

# Critical reassessment of recent developments in functional theory: From Hartree-Fock to i-DMFT

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arXiv:2202.05532

# What is i-DMFT

J. Wang, E. Baerends, PRL, 128, 013001 (2022)

- 1RDM ( $\gamma = (\mathbf{n}, \{\phi_i\})$ ) functional theory for  $H(h) = h + W$

$$E(\gamma) = \text{Tr}[h\gamma] + \mathcal{F}_W(\gamma) \leftarrow \text{universal functional}$$

D. Collins, Z. Naturforsch. Teil A 48, 68 (1993)

- Premise - **modified Collin's conjecture**

$$E_{\text{cum}} = \mathcal{F}_W^{\text{HF}} - \mathcal{F}_W \approx \kappa S^{\text{ph}}(\mathbf{n}) + b, \text{ with}$$

$$S^{\text{ph}}(\mathbf{n}) = - \sum_i n_i \log(n_i) + (1 - n_i) \log(1 - n_i)$$

**Particle Correlation**



**Correlation Energy**

- Optimal  $\gamma$  is obtained self-consistently, with  $(\kappa, b)$  pre-fit by two FCI solutions

# Capability of *i*-DMFT

J. Wang, E. Baerends, PRL, 128, 013001 (2022)

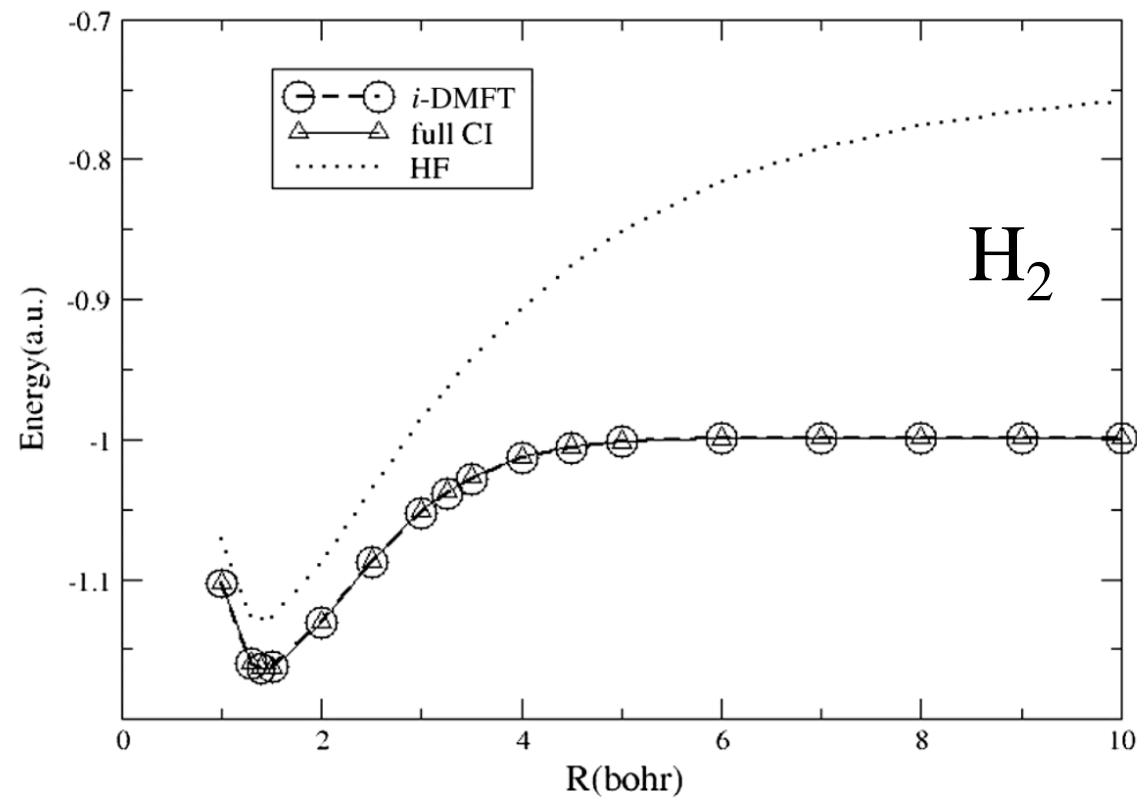


FIG. 1. Total energy curves for HF, full CI, and *i*-DMFT calculations along the dissociation coordinate of H<sub>2</sub>. Basis set cc-pVDZ,  $\kappa = 0.094681$ , and  $b = 0.0286186$ .

