

## Radiative corrections and form factors (not only) in Dalitz decays of lightest mesons

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## Hadronic parameters from experiment

- QED corrections left out entirely
- some of the relevant terms neglected
- approximative results (leading logs, soft-photon approximation)

## Artificial discrepancies between theory and experiment

↪ measured observables or related hadronic parameters may include unsubtracted QED part

## QED radiative corrections in the low-energy QCD sector

↪  $\pi^0, \eta^{(\prime)}, \Sigma^0, K^+, \dots$

↪ direct application to experiment

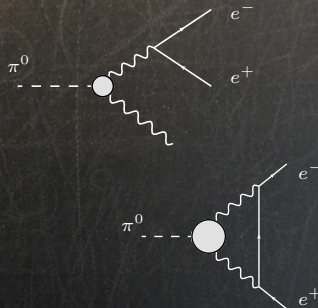
## Radiative corrections for $\pi^0$ decays

# Rare decay of $\pi^0$

Introduction

Decay modes of the neutral pion:

Process	Branching ratio
$\pi^0 \rightarrow \gamma\gamma$	$(98.823 \pm 0.034) \%$
$\pi^0 \rightarrow e^+e^-\gamma$	$(1.174 \pm 0.035) \%$
$\pi^0 \rightarrow e^+e^+e^-e^-$	$(3.34 \pm 0.16) \times 10^{-5}$
$\pi^0 \rightarrow e^+e^-$	$(6.46 \pm 0.33) \times 10^{-8}$



Rare decay  $\pi^0 \rightarrow e^+e^-$

- interesting way to study low-energy (long-distance) dynamics in the SM
- systematic theoretical treatment dates back to [Drell, NC \(1959\)](#)
- suppressed compared to the decay  $\pi^0 \rightarrow \gamma\gamma$  by a factor of  $2(\alpha m_e/M_\pi)^2$ 
  - ↪ one-loop structure + helicity suppression
  - ↪ may be sensitive to possible effects of new physics

# Rare decay of $\pi^0$

KTeV measurement

KTeV-E799-II experiment at Fermilab (*Abouzaid et al., PRD 75 (2007)*)

↪ precise measurements of branching ratio  $\pi^0 \rightarrow e^+e^-$  (794 candidates)

$$\frac{\Gamma(\pi^0 \rightarrow e^+e^-(\gamma), x > 0.95)}{\Gamma(\pi^0 \rightarrow e^+e^-\gamma, x > 0.232)} = (1.685 \pm 0.064 \pm 0.027) \times 10^{-4}$$

↪ final result for lowest order (no final state radiation)

$$B_{\text{KTeV}}^{\text{no-rad}}(\pi^0 \rightarrow e^+e^-) = (7.48 \pm 0.29 \pm 0.25) \times 10^{-8}$$

Comparison with SM prediction (*Dorokhov and Ivanov, PRD 75 (2007)*)

$$B_{\text{SM}}^{\text{no-rad}}(\pi^0 \rightarrow e^+e^-) = (6.23 \pm 0.09) \times 10^{-8}$$

↪ interpreted as 3.3  $\sigma$  discrepancy between theory and experiment

# Radiative corrections for $\pi^0 \rightarrow e^+e^-$

Final results



Vaško and Novotný, JHEP 1110 (2011)

Size of the radiative corrections (newly calculated)

$$\delta^{\text{NLO}}(0.95) \equiv \delta^{\text{virt.}} + \delta^{\text{BS}}(0.95) = (-5.5 \pm 0.2) \%$$

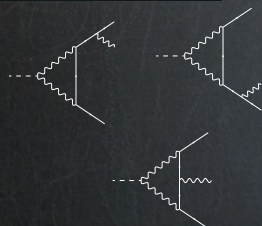
- can be thought as model-independent
- differs significantly from previous approximate calculations

*Bergström, Z.Ph.C 20 (1983):  $\delta(0.95) = -13.8 \%$*   
*Dorokhov et al., EPJC 55 (2008):  $\delta(0.95) = -13.3 \%$*

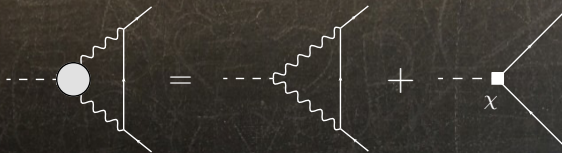
- original KTeV vs. SM discrepancy reduced to the  $2\sigma$  level
- contact interaction coupling finite part set to

$$\chi_{\text{LMD}}^{(r)}(M_\rho) = 2.2 \pm 0.9$$

TH, Kampf and Novotný,  
EPJC 74 (2014)



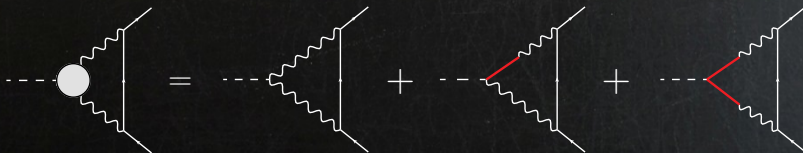
# Resonances



Chiral Perturbation Theory ( $\chi$ PT)



Resonance Chiral Theory ( $R\chi$ T)



# THS model for $PVV$ correlator

## 1) Ansatz for pseudoscalar-vector-vector ( $PVV$ ) correlator

- Two-hadron saturation (THS) - 2 meson multiplets per channel

$$\Pi^{\text{THS}}(r^2; p^2, q^2) \sim \frac{1}{r^2(r^2 - M_P^2)} \frac{P(r^2; p^2, q^2)}{(p^2 - M_{V_1}^2)(p^2 - M_{V_2}^2)(q^2 - M_{V_1}^2)(q^2 - M_{V_2}^2)}$$

- numerator: general polynomial symmetrical in  $p^2$  and  $q^2$ 
  - ↪ simple requirement: correlator must drop at large momenta
  - ↪ 22 free parameters

$$P(r^2; p^2, q^2) = c_0 p^2 q^2 + c_1 [(p^2)^3 q^2 + (q^2)^3 p^2] + c_2 (r^2)^2 p^2 q^2 + \dots$$

