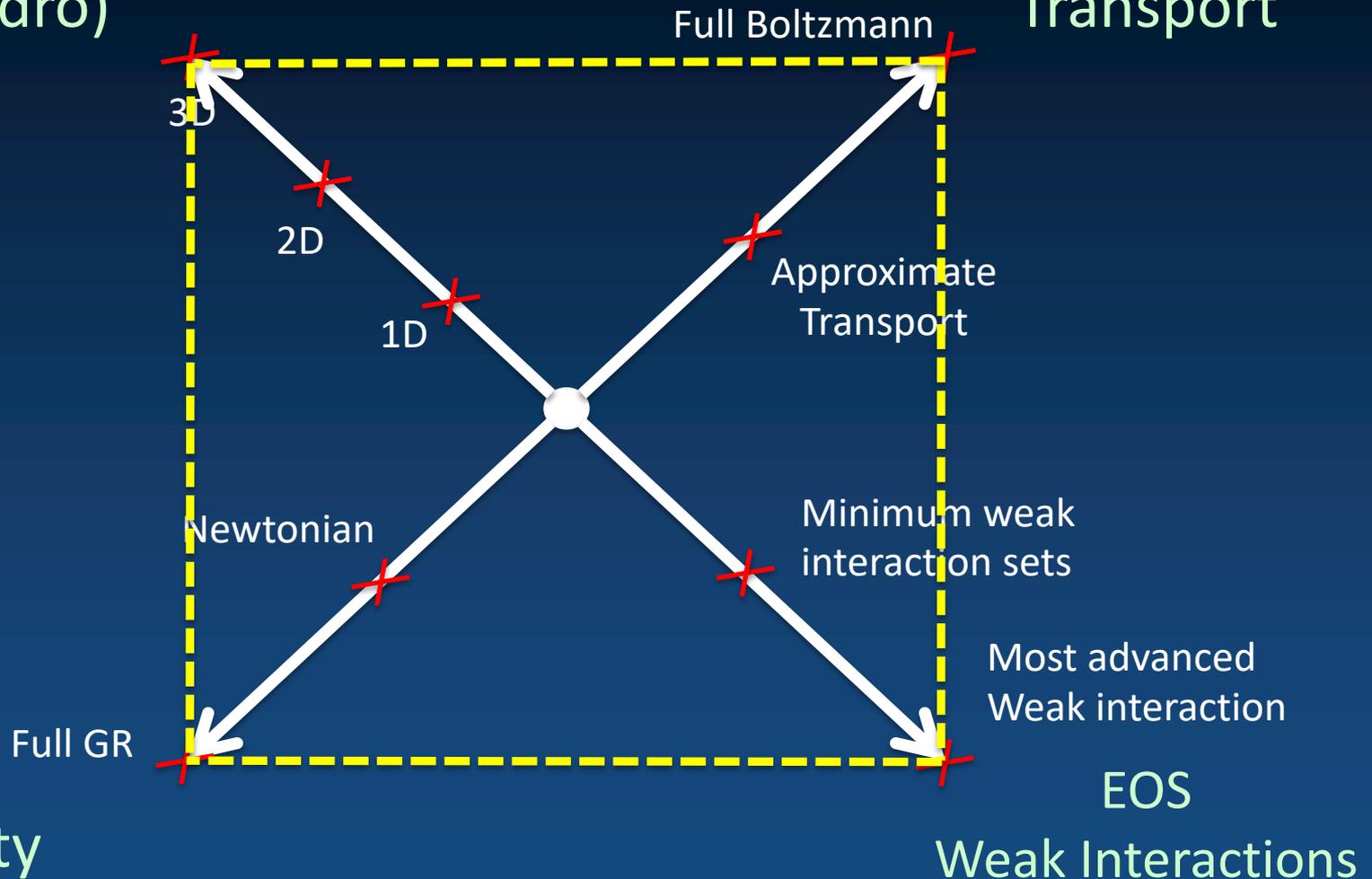
- ✓ Major production channel of muon- and tau- neutrinos
- ✓ Major role in proto-neutron star cooling phase



# - Towards first-principles CCSN simulations

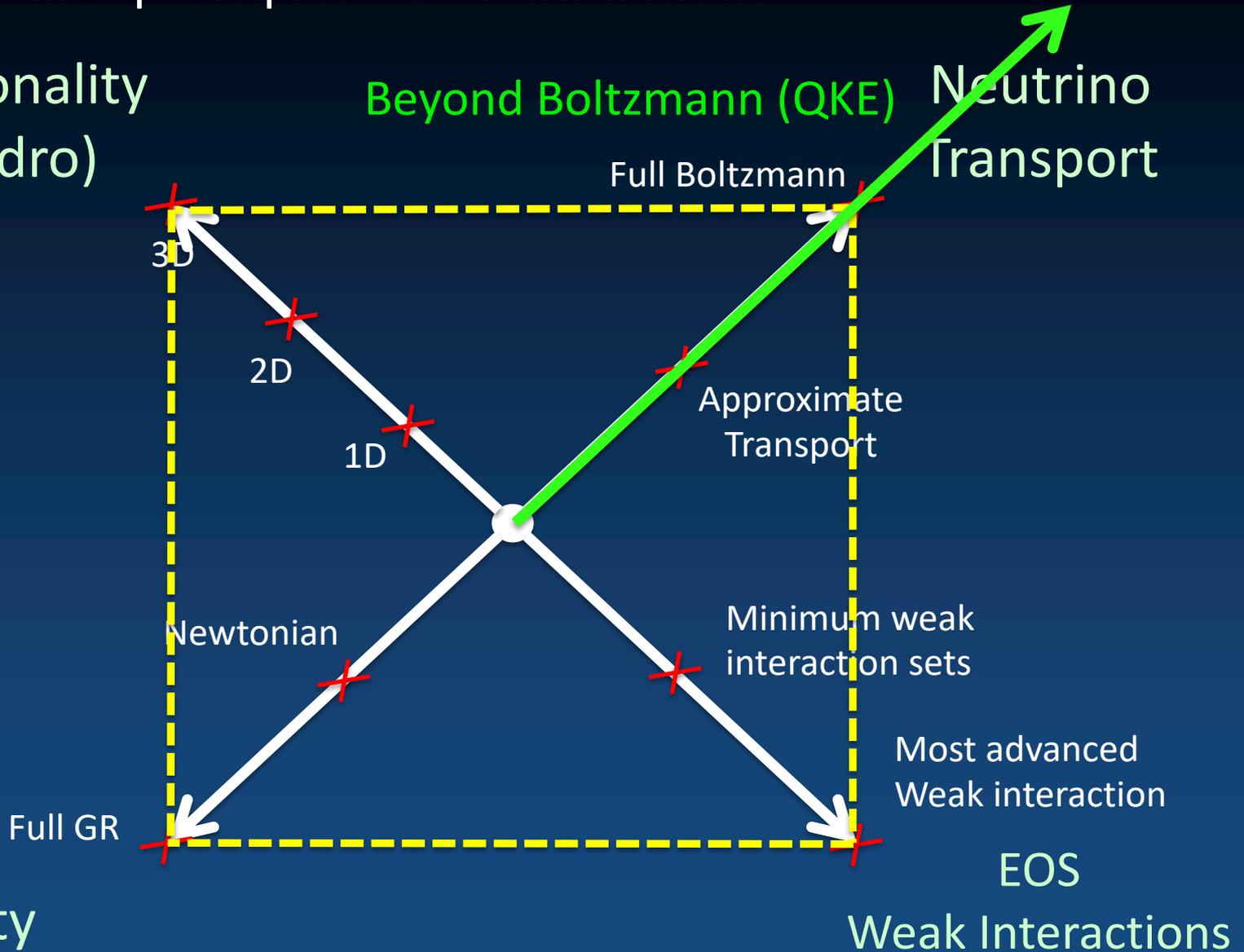
Dimensionality  
(for Hydro)

Neutrino  
Transport



# - Towards first-principles CCSN simulations

Dimensionality  
(for Hydro)



# Neutrino self-interactions can induce flavor-conversion instabilities

See also talk by Gail McLaughlin on Thursday



## Quantum Kinetics neutrino transport

Vlasenko et al. 2014, Volpe 2015,  
Blaschke et al. 2016, Richers et al. 2019

$$p^\mu \frac{\partial f^{(-)}}{\partial x^\mu} + \frac{dp^i}{d\tau} \frac{\partial f^{(-)}}{\partial p^i} = -p^\mu u_\mu S_{\text{col}}^{(-)} + \underbrace{ip^\mu n_\mu [H, f]^{(-)}}_{\text{Oscillation term}},$$

