Tomography of light nuclei at an EIC

9-10 November 2022 ECT* (ITALY)

Deeply Virtual Compton Scattering off light nuclei: where do we stand?

Sara Fucini

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The EMC effect

The nuclear medium modifies the structure of bound nucleons



Collinear information led to many models but not yet to a complete explanation (e.g., see Cloët et al. JPG (2018), for a recent report)

Q2: 50 Gev2

Deeply Virtual Compton Scattering off nuclei

• Exclusive electro-production of a real photon \rightarrow clean access to Generalized Parton Distributions





- Two DVCS channels in nuclei:
- ▶ Coherent channel → GPDs of the whole nucleus
- ▶ Incoherent channel → GPDs of the bound nucleon



GPDs from lattice QCD

- Transversity GPDs of the proton from lattice QCD (Alexandrou at al., arXiv:2108.10789 (2021))
- Pion generalized parton distribution from lattice QCD (Chen et al., NPB, 952 (2020))

Deconvolution from CFFs

- The deconvolution problem of DVCS, see Bertone et al., PRD 103 (2021)
- Global fit (Guidal et al., Rept. Prog. Phys. 76 (2013)) and local fit (Dupré et al., PRD 95.1 (2017))

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All this concerns the free proton!

What about nuclei?

- Position-momentum structure of SRC for ⁴⁰Ca, ⁴⁸Ca, ¹²C (Cosyn et al. ,PLB 820 (2021))
- FSI in DIS off deuteron (Strikman et Weiss, PRC 57 (2018), transversity GPDs (Cosyn et al., PRD 98 (2018))
- NPLQCD collaboration (e.g., PRL 120 (2018)): nuclei with A<5 but unphysical q masses
- Phenomenological models for helium targets (ours or e.g. Liuti et al., PRC 72 (2005))

Incoherent DVCS off light nuclei

Incoherent DVCS off 4 He: S.F., S. Scopetta, M. Viviani, PRC(2021)-PRD(2021)



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In impulse approximation, for the cross section we get

$$d\sigma_{Incoh}^{\pm} = \int_{exp} dE d\vec{p} \frac{p \cdot k}{p_0 |\vec{k}|} P^{4He}(\vec{p}, E) \qquad d\sigma_b^{\pm}(\vec{p}, E, K)$$

the DVCS cross section off a bound proton_

For the **BSA**

$$\begin{split} A_{LU}^{Incoh}(K) &= \frac{d^4\sigma^+ - d^4\sigma^-}{d^4\sigma^+ + d^4\sigma^-} \approx \frac{\mathcal{I}^{4He}(K)}{T_{BH}^{24He}(K)} = \frac{\int_{\tilde{K}} dE \, d\vec{p} \, P^{4He}(\vec{p}, E) \, g(\vec{p}, E, K) \, \mathcal{I}(\vec{p}, E, K)}{\int_{\tilde{K}} dE \, d\vec{p} \, P^{4He}(\vec{p}, E) \, g(\vec{p}, E, K) T_{BH}^2(\vec{p}, E, K)} \\ \bullet \, \mathcal{I}(\vec{p}, E, K) \propto \Im \mathcal{M}\mathcal{H}(\xi', \Delta^2) = H(\xi', \xi', \Delta^2) - H(-\xi', \xi', \Delta^2), \\ \text{the nucleon GPD } H \text{ is evaluated for } \xi' = \frac{\mathbf{Q}^2}{(\mathbf{p} + \mathbf{p}')(\mathbf{q_1} + \mathbf{q_2})} \end{split}$$

The nuclear ingredient

$$\begin{split} P_{N}^{A}(\vec{p},E) &= \sum_{f_{A-1}} \langle {}^{4}He|f_{A-1}; N\vec{p}\rangle \langle f_{A-1}; N\vec{p}|^{4}He\rangle \delta(E-E_{min}-\epsilon_{A-1}^{*}) \\ \xrightarrow{f_{A-1}} & P_{0}(\vec{p},E) + P_{1}(\vec{p},E) \approx n_{0}(\vec{p})\delta(E-E_{min}) + n_{1}(\vec{p})\delta(E-\bar{E}) \\ \xrightarrow{p} & A_{1} \\ \xrightarrow{f_{He}} & A_{1} \\ \xrightarrow{f_{He}} & A_{1} \\ \xrightarrow{f_{He}} & A_{1} \\ \xrightarrow{f_{He}} & A_{2} \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\$$