## 2) High- and low-energy limits to constrain the parameters $c_i$

- operator-product expansion (OPE)
- Brodsky–Lepage (BL) scaling limit
- chiral anomaly

*TH and S. Leupold, EPJC 75 (2015)*



# THS and $\mathcal{F}_{\pi^0\gamma^*\gamma^*}$ form factor

Form factor is in general related to  $PVV$  correlator as

$$\mathcal{F}_{\pi^0\gamma^*\gamma^*}(p^2, q^2) \sim \lim_{r^2 \rightarrow 0} r^2 \Pi(r^2; p^2, q^2)$$

↪ in our case complicated, but with only **one** free parameter

$$\mathcal{F}_{\pi^0\gamma^*\gamma^*}^{\text{THS}}(p^2, q^2) = -\frac{N_c}{12\pi^2 F} \left[ \frac{M_{V_1}^4 M_{V_2}^4}{(p^2 - M_{V_1}^2)(p^2 - M_{V_2}^2)(q^2 - M_{V_1}^2)(q^2 - M_{V_2}^2)} \right] \\ \times \left\{ 1 + \frac{\kappa}{2N_c} \frac{p^2 q^2}{(4\pi F)^4} - \frac{4\pi^2 F^2 (p^2 + q^2)}{N_c M_{V_1}^2 M_{V_2}^2} \left[ 6 + \frac{p^2 q^2}{M_{V_1}^2 M_{V_2}^2} \right] \right\}$$

$\kappa$  determined from fit to  $\omega$ - $\pi$  transition form factor measurements

↪ **NA60, PLB 677 (2009)**

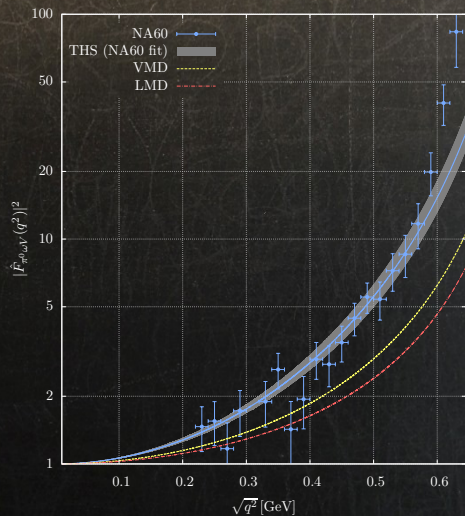
$$\kappa = 21 \pm 3$$

$M_{V_1} \sim \rho, \omega$  vector-meson mass

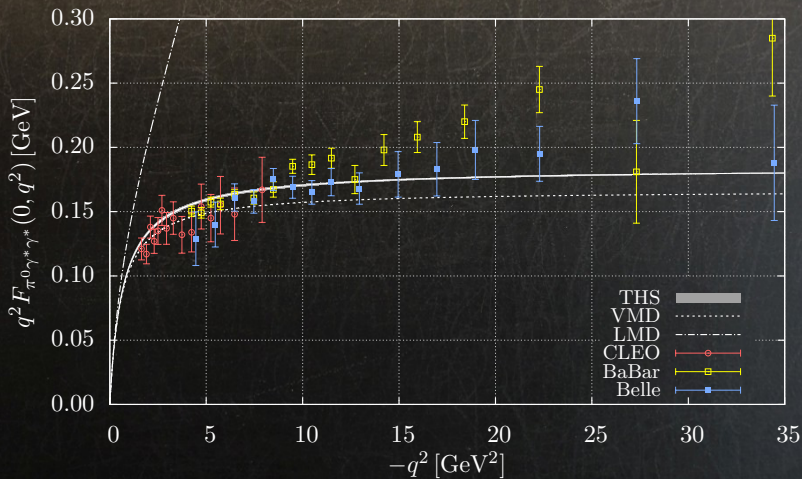
$M_{V_2} \sim$  between physical masses of first and second vector-meson excitations

$$M_{V_2} \in [1400, 1740] \text{ MeV}$$

# Fit of $\kappa$ to NA60 data



# Form-factor data



## Theoretical prediction within THS model

$$B^{\text{THS}}(\pi^0 \rightarrow e^+e^-\gamma), x_D > 0.95) = (5.76 \pm 0.7 \pm 2.0) \times 10^{-8}$$

$$\Updownarrow$$

$$B_{\text{THS}}^{\text{no-rad}}(\pi^0 \rightarrow e^+e^-) = 6.11(7)(20) \times 10^{-8}$$

- experimental value:  $B^{\text{KTeV}} = 6.44(33) \times 10^{-8} \Leftrightarrow B_{\text{KTeV}}^{\text{no-rad}} = 6.82(35) \times 10^{-8}$   
 $\hookrightarrow$  disagreement at the level of only **1.8  $\sigma$**
- matching on LO  $\chi$ PT gives  $\chi_{\text{THS}}^{(r)}(M_\rho) = 2.2 \pm 0.7$
- if KTeV result confirmed (e.g. by NA62)  $\rightarrow$  two scenarios are conceivable:
  - a) some aspects of the THS approach not well-suited for  $\pi^0 \rightarrow e^+e^-$
  - b) beyond-Standard Model physics influences the rare pion decay significantly
- under the present circumstances the current discrepancy is **inconclusive**

Extract  $\chi^{(r)}(M_\rho)$  from various models and measurements

VMD	LMD	THS	CLEO-based	KTeV (old)	KTeV (upd.)
2.87	2.29( $\pm 0.9$ )	2.2 $\pm$ 0.7	2.6 $\pm$ 0.3	6.0 $\pm$ 1.0	4.5 $\pm$ 1.0