Accurate ground state energy for all geometry (with two data points as input)

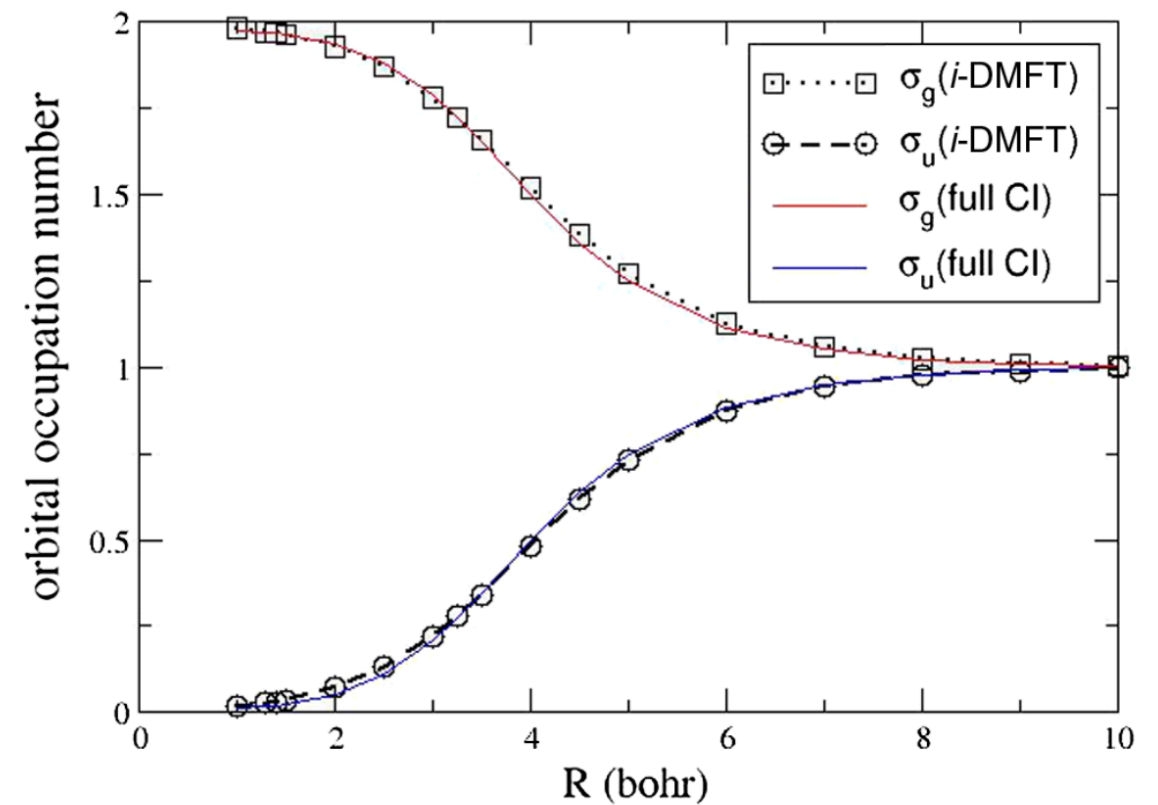


FIG. 2. Occupation numbers of the  $\sigma_g$  and  $\sigma_u$  orbitals in H<sub>2</sub> as a function of the nuclear separation  $R$ . The data are calculated with the basis set cc-pVDZ,  $\kappa = 0.094681$ , and  $b = 0.0286186$ .

