

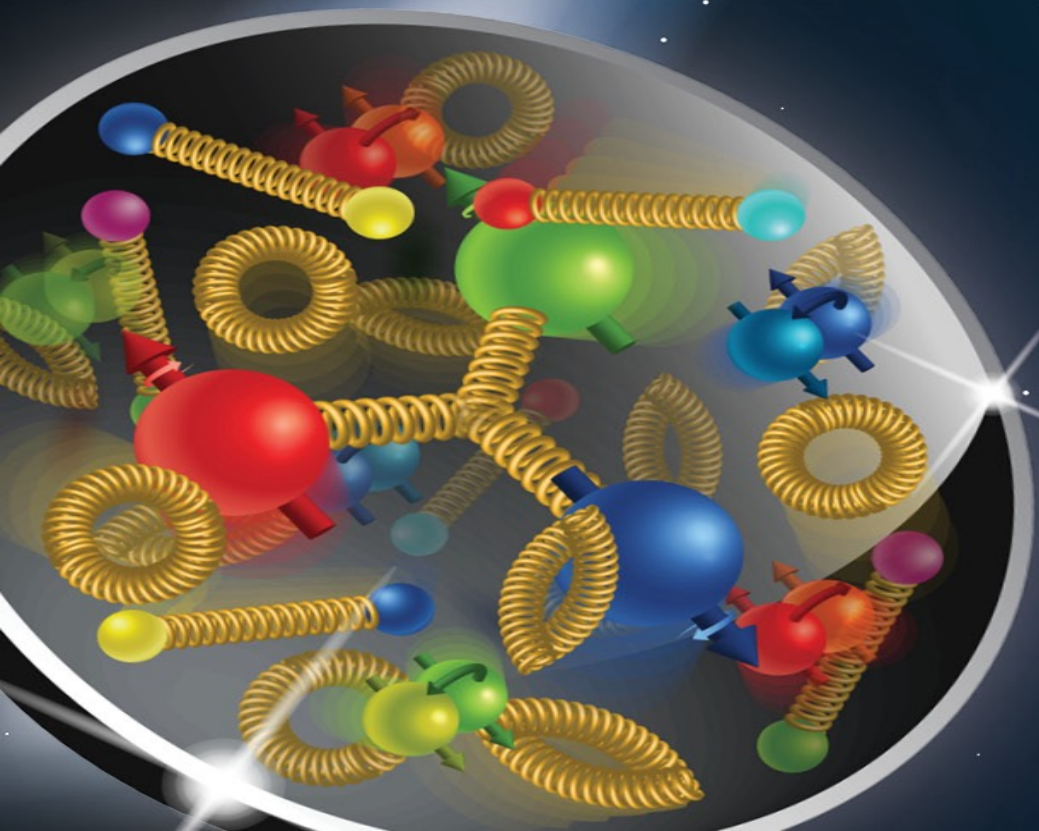
(some) Exclusive Physics Opportunities with Far-Forward Detectors at the EIC

Alex Jentsch, *Brookhaven National Lab*
ajentsch@bnl.gov

ECT*-APCTP Joint Workshop: Exploring resonance structure with transition GPDs

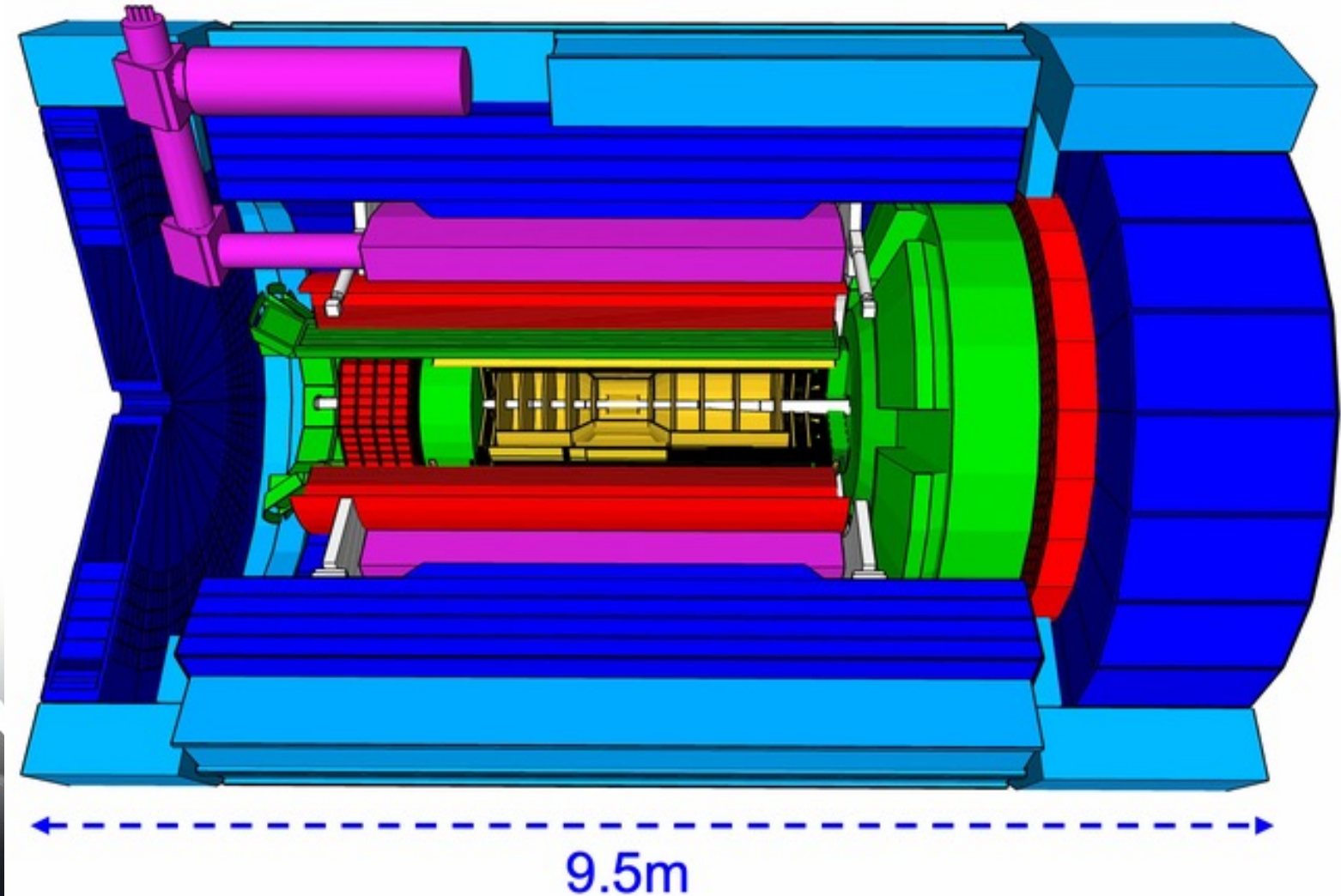
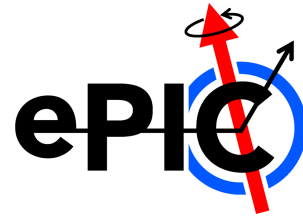
August 21st- 25th, 2023

ECT*: Trento, Italy

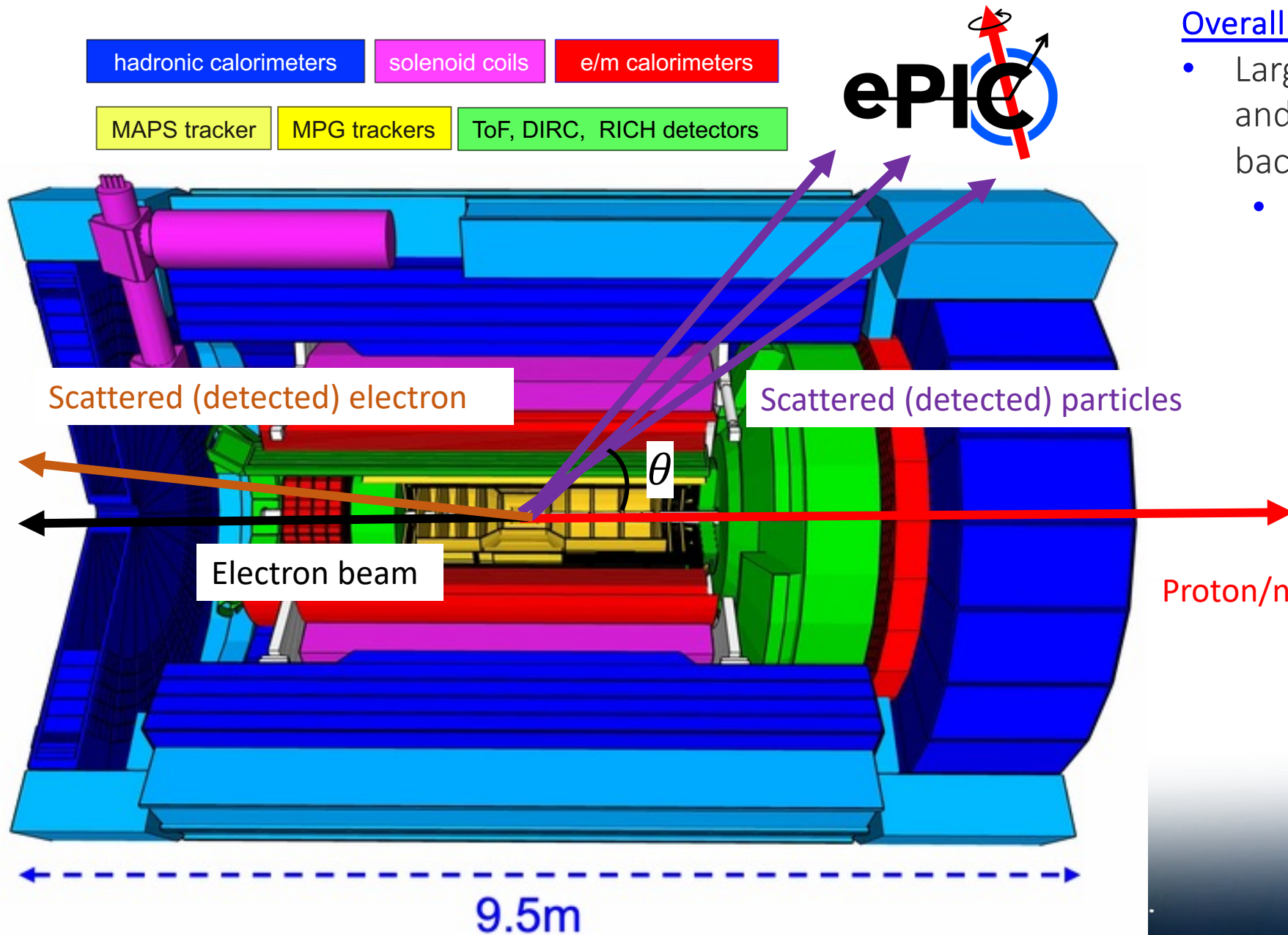


What is meant by Far-Forward?

- hadronic calorimeters
- solenoid coils
- e/m calorimeters
- MAPS tracker
- MPG trackers
- ToF, DIRC, RICH detectors



What is meant by Far-Forward?



Overall detector requirements:

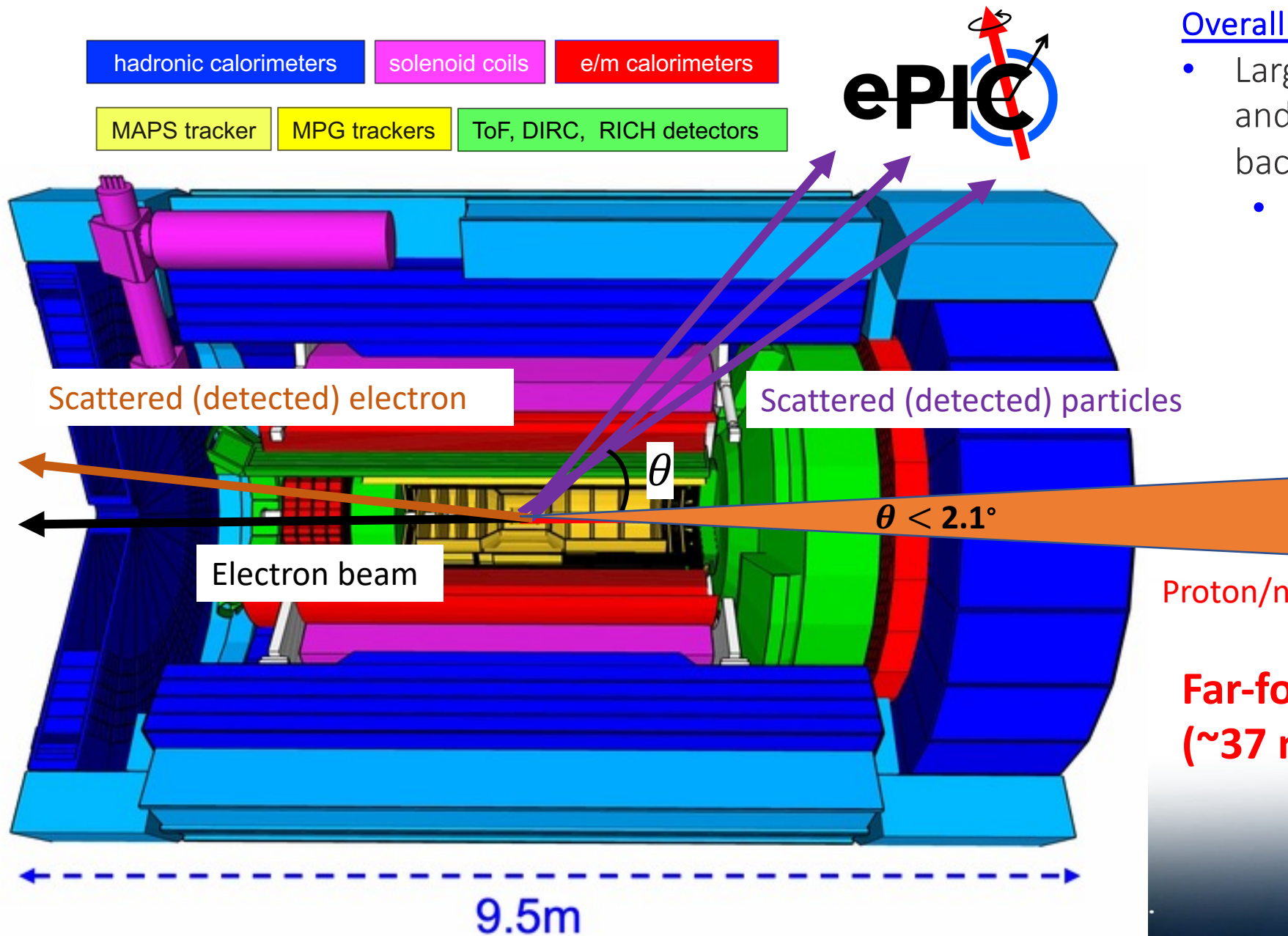
- Large rapidity ($-4 < \eta < 4$) coverage; and far beyond in far-forward/far-backward detector regions

- Rapidity is related to the polar angle $\rightarrow 0 < \eta < 4$ equates to $2.1^\circ < \theta < 90^\circ$

$$\eta = -\ln(\tan(\theta/2))$$

pseudorapidity

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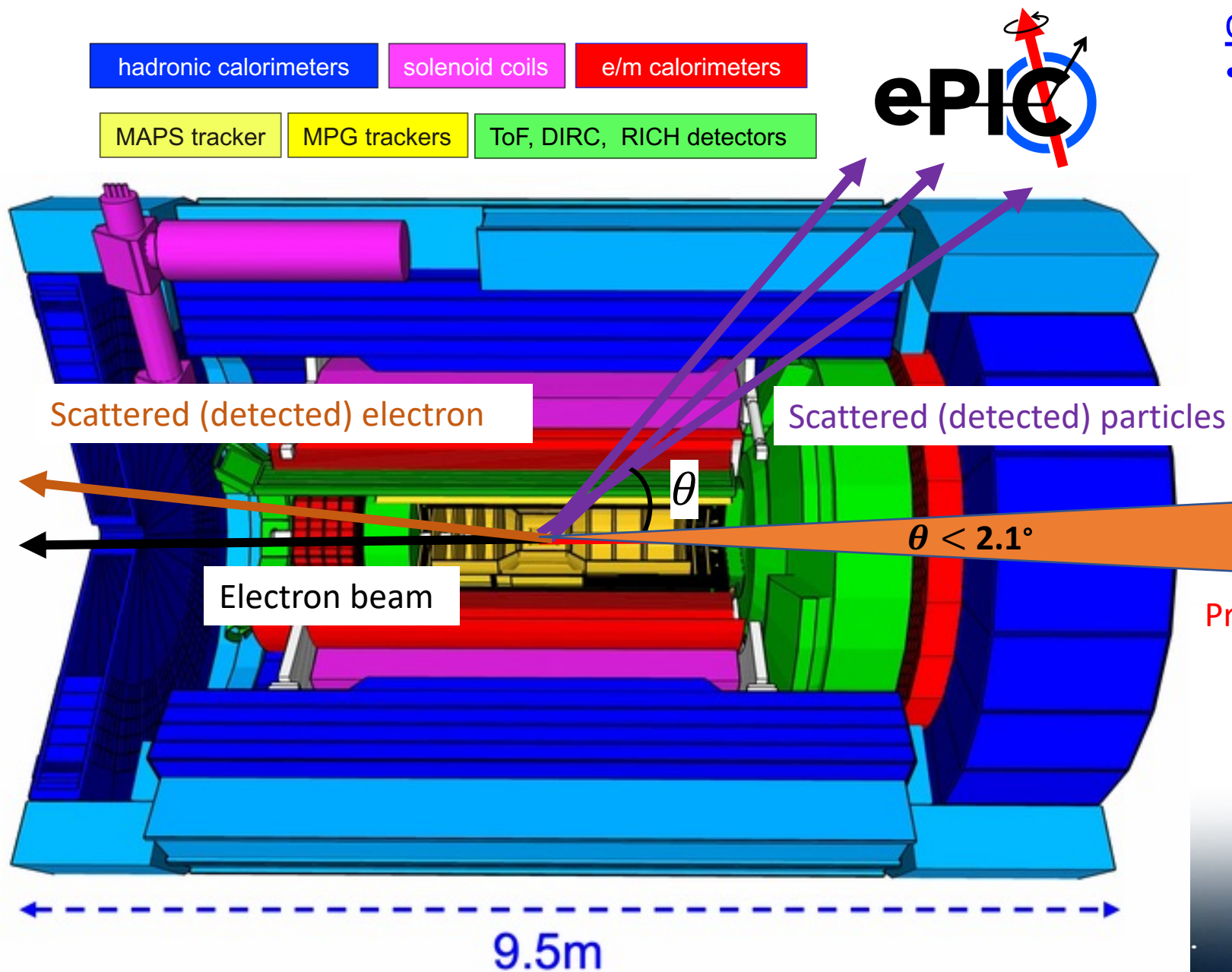
$$\eta = -\ln(\tan(\theta/2))$$

pseudorapidity

Proton/nucleus beam

Far-forward here means $\theta < 2.1^\circ$ (~37 mrad)

What is meant by Far-Forward?



Overall detector requirements:

- Large rapidity ($-4 < \eta < 4$) coverage; and far beyond in far-forward/far-backward detector regions

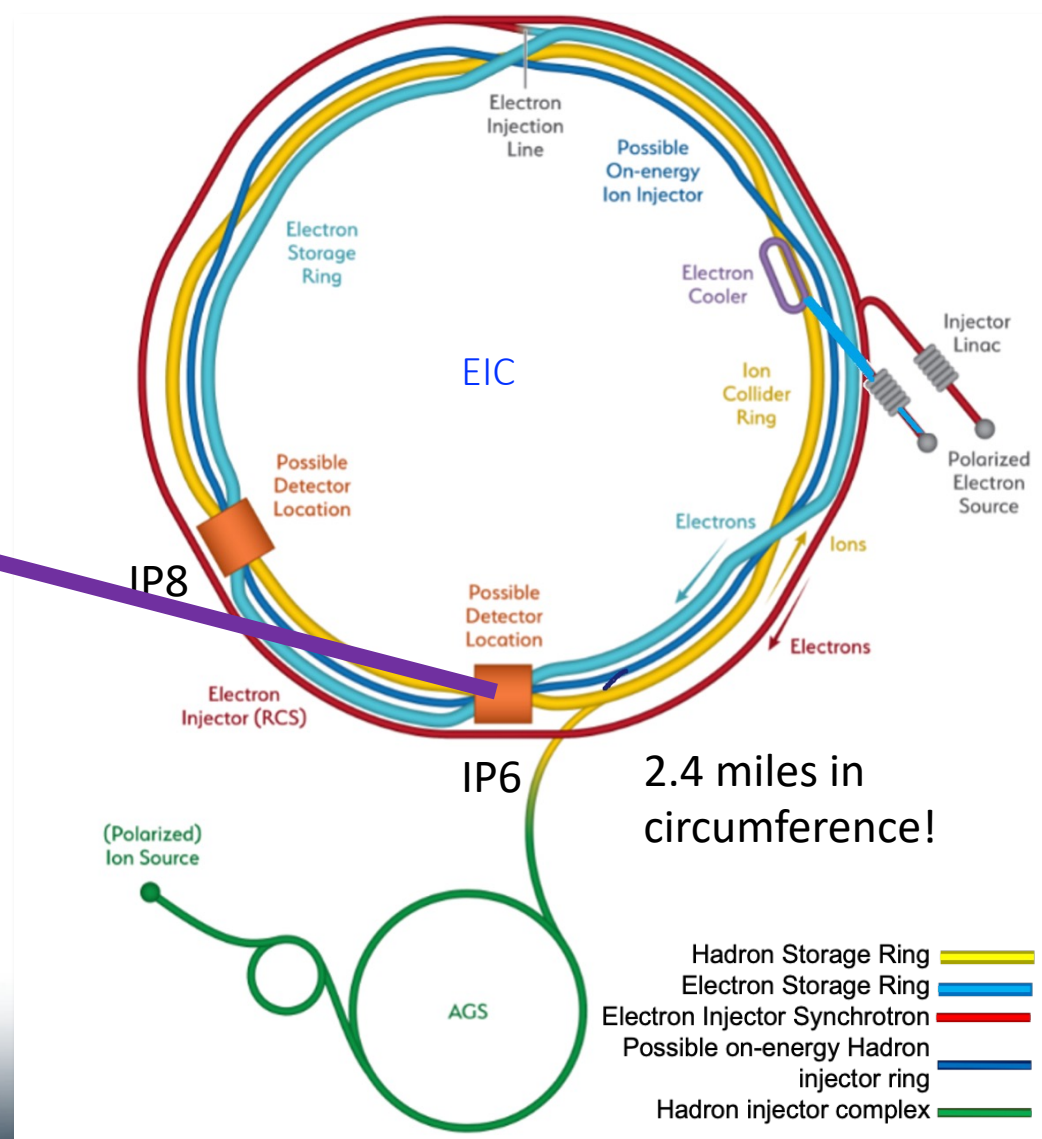
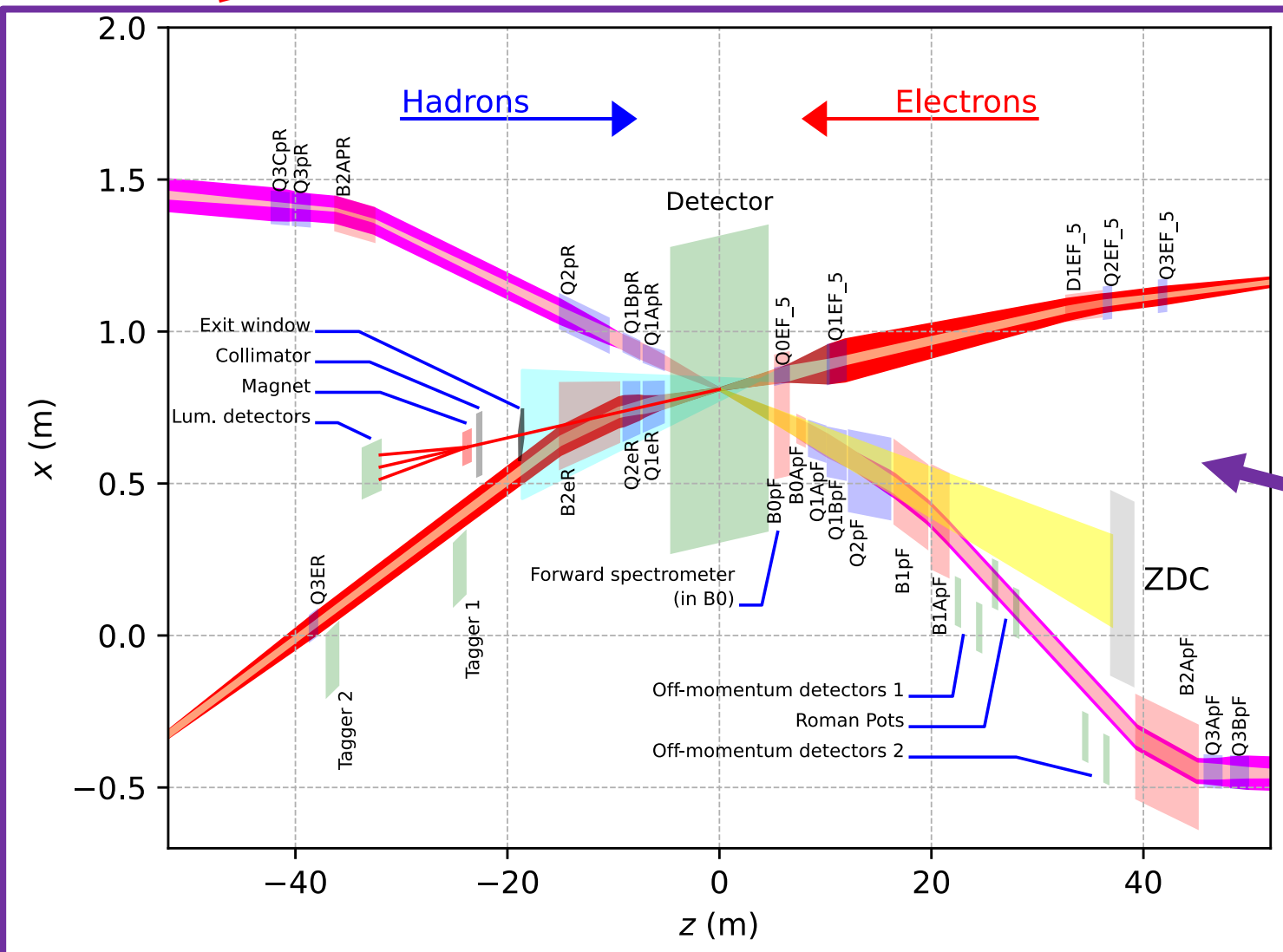
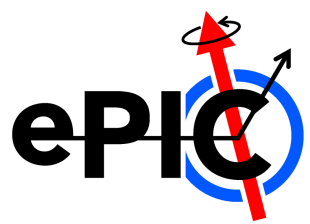
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pseudorapidity

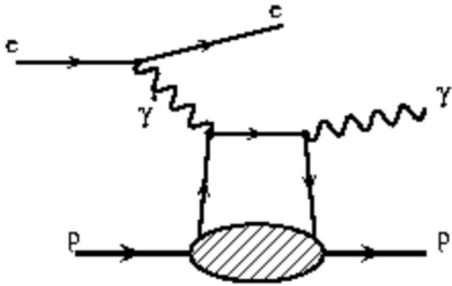


Need detectors here!!

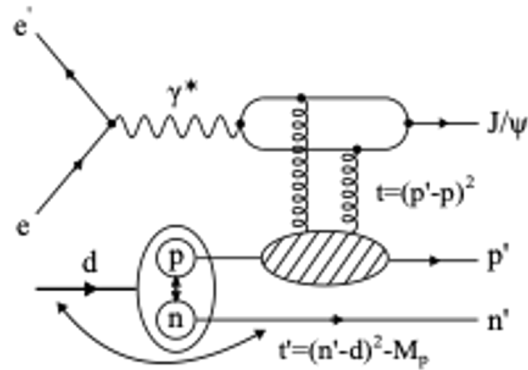


(some) Far-Forward Processes at the EIC

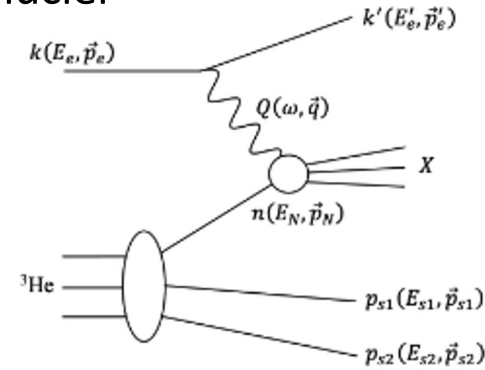
e+p DVCS



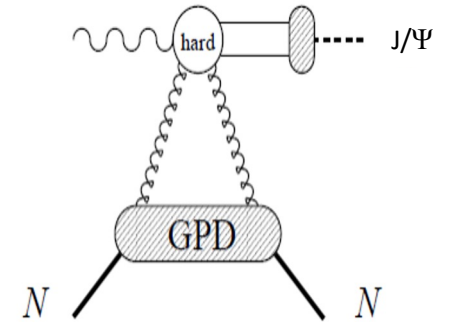
e+d exclusive J/Psi with p/n tagging



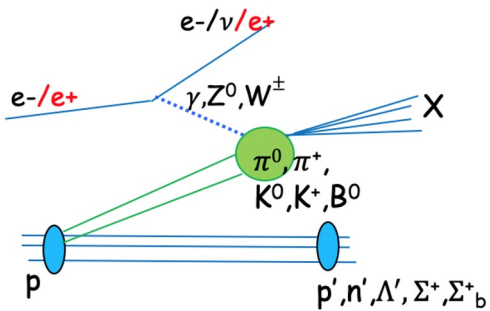
spectator tagging in light nuclei



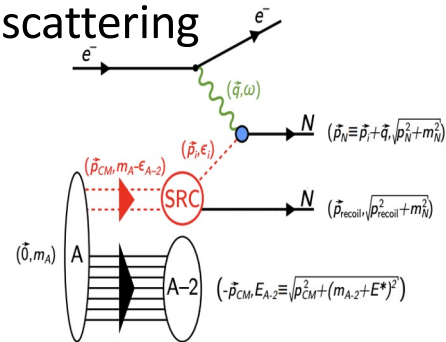
coherent/incoherent J/ψ production in e+A



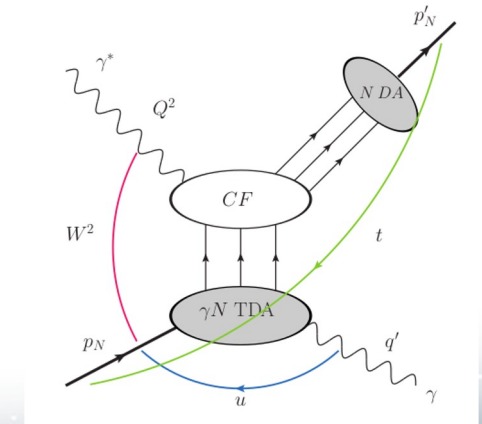
Sullivan process



Quasi-elastic electron scattering



u-channel backward exclusive electroproduction

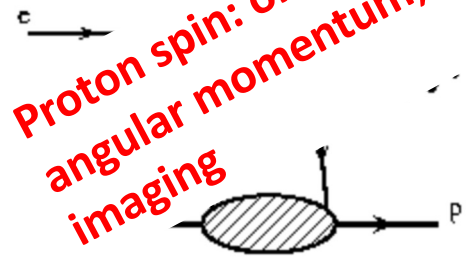


...and MANY more!

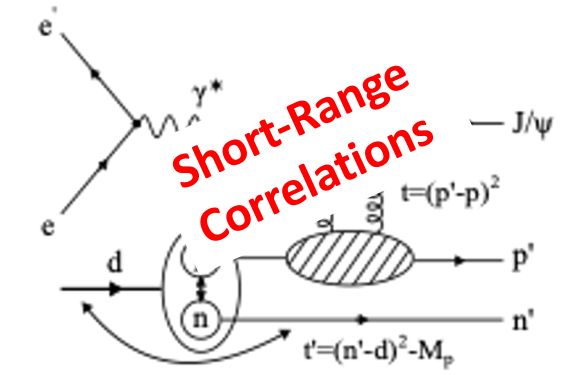
(some) Far-Forward **Physics** at the EIC

e+p DVCS

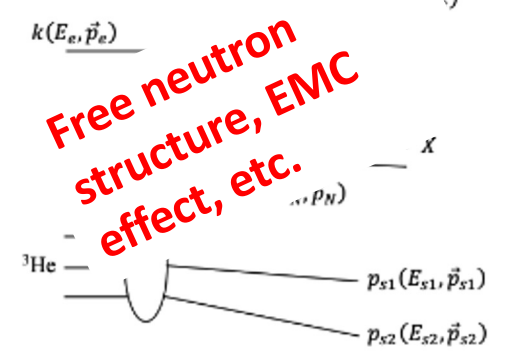
Proton spin: orbital angular momentum; imaging



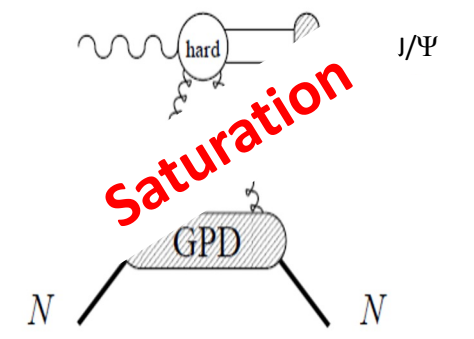
e+d exclusive J/Psi with p/n tagging



spectator tagging in light nuclei

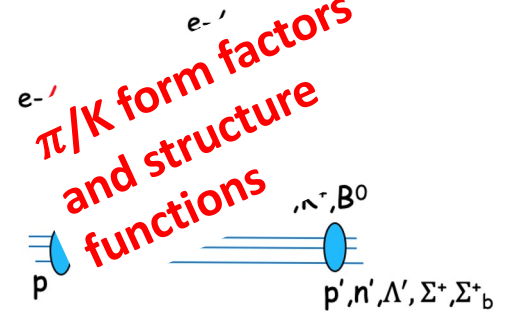


coherent/incoherent J/psi production in e+A

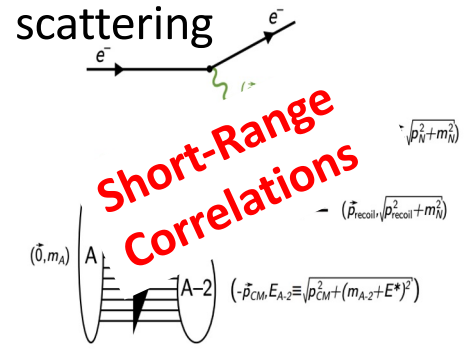


Sullivan process

pi/K form factors and structure functions

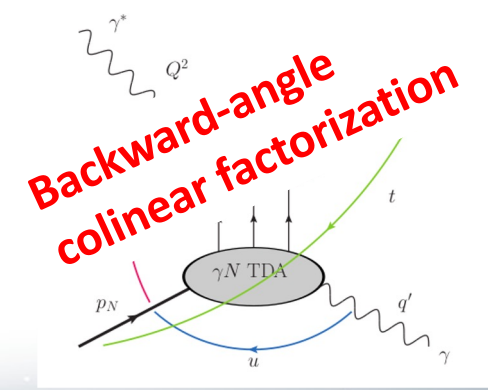


Quasi-elastic electron scattering



[1] Z. Tu, A. Jentsch, et al., Physics Letters B, (2020)
 [2] I. Friscic, D. Nguyen, J. R. Pybus, A. Jentsch, et al., Phys. Lett. B, **Volume 823**, 136726 (2021)
 [3] W. Chang, E.C. Aschenauer, M. D. Baker, A. Jentsch, J.H. Lee, Z. Tu, Z. Yin, and L.Zheng, Phys. Rev. D **104**, 114030 (2021)
 [4] A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104**, 065205, (2021) (**Editor's Suggestion**)

u-channel backward exclusive electroproduction

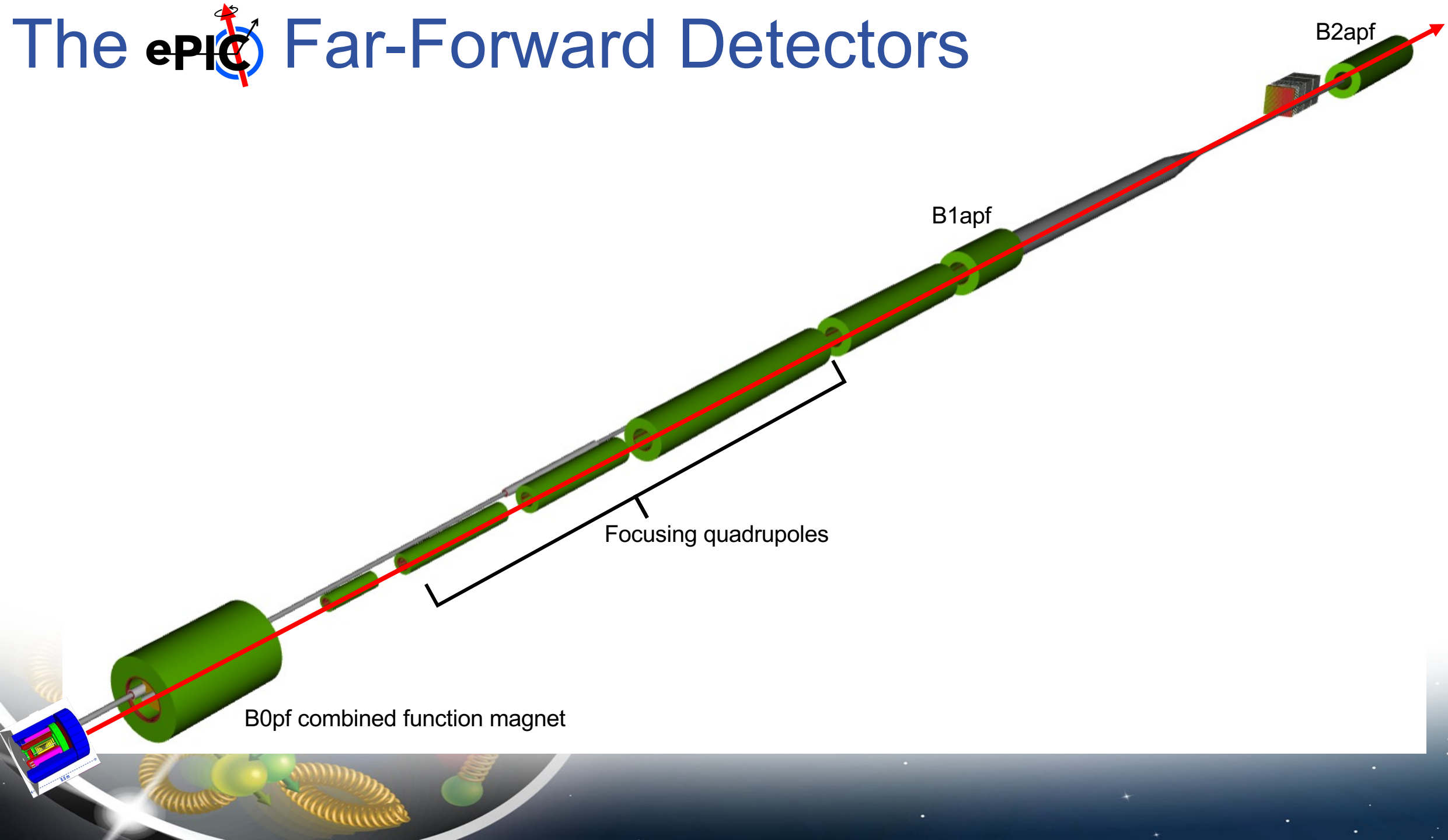


...and **MANY** more!

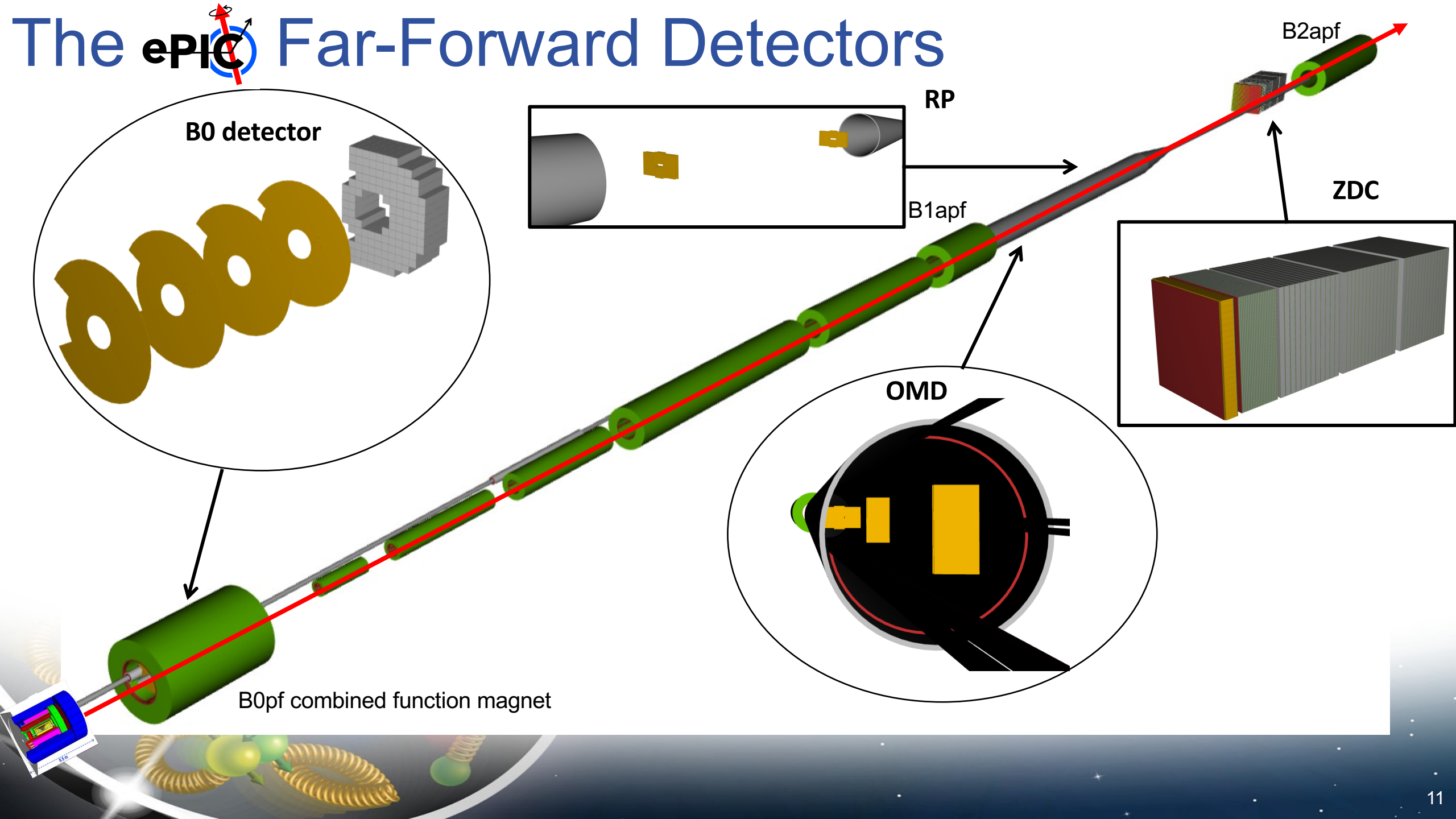
(some) Far-Forward **Physics** at the EIC

- Physics channels require tagging of **charged hadrons** (protons, pions) or **neutral particles** (neutrons, photons) at **very-forward rapidities** ($\eta > 4.5$).
- Different final states \rightarrow tailored detector subsystems.
- Various beams and energies (h: 41, 100-275 GeV, e: 5-18 GeV; e+p, e+d, e+Au, etc.).
- Placing and operation of far-forward detectors uniquely challenging due to integration with accelerator.

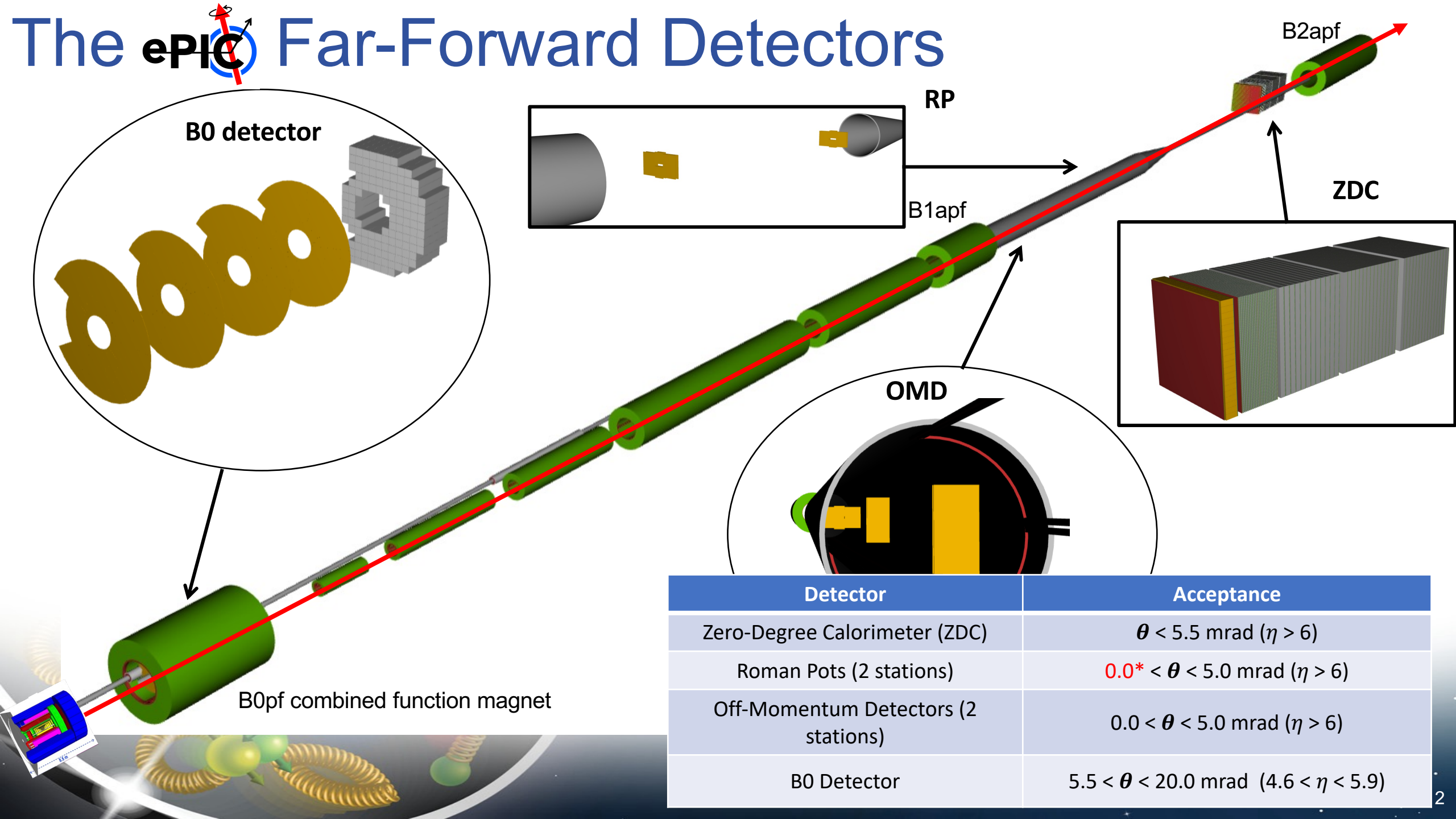
The ePIC Far-Forward Detectors



The **ePIC** Far-Forward Detectors

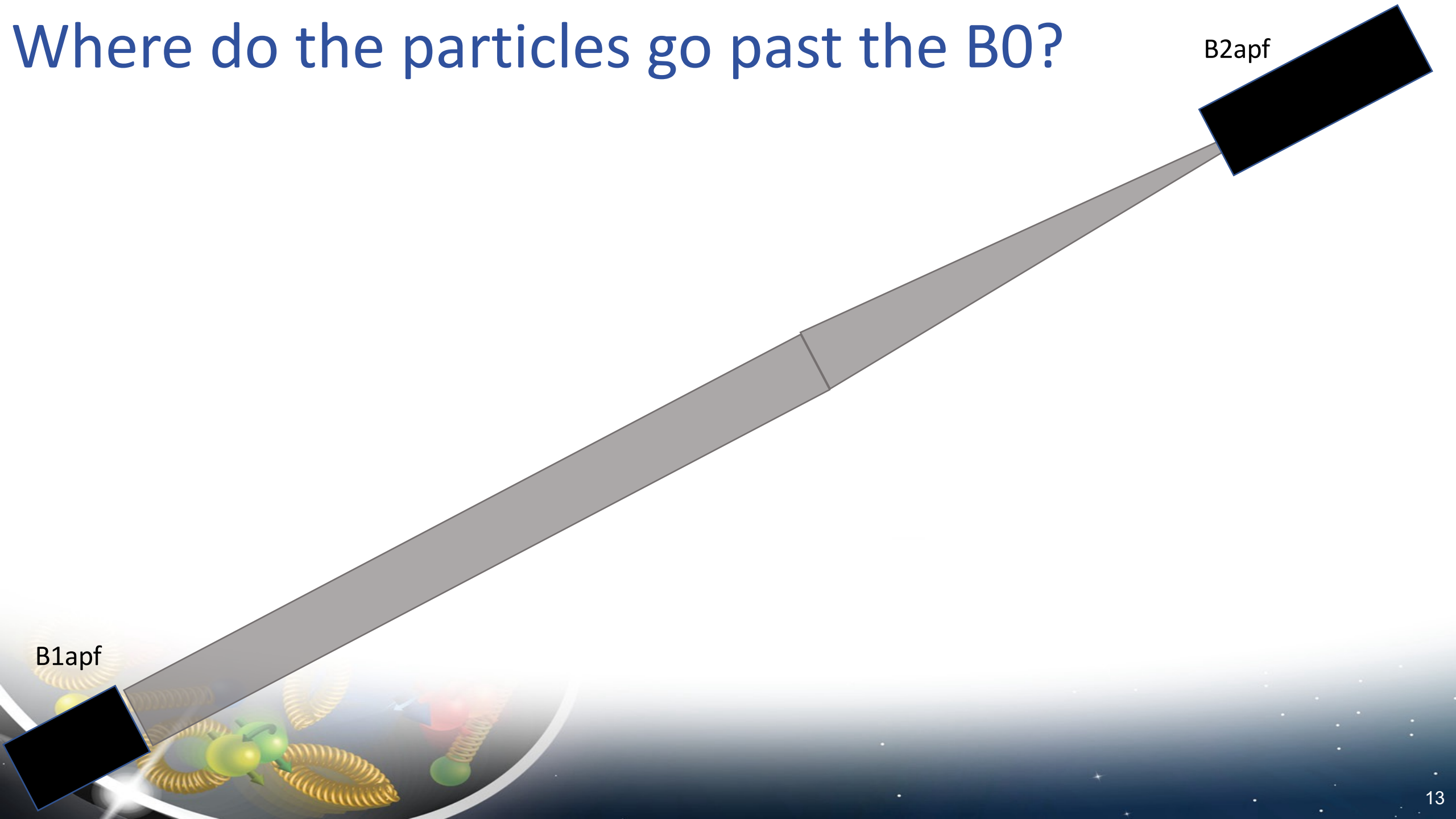


The ePIC Far-Forward Detectors



Detector	Acceptance
Zero-Degree Calorimeter (ZDC)	$\theta < 5.5 \text{ mrad}$ ($\eta > 6$)
Roman Pots (2 stations)	$0.0^* < \theta < 5.0 \text{ mrad}$ ($\eta > 6$)
Off-Momentum Detectors (2 stations)	$0.0 < \theta < 5.0 \text{ mrad}$ ($\eta > 6$)
B0 Detector	$5.5 < \theta < 20.0 \text{ mrad}$ ($4.6 < \eta < 5.9$)

Where do the particles go past the B0?

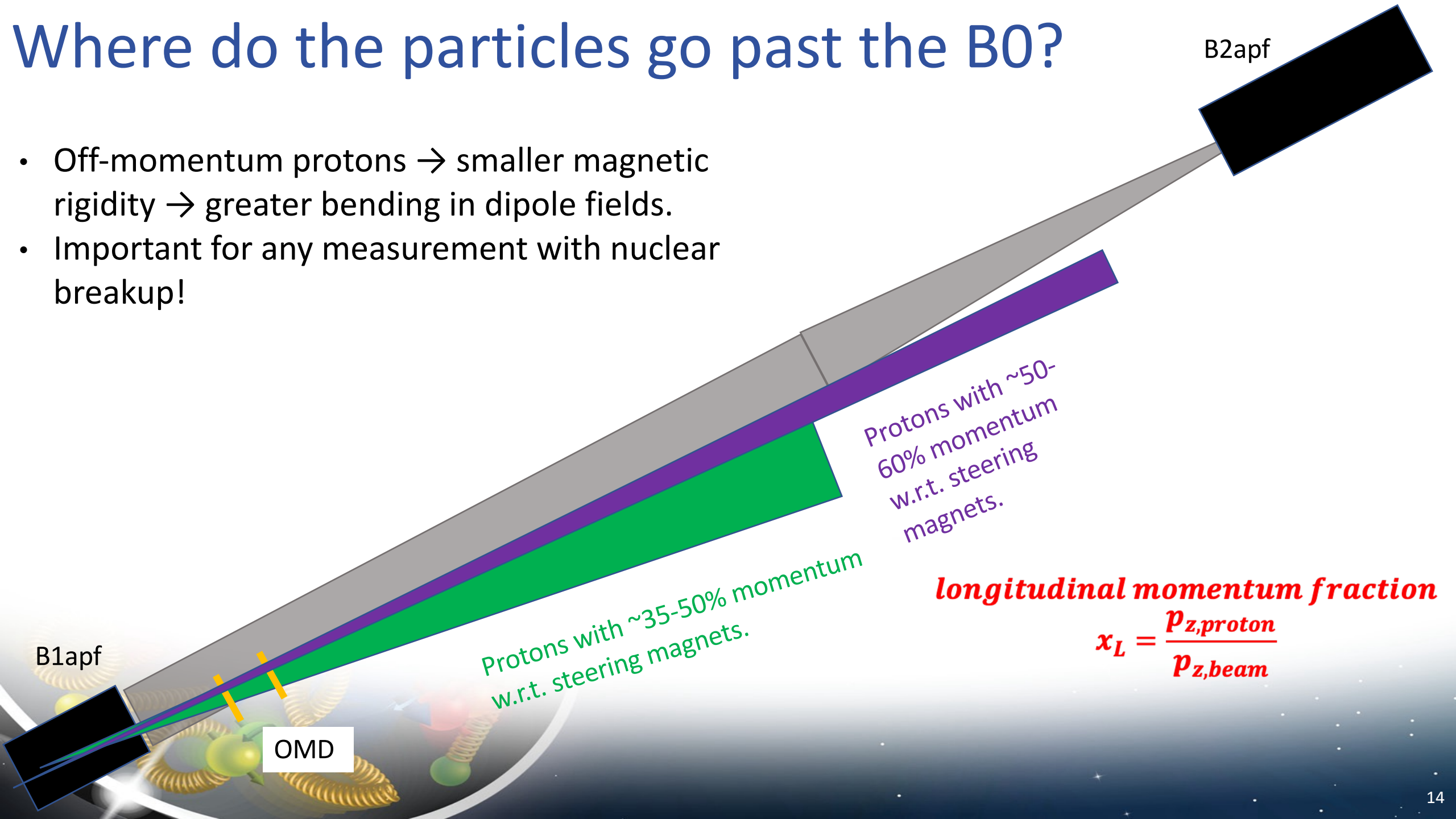


B1apf

B2apf

Where do the particles go past the B0?

- Off-momentum protons → smaller magnetic rigidity → greater bending in dipole fields.
- Important for any measurement with nuclear breakup!

