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K⁻p correlation function and chiral SU(3) dynamics

"EXOTICO: EXOTIc atoms meet nuclear COllisions for a new frontier precision era in low-energy strangeness nuclear physics" @ ECT* Trento, 2022/10/18

High energy nuclear collision and FSI



Hadron-hadron correlation

$$C_{12}(k_1, k_2) = \frac{N_{12}(k_1, k_2)}{N_1(k_1)N_2(k_2)}$$

=
$$\begin{cases} 1 & (\text{w/o correlation}) \\ \text{Others (w/ correlation)} \end{cases}$$

High energy nuclear collision and FSI



Hadron-hadron correlation

• Koonin-Pratt formula : S.E. Koonin, PLB 70 (1977) S. Pratt et. al. PRC 42 (1990) $C(\mathbf{q}) \simeq \int d^3 \mathbf{r} \ S(\mathbf{r}) | \varphi^{(-)}(\mathbf{q}, \mathbf{r}) |^2_{\mathbf{q} = (m_2 \mathbf{k}_1 - m_1 \mathbf{k}_2)/(m_1 + m_2)}$ $S(\mathbf{r}) \quad : \text{Source function}$

 $\varphi^{(-)}(\mathbf{q},\mathbf{r})$: Relative wave function

High energy nuclear collision and FSI



• High energy nuclear collision and FSI A_2 Final State Interaction (FSI)

Hadronization

Hadron-hadron correlation

A

- Koonin-Pratt formula : $\underset{S.E. \text{ Koonin, PLB 70 (1977)}}{\text{S. Pratt et. al. PRC 42 (1990)}}$ $C(\mathbf{q}) \simeq \int d^3 \mathbf{r} S(\mathbf{r}) | \varphi^{(-)}(\mathbf{q}, \mathbf{r}) |^2_{\mathbf{q} = (m_2 \mathbf{k}_1 - m_1 \mathbf{k}_2)/(m_1 + m_2)}$ $S(\mathbf{r})$: Source function $\varphi^{(-)}(\mathbf{q}, \mathbf{r})$: Relative wave function
- Depends on ...

Interaction (strong and Coulomb)

mmm

quantum statistics (Fermion, boson)

• Simple model: Lednicky-Lyuboshits $\begin{pmatrix} 0 \\ LL \end{pmatrix}$ + Gaussian source with radius <math>R• $C(\mathbf{q}) \simeq \int d^3\mathbf{r} \ S(\mathbf{r}) | \varphi^{(-)}(\mathbf{q}, \mathbf{r}) |^2 = C(qR, R/a_0)$ • $\mathcal{F}(q) = [-1/a_0 - iq]^{-1}$ with scat. length a_0



Y. Kamiya, K. Sasaki, T. Fukui, K. Morita, K. Ogata, A. Ohnishi, T. Hatsuda, *Phys.Rev.C* 105 (2022) 1,014915

- Clear relation between C(q) and interaction
- Sensitive to (non)existence of bound state



Hadron correlation in high Simple model: Lednicky-Lyuboshits (LL) formula 0 2 4 $C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\varphi^{(-)}(\mathbf{q}, \mathbf{r})|^2 = C(qR, R/a_0)$ energy nuclear collisionUn-bound Unitary Bound<math>0 2 4 $Gaussian source R/a_0$ $F(q) = [-1/a_0 - iq]^{-1}$ with scat. length a_0



Y. Kamiya, K. Sasaki, T. Fukui, K. Morita, K. Ogata, A. Ohnishi, T. Hatsuda, *Phys.Rev.C* 105 (2022) 1,014915

- Clear relation between C(q) and interaction
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$\bar{K}N$ interaction and $\bar{K}p$ correlation

SIDDHARTA

 $\Lambda(1405)$

constraint on $a_0^{K^-p}$

• $\bar{K}(s\bar{l})N$ interaction and $\Lambda(1405)$

- Coupled-channel system of $\pi\Sigma$ - $\pi\Lambda$ - $\bar{K}N$
- Strong attraction reproducing quasi-bound state $\Lambda(1405)$
- Strong constraint on $a_0^{K^-p}$ by SIDDHARTA experiment of Kaonic hydrogen M. Bazzi, et al.. PLB 704 (2011)
- Structure of $\Lambda(1405)$
 - two pole structure J. A. Oller and U. G. Meißner, PLB500, 263 (2001)
 - *K̄N* molecular picture (high-mass pole) R.H. Dalitz, S.F. Tuan, PRL 425 (1959).
- Chiral SU(3) based $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$ potential
 - Constructed based on the amplitude with NLO chiral SU(3) dynamics $< -a_0^{K^-p}$, σ fitted Ikeda, Hyodo, Weise, NPA881 (2012)

