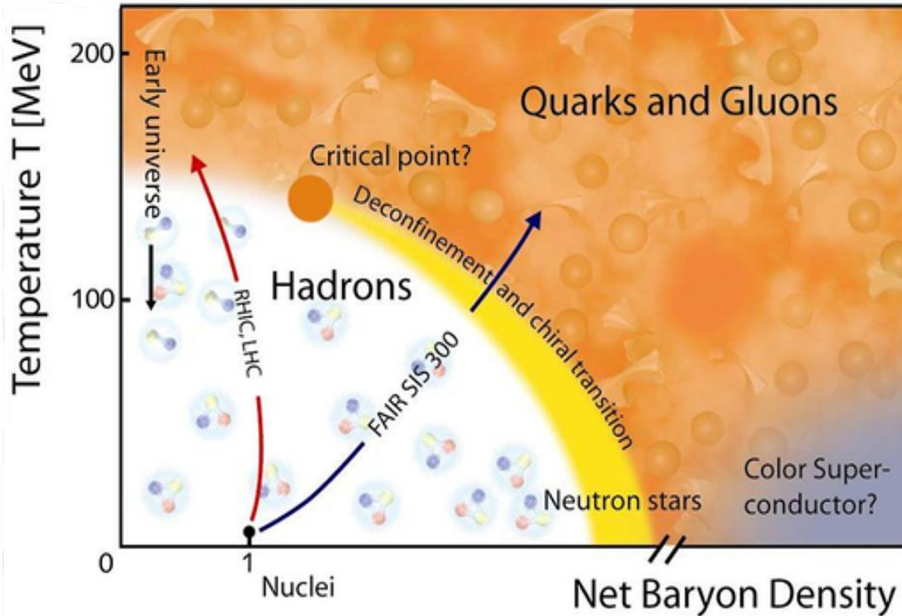


Towards the QCD Phase Structure with the fRG

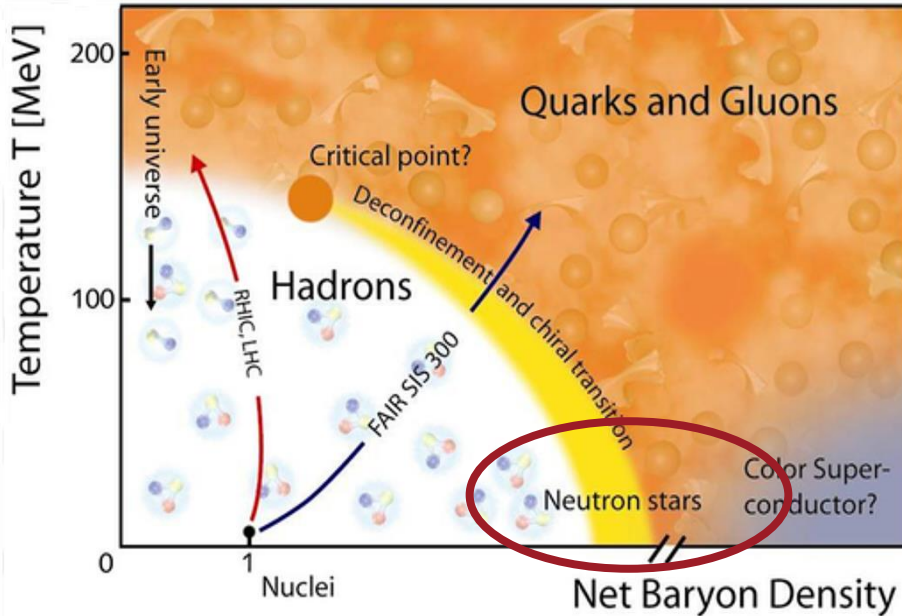
A Comprehensive Analysis

Nicolas Wink

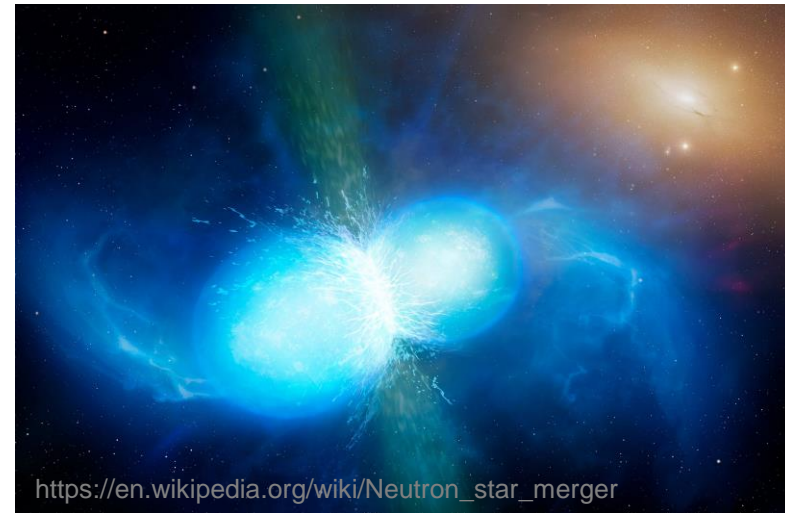
Schematic Phase Structure



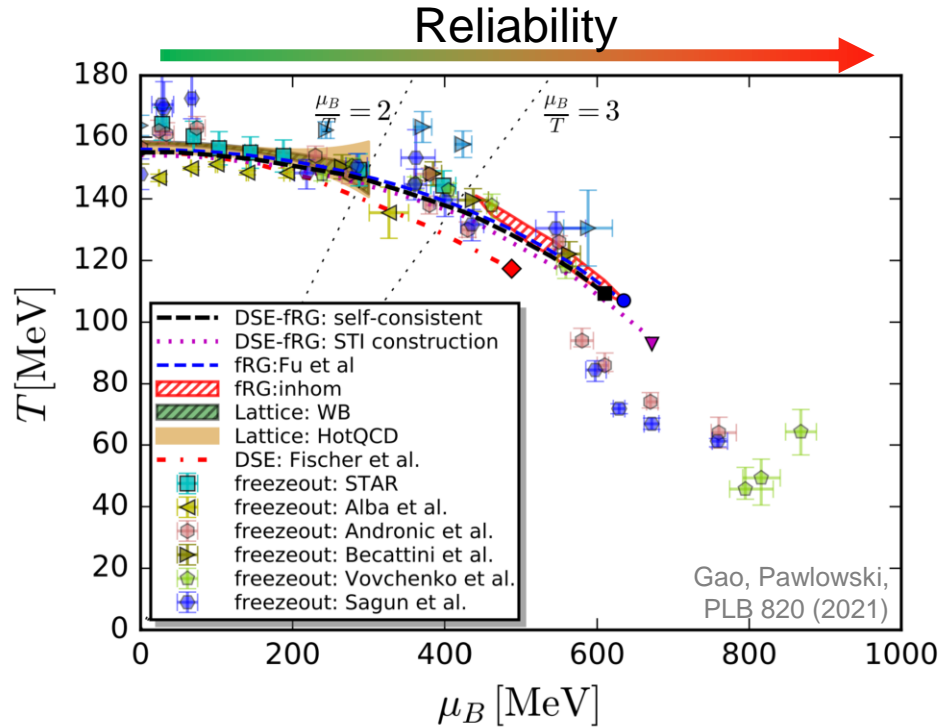
Schematic Phase Structure



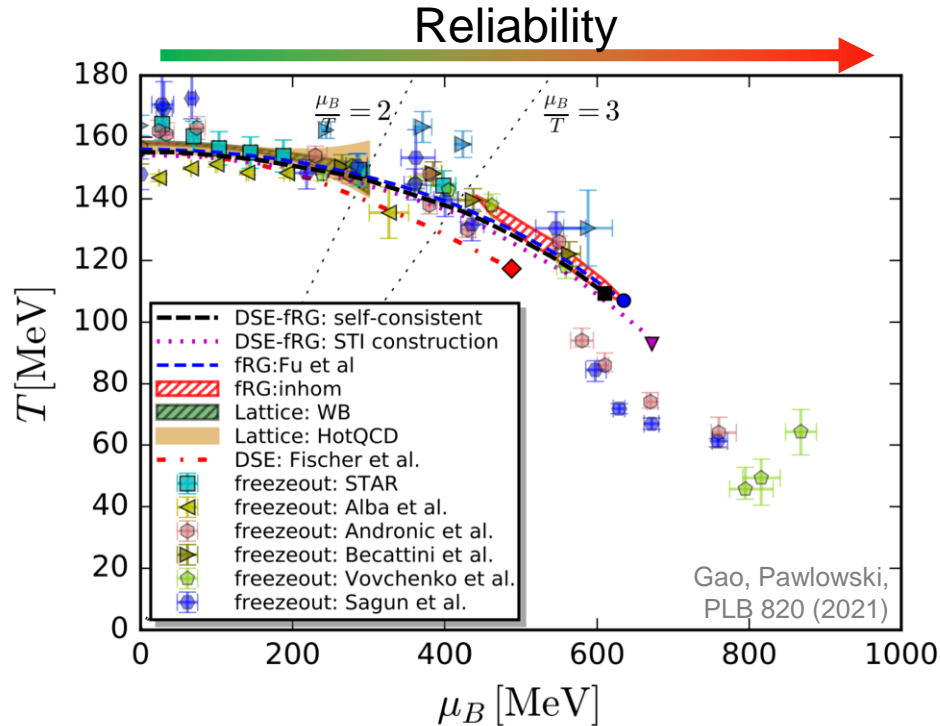
Neutron star mergers



Current state



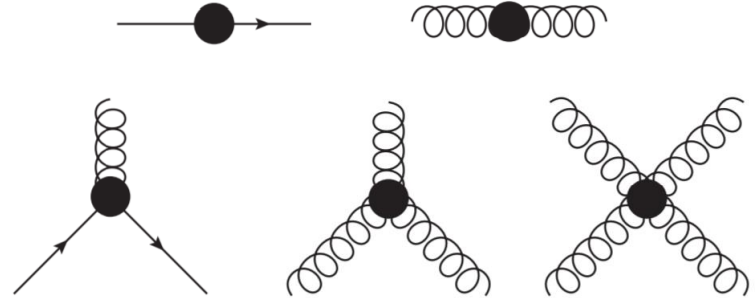
Current state



Focus on functional
Renormalization Group
approach

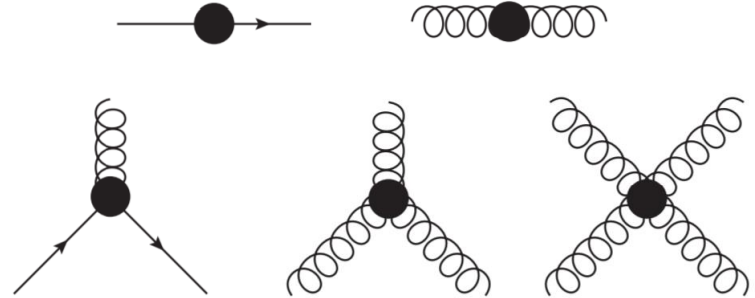
Described by Action

$$S_{\text{QCD}} = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \int_x \bar{q} \not{D} q$$



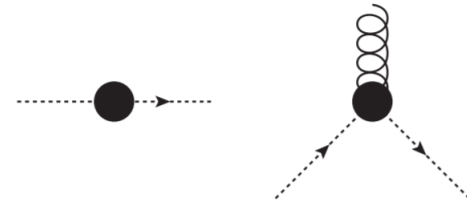
Described by Action

$$S_{\text{QCD}} = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \int_x \bar{q} \not{D} q$$



Functional Methods require gauge fixing

➔ Introduces ghosts (additional terms)



Quantum Effective Action

Generating functional

$$\mathcal{Z}[J] = \frac{1}{\mathcal{N}} \int [d\varphi]_{\text{ren}} \exp \{ -S[\varphi] + J \cdot \varphi \}$$

