

Chiral Anomaly in Effective Open Quantum Systems

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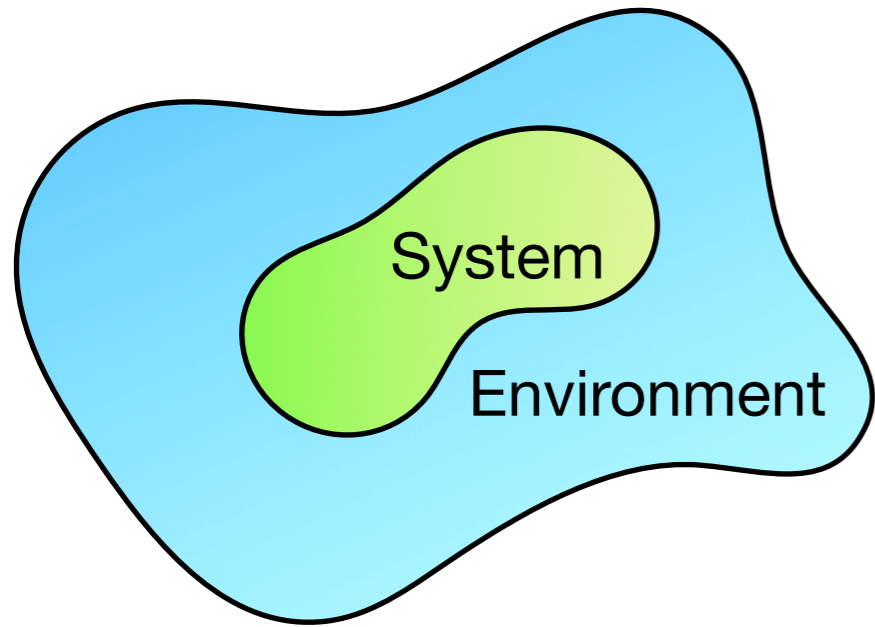
In collaboration with:

Julia Hannukainen (KTH, Sweden)

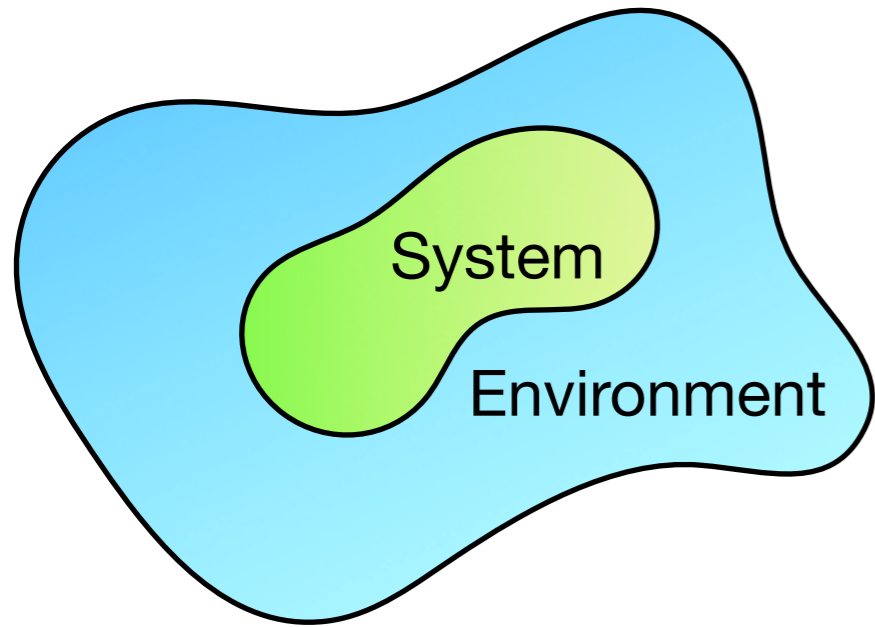
Adolfo Grushin (Institut Néel, France)

S.Sayyad, et al., Phys. Rev. Research 4, L042004 (2022)

#Open quantum systems



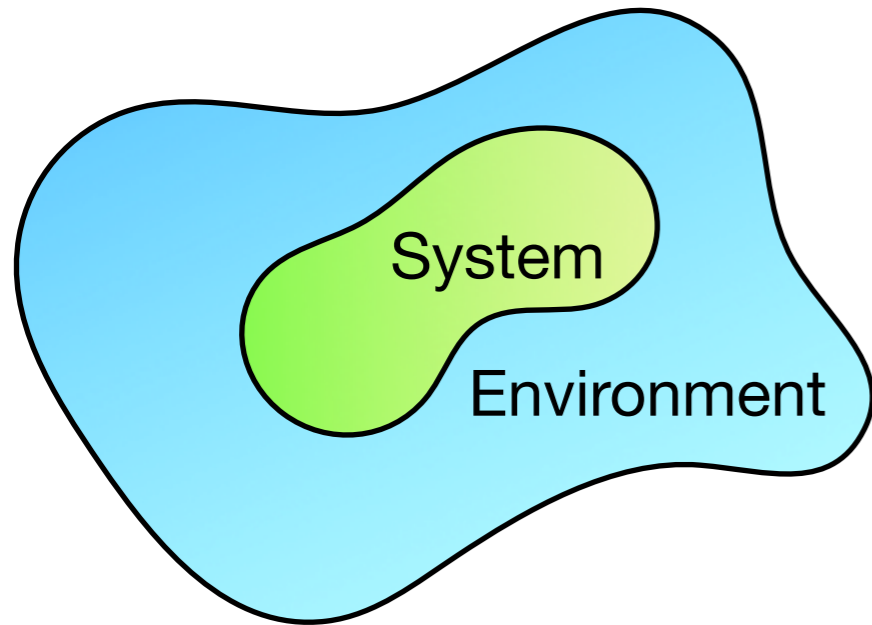
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Closed quantum systems



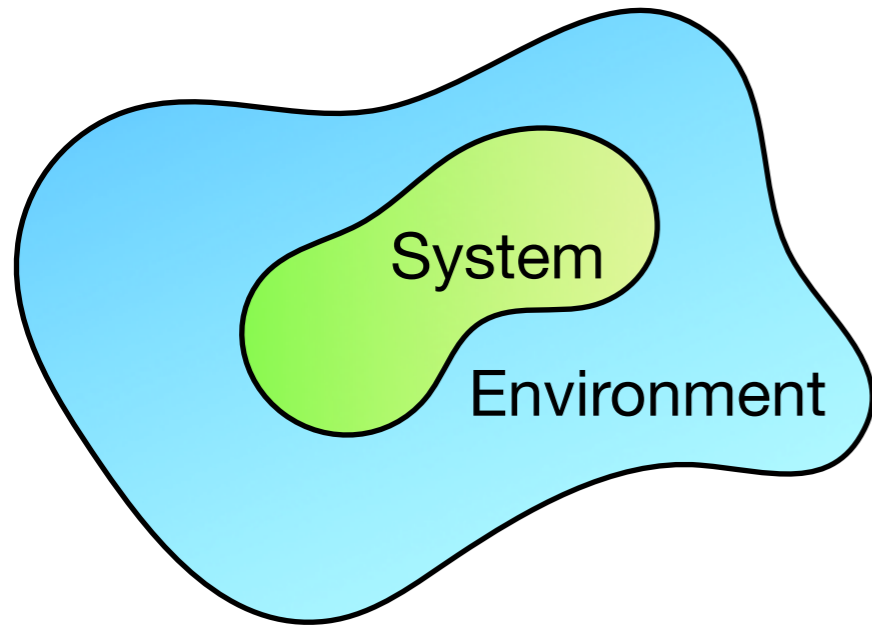
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$$\mathcal{H}_{\text{tot}} = \mathcal{H}_{\text{sys}} \otimes 1_{\text{env}} + 1_{\text{sys}} \otimes \mathcal{H}_{\text{env}} + \mathcal{H}_{\text{sys-env}}$$