# $\pi^0 \rightarrow e^+e^-$ branching ratio

Standard Model prediction

$$\frac{B(\pi^0 \rightarrow e^+e^-)}{B(\pi^0 \rightarrow \gamma\gamma)} = 2\beta \left(\frac{\alpha m}{\pi M}\right)^2 \left\{ \left[ -\frac{5}{2} + \chi^{(r)}(\mu) + \frac{3}{2} \log \frac{m^2}{\mu^2} + \frac{1}{2\beta} \left( \text{Li}_2 z - \text{Li}_2 \frac{1}{z} \right) \right]^2 + \left[ \frac{\pi}{2\beta} \log(-z) \right]^2 \right\}$$
$$z = -\frac{1-\beta}{1+\beta}, \quad \beta = \sqrt{1 - \frac{4m^2}{M^2}}$$

$$B(\pi^0 \rightarrow e^+e^-) \approx (6.21 + 0.15\tilde{\chi}) \times 10^{-8}, \quad \tilde{\chi} \equiv 2 \left[ \chi^{(r)}(770 \text{ MeV}) - \frac{5}{2} \right] \in (-1, 1)$$

$$B_{\text{KTeV}}^{\text{no-rad}} = 6.82(35) \times 10^{-8}$$

Theoretical predictions and models suggest  $\chi^{(r)}(770 \text{ MeV}) \sim 2-3$

*Knecht et al.*, PRL 83 (1999)

*Dorokhov and Ivanov*, PRD 75 (2007)

*TH and Leupold*, EPJC 75 (2015)

*Hoferichter et al.*, PRL 128 (2022)

# Dalitz decay of the neutral pion

## Introduction

Quantity **really** measured by KTeV

$$\frac{\Gamma(\pi^0 \rightarrow e^+e^-\gamma), x > 0.95}{\Gamma(\pi^0 \rightarrow e^+e^-\gamma), x > 0.2319} \Bigg|_{\text{KTeV}} = (1.685 \pm 0.064 \pm 0.027) \times 10^{-4}$$

↪ Dalitz decay comes into play

- **second** most important decay channel of the neutral pion  
↪ branching ratio  $(1.174 \pm 0.035) \%$
- first studied by **Richard H. Dalitz, PPSA 64 (1951)**
- experimental data of this process provide the information about **singly-virtual** pion transition form factor  $\mathcal{F}_{\pi^0\gamma^*\gamma^*}(0, q^2)$   
↪ in particular about its **slope** parameter  $a_\pi$

$$\frac{\mathcal{F}_{\pi^0\gamma^*\gamma^*}(0, M^2x)}{\mathcal{F}_{\pi^0\gamma^*\gamma^*}(0, 0)} \simeq 1 + a_\pi x$$

$$x = \frac{(p_{e^+} + p_{e^-})^2}{M_{\pi^0}^2}, \quad y = -\frac{2}{M_{\pi^0}^2} \frac{p_\gamma \cdot (p_{e^-} - p_{e^+})}{1 - x}$$

# Radiative corrections for $\pi^0 \rightarrow e^+e^-\gamma$

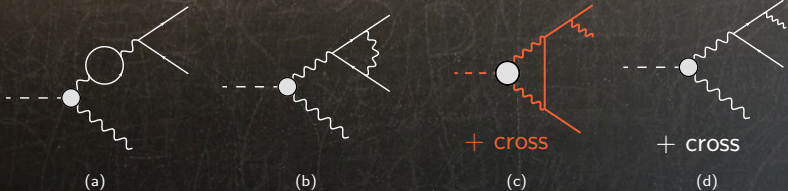
## Introduction

- radiative corrections to the **total** decay rate of the Dalitz decay  
↪ first addressed (numerically) by *Joseph, NC 16 (1960)*
- pioneering study of corrections to the **differential** decay rate  
↪ *Lautrup and Smith, PRD 3 (1971)*  
↪ soft-photon approximation
- extended by *Mikaelian and Smith, PRD 5 (1972)*  
↪ hard-photon corrections  
↪ **whole** range of bremsstrahlung photon energy  
↪ table of values



# Radiative corrections for $\pi^0 \rightarrow e^+e^-\gamma$

Novel approach



- new calculations motivated by needs of NA48/NA62 experiments at CERN  
↪ measure the slope of  $\mathcal{F}_{\pi^0\gamma^*\gamma^*}(0, q^2)$ : [Lazzeroni et al., PLB 768 \(2017\)](#)

$$a_{\pi}^{\text{NA62}} = 3.68(57) \%$$

- unlike before **no approximation** was used + **1 $\gamma$ IR correction**
- C++ code returns the correction for any given  $x$  and  $y$   
↪ propagated into **MC generator** of NA62 experiment
- [TH, Kampf and Novotný, PRD 92 \(2015\)](#)



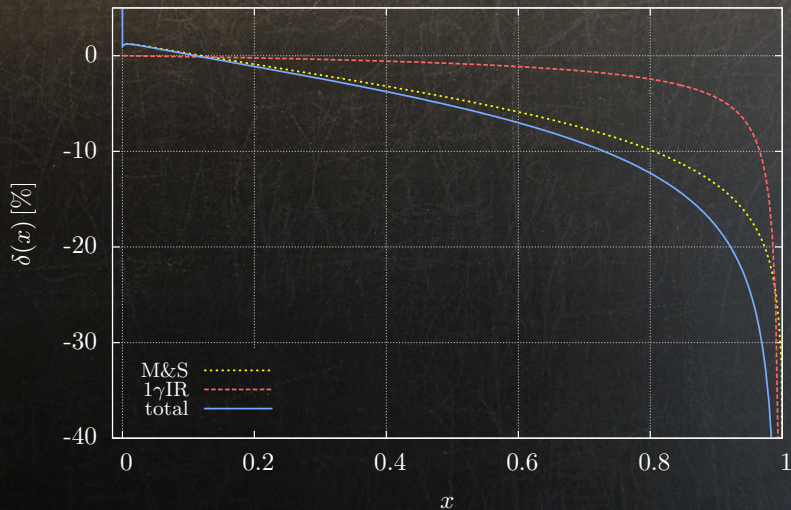
# Radiative corrections for $\pi^0 \rightarrow e^+e^-\gamma$