### Density matrix

$$f^{(-)} = \begin{bmatrix} f_{ee}^{(-)} & f_{e\mu}^{(-)} & f_{e\tau}^{(-)} \\ f_{\mu e}^{(-)} & f_{\mu\mu}^{(-)} & f_{\mu\tau}^{(-)} \\ f_{\tau e}^{(-)} & f_{\tau\mu}^{(-)} & f_{\tau\tau}^{(-)} \end{bmatrix}$$

### Hamiltonian

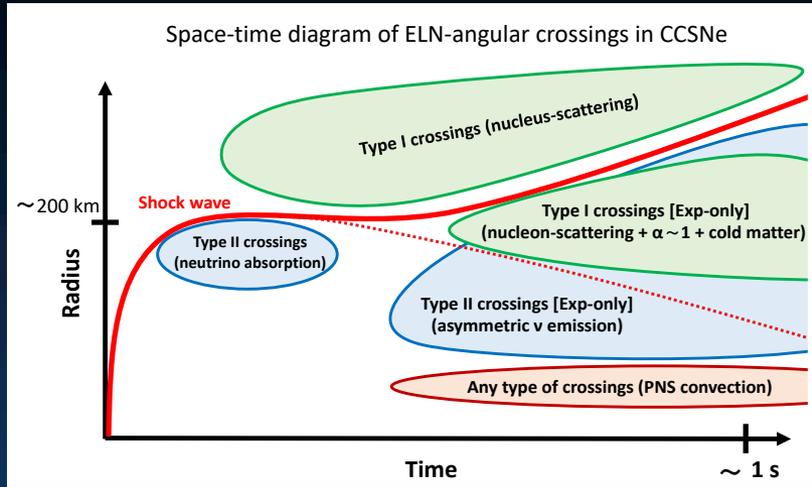
$$\bar{H}^{(-)} = \bar{H}_{\text{vac}}^{(-)} + \bar{H}_{\text{mat}}^{(-)} + \bar{H}_{\nu\nu}^{(-)},$$

Self-interaction

$$H_{\nu\nu} = \sqrt{2}G_F \int \frac{d^3q'}{(2\pi)^3} \left(1 - \sum_{i=1}^3 \ell'_{(i)} \ell_{(i)}\right) (f(q') - \bar{f}^*(q')),$$

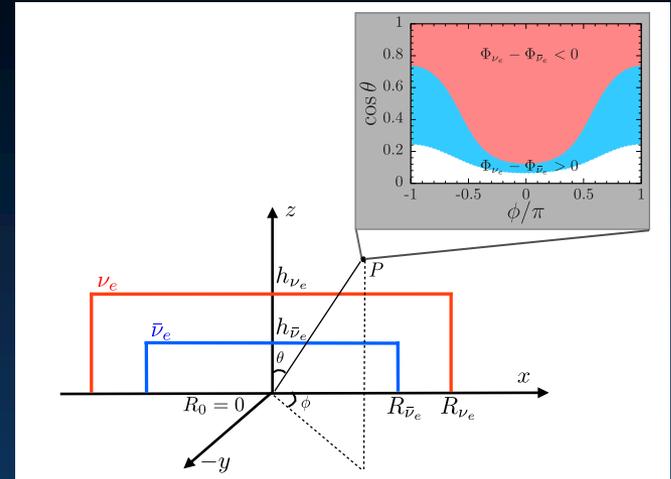
- Fast neutrino-flavor conversion (FFC)

CCSN



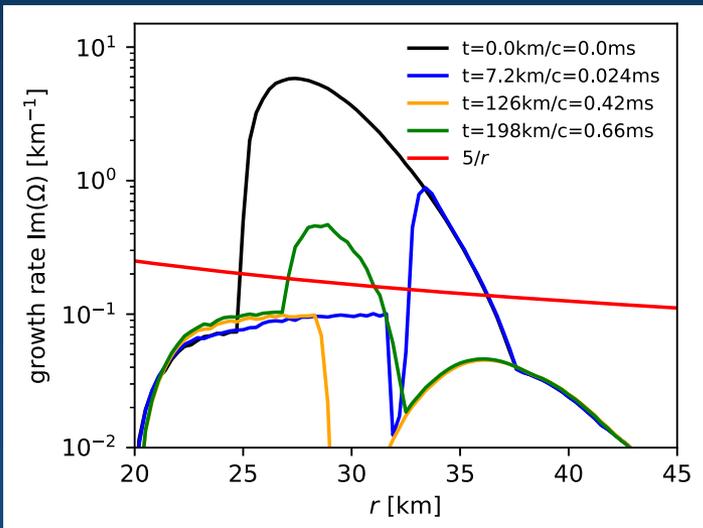
Nagakura et al. 2021

Binary neutron star merger (BNSM)

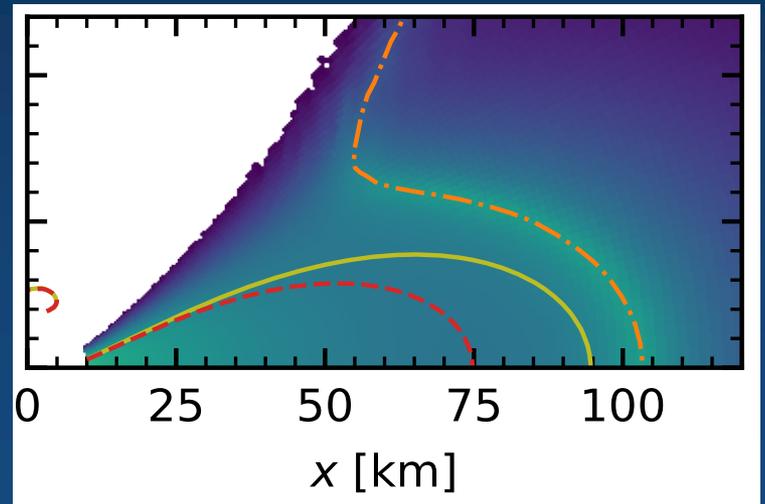


Wu and Tamborra 2017

- Collisional instability



Xiong et al. 2023



Xiong et al. 2022 16

## - Global Simulations: code development

### General-relativistic quantum-kinetic neutrino transport (GRQKNT)

Nagakura 2022

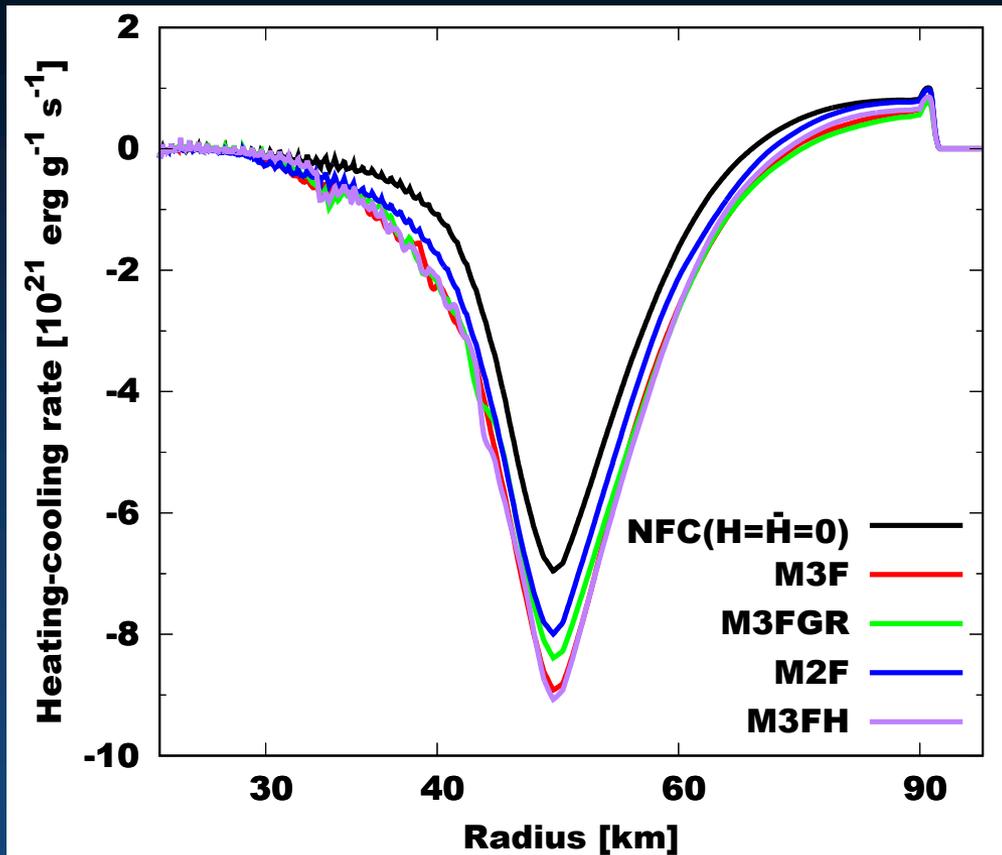
$$p^\mu \frac{\partial f^{(-)}}{\partial x^\mu} + \frac{dp^i}{d\tau} \frac{\partial f^{(-)}}{\partial p^i} = -p^\mu u_\mu \overset{(-)}{S}_{\text{col}} + ip^\mu n_\mu [ \overset{(-)}{H}, f^{(-)} ],$$

