

Transport of unpaired quark matter in neutron star densities

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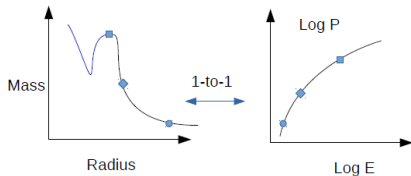
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Holographic perspectives on chiral transport

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- EoS (pressure \leftrightarrow energy density) determines most important properties of stars,

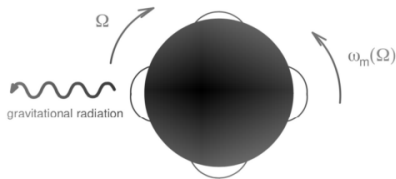


but alone is insufficient to address if \exists quark matter.

- No symmetry arguments (quark-hadron continuity).
- Sharp deconfinement (Occam's razor?) phase transition may lead to distinct signals
- Root for transport: quarks are relativistic unlike hadrons

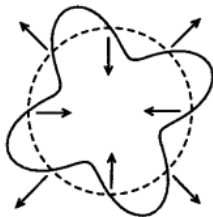
r-mode instability window

- Non-radial modes are driven unstable if the star rotates fast enough.
- These modes are strongly sheared and damped.
- ~~full~~ GR study, estimate:
[Andersson-Kokkotas gr-qc/0010102]



[Figure: Bratton et al.'22]

$$\frac{dE}{dt} = -2 \frac{E}{\tau}$$
$$\propto \exp(i\omega_m(\Omega)t + im\phi - t/\tau_m(\Omega))$$



- Maximum stable frequency for the star:

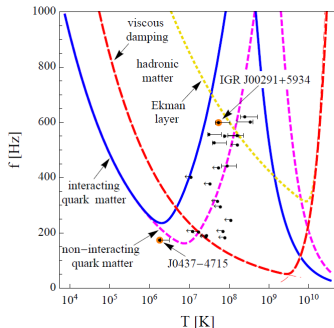
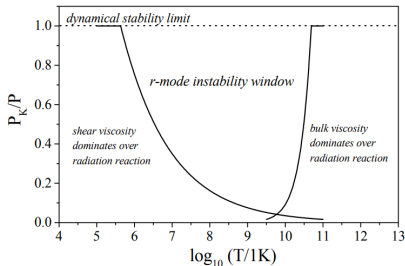
$$0 = 1/\tau_m(\Omega) = -1/\tau_{\text{GW}} + 1/\tau_{\eta} + 1/\tau_{\zeta}$$

[Figure: Lindblom'98]

r-mode instability window

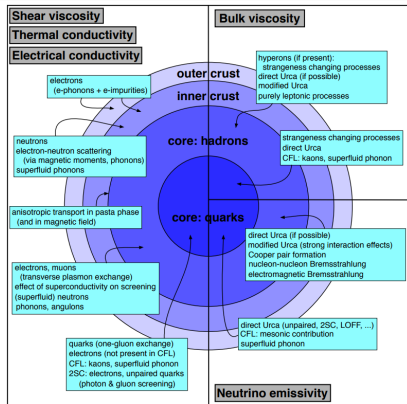
- Indirect evidence for quark matter?

[Figures: Andersson-Kokkotas gr-qc/0010102, Alford-Schwenzer 1310.3524]



- Bulk viscosity ζ seems interesting

Relevance of transport in stars



[Schmitt-Shternin 1711.06520]

- Damping of oscillations: shear η and bulk ζ viscosities
- Cooling: thermal conductivity κ , neutrino emissivities
- Magnetic fields: electric σ and thermal conductivities
- Mergers: viscosities, conductivities, evolution far from equilibrium

- Computing conductivities need care since spatial homogeneity
 - Lorentz invariance: only one hydro transport coeff σ
 - No force condition:

$$\rho_i E - s \nabla_i T = 0$$

- Conductivities

$$J^i = \sigma^{ij} E_j \quad , \quad \sigma^{ij} = \sigma \frac{\epsilon + p}{Ts} \delta^{ij}$$

$$Q^i = -\kappa^{ij} \nabla_j T \quad , \quad \kappa^{ij} = \frac{\mu}{T} \sigma \frac{\epsilon + p}{\rho} \delta^{ij}$$

- Quiescent stars

- Local charge neutrality: $\frac{2}{3}n_u - \frac{1}{3}n_d - \frac{1}{3}n_s - n_e = 0$
- Beta equilibrium: $\mu_s = \mu_d, \mu_u = \mu_d - \mu_e$

Holographic models encompassed

Action $S = S_{\text{glue}} + S_{\text{flavor}}$ w/ $\kappa_5^2 \sim 1/N_c^2, \mathcal{T}_b \sim N_f/N_c$:

$$S_{\text{glue}} = \frac{1}{2\kappa_5^2} \int d^5x \sqrt{-g} \left(R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi, \chi) \right)$$

$$S_{\text{flavor}} = -\frac{\mathcal{T}_b}{2\kappa_5^2} \int d^5x \mathcal{Z}(\phi, \chi) \sqrt{-\det(g_{\mu\nu} + \kappa(\phi, \chi) \partial_\mu \chi \partial_\nu \chi + \mathcal{W}(\phi, \chi) F_{\mu\nu})}$$

