

# **1-D Steps Toward Digital Quantum** Simulations of Standard Model Physics - some of our recent results

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ECT\* June 6, 2023





Beimaweres . 1996.

#### **Simulating Physics with Computers**

Department of Physics, California Institute of Technology, Pasadena, California 91107





#### **Richard P. Feynman**

Received May 7, 1981



#### Simulation Objectives for the Standard Model and Beyond **Gauge Theories and Descendent Effective Field Theories and Models**



Real-time dynamics particle production, fragmentation vacuum and in medium

Low-energy reactions

Electroweak processes (e.g., nu-A)

BQP

Neutrino dynamics

Matter-antimatter asymmetry



Equation of state of dense hot matter and dynamics viscosity, etc

Conquering some "sign problems"

The early universe

Supernova/Neutron stars

Precision structure and interactions of nuclei

Many-body systems

Rare processes, double-beta decay

- symmetries





#### **Physical Systems in Multi-Hilbert Space, Hybrid Devices**

#### Map scalar, fermion and vector systems

Optimize for target observables

Minimize time-to-solution within a specified error



Bauer, Davoudi, Klco, Savage







## **Digital To-Do List for Quantum Chromodynamics**

- 1. Map quarks and gluons on a quantum register of qubits, qutrits, ...
- 2. Develop unitary operators to evolve initial wavefunctions forward in time
- **3. Develop observables**









## **Real-Time Dynamics and Improved modeling of Reaction Pathways**



J. Phys. Chem. B 2013, 117, 49, 15894-15902

#### Femto-second chemistry reveals reaction mechanisms Quantum simulations will reveal the reactions pathways of QCD





#### **Gold-Standard for QFT Can entanglement be used more strategically?**

#### Jordan, Lee, Preskill Scalar field theory is BQP-complete

Parallelizes easily at the circuit level - dual layer application per Trotter step X

Double exponential convergence of field digitization

- Nyquist-Shannon JLP, FNAL, UW
- QFT and exact conjugate-momentum space operator





**Could it be done better ? Can entanglement be used strategically?** 







#### Lattice Gauge Field Theories and the Standard Model

Hamiltonian: Kogut-Susskind 1970's

Yang-Mills: Byrnes-Yamamoto 2005

SU(N): Zohar et al (2013)

QLM: Banerjee et al Tagliacozzo et al (2013)

First Quantum Simulation: Innesbruck, 4 Trapped Ions (2016)





## **Yang-Mills Byrnes-Yamamoto – Kogut-Susskind**

Many valid ways to distribute fields in the UV with same IR physics e.g., Kogut-Susskind basis = electric basis



Magnetic Field operator Off-diagonal on electric basis

> SU(N) Gauge invariant Hilbert space

Truncate in Casimir = dimensionality of irrep

Continuum limit



![](_page_8_Figure_10.jpeg)

![](_page_8_Picture_11.jpeg)

![](_page_9_Picture_0.jpeg)

#### Lattice Hamiltonian for two-flavor QCD in 1+1D A<sub>0</sub>=0 Weyl Gauge

Formal : Banuls, Cirac, Jansen, ....

![](_page_9_Figure_3.jpeg)

Staggered Lattice of size 2L with (anti)quarks on (odd) even numbered sites

$$H_{\rm KS} = \sum_{f=u,d} \left[ \frac{1}{2a} \sum_{n=0}^{2L-2} \left( \psi_n^{(f)\dagger} U_n \psi_{n+1}^{(f)} + \text{h.c.} \right) \right]$$

Quark Kinetic Term (Hopping)

Explicit degrees of freedom (qubits) for gauge fields - locally constrained by Gauss's law

Lattice spacing = a

![](_page_9_Figure_9.jpeg)

![](_page_9_Figure_11.jpeg)

![](_page_9_Picture_12.jpeg)

![](_page_10_Picture_0.jpeg)