- ✓ Fully general relativistic (3+1 formalism) neutrino transport
- ✓ Multi-Dimension (6-dimensional phase space)
- ✓ Neutrino matter interactions (emission, absorption, and scatterings)
- ✓ Neutrino Hamiltonian potential of vacuum, matter, and self-interaction
- ✓ 3 flavors + their anti-neutrinos
- ✓ Solving the equation with Sn method (explicit evolution: WENO-5th order)
- ✓ Hybrid OpenMP/MPI parallelization

# Global simulations of FFC in a CCSN environment

Nagakura PRL 2023

## Neutrino heating/cooling



### Numerical setup:

Collision terms are switched on.

Fluid-profiles are taken from a CCSN simulation.

General relativistic effects are taken into account.

A wide spatial region is covered.

Three-flavor framework

Neutrino-cooling is enhanced by FFCs  
Neutrino-heating is suppressed by FFCs



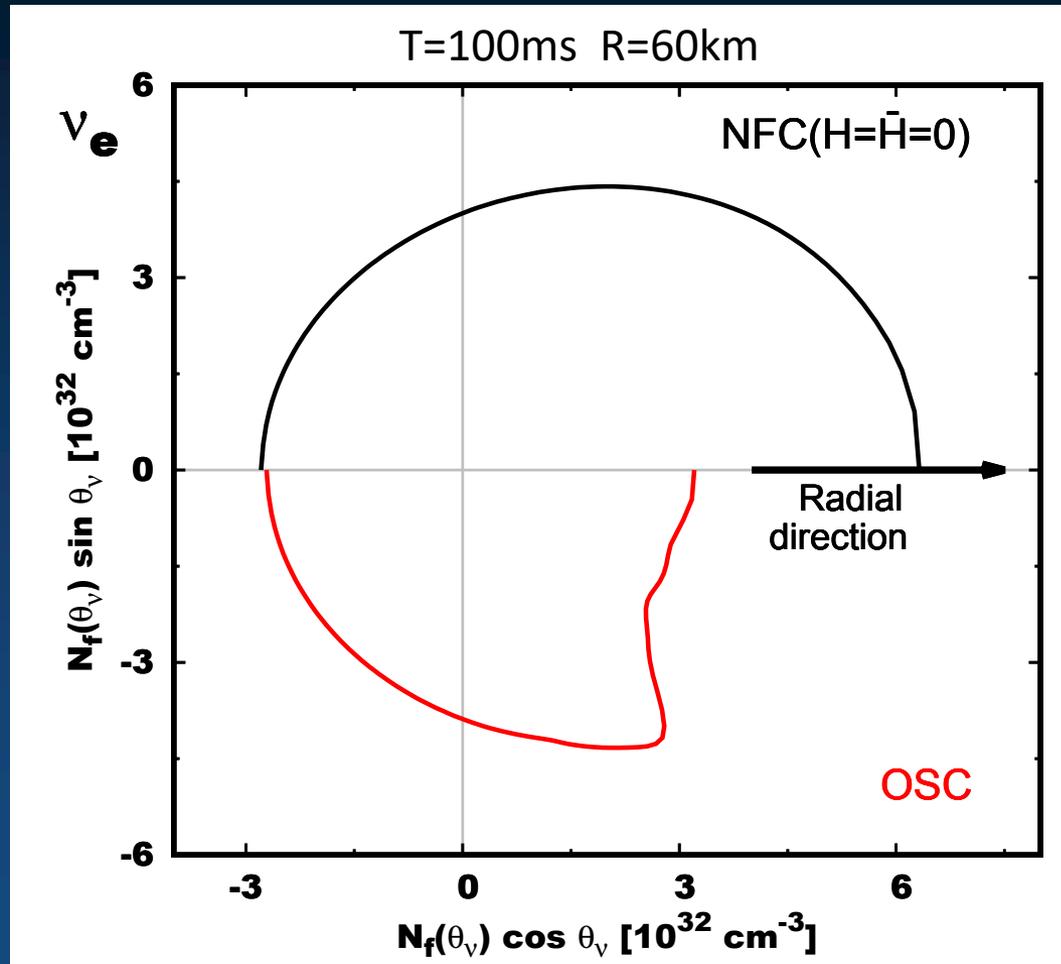
Impacts on CCSN explosion !!

# Global simulations of FFC in a CCSN environment

Nagakura and Zaizen (arXiv:2308.14800)

$K^{rr}$  (r-r component of Eddington tensor) becomes **less than 1/3**.