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Quantum Effective Action

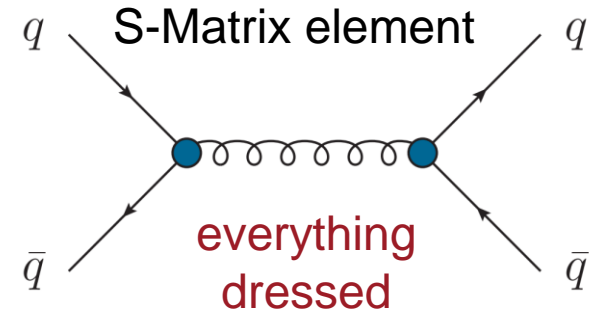
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➔ Encode full information of the QFT



Quantum Effective Action

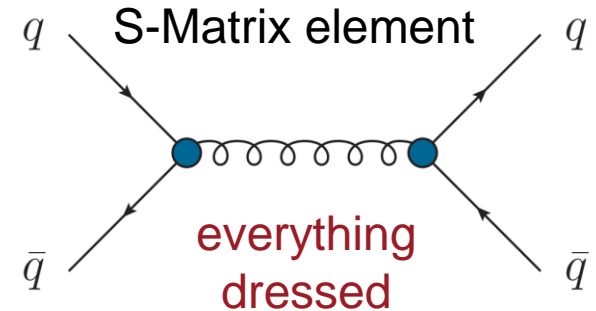
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- ➔ Encode full information of the QFT
- ➔ Think in terms of **tensor structures** (basis of amplitudes)



Wetterich equation

$$\partial_k \Gamma_k[\Phi] = \frac{1}{2} \text{Tr} \left\{ \frac{1}{\Gamma^{(2)}[\phi] + R_k} \partial_k R_k \right\}$$

1-loop exact equation

Wetterich equation

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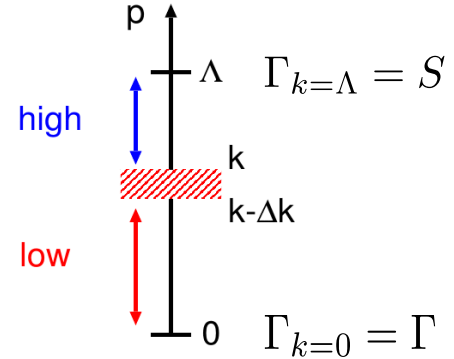
1-loop exact equation

➔ Renormalization scale k and regulator $R_k \sim k^2$ for $p \rightarrow 0$

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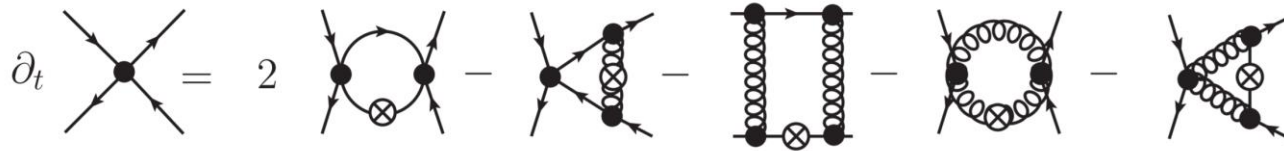


➔ Renormalization scale k and regulator $R_k \sim k^2$ for $p \rightarrow 0$

➔ Γ_k interpolates between the classical action $S = \Gamma_{k=\Lambda}$ and the full QEA $\Gamma = \Gamma_{k=0}$