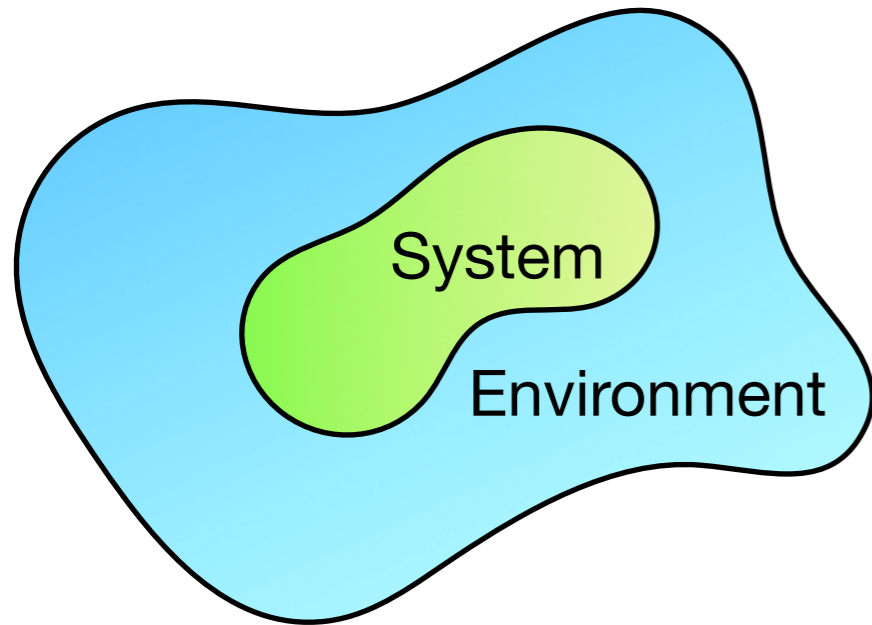
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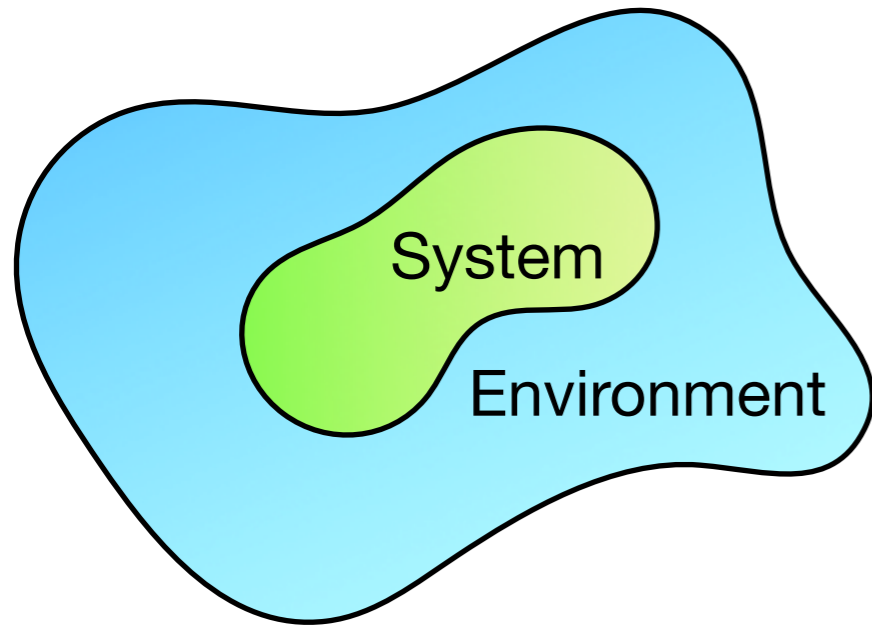
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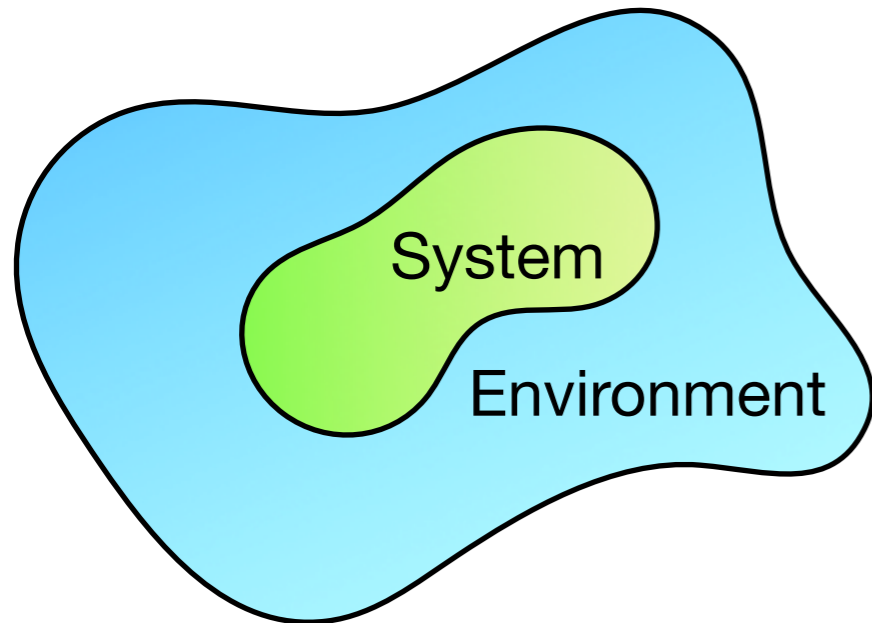
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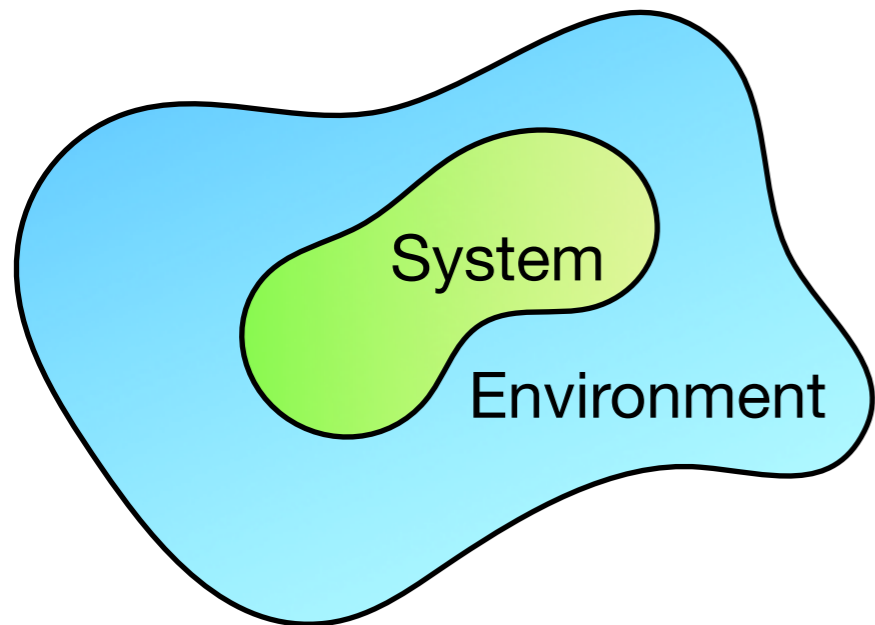
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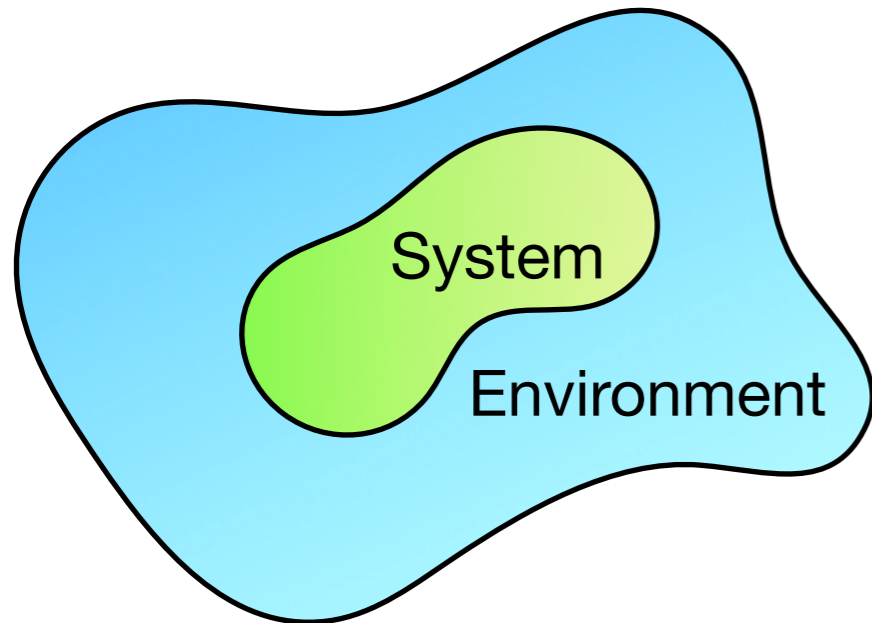
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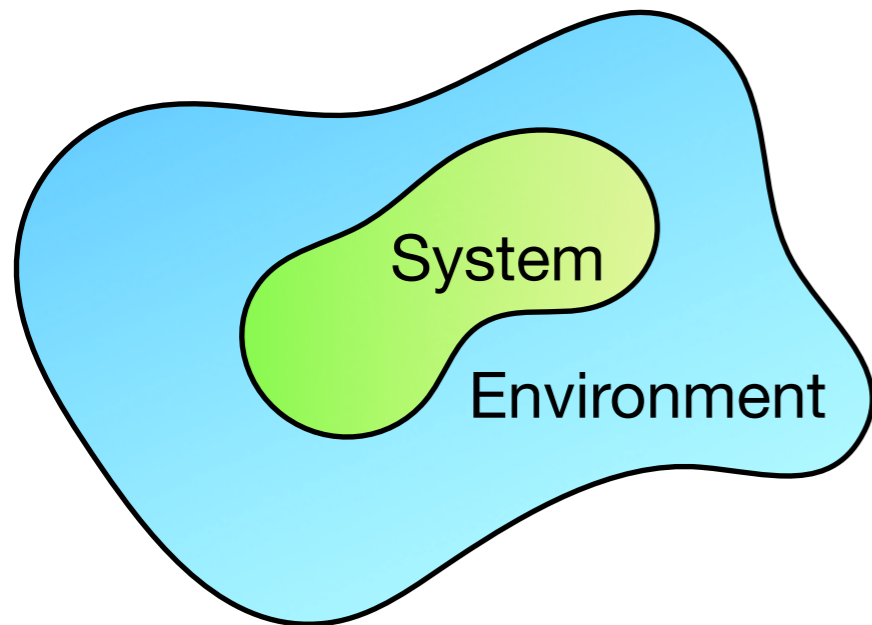
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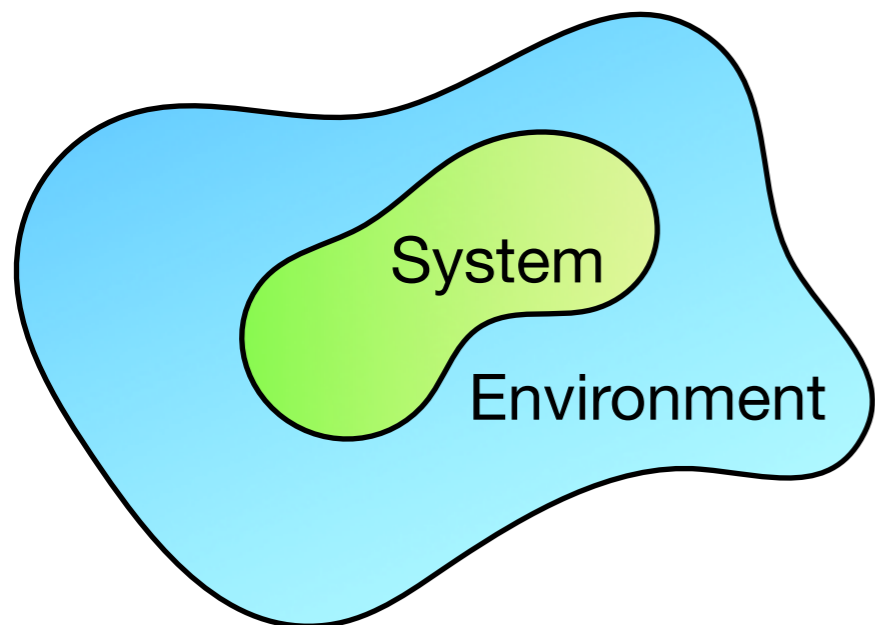
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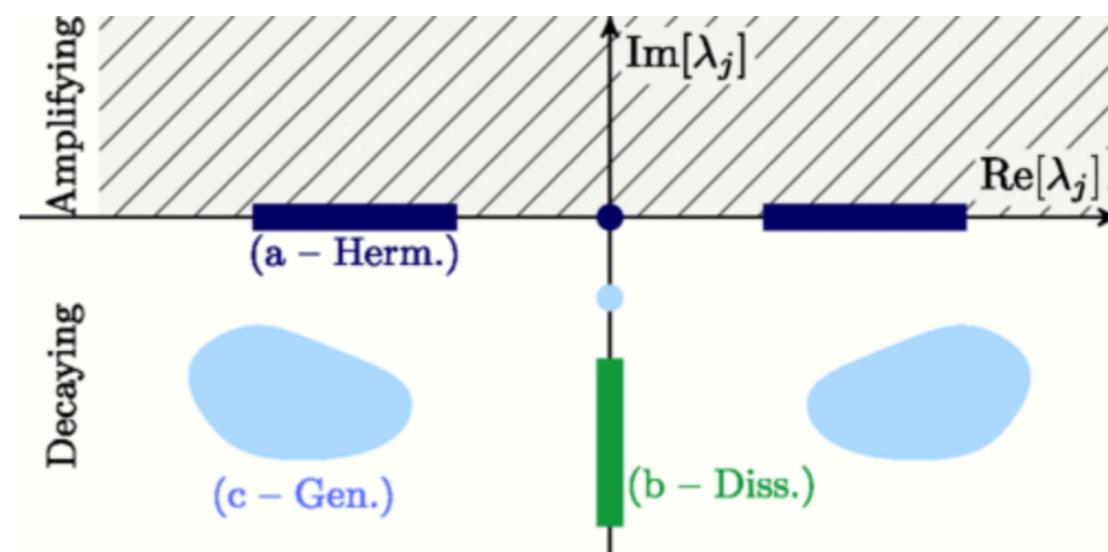
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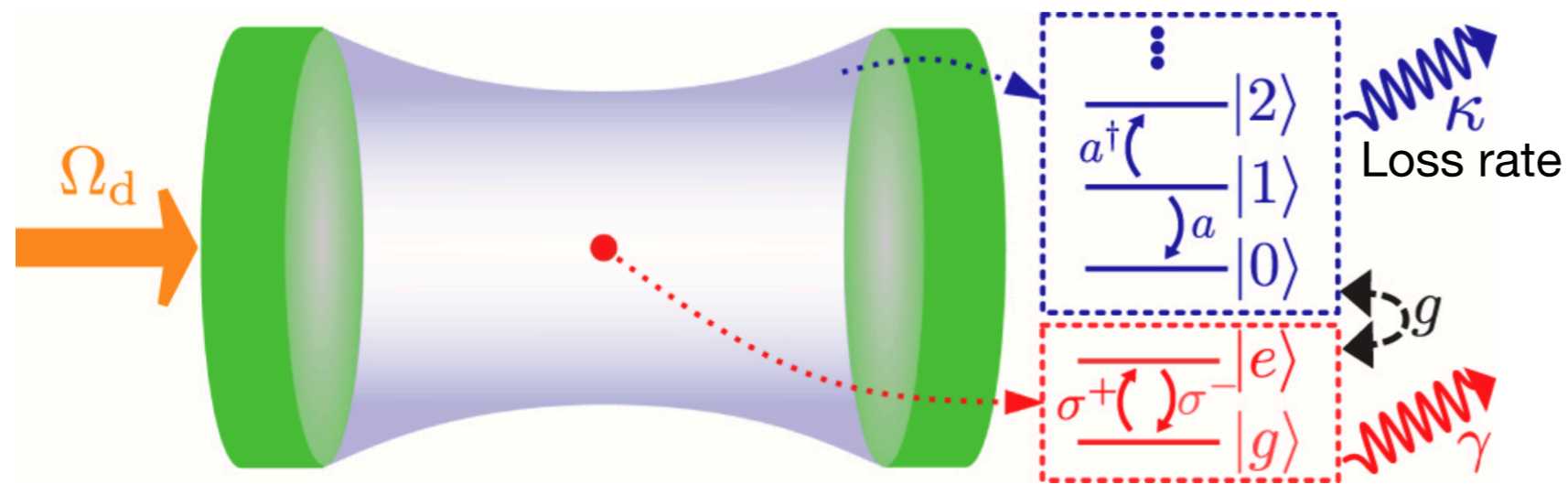
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S. Lieu, et al. PRL(2020)

#Open quantum systems in transport experiments

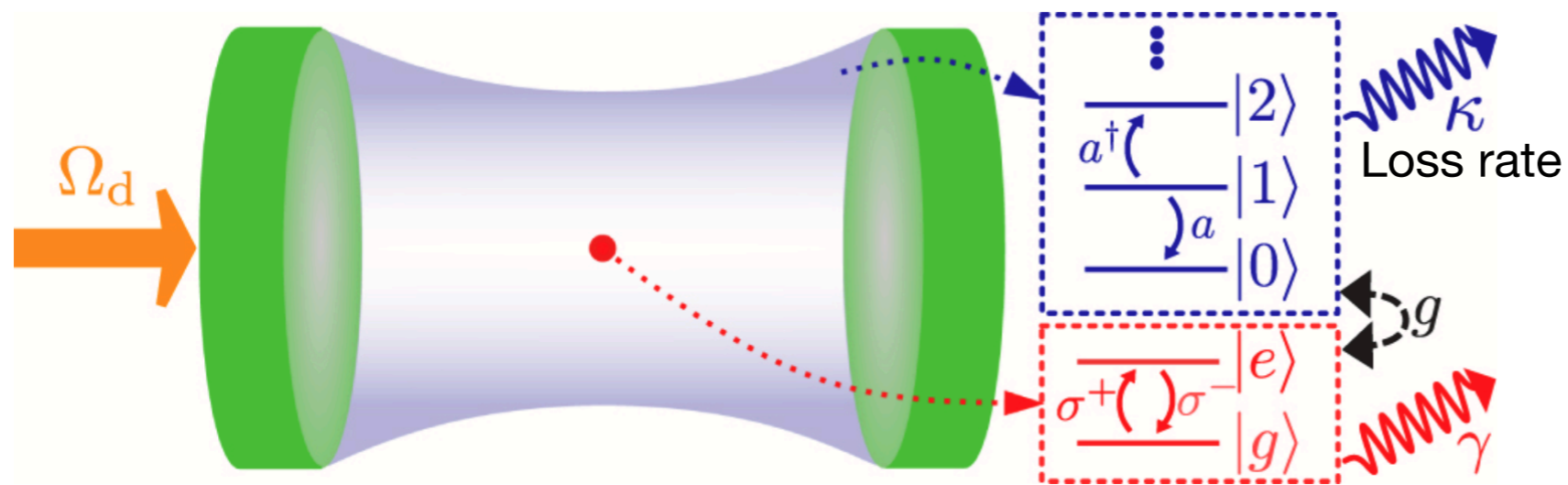
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Quantum optics
Strong light-matter interaction