## Lattice Hamiltonian for two-flavor QCD in 1+1D A<sub>x</sub>=0 Axial Gauge

 $V_{\text{QCD}} \sim g^2 \sum_{a=1}^{8} Q_0^{(a)} Q_2^{(a)}$ 

Staggered Lattice of size 2L with (anti)quarks on (odd) even numbered sites

$$H_{\rm KS} = \sum_{f=u,d} \left[ \frac{1}{2} \sum_{n=0}^{2L-2} \left( \psi_n^{(f)\dagger} \psi_{n+1}^{(f)} + \text{h.c.} \right) + \eta \right]$$

 $m_f \sum_{n=0}^{2L-1} (-1)^n \psi_n^{(f)\dagger} \psi_n^{(f)} \bigg] + \frac{g^2}{2} \sum_{n=0}^{2L-2} \sum_{a=1}^8 \left( \sum_{m \le n} Q_m^{(a)} \right) \bigg]$ Chromo-electric energy Quark Kinetic Term (Hopping) Quark Mass Term

 $\mathcal{U}_{\mathcal{I}}$ 

#### Q<sup>(a)</sup> have diagonal and also off-diagonal action in color space Entangles in color space - distinct from QED

![](_page_10_Picture_7.jpeg)

![](_page_11_Picture_0.jpeg)

## **Simulations using IBM's Quantum Computers** 1+1D QCD

![](_page_11_Figure_2.jpeg)

![](_page_11_Figure_3.jpeg)

![](_page_12_Picture_0.jpeg)

#### **Color Edge States**

![](_page_12_Figure_2.jpeg)

$$H_{1} = \frac{h^{2}}{2} \sum_{n=0}^{2L-1} \left( \sum_{f=0}^{1} Q_{n,f}^{(a)} Q_{n,f}^{(a)} + 2Q_{n,0}^{(a)} \right)$$

![](_page_13_Picture_0.jpeg)

![](_page_13_Figure_2.jpeg)

#### Low-Lying Spectra

![](_page_13_Figure_4.jpeg)

![](_page_14_Picture_0.jpeg)

## Simulations using IBM's Quantum Computers 1-site, 3 colors, 1 flavor: m=g=L=1

#### Circuits to implement Trotterized Gauge term

![](_page_14_Figure_3.jpeg)

![](_page_14_Figure_4.jpeg)

![](_page_15_Picture_0.jpeg)

## Simulations using IBM's Quantum Computers 1-site, 3 colors, 1 flavor: m=g=L=1

![](_page_15_Picture_2.jpeg)

#### Trivial Vacuum-to-Vacuum

![](_page_15_Figure_4.jpeg)

#### IBM 7 qubit Perth and Jakarta

34 CNOTs per step 447 Pauli-Twirled circuits 1000 shots per circuits

Dynamic Decoupling Pauli-Twirling Post selection **De-coherence** renormalization

N	umber of CNOT gates for one Trotter step of $SU($		
L	$N_f = 1$	$N_f = 2$	$N_f = 3$
1	30	114	242
2	228	878	$1,\!940$
5	$1,\!926$	$7,\!586$	$16,\!970$
10	$8,\!436$	$33,\!486$	$75,\!140$
100	912,216	$3,\!646,\!086$	8,201,600

![](_page_15_Picture_9.jpeg)

![](_page_15_Figure_10.jpeg)

![](_page_15_Figure_11.jpeg)

![](_page_16_Picture_0.jpeg)

#### The Difference 5 Years Makes

2017-8

![](_page_16_Figure_3.jpeg)

#### 2022

![](_page_16_Figure_5.jpeg)

![](_page_16_Picture_6.jpeg)

![](_page_16_Picture_7.jpeg)

![](_page_17_Picture_0.jpeg)

#### **Error Mitigation - NISQ-Life**

![](_page_17_Picture_2.jpeg)

#### Coherent errors from errors on control and target transformed to incoherent, averaged over all channels

Mitigating depolarizing noise on quantum computers with noise-estimation circuits

Miroslav Urbanek,<sup>1,\*</sup> Benjamin Nachman,<sup>2</sup> Vincent R. Pascuzzi,<sup>2</sup> Andre He,<sup>2,  $\dagger$ </sup> Christian W. Bauer,<sup>2</sup> and Wibe A. de Jong<sup>1,  $\ddagger$ </sup>

#### **Post-Selection**

Pauli-Twirling

Device Hilbert space >>> Physical Hilbert space Eliminates order-p errors.