Recovering static correlation and NONs at dissociation

# How useful/good/accurate is the modified Collin's Conjecture?

$$E_{\text{cum}} \approx \kappa S^{\text{ph}}(\mathbf{n}) + b$$

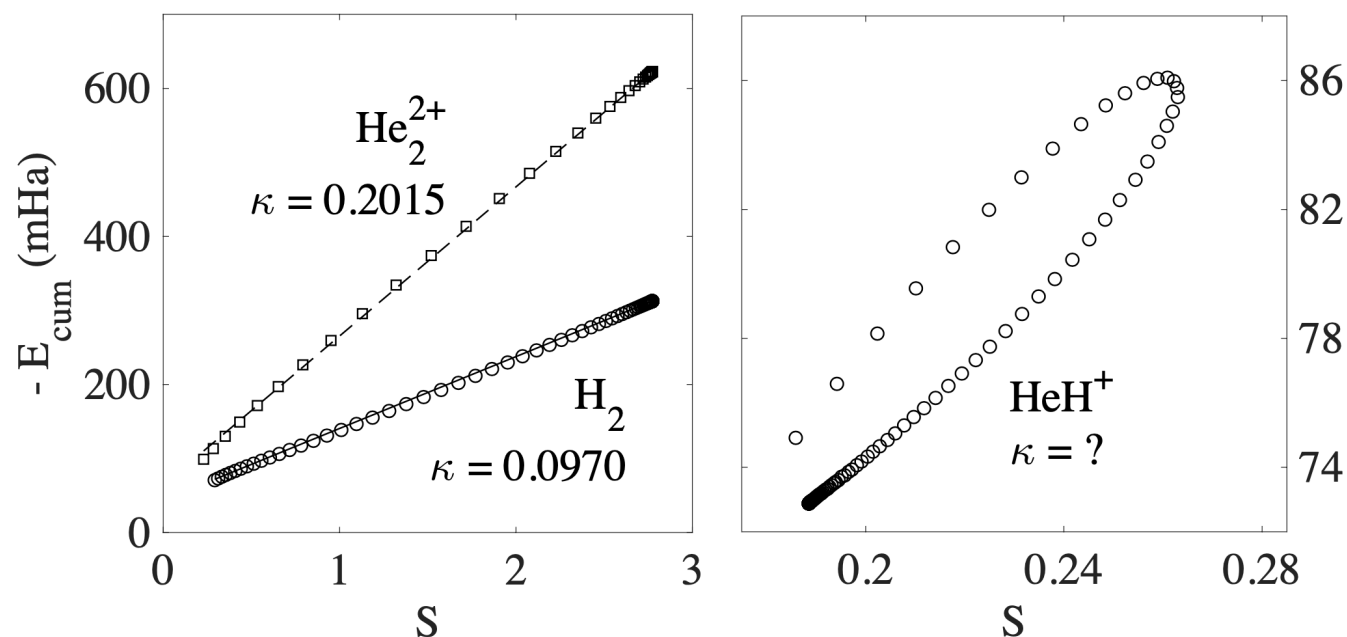


FIG. 1. Cumulant energy  $E_{\text{cum}}(R)$  versus particle-hole symmetric entropy  $S(\mathbf{n}(R))$  for the ground states of  $\text{H}_2$ ,  $\text{He}_2^{2+}$  (left) and  $\text{HeH}^+$  (right) for different internuclear separations  $R$ . The basis set cc-pVDZ was used.

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- For fixed  $N_e$  and nuclear charge, sometimes it is extremely good! Why?
- Different  $(\kappa, b)$  for same  $N_e$  but different nuclear charge  
 $\Rightarrow$  give up universality entirely
- $(\kappa, b)$  depends on data points chosen for fitting
- $E_{\text{cum}}(S)$  is not always a function