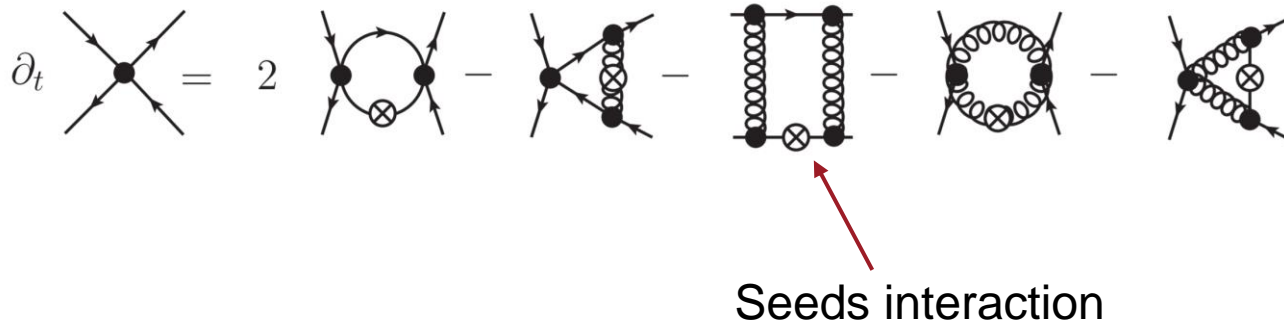
Dynamical Hadronization

RG flow of the four-quark interaction



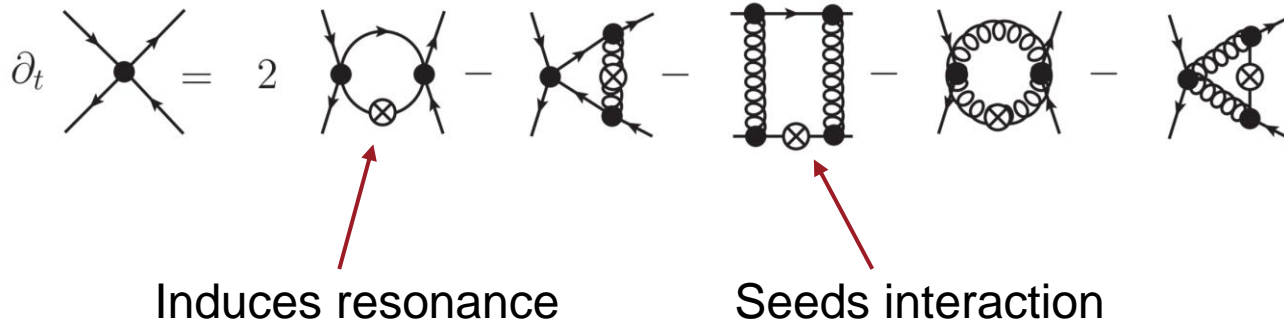
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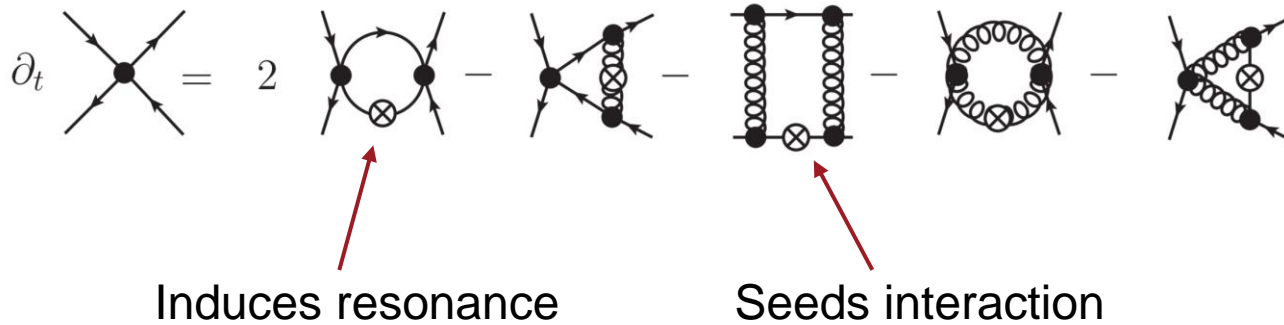


Dynamical Hadronization

RG flow of the four-quark interaction



RG flow of the four-quark interaction



Absorb resonance in effective field

$$\phi \propto \bar{q}(T_f^0, i\gamma_5 T_f^a)q$$

Introduces correction terms

Exact transformation

$$\begin{aligned} & \partial_t \Gamma_k[\Phi] + \int \langle \partial_t \hat{\phi}_{k,i} \rangle \left(\frac{\delta \Gamma_k[\Phi]}{\delta \phi_i} + c_\sigma \delta_{i\sigma} \right) \\ &= \frac{1}{2} \text{Tr}(G_k[\Phi] \partial_t R_k) + \frac{1}{2} \text{Tr} \left(G_{\phi\Phi_j}[\Phi] \frac{\delta \langle \partial_t \hat{\phi}_k \rangle}{\delta \Phi_j} R_\phi \right) \end{aligned}$$

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Introduced scalar field is O(4)-symmetric

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Introduced scalar field is O(4)-symmetric

- ➡ Components typically identified as pion and sigma meson
- ➡ Breaking of O(4) symmetry ↔ Chiral symmetry breaking

Dynamic emergence of LEFT

Fields decouple at their mass scale

Regularized propagator

$$G_k \sim \frac{1}{k^2 + m^2}$$

Dynamic emergence of LEFT

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Effective decoupling

Gluons

Quarks

Sigma/Pions

RG scale



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Natural to work in terms of
Low Energy Effective Theories

Current limitations to the
reliability at higher densities

Truncation used in Fu, Pawłowski, Rennecke, PRD 101 (2020)

$$\begin{aligned}\Gamma_k = \int_x \left\{ \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + Z_c (\partial_\mu \bar{c}^a) D_\mu^{ab} c^b + \frac{1}{2\xi} (\partial_\mu A_\mu^a)^2 \right. \\ + \frac{1}{2} \int_p A_\mu^a(-p) \left(\Gamma_{AA\mu\nu}^{(2)ab}(p) - Z_A \Pi_{\mu\nu}^\perp \delta^{ab} p^2 \right) A_\nu^b(p) \\ + \bar{q} \left[Z_q (\gamma_\mu D_\mu - \gamma_0 \hat{m}) + m_s(\sigma_s) \right] q \\ - \lambda_q \left[(\bar{q} \tau^0 q)^2 + (\bar{q} \boldsymbol{\tau} q)^2 \right] + h \bar{q} (\tau^0 \sigma + \boldsymbol{\tau} \cdot \boldsymbol{\pi}) q \\ \left. + \frac{1}{2} Z_\phi (\partial_\mu \phi)^2 + V_k(\rho, A_0) - c_\sigma \sigma - \frac{1}{\sqrt{2}} c_{\sigma_s} \sigma_s \right\},\end{aligned}$$

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Current limitations to the reliability at higher densities

➔ Competing order effects in 4-quark interaction

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Current limitations to the reliability at higher densities

➔ Competing order effects in 4-quark interaction

➔ Proper resolution of field dependence

Truncation used in Fu, Pawłowski, Rennecke, PRD 101 (2020)

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Competing order in the four-quark sector



TECHNISCHE
UNIVERSITÄT
DARMSTADT

Four-quark interaction dominates

➔ Good arguments that the four-quark interaction dominates

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Small to moderate couplings:

Vertex expansion in the fermionic sector converges + four-quark dominates

Metzner, Salmhofer, Honerkamp, Meden, Schönhammer, *Reviews of Modern Physics* 84 (2012)

Four-quark interaction dominates

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Do gluons change this? Most likely not

Genuine three-body forces sub-leading in Baryon masses

e.g. Eichmann, Few Body Syst 63 (2022)

Four-quark sector

A lot of different channels

Four-quark sector

A lot of different channels

➔ Which are important?

Four-quark sector

A lot of different channels

➔ Which are important?

At chemical potentials $\frac{\mu_B}{T} \lesssim 2-3$ well known

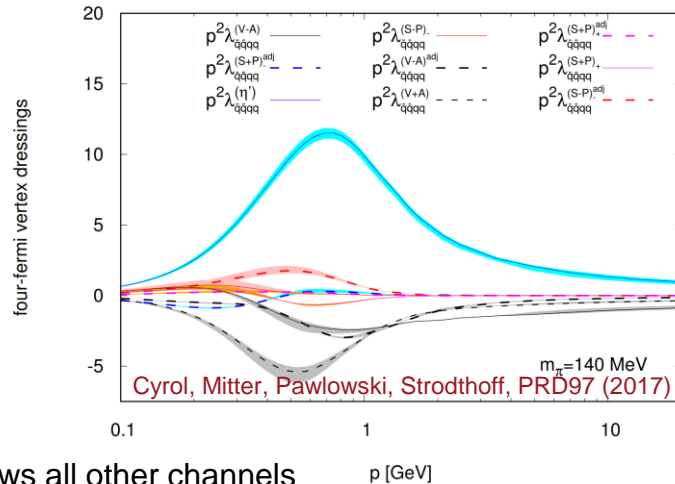
Scalar Pseudo-Scalar channel dominates

Four-quark sector

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Scalar Pseudo-Scalar channel dominates

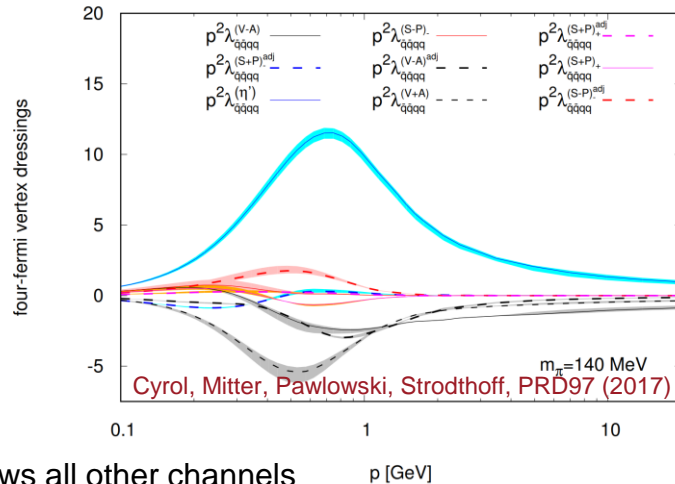
Shows all other channels

Four-quark sector

A lot of different channels

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See also

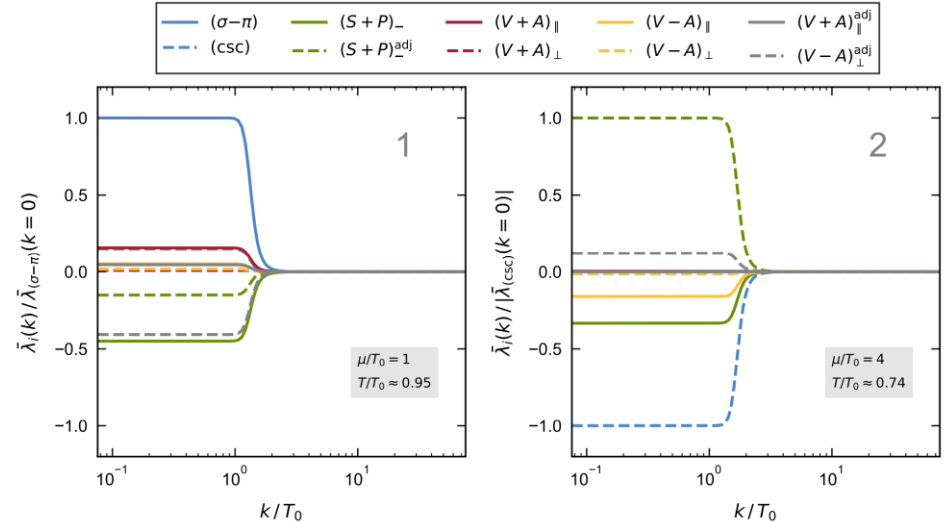
Braun, Leonhardt, Pospiech 2017-2020

Fu, Huang, Pawlowski, Tan, SciPost 14 (2023)

