

Hadronic resonances from Lattice QCD

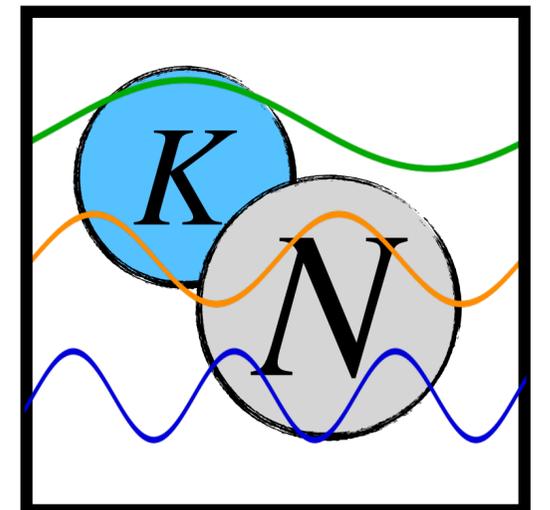
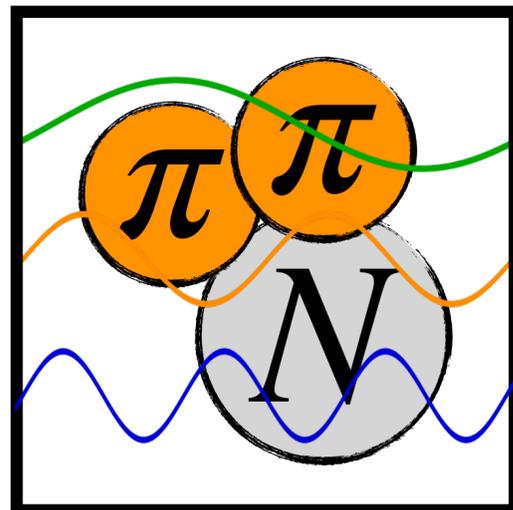
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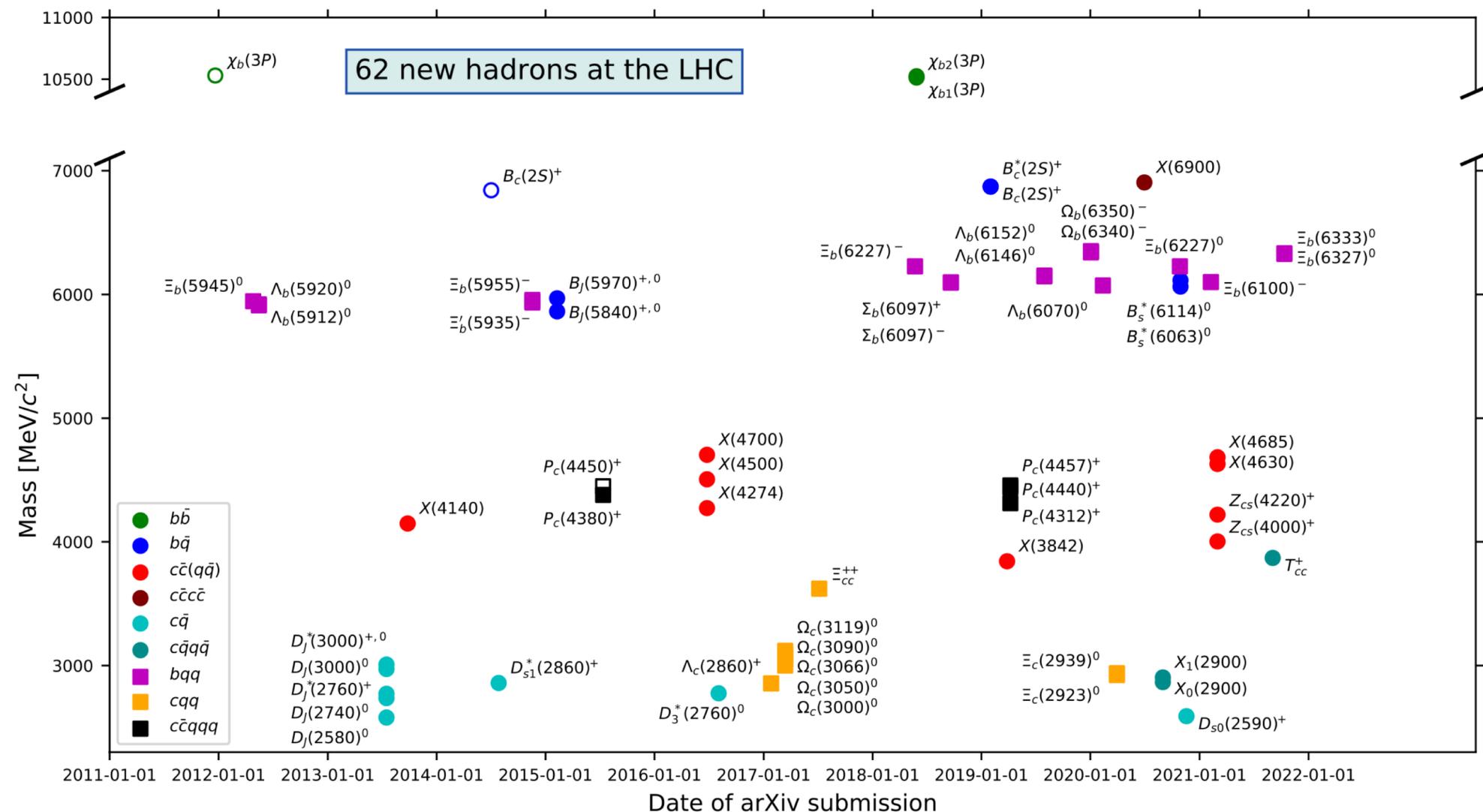
from first-principles QCD to experiment @ ECT*

May 24th



The Hadron Spectrum

Lattice QCD offers a systematically improvable approach to investigate the hadron spectrum

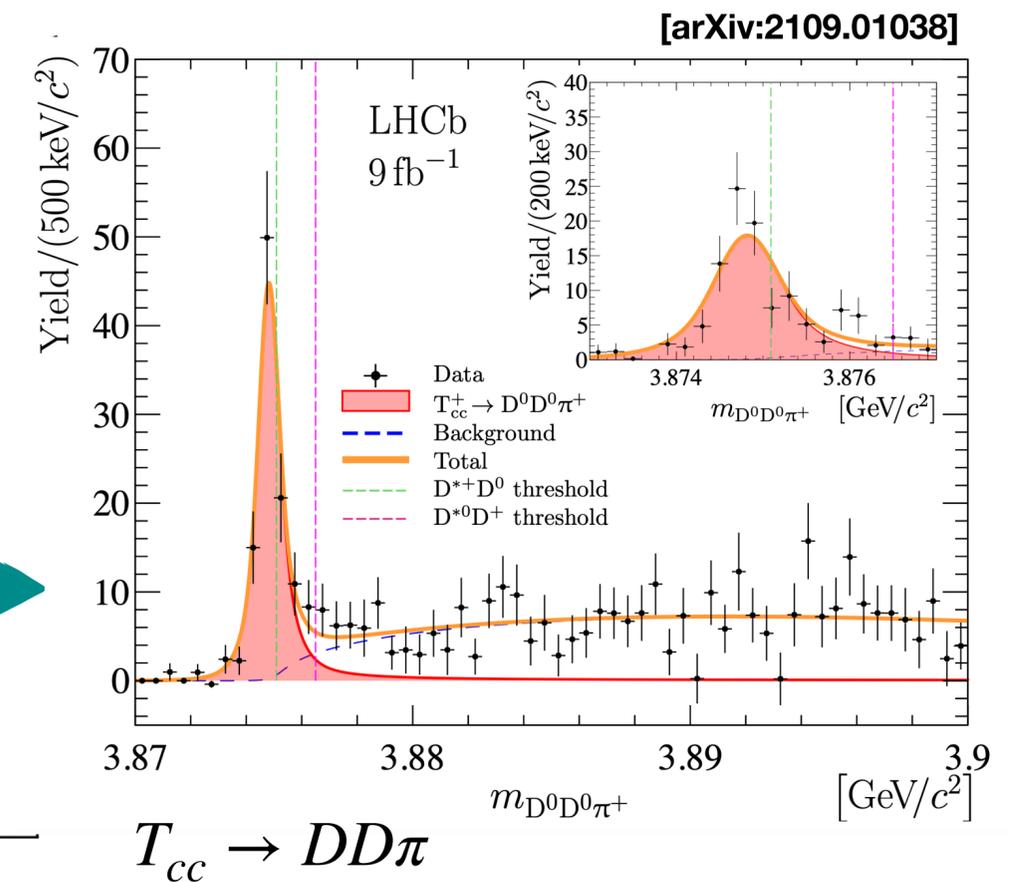
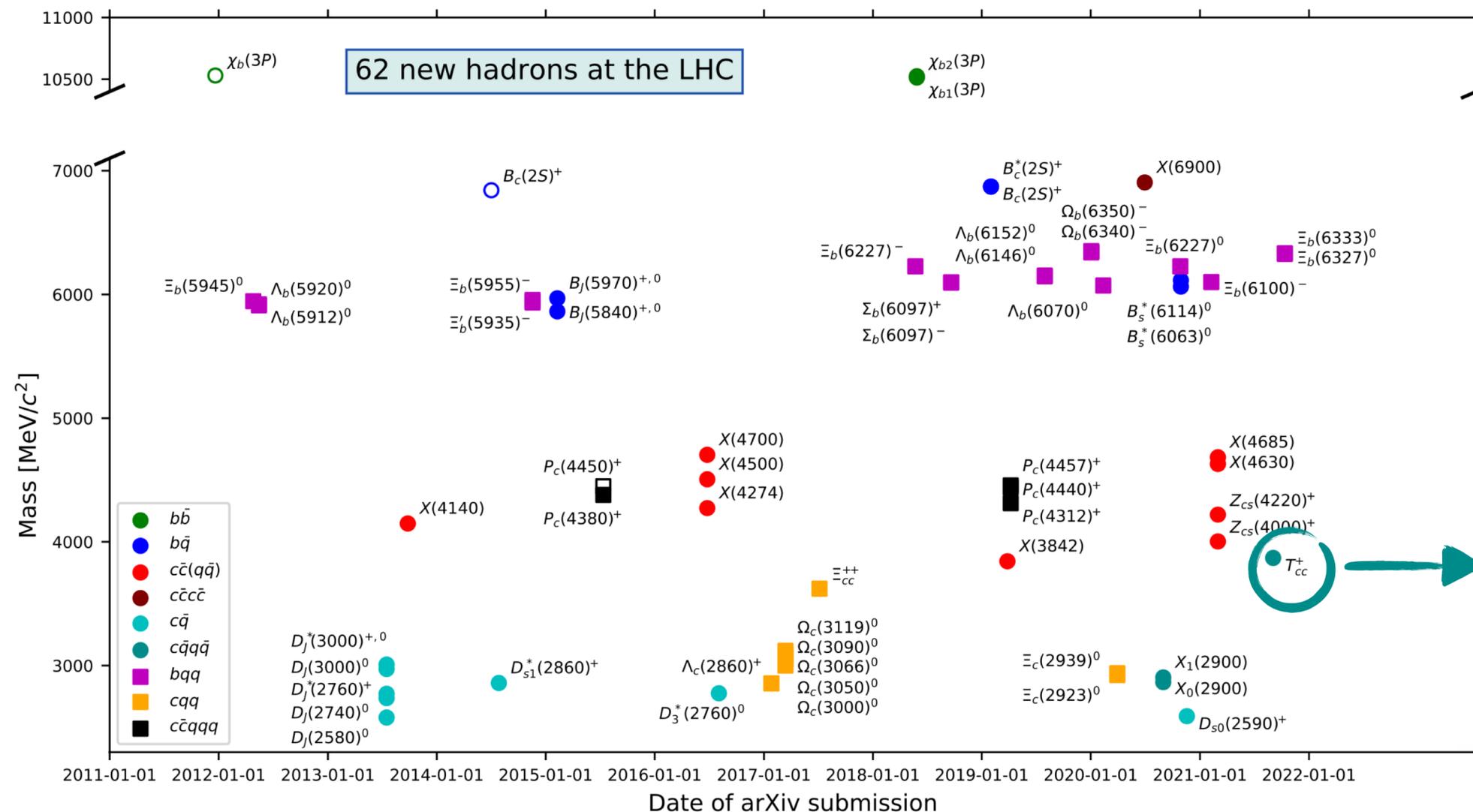


[I. Danilkin, talk@INT, March 2023]

(+ Babar, Belle, COMPASS)

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Multi-hadron interactions

Beyond QCD, understanding resonance properties is important for new physics searches

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► Tests of the Standard Model in meson weak decays

CP violation in $K \rightarrow \pi\pi$ weak decays

$$(\varepsilon'/\varepsilon)_{\text{exp}} = (16.6 \pm 2.3) \times 10^{-4}$$

[NA48 & KTeV, 2002 & 2009]

σ resonance

CP violation in $D^0 \rightarrow K^+K^-/\pi^+\pi^-$ decays

$$\Delta a_{CP}^{\text{dir}} = (-15.7 \pm 2.9) \times 10^{-4}$$

[LHCb, 2019]

$f_0(1710)$
resonance

Multi-hadron interactions

Beyond QCD, understanding resonance properties is important for new physics searches

► Tests of the Standard Model in meson weak decays

► Neutrino-nucleus scattering (DUNE, Hyper-K)

$$\nu N \rightarrow \ell N \pi$$

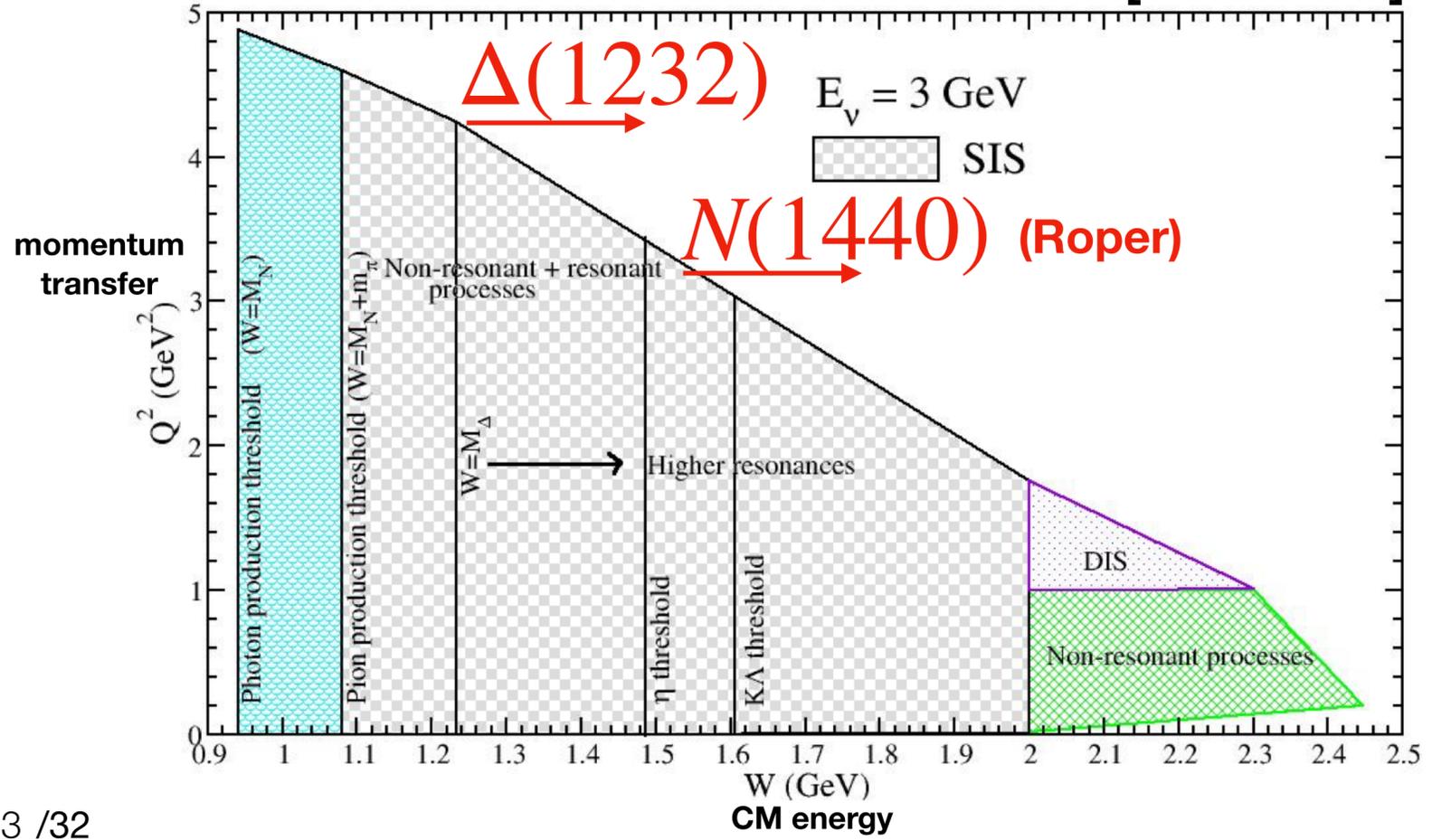
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[2203.09030]

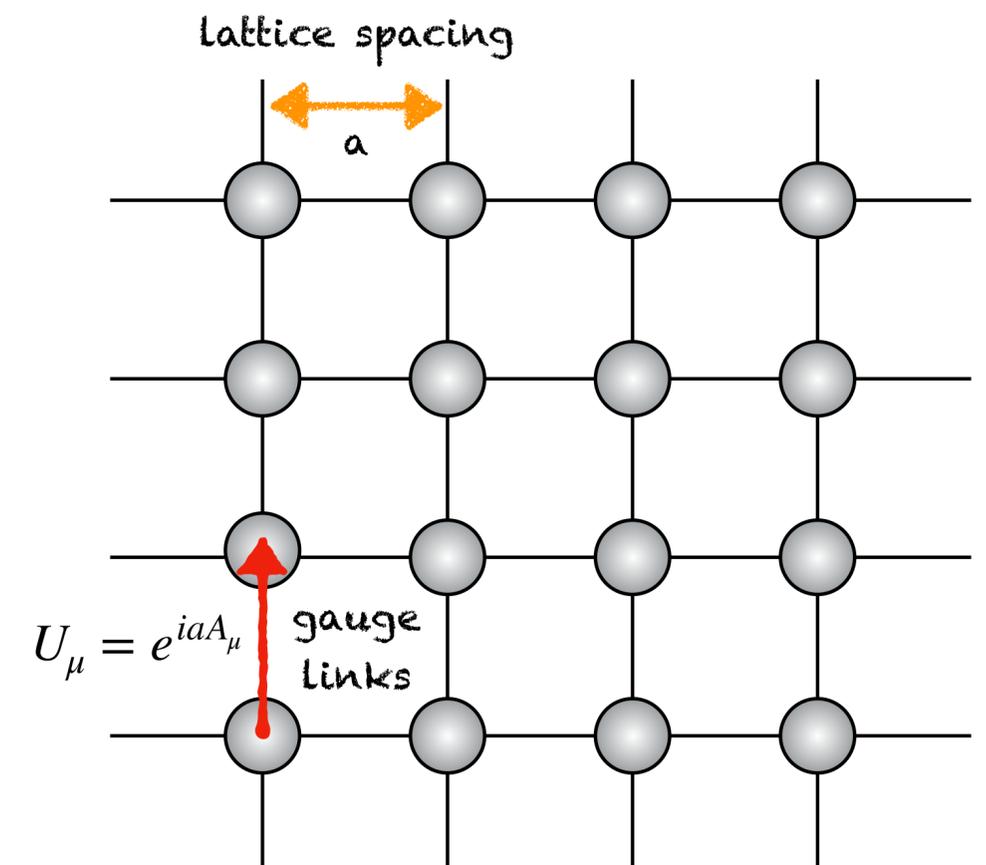


Lattice QCD

- Lattice QCD is a first-principles numerical approach to the strong interaction

$$\langle \mathcal{O}(t)\mathcal{O}(0) \rangle = \frac{1}{\mathcal{Z}} \int D\psi D\bar{\psi} DA \mathcal{O}(t)\mathcal{O}(0) e^{-S_E(\psi, \bar{\psi}, A_\mu)}$$

Euclidean action



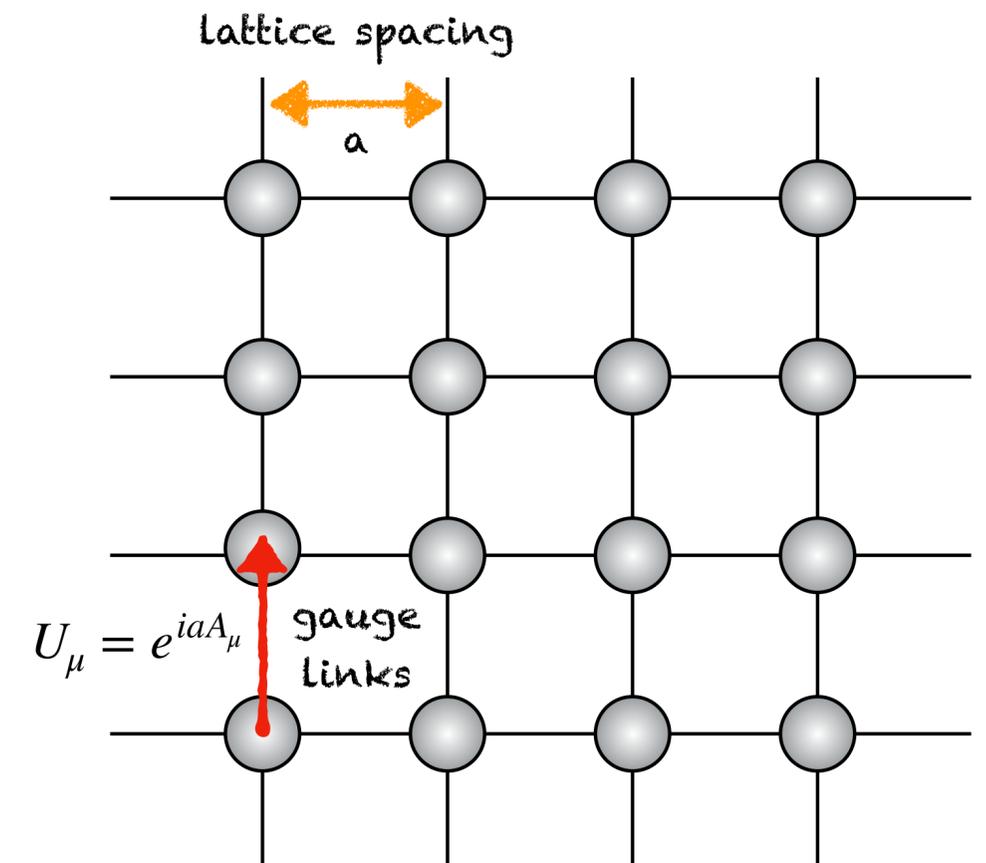
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Euclidean action

Can we obtain resonance properties from Euclidean correlation functions?



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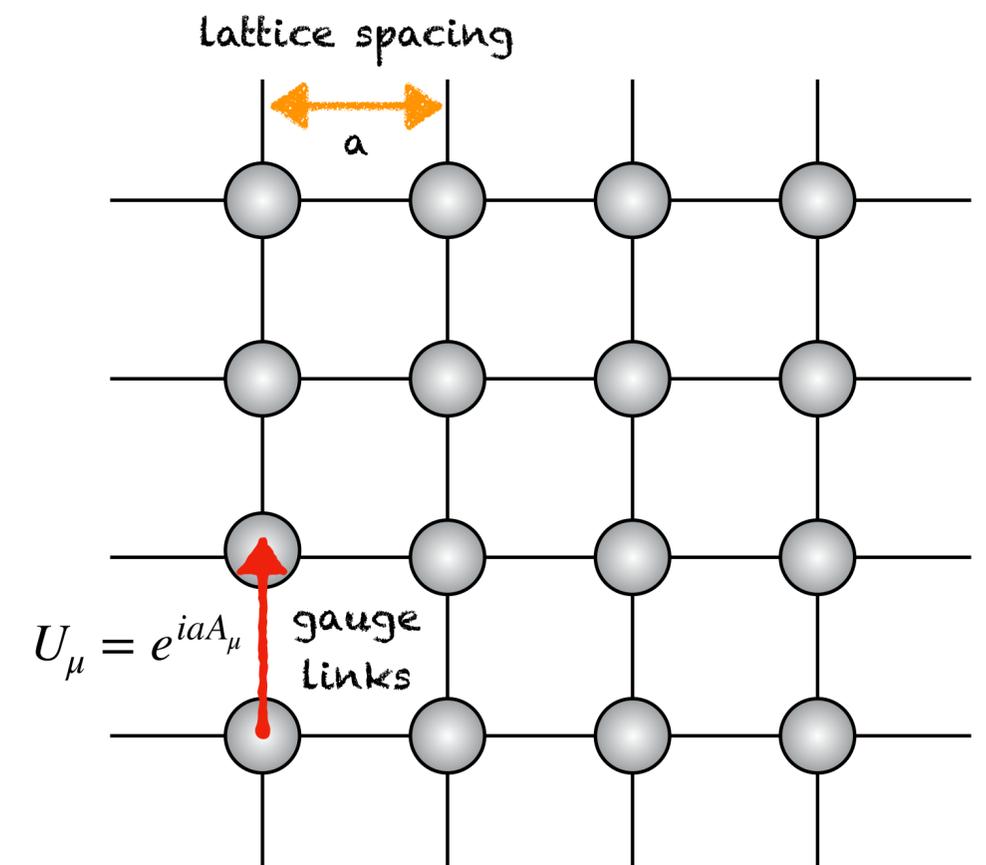
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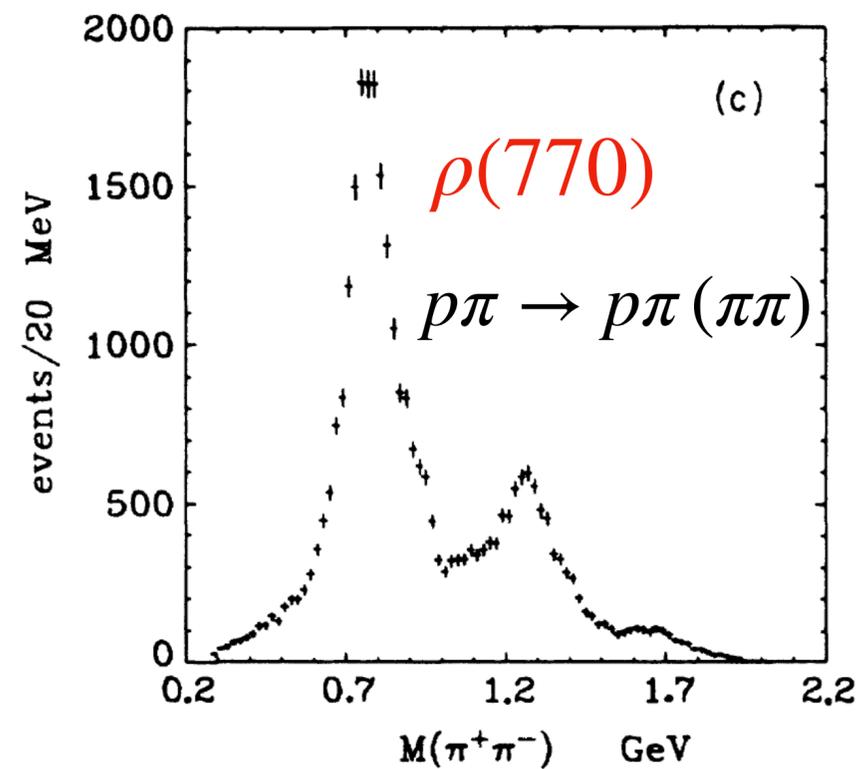
Yes, but not that simple!



Infinite vs finite volume

Experiments

- Asymptotic states
- Direct access to scattering amplitudes

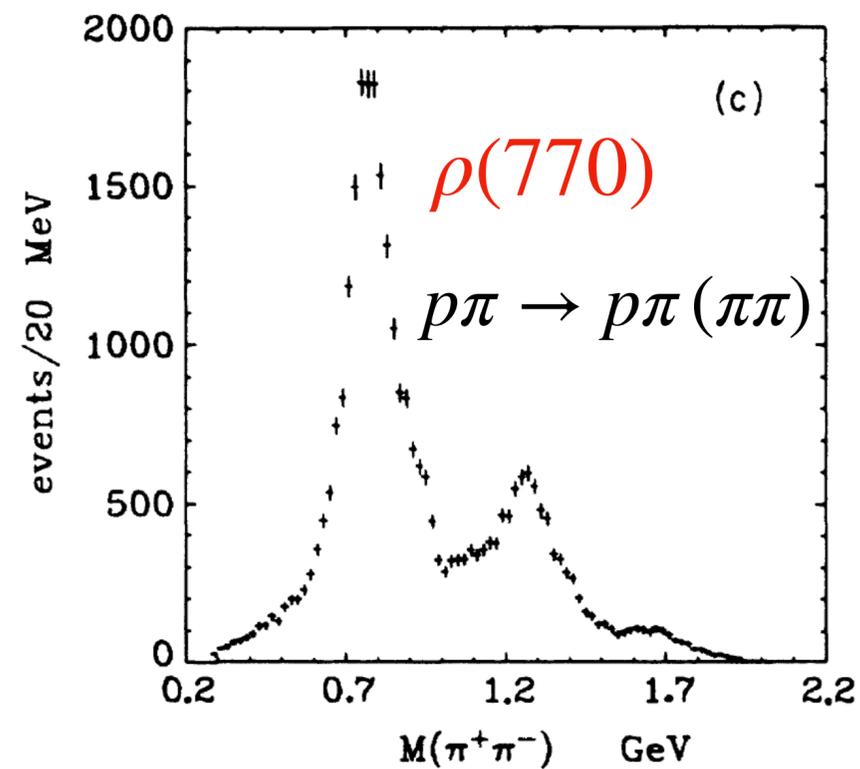


[Protopescu et al, PRD7 1973]

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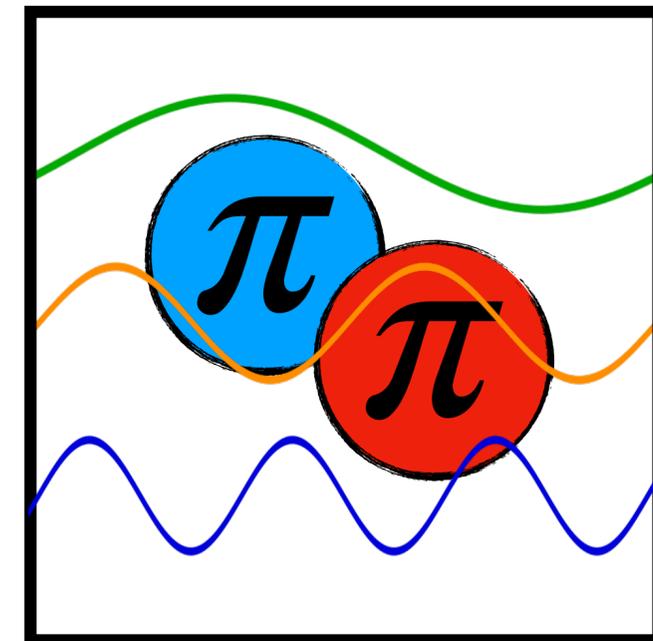
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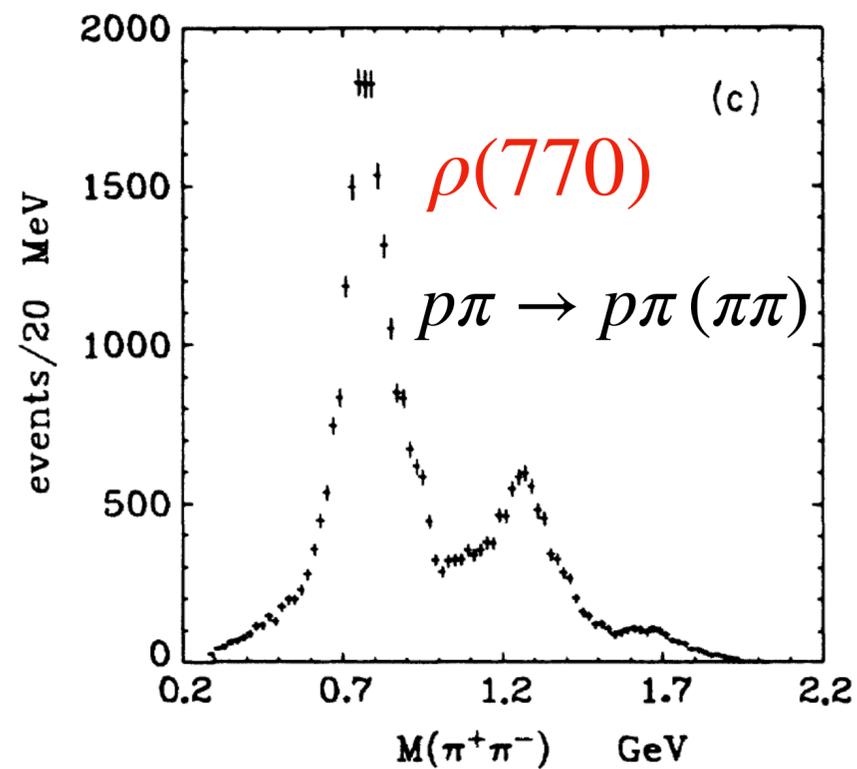
- Euclidean time
- Stationary states in a box