 $\pi\Sigma$

• Coupled-channel, energy dependent as

 $V_{ij}^{\text{strong}}(r, E) = e^{-(b_i/2 + b_j/2)r^2} \sum_{\alpha=0}^{\alpha_{\text{max}}} K_{\alpha,ij} (E/100 \text{ MeV})^{\alpha}$

• Constructed to reproduce the chiral SU(3) amplitude around the $\overline{K}N$ sub-threshold region

 $\sigma_{K^-p\to K^-p}$

 $\sigma_{K^-p \to \bar{K}^0 n}$

 $\bar{K}^0 n$

 K^-p

Miyahara, Hyodo, Weise, PRC 98 (2018)

Re \sqrt{s}

 K^-p correlation

Coupled-channel effect

Koonin-Pratt-Lednicky-Lyuboshits-Lyuboshits (KPLLL) formula

 $C(\mathbf{q}) = \int d^3 \mathbf{r} \, S(\mathbf{r}) \, |\psi^{(-)}(q;r)|^2 + \sum_{j \neq i} \omega_j \int d^3 \mathbf{r} \, S_j(\mathbf{r}) \, |\psi_j^{(-)}(q;r)|^2$

S.E. Koonin, PLB 70 (1977)S. Pratt et. al. PRC 42 (1990)R. Lednicky, et.al. Phys. At. Nucl. 61(1998)

Contribution from Coupledchannel Source

- Coupled-channel wave function $\psi_i \rightarrow (\text{out-going wave}) + S^{\dagger} (\text{incoming wave})$
 - $|S_{ij}| < 1$ —> Decrease the correlation
 - At channel threshold —> Cusp structure
 - ψ_i : obtained by solving the c.c. Schrödinger eq.

$$\begin{array}{cccc} -\frac{\nabla^2}{2\mu_1} + V_{11}(r) & V_{12}(r) & \cdots & V_{1n}(r) \\ V_{21}(r) & -\frac{\nabla^2}{2\mu_2} + V_{22}(r) + \Delta_2 & \cdots & V_{2n}(r) \\ \vdots & \vdots & \ddots & \vdots \\ V_{n1}(r) & V_{n2}(r) & \cdots & -\frac{\nabla^2}{2\mu_n} + V_{nn}(r) + \Delta_n \end{array} \right)$$

 $V_{ij} = V_{ij}^{\text{strong}} (+V^{\text{Coulomb}}) \quad \Delta_i$; threshold energy diff.

• Contribution from coupled-channel source

$$K^-p, \bar{K}^0n, \pi^0\Sigma^0, \pi^+\Sigma^-, \pi^-\Sigma^+, \pi^0\Lambda$$

$$\begin{array}{c} FSI \\ FSI \\ p \end{array} \begin{array}{c} K^{-} \\ p \end{array} \begin{array}{c} C_{K^{-}p} \end{array}$$

 $\Psi(q_1, r) = E\Psi(q_1, r), \bullet$ Enhance C(q)

- Enhance cusp structure
- ω_i : production rate (compared to measured channel)

Coupled-channel effect

Coupled-channel source effect



- Strong source size dependence < == Due to the near-threshold $\Lambda(1405)$ pole
- Coupled—channel effect
 - Enhance the correlation
 - Enhance the cusp structure
 - $\pi\Sigma$ and $\bar{K}^0 n$ w.f. components are significant

Kamiya, Hyodo, Morita, Ohnishi, Weise, PRL 124 (2020) 13, 132501

Coupled-channel effect

Source size dependence of coupled-channel effect



- Strong source size dependence < == Due to the near-threshold $\Lambda(1405)$ pole
- C(q) with large source
 - Less prominent cusp structure
 - Weaker coupled-channel source contribution



Source size dependence

Coupled-channel wave function

 $C_{K^{-}p}(\mathbf{q}) = \int d^3 \mathbf{r} \ S_{K^{-}p}(\mathbf{r}) \left| \psi_{K^{-}p}^{C,(-)}(q;r) \right|^2 + \sum_{i \neq i} \omega_j \int d^3 \mathbf{r} \ S_j(\mathbf{r}) \left| \psi_j^{C,(-)}(q;r) \right|^2$



Coupled-channel wave function $\bar{K}^0 n$, $\pi^0 \Sigma^0$, ...