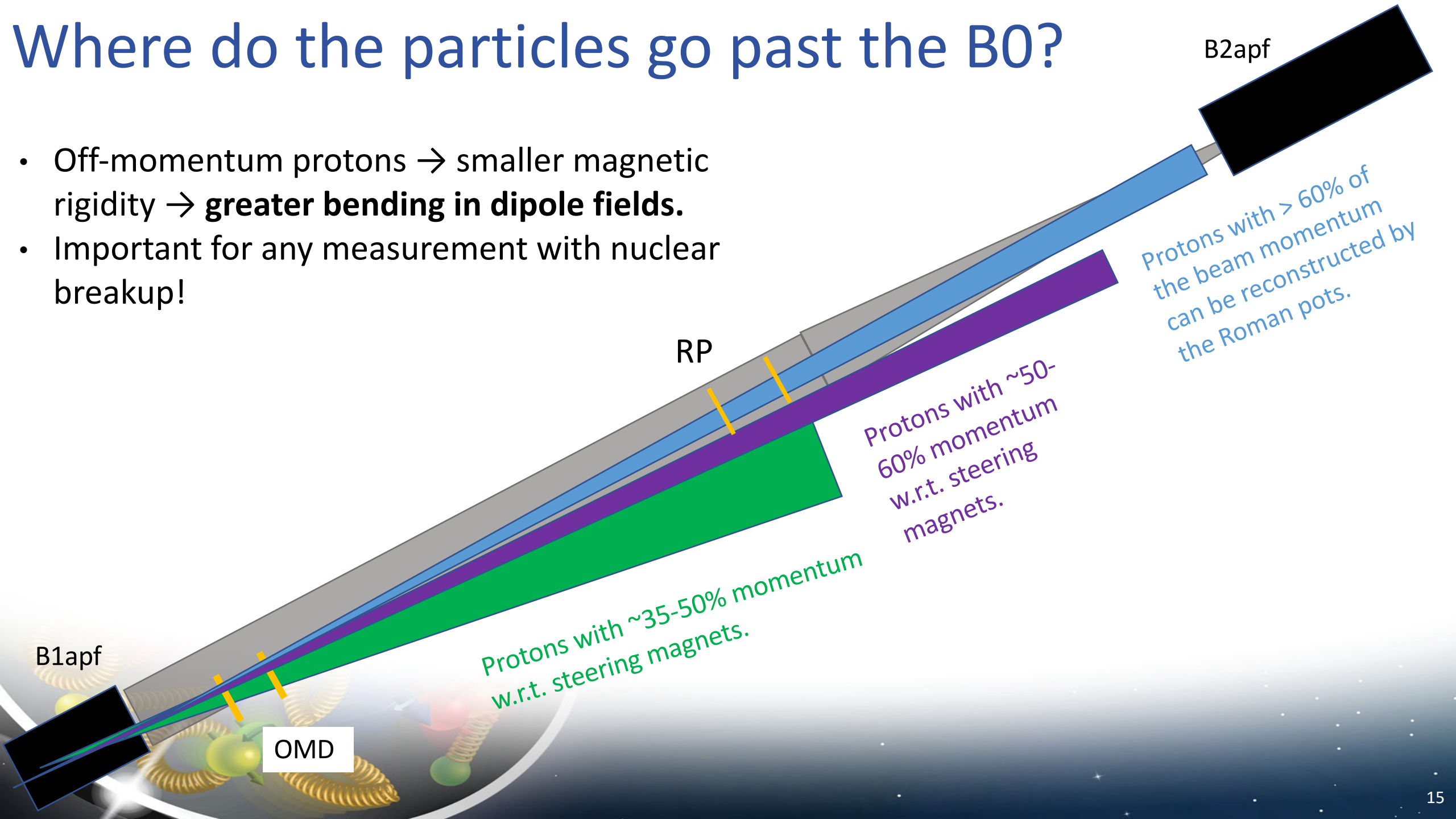


longitudinal momentum fraction

$$x_L = \frac{p_{z,\text{proton}}}{p_{z,\text{beam}}}$$

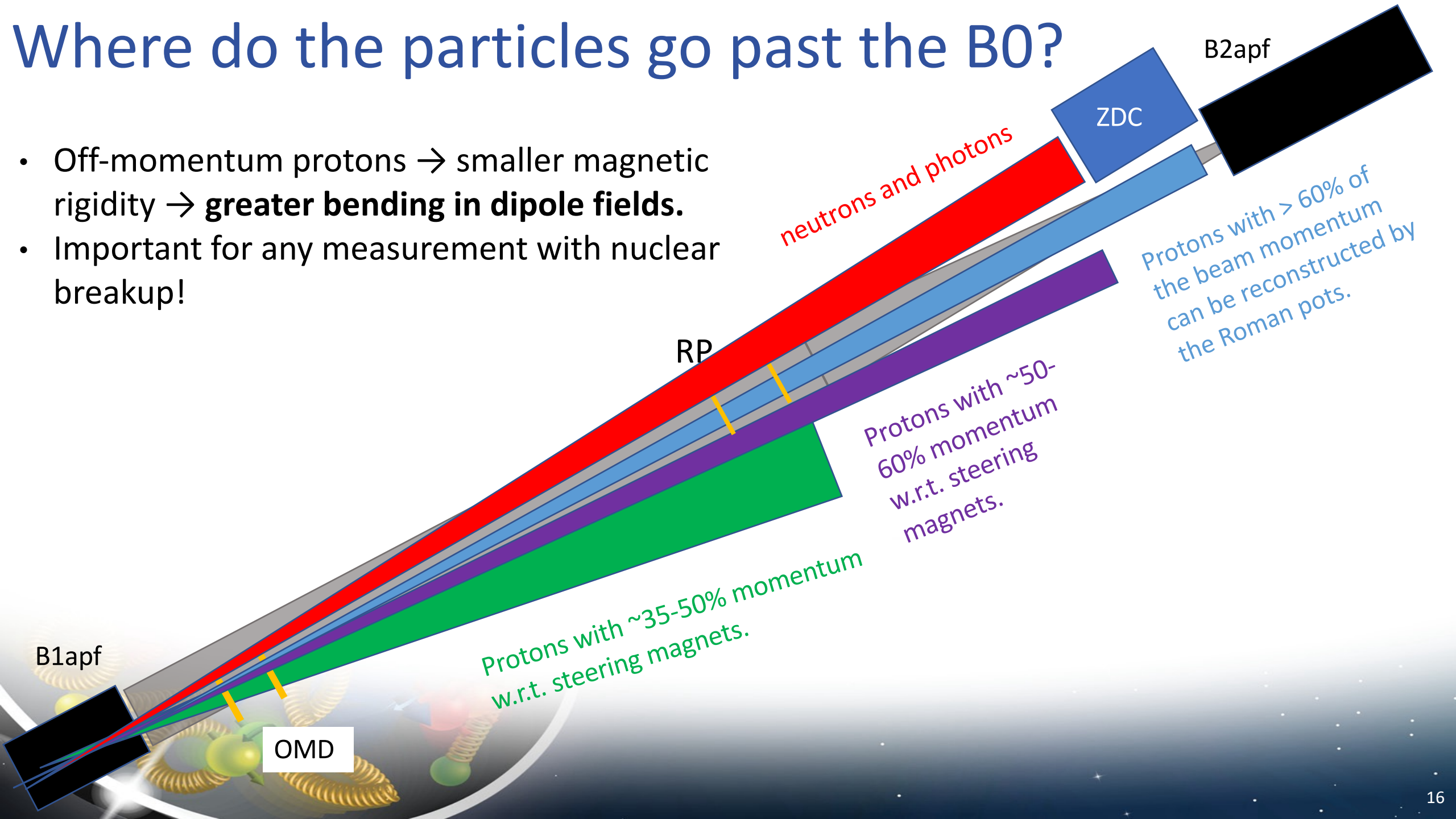
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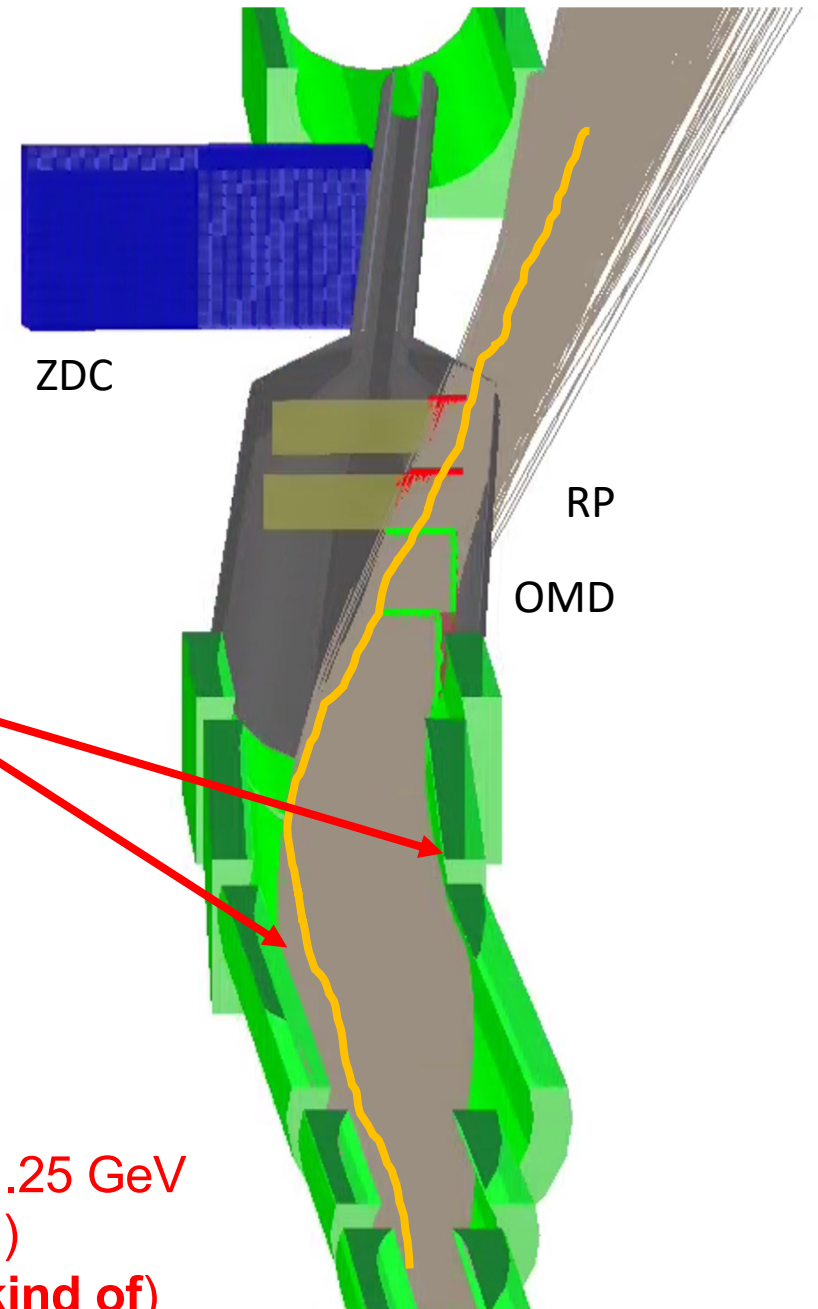
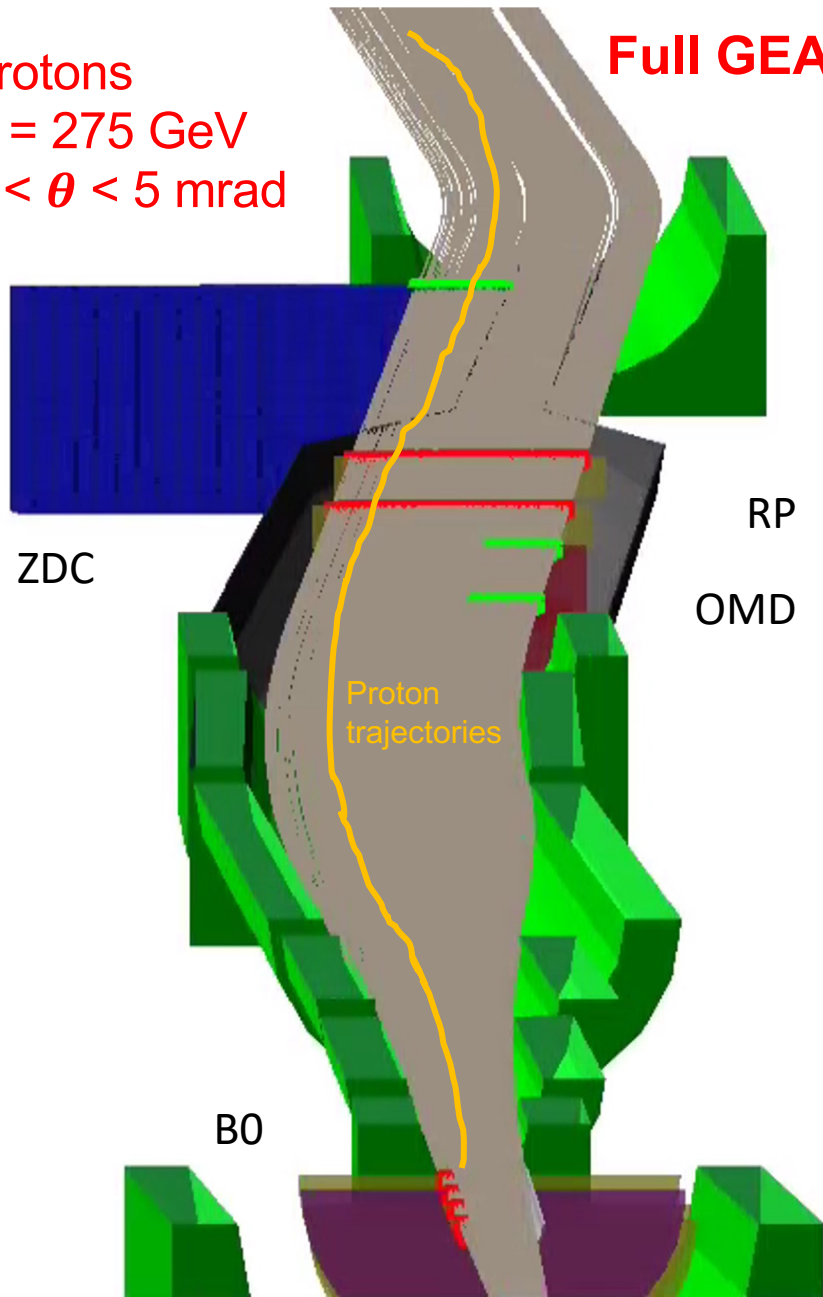
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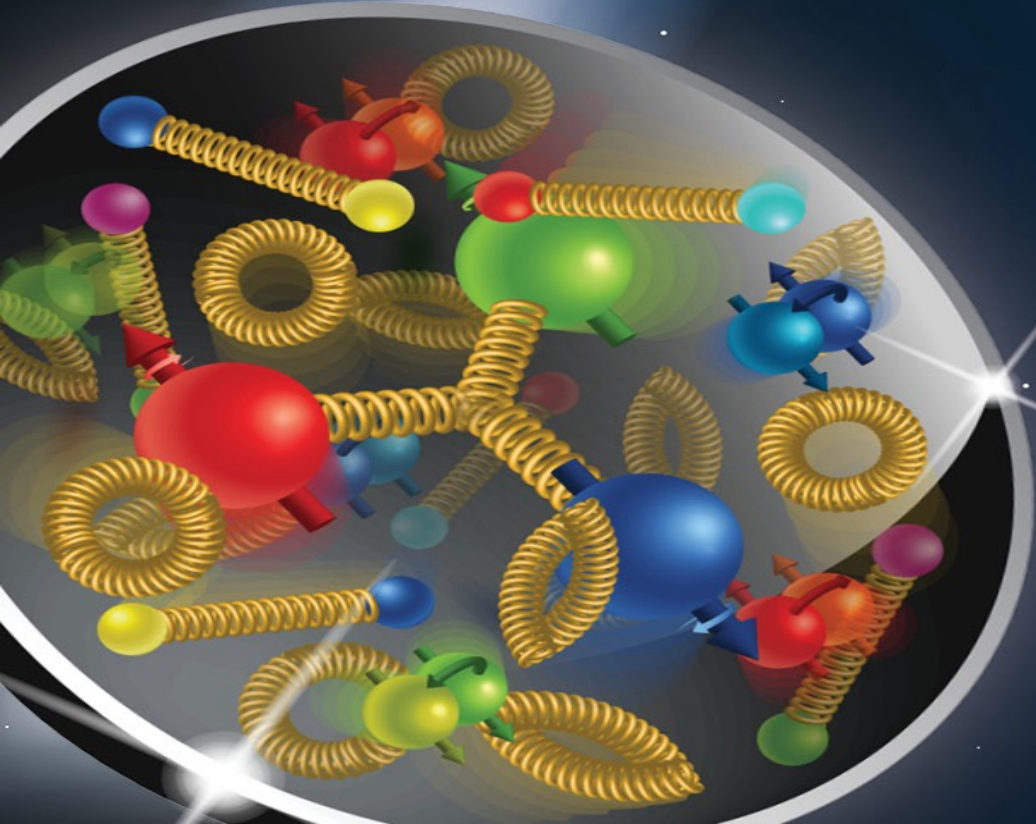
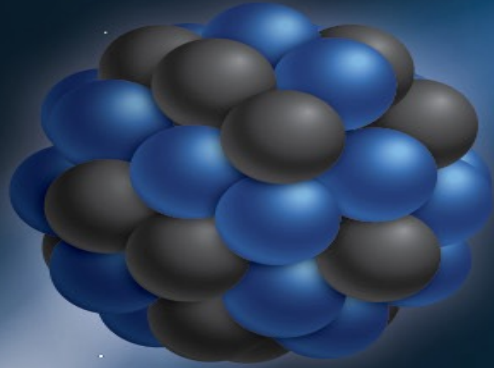
Roman Pots and OMD

Protons
 $E = 275 \text{ GeV}$
 $0 < \theta < 5 \text{ mrad}$

Full GEANT4 simulation.



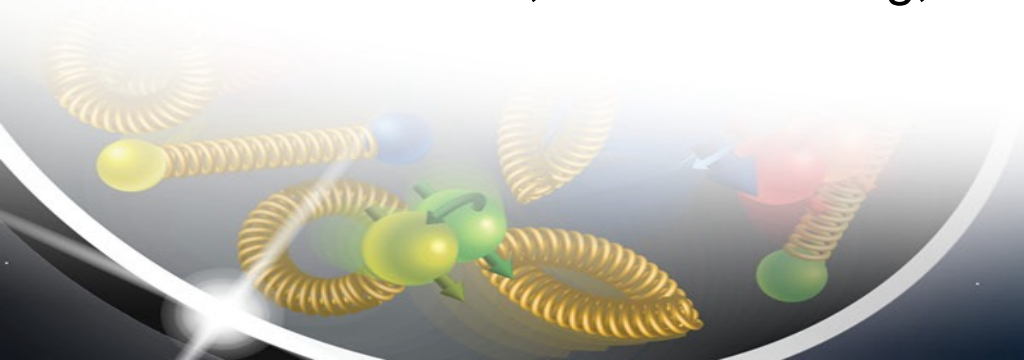
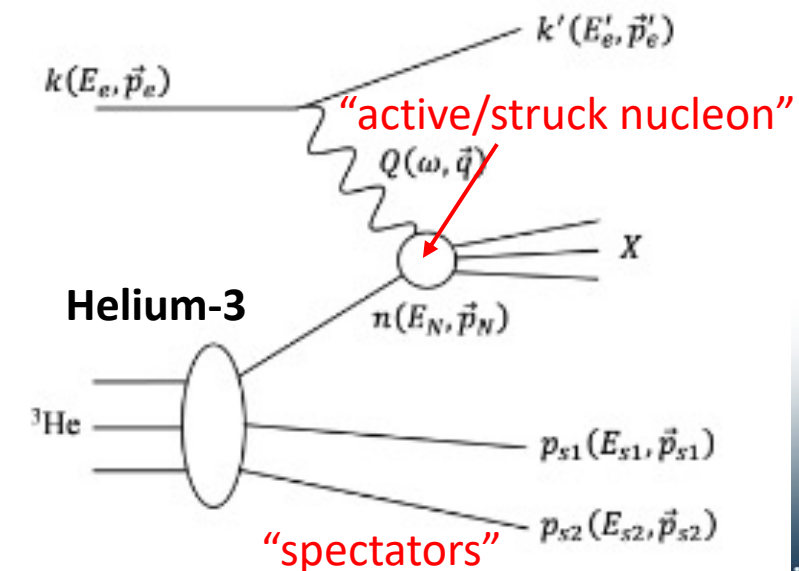
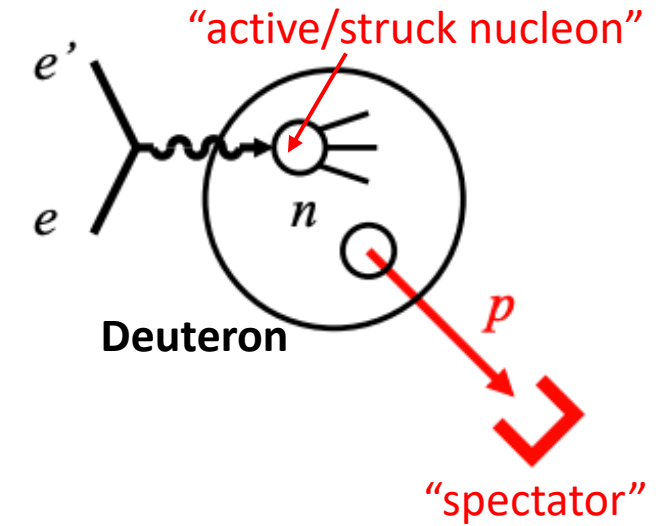
Protons
 $123.75 < E < 151.25 \text{ GeV}$
 $(45\% < x_L < 55\%)$
 $0 < \theta < 5 \text{ mrad (kind of)}$



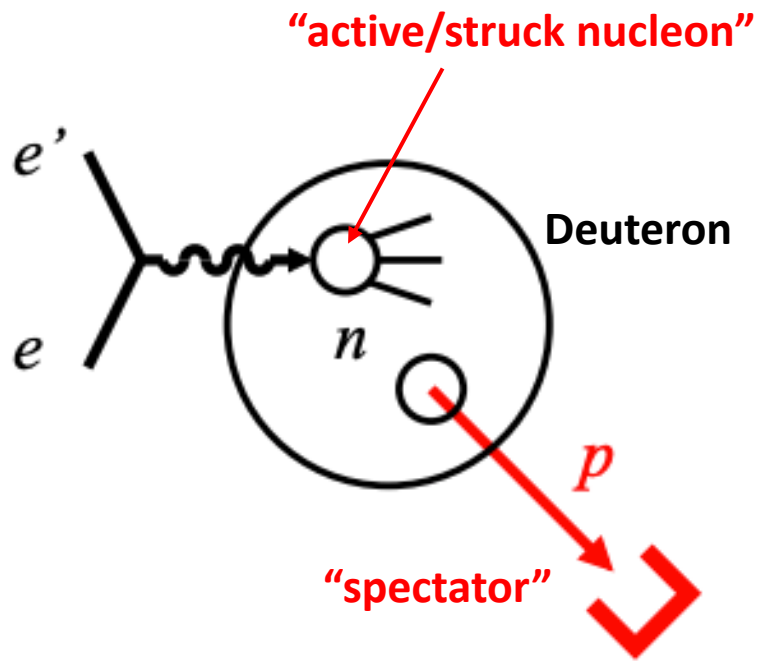
My Interest: Deuterons as an unexpected QCD laboratory at the EIC.

Deuteron tagged DIS as a tool at the EIC

- **Tagged DIS** measurements on light nuclei → "tag" (generally) far-forward particles in final state for useful kinematic information!
 - Provides more information than inclusive cross sections!
- Lots of topics!
 - Short-range correlations.
 - Gluon distributions in nuclei.
 - Free neutron structure functions.
 - Nuclear modifications of nucleons in light nuclei.
 - EMC effect, anti-shadowing, etc.



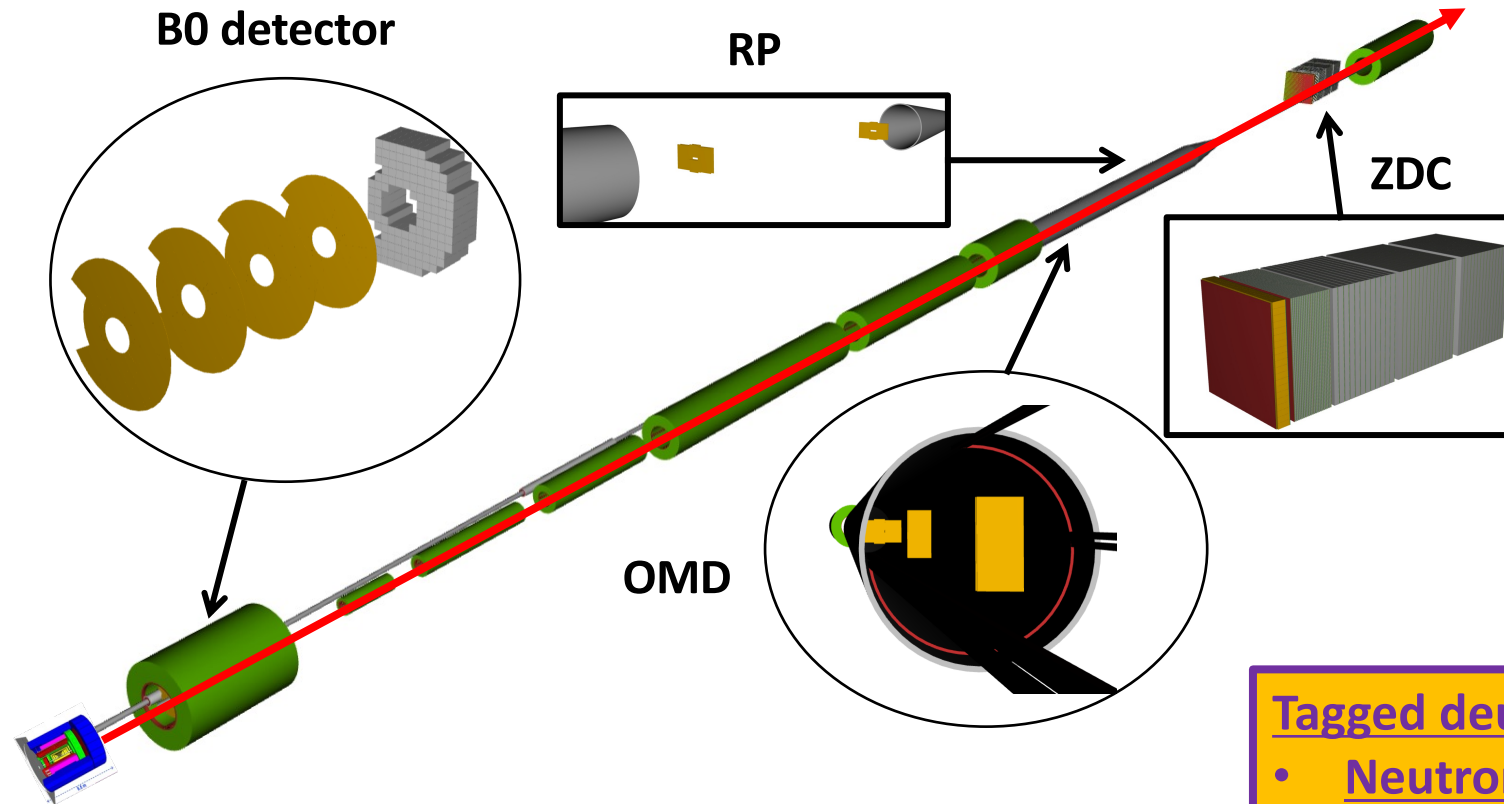
Tagged DIS with deuterons



- Spectator kinematics \rightarrow determines nuclear configuration.
 - Loosely bound configuration – enables extraction of free nucleon structure via pole extrapolation (previous study²).
 - Configuration with strongly-interacting nucleons – opens up study of nuclear modifications.
 - Differential study of transition region where nuclear effects manifest!

Tagged DIS on the deuteron enables study of free and modified nuclear structure in a single nucleus!

Full Detector Simulations – Tagged Spectators



Tagged deuteron spectators

- Neutrons: reconstructed in ZDC ($\theta < 5$ mrad acceptance).
- Protons: reconstructed in B0 tracker ($6 < \theta < 20$ mrad) and off-momentum detectors ($\theta < 5$ mrad).

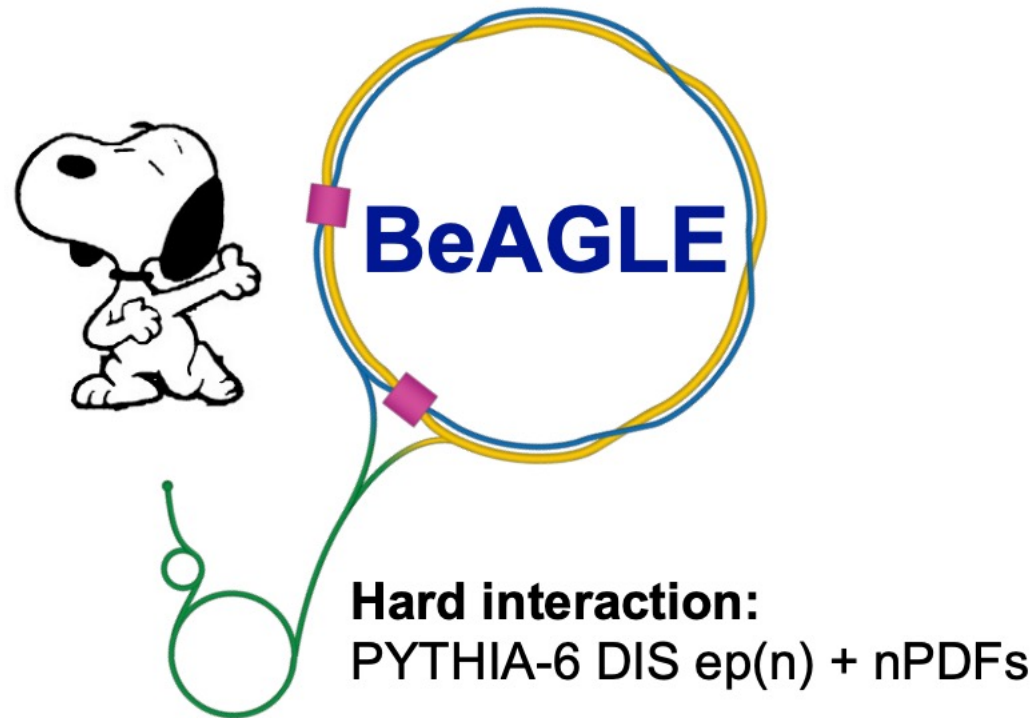


Deuterons: Gluons and Short-Range Correlations

Monte Carlo for all e+d studies presented here

General-purpose eA DIS MC generator

<https://eic.github.io/software/beagle.html>

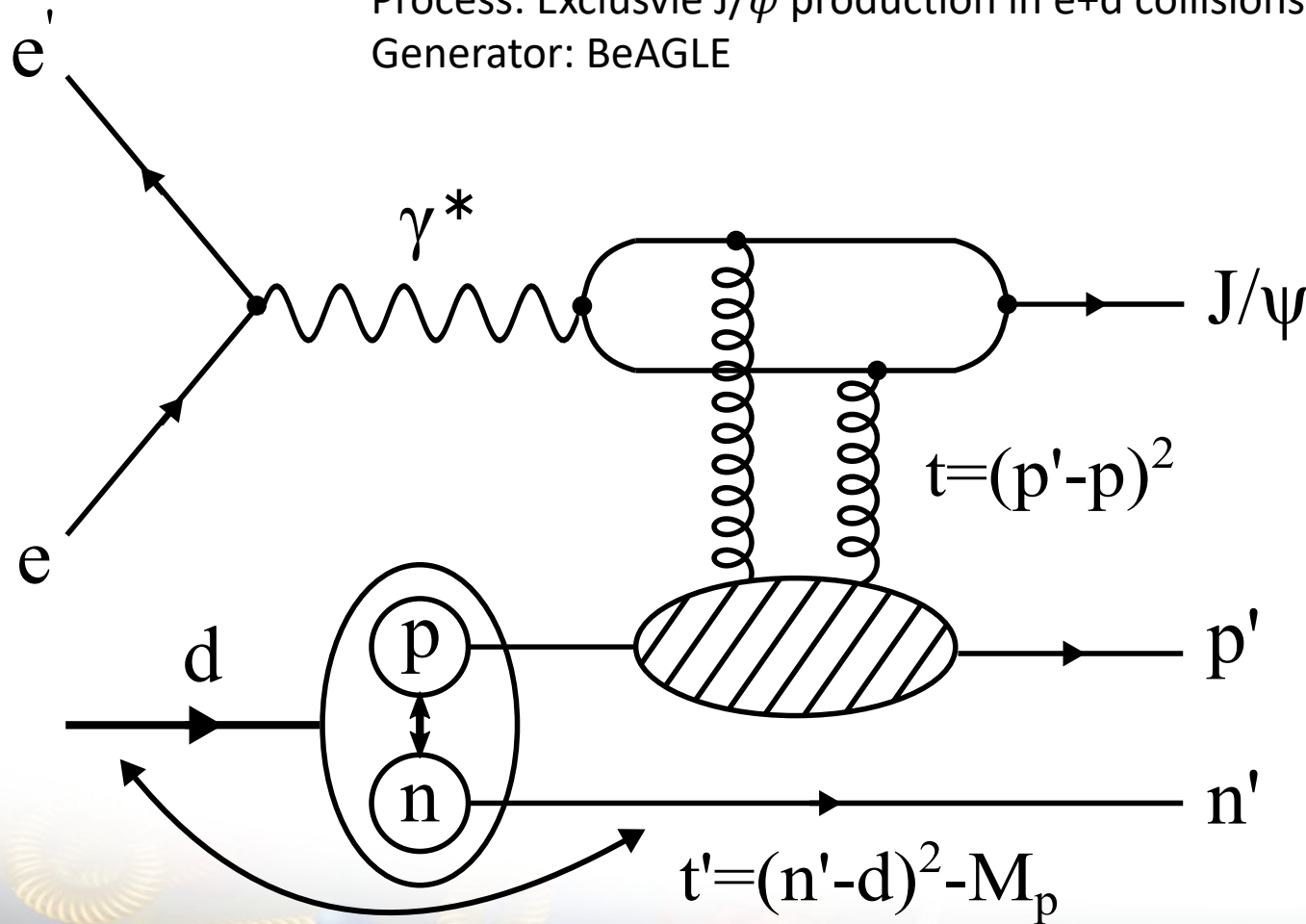


- Use BeAGLE to simulate the hard $e +$ (active) nucleon scattering and primary process (e.g. J/ψ production, DIS, etc.)
 - **For heavy A:** DPMJET and FLUKA
 - **For deuteron:** Spectator momentum spectra calculated via deuteron spectral function, using parametrization of Ciofi and Simula.
 - C. Ciofi degli Atti and S. Simula, Phys. Rev. C 53, 1689 (1996)
- BeAGLE MC samples passed through full detector simulations, including beam effects to study prospects for future analysis!

Wan Chang, Elke-Caroline Aschenauer, Mark D. Baker, Alexander Jentsch, Jeong-Hun Lee, Zhoudunming Tu, Zhongbao Yin, and Liang Zheng
Phys. Rev. D **106**, 012007 (2022)

Short-Range Correlations in Deuterons

Process: Exclusive J/ψ production in $e+d$ collisions.
Generator: BeAGLE

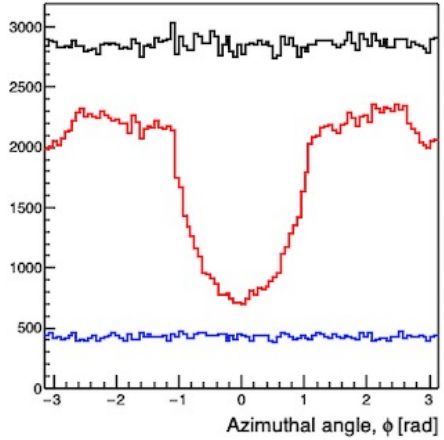


- J/ψ produced at mid-rapidity.
 - **Sensitive to gluons!**
- Tagging active and spectator nucleons allow for experimental control of nuclear configuration \rightarrow study transition into SRC region (e.g. where nuclear effects become larger).
- Tagging **both** nucleons allows for full reconstruction of momentum transfer!

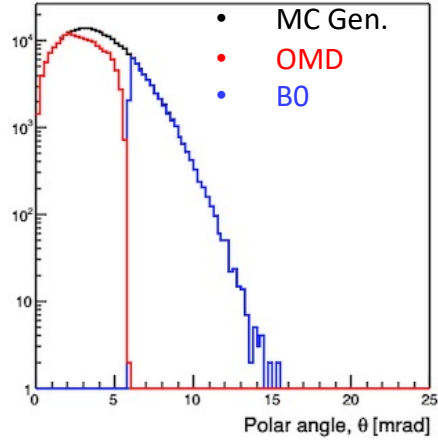
Short-Range Correlations in Deuterons

Z. Tu, A. Jentsch *et al.*, Phys. Lett. B, **811** (2020)

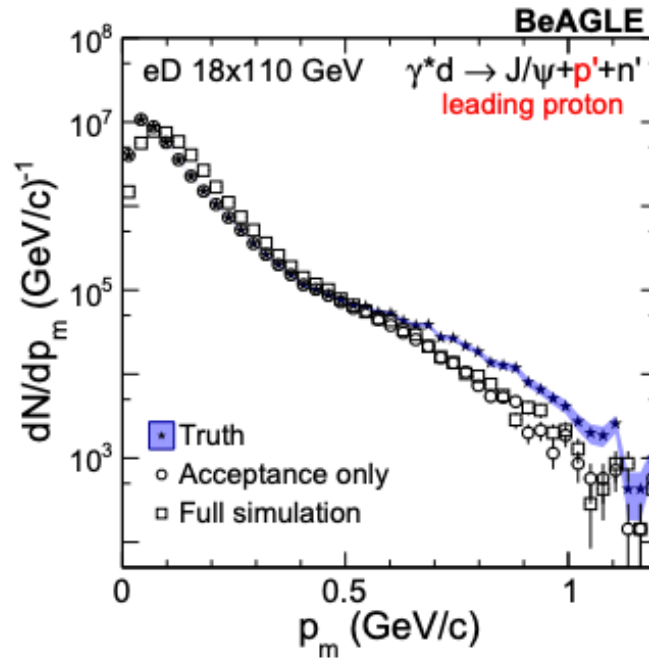
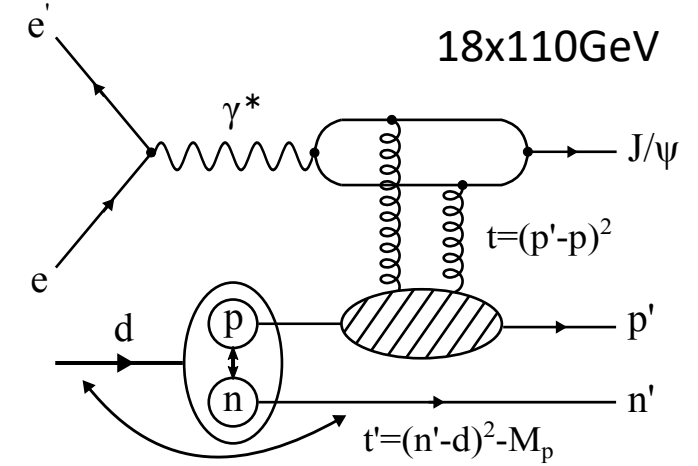
“active” protons



“active” protons



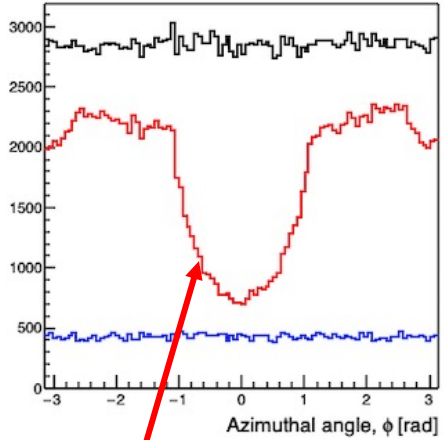
Neutron “spectator” case.



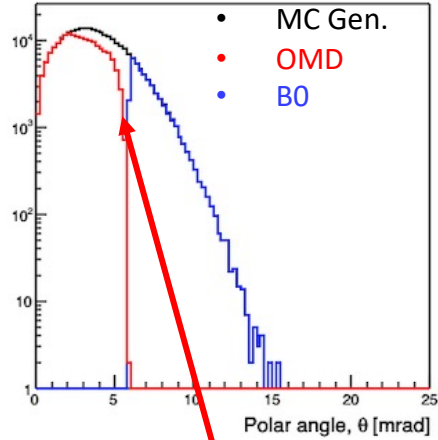
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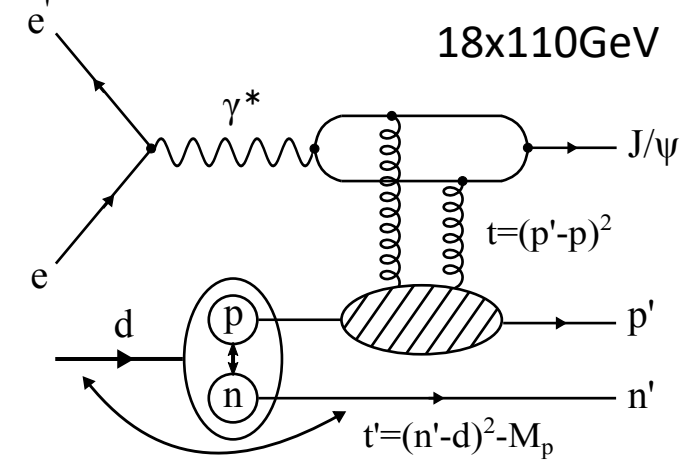
“active” protons



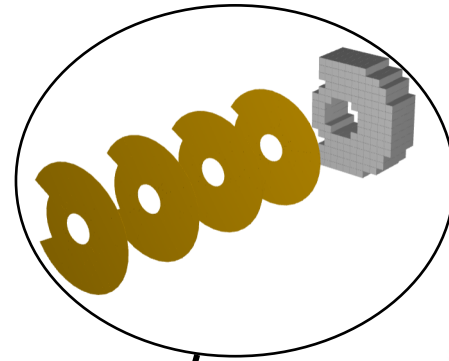
“active” protons



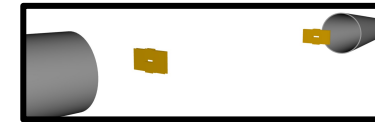
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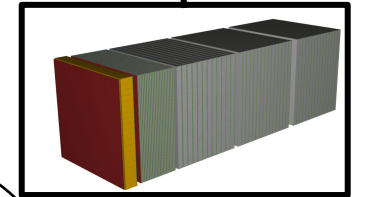
B0 detector



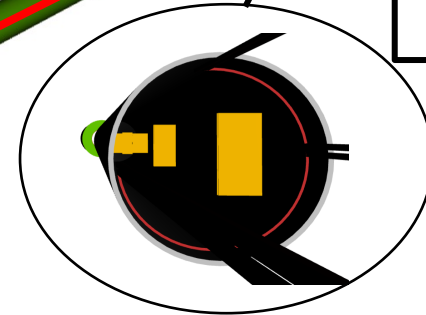
RP



ZDC

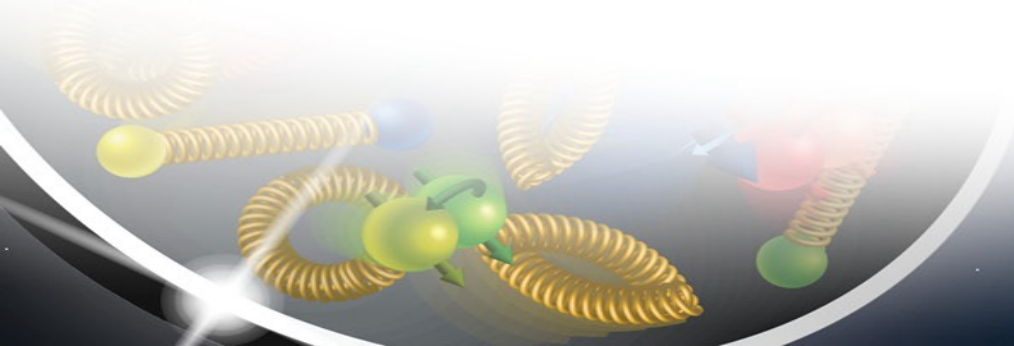


OMD



Protons lost in transition between very far-forward detectors and B0 spectrometer.

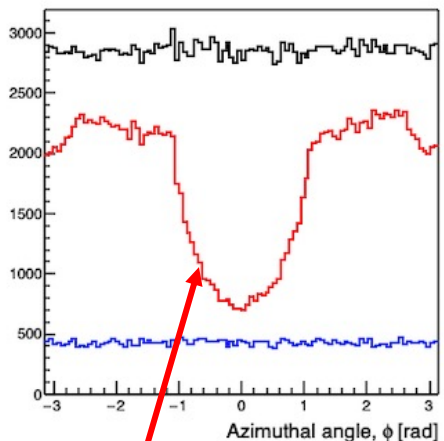
Off-momentum protons lost in quadrupole magnets.



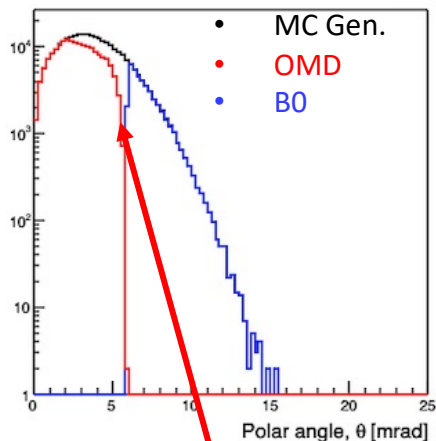
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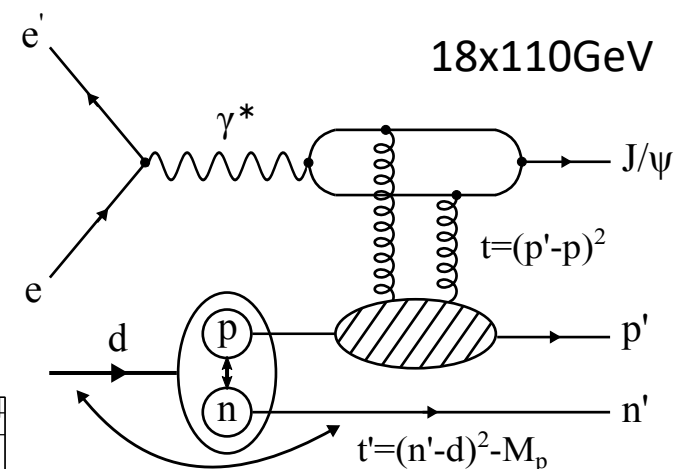
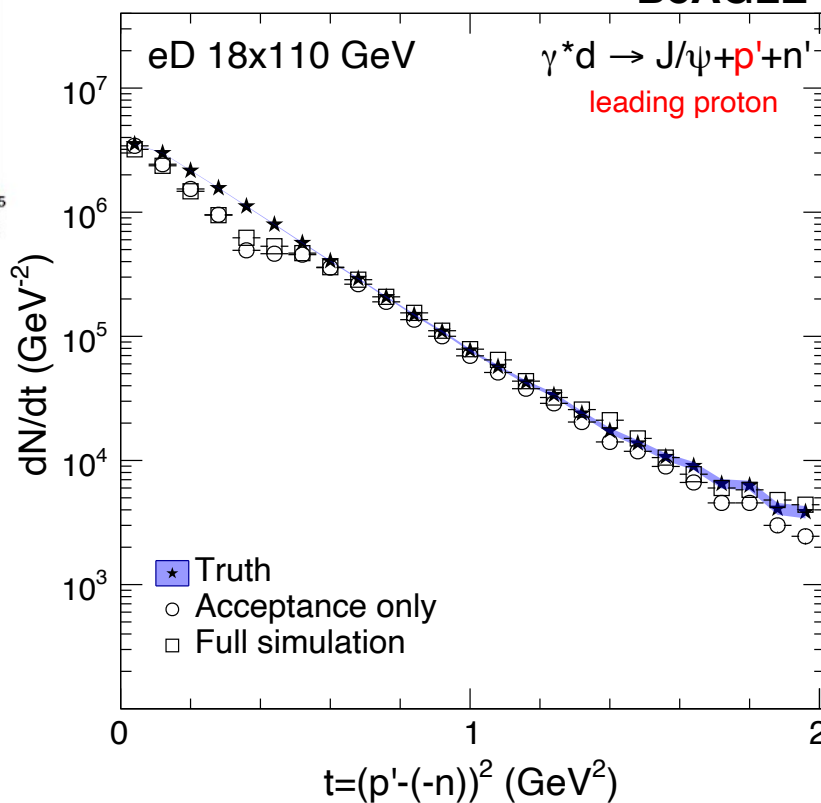


“active” protons



Neutron “spectator” case.

BeAGLE

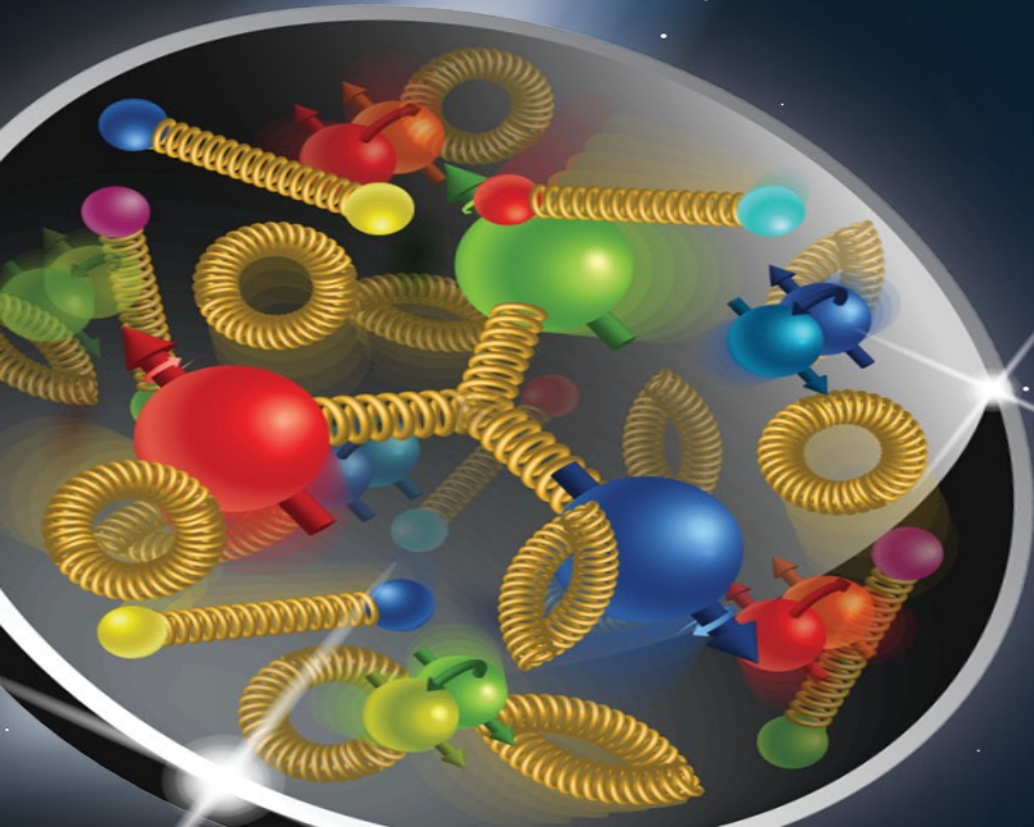
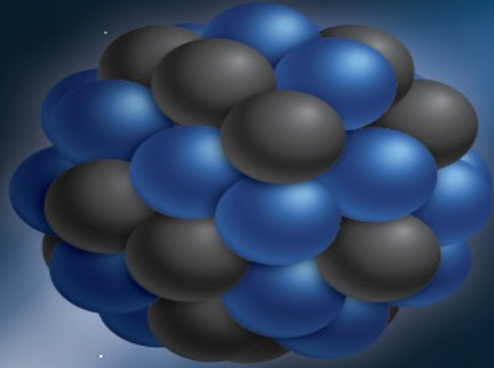


t-reconstruction using double-tagging (both proton and neutron reconstructed).

Protons lost in transition between very far-forward detectors and BO spectrometer.

Off-momentum protons lost in quadrupole magnets.

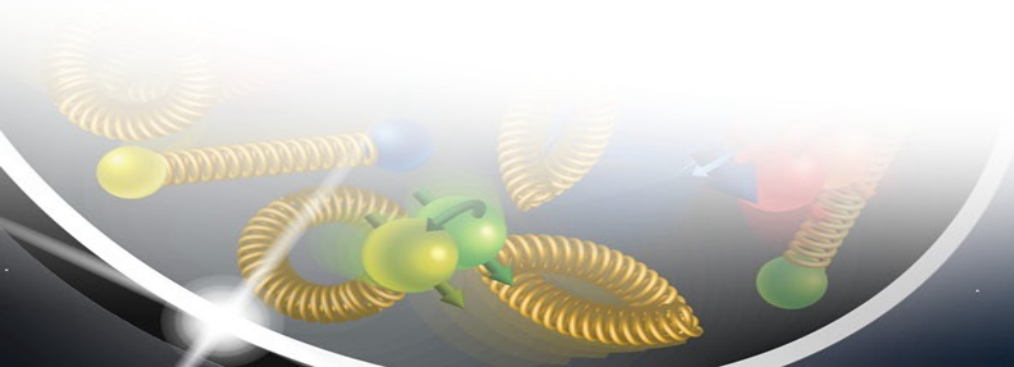
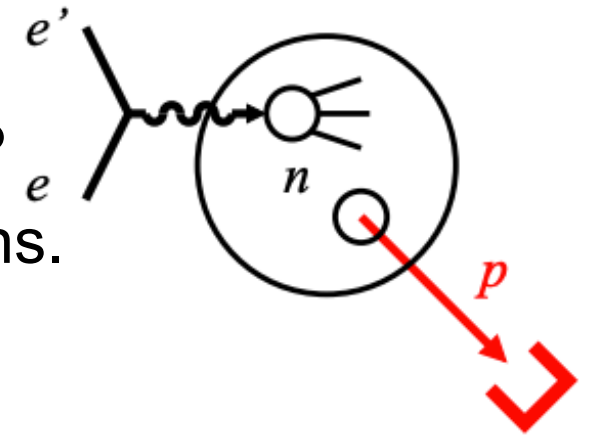
➤ Spectator information is the “dial” for the SRC region.



Deuterons: Free Neutron Structure

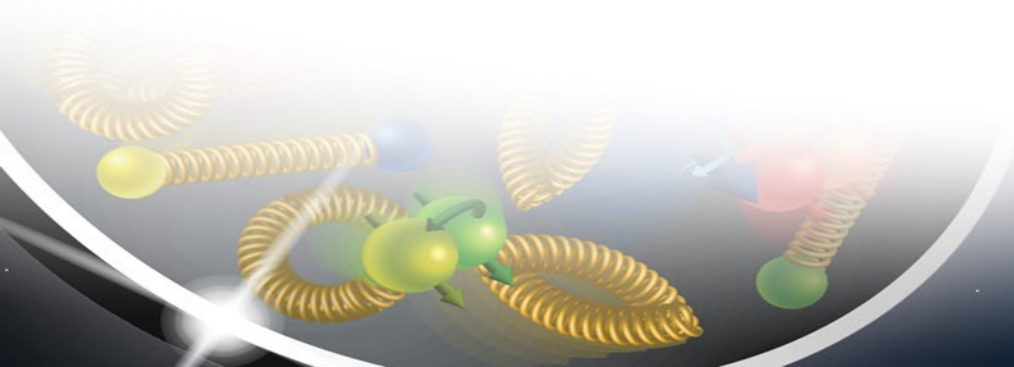
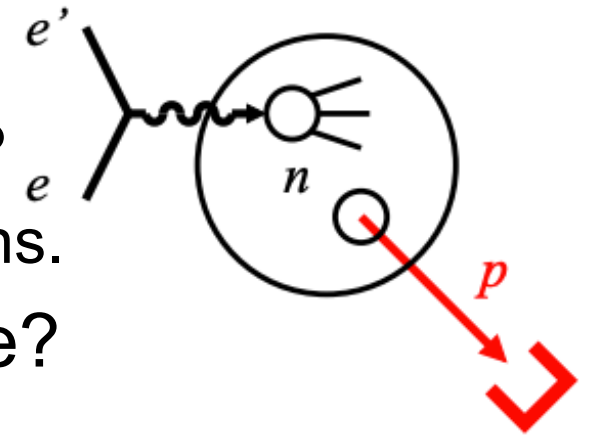
Neutron Structure

- Protons well-studied at HERA -> So...why the neutron?
 - Flavor separation, baseline for studies of nuclear modifications.



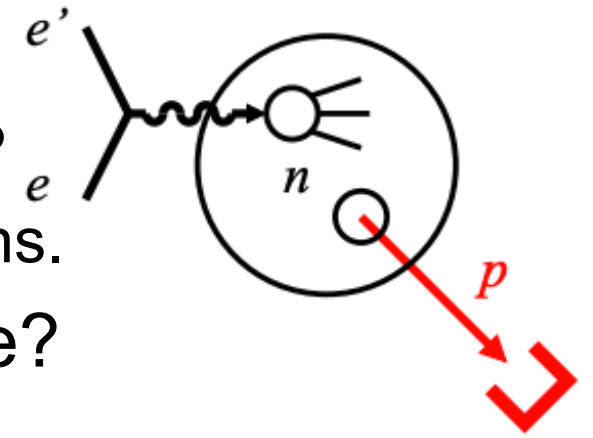
Neutron Structure

- Protons well-studied at HERA -> So...why the neutron?
 - Flavor separation, baseline for studies of nuclear modifications.
- What makes the free neutron structure hard to measure?
 - Can only access neutrons *in a nucleus*.
 - Includes nuclear binding effects, Fermi motion, etc.



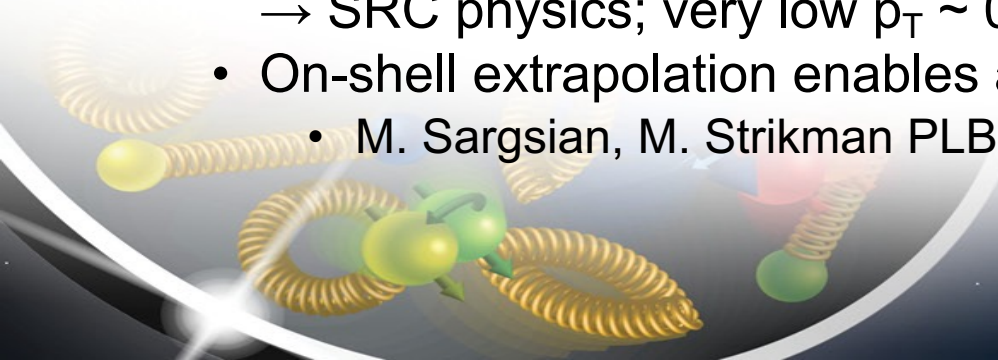
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- What makes the free neutron structure hard to measure?
 - Can only access neutrons *in a nucleus*.
 - Includes nuclear binding effects, Fermi motion, etc.



- Two options:

1. Inclusive measurements → Average over all nuclear configurations, use theory input to correct for nuclear binding effects.
2. Tagged measurements → Select nuclear configuration via spectator kinematics, allows for differential study.
 - Spectator kinematics provide a knob to dial in different regions of interest for study (i.e. high p_T → SRC physics; very low $p_T \sim 0$ GeV/c yields access to on-shell extrapolation).
 - On-shell extrapolation enables access to **free** nucleon structure.
 - M. Sargsian, M. Strikman PLB **639** (iss. 3-4) 223231 (2006)

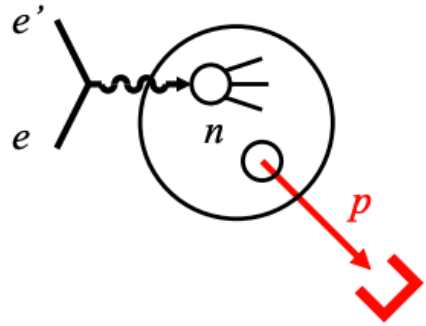


Neutron Structure

- Previous fixed target experiments with tagging have measured the neutron F_2 at high- x .
 - CLAS - Phys. Rev. Lett. **108**, 199902 (2012)
 - CLAS + BONUS - Phys. Rev. C 89, 045206 (2014)
 - measurement had a lower p_T cutoff ~ 70 MeV/c.
- Future JLAB 12 GeV studies planned.
 - ALERT - <https://arxiv.org/abs/1708.00891>
 - CLAS - https://www.jlab.org/exp_prog/proposals/10/PR12-06-113-pac36.pdf
- **Tagged DIS @ the EIC:**
 - In a collider, can tag spectators down to $p_T \sim 0$ MeV/c \rightarrow Enables extraction of free neutron structure function via pole extrapolation.
 - Can extend tagged DIS measurement to $x \lesssim 0.1$.



