

Computation of Relativistic Corrections to the Static Potential

from Generalized Wilson Loops at Finite Flow Time

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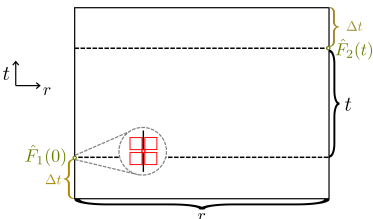
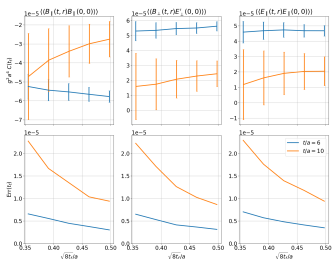
Trento, 24.03.2023

Introduction

- Non-relativistic quarkonia can be described via [static potentials](#) + [rel. corrections](#) [[2212.11107](#)].
- Well understood for $t_f = 0$ to high order in $\bar{m}s$, but only *few results* from lattice QCD.
- Correlators **strongly affected** by *UV fluctuations* and need *renormalization* [[hep-lat/9703019](#)].
- Gradient flow expected to effectively solve these problems (cf. talk by Julian Mayer-Steudte).
- The types of Wilson loops used here are great to study *properties of colour-fields* in gradient flow, e.g. t_f -dependence and matching (cf. talks by Viljami Leino, Nora Brambilla).

Outline

- Overview of the potentials.
- How to compute Correlators.
- Preliminary results for correlators.
- Where to improve and what can be learned.



Potentials

$$V(r) = V^{(0)}(r) + \frac{1}{m} V^{(1)}(r) + \frac{1}{m^2} (V_{SD}(r) + V_{SI}(r)) + \mathcal{O}(1/m^3)$$

Pineda, Vairo, 2001

- $V^{(0)}$: **Static** $\bar{q}q$ -potential. Rest: **Relativistic** correction terms.
- $V^{(0)}$ and corrections well understood in *perturbation theory*, but results from lattice QCD **incomplete**. [Bali, 1997](#), [Koma, Koma, 2007](#)
- **Precise** long-range expressions relevant for phenomenology, e.g. determination of quarkonia masses and related areas.
- **Properties** of colour-field insertions in gradient flow (flow dependence, matching, ...) can be studied here **independently** \Rightarrow Relevant for e.g. computation of **HQ diffusion coefficients**.

Potentials

Potentials are computed using correlators $\langle\langle \dots \rangle\rangle = \langle \dots \rangle_W / \langle 1 \rangle_W \Rightarrow V_0$ (in particular $V_{\text{self}}(\mu)$) **does not appear** in resulting spectrum.

$$V_{\mathbf{p}^2} = \frac{1}{2} \left\{ \mathbf{p}^2, (\mathcal{I}(E_{\parallel}(t, 0)E_{\parallel}(0, 0)) + \mathcal{I}(E_{\parallel}(t, r)E_{\parallel}(0, 0))) \right\},$$

$$V_{\text{LS}} = \epsilon_{ij\parallel} \frac{c_F(\mu)}{r} (\mathcal{I}(B_i(t, 0)E_j(0, 0)) + \mathcal{I}(B_i(t, r)E_j(0, 0))) \hat{\mathbf{L}} \hat{\mathbf{S}},$$

$$V_{\text{S}^2} = \frac{2c_F^2(\mu)}{3} \sum_i (\mathcal{I}(B_i(t, r)B_i(0, 0))) (\hat{\mathbf{S}}_1 \hat{\mathbf{S}}_2),$$

...

$$\text{with } \mathcal{I}(F_i^1(t_f, r_2)F_j^2(t_i, r_1)) = \lim_{T \rightarrow \infty} \int_0^T dt \, t^s \langle\langle g^2 F_i^1(t_f, r_2)F_j^2(t_i, r_1) \rangle\rangle_{(c)}$$

$c_F(\mu)$: Matching coefficient for B -insertions.

Indices \parallel, \perp, \perp' : Paralell or orthogonal directions, relative to $\bar{q}q$ -separation axis.

Differences for gradient flow

- Insertions without gradient flow:
 - Small signal-to-noise ratio due to **UV-fluctuations**.
 - Sizeable **discretization errors** with poor convergence .
 - B -insertions **log-divergent** in μ , lattice artefacts difficult to remove.
- In gradient flow:
 - + **UV-fluctuations** effectively suppressed.
 - + **Additional renormalization** of insertions not necessary $r_f/a > 1$.
 - + **Divergencies** regulated by t_f .
 - Effect of r_f larger than loop separations unknown and may be complicated to handle.
- Find the best compromise: Use flow times *large enough* to make precise measurements, while keeping r_f *small compared to separations*.

