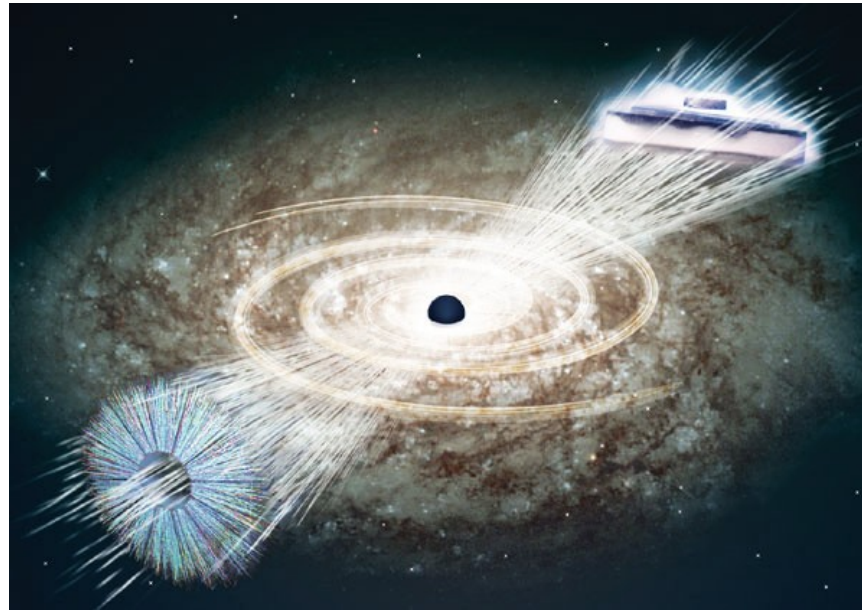


Pseudo-spontaneous symmetry breaking in hydrodynamics and holography



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Based on work with

Daniel Areán, Matteo Baggioli, Seán Gray and Sebastian Grieneringer

*Pseudo-spontaneous U(1) Symmetry Breaking
in Hydrodynamics and Holography*

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What is hydrodynamics?

- EFT of long-lived, long-wavelength excitations
- Equations of motion are conservation laws

$$\partial_{\mu} \langle T^{\mu\nu} \rangle = 0$$

- Recent interest in weakly broken symmetries:
 - pseudo-spontaneous breaking of spatial translations

see e.g. [Amoretti, Areán, Goutéraux, Musso, '16], [Ammon, Baggioli, Jiménez-Alba, '19], [Baggioli and Grienerger, '19], [Donos, Martin, Pantelidou, Ziogas '19]

potentially interesting for physics of strange metals and their charge-density fluctuations

[Delacretaz, Gouteraux, Hartnoll, Karlsson, '16], [Amoretti, Arian, Gouteraux, Musso, '18]

- pseudo-spontaneous breaking of U(1) superfluid *In this talk!*

[MA, Areán, Baggioli, Gray, Grienerger, '21], see also [Donos, Kailidis, Pantelidou, '21]
[Delacretaz, Gouteraux, Ziogas, '21], [Armas, Jain, Lier, '21]

Why is holography useful?

- Holography geometrises the hydrodynamic (derivative) expansion

Fluid/Gravity correspondence

- Thermodynamic and hydrodynamic (transport) coefficients can be determined

Are those non-zero?

- Universal relations may be detected

If you are lucky

Find a field-theoretic argument ...