The overall NLO correction  $\delta(x, y)$  given in percent (Dalitz-plot corrections)

$x \backslash y$	0.00	0.10	0.20	0.30	0.40	0.50	0.60	0.70	0.80	0.90	0.99
0.01	2.761	2.714	2.599	2.449	2.273	2.061	1.786	1.402	0.803	-0.357	-5.657
0.02	2.756	2.720	2.622	2.480	2.300	2.073	1.774	1.355	0.703	-0.546	-5.859
0.03	2.669	2.639	2.552	2.419	2.242	2.012	1.704	1.267	0.586	-0.716	-6.125
0.04	2.558	2.531	2.452	2.327	2.155	1.925	1.611	1.164	0.464	-0.874	-6.372
0.05	2.437	2.412	2.340	2.221	2.053	1.824	1.509	1.054	0.341	-1.025	-6.601
0.06	2.311	2.288	2.221	2.108	1.944	1.717	1.400	0.940	0.216	-1.172	-6.815
0.07	2.184	2.163	2.099	1.990	1.830	1.605	1.288	0.824	0.092	-1.315	-7.017
0.08	2.056	2.036	1.975	1.870	1.714	1.491	1.173	0.707	-0.033	-1.455	-7.211
0.09	1.928	1.909	1.851	1.749	1.596	1.374	1.057	0.588	-0.157	-1.593	-7.397
0.10	1.801	1.783	1.726	1.628	1.477	1.257	0.940	0.469	-0.281	-1.729	-7.578
0.15	1.170	1.154	1.105	1.016	0.874	0.661	0.345	-0.131	-0.900	-2.394	-8.424
0.20	0.546	0.532	0.486	0.402	0.266	0.057	-0.258	-0.738	-1.520	-3.048	-9.219
0.25	-0.079	-0.092	-0.135	-0.217	-0.350	-0.556	-0.871	-1.355	-2.148	-3.704	-9.995
0.30	-0.713	-0.726	-0.768	-0.847	-0.978	-1.184	-1.499	-1.988	-2.790	-4.372	-10.77
0.35	-1.366	-1.378	-1.419	-1.497	-1.627	-1.833	-2.149	-2.641	-3.454	-5.058	-11.56
0.40	-2.044	-2.056	-2.097	-2.174	-2.304	-2.509	-2.827	-3.324	-4.146	-5.773	-12.37
0.45	-2.759	-2.771	-2.811	-2.887	-3.017	-3.222	-3.543	-4.044	-4.875	-6.525	-13.22
0.50	-3.521	-3.533	-3.572	-3.648	-3.777	-3.983	-4.306	-4.811	-5.653	-7.324	-14.12
0.55	-4.344	-4.356	-4.395	-4.470	-4.599	-4.806	-5.130	-5.640	-6.492	-8.186	-15.08
0.60	-5.249	-5.261	-5.299	-5.373	-5.501	-5.708	-6.034	-6.549	-7.410	-9.128	-16.12
0.65	-6.262	-6.273	-6.310	-6.383	-6.510	-6.717	-7.044	-7.563	-8.435	-10.18	-17.28
0.70	-7.425	-7.435	-7.470	-7.541	-7.666	-7.871	-8.198	-8.721	-9.603	-11.37	-18.60
0.75	-8.802	-8.811	-8.844	-8.910	-9.031	-9.232	-9.558	-10.08	-10.98	-12.77	-20.14
0.80	-10.51	-10.52	-10.54	-10.60	-10.72	-10.91	-11.23	-11.76	-12.66	-14.49	-22.02
0.85	-12.78	-12.78	-12.80	-12.85	-12.95	-13.13	-13.44	-13.96	-14.86	-16.72	-24.47
0.90	-16.21	-16.21	-16.21	-16.23	-16.29	-16.43	-16.71	-17.21	-18.11	-20.00	-28.00
0.95	-23.17	-23.14	-23.08	-23.01	-22.96	-22.98	-23.14	-23.53	-24.36	-26.26	-34.45
0.99	-54.29	-54.07	-53.44	-52.50	-51.35	-50.15	-49.03	-48.16	-47.76	-48.47	-55.83

# Radiative corrections for $\pi^0 \rightarrow e^+e^-\gamma$

Results



# Radiative corrections for $\pi^0 \rightarrow e^+e^-\gamma$

Determination of ratio  $R$

Precise and reliable determination of  $R \equiv \frac{\Gamma(\pi^0 \rightarrow e^+e^-\gamma)}{\Gamma(\pi^0 \rightarrow \gamma\gamma)}$

$\leftrightarrow$  for small slope and up to NLO radiative corrections

$$R \simeq \frac{\alpha}{\pi} \iint (1 + a_\pi x)^2 (1 + \delta(x, y)) \frac{(1-x)^3}{4x} \left[ 1 + y^2 + \frac{4m_e^2}{M_\pi^2 x} \right] dx dy$$

Conservative estimate for uncertainty ( $a_\pi$ , NNLO):  $R = 1.1978(5)(3) \%$

$\leftrightarrow$  chosen  $a_\pi^{\text{univ}} = 3.55(70) \%$ , covers

source	VMD	LMD	THS	dispers.	Padé aps.	NA62	A2	PDG
$a_\pi$ [%]	3.00	2.45	2.92(4)	3.15(9)	3.21(19)	3.68(57)	3.0(1.0)	3.35(31)
$b_\pi$ [ $10^{-3}$ ]	0.90	0.74	0.87(2)	1.14(4)	1.04(22)	×	×	×

Constraint:  $1 \simeq \mathcal{B}(\pi^0 \rightarrow \gamma\gamma) + \mathcal{B}(\pi^0 \rightarrow e^+e^-\gamma(\gamma)) + \mathcal{B}(\pi^0 \rightarrow e^+e^-e^+e^-)$

$$\mathcal{B}(\pi^0 \rightarrow \gamma\gamma) = 98.8131(6) \%, \quad \mathcal{B}(\pi^0 \rightarrow e^+e^-\gamma(\gamma)) = 1.1836(6) \%$$

