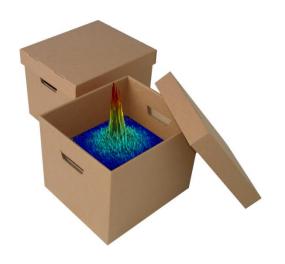






Turbulence in a Quantum Gas



Nir Navon



Experimentalists: Christoph Eigen, Jinyi Zhang, Alex Gaunt, Raphael Lopes, Sam Garratt,

Patrik Turzak, Rob Smith, Zoran Hadzibabic

Theorists: Kazuya Fujimoto, Makoto Tsubota







Outline

What is turbulence, and why is it interesting?

The optical box trap: a new tool for ultracold atoms

Low-energy excitations of a box-trapped BEC

A turbulent steady-state

A synthetic dissipation scale and turbulent-cascade fluxes

Hierarchy in complexity

Equilibrium



Overarching principle: maximum entropy
Many methods to efficiently calculate static quantities
Phase transitions, etc.

Near equilibrium



Linear response theory Universality of low-energy excitations Transport coefficients, etc.

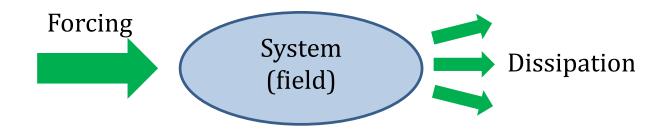
Far from equilibrium



Nothing like equilibrium states Existence of an overarching principle? Even classical problem intractable

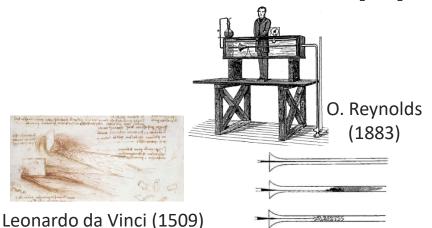
What is Turbulence?

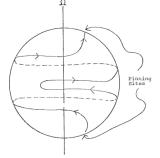
Quintessential phenomenon of out-of-equilibrium physics



Far-from-equilibrium states that (usually) involves

- drive + dissipation
- steady in a statistical sense
- local restoration of symmetries (isotropy/homogeneity)
- many interacting degrees of freedom
- chaotic properties





SF turbulence neutron stars / glitches? (1969)

And ongoing...

The Onset of Turbulence in Pipe Flow

Kerstin Avila, 1+ David Moxey, 2 Alberto de Lozar, 1 Marc Avila, 1 Dwight Barkley, 2,3 Björn Hof1+

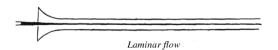
Shear flows undergo a sudden transition from laminar to turbulent motion as the velocity increases, and the onset of turbulence radically changes transport efficiency and mixing properties. Even for the well-studied case of pipe flow, it has not been possible to determine at what Reynolds number the motion will be either persistently turbulent or ultimately laminar. We show that in pipes, turbulence that is transient at low Reynolds numbers becomes sustained at a distinct critical point. Through extensive experiments and computer simulations, we were able to identify and characterise the processes ultimately responsible for sustaining turbulence. In

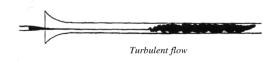
Science (2011)

Hydrodynamic vs. Wave Turbulence

Hydrodynamic (HD) equations v(r,t)

- Static solutions (e.g. Poiseuille)
- · Static solutions can be unstable
 - → time invariance is broken
 - → chaotic, turbulent state can emerge





Eq of motion for wave occupations e.g. $\mathbf{n}(\mathbf{k},t)$

Nonlinearity: waves eigenmodes
Analytic insights for weak interactions
(weak wave turbulence)





Turbulence is complicated because:

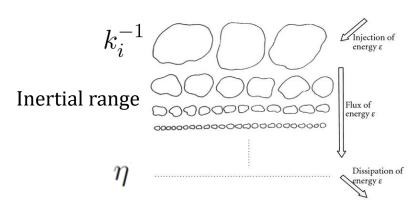
- nonlinearity (e.g. inertial term in NS)
- no perturbation theory, far beyond linear response
- many length and time scales simultaneously involved

Even the fully classical problem is numerically untractable for many realistic flows

How can one say anything general?