PART I

Hydrodynamics for U(1) pseudo-spontaneous symmetry breaking

Hydrodynamics for a $U(1)$ superfluid

for simplicity: “probe limit”

- Goldstone boson φ with gauge invariant combination:

$$\xi_\mu \equiv \partial_\mu \varphi - A_\mu$$

- Equations: $\partial_\mu \langle J^\mu \rangle = 0$ and $u^\mu \xi_\mu = -\mu$

- Constitutive and Josephson relation to first order in derivatives:

$$\langle J^\mu \rangle = \rho_t u^\mu + \rho_s n^\mu - \sigma_0 T \Delta^{\mu\nu} \partial_\nu \left(\frac{\mu}{T} \right) + \mathcal{O}(\partial^2)$$

$$u^\mu \partial_\mu (\Delta_{\nu}^{\rho} \xi_\rho) = \Delta_{\nu}^{\rho} [-\partial_\rho \mu + \zeta_3 \partial_\rho \partial_\mu \rho_s n^\mu] + \mathcal{O}(\partial^2)$$

with superfluid bulk viscosity $\zeta_3 = \lim_{\omega \rightarrow 0} \omega \text{Im} G_{\varphi\varphi}^{\text{R}}(\omega, 0)$

and superfluid velocity defined by $\mu n^\mu \equiv \Delta^{\mu\nu} \xi_\nu$ $\Delta_{\mu\nu} = u_\mu u_\nu + \eta_{\mu\nu}$

Hydrodynamics for a $U(1)$ superfluid (II)

- Dispersion relation (of fourth sound)

$$\omega(k) = \pm v_s k - \frac{i}{2} \Gamma_s k^2 + \dots$$

where

$$v_s^2 = \frac{\rho_s}{\mu(\partial\rho_t/\partial\mu)} \quad \Gamma_s = \frac{\sigma_0}{(\partial\rho_t/\partial\mu)} + \zeta_3 \frac{\rho_s}{\mu}$$

Pseudo-spontaneous hydrodynamics

for simplicity: “probe limit”

- Goldstone boson acquires a mass, shift symmetry is broken

- Conservation law:
$$\partial_\mu \langle J^\mu \rangle = \Gamma u_\mu \langle J^\mu \rangle + m \varphi$$

with charge relaxation rate Γ and m related to mass of pseudo-Goldstone boson

Assumption: Γ and m are tiny

- Modified Josephson relation to first order in derivatives:

$$(u^\mu \partial_\mu + \Omega) (\Delta_\nu^\rho \xi_\rho) = \Delta_\nu^\rho [-\partial_\rho \mu + \zeta_3 \partial_\rho \partial_\mu \rho_s n^\mu] + \mathcal{O}(\partial^2)$$

with phase relaxation rate Ω

Pseudo-spontaneous hydrodynamics (II)

for simplicity: “probe limit”

- Dispersion relation $\omega(k) = \alpha_{\pm} - iD_{\pm}k^2 + \dots$

where
$$\alpha_{\pm} = -\frac{i}{2}(\Gamma + \Omega) \pm \sqrt{\omega_0^2 - \frac{(\Gamma - \Omega)^2}{4}} \quad \omega_0^2 \equiv \frac{m}{(\partial\rho_t/\partial\mu)}$$

$$D_{\pm} = \frac{1}{2} \frac{\sigma_0}{(\partial\rho_t/\partial\mu)} + \zeta_3 \frac{\rho_s}{\mu} \pm \frac{i}{2} \frac{\zeta_3 \rho_s (\partial\rho_t/\partial\mu) (\Gamma - \Omega) + 2\rho_s - \sigma_0 \mu (\Gamma - \Omega)}{\mu \sqrt{4m(\partial\rho_t/\partial\mu) - (\partial\rho_t/\partial\mu)^2 (\Gamma - \Omega)^2}}$$

- ‘sum relation’ $D_+ + D_- = \Gamma_s$

also holds when taking temperature and fluid velocity fluctuations into account.

