N* Physics with CLAS

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ECT*-APCTP Joint Workshop: Exploring Exploring Resonance Structure with Transition GPDs August 21-25, 2023 Trento, Italy





CLAS EXPERIMENT 1997-2012





CLAS12 EXPERIMENT 2018 ~





N* Program

- 1. Spectroscopy (mainly driven by real photon scattering)
 - Measure the N* spectrum more precisely and more completely.
 - The impact of nucleon spectroscopy for QCD compared with that of atomic spectroscopy for QED.
- 2. Structure of excited baryons (mainly driven by electron scattering)
 - Studying underlying symmetries of hadron system
 - Understanding the effective degrees of freedom



Quark Model Classification of N*



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Search for excited baryons in 2-body channels

Lab

data ac	cqui	irec	d		- analyzed/published									CEBAF Large A			
Observable	σ	Σ	Т		Р	E	F	G	Н	T _x	Tz	L _x	Lz	O _x	O _z	C _x	Cz
0	\checkmark	\checkmark	1			1	1	1		1							
ρπ ⁺				-						+							
				-						+				N			
pri pri										+			γp-	→ X			
κ+ ν				/	\checkmark					1	1	1	1	V	\checkmark	V	V
Κ+Σ 0	V	v	v	/	\checkmark	 √								√	√	√	√
ρω/Φ	√	√	√	/		√	√	✓	√				✓ SD	ME			
K ^{+*} Λ	√				\checkmark								✓ SD	ME			
K ^{0*} Σ ⁺	v	 ✓ 	+									 ✓ 	✓		SD	OME	
рπ⁻	\checkmark	\checkmark				\checkmark	✓	\checkmark					yn-	→x			
pρ⁻	\checkmark	\checkmark				\checkmark	\checkmark	\checkmark									
Κ -Σ+	\checkmark	\checkmark				\checkmark	✓	 ✓ 									
					1		1	1		1	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark	1
K ⁰ Λ		\checkmark			v	v	×										
Κ⁰Λ		✓ ✓	✓ ✓		✓ ✓		 ✓ 	✓ ✓		√ 	✓	 ✓ 	✓	✓	 ✓ 	 ✓ 	 ✓

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Establishing the N* spectrum

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Establishing the N* spectrum 0.2 Hyperon photoproduction $\gamma p \rightarrow K^+ \Lambda \rightarrow K^+ p \pi^ \frac{P}{105}$ $\frac{1}{105}$ $\frac{1$

Target T and beam-target E spin asymmetries

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Establishing the N* and Δ^* Spectrum

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Quark Model Classification of N*

Quark Model Classification of N*

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BnGa energy-dependent coupled-channel PWA of CLAS $K^+\Lambda$ and other data

New evidence for excited nucleons

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State N((mass)J ^P	PDG pre 2012	PDG 2018 evidence	Mass (Pole)	
N(1710)1/2+	***	***	1700	
N(1880)1/2+		***	1860	
N(2100)1/2+	*	***	2100	*
N(1895)1/2 ⁻		* * * *	1910	*
N(1900)3/2+	**	***	1920	*
N(1875)3/2 ⁻		* * *	1900	
N(2120)3/2-		***	2100	
N(2060)5/2-		***	2070	
∆(2200)7/2 ⁻	*	***	2150	

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****	- existence is certain
***	- existence is likely
**	- evidence of existence is fair

- evidence of existence is poor

Photoproduction data led to the discovery of new states and fully established poorly known states.

Structure of Excited Baryons with Electromagnetic Probes

Central question in hadron physics What are the effective degrees of freedom at varying distance scale?

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Measure the strength of resonance excitations versus distance scale in meson electroproduction

Inclusive electron scattering

p(e,e')X

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 In contrast to elastic scattering, resonances cannot be uniquely separated in inclusive scattering → need to measure exclusive processes.

Inclusive vs. Exclusive electron scattering

• Resonance states can be better isolated by exclusive processes.

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Missing Mass² vs W for $ep \rightarrow e'pX$

Final data sample from 6 GeV run (e16)

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W-Dependence of Selected Exclusive Channels

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Integrated cross section in the resonance region

 $\gamma^* p \rightarrow \pi^+ n$

UCONN | UNIVERSITY OF CONNECTICUT K. Park et al., PR C77 (2008) 015208; PR C91 (2015) 045203

Why study more than one resonance, aren't they all the same?