Tagged Deuteron Cross Section



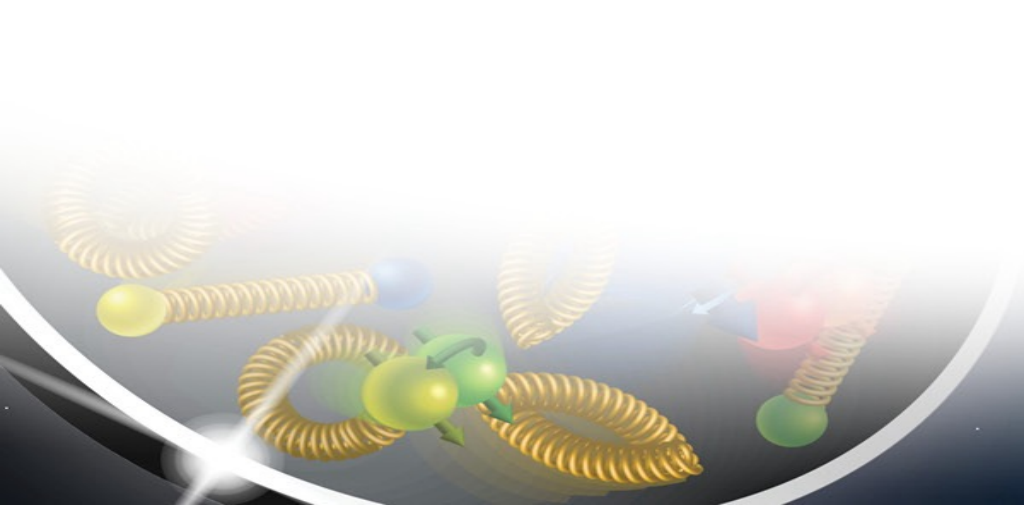
spectator nucleon (p_{pT}, α_p)

Total cross section $d\sigma = Flux(x, Q^2) \times \sigma_{red,d} \times \frac{dx}{2} dQ^2 \frac{d\phi_{e'}}{2\pi} [2(2\pi)^3]^{-1} \frac{d\alpha_p}{\alpha_p} \frac{dp_{pT}^2}{2} d\phi_p$

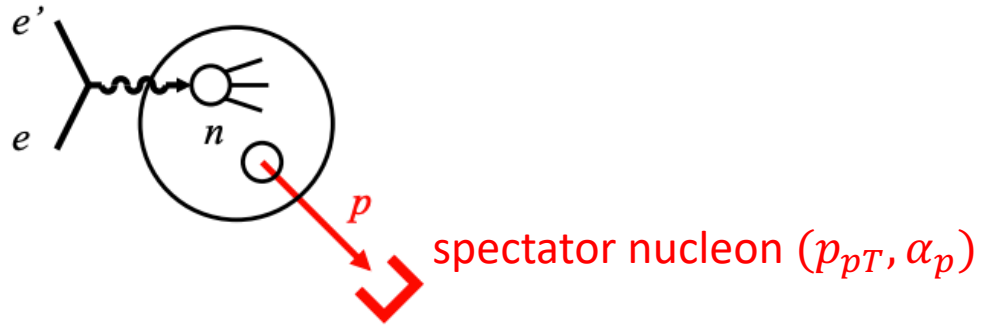
α_p : light-cone momentum fraction

$$\alpha_p \equiv \frac{2p_p^+}{p_d^+} = \frac{2(E_p + p_{z,p})}{M_d}$$

S_d : deuteron spectral function pole



Tagged Deuteron Cross Section



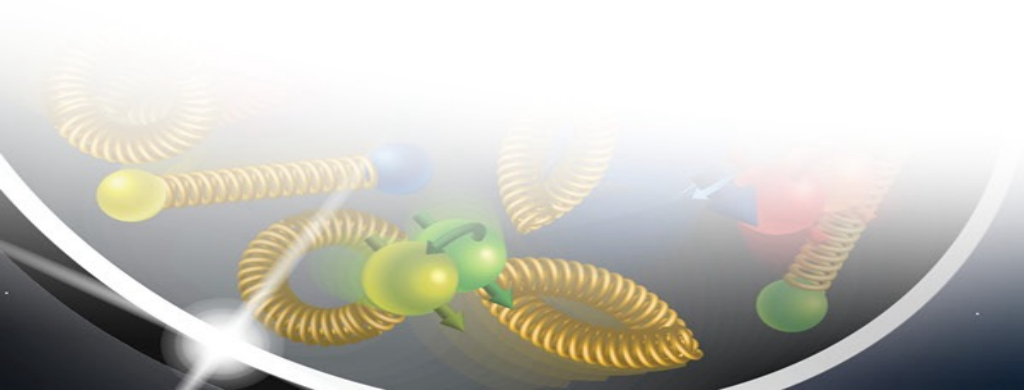
Total cross section
$$d\sigma = Flux(x, Q^2) \times \sigma_{red,d} \times \frac{dx}{2} dQ^2 \frac{d\phi_{e'}}{2\pi} [2(2\pi)^3]^{-1} \frac{d\alpha_p}{\alpha_p} \frac{dp_{pT}^2}{2} d\phi_p$$

α_p : light-cone momentum fraction

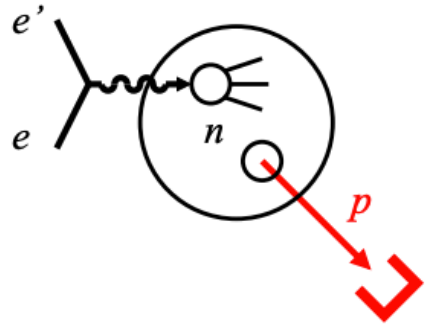
$$\alpha_p \equiv \frac{2p_p^+}{p_d^+} = \frac{2(E_p + p_{z,p})}{M_d}$$

S_d : deuteron spectral function pole

- Measure the cross-section differential on the spectator kinematics.
 - Spectator kinematics provide control knob on the nuclear configuration.
- Solve for the deuteron reduced cross section.



Tagged Deuteron Cross Section



spectator nucleon (p_{pT}, α_p)

Total cross section
$$d\sigma = Flux(x, Q^2) \times \sigma_{red,d} \times \frac{dx}{2} dQ^2 \frac{d\phi_{e'}}{2\pi} [2(2\pi)^3]^{-1} \frac{d\alpha_p}{\alpha_p} \frac{dp_{pT}^2}{2} d\phi_p$$

α_p : light-cone momentum fraction

$$\alpha_p \equiv \frac{2p_p^+}{p_d^+} = \frac{2(E_p + p_{z,p})}{M_d}$$

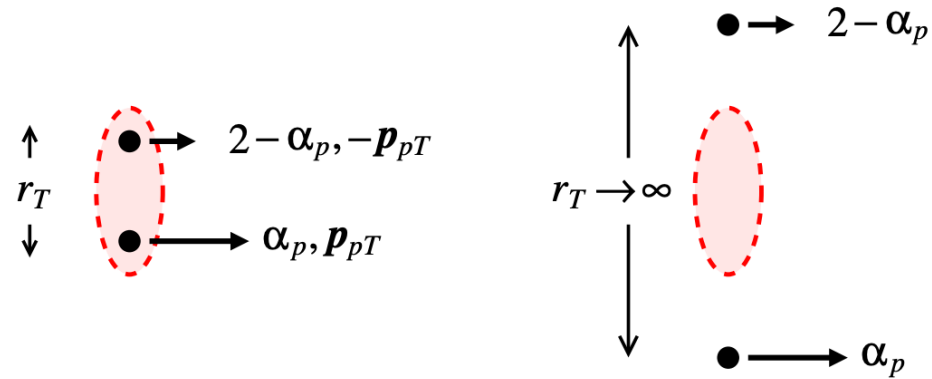
S_d : deuteron spectral function pole

- Measure the cross-section differential on the spectator kinematics.
 - Spectator kinematics provide control knob on the nuclear configuration.
- Solve for the deuteron reduced cross section.
- Deuteron reduced cross section related to the struck nucleon reduced cross section via the deuteron spectral function.

$$\sigma_{red,d}(x, Q^2; p_{pT}, \alpha_p) = [2(2\pi)^3] \times S_d(p_{pT}, \alpha_p) [pole] \times \sigma_{red,n}(x, Q^2)$$

Measurement of the deuteron reduced cross section yields access to the struck nucleon structure via the tagged spectator!

Pole Extrapolation



- Divide by deuteron spectral function (nucleon pole).
 - The resulting distribution is the active nucleon reduced cross section as a function of p_{pT}^2 .

$$\sigma_{red,n}(x, Q^2) = \frac{\sigma_{red,d}(x, Q^2; p_{pT}, \alpha_p)}{[2(2\pi)^3]S_d(p_{pT}, \alpha_p)[pole]}$$

$$S_d(p_{pT}, \alpha_p)[pole] = \frac{R}{(p_{pT}^2 + a_T^2)^2} \quad \text{Deuteron spectral function}$$

$p_{pT}^2 > 0$
physical region

$p_{pT}^2 \rightarrow -a_T^2$
pole extrapolation

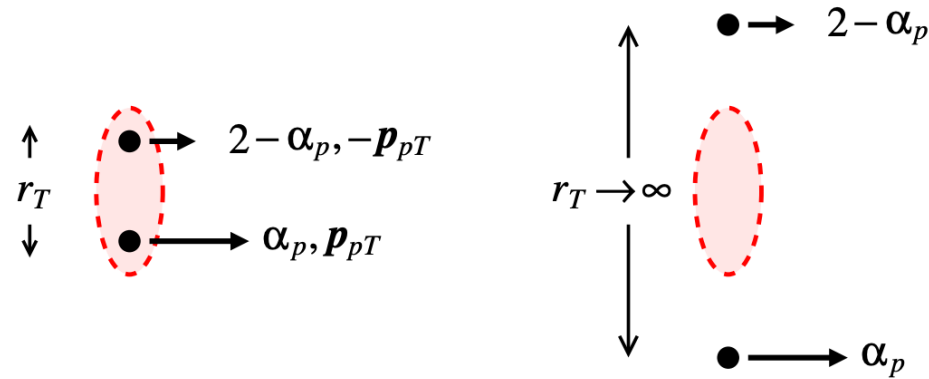
$$R = 2\alpha_p^2 m_N \Gamma^2 (2 - \alpha_p)$$

$$a_T^2 = m_N^2 - \alpha_p(2 - \alpha_p) \frac{M_d^2}{4}$$

$R = \text{residue of spectral function}$

$a_T^2 = \text{position of pole}$

Pole Extrapolation



- Divide by deuteron spectral function (nucleon pole).
 - The resulting distribution is the active nucleon reduced cross section as a function of p_{pT}^2 .

$$\sigma_{red,n}(x, Q^2) = \frac{\sigma_{red,d}(x, Q^2; p_{pT}, \alpha_p)}{[2(2\pi)^3]S_d(p_{pT}, \alpha_p)[pole]}$$

$$S_d(p_{pT}, \alpha_p)[pole] = \frac{R}{(p_{pT}^2 + a_T^2)^2} \quad \text{Deuteron spectral function}$$

- Extrapolate to $p_{pT}^2 \rightarrow -a_T^2$ to extract F_2 to extract free nucleon F_2 .
 - Pole extrapolation selects large-size pn configurations where nuclear binding and FSI are absent.

$$R = 2\alpha_p^2 m_N \Gamma^2 (2 - \alpha_p)$$

$$a_T^2 = m_N^2 - \alpha_p(2 - \alpha_p) \frac{M_d^2}{4}$$

$R =$ residue of spectral function

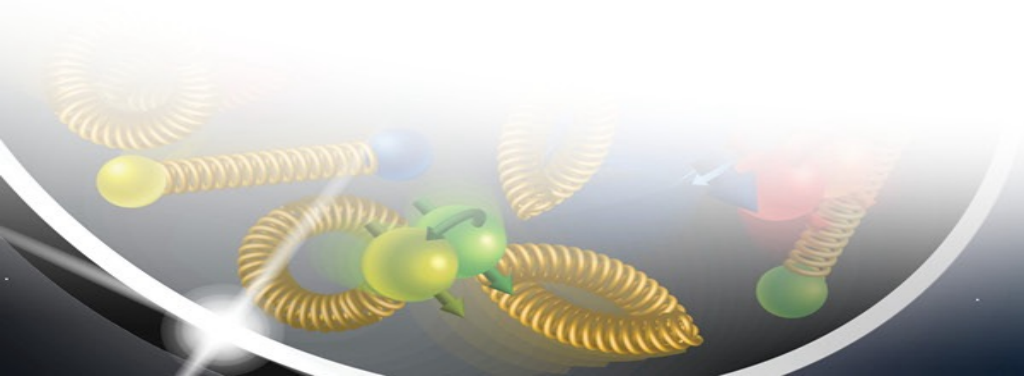
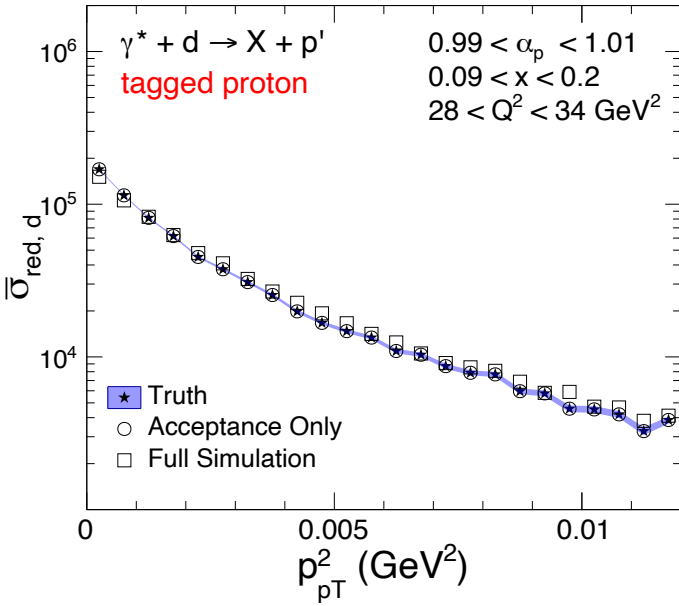
$a_T^2 =$ position of pole

Free Neutron F_2 Extraction

A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104**, 065205, (2021) (Editor's Suggestion)

- Start with the deuteron reduced cross section \rightarrow **direct measurement!**

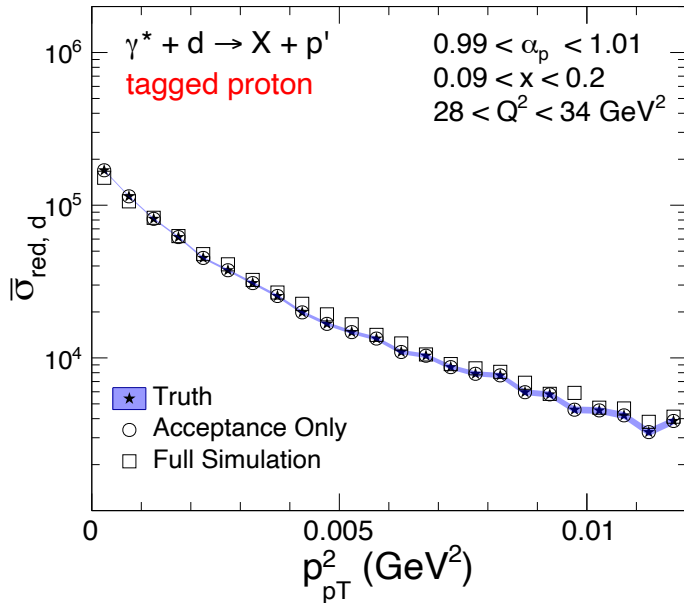
(deuteron reduced cross section)
eD 18 x 110 GeV² BeAGLE



Free Neutron F_2 Extraction

A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104**, 065205, (2021) (Editor's Suggestion)

(deuteron reduced cross section)
eD 18 x 110 GeV² BeAGLE

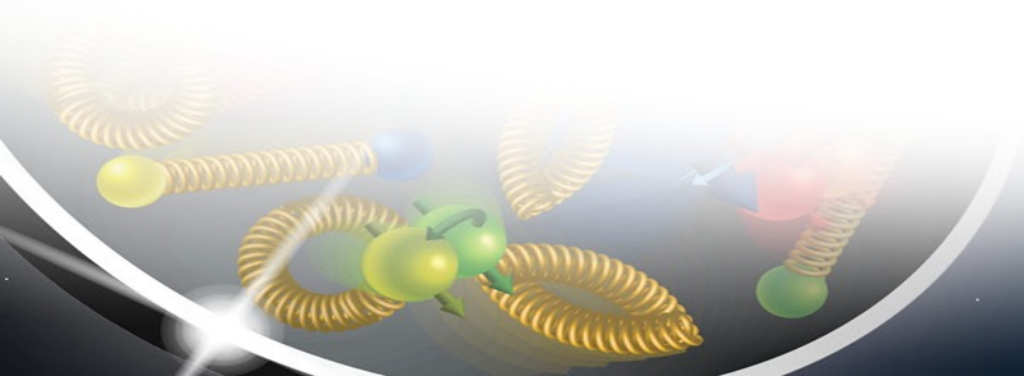


- Start with the deuteron reduced cross section \rightarrow direct measurement!
- Multiply by the inverse of the deuteron spectral function pole.



$$\frac{1}{S_d(p_{pT}, \alpha_p)[pole]}$$

(inverse pole of deuteron spectral function)



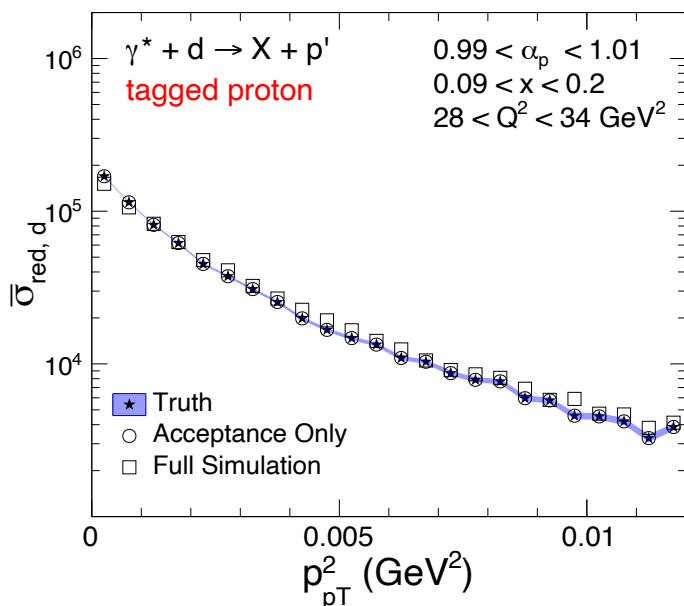
Free Neutron F_2 Extraction

A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104**, 065205, (2021) (Editor's Suggestion)

(deuteron reduced cross section)

eD 18 x 110 GeV²

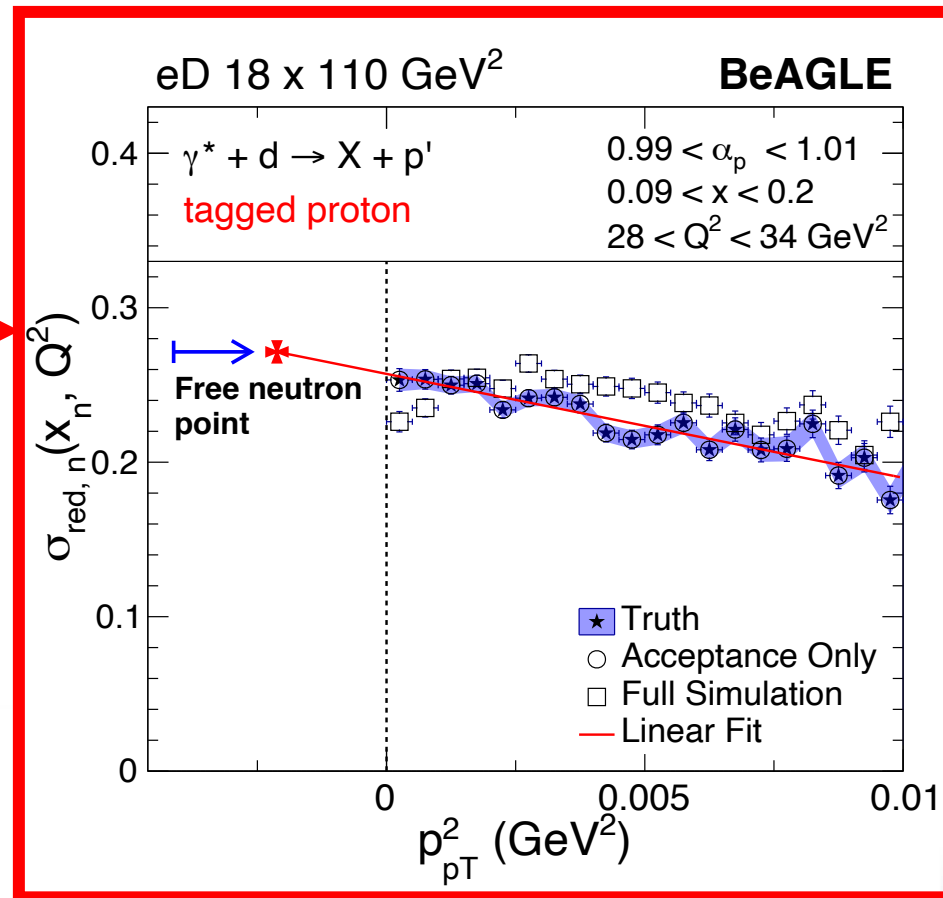
BeAGLE



RESULT: Reduced cross section on the active nucleon.

$$\frac{1}{S_d(p_{pT}, \alpha_p)[pole]}$$

(inverse pole of deuteron spectral function)



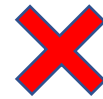
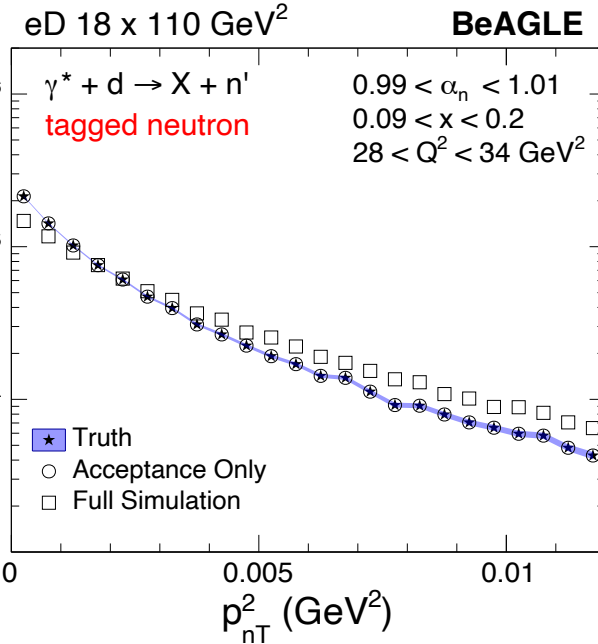
(Active nucleon reduced cross section)

$$\sigma_{red,n}(x, Q^2) = \frac{\sigma_{red,d}}{[2(2\pi)^3]S_d(p_{pT}, \alpha_p)}$$

Free Proton F_2 Extraction

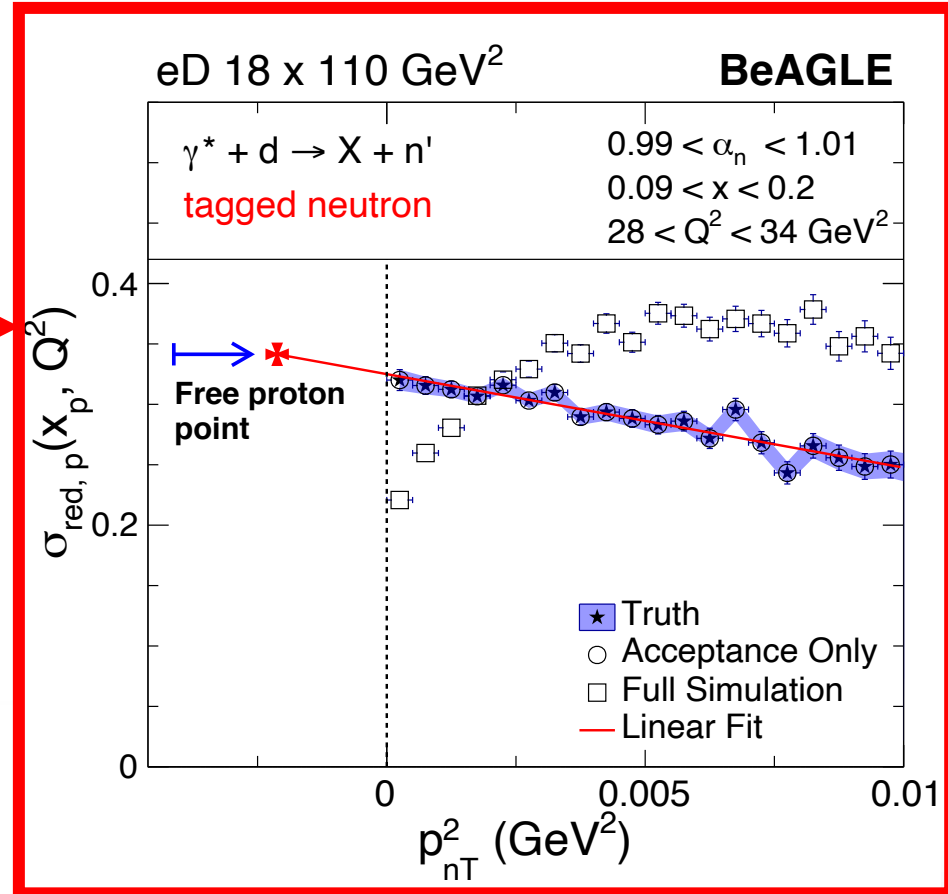
A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104**, 065205, (2021) (Editor's Suggestion)

(deuteron reduced cross section)



$$\frac{1}{S_d(p_{pT}, \alpha_p)[pole]}$$

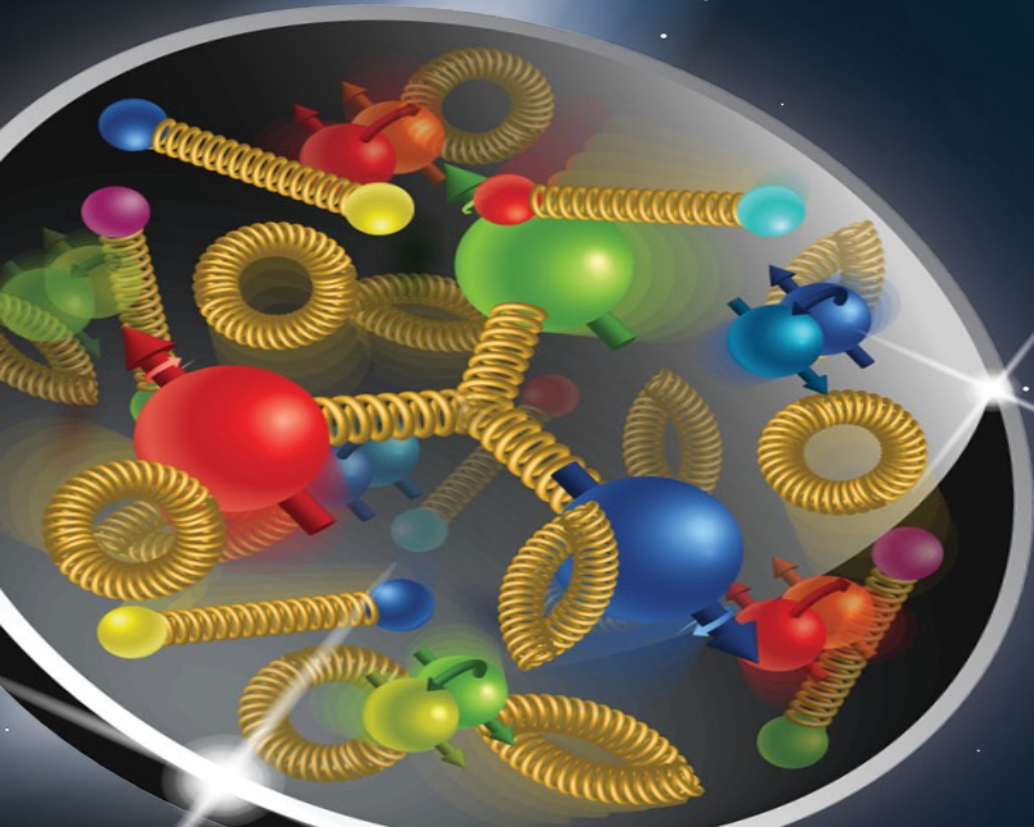
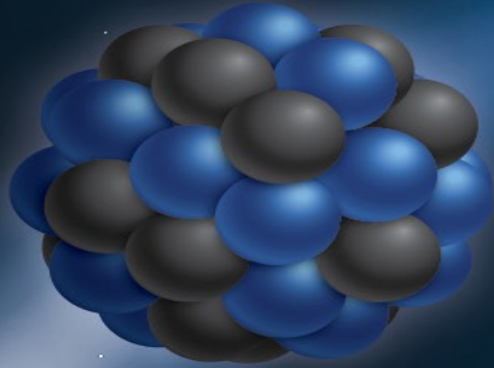
(inverse pole of deuteron spectral function)



(Active nucleon reduced cross section)

Measurement of proton F_2 using this method provides ability to directly estimate systematics for extrapolation procedure, since proton F_2 directly measurable in e+p scattering!

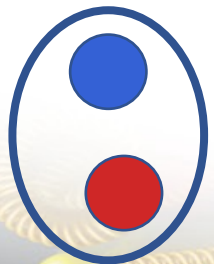
$$\sigma_{\text{red},p}(x, Q^2) = \frac{\sigma_{\text{red},d}}{[2(2\pi)^3]S_d(p_{nT}, \alpha_n)}$$



Deuterons:
The EMC Effect
(on-going study)

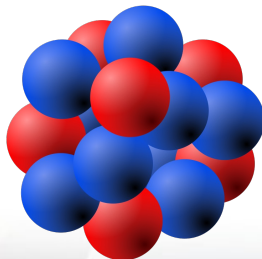
The EMC Effect

- Discovered by the European Muon Collaboration ~40 years ago.
 - Puzzle: why the dip?
- Still an unanswered question, and one we hope the EIC can aid in answering.
- Established via measurements with **different nuclear targets!**



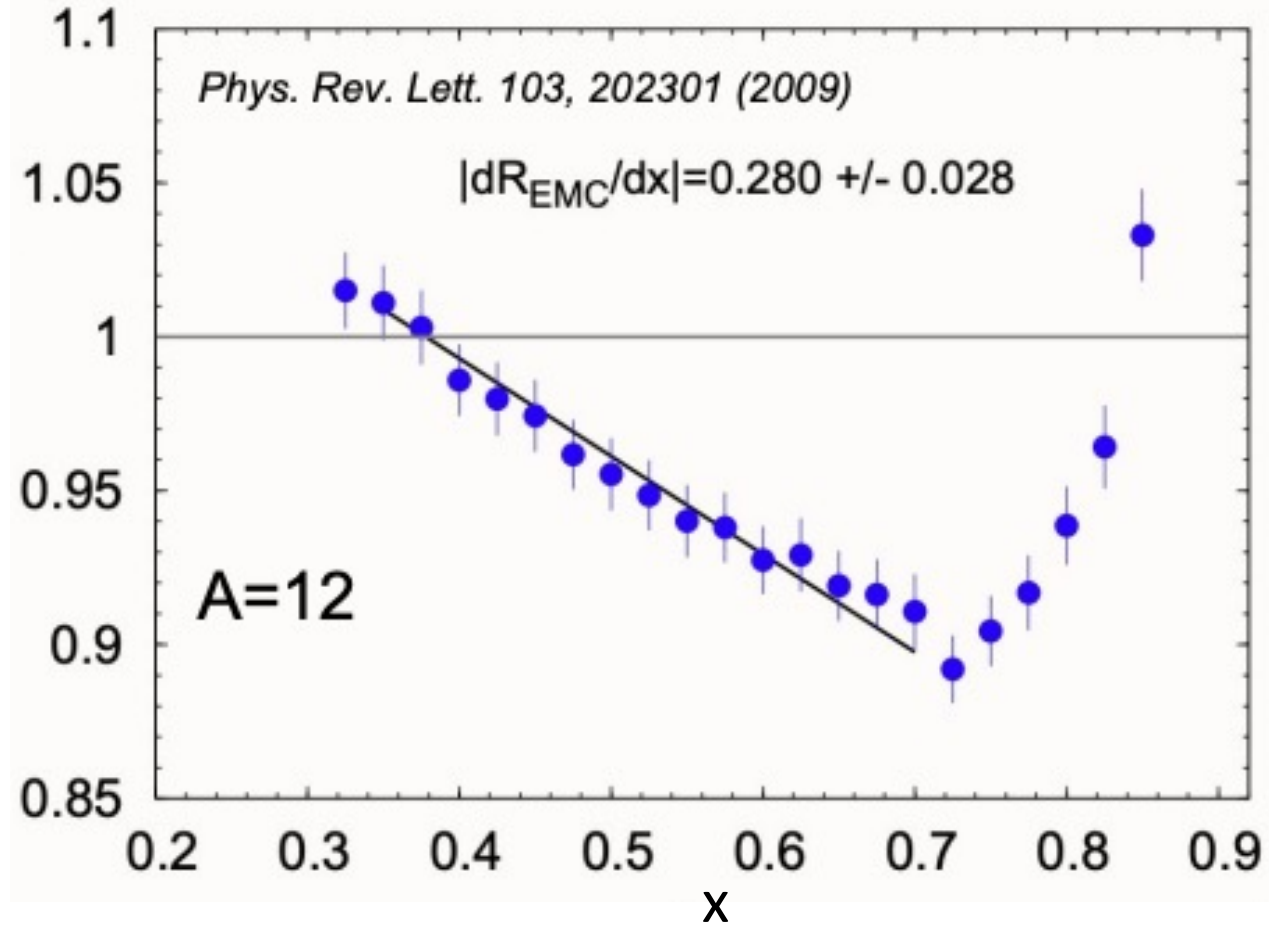
deuteron

Nuclear effects
modify nucleon
structure? How?



Heavier nucleus (A > 2)

$$R_{EMC} = (\sigma_A/A)/(\sigma_D/2)$$



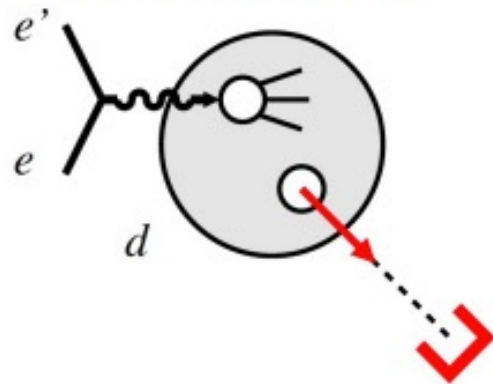
Understanding the origin of the EMC effect and nuclear modifications of prime interest in nuclear physics!

The Deuteron – a stand-alone lab for nuclear physics

- **Off-shellness in deuterons as a probe of nuclear effects.**

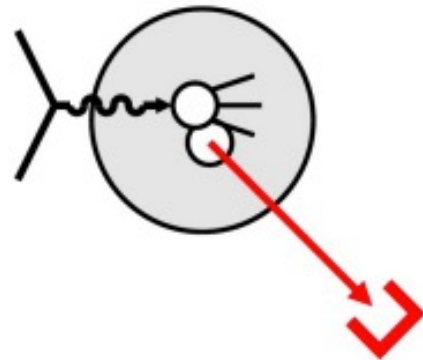
Tagged DIS Process: $e + d \rightarrow e' + X + p' \text{ or } n'$

Low off-shellness



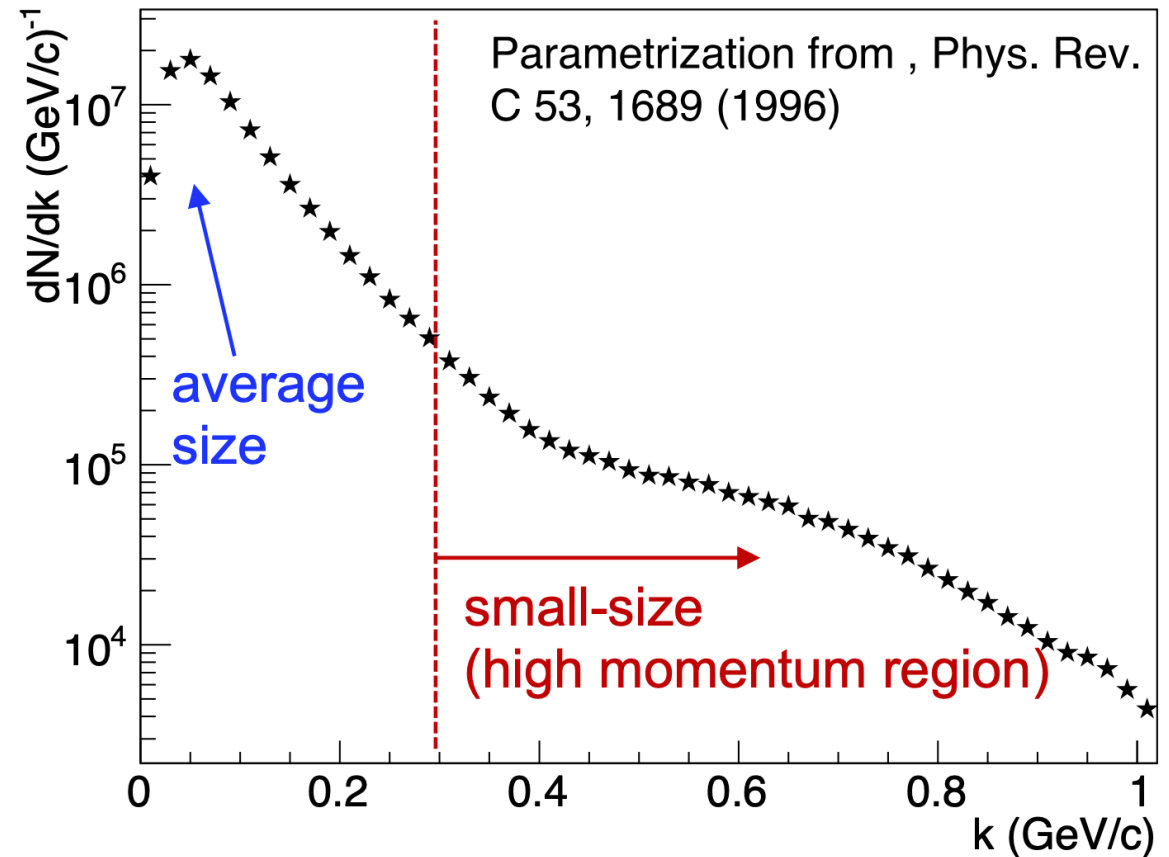
average-size
more-likely

High off-shellness



small-size
less-likely

Deuteron: nucleon internal momentum

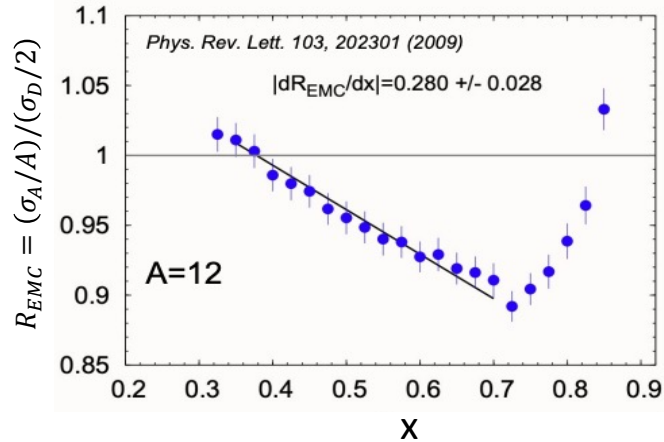


$$-t'^2 = M_N^2 - (p_d - p_p)^2$$

Virtuality/off-shellness in the deuteron

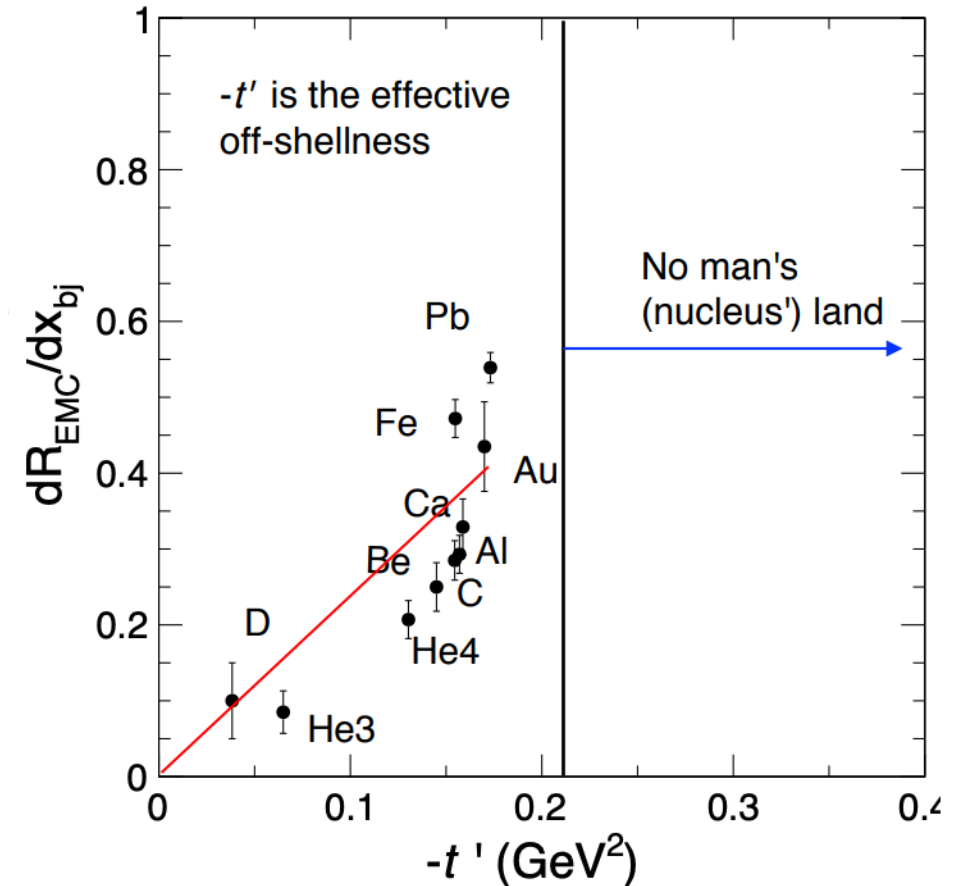
Question: can the EMC effect be controlled via the off-shellness **without altering the nuclear species?**

Simulating the EMC Effect in BeAGLE



Use EMC effect slope measurements from data with different nuclear targets.

*Data from J. Seely *et al.* Phys. Rev. Lett. **103**, 202301 (2009)

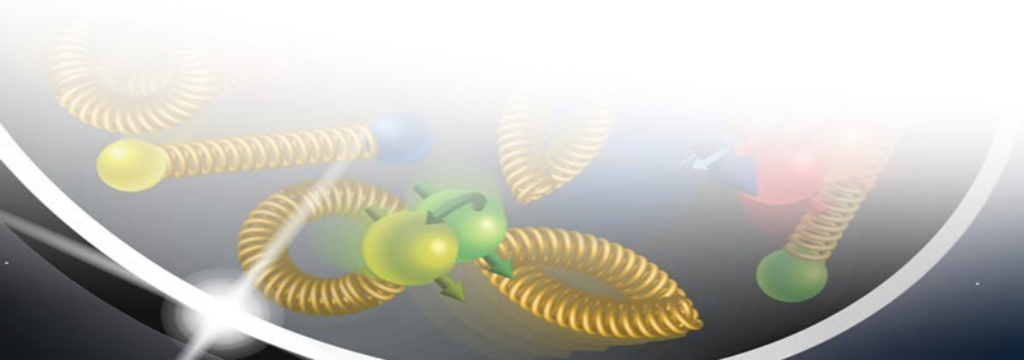


Linear fit to virtuality dependence \rightarrow Minimal parametrization:

Frankfurt and Strikman, Nuc. Phys. B **250** (1985)

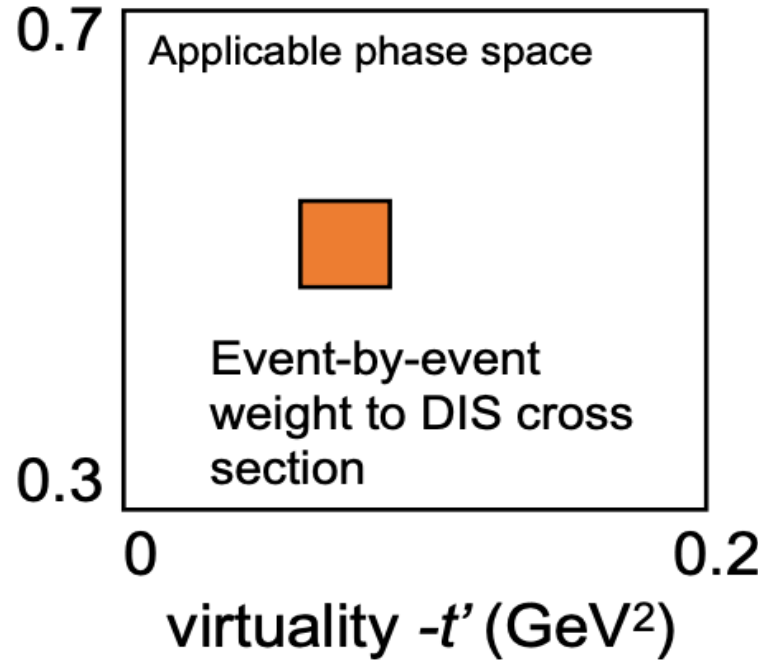
C. Ciofi *et al.*, Phys. Rev. C **76**, 055206 (2007)

And others...



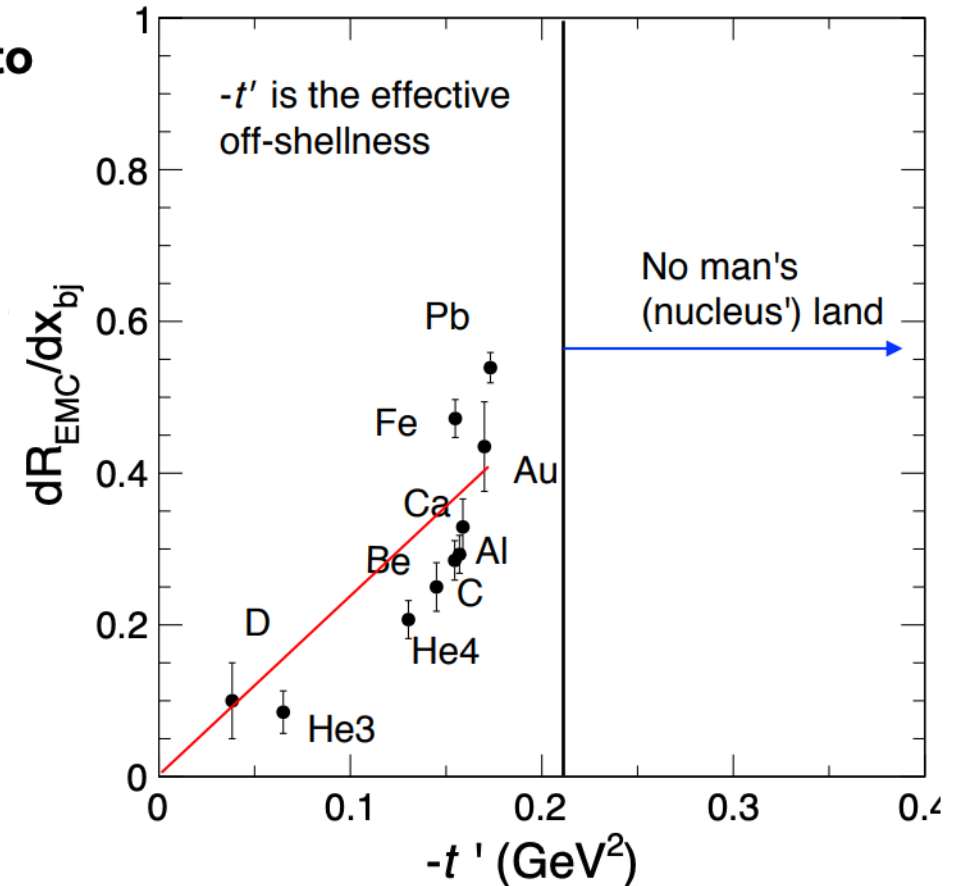
Simulating the EMC Effect in BeAGLE

BeAGLE



- Only apply to $0.3 < x_{bj} < 0.7$
- Q^2 independent
- Weight = $F_2(\text{bound}) / F_2(\text{free})$

Add EMC effect according to the linear parametrization



Linear fit to virtuality dependence → Minimal parametrization:

Frankfurt and Strikman, Nuc. Phys. B **250** (1985)

C. Ciofi *et al.*, Phys. Rev. C **76**, 055206 (2007)

And others...

The EMC Effect @ the EIC

- Approach:

- Measure deuteron reduced cross-section σ_D , with and without the off-shell effects included.
 - No FSI included.
- Ratio of σ_D **inside and outside the EMC region** (e.g. $x \sim 0.5$ and $x \sim 0.2$)

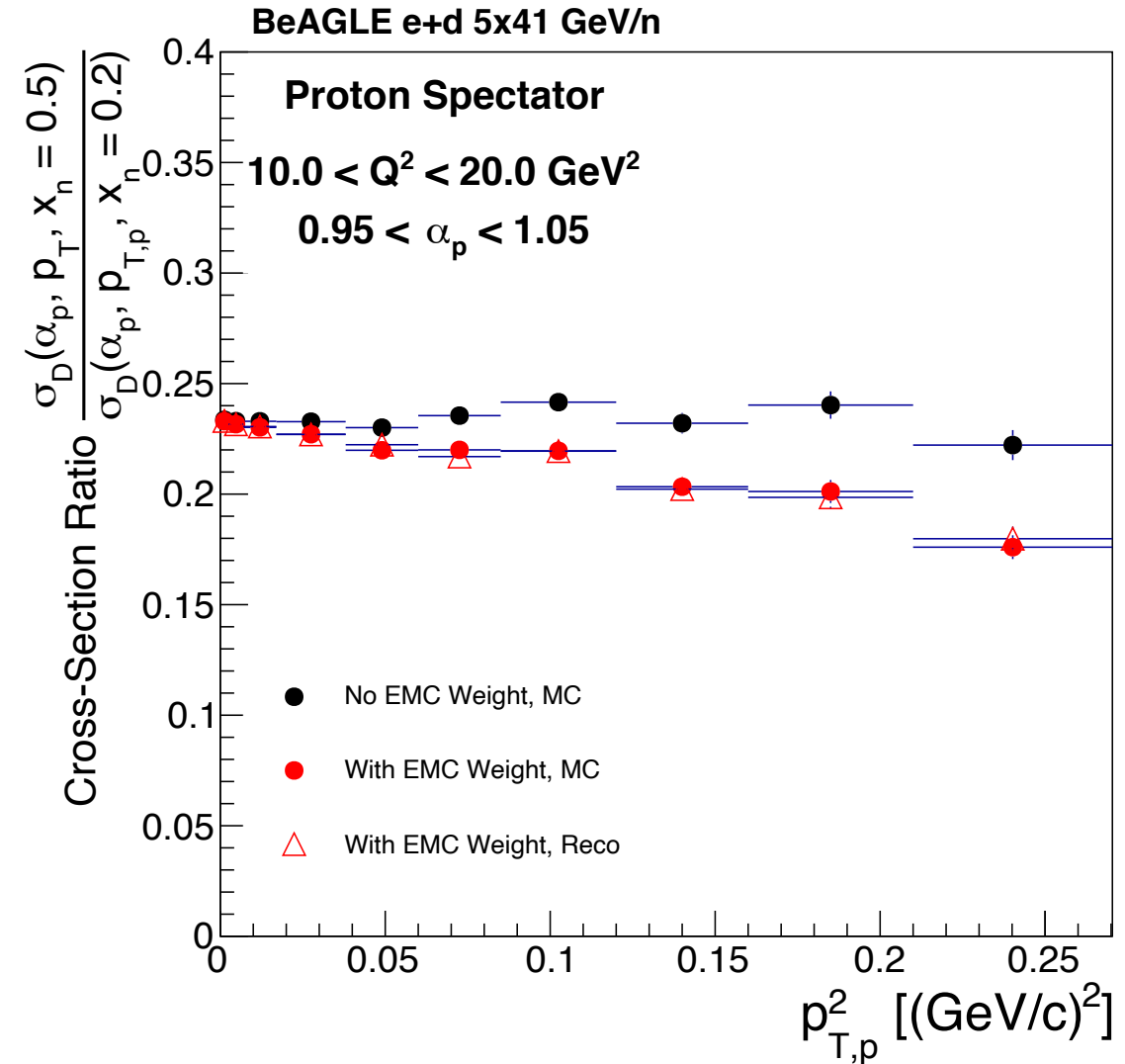
➤ Quantity allows direct comparison of cross section with and without EMC weight ($x \sim 0.2$ chosen to avoid anti-shadowing region).

$$\frac{\sigma_D(\alpha_p, p_{T,p}, x_n = 0.5)}{\sigma_D(\alpha_p, p_{T,p}, x_n = 0.2)}$$

The EMC Effect @ the EIC

5x41 GeV/n Integrated Luminosity $\sim 25 \text{ fb}^{-1}$

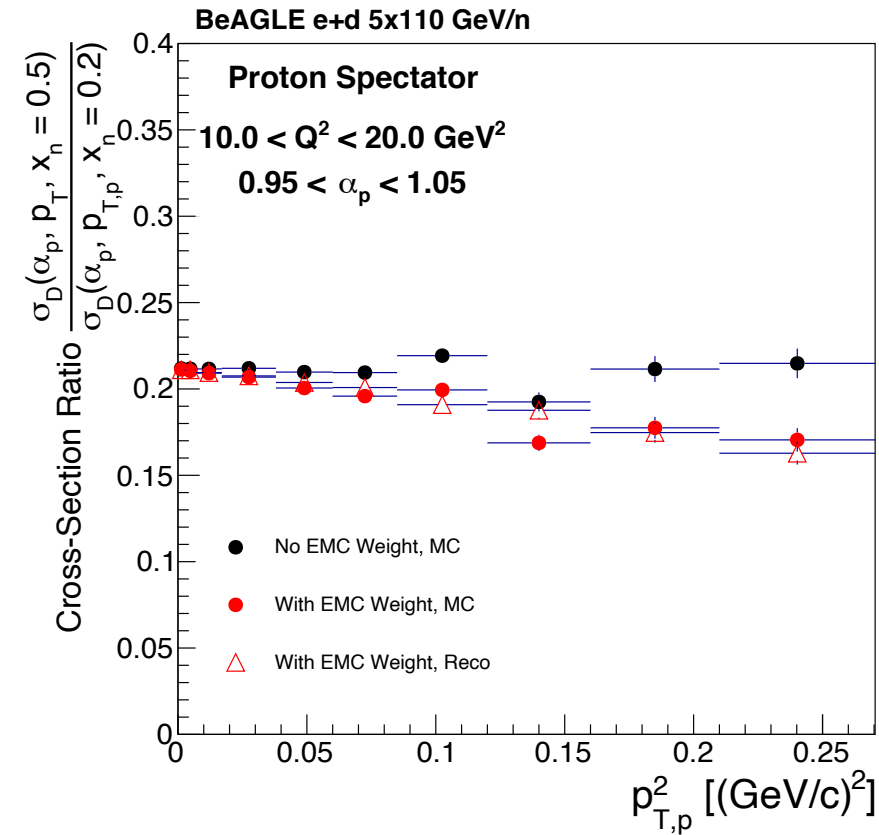
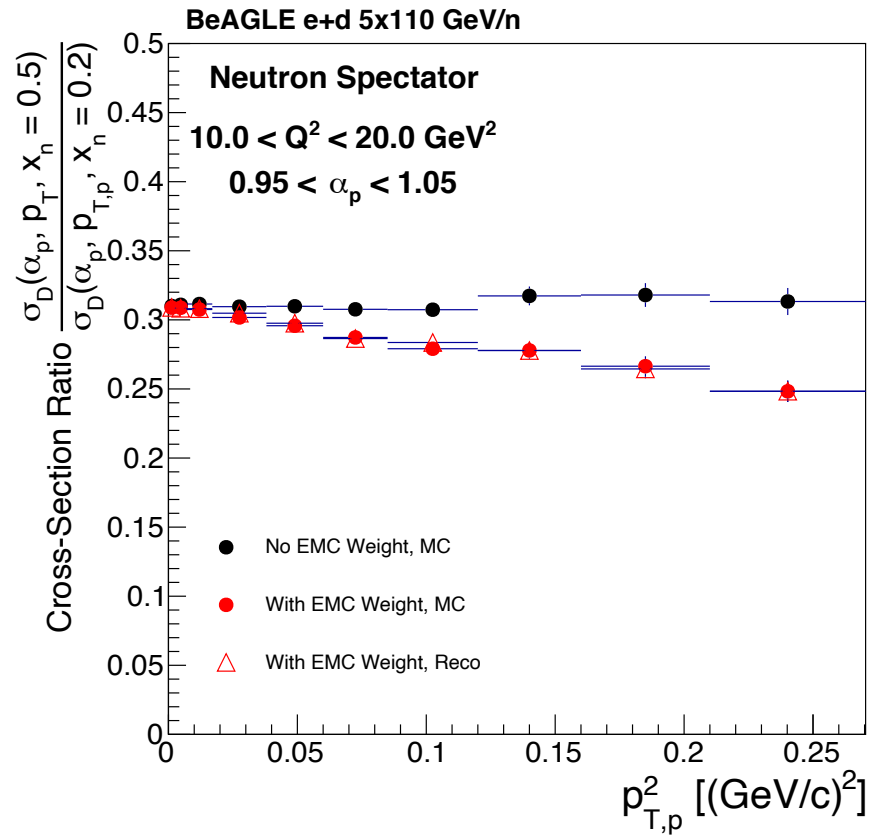
- Approach:
 - Measure deuteron reduced cross-section σ_D , with and without the off-shell effects included.
 - No FSI included.
 - Ratio of σ_D **inside and outside the EMC region** (e.g. $x \sim 0.5$ and $x \sim 0.2$)
 - Establish required integrated luminosity.
 - **Challenging measurement** \rightarrow high- x + low probability nuclear configuration + lower beam energies.
 - **Neutron spectator not possible in 5x41 GeV/n due to detector acceptance.**



The EMC Effect @ the EIC

5x110 GeV/n Integrated Luminosity $\sim 16 \text{ fb}^{-1}$

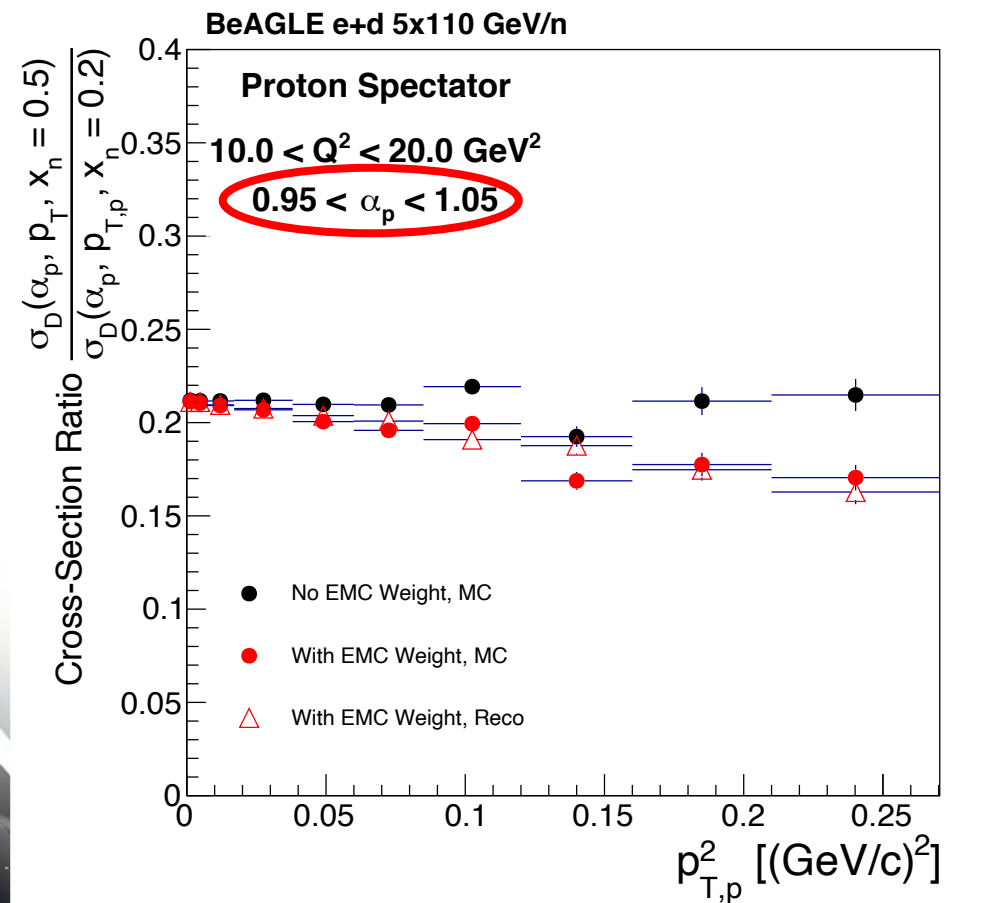
- EIC versatility \rightarrow different beam energy configurations!