# Errors swept under the rug... why do they almost "cancel"?

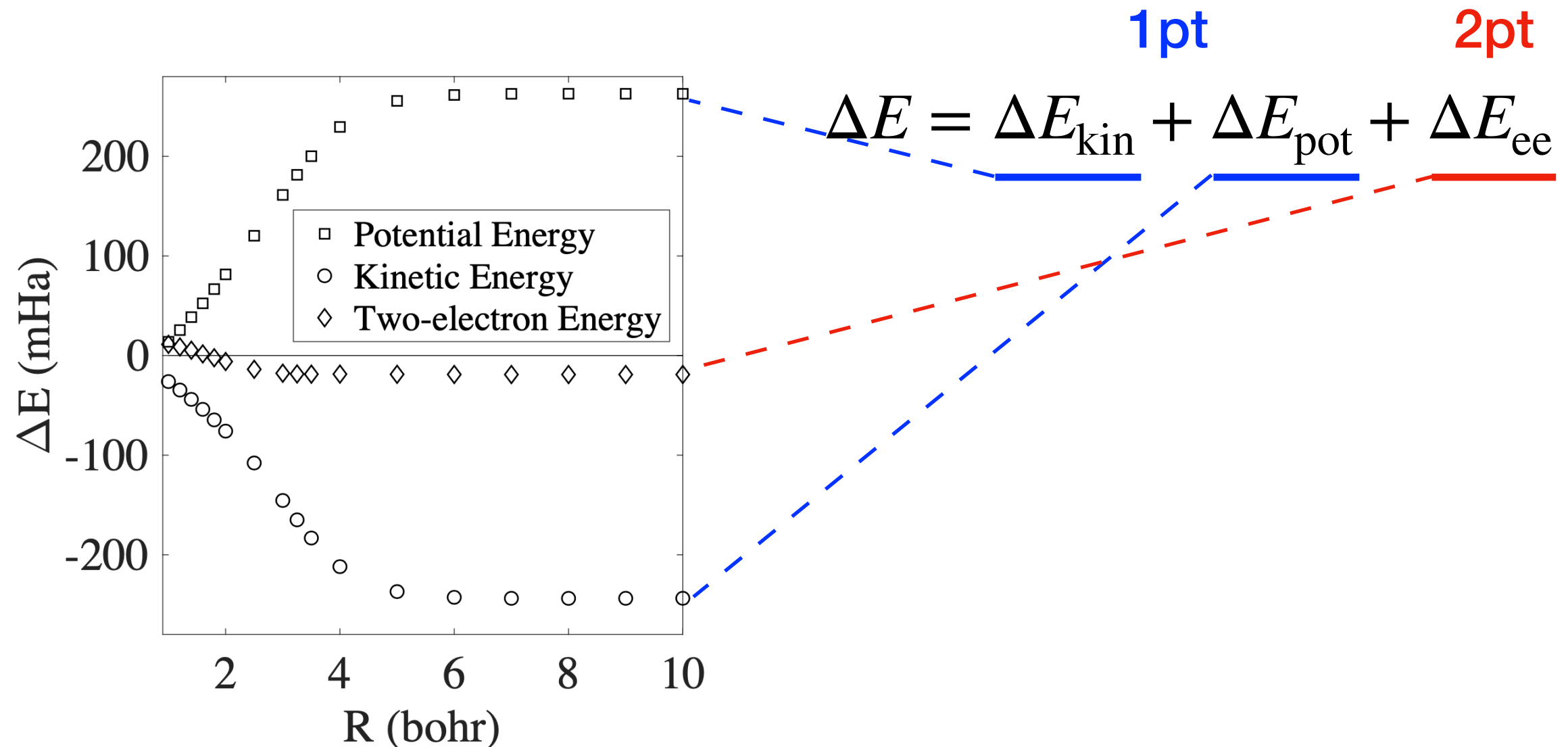
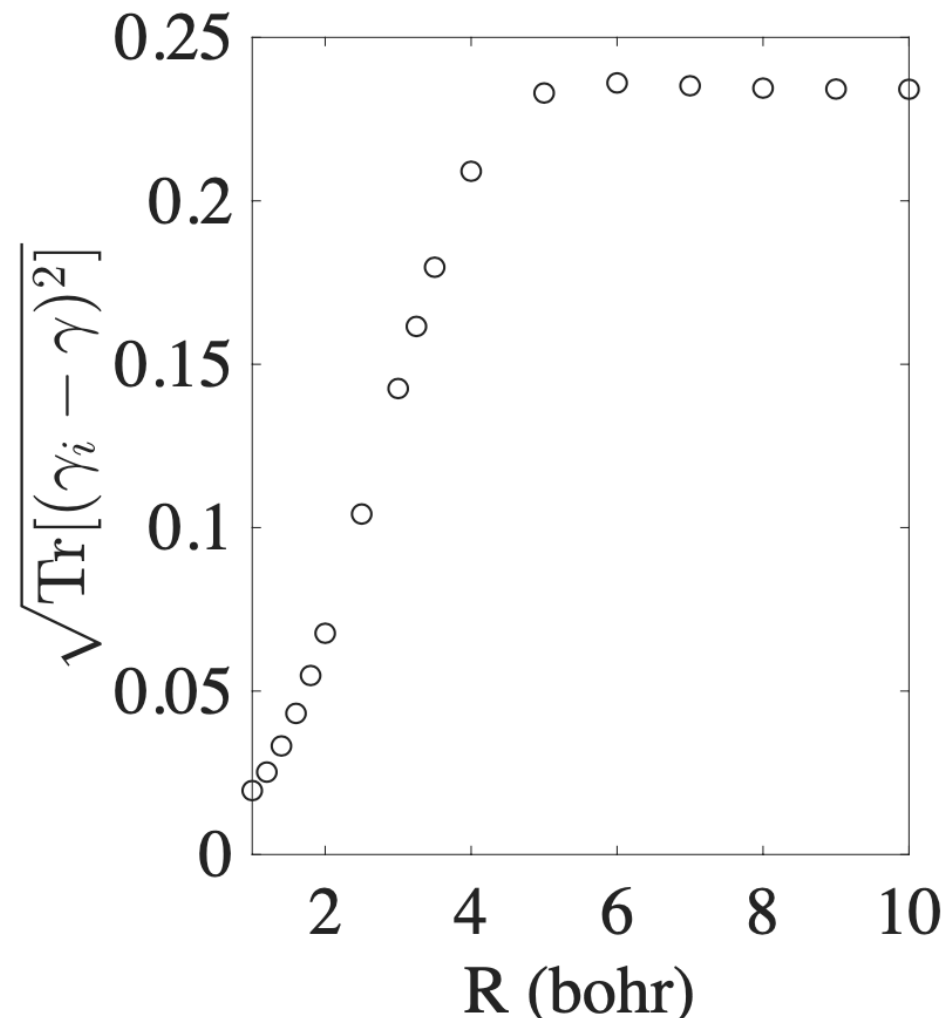