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Correlators

- **Integration** of correlation functions: **Fit ansatz** and integrate analytically.
- Motivate ansatz from **spectral decomposition**:

$$\langle\langle F^2(t_2, r_2) F^1(t_1, r_1) \rangle\rangle = \sum_m D_m^{12} e^{-\Delta V_m t} (1 + \dots)$$
with $\Delta V_m = V_m - V_0$.
- $(1 + \dots) \rightarrow 1$ for $\Delta t \rightarrow \infty$, or $|\langle 0|n, r\rangle|^2 / |\langle 0|0, r\rangle|^2 \Rightarrow$ use additional **ground state enhancement**.
- Here: APE-smearing¹ with N_{APE} adjusted such that $\ln(W(t_W, r)/W(t_W + a, r))$ forms plateau for small t_W .
- ΔV_m 's can be extracted from **combined** fits for multiple correlators at fixed t_f, r .
- Number of exponentials m_{max} and range for integration must be chosen carefully!

¹Spatial HYP- or gradient-flow-smearing planned for future computations.

Noise reduction for colour fields

Preliminary

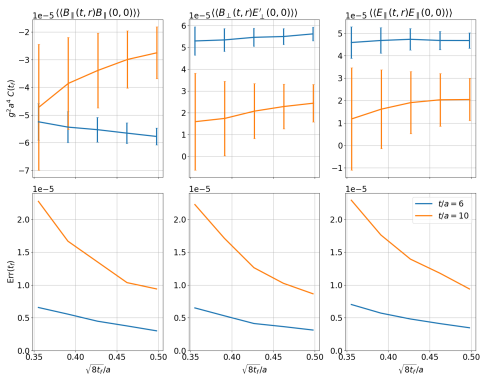


Figure: Noise reduction for different correlation functions for $a = 0.06$ fm at $r/a = 8$ as function of t_f . **Top** row: Data points with error bars. **Bottom** row: Statistical errors. **The effect is most noticeable**, if t_f is fixed in physical units when taking the continuum limits, as this goes as function of t_f/a^2 .

Gradient flow vs. smearing

Treat correlators with gradient flow, or smearing? (talk by Rainer Sommer)

- Smearing possible, if correlator contains E -field insertions only. B -field insertions are log-divergent and correlators **cannot** be continuum extrapolated.
- $t_f/a^2 = \text{const.}$ for smearing, amount of **noise reduction remains similar** for ensembles at different lattice spacings.
- For physical gradient flow, larger t_f/a^2 are accessed for decreasing lattice spacings, giving **more noise reduction**.
- Continuum extrapolated result at finite t_f has to be related to $t_f = 0$, either by **extrapolation**, or **matching** (e.g. with B -fields).
→ Matching coefficients (without GF) known in $\overline{\text{ms}}$, computation in gradient flow might be similar.

⇒ Gradient flow seems to be the most sensible choice here, but matching to $t_f = 0$ has to be worked out.

Lattice ensembles

- Simulation: $SU(3)$ heatbath with CL2QCD.
- Wilson flow with adaptive step size.
- Whole Wilson loop, including insertions, flowed.
- Statistical errors propagated via bootstrap samples.

Table: Lattice ensembles shown in the following (black), used, but not shown (dark grey), or planned (light grey).

β	T/a	L/a	a [fm]	# confs.
6.091	36	18	0.08	20000
6.284	48	24	0.06	10000
6.451	60	30	0.048	—
6.594	36	18	0.04	—

Static potential

Preliminary

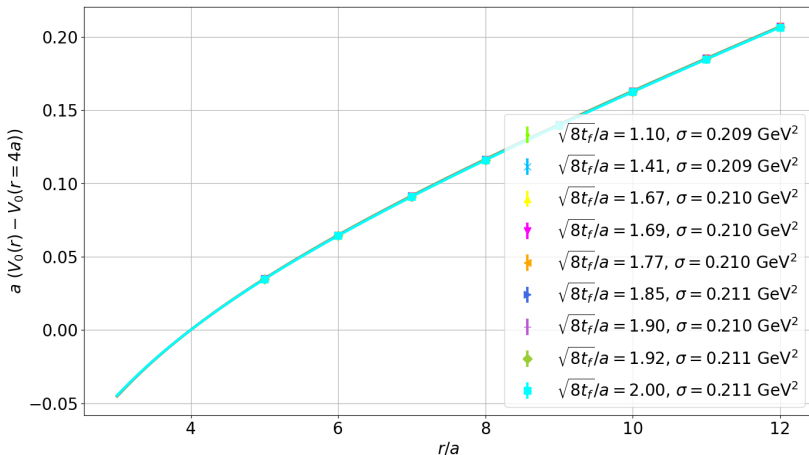


Figure: Static potential for $a = 0.06$ fm, normalized with $V_0(r/a = 4)$, at different flow times. The physical scale has been fixed with $r_0 = 0.5$ fm.