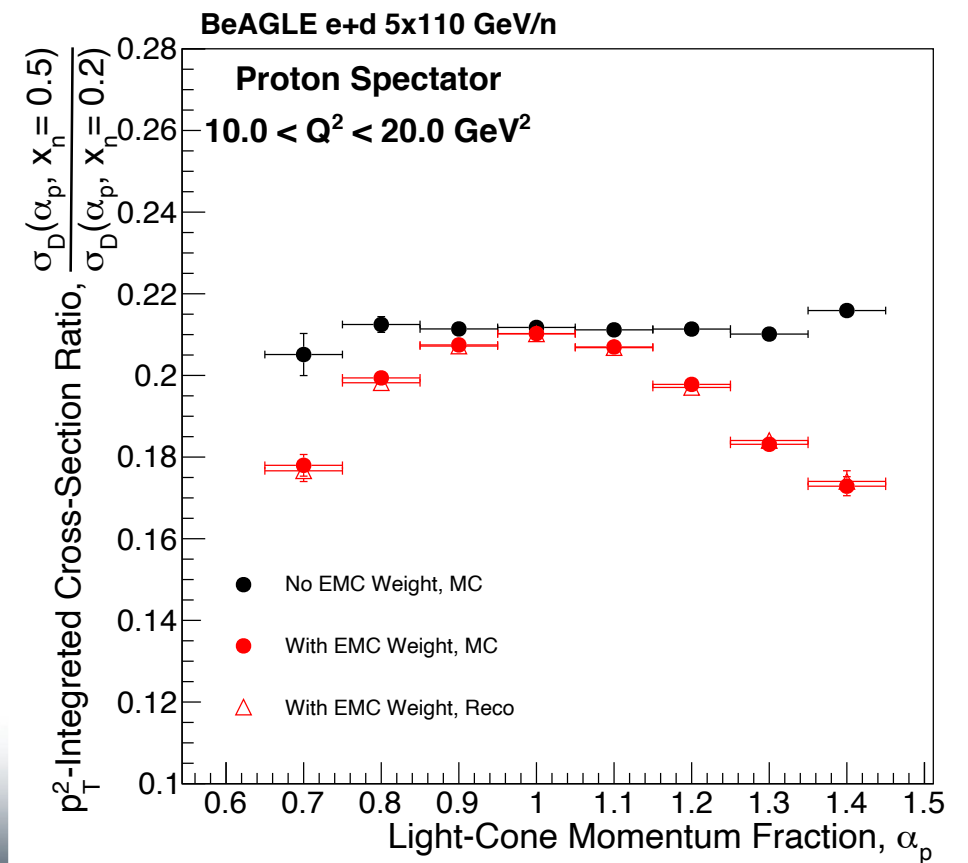
- Higher energy configuration (5x110 GeV/n).
- **More favorable detector acceptance \rightarrow study of proton *and* neutron spectators with same beam configuration.**
- Measurement of same observable with different beam energies/spectator reconstruction enables better understanding of experimental systematics.

Different nuclear configurations

- EIC kinematic coverage enables broad, differential study of effects.
 - Spectator kinematic coverage \rightarrow varied deuteron nuclear configurations.

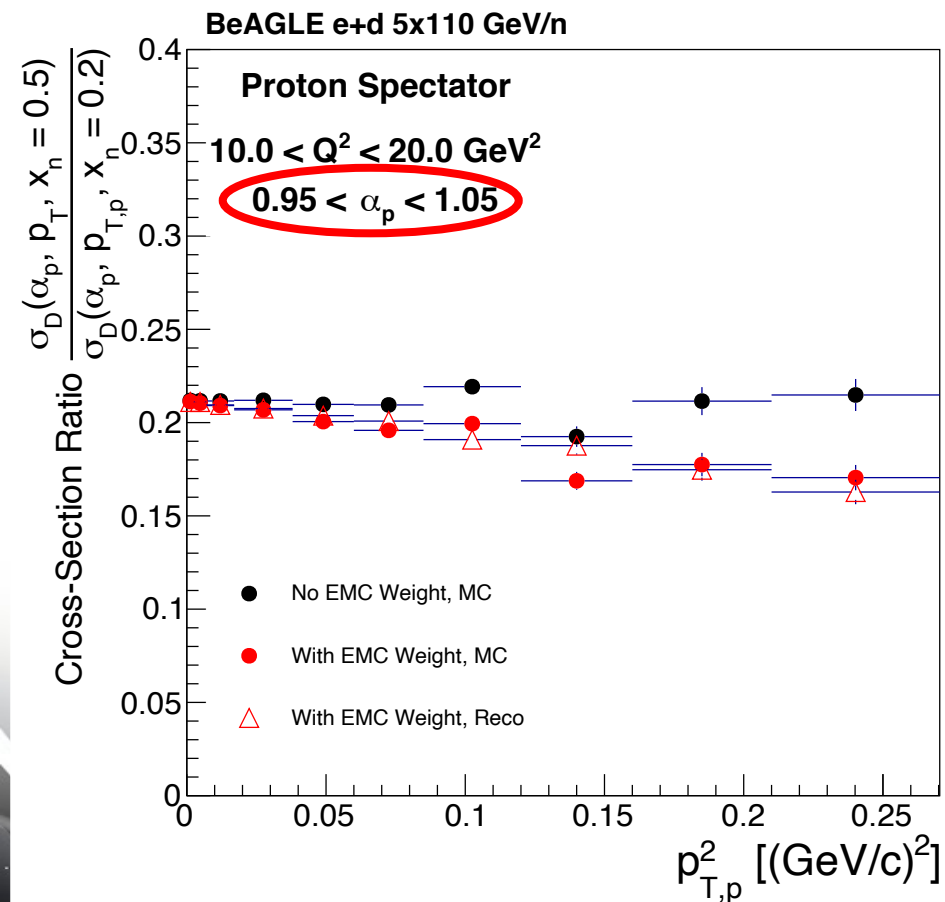


Integrate cross section over $p_{T,p}^2$ in each α bin.

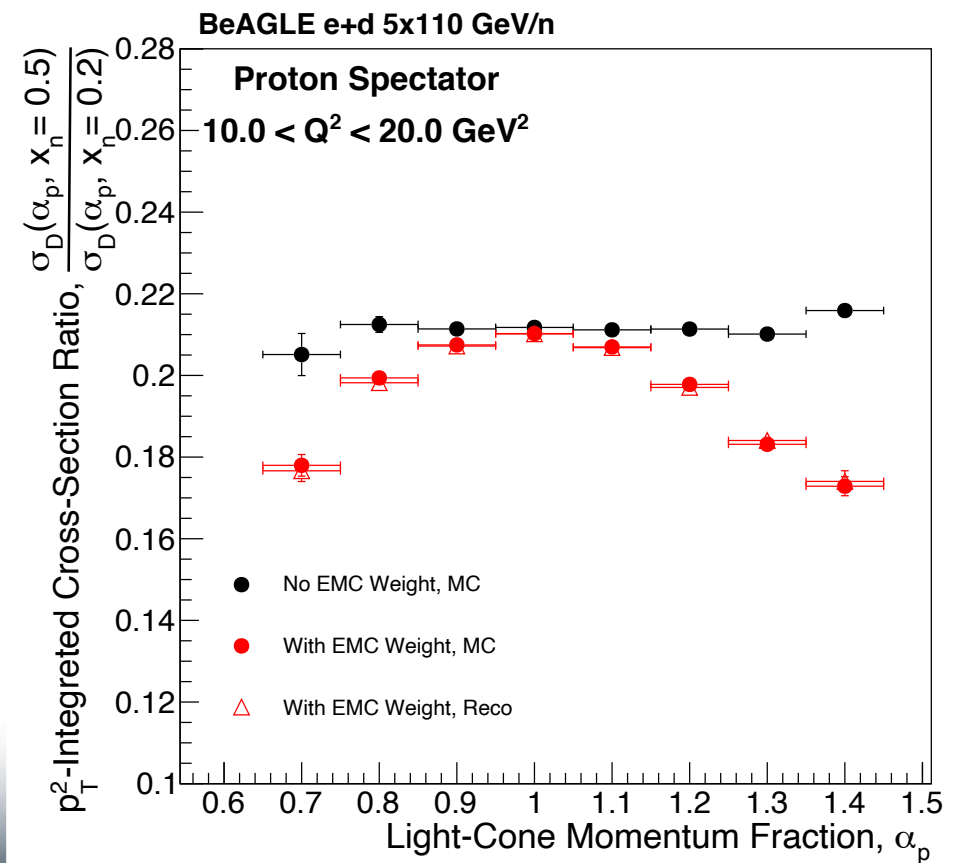


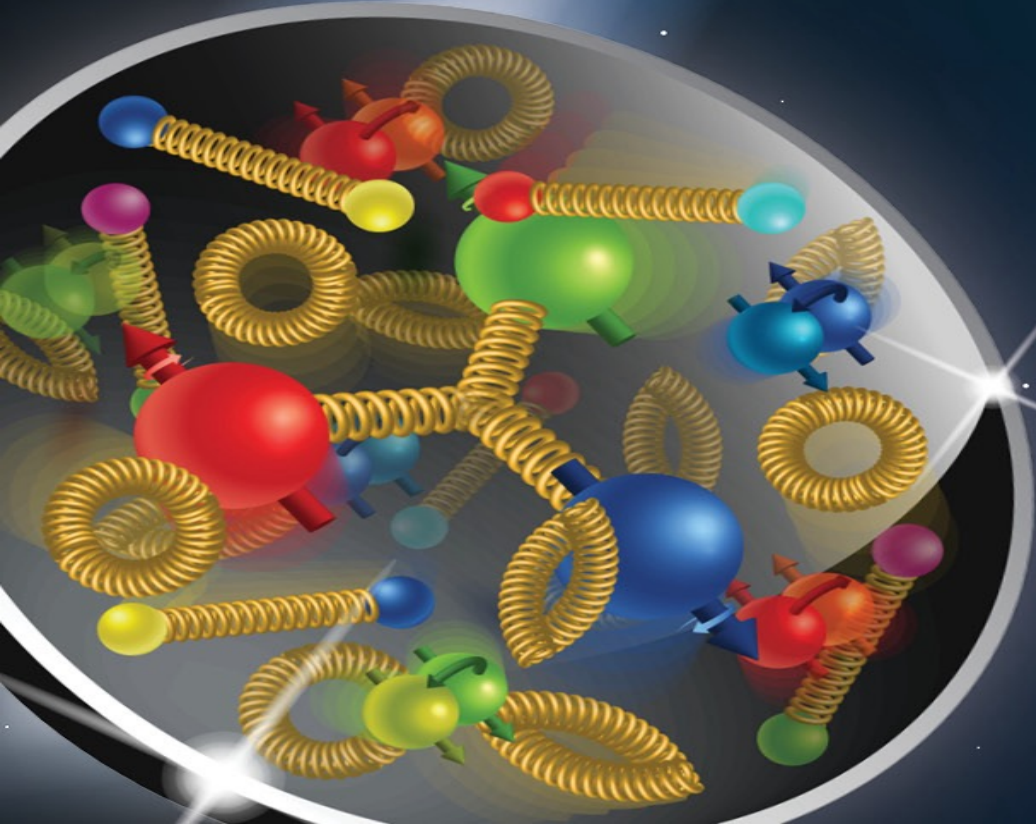
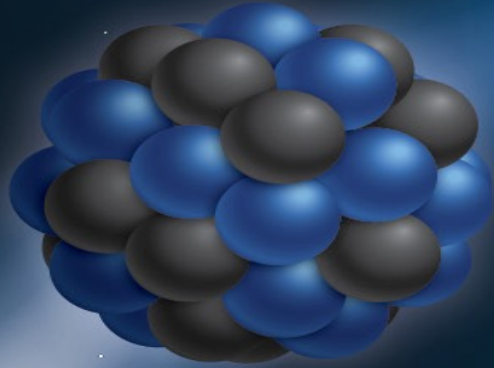
Different nuclear configurations

Study of FSI and comparisons in-progress (see backup).



Integrate cross section
over $p_{T,p}^2$ in each α bin.





What about $e + (\text{heavy})A$?

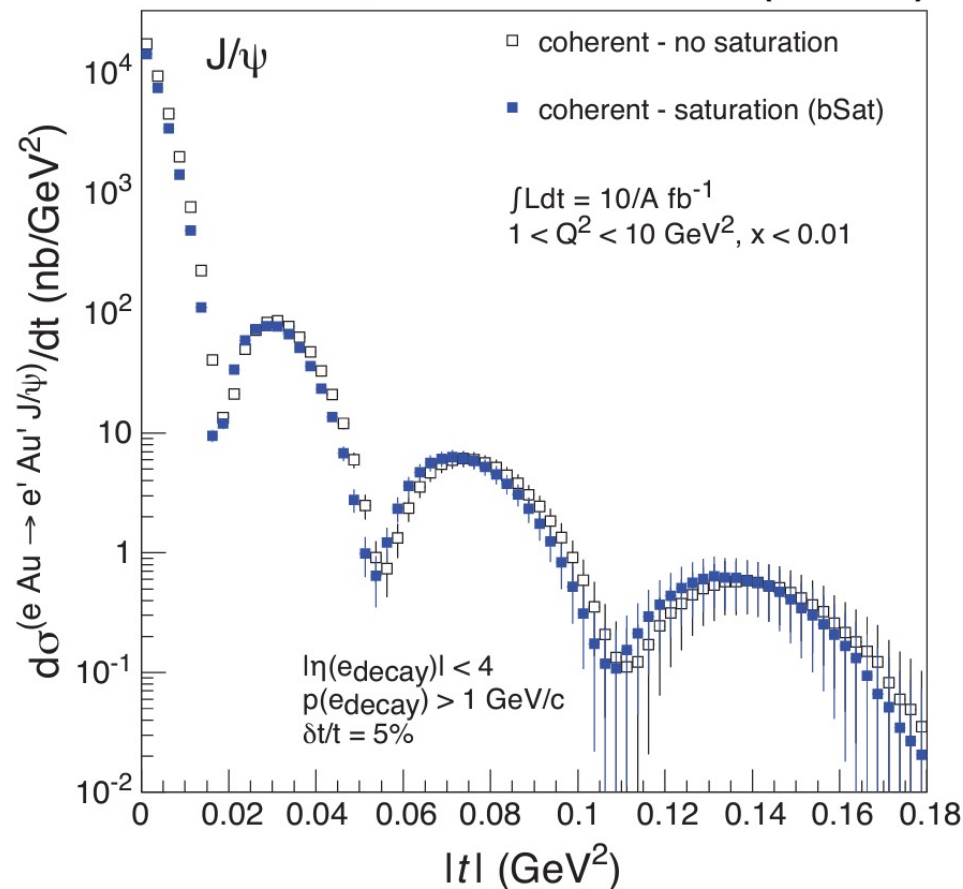
J/ψ in $e+A$ as a tool for nuclear imaging

EIC White Paper: golden channel

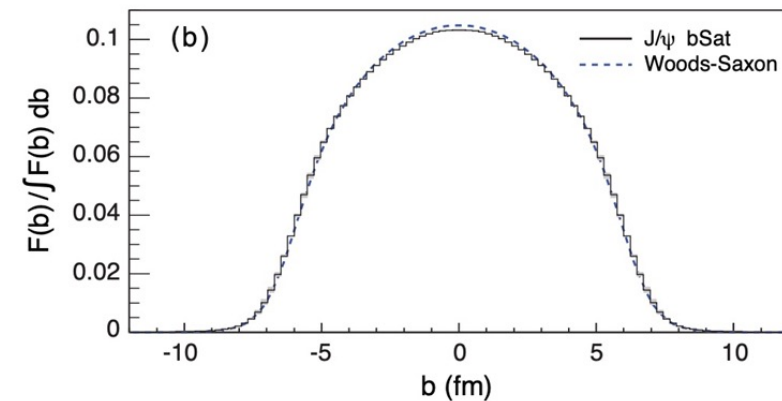
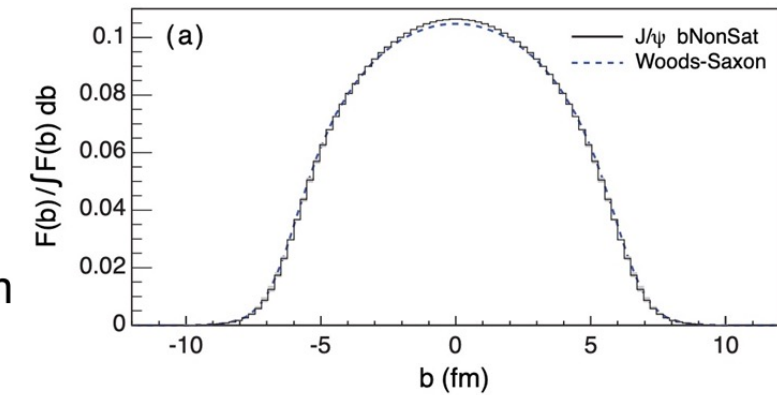
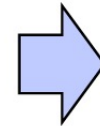
From K. Tu @ DIS 2023:

<https://indico.cern.ch/event/1199314/contributions/5189840/attachments/2621029/4531556/ePIC-exclusive-slides-Tu-v3.pdf>

EIC WP, Toll & Ullrich (2012)

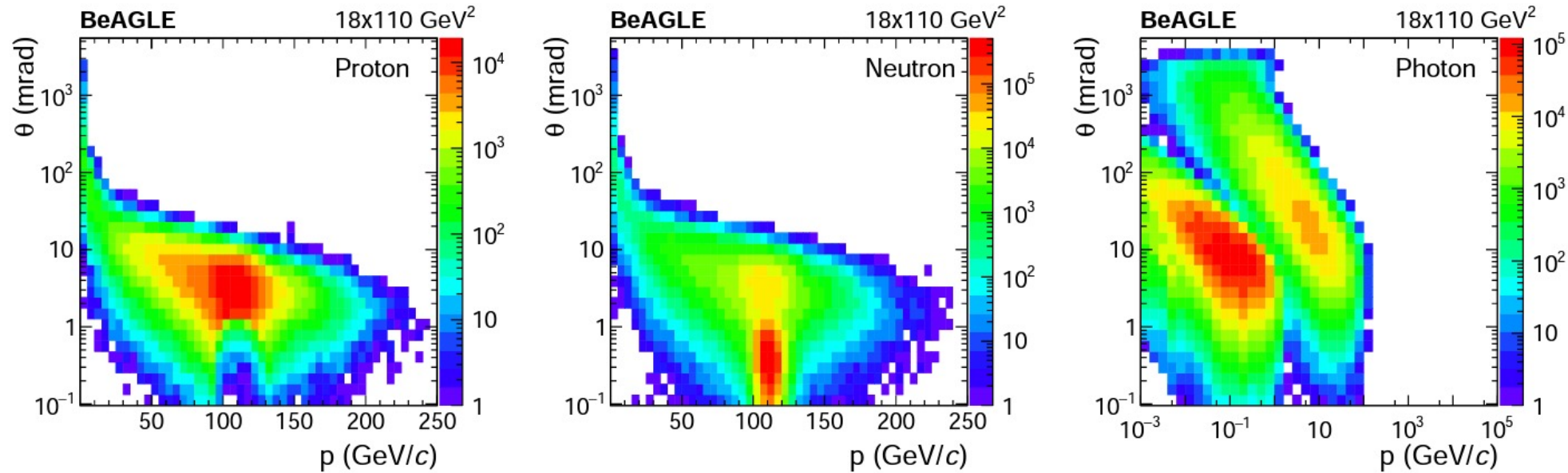


Fourier Transform



Vetoing Incoherent Background Using FF

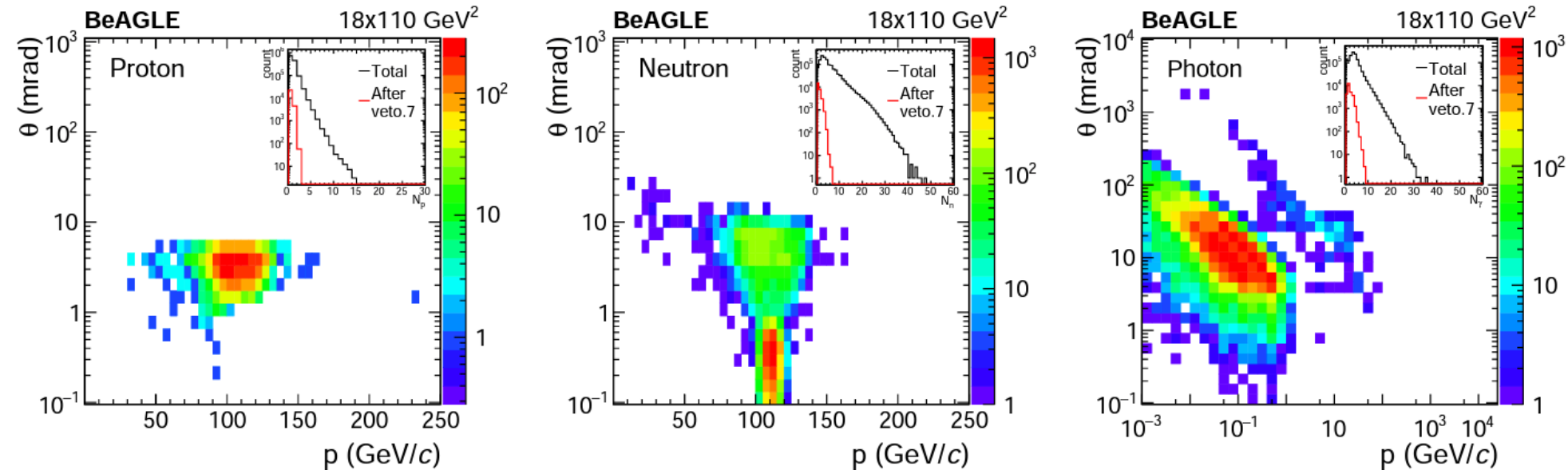
Process in BeAGLE: incoherent J/ψ production $e+Pb \rightarrow e + J/\psi + X$ (18x110 GeV/n)



Production of protons, neutrons, and photons from breakup of lead nucleus using BeAGLE.

Vetoing Incoherent Background Using FF

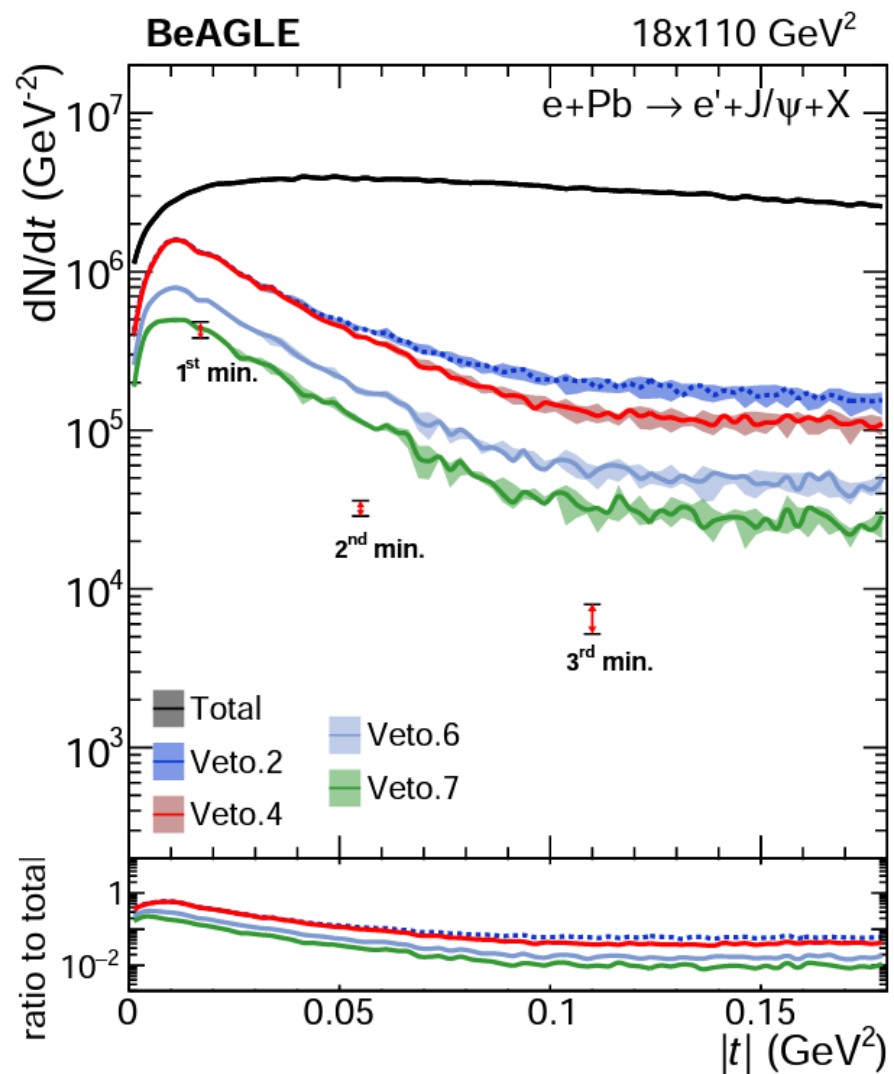
Process in BeAGLE: incoherent J/ψ production $e+Pb \rightarrow e + J/\psi + X$ (18x110 GeV/n)



Tagging of protons, neutrons, and photons from breakup of lead nucleus using BeAGLE in **FF detectors**.

Particles tagged in all 4 far-forward detectors.

Vetoing Incoherent Background Using FF



Different t-distribution curves are shown with different FF detectors vetoing particles. All vetoes shown here assume an exclusivity cut in the main detector (scattered electron + J/psi).

ZDC neutron only

ZDC neutron + RP proton + OMD proton

ZDC neutron + RP proton + OMD proton + B0 proton and photon

ZDC neutron + RP proton + OMD proton + B0 proton and photon + ZDC photon (> 50 MeV)

The “best” curve represents the FULL far-forward acceptance, but also includes an unoptimized beam pipe which can reduce the efficiency for the ZDC (especially low E photons).

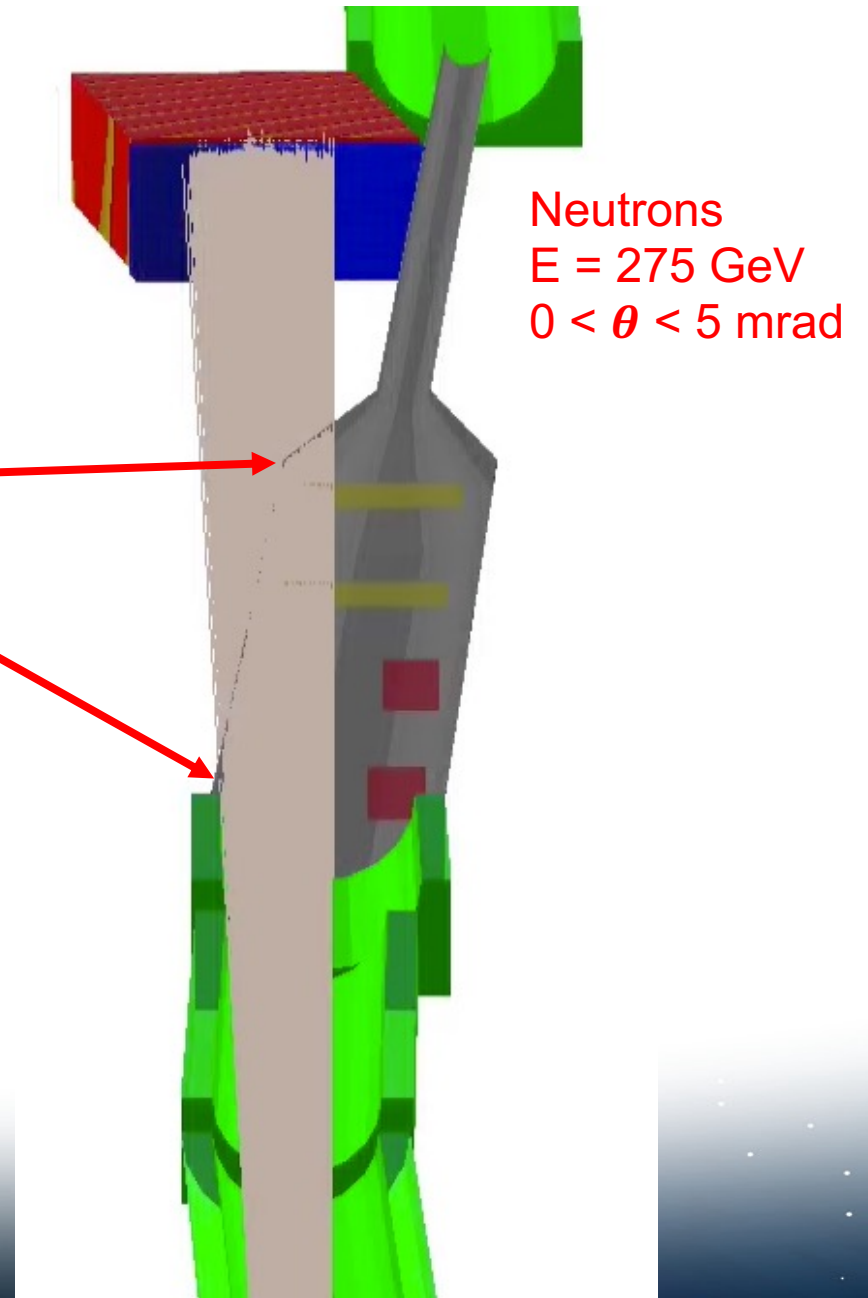
ZDC & neutral particle exit

Want to have as large an incident angle with the beam pipe as possible.

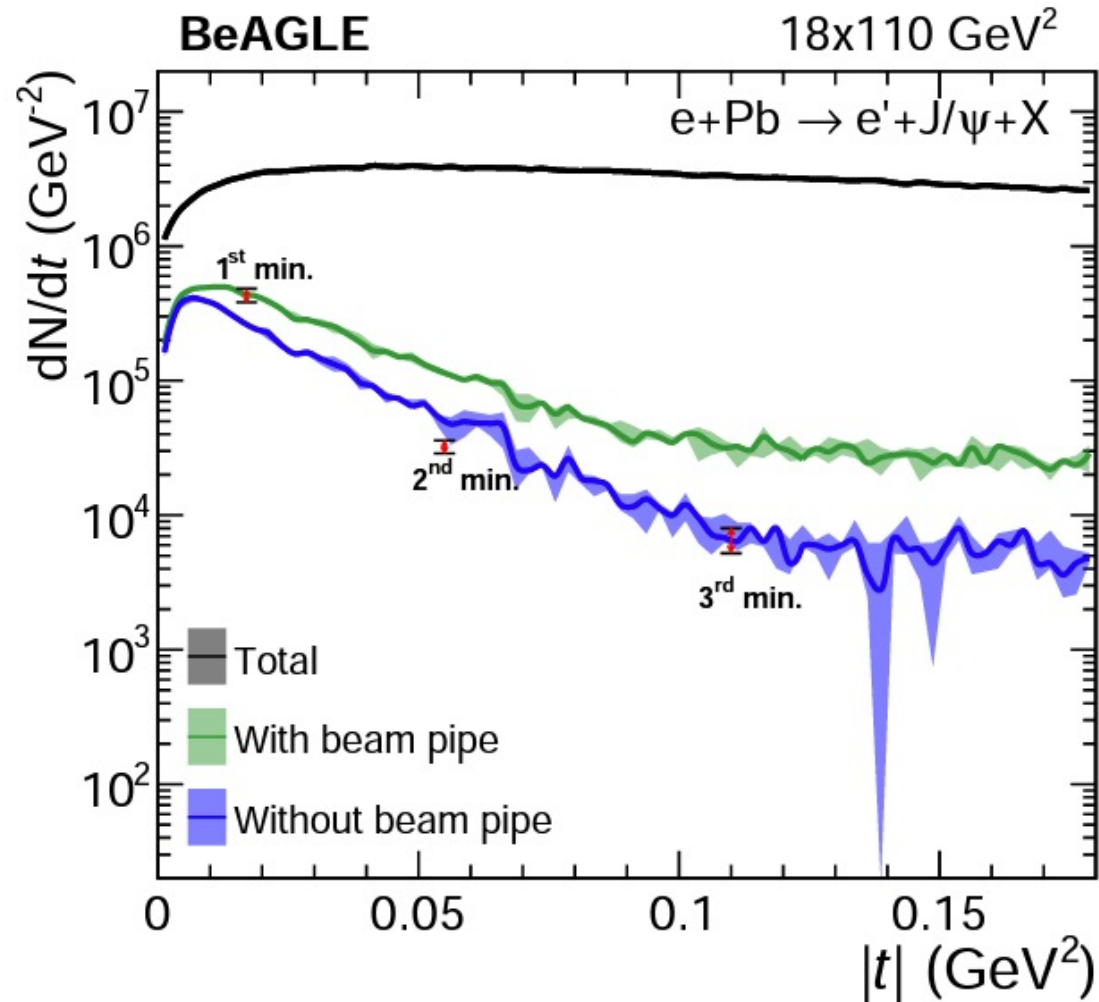
This is the problem area → shallow incident angle can increase effective material thickness by ~ factor of 10!!

This will reduce our detection efficiency beyond just the aperture limit!

➤ Updated design in-production.



Vetoing Incoherent Background Using FF

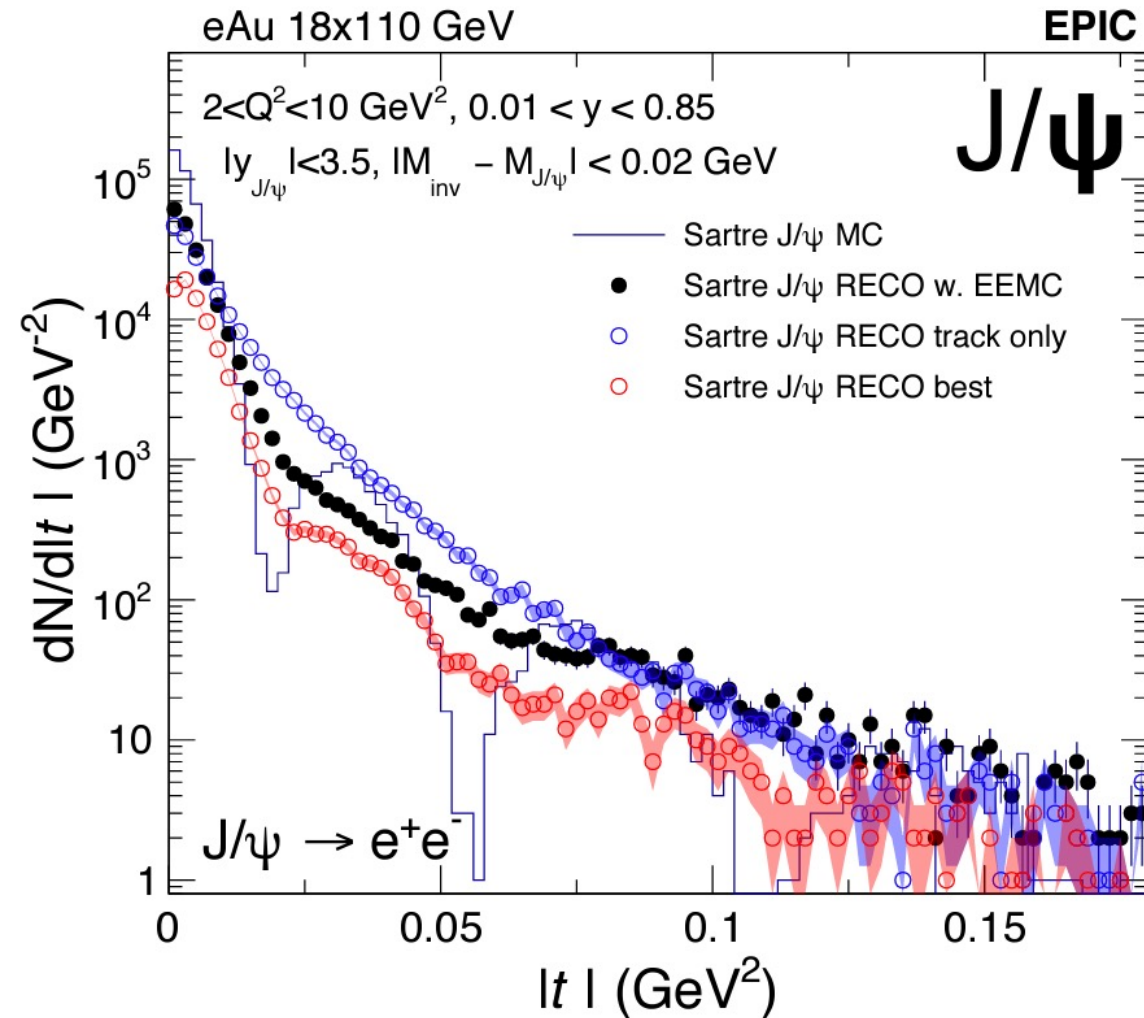


Optimizing the beam pipe exit for neutral particles will have a major impact on the vetoing efficiency (in-progress now!).

Measuring t-distribution \rightarrow Full ePIC simulations

From K. Tu @ DIS 2023:

<https://indico.cern.ch/event/1199314/contributions/5189840/attachments/2621029/4531556/ePIC-exclusive-slides-Tu-v3.pdf>

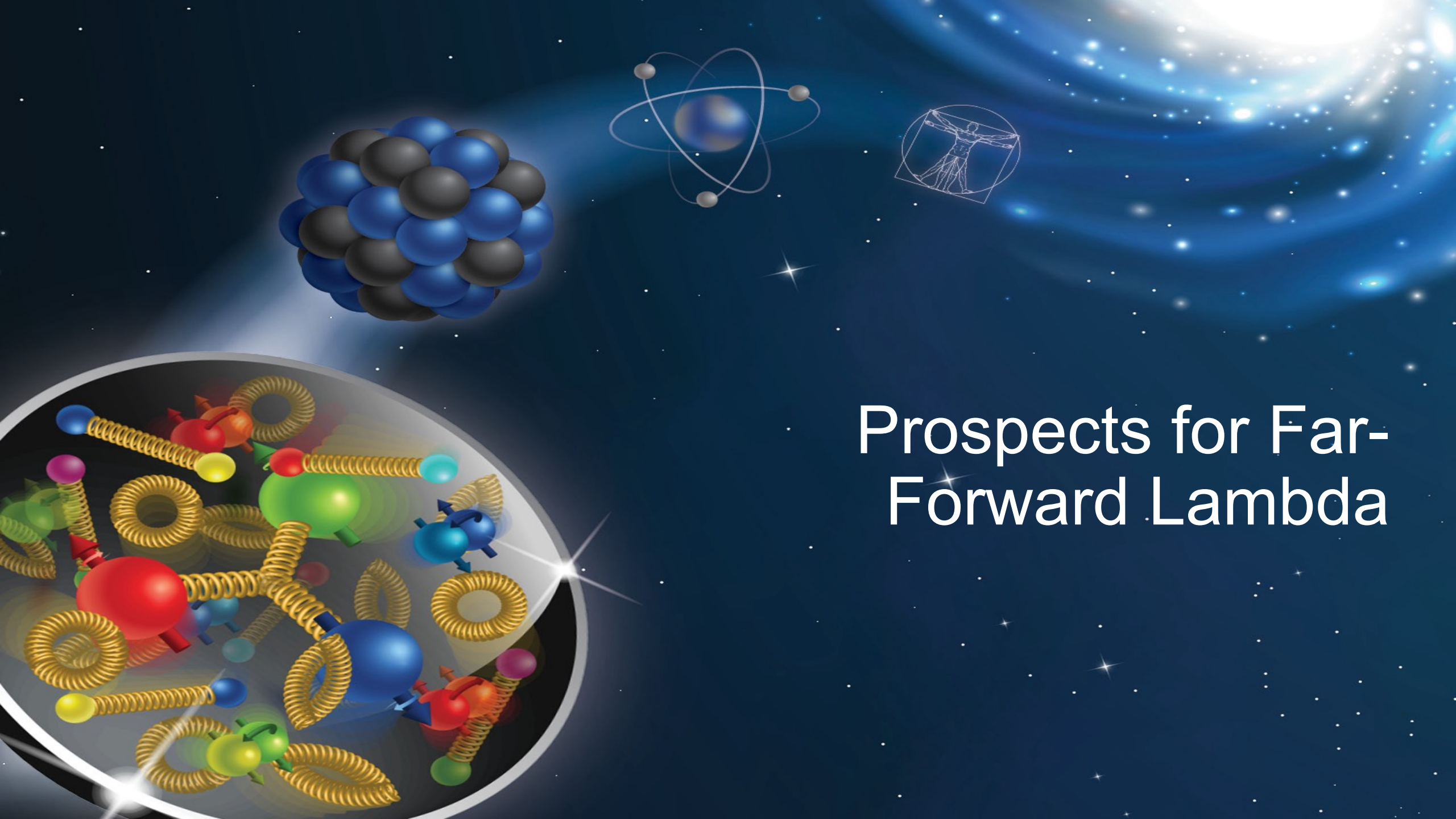


Legend details:

- w. EEMC: electron energy from EEMC, electron mass (PDG), angle (eta, phi) from tracking; $\varphi \rightarrow KK$ from tracking.
- Track only: e' , $\varphi \rightarrow KK$, all from tracking
- Best: average of the above 2 E-by-E.

Improvements from *algorithm*:

- The two methods can be used together to further improve the $|t|$ resolution.

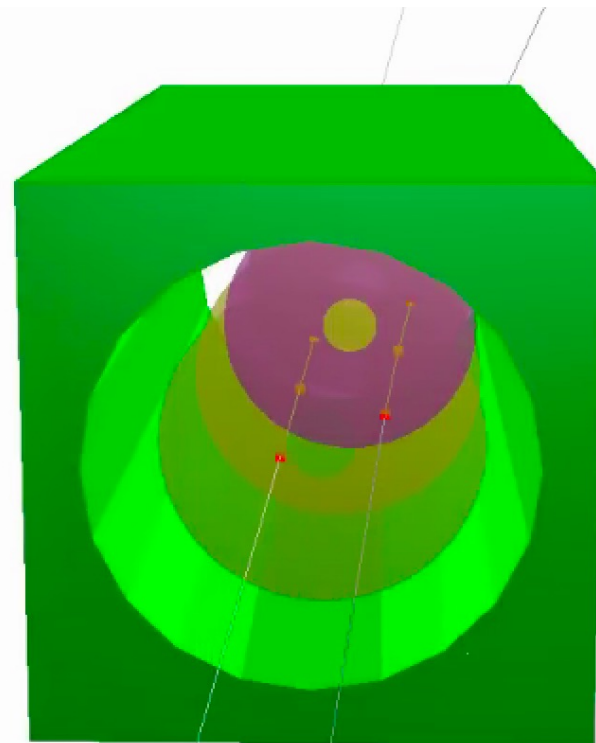
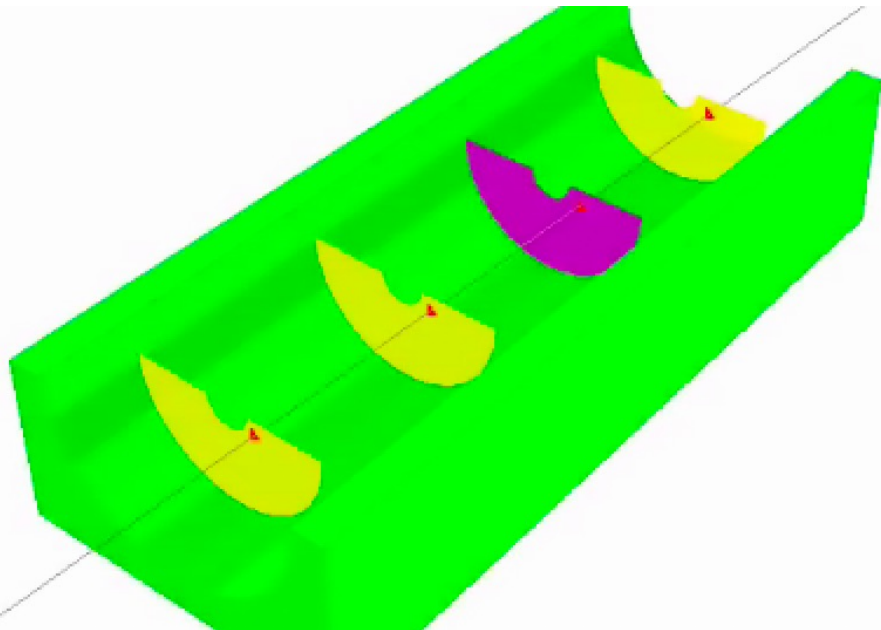


Prospects for Far-Forward Lambda

The importance of the B0 for the meson program

- Needed for measuring final states with $\theta > 5.5$ mrad.
 - Especially important at medium and low hadron beam energies at the EIC.
- Important for incoherent vetoing in e+A (heavy nuclear) collisions.
 - Charged particles and photons.
- The B0 tracking system behaves like a normal spectrometer, so anything which decays with particles in its acceptance can be reconstructed just like in the forward tracking disks!

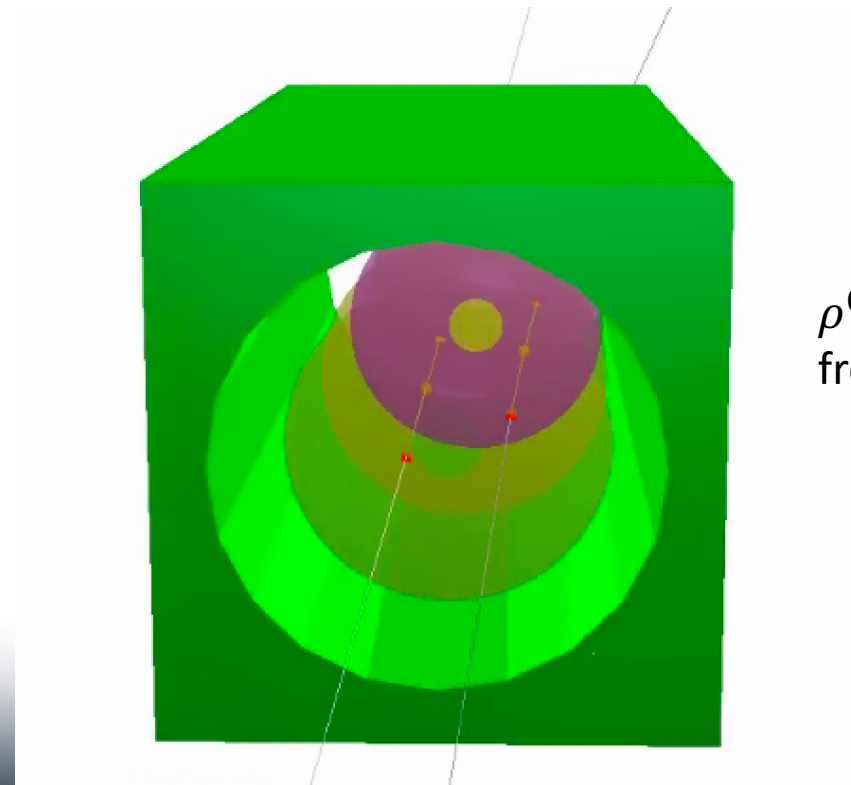
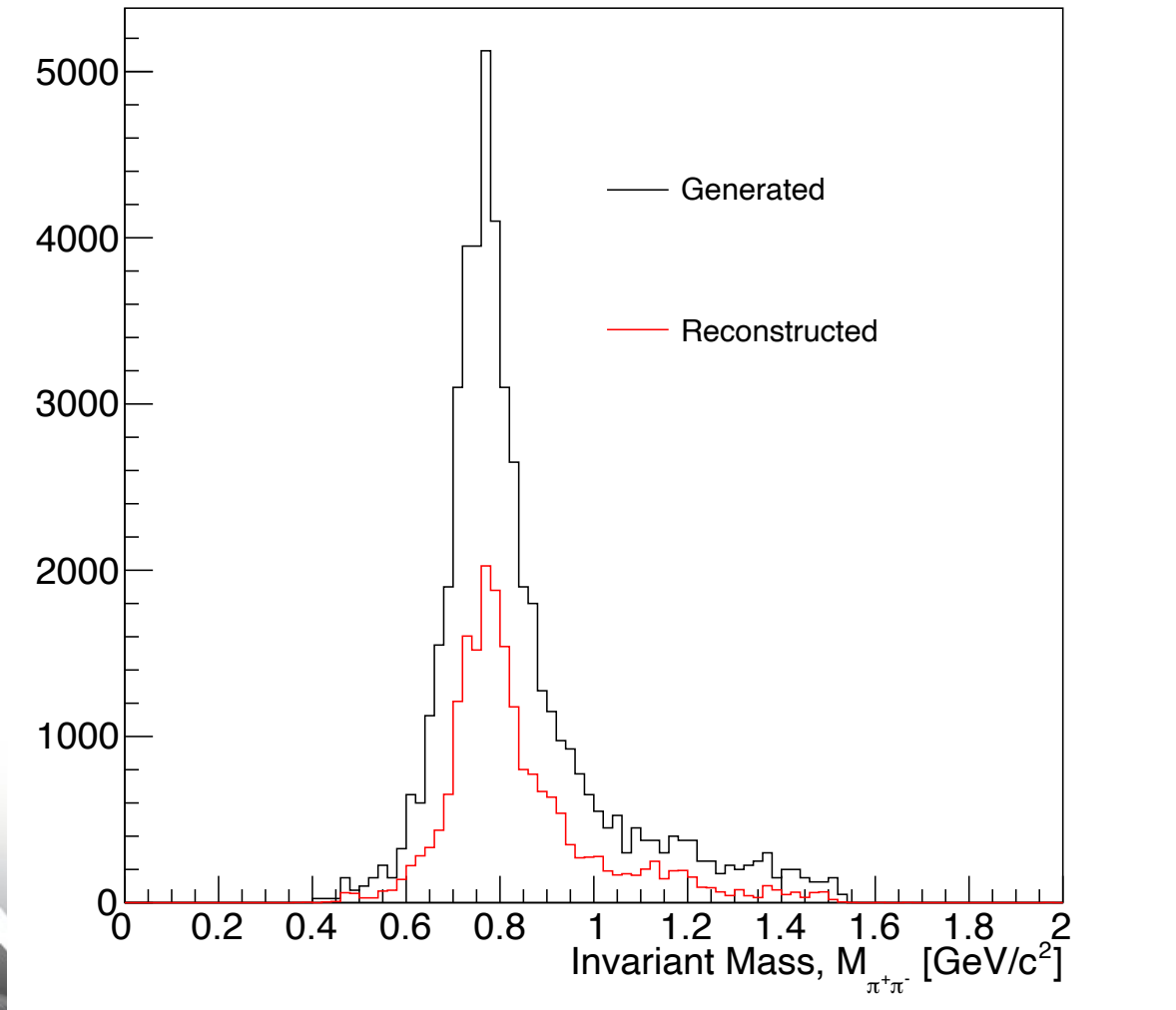
GEANT simulation: 100 GeV proton



$\rho^0 \rightarrow \pi^+ \pi^-$ decay
from u-channel production

The importance of the B0 for the meson program

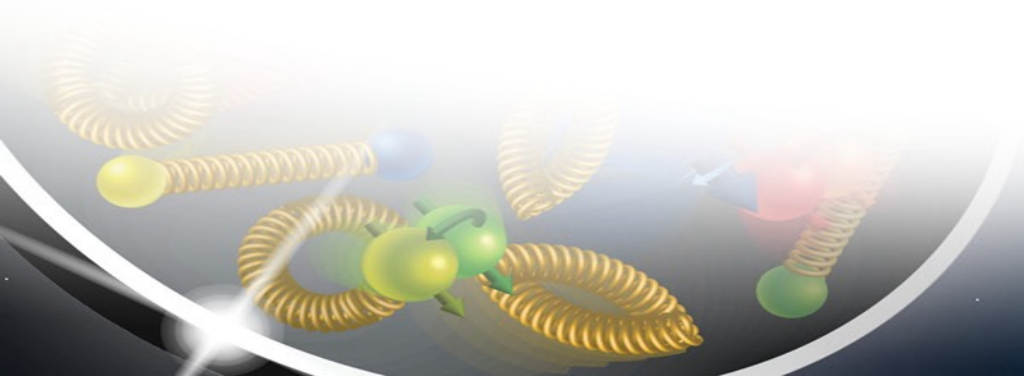
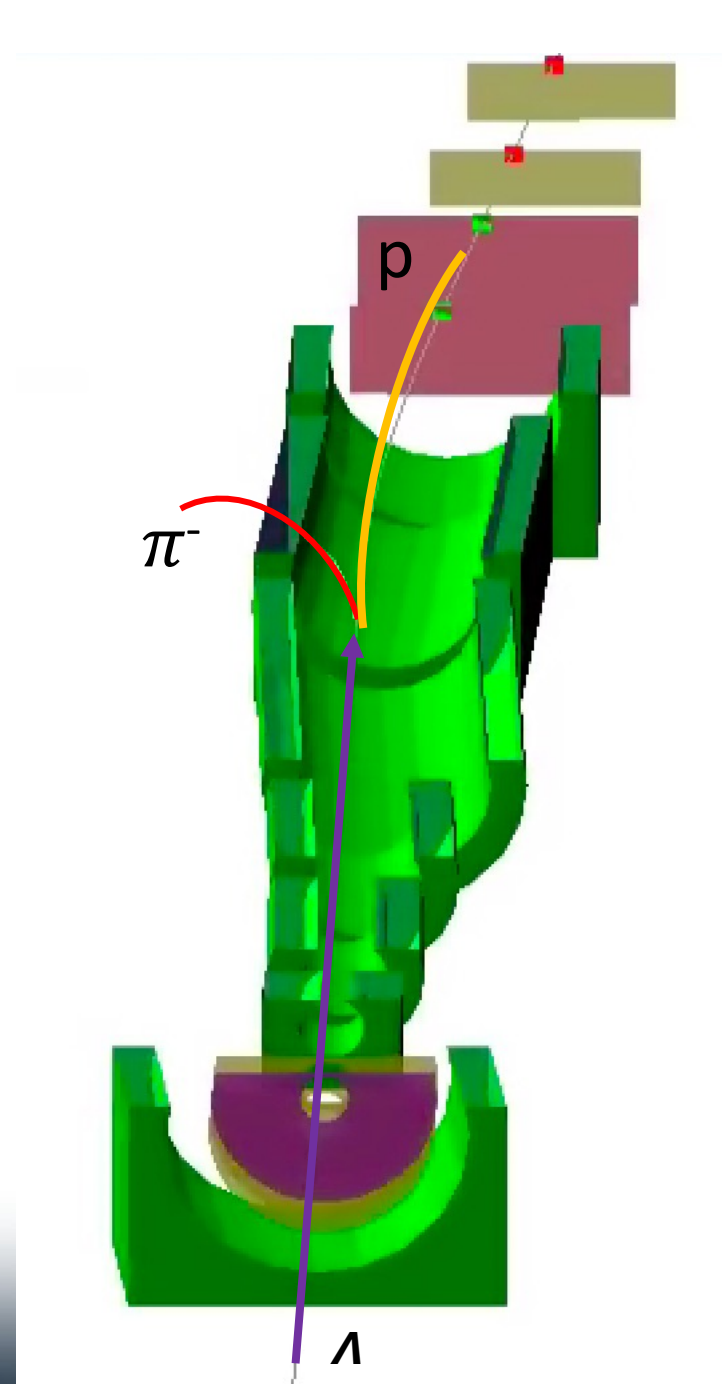
- $\rho^0 \rightarrow \pi^+ \pi^-$ decay studied with eSTARLight 5x41 events (generated by Zach Sweger).
- Reconstruction performed with EicRoot.