Talk by *Wei-Jie Fu*

Four-quark at larger densities

Other channels become important

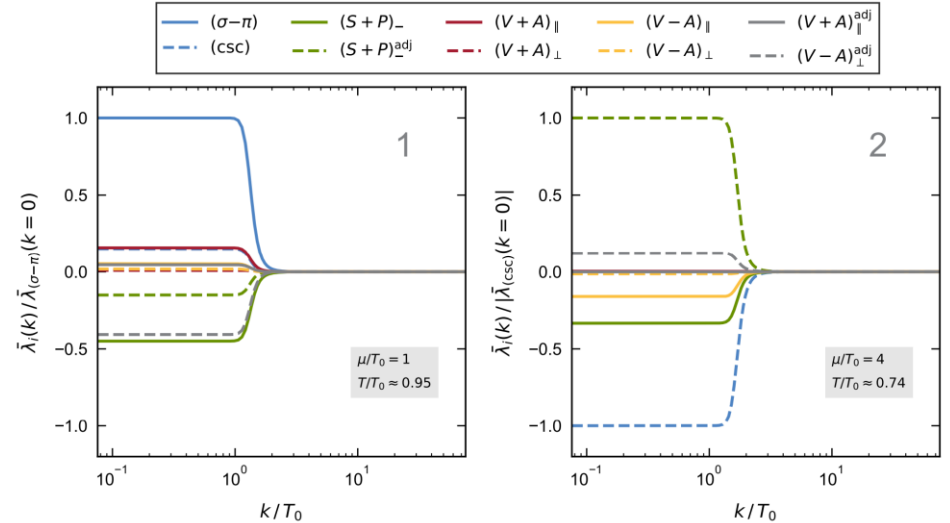


Braun, Leonhardt, Pospiech, PRD 101 (2020)

Four-quark at larger densities

Other channels become important

➔ Diquarks



Braun, Leonhardt, Pospiech, PRD 101 (2020)

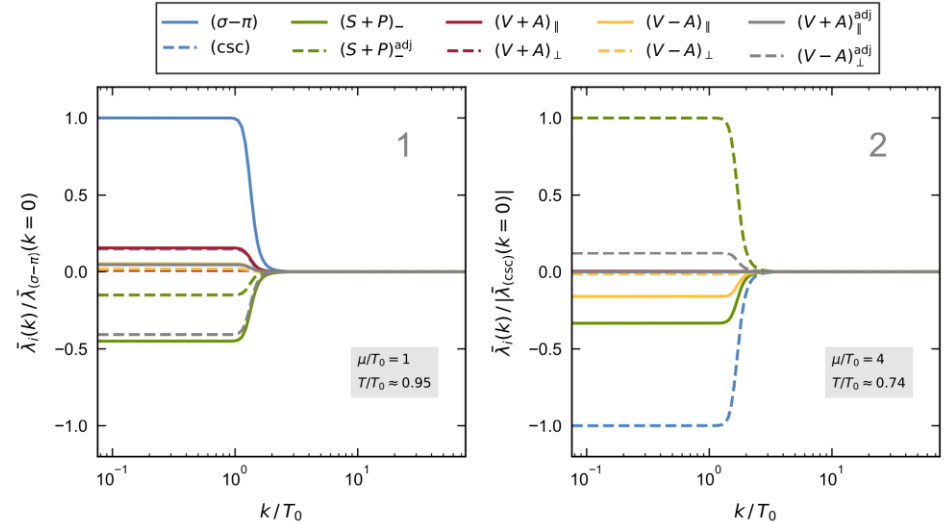
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Other potentially important effects

- Anomaly
- Vector channel
- ...



Braun, Leonhardt, Pospiech, PRD 101 (2020)

Four-quark at larger densities

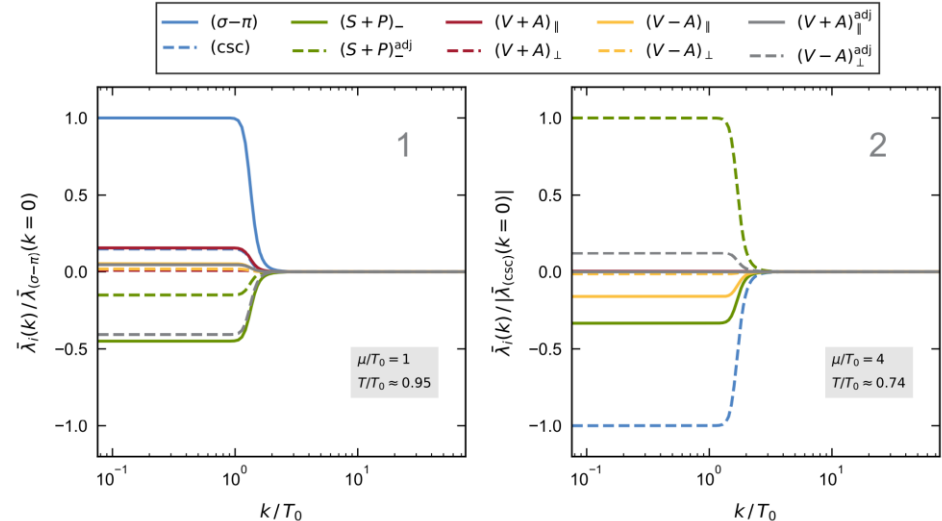
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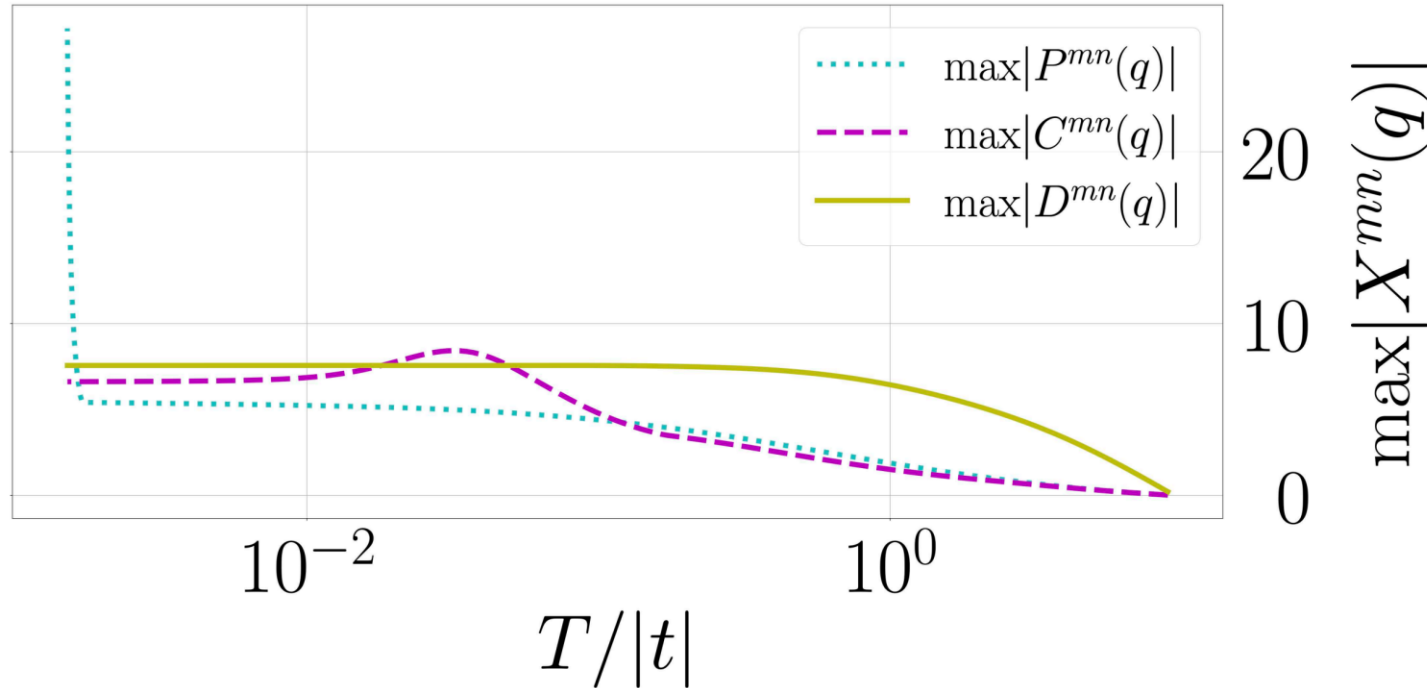
- Anomaly
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- ...