![](_page_17_Picture_8.jpeg)

#### Self-mitigating Trotter circuits for SU(2) lattice gauge theory on a quantum

Sarmed <u>A Rahman</u>, Randy <u>Lewis</u>, Emanuele <u>Mendicelli</u>, and Sarah <u>Powell</u> Department of Physics and Astronomy, York University, Toronto, Ontario, Canada, M3J 1P3 (Dated: May 18, 2022)

# Select only members of measurement ensemble entirely in Physical space

![](_page_17_Picture_12.jpeg)

com	puter

![](_page_18_Picture_0.jpeg)

#### **Decoherence Renormalization** The Difference 1 Year Can Make!

#### Self-mitigating Trotter circuits for SU(2) lattice gauge theory on a quantum computer

Sarmed <u>A Rahman</u>, Randy <u>Lewis</u>, Emanuele <u>Mendicelli</u>, and Sarah <u>Powell</u> Department of Physics and Astronomy, York University, Toronto, Ontario, Canada, M3J 1P3

(Dated: May 2022. Updated: October 2022.)

![](_page_18_Figure_5.jpeg)

FIG. 3. Time evolution by self-mitigation on a two-plaquette lattice from the initial state of Fig. 1 with gauge coupling x = 2.0 and time step dt = 0.08. In both panels, the red solid (blue dashed) curve is the exact probability of the left (right) plaquette being measured to have  $j = \frac{1}{2}$ . Upper panel: The red left-pointing (blue right-pointing) triangles are the physics data computed from the ibm\_lagos quantum processor. The red (blue) error bars without symbols are the mitigation data computed on ibm\_lagos from the same circuit but with half the steps forward in time and then half backward in time. Lower panel: The triangles are the physics results obtained by applying Eq. (8) to the data from the upper panel.

![](_page_18_Figure_7.jpeg)

![](_page_19_Picture_1.jpeg)

![](_page_19_Figure_2.jpeg)

#### **Twirling and Decoherence Mitigation**

![](_page_19_Figure_4.jpeg)

![](_page_20_Picture_0.jpeg)

#### Entanglement structure in the mesons for L=2

![](_page_20_Figure_2.jpeg)

Peak in entanglement coincides with transition from quark-antiquark to baryon-anti-baryon structure

![](_page_20_Figure_4.jpeg)

![](_page_20_Picture_7.jpeg)

![](_page_21_Picture_0.jpeg)

## **Semileptonic Weak Decays : L=1, Nf=2**

![](_page_21_Figure_2.jpeg)

nQ=16 qubit JW mapping, with leptons at the end of the site. - 6 quarks, 6 antiquarks and 4 leptons and anti-leptons

![](_page_21_Picture_5.jpeg)

![](_page_21_Picture_6.jpeg)

![](_page_21_Picture_7.jpeg)

![](_page_21_Picture_8.jpeg)

## **Semileptonic Decays : Expectations — Recovering Real-Time Exponential-Decay**

![](_page_22_Picture_1.jpeg)

![](_page_22_Figure_2.jpeg)

One available final state

![](_page_22_Figure_4.jpeg)

![](_page_22_Figure_5.jpeg)

Y<sub>F</sub> available final states

![](_page_22_Picture_7.jpeg)

![](_page_23_Picture_0.jpeg)

## Semileptonic Decays : Real-time Baryon Decay Quantum Simulation using Quantinuum, 16-qubits

![](_page_23_Figure_2.jpeg)

FIG. 9. A quantum circuit that provides the time evolution Hamiltonian, with  $\alpha = \sqrt{2}Gt/8$ .