- the total momentum distribution is $n(p) \propto \int d\vec{r_1} d\vec{r'_1} e^{i\vec{p}\cdot(\vec{r_1}-\vec{r'_1})} \rho_1(\vec{r_1},\vec{r'_1})$
- the ground momentum distribution is $n_0(|\vec{p}|) = |a_0(|\vec{p}|)|^2$ with

$$a_0(|\vec{p}|) \approx \langle \Phi_{^3He/^3H} | \Phi_{^4He} \rangle .$$

- the excited momentum distribution is $\mathbf{n_1}(|\vec{p}|) = n(|\vec{p}|) n_0(|\vec{p}|)$
- n(p), $n_0(p)$ can be evaluated within the Av18 NN interaction (Wiringa et al., PRC (1995)) + UIX 3-body forces (Pudliner et al., PRL (1995))
- $P_1^{\text{our model}}(\vec{p}, E) = N(p) P_{exc}^{\text{Ciofi's model (PRC(1996))}}(\vec{p}, E)$

MESSAGE TO TAKE HOME

- Realistic calculations for light nuclei $A \leq 6$
- Many body calculation accounting for mean field potential for heavier nuclei

Incoherent DVCS: results

Our results compared with the data from EG6 collaboration at JLab (PRL 123 (2019)).



Nuclear effects in A^{Incoh}: S.F., S. Scopetta, M. Viviani PRC(2021)

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What kind of nuclear effects are we describing? Let us consider the super ratio

$$A_{LU}^{Incoh}/A_{LU}^p = \frac{\mathcal{I}^{4}{}^{He}}{\mathcal{I}{}^{p}} \frac{T_{BH}^{2}}{T_{BH}^{2}} = \frac{R_{\mathcal{I}}}{R_{BH}} \propto \frac{(nucl.eff.)_{\mathcal{I}}}{(nucl.eff.)_{BH}} \,,$$



Is this behaviour due to a modification of the parton structure?

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Is this behaviour due to a modification of the parton structure?

- the ratio A_{LU}^{Incoh}/A_{LU}^p for "pointlike" protons
- the "EMC-like" trend

$$R_{EMC-like} = \frac{1}{\mathcal{N}} \frac{\int_{\tilde{K}} dE \, d\vec{p} \, P^{^{4}He}(\vec{p}, E) \, \Im m \, \mathcal{H}(\xi', \Delta^2)}{\Im m \, \mathcal{H}(\xi, \Delta^2)}$$



Incoherent DVCS off unpolarized deuteron

• The nuclear ingredient is easier than for ⁴He: just **momentum distribution** (totally realistic within AV18 potential!)

•
$$\Delta^2_{vertex} = (p_{final} - p_{inner})^2 \implies p_{final} \text{ fixed with } \Delta^2_{exp} \Big|_{\substack{p_{initial}}^{rest}} \dot{a} \text{ la CLAS}$$

•
$$\Delta^2_{vertex} = (p_{final} - p_{rest})^2 \implies p_{final}$$
 fixed with $\Delta^2_{exp} \Big|_{photons}$ à la HERMES

• Experimental data for pDVCS and nDVCS are coming out at JLab using a 12 GeV electron beam.

Analysis under review, see e.g.

https://indico.cern.ch/event/1104299/contributions/5055280/ attachments/2536704/4365938/EuNPC2022_ajh.pdf.

Within our model we can deliver

- Predictions for pDVCS
- Preliminary results for nDVCS

Stay tuned for the comparison with CLAS data!

nDVCS: preliminary results

$$\mathcal{I}(\vec{p}, E, K) \propto Im \left[F_1(\Delta^2) \mathcal{H}(\xi', \Delta^2) - F_2(\Delta^2) \mathcal{E}(\xi', \Delta^2) \left(\frac{\Delta^2}{4M^2} + \frac{\xi'(\Delta^2 - 2M^2 + 2p \cdot p'))}{4M^2} \right) \right]$$



Considering the DD formalism for the GPD E from GK EPJ (2008)

$$e_{val}(x) \propto B(\beta_{val})(1-x)^{\beta_{val}}$$

 $e_s(x) \propto N_s(1-x)^{\beta_s}$

In variant 1-6 $\beta_{val, s}$ and N_s are varied to have still a reasonabe fit to the Pauli FF.

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Incoherent on the deuteron: preliminary results



PDVCS



- Mild nuclear effects \implies scan in x_B for Δ^2 fixed and viceversa
- The contribution $\propto F_2 \mathcal{E}$ is crucial in nDVCS
- Better understanding of the flipped sign for pDVCS and nDVCS \implies insights on the the value of $J_{u,\,d}$
- Include possible FSI