Infinite vs finite volume

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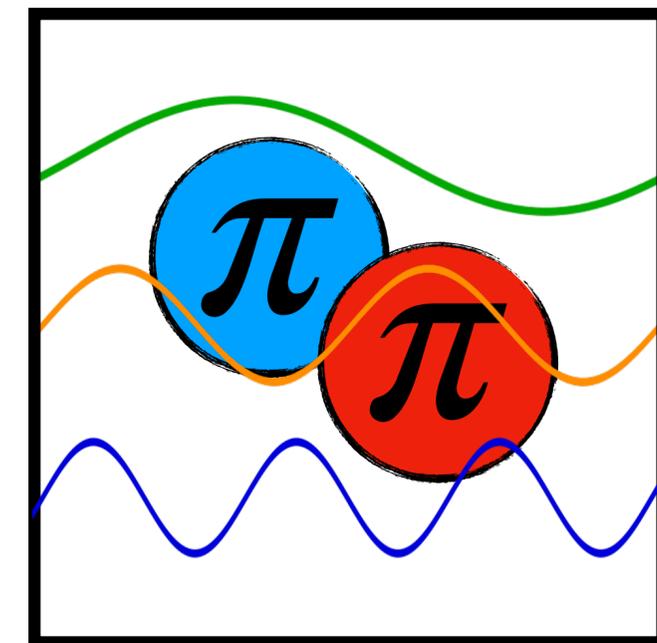
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Lattice QCD

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Towards the QCD S-Matrix

The S-Matrix contains the physical information of the theory:

$$S_{ab}(E) \equiv \langle \text{out} | \mathcal{S} | \text{in} \rangle$$

Lattice QCD \longrightarrow QCD S-matrix

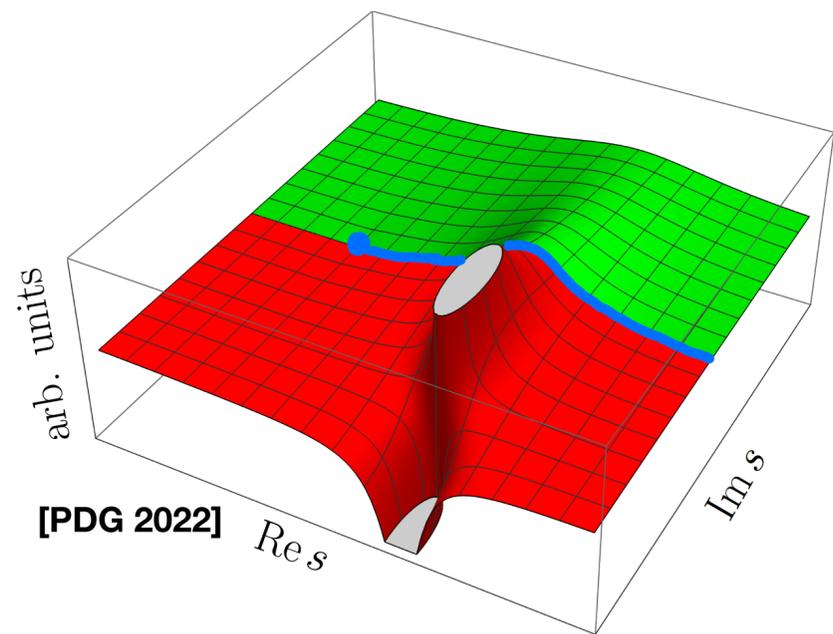
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► Resonances as poles in the S-matrix (or scattering amplitude)



$$\sim \frac{g}{E^2 - E_R^2}$$

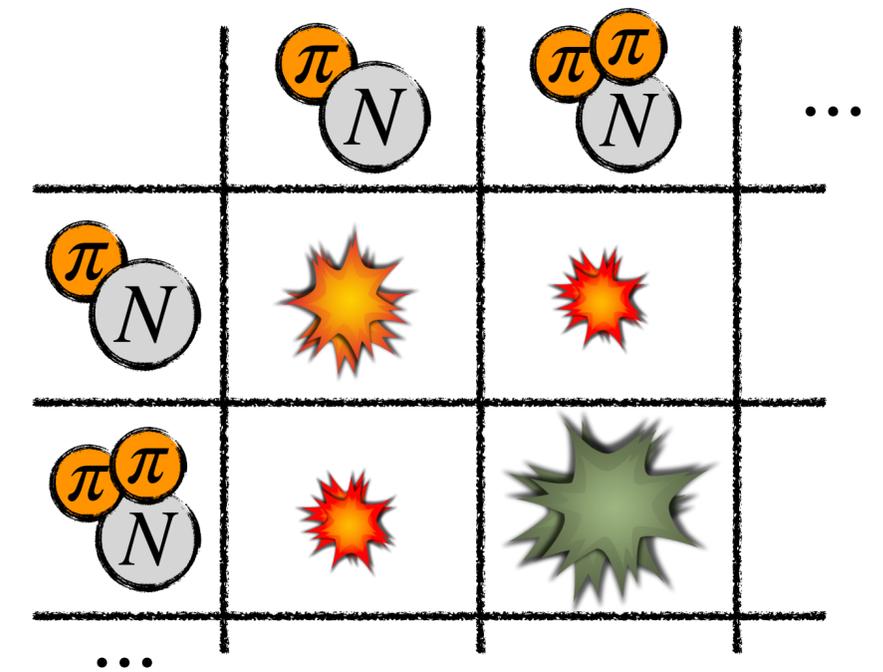
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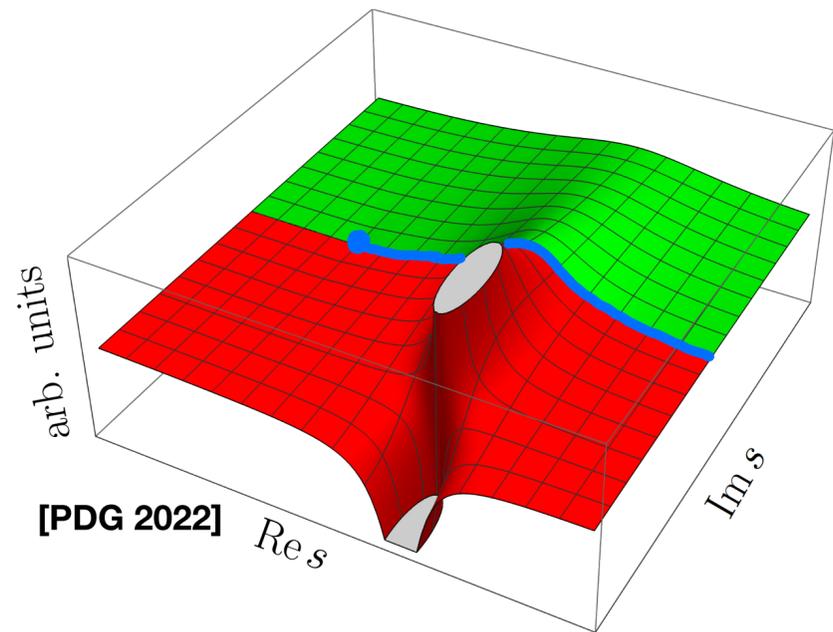
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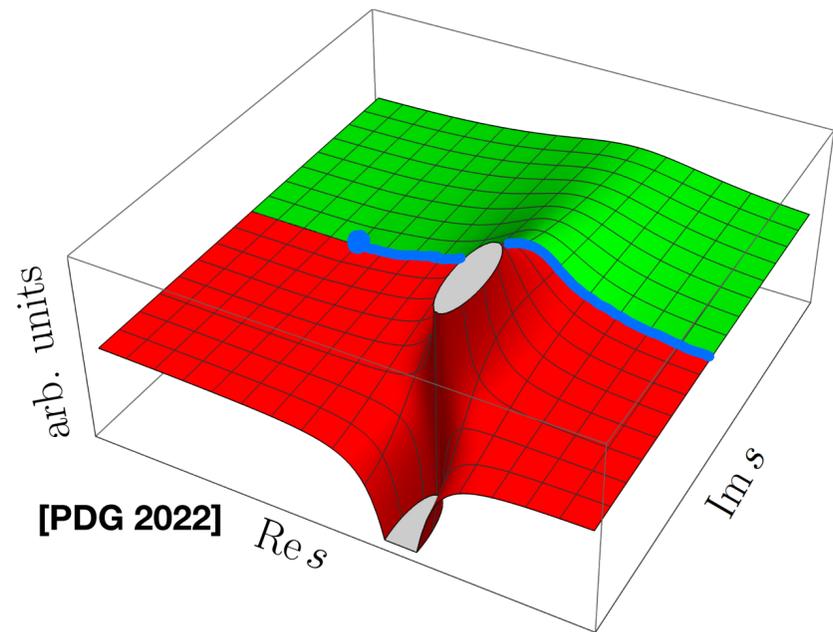
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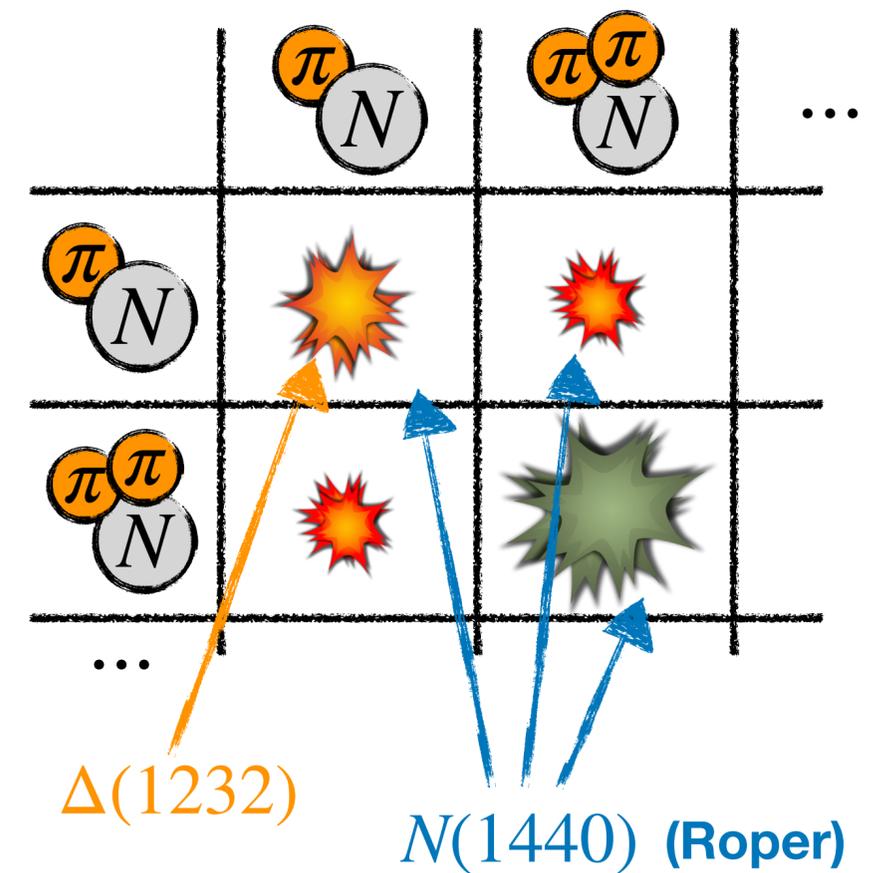
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Outline

1. Scattering amplitudes from lattice QCD
2. Meson-Baryon scattering: $\Delta(1232)$ and $\Lambda(1405)$
3. Three-particle systems

Scattering amplitudes from lattice QCD

Energy levels from Lattice QCD

- The **energy levels** of the theory are measured from Euclidean correlation functions

$$C(t) = \langle \mathcal{O}(t)\mathcal{O}(0) \rangle = \sum_n \left| \langle 0 | \mathcal{O}(0) | n \rangle \right|^2 e^{-E_n t}$$

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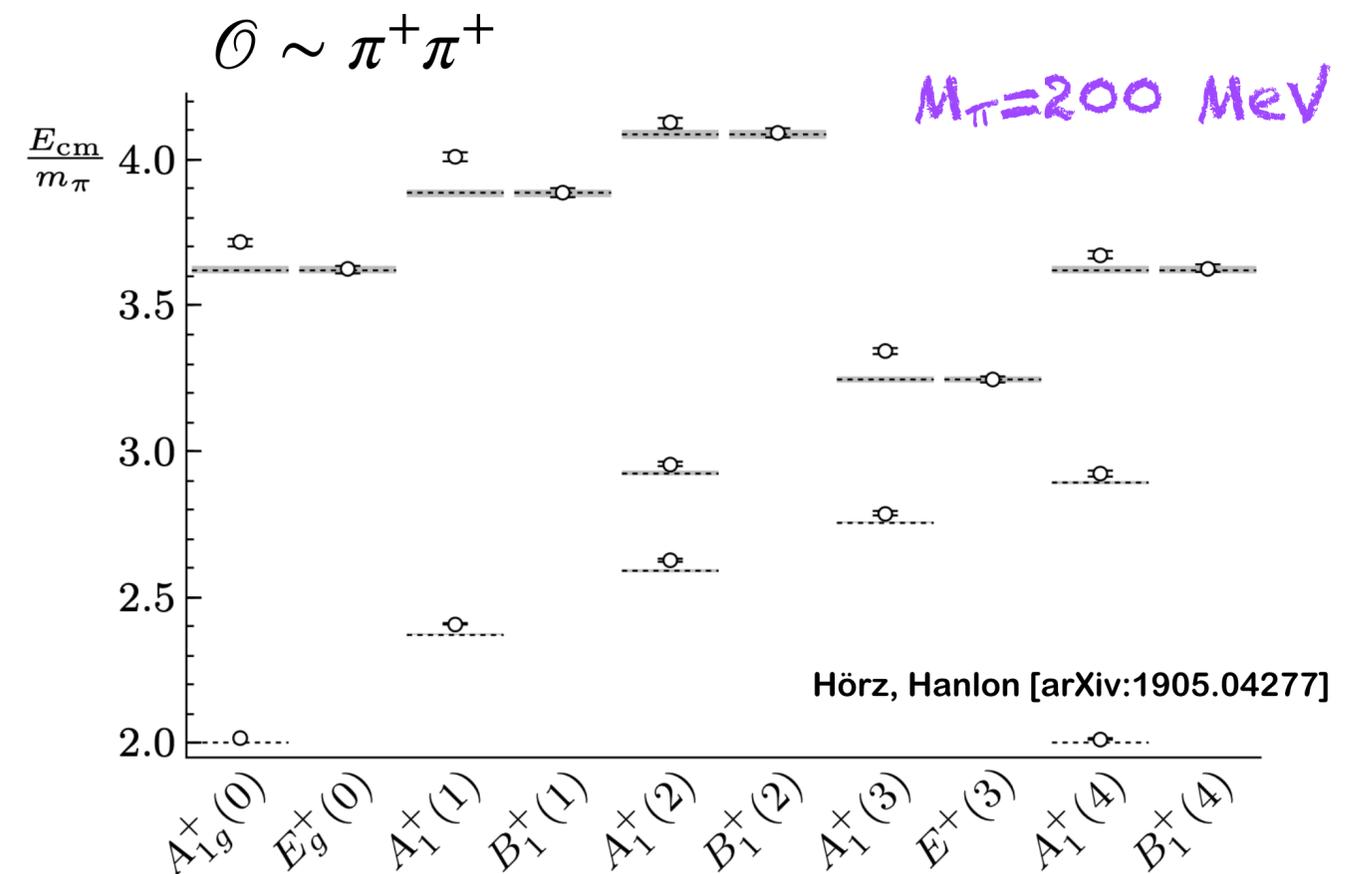
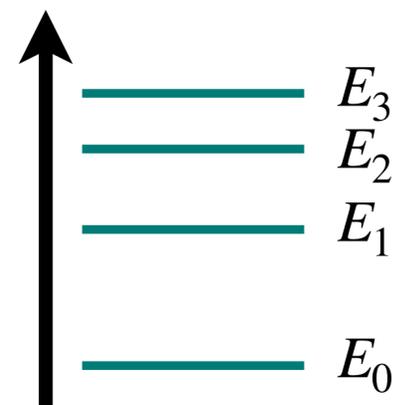
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- Multiple operators with the same quantum names to obtain several energy levels

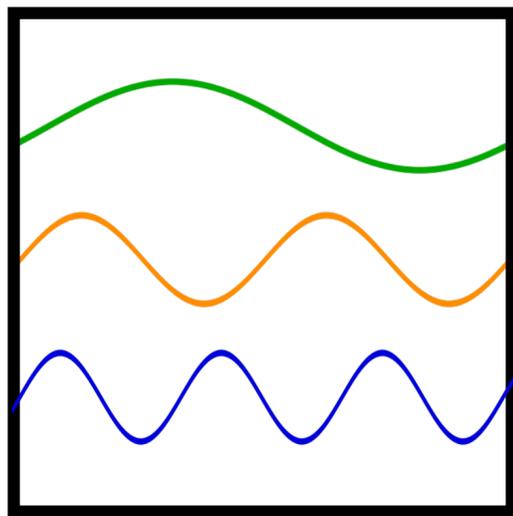
► Variational techniques
(Generalized EigenValue Problem, GEVP)

The Spectrum



Finite-Volume Spectrum

Free scalar particles in finite volume
with periodic boundaries

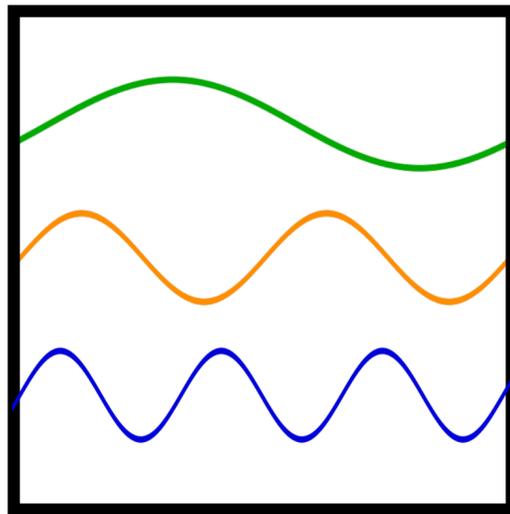


$$\vec{p} = \frac{2\pi}{L}(n_x, n_y, n_z)$$

Two particles: $E = 2\sqrt{m^2 + \frac{4\pi^2}{L^2}\vec{n}^2}$

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Interactions change the spectrum:
it can be treated as a perturbation

Ground state to leading order

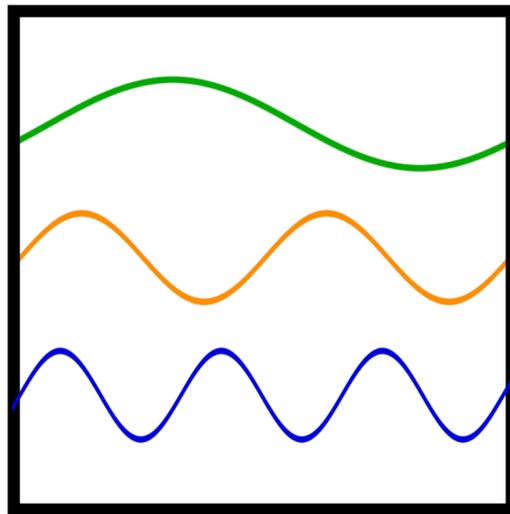
$$E_2 - 2m = \langle \phi(\vec{0})\phi(\vec{0}) | \mathbf{H}_I | \phi(\vec{0})\phi(\vec{0}) \rangle$$

$$\Delta E_2 = \frac{\mathcal{M}_2(E = 2m)}{8m^2L^3} + O(L^{-4})$$

[Huang, Yang, 1958]

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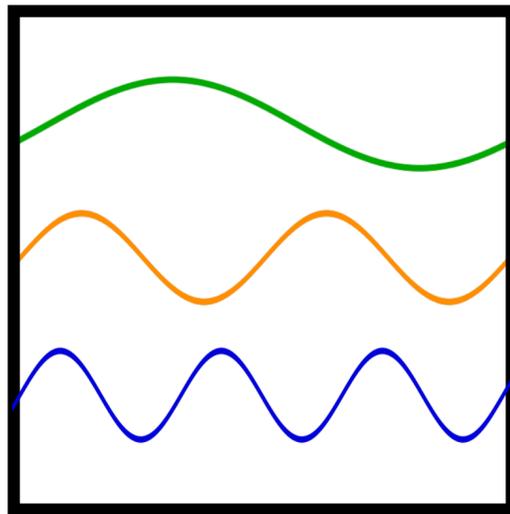
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The **energy shift** of the two-particle ground state is related to the $2 \rightarrow 2$ **scattering amplitude**

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In general a problem of Quantum Field Theory in finite volume

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Finite-Volume formalism

- Indirect connection between the spectrum and the two-particle scattering amplitude [Lüscher 89']

Two-particle Quantization Condition

$$\det_{\ell m} \left[\mathcal{K}_2(E) + F^{-1}(E, \vec{P}, L) \right] \Big|_{E=E_n} = 0$$

Scattering K-Matrix Known kinematic function

"QC2"

! Note: only valid for two particles below inelastic thresholds.

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K-matrix parametrized
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Finite-volume information

$$F_{00}(q^2) \sim \left[\frac{1}{L^3} \sum_{\vec{k}} - \int \frac{d^3k}{(2\pi)^3} \right] \frac{1}{k^2 - q^2}$$

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Example $I=2$ $\pi\pi$ scattering