TH, Goudzovski and Kampf, PRL 122 (2018)

$\leftrightarrow$  could be used in future exp. analysis of  $K^+ \rightarrow \pi^+e^+e^-$

PDG

$$R = 1.188(35) \%, \quad \mathcal{B}(\pi^0 \rightarrow \gamma\gamma) = 98.823(34) \%, \quad \mathcal{B}(\pi^0 \rightarrow e^+e^-\gamma) = 1.174(35) \%$$

# 'Recent' KTeV measurement

*Abouzaid et al.*, PRD 100 (2019)

↪ based on 1999 data and *E. Abouzaid*, Ph.D. thesis (2007)

$$\frac{\Gamma(\pi^0 \rightarrow e^+e^-\gamma)}{\Gamma(\pi^0 \rightarrow \gamma\gamma)} = (1.1559 \pm 0.0046 \pm 0.0106) \%$$

*TH, Goudzovski and Kampf*, PRL 122 (2018)

Conservative estimate for uncertainty ( $a_\pi$ , NNLO):  $R = 1.1978(5)(3) \%$

↪ chosen  $a_\pi^{\text{univ}} = 3.55(70) \%$

⇒ 3.6  $\sigma$  discrepancy between theory and experiment

PDG average:  $R = 1.188(35) \%$

↪ most recent (archived ALEPH data) *Beddall and Beddall*, EPJC 54 (2008)

Need for new measurements

↪  $R$ , improvement on  $a_\pi$  (maybe  $b_\pi$ ) welcome

## Radiative corrections for $\eta^{(\prime)}$ Dalitz decays

# Radiative corrections for $\eta^{(\prime)} \rightarrow \ell^+ \ell^- \gamma$ decays

What should be taken into account?

naïve rad. corrections for  $\eta \rightarrow e^+ e^- \gamma$ : *Mikaelian and Smith, PRD 5 2890 (1972)*

- numerical values correspond to simple change  $M_{\pi^0} \rightarrow M_\eta$   
 $\hookrightarrow \pi^0$  case: *Mikaelian and Smith, PRD 5 1763 (1972)*

$\eta^{(\prime)}$  case compared to  $\pi^0$

- larger rest mass  
 $\hookrightarrow M_\eta$  above muon-pair threshold:  $M_\eta > 2m_\mu$   
 $\hookrightarrow M_{\eta'}$  above lowest-lying resonances:  $M_{\eta'} > M_\rho, M_\omega$   
 $\hookrightarrow$  sensitive to the **widths** of resonances  
 $\hookrightarrow \omega$  narrow,  $\rho$  **broad** resonance in  $\pi\pi$  scattering
- strange-flavor content**  
 $\hookrightarrow$  quark-flavor basis  
*Feldmann et al., PLB 449 (1999), Escribano et al., JHEP 06 (2005)*

$$j^\ell \equiv \frac{i}{2} [\bar{u}\gamma_5 u + \bar{d}\gamma_5 d], \quad j^s \equiv \frac{i}{\sqrt{2}} [\bar{s}\gamma_5 s]$$

- $\eta$ - $\eta'$  **mixing**:  $\langle 0 | j^A | \eta^B \rangle = B_0 F_\pi f_A \delta^{AB}$ ,  $\langle \eta^A | \eta^B \rangle = \delta^{AB}$ ,  $A, B \in \{\ell, s\}$

$$|\eta\rangle = \cos \phi |\eta^\ell\rangle - \sin \phi |\eta^s\rangle$$

$$|\eta'\rangle = \sin \phi |\eta^\ell\rangle + \cos \phi |\eta^s\rangle$$

# Radiative corrections for $\eta^{(\prime)} \rightarrow l^+l^-\gamma$ decays

Novel approach

Full set of NLO QED radiative corrections:

*TH, Kampf, Leupold and Novotný, PRD 97 (2018)*

- compared to previous approach:
  - ↪ muon loops + **hadronic** VP
  - ↪ **1 $\gamma$ IR** at one-loop level
  - ↪ **form-factor** effects (also in BS)
  - ↪ higher orders in the final-state-lepton mass **not** neglected
- general framework: **three** additional processes
  - ↪ also muon decay modes

$\eta$  case: **most** of the ingredients in *TH, Kampf and Novotný, PRD 92 (2015)*

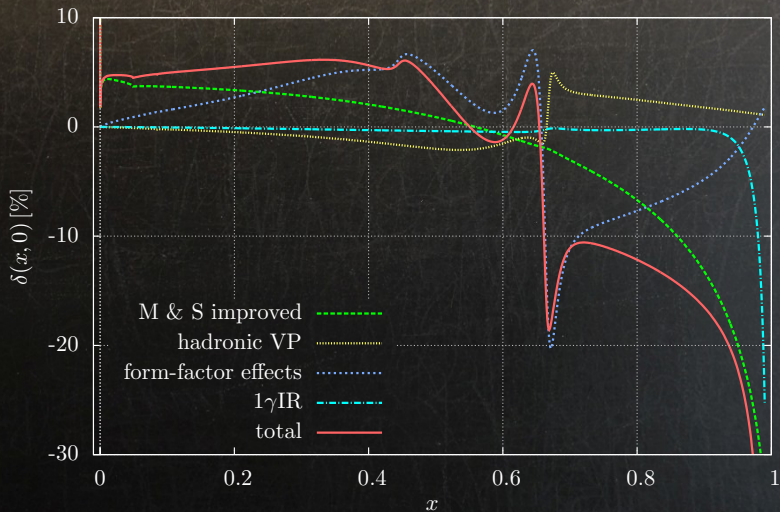
$\eta'$  case: real challenge

- ↪ resulting framework also **applicable** to the  $\pi^0$  case (numerically compatible)
  - ↪ overkill (correction to the correction of order 1%)