FIG. 3. *i*-DMFT (with basis set cc-pVDZ) applied to H<sub>2</sub> at different internuclear separation *R*: Error in the 1RDM  $\gamma_i$  measured by the Frobenius norm  $\sqrt{\text{Tr}[(\gamma_i - \gamma)^2]}$  with  $\gamma$  the exact/FCI 1RDM (left), error  $\Delta E$  in the kinetic energy, potential energy and interaction energy (right).

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# Does i-DMFT recover the correct 1RDM?



Similarity between matrices

$$\|\rho - \sigma\|_F = \sqrt{(\text{Tr}(\rho - \sigma)^\dagger (\rho - \sigma))}$$

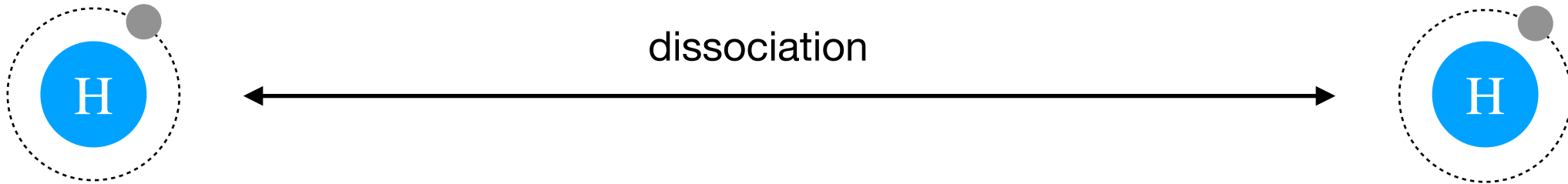
Bounds discrepancy in observables

$$\|(\rho - \sigma)O\|_F \leq \|O\|_{\text{op}} \|\rho - \sigma\|_F$$

**Bad orbitals  $\longrightarrow$  Bad 1pt quantities**

FIG. 3. *i*-DMFT (with basis set cc-pVDZ) applied to  $\text{H}_2$  at different internuclear separation  $R$ : Error in the 1RDM  $\gamma_i$  measured by the Frobenius norm  $\sqrt{\text{Tr}[(\gamma_i - \gamma)^2]}$  with  $\gamma$  the exact/FCI 1RDM (left), error  $\Delta E$  in the kinetic energy, potential energy and interaction energy (right).