Two pions in s-wave

$$\mathcal{K}_2^{s\text{-wave}}(E_n) = \frac{-1}{F_{00}(E_n, \vec{P}, L)}$$

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energy
level
 E_n



phase
shift
 $\delta_0(E_n)$

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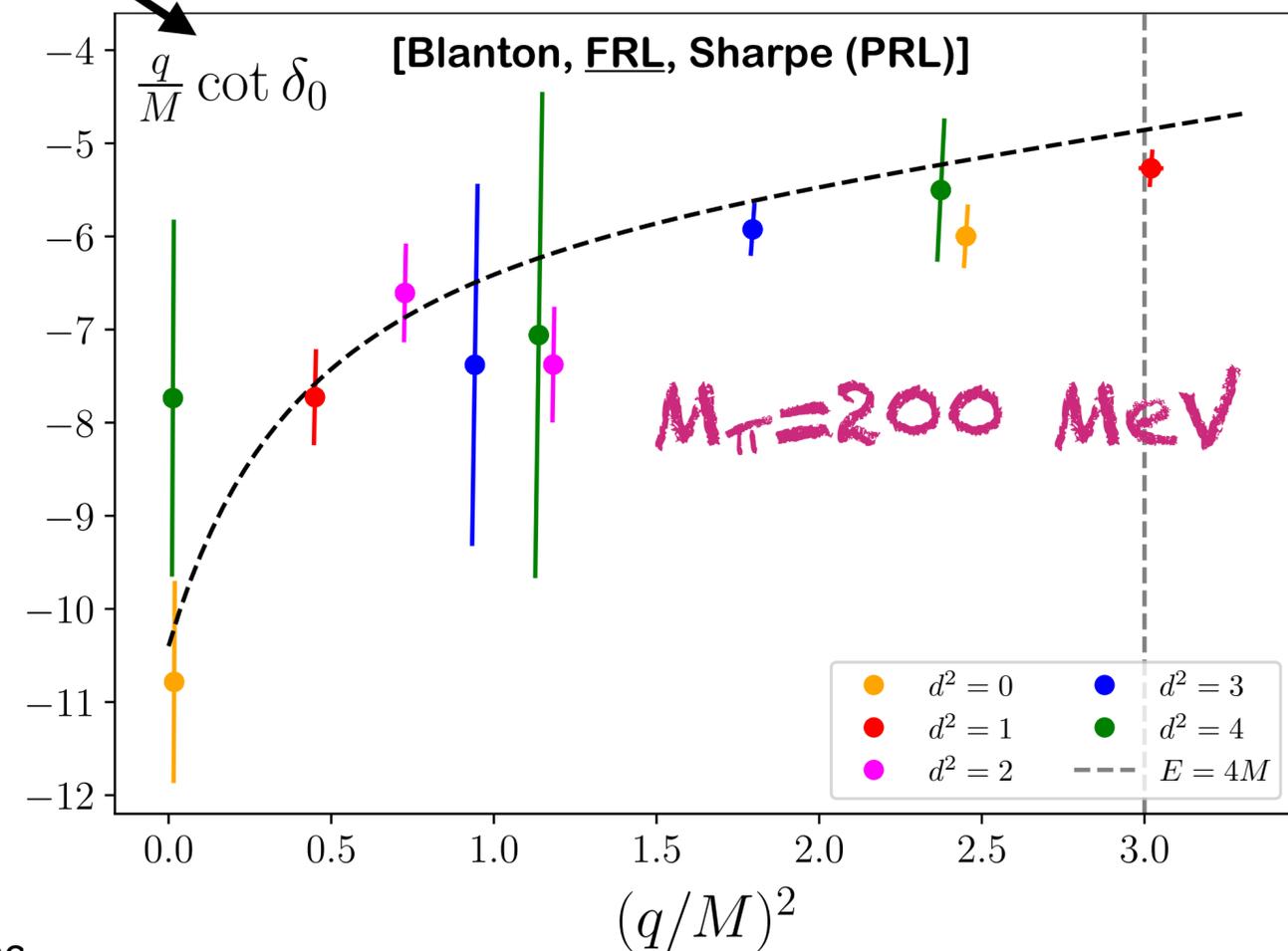
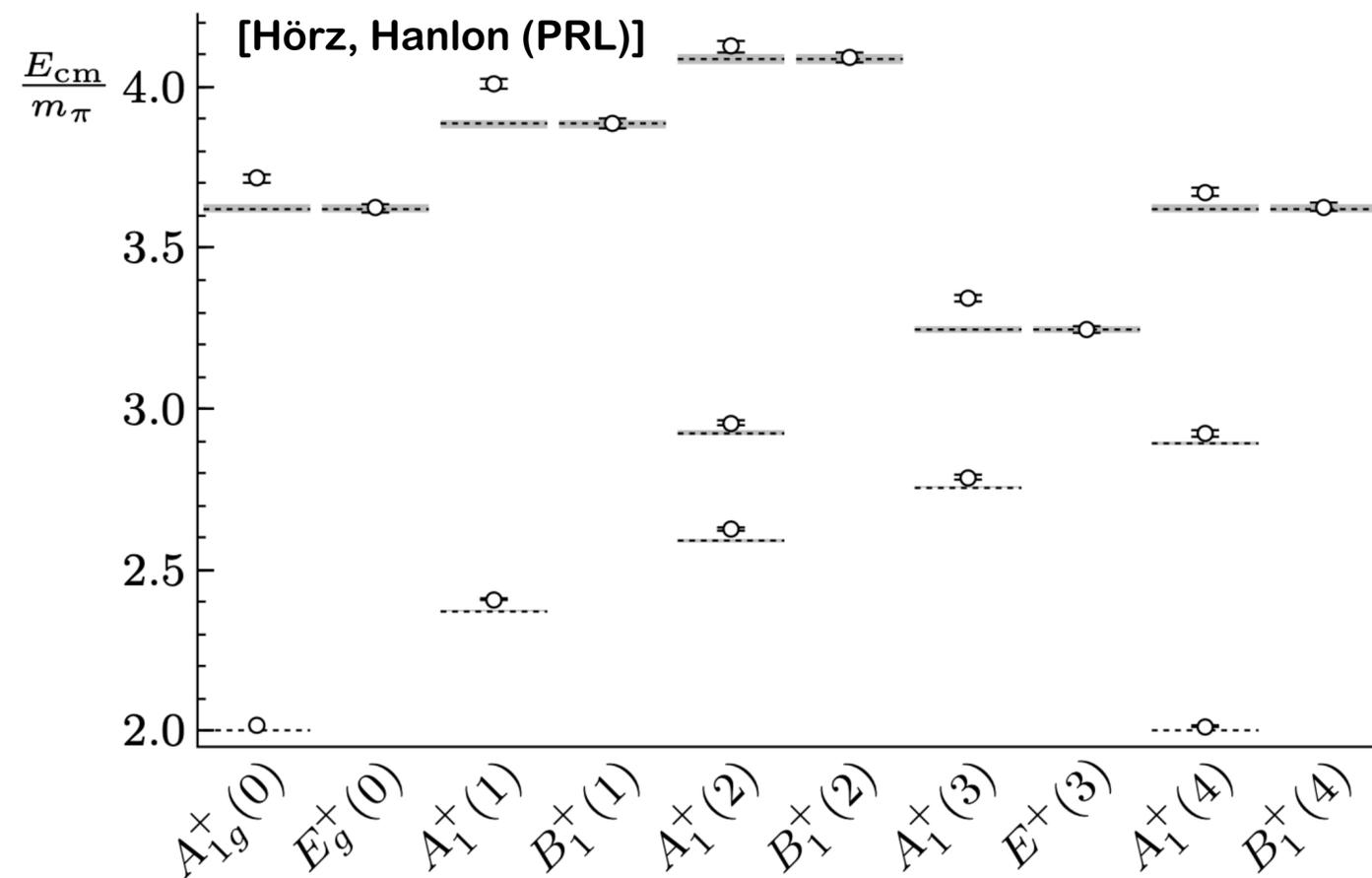
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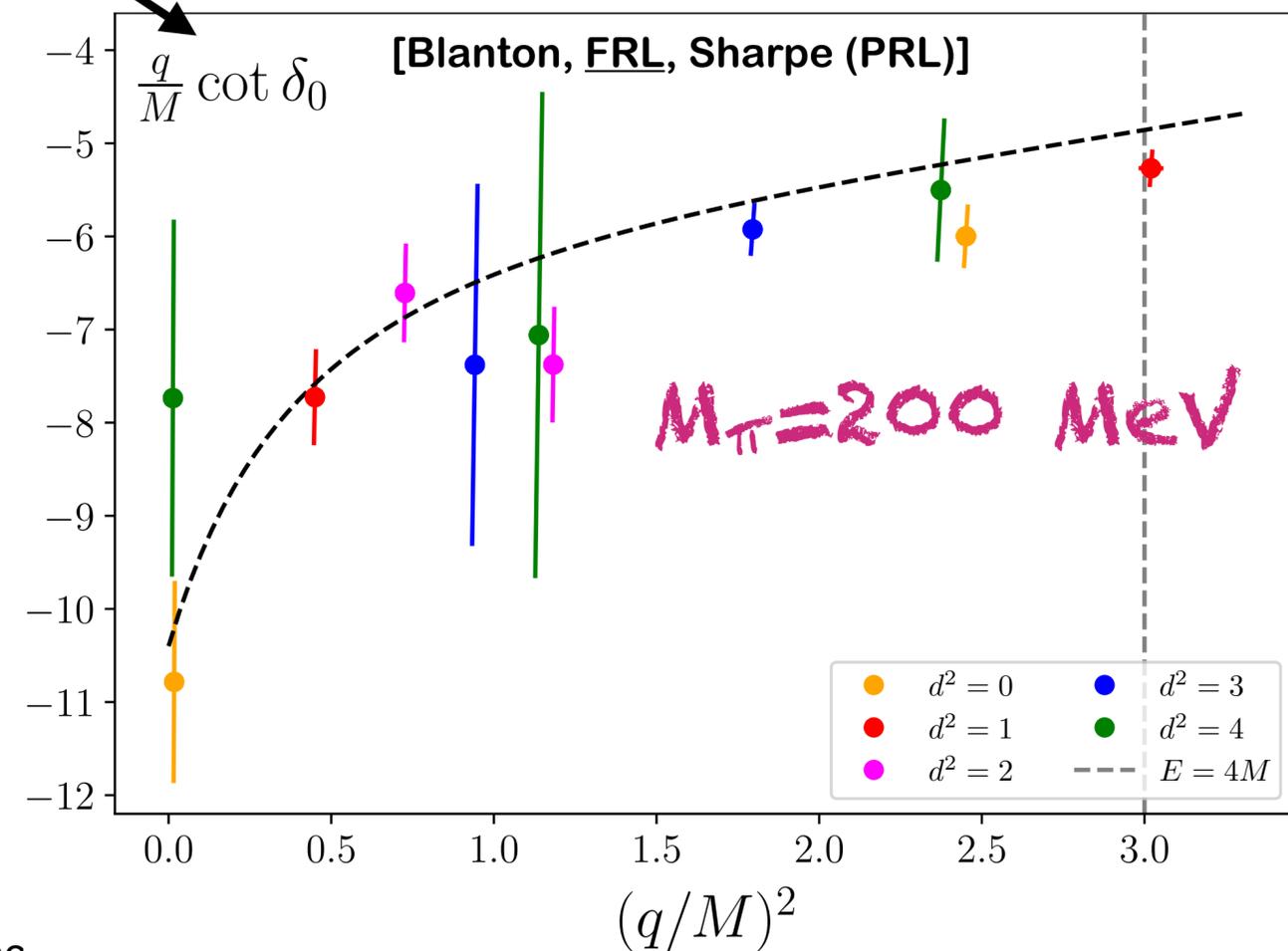
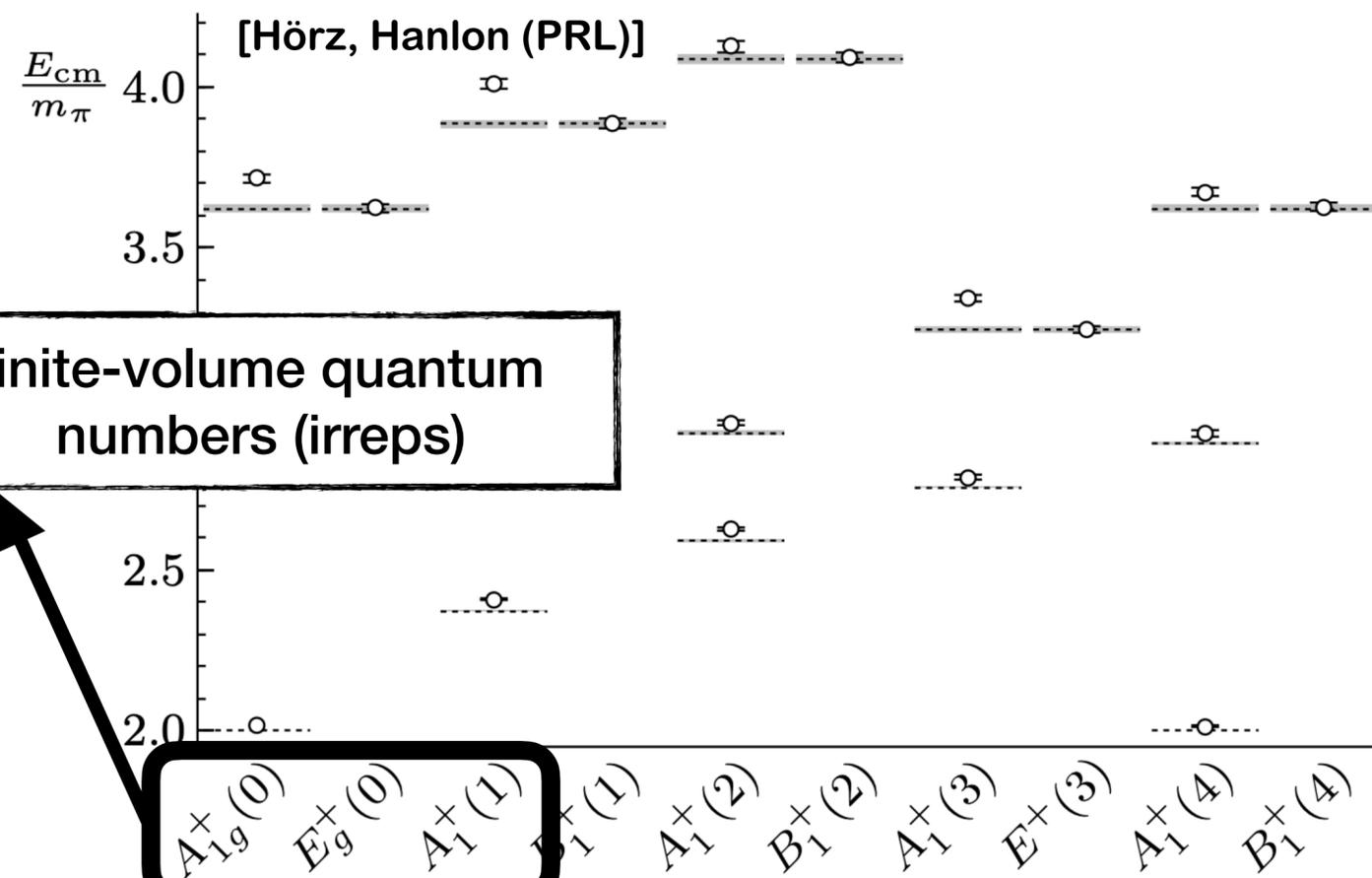
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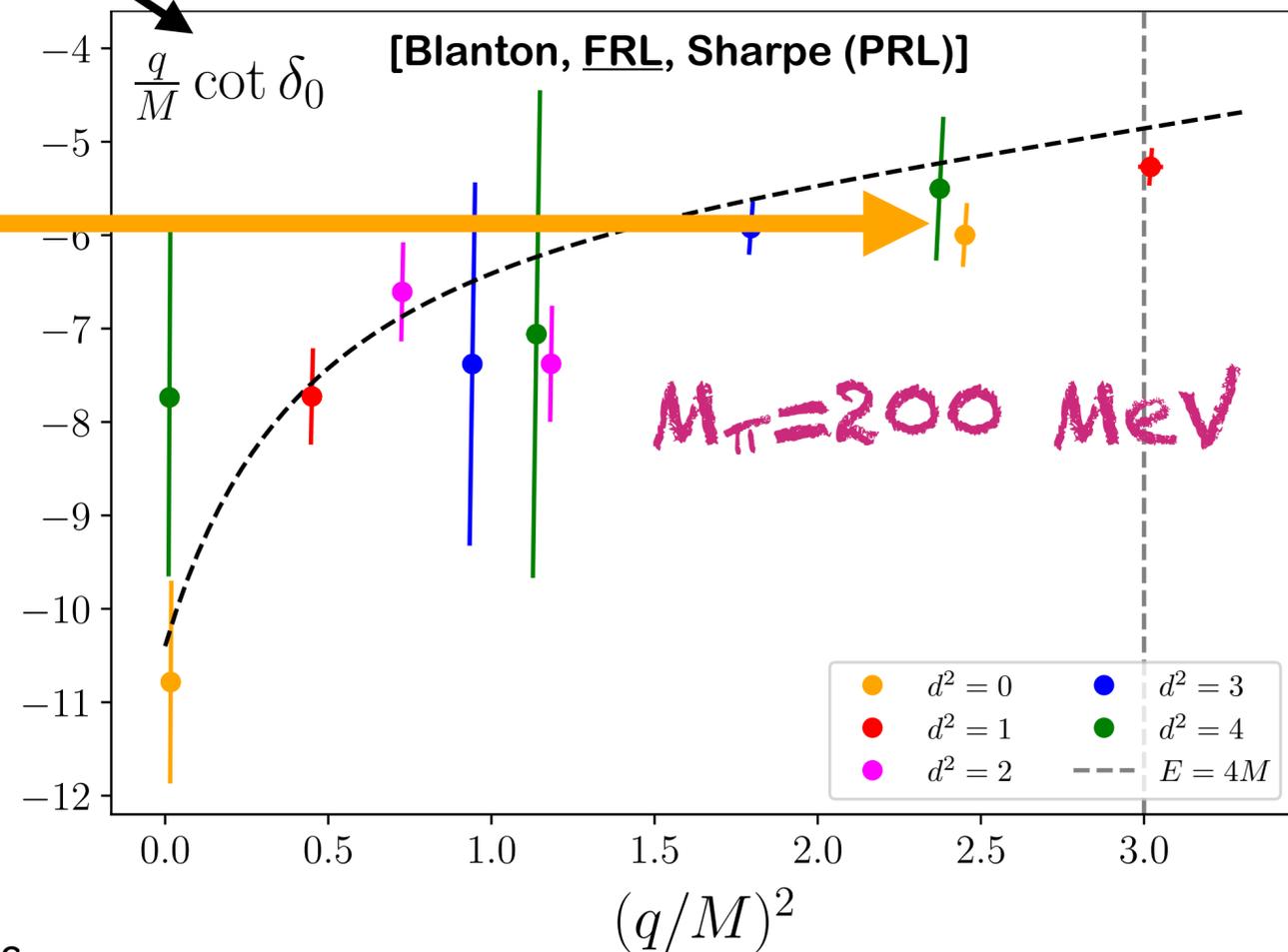
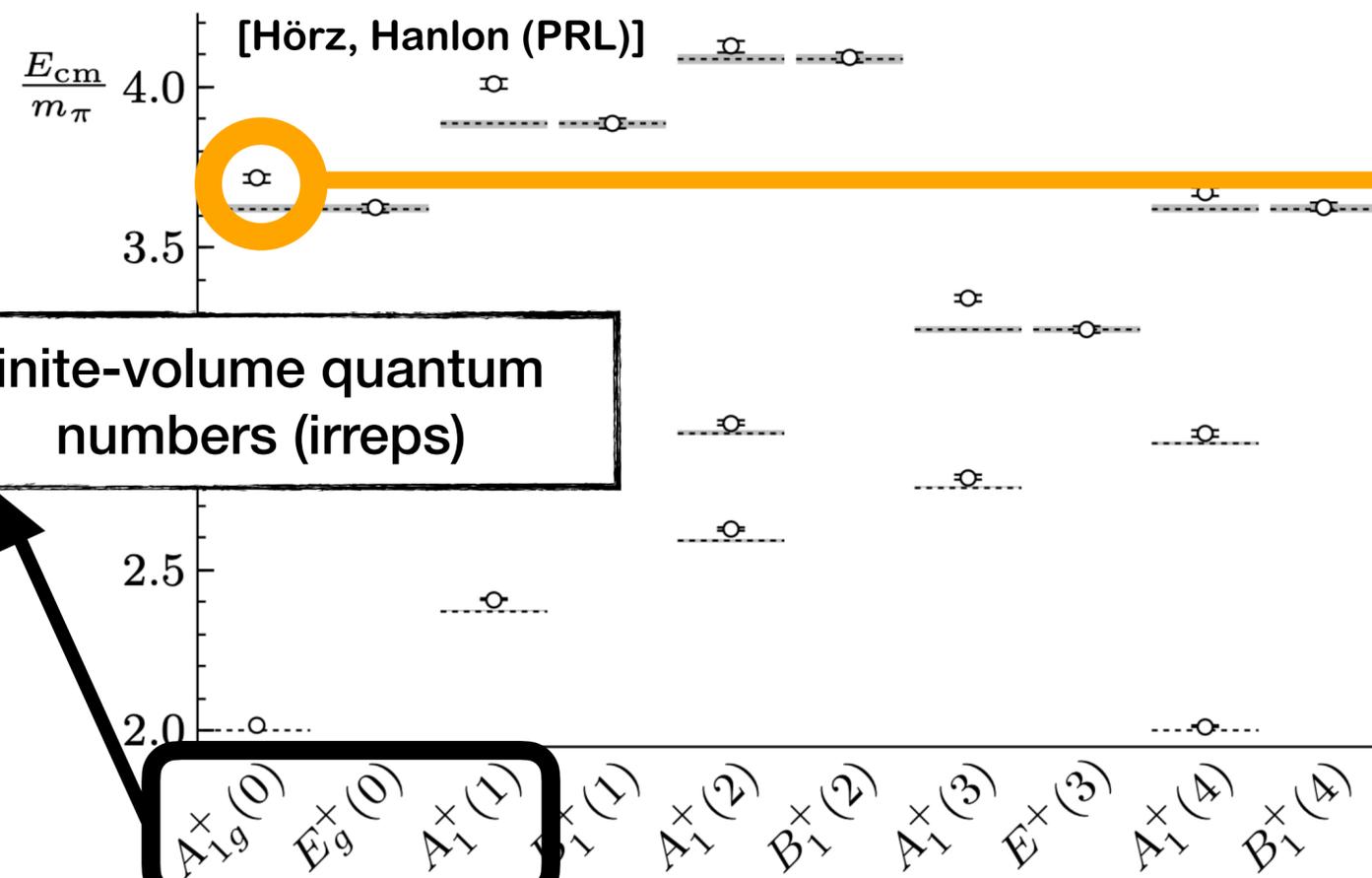
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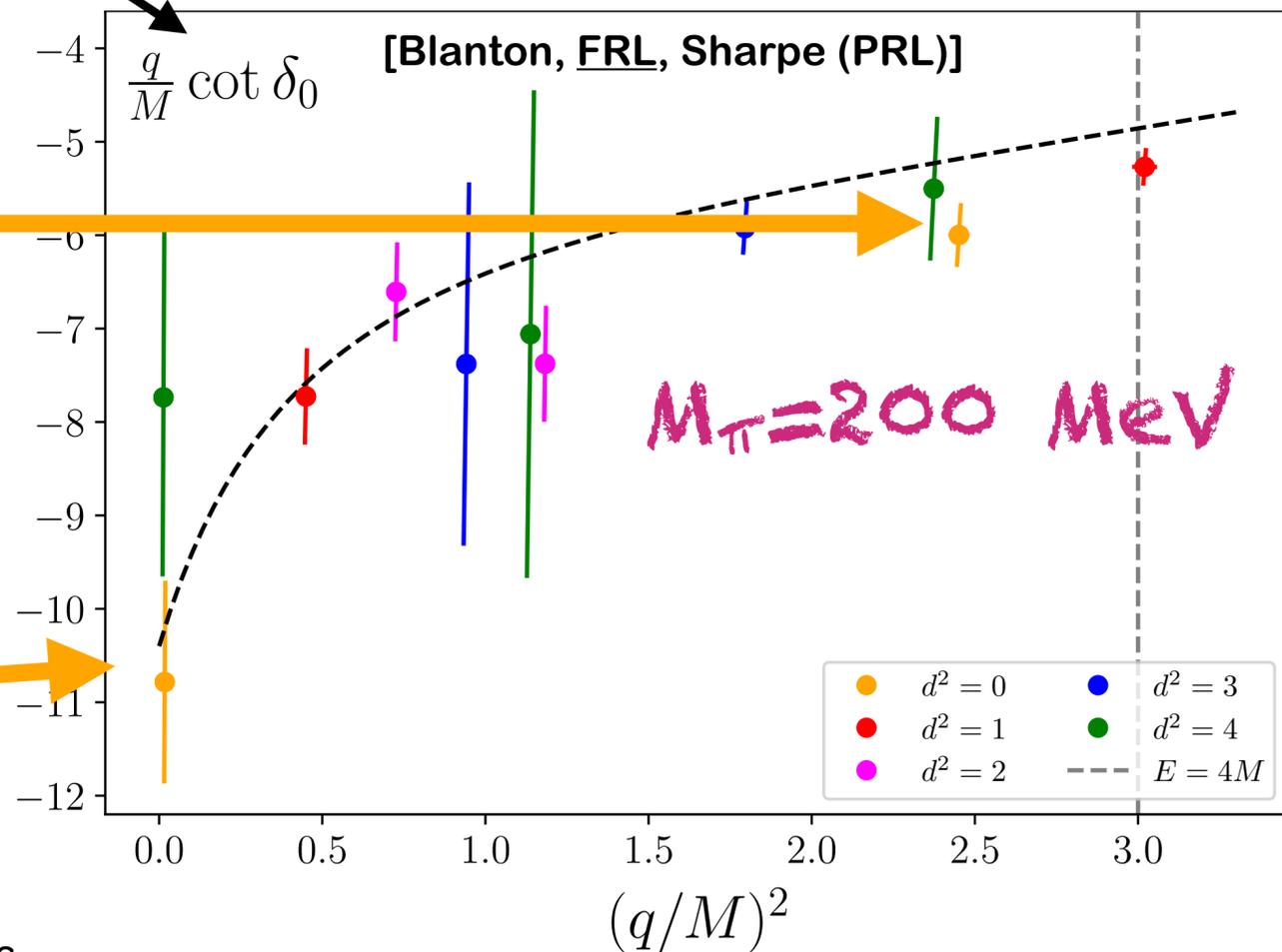
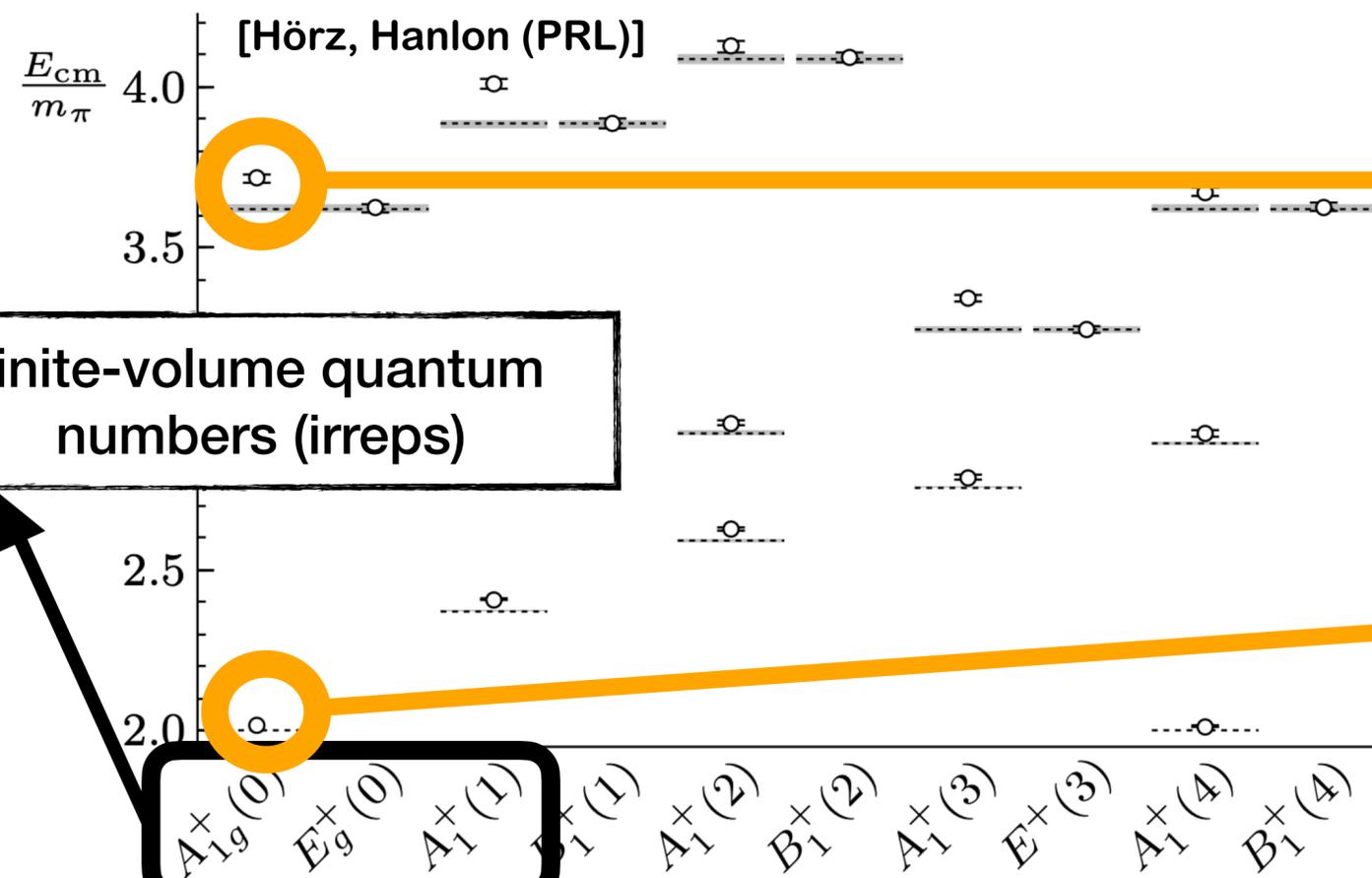
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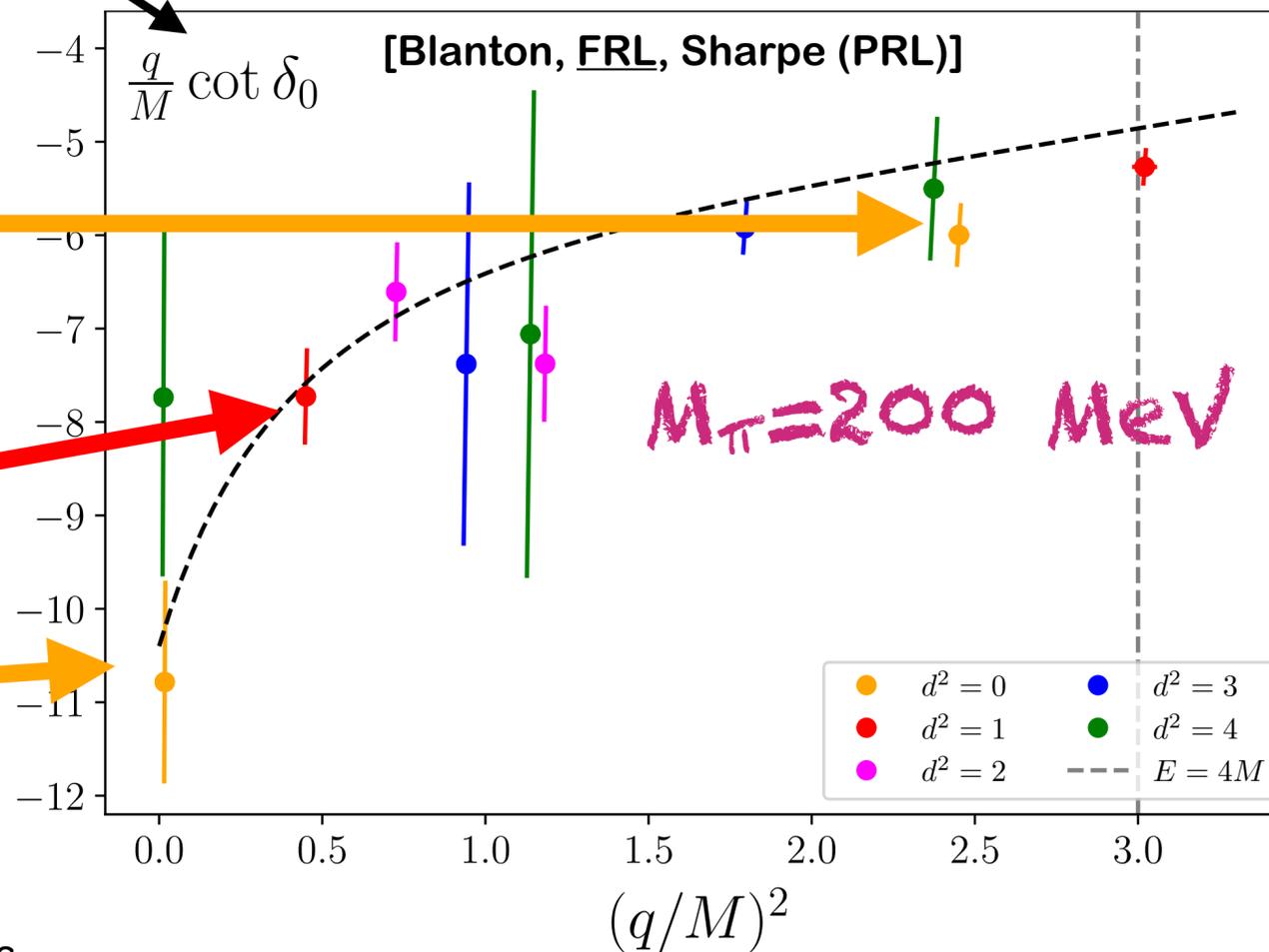
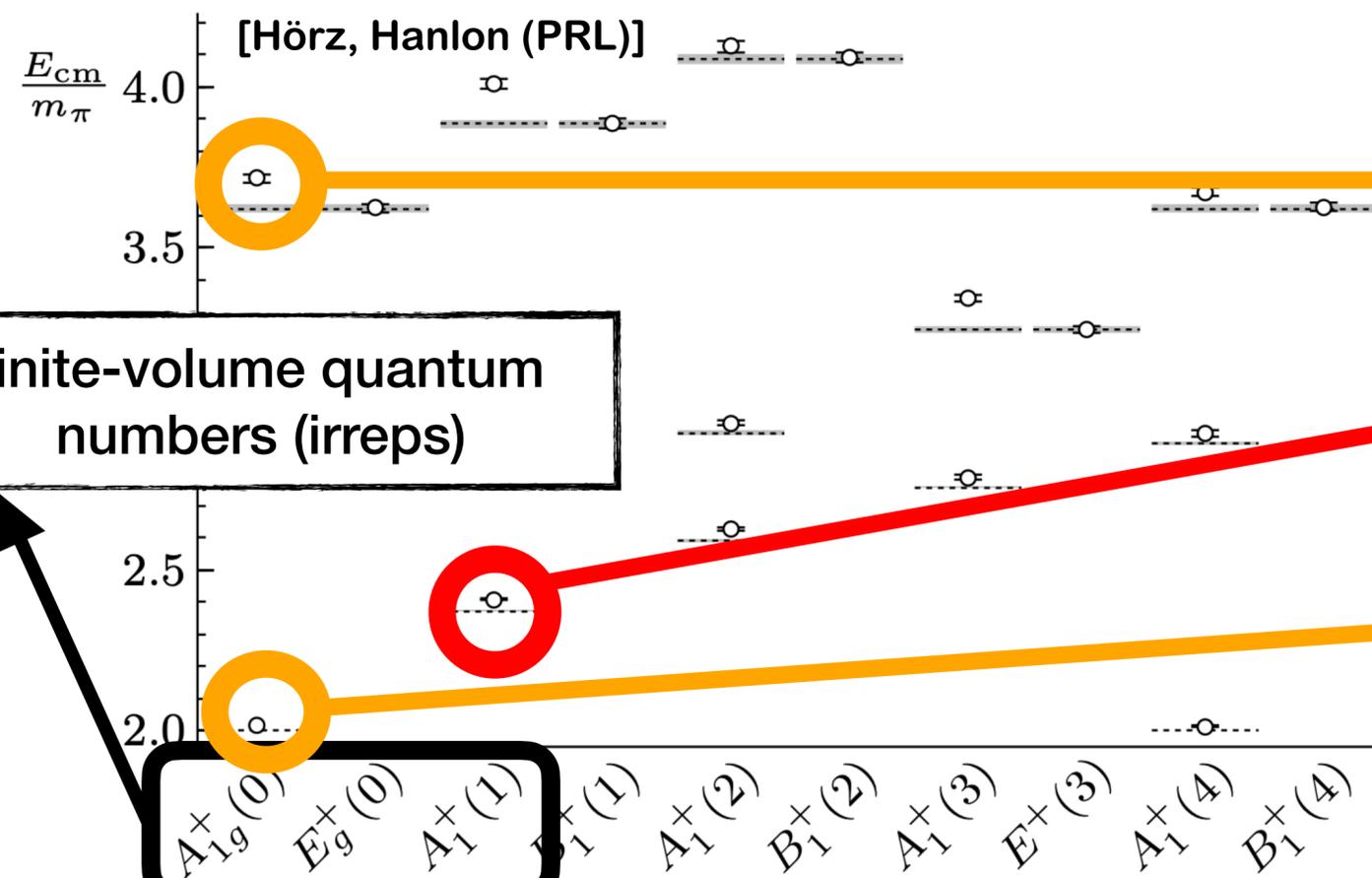
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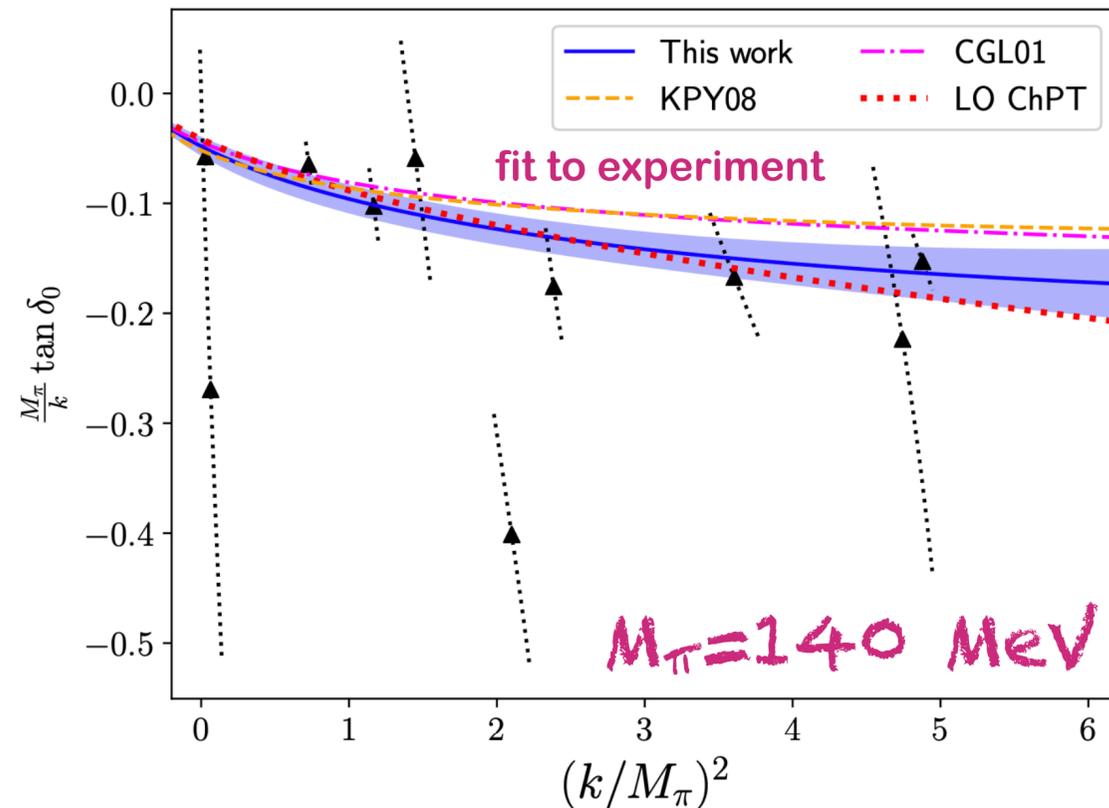
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Towards the physical point

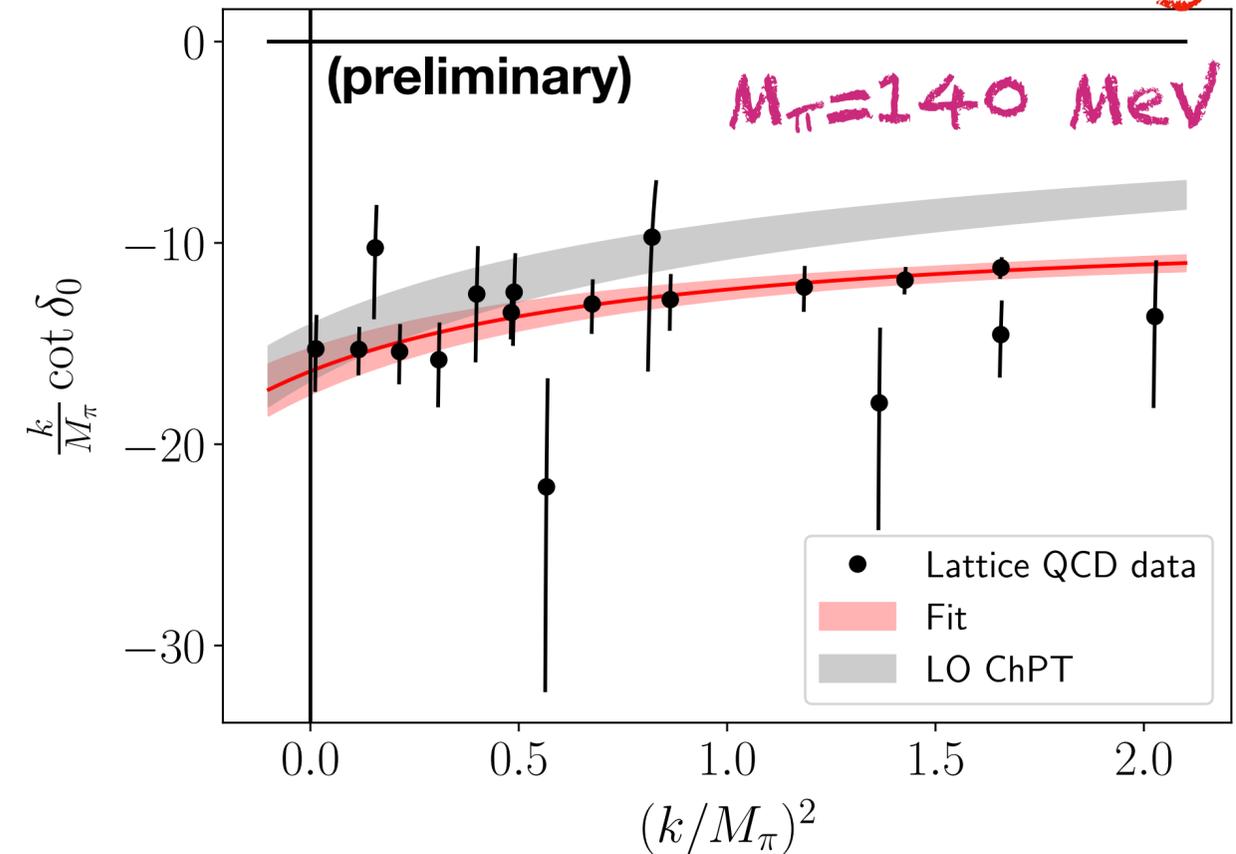
- Some systems already being studied at the **physical point!**

I=2 $\pi\pi$ scattering



[Fischer, Kostrzewa, Liu, FRL, Ueding, Urbach (ETMC), EPJC 2021]