FIG. 9. A quantum circuit that provides the time evolution associated with the  $\sigma_{\overline{\nu}}^- \sigma_e^+ \sigma_{d,r}^- Z_{u,b} Z_{u,g} \sigma_{u,r}^+$  operator in the  $\beta$ -decay

![](_page_24_Picture_0.jpeg)

## Semileptonic Decays : Real-time Baryon Decay Quantum Simulation using Quantinuum, 16-qubits

![](_page_24_Figure_2.jpeg)

1-step: 59 ZZ

![](_page_24_Figure_4.jpeg)

2-steps: 212 ZZ gates

![](_page_25_Picture_0.jpeg)

#### For 1 Trotter step

L	# of qubits	CNOTs
5	80	9,874 <b>2.5</b> K
10	160	38,074 <b>10</b> K
50	800	926,074 25K
100	1600	3,692,074 900

- $H_{\rm glue} \sim L^2$

#### **Resource Requirements**

• BUT this is too naive — in reality it will be ~ L R, i.e. much fewer gates • R = confinement scale (an emergent scale in the simulation) Now consistent with Feynman's criterion for simulation

![](_page_25_Picture_11.jpeg)

## **State Preparation with Localizable** or Physics-Aware Quantum Circuits

![](_page_26_Picture_1.jpeg)

![](_page_26_Figure_2.jpeg)

![](_page_26_Figure_3.jpeg)

Correlation length allows for fixed-point angles to be determined exponentially well with small-scale simulations

#### Systematically Localizable Operators for Quantum Simulations of Quantum Field Theories

Natalie Klco<sup>\*</sup> and Martin J. Savage<sup> $\dagger$ </sup>

![](_page_26_Figure_7.jpeg)

![](_page_27_Picture_0.jpeg)

#### **Doublers are less of a problem**

From discussions with Anthony Ciavarella

# One d-dim Kogut-Susskind fermion has 2<sup>d</sup>-1 doublers

4+0 D: 1 fermion -> 16 massless fermions, 4 tastes of 4 component Dirac spinors

3+1D: 1 fermion -> 8 massless fermions, 2 tastes of 4 component Dirac spinors

![](_page_28_Picture_0.jpeg)

- Quantum Volume in Simulations and Errors
- Designing quantum circuits that scale efficiently
- (Sufficient) Access to devices TI Vs SC
- 2+1D cold-atom systems
- Integration of HPC and Quantum algorithms
- Can quantum simulations crack chiral gauge theories ?

#### Some of the Challenges

![](_page_29_Picture_0.jpeg)

## **A High-Level Comment and a Conjecture**

We are likely missing an important ingredient so far:

- all of the "power" of computation the gates are being applied at the scale of the lattice spacing
- this becomes increasingly disparate from the scale of physics in the continuum limit
- Seek physics and circuit scalings to move away from the UV

I conjecture that efficient digital quantum circuits exist for Standard Model lattice field theory simulations where the gate-structure, or power, is dominantly focused at the scale of the physics/ observable(s). i.e. EFTs can manifest at the quantum circuit level.

![](_page_29_Picture_8.jpeg)

![](_page_29_Picture_9.jpeg)

# from the Standard Model 1+1D heading toward 2+1D

Advances in computers and simulators

Quantum simulations of Standard-Model physics face challenges

## Summary

• Near-term : Great progress toward dynamic properties of matter

![](_page_30_Picture_7.jpeg)

![](_page_31_Picture_1.jpeg)

## Recent IQuS Workshops – 2022

#### Next-Generation Computing for Nuclear Physics August 2022

![](_page_32_Picture_2.jpeg)

![](_page_32_Picture_3.jpeg)

![](_page_32_Picture_4.jpeg)

#### At the Interface of Quantum Sensors and Quantum Simulations (22-3b)

IQUS InQubator for Quantum Simulation

Organizers: Doug Beck (UIUC), Natalie Klco (Caltech), Crystal Noel (UMD) and Joel Ullom (NIST)

![](_page_32_Picture_8.jpeg)

![](_page_32_Picture_9.jpeg)

Thank you !!

![](_page_32_Picture_11.jpeg)

Institute for People and 158

![](_page_32_Picture_13.jpeg)

IQuS University of W Building, Scattle

![](_page_32_Picture_15.jpeg)

https://iqus.uw.edu/events/iqus-workshop-22-3b/

# 022 Udear Theory

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![](_page_32_Picture_19.jpeg)

w.edu