

Multigluon Correlation Functions from lattice QCD

The four gluon vertex

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From first-principles QCD to experiments, Trento, Italy, May 2023

Motivation

QCD Dynamics at a fundamental level

importance of higher order corrections (two loop contributions to gluon DSE, etc.)

Test our understanding of QCD Green functions (ghost dominance at IR)

Effective charge $\frac{g^2}{4\pi} \Gamma^{(4g)}(p^2) p^2 D(p^2)$

Landau gauge

(pure gauge) Lattice point of view

Lattice Approach

Importance Sampling **allows to access the gluon field** $A_{\mu}^a(x)$

$$\mathcal{G}^{(n)}(x_1, \dots, x_n) = \langle 0 | T \left(A_{\mu_1}^{a_1}(x_1) \cdots A_{\mu_n}^{a_n}(x_n) \right) | 0 \rangle$$

Larger n:

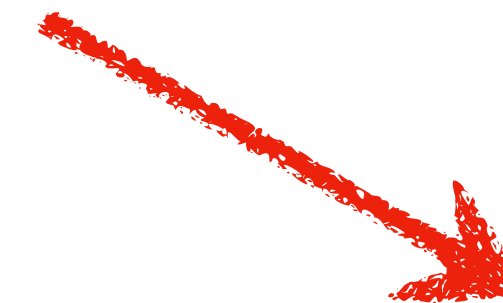


Noisy

Disconnected parts



Larger sets of configurations

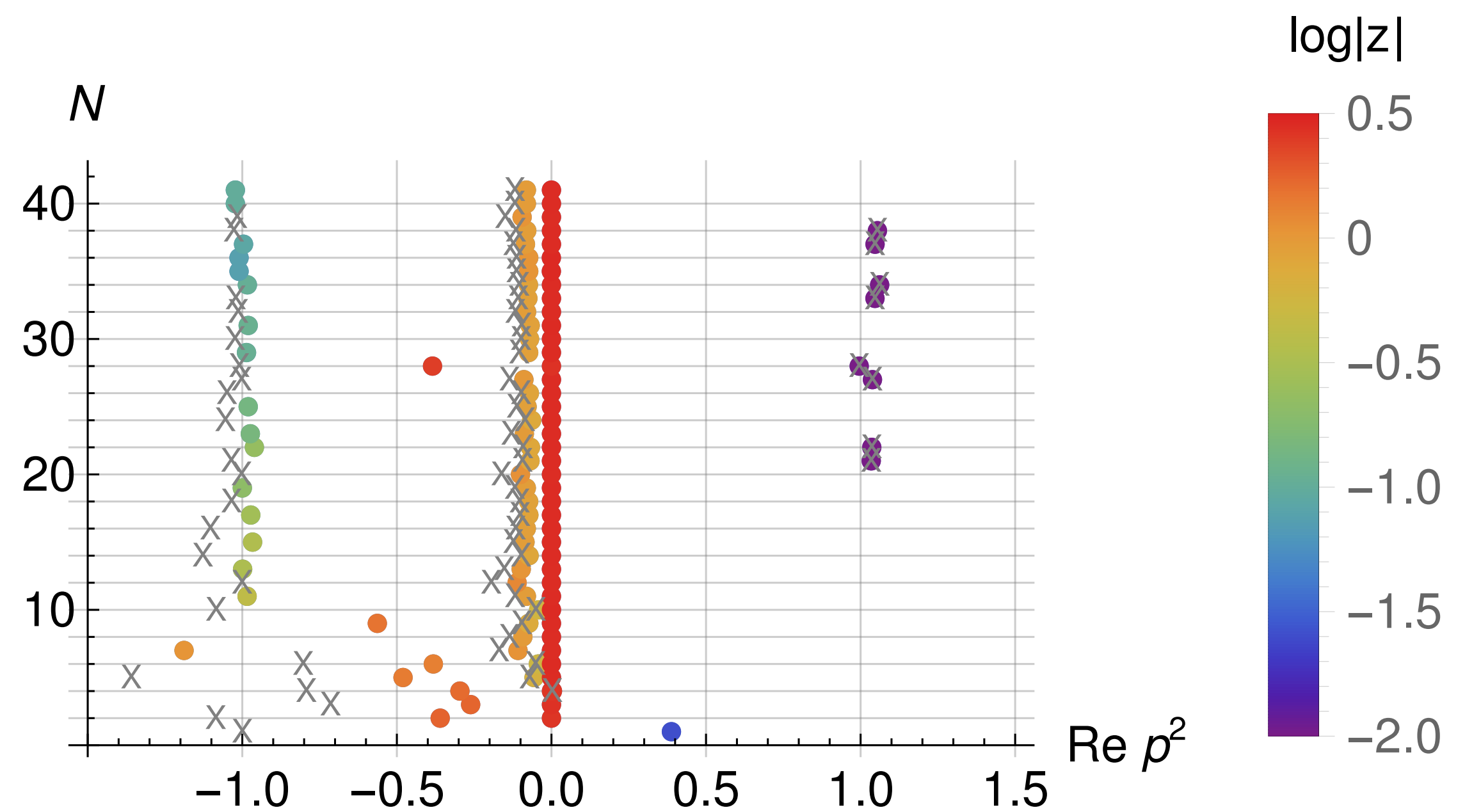
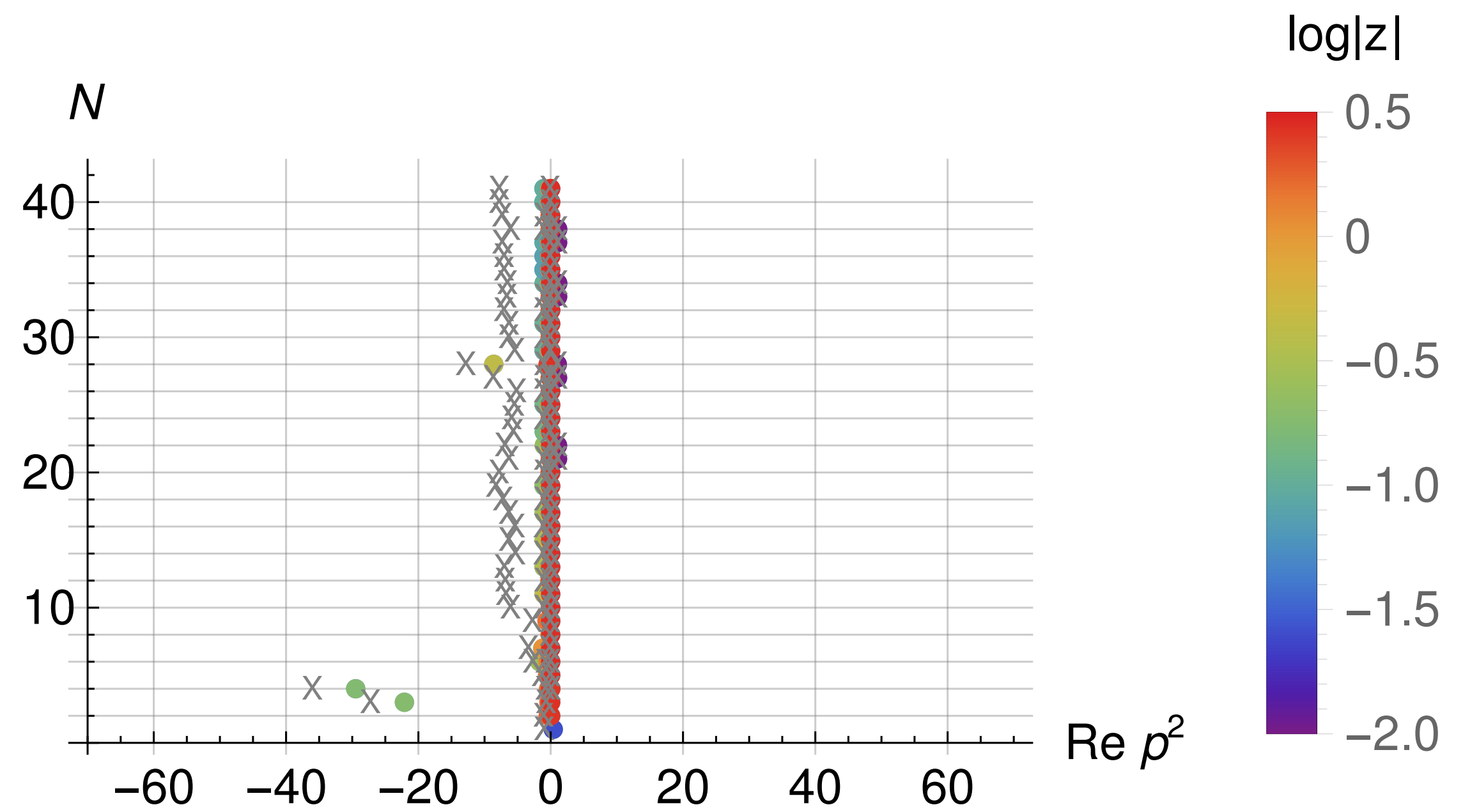


kinematical configurations
or
Combinations of components

Massless Ghost

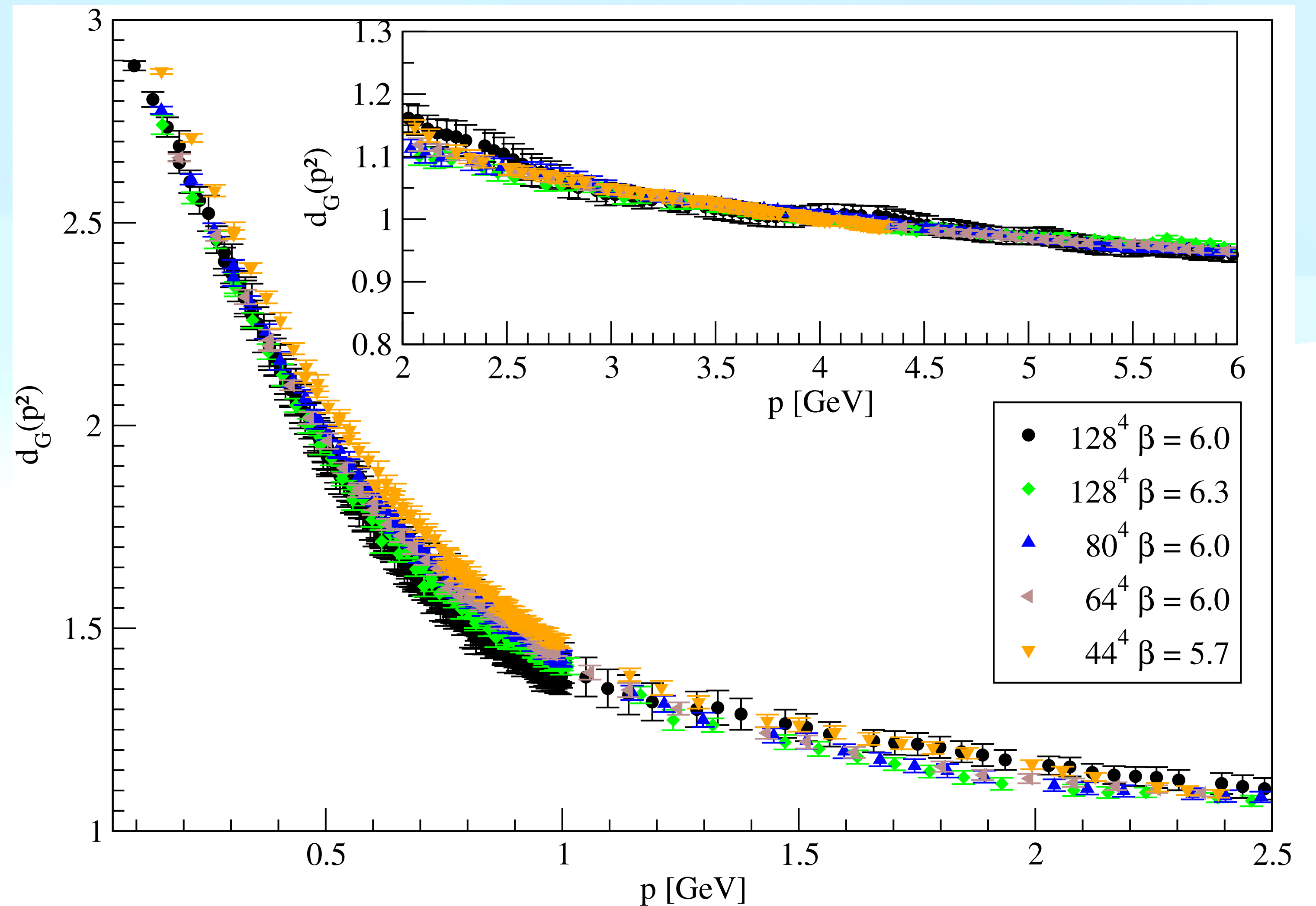
$$D^{ab}(p^2) = -\delta^{ab} \frac{d_G(p^2)}{p^2}$$

lattice studies support a massless ghost



IR Ghost Dominance / Log Divergence

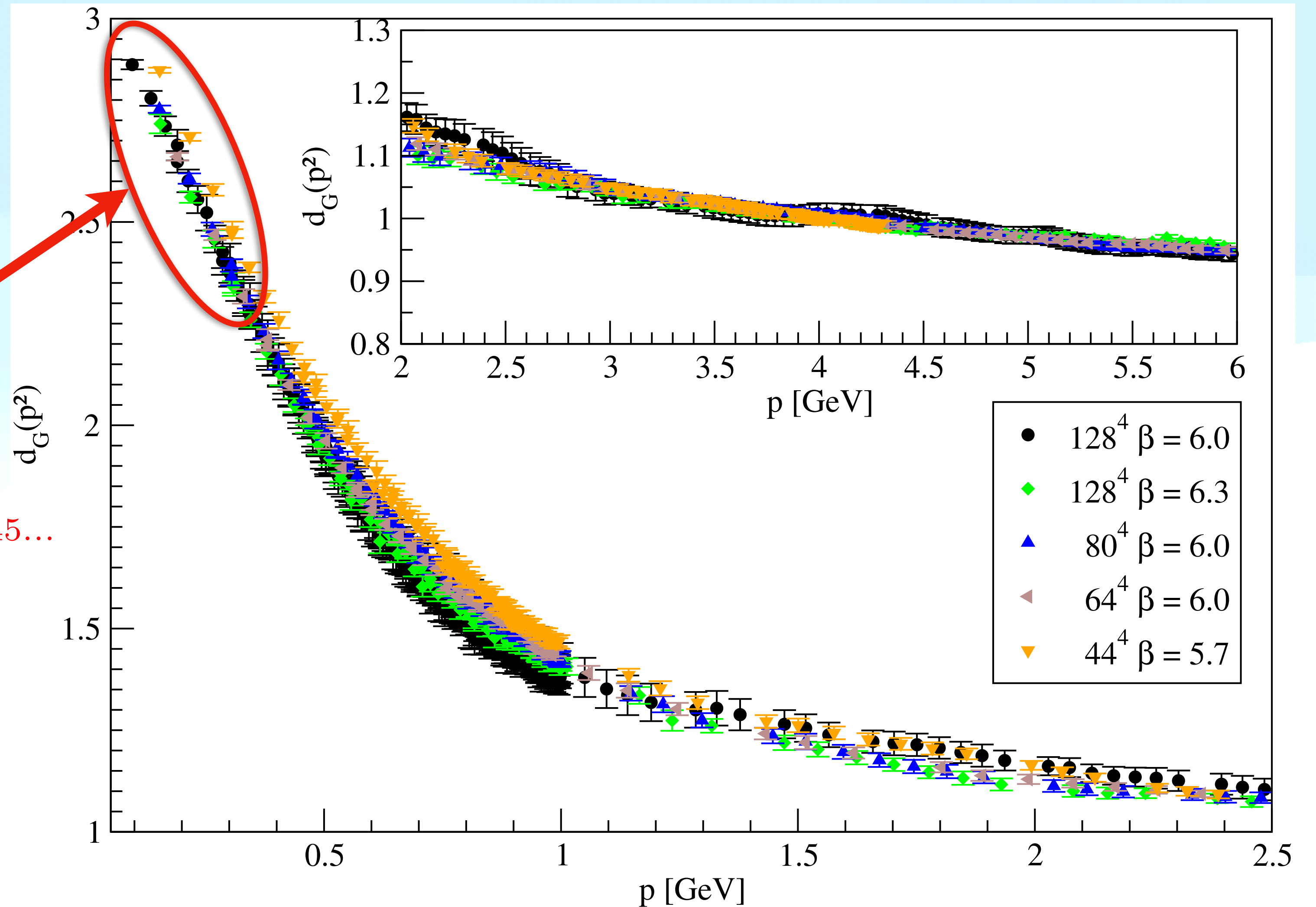
$$D^{ab}(p^2) = -\delta^{ab} \frac{d_G(p^2)}{p^2}$$