Pseudo-spontaneous hydrodynamics (III)

for simplicity: “probe limit”, $\Gamma = 0$

- Static Correlators:

$$G_{\varphi\varphi}(0, k) = \frac{\chi_{\xi\xi}}{k^2 + m^2} \quad m^2 = \chi_{\xi\xi} m, \quad \chi_{\xi\xi} = \frac{\mu}{\rho_s}$$

‘Gell-Mann-Oakes-Renner’ relation

- Kubo-Formulae for Ω and ζ_3

$$\Omega = \lim_{\omega \rightarrow 0} \frac{m}{\omega \chi_{\rho\rho}^2} \text{Im} G_{J^t J^t}^{\text{R}}(\omega, 0)$$

$$\lim_{\omega \rightarrow 0} \frac{1}{\omega} \lim_{k \rightarrow 0} \frac{\partial^2}{\partial k^2} \text{Im} G_{J^t J^t}^{\text{R}}(\omega, k) = -2 \frac{\chi_{\rho\rho}^2}{m^2 \mu} (\zeta_3 m \rho_s - \rho_s \Omega - \mu \sigma_0 \Omega^2)$$

- AC-conductivity

$$\sigma(\omega) = \lim_{k \rightarrow 0} \frac{i}{\omega} [G_{J^x J^x}^{\text{R}}(\omega, k) - G_{J^x J^x}^{\text{R}}(0, k)] = \sigma_0.$$

PART II

Holographic models for pseudo-spontaneous symmetry breaking

Holographic model: s-wave superfluid

[Hartnoll, Herzog, Horowitz, '08]

- Action

$$S = \int d^4x \sqrt{-g} \left[-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - |D\psi|^2 - M^2 |\psi|^2 \right]$$

with background metric

$$ds^2 = \frac{1}{z^2} \left[-u(z) dt^2 - 2 dt dz + dx^2 + dy^2 \right] \quad u(z) = 1 - z^3$$

and $M^2 = -2$. Hence

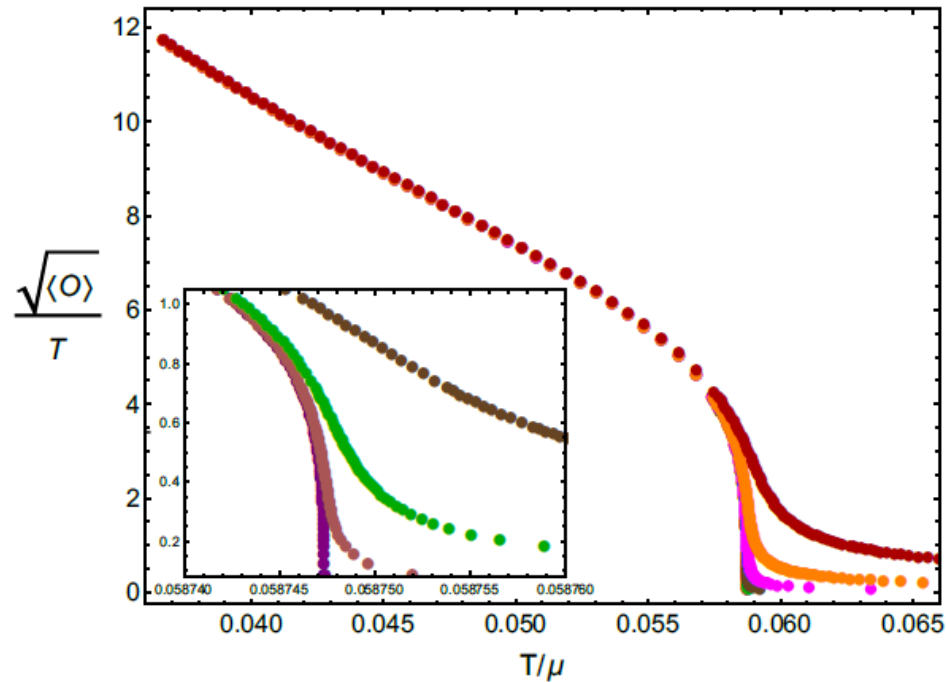
$$\psi(z) = \psi^{(l)} z + \psi^{(s)} z^2 + \mathcal{O}(z^3)$$

- **Model I:** tiny source for the scalar field $\psi^{(l)} = \lambda$

Identification of parameter: $m = q \lambda \langle \mathcal{O} \rangle_{\text{eq}} \quad \Gamma = 0$