➔ Work in Fierz complete Basis



Braun, Leonhardt, Pospiech, PRD 101 (2020)

Competing order



Gneist, Classen, Scherer PRB 106 (2022)

Resolving field dependencies

Why field dependencies matter

“Field dependencies = Expectation values dependencies”

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➔ Derivatives encodes higher order mesonic scatterings

Why field dependencies matter

“Field dependencies = Expectation values dependencies”

- ➔ Derivatives encodes higher order mesonic scatterings
- ➔ In the vicinity of first order transitions qualitatively important
- ➔ Formation of discontinuities

Why field dependencies matter

“Field dependencies = Expectation values dependencies”

- ➔ Derivatives encodes higher order mesonic scatterings
- ➔ In the vicinity of first order transitions qualitatively important
- ➔ Formation of discontinuities

Importance starts to increase significantly at large chemical potential

RG-flows & convection

Subsequently integrate out momentum shells

RG-flows & convection

Subsequently integrate out momentum shells

➔ Introduces strong notion of directionality

RG-flows & convection

Subsequently integrate out momentum shells

➔ Introduces strong notion of directionality

➔ **convection dominated**

Subsequently integrate out momentum shells

➔ Introduces strong notion of directionality

➔ convection dominated

Numerical scheme **must** reflect this

➔ Discontinuous Galerkin

➔ Finite Volume

➔ Upwind Finite Difference

Subsequently integrate out momentum shells

➡ Introduces strong notion of directionality

➡ convection dominated

Numerical scheme **must** reflect this

➡ Discontinuous Galerkin

➡ Finite Volume

➡ Upwind Finite Difference

Grossi, NW, arxiv:1903.09503

Grossi, Ihssen, Pawlowski, NW, PRD 104 (2021)

Koenigstein, Steil, NW, Grossi, Braun, Buballa, Rischke, PRD 106 (2022)

Koenigstein, Steil, NW, Grossi, Braun, PRD 106 (2022)

Steil, Koenigstein, PRD 106 (2022)

Stoll, Zorbach, Koenigstein, Steil, Rechenberger arxiv: 2108.10616

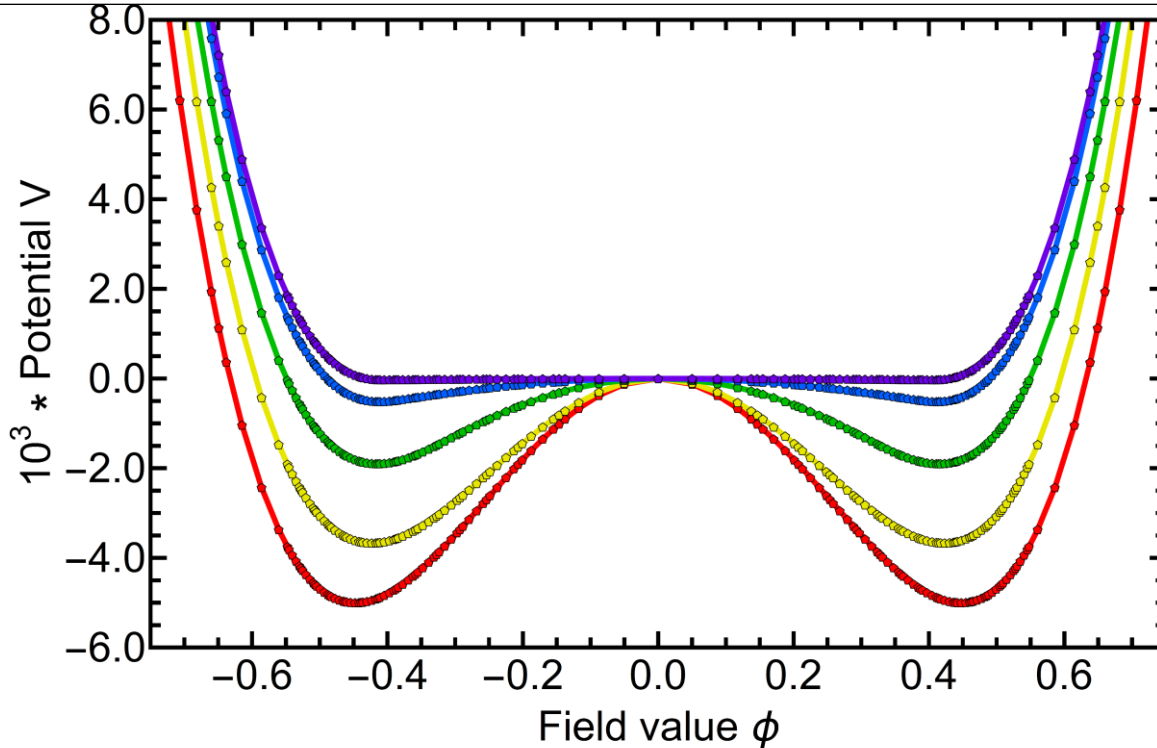
Ihssen, Pawlowski, arxiv:2207.10057

Ihssen, Pawlowski, Sattler, NW arxiv:2207.12266

Ihssen, Sattler, NW arxiv:2302.04736 (to appear in PRD)

Murgana, Koenigstein, Rischke arXiv:2303.16838

Illustration – Effective potential

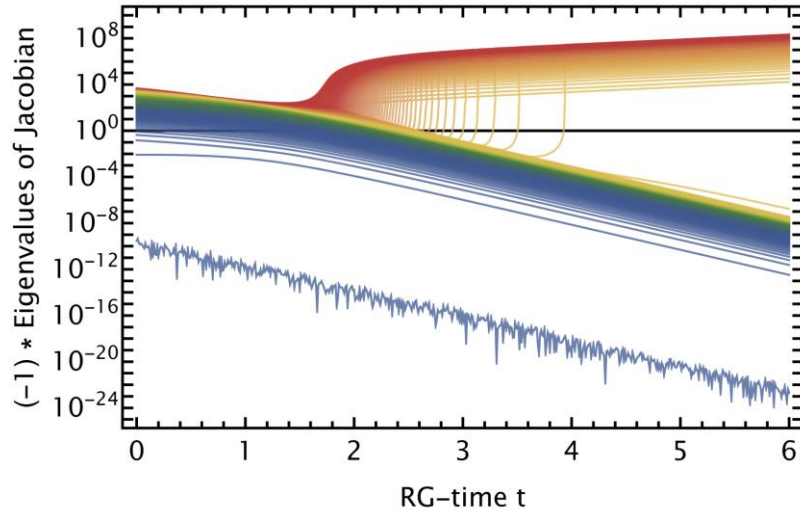


➔ Broken phase

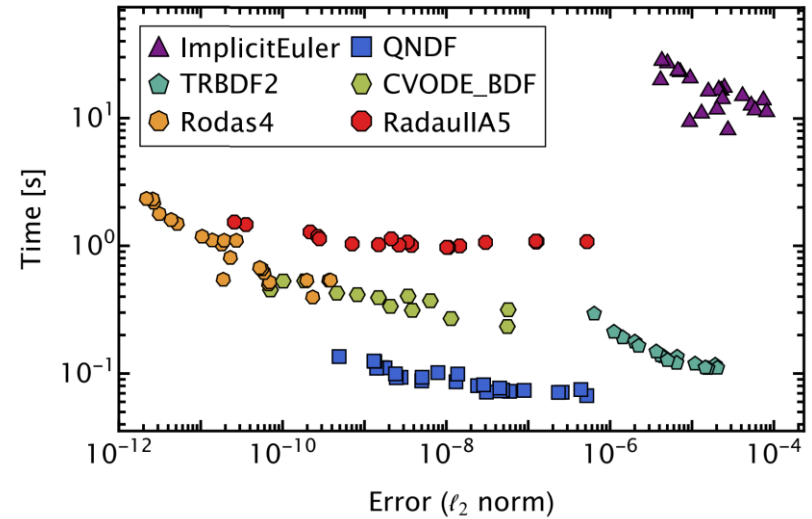
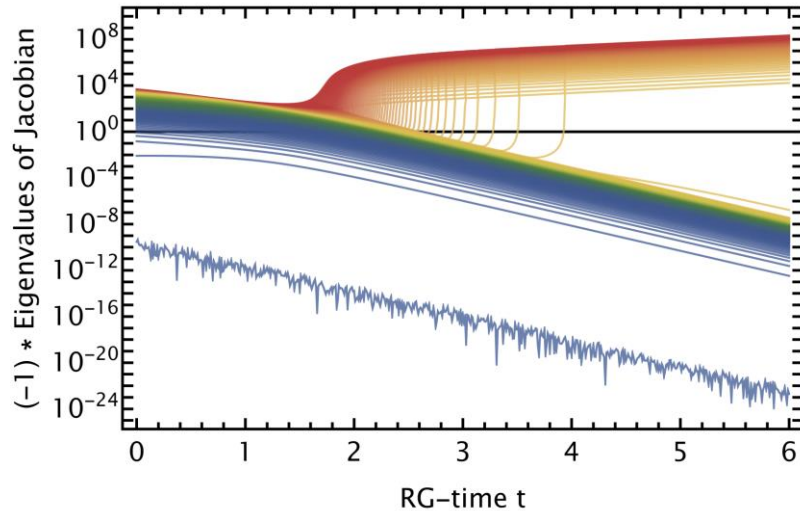
➔ Convexity restoration