IR Ghost Dominance / Log Divergence

$$D^{ab}(p^2) = -\delta^{ab} \frac{d_G(p^2)}{p^2}$$

$$d_G(p^2) \propto \left(\ln \frac{p^2 + (1.7 \text{ GeV})^2}{(0.4 \text{ GeV})^2} \right)^{-0.2045\dots}$$




IR Ghost Dominance / Log Divergence

Ward identities/Slavnov-Taylor identities

See Ferreira talk

Ward identity displacement in QCD

$$q^\alpha \Pi_{\alpha\mu\nu}(q, r, p) = F(q^2) [\Delta^{-1}(p^2) P_\nu^\sigma(p) H_{\sigma\mu}(p, q, r) - \Delta^{-1}(r^2) P_\mu^\sigma(r) H_{\sigma\nu}(r, q, p)]$$

$q \rightarrow 0$ 

Ward identity

$$L_{\text{sg}}(r^2) = F(0) \left[\frac{\mathcal{W}(r^2)}{r^2} \Delta^{-1}(r^2) + \frac{\partial \Delta^{-1}(r^2)}{\partial r^2} \right] + \mathbb{C}(r^2)$$

Displacement = BS amplitude

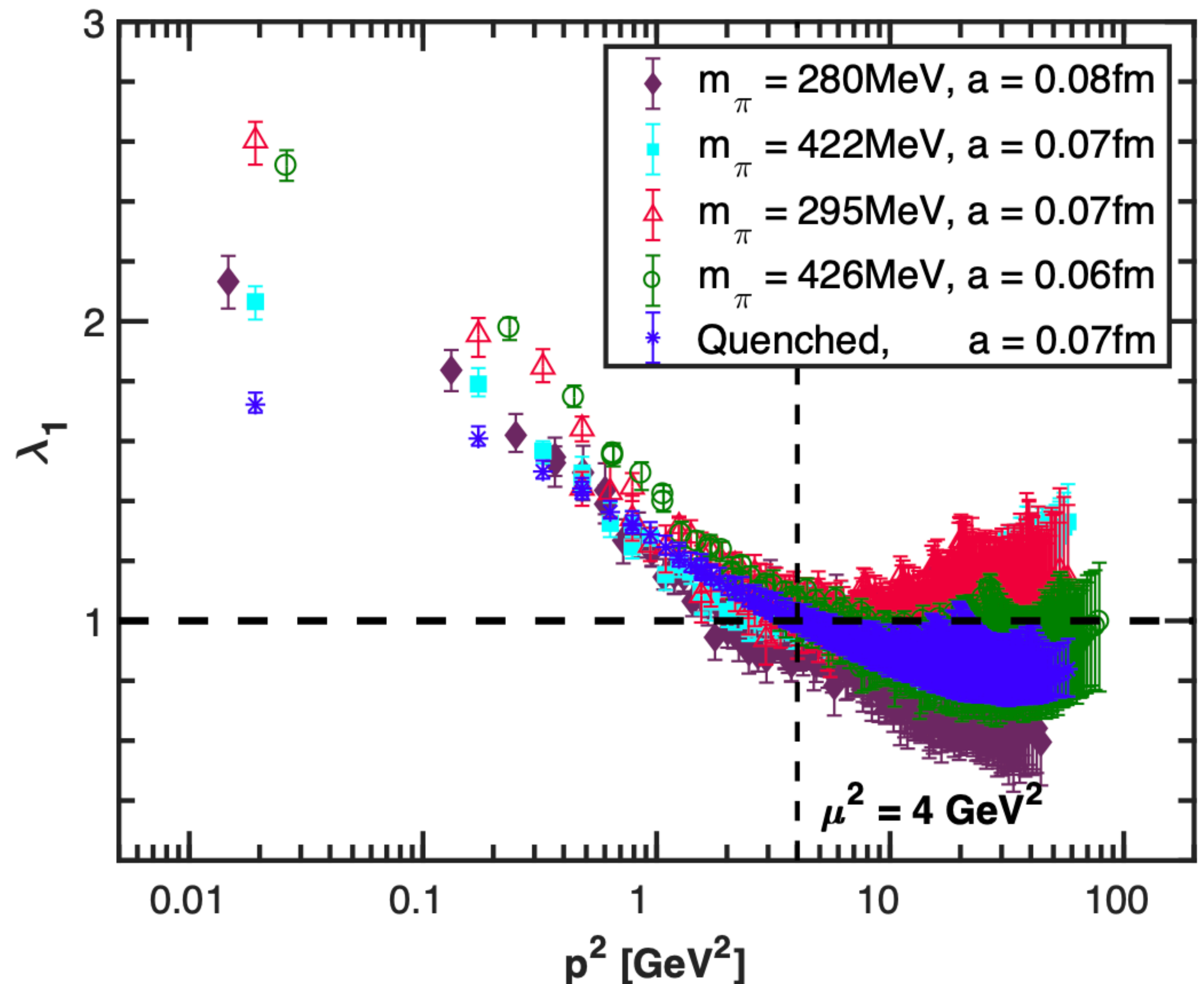
IR Ghost Dominance / Log Divergence

Full QCD

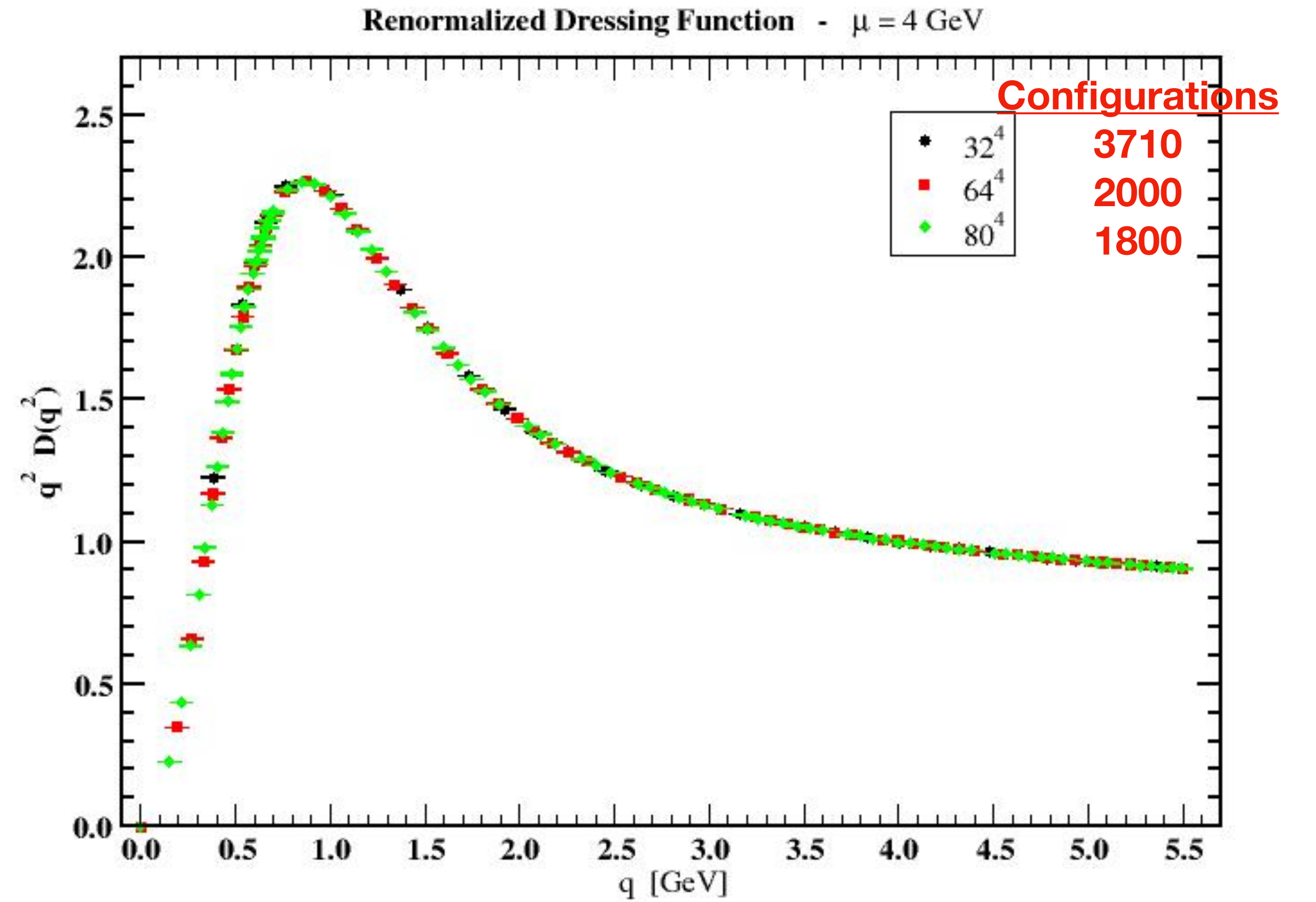
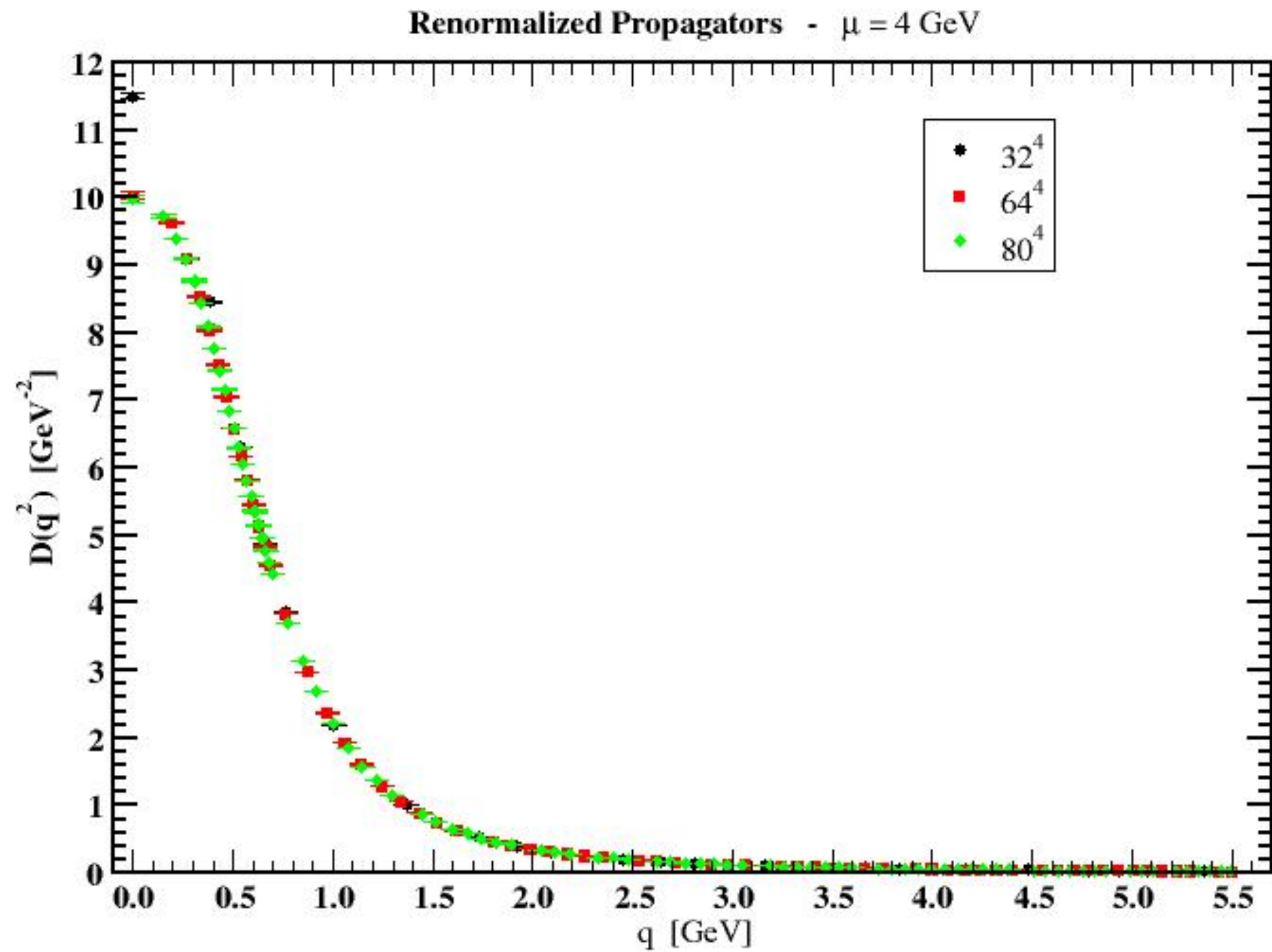
$$\lambda_1(p_1, p_2, p_3) = \frac{F(p_3^2)}{2} \left\{ \begin{aligned} &A(p_1^2) \left[X_0 + (p_1^2 - (p_1 p_2)) X_3 \right] + \\ &A(p_2^2) \left[\bar{X}_0 + (p_2^2 - (p_1 p_2)) \bar{X}_3 \right] + \\ &B(p_1^2) \left[X_1 + X_2 \right] + \\ &B(p_2^2) \left[\bar{X}_1 + \bar{X}_2 \right] \end{aligned} \right\}$$