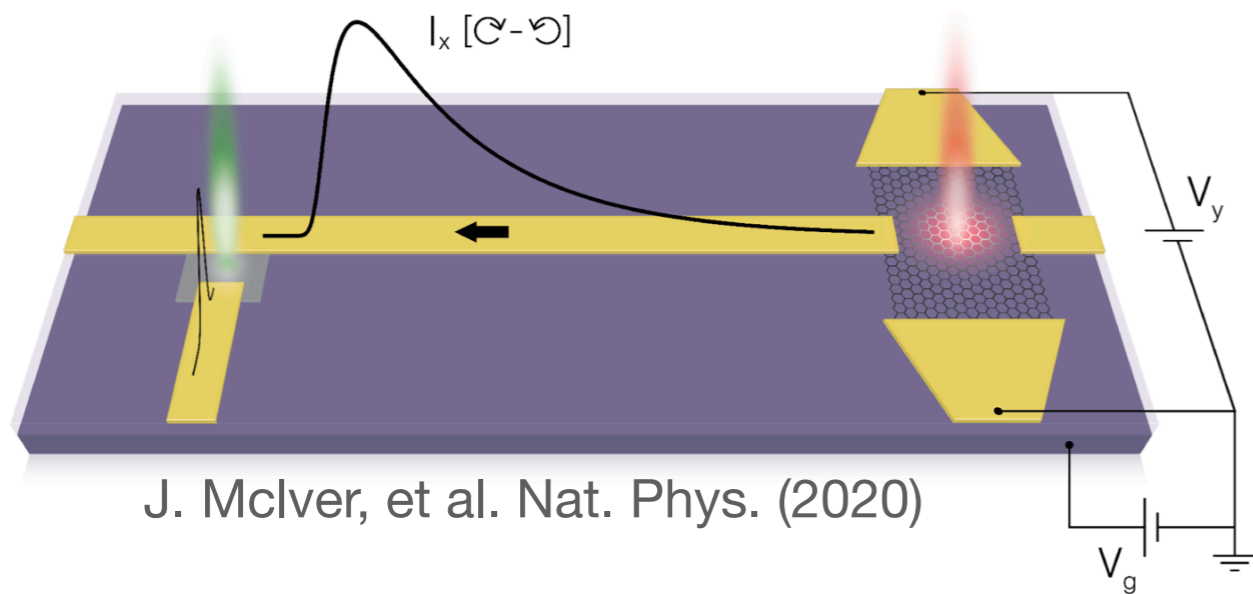
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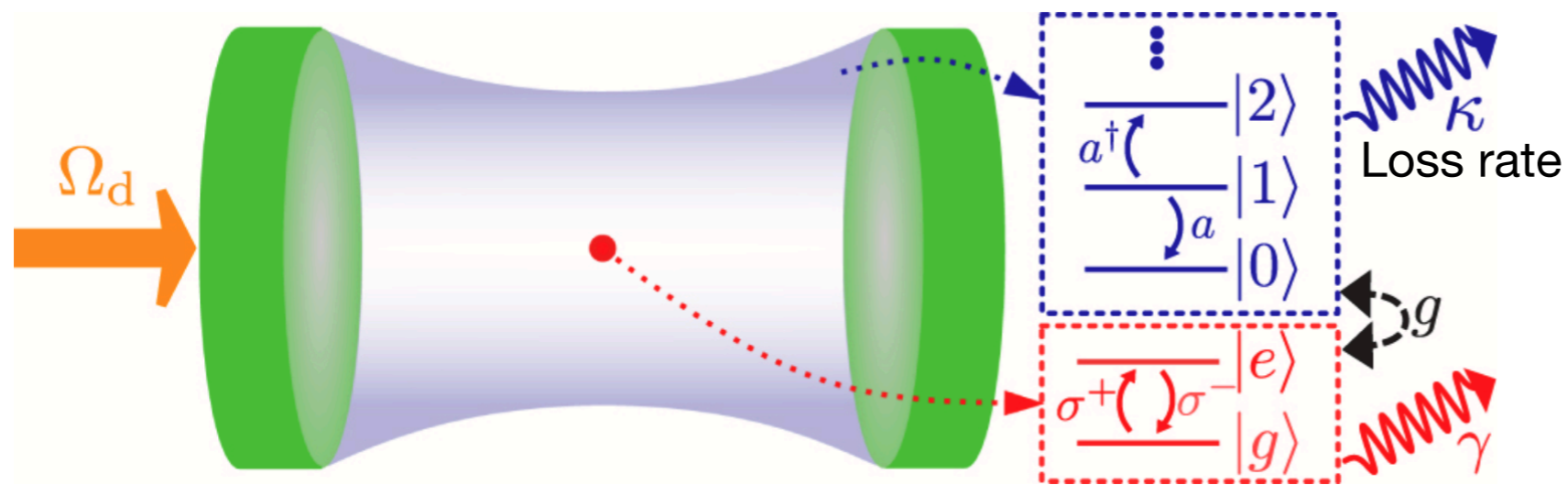
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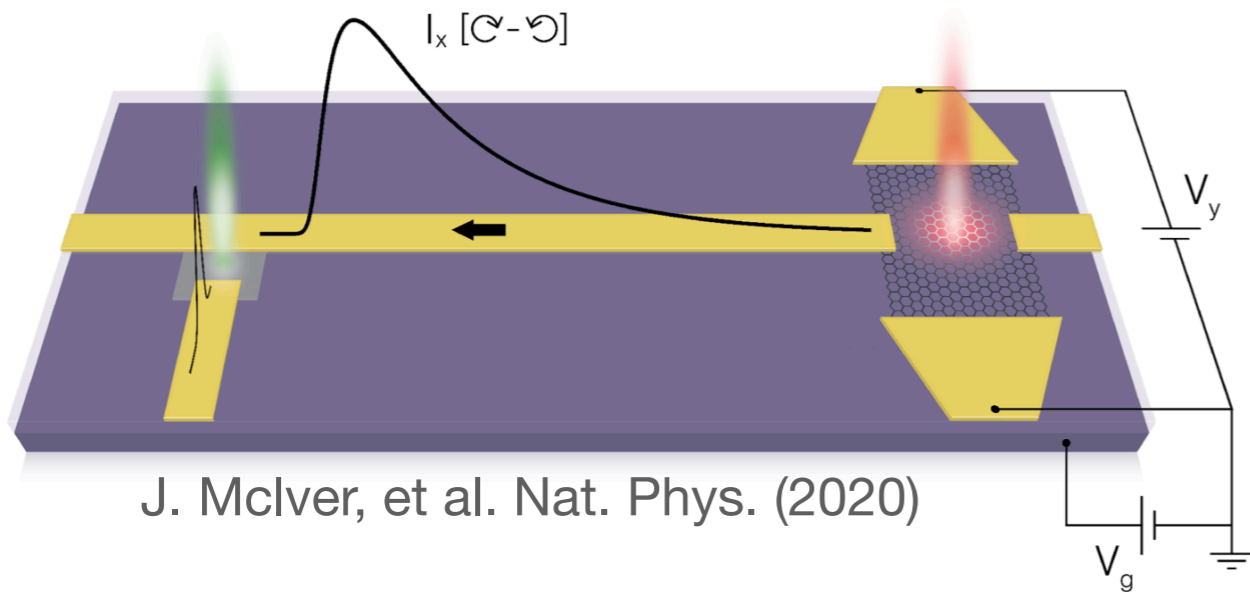
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Light-driven systems
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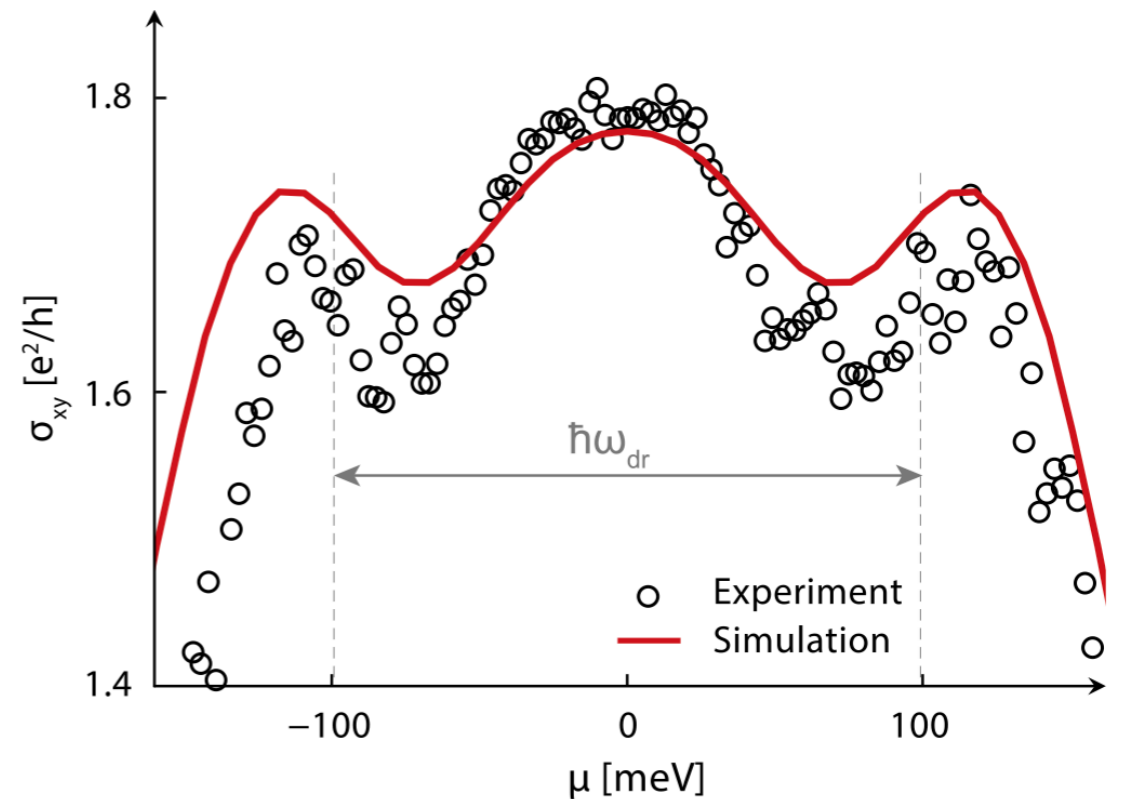
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M. Nuske, et al. *PRR* (2020)

#Effective description of open quantum systems

$$\frac{d\rho_{\text{sys}}}{dt} = -i[\mathcal{H}_{\text{sys}}, \rho_{\text{sys}}] + \sum_i (2L_i^\dagger \rho_{\text{sys}} L_i - \{L_i^\dagger L_i, \rho_{\text{sys}}\}) \equiv \hat{\mathcal{L}}[\rho_{\text{sys}}]$$

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Generating effective **non-Hermitian** Hamiltonians

Momentum-based effective Hamiltonian

Third quantization

Vectorized form of Lindbladian

Effective recycled Hamiltonian

Effective Hamiltonian

Dynamical matrix

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$$\dot{\rho}_{\text{sys}} = -i(H_{\text{eff}}\rho_{\text{sys}} - \rho_{\text{sys}}H_{\text{eff}}^\dagger) + \sum_i 2L_i\rho_{\text{sys}}L_i^\dagger$$
$$H_{\text{eff}} = H_{\text{sys}} - i \sum_l L_l^\dagger L_l$$

#Hermitian vs. non-Hermitian systems

Hermitian QM

Non-Hermitian QM

Eigenvectors

Spectrum

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Orthogonal Eigenvectors



Bi-orthogonal/Self-orthogonal
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Localization of bulk eigenvectors
close to boundaries
(NH skin effect)

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Real-valued Eigenvalues

Complex-valued Eigenvalues

Degenerate eigenvalues

Degenerate eigenvalues/
Coalescence of eigenvalues and eigenvectors

Real-line gaps

Real/imaginary-line gaps
& point gaps

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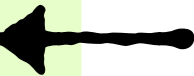
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Non-defective degeneracy

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Both eigenvalues and eigenvectors coalesce

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$$\mathcal{H} = \begin{pmatrix} 0 & \alpha \\ 1 & 0 \end{pmatrix} \quad E_{\pm} = \pm \sqrt{\alpha}$$

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Defective degeneracy at $\alpha = 0$
& self-orthogonal basis $\langle \psi_L | \psi_R \rangle = 0$

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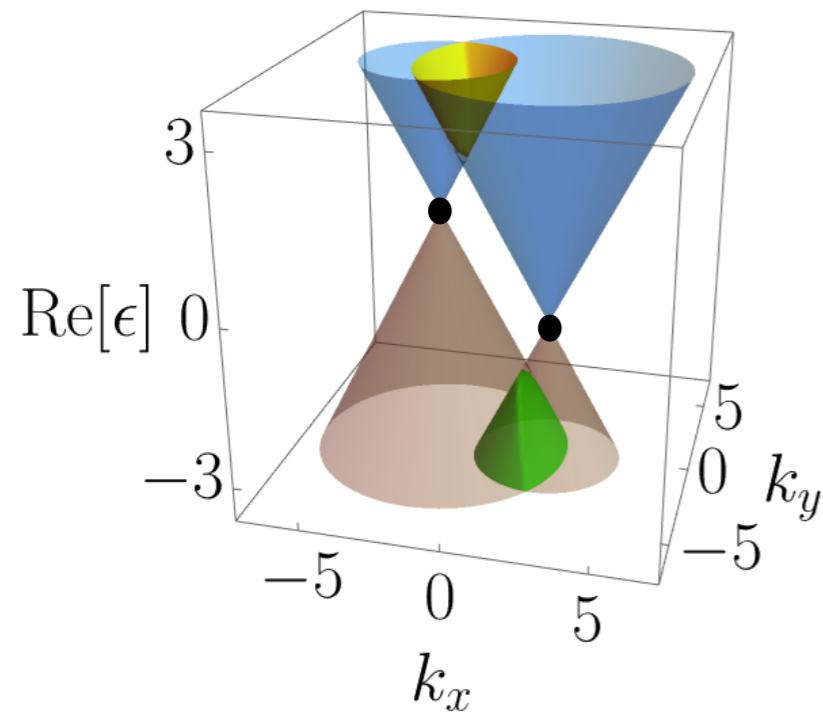
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$$H_{\text{sys}} = \pm [b_0 + M_i^j (k_j \mp b_j) \sigma^i]$$