# Radiative corrections for $\eta' \rightarrow e^+e^-\gamma$ decays

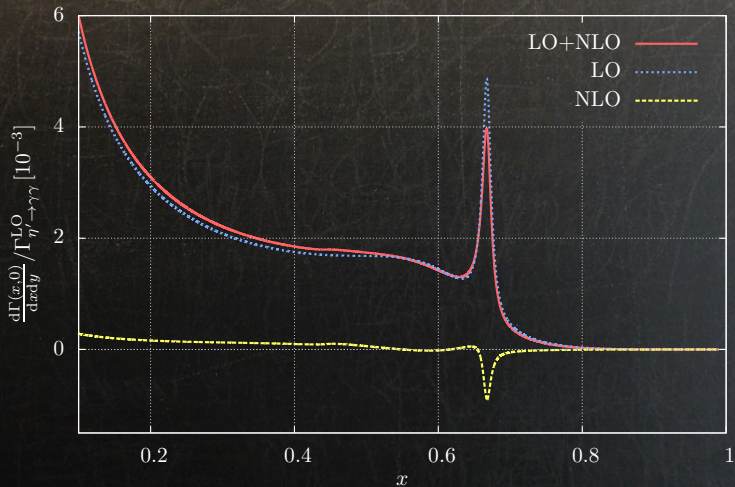
The overall NLO correction  $\delta(x, 0)$  in comparison to its constituents





# Radiative corrections for $\eta' \rightarrow e^+e^-\gamma$ decays

The two-fold differential decay width  $d\Gamma(x, 0)$  at NLO



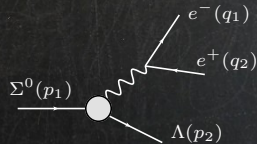
## Radiative corrections for $\Sigma^0$ Dalitz decay

# Dalitz decay of $\Sigma^0$

## Introduction

Dalitz decay  $\Sigma^0 \rightarrow \Lambda e^+ e^-$

- $e^+e^-$  invariant mass only up to  $M_\Sigma - M_\Lambda \simeq 77 \text{ MeV}$
- provides electric and magnetic transition form factors of  $\Sigma^0 \rightarrow \Lambda$  transition
- extracting transition radii challenging
  - high-precision measurement required
  - competing with QED radiative corrections



Predictions of electric and magnetic radii

- *Kubis and Meissner, EPJC 18 (2001)*
- *Granados, Leupold and Perotti, EPJA 53 (2017)*

# Radiative corrections for $\Sigma^0 \rightarrow \Lambda e^+ e^-$

Introduction

Radiative corrections to the differential decay width in **soft-photon** approximation

↔ *Sidhu and Smith, PRD 4 3344 (1971)*

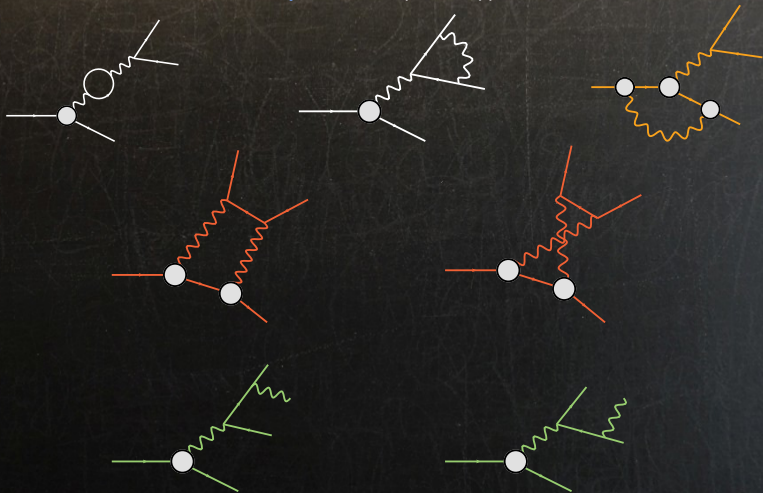


# Radiative corrections for $\Sigma^0 \rightarrow \Lambda e^+ e^-$

Diagrams

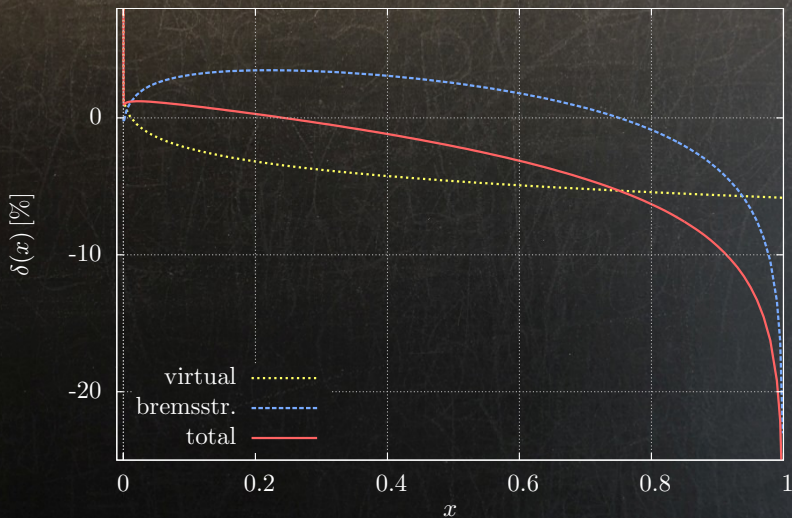
*TH and Leupold, EPJC 80 (2020)*

↪ inclusive radiative corrections **beyond** the soft-photon approximation



# Radiative corrections for $\Sigma^0 \rightarrow \Lambda e^+ e^-$

Results



# Radiative corrections for $\Sigma^0 \rightarrow \Lambda e^+ e^-$

Correction to the form-factor slope

Estimate of size of correction to the magnetic form-factor slope  $a$

- take half of the slope of the curve in the **low- $x$**  region
- **farther** from the threshold ( $\nu^2 \ll x_0 \ll 1$ ):

$$\Delta a \equiv a_{(+\text{QED})} - a \simeq \frac{1}{2} \left. \frac{d\delta(x)}{dx} \right|_{x=x_0}$$