IR Ghost Dominance / Log Divergence

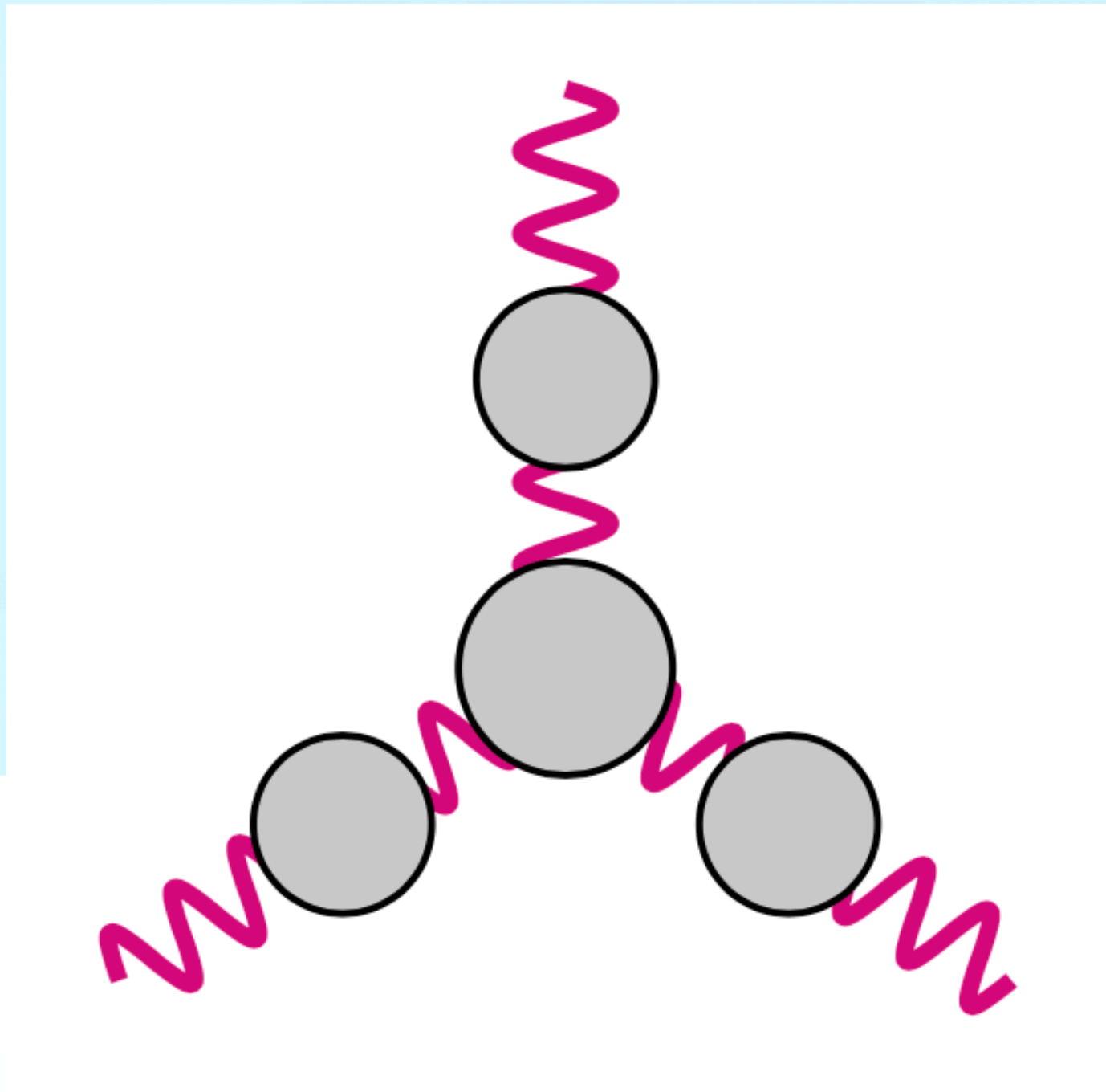
Full QCD $\lambda_1(p_1, p_2, p_3) =$



Gluon Propagator



3-point function



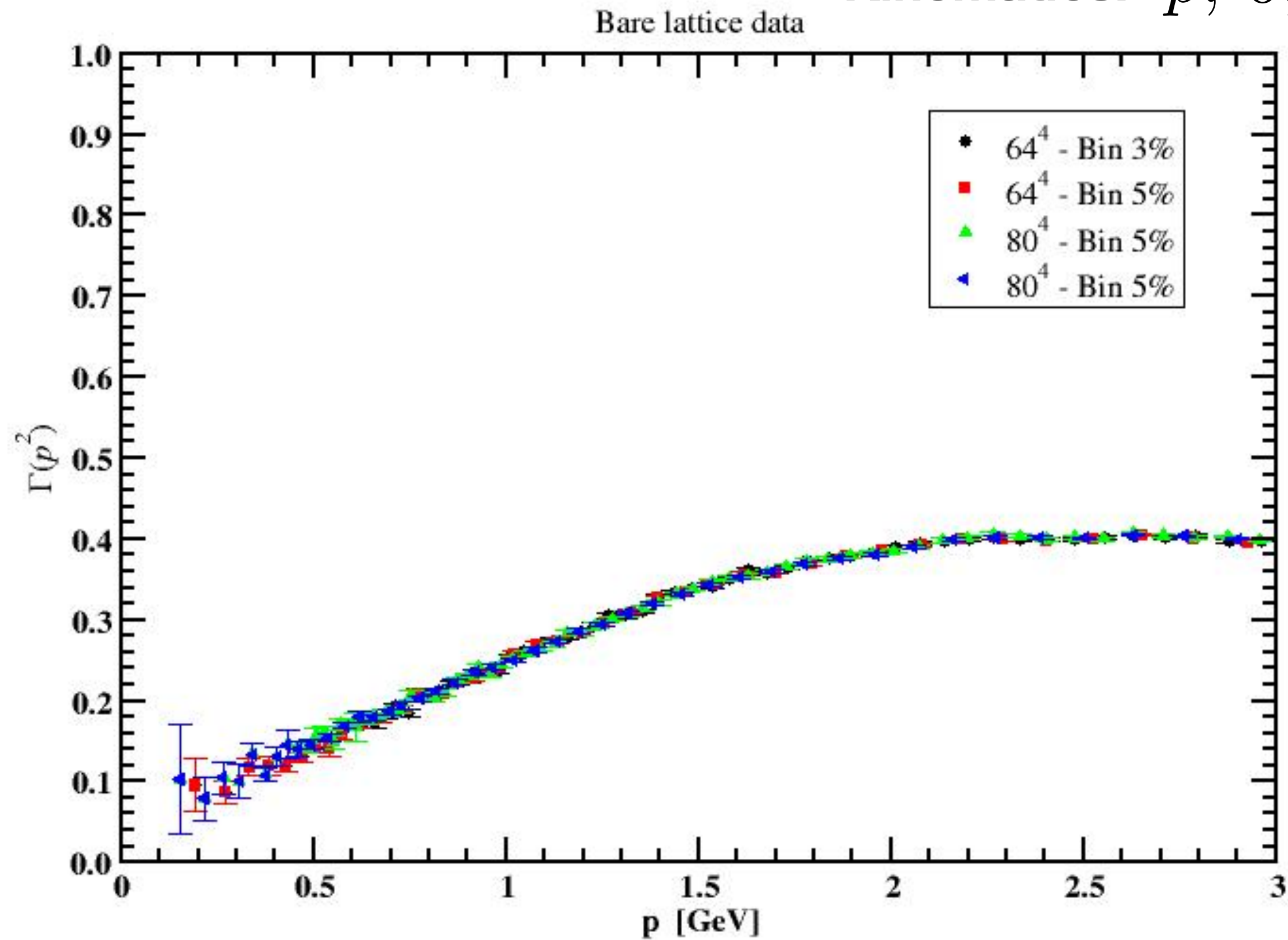
$$\mathcal{G}^{(3)}{}_{\alpha\beta\zeta}{}^{abc}(p_1, p_2, p_3) = D_{\alpha\alpha'}{}^{aa'}(p_1) D_{\beta\beta'}{}^{bb'}(p_2) D_{\zeta\zeta'}{}^{cc'}(p_3) \underbrace{\Gamma_{\alpha'\beta'\zeta'}{}^{a'b'c'}(p_1, p_2, p_3)}_{f_{a'b'c'} \Gamma_{\alpha'\beta'\zeta'}(p_1, p_2, p_3)}$$

Six form factors

J S Ball, T.-W. Chiu, PRD22, 2550 (1980)

Landau gauge: only the transverse form factors

Kinematics: $\vec{p}, \vec{0}, -\vec{p}$



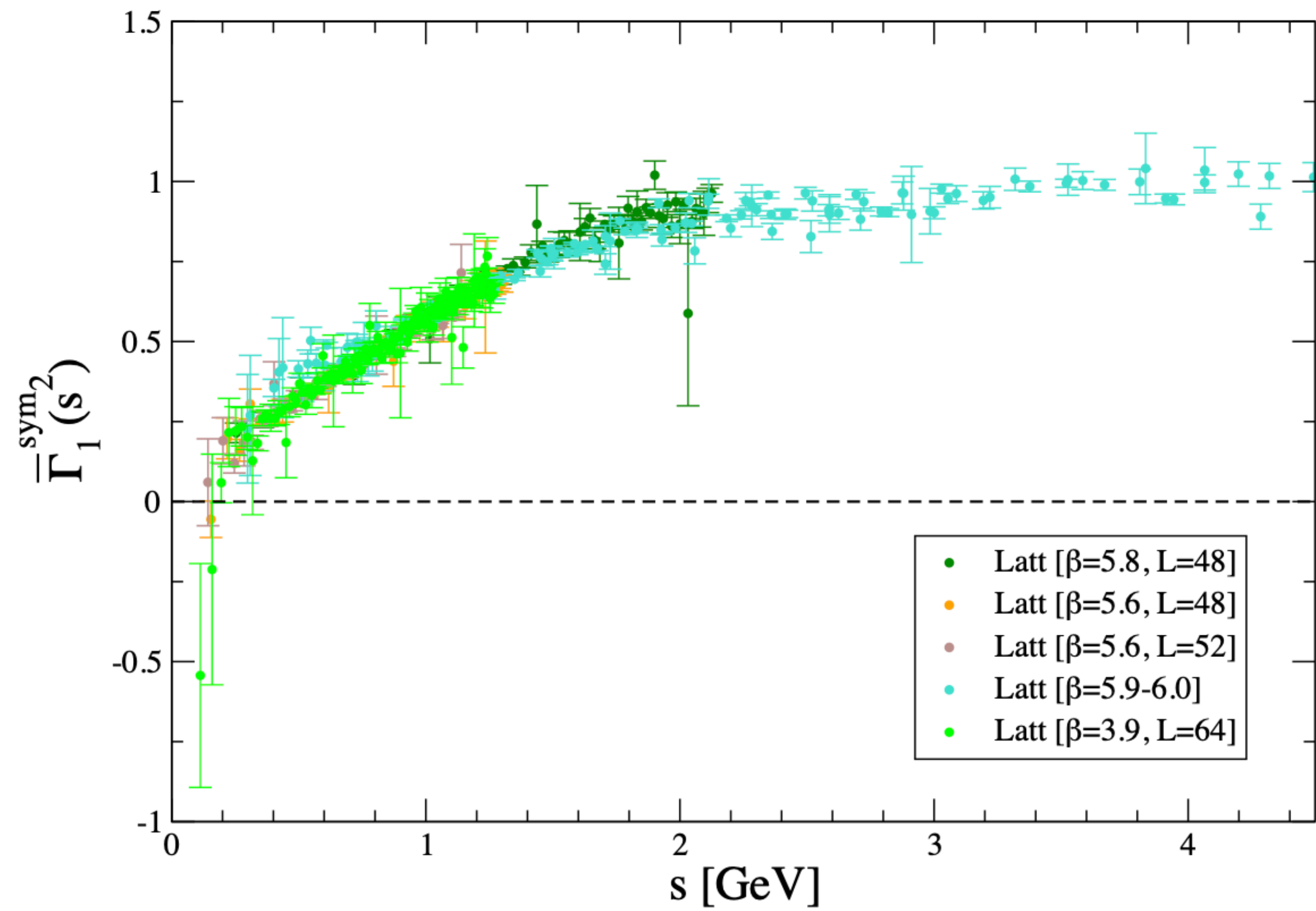
G. T. R. Catumba, O. Oliveira, P. J. Silva, EPJ Web Conf. **258** (2022) 02008

A G Duarte, O. Oliveira, Phys Rev D **94** (2016) 074502

A. Cucchieri, A. Maas, T. Mendes, Phys Rev D **77** (2008) 0945001
and many other works by several groups

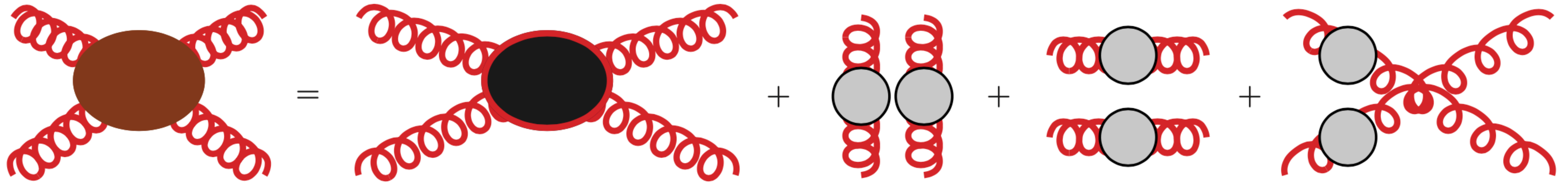
See all talks by:
Feliciano de Soto

Typically more than ~1000 configurations

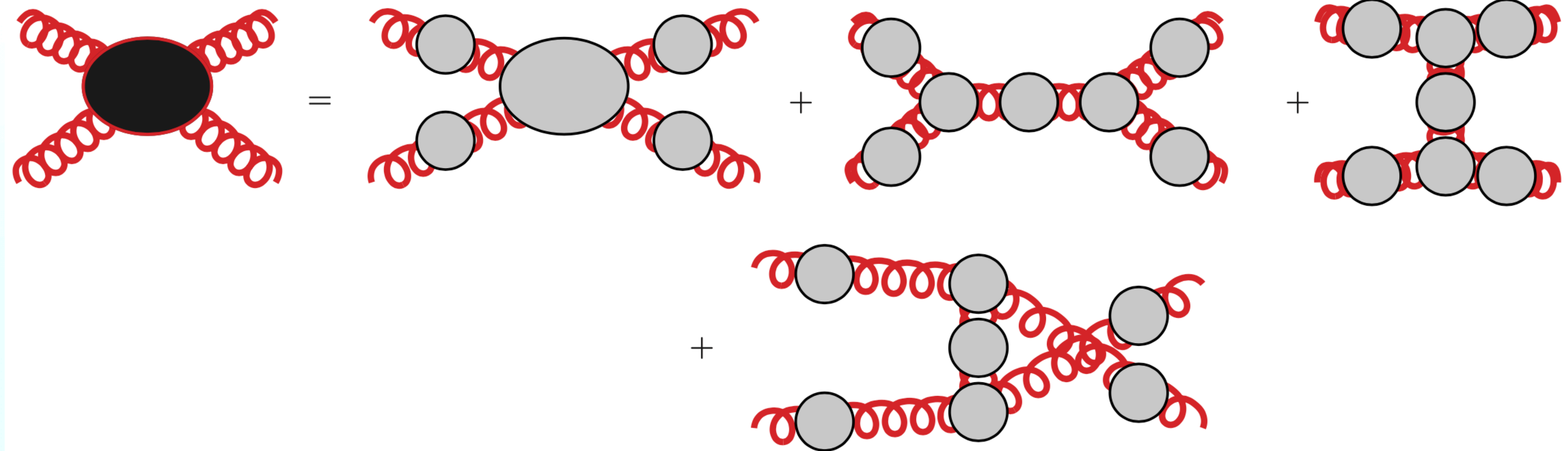


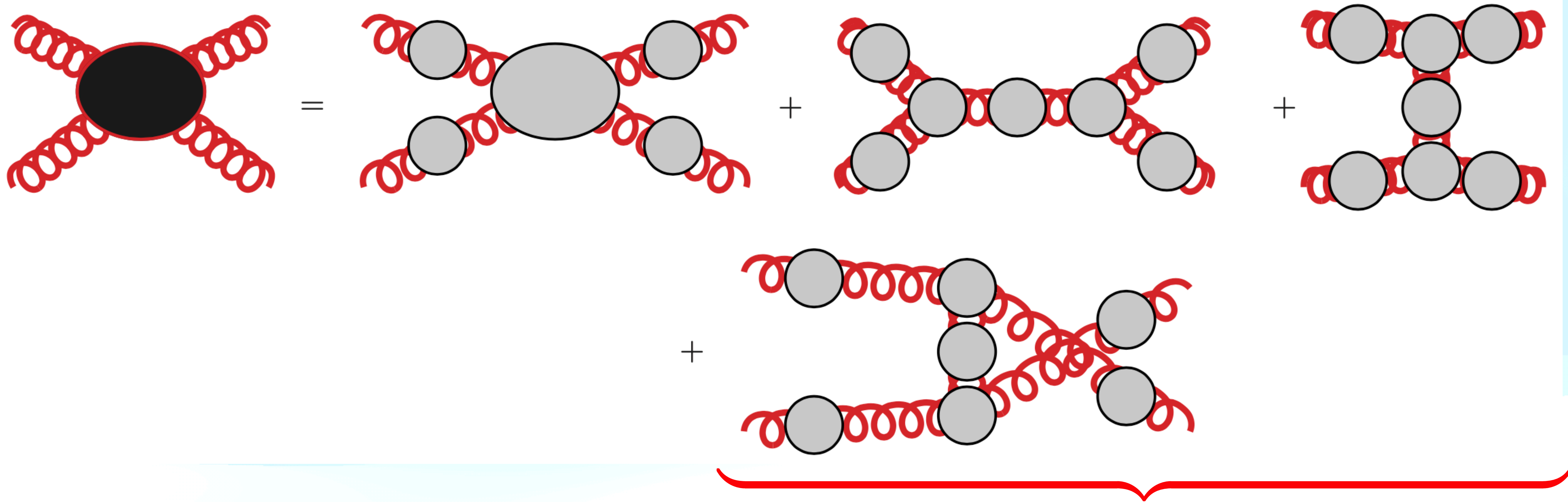
A. C. Aguilar et al, *Phys. Lett. B* **818** (2021) 136352

4-gluon correlation function



$\vec{p}_i \neq \vec{p}_j$





Single momentum scale

$$\vec{p}_i \propto \vec{p}$$

In **Landau gauge** simplifies the tensor analysis

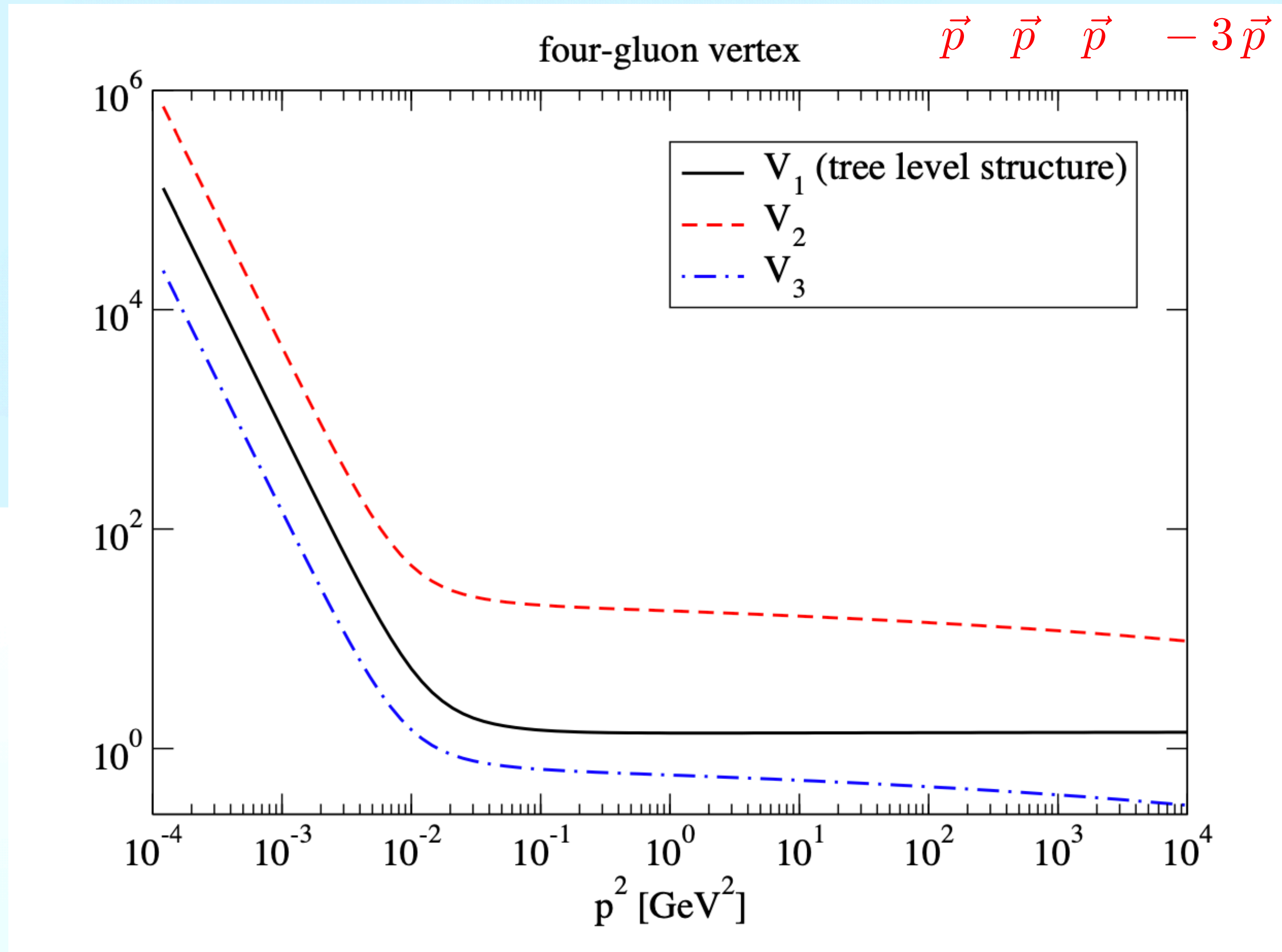
Contributions only from the tensors proportional to $\delta_{\mu\nu}$