Angular distribution

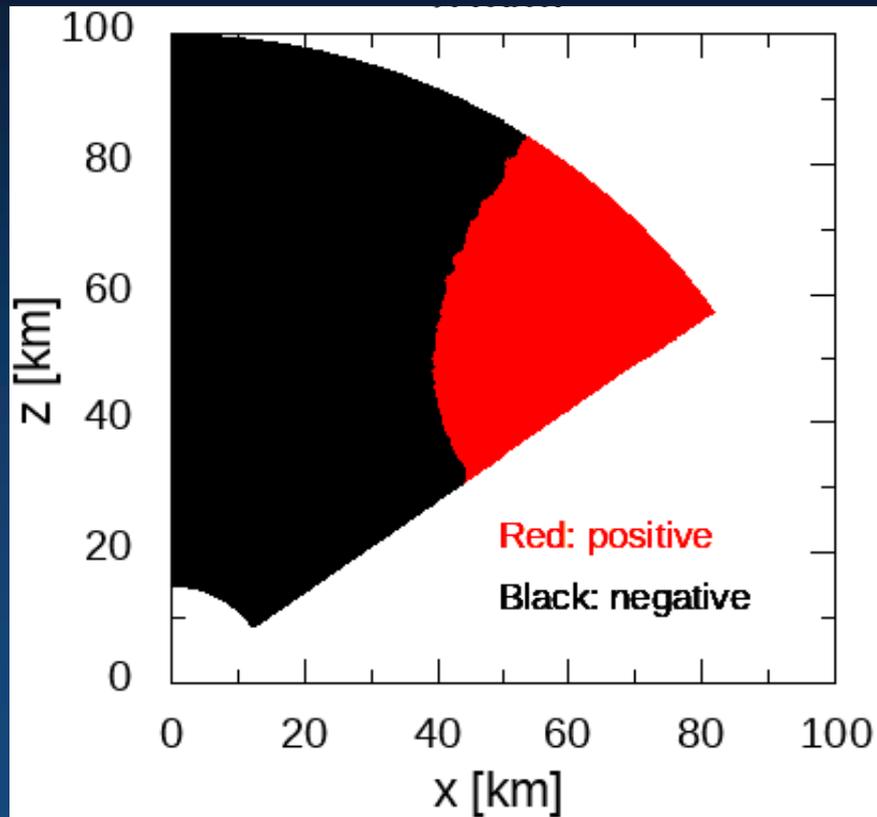


# Global Simulations of FFC in binary neutron star merger remnant

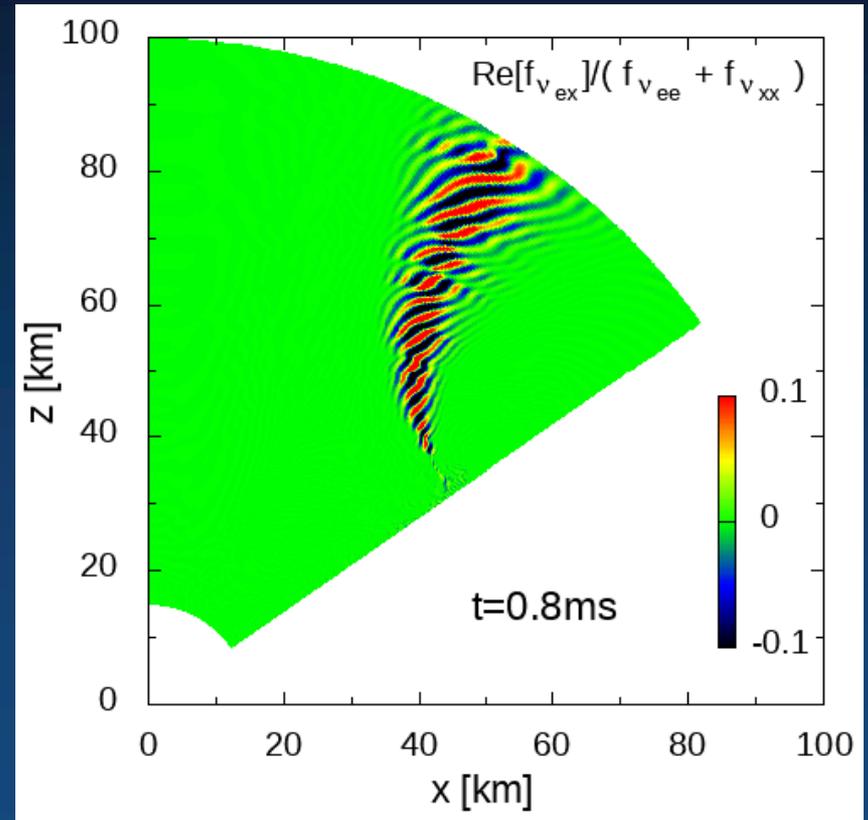
Nagakura (arXiv:2306.10108)

✓ EXZS (ELN-XLN Zero Surface):

ELN - XLN

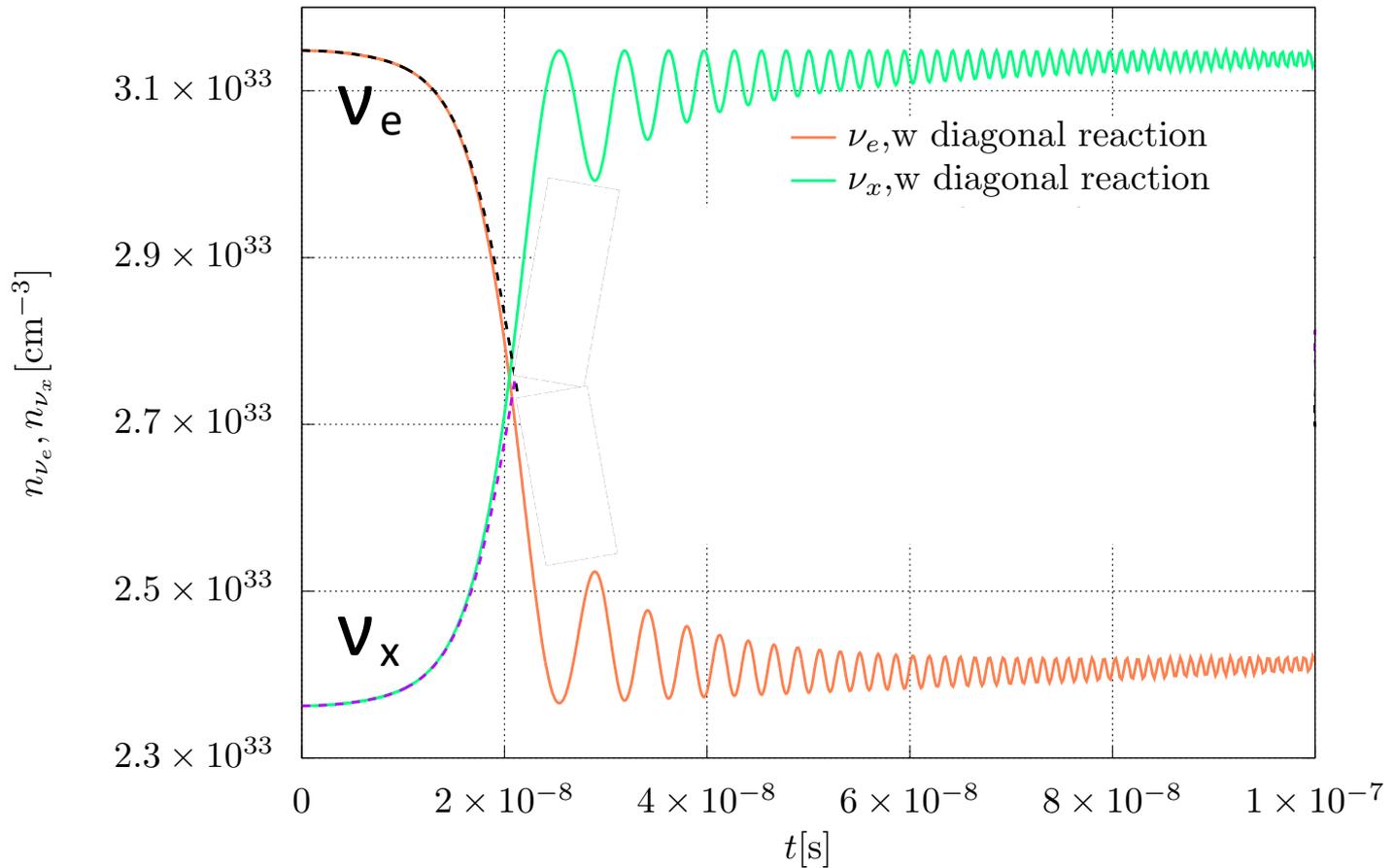


Flavor coherency



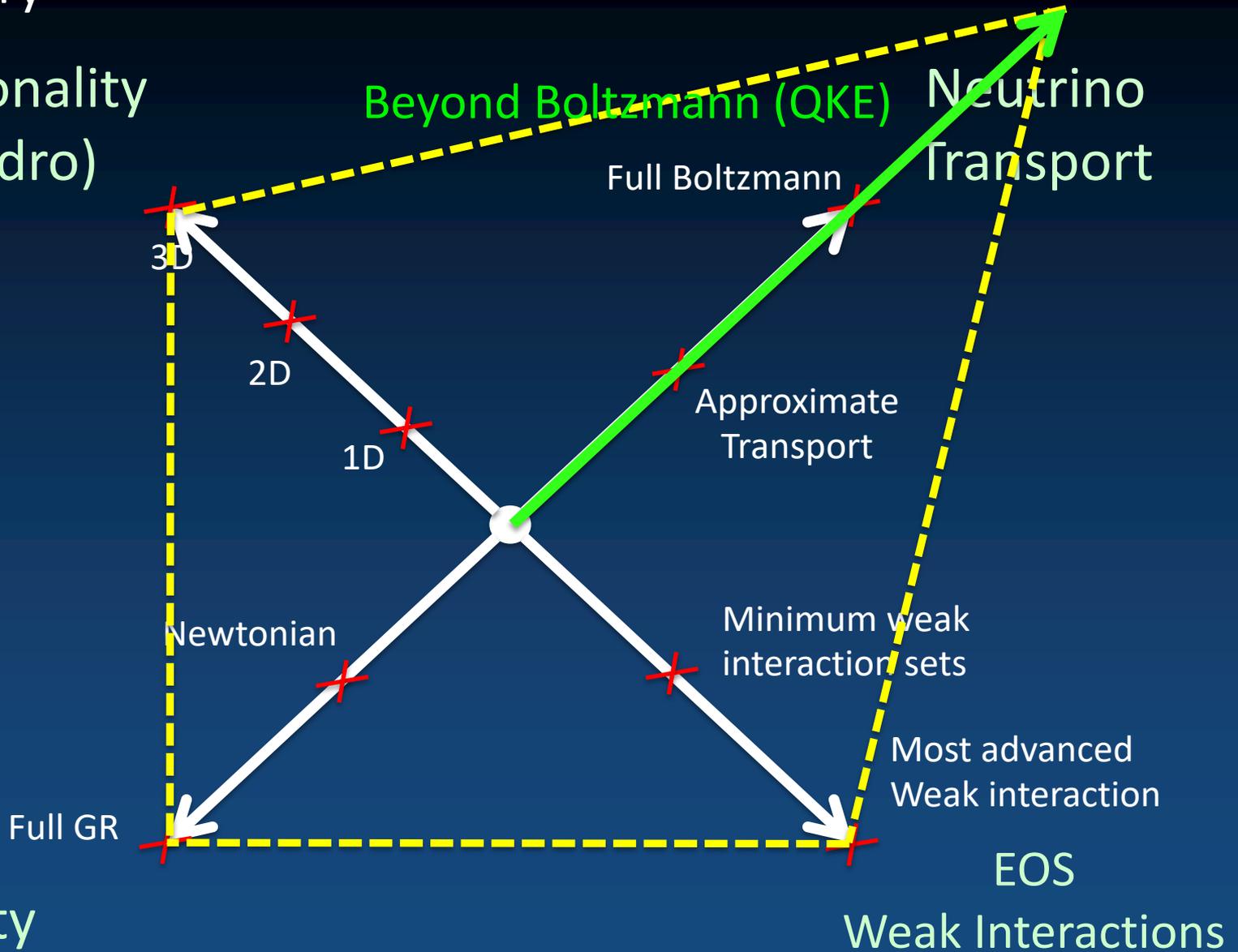
# Collisional flavor swap (associated with collisional instability)