$\rho^0 \rightarrow \pi^+ \pi^-$ decay
from u-channel production

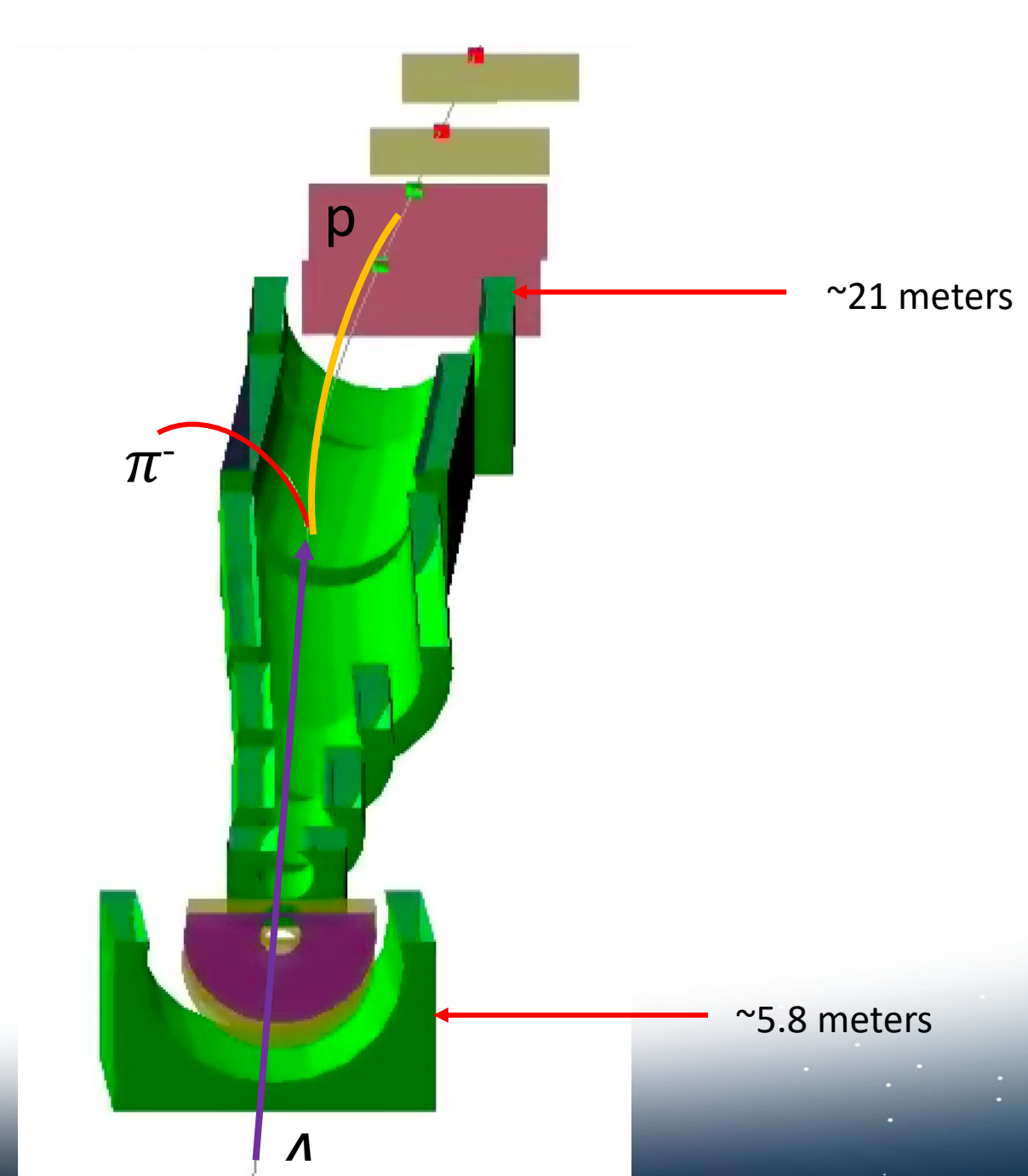
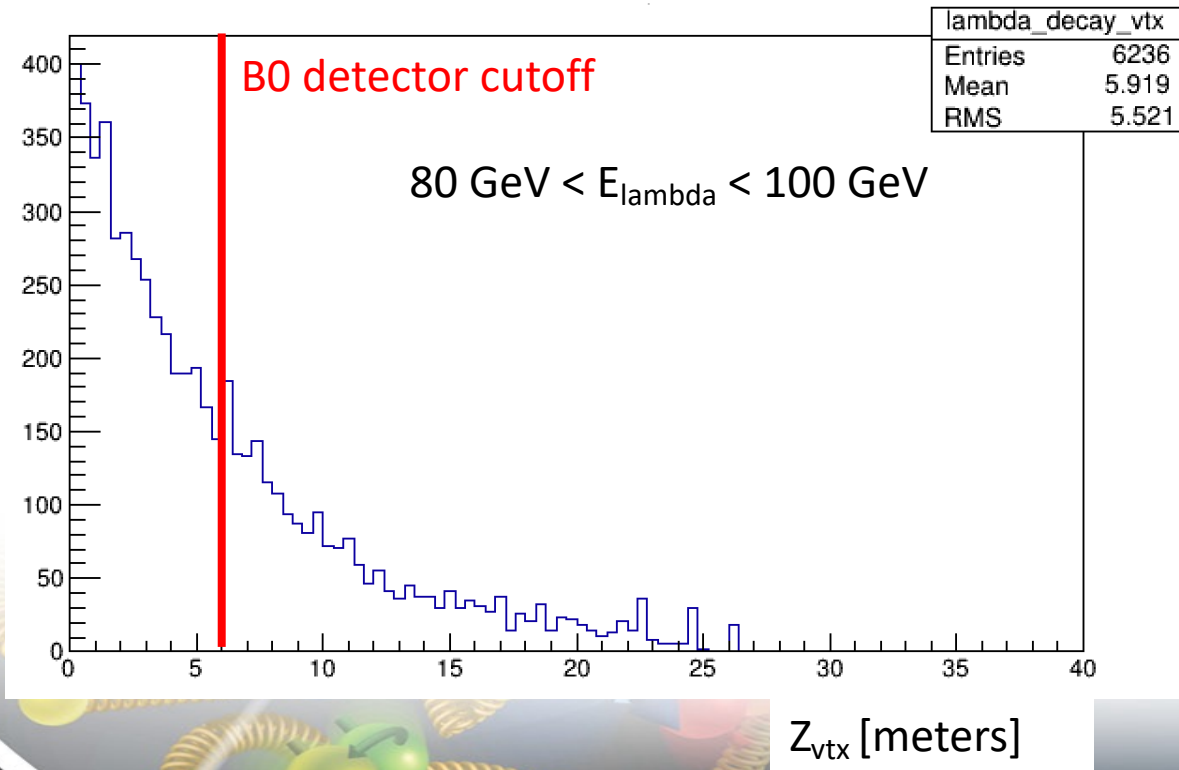
Lambda Decay ($p + \pi^-$)

- Boost causes the lambda to be able to decay 10s of meters from the IP.
 - Significant problem since reconstruction of this displaced secondary vertex within the hadron magnets is very challenging.



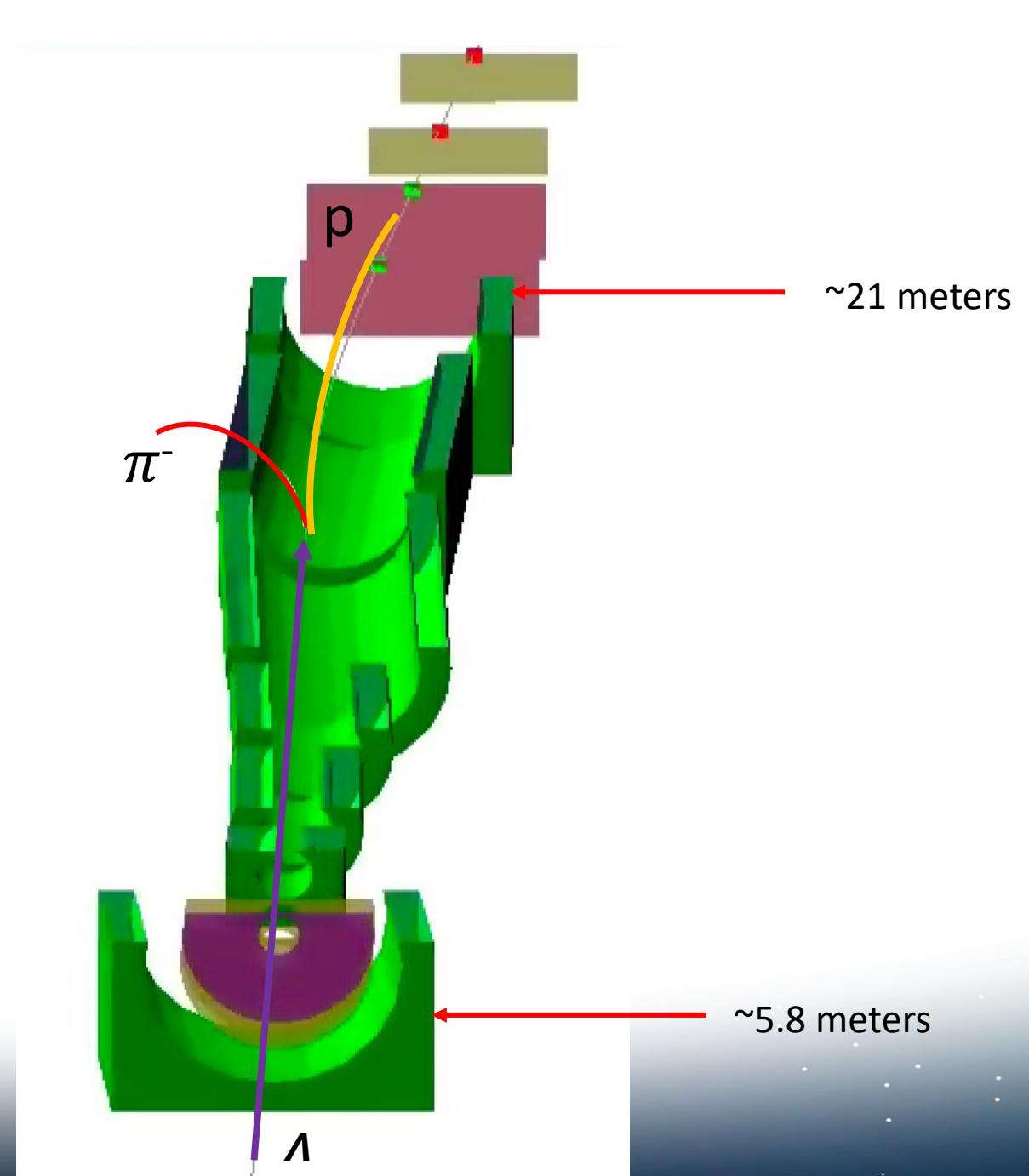
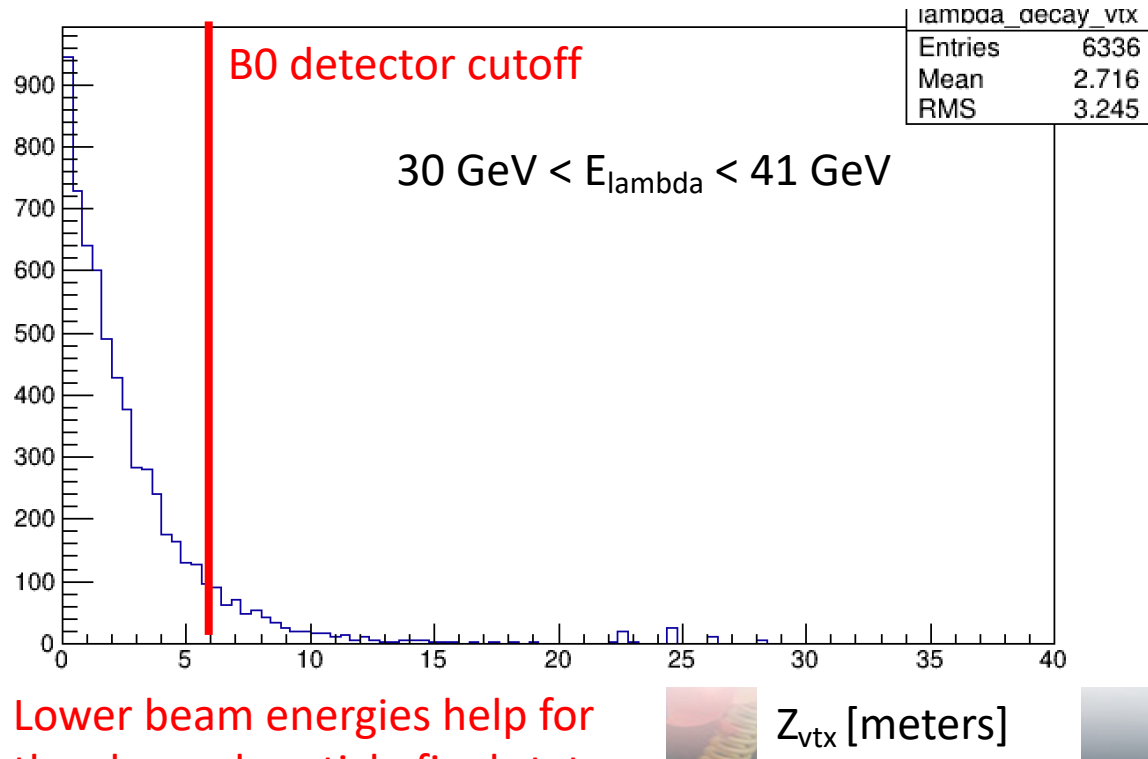
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Lambda Decay ($p + \pi^-$)

- Boost causes the lambda to be able to decay 10s of meters from the IP.
 - Significant problem since reconstruction of this displaced secondary vertex within the hadron magnets is very challenging.



(some) Far-Forward **Physics** at the EIC

e+p DVCS

Proton spin: orbital angular momentum; imaging

e+d exclusive J/Psi with p/n tagging

Short-Range Correlations

spectator tagging in light nuclei

Free neutron structure, EMC effect, etc.

coherent/incoherent J/psi production in e+A

Saturation

Sullivan process

pi/K form factors and structure functions

Quasi-elastic electron scattering

Short-Range Correlations

[1] Z. Tu, A. Jentsch, et al., Physics Letters B, (2020)
 [2] I. Friscic, D. Nguyen, J. R. Pybus, A. Jentsch, et al., Phys. Lett. B, **Volume 823**, 136726 (2021)
 [3] W. Chang, E.C. Aschenauer, M. D. Baker, A. Jentsch, J.H. Lee, Z. Tu, Z. Yin, and L.Zheng, Phys. Rev. D **104**, 114030 (2021)
 [4] A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104**, 065205, (2021) (Editor's Suggestion)

u-channel backward exclusive electroproduction

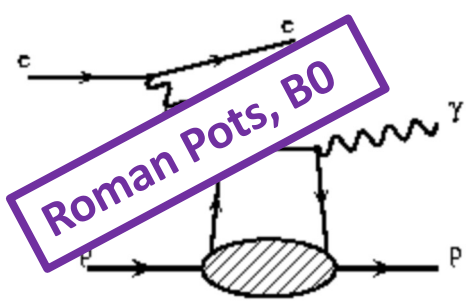
Backward-angle colinear factorization

...and **MANY** more!



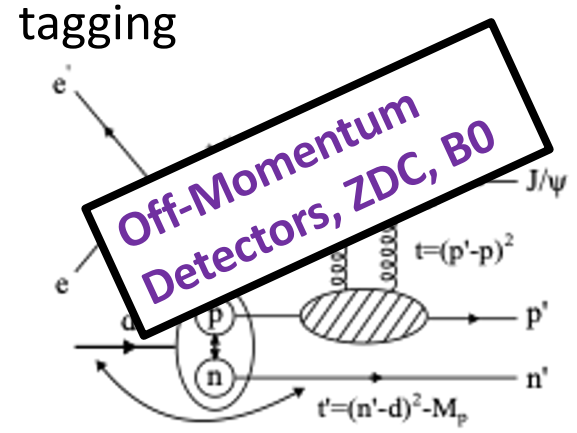
(some) Far-Forward **Physics** at the EIC

e+p DVCS



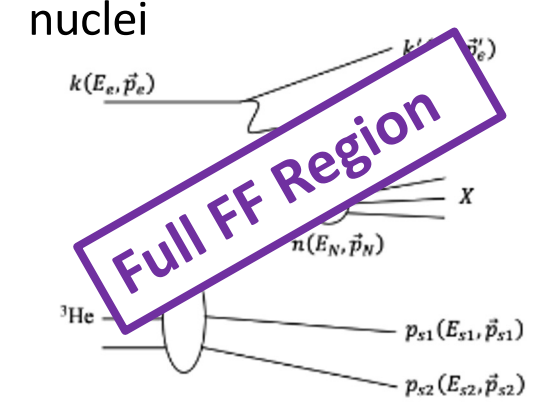
Roman Pots, B0

e+d exclusive J/Psi with p/n tagging



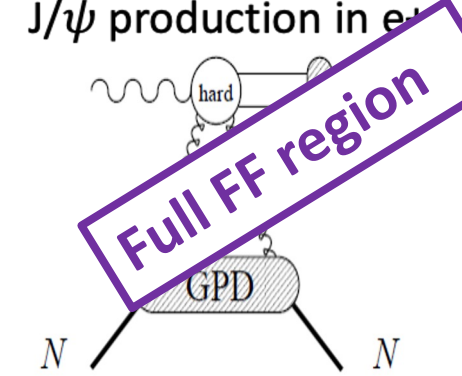
Off-Momentum Detectors, ZDC, B0

spectator tagging in light nuclei



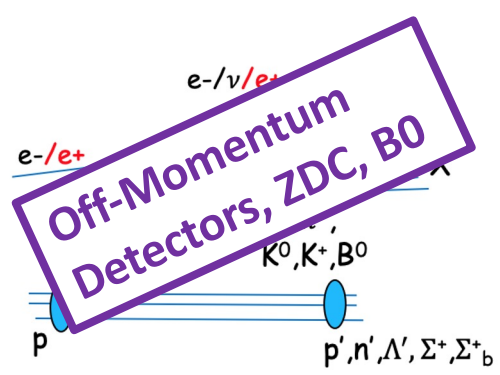
Full FF Region

coherent/incoherent J/psi production in e+



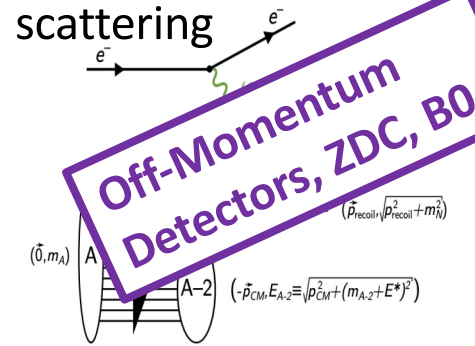
Full FF region

Sullivan process



Off-Momentum Detectors, ZDC, B0

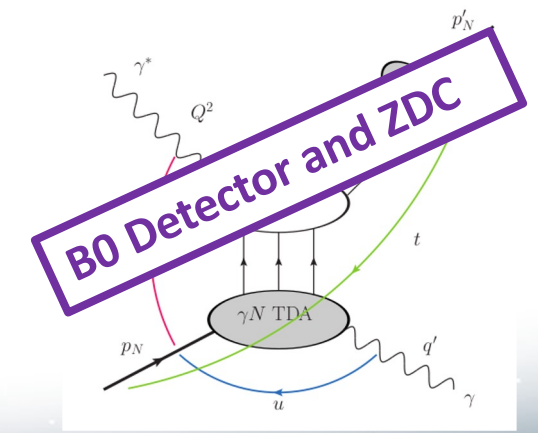
Quasi-elastic electron scattering



Off-Momentum Detectors, ZDC, B0

- [1] Z. Tu, A. Jentsch, et al., Physics Letters B, (2020)
- [2] I. Friscic, D. Nguyen, J. R. Pybus, A. Jentsch, et al., Phys. Lett. B, **Volume 823**, 136726 (2021)
- [3] W. Chang, E.C. Aschenauer, M. D. Baker, A. Jentsch, J.H. Lee, Z. Tu, Z. Yin, and L.Zheng, Phys. Rev. D **104**, 114030 (2021)
- [4] A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104**, 065205, (2021) (**Editor's Suggestion**)

u-channel backward exclusive electroproduction



B0 Detector and ZDC

...and **MANY** more!

Summary and Takeaways

- Far-forward physics characterized by exclusive+diffractive final states.
 - Lots to unpack! – proton spin, neutron structure, saturation, partonic imaging, meson structure, etc.
- There is lots of interest in the EIC community in studying this physics via these final states!
 - Exciting time to get involved!!

Email me if you have any questions: ajentsch@bnl.gov

Want to get involved?? Join our meetings and learn how!

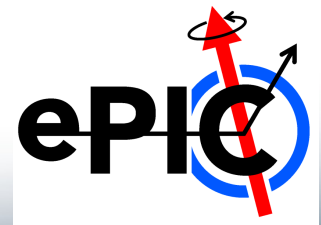
Meeting time: Tuesdays @ 9am EDT (bi-weekly, or weekly, as needed)

Indico: <https://indico.bnl.gov/category/407/>

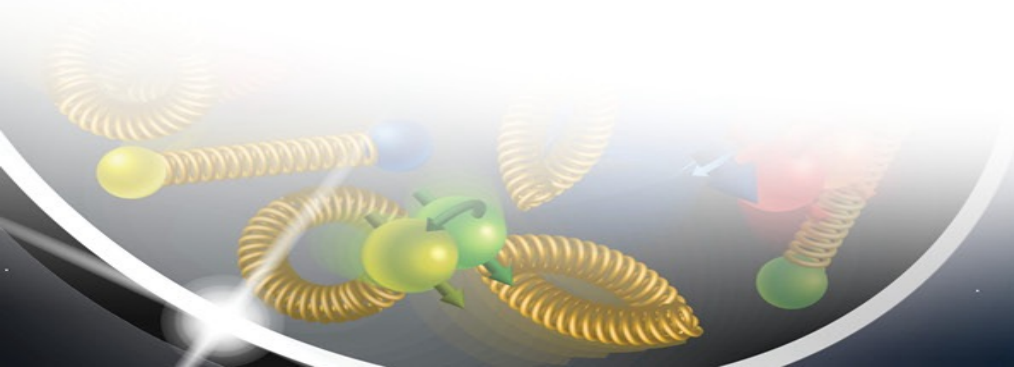
Wiki: <https://wiki.bnl.gov/eic-project-detector/index.php?title=Collaboration>

Email-list: eic-projdet-FarForw-l@lists.bnl.gov

Subscribe to mailing list through: <https://lists.bnl.gov/mailman/listinfo/eic-projdet-farforw-l>

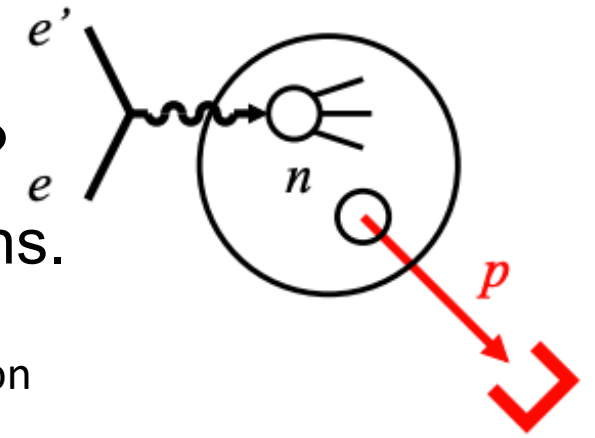


Backup



Neutron Structure

- Protons well-studied at HERA -> So...why the neutron?
 - Flavor separation, baseline for studies of nuclear modifications.

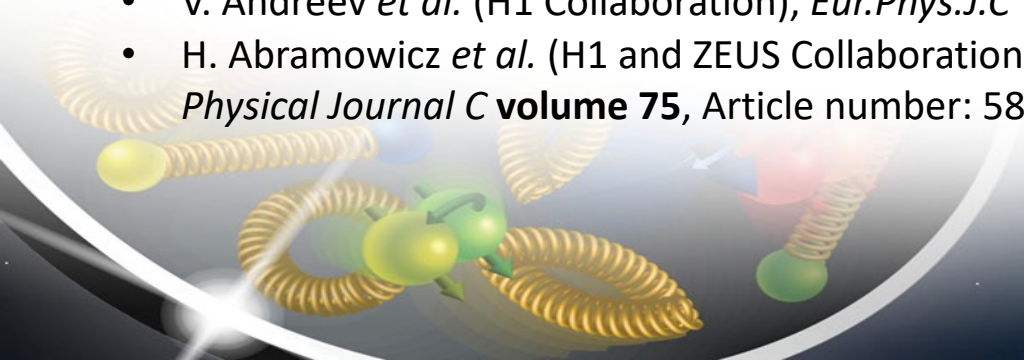
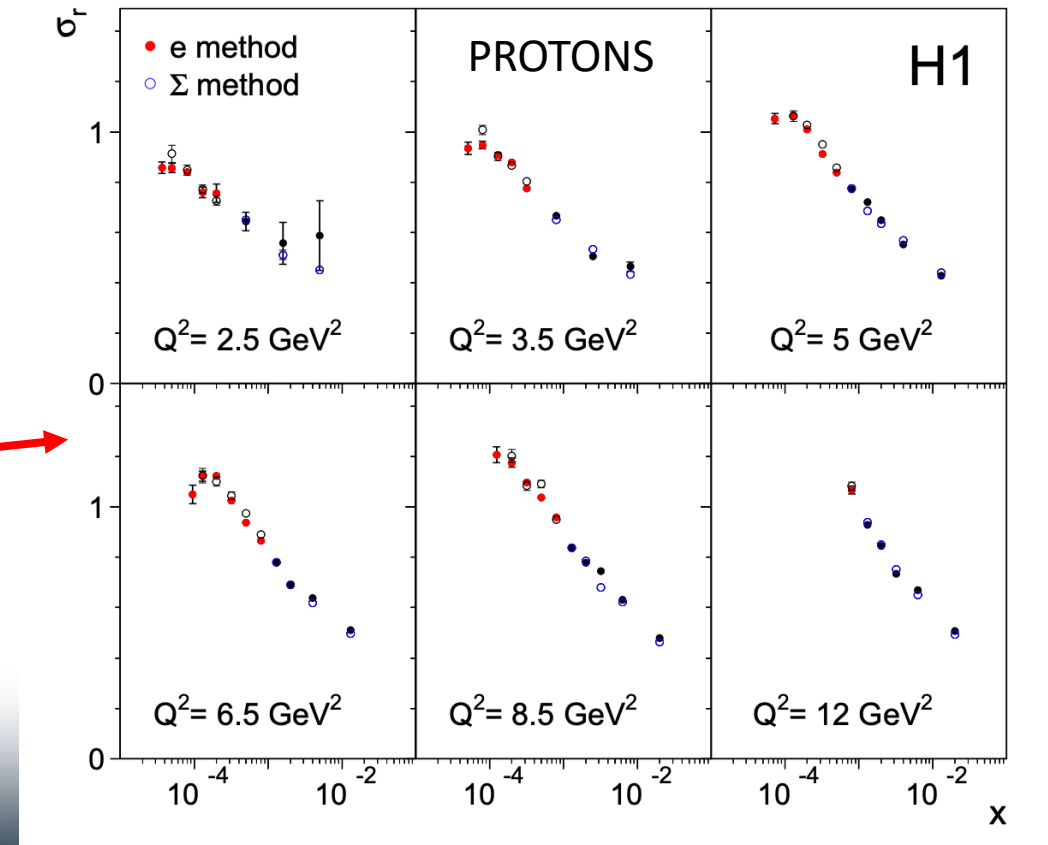


$$\sigma_r = \underbrace{\frac{Q^4 x}{2\pi\alpha^2 [1 + (1-y)^2]}}_{\text{"Flux factor"}} \cdot \underbrace{\frac{d^2\sigma}{dx dQ^2}}_{\text{Differential cross section}} = \underbrace{F_2(x, Q^2) - f(y) \cdot F_L(x, Q^2)}_{\text{Structure functions}}$$

Reduced cross section

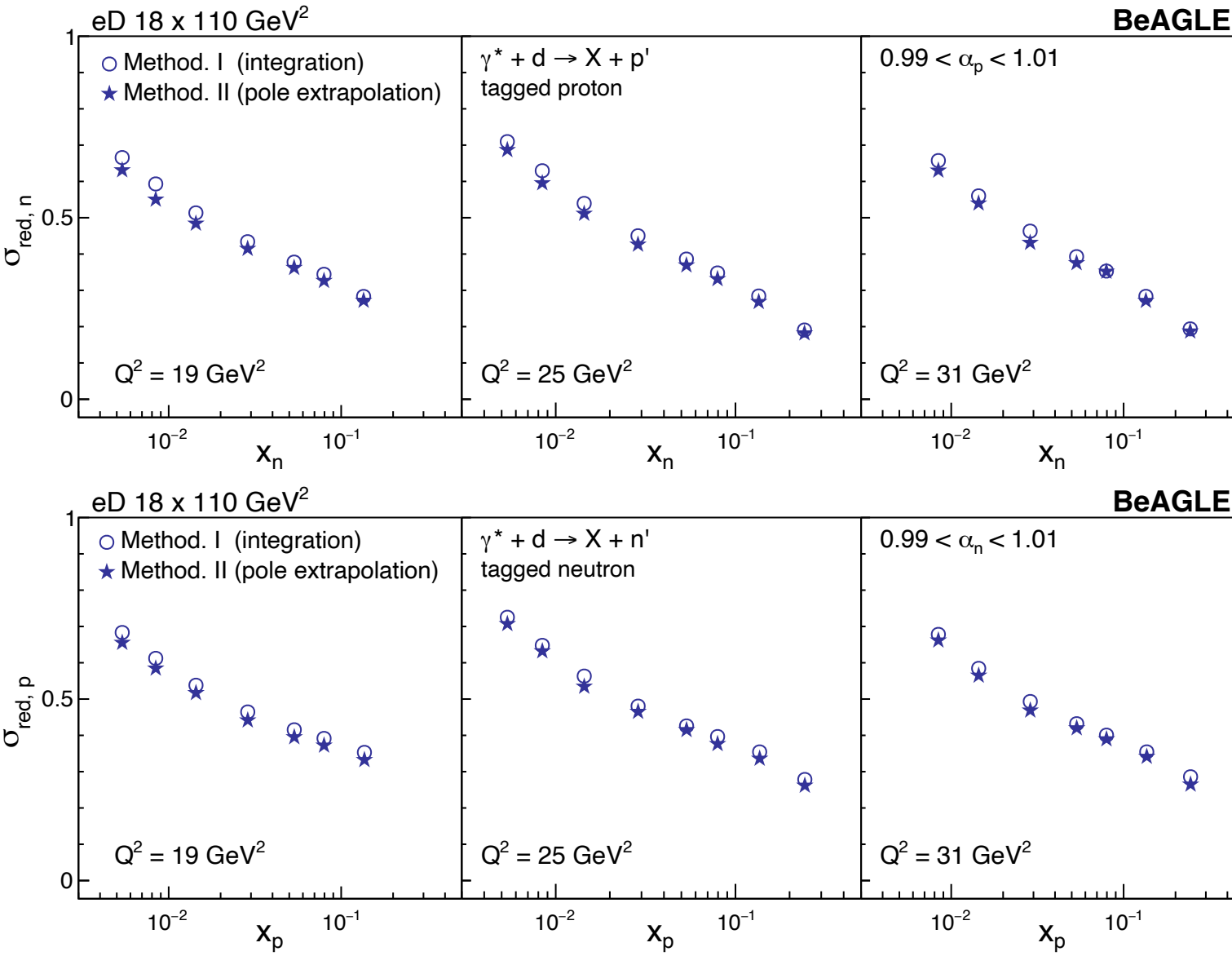
Some useful HERA references for measurements on proton

- F. Aaron *et al.* (H1 Collaboration), *The European Physical Journal C* volume 63, Article number: 625 (2009)
- V. Andreev *et al.* (H1 Collaboration), *Eur.Phys.J.C* 74 (2014) 4, 2814
- H. Abramowicz *et al.* (H1 and ZEUS Collaborations) *The European Physical Journal C* volume 75, Article number: 580 (2015)



Free Nucleon Structure

A. Jentsch, Z. Tu, and C. Weiss, Phys. Rev. C **104**, 065205, (2021) (Editor's Suggestion)



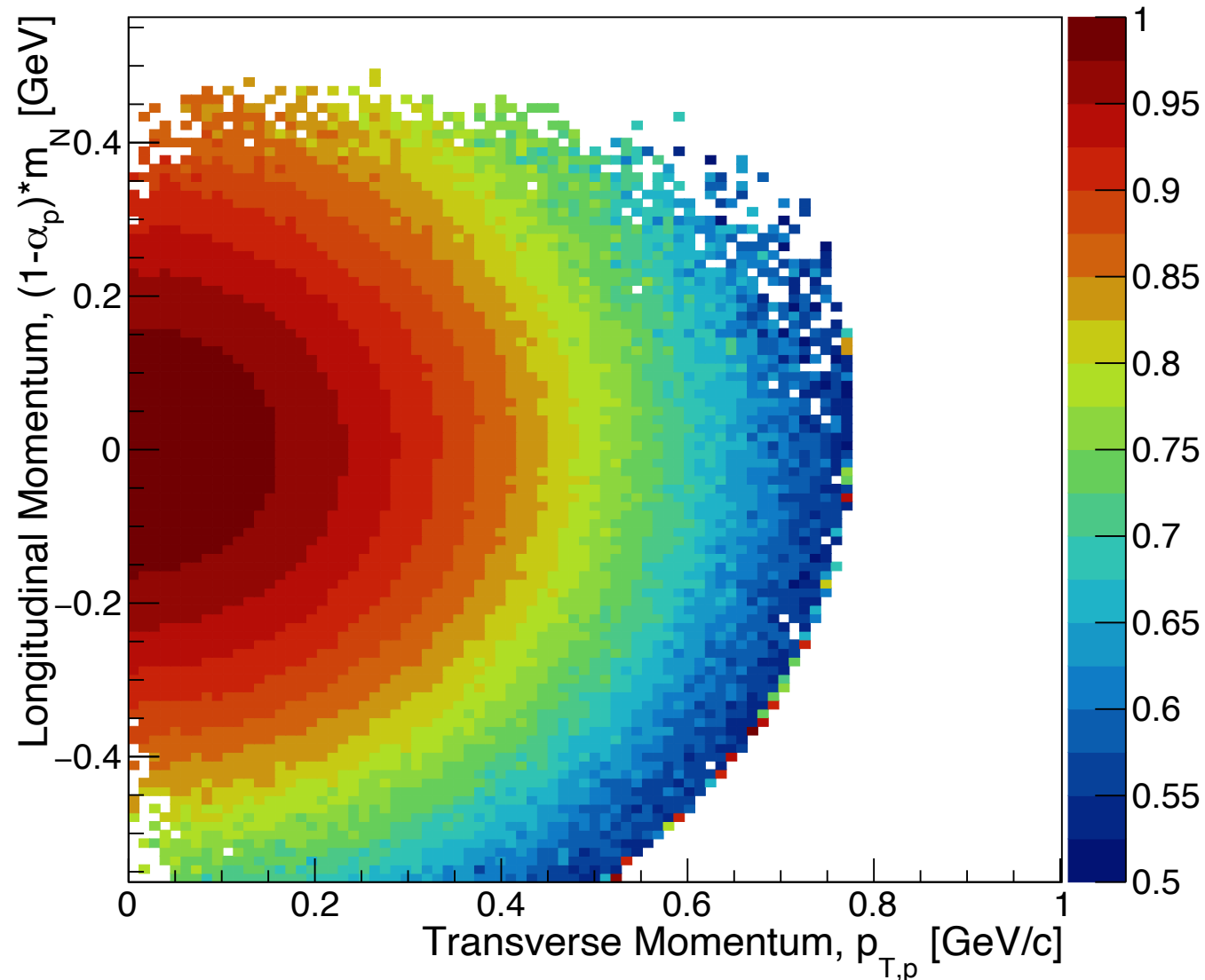
Open circles: “inclusive” measurement.
Stars: pole extrapolation procedure.

Differences driven by evaluation of pole (average in bin, vs. event-by-event).

- Similar kinds of high-precision results achievable as was done for proton F_2 at HERA!

Simulating the EMC Effect in BeAGLE

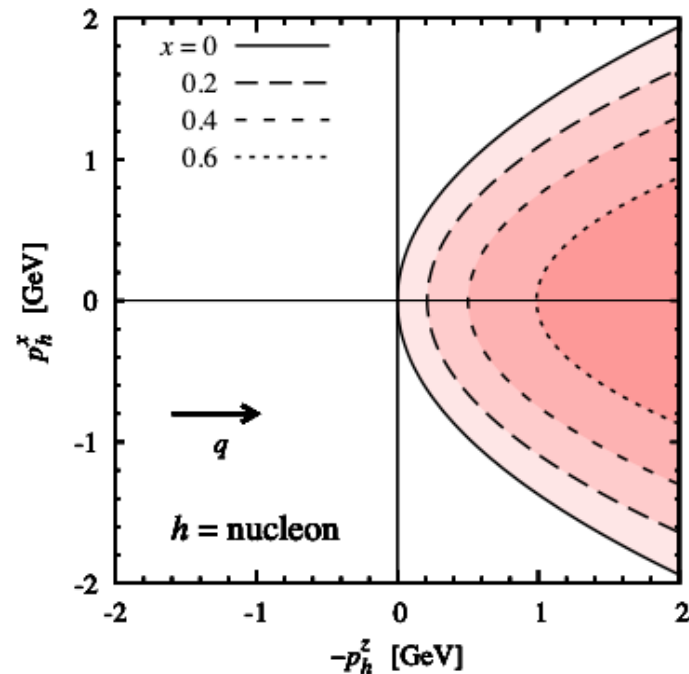
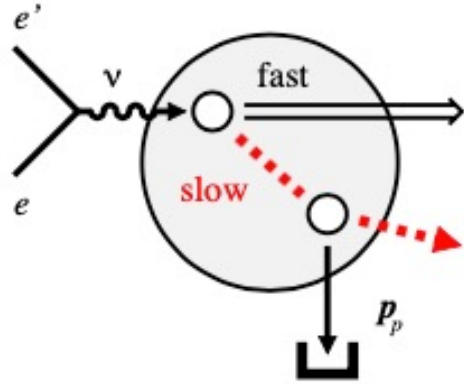
EMC Weight Distribution, $0.45 < x_n < 0.55$



Result → EMC Weight in BeAGLE

- Weight factor simulates the EMC effect from the *virtuality* in the deuteron.
- Applied event-by-event to compare **with and without weight** → enables study of sensitivity to EMC effect in various observables.

Final-State Interaction: Physical Picture



Space-time picture in deuteron rest-frame

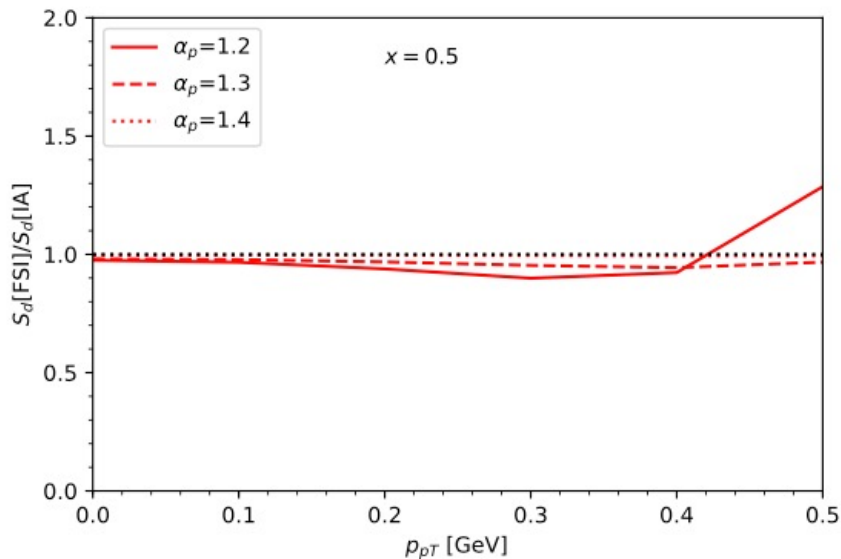
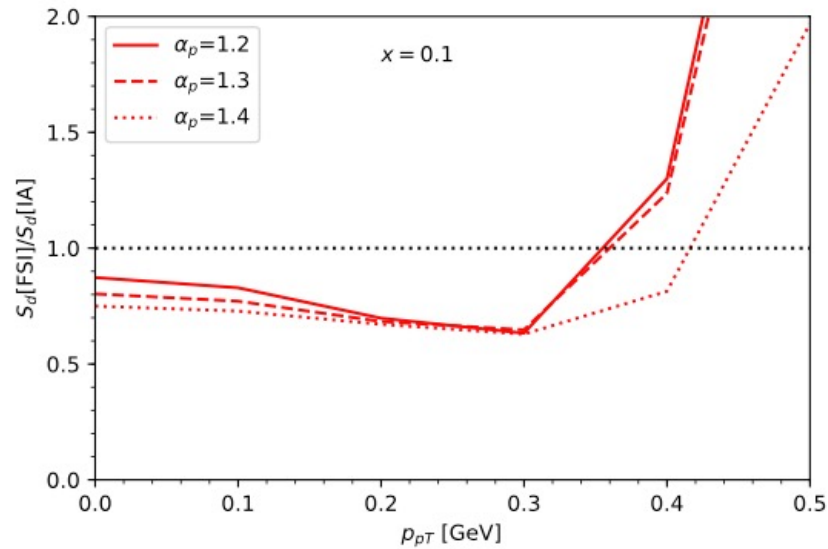
- $\nu \gg$ hadronic scale: large phase space for hadron production.
- “Fast” hadrons $E_h = \mathcal{O}(\nu) \rightarrow$ current fragmentation region: Formed outside the nucleus, interaction with the spectator suppressed.
- “Slow” hadrons $E_h = \mathcal{O}(1 \text{ GeV}) \rightarrow$ target fragmentation region: Formed inside the nucleus, interact with hadronic cross sections.
 - Source of FSI in tagged DIS!

Implementation

- Distributions of slow hadrons in DIS on nucleon: kinematic dependence, empirical distributions
- Hadron-nucleon scattering amplitudes: Re/Im
- Calculation of rescattering process: phase space integral
- Study kinematic dependences: x, α_p, p_{pT}

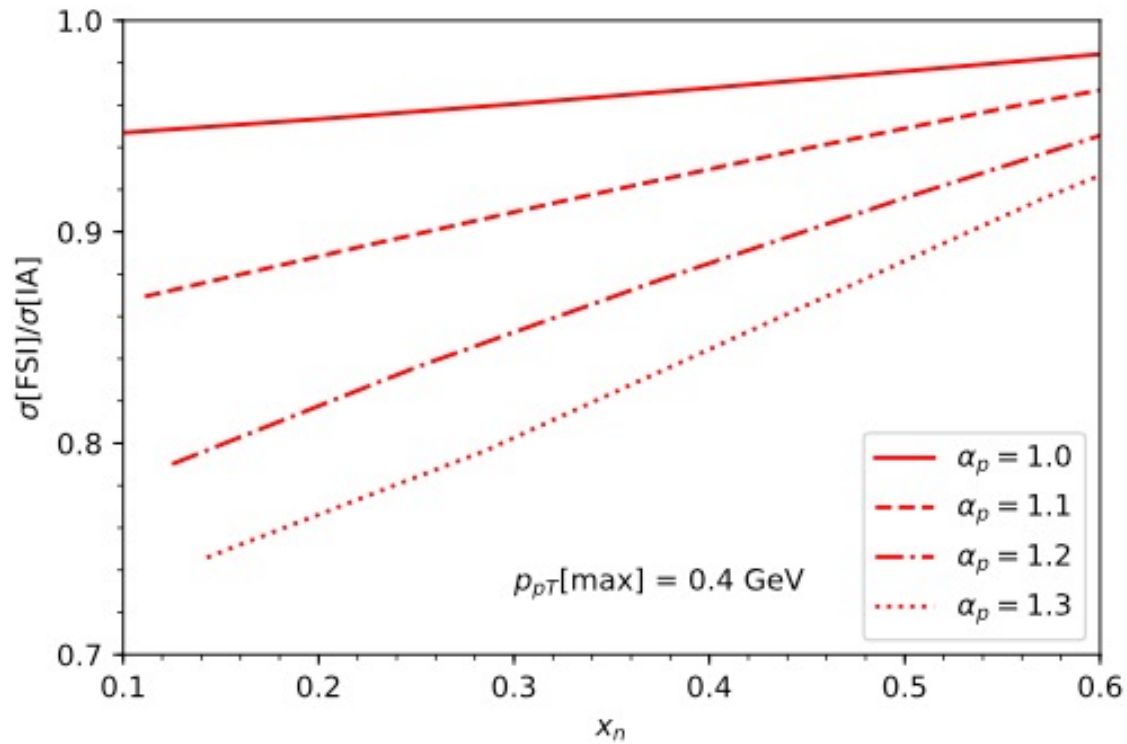
Momentum distribution of slow hadrons in nucleon rest frame: Cone in virtual photon direction.

FSI: Kinematic Dependence



- FSI Ratio $S_d[\text{FSI}]/S_d[\text{IA}]$
- p_{pT} dependence: weak up to ~ 0.3 GeV, strong rise above
- α_p dependence: FSI increases with $\alpha_p - 1$ at small p_{pT}
- x dependence: FSI decreases with increasing x due to depletion of slow hadrons

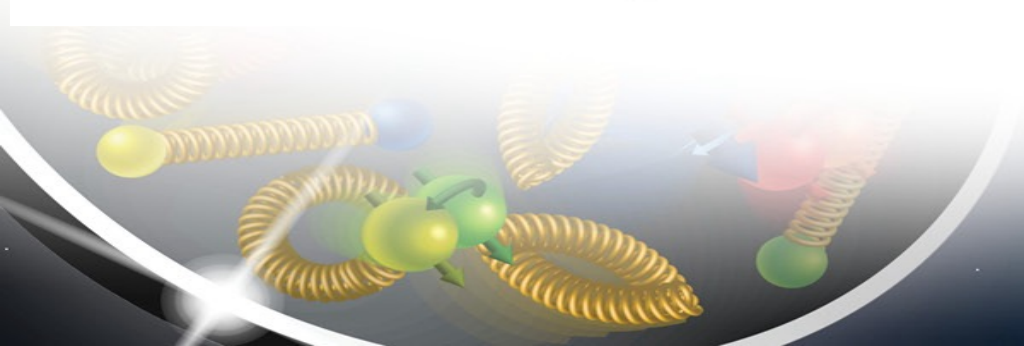
FSI: pT-integrated cross-section



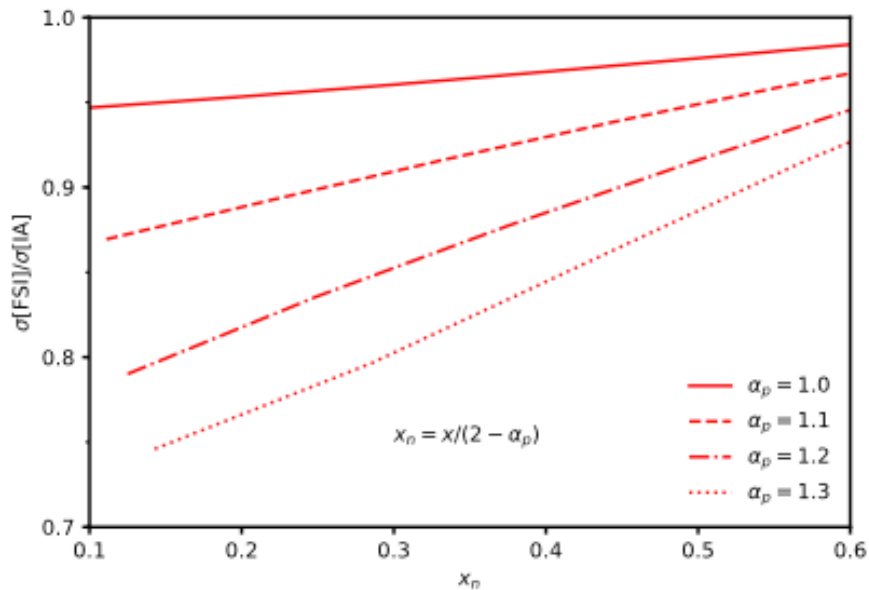
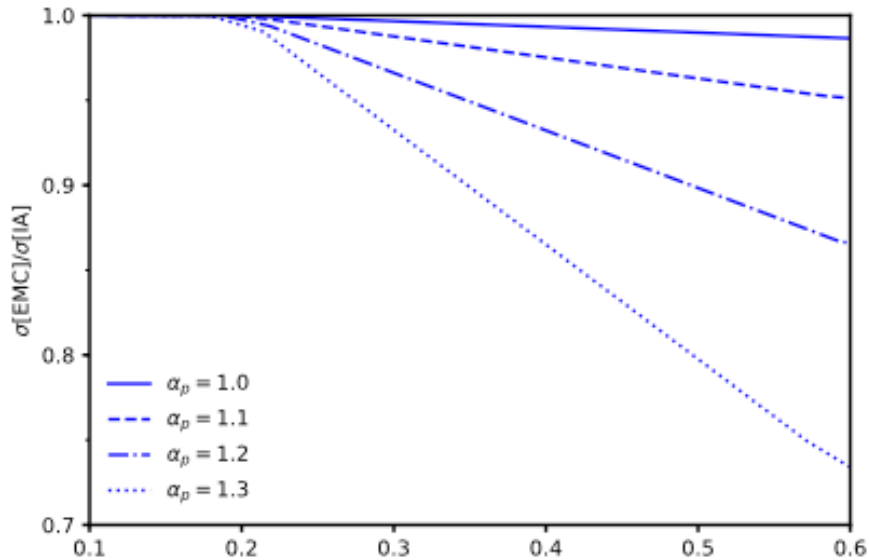
- p_{pT} - integrated cross section:

$$\sigma = \int_{p_{pT}[\max]} d^2 p_{pT} S_d(\alpha_p, p_{pT}) \sigma_n(x_n)$$

- Here: Plotted as a function of $x_n = x/(2 - \alpha_p)$
- Simple dependence of α_p and x_n .
- FSI effect typically 10-20%



FSI: Initial state vs. final-state modification



- Here: p_{pT} - integrated cross section, $p_{pT} [max] = 0.4$ GeV
- EMC Effect: virtuality-dependent model

$$\frac{\sigma_n[bound]}{\sigma_n[free]} = 1 + \frac{t}{\langle t \rangle} f_{EMC}(x_n)$$

$$t = t(\alpha_p, p_{pT})$$

- Compare EMC and FSI

→ Currently in-progress!

Measuring t-distribution

From K. Tu @ DIS 2023:

<https://indico.cern.ch/event/1199314/contributions/5189840/attachments/2621029/4531556/ePIC-exclusive-slides-Tu-v3.pdf>

Reconstruction method of $-t$

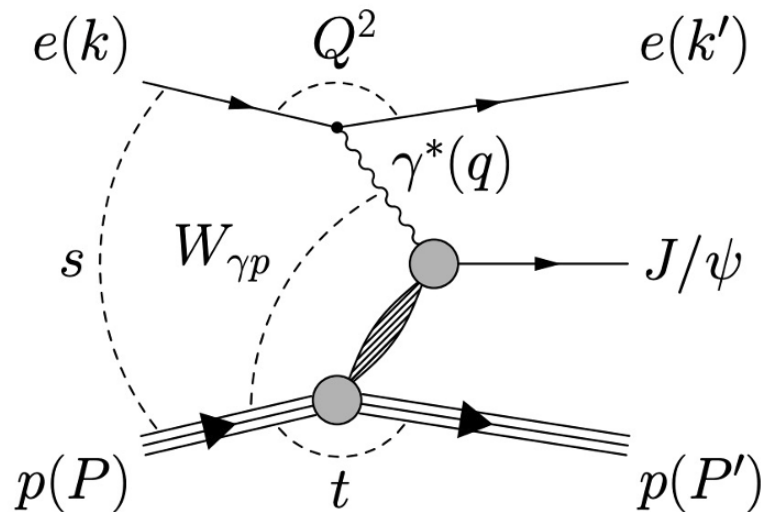
- Method Exact (E):
- Method Approximate (A) (UPCs)
- Method with **exclusivity corrected** (L):

$$-t = -(\mathbf{p}_e - \mathbf{p}_{e'}, -\mathbf{p}_{VM})^2 = -(\mathbf{p}_{A'} - \mathbf{p}_A)^2$$

$$-t = (p_{T,e} + p_{T,VM})^2$$

$$-t = -(\mathbf{p}_{A',\text{corr}} - \mathbf{p}_A)^2,$$

where $\mathbf{p}_{A',\text{corr}}$ is constrained by exclusive reaction.



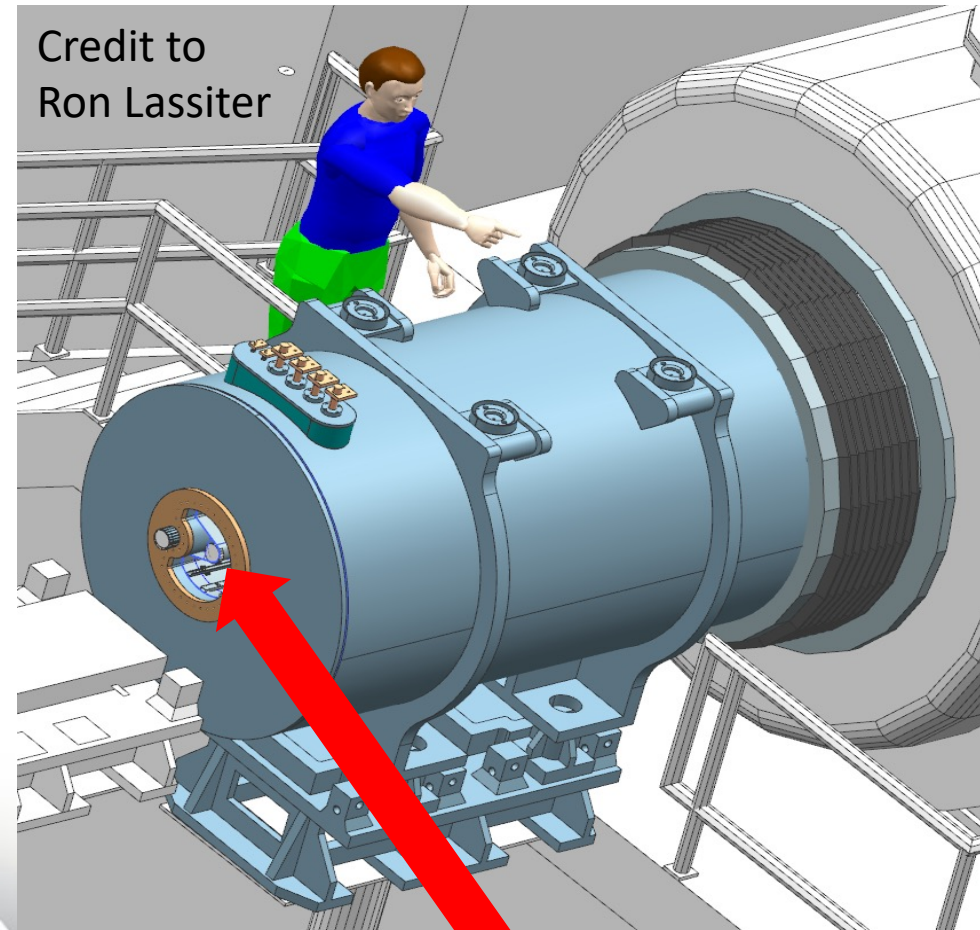
Best method concluded from the EIC Yellow Report* is with **exclusivity corrected**:

- Insensitive to beam effects, e.g., angular divergence and momentum spread.
- More precise than Method A for electroproduction

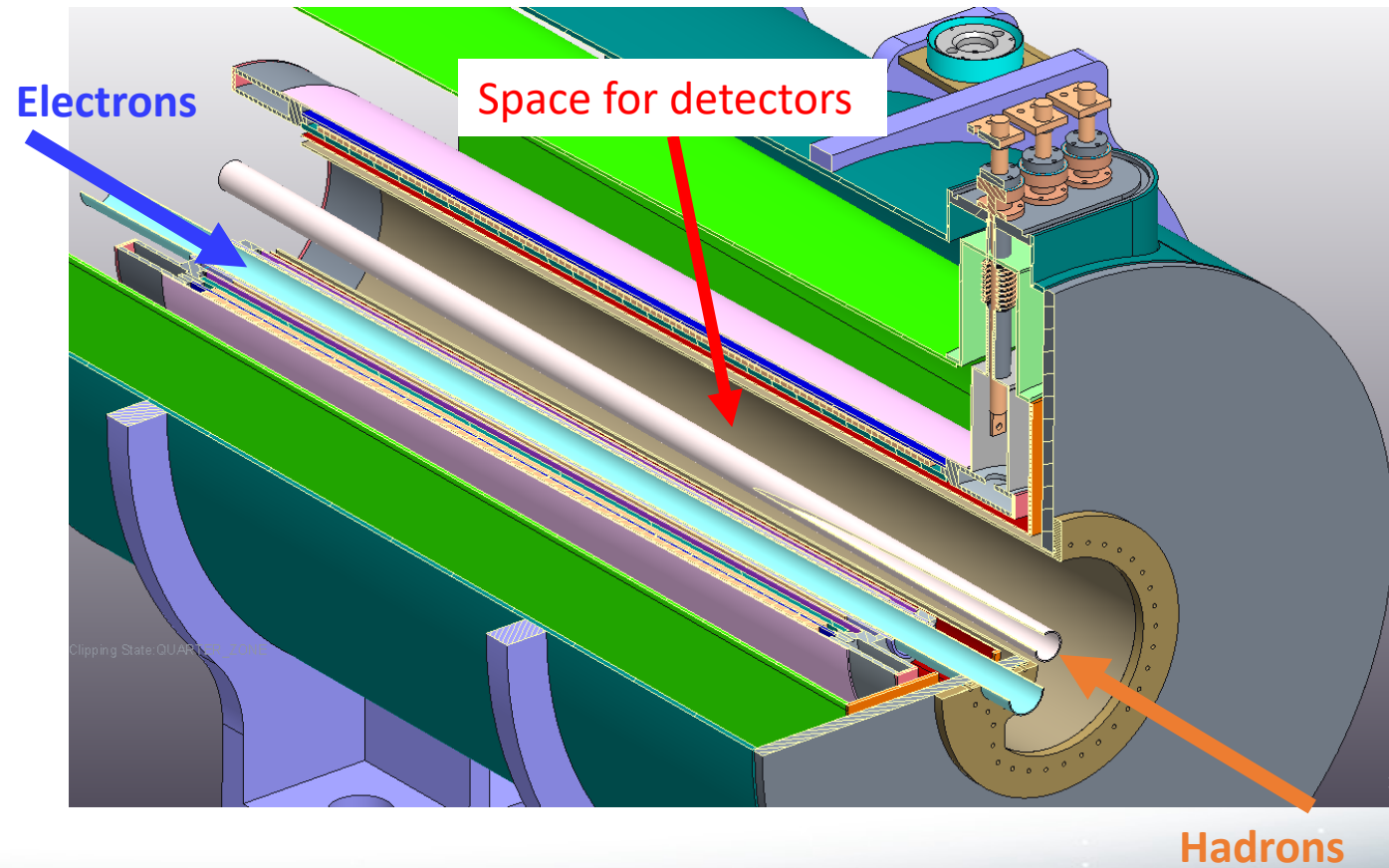
* also known as **Method L** in the Yellow Report

B0 Detectors

- Detector subsystem embedded in an accelerator magnet.