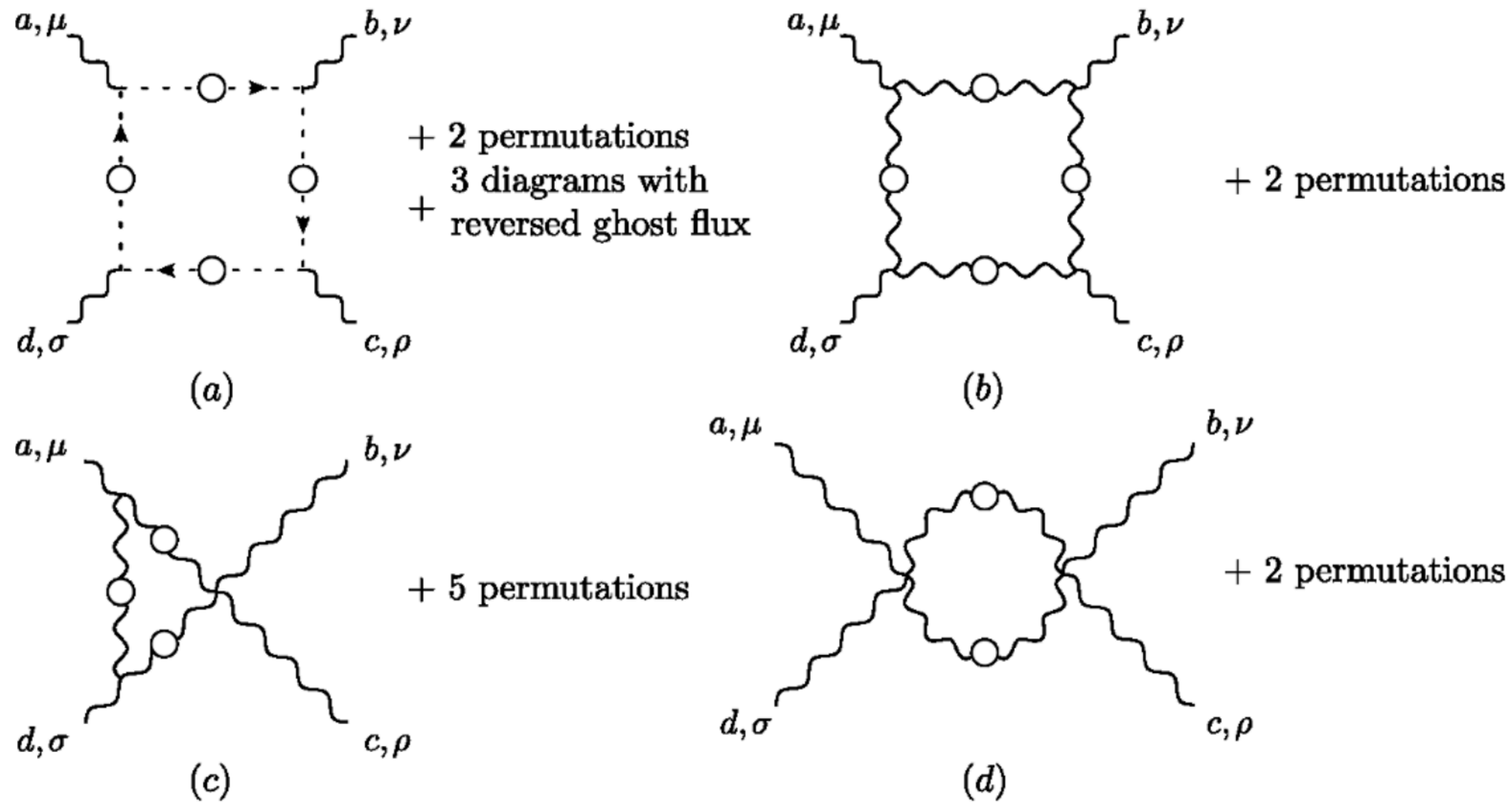
Recent Continuum Calculations

C Kellermann, C S Fischer, *Phys. Rev. D* **78** (2008) 025015

D Binosi, D Ibañez, J Papavassiliou, *JHEP* **1409** (2014) 059

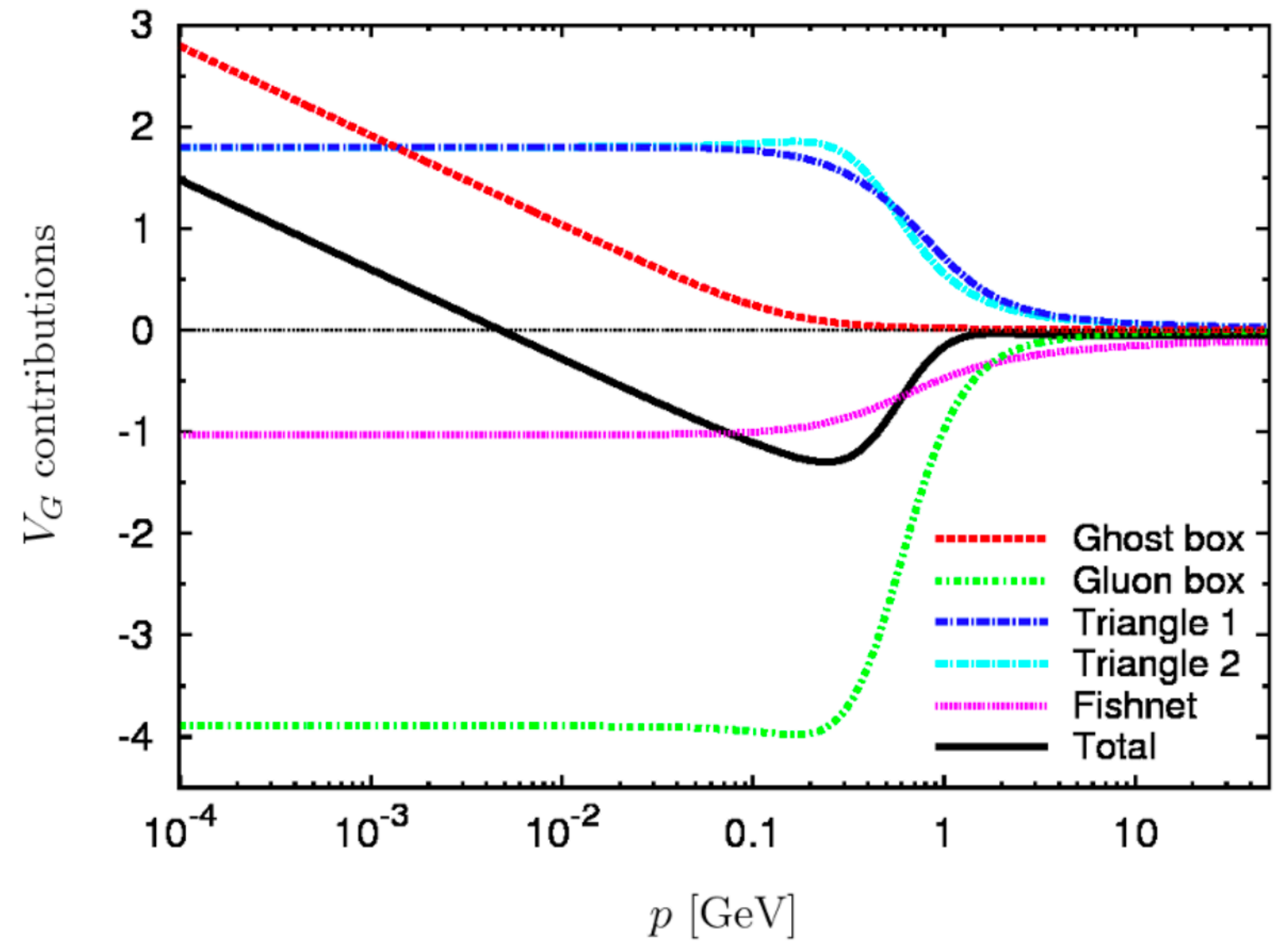
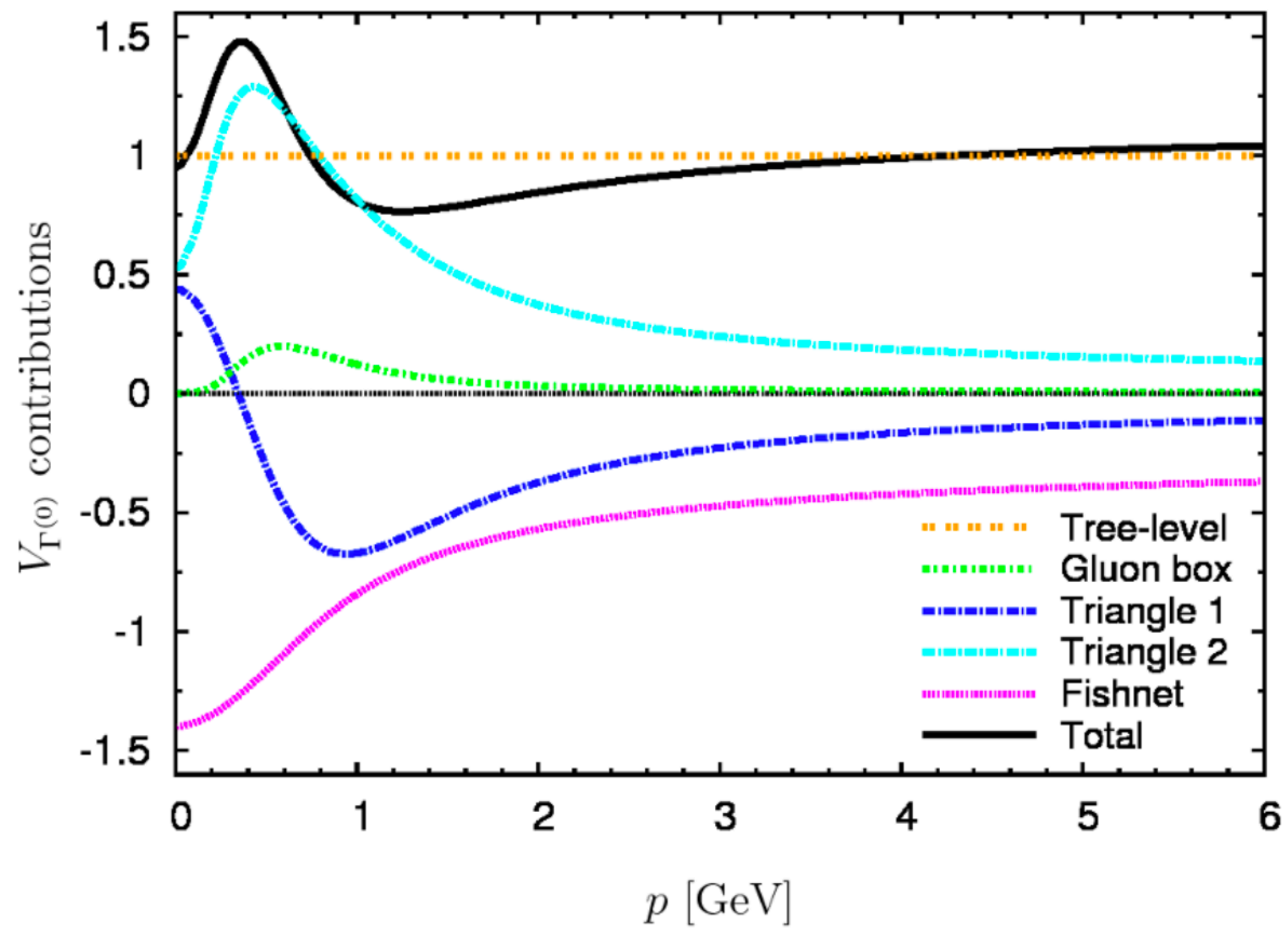
A K Carol, M Q Huber, L von Smekal, *Eur Phys J C* **75** (2015) 102