Kato, Nagakura, and Johns (arXiv:2309.02619)



# - Summary

Dimensionality  
(for Hydro)



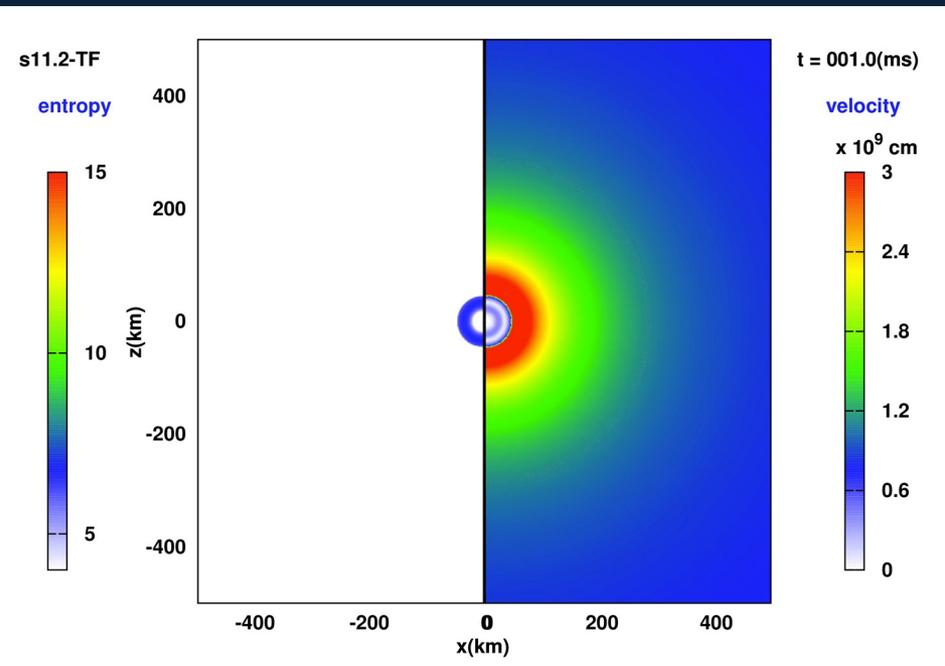
# Backup

# Multi-dimensional (or alternative) CCSN simulations

See also other talks:

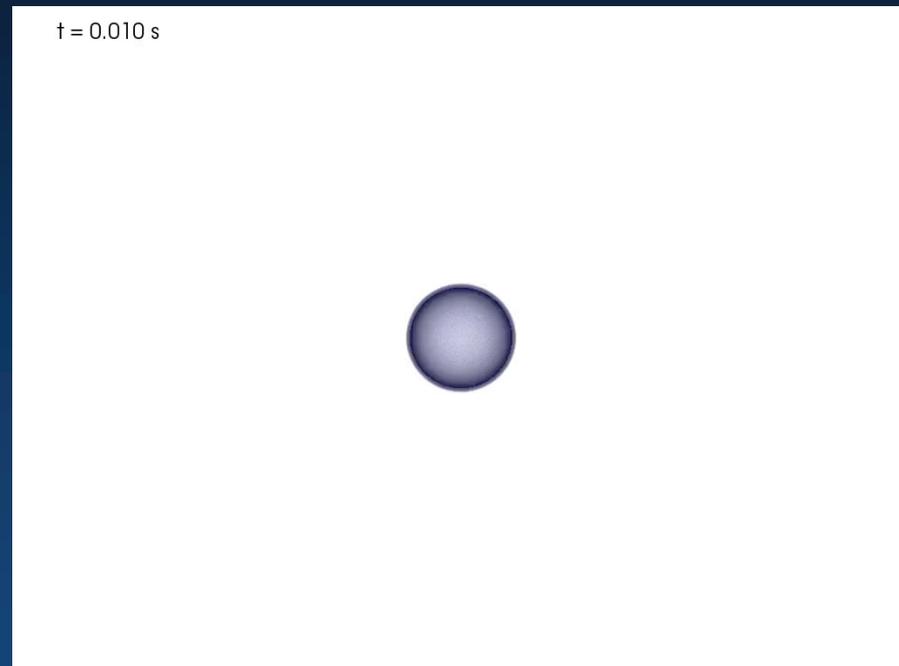
H. Andresen, Boccioli, M. Mori, Gogilashvili, Dunham, Pajkos, O. Andersen, Endeve, Akaho, Betranhandy, Yeow, K. Mori

CCSN simulations with full Boltzmann transport



Nagakura et al. 2019

CCSN simulations with two-moment method

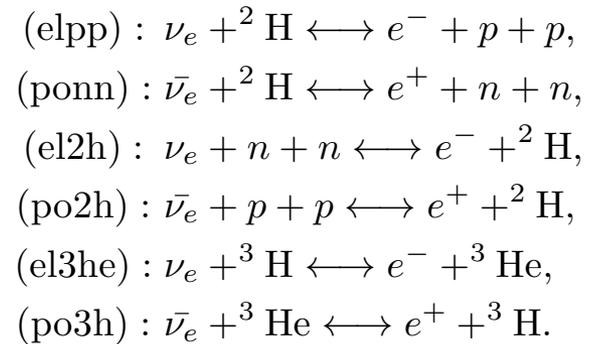
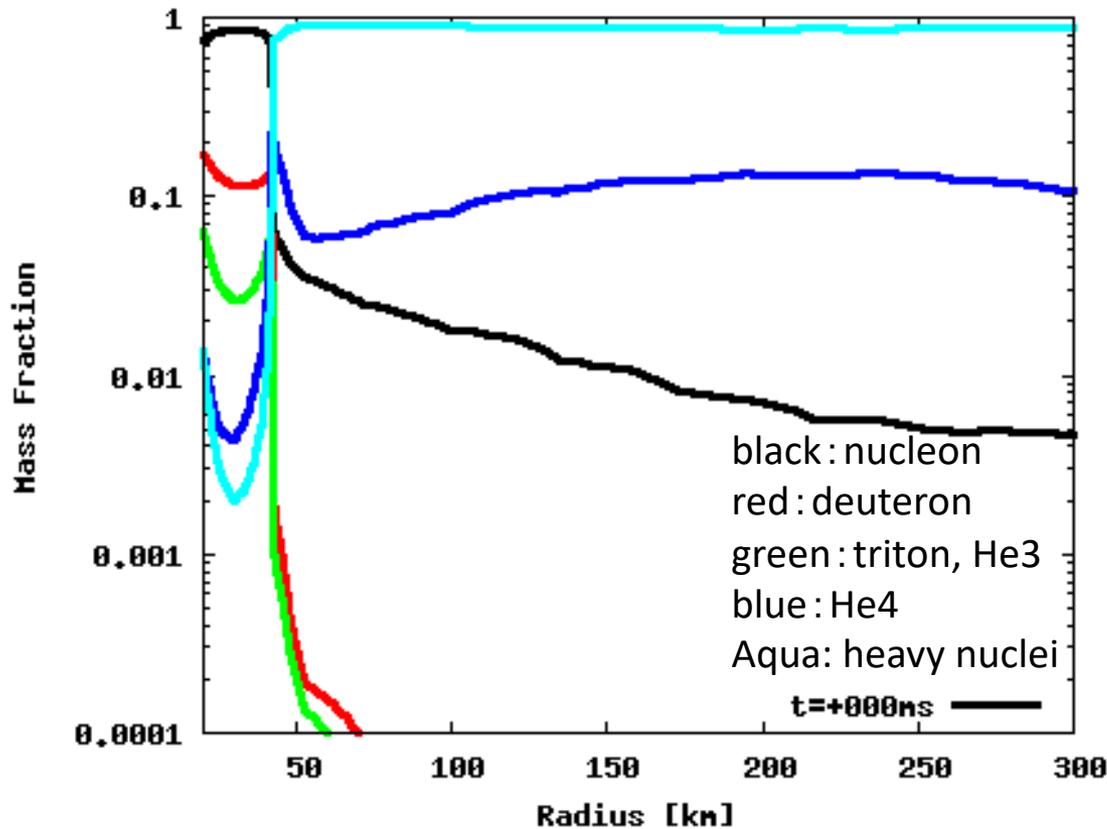