Turbulent cascades

A central phenomenology of turbulence: cascades of *something*, *gradually* transported across different length scales

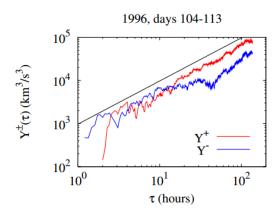


The most famous example: "Kolmogorov 5/3" (hydro)

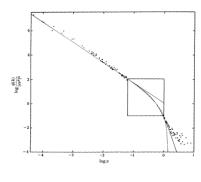
kinetic energy cascade of incompressible flow

$$E_k = C_K \epsilon^{2/3} k^{-5/3}$$

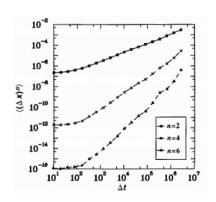
Turbulent cascades are everywhere



Interplanetary PlasmaSorriso-Valvo et al, PRL 2007



Water in a tidal channel Grant et al, J. Fluid Mech. 1962



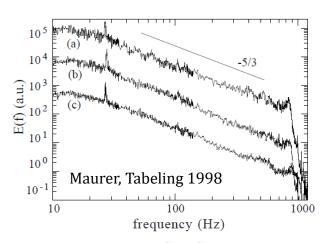
Foreign exchange markets Ghashgale et al, Nature 1996

Turbulence in quantum fluids

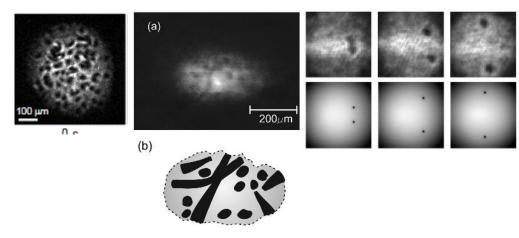
Interesting twists:

- no (intrinsic) dissipation scale (at low T)
- strongly restricted flow (irrotational)

For a long time, only SF helium, but no 1st-principle description (neutron stars are conjecured to be in a SF turbulent state, possible



Groups in Maryland, Lancaster, Saclay, Lyon, Grenoble, ...

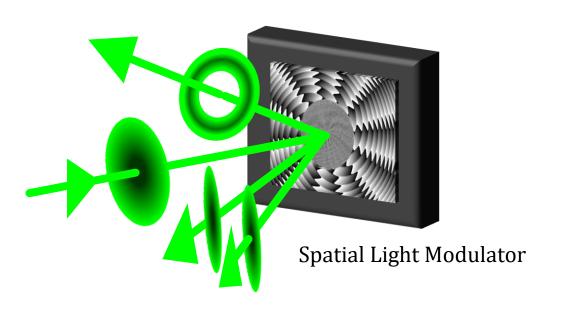


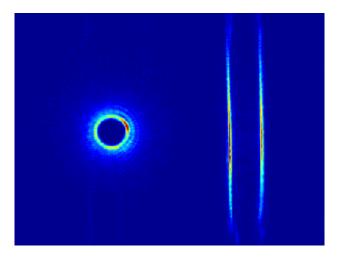
V. Bagnato (Brazil), B. P. Anderson (US), Y. Shin (South Korea), P. Engels (Washington), K. Helmerson and T. Neely (Australia)

Nonlinear Schrodinger Equation $i\hbar \frac{\partial \psi}{\partial t} = \left(-\frac{\hbar^2}{2m}\Delta + g|\psi|^2\right)\psi$

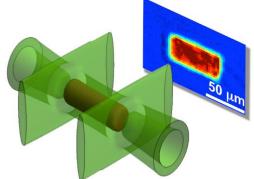
Quantum-gas fluid than can be simulated on all relevant scales, and we have new knobs!

The Box Potential





With (bosonic) ⁸⁷Rb atoms in:



A. Gaunt et al., PRL 110, 200406 (2013)

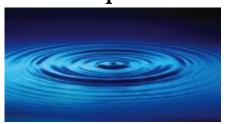
Hierarchy in complexity

Equilibrium



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Many methods to efficiently calculate static quantities
Phase transitions, etc.

Near equilibrium



Linear response theory Universality of low-energy excitations Transport coefficients, etc.