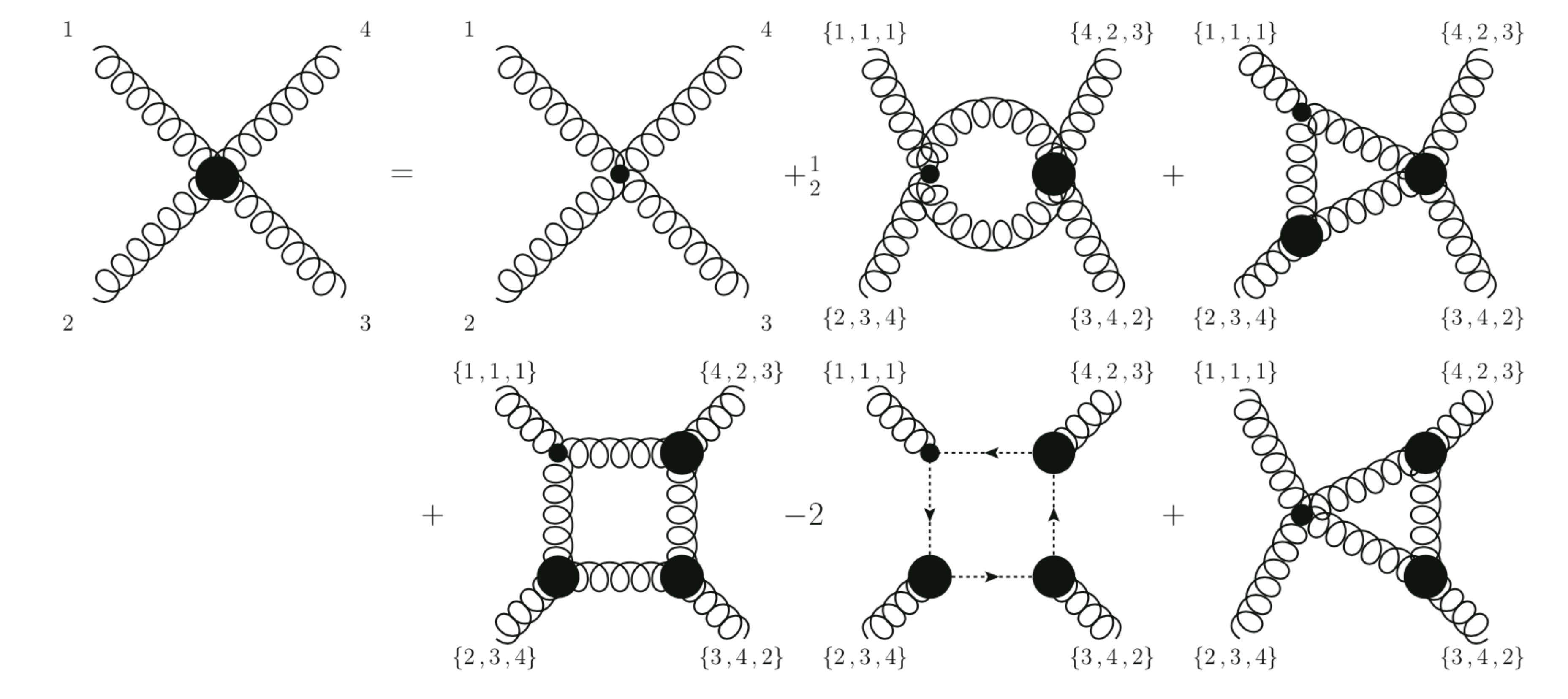


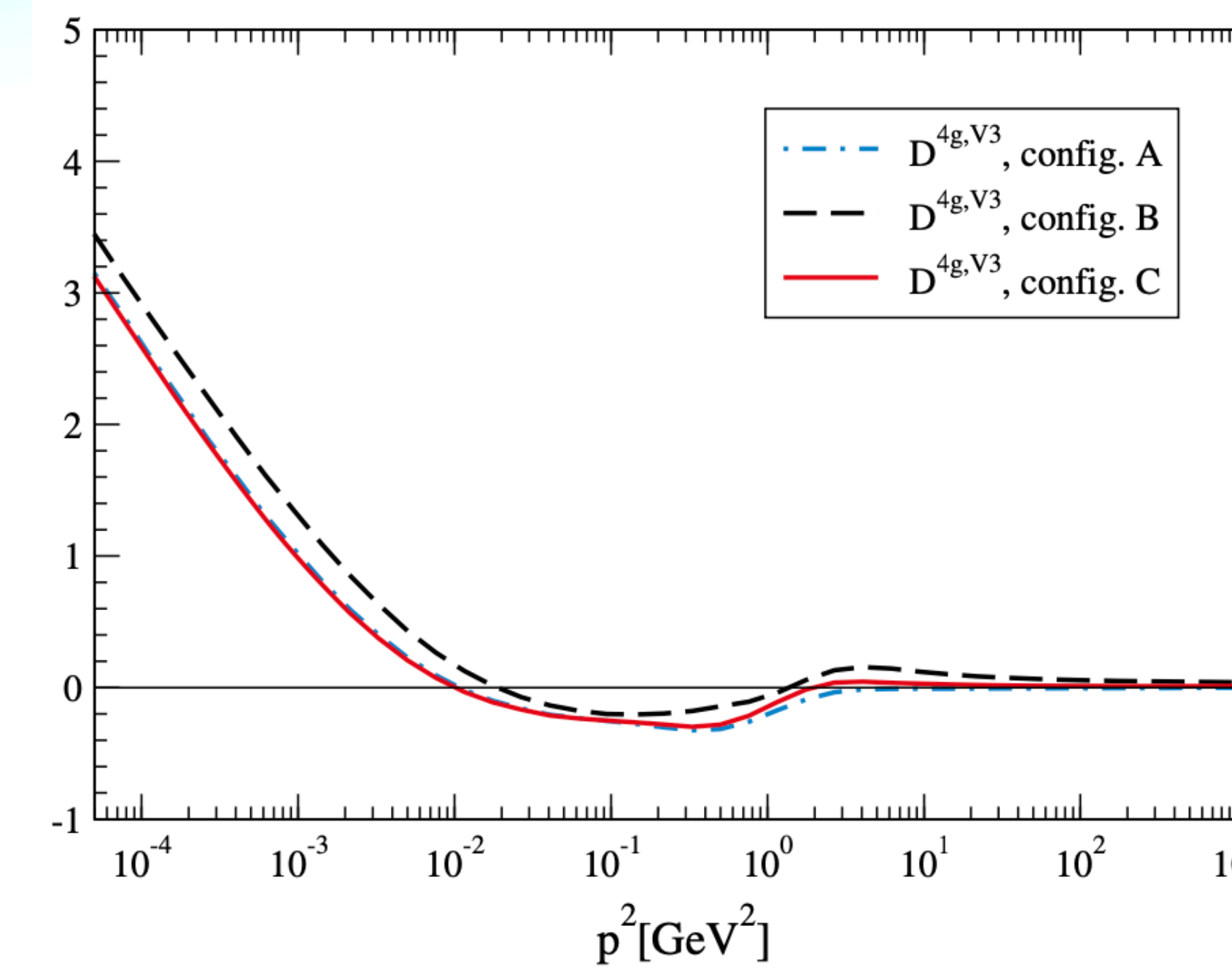
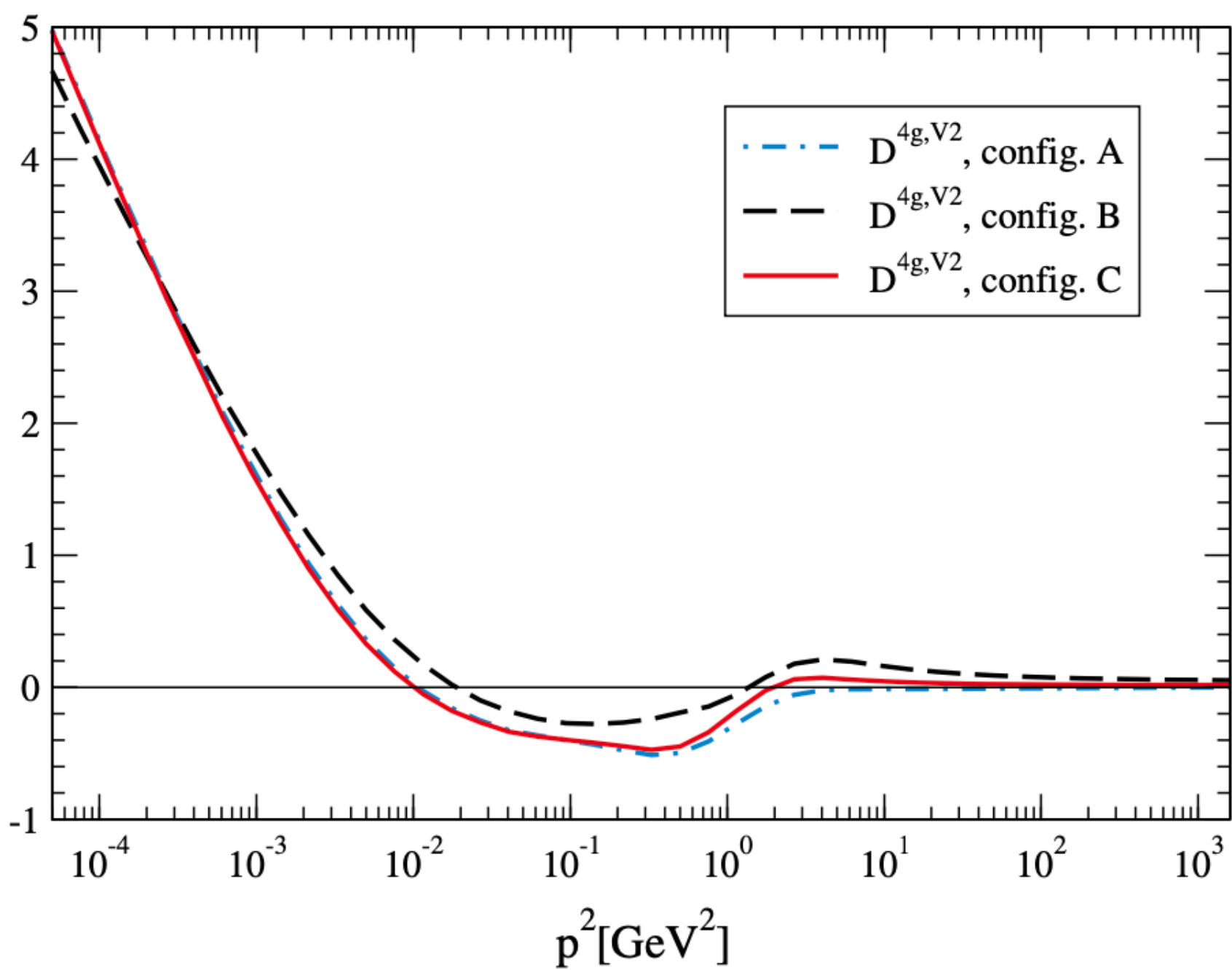
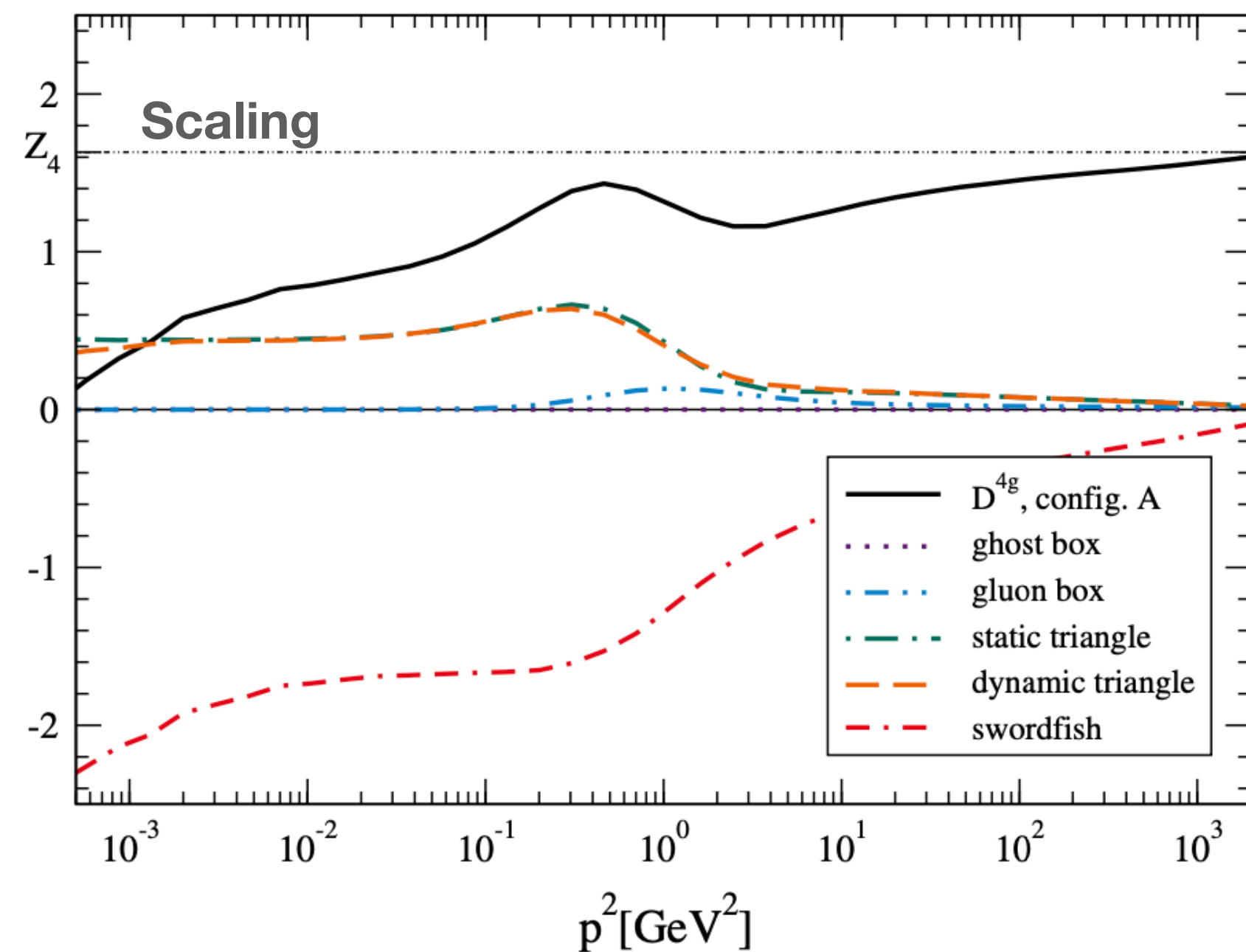
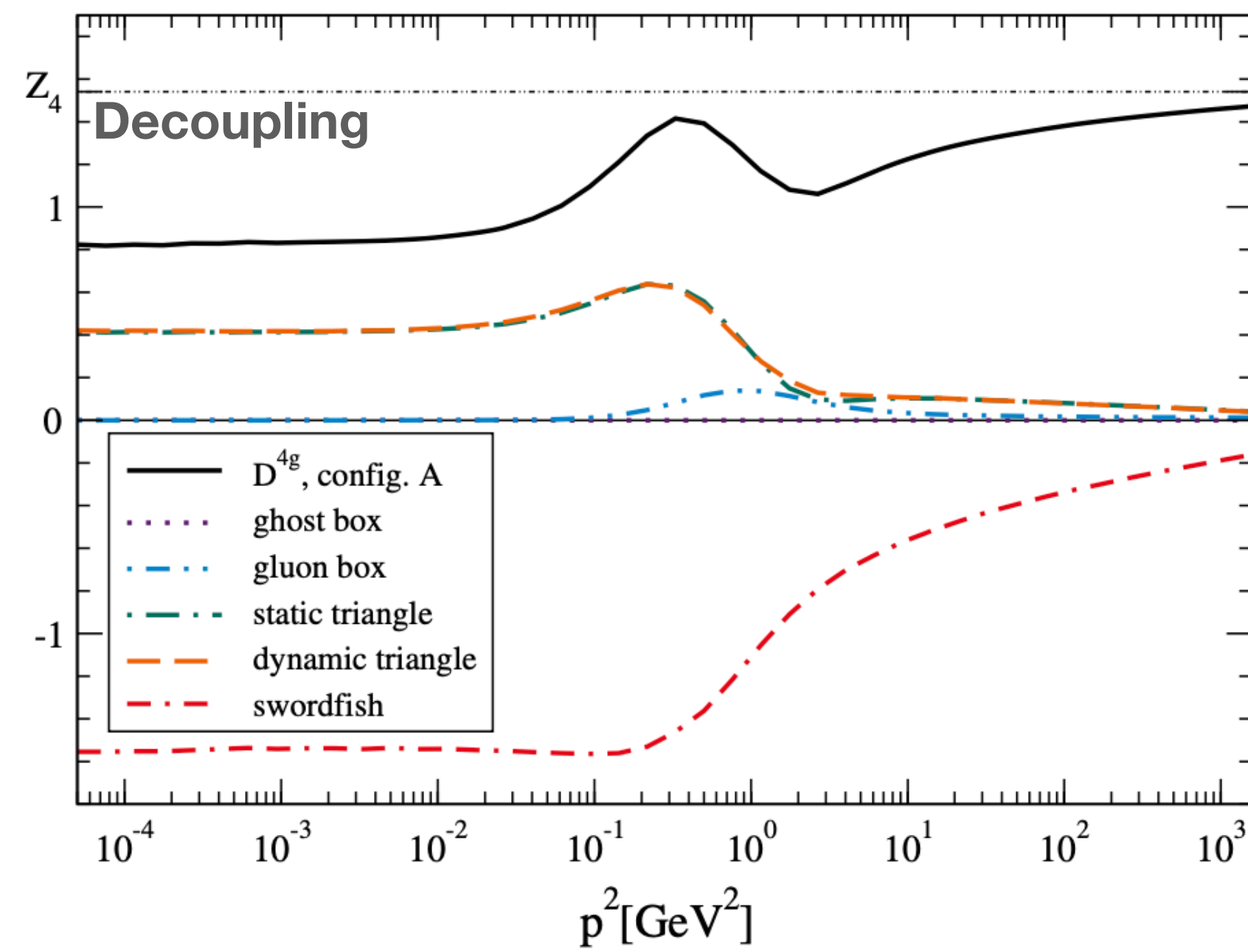
$$\Gamma_{\mu\nu\rho\sigma}^{abcd}(p, p, p, -3p) \Big|_{gg} = V_{\Gamma^{(0)}}(p^2) \Gamma_{\mu\nu\rho\sigma}^{abcd(0)} + V_G(p^2) G_{\mu\nu\rho\sigma}^{abcd},$$

$$G_{\mu\nu\rho\sigma}^{abcd} = (\delta^{ab} \delta^{cd} + \delta^{ac} \delta^{bd} + \delta^{ad} \delta^{bc}) \underbrace{(g_{\mu\nu} g_{\rho\sigma} + g_{\mu\rho} g_{\nu\sigma} + g_{\mu\sigma} g_{\nu\rho})}_{R_{\mu\nu\rho\sigma}}$$