Nagakura et al. 2019

Neutrino kinetics (transport, neutrino-matter collisions, and oscillation) plays key roles on CCSN dynamics

## - Weak reactions with light nuclei

Nagakura et al. 2019, Furusawa and Nagakura 2022



Multi-nuclear treatments of EOS are mandatory for accurate computations of nuclear-weak reaction rates

Hempel et al. 2011, Steiner et al. 2013, Furusawa and Nagakura et al. 2017

# Various Approximations for Multi-D Neutrino Transfer

See a review by Mezzacappa et al. 2020

## ✓ Ray-by-Ray Approach (UTK-Oak Ridge, MPA)

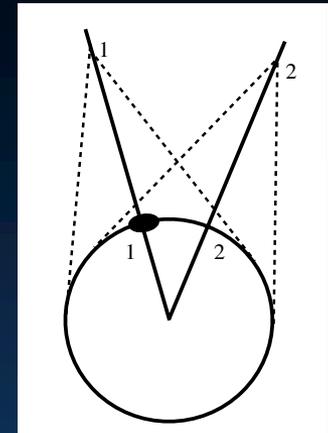
Neutrino-transport is essentially same as spherical symmetry.

## ✓ Isotropic Diffusion Source Approximation (IDSA) (Basel, Japan)

Neutrinos are decomposed into trapped and streaming parts.

Two reduced equations are coupled by each source term, which is approximately described under diffusion treatment.

(See e.g., Berninger et al. 2013)



Schematic picture of ray-by-ray approach (Lentz et al. 2012)

## ✓ Moment method (Many groups....)

Neutrino angular direction is integrated. The so-called “closure relation” is imposed in the higher moment.

$$M_{(\nu)}^{\alpha_1 \alpha_2 \dots \alpha_k}(x^\beta) = \int \frac{f(p'^\alpha, x^\beta) \delta(\nu - \nu')}{\nu'^{k-2}} p'^{\alpha_1} p'^{\alpha_2} \dots p'^{\alpha_k} dV'_p,$$

Shibata et al. 2011

## ✓ Multi-Group Flux-Limited-Diffusion (MGFLD) (UTK-Oak Ridge)

Neutrino Transports are treated as the Energy-Dependent Diffusion Equation.

# Numerical methods of Boltzmann solver (Sn method)

Large-matrix Inversion is required.

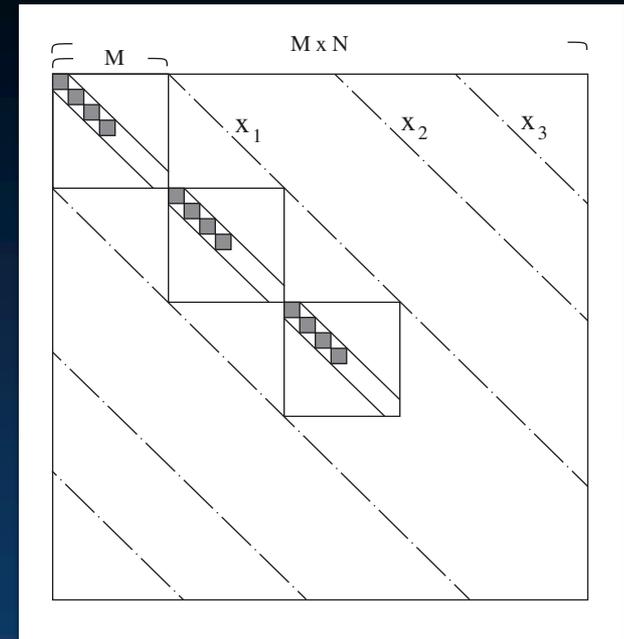
$$\begin{aligned} \frac{\partial f}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} (f \cos \bar{\theta} r^2) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (f \sin \theta \sin \bar{\theta} \cos \bar{\varphi}) \\ + \frac{1}{r \sin \theta} \frac{\partial}{\partial \varphi} (f \sin \bar{\theta} \sin \bar{\varphi}) - \frac{1}{r \sin \bar{\theta}} \frac{\partial}{\partial \bar{\theta}} (f \sin^2 \bar{\theta}) \\ - \frac{\partial}{\partial \bar{\varphi}} \left( f \frac{\cot \theta}{r} \sin \bar{\theta} \sin \bar{\varphi} \right) = S_{\text{rad}}. \end{aligned}$$



Implicit time evolution

$$\mathbf{Ax} = \mathbf{b} \quad (\text{Matrix Equation}) \quad a_j^i f_i^{(n+1)} = b_j(f^{(n)})$$

$\mathbf{A} =$



Block-diagonal sparse matrix

✓ Solved by BiCGSTAB with Damped Jacobi-type Preconditioner

(Imakura et al. 2012)

✓ Scale of axisymmetric simulations

Memory:  $\sim 2$  TB,

Operation: 20TFlops  $\times$  2000 hours

We achieve  $\sim 10\%$  performance on “K” and “Fugaku” supercomputers

Full 7D simulation needs **100 times** computational resources are necessary.

# Rich flavor-conversion phenomena driven by neutrino-neutrino self-interactions

## - Slow-mode (Duan et al. 2010)

- Energy-dependent flavor conversion occurs.
- The frequency of the flavor conversion is proportional to

$$\sqrt{\omega\mu}$$

Vacuum:	$\omega = \frac{\Delta m^2}{2E_\nu}$ ,
Matter:	$\lambda = \sqrt{2}G_F n_e$ ,
Self-int:	$\mu = \sqrt{2}G_F n_\nu$ ,

## - Fast-mode (FFC) (Sawyer 2005)

- Collective neutrino oscillation in the limit of  $\omega \rightarrow 0$ .
- The frequency of the flavor conversion is proportional to
- Anisotropy of neutrino angular distributions drives FFCs.

$$\mu$$

## - Collisional instability (Johns 2021)

- Asymmetries of matter interactions between neutrinos and anti-neutrinos drive flavor conversion.

$$\text{Im } \Omega \cong \pm \frac{\Gamma - \bar{\Gamma}}{2} \frac{\mu S}{\sqrt{(\mu D)^2 + 4\omega\mu S}} - \frac{\Gamma + \bar{\Gamma}}{2}$$