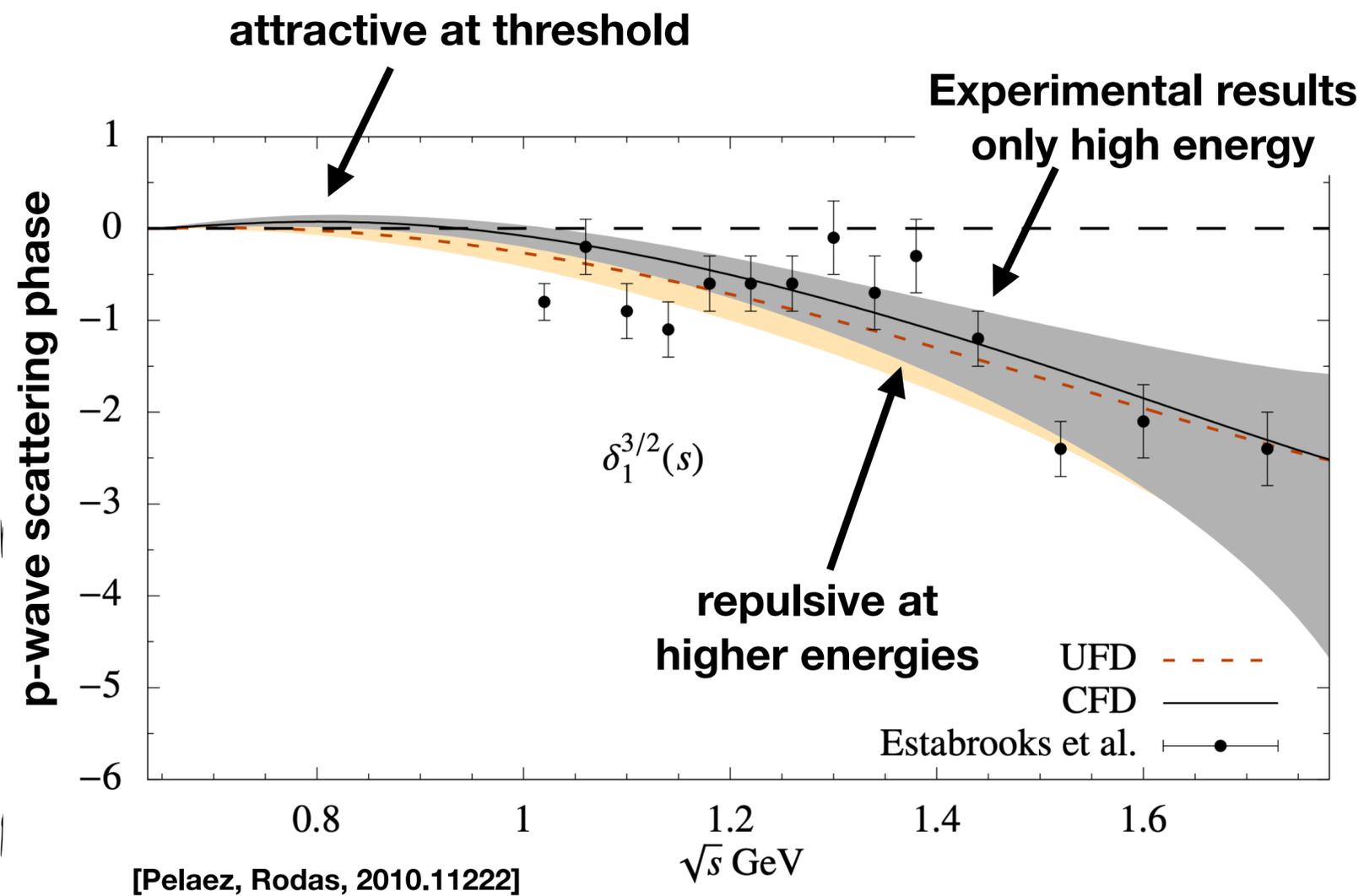
I=3/2 πK scattering



[Draper, Hanlon, Hörz, Morningstar, FRL, Sharpe (in preparation)]

p-wave π^+K^+ scattering

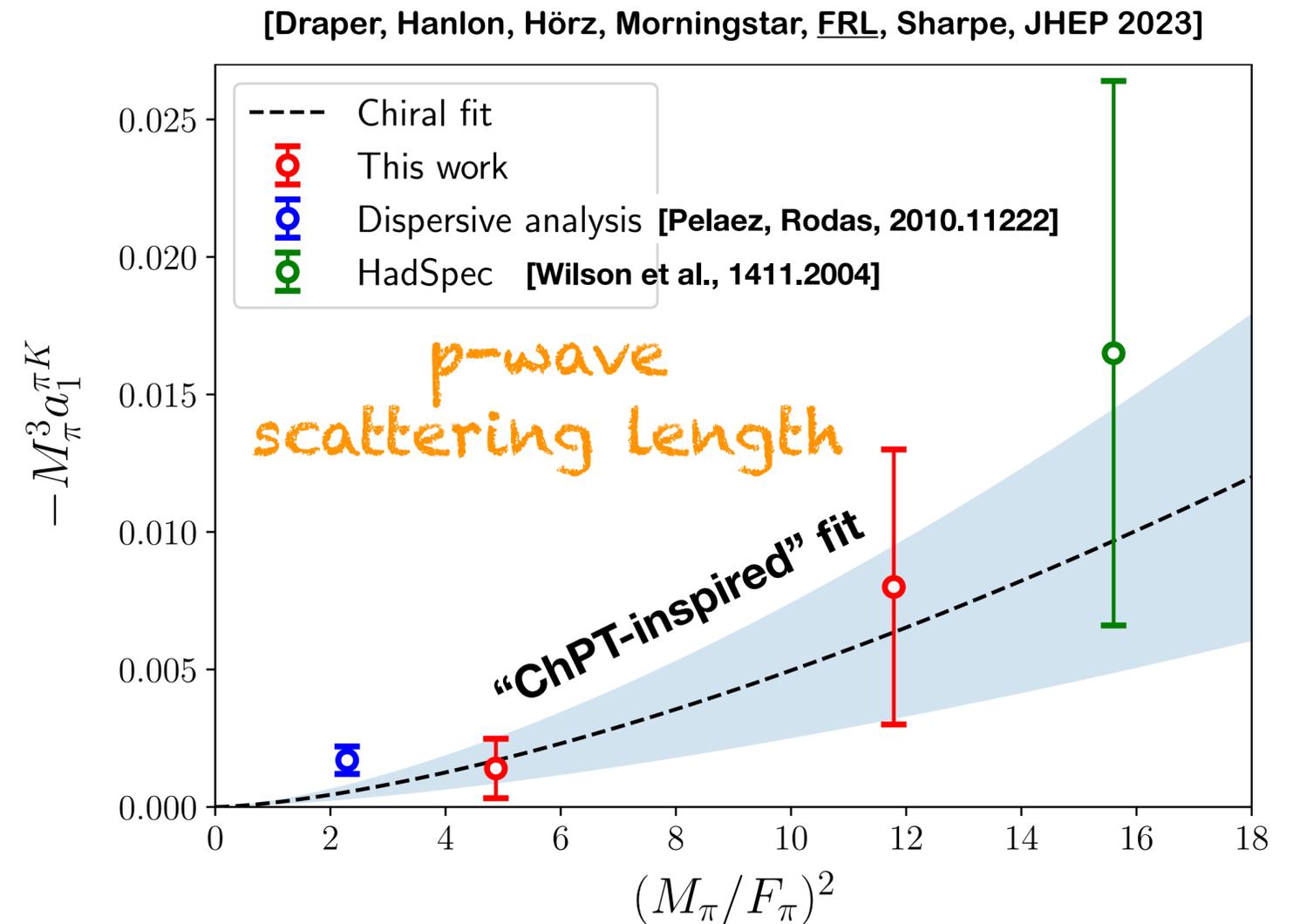
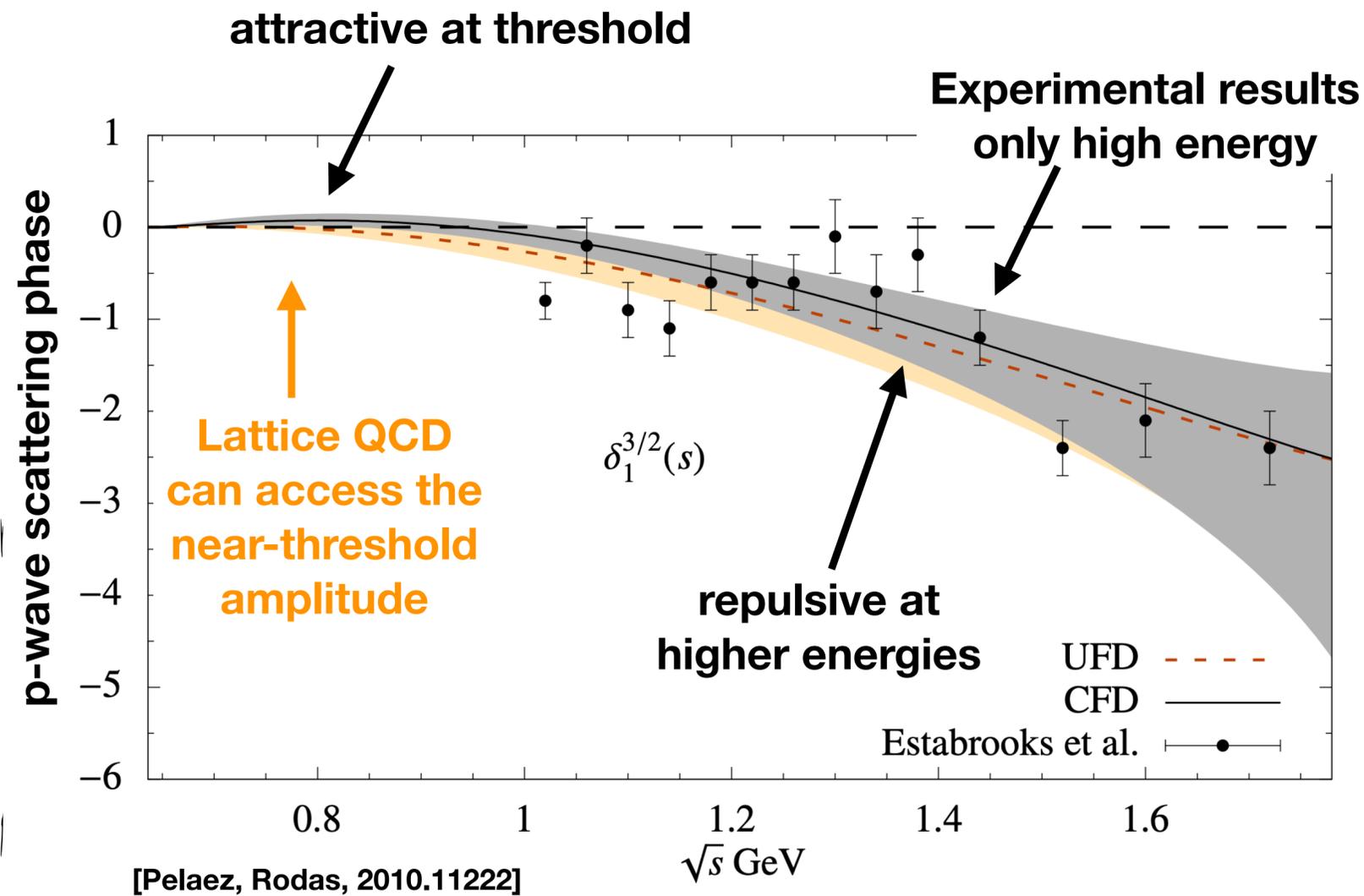
○ Lattice QCD results are complementary to experiment!



[Pelaez, Rodas, 2010.11222]

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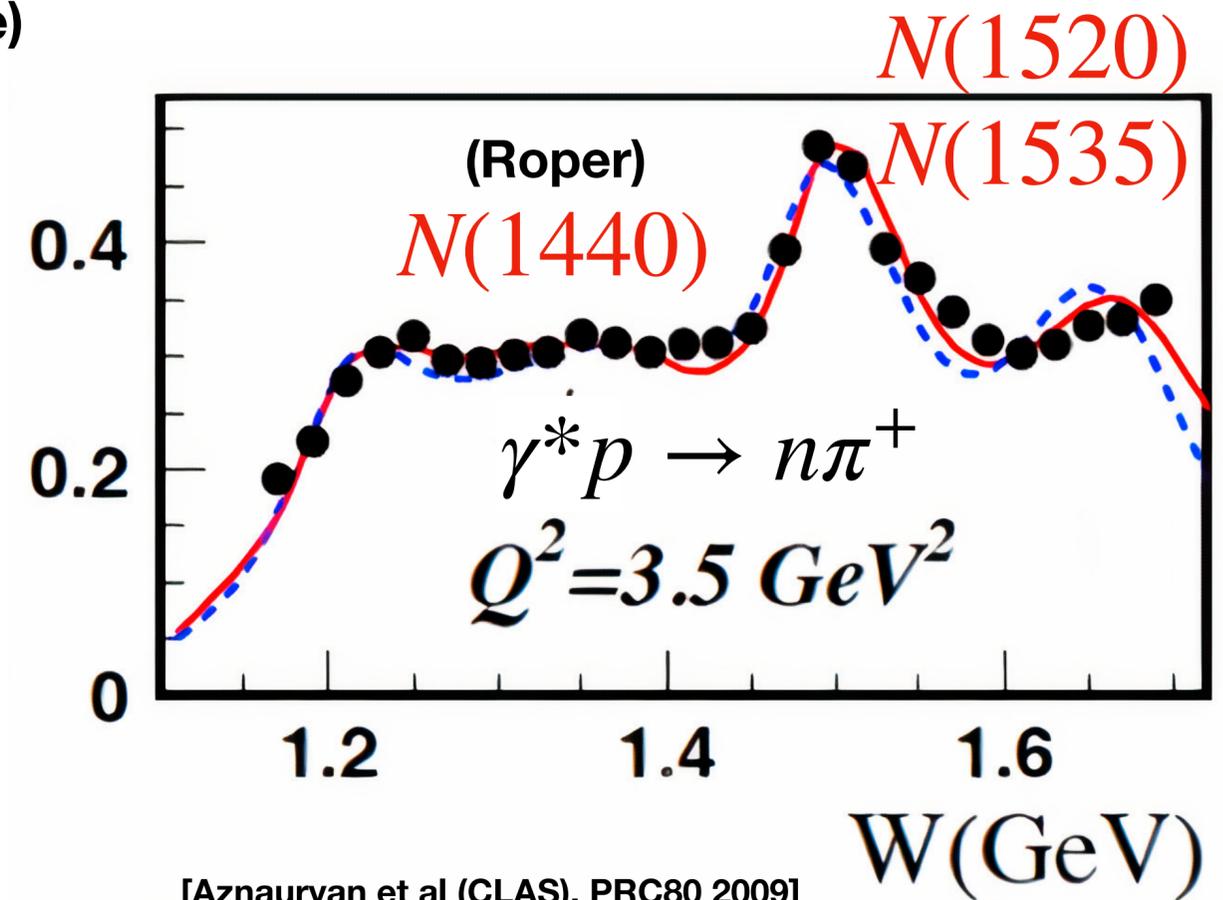
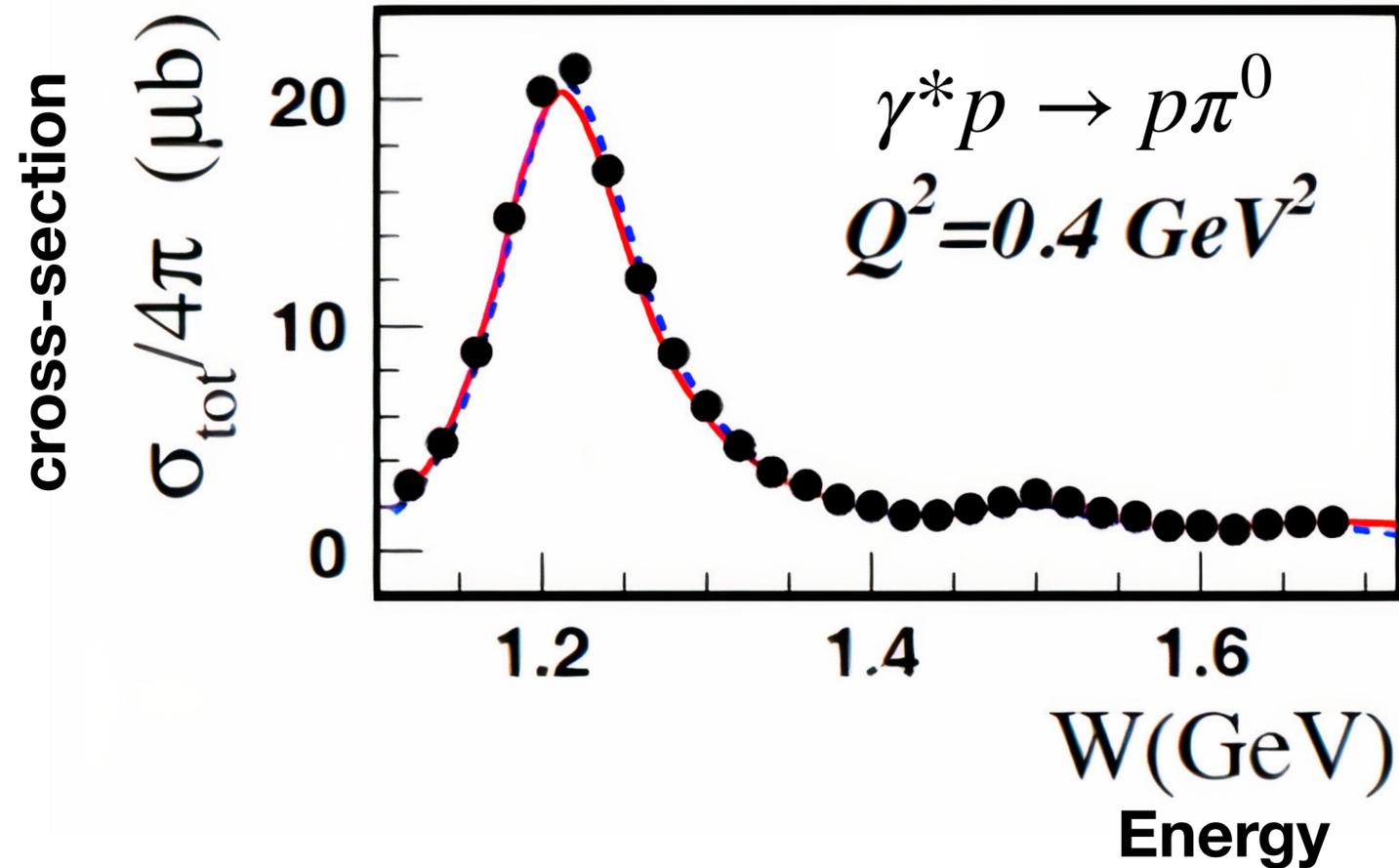


The πN scattering amplitude

πN scattering

- Pion-nucleon scattering is an important process in QCD

$\Delta(1232)$ (lowest-lying baryon resonance)

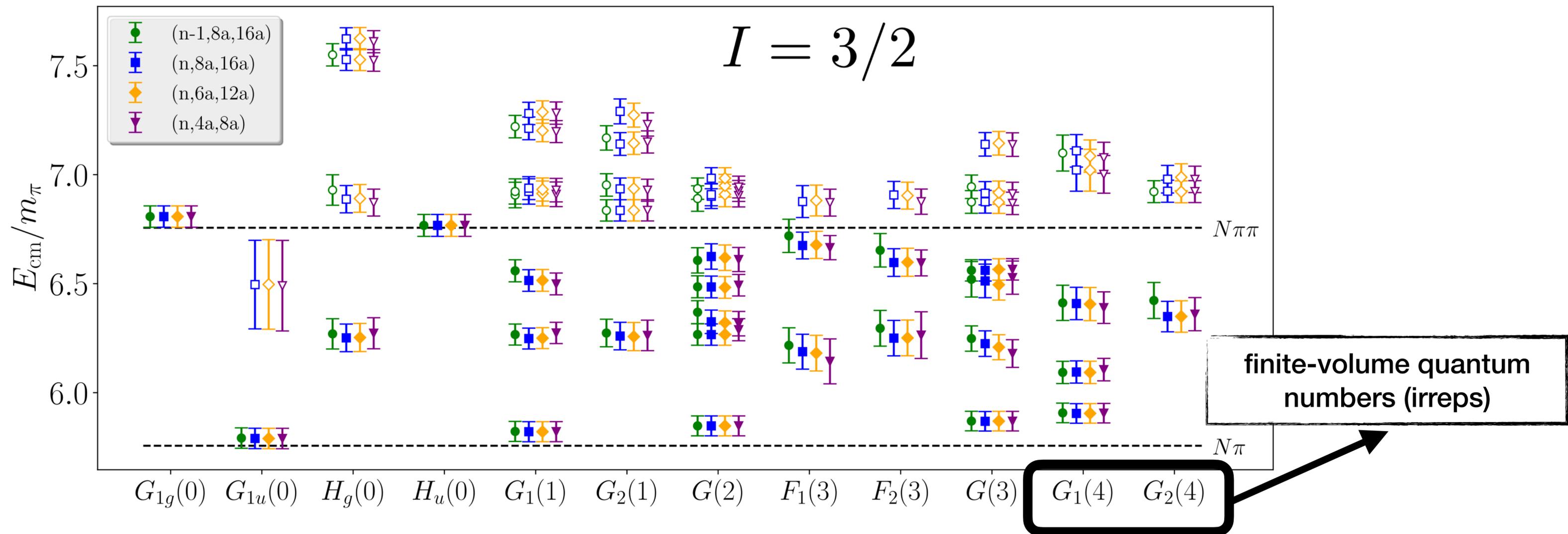


[Aznauryan et al (CLAS), PRC80 2009]

[Burkert, Roberts, 1710.02549]

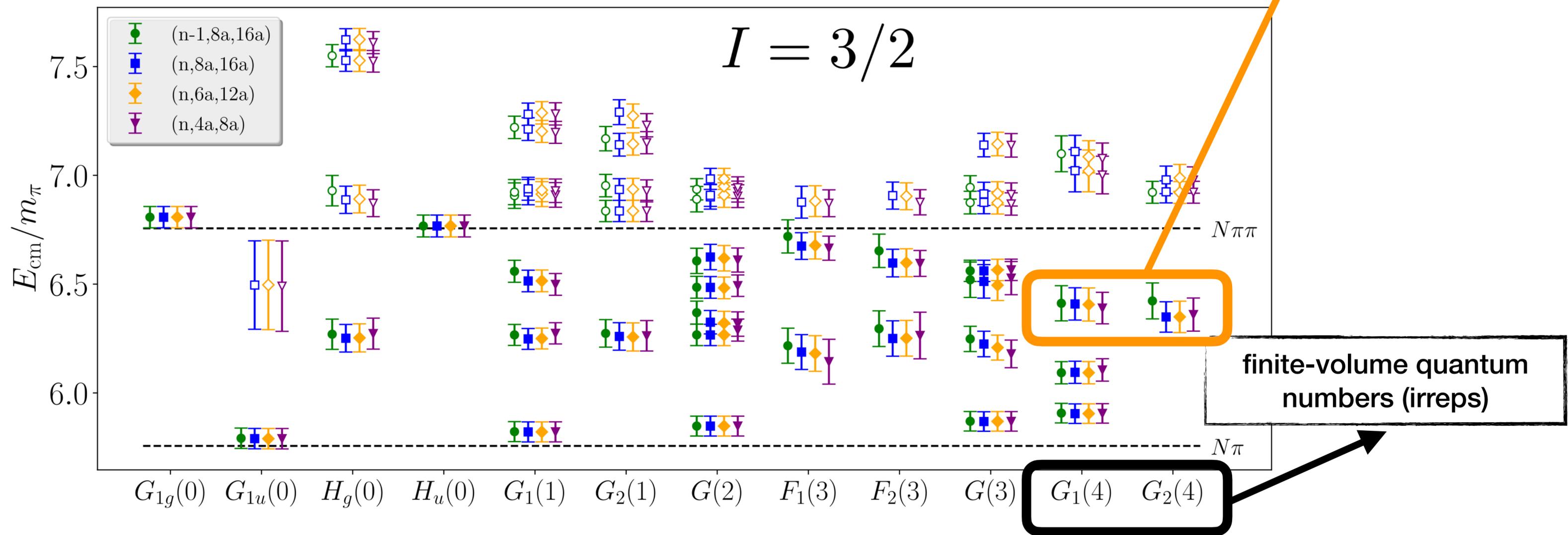
Lattice QCD πN spectrum

Key ingredient: reliable variational extractions of the lattice QCD energy levels: **GEVP + stability**



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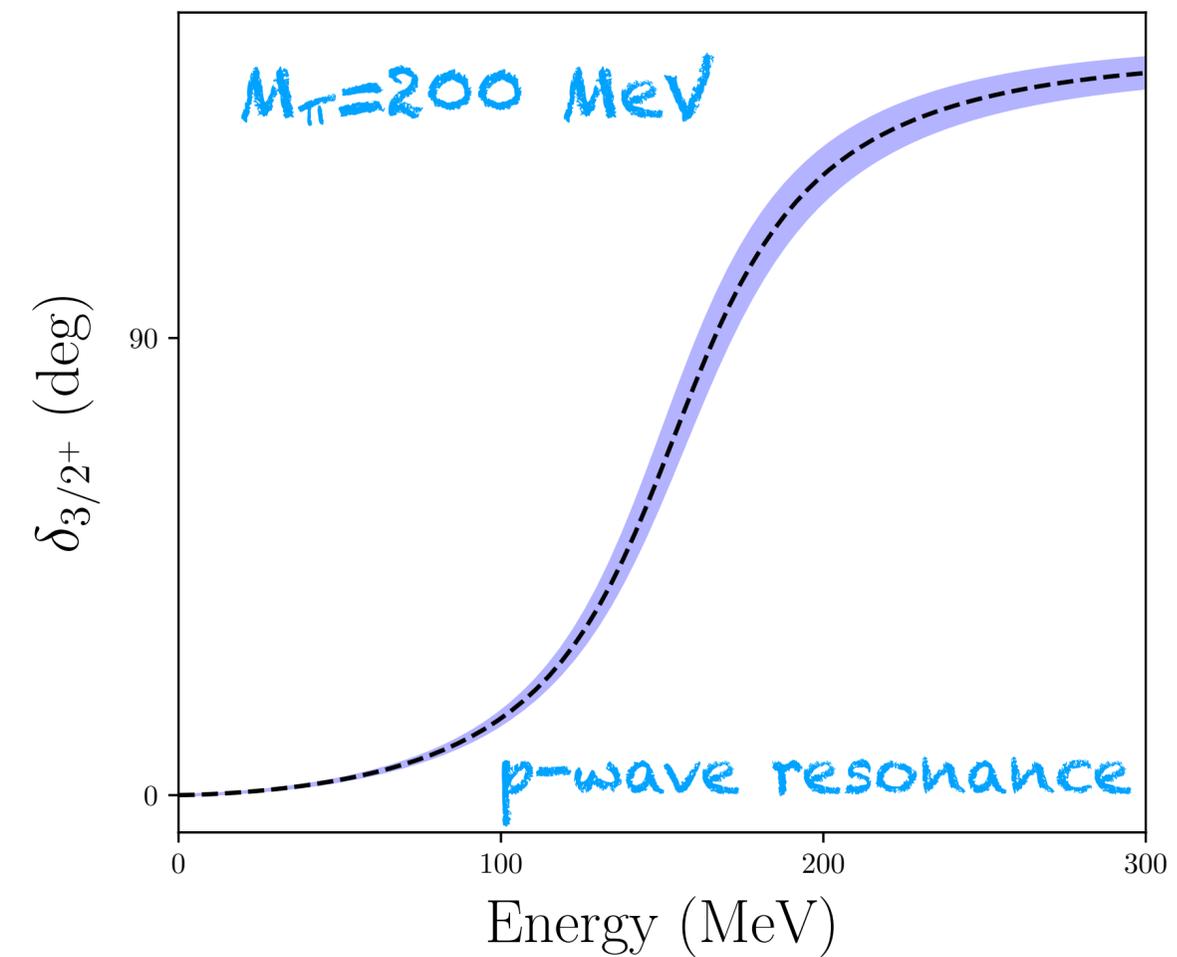
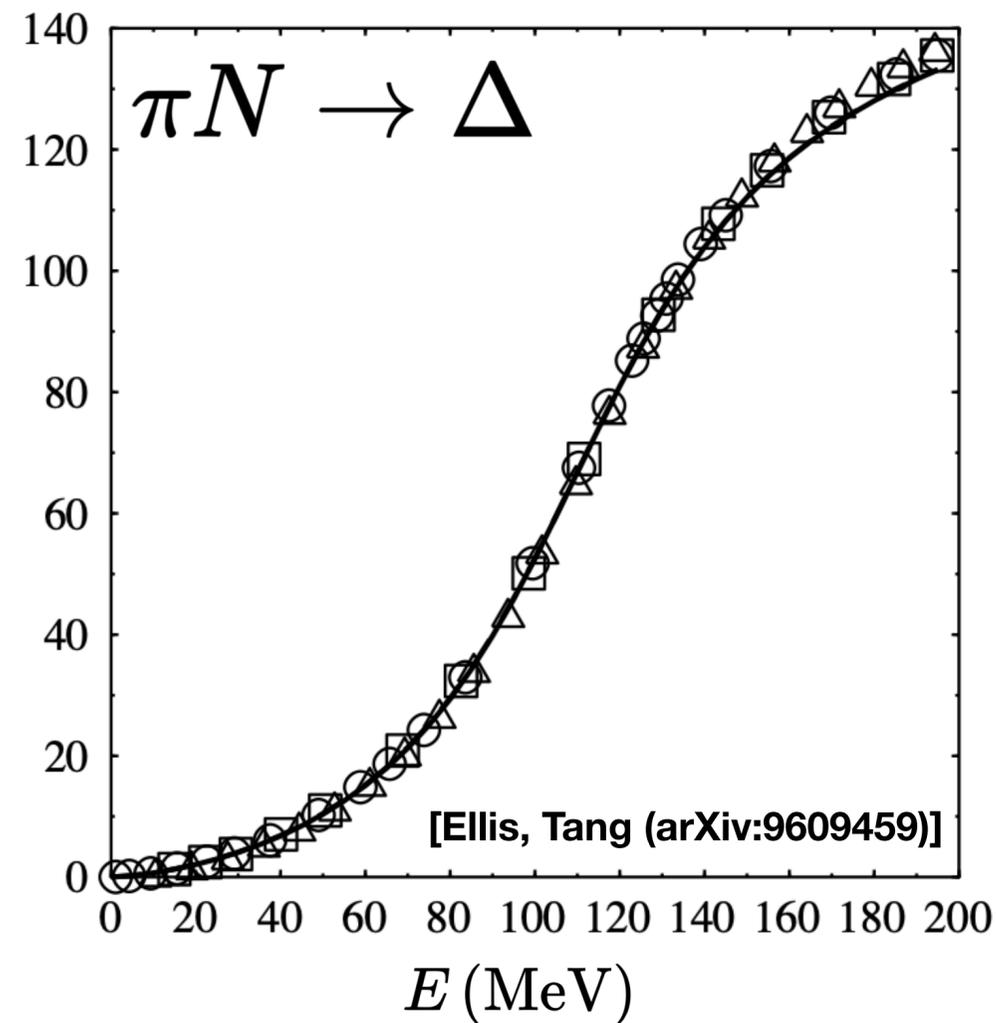


The $\Delta(1232)$ resonance

Experiment

Lattice QCD

[Bulava, Hanlon, Hörz, Morningstar, Nicholson, FRL, Skinner, Vranas, Walker-Loud, 2208.03867]

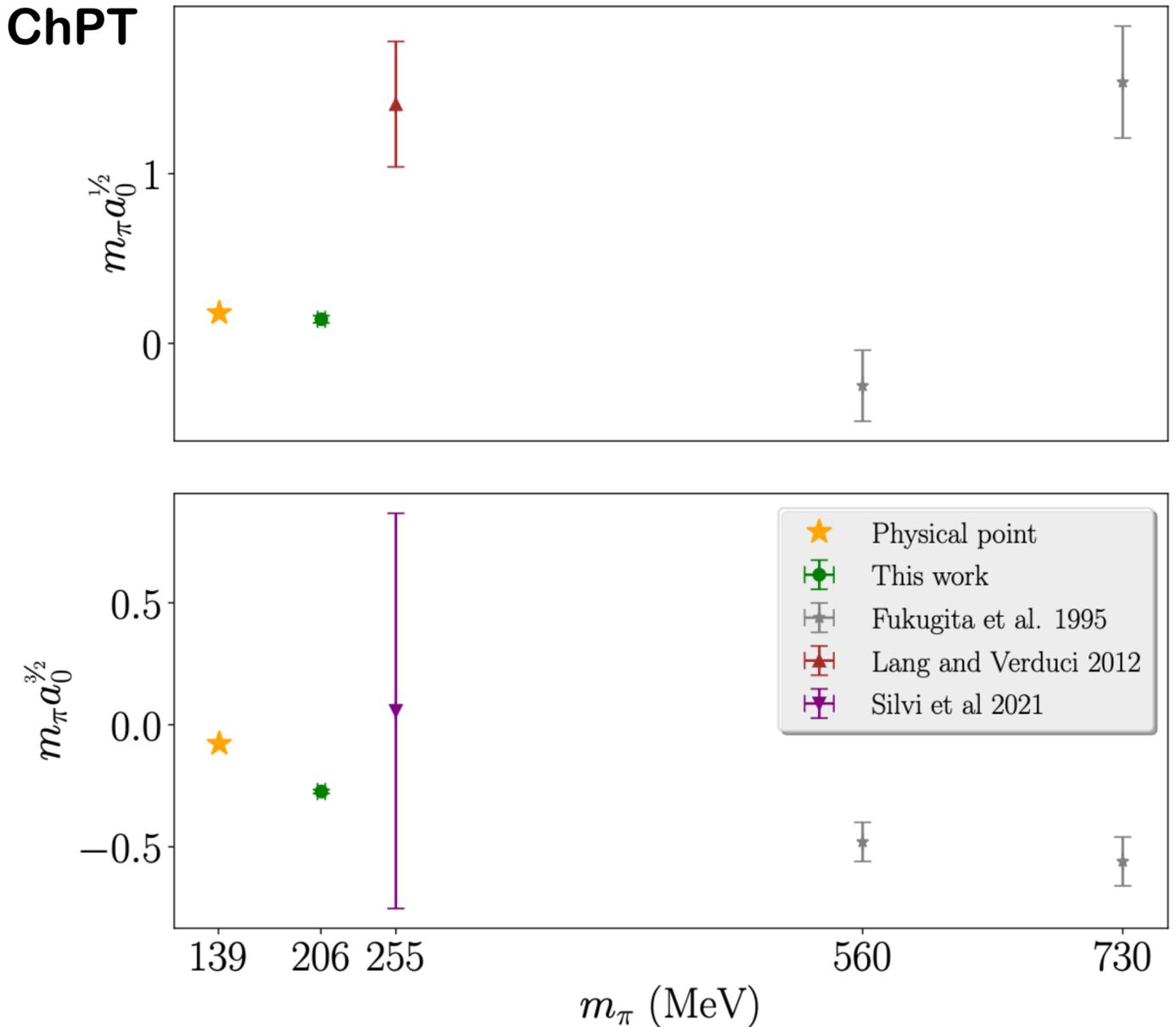


The πN scattering lengths

[Bulava, Hanlon, Hörz, Morningstar, Nicholson, [FRL](#),
Skinner, Vranas, Walker-Loud, 2208.03867]