A K Cyrol, M Q Huber, L von Smekal, *Eur Phys J C* **75** (2015) 102





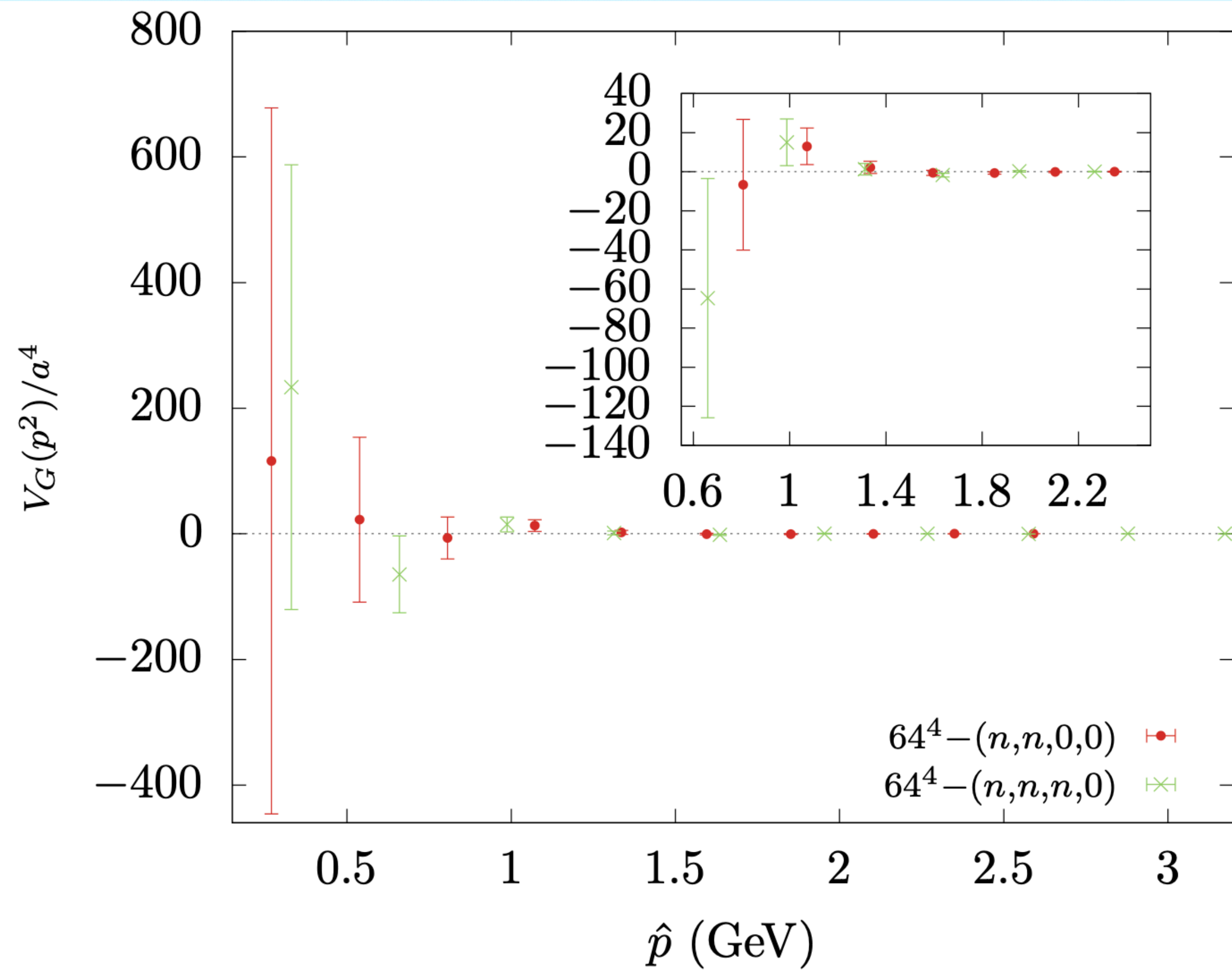
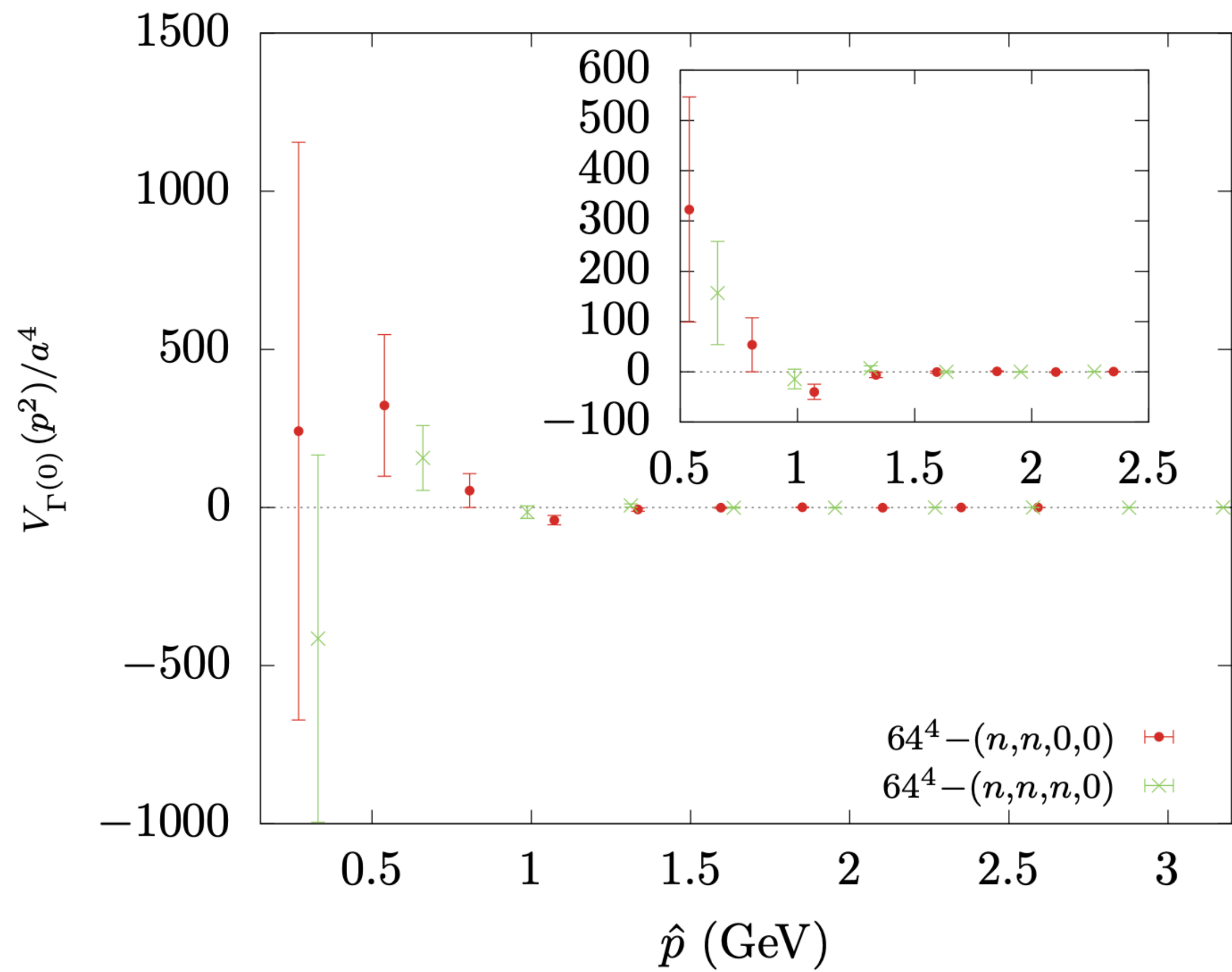
First tentative of a lattice calculation can be found in
G T R Catumba Master Thesis [ArXiv: 2101.06074](https://arxiv.org/abs/2101.06074)

D. Binosi, D. Ibañez, J. Papavassiliou JHEP 9, 059 (2014) arXiv:1407.3677

$p, p, p, -3p$ Two tensor structures

$$\tilde{\Gamma}_{\mu\nu\eta\zeta}^{(0)abcd} = f_{abr}f_{cdr}(\delta_{\mu\eta}\delta_{\nu\zeta} - \delta_{\mu\zeta}\delta_{\nu\eta}) + f_{acr}f_{bdr}(\delta_{\mu\nu}\delta_{\eta\zeta} - \delta_{\mu\zeta}\delta_{\nu\eta}) + f_{adr}f_{bcr}(\delta_{\mu\nu}\delta_{\eta\zeta} - \delta_{\mu\eta}\delta_{\nu\zeta})$$

$$G_{\mu\nu\eta\zeta}^{abcd} = (\delta^{ab}\delta^{cd} + \delta^{ac}\delta^{bd} + \delta^{ad}\delta^{bc})(\delta_{\mu\nu}\delta_{\eta\zeta} + \delta_{\mu\eta}\delta_{\nu\zeta} + \delta_{\mu\zeta}\delta_{\nu\eta})$$



For **tensor analysis** see:

J A Gracey, Phys Rev D90, 025011 (2014) arXiv: 1406.1618

G Eichmann, C S Fischer, W Heupel, Phys Rev D92, 056006 (2015) arXiv: 1505.06336

$$\tilde{\Gamma}^{(0)}{}_{\mu\nu\eta\zeta}{}^{abcd} = f_{abr}f_{cdr}(\delta_{\mu\eta}\delta_{\nu\zeta} - \delta_{\mu\zeta}\delta_{\nu\eta}) + f_{acr}f_{bdr}(\delta_{\mu\nu}\delta_{\eta\zeta} - \delta_{\mu\zeta}\delta_{\nu\eta}) + f_{adr}f_{bcr}(\delta_{\mu\nu}\delta_{\eta\zeta} - \delta_{\mu\eta}\delta_{\nu\zeta})$$

$$\tilde{\Gamma}^{(1)}{}_{\mu\nu\eta\zeta}{}^{abcd} = d_{abr}d_{cdr}(g_{\mu\eta}g_{\nu\zeta} + g_{\mu\zeta}g_{\nu\eta}) + d_{acr}d_{bdr}(g_{\mu\zeta}g_{\nu\eta} + g_{\mu\nu}g_{\eta\zeta}) + d_{adr}d_{bcr}(g_{\mu\nu}g_{\eta\zeta} + g_{\mu\eta}g_{\nu\zeta})$$

$$\tilde{\Gamma} = F(p^2)\tilde{\Gamma}^{(0)} + G(p^2)\tilde{\Gamma}^{(1)} + H(p^2)\tilde{\Gamma}^{(2)}$$

Not an orthogonal basis

$$\mathcal{G}^{(4)} = \tilde{\Gamma} \left(P^\perp(p) D(p^2) \right)^3 \left(P^\perp(3p) D(9p^2) \right)$$

p, p, p, -3p

Kinematics

0, p, p, -2p

Measure the two form factors

$$\beta = 6.0 \quad a = 0.102 \text{ fm} \quad a^{-1} = 1.943 \text{ GeV}$$

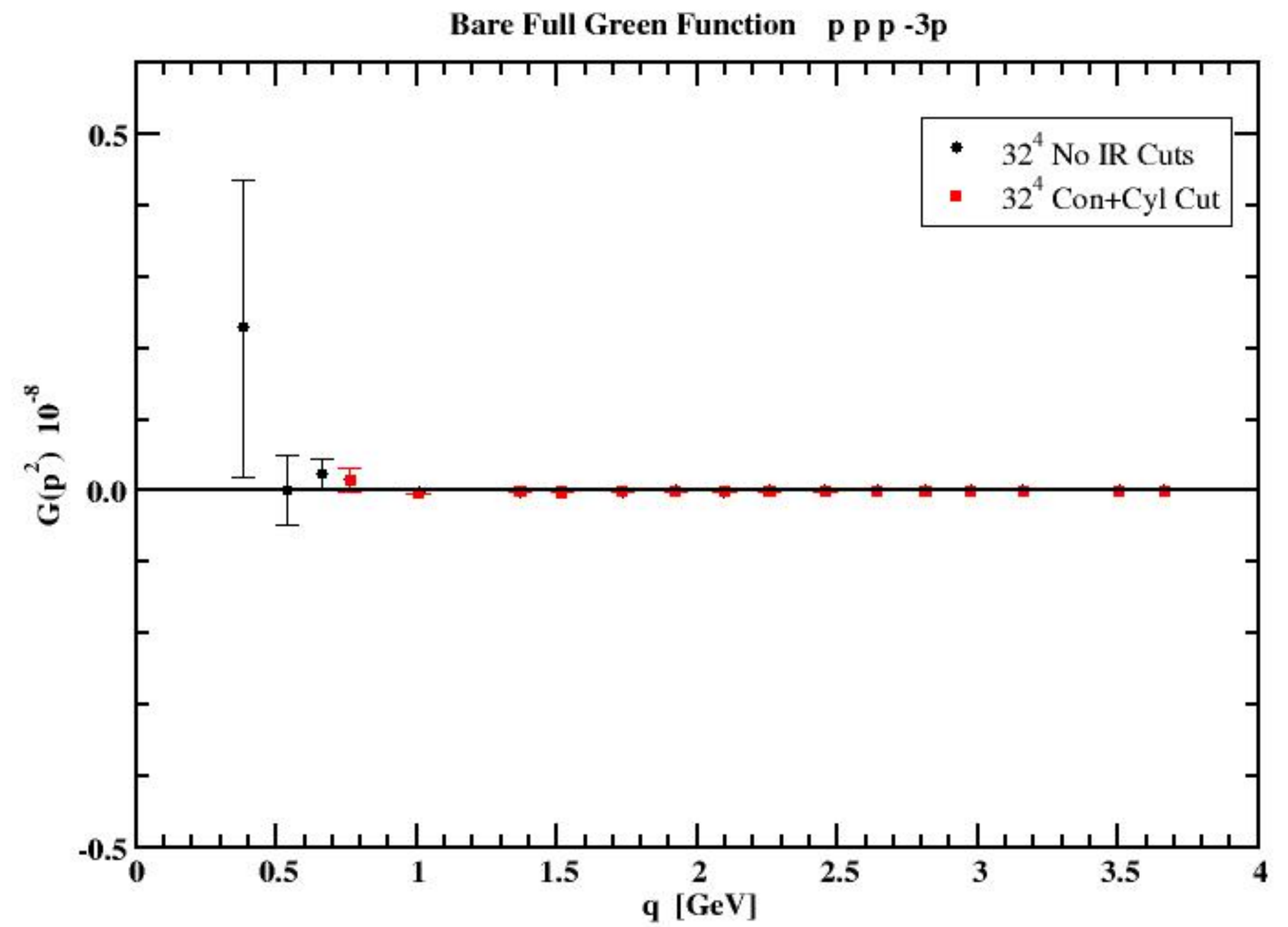
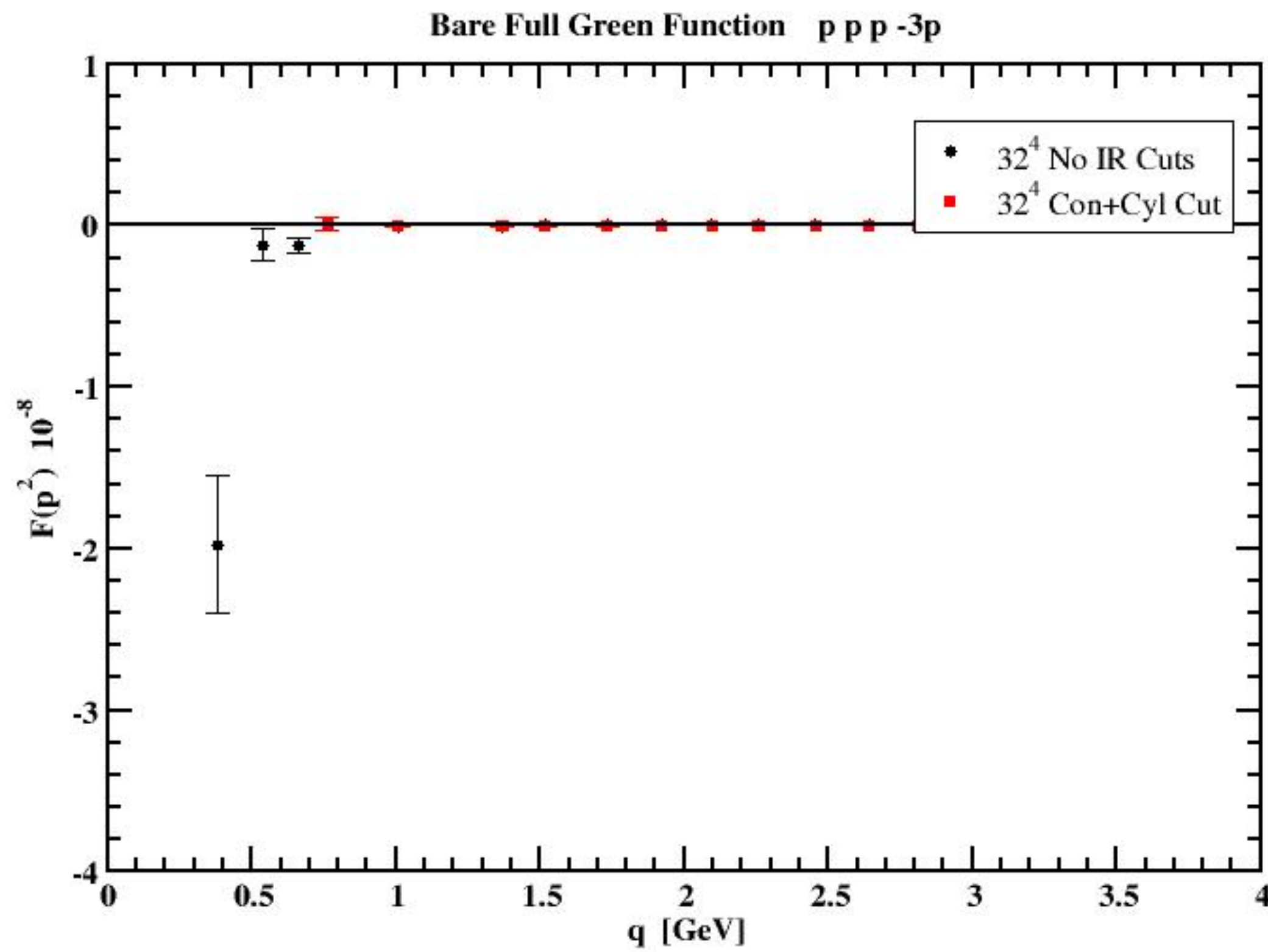
Lattice	Configs	
32^4	3710	381 MeV
64^4	2000	191 MeV
80^4	1801	153 MeV

Averaged over equivalent momenta, including the negative momenta !