# From Hartree-Fock to i-DMFT: conceptual flaw



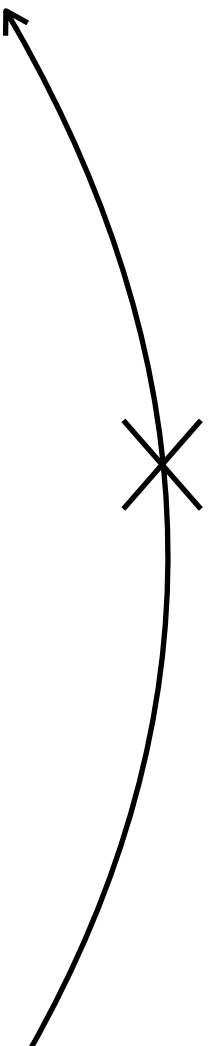
Suppose at dissociation limit for  $\text{H}_2$ , i-DMFT obtained the correct 1RDM

$$\bar{\gamma} = \frac{1}{2} \sum_{i=L/R, \tau=\uparrow, \downarrow} |1s_i\tau\rangle\langle 1s_i\tau|$$

Corresponding HF 2RDM  $\Gamma_{\text{HF}} = \frac{1}{2}(1 - \text{Ex})\gamma \otimes \gamma$

$$Y = \text{Tr}[W\Gamma_{\text{HF}}] = \frac{1}{4} \sum_{\tau, \tau'=\uparrow, \downarrow} \langle 1s_L\tau, 1s_R\tau' | W | 1s_L\tau, 1s_R\tau' \rangle \quad \text{Inter-nuclear}$$

$$+ \frac{1}{4} \sum_{i=L/R} \langle 1s_i\uparrow, 1s_i\downarrow | W | 1s_i\uparrow, 1s_i\downarrow \rangle \quad \text{Intra-nuclear}$$



# From Hartree-Fock to i-DMFT: conceptual flaw

$$\frac{1}{4} \sum_{i=L/R} \langle 1s_i \uparrow, 1s_i \downarrow | W | 1s_i \uparrow, 1s_i \downarrow \rangle$$

This intra-nuclear energy is unlikely to be accommodated by 1pt quantities such as  $S(\mathbf{n})$