- metric $g_{\mu\nu} \leftrightarrow$ energy-momentum tensor
- flavor gauge field $A_\mu \leftrightarrow$ baryon charge
- “dilaton” $\phi \leftrightarrow$ gauge coupling
- “tachyon” \leftrightarrow quark masses

Our examples include:

- bottom-up V-QCD model
- top-down D3-D7 model in the quenched approximation

Transport for flavor independent masses

- Transport determined at the black hole horizon
[Hoyos-NJ-Järvinen-Subils-Tarrio-Vuorinen 2005.14205,2109.12122]

T (surface gravity) , s (area) , ρ (electric flux)

$$\phi_H = \phi(r_H) , \chi_H = \chi(r_H)$$

- Boundary values determine thermo: ϵ, p, μ
- Thermal and electric conductivities

$$\sigma = \frac{\mathcal{W}_H}{2\kappa_5^2 g_{xx}^H} \sqrt{(2\kappa_5^2 \rho)^2 + (g_{xx}^H)^3 \mathcal{T}_b^2 \mathcal{W}_H^2 \mathcal{Z}_H^2}$$

- Shear viscosity

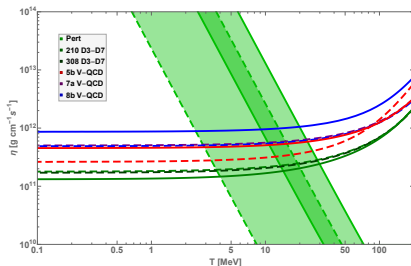
$$\eta = \frac{s}{4\pi} = \frac{(g_{xx}^H)^{3/2}}{2\kappa_5^2} + \frac{s_{\text{flavor}}}{4\pi}$$

- Bulk viscosity

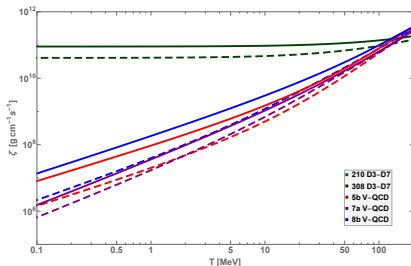
$$\frac{\zeta}{\eta} = (s\partial_s\phi_H + \rho\partial_\rho\phi_H)^2 + \frac{2\kappa_5^2}{(g_{xx}^H)^{1/2}} \frac{\kappa_H}{\mathcal{W}_H^2} (s\partial_s\chi_H + \rho\partial_\rho\chi_H)^2$$

Transport of cool quark matter

$\log \eta$ vs. $\log T$



$\log \zeta$ vs. $\log T$



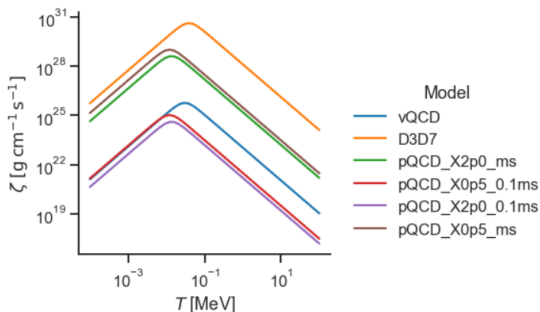
- Predictions for viscosities for unpaired quark matter (dashed $\mu = 450 \text{ MeV}$, solid $\mu = 600 \text{ MeV}$)
- Large deviation from perturbative QCD results
[Heiselberg-Pethick'93, Arnold-Dogan-Moore'06]
- Small results due to “idealized” case. Flavor-independent masses so get only QCD contributions, no weak interactions or electrons

Flavor dependence

- Bulk viscosity can be enhanced by resonant EW processes
[Alford-Mahmoodifar-Schwenzer '12]

$$u + d \leftrightarrow u + s$$

- Relevant for **unequal**, flavor-dependent, quark masses
[Gorda-Hoyos-NJ-Järvinen-Kurkela-Paatelainen-Rojas-Säppi-Vuorinen]



$$\zeta_{\text{D3-D7}} \Big|_{T=0} = \frac{4\lambda_1 \mu_d^6 (m_s^2 - m_d^2)}{\pi^4 \lambda_1^2 (-6\mu_d^2 + m_d^2 + m_s^2)^2 + \omega^2 (m_d^2 - 3\mu_d^2)^2 (m_s^2 - 3\mu_d^2)^2}$$

$$\lambda_1 = \frac{64 G_F^2 \sin^2(\theta_c) \cos^2(\theta_c)}{5\pi^3} \mu_d^5 T^2$$

- Observable effects in neutron star physics?
- Extensions:
 - neutrino emissivity
 - magnetic field
 - anisotropic equation of state
 - inhomogeneous phases
 - quark pairing (color “superconductivity”)

Grazie a tutti!