Noise reduction for colour fields

Preliminary

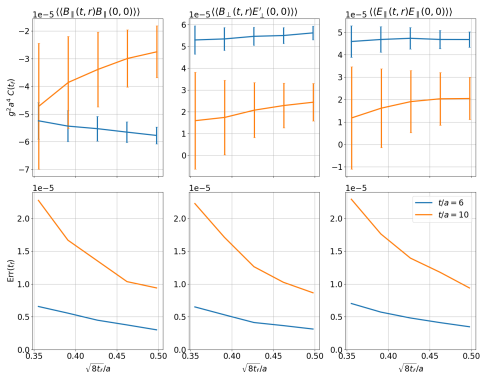


Figure: Noise reduction for different correlation functions for $a = 0.06$ fm at $r/a = 8$ as function of t_f . **Top** row: Data points with error bars. **Bottom** row: Statistical errors.

Fitting of correlation functions I

Preliminary

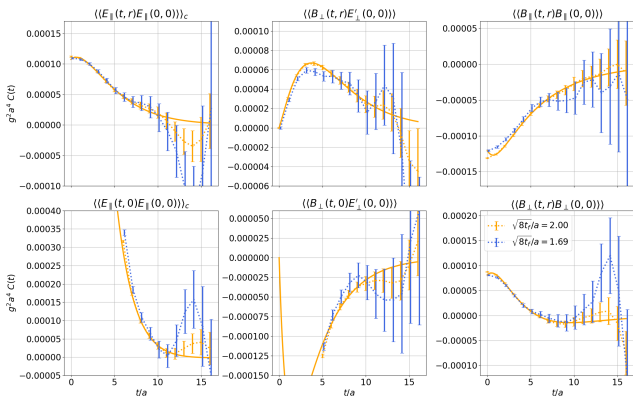


Figure: (Unweighted) correlators at $r/a = 8$ for $a = 0.06$ fm as function of euclidean time t . For $\sqrt{8t_f}/a = 2$ (orange), the fit result is drawn. The first data point was left out of the fitting procedure, for demonstration purposes.

Fitting of correlation functions II

Preliminary

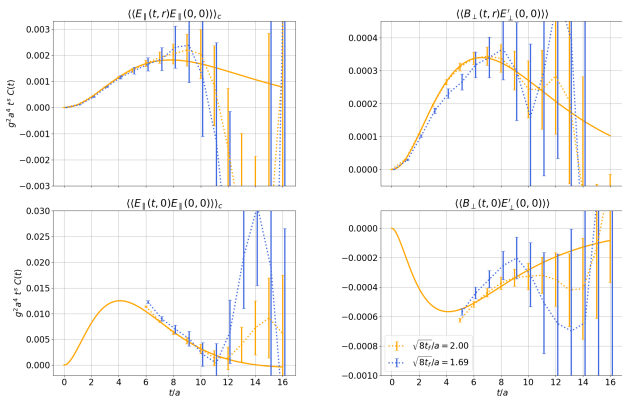


Figure: Weighted correlators at $r/a = 8$ for $a = 0.06$ fm as function of euclidean time t . For $\sqrt{8t_f}/a = 2$ (orange), the fit result is drawn. The first data point was left out of the fitting procedure, for demonstration purposes.

Conclusions

Summary

- $\mathcal{O}(1/m^2)$ -corrections to the static $\bar{q}q$ -potential have been computed at $a = 0.06$ fm and several values of t_f for distances up to ~ 0.7 fm.
- *Effective reduction of noise* in colour-field correlators.
- Good ground to *test effect of gradient flow* on colour-fields.
- Data at small separations not accessible due to $r_f/a > 1$, *difficult to get robust fits*.

Outlook

- Do *continuum* extrapolation and relate it to $t_f = 0$, e.g. via matching (perturbative calculation required).
- Improve *robustness* of correlator fits, e.g. *combined fitting* for multiple correlators.
- Access *smaller* min. separations < 0.2 fm.