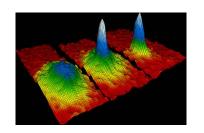
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Harmonic vs. Uniform BEC

BEC occurs both in real and momentum space



Time of Flight (JILA)

In-situ (MIT)



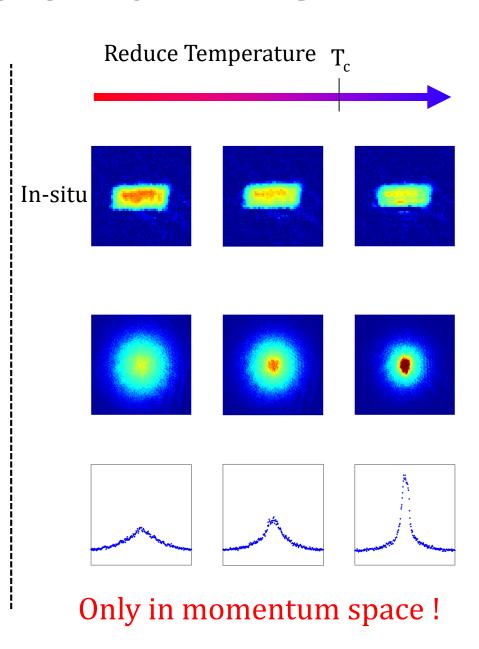






C. Wieman, E. Cornell

W. Ketterle

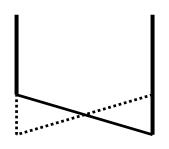


Shaken, not stirred

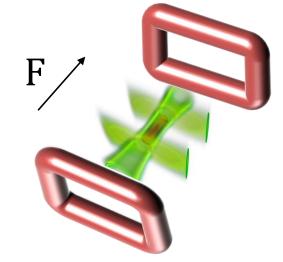
Pump energy at the largest scale, see what happens...



Easily realisable perturbation potential: time-dependent gradient



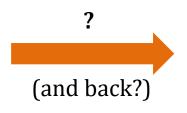
$$V(\mathbf{r},t) = \alpha \sin(\Omega t)x$$



4 Stages:

- Linear response
- Nonlinear response
- Turbulence
- Relaxation (ongoing)







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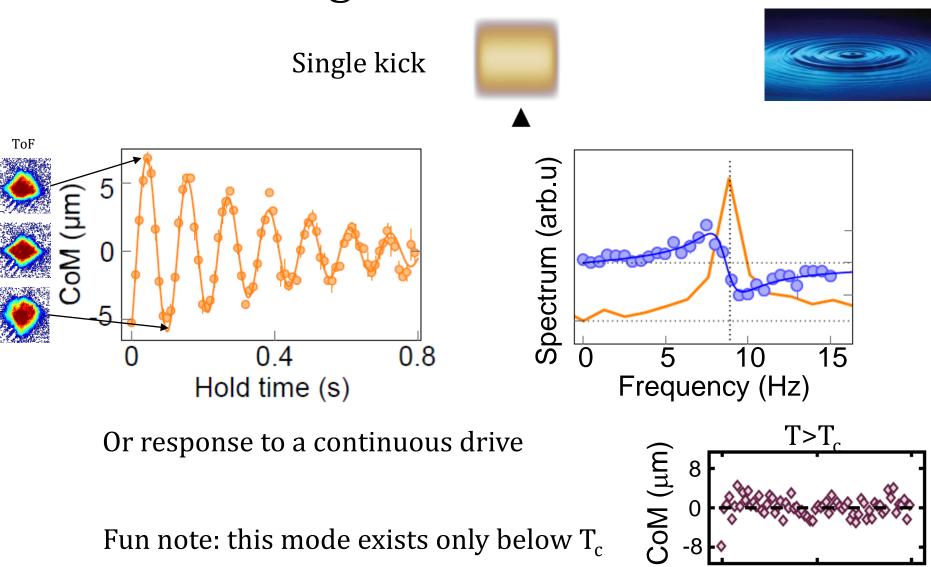
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Linear regime: collective modes

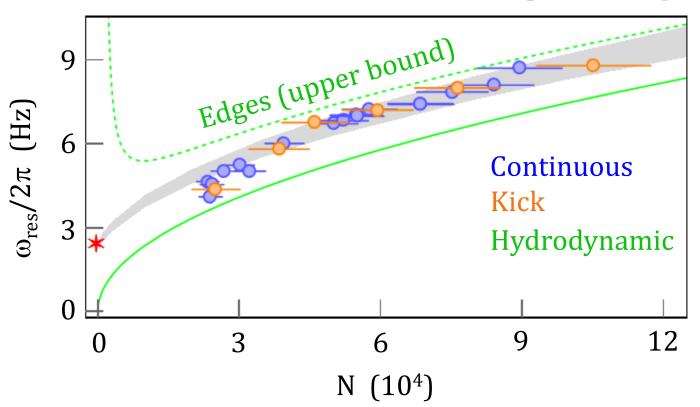


0.6

Hold time (s)