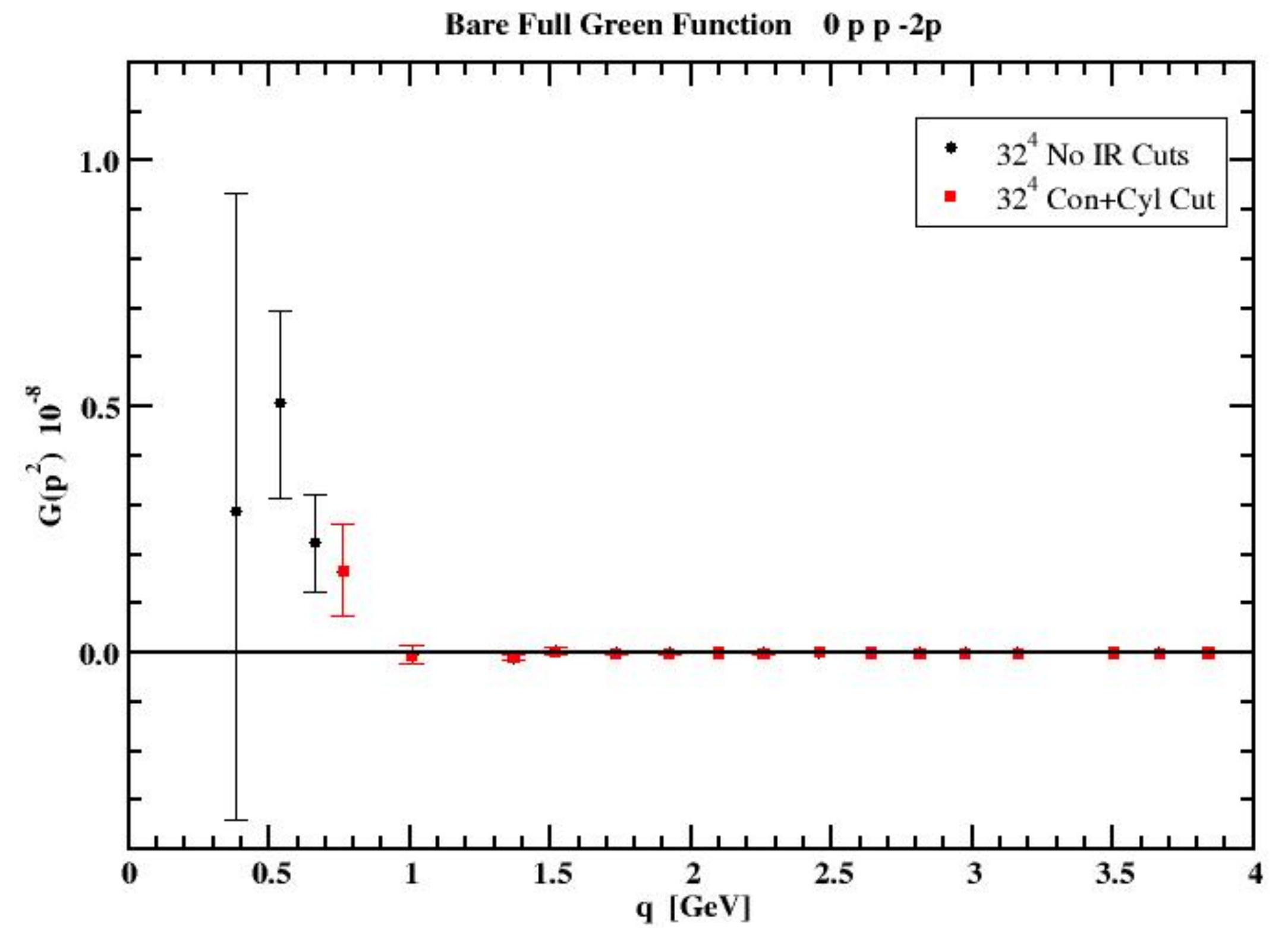
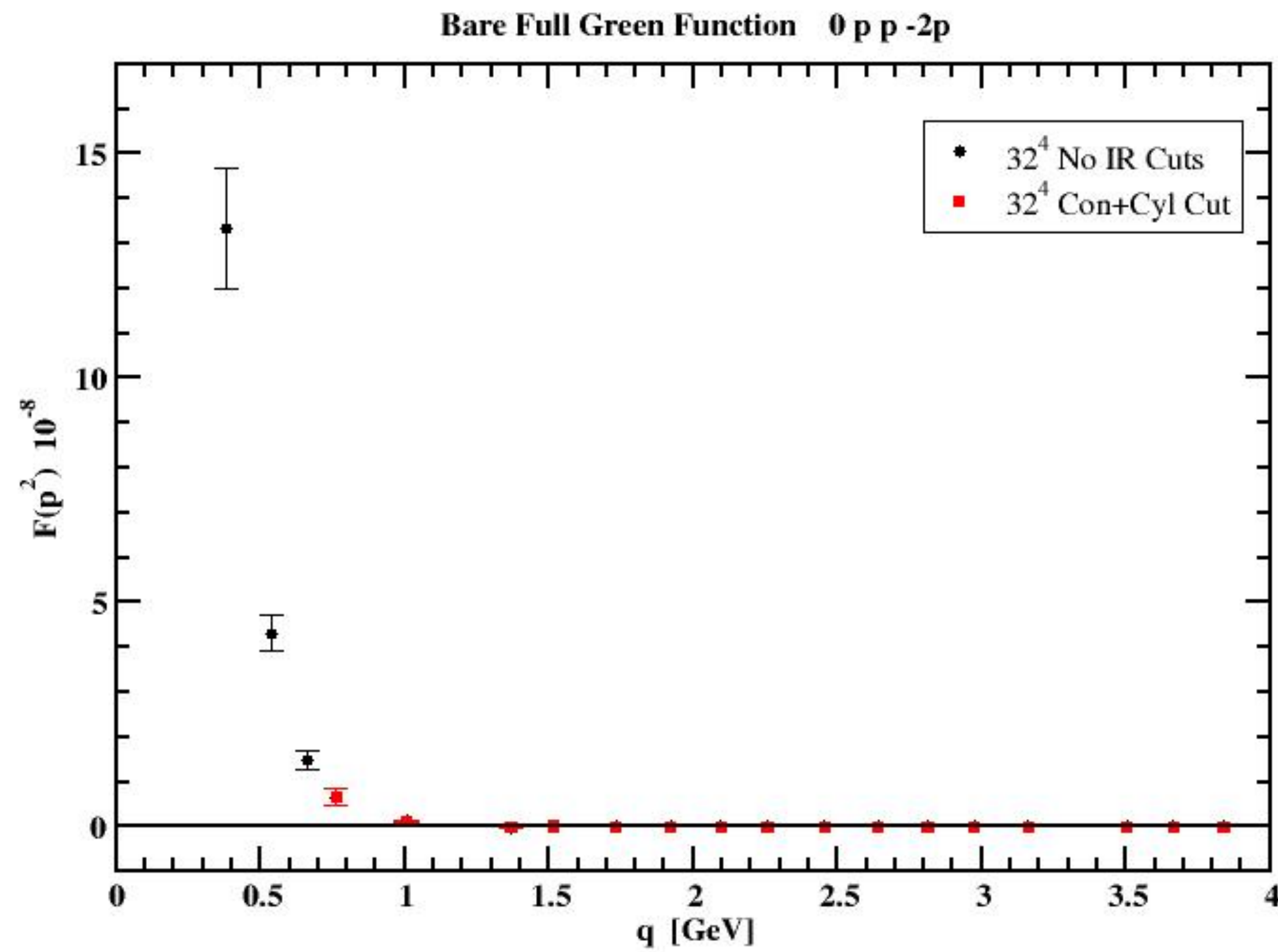
(1 1 1 1)

(-1 1 1 1) (1 -1 1 1) (1 1 -1 1) (1 1 1 -1)
(-1 -1 1 1) (-1 1 -1 1) (-1 1 1 -1) (1 -1 -1 1) (1 -1 1 -1) (1 1 -1 -1)
(1 -1 -1 -1) (-1 1 -1 -1) (-1 -1 1 -1) (-1 -1 -1 1)
(-1 -1 -1 -1)

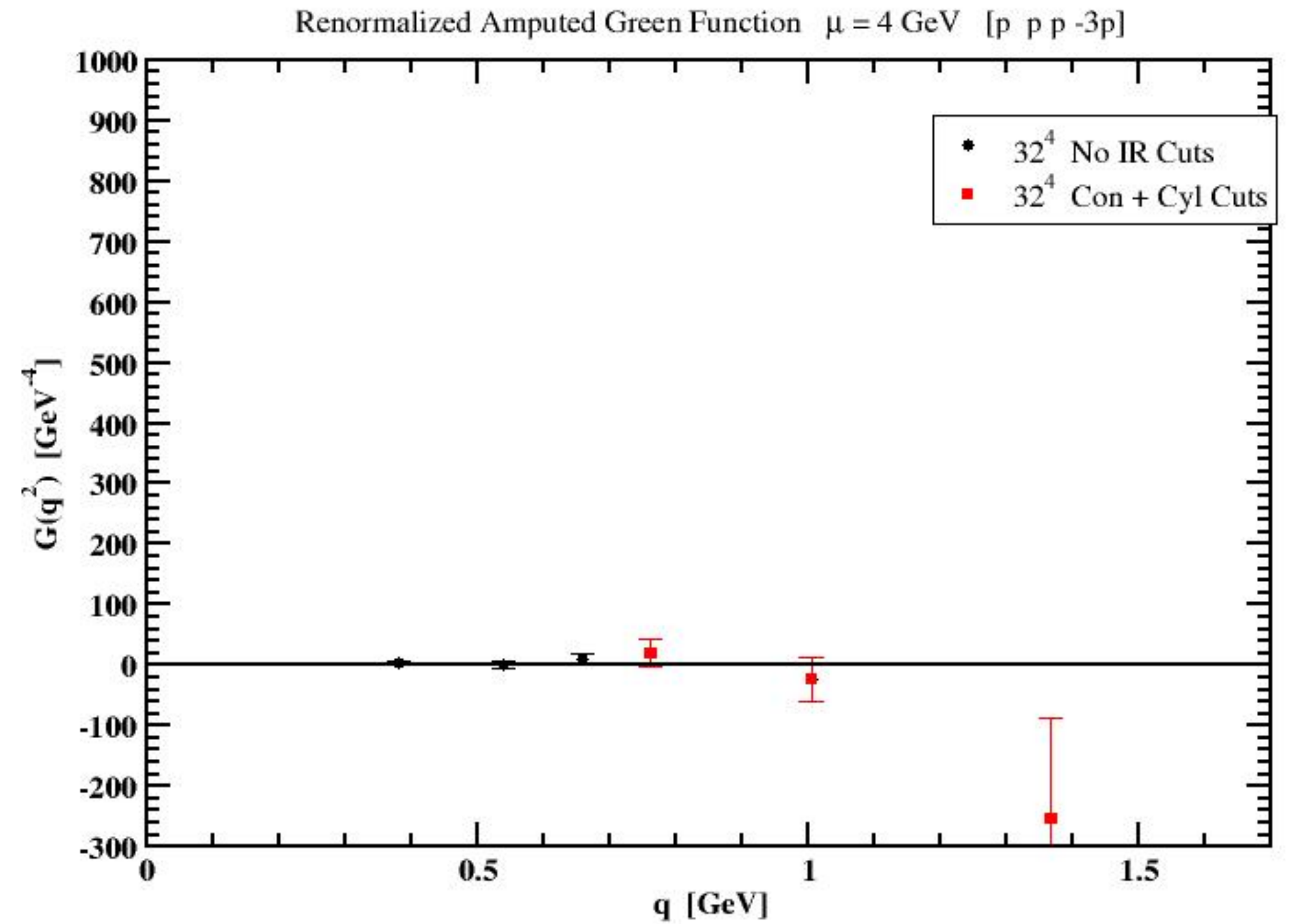
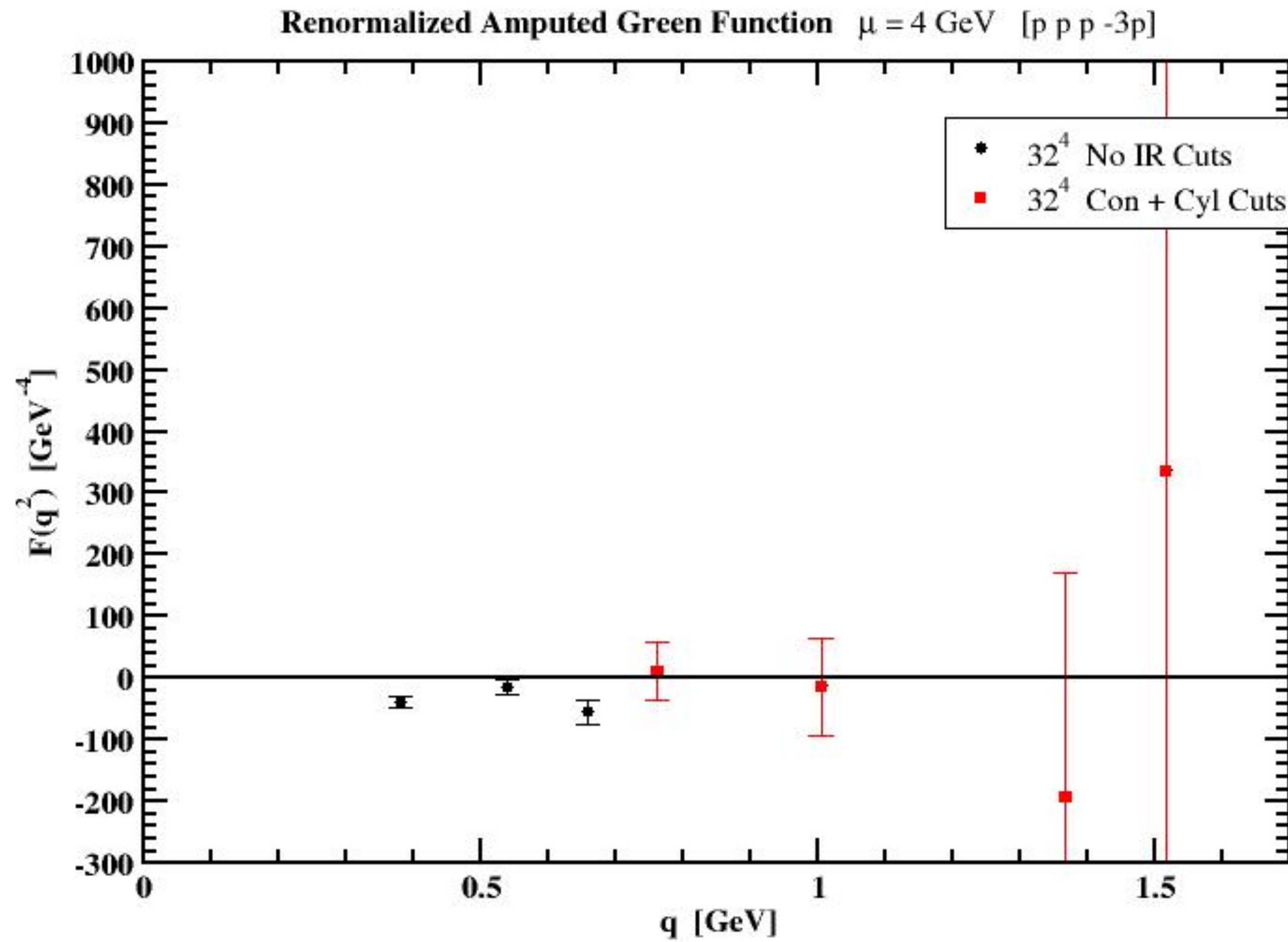
Bare Green Function for kinematics: $p p p -3p$