Grossi, NW, arxiv:1903.09503

Numerical aspects of the RG-scale evolution



Numerical aspects of the RG-scale evolution



Ihssen, Sattler, NW arxiv:2302.04736

Low energy effective theories

Parallel development of technical advances in LEFTs and QCD

Low energy effective theories

Parallel development of technical advances in LEFTs and QCD

➔ Isolate (and separate) technical and conceptual problems

Low energy effective theories

Parallel development of technical advances in LEFTs and QCD

➔ Isolate (and separate) technical and conceptual problems

Current construction sites:

➔ Higher order scattering

➔ Diquarks

➔ Vector mesons

Simple model only including quarks and mesons

➔ Ideal for technical developments

$$\Gamma_k[\bar{\psi}, \psi, \phi] = \int_x \left\{ i\bar{q}(\not{\partial} + \mu\gamma_0)q + \frac{1}{2}(\partial_\mu\phi)^2 + h_k(\rho)\bar{q}(\tau_0\sigma + \boldsymbol{\tau}\boldsymbol{\pi})q + V_k(\rho) - c_\sigma\sigma \right\}$$

Simple model only including quarks and mesons

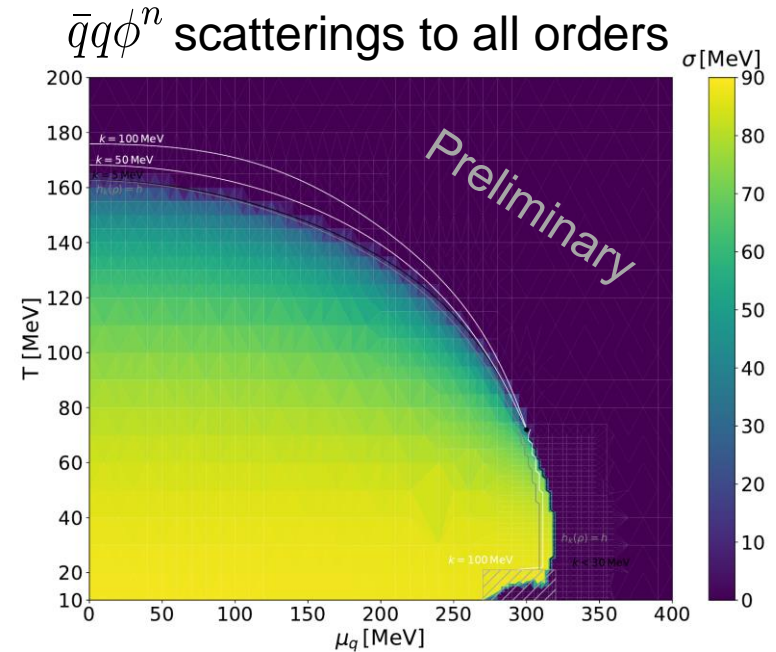
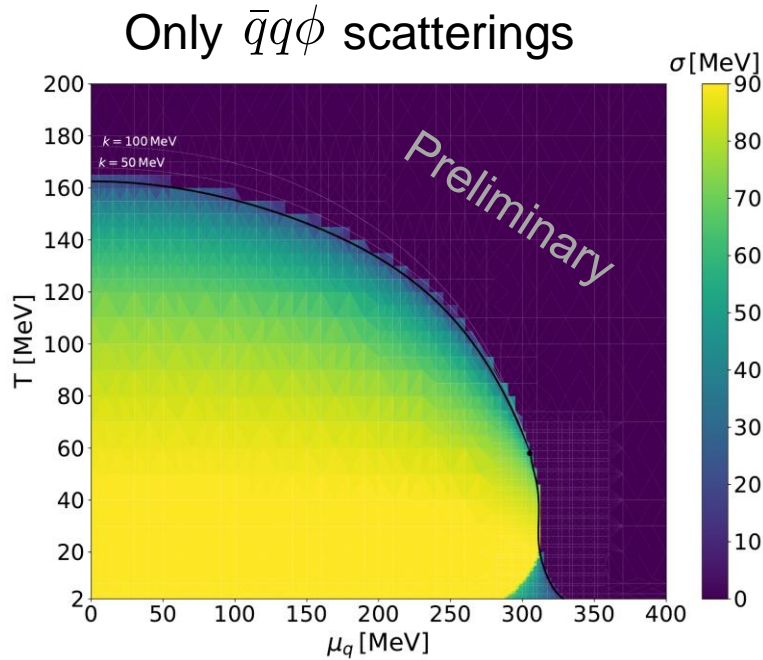
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➔ Include n-meson—two-quark scatterings

$$\left. + h_k(\rho)\bar{q}(\tau_0\sigma + \boldsymbol{\tau}\boldsymbol{\pi})q + V_k(\rho) - c_\sigma\sigma \right\}$$

Quark-Meson model



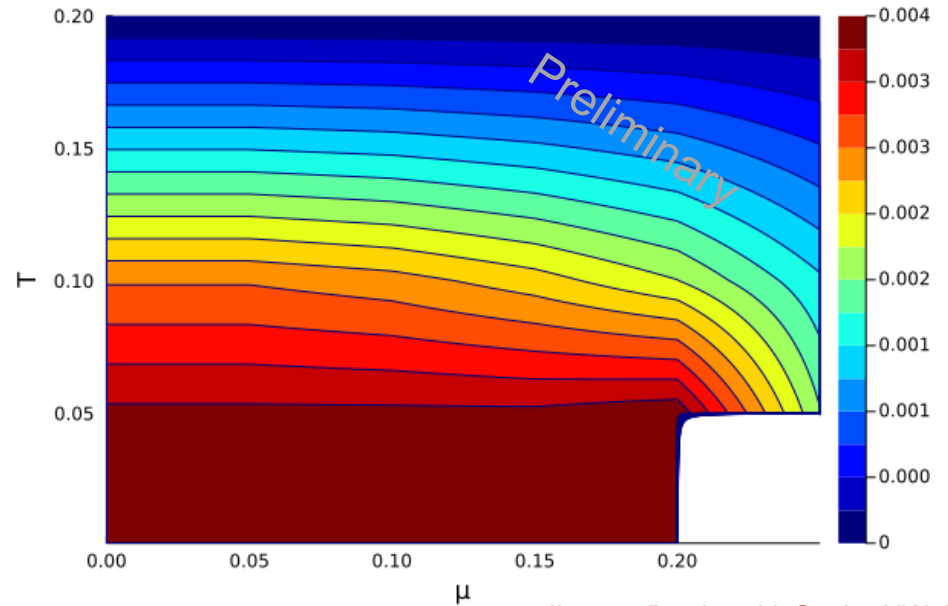
Ihssen, Pawlowski, Sattler NW, in prep

Diquarks

Diquark channel important at large densities

➔ Inclusion important to go beyond critical endpoint

cf. talk by **Ugo Mire**



Ihssen, Pawlowski, Sattler NW, in prep

Including Vector exchange channels

Dynamical Hadronization of the Vector channel in the four-quark interaction

Rennecke, Phys.Rev.D 92 (2015)

Fukushima, Pawlowski, Strodthoff, Annals Phys. 446 (2022)

Including Vector exchange channels

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➔ Important to describe Liquid-Gas transition

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➔ Important to describe Liquid-Gas transition

➔ Additional “spatial” direction in resulting PDEs

$$\Gamma_k = \int_p \left(\bar{q} (i\not{p} - \mu_q \gamma_0) q + \frac{1}{2} \phi p^2 \phi + \frac{1}{2} \omega_\mu p^2 \omega_\mu + h_\omega \bar{q} \psi q + h_\phi \bar{q} (\sigma + i\gamma_5 \vec{\tau} \vec{\pi}) q + V(\rho, \omega^2) \right)$$