• Coupled-channel wave function satisfies the out-going boundary condition

• Measured channel (*K*⁻*p*) has out going wave

• Coupled-channel w.f. emerges only in int. region

- Small source ==> W.F. of Coupled-channels counts
- Large source ==> Measured channel contribution dominant

Comparison with ALICE data

Comparison to the Exp. data



- Strong Coulomb enhancement at small q
- $\Lambda(1520)$ peak
- Comparison with Chiral model 0
 - data well reproduced with the reasonable values of parameter
 - Sizable $\pi\Sigma$ source contribution



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$\bar{K}N$ interaction and $\bar{K}p$ correlation

- Source size dependence of K^-p data
 - ALICE data PbPb collisions data ALICE PLB 822 (2021) 136708
 - Large source —> weaker coupled-channel effect
 - —> more direct approach to interaction of the measured channel
 - Extraction of the K^-p scattering length from correlation function * Fitting with 1 channel LL model with Gaussian source





Latest K^-p correlation results

3.5) C(k*)

3.0

2.5

2.0F

1.5

1.0

n_{o stat}

ALICE [2205.10258

- *p*Pb : 0-20%, 20-40% 40-100%
- PbPb : 60-70%, 70-80% 80-90%
- Discrepancy around $\bar{K}^0 n$ threshold between chiral SU(3) model and exp. data for small source data

ALICE *p*Pb

0-20%

150

100

 \mathbf{U}

²/NDF = 11.13

200

50

250

100

150

k* (MeV/c)



50

KN interaction and K^-p correlation

Latest K^-p correlation results

ALICE [2205.10258]

- *p*Pb : 0-20%, 20-40% 40-100%
- PbPb : 60-70%, 70-80% 80-90%
- Discrepancy around $\bar{K}^0 n$ threshold between chiral SU(3) model and exp. data for small source data



Analysis with the effective Kyoto potential Vary the interaction with the factor β as

 $V_{\bar{K}N,I=1} \to \beta_{I=0} V_{\bar{K}N,I=1}$

Check the consistency buits NR0 SIDD 02 RT eVan 20-40

6

5

150

 $0.7 < S_{T}$

200

= 2.2

 $\underbrace{\overset{*}{\underset{N}{\overset{*}{\underset{N}}}}^{3.5}}_{\bar{K}N,I=0} \xrightarrow{\gamma}_{A} \overset{\beta}{\underset{K}{\overset{K}{\underset{N}}}^{I=0}} \overset{K}{\underset{N}{\overset{K}{\underset{N}}}^{I=1}} \overset{K}{\underset{N}{\underset{N}}} \overset{K}{\underset{N}{\underset{N}}} \overset{K}{\underset{N}{\underset{N}}} \overset{K}{\underset{N}} \overset{K}{\underset{N}} \overset{K}{\underset{N}} \overset{K}{\underset{N}{\underset{N}}} \overset{K}{\underset{N}} \overset{K}{\underset{N}} \overset{K}{\underset{N}} \overset{K}{\underset{N}} \overset{K}{\underset{N}} \overset{K}{\underset{N}} \overset{K}{\underset{N}} \overset{K}{\underset{N}{\underset{N}}} \overset{K}{\underset{N}} \overset$

Further constraint on *KN* interaction?



Cieply and Mai, EPJ Web Conf. 130, 02001 (2016)

• Can we constrain $\overline{K}NI = 1$ interaction / amplitude from femtoscopy?

$\bar{K}N$ interaction from $K_S^0 p$ correlation 0.7550 100



- Chiral amplitude 1600 1650 1700 1750 Ikeda, Hyodo, Weise, NPA881 (2012)
- Effective potential
- Miyahara, Hyodo, Weise, PRC 98 (2018)



- Well determined with scat. exp.
- Chiral amplitude
 - K. Aoki and D. Jido, PTEP (2019)
- Effective potential
 - Constructed from chiral amp.



Y. Kamiya, T. Hyodo, A. Ohnishi. in preparation

250

200

 $q \, [\text{MeV}/c]$

 $\overline{300}$

 $C_{K_{s}^{0}p} = [C_{\bar{K}^{0}p} + C_{K^{0}p}]/2$



Enhancement by $\overline{K}^0 p(\overline{K}N I = 1)$ is sizable.



Summary

- Femtoscopic correlation function in high energy nuclear collisions is a powerful tool to investigate the hadron-hadron interaction.
- Chiral SU(3) based effective potential model reproduces the ALICE K^-p correlation data from pp collisions with the reasonable by including the coupled-channel source effect.
- Detailed source size dependence can be investigated with the latest data from the different collision experiments.
 - Large source: Good agreement with the chiral model of K^-p channel
 - Small source: Finite deviation indicates the need for the modification of the coupled-channel interaction.
- $K_s^0 p$ correlation is useful to directly see the $I = 1 \ \overline{K}N$ interaction.

Thank you for your attention!



Thank you!