 \rightarrow States with different quantum numbers respond differently to increase in Q².

Exclusive Electroproduction x-section and BSA

$$\frac{Cross \ section}{dt} (\text{longitudinally pol. beam and unpol. target} \left[2\pi \frac{d^2\sigma}{dtd\phi} = \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cdot \cos(2\phi) \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cdot \cos(\phi) \frac{d\sigma_{LT'}}{dt} + h \cdot \sqrt{2\epsilon(1-\epsilon)} \cdot \sin(\phi) \frac{d\sigma_{LT'}}{dt} \right]$$

$$\frac{Beam \ Spin \ Asymmetry:}{BSA(t, \phi, x_B, Q^2) = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-}} = \frac{\sqrt{2\epsilon(1-\epsilon)} \frac{\sigma_{LT'}}{\sigma_0}}{1 + \sqrt{2\epsilon(1+\epsilon)} \frac{\sigma_{LT}}{\sigma_0}} \cos \phi + \epsilon \frac{\sigma_{TT}}{\sigma_0} \cos 2\phi}$$

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Power of Interference

• Unpolarized structure function

 $\sigma_{LT} \sim Re(L*T) = Re(L)Re(T) + Im(L)Im(T)$

- Amplify small resonance multipole by an interfering larger resonance multipole
- Polarized structure function
 - $\sigma_{LT'} \sim Im(L*T) \\ = Re(L)Im(T) + Im(L)Re(T)$
 - Amplify resonance multipole by a large background amplitude

K. Joo et al. (CLAS Collaboration), Phys. Rev. Lett. 88, 122001 (2002).

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K. Joo et al. (CLAS Collaboration), Phys. Rev. C 72, 058202 (2005) K. Joo et al. (CLAS Collaboration), Phys. Rev. C 70, 042201 (2004) K. Joo et al. (CLAS Collaboration), Phys. Rev. C 68, 032201 (2003)

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Power of Interference (sensitivity to $P_{11}(1440)$)

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 $\sigma_{LT'}$ structure function is sensitive to imaginary part of P₁₁(1440) through interference with real non resonant background.

K. Joo et al. (CLAS Collaboration), Phys. Rev. C 72, 058202 (2005)
K. Joo et al. (CLAS Collaboration), Phys. Rev. C 70, 042201 (2004)
K. Joo et al. (CLAS Collaboration), Phys. Rev. C 68, 032201 (2003)

Solving the Roper N(1440)1/2⁺ Puzzle

Importance of MB at Q² < 1.5GeV². Quark core contributions dominate at Q² > 2 GeV²

The 1st radial excitation of the 3-quark core seen when the probe penetrates the MB cloud.

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q^3 and MB contributions in N Δ (1232)

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Exclusive Electroproduction x-section and BSA

$$\frac{\operatorname{Cross\ section}}{2\pi \frac{d^2\sigma}{dtd\phi} = \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cdot \cos(2\phi) \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cdot \cos(\phi) \frac{d\sigma_{LT}}{dt} + h \cdot \sqrt{2\epsilon(1-\epsilon)} \cdot \sin(\phi) \frac{d\sigma_{LT'}}{dt}$$

$$\frac{\operatorname{Beam\ Spin\ Asymmetry:}}{BSA(t, \phi, x_B, Q^2) = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-} = \frac{\sqrt{2\epsilon(1-\epsilon)} \frac{\sigma_{LT'}}{\sigma_0} \sin\phi}{1 + \sqrt{2\epsilon(1+\epsilon)} \frac{\sigma_{LT}}{\sigma_0} \cos\phi + \epsilon \frac{\sigma_{TT}}{\sigma_0} \cos 2\phi}$$

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Multipoles, CGLN Amplitudes and Structure Functions

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G. Knöchlein, D. Drechsel, L. Tiator, Z.Phys.A352:327-343,1995; arXiv:nucl-th/9506029v1