$a_{(+\text{QED})}$ : measured value implicitly containing the QED radiative correction

$$\frac{1}{2} \left. \frac{d\delta(x)}{dx} \right|_{x=x_0} \approx -3.5\%$$

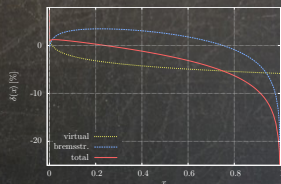
- **bigger** than the estimate on the slope  $a \equiv \frac{1}{6} \langle r_M^2 \rangle \Delta_M^2$  itself ( $a \approx 1.8(3)\%$ )

Using **no** radiative corrections in the experimental analysis

- one expects “measured” radius  $\langle r_M^2 \rangle_{(+\text{QED})}$  to be **negative**

$$\langle r_M^2 \rangle_{(+\text{QED})} = \langle r_M^2 \rangle + \frac{6}{\Delta_M^2} \Delta a, \quad \text{with } \frac{6}{\Delta_M^2} \Delta a \approx -35 \text{ GeV}^{-2}$$

( $\chi\text{PT}$ :  $\langle r_M^2 \rangle = 18.5(2.6) \text{ GeV}^{-2}$ ); in general for hadronic radii:  $\langle r^2 \rangle \leq (1 \text{ fm})^2 \approx 25 \text{ GeV}^{-2}$



# Radiative corrections for $\Sigma^0 \rightarrow \Lambda e^+ e^-$

Integrated decay width

Integrate over the Dalitz plot (values from *Kubis and Meissner, EPJC 18 (2001)*)

$$R \equiv \frac{\Gamma(\Sigma^0 \rightarrow \Lambda e^+ e^-)}{\Gamma(\Sigma^0 \rightarrow \Lambda \gamma)} = 5.541(2) \times 10^{-3}$$

$\hookrightarrow$  neglect **electric** form factor  $\rightarrow$  **expansion** in **magnetic** form-factor slope  $a \equiv \frac{1}{6} \langle r_M^2 \rangle \Delta_M^2$

$$R = R_0 + a R_1 + \mathcal{O}(a^2)$$

+ higher order corrections as additional uncertainty

$$R = [5.530(3) + 0.626(2)a] \times 10^{-3}$$

$\hookrightarrow$  **consistent** with the **NLO** result in *Sidhu and Smith, PRD 4 3344 (1971)*

$$R_{S\&S} = (5.532 + 0.627a) \times 10^{-3} [\approx 5.544 \times 10^{-3}]$$

From  $\mathcal{B}(\Sigma^0 \rightarrow \Lambda \gamma) + \mathcal{B}(\Sigma^0 \rightarrow \Lambda e^+ e^-) + \mathcal{B}(\Sigma^0 \rightarrow \Lambda \gamma \gamma) \simeq 1$

$$\mathcal{B}(\Sigma^0 \rightarrow \Lambda \gamma) = [99.4501(3) - 0.0619(2)a] \%$$

$$\mathcal{B}(\Sigma^0 \rightarrow \Lambda e^+ e^-) = [0.5499(3) + 0.0619(2)a] \%$$

$a = 0.02(2)$ :  $\mathcal{B}(\Sigma^0 \rightarrow \Lambda \gamma) = 99.449(2) \%$ ,  $\mathcal{B}(\Sigma^0 \rightarrow \Lambda e^+ e^-) = 0.551(2) \%$

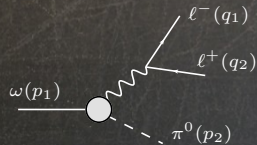


## Radiative corrections for $\omega \rightarrow \pi^0 \ell^+ \ell^-$ decays

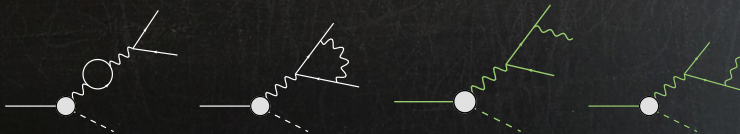
# Radiative corrections for $\omega \rightarrow \pi^0 \ell^+ \ell^-$ decays

Dalitz decay  $\omega \rightarrow \pi^0 \ell^+ \ell^-$

- $\pi\omega V$  correlator measured in the past
  - ↪ NA60: [Arnaldi et al., PLB 677 \(2009\)](#)
  - ↪ A2: [Adlarson et al., PRC 95 \(2017\)](#)
- radiative corrections not included (not available)



Inclusive NLO QED radiative corrections **beyond** the soft-photon approximation



$$|\mathcal{F}_{(+QED)}(q^2)|^2 = |\mathcal{F}(q^2)|^2 [1 + \delta(x)]$$

Bachelor project: [Jacob Lindahl \(2021\)](#)

## Radiative corrections for $K^+ \rightarrow \pi^+ \ell^+ \ell^-$ decays

# Radiative decays $K^+ \rightarrow \pi^+ \ell^+ \ell^-$

Form-factor parametrization

LO appears at  $\mathcal{O}(p^4)$  + unitarity loop correction from  $\pi\pi$  rescattering ↓

↪ universally used parametrization for the fit:  $V_+(x) = a_+ + b_+x + V_+^{\pi\pi}(x)$

↪ *Ecker et al., NPB 291 (1987), D'Ambrosio et al., JHEP 08 (1998)*

$$\frac{d\Gamma_+}{dx} = \frac{G_F^2 \alpha^2 M_K^5}{3(4\pi)^5} \lambda^{3/2}(x) \sqrt{1 - \frac{4r_\ell^2}{x}} \left(1 + \frac{2r_\ell^2}{x}\right) |V_+(x)|^2$$

LFU

↪  $a_+$ s and  $b_+$ s should be the same for both ( $e$  and  $\mu$ ) channels

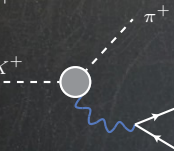
↪ discrepancy due to NP via SD effects

Moreover, the ratio deviates significantly from the VMD ansatz

$$\text{VMD: } \frac{b_+}{a_+} = \frac{M_K^2}{M_\rho^2} \approx 0.4, \quad \text{exp.: } \frac{b_+}{a_+} \approx 1.25$$