- Our results can be used to test the convergence of baryon ChPT

	m_π (MeV)	$m_\pi a_0^{1/2}$	$m_\pi a_0^{3/2}$
This work	200	0.142(22)	-0.2735(81)
LO χ PT	200	0.321(04)(57)	-0.161(02)(28)

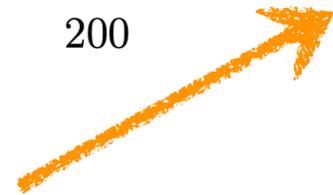


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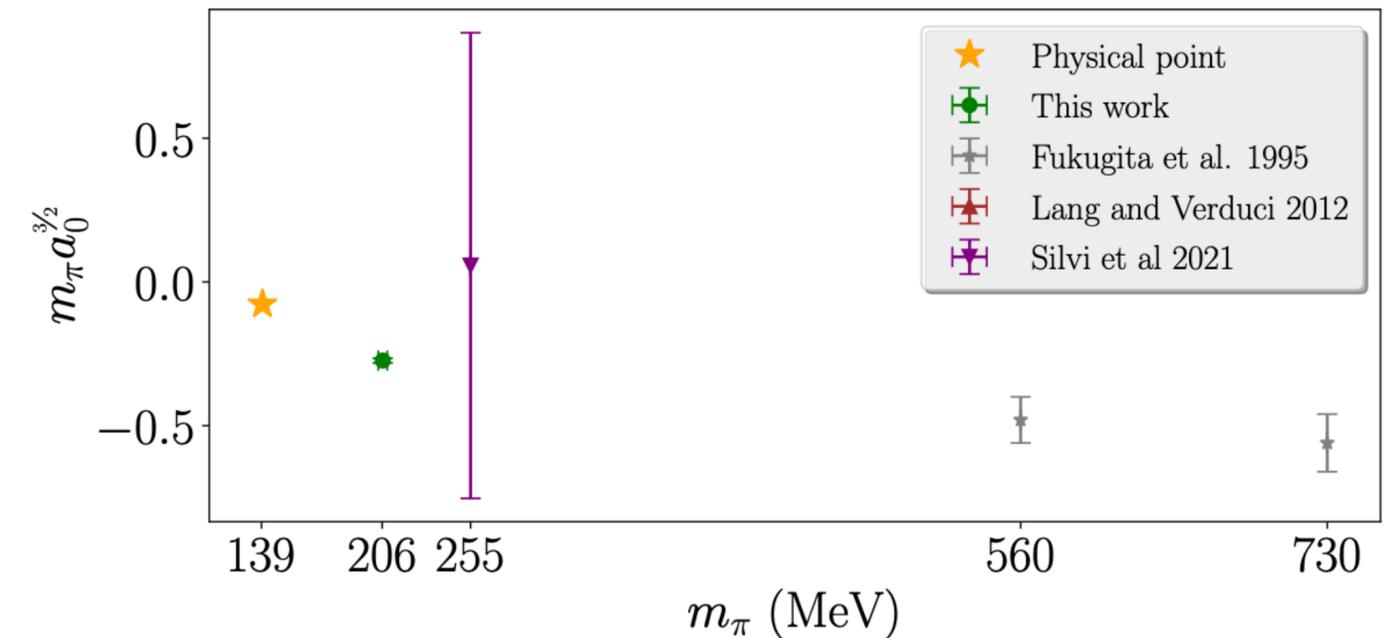
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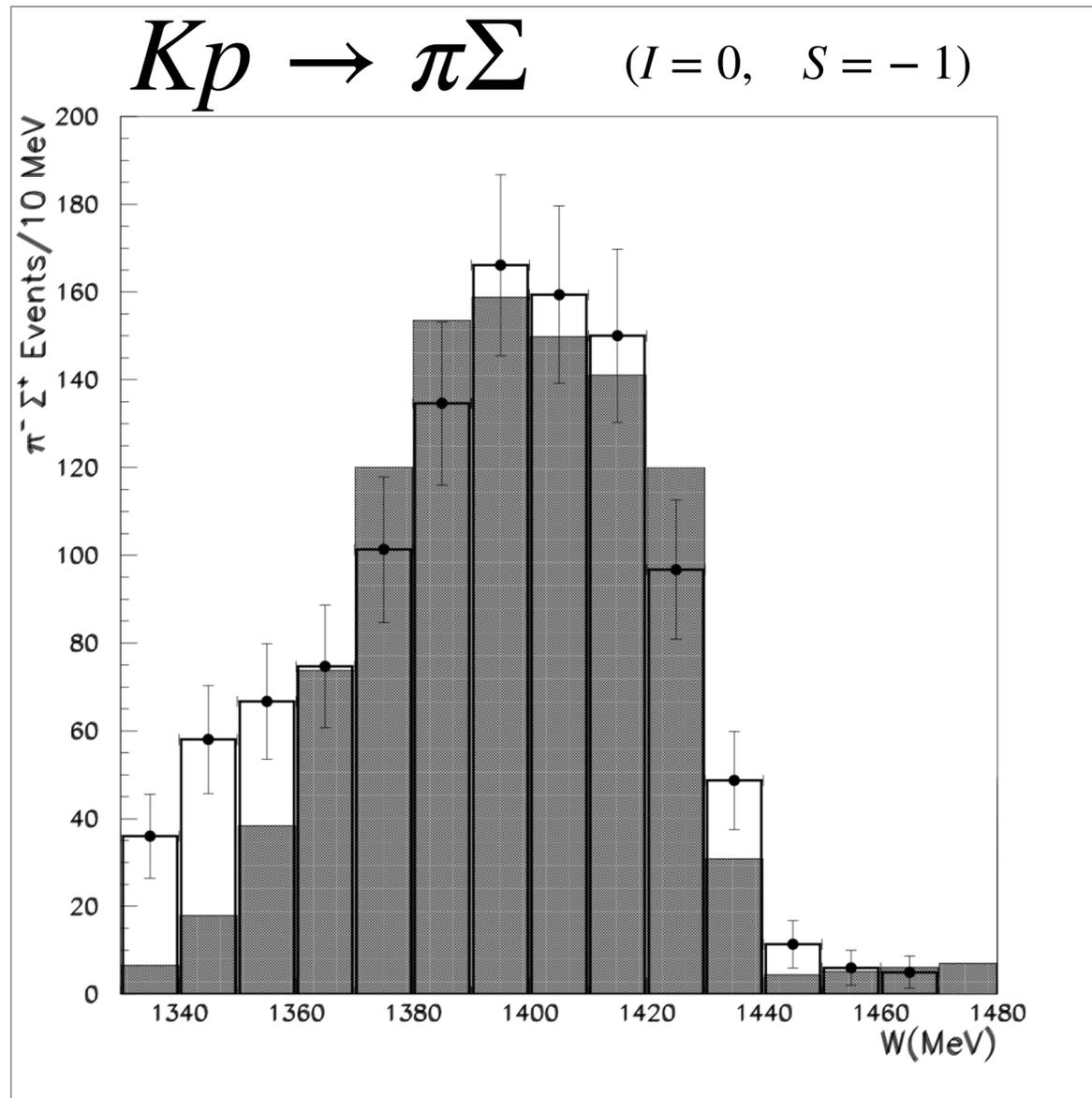
We find poor convergence at $M_\pi = 200$ MeV

- Additional values of the pion mass are needed!



$\wedge(1405)$

The $\Lambda(1405)$ resonance



[Oller, Meißner, 0011146], [Hemingway, NPB 1985]

PDG (4 star resonance):

$\Lambda(1405) 1/2^-$

$$I(J^P) = 0(\frac{1}{2}^-)$$

$$\text{Mass } m = 1405.1^{+1.3}_{-1.0} \text{ MeV}$$

$$\text{Full width } \Gamma = 50.5 \pm 2.0 \text{ MeV}$$

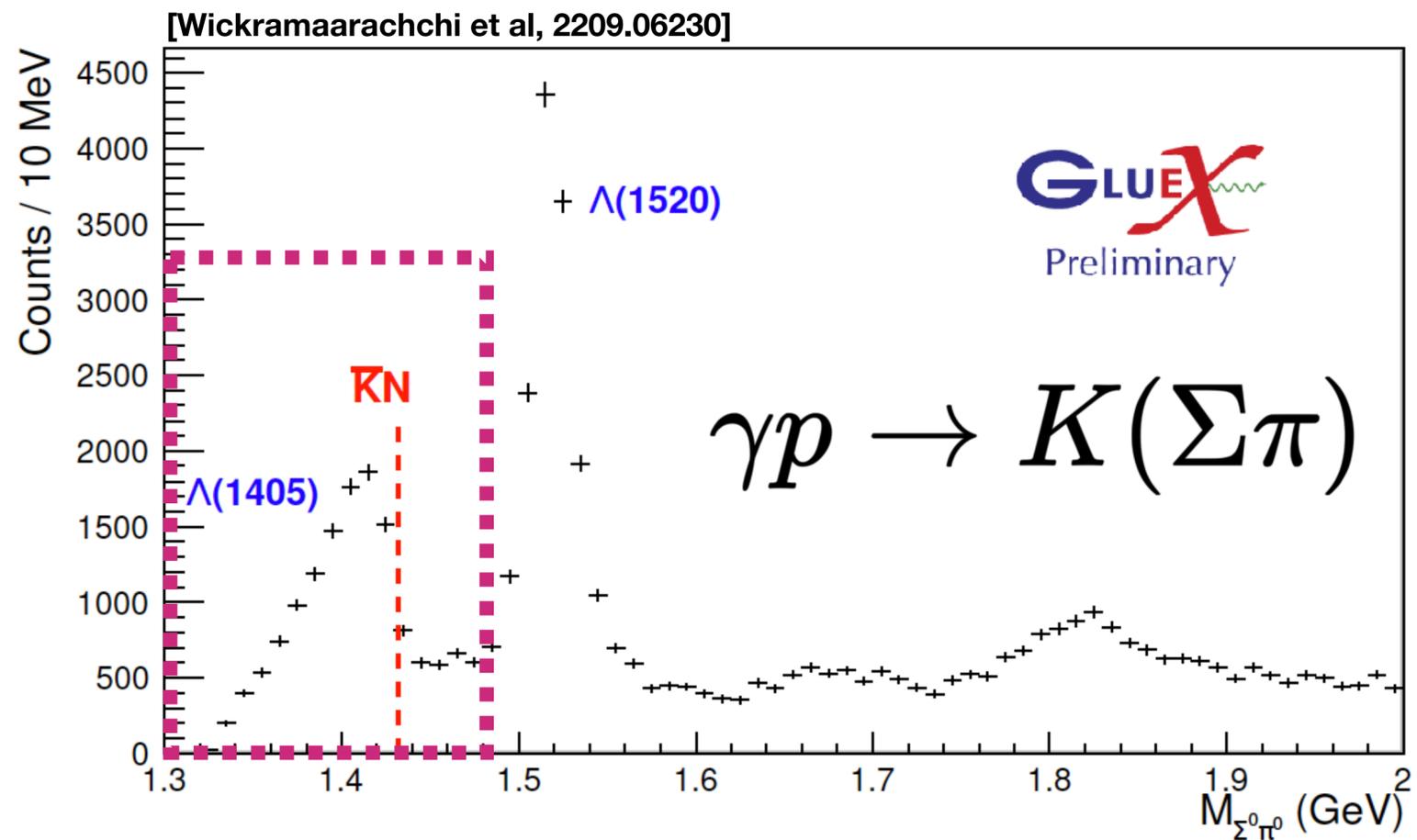
Below $\bar{K}N$ threshold

$\Lambda(1405)$ DECAY MODES	Fraction (Γ_i/Γ)	p (MeV/c)
$\Sigma \pi$	100 %	155

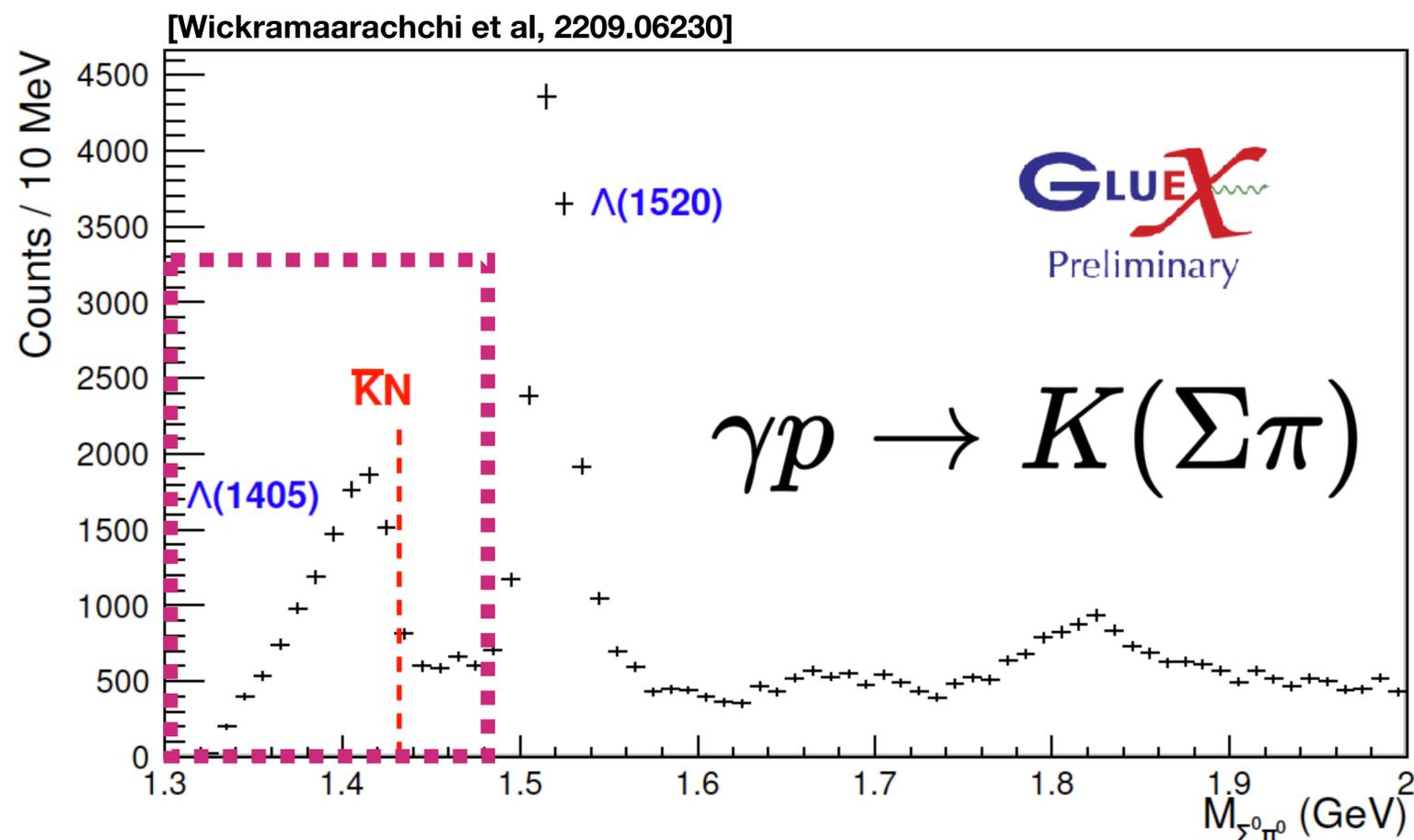
Resonance appearing in multi-channel scattering:

$$\begin{pmatrix} \pi\Sigma \rightarrow \pi\Sigma & \pi\Sigma \rightarrow Kp \\ Kp \rightarrow \pi\Sigma & Kp \rightarrow Kp \end{pmatrix}$$

The $\Lambda(1405)$ resonance



The $\Lambda(1405)$ resonance



○ Long-standing puzzle about its fundamental nature:

➔ One or two resonance poles?

➔ Chiral EFT analysis indicate two poles

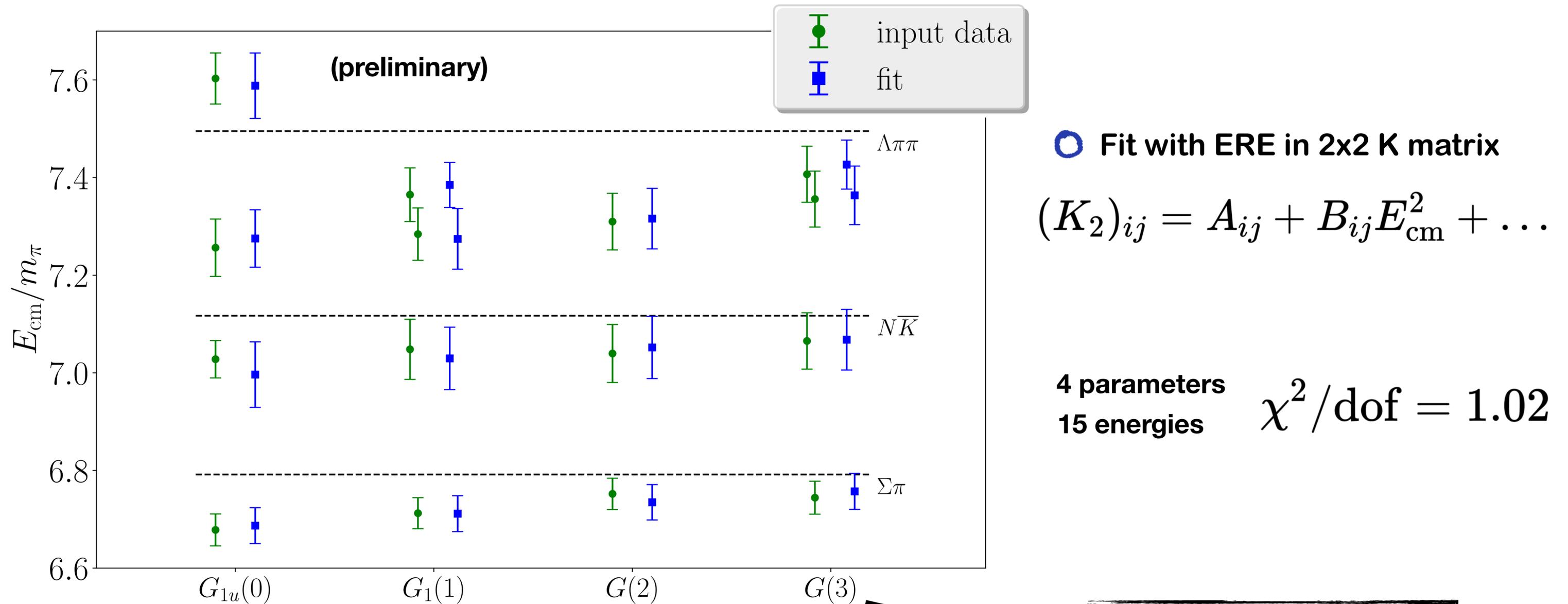
approach	[PDG 2022]pole 1 [MeV]	pole 2 [MeV]	
Refs. [14, 15], NLO	$1424_{-23}^{+7} - i 26_{-14}^{+3}$	$1381_{-6}^{+18} - i 81_{-8}^{+19}$	[1109.3005, 1201.6549]
Ref. [17], Fit II	$1421_{-2}^{+3} - i 19_{-5}^{+8}$	$1388_{-9}^{+9} - i 114_{-25}^{+24}$	[1210.3485]
Ref. [18], solution #2	$1434_{-2}^{+2} - i 10_{-1}^{+2}$	$1330_{-5}^{+4} - i 56_{-11}^{+17}$	[1411.7884]
Ref. [18], solution #4	$1429_{-7}^{+8} - i 12_{-3}^{+2}$	$1325_{-15}^{+15} - i 90_{-18}^{+12}$	

well resolved

more uncertain

○ Lattice QCD: essential non-perturbative verification

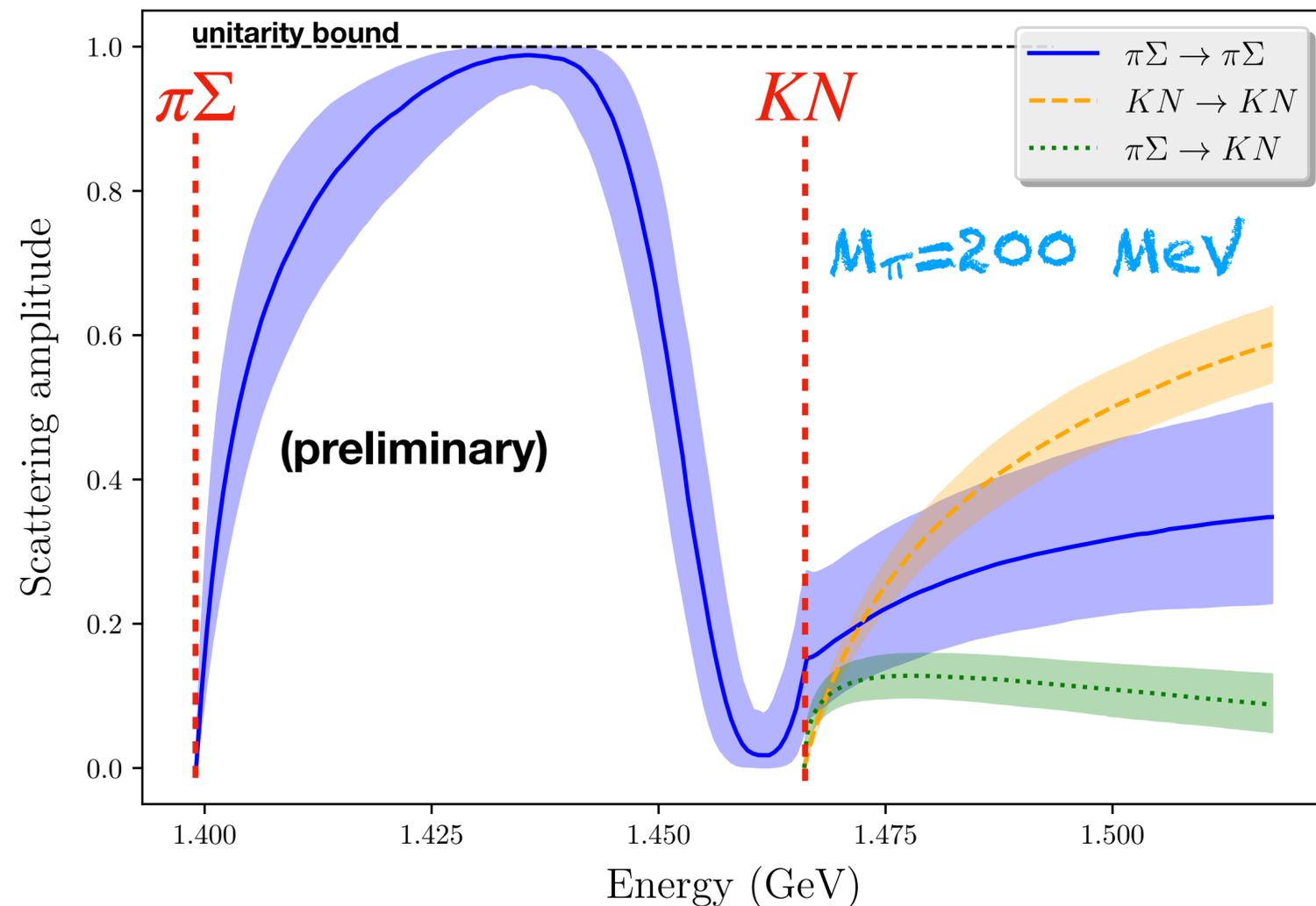
The $\Lambda(1405)$ in Lattice QCD



finite-volume quantum numbers (irreps)

The $\Lambda(1405)$ in Lattice QCD

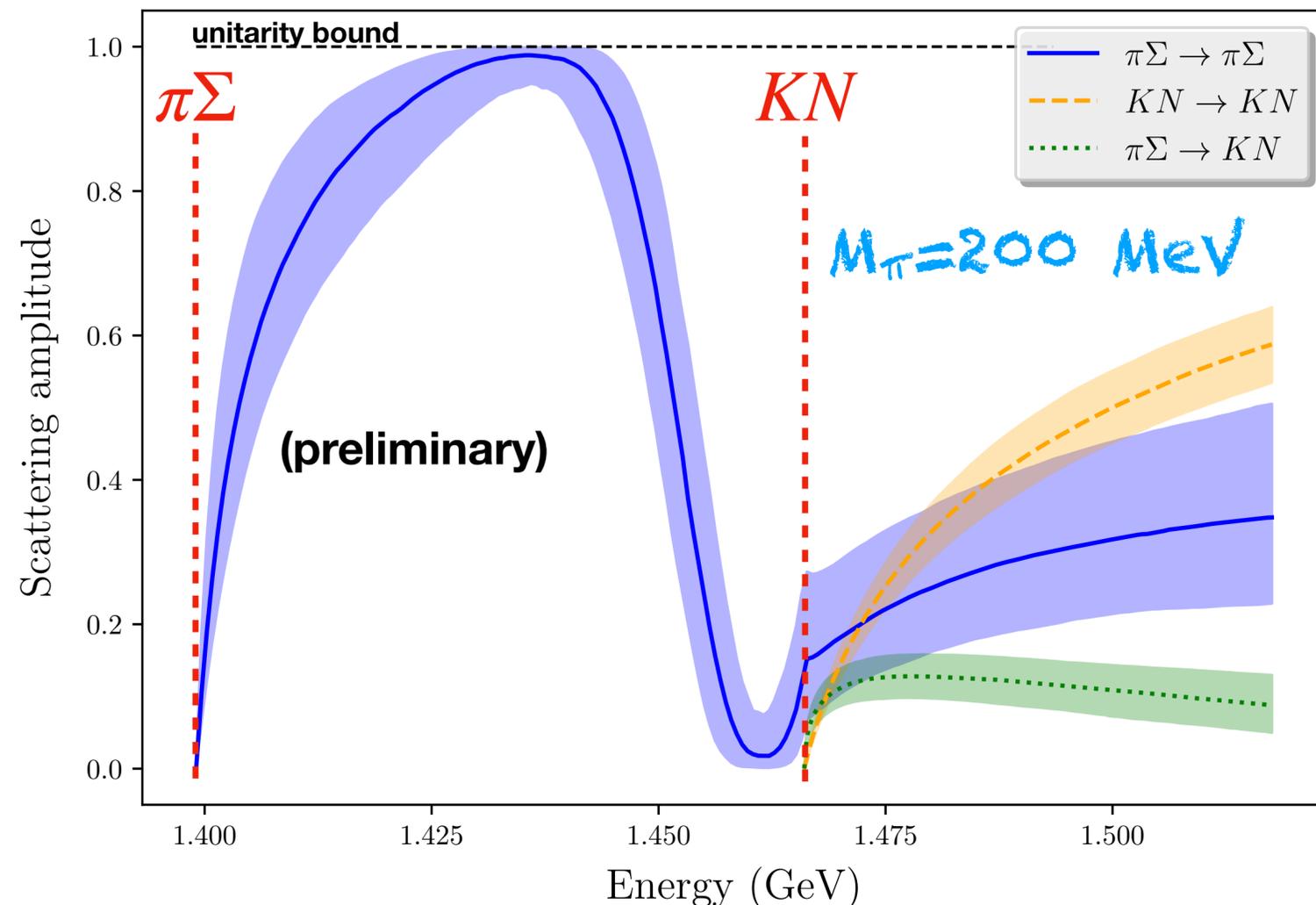
- ✓ First lattice QCD study of its analytic structure using multi-channel scattering



[Bulava, Cid Mora, Hanlon, Hörz, Mohler, Morningstar,
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The $\Lambda(1405)$ in Lattice QCD

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- ✓ We find two resonance poles (preliminary)

$$\sqrt{s}_{\text{pole 1}} = 1456(21) - i 12(6) \text{ MeV} \quad \sqrt{s}_{\text{pole 2}} = 1395(18) \text{ MeV}$$

(qualitative agreement with chiral EFTs)

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- ✓ Qualitative picture stable under model variations

- ✓ Results close to the physical point

➔ Work in progress for a physical point calculation

Three-particle processes

Why three particles?

○ The two-body formalism is restricted to few interesting resonances

▶ Exotics: $T_{cc} \rightarrow DD^*, DD\pi$

▶ Roper: $N(1440) \rightarrow \Delta\pi \rightarrow N\pi\pi$

Resonance	$I_{\pi\pi\pi}$	J^P
$\omega(782)$	0	1^-
$h_1(1170)$	0	1^+
$\omega_3(1670)$	0	3^-
$\pi(1300)$	1	0^-
$a_1(1260)$	1	1^+
$\pi_1(1400)$	1	1^-
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$a_2(1320)$	1	2^+
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(with $\geq 3\pi$ decay modes)

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- ☑ Major developments in the three-particle finite-volume formalism

[Hansen, Sharpe, PRD 2014 & 2015], [Hammer, Pang, Rusetsky, JHEP 2017] x 2

[Mai, Döring, EPJA 2017]

[...]