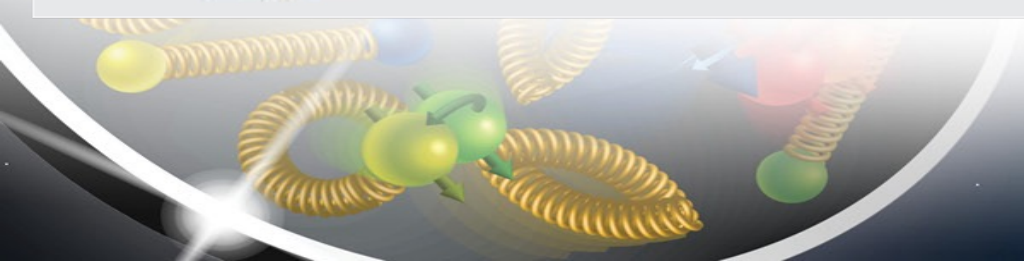
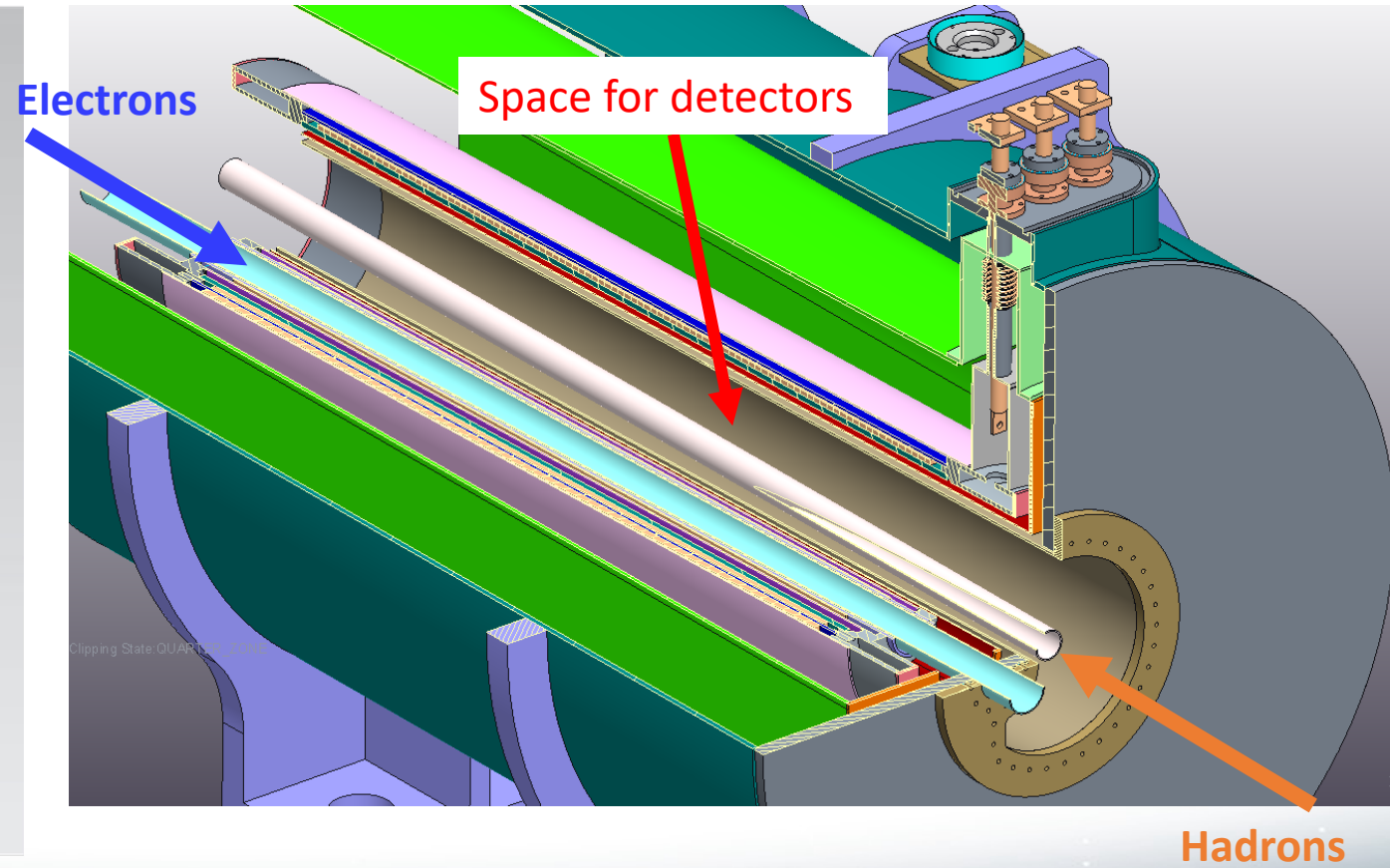
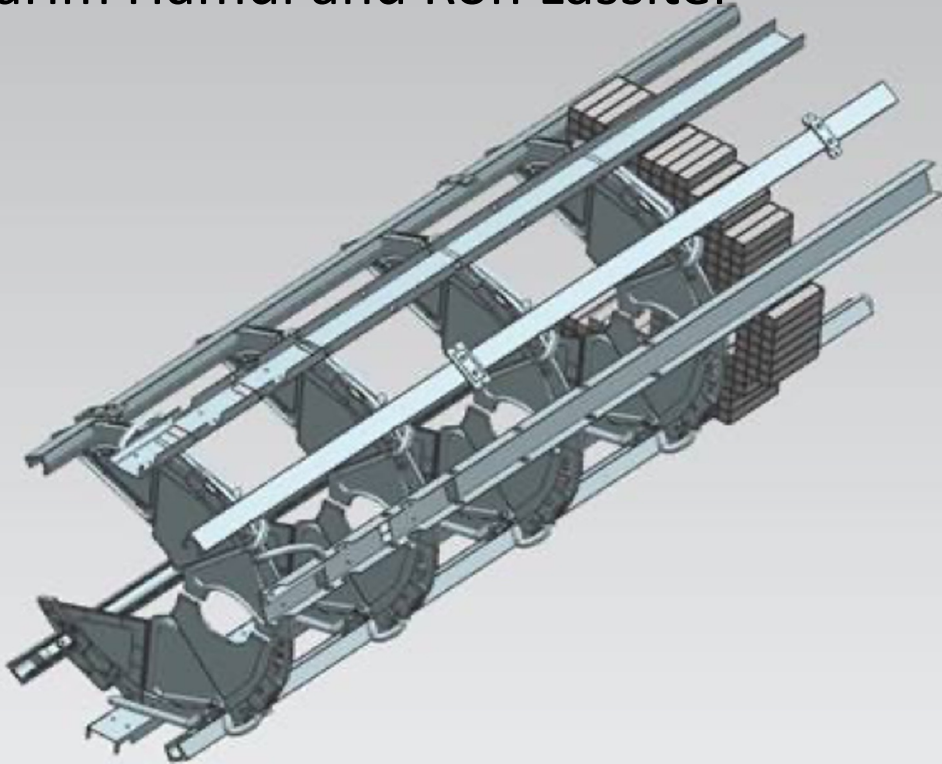
This is the opening where the detector planes will be inserted



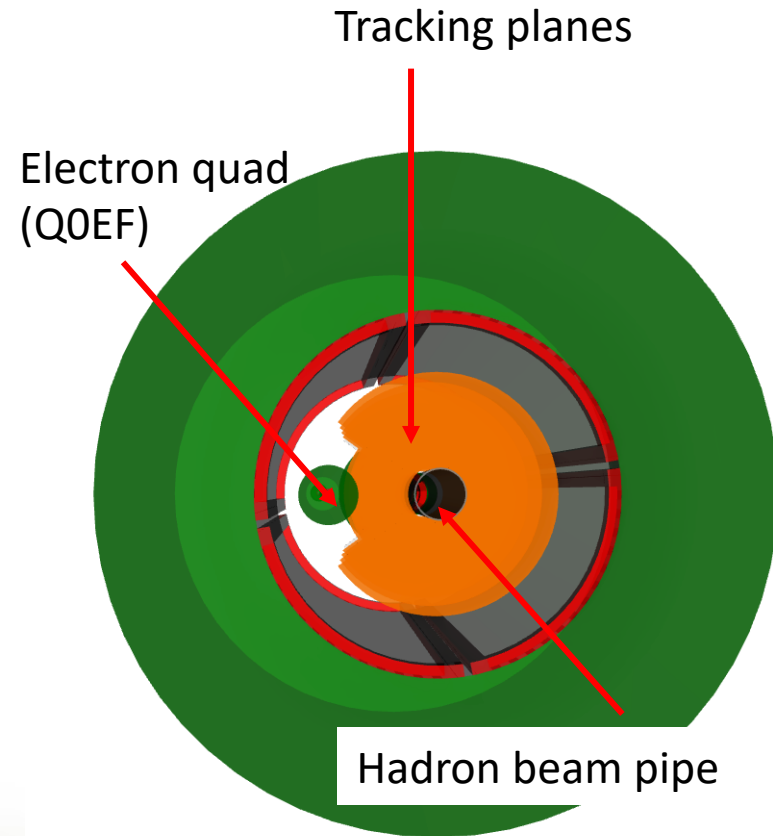
B0 Detectors

- Detector subsystem embedded in an accelerator magnet.

Karim Hamdi and Ron Lassiter



B0 Tracking and EMCAL Detectors



ePIC DD4HEP Simulation



PbWO₄/LYSO
EMCAL (behind
tracker)

- Technology choices:
 - Tracking: 4 layers AC-LGADs
 - PbWO₄ or LYSO EMCAL.

➤ Status

- ✓ Used to reconstruct charged particles and photons.
 - ✓ Acceptance: $5.5 < \theta < 20.0$ mrad on one side, up to 13mrad on the other.
 - ✓ Focus now is on readout, new tracking software, and engineering support structure.
- ✓ Stand-alone simulations have demonstrated tracking resolution.
 - <https://indico.bnl.gov/event/17905/>
 - <https://indico.bnl.gov/event/17622/>

Bee Detectors

Design for two detectors is converging:

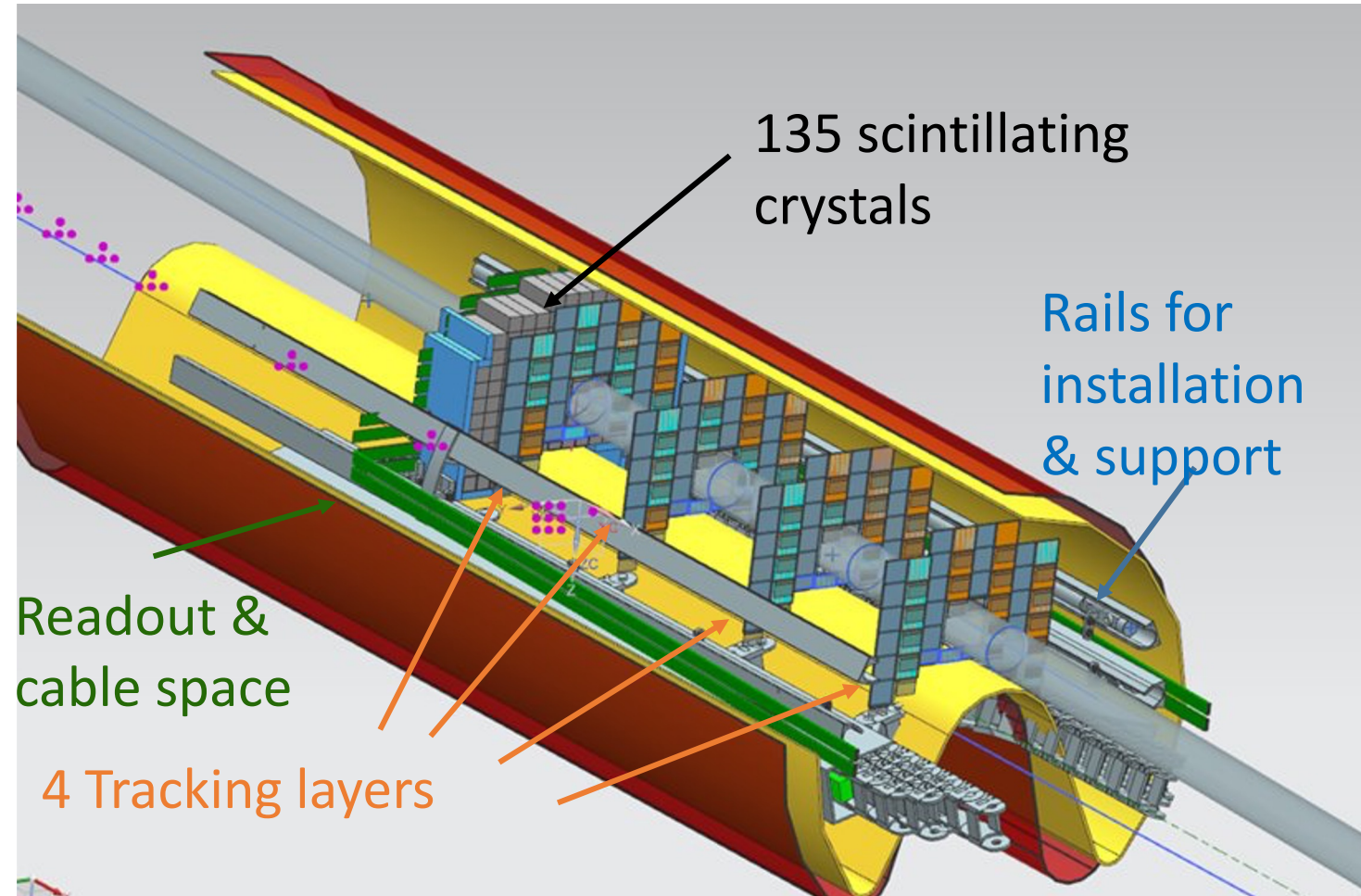
Si Tracker:

- 4 Layers of AC-LGAD → provide ~20um spatial resolution (with charge sharing) and 20-40ps timing resolution.
- Technology overlap w/ Roman pots

EM Calorimeter:

- 135 2x2x7* cm³ LYSO crystals
- Good timing and position resolution
- Technology overlap with ZDC

CAD Look credit: Jonathan Smith



* ZDC wants slightly longer crystals, ideally, we will use the same length in both detectors

Detectors - Simulation Studies

Si Tracker:

- Resolution plots made by Alex Jentsch with standalone setup (more [here](#) and [here](#))
- ACTS Tracking (a long-standing problem) was recently solved and is implemented in the simulation (see recent Sakib R [slides](#)), we expect more results soon

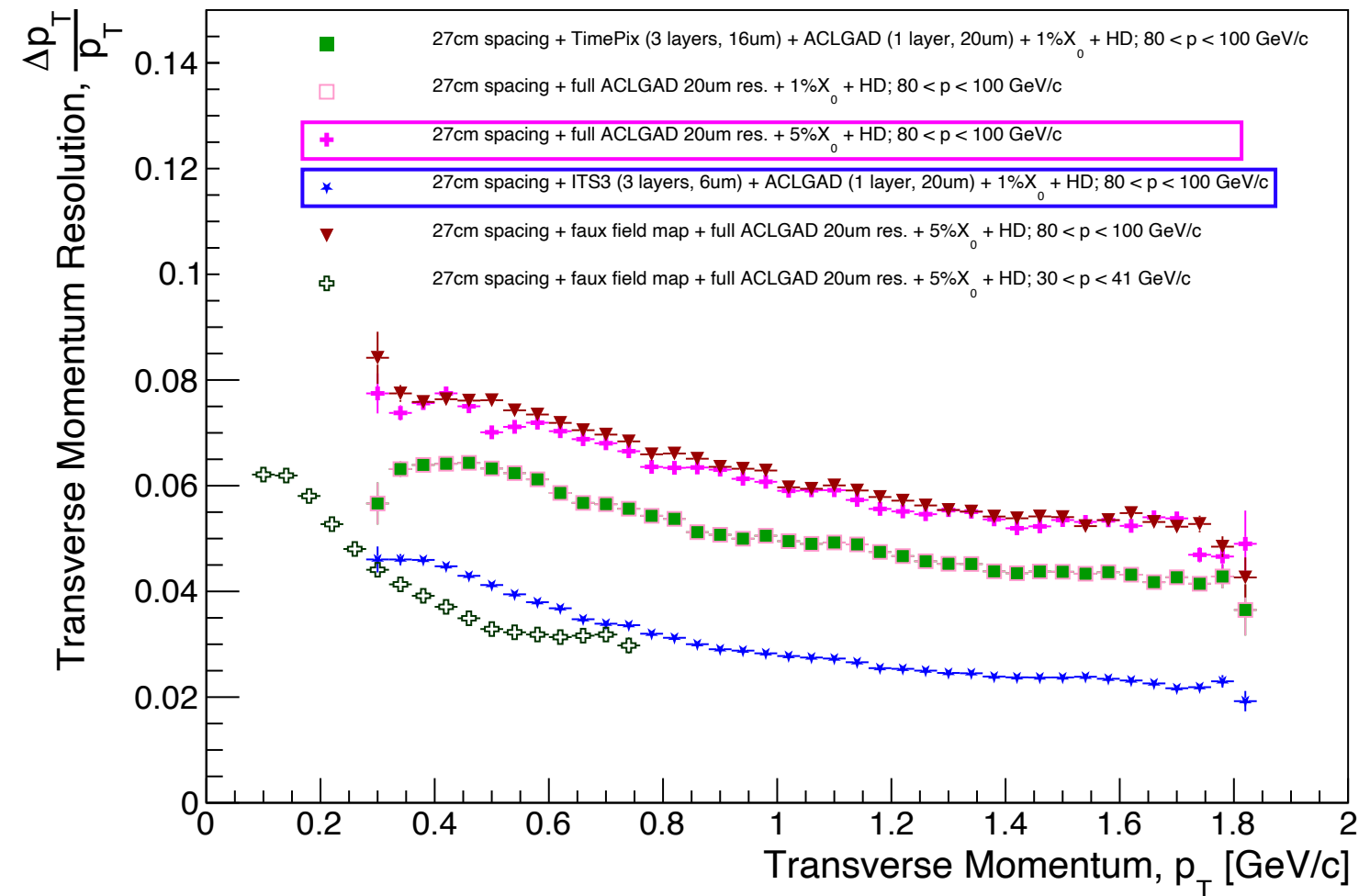
EM Calorimeter:

- Caveat - studies performed with PbWO4 crystals, LYSO crystals still to be implemented in the simulation.
- General performance studies by Michael Pitt (more in [FF weekly meeting](#))
- Sensitivity to soft photons (see Eden Mautner [talk](#) at the EICUG EC workshop early this week)



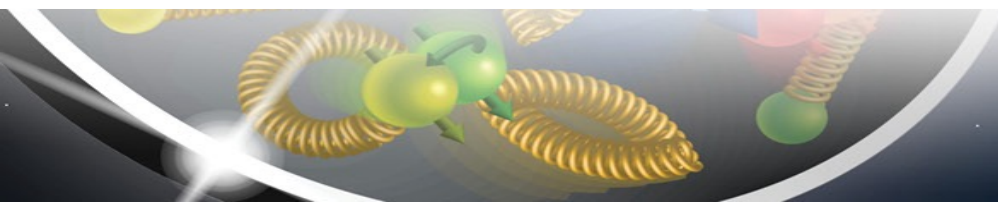


Tracking - Performance



- 27cm spacing with fully AC-LGAD system and 5% radiation length may be the most-realistic option.
- Reduced spacing (from 30cm) to make room for EMCAL.
- Needs to be looked at with proper field map and layout.
- Resolution impact on physics still being evaluated.

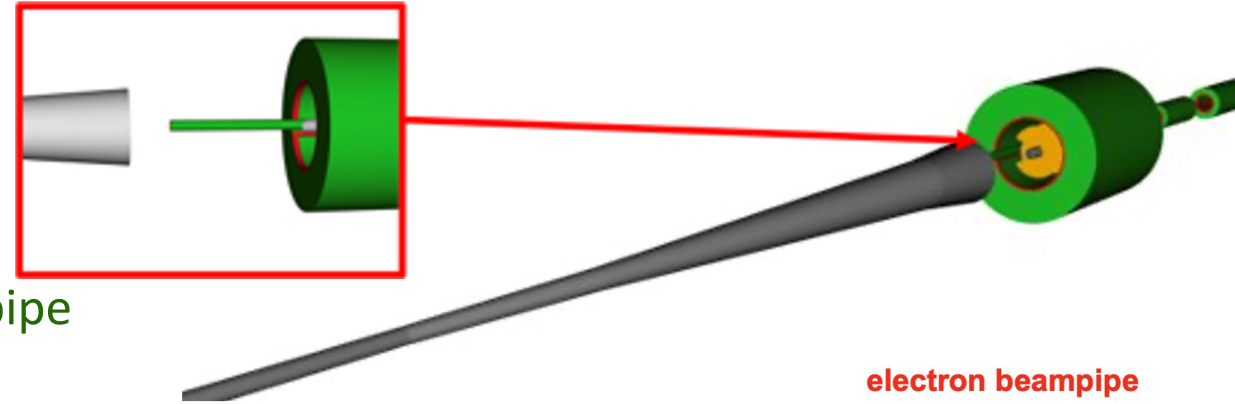
Note: momentum resolution (dp/p) is $\sim 2-4\%$, depending on configuration.



BEEEMCal - Performance

- Acceptance $5.5 < \theta < 23$ mrad
- Very low material budget in $5 < \eta < 5.5$

Particles within $5.5 < \theta < 15$ mrad don't cross the beampipe

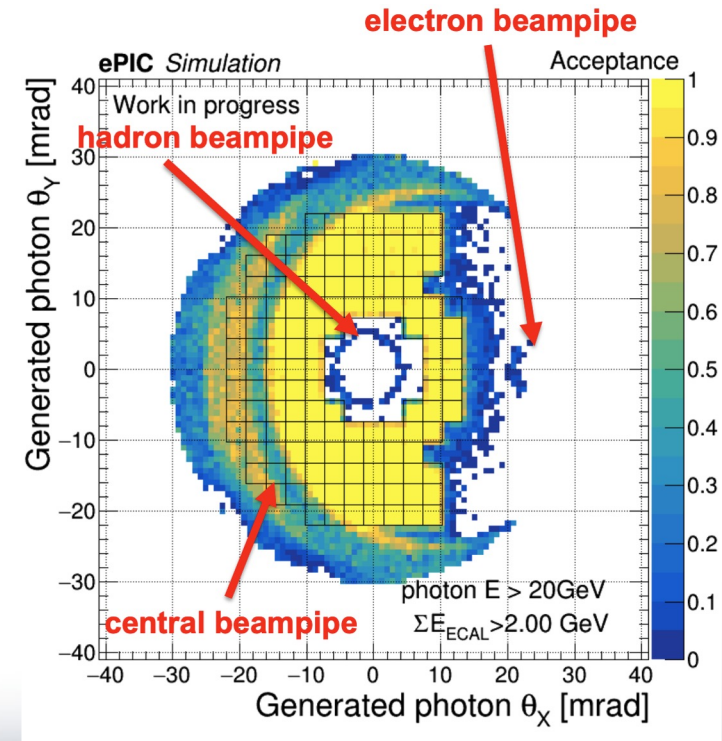
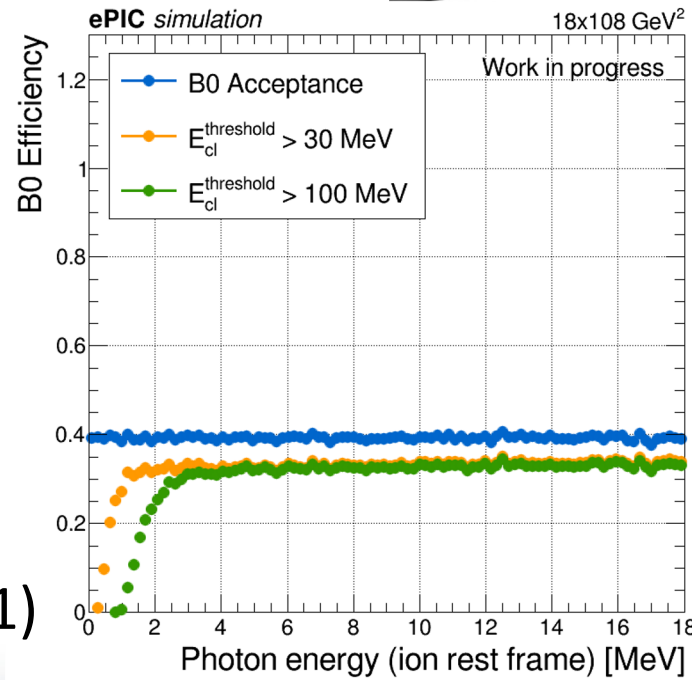


Photons:

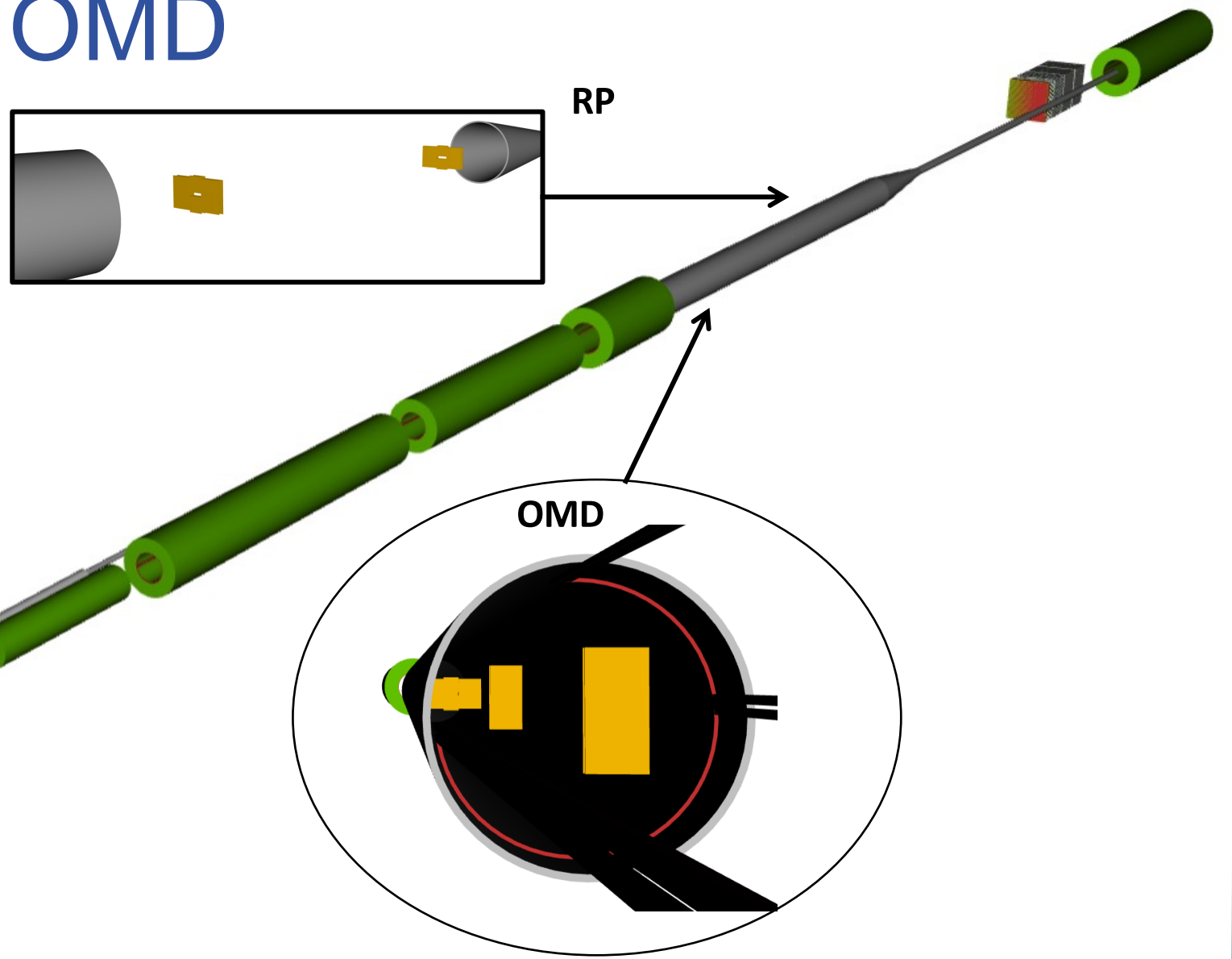
- High acceptance in a broad energy range (> 100 s MeV), including \sim MeV de-excitation photons
- Energy resolution of 6-7%
- Position resolution of ~ 3 mm

Neutrons:

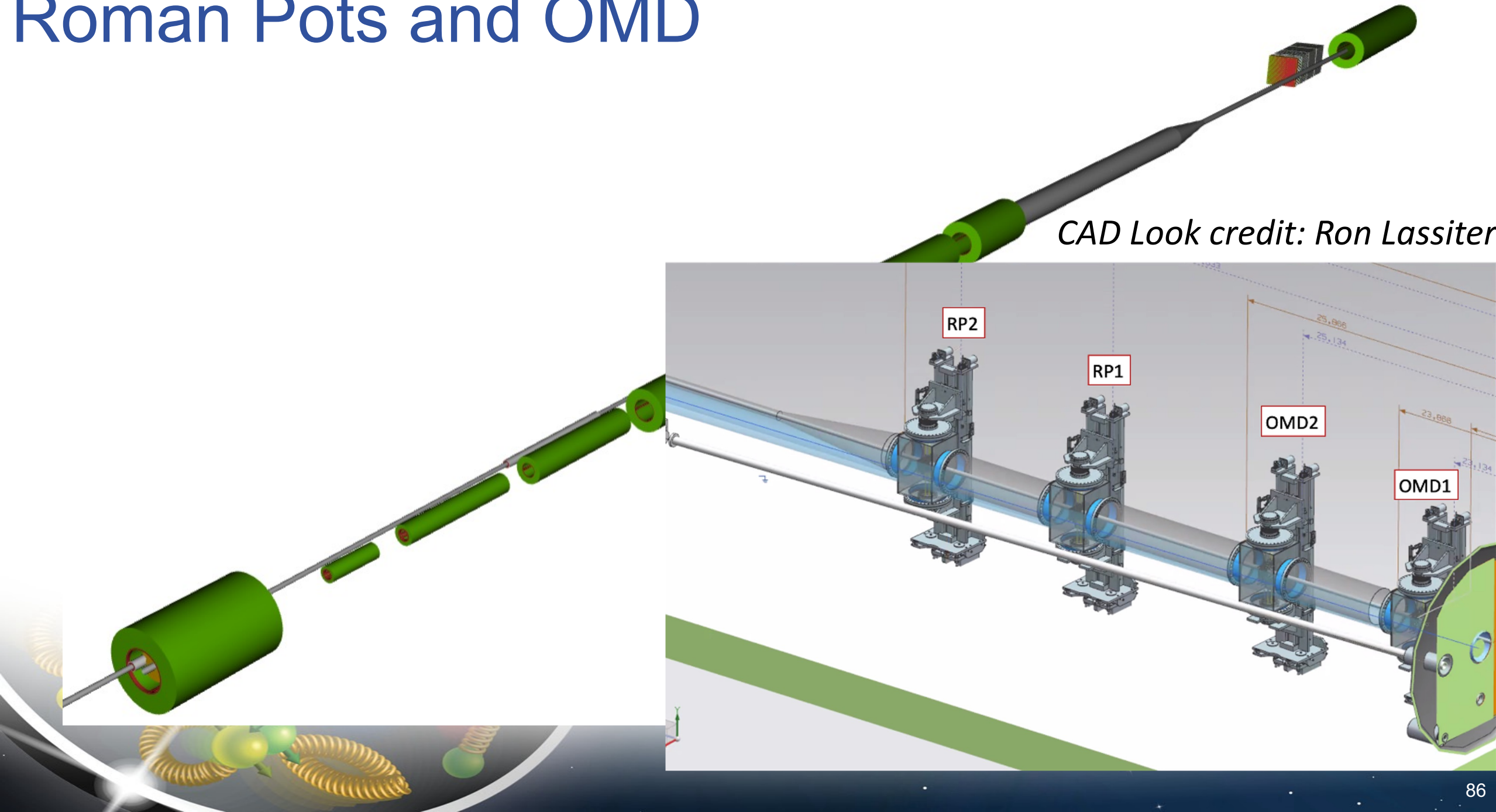
- 50% detection efficiency (λ is almost 1)



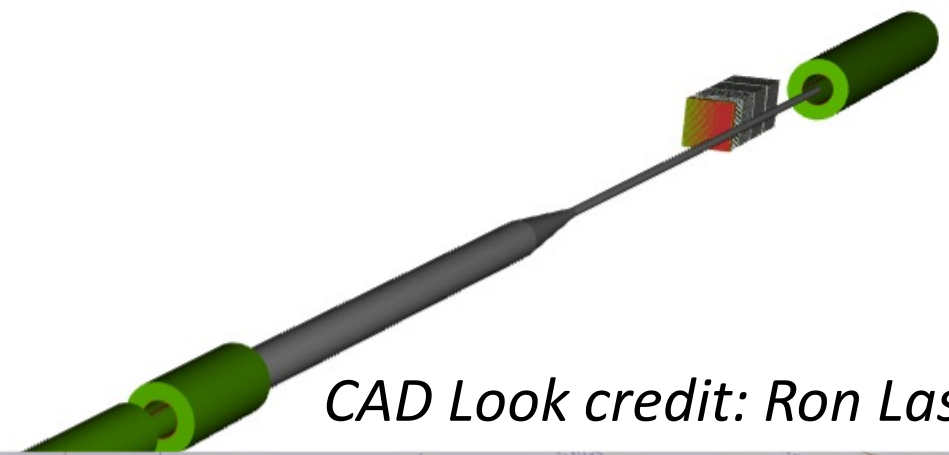
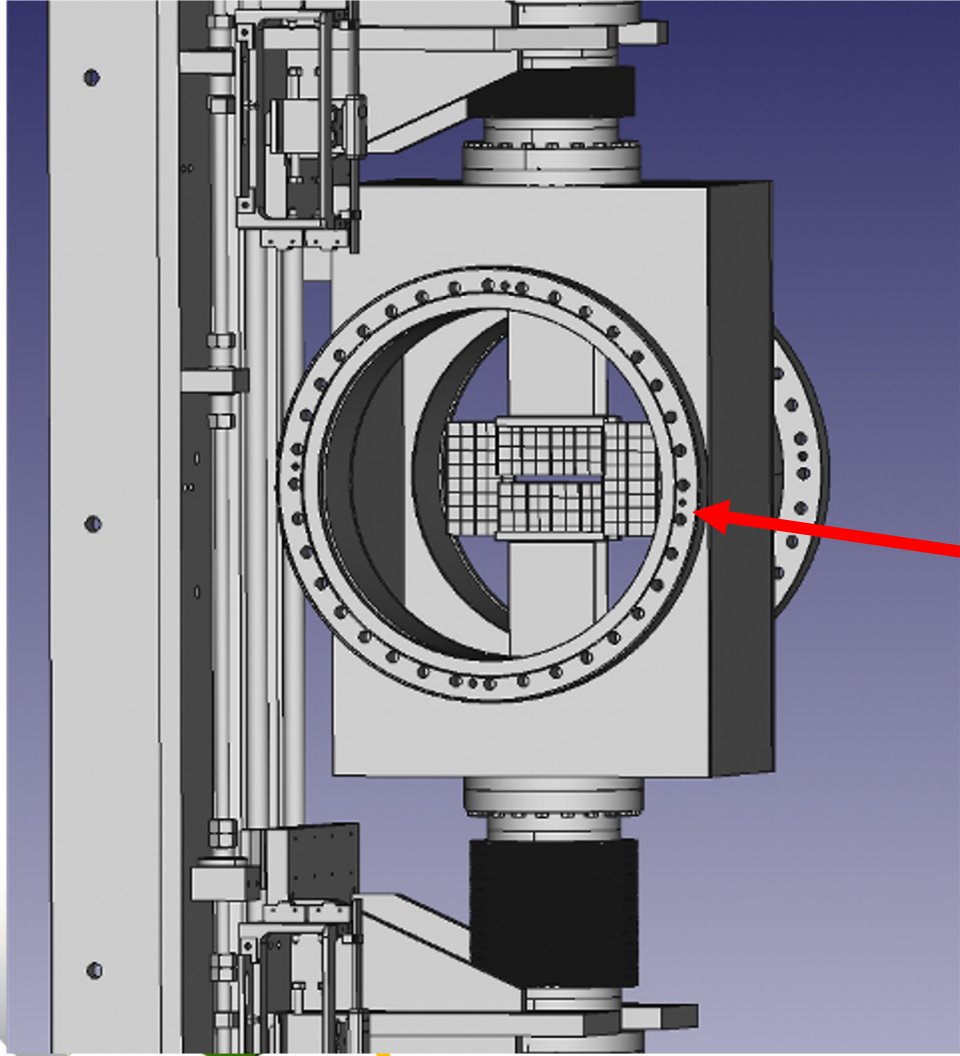
Roman Pots and OMD



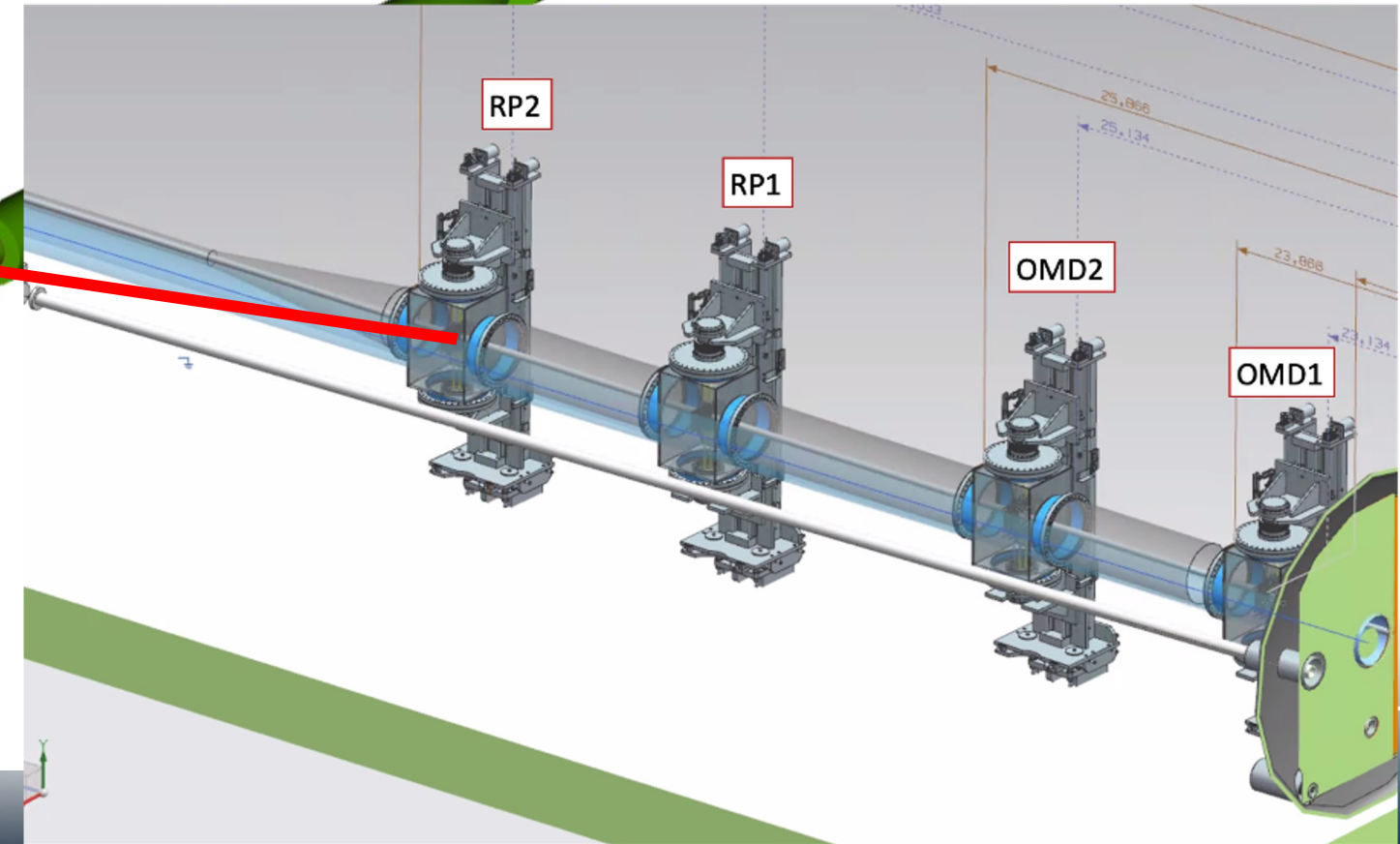
Roman Pots and OMD



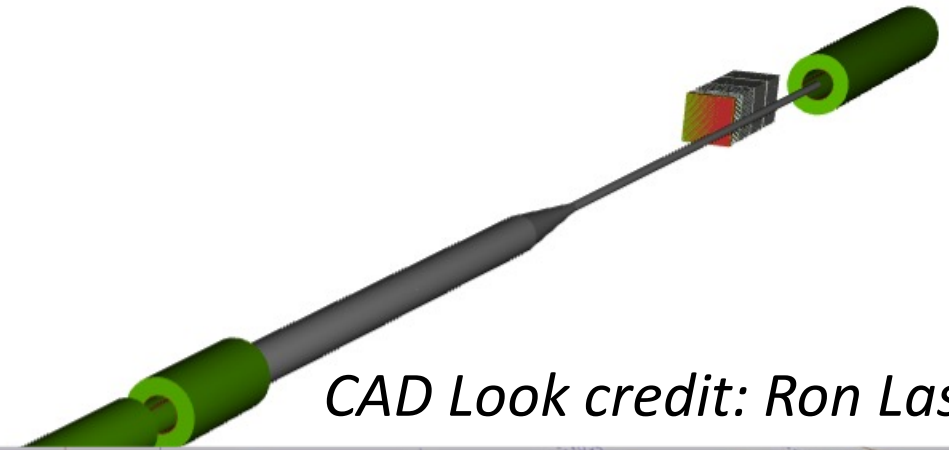
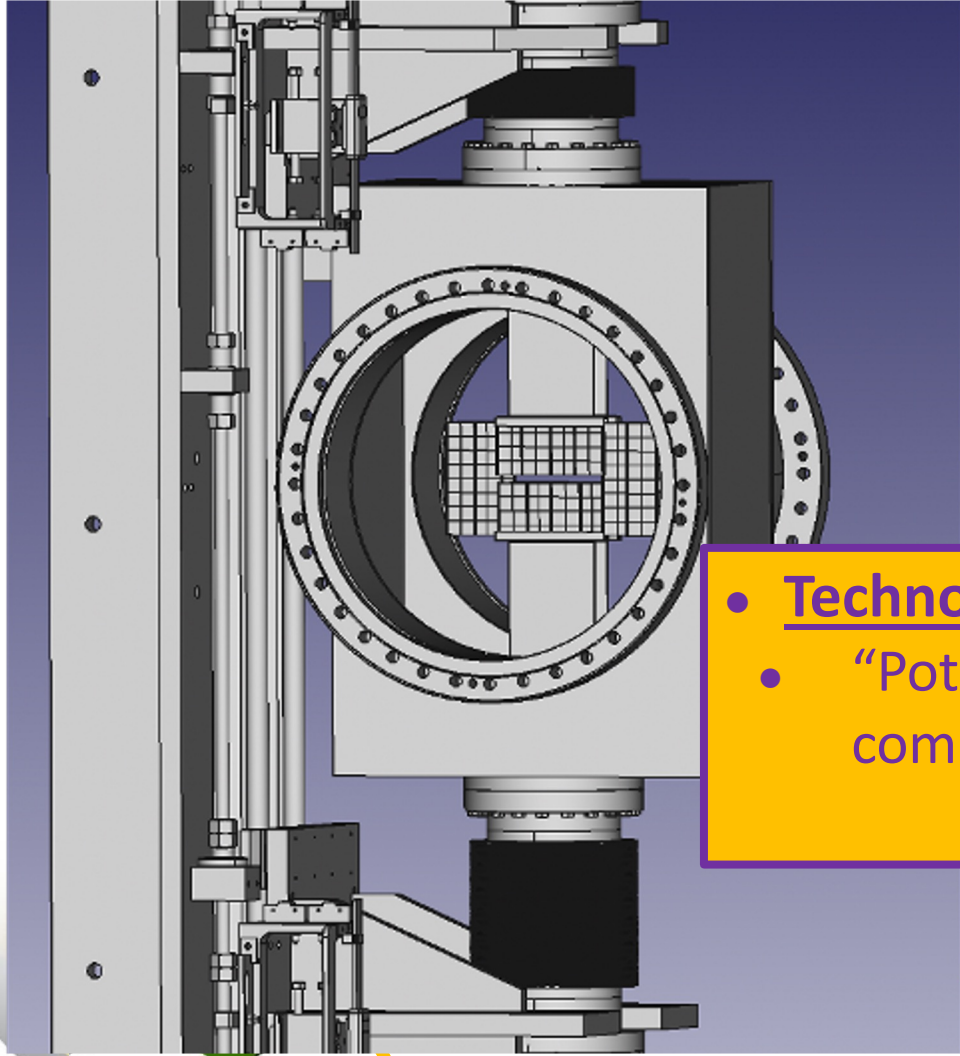
Roman Pots and OMD



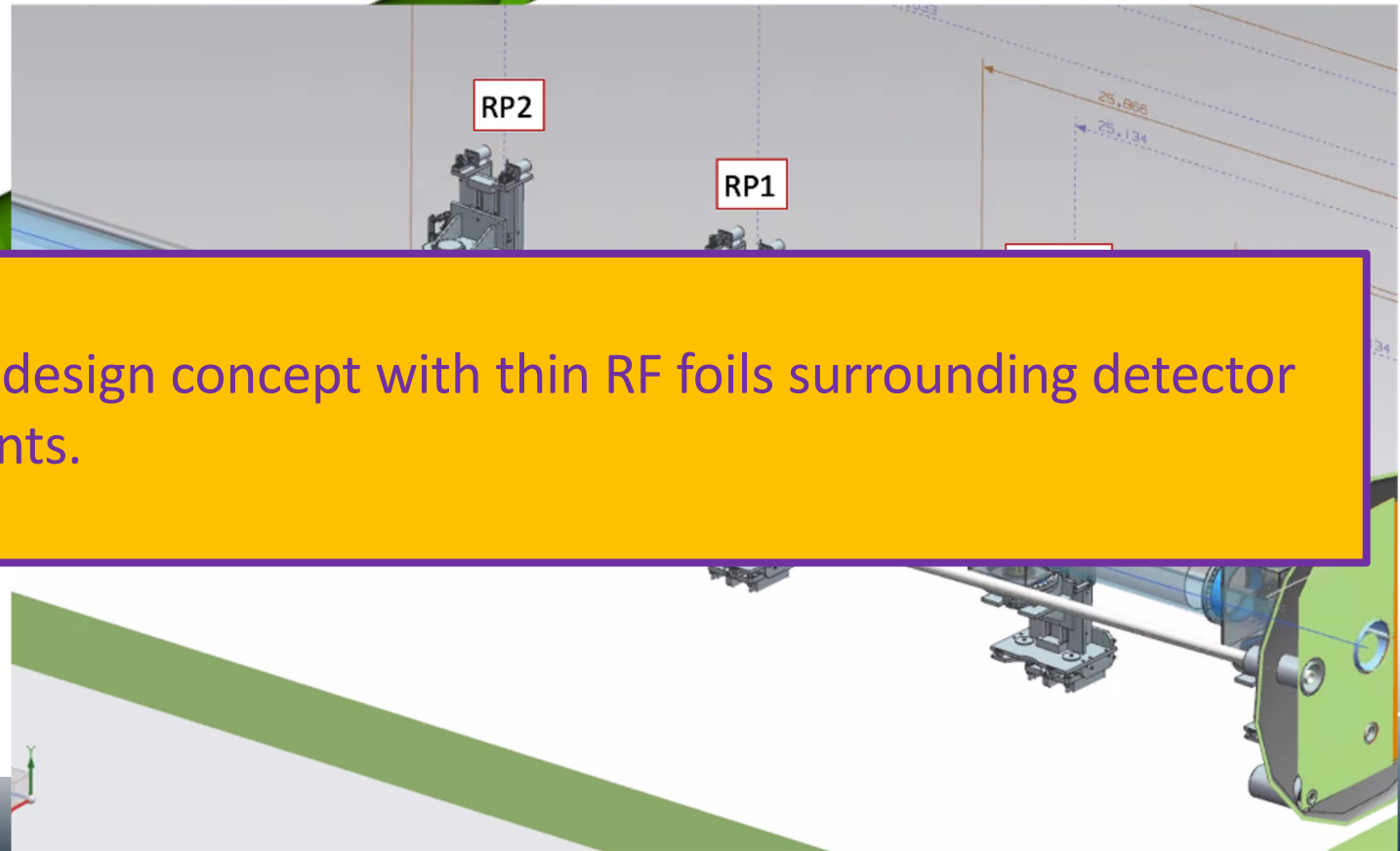
CAD Look credit: Ron Lassiter



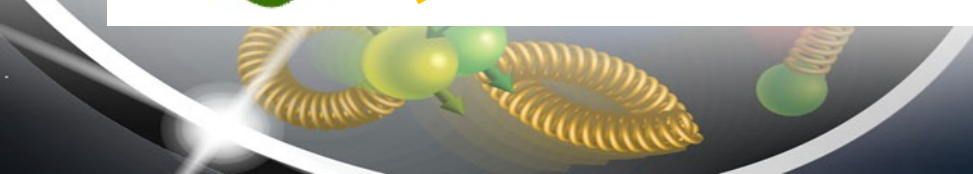
Roman Pots and OMD



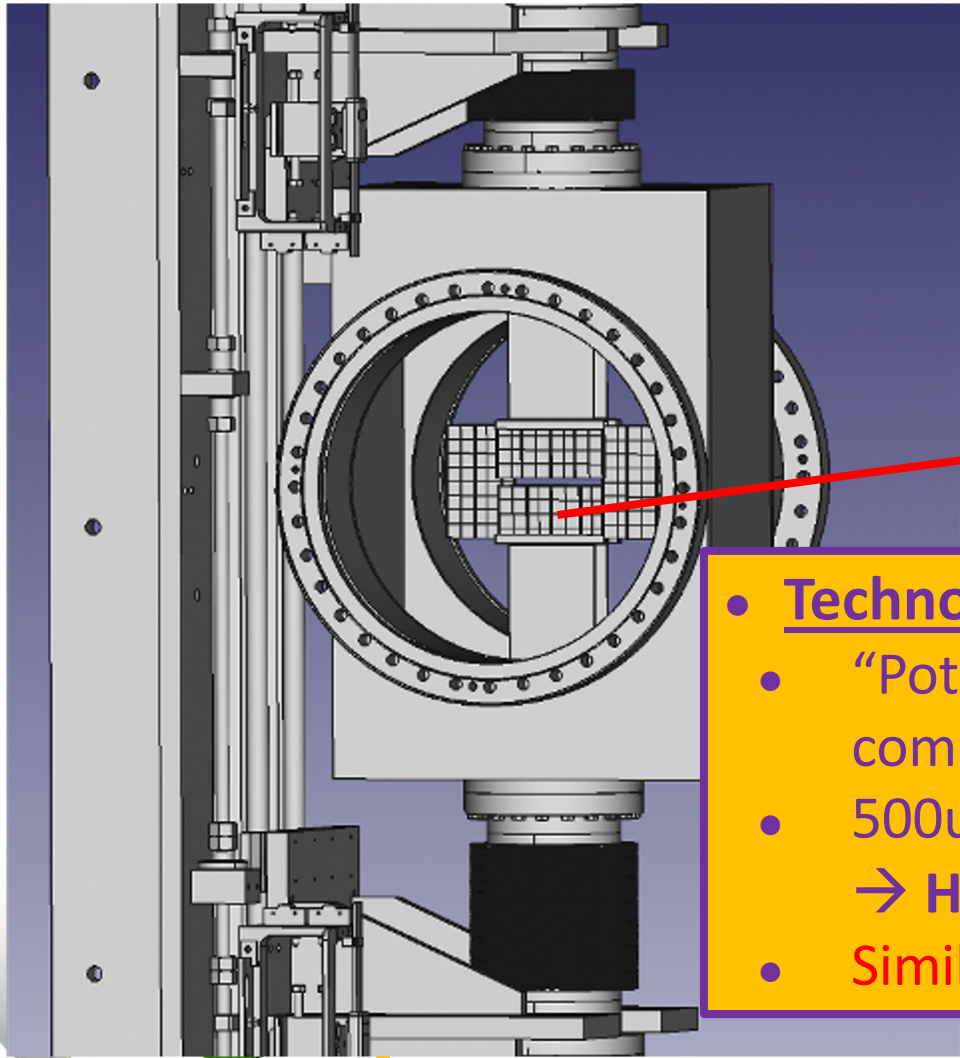
CAD Look credit: Ron Lassiter



- Technology
 - “Potless” design concept with thin RF foils surrounding detector components.