Linear regime: collective modes

Bogoliubov equations



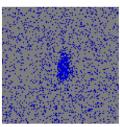
Weak perturbation: Bogoliubov equations from linearised GPE

$$\psi(\mathbf{r},t) = e^{-i\mu t/\hbar} (\psi_0(\mathbf{r}) + u(\mathbf{r})e^{-i\omega t} - v^*(\mathbf{r})e^{i\omega t})$$

Non-trivial modes in the box (with fixed boundary conditions)

Linear regime: collective modes

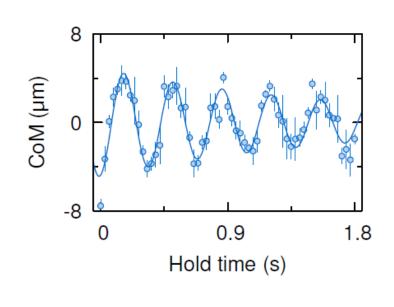
The single-particle limit

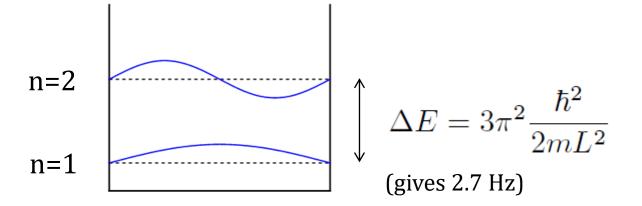


 $N \approx 10^3$ (with 140ms ToF)

We measure

$$\frac{\hbar\omega_{\rm res}}{2\pi}\approx 2.6~{\rm Hz}$$



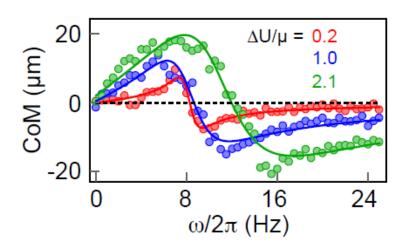


Good thing we have a BEC... $E_{n=1} \approx k_B \times 40 \text{ pK}$

$$E_{n=1} \approx k_{\rm B} \times 40 \text{ pK}$$

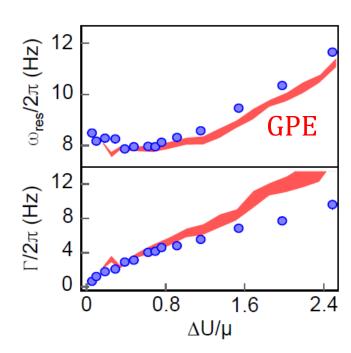
but $\mu \ll E_{n=1}(0.3 \text{ Hz} \ll 0.8 \text{ Hz})$

Nonlinear Response



Beyond Bogoliubov necessary

No "leader-order" frequency shift Significant "leading-order" broadening Energy has to go somewhere... everywhere?



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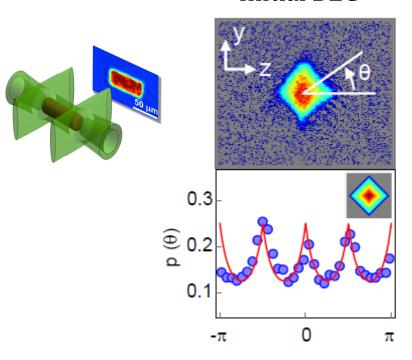
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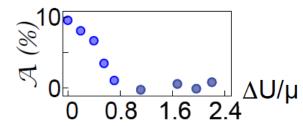
Nothing like equilibrium states Existence of an overarching principle? Even classical problem intractable

Emergence of Turbulence (in the experiment)

Initial BEC



Anistropy parameter
$$\mathcal{A} = \frac{1}{2} \int d\theta \, \left| p(\theta) - \frac{1}{2\pi} \right|$$