[Blanton, FRL, Sharpe, JHEP 2019], [Hansen, FRL, Sharpe, JHEP 2020]

[Hansen, FRL, Sharpe, JHEP 2021], [Blanton, FRL, Sharpe, JHEP 2022]

[Draper, Hansen, FRL, Sharpe, JHEP 2023]

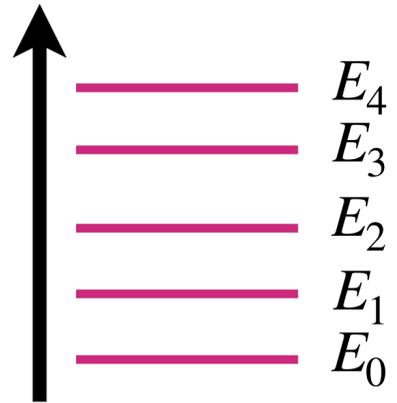
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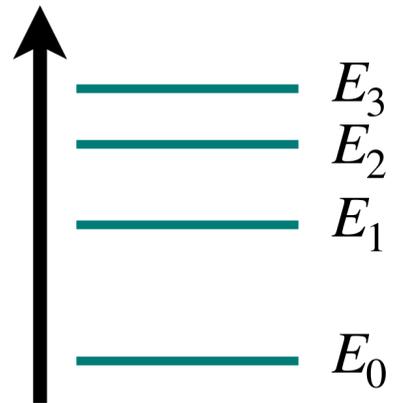
Relativistic three-particle formalism for identical particles

[Hansen, Sharpe, PRD 2014 & 2015]

2π Spectrum



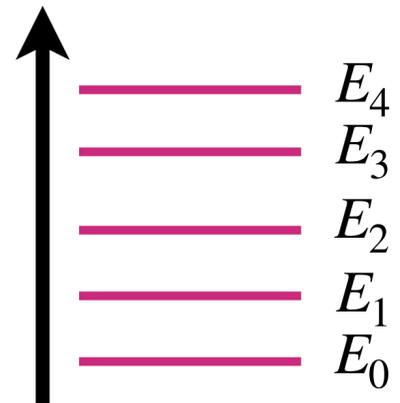
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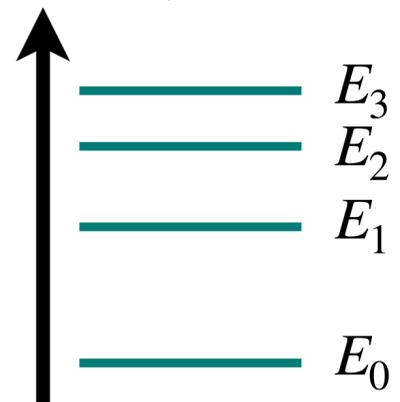
2π Spectrum



Quantization
conditions

$$\det_{\ell m} [\mathcal{K}_2 + F_2^{-1}] = 0$$

3π Spectrum



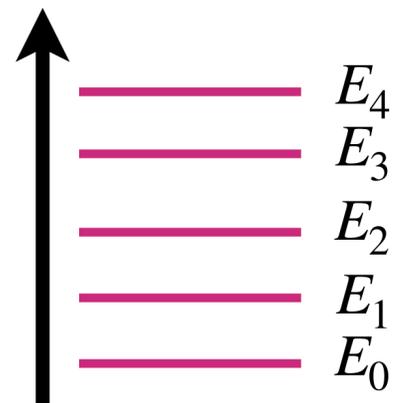
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Matrix indices describe
three on-shell particles

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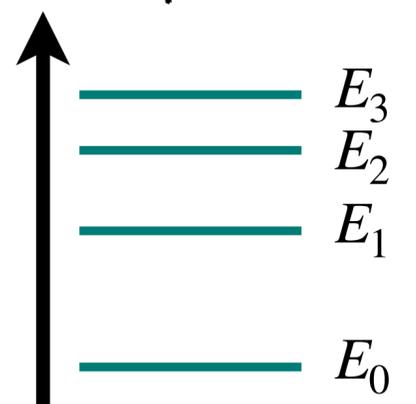


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K-matrices

\mathcal{K}_2

$\mathcal{K}_{df,3}$

Fit

Parametrize:

$$\mathcal{K}_2 = c_0 + c_1 k^2 + \dots$$

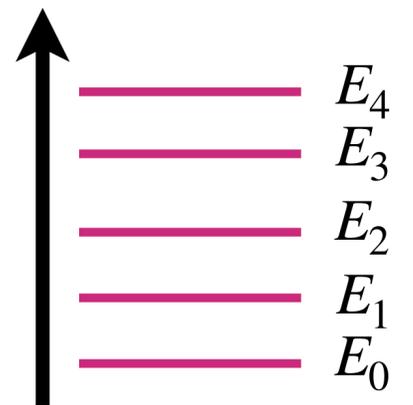
$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{\text{iso},0} + \mathcal{K}_{df,3}^{\text{iso},1} \left(\frac{s - 9m^2}{9m^2} \right) + \dots$$

[Blanton, FRL, Sharpe, JHEP 2019]

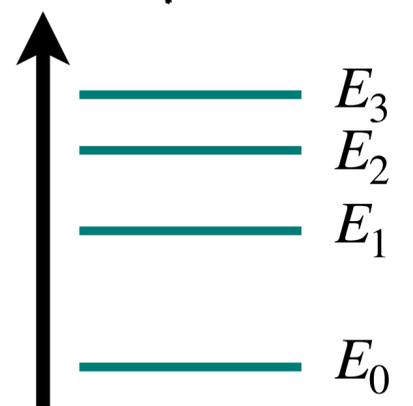
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Unitarity relations

Scattering amplitudes

$$\mathcal{M}_2$$

Integral equations

$$\mathcal{M}_3$$

[Briceño et al., PRD 2018]
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[Dawid et al., 2303.04394]

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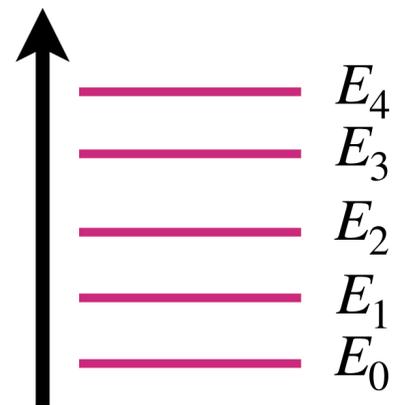
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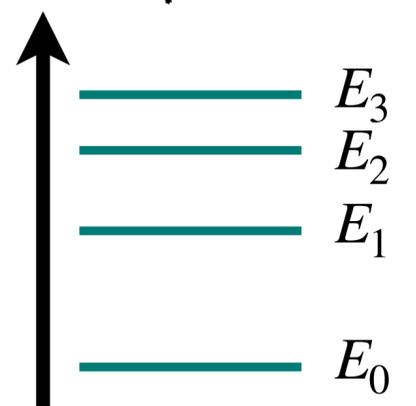
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Experiments

⊗

resonance properties

Three-meson systems

○ Three-particle formalism applied to weakly-interacting (non-resonant) systems: $\pi^+\pi^+\pi^+$, $\pi^+\pi^+K^+$

[Blanton ... [FRL...](#) et al., PRL 2020 & JHEP 2021], [Draper ... [FRL...](#) et al., JHEP 2023], [Fischer ... [FRL...](#) et al, EPJC 2021]

[Alexandrou et al, Brett et al, Culver et al, Hansen et al, Mai et al]

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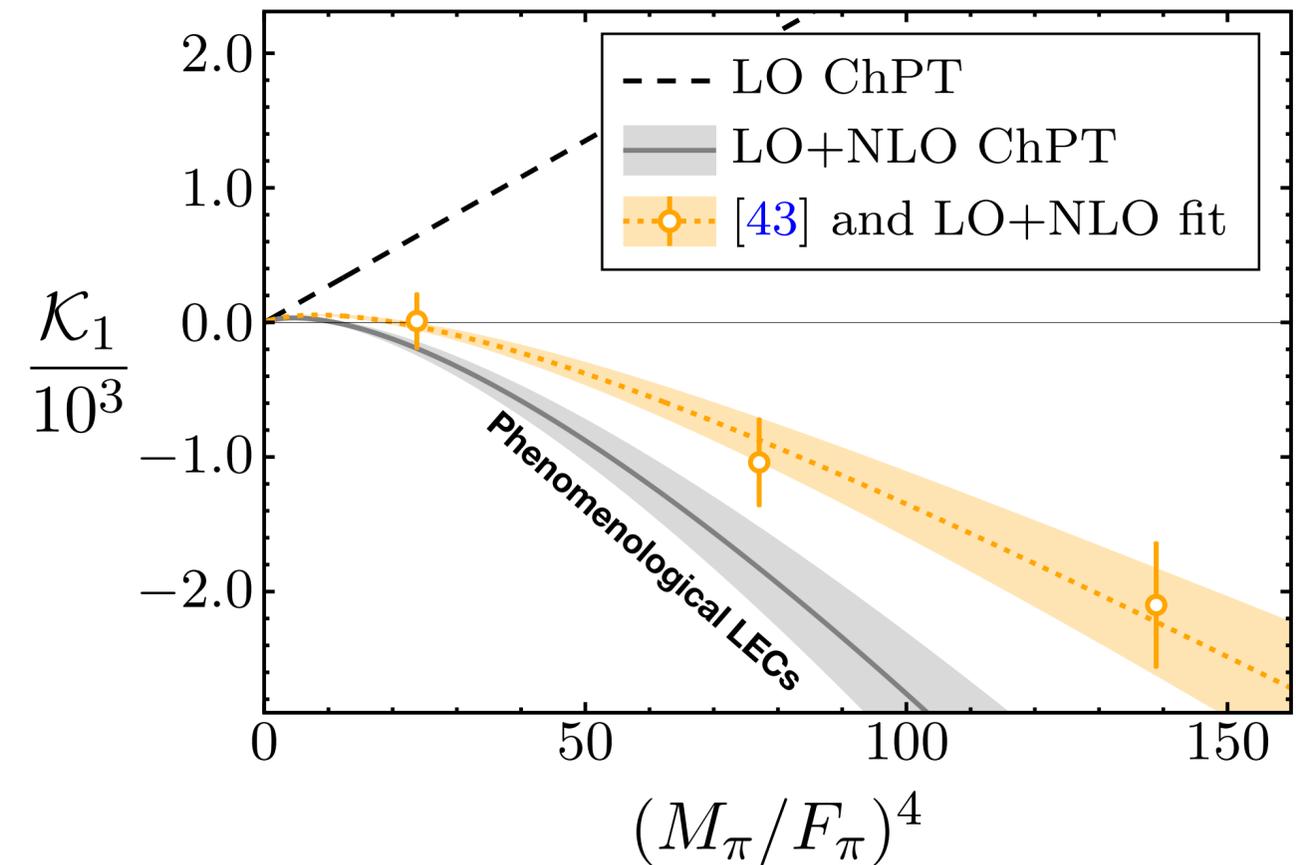
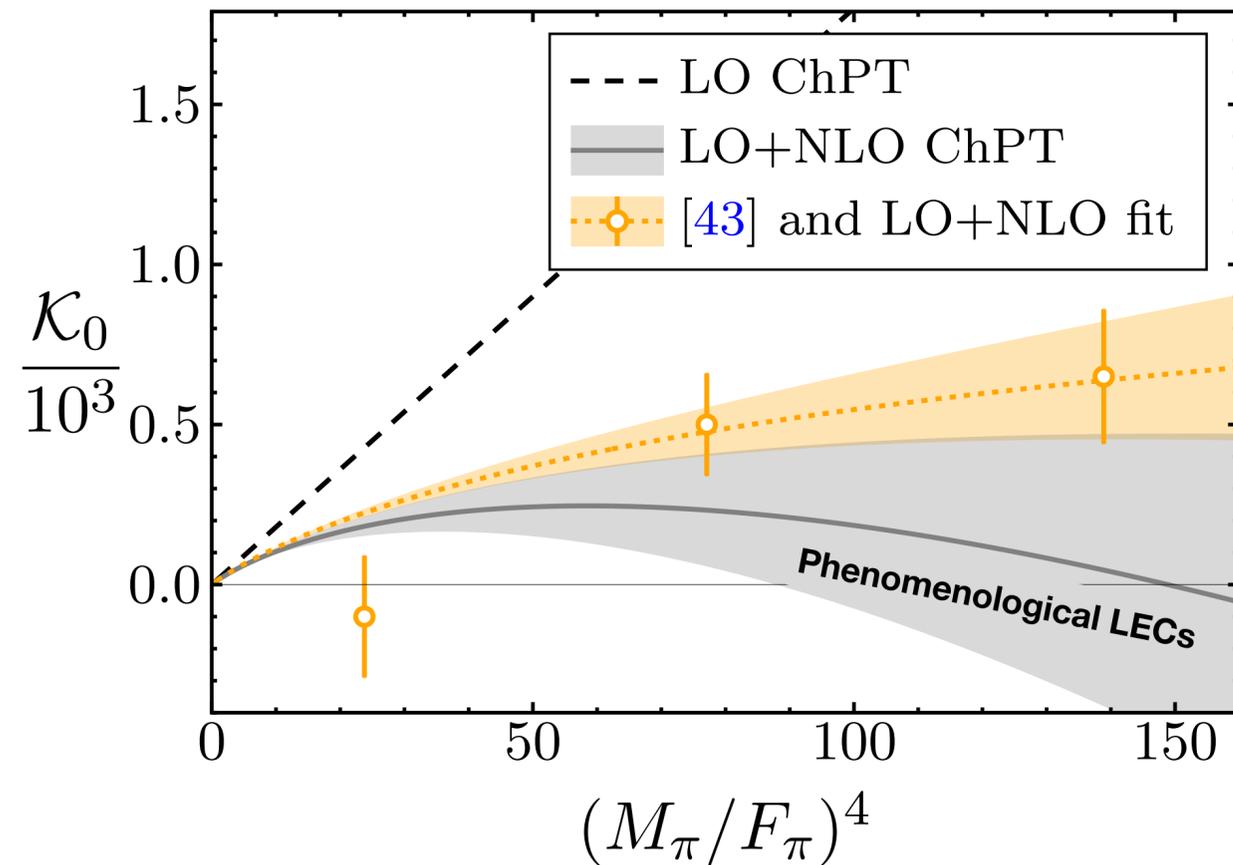
Example: $\pi^+\pi^+\pi^+$ scattering

Lattice data: [Blanton, Hanlon, Hörz, Morningstar, [FRL](#), Sharpe, JHEP 2021]

NLO ChPT: [Baeza-Ballesteros, Bijens, Husek, [FRL](#), Sharpe, Sjö, JHEP 2023]

parametrized by the three-particle K-matrix

$$\mathcal{K}_{\text{df},3} = \mathcal{K}_0 + \mathcal{K}_1 \left(\frac{s - 9M_\pi^2}{9M_\pi^2} \right) + \dots$$



Three-body resonances

- Formalism to study relevant three-pion resonances is available
[Blanton, Sharpe, PRD 2021] x 2, [Hansen, [FRL](#), Sharpe, JHEP 2021], [Mai et al (GWQCD), PRL 2021]
 - ▶ Preparing formalism for T_{cc} [Draper, Hansen, [FRL](#), Sharpe, (in prep)]
- Extensive lattice QCD data is still not available

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- Extract resonance properties on a **toy model** [Garofalo, Mai, FRL, Rusetsky, Urbach (2211.05605)]

$$\mathcal{L} = \sum_{i=0,1} \left(\partial_\mu \phi_i^\dagger \partial_\mu \phi_i + m_i \phi_i^\dagger \phi_i + \lambda_i \phi_i^4 \right) + \frac{g}{2} \phi_1^\dagger \phi_0^3 + h.c.$$

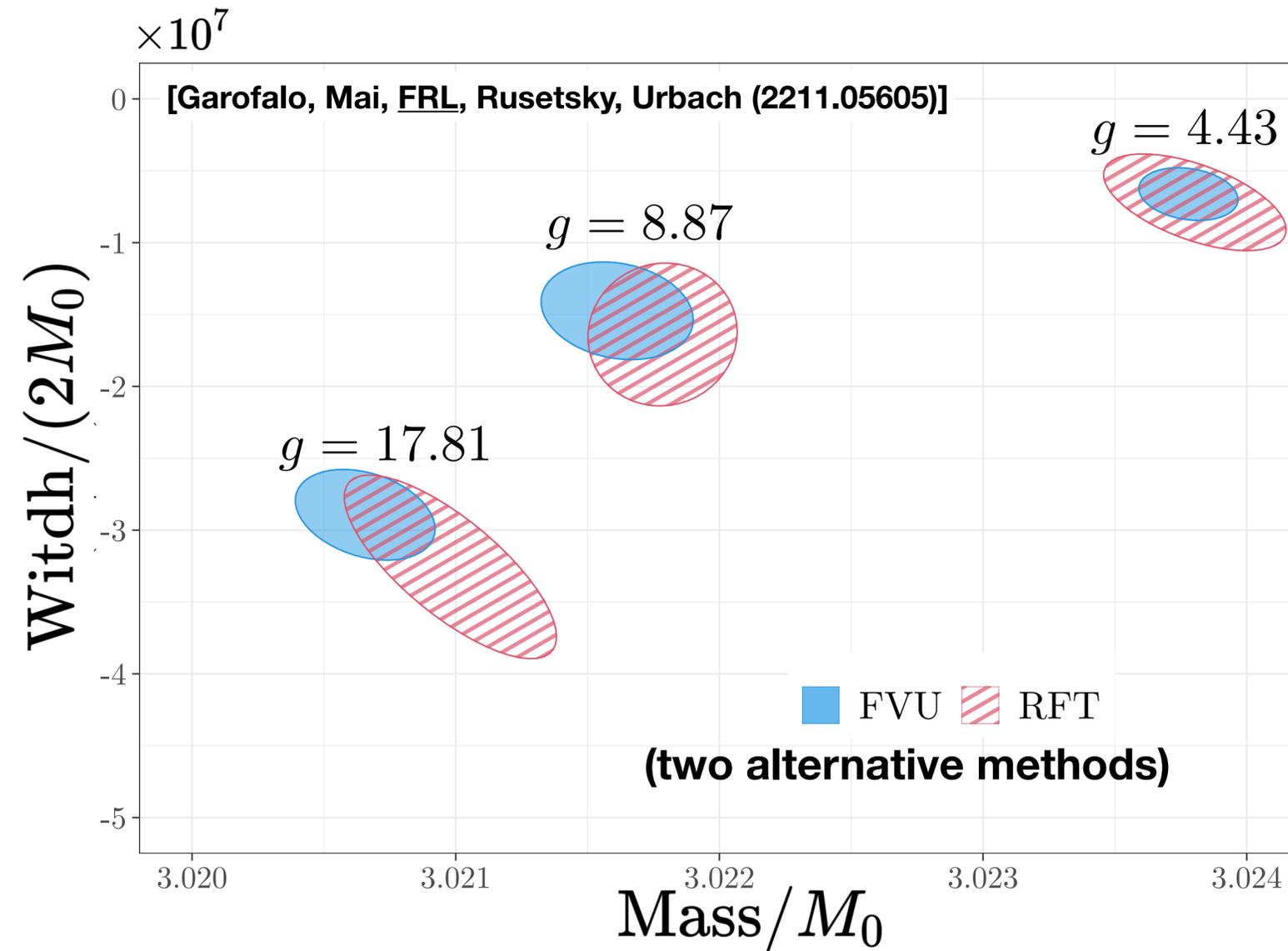
- ▶ Test formalism in a controlled setup
- ▶ Computationally cheaper
- ▶ Allows access to several volumes & energy levels

Induces transitions:

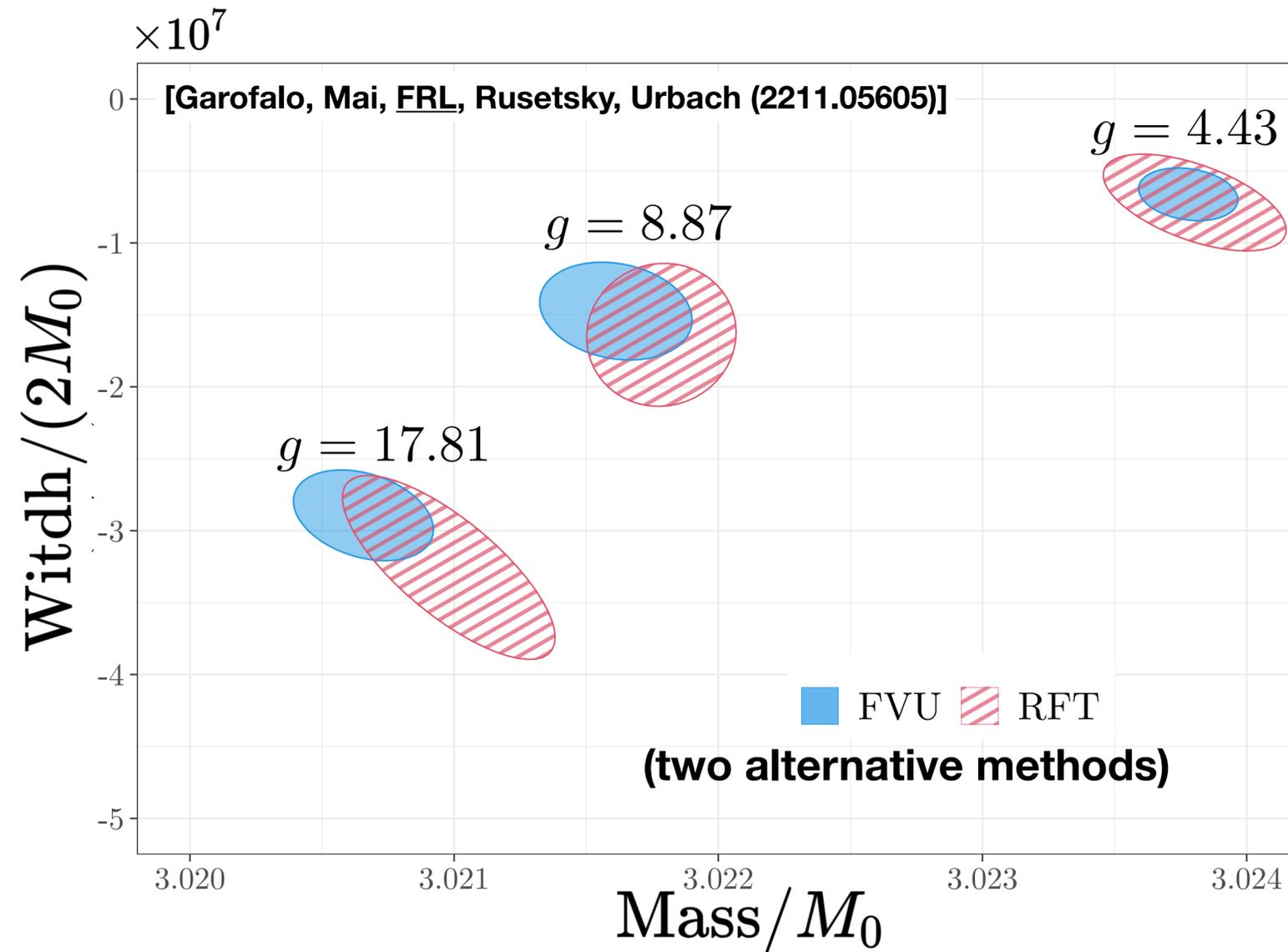
$$\phi_1 \rightarrow 3\phi_0$$

$(M_1 > 3M_0)$

A toy three-body resonance



A boy three-body resonance



Successfully determined properties of three-particle resonance for the first time!

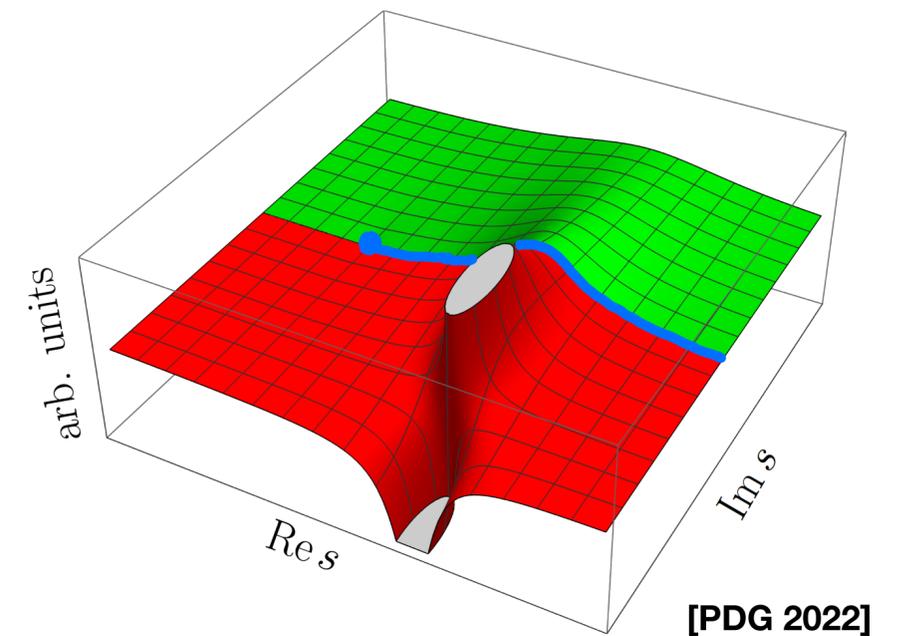
Good agreement between methods

Next: QCD resonances.
(e.g. $h_1(1170)$, T_{cc}^+ and Roper)

Summary & Outlook

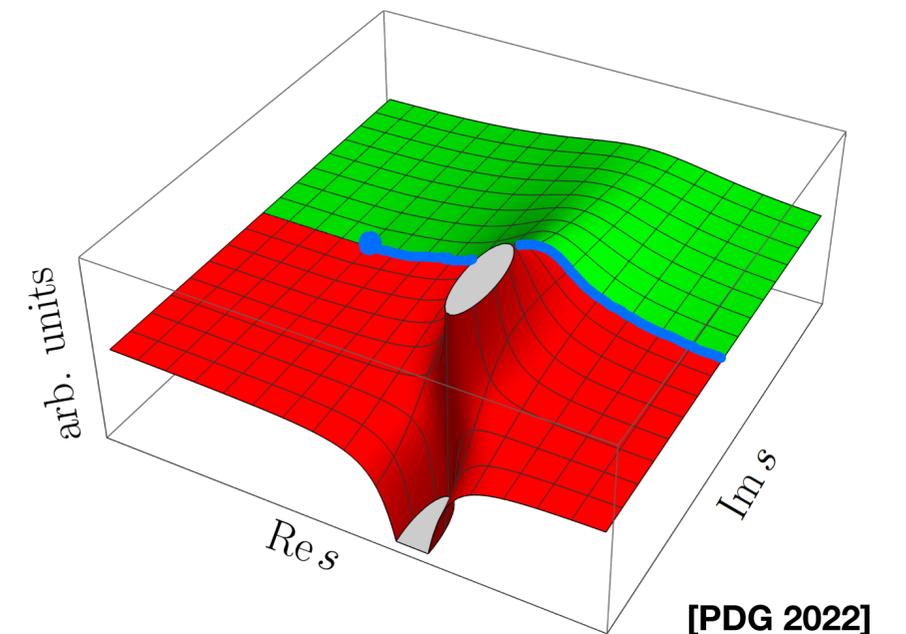
Summary and outlook

- ☑ Lattice QCD provides a first-principle tool to investigate the hadron spectrum
- ☑ Access to quantities that are hard experimentally, e.g. hyperons
- ☑ Several studies of scattering lattice QCD: $\Delta(1232)$, $\Lambda(1405)$
- ☑ First results on three-particle resonances on a toy model



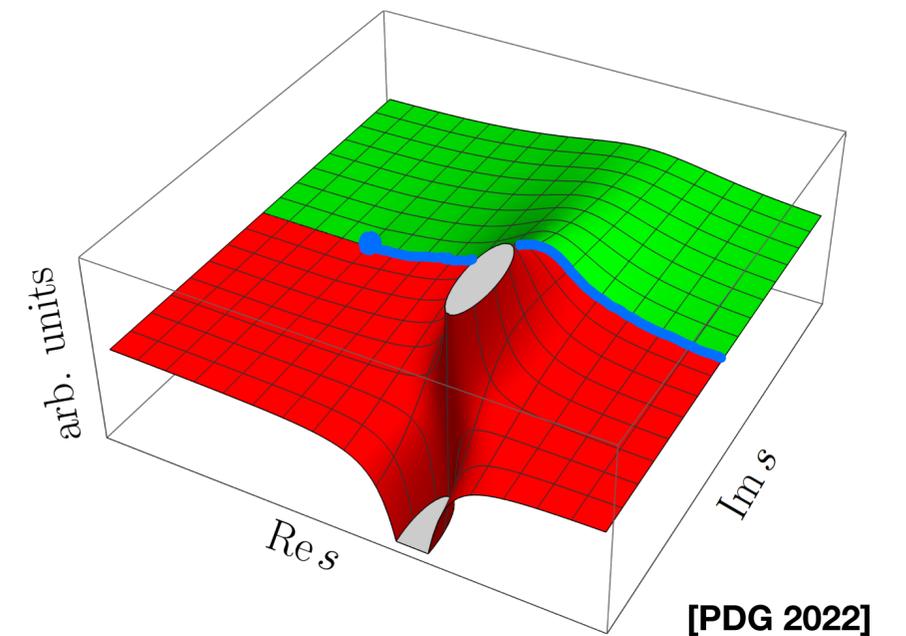
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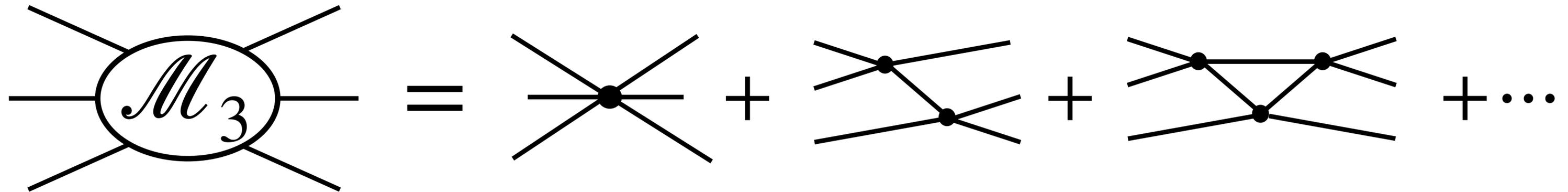


Thanks!

Backup

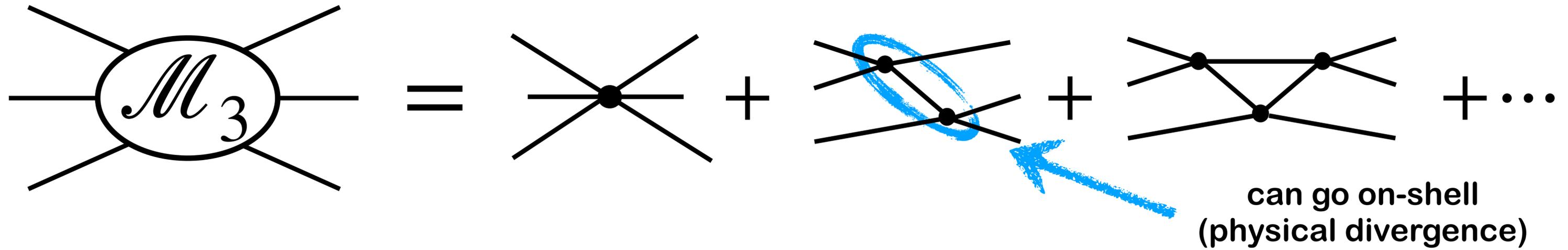
Three-particle scattering

Qualitatively more complicated than the two-particle case!