Roman Pots and OMD



12.8 cm

25.6 cm

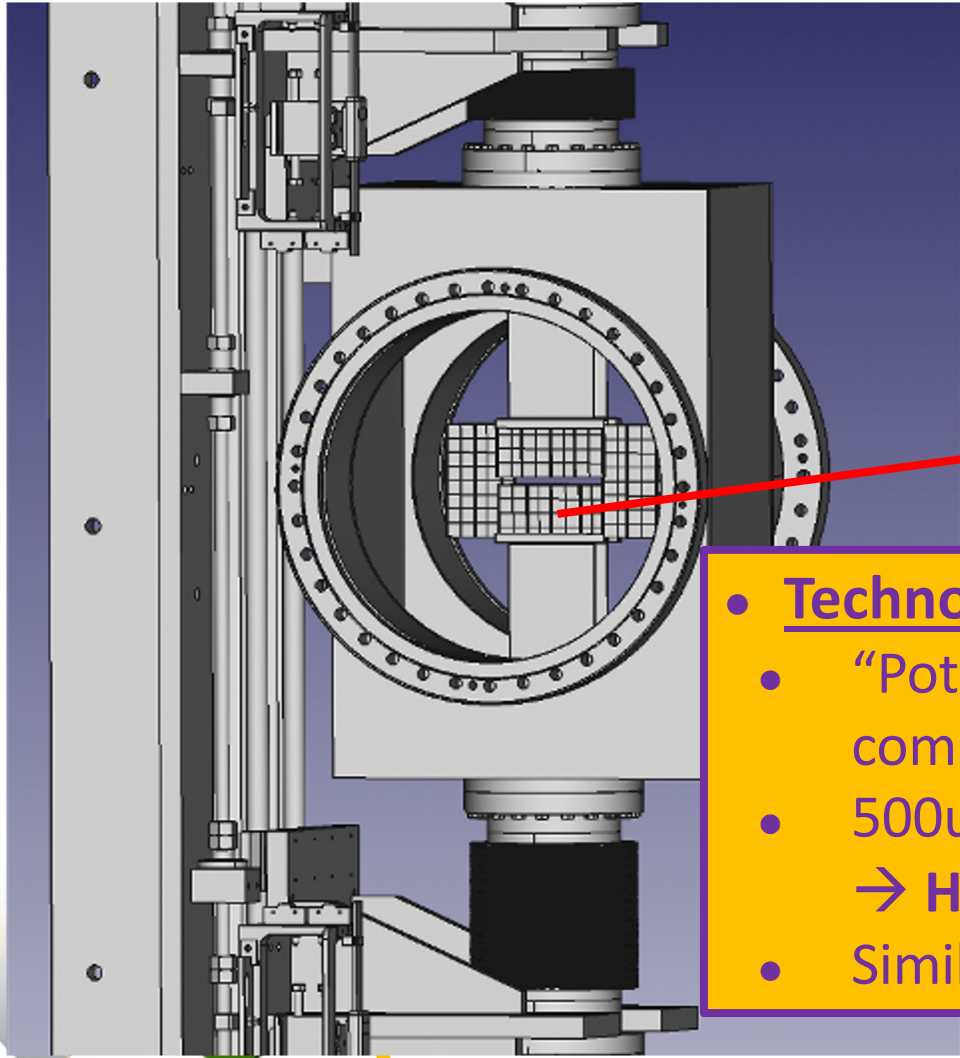


site

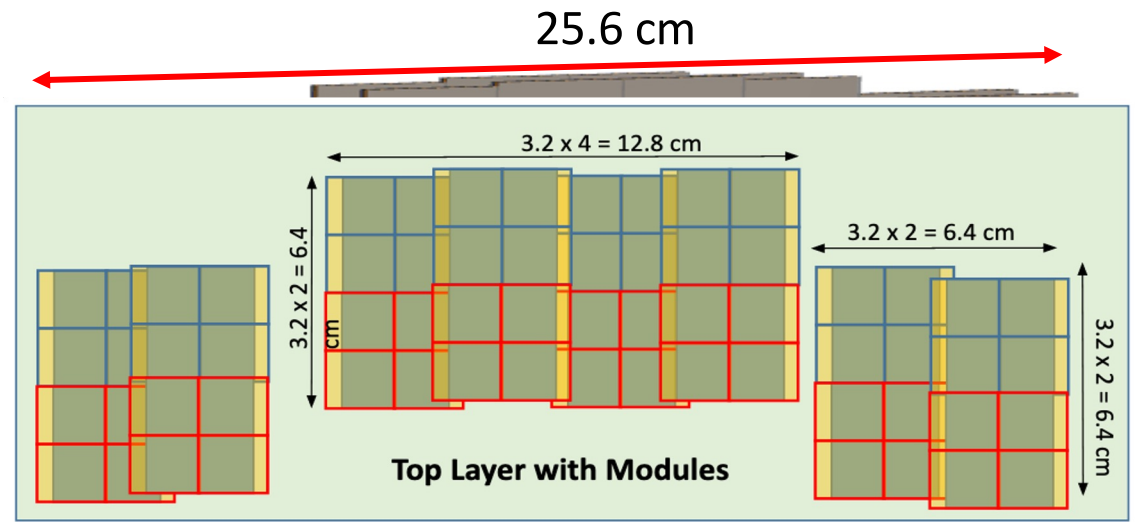
- Technology

- “Potless” design concept with thin RF foils surrounding detector components.
- 500um, **pixilated AC-LGAD sensor**, with 30-40ps timing resolution
→ **High-precision space and time information!**
- Similar concept for the OMD, just different active area and shape.

Roman Pots and OMD



12.8 cm



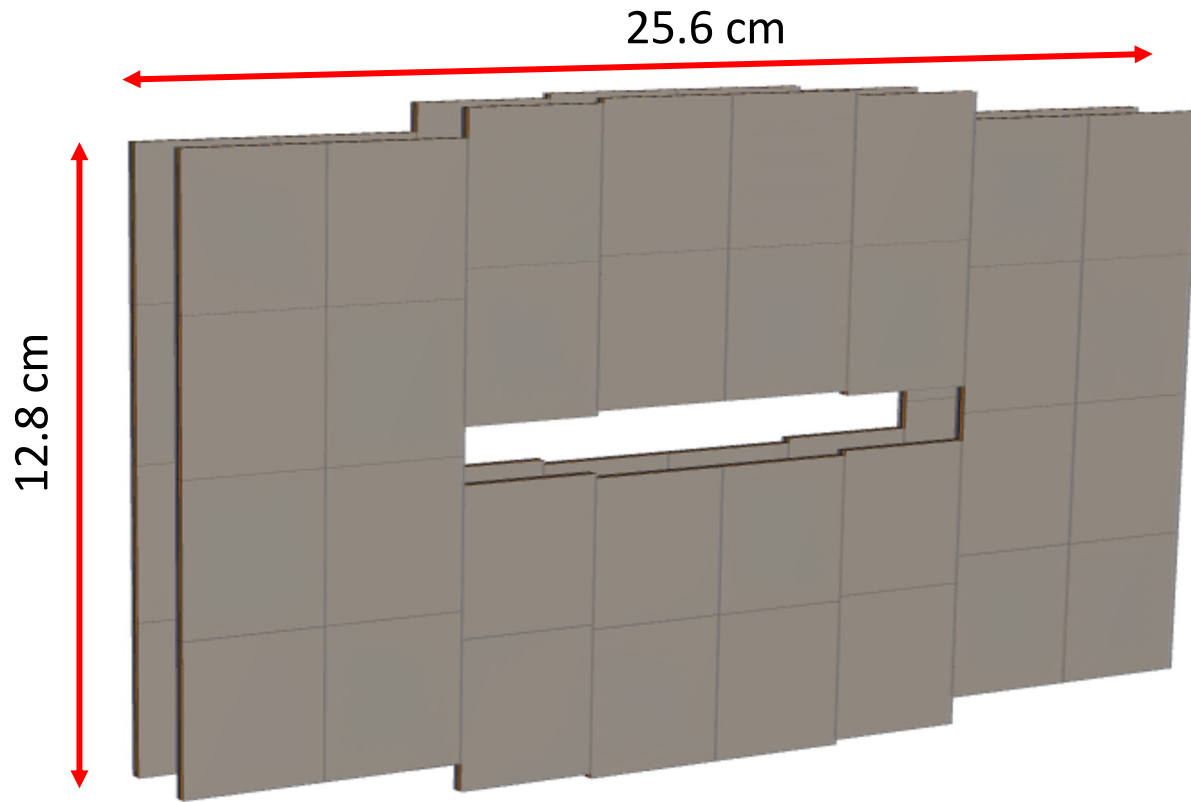
- Technology

- “Potless” design concept with thin RF foils surrounding detector components.
- 500um, **pixilated AC-LGAD sensor**, with 30-40ps timing resolution → **High-precision space and time information!**
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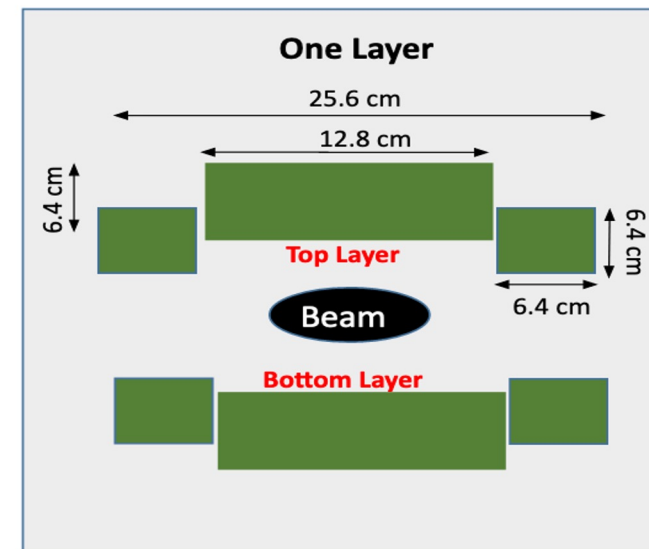
More engineering work is currently underway to optimize the layout, support structure, cooling, and movement systems for inserting the detectors into the beamline.

Roman "Pots" @ the EIC

$\sigma(z)$ is the Gaussian width of the beam, $\beta(z)$ is the RMS transverse beam size, ϵ is the beam emittance, and D is the momentum dispersion.



$$\sigma_{x,y} = \sqrt{\beta(z)_{x,y} \epsilon_{x,y} + \left(D_{x,y} \frac{\Delta p}{p}\right)^2}$$

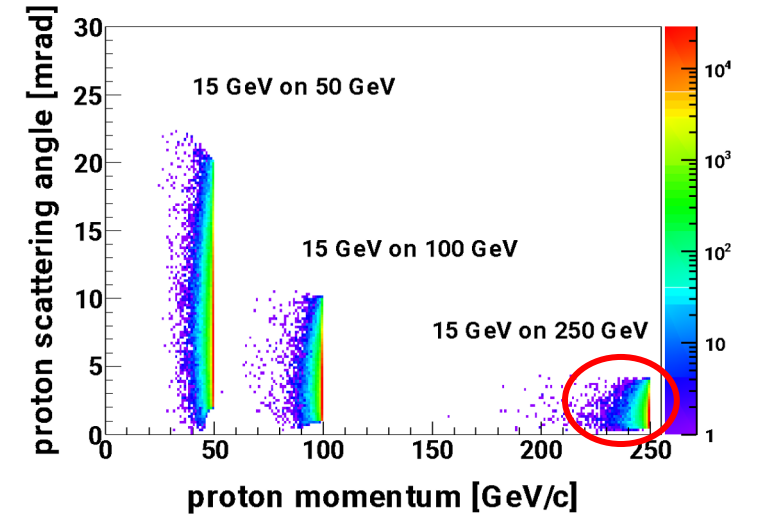
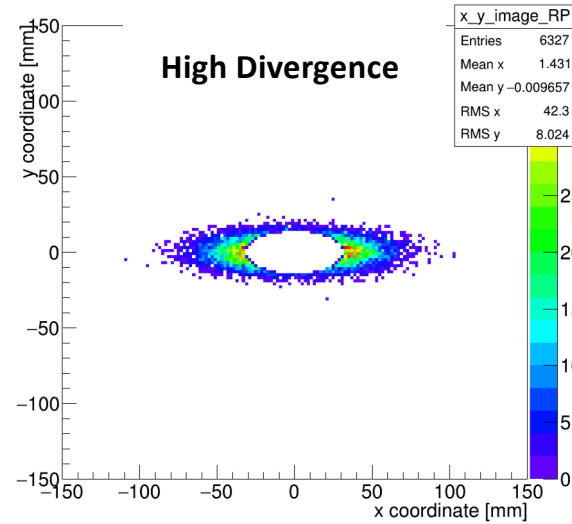
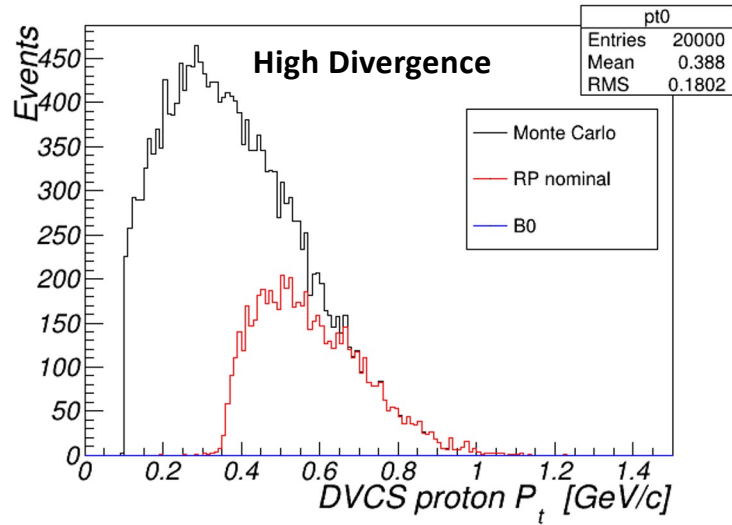


DD4HEP
Simulation

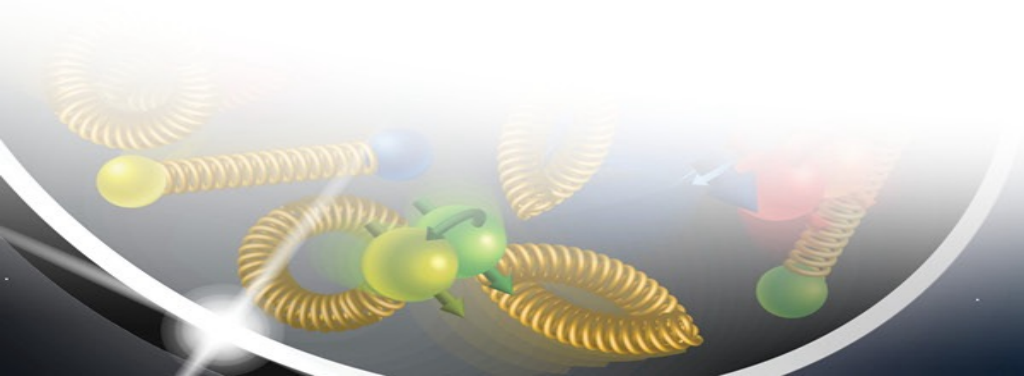
- Low-pT cutoff determined by beam optics.
 - The safe distance is $\sim 10\sigma$ from the beam center.
 - $1\sigma \sim 1\text{mm}$
- These optics choices change with energy, but can also be changed within a single energy to maximize *either acceptance at the RP, or the luminosity.*

Digression: Machine Optics (IP6)

275 GeV DVCS Proton Acceptance

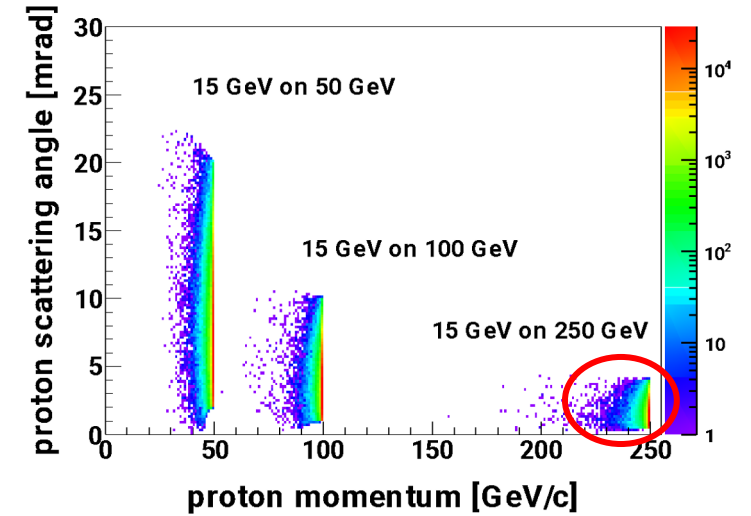
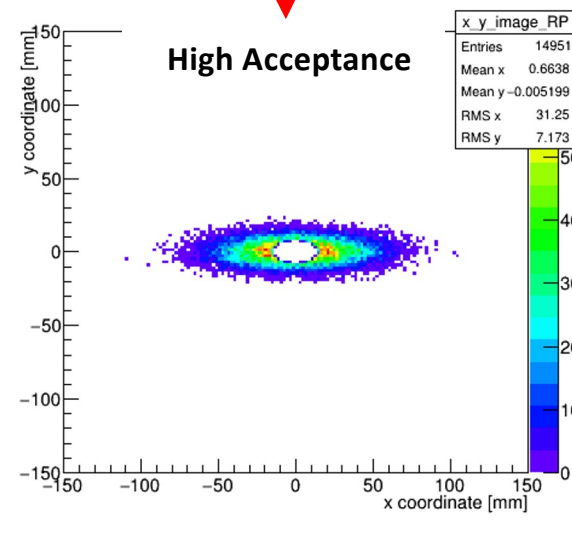
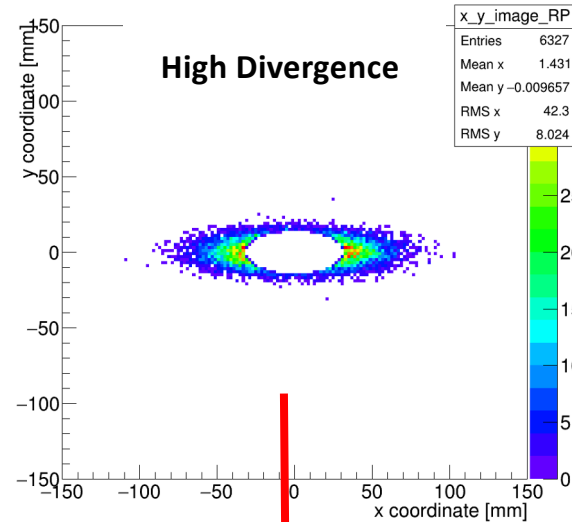
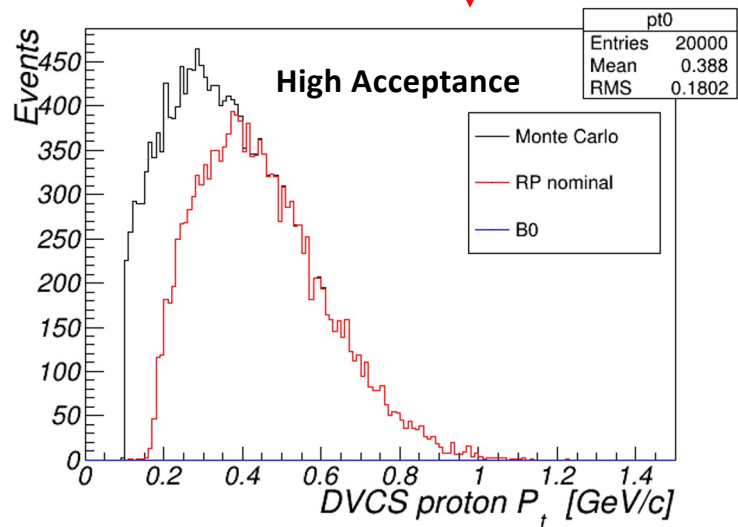
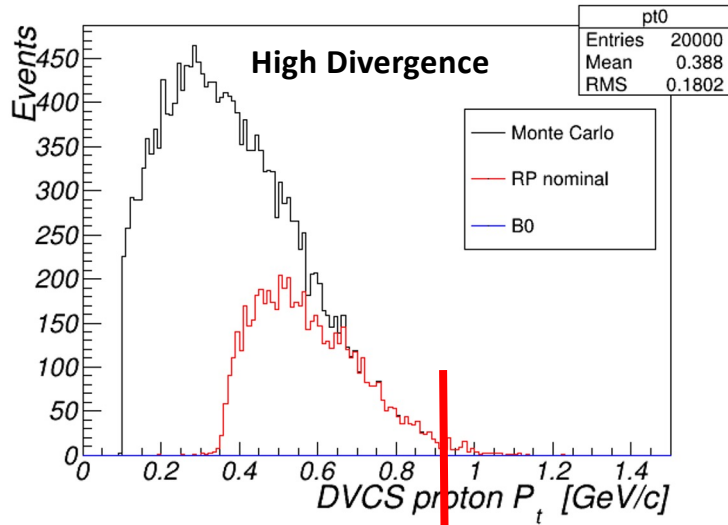


High Divergence: smaller β^* at IP, but bigger $\beta(z = 30m)$ -> higher lumi., larger beam at RP



Digression: Machine Optics (IP6)

275 GeV DVCS Proton Acceptance

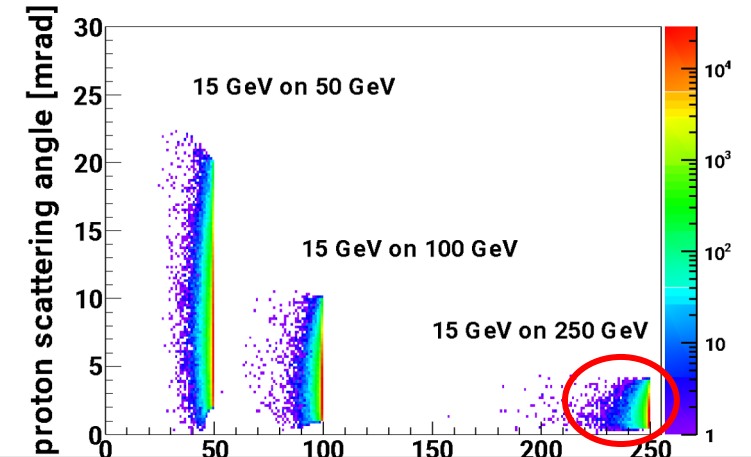
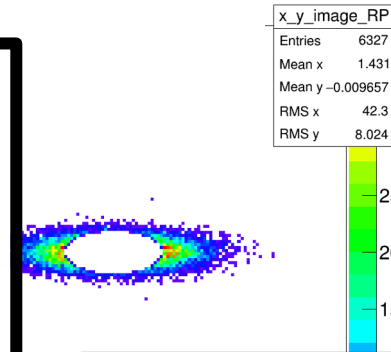
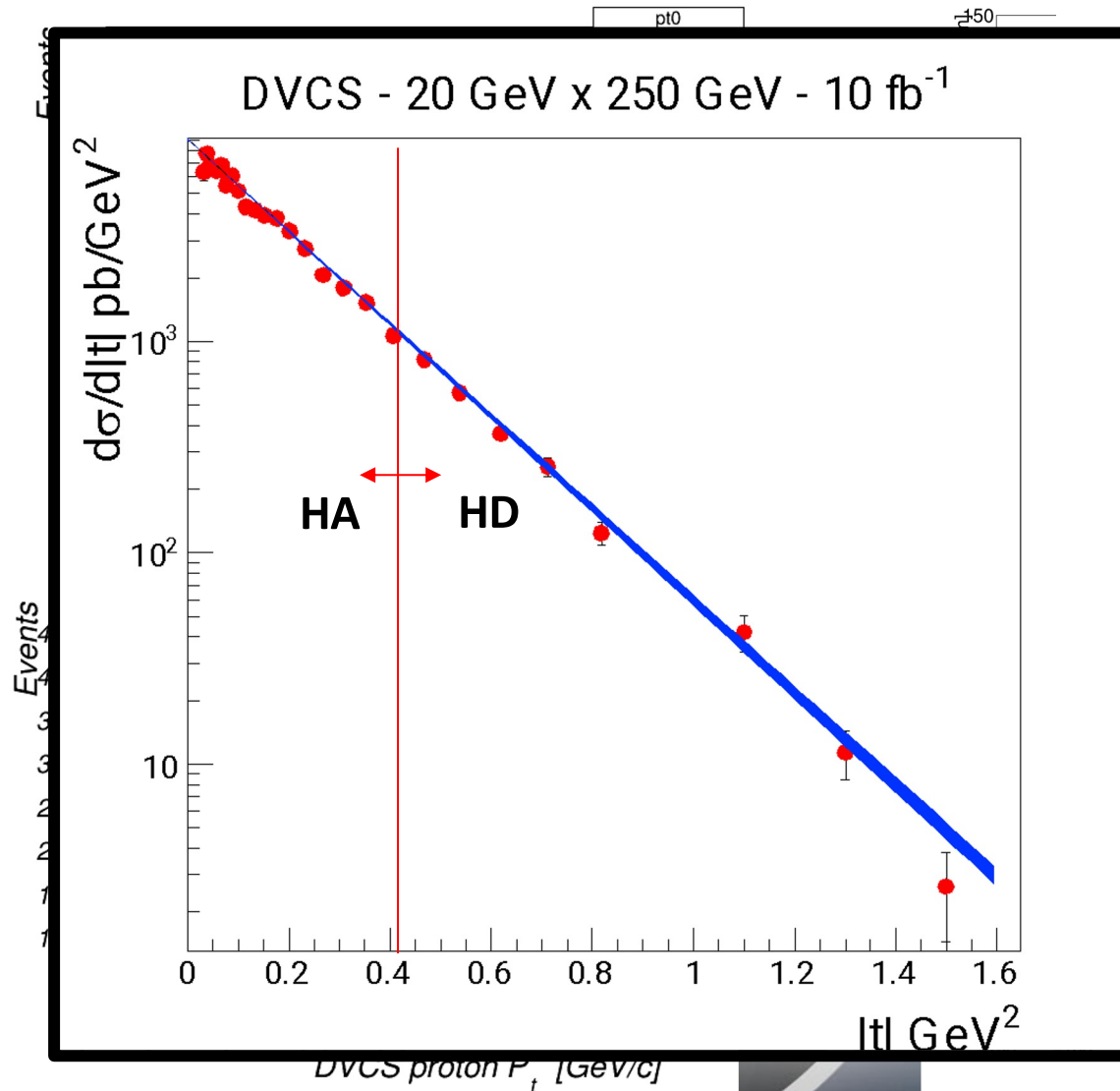


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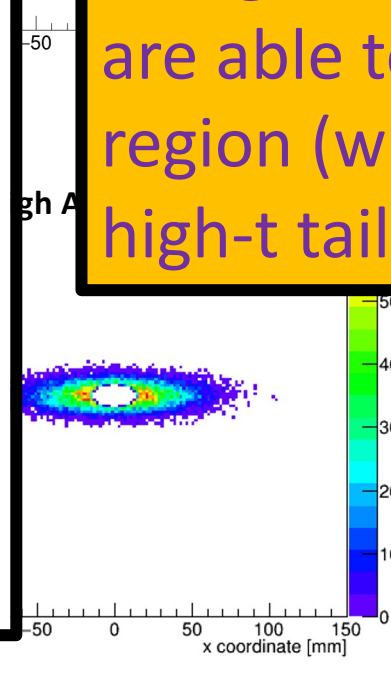
High Acceptance: larger β^* at IP, smaller $\beta(z = 30m)$ -> lower lumi., smaller beam at RP

Digression: Machine Optics (IP6)

275 GeV DVCS Proton Acceptance



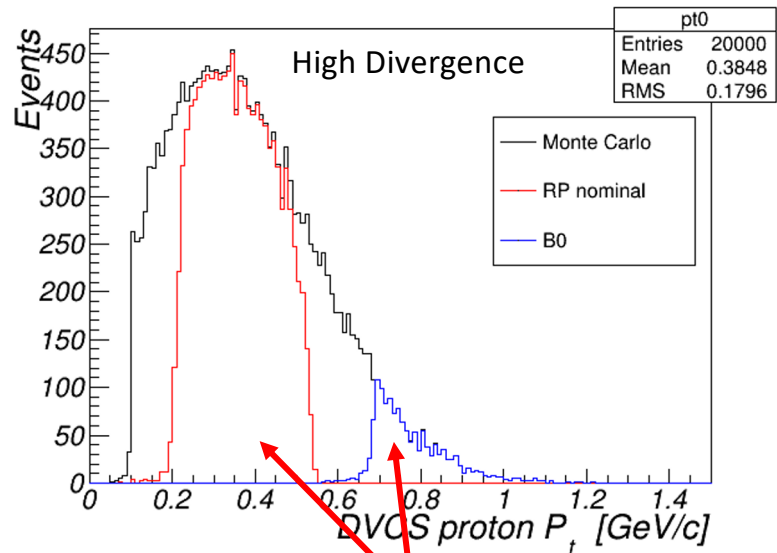
Using the two configurations, we are able to measure the low-t region (with better acceptance) and high-t tail (with higher luminosity).



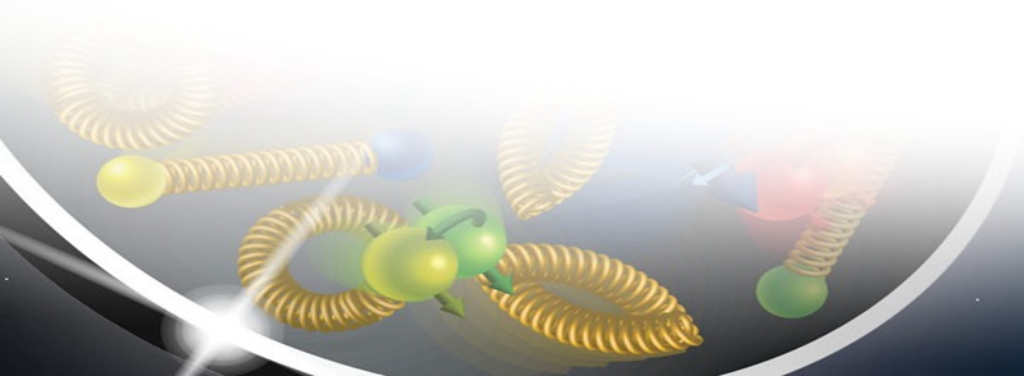
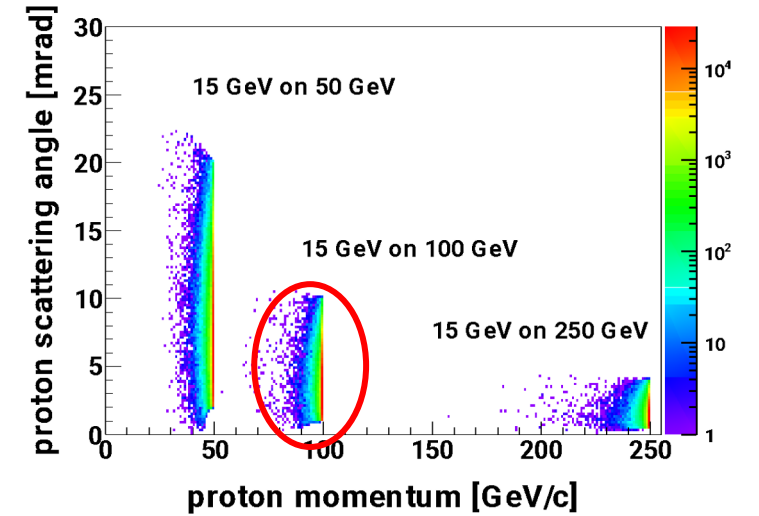
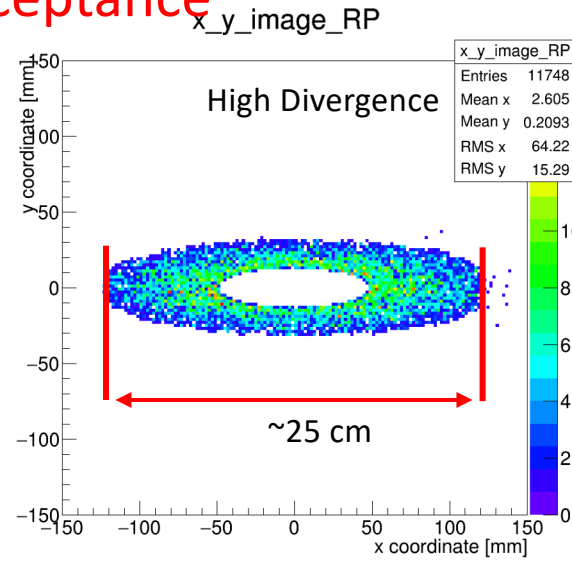
High Acceptance: larger β^* at IP, smaller $\beta(z = 30m)$ -> lower lumi., smaller beam at RP

Digression: Machine Optics (IP6)

100 GeV DVCS Proton Acceptance

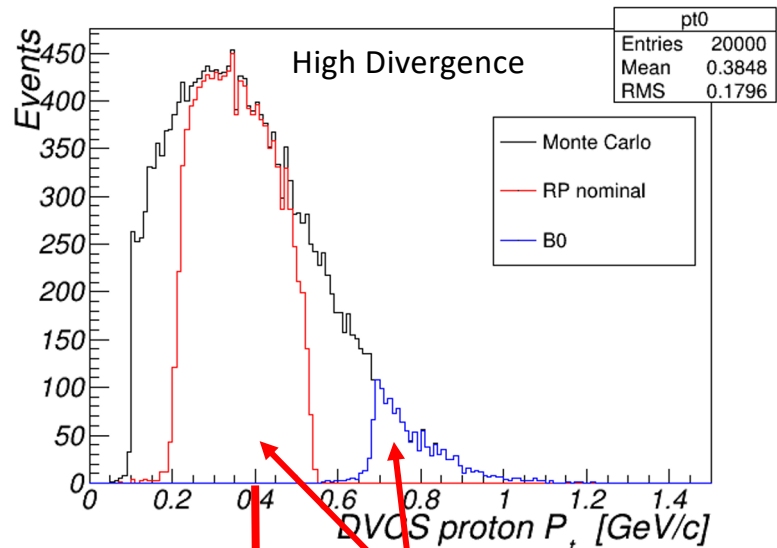


Need both detector systems together here!

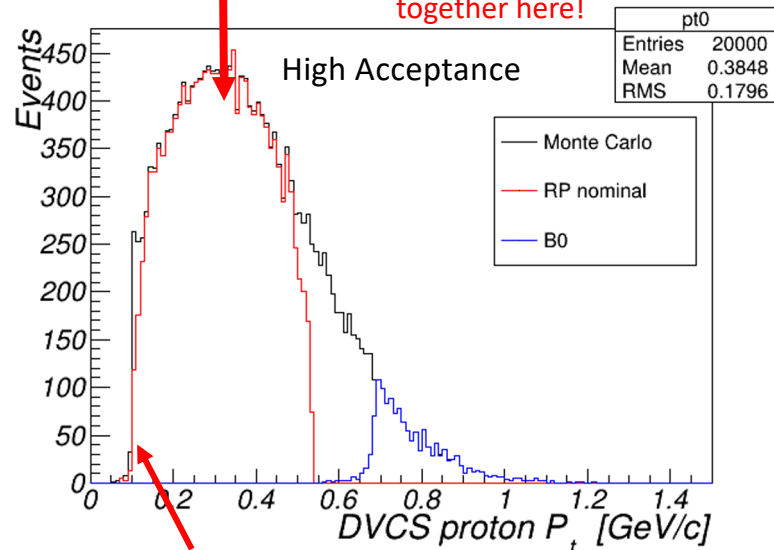
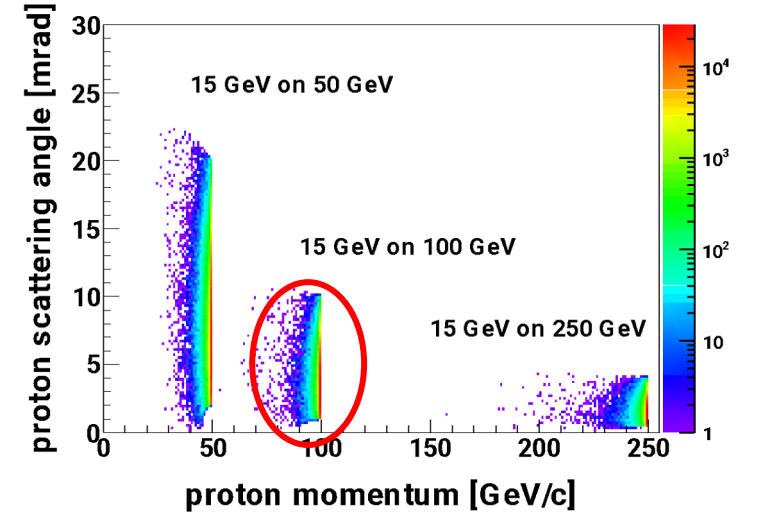
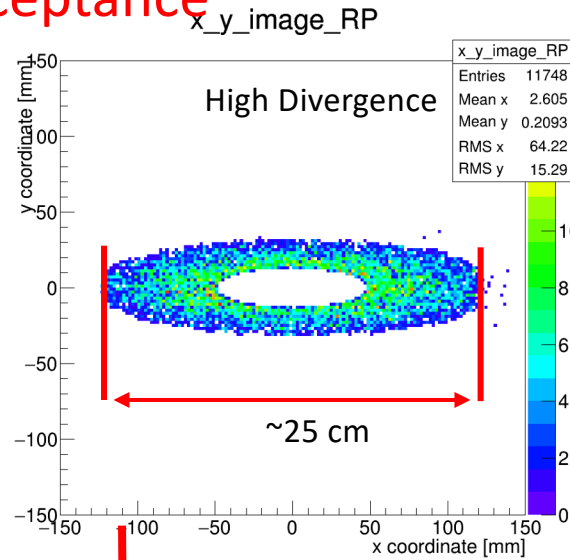


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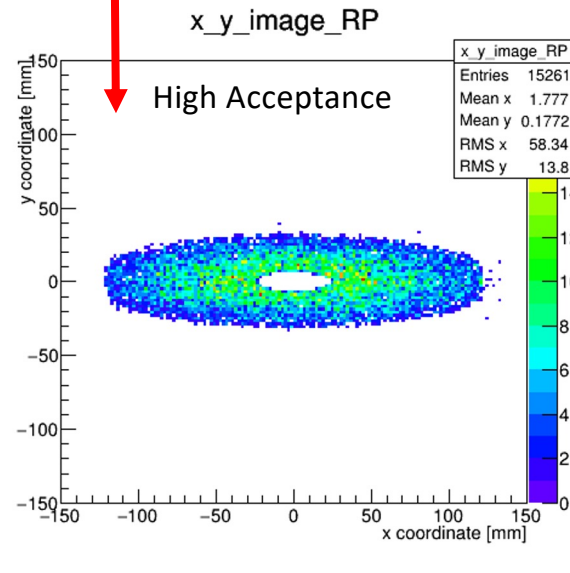
100 GeV DVCS Proton Acceptance



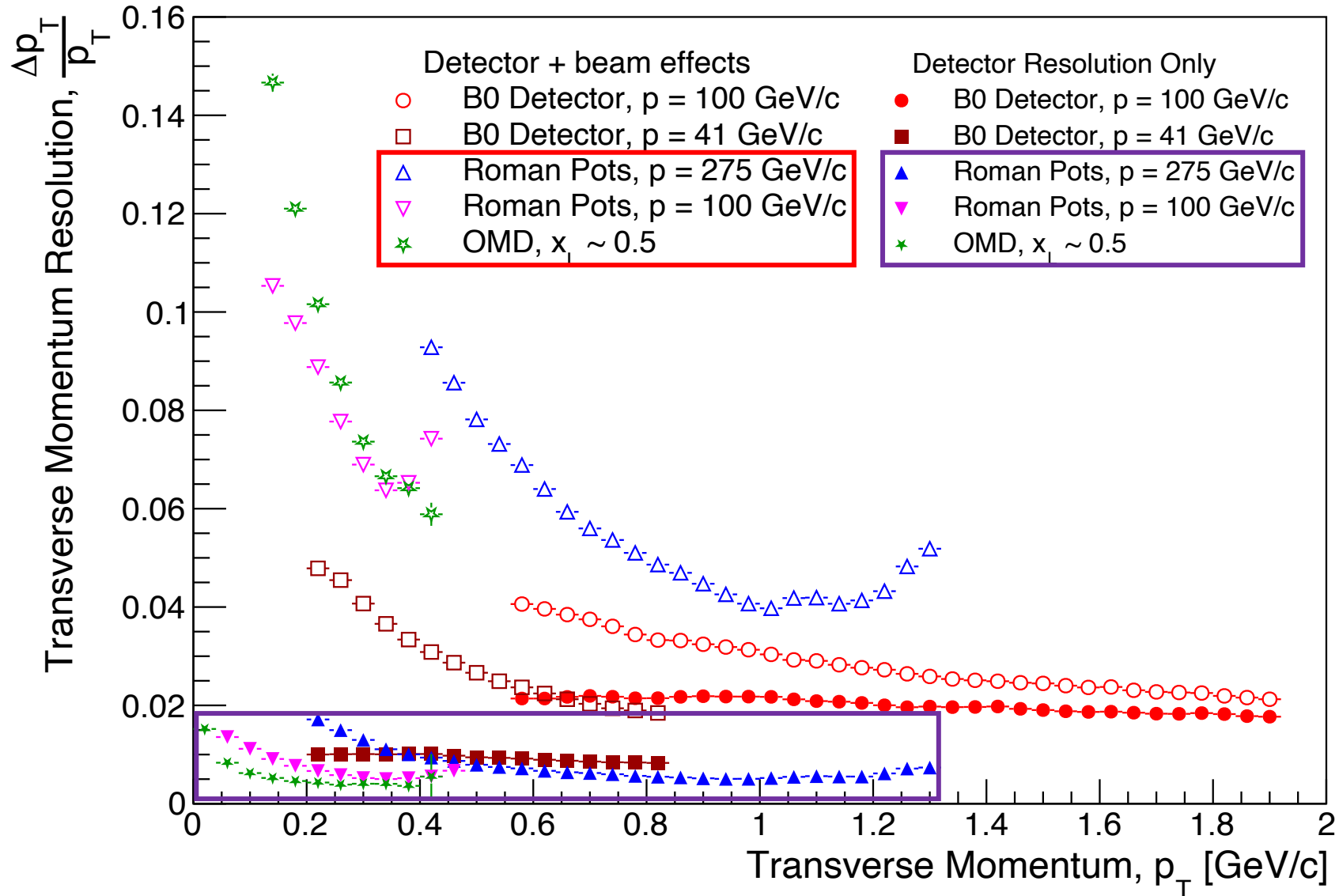
Need both detector systems together here!



Improves low p_t acceptance.



Summary of Detector Performance

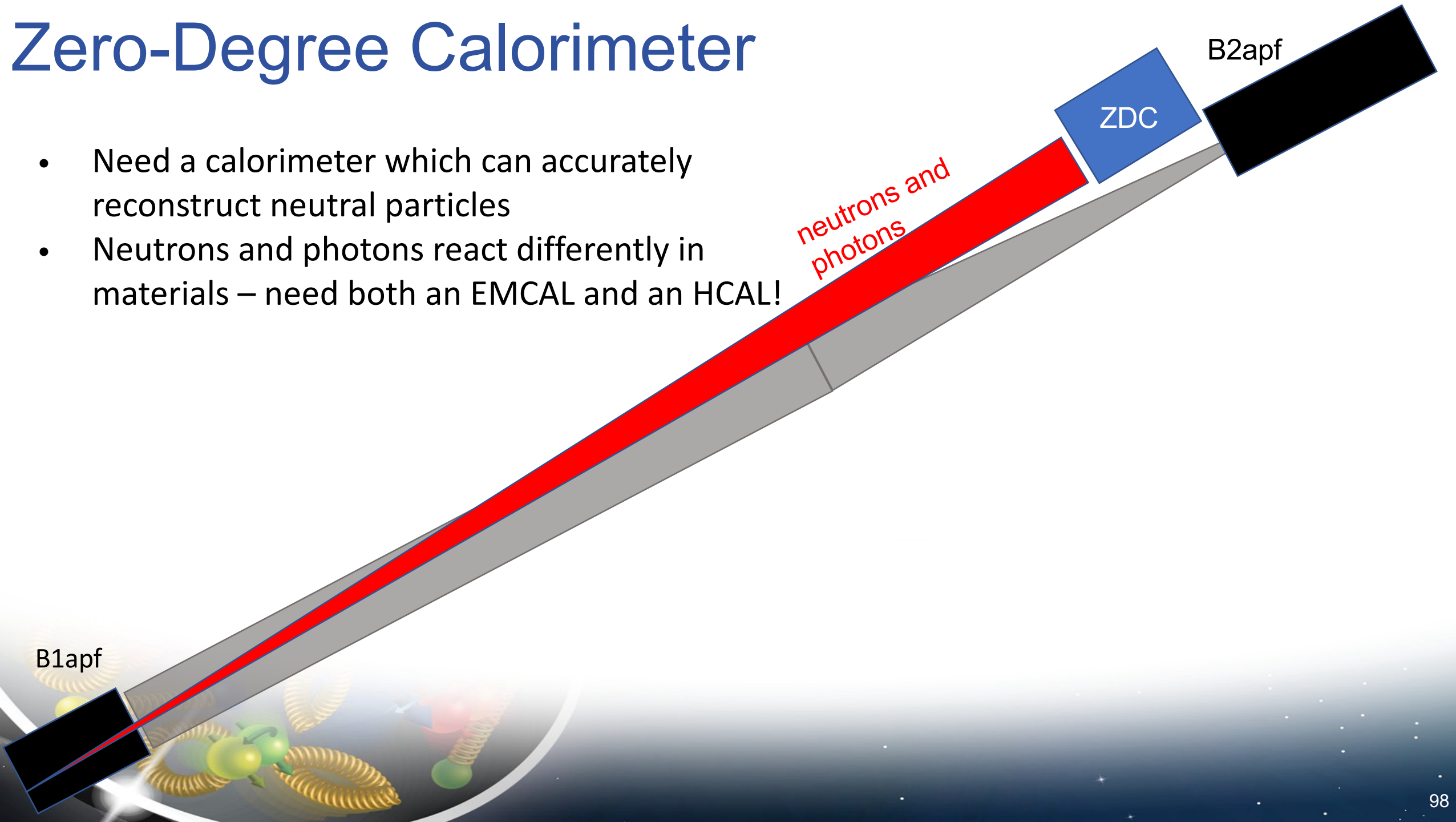


- All beam effects included!
 - Angular divergence.
 - Crossing angle.
 - Crab rotation/vertex smearing.

Beam effects the dominant source of momentum smearing!

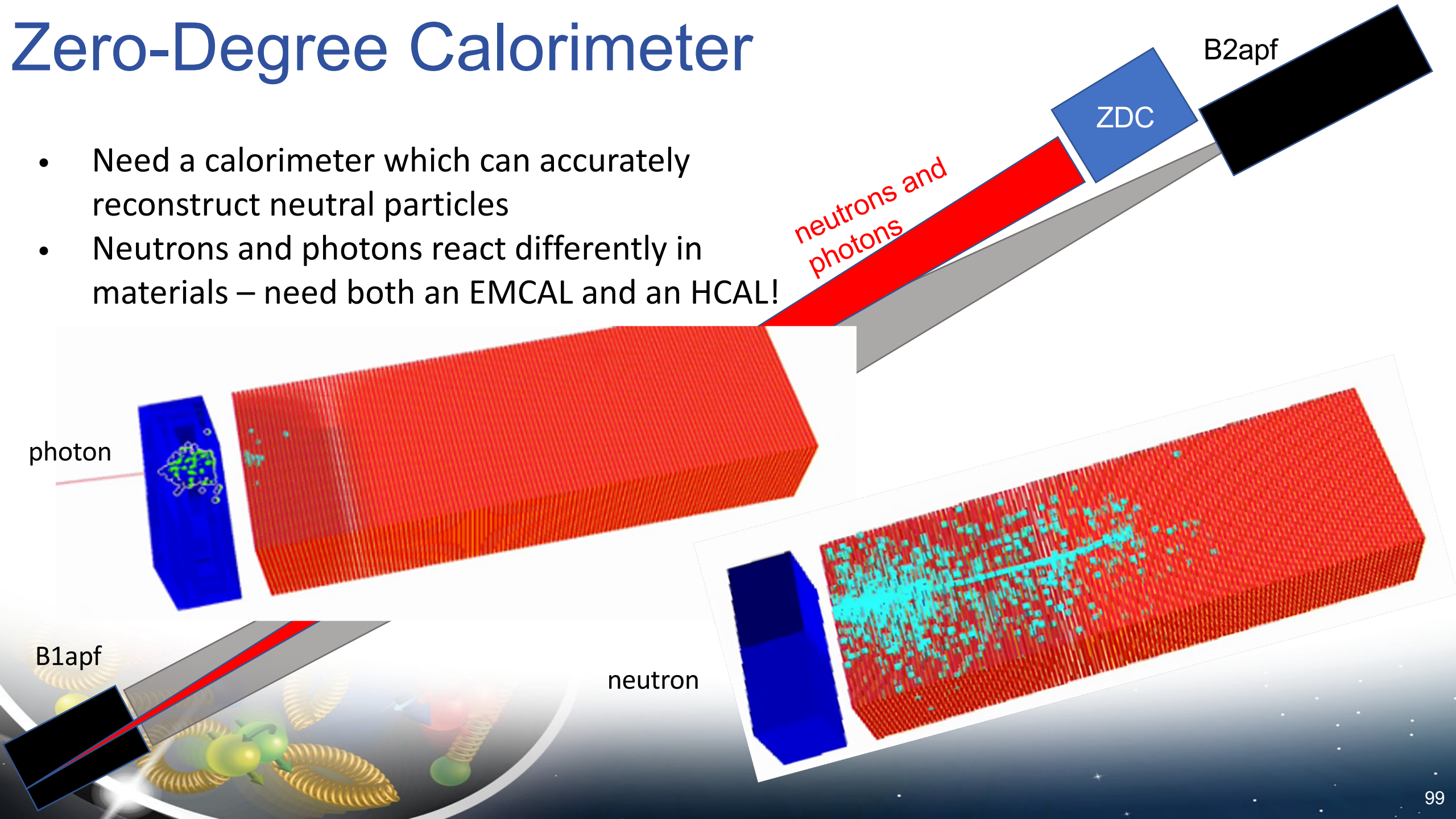
Zero-Degree Calorimeter

- Need a calorimeter which can accurately reconstruct neutral particles
- Neutrons and photons react differently in materials – need both an EMCAL and an HCAL!



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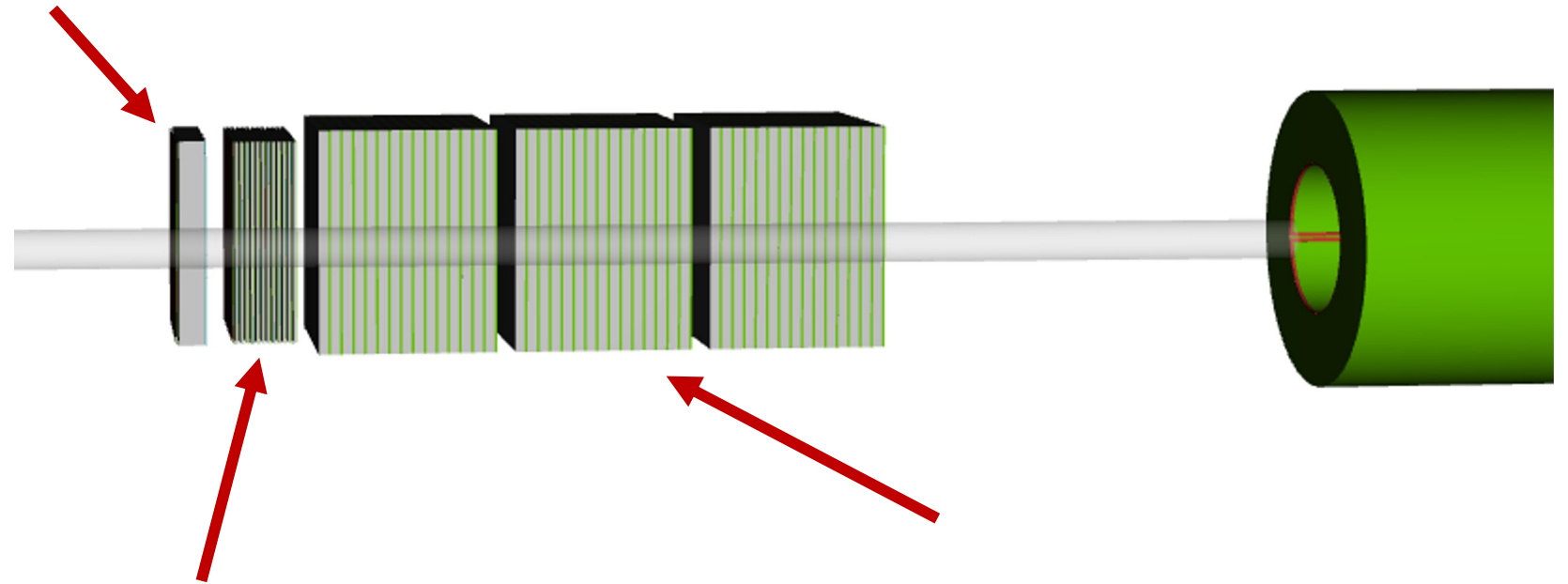
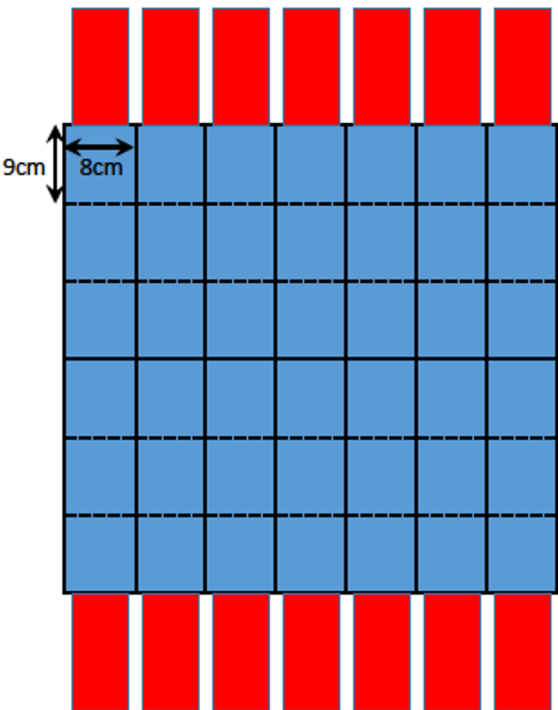


ZDC - What's New

- 1st Silicon & crystal calorimeter (PbWO₄ or LYSO):
 - **Smaller lateral dimension** (x, y) = (56, 54) cm.