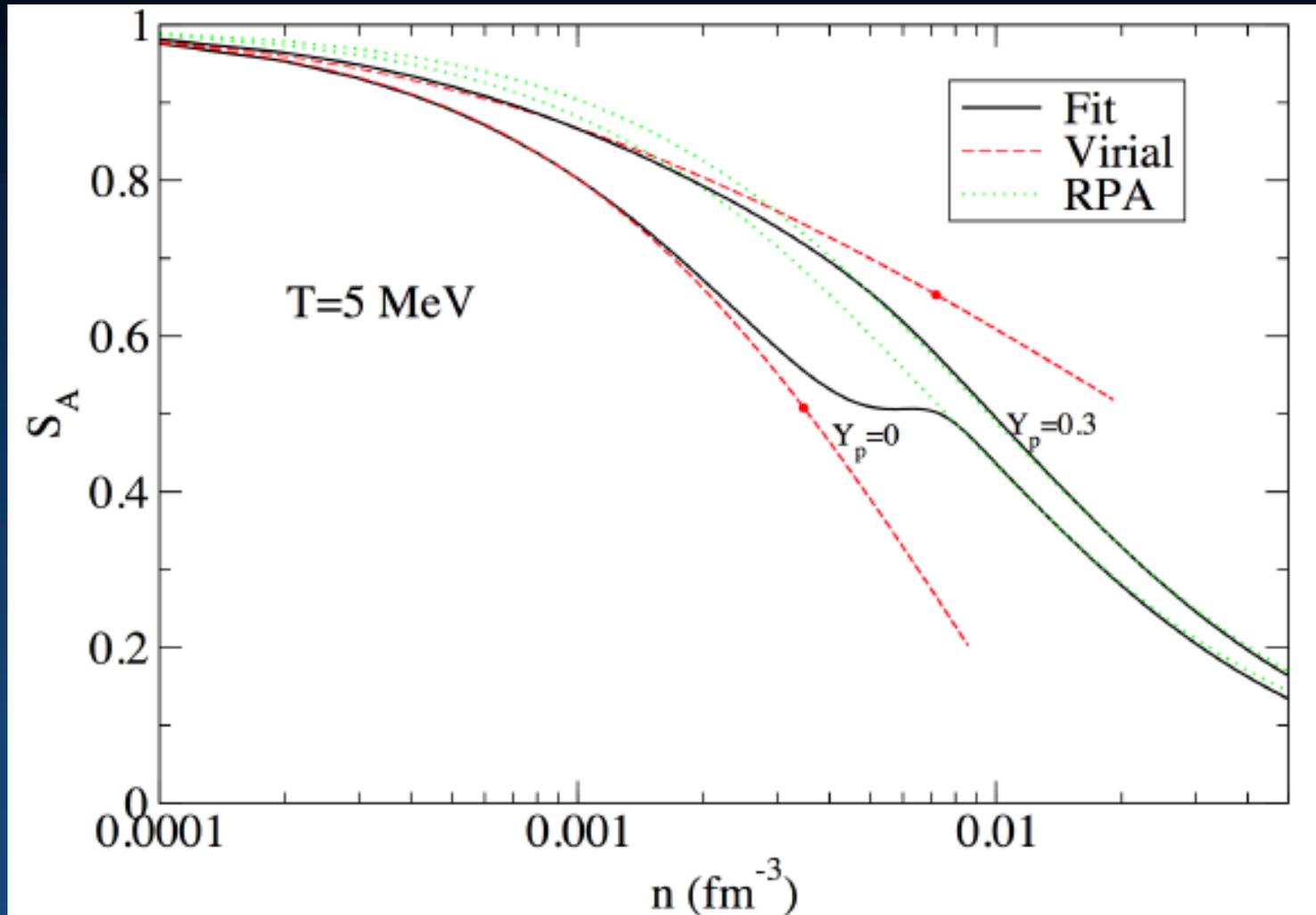
$\Gamma$ : Matter-interaction rate

## - Matter-neutrino resonance (Malkus et al. 2012)

- The resonance potentially occur in BNSM/Collapsar environment (but not in CCSN).
- Essentially the same mechanism as MSW resonance.

$$|\lambda + \mu| \sim |\omega|$$

# The spin (axial) $S_A$ response

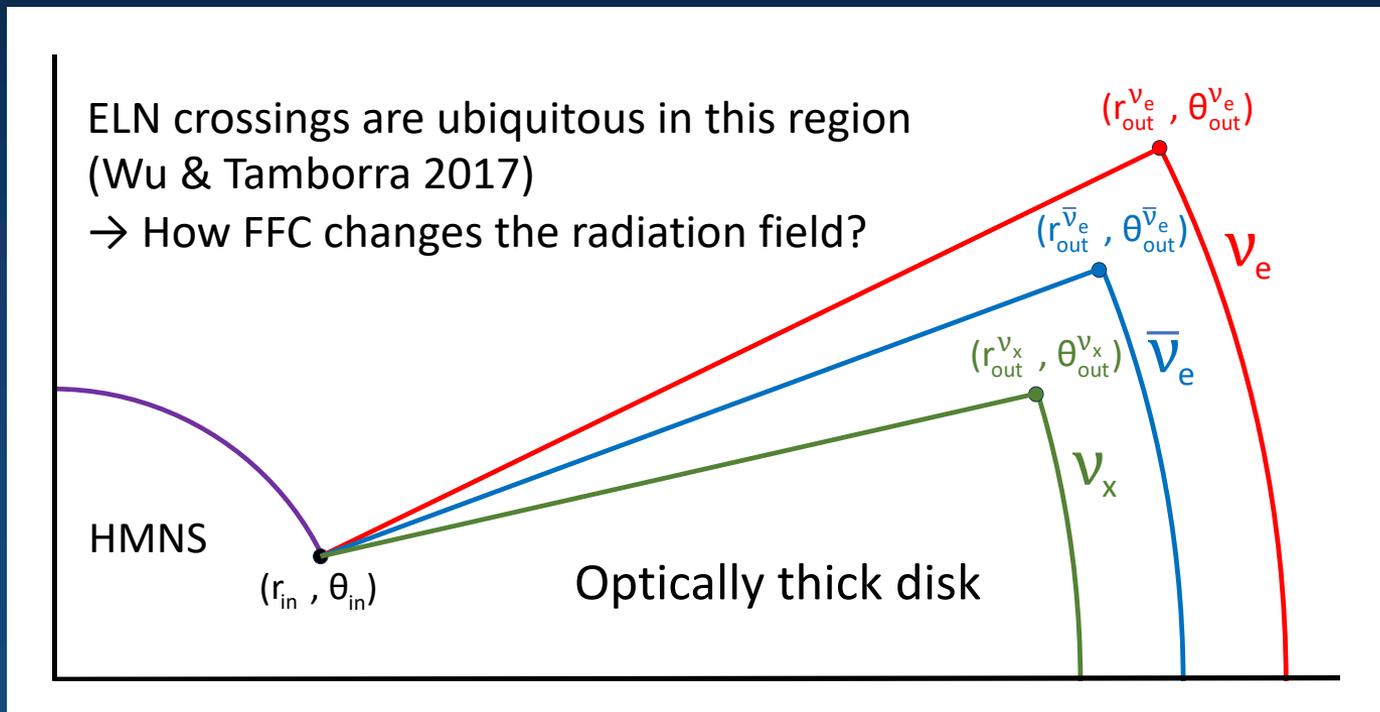


# Global Simulations of FFC in **binary neutron star merger**

Nagakura (arXiv:2306.10108)

## ✓ Setup:

- Hypermassive neutron star (HMNS) + disk geometry
- Thermal emission on the neutrino sphere
- QKE (FFC) simulations in axisymmetry
- Resolutions:  $1152 (r) \times 384 (\theta) \times 98 (\theta_v) \times 48 (\phi_v)$

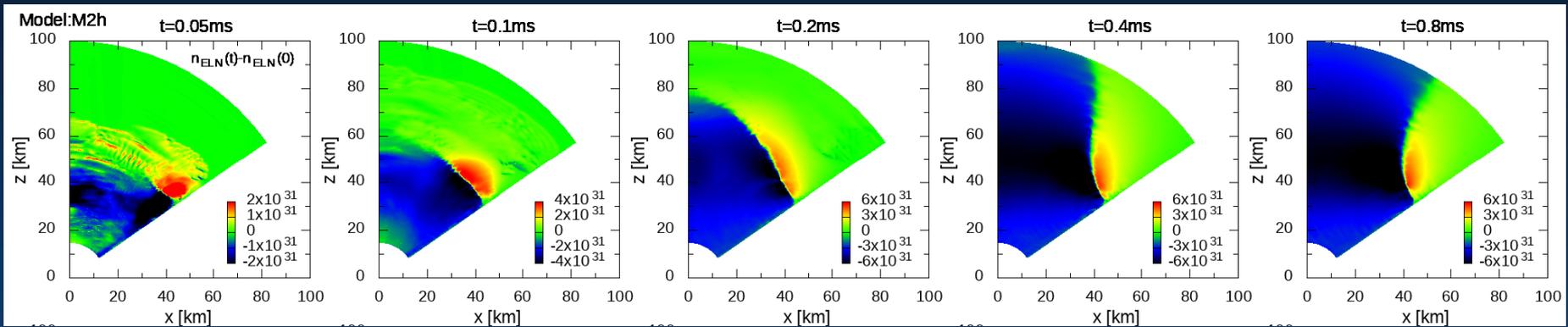


# Global Simulations of FFC in a BNSM environment

Nagakura (arXiv:2306.10108)