Real Fermi velocity: $M = \text{diag}[v_1, v_2, v_3]$

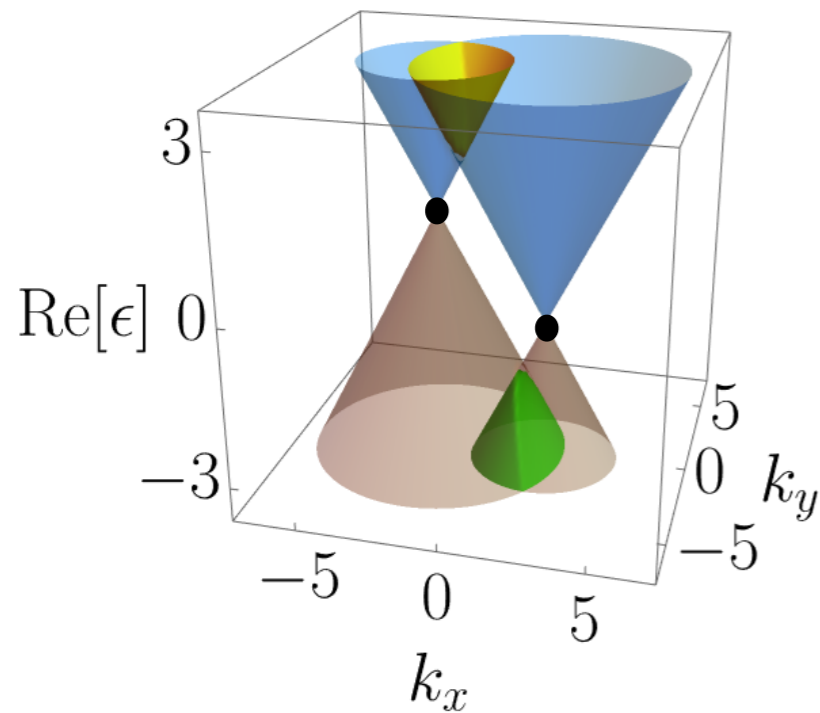


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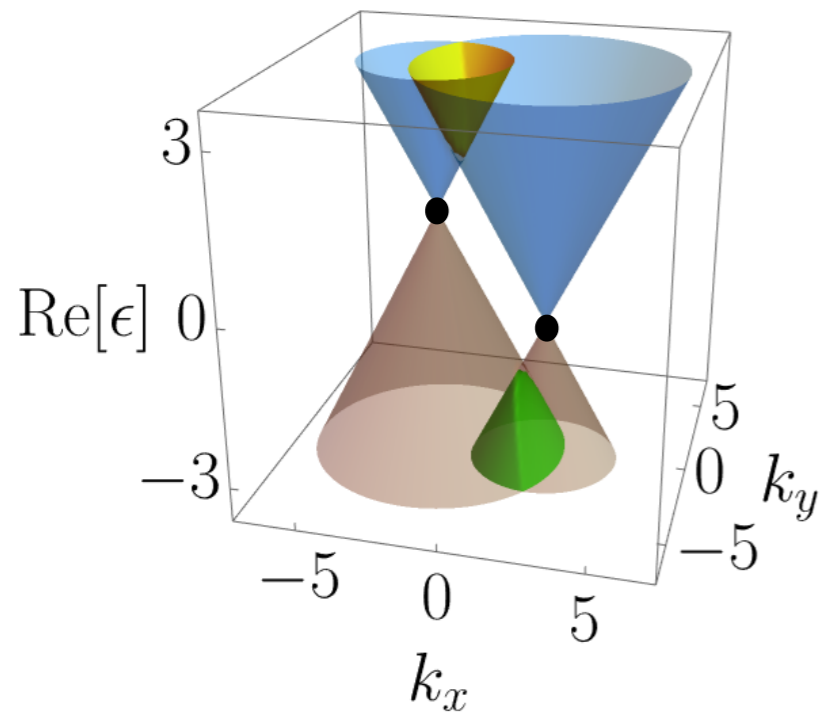
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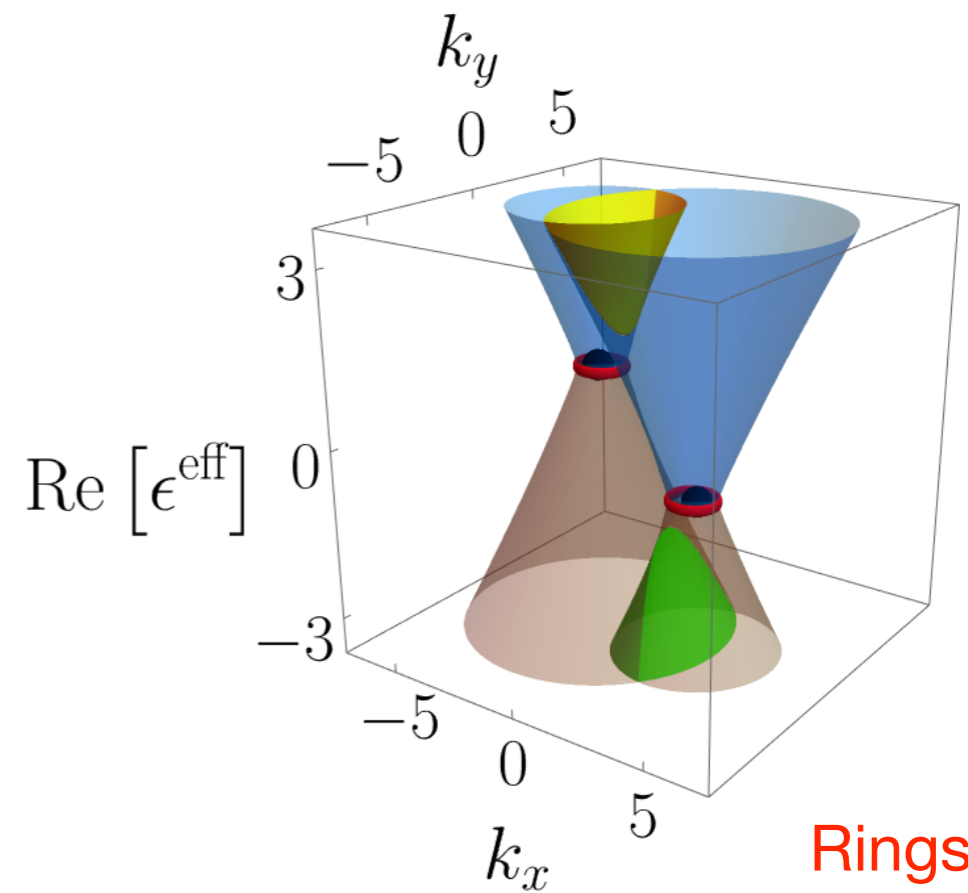
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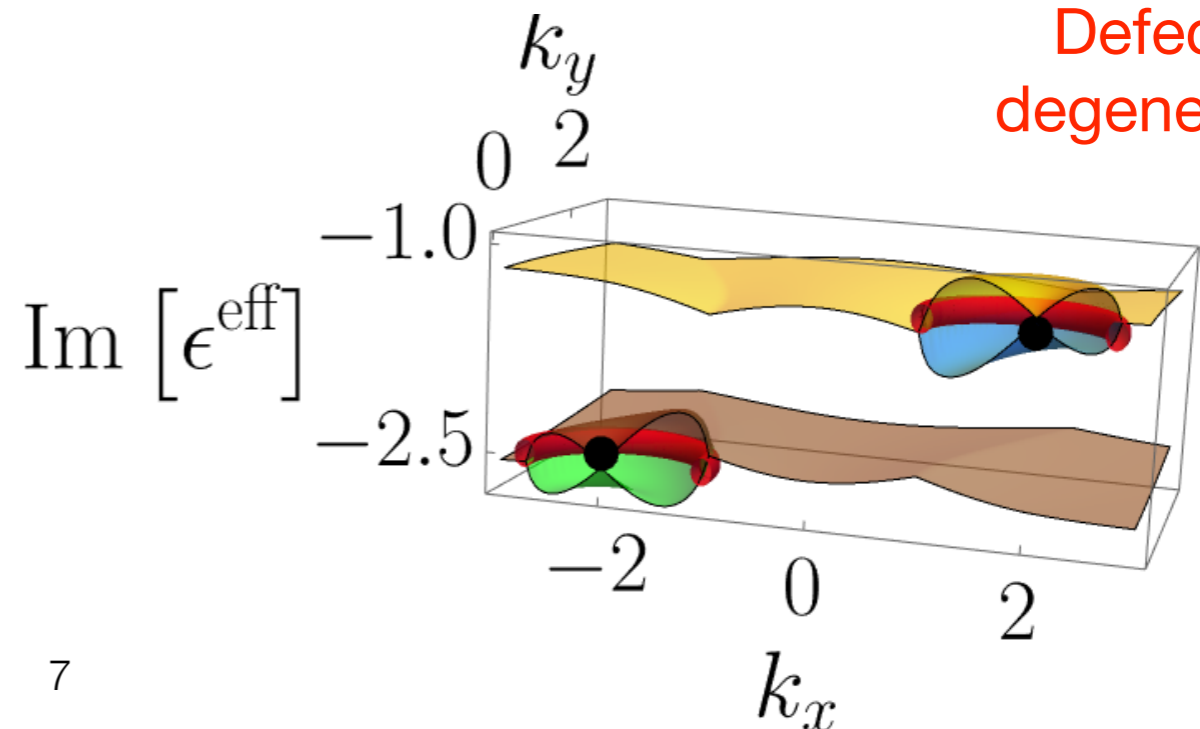
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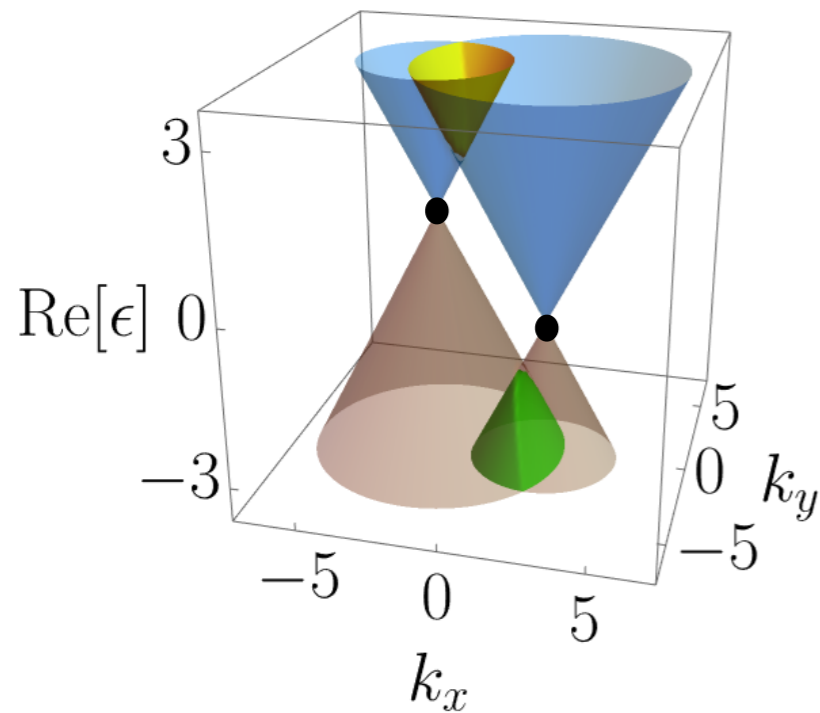
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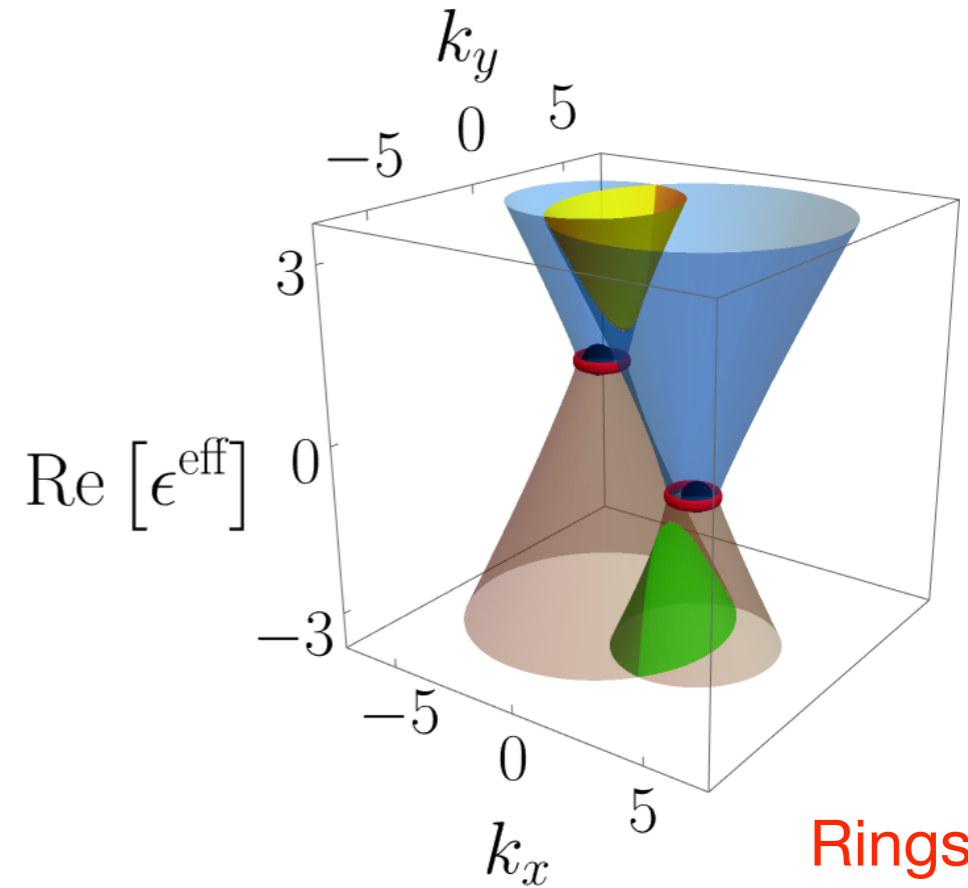
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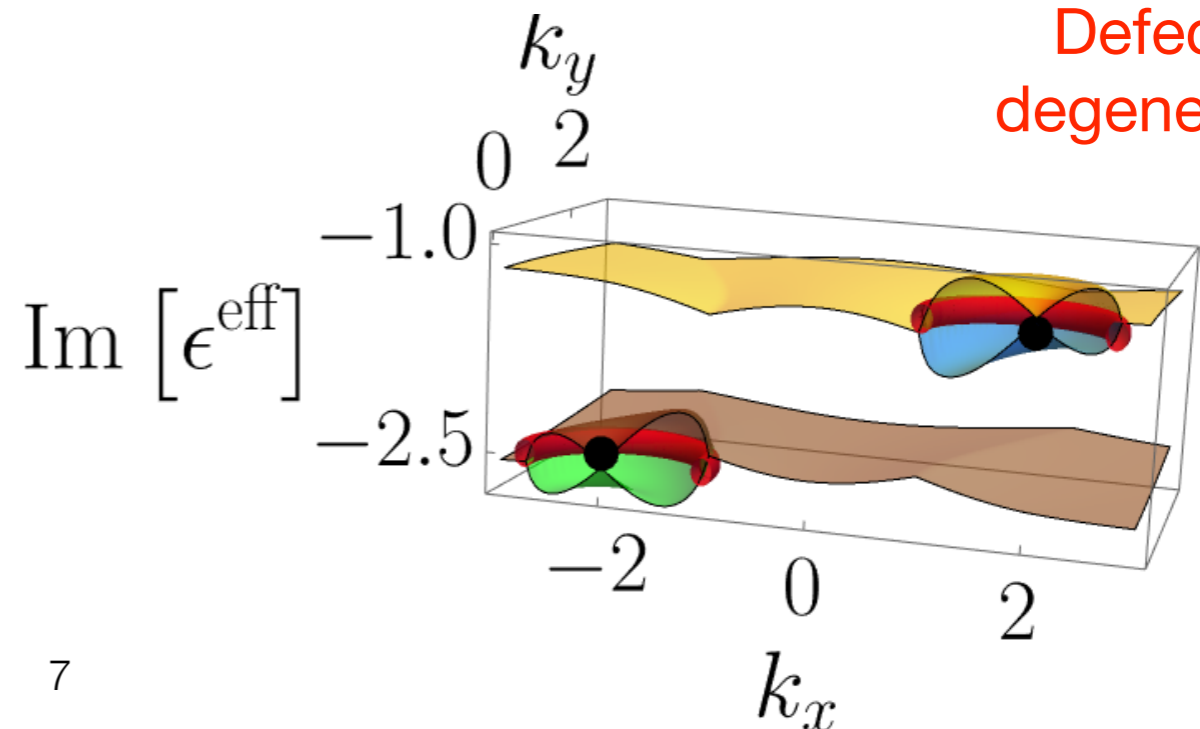


Rings:
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In the vicinity of
non-defective degeneracies

$$H_{\text{eff}}^{\text{low-en}} = \pm [ib_0 + iM_i^j (k_j \mp b_j) \sigma^i]$$

Imaginary Fermi velocity



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Hermitian

$$\mathcal{D}_{\text{herm}}^2 |\psi_n\rangle = \lambda_n^2 |\psi_n\rangle$$

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$$\mathcal{D}[A^c, A^{5c}] \neq \mathcal{D}^\dagger[A^c, A^{5c}]$$

Conserved classical currents

$$\partial_\mu j^\mu = 0 \quad \partial_\mu j_5^\mu = 0$$

How about divergence of quantum currents?