Holographic model: s-wave superfluid

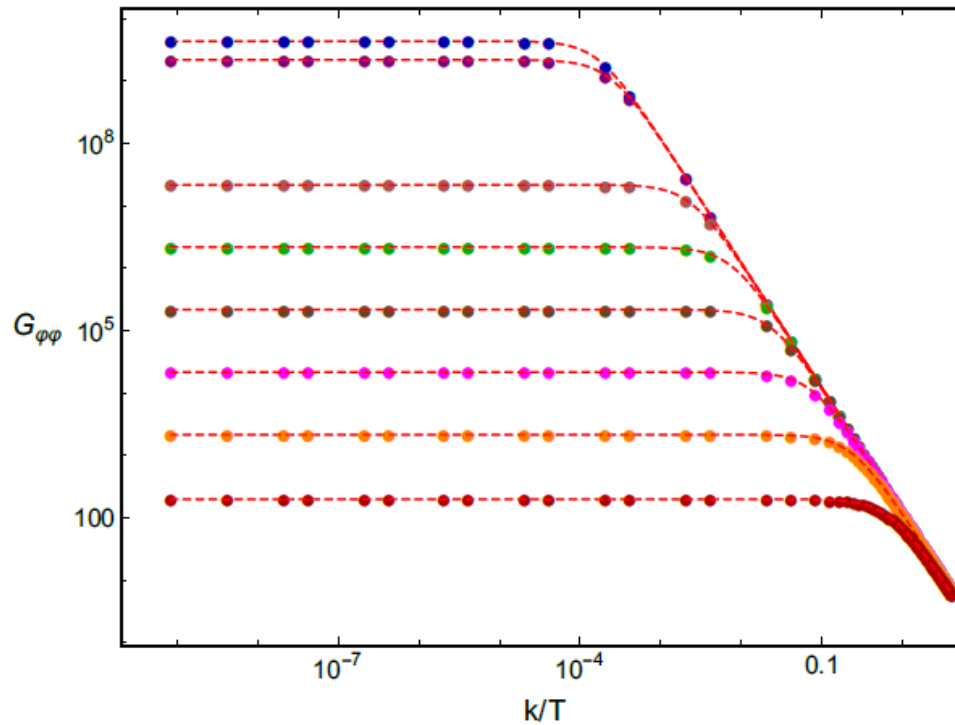
- Results:



Values for λ/T : blue $5 \cdot 10^{-10}$, purple 10^{-9} , light brown 10^{-7} , green 10^{-6} , dark brown 10^{-5} , magenta 10^{-4} , orange 10^{-3} , red 10^{-2} .

Holographic model: s-wave superfluid

- Results:



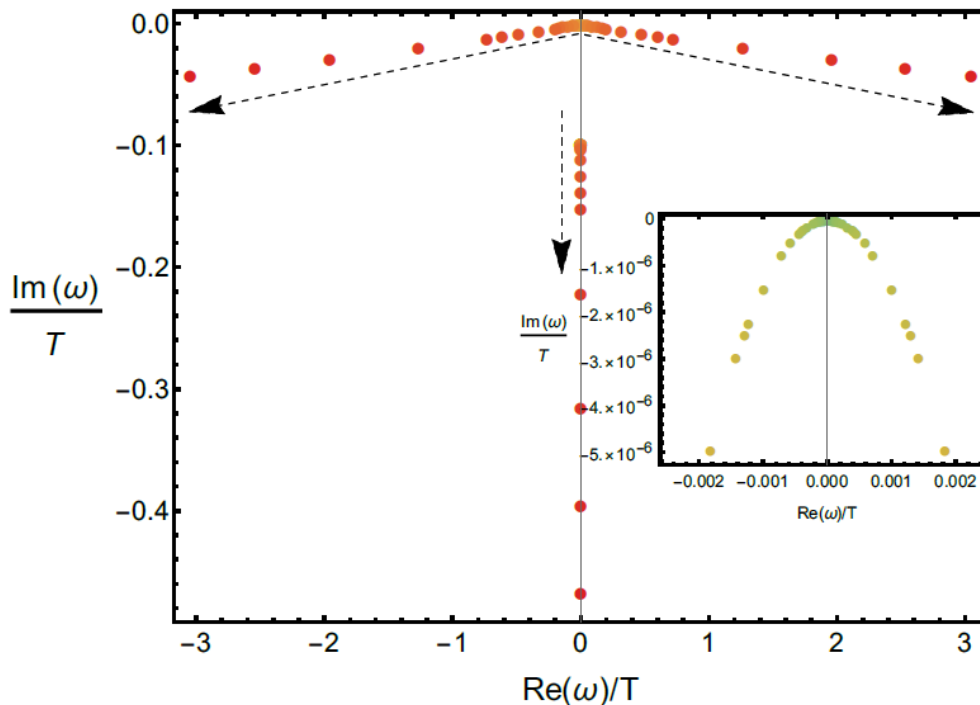
$$T/\mu = 0.0575$$

Values for λ/T : blue $5 \cdot 10^{-10}$, purple 10^{-9} , light brown 10^{-7} , green 10^{-6} , dark brown 10^{-5} , magenta 10^{-4} , orange 10^{-3} , red 10^{-2} .

Holographic model: s-wave superfluid

- Results: Quasi-normal Modes at zero momentum

$$\omega(k=0) = -\frac{i}{2}(\Gamma + \Omega) \pm \sqrt{\frac{m}{(\partial\rho_t/\partial\mu)} - \frac{(\Gamma - \Omega)^2}{4}}$$



Values: $\lambda/T \in [10^{-16}, 0.1]$

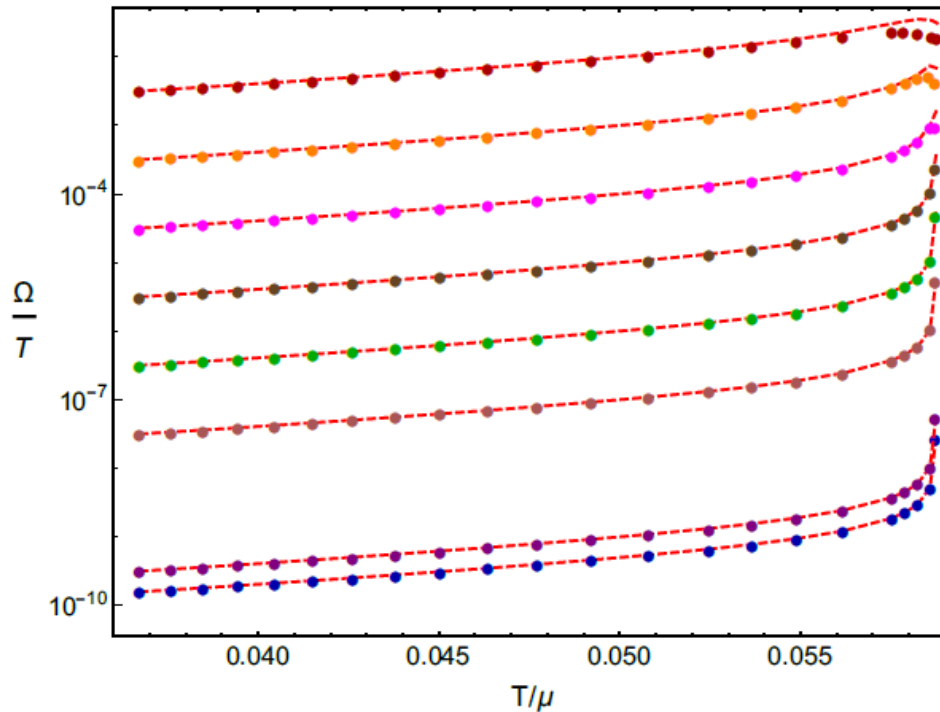
$T/\mu = 0.0582$

Conclusion:

$\Omega \neq 0$ but $\Gamma = 0$

Holographic model: s-wave superfluid

- Results:



Values for λ/T : blue $5 \cdot 10^{-10}$, purple 10^{-9} , light brown 10^{-7} , green 10^{-6} , dark brown 10^{-5} , magenta 10^{-4} , orange 10^{-3} , red 10^{-2} .