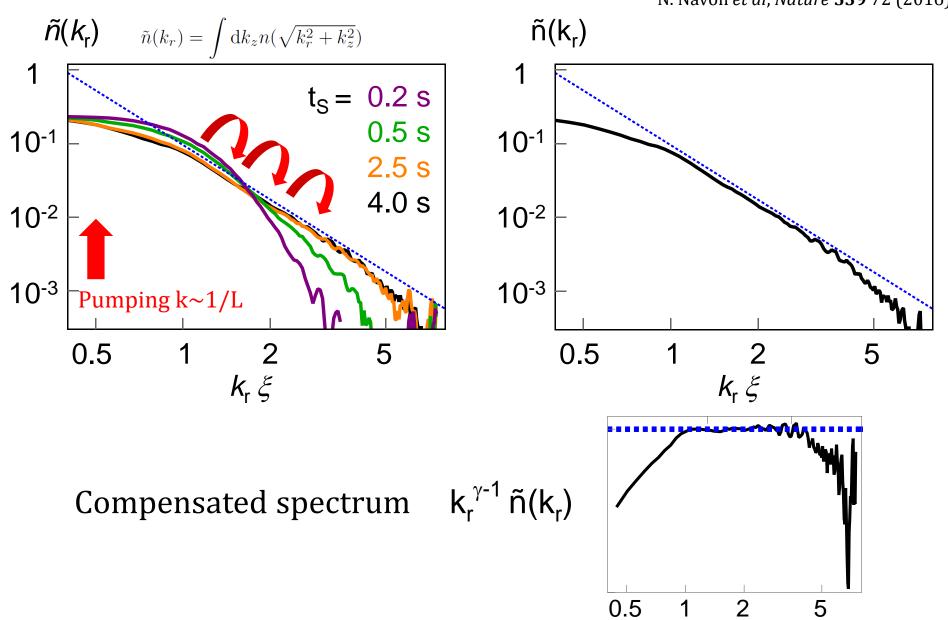


Isotropic state yet highly degenerate (not thermal!)

Is ToF a sane technique to use for a highly degenerate gas?

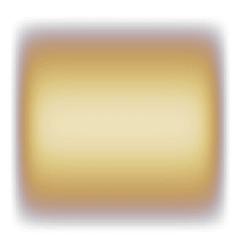
Turbulent Cascade

N. Navon et al, Nature **539** 72 (2016)



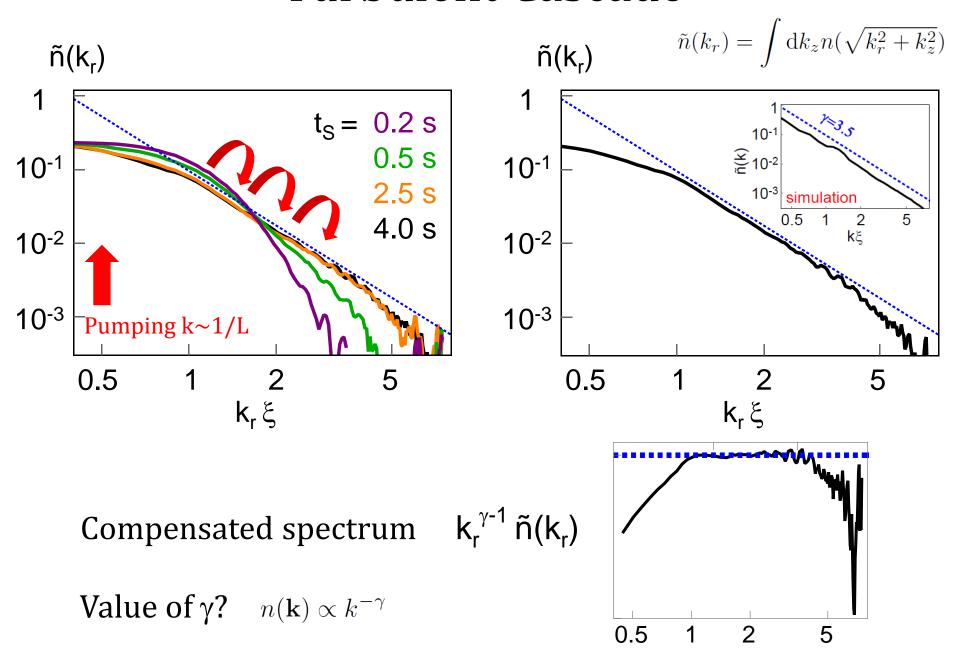
Emergence of Turbulence

(numerical simulations)



In situ density GPE in 256^3 grid with $10~\mu s$ steps for up to 4s (GPU) 30~min computation time per 1~sec experimental time

Turbulent Cascade



Summary/Outlook

- Even one of the simplest quantum system imaginable exhibits spectacularly rich physics