Measurement of quadratic term  $c_+x^2$  may further test the VMD hypothesis

$\ell$	$a_+$	$b_+$	exp.
$e$	-0.587(10)	-0.655(44)	E865
$e$	-0.578(16)	-0.779(66)	NA48/2
$\mu$	-0.575(39)	-0.813(145)	NA48/2
$\mu$	-0.575(13)	-0.722(43)	NA62(2022)



← *JHEP 11 (2022) 011*

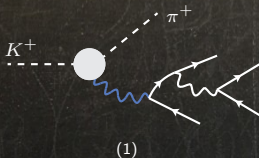
improve precision → radiative corrections, studied earlier: *Kubis and Schmidt, EPJC 70 (2010)*

$$K^+ \rightarrow \pi^+ 4e \text{ decay}$$

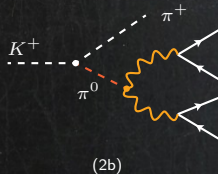
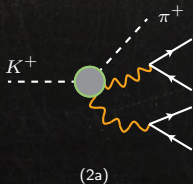
$$K^+ \rightarrow \pi^+ e^+ e^- e^+ e^-$$

Standard Model prediction: Topologies

### One-photon-exchange topology

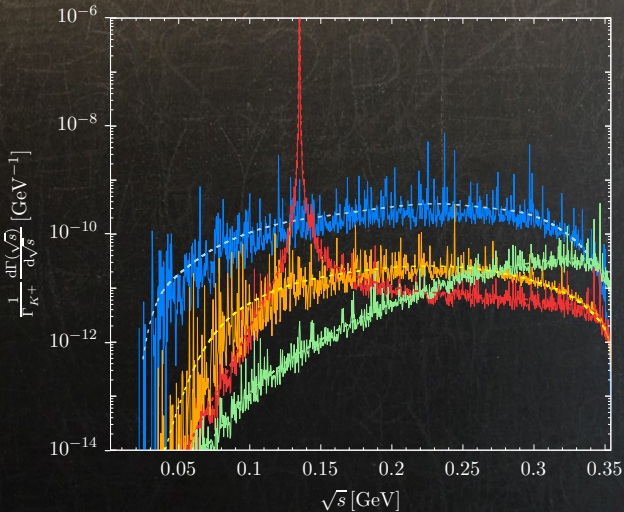


### Two-photon-exchange topology



TH, PRD 106 (2022)

$K^+ \rightarrow \pi^+ e^+ e^- e^+ e^-$   
Contributions to the branching ratio



[large MC samples generated by A. Shaikhiev, E. Goudzovski]

$$K^+ \rightarrow \pi^+ e^+ e^- e^+ e^-$$

Contributions to the branching ratio

Branching ratio calculated using Monte Carlo event generator technique:

$$B = \frac{1}{\Gamma_0} \frac{1}{4} \frac{1}{2M_K} \Phi_5 \frac{1}{N} \sum_{N \text{ events}} |\mathcal{M}|^2$$

	$B(\sqrt{s} < 120 \text{ MeV})$	$B(\sqrt{s} > 150 \text{ MeV})$	$B$
(1)	$5.60 \times 10^{-12}$	$5.44 \times 10^{-11}$	$6.70 \times 10^{-11}$
(2a)	$3.11 \times 10^{-13}$	$3.85 \times 10^{-12}$	$4.60 \times 10^{-12}$
(2b)	$1.40 \times 10^{-13}$	$1.97 \times 10^{-12}$	$7.0(3) \times 10^{-6}$
$\kappa$	$7.08 \times 10^{-15}$	$3.69 \times 10^{-12}$	$3.72 \times 10^{-12}$
$\Sigma$	$6.1(4) \times 10^{-12}$	$6.0(6) \times 10^{-11}$	$7.2(7) \times 10^{-11}$

$$B(K^+ \rightarrow \pi^+ 4e) \simeq B(K^+ \rightarrow \pi^+ \pi^0) B(\pi^0 \rightarrow 4e)$$

$$\Leftrightarrow B(K^+ \rightarrow \pi^+ \pi^0) = 20.67(8) \% \text{ and } B(\pi^0 \rightarrow 4e) = 3.38(16) \times 10^{-5}$$

TH, PRD 106 (2022)



# Summary

NLO QED radiative corrections for discussed decays are now available

Meson sector

- $\pi^0 \rightarrow e^+e^-$   
*Vaško and Novotný, JHEP 1110 (2011)*  
*TH, Kampf and Novotný, EPJC 74 (2014)*  
↪ THS model: *TH and S. Leupold, EPJC 75 (2015)*  
↪ measure  $B(\pi^0 \rightarrow e^+e^-)$ , extract  $\chi^{(r)}(M_\rho)$
- $\pi^0 \rightarrow e^+e^-\gamma$   
*TH, Kampf and Novotný, PRD 92 (2015)*  
↪ precise determination of  $R$ : *TH, Goudzovski and Kampf, PRL 122 (2018)*  
↪ could be used in future exp. analysis of  $K^+ \rightarrow \pi^+e^+e^-$
- $\eta^{(\prime)} \rightarrow \ell^+\ell^-\gamma$   
*TH, Kampf, Leupold and Novotný, PRD 97 (2018)*
- $K^+ \rightarrow \pi^+\ell^+\ell^-$   
*TH, in preparation, used in recent NA62 analysis → JHEP 11 (2022) 011*  
↪ SM estimate of  $B(K^+ \rightarrow \pi^+4e)$ : *TH, PRD 106 (2022) 7*

Baryon sector

- $\Sigma^0 \rightarrow \Lambda e^+e^-$   
*TH and Leupold, EPJC 80 (2020)*

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NLO QED radiative corrections for discussed decays are now available

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# Thank you for listening!