Fujikawa method + Heat-kernel regularization

Hermitian

$$\mathcal{D}_{\text{herm}}^2 |\psi_n\rangle = \lambda_n^2 |\psi_n\rangle$$

Non-Hermitian

$$\begin{aligned} \tilde{\mathcal{D}} \tilde{\mathcal{D}}^\dagger |\eta_n\rangle &= |\lambda_n|^2 |\eta_n\rangle, & \tilde{\mathcal{D}}^\dagger \tilde{\mathcal{D}} |\xi_n\rangle &= |\lambda_n|^2 |\xi_n\rangle, \\ \tilde{\mathcal{D}}^\dagger |\eta_n\rangle &= \lambda_n^* |\xi_n\rangle, & \tilde{\mathcal{D}} |\xi_n\rangle &= \lambda_n |\eta_n\rangle. \end{aligned}$$



#Chiral anomaly in non-Hermitian systems

$$\mathcal{Z} \propto \int \mathcal{D}\Psi \mathcal{D}\bar{\Psi} e^{-\int d^d x \left[\bar{\Psi} \gamma^\mu \left(M_\mu^{c\nu} \partial_\nu - i M_\mu^{c\nu} A_\nu^c + i \gamma_5 M_\mu^{c\nu} A_\nu^{5c} \right) \Psi \right]}$$

$M^c = \text{diag}[1, \nu_1, \nu_2, \nu_3]$

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Divergence of covariant quantum currents:

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Divergence of covariant quantum currents:

$$\tilde{d}_\mu j^{5,\mu} \propto v_1 \tilde{E}_1^\dagger + v_1^* \tilde{E}_1 \quad \text{in } d = 2,$$

$$\begin{aligned} \tilde{d}_\mu j^{5,\mu} \propto v_1 v_2 v_3 (\tilde{E}^\dagger \cdot \tilde{B}^\dagger + \tilde{E}^{5\dagger} \cdot \tilde{B}^{5\dagger}) \\ + v_1^* v_2^* v_3^* (\tilde{E} \cdot \tilde{B} + \tilde{E}^5 \cdot \tilde{B}^5) \quad \text{in } d = 4. \end{aligned}$$

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$$+ v_1^* v_2^* v_3^* (\tilde{E} \cdot \tilde{B} + \tilde{E}^5 \cdot \tilde{B}^5) \quad \text{in } d = 4.$$

Electric field

$$\tilde{E}_j = e^{2i\phi_j} \partial_t A_j^c - \partial_j A_t^c$$

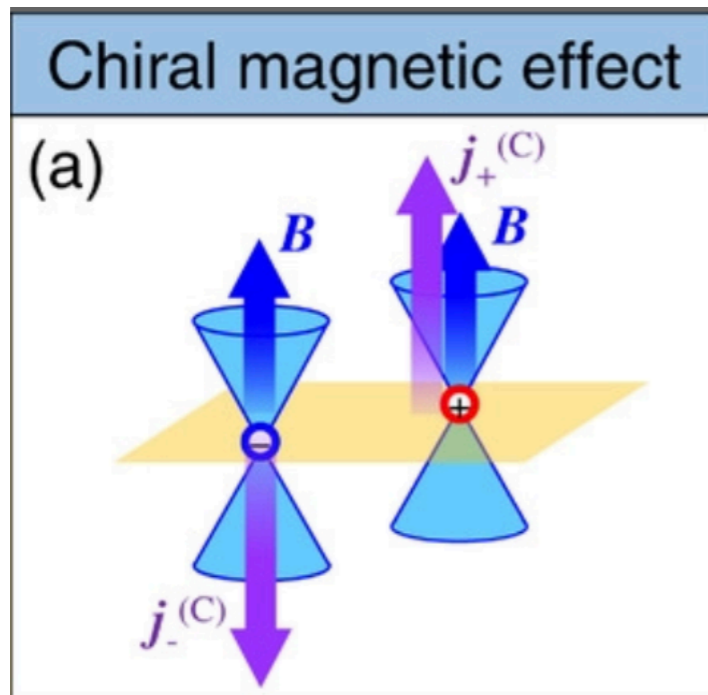
Phase of complex Fermi velocity

$$\frac{v_j}{v_j^*} = e^{i2\phi_j}$$

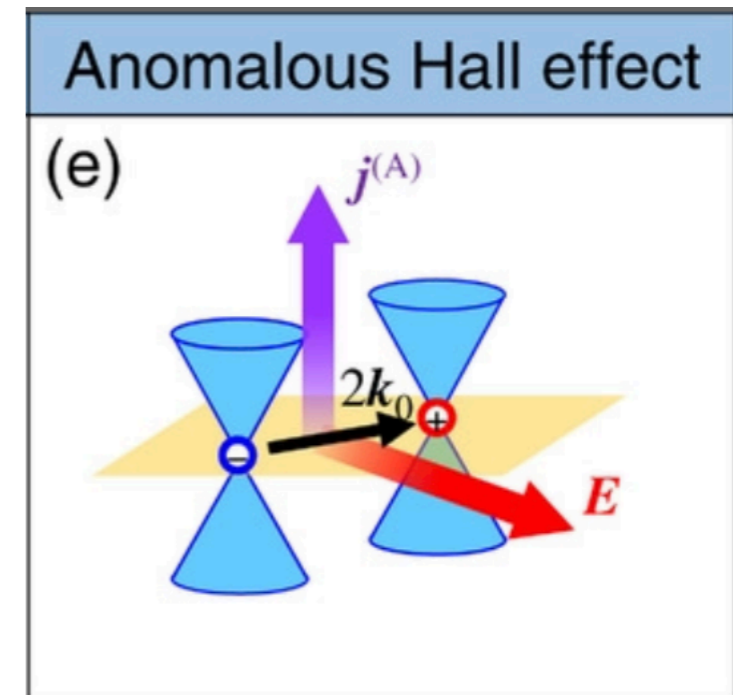
Magnetic field

$$\tilde{B}_i = \varepsilon^{ijk} \tilde{B}_{jk} = \varepsilon^{ijk} (e^{2i\phi_k} \partial_j A_k^c - e^{2i\phi_j} \partial_k A_j^c)$$

#Physical consequences of non-Hermitian chiral anomaly



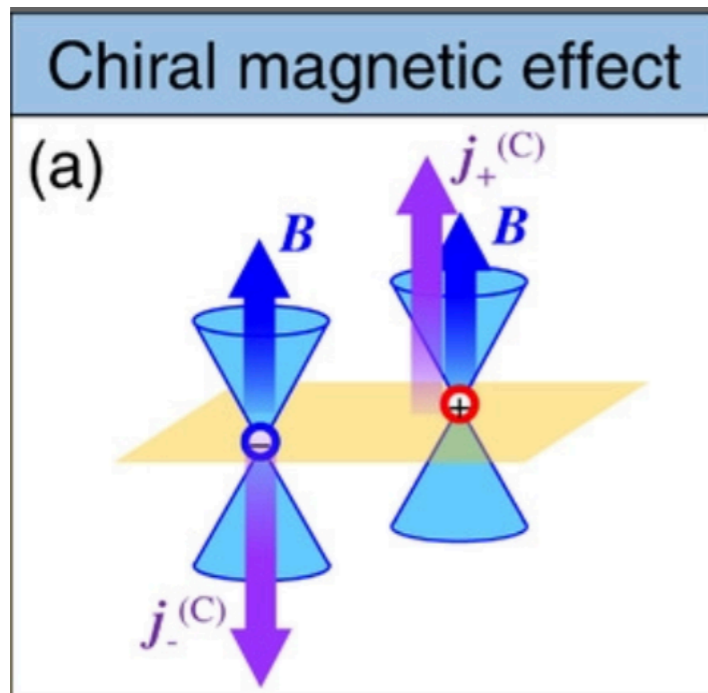
For each Weyl node $j^{(c)} \propto B$



$$j^{(A)} \propto \vec{k}_0 \times E$$

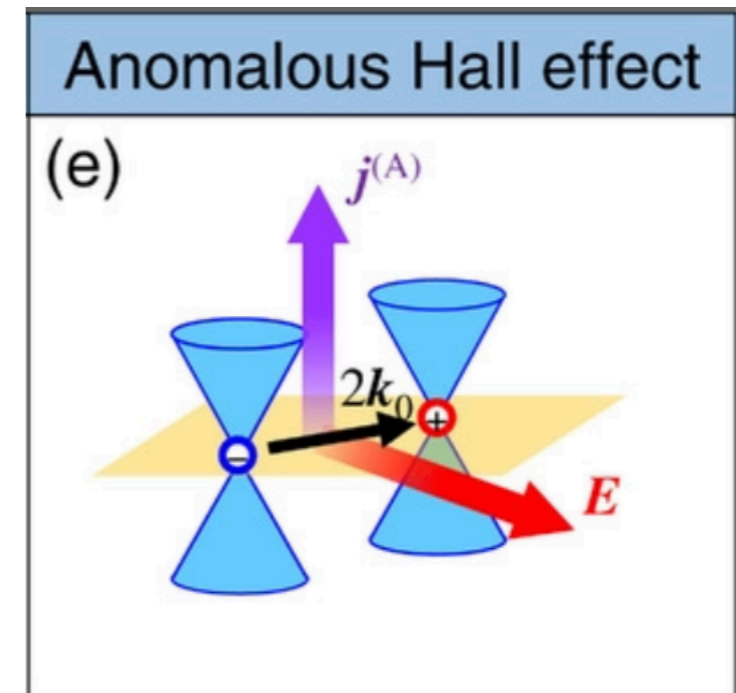
Y. Araki, Annalen der Physik (2020)

#Physical consequences of non-Hermitian chiral anomaly



For each Weyl node $j^{(c)} \propto B$

$$M_\nu^\alpha j^\nu \propto \varepsilon^{0\nu\eta\xi} \text{Re}[M_\nu^{\alpha*} M_\eta^\nu M_\xi^\rho \tilde{B}_{\rho}^\dagger]$$



$$j^{(A)} \propto \vec{k}_0 \times E$$

Y. Araki, Annalen der Physik (2020)

$$M_\nu^\alpha j^\nu \propto \varepsilon^{\mu\nu\eta\xi} \tilde{k}_{0\delta} \text{Re}[M_\nu^{\alpha*} M_\mu^\delta M_\xi^\rho \tilde{E}_\rho^\dagger]$$

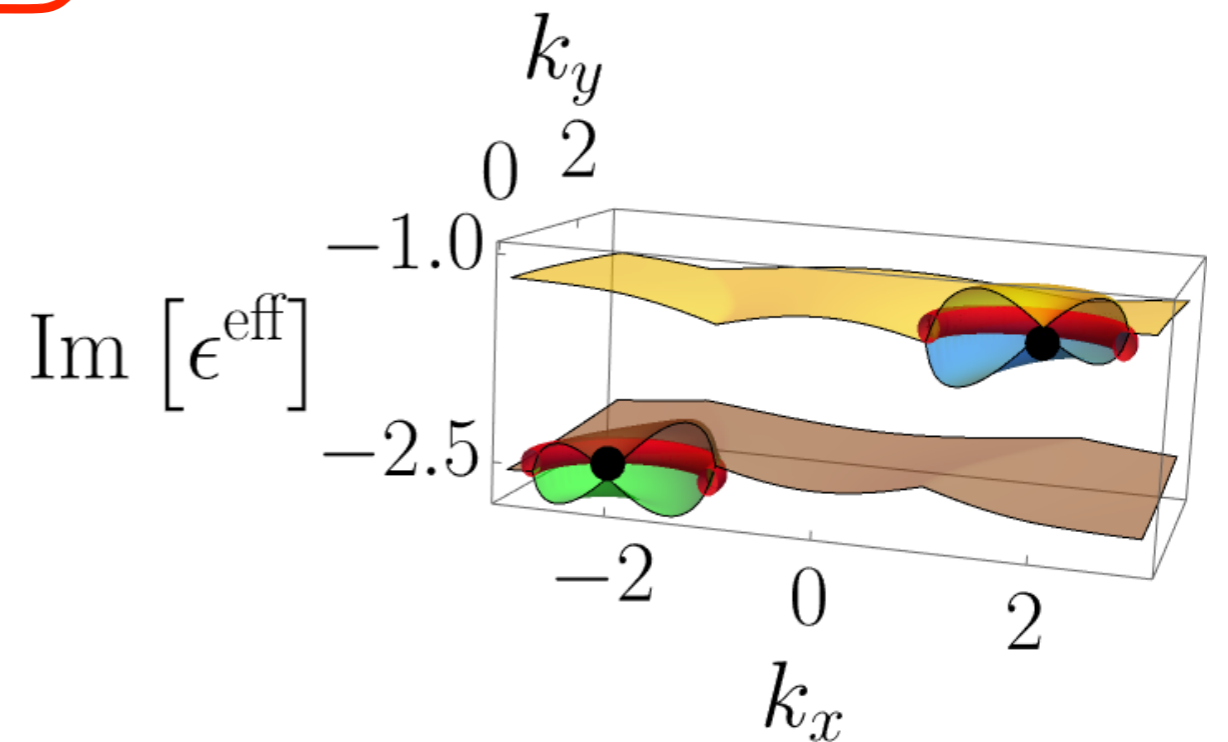
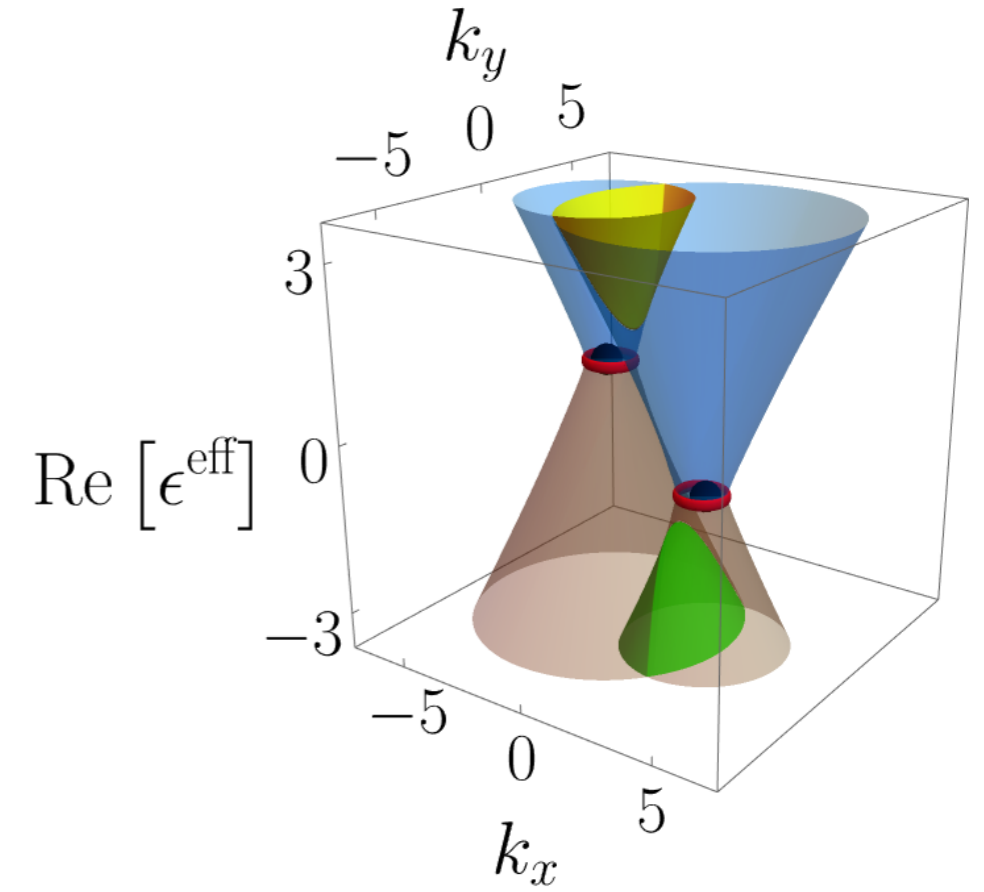
S.Sayyad, et al., PRR (2022)

#Conclusion

$$\tilde{d}_\mu j^{5,\mu} \propto v_1 \tilde{E}_1^\dagger + v_1^* \tilde{E}_1 \quad \text{in } d = 2,$$

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Effective NH Weyl semimetals