- New system and new (AMO-specific) tools to study turbulence

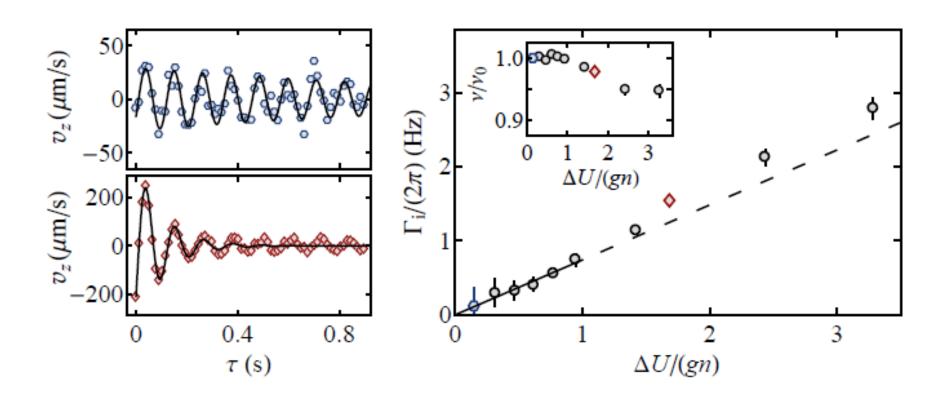
- GPE simulations confirm experiment (a surprise in and of itself)

- Many questions: (i) Cascade exponent not understood yet (related: Quantify the deviation from weak wave turbulence theory)

(ii) Establish the validity of classical field methods for this system

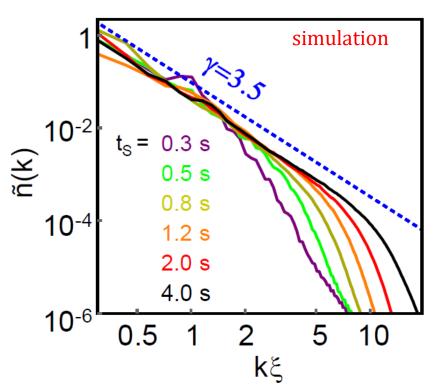
(iii) Interplay play of compressible and incompressible-fluid turbulence

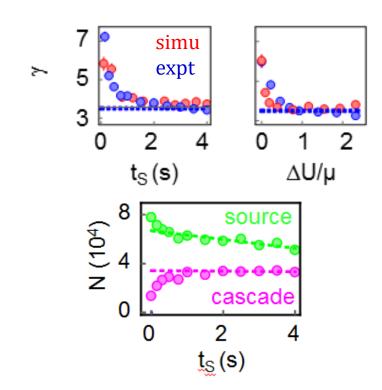
Nonlinear behavior in real time



Dynamics of the Cascade

$$n(\mathbf{k}) \propto k^{-\gamma}$$
 $\tilde{n}(k_r) = \int \mathrm{d}k_z n(\sqrt{k_r^2 + k_z^2})$





What type of cascade is it?

Two broad classes: - vortex turbulence

(quantized circulation / incompressible flow)

wave turbulence
 (density fluctuations / compressible flow)

Energy spectra

(with K. Fujimoto and M. Tsubota)

$$E[\psi] = \int d\mathbf{r} \ \psi^* \left(-\nabla^2 + \frac{g}{2} |\psi|^2 \right) \psi$$

Define flow field

$$\mathbf{w}(\mathbf{r}) = \sqrt{n}\nabla\phi$$

$$E_q = \int d\mathbf{r} (\nabla \sqrt{n})^2$$

$$E_{\text{kin}} = \int d\mathbf{r} n (\nabla \phi)^2$$

$$E_{\text{int}} = \frac{g}{2} \int d\mathbf{r} |\psi|^4$$

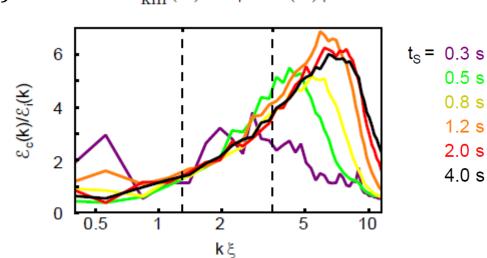
Helmholtz decomposition $\mathbf{w}(\mathbf{r}) = \mathbf{w}^c(\mathbf{r}) + \mathbf{w}^i(\mathbf{r})$

$$\mathbf{w}(\mathbf{r}) = \mathbf{w}^c(\mathbf{r}) + \mathbf{w}^i(\mathbf{r})$$

$$\nabla \times \mathbf{w}^c = 0 \text{ and } \nabla \cdot \mathbf{w}^i = 0$$

Then (Fourier)

$$E_{\rm kin}^{i,c}(\mathbf{k}) = |\tilde{\mathbf{w}}^{i,c}(\mathbf{k})|^2$$



Compressible-flow energy dominates over incompressible-flow one