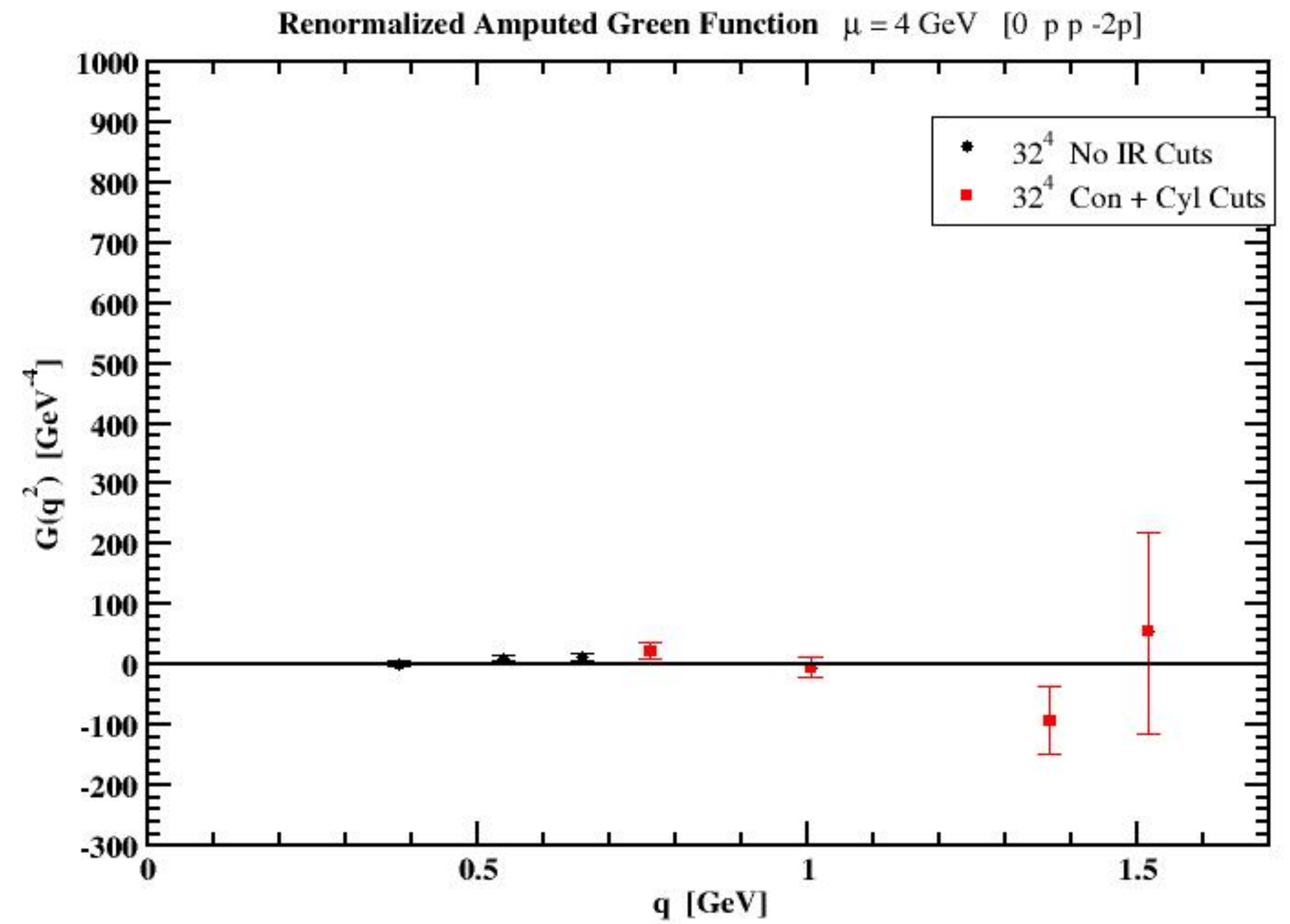
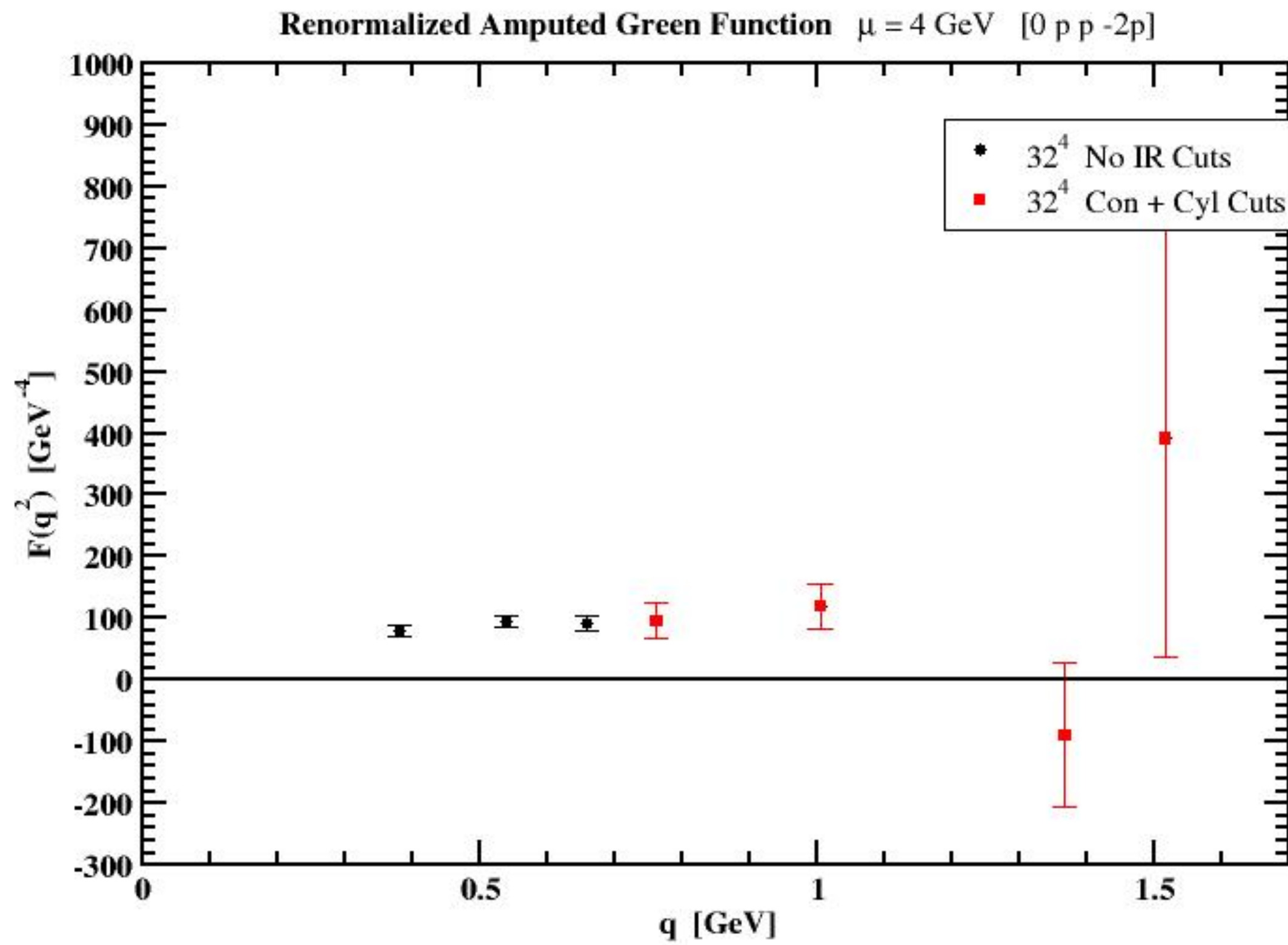
Bare Green Function for kinematics: 0 p p -2p



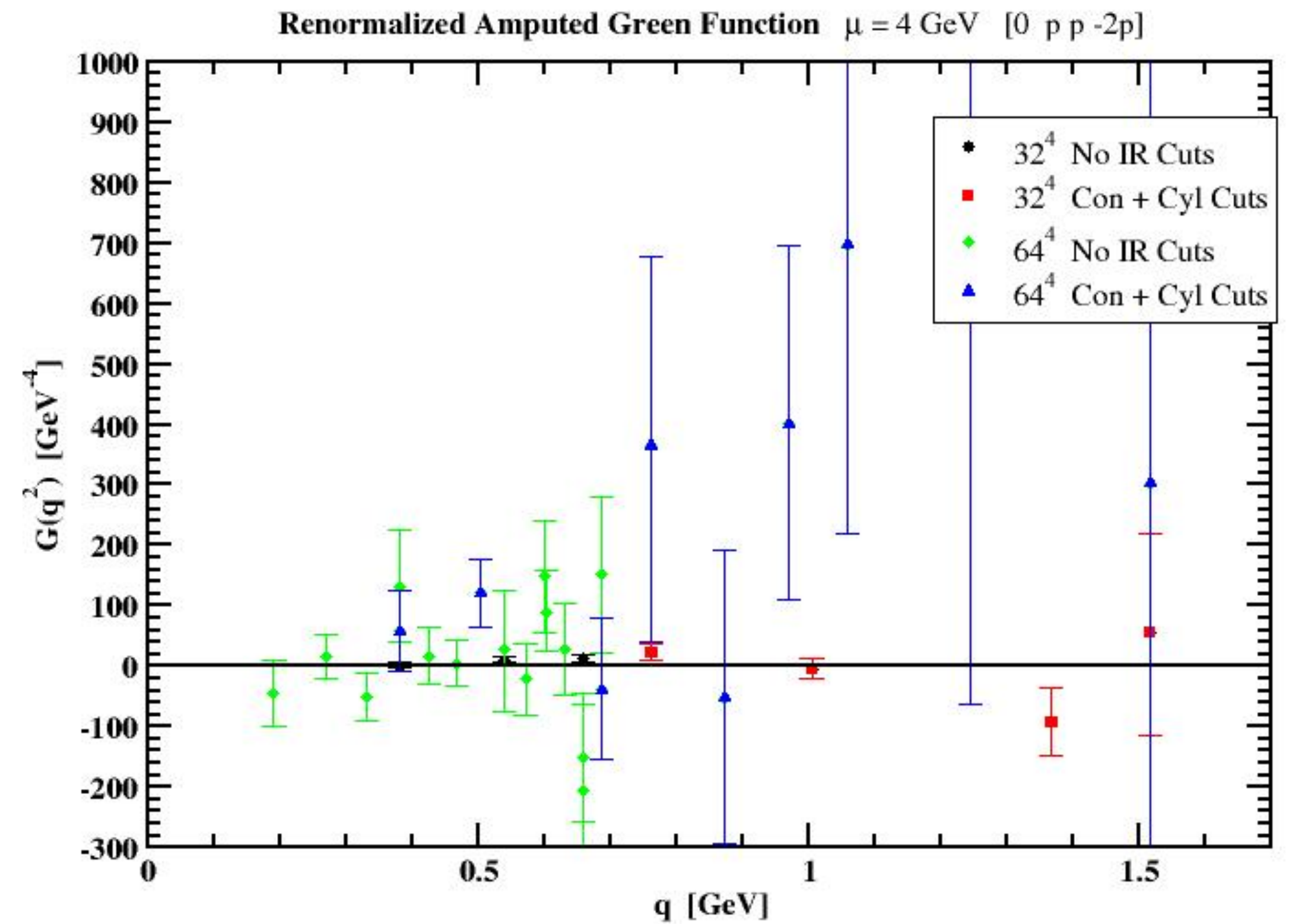
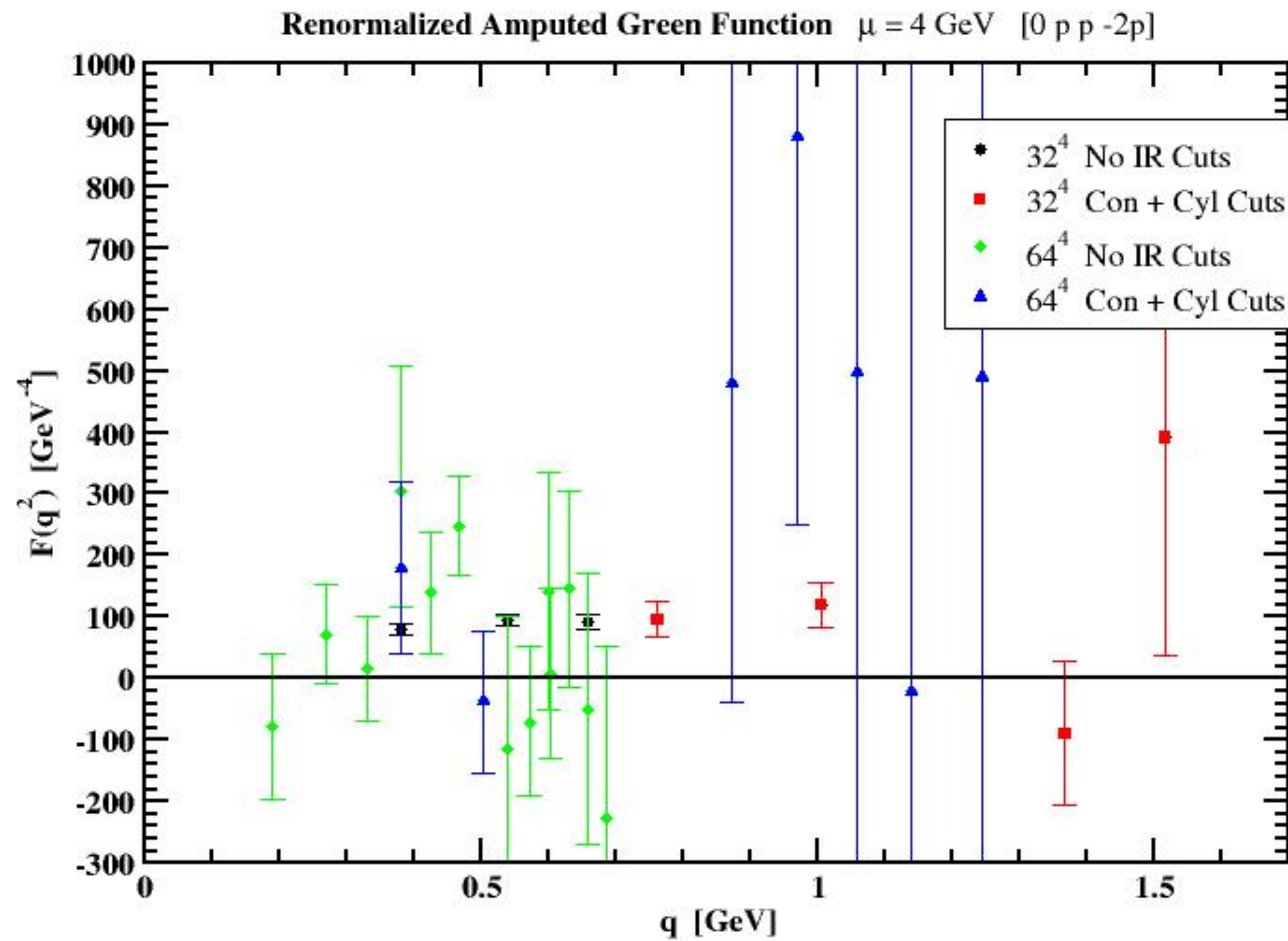
Vertex Form Factors Renormalized @ $\mu = 4 \text{ GeV}$



Vertex Form Factors Renormalized @ $\mu = 4 \text{ GeV}$



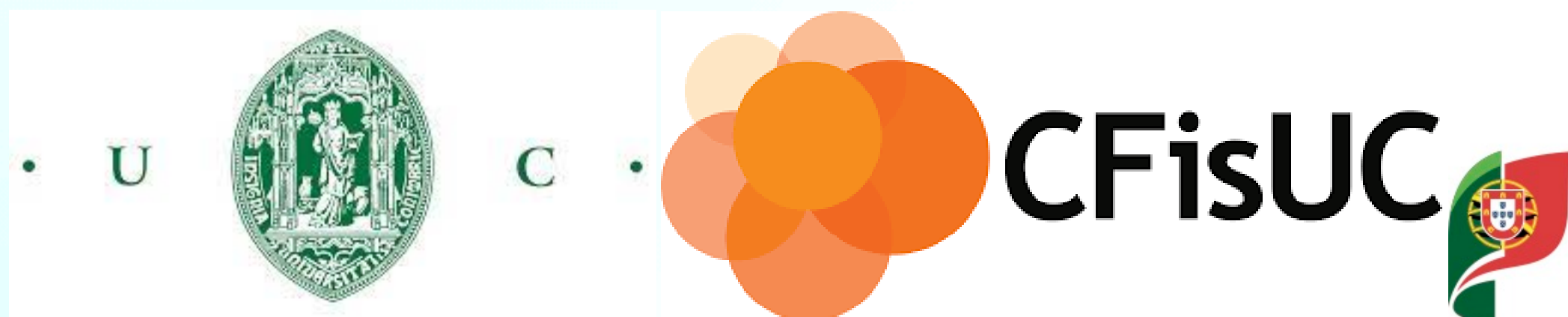
Vertex Form Factors Renormalized @ $\mu = 4$ GeV



Summary and Conclusions

- gluon correlation functions up to 4-external legs are possible with standard lattice methods
- need to increase the statistics and improve the tensor analysis
- It would be great to explore further kinematics (need to be studied)
- How low can we go in momentum?

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