Holographic model: s-wave superfluid

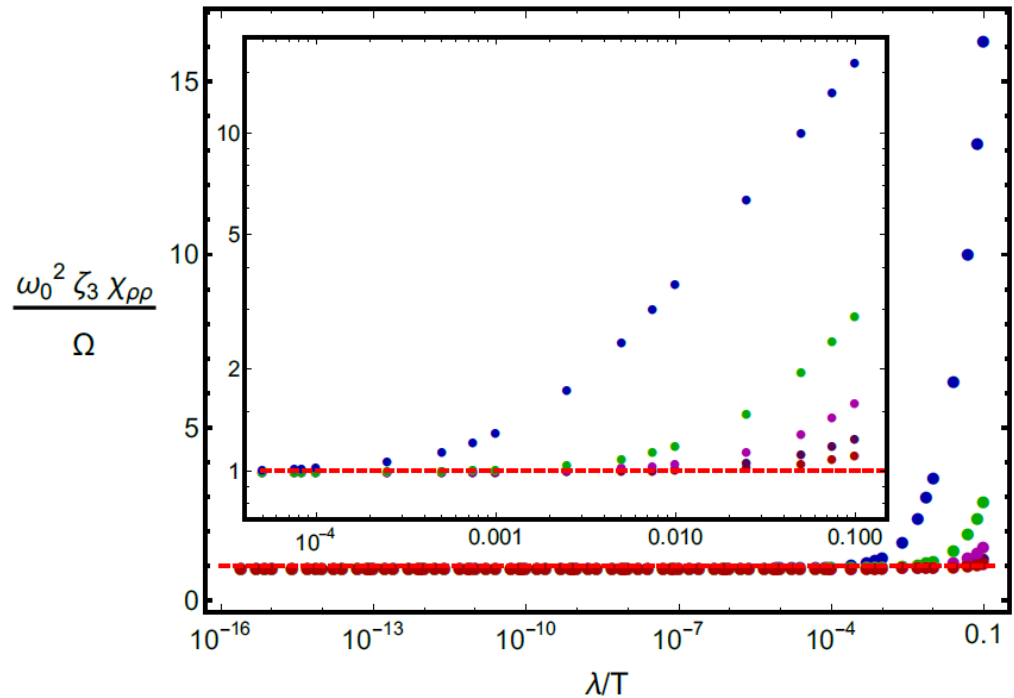
- Results: for small λ/T

$$\Omega = m \zeta_3$$

Universal relationship?

see also

[Delacretaz, Gouteraux, Ziogas, '21],
[Armas, Jain, Lier, '21]
[Grossi, Soloviev, Teaney, Yan, '20]



- Electrical conductivity:
finite DC conductivity despite no Drude-like behaviour

Holographic model: s-wave superfluid

- Interpretation of results:

pseudo-Goldstone bosons
as well as
(induced) effective phase relaxation

- We considered a *second model* based on massive gauge fields
 - induces effectively charge relaxation
 - breaks U(1) symmetry *but not* shift symmetry

Summary & Outlook

- ‘Hydrodynamics’ for pseudo-spontaneous breaking of U(1) symmetry
 - ‘sum relation’
- Two holographic models with distinct symmetry breaking behaviour
 - ‘universal relation’
- Outlook: how general is the ‘universal relation’?
 - type II (pseudo-)Goldstone modes?
 - light dilatons arising from the pseudo-spontaneous breaking of scale invariance

**Thank you very much
for your attention!**



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