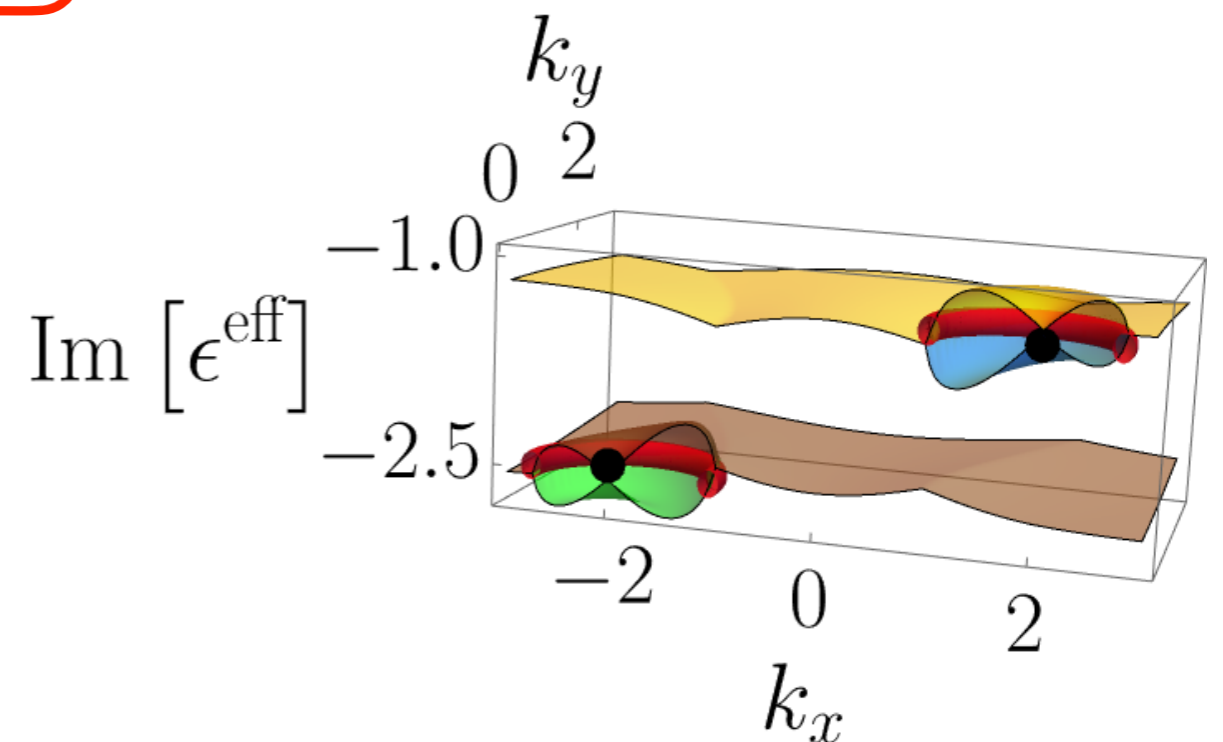
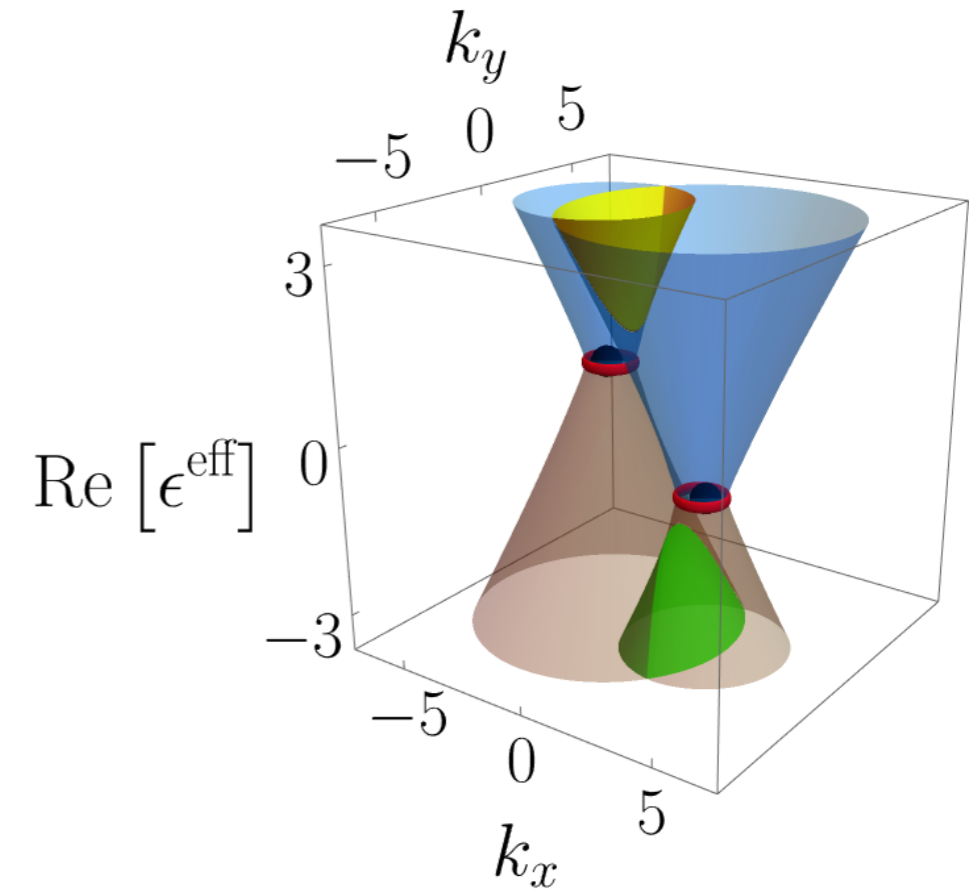
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Thank you for your attention!

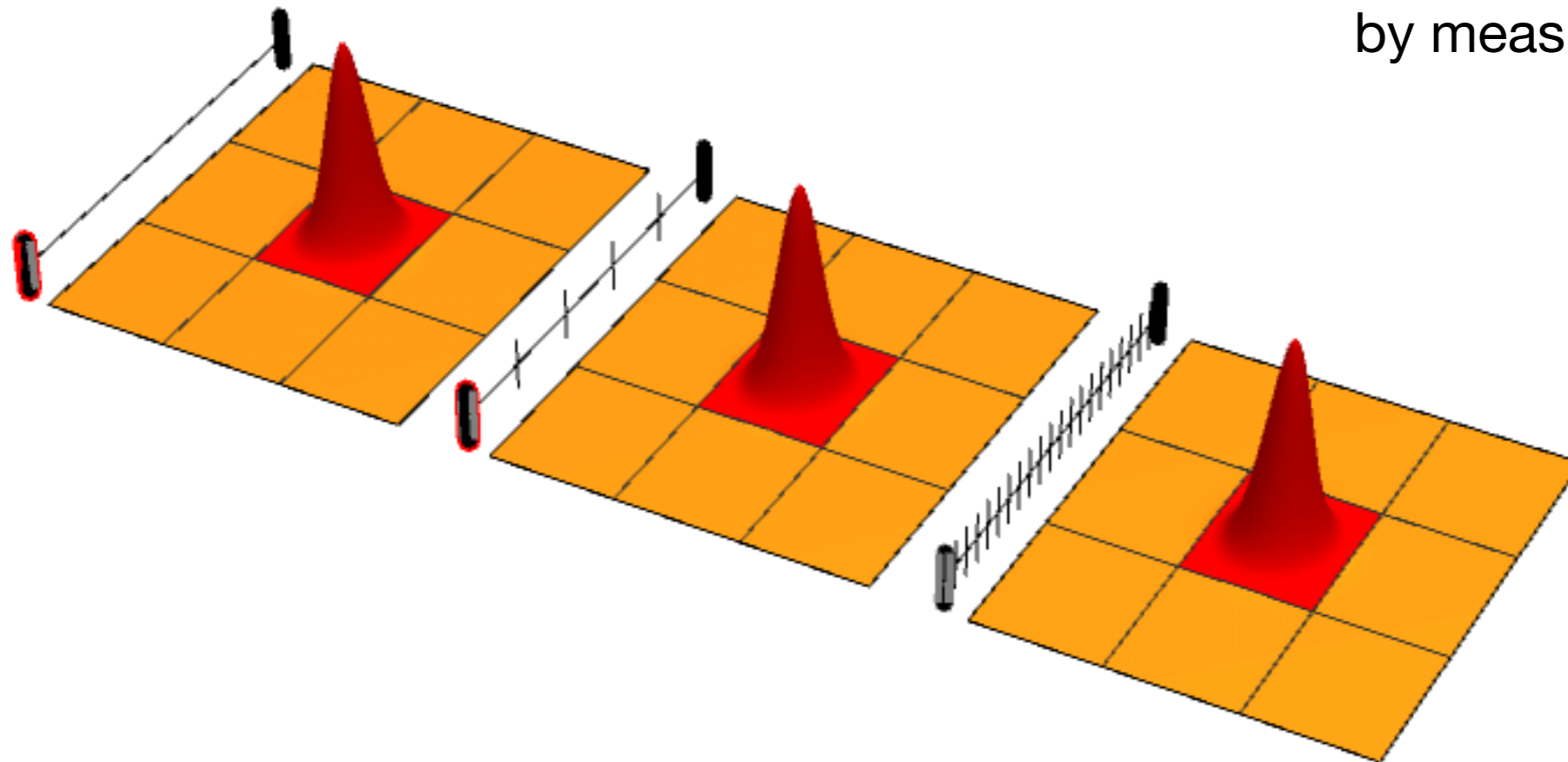
Effective NH Weyl semimetals



#Open quantum systems in transport experiments

Quantum Zeno effect

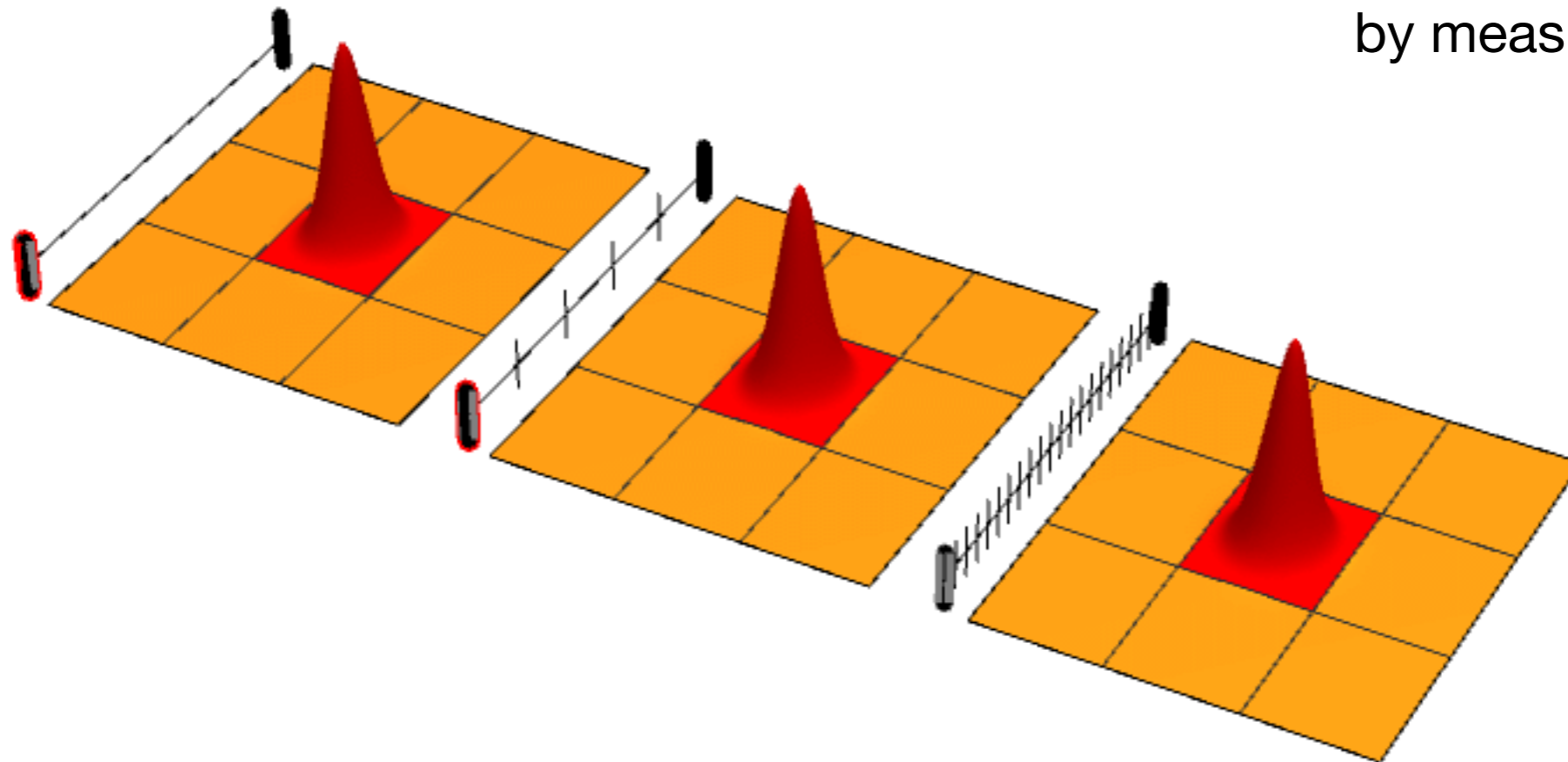
Slowing down the time-evolution
by measurements



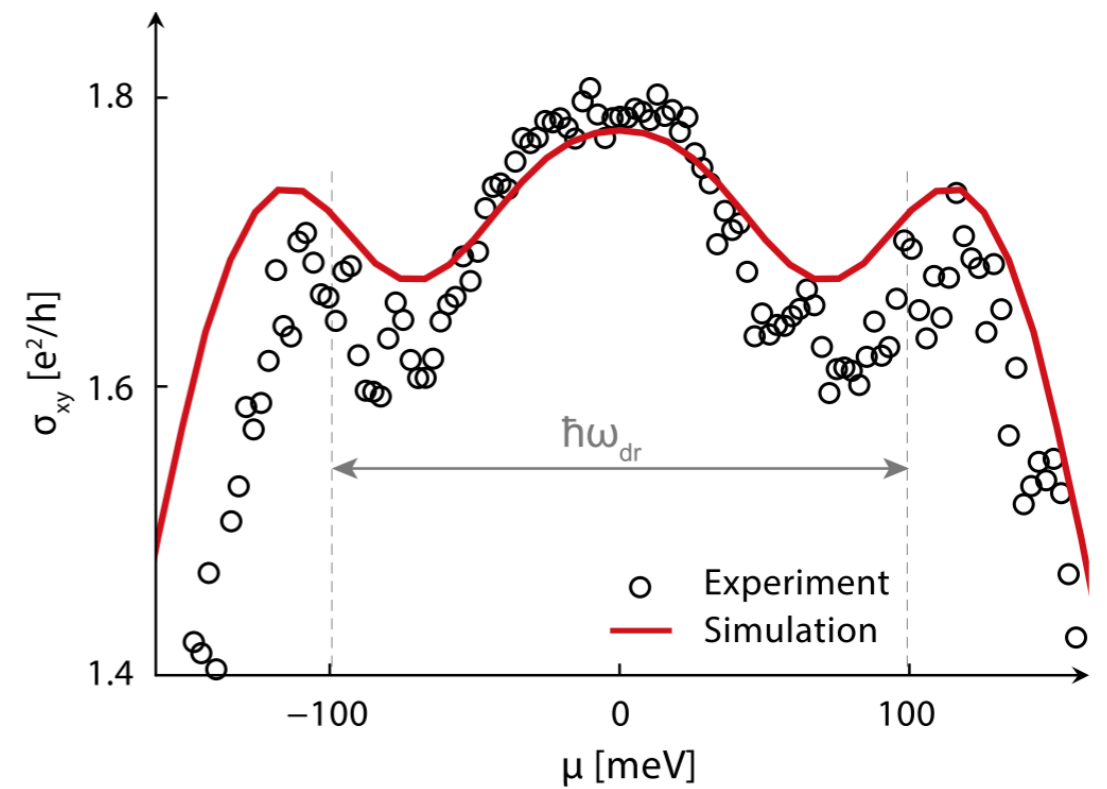
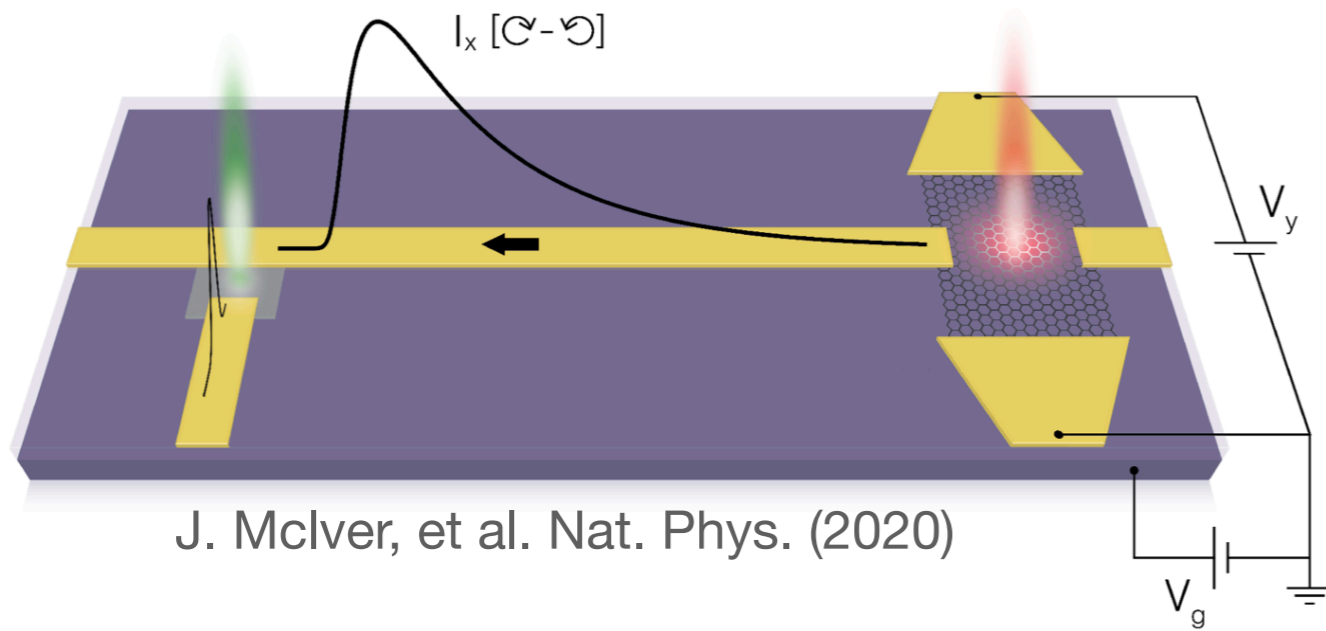
#Open quantum systems in transport experiments