## Temporal evolution of FFCs in global scale:

$ELN(t) - ELN(0)$



Time

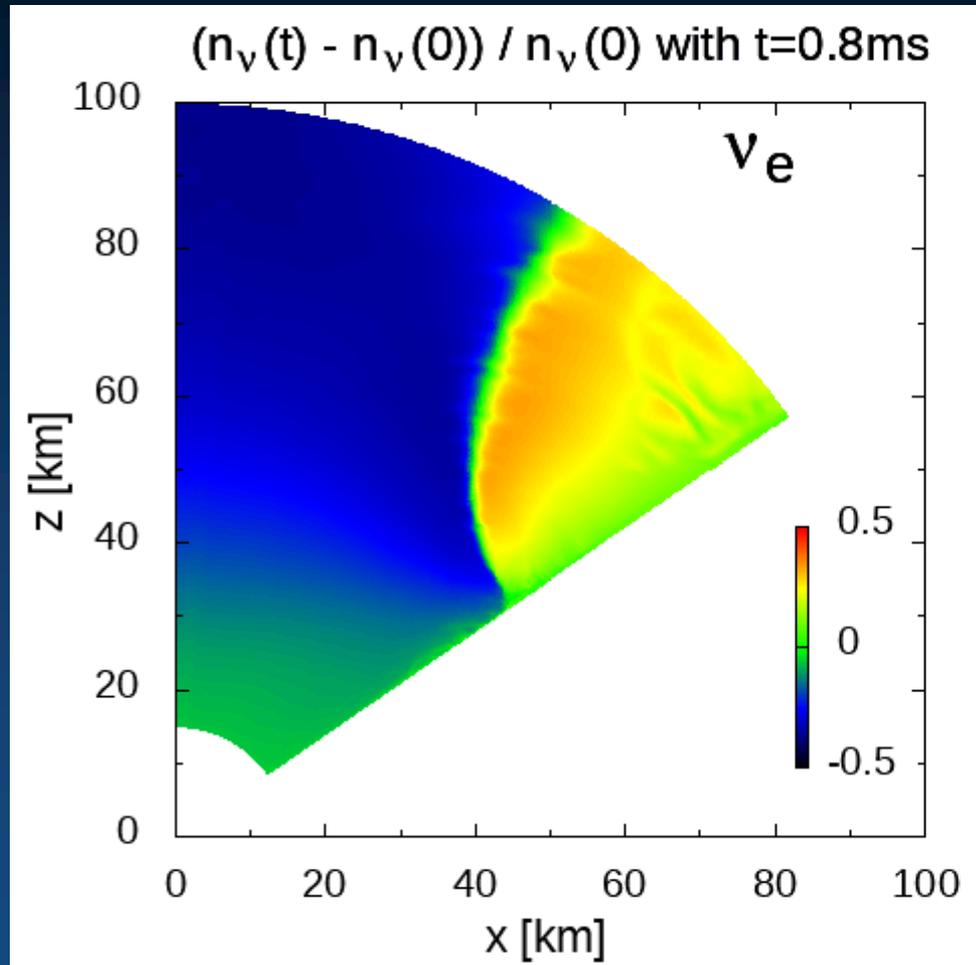
Take-home message 1

Non-conservations of ELN (and XLN) number density  
represent the importance of global advection of neutrinos in space!

# Global Simulations of FFC in **binary neutron star merger**

Nagakura (arXiv:2306.10108)

✓ Substantial change of neutrino radiation field:



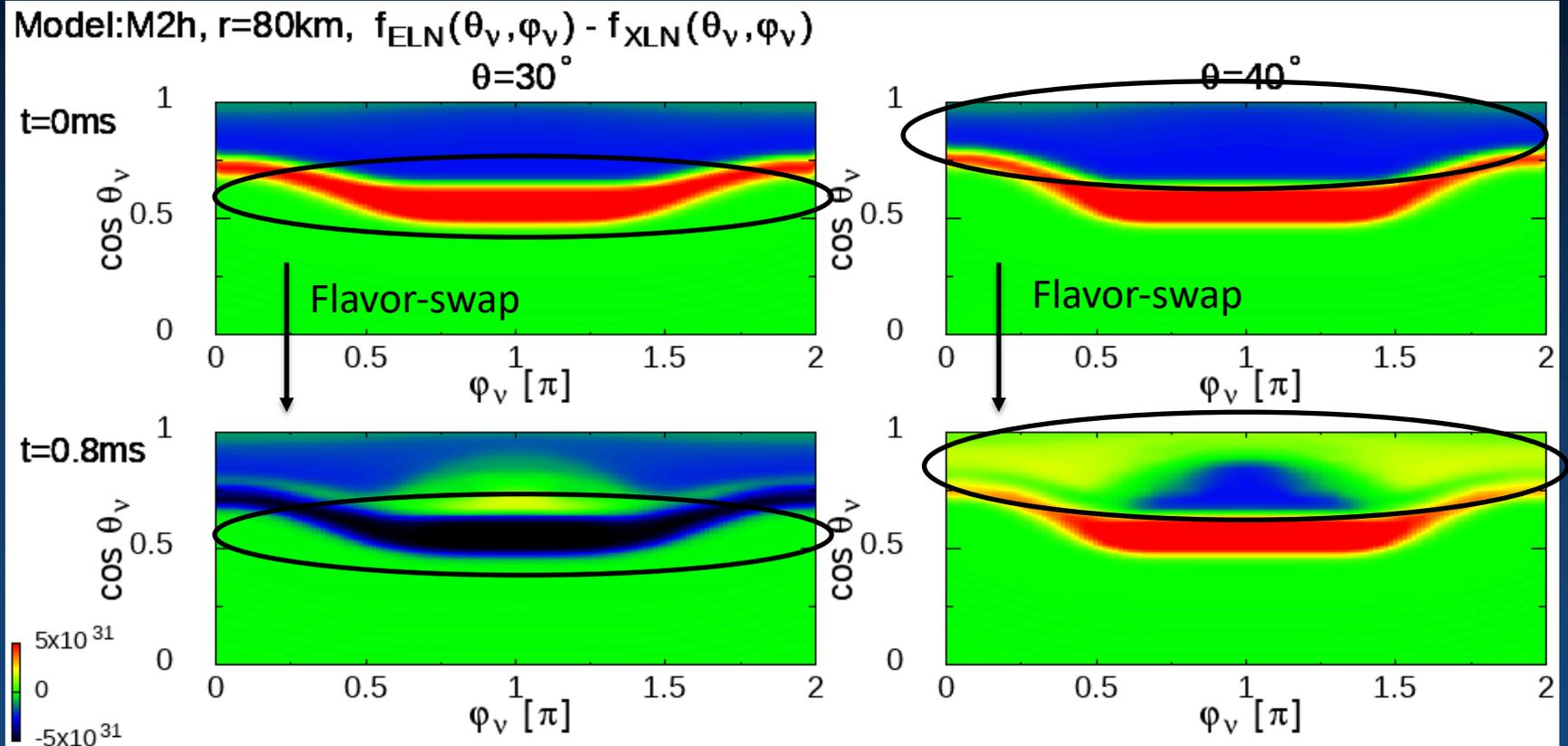
Note: Increase or decrease of electron-type neutrinos hinge on heavy-leptonic neutrinos

More detailed study is required!!

# Global Simulations of FFC in a BNSM environment

Nagakura (arXiv:2306.10108)

## ✓ Flavor swap between electron- and heavy-leptonic neutrinos:



# Global simulations of FFC in a CCSN environment

Nagakura and Zaizen (arXiv:2308.14800)

- Eddington tensor (and comparing to analytic closure relations)