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$$\label{eq:relation} \begin{split} {}^{c}\!R_{TL}^{00} &= \sin \Theta \mathrm{Re} \left\{ -F_{2}^{*}F_{5} - F_{3}^{*}F_{5} - F_{1}^{*}F_{6} - F_{4}^{*}F_{6} \right. \\ &- \cos \Theta \left(F_{4}^{*}F_{5} + F_{3}^{*}F_{6} \right) \right\}, \\ {}^{s}\!R_{TL'}^{00} &= -\sin \Theta \mathrm{Im} \left\{ F_{2}^{*}F_{5} + F_{3}^{*}F_{5} + F_{1}^{*}F_{6} + F_{4}^{*}F_{6} \right. \\ &+ \cos \Theta \left(F_{4}^{*}F_{5} + F_{3}^{*}F_{6} \right) \right\}, \end{split}$$

$$\begin{split} E_{l+} &= \frac{1}{l+1} A_{l+} + \frac{l}{2(l+1)} B_{l+} ,\\ E_{l-} &= -\frac{1}{l} A_{l-} + \frac{l+1}{2l} B_{l-} ,\\ M_{l+} &= \frac{1}{l+1} A_{l+} - \frac{l+2}{2(l+1)} B_{l+} ,\\ M_{l-} &= \frac{1}{l} A_{l-} + \frac{l-1}{2l} B_{l-} ,\\ S_{l+} &= \frac{1}{l+1} C_{l+} ,\\ S_{l-} &= -\frac{1}{l} C_{l-} .\\ \end{split}$$

$$\begin{split} A_{l\pm} &= \pm K A_{1/2}^{N},\\ B_{l\pm} &= \mp K \sqrt{\frac{4}{l(l+2)}} A_{3/2}^{N},\\ C_{l\pm} &= \pm K C_{1/2}^{N}.\\ \end{split}$$

$$K &= \sqrt{\frac{k_{\gamma}^{cm}}{|\vec{k}|} \frac{m}{W} \frac{\Gamma_{\eta}}{\pi (2J+1)}} \frac{M^{*}}{M^{*2} - W^{2} - iW\Gamma}$$

Data Analysis Frame Work

Inclusive Data Analysis from 10.6 GeV with CLAS12

W yield in Q² bins, Normalization to integral

Orange – Data

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• Black Dashed Line – MC no smearing

Inclusive Cross Sectioins from 10.6 GeV with CLAS12

- Preliminary CLAS12 measurements.
- CLAS data (after interpolation into the grid of our experiment), Phys. Rev. D67, 092001 (2003).

CLAS12 N* Program

Extend the Q2 to a higher range.

Study the hyper-baryon program.

 Map out the transition of strong QCD to pertubative QCD

Extend to Transition GPDs

ep->enπ⁺

VS.

ep->eγnπ⁺

Electron scattering binning scheme

Resonance Region DIS Region Q², W Q^2 , X_B **Inclusive Scattering** Exclusive Process (γ , π , ρ , ϕ , ...) Q², W, $\cos\theta^*$, ϕ^* Q², X_B, -t, ϕ^* **Off-diagonal DVCS or DVMP** Q², x_{B} , -t, ϕ^* , $M_{\pi N}$, $\cos\theta_{cm}$, ϕ_{cm}

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Summary

- First high precision photo- and electroproduction data have become available and led to a new wave of significant developments in N* program.
- Large amounts of high precision data and multi-channel PWA have been essential in the discovery of new N* resonances.
- Electroexcitation of nucleon resonances are sensitive to the effective degrees of freedom versus distance scale.

Transition Form Factors (N* Physics) at 6 GeV JLab Era Transition GPDs (3D N* Physics) at 12-22 GeV JLab Era

Establishing the N* spectrum

Photon			Target			Recoil		Target + Recoil				
	_	_	_	_	x'	y'	z'	x'	x'	z'	z'	
	_	\boldsymbol{x}	\boldsymbol{y}	z	_	_	_	x	z	x	z	
transverse unpolarized (T)	σ_0	0	T	0	0	P	0	$T_{x'}$	$-L_{x'}$	$T_{z'}$	$L_{z'}$	
transverse (linear pol.) (TT)	$-\Sigma$	H	(-P)	-G	$O_{x'}$	(-T)	$O_{z'}$	$(-L_{z'})$	$(T_{z'})$	$(-L_{x'})$	$(-T_{x'})$	
transverse (circular pol.) (TT')	0	F	0	-E	$-C_{x'}$	0	$-C_{z'}$	0	0	0	0	

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