Quantum Zeno effect

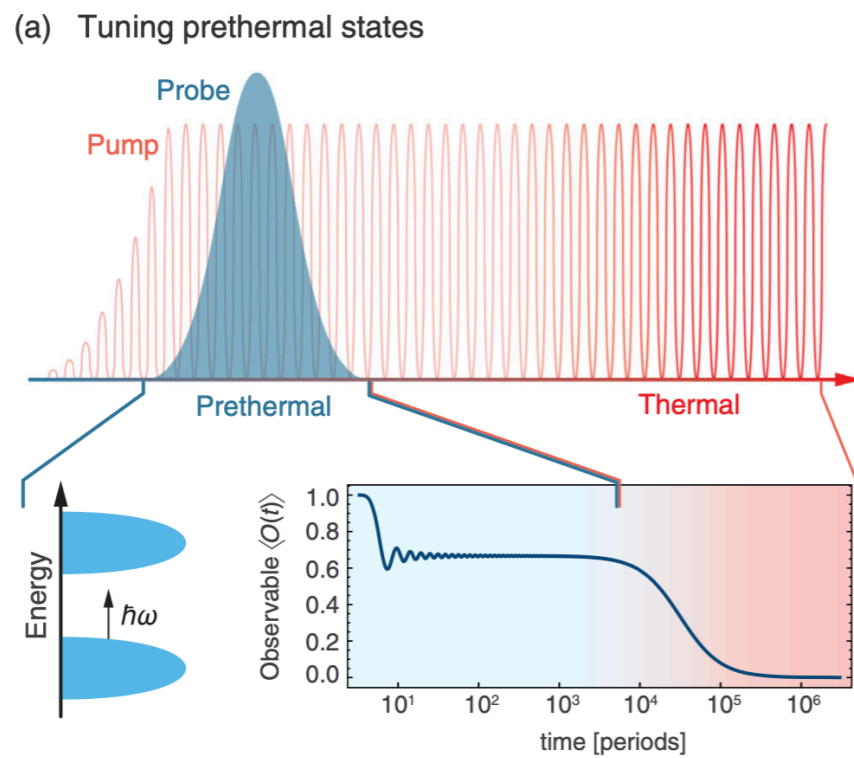
Slowing down the time-evolution
by measurements



#Open Quantum Systems in Transport Experiments



M. Nuske, et al. PRR (2020)



A. De la Torre, et al RMP (2021)

#Chiral anomaly in Non-Hermitian systems

Non-Hermitian effective model

$$H_{\text{eff}} \approx \pm [A_0^{5c} + M_i^{cj}(\vec{k} \mp \vec{A}^{5c}) \cdot \sigma]$$

Complex-valued

$$\mathcal{L} \propto \int \mathcal{D}\Psi \mathcal{D}\bar{\Psi} e^{-\int d^d x \left[\bar{\Psi} \gamma^\mu (M_\mu^{c\nu} \partial_\nu - i M_\mu^{c\nu} A_\nu^c + i \gamma_5 M_\mu^{c\nu} A_\nu^{5c}) \Psi \right]}$$

$$\mathcal{D}[A^c, A^{5c}] \neq \mathcal{D}^\dagger[A^c, A^{5c}]$$

Conserved classical currents

$$\partial_\mu j^\mu = 0 \quad \partial_\mu j_5^\mu = 0$$

Divergence of quantum currents:

$$[M_\mu^{c\nu} \partial_\nu j^\mu]_{\text{cov}} \propto \frac{1}{32\pi^2 \det[B]} \varepsilon^{\mu\nu\rho\lambda} \left(\tilde{F}_{\mu\nu}[A^{c\dagger}] \tilde{F}_{\rho\lambda}[A^{5c\dagger}] + \tilde{F}_{\mu\nu}^\dagger[A^c] \tilde{F}_{\rho\lambda}^\dagger[A^{5c}] \right)$$

$$[M_\mu^{c\nu} \partial_\nu j_5^\mu]_{\text{cov}} \propto \frac{1}{32\pi^2 \det[B]} \varepsilon^{\mu\nu\rho\lambda} \left(\tilde{F}_{\mu\nu}[A^{c\dagger}] \tilde{F}_{\rho\lambda}[A^{c\dagger}] + \tilde{F}_{\mu\nu}^\dagger[A^c] \tilde{F}_{\rho\lambda}^\dagger[A^c] + \tilde{F}_{\mu\nu}[A^{5c\dagger}] \tilde{F}_{\rho\lambda}[A^{5c\dagger}] + \tilde{F}_{\mu\nu}^\dagger[A^{5c}] \tilde{F}_{\rho\lambda}^\dagger[A^{5c}] \right)$$

Related to M^c

$$\tilde{F}_{\mu\nu}[V^\dagger] = M_\mu^{c\eta} M_\nu^{c\beta*} (\partial_\eta V_\beta^* - \partial_\beta V_\eta^*)$$

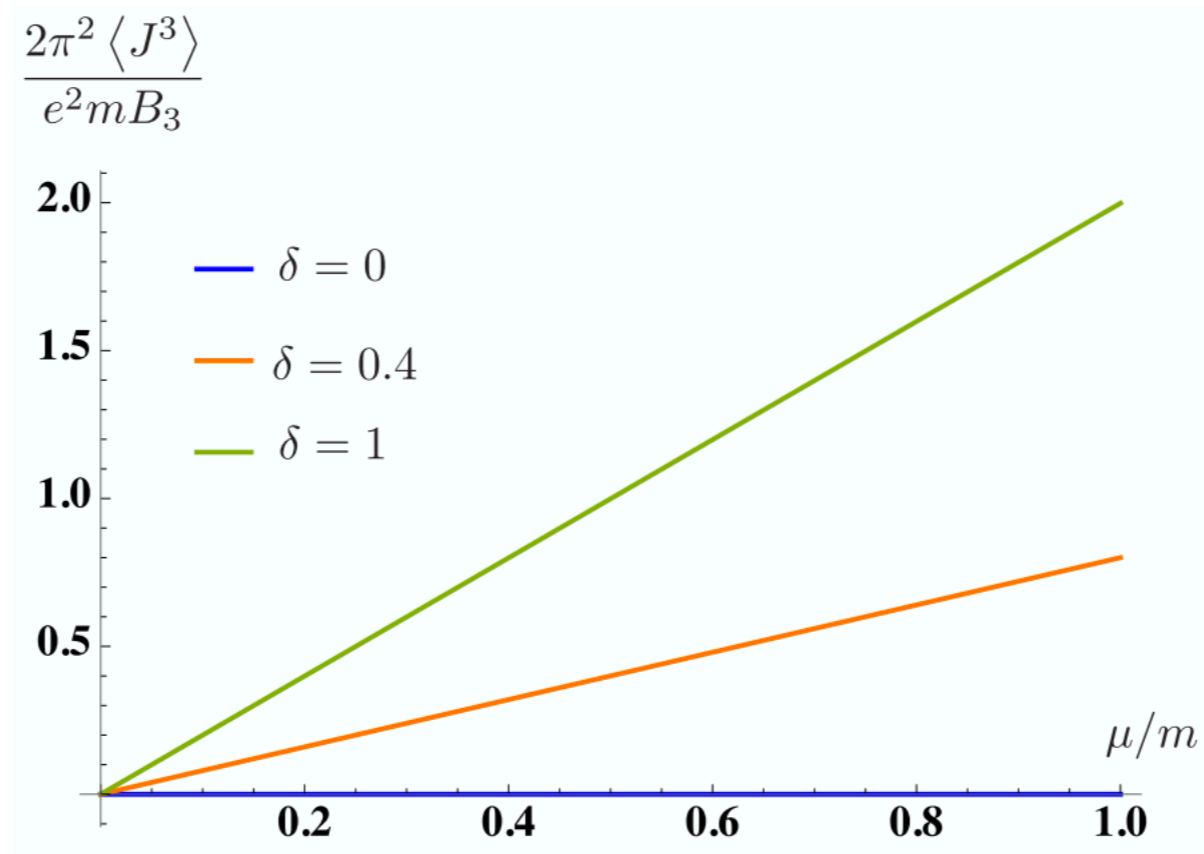
$$\tilde{F}_{\mu\nu}^\dagger[V] = M_\mu^{c\eta*} M_\nu^{c\beta} (\partial_\eta V_\beta - \partial_\beta V_\eta)$$

#Chiral anomaly in \mathcal{PT} -symmetric systems

$$H_{\mathcal{PT}} = \gamma_0 \boldsymbol{\gamma} \cdot \mathbf{k} + m\gamma_0 + m_5 \gamma_0 \gamma_5 - \mu \left(I - \frac{m_5}{m} \gamma_5 \right)$$

$$H_{\text{herm}} = S H_{\mathcal{PT}} S^{-1}$$

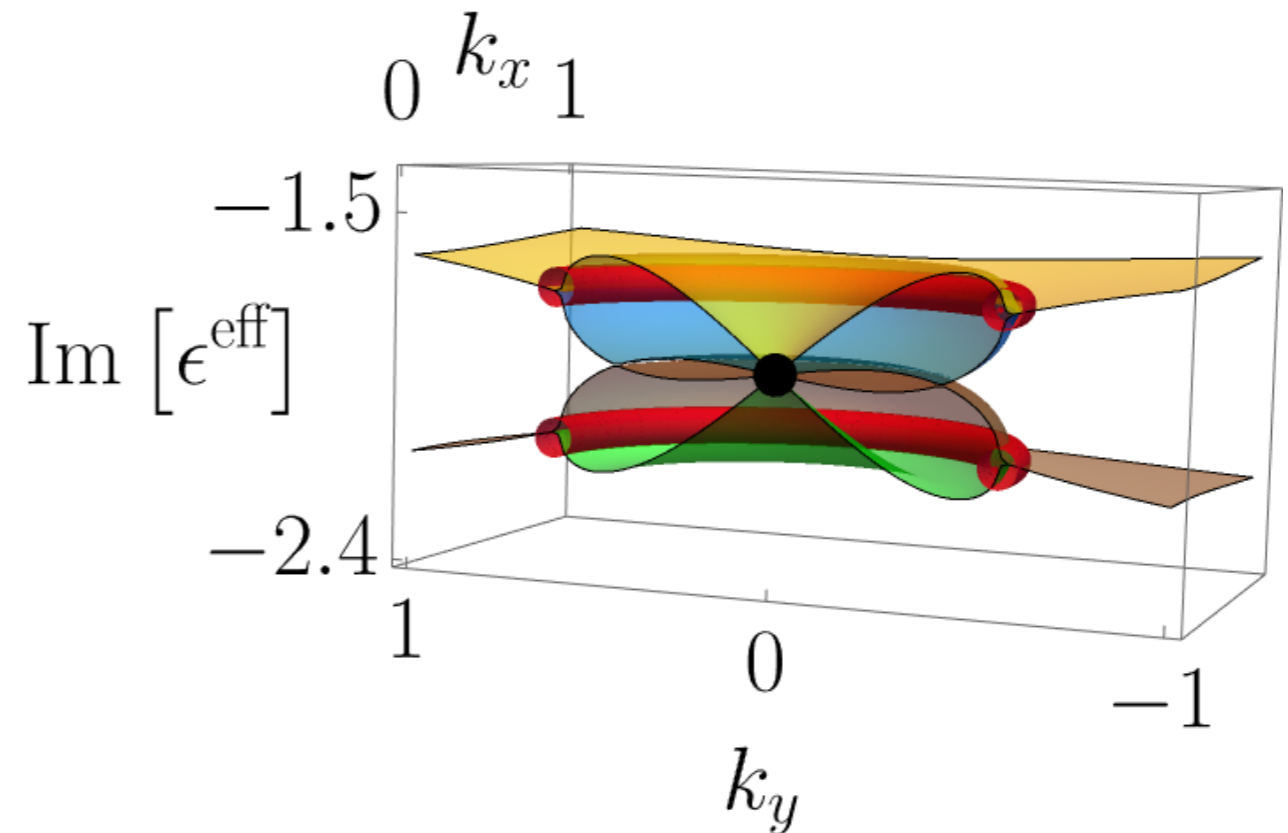
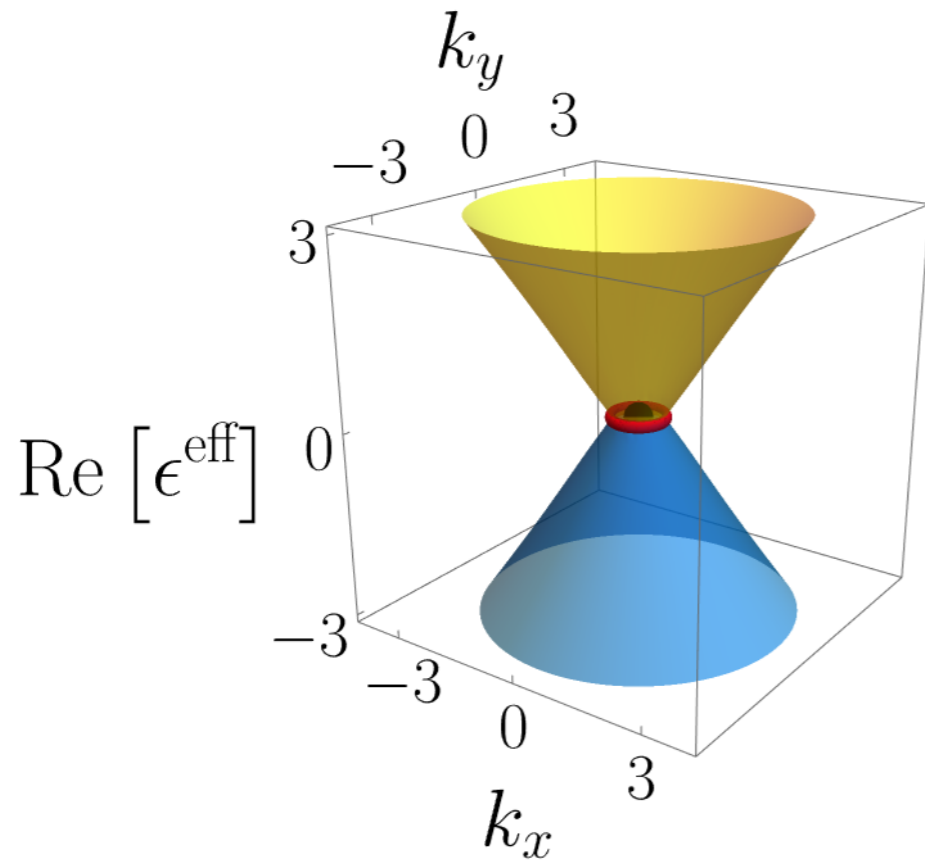
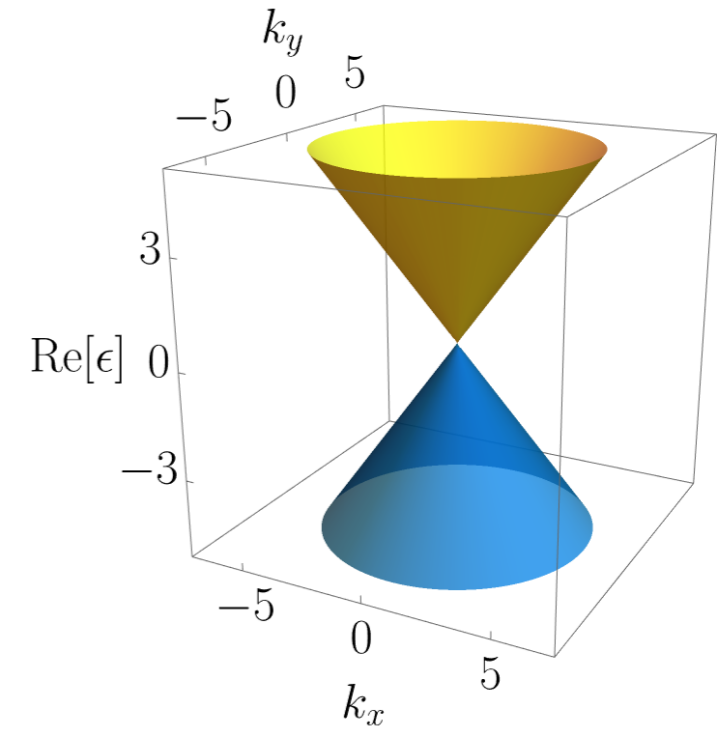
$$\langle J^3 \rangle = \frac{e^2 B_3}{2\pi^2} \frac{m_5}{m} \mu.$$