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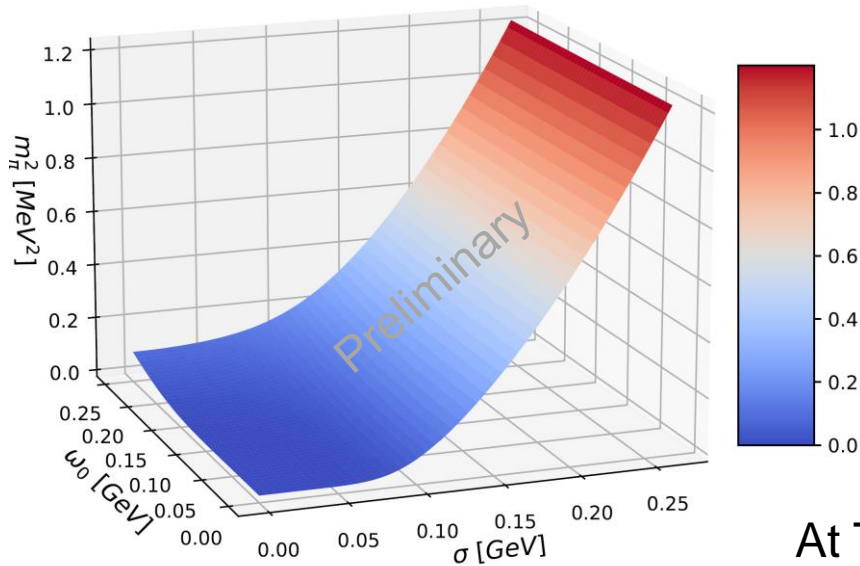
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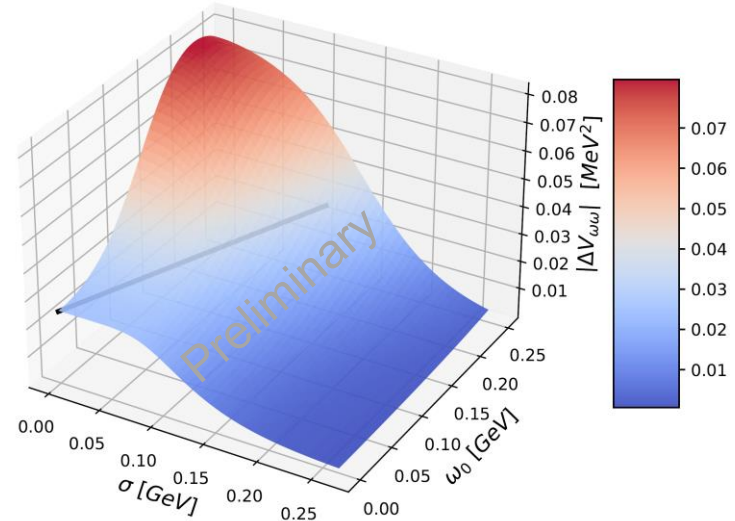
➔ Possible to absorb chemical potential into ω_0

Including Vector exchange channels

Pion mass



ω_0 mass correction



At T=150MeV

Master Thesis N. Hendricks

Technical progress included in QCD

Technical progress included in QCD

→ Include all order
mesonic scatterings

Technical progress included in QCD



Include all order
mesonic scatterings

Improved truncation

$$\begin{aligned}
 \Gamma_k = & \int_x \left\{ \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \frac{1}{2\zeta} (\partial_\mu A_\mu^a)^2 \right. \\
 & + \frac{1}{2} \int_p A_\mu^a(-p) \left(\Gamma_{AA\mu\nu}^{(2)ab}(p) - Z_A \Pi_{\mu\nu}^\perp \delta^{ab} p^2 \right) A_\nu^b(p) \\
 & + Z_c (\partial_\mu \bar{c}^a) D_\mu^{ab} c^b + \bar{q} Z_q (\gamma_\mu D_\mu - \gamma_0 \hat{m}_q) q \\
 & - \underline{\lambda_q(\rho)} \left[(\bar{q} \tau^0 q)^2 + (\bar{q} \boldsymbol{\tau} q)^2 \right] \\
 & \left. + \underline{h(\rho)} \bar{q} (\tau^0 \sigma + \boldsymbol{\tau} \cdot \boldsymbol{\pi}) q + \frac{1}{2} Z_\phi (\partial_\mu \phi)^2 + V(\rho) - c_\sigma \sigma \right\}
 \end{aligned}$$

Technical progress included in QCD

➔ Include all order mesonic scatterings

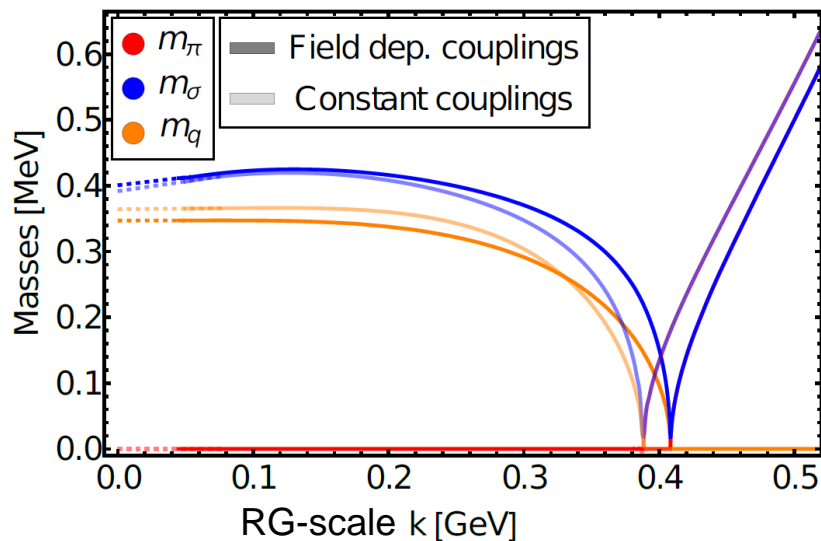
Ideal outcome:

➔ No/very small changes at vanishing/small chemical potential

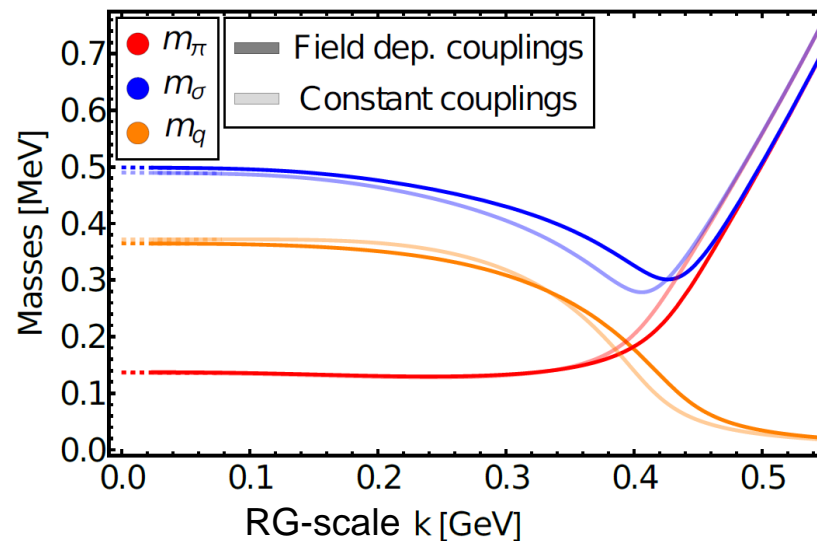
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 \end{aligned}$$

Chiral limit

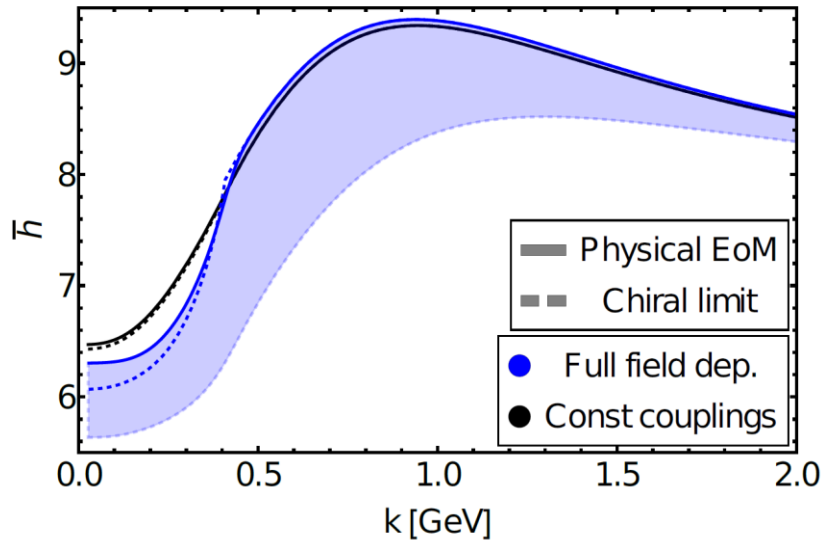


Physical point

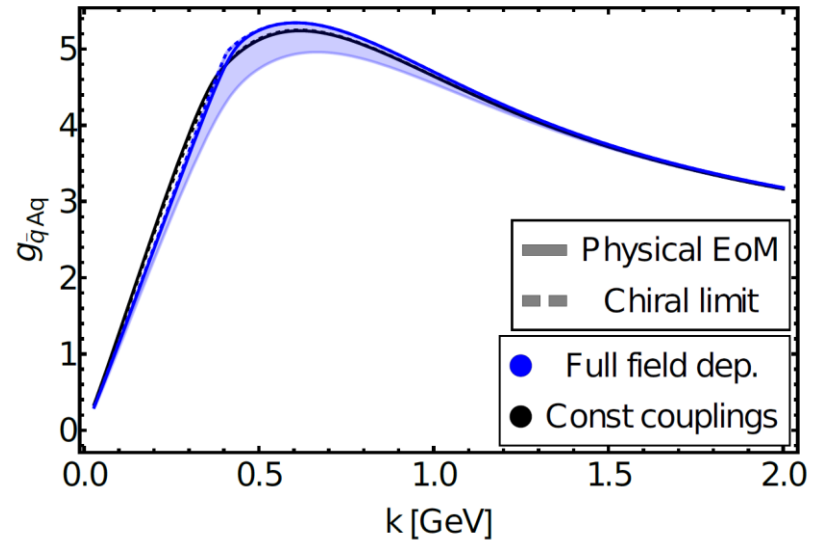


Ihssen, Pawłowski, Sattler NW, in prep

Yukawa coupling



Quark-Gluon coupling



Ihssen, Pawłowski, Sattler NW, in prep

Connect to phenomenology

Connect fRG to Transport/Hydro

Reduce uncertainties from QCD in phenomenology

Batini, Grossi, NW, in prep

Connect fRG to Transport/Hydro

Reduce uncertainties from QCD in phenomenology

➔ Use fRG input for critical mode in Transport/Hydro evolution

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Connect fRG to Transport/Hydro