Toward the tomography of ${}^{4}\mathrm{He}$ at the EIC

TOPEG: a Monte Carlo event generator for DVCS off light nuclei



x-section of coherent DVCS off ⁴He (S. F., S.Scopetta, M. Viviani, PRC 98 (2018))

$$\frac{d^4\sigma^{\lambda=\pm}}{dx_A dt dQ^2 d\phi} = \frac{\alpha^3 x_A y^2}{8\pi Q^4 \sqrt{1+\epsilon^2}} \frac{|T_{BH}|^2 + |T_{DVCS}|^2 + I^{\lambda}_{BH-DVCS}}{e^6}$$

$$T_{BH}^2 \propto F_A^2(t); \ T_{DVCS}^2 \propto \Im m \mathcal{H}^2 + \Re e \mathcal{H}^2; \ I_{BH-DVCS}^\lambda \propto F_A(t) \Im m \mathcal{H}$$
$$\mathcal{H}_q(\xi, t) = \int_0^1 dx \left(\frac{1}{x+\xi} + \frac{1}{x-\xi}\right) \left(\mathbf{H}_{\mathbf{q}}^{\mathbf{A}}(\mathbf{x}, \xi, \mathbf{t}) - \mathbf{H}_{\mathbf{q}}^{\mathbf{A}}(\mathbf{x}, -\xi, \mathbf{t})\right)$$

$$\mathbf{H}_{\mathbf{q}}^{\mathbf{A}}(\mathbf{x},\xi,\boldsymbol{\Delta^{2}}) \approx \sum_{N} \int \frac{dz}{z} \int dE d\vec{p} P_{N}^{A}(\vec{p},\vec{p}+\vec{\Delta},E) H_{q}^{N}\left(\frac{x}{z},\frac{\xi}{z},\Delta^{2}\right) \delta\left(z-\frac{\vec{p}^{+}}{\vec{p}^{+}}\right)$$
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Version 1.0 released:

JLab

- · Check for the events generated at the kinematics with 6 GeV electron beam
- Good also for CLAS 12 GeV

► EIC

- · We generated events for the three electron helium-4 beam energy configurations
 - (5x41) GeV
 - (10x110) GeV
 - (18x110) GeV
- These latter results are included in the EIC Yellow Report (Nucl.Phys.A 1026 (2022))
 - the NUCLEAR DVCS can be observed at the EIC
 - TOPEG is a flexible tool to do the GPDs phenomenology
 - Soon arriving the version 1.1

Toward the nuclear tomography

Promising results!!



Our assumptions in doing the fit of the pseudo-data generated with TOPEG

- using the leading order formalism
- 3 different minimum transverse momenta for the Roman pots
- 10 fb⁻¹ integrated luminosity

We conclude that

- · the error is highly correlated to the measurement threshold of the Roman Pots
- · the density profile extraction is anyway doable

► Incoherent DVCS off ⁴He and ²H

- New formalism for ⁴He and the deuteron (in progress)
- Insights for the generalized EMC effect
- Introduction of some final state interaction effects (TBD)
- Study of the A- **dependence** of the average BSA for light nuclei (see *Dupré's talk*)

► TOPEG

- For the **coherent DVCS** off ⁴He
 - Nuclear DVCS can be performed at the EIC: toward the 3D imaging of nuclei
 - Constrain the calculation of the nuclear spectral function with the help of tagged measuments
- TOPEG is a suitable phenomenological tool to study light nuclei at the EIC.

Backup slides

Incoherent channel

- · Nuclear part: momentum distribution (it is exact: instant form or light front)
- · Key study also for heavier nuclei

Coherent channel

- 9 quark GPDs
- Formalism already developed and established (see Cano, Pire EPJA (2004))
- there is a connection between the light-cone wave function of the deuteron (helicity amplitudes → GPDs) in terms of light-cone coordinates and the ordinary (instant-form) relativistic wave function that fulfills a Schrödinger type equation (we can update the potential)
- · we can compute

$$\chi(\vec{k};\mu_{1},\mu_{2}) = \sum_{L;m_{L};m_{S}} \langle \frac{1}{2} \frac{1}{2} 1 | \mu_{1},\mu_{2},m_{S} \rangle \langle L11 | m_{L}m_{S}\lambda \rangle Y_{L,M_{L}}(\hat{k}) u_{L}(k)$$

with AV18 and perform a Melosh rotation to relate the spin in the light-front with the spin in the instant-form frame of the dynamics

(18 x 110) GeV: analysis

Is it possible to study the region around the first diffraction minimum in the ^4He FF (t_{dif.min} = -0.48 GeV²)? YES, we can!

- 99%+ electrons and photons are in the acceptance of the detector matrix
- · This is true for all energy configurations

Electrons and photons appear in easily accessible kinematics according to the detector matrix requirements (exceptions for small angles photons)

- · Acceptance at low -t will be cut passing through the detectors
 - t_{min} is set by the detector features
 - t_{max} is fixed by the luminosity (billion of events to generate)

From left to right, the kinematical distributions of the final particles: electron, photon and ⁴He





From Hobart's talk

https://indico.cern.ch/event/1104299/contributions/5055280/ attachments/2536704/4365938/EuNPC2022_ajh.pdf