Overall length within 2m limit

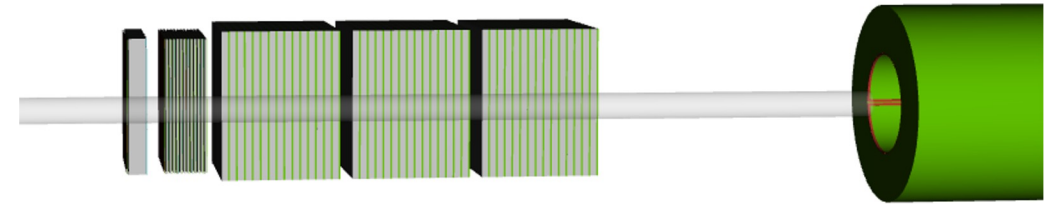
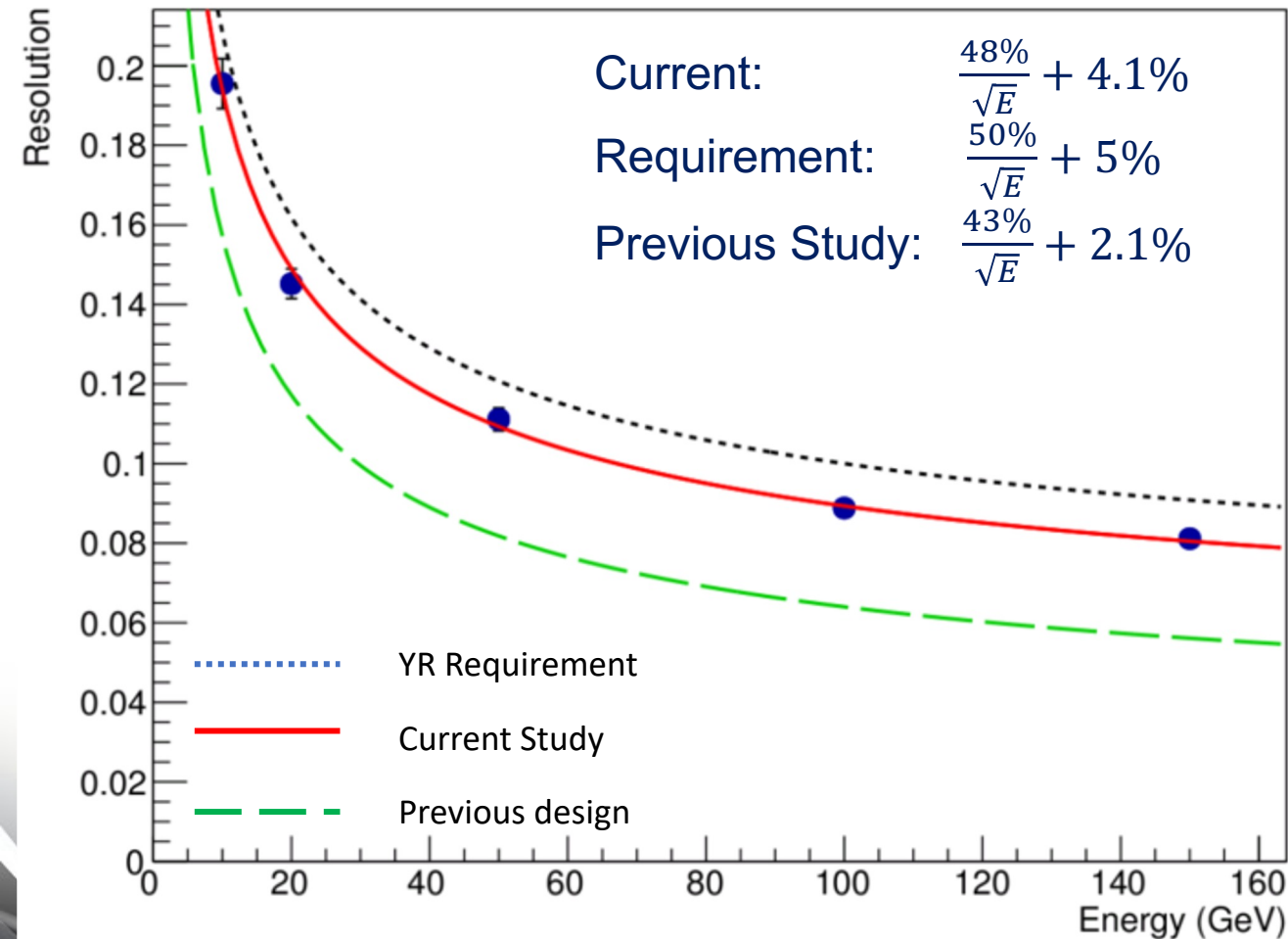
Readout setup
from top & bottom



- W/Silicon Imaging EMCAL
 - Transverse size (x,y) = (56, 54) cm
 - 12 layers ($\sim 24\chi_0$)
- Pb-Scintillator (+ fused silica)
 - Towers of 10cm x 10cm x 48cm, each module 60cm x 60cm x 48cm
 - 3 modules

ZDC - Performance

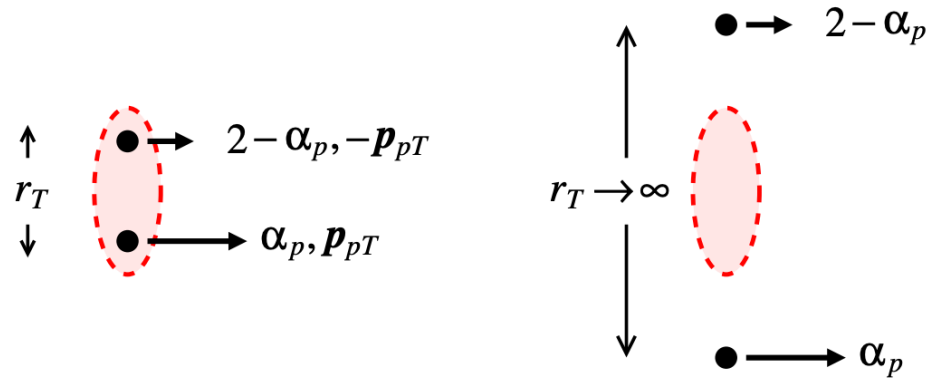
Neutron Energy Resolution



- Energy resolution in the new design acceptable → Optimization, test of different ideas within the size limit.
- **Next steps:**
 - Implementation of reconstruction
 - Position resolution & shower development study ongoing for the imaging part of HCAL

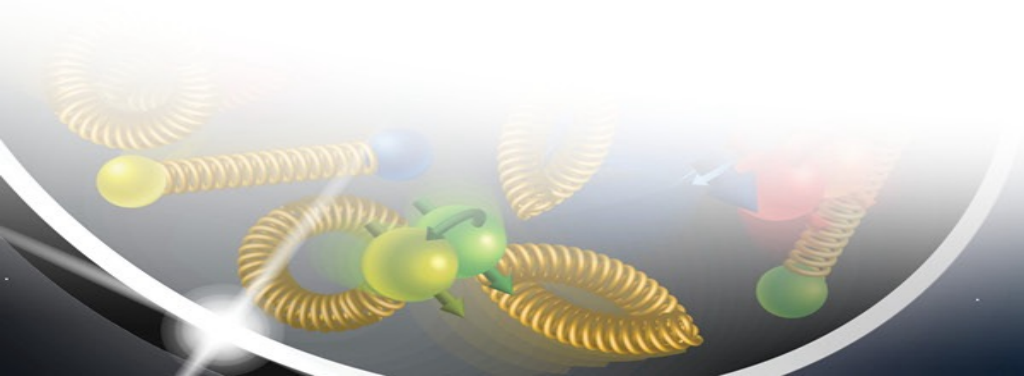
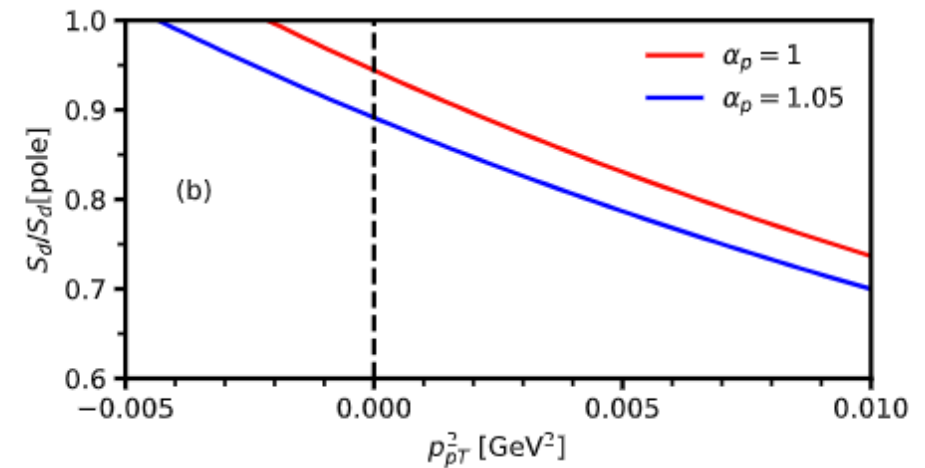
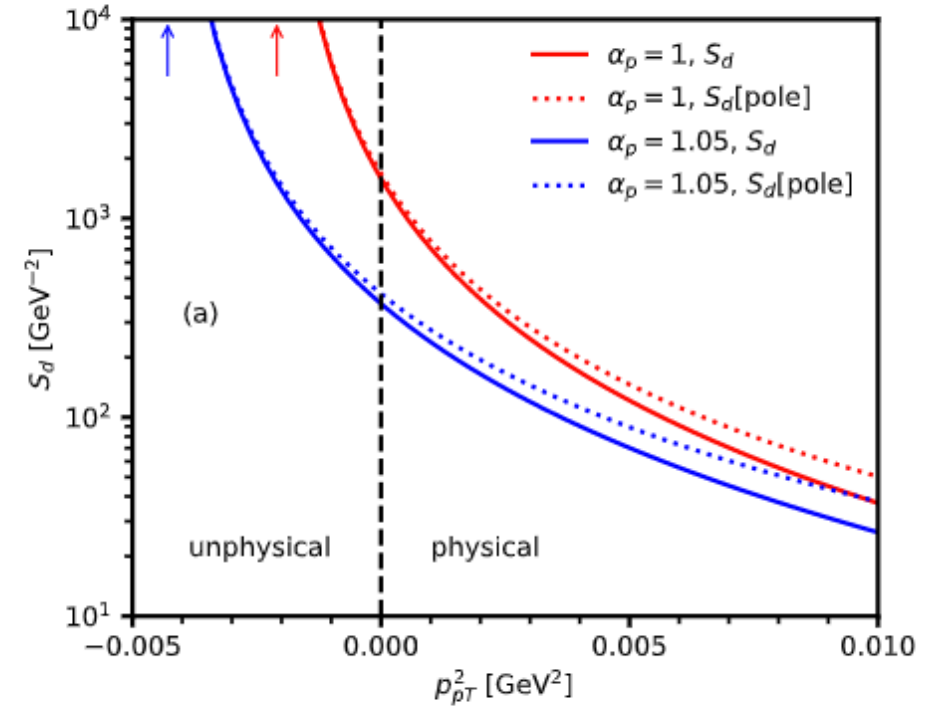
Pole Extrapolation

C. Weiss and W. Cosyn
Phys. Rev. C **102**, 065204 (2020)



$p_{pT}^2 > 0$
physical region

$p_{pT}^2 \rightarrow -a_T^2$
pole extrapolation

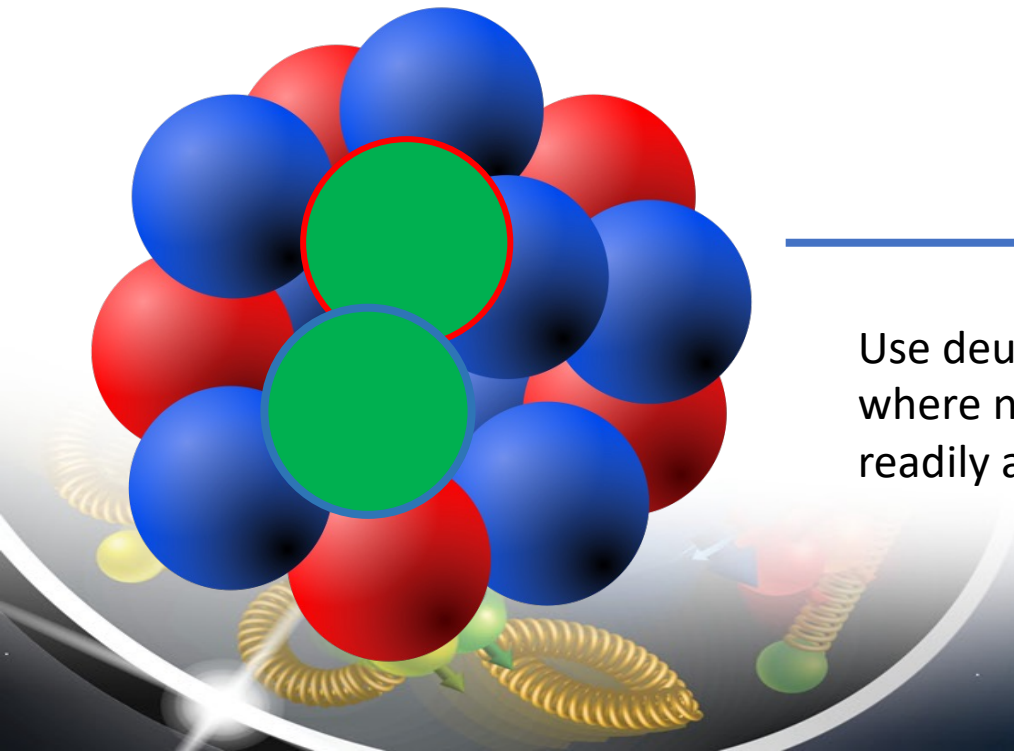


Short-Range Correlations

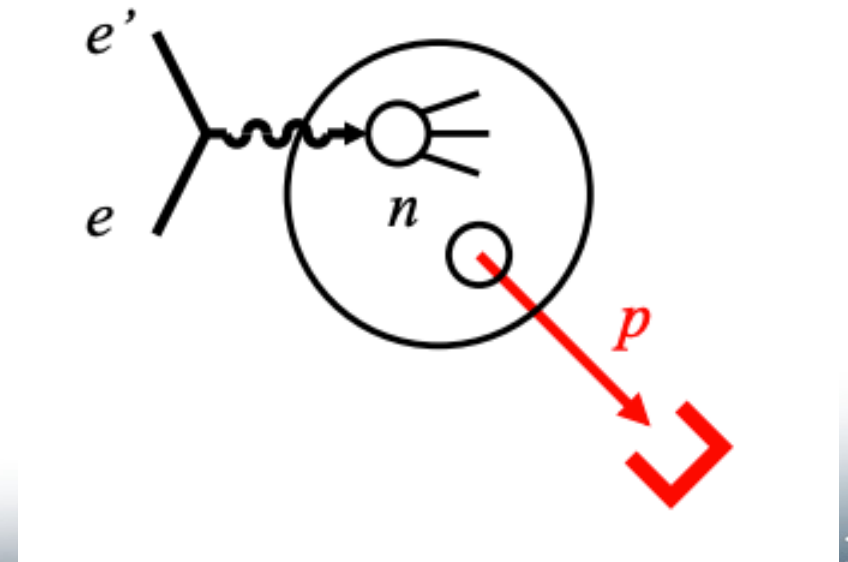
“The nucleus can often be approximated as an independent collection of protons and neutrons confined in a volume, but for short periods of time, the nucleons in the nucleus can strongly overlap. This quantum mechanical overlapping, known as a nucleon-nucleon short-range correlation, is a manifestation of the nuclear strong force, which produces not only the long-range attraction that holds matter together, but also the short-range repulsion that keeps it from collapsing.”

Excerpt from: https://www.jlab.org/research/nucleon_nucleon

Lots of SRC pairs!!! -> Really tough!



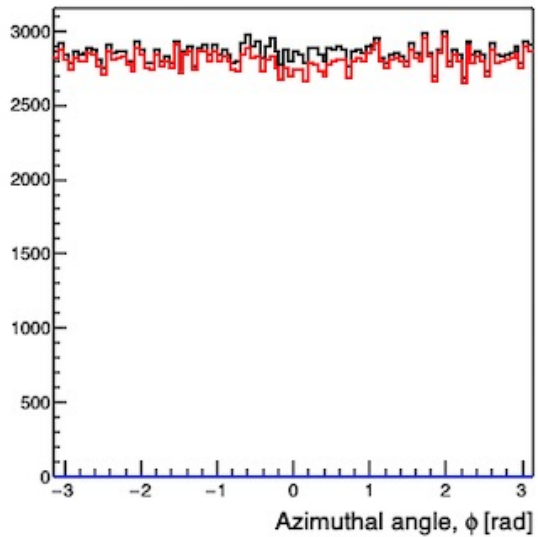
Use deuteron as “SRC laboratory”, where nucleon kinematics are readily accessible.



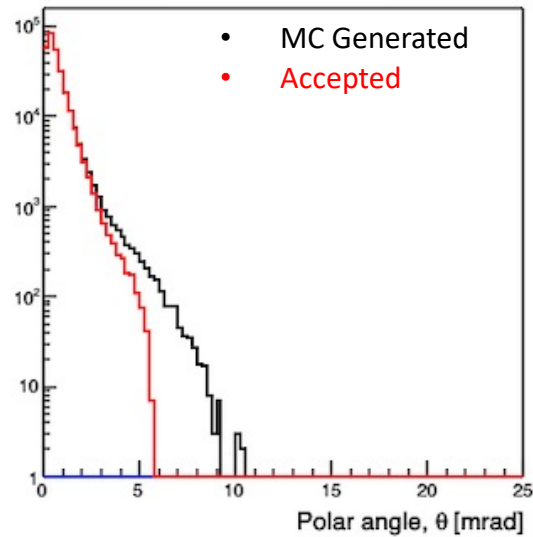
Short-Range Correlations in Deuterons

Z. Tu, A. Jentsch *et al.*, Phys. Lett. B, **811** (2020)

Protons

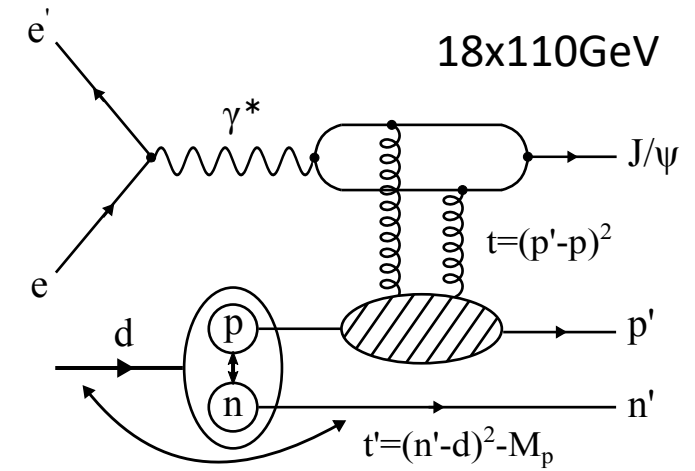


Protons

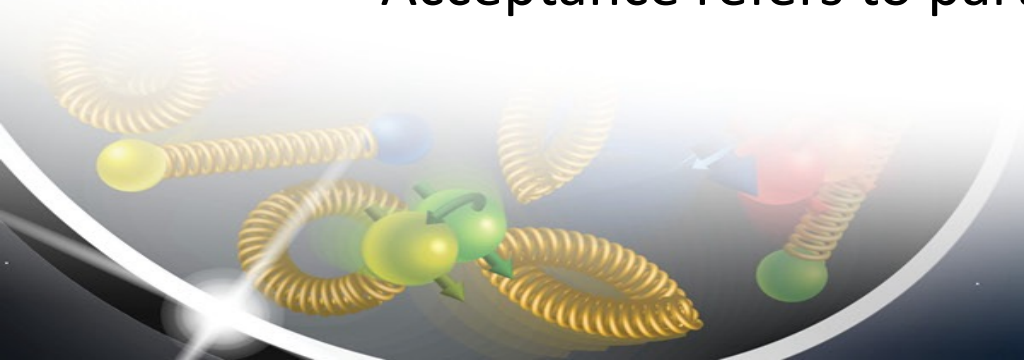


Proton "spectator" case.

Particular process in BeAGLE:
incoherent diffractive J/ψ
production off bounded nucleons.

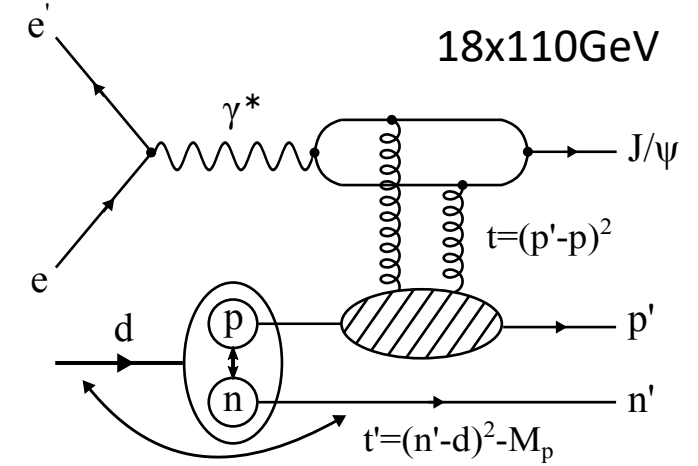


MC generated events shown in black – “accepted” protons in red.
Acceptance refers to particles which are actually captured by the detector.



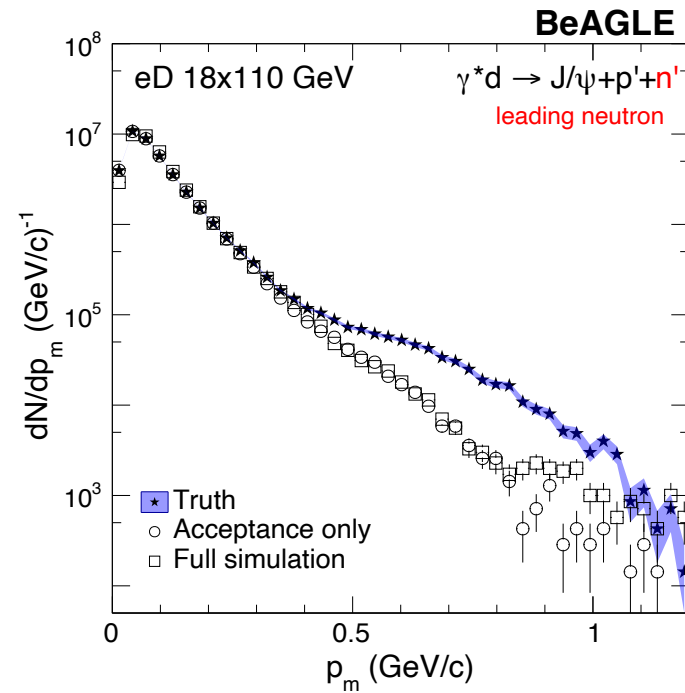
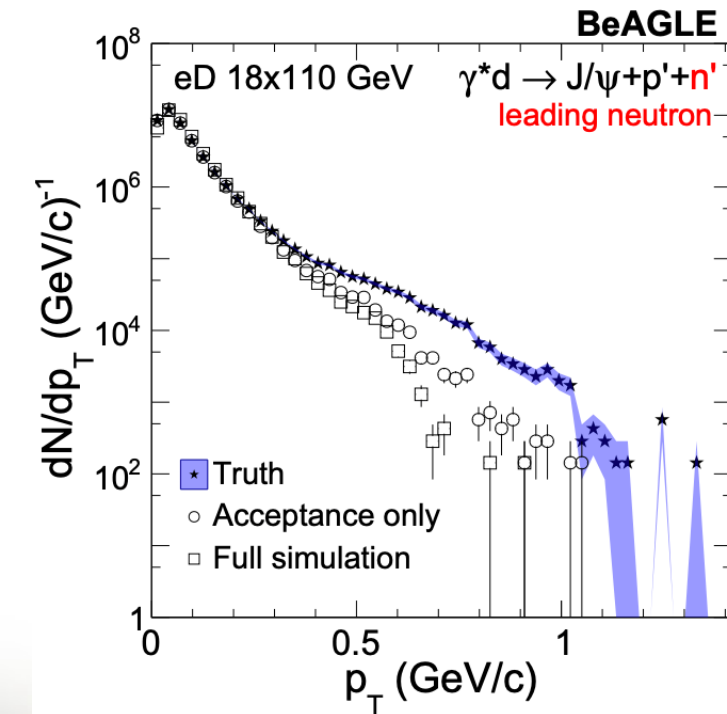
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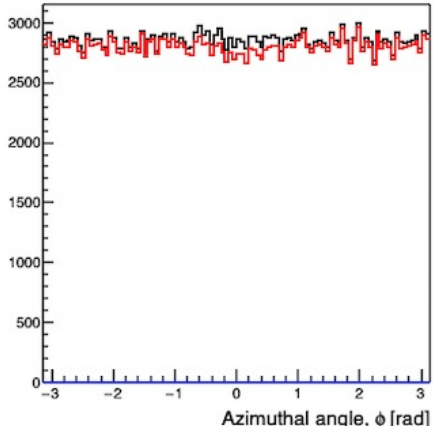
- Spectator kinematic variables reconstructed over a broad range.
- **All detector and beam effects included in the full GEANT simulations!**
 - Bin migration is observed due to smearing in the reconstruction.



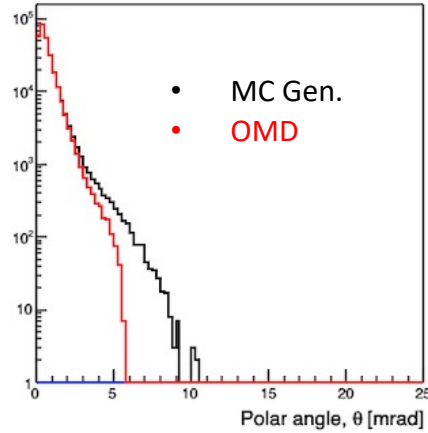
- In the proton spectator case, essentially all spectators tagged up to $p_T \sim 600 \text{ MeV/c}$.
- Active neutrons only tagged up to 4.5 mrad \rightarrow double-tagging efficiency very low.

e+d Spectator Tagging

Protons

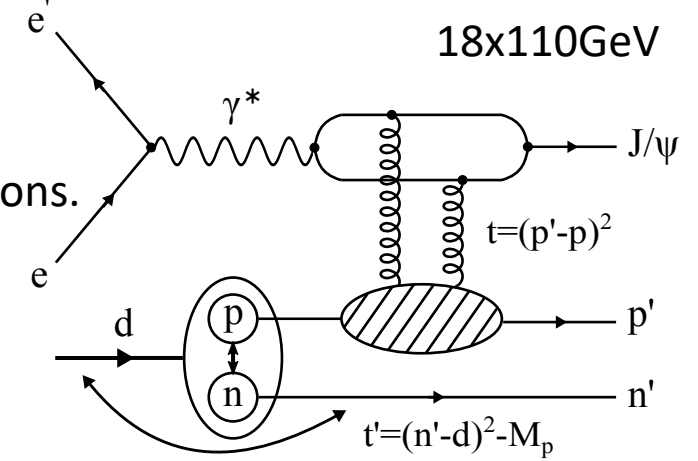


Protons

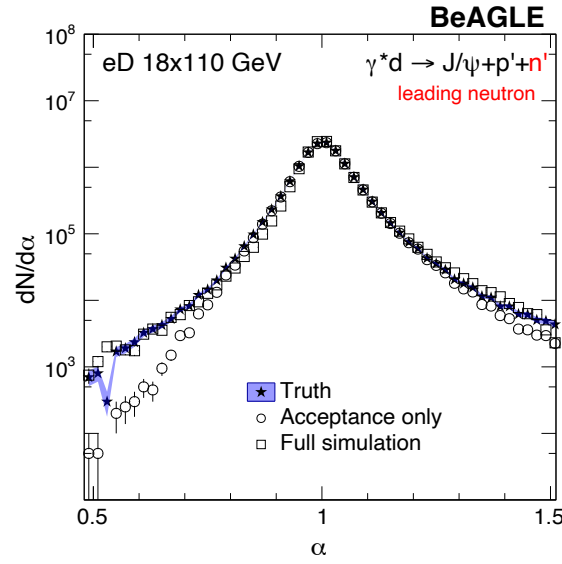
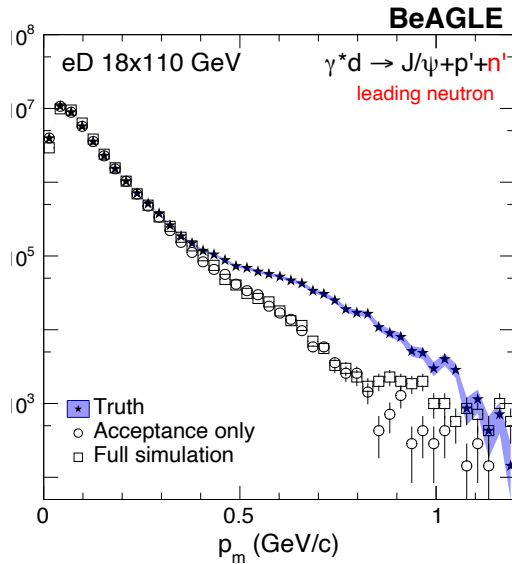
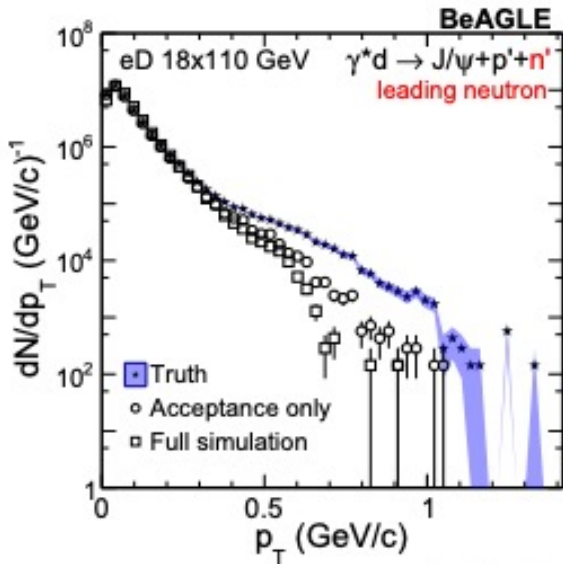


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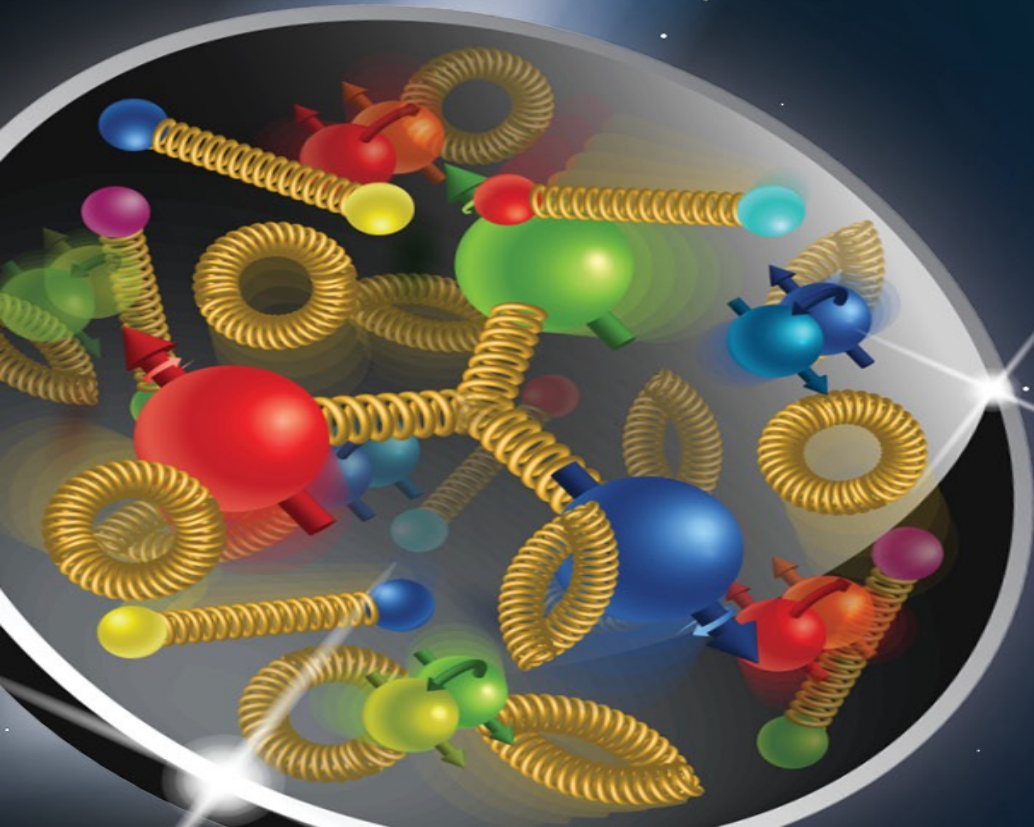
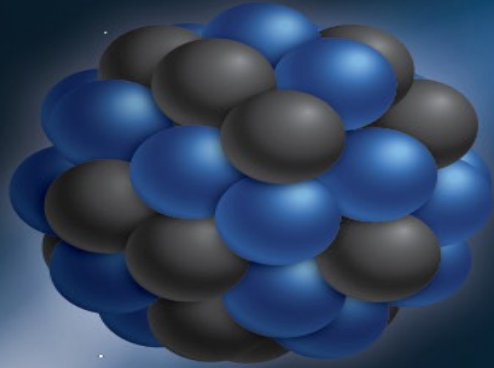
Particular process in BeAGLE:
incoherent diffractive J/psi
production off bounded nucleons.



Spectator kinematic variables reconstructed over a broad range. Bin migration is observed due to smearing in the reconstruction. Each plot shows the MC (closed circles), acceptance effects only (open circles), and full reconstruction (open squares).



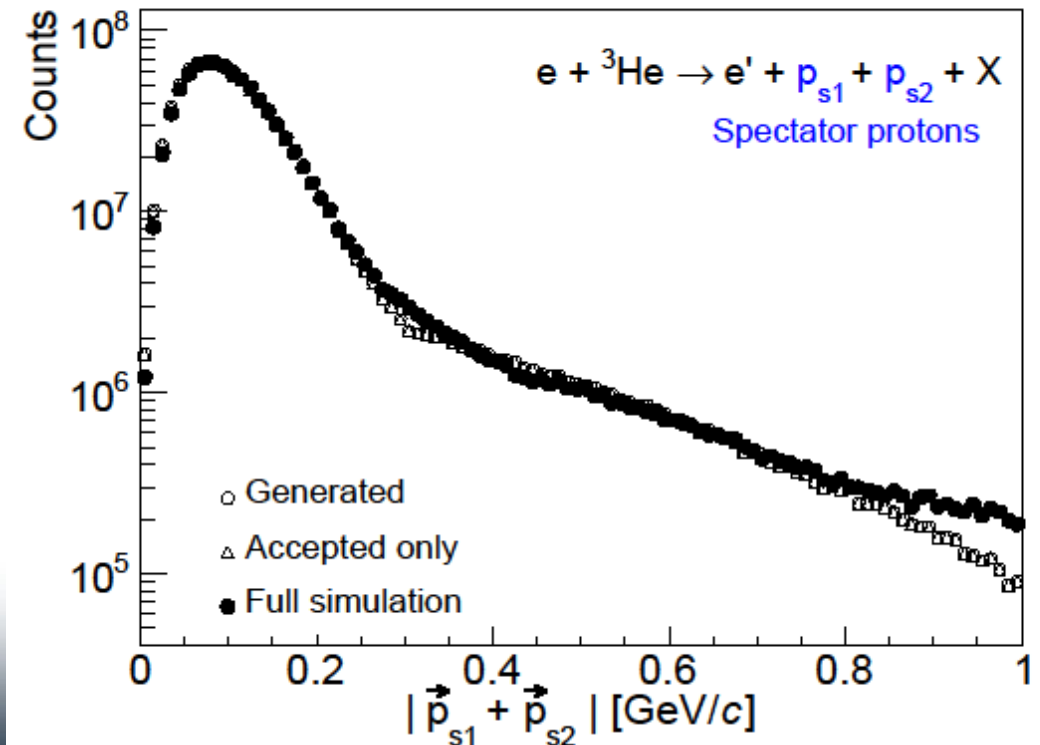
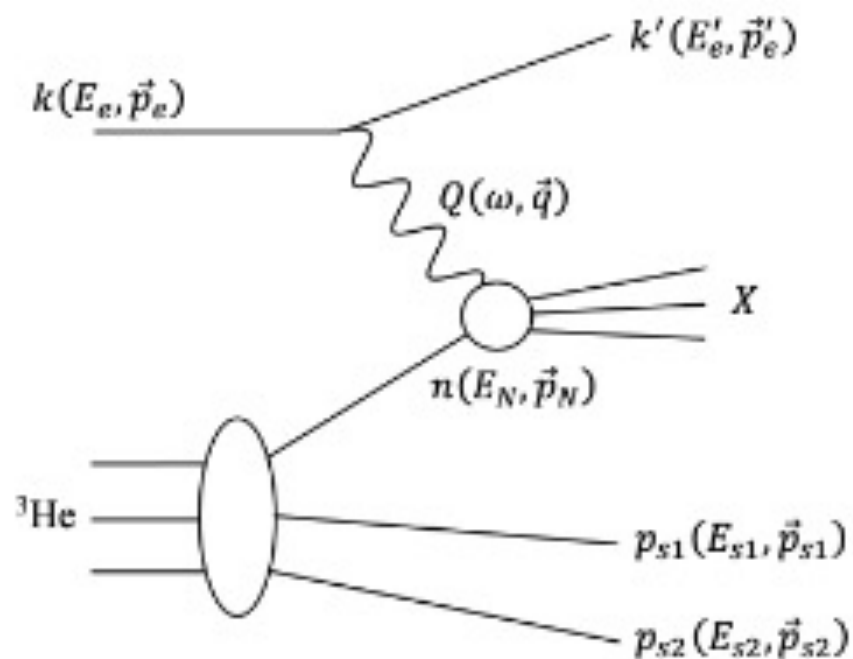
- In the proton spectator case, essentially all spectators tagged.
- Active neutrons only tagged up to 4.5 mrad.



Light nuclei – Helium-3: Neutron Spin Structure

Neutron Spin Structure in He3

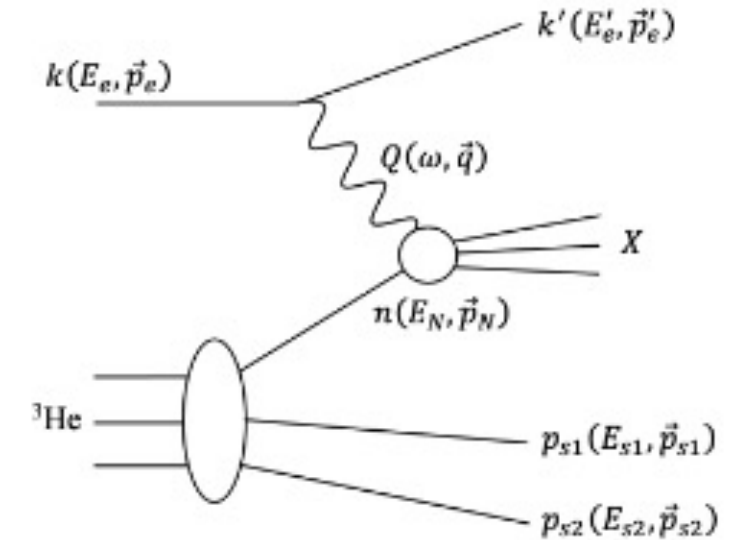
- Studies of neutron structure with a *polarized* neutron.
- More challenging final state tagging since *both* protons must be tagged.
- MC events generated with CLASDIS in fixed-target frame, and then boosted to collider frame.



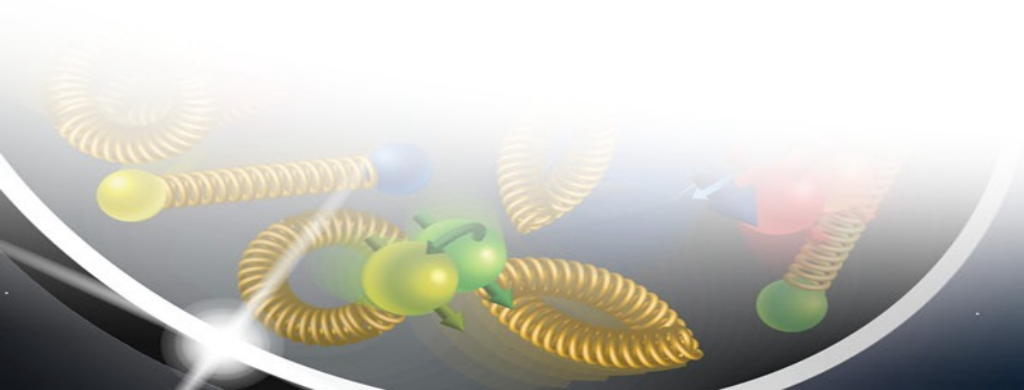
Neutron Spin Structure in He3

- Spin structure probed via spin asymmetries!

$$A_1^{3\text{He}} = \underbrace{P_n \frac{F_2^n}{F_2^{3\text{He}}} A_1^n}_{\text{Neutron}} + \underbrace{2P_p \frac{F_2^p}{F_2^{3\text{He}}} A_1^p}_{\text{Protons}}$$



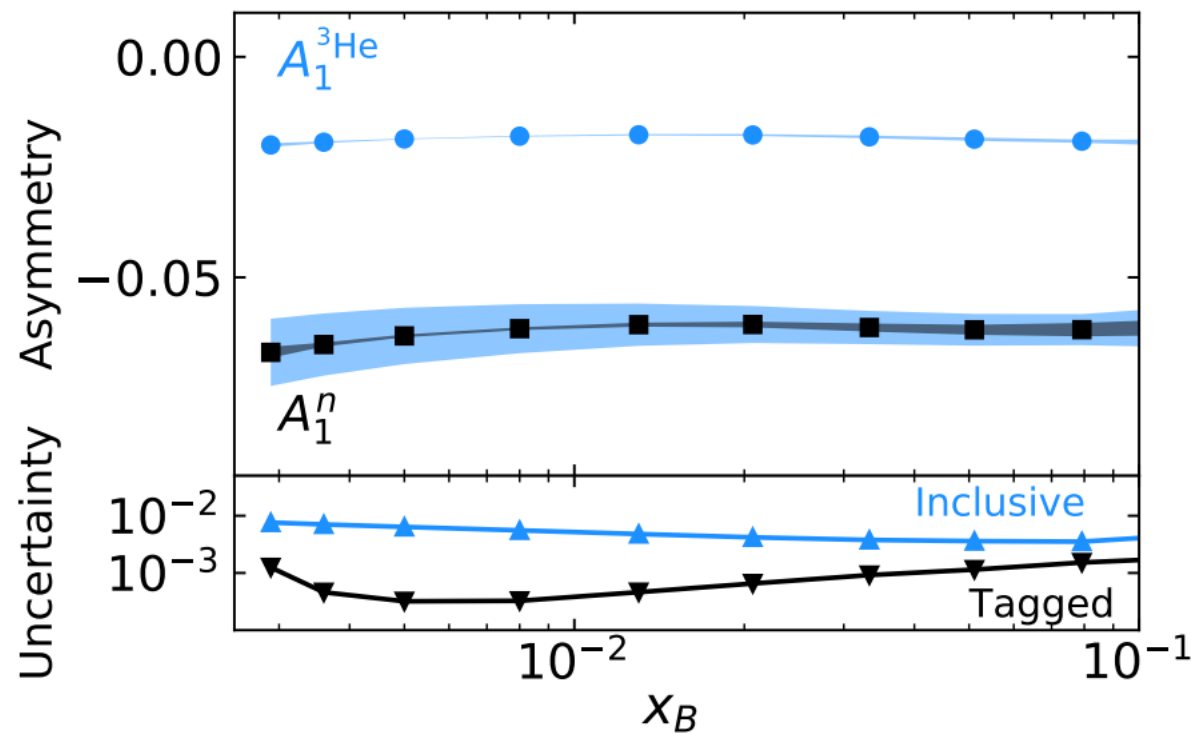
- (double) Tagged DIS measurement capable of measuring A_1^n directly!
- Complementary to measurements at JLAB.



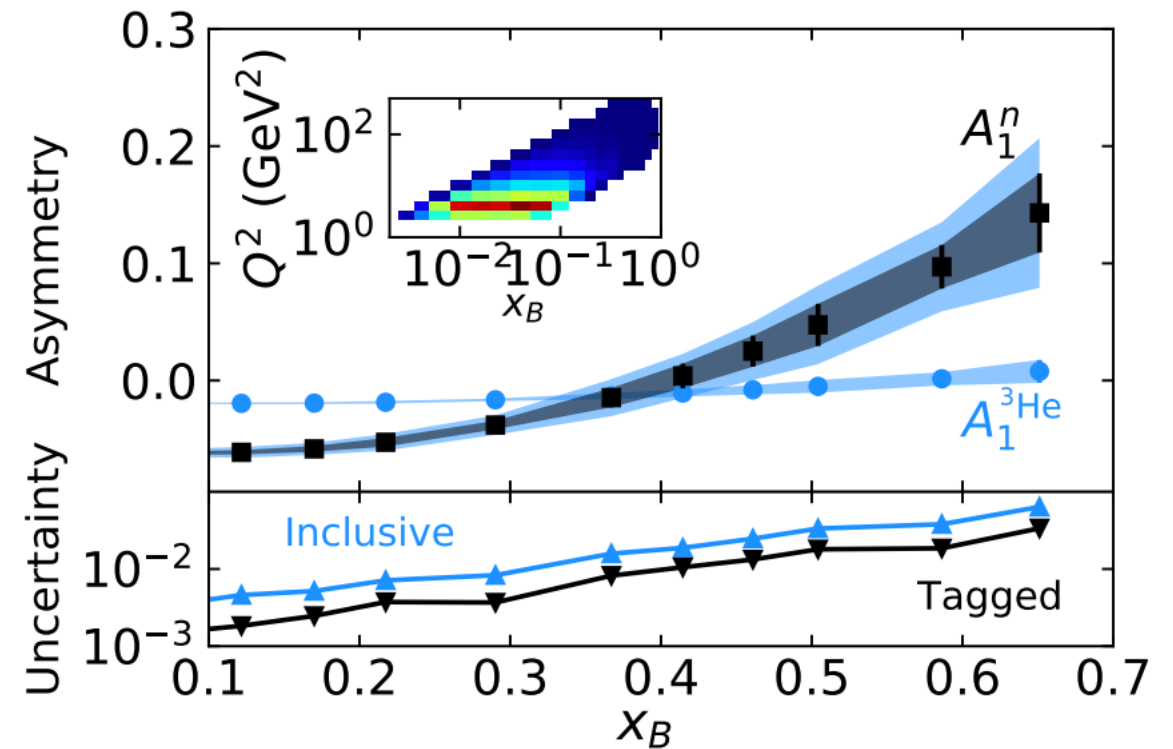
Neutron Spin Structure in He3

- Neutron spin asymmetries can be measured from kinematics of the tagged protons.
- EIC can build upon measurements at JLAB by reducing polarization uncertainties, and opening a broader Q^2 range for study.
- Can aid in our understanding of quark orbital angular momentum in nucleons.

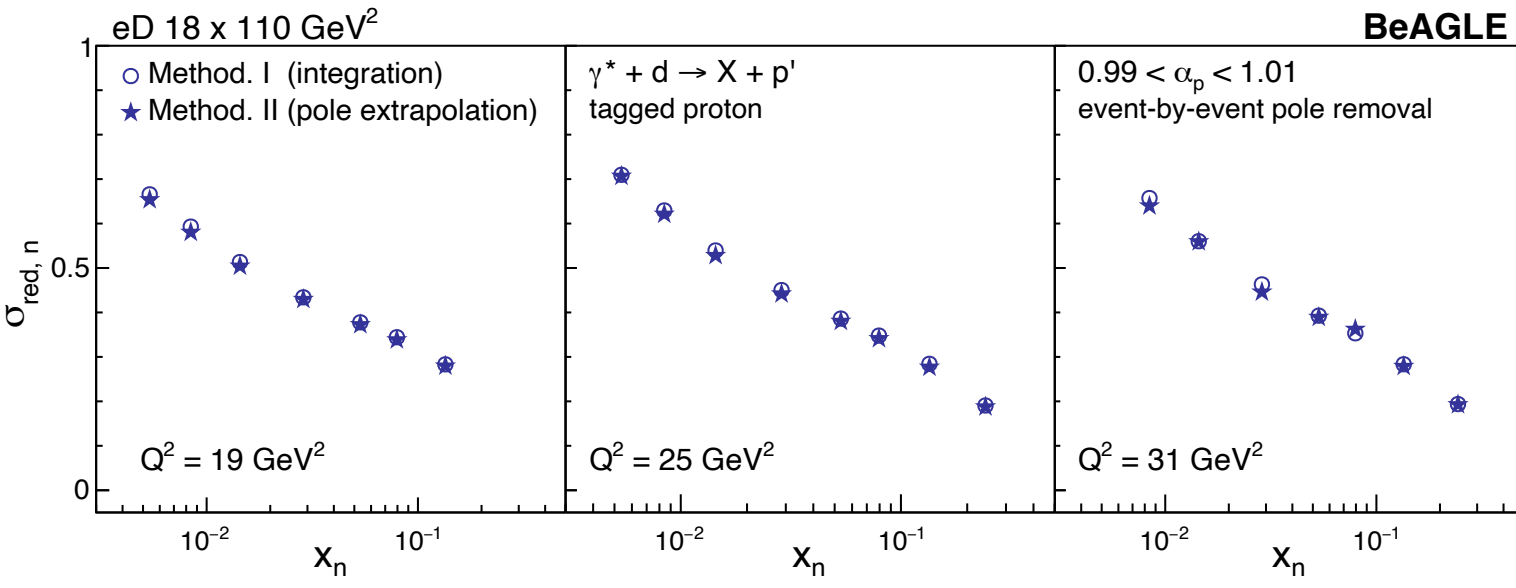
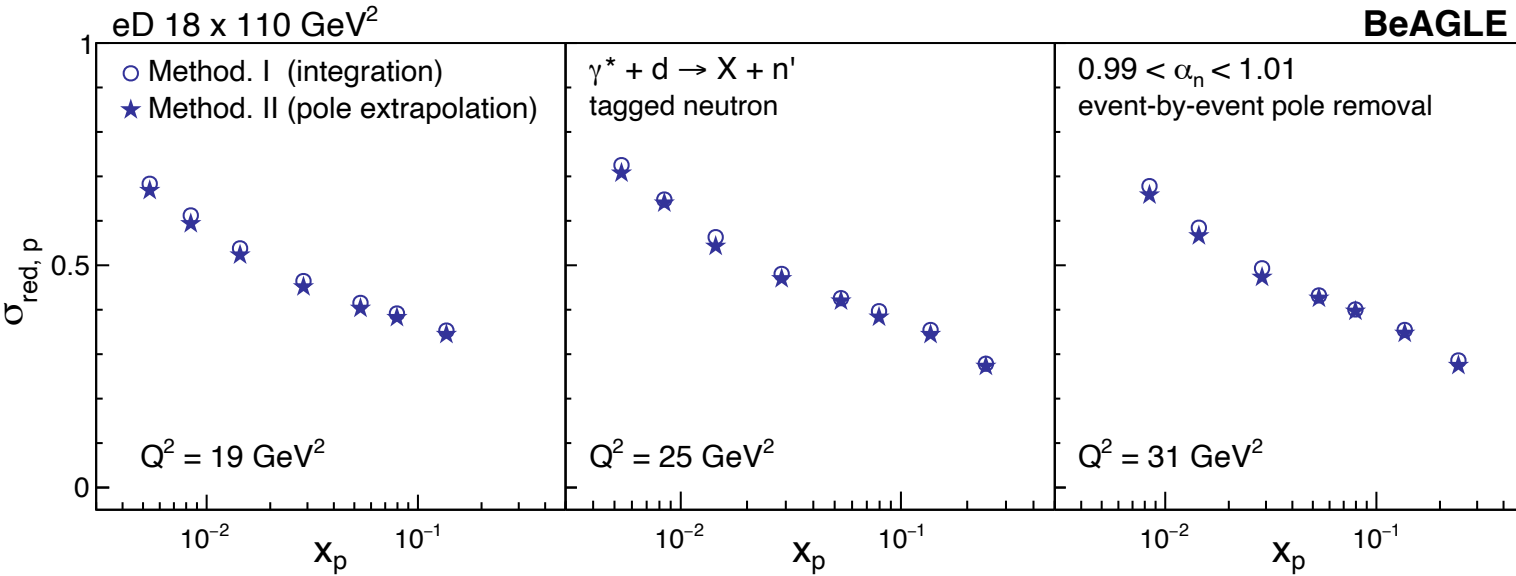
Low-x



High-x

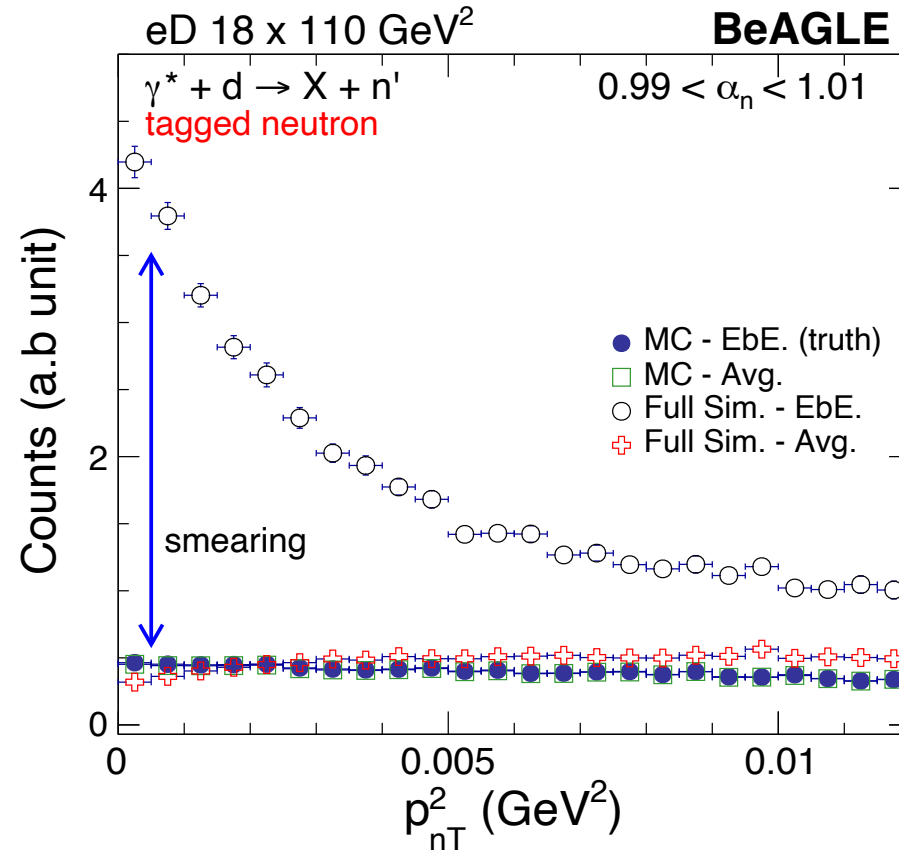
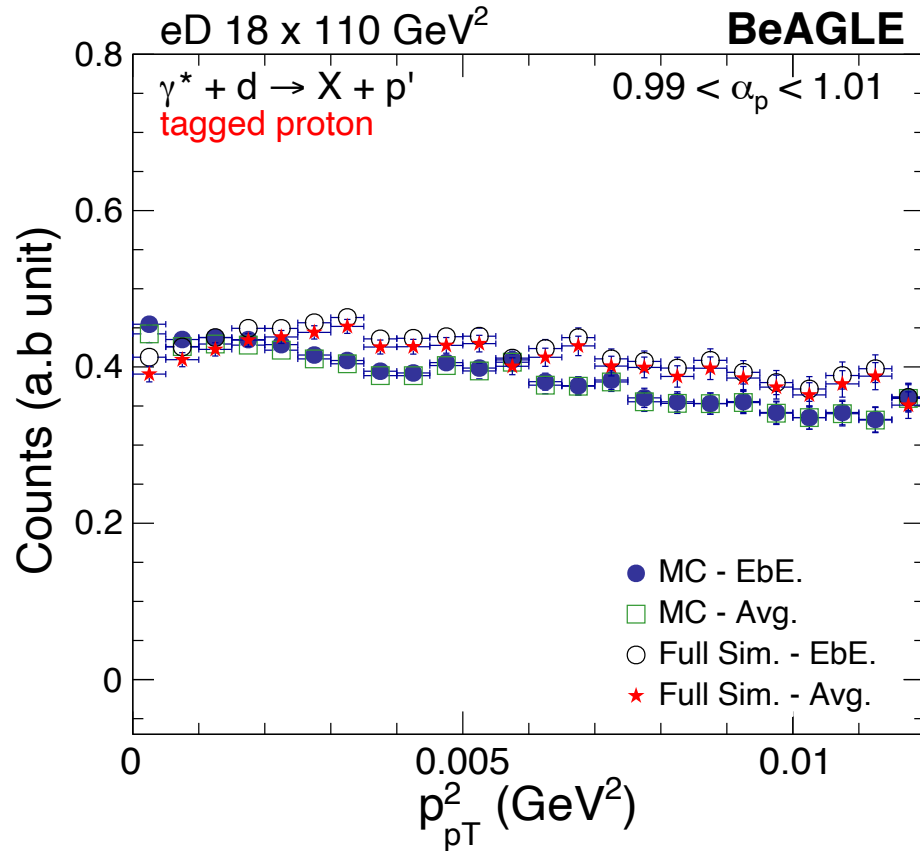


Closure Test – Event by Event Pole Removal



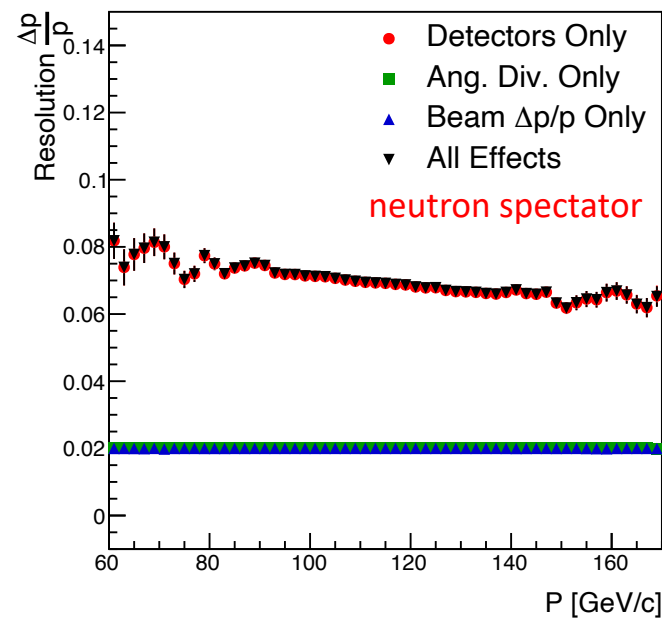
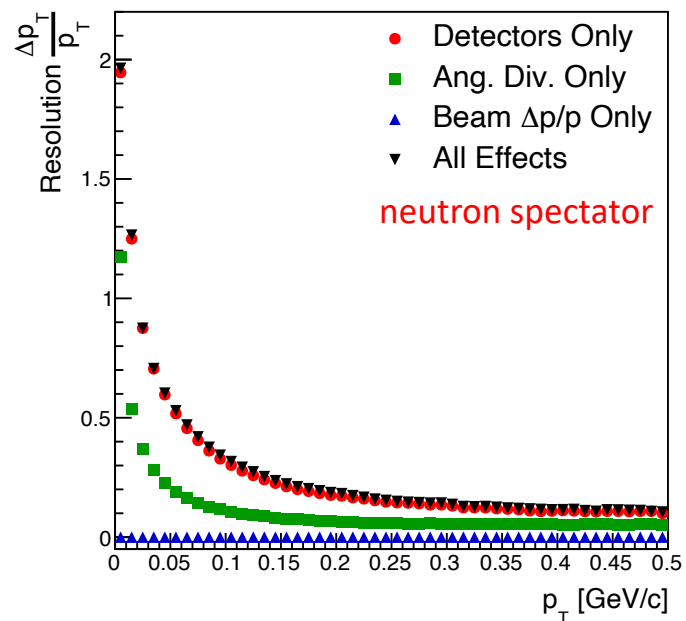
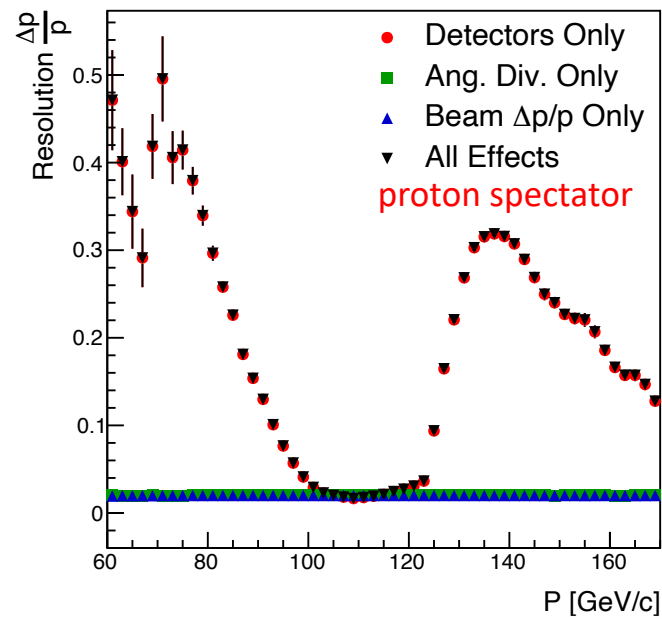
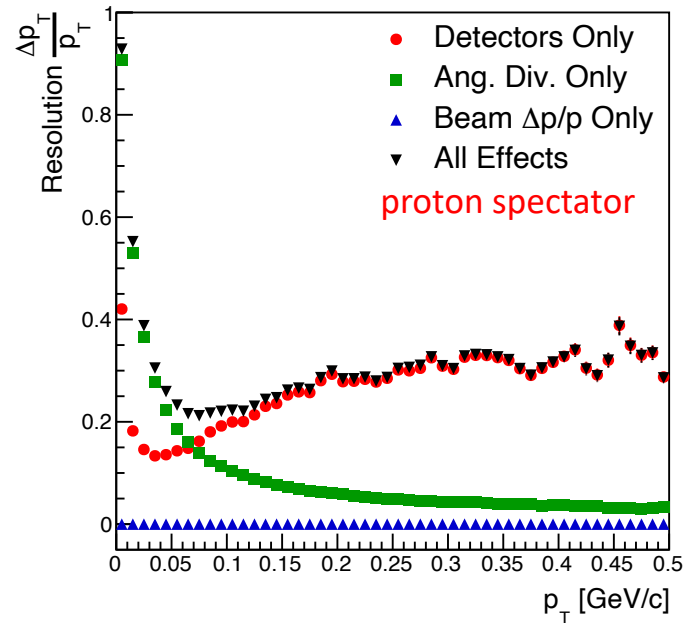
- Pole factor removed using “**event by event (EbE)**” approach.
 - Pole factor calculated and applied for each event (i.e. pole factor calculated for each exact nuclear configuration).
 - Solved discrepancy at generator level.
 - This was also checked using an independent toy MC to confirm there was nothing related to our analysis code causing an issue.
- Remaining differences due to fitting and statistics.

Effects of momentum smearing on pole factor



- Detector smearing has a drastic impact when the EbE method is used.
 - If you calculate the pole factor on an EbE basis with *smeared* spectator kinematic values, you now remove the pole factor for the wrong nuclear configuration!

Kinematic Distributions and Smearing



- Event sub-sample passed through full GEANT4 simulations.
 - Smearing parametrizations extracted for (p_x, p_y, p_z, E) .
- Larger overall smearing observed for neutrons, consistent with previous study.
- Anomalous proton smearing at high p_T and $p > 120$ GeV/c and $p < 100$ GeV/c due to linear transfer matrix assumption.
 - Will be fixed in the future for TDR studies.