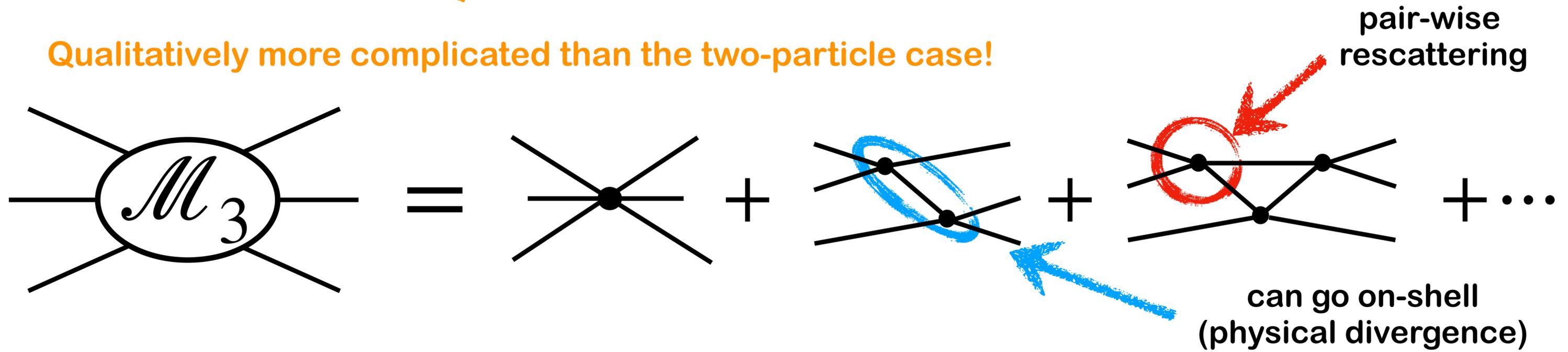
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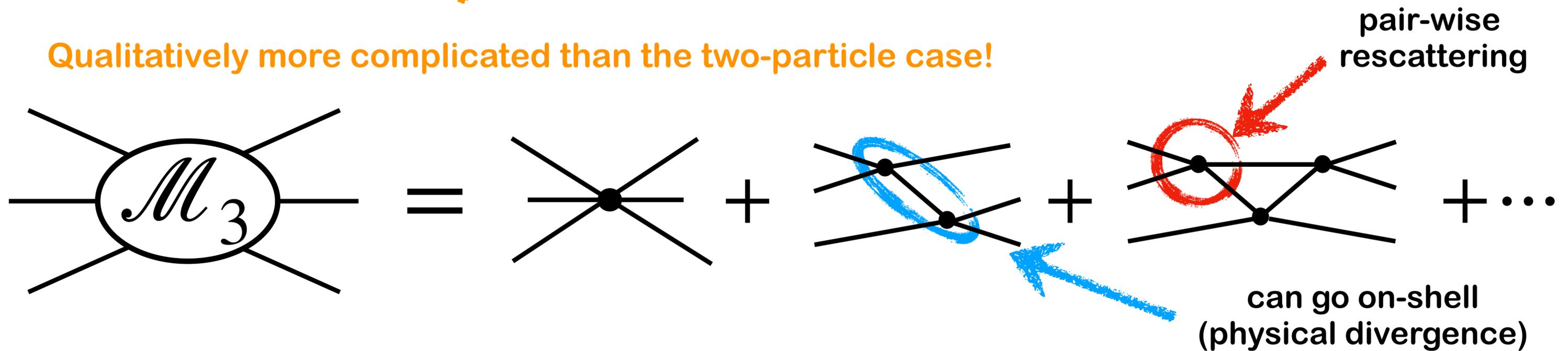
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Finite-volume formalism has been developed independently by three groups

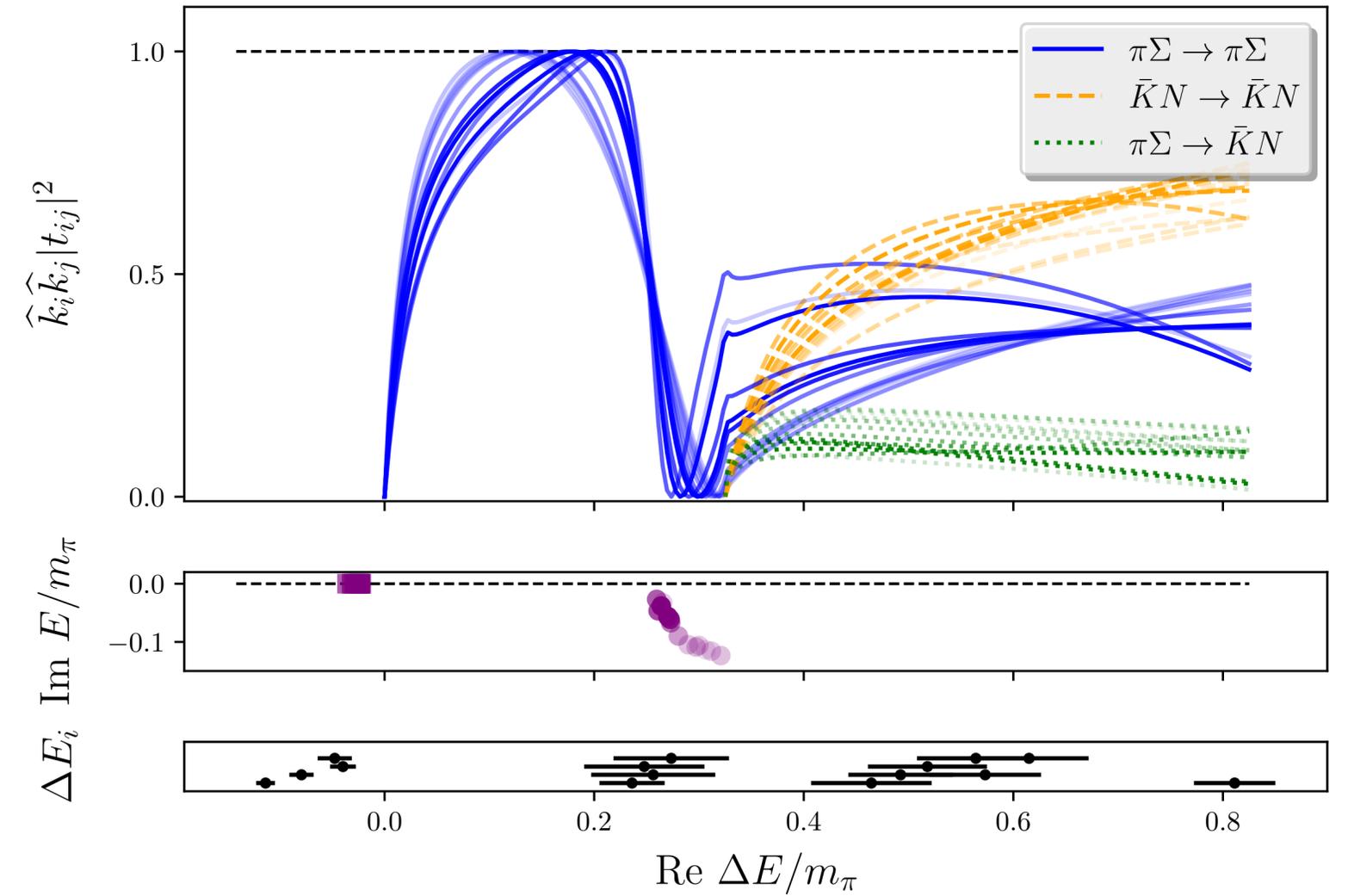
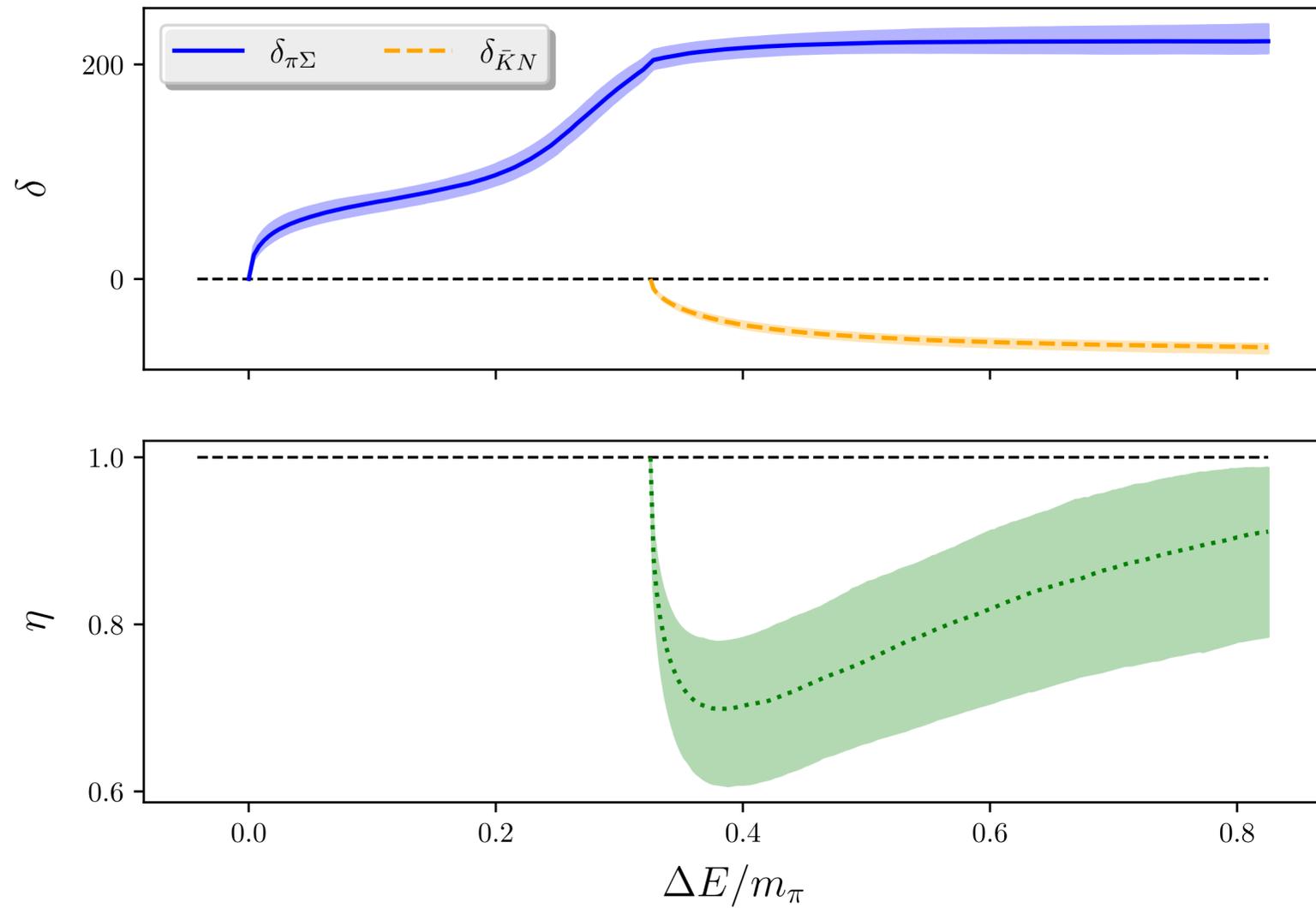
● **Generic Relativistic Field Theory (RFT)** [Hansen, Sharpe, PRD 2014 & 2015]

● **Non-Relativistic EFT (NREFT)** [Hammer, Pang, Rusetsky, JHEP 2017] x 2

● **Finite-Volume Unitarity (FVU)** [Mai, Döring, EPJA 2017]

Equivalence has been established
[Jackura et al. PRD 2019], [Blanton, Sharpe, PRD 2020],
[Jackura, 2208.10587]

The $\Lambda(1405)$ resonance



Poles

$$\mathcal{M}_2(s) = \frac{16\pi\sqrt{s}}{k \cot \delta(k) - ik},$$
$$k = \pm \sqrt{k^2}$$

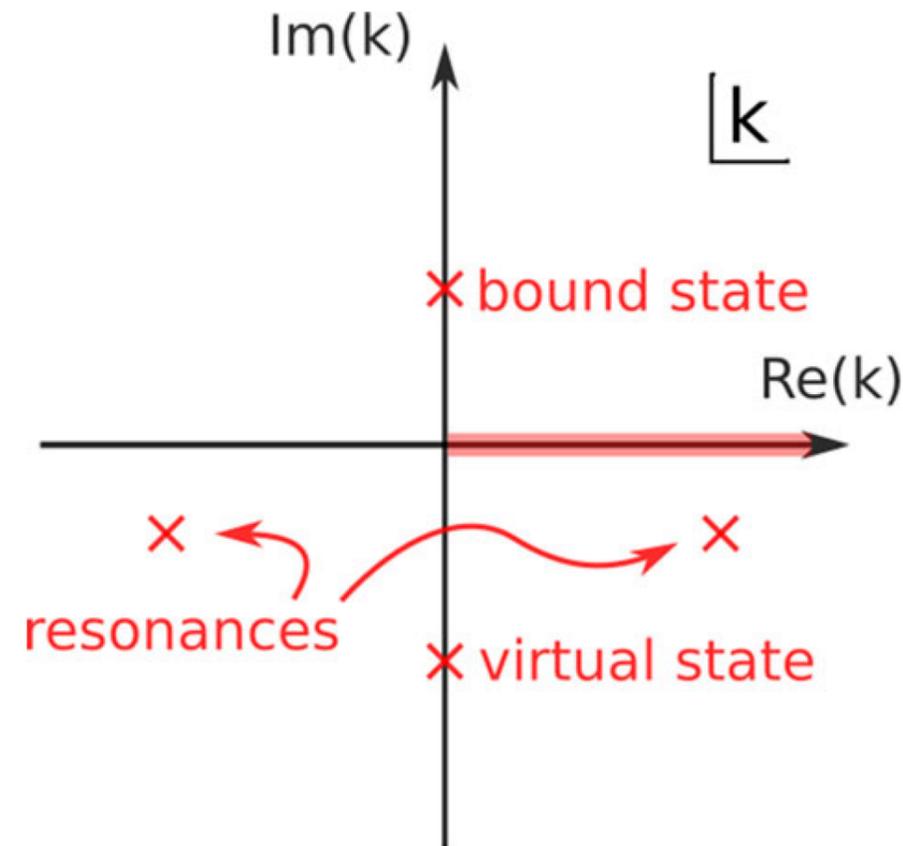


Fig. 1 Naming convention for the poles in the k -plane. The thick red line for positive real valued k marks the physical momenta in the scattering regime

[Matuschek et al, EPJA 2021]

Applying the three-body formalism

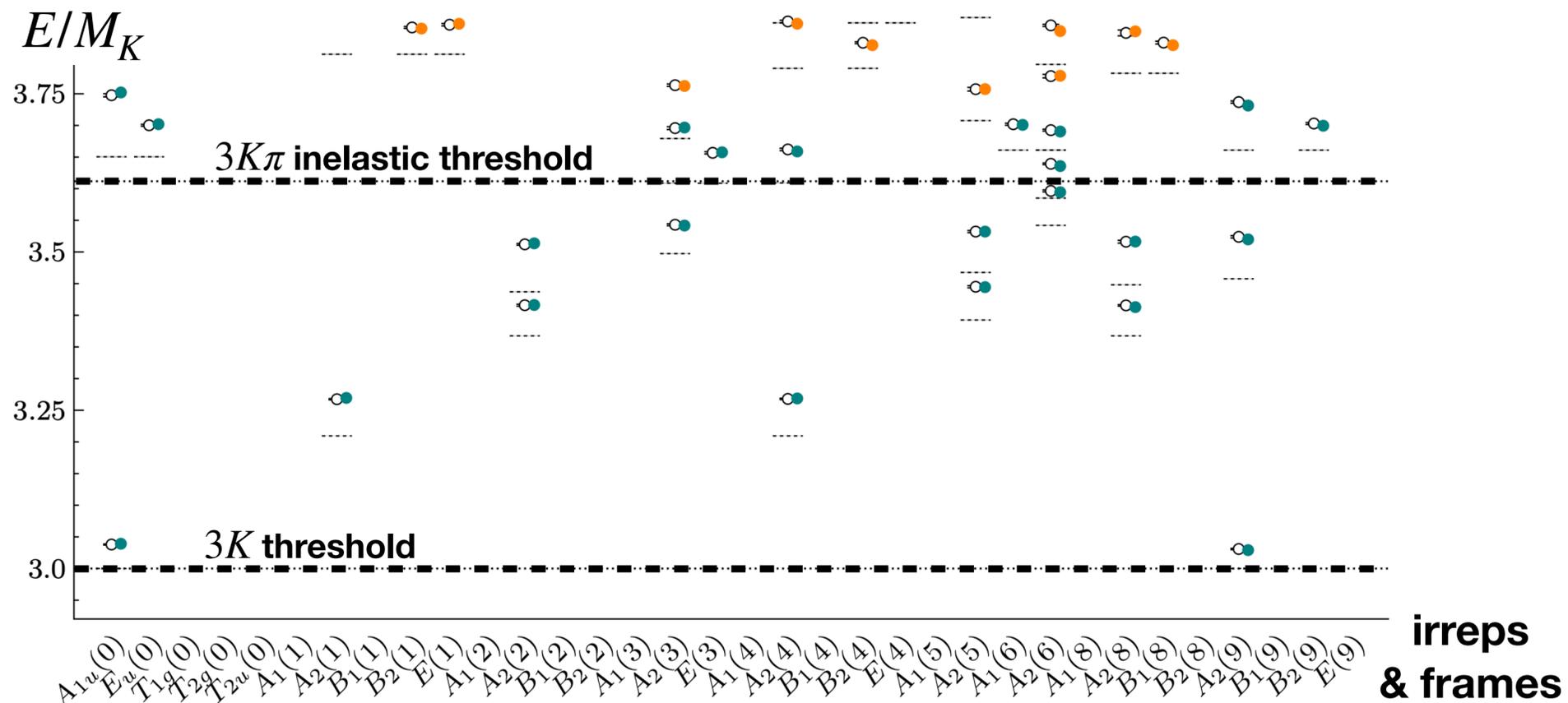
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- Requires large sets of energy levels obtained using variational techniques

Using stochastic LapH method [Morningstar et al, PRD 2011]

3K⁺ energy levels



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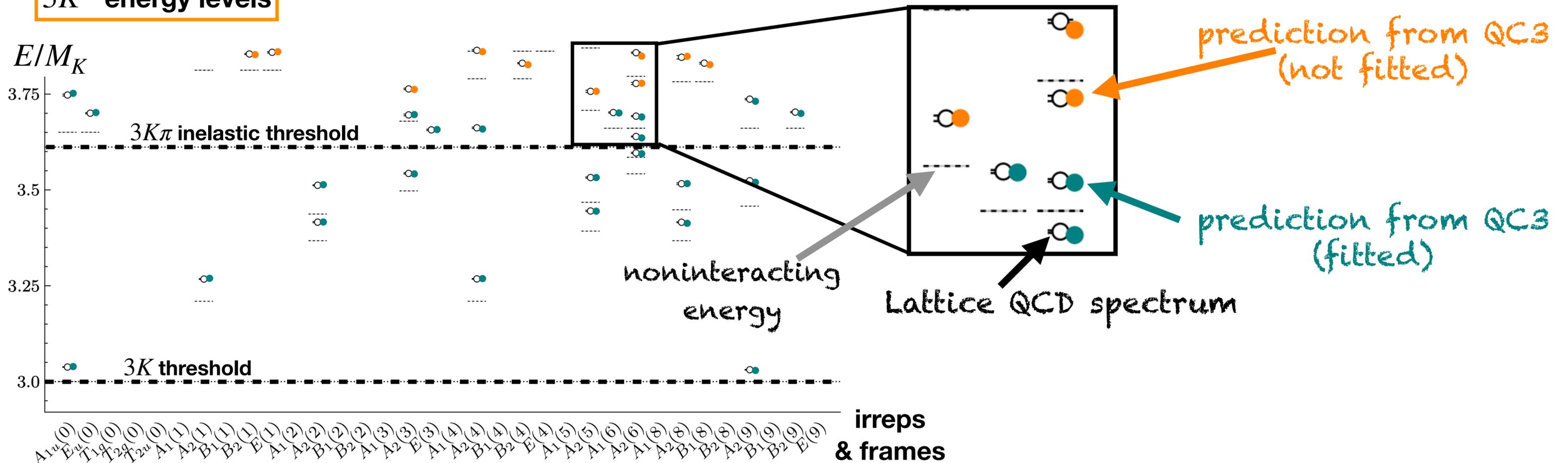
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[Alexandrou et al, PRD 2020], [Blanton et al., PRL 2020 & JHEP 2021], [Brett et al, PRD 2021], [Culver et al, PRD 2021], [Fischer et al, EPJC 2021], [Hansen et al, PRL 2021], [Mai et a PRL 2019 & 2021]

- Requires large sets of energy levels obtained using variational techniques

Using stochastic LapH method [Morningstar et al, PRD 2011]

3K⁺ energy levels



Applying the three-body formalism

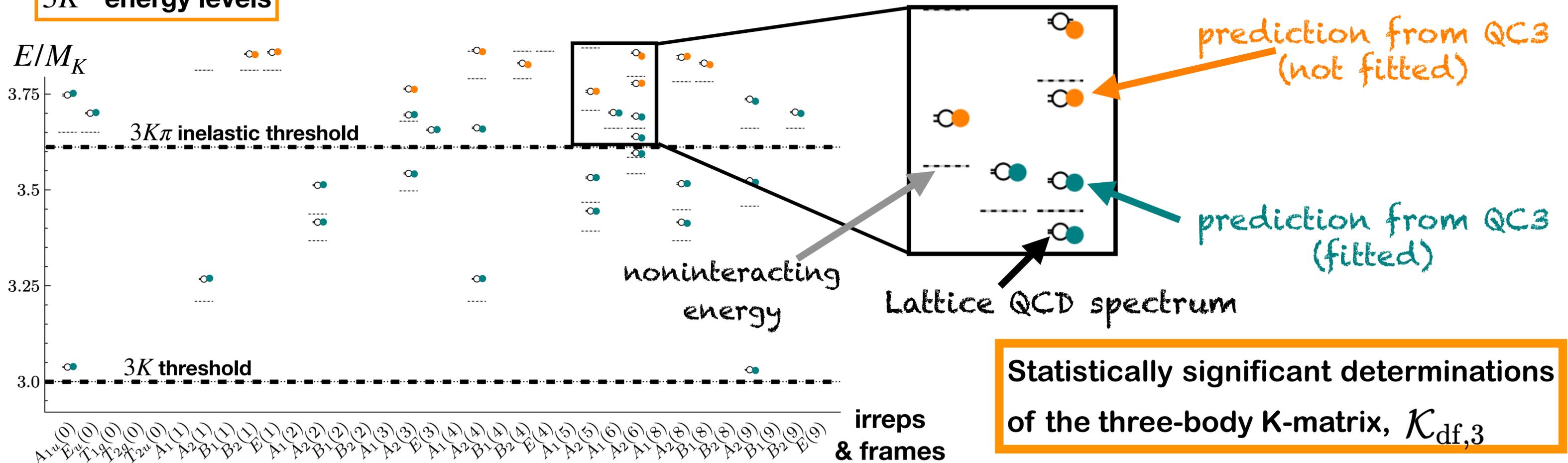
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Statistically significant determinations of the three-body K-matrix, $\mathcal{K}_{df,3}$

$\mathcal{K}_{df,3}$ from Lattice QCD

$$\mathcal{K}_{df,3} = \underbrace{\mathcal{K}_{df,3}^{\text{iso},0} + \mathcal{K}_{df,3}^{\text{iso},1} \Delta + \mathcal{K}_{df,3}^{\text{iso},2} \Delta^2}_{\text{Depend of CM energy}} + \underbrace{\mathcal{K}_A \Delta_A + \mathcal{K}_B \Delta_B}_{\text{Angular dependence}},$$

$\Delta \equiv \frac{s - 9m^2}{9m^2}$

$\Delta_B = \sum_{i,j=1}^3 \tilde{t}_{ij}^2 - \Delta^2,$

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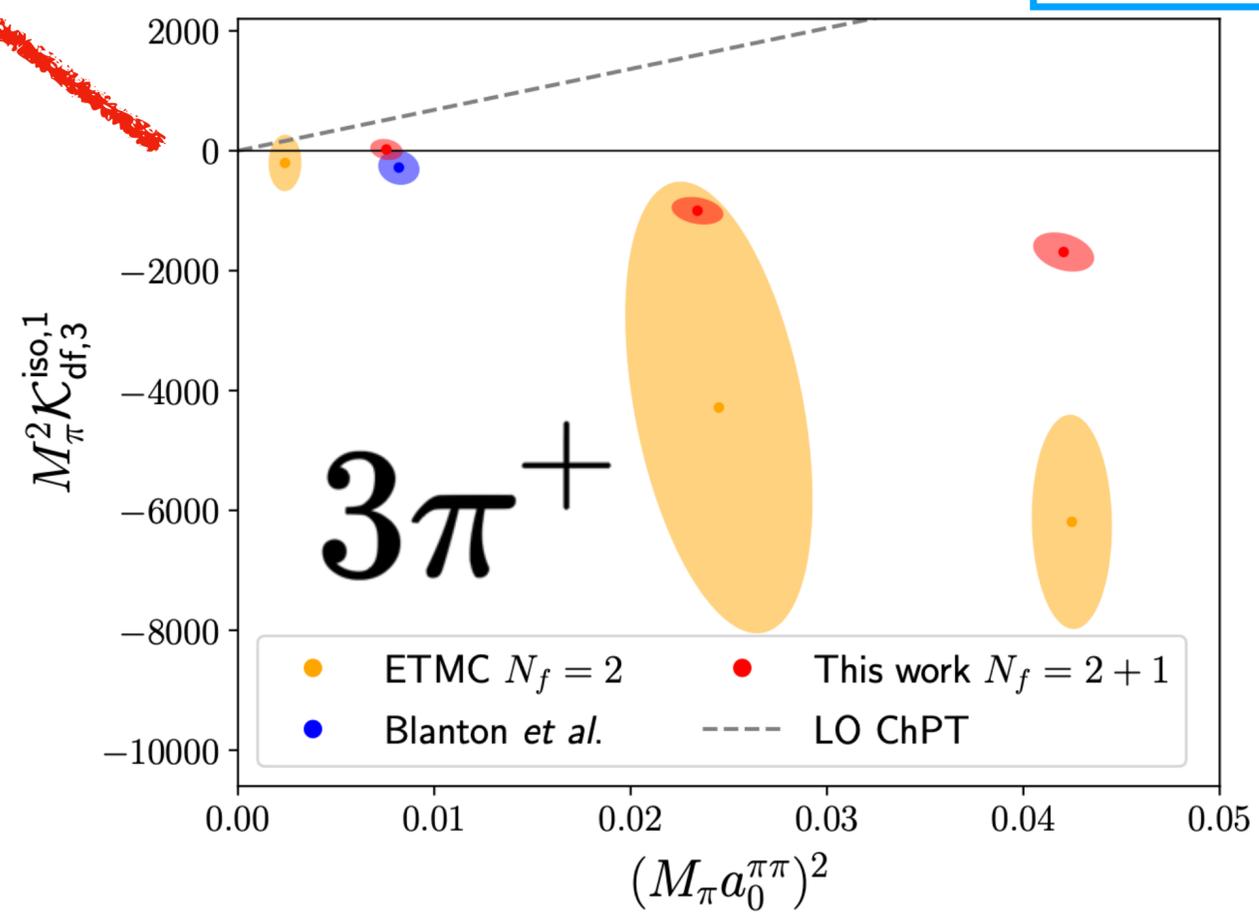
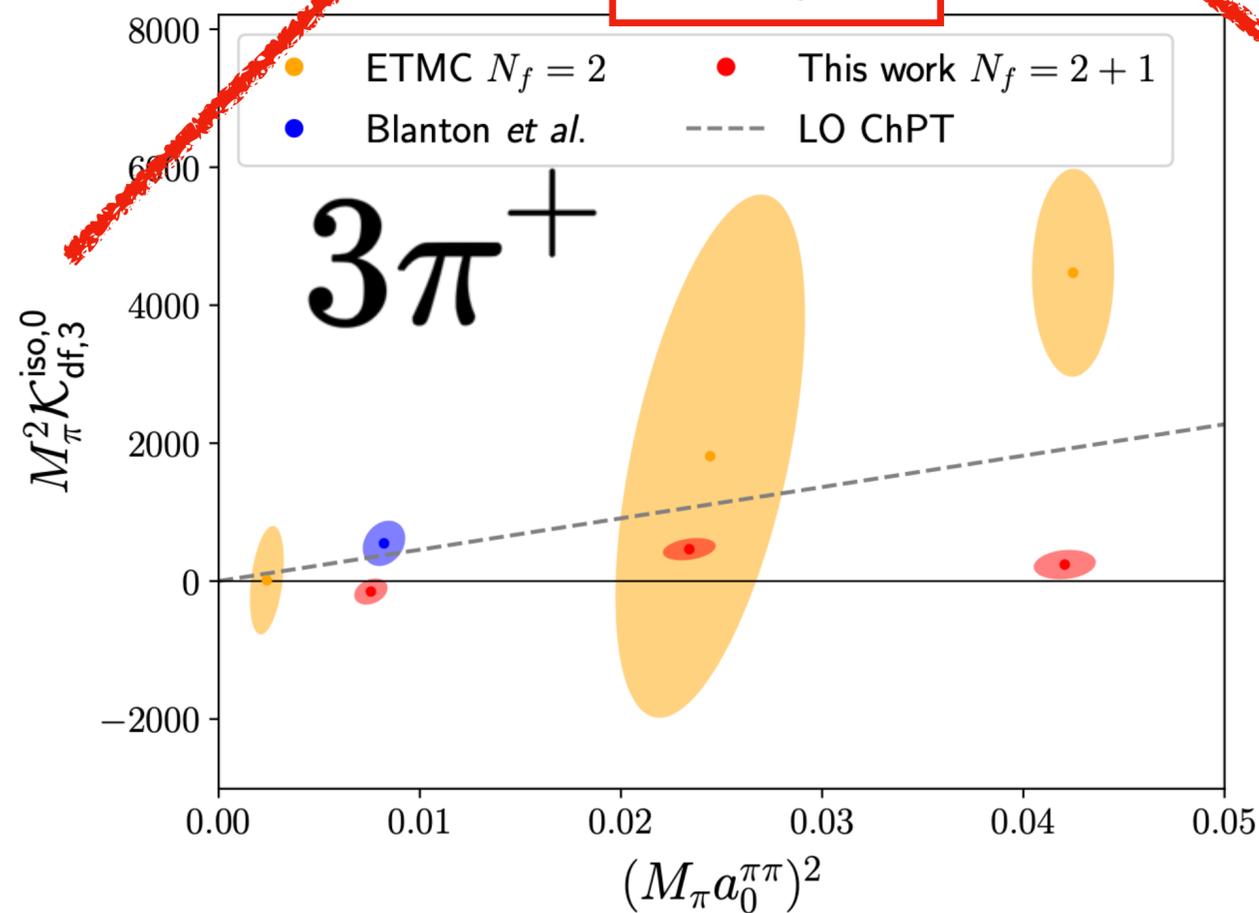
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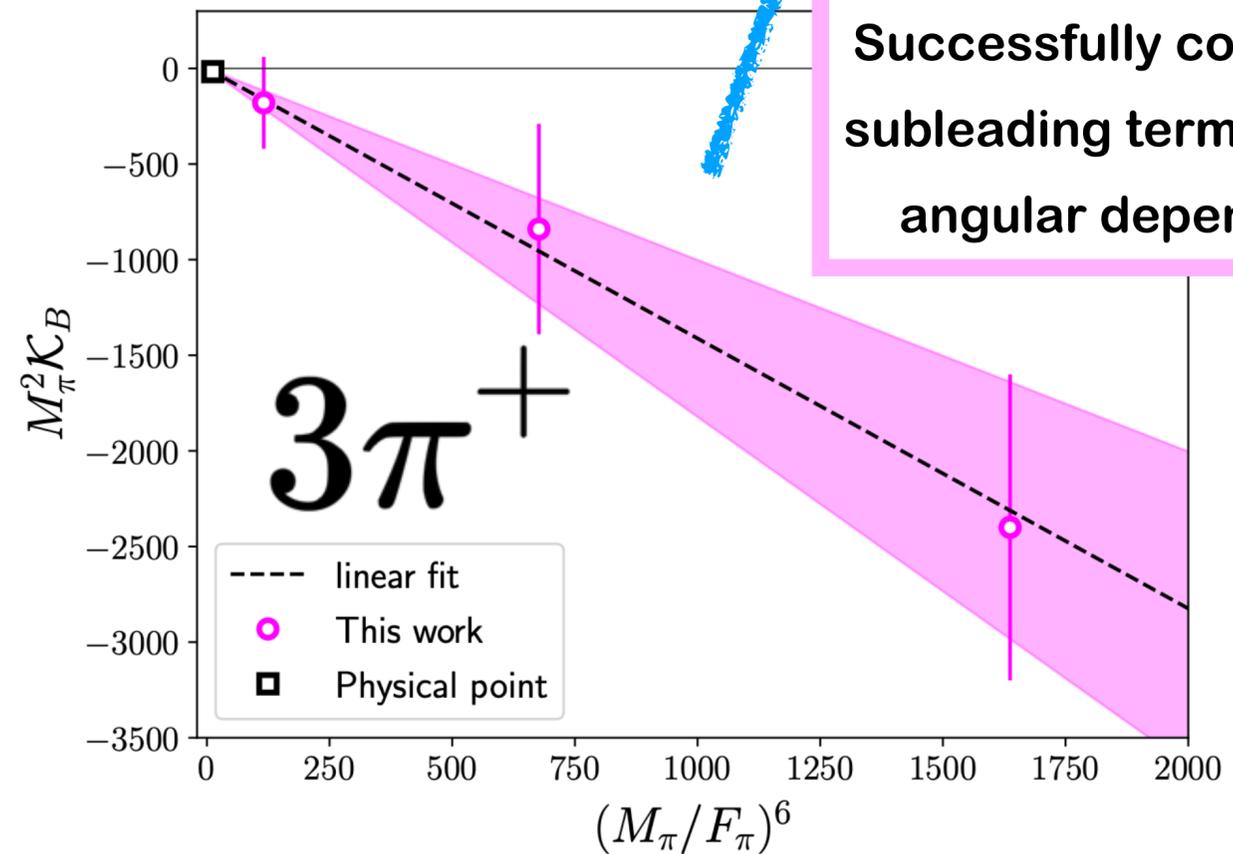
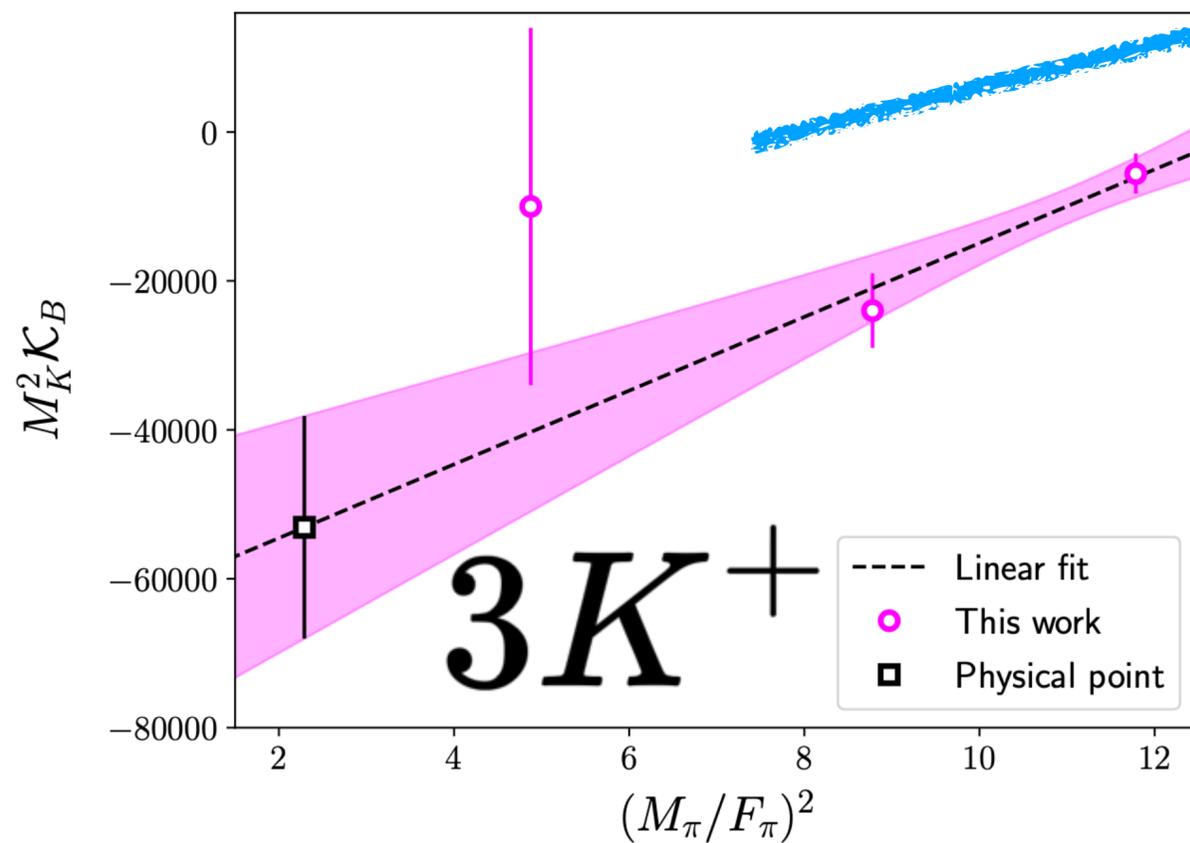


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Successfully constrained subleading term including angular dependence!