Reduce uncertainties from QCD in phenomenology

- ➔ Use fRG input for critical mode in Transport/Hydro evolution
- ➔ Leading order contribution: Effective potential

Connect fRG to Transport/Hydro

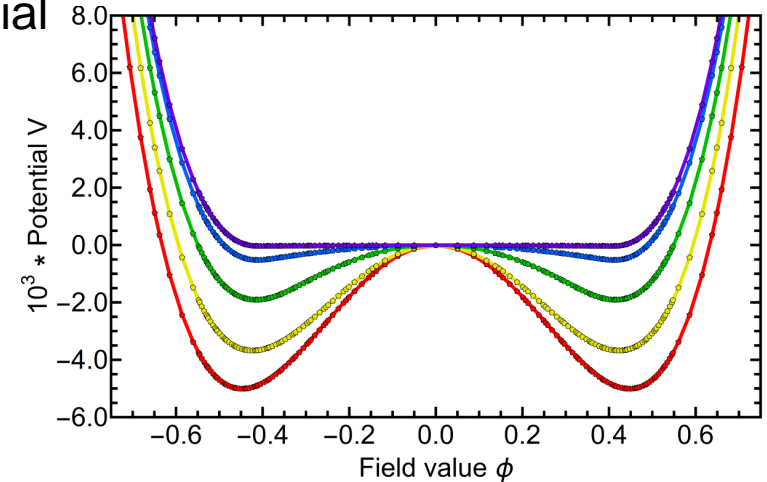
Reduce uncertainties from QCD in phenomenology

➔ Use fRG input for critical mode in Transport/Hydro evolution

➔ Leading order contribution: Effective potential

... Full effective potential is convex

➔ How to resolve in practice?



Batini, Grossi, NW, in prep

Relaxation rate

Simplest model to study

➔ Model A

Relaxation rate

Simplest model to study

➔ Model A

Consider evolution of scalar field in a hydro setting

$$u^\mu \partial_\mu \phi = \Gamma \left[\partial_\perp^2 \phi - \frac{\partial U}{\partial \phi} \right] + \zeta^{(1)} \phi \partial_\mu u^\mu$$

Batini, Grossi, NW, in prep

Relaxation rate

Simplest model to study

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Need this

Batini, Grossi, NW, in prep

Relaxation rate

Simplest model to study

➔ Model A

Consider evolution of scalar field in a hydro setting

Scalar field on Keldysh contour

$$u^\mu \partial_\mu \phi = \Gamma \left[\partial_\perp^2 \phi - \frac{\partial U}{\partial \phi} \right] + \zeta^{(1)} \phi \partial_\mu u^\mu$$

Need this

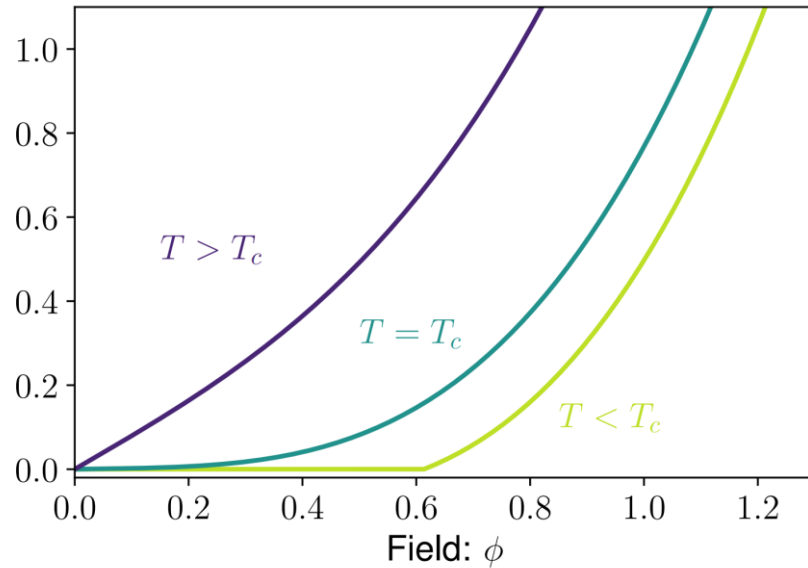
Relaxation rate $X = \frac{1}{\Gamma}$

$$\Gamma_k[\Phi] = \int_{t, \mathbf{x}} \phi_a \left(-X_k(\phi_r) (\partial_t \phi_r - i\phi_a) - \nabla^2 \phi_r + U_k^{(1)}(\phi_r) \right)$$

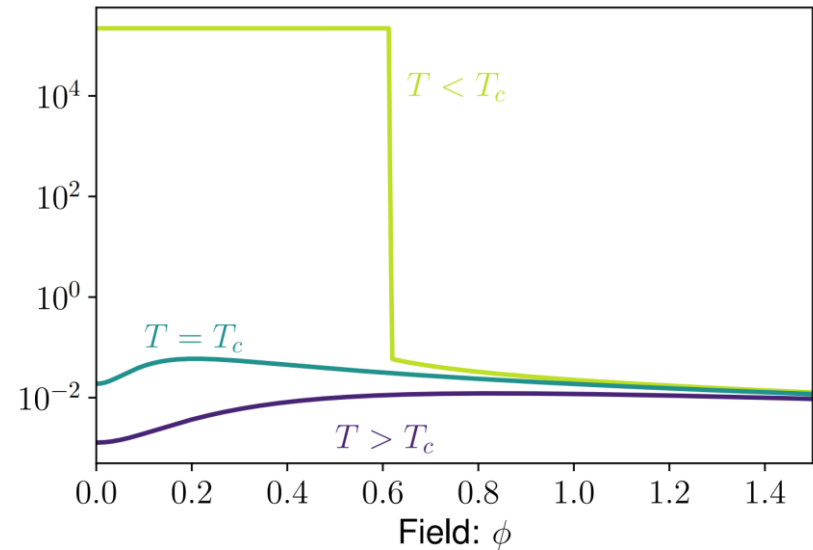
Batini, Grossi, NW, in prep

Relaxation rate

Derivative effective potential

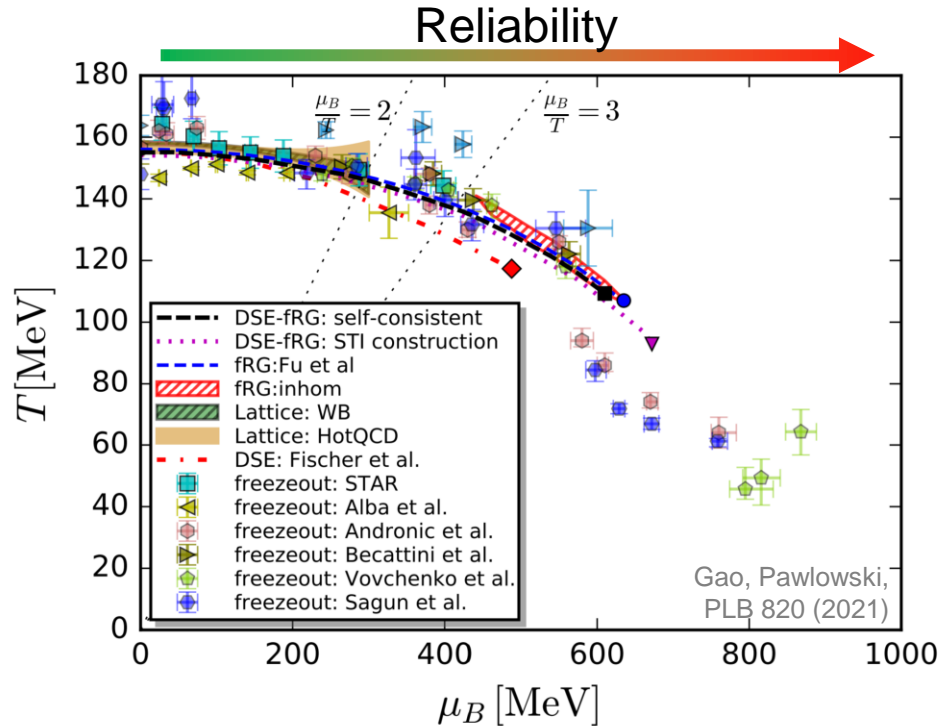


Relaxation rate



Batini, Grossi, NW, in prep

Summary



➔ Why it works up to $\frac{\mu_B}{T} \lesssim 3$

➔ Efforts and status to increase chemical potential

➔ Connection to dynamic evolution