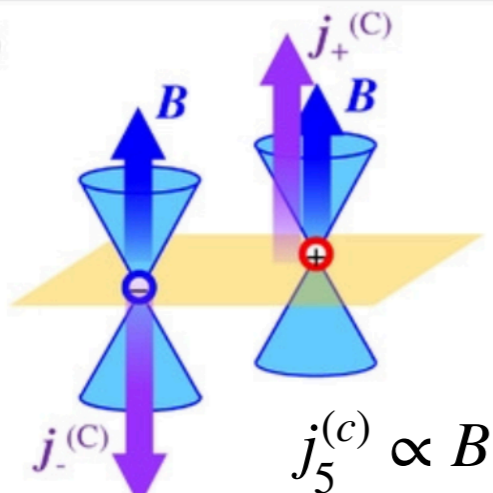
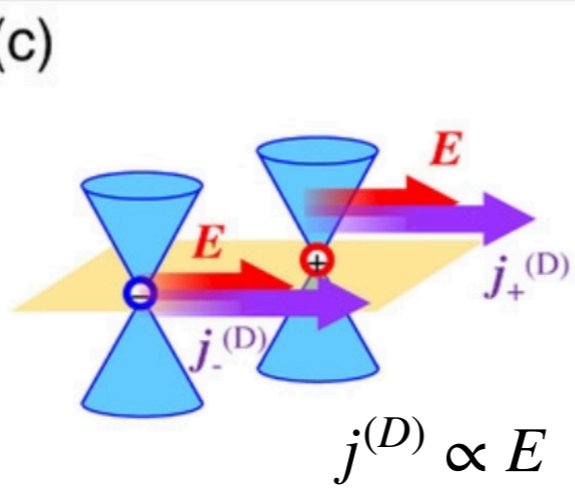
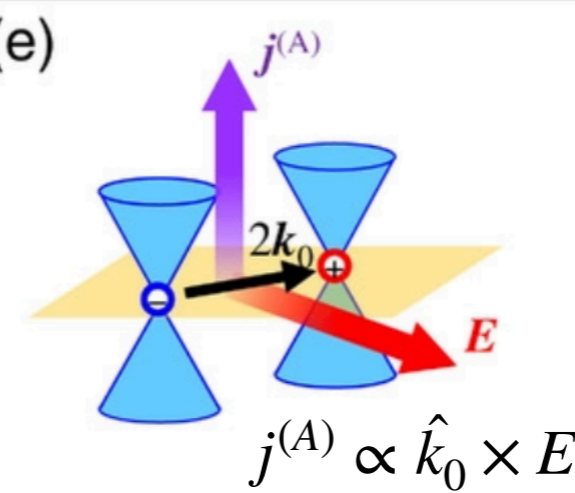
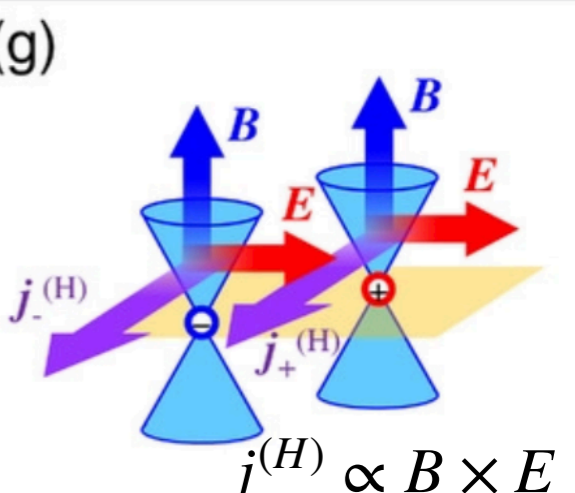
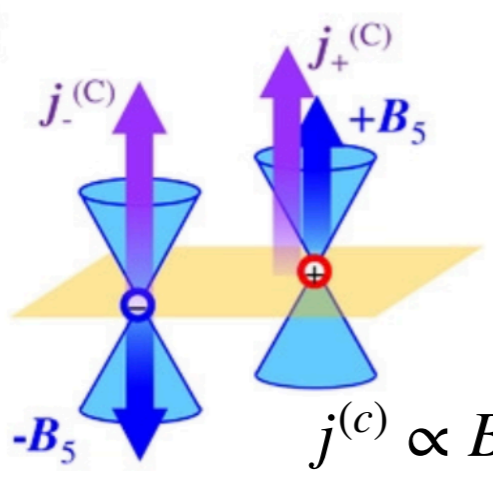
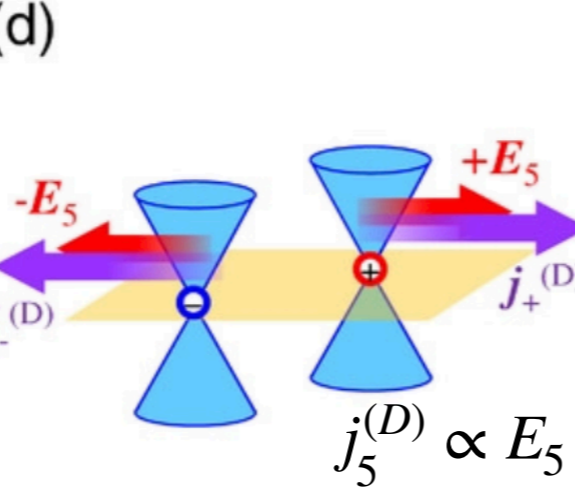
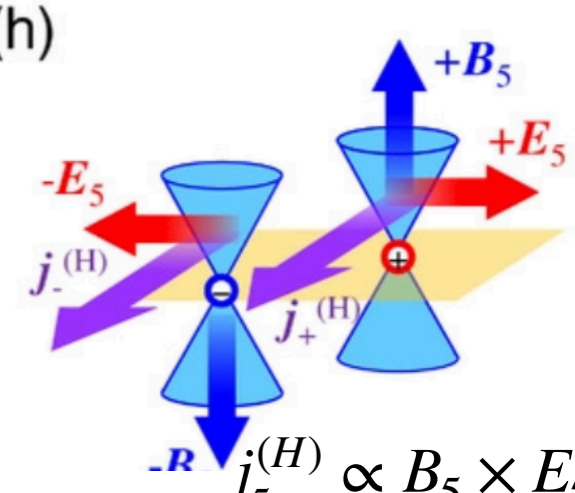


#Effective Description of Open Quantum Systems

$$\dot{\rho}_{\text{sys}} = -i(H_{\text{eff}}\rho_{\text{sys}} - \rho_{\text{sys}}H_{\text{eff}}^\dagger) + \sum_i 2L_i\rho_{\text{sys}}L_i^\dagger$$

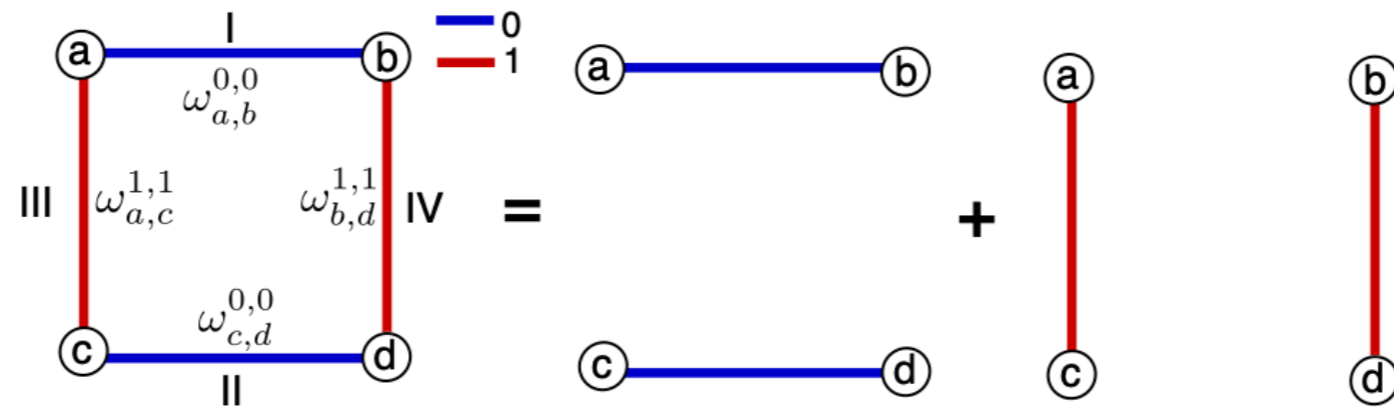
$$H_{\text{eff}} = H_{\text{sys}} - i \sum_l L_l^\dagger L_l$$



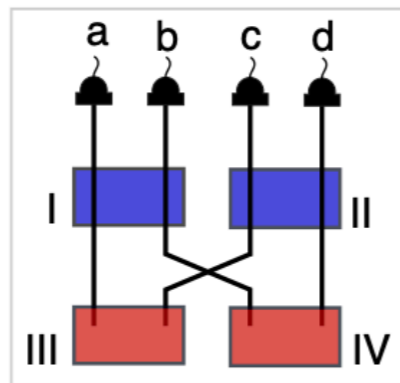
	Chiral magnetic effect	Drift effect	Anomalous Hall effect	Regular Hall effect
(E, B)	<p>(a)  $j_5^{(c)} \propto B$</p>	<p>(c)  $j^{(D)} \propto E$</p>	<p>(e)  $j^{(A)} \propto \hat{k}_0 \times E$</p>	<p>(g)  $j^{(H)} \propto B \times E$</p>
(E_5, B_5)	<p>(b)  $j^{(c)} \propto B_5$</p>	<p>(d)  $j_5^{(D)} \propto E_5$</p>	<p>(f) not applicable.</p>	<p>(h)  $j_5^{(H)} \propto B_5 \times E_5$</p>

Y. Araki, Annalen der Physik (2020)

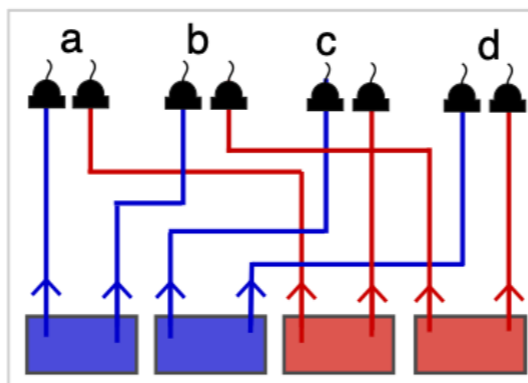
#Computer-discovery of Optical experiments



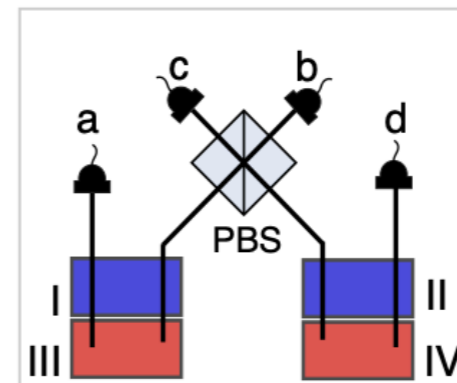
Setups for the Graph:



Entanglement by path identity



Integrated photonics i.e., path encoding



Bulk optics

$$|\psi^{(3)}\rangle = |0, -1, +1\rangle - |-1, 0, +1\rangle + |-1, +1, 0\rangle - |0, +1, -1\rangle + |+1, 0, -1\rangle - |+1, -1, 0\rangle, \quad (85)$$

$$|\psi^{(4)}\rangle = 2(|-1, 1, -1, 1\rangle + |1, -1, 1, -1\rangle) + |1, 0, -1, 0\rangle - |1, -1, 0, 0\rangle + |0, 1, 0, -1\rangle - |0, 1, -1, 0\rangle - |0, 0, 1, -1\rangle + |0, 0, 0, 0\rangle - |0, 0, -1, 1\rangle - |0, -1, 1, 0\rangle + |0, -1, 0, 1\rangle - |-1, 1, 0, 0\rangle - |1, 0, 0, -1\rangle + |-1, 0, 1, 0\rangle - |-1, 0, 0, 1\rangle. \quad (86)$$