Nondegenerate systems

○ Relevant three-body systems involve nonidentical particles ($\pi\pi N$)

○ First step: formalism for three nonidentical scalars

e.g. $\pi^+\pi^0\pi^-$, $K^+K^+\pi^+$, $D_s^+D^0\pi^-$

[Blanton, Sharpe, PRD 2021] x 2, [Hansen, FRL, Sharpe, JHEP 2021], [Mai et al (GWQCD), PRL 2021]

$$\det_{k,\ell,m,\mathbf{f}} [1 - \mathbf{K}_{\text{df},3}(E^*)\mathbf{F}_3(E, \mathbf{P}, L)] = 0$$

determinant runs over an additional “flavor” index

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Example:
 $\pi^+\pi^+K^+$ scattering

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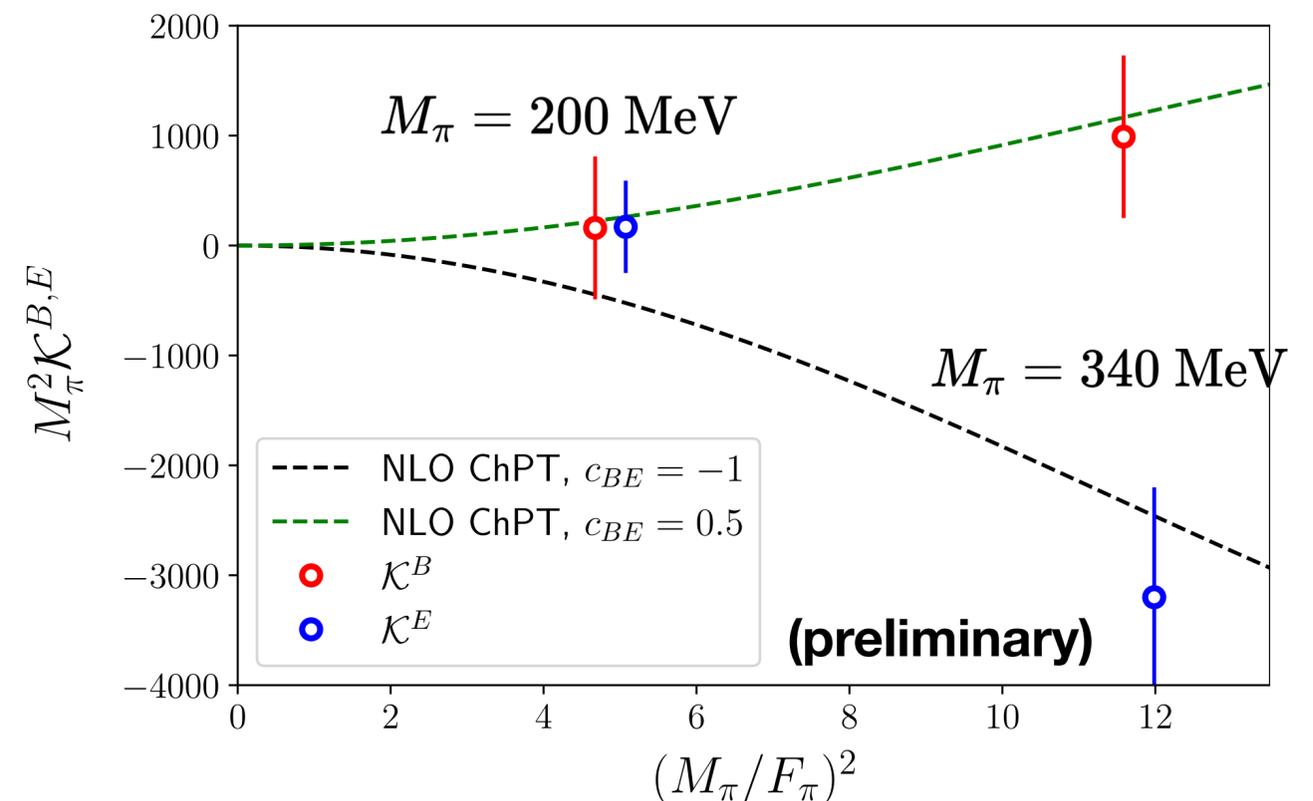
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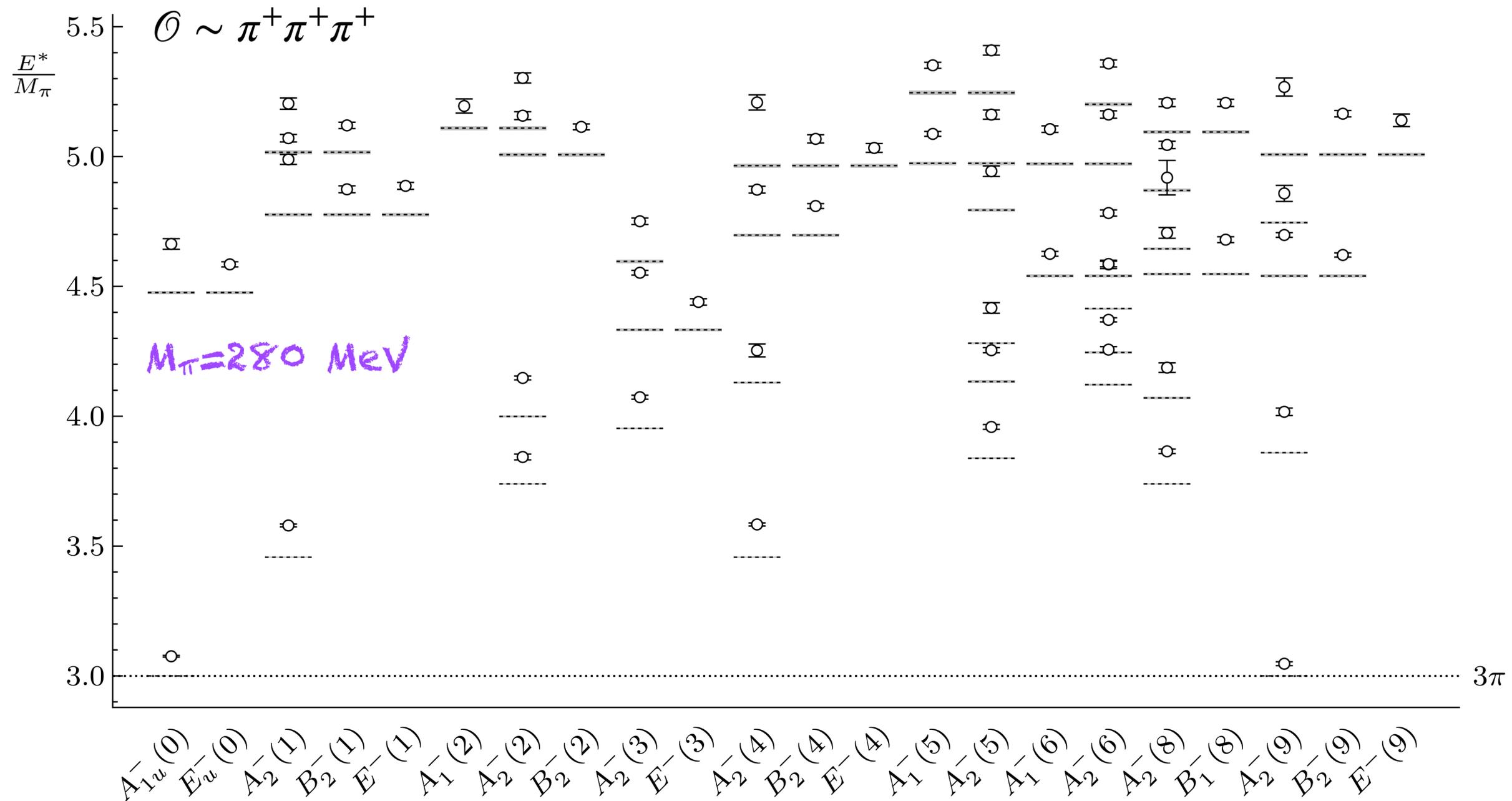
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[Draper, Hanlon, Hörz, Morningstar, FRL, Sharpe (in preparation)]

[Talk by S. Sharpe]

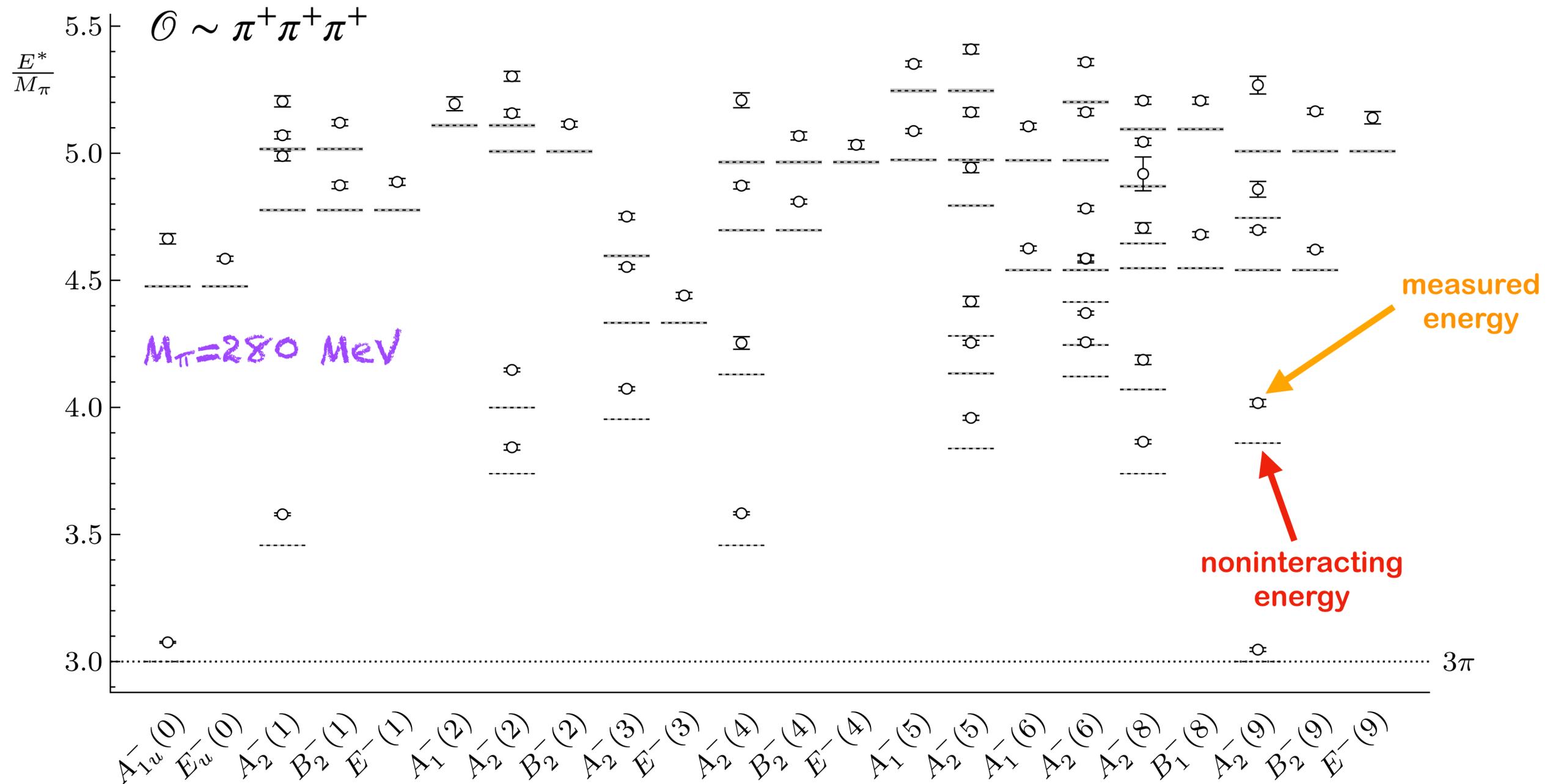


3π⁺ energy levels



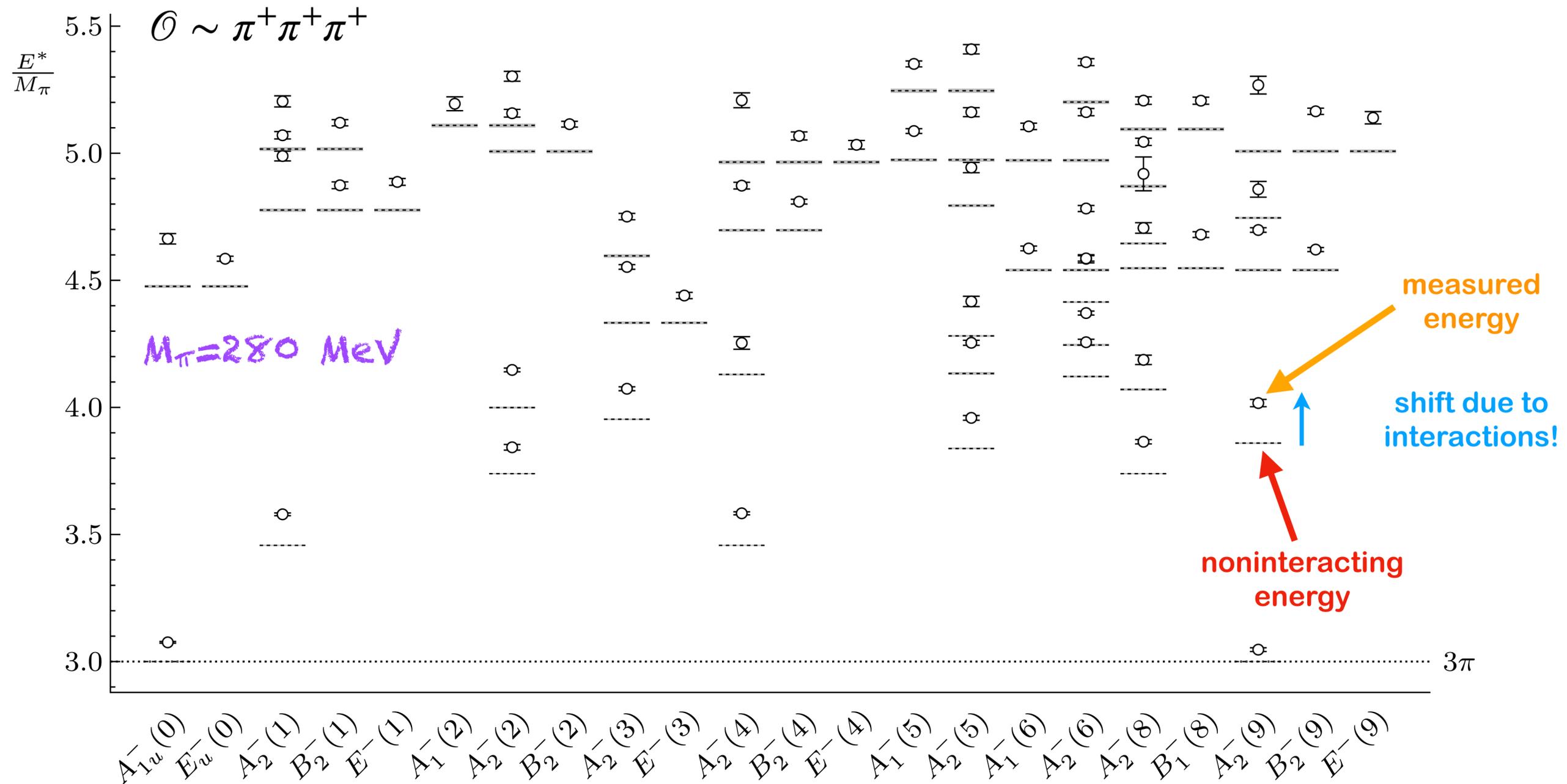
[Blanton, Hanlon, Hörz, Morningstar, FRL, Sharpe]

$3\pi^+$ energy levels



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3π⁺ energy levels



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Integral equations (RFT)

Final step

Physical 3- \rightarrow 3
amplitude

$\mathcal{K}_2, \mathcal{K}_{df,3}$



Integral
equations

\mathcal{M}_3

Integral equations (RFT)

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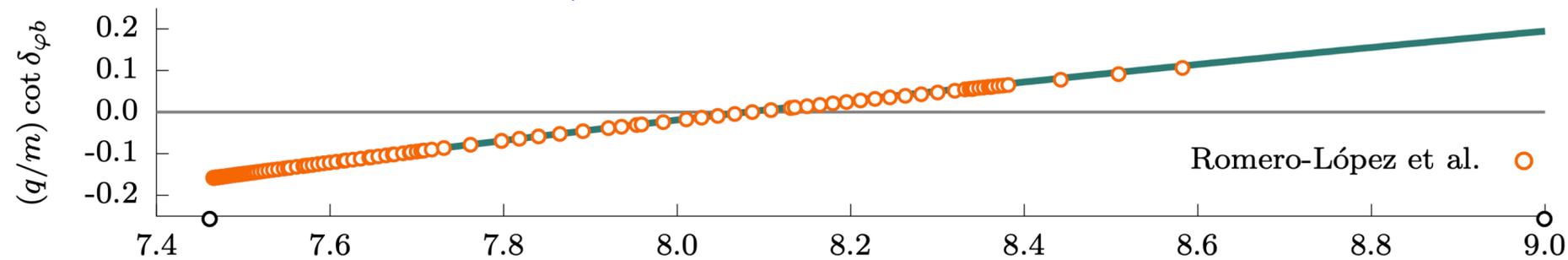
$\mathcal{K}_2, \mathcal{K}_{df,3}$



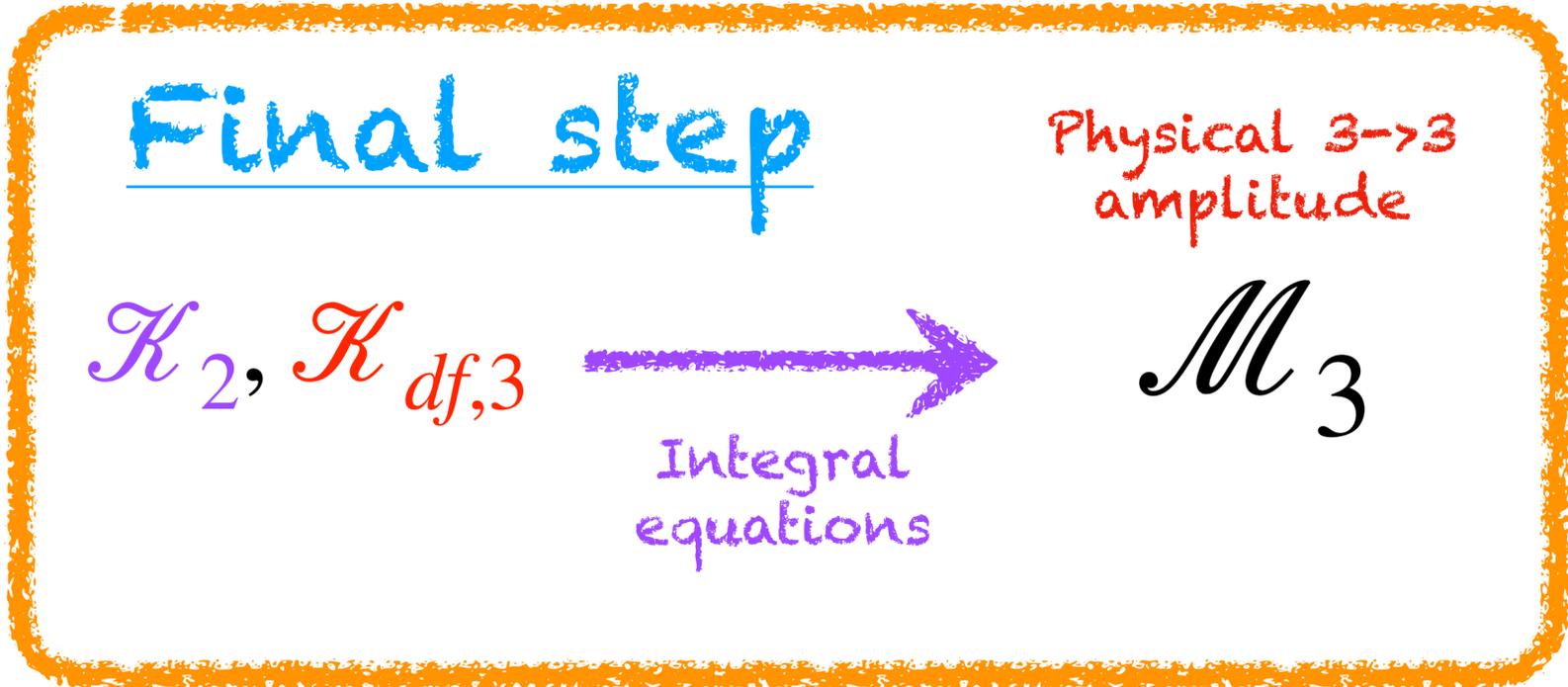
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Integral
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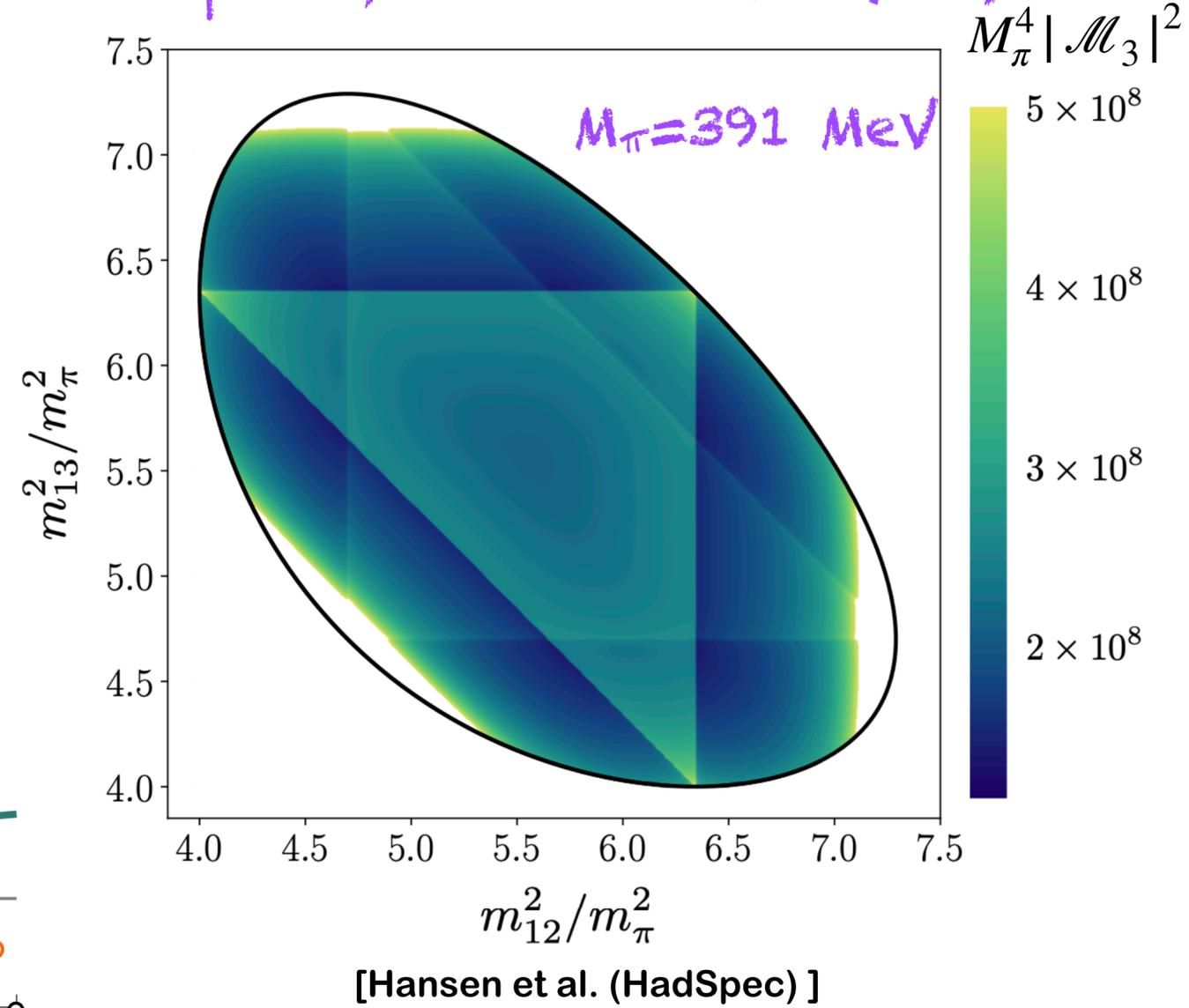
Particle-Dimer phase shift [Jackura et al.]



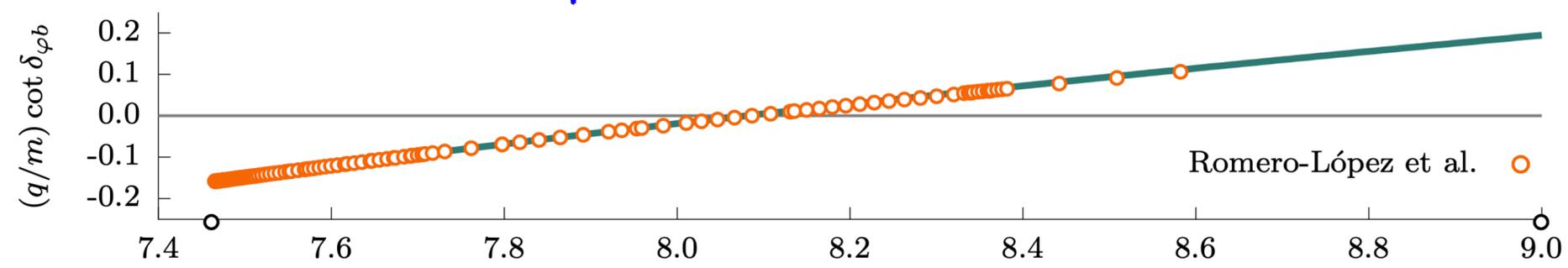
Integral equations (RFT)



Dalitz plots from lattice QCD ($3\pi^+$)



Particle-Dimer phase shift [Jackura et al.]

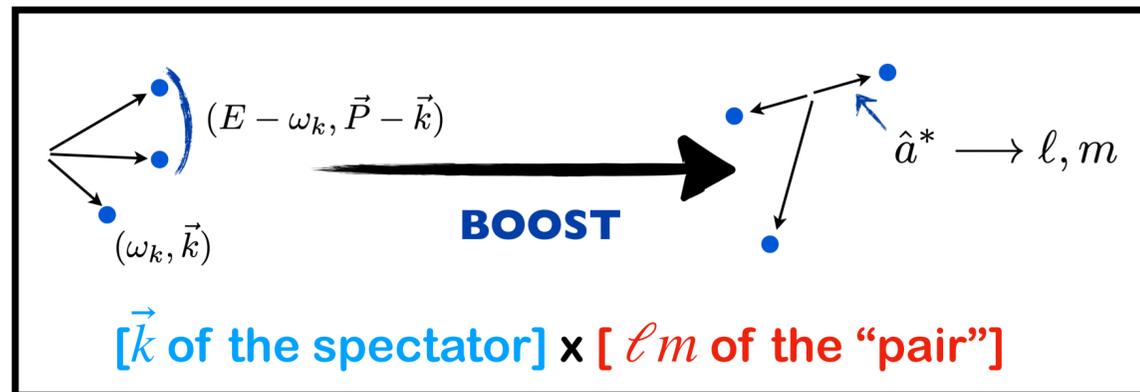


Quantization Condition

$$\det_{klm} [\mathcal{K}_{df,3} + F_3^{-1}] = 0$$

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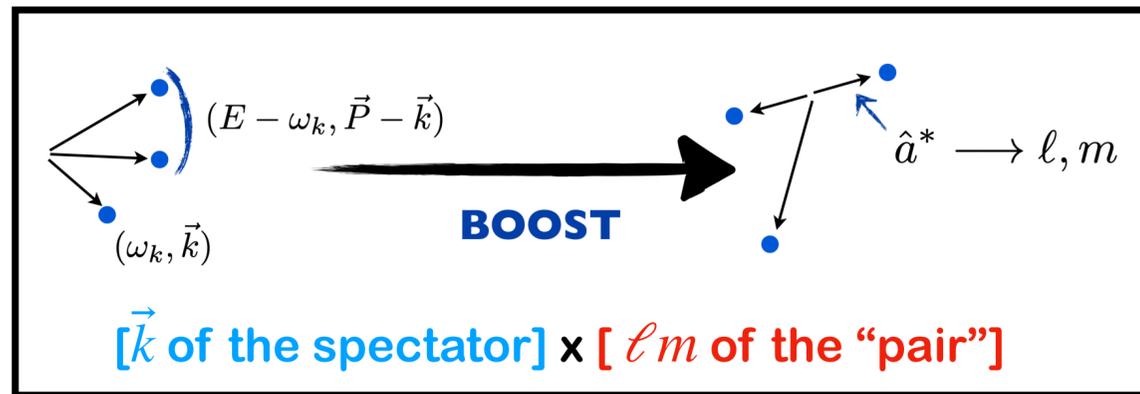
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Quantization Condition

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Finite-volume information & two-body interactions



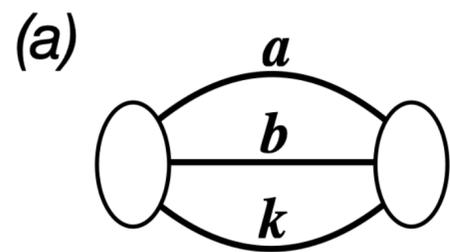
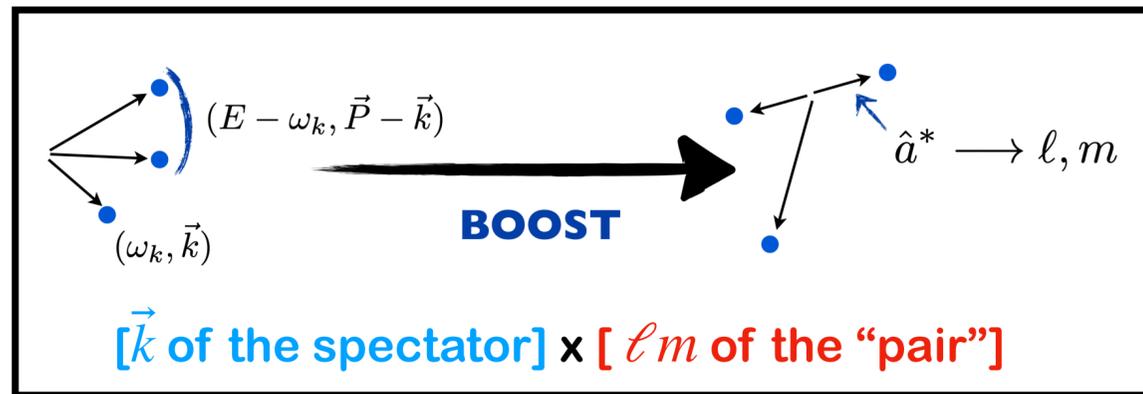
$$F_3 = \frac{1}{L^3} \left[\frac{F}{3} - F \frac{1}{(\mathcal{K}_2)^{-1} + F + G} F \right]$$

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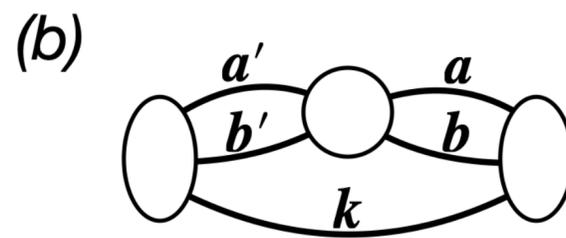
$$\det_{klm} [\mathcal{K}_{\text{df},3} - F_3^{-1}] = 0$$

Finite-volume information & two-body interactions

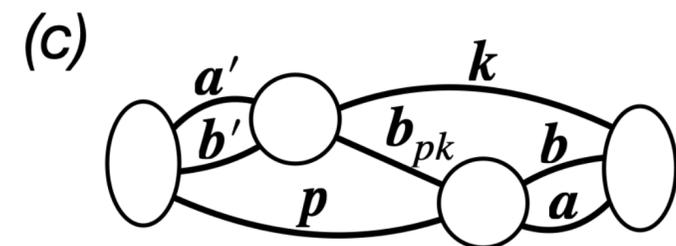
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F



\mathcal{K}_2



G

$$F_{00}(q^2) \sim \left[\frac{1}{L^3} \sum_{\vec{k}} - \int \frac{d^3k}{(2\pi)^3} \right] \frac{1}{k^2 - q^2}$$

$$\mathcal{K}_2^\ell = \frac{16\pi\sqrt{s}}{q^{2\ell+1} \cot \delta_\ell}$$

$$G_{p00;k00} \equiv \frac{1}{L^3} \frac{1}{2\omega_p} \frac{H(\vec{p})H(\vec{k})}{b_{pk}^2 - m^2} \frac{1}{2\omega_k}$$

Including isospin

○ Relevant three-body systems involve nonidentical particles ($\pi\pi N$)

○ Let us consider mass-degenerate pions with different flavor e.g. $\pi^+\pi^0\pi^-$

[Hansen, FRL, Sharpe, JHEP 2020]

▶ All pions have the same mass

▶ Overall isospin is conserved

▶ Presence of resonances

▶ Example of multi-channel scattering