J. Wang, E. Baerends, arxiv:2207.02616v1 (2022)

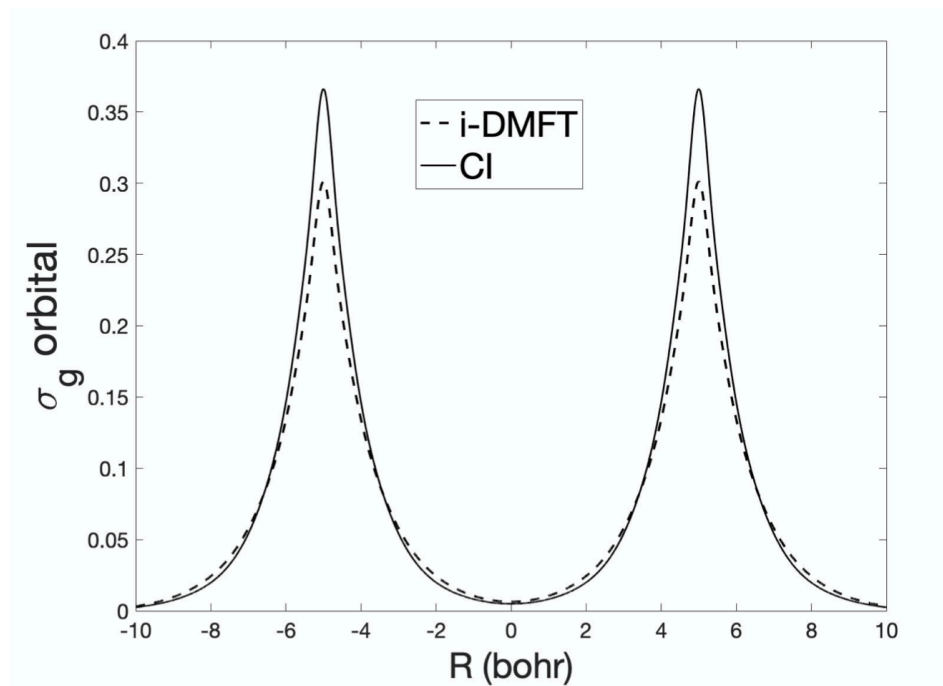


FIG. 1. The  $\sigma_g$  orbital from the  $i$ -DMFT calculation and the  $\sigma_g$  NO from a CI calculation at  $R=10$  bohr with the basis set cc-pVDZ. The  $i$ -DMFT parameters are  $\kappa=0.094681$  and  $b=0.0286189$  (in a.u.).

Adjustment by Wang & Baerends

$$E_{cum} = AE_x S - b \quad E_x = -\frac{1}{2} \sum_{ij} n_i n_j \langle ij | ji \rangle.$$



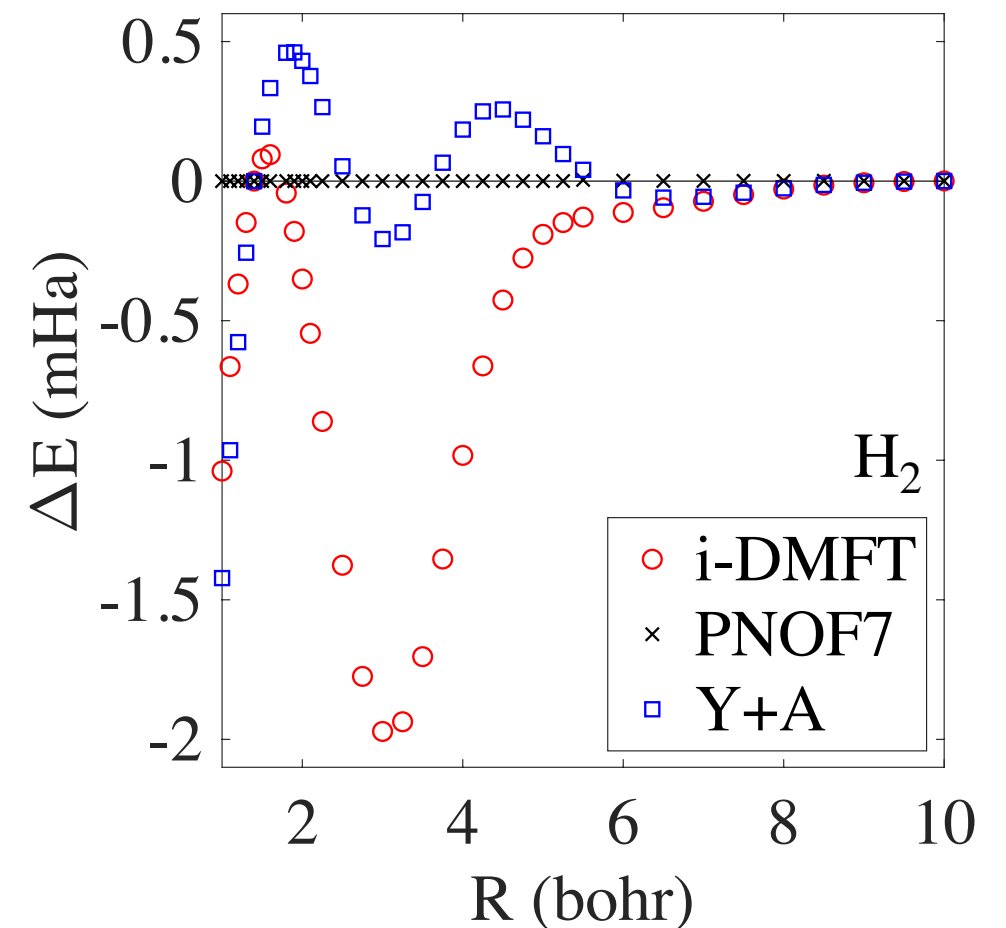
Collin's conjecture is perfect for  $H_2$

$$E_{cum} \approx \kappa S^{ph}(\mathbf{n}) + b$$

# Open questions

- When and why is the Collin's conjecture correct?
- Is convergence of i-DMFT guaranteed?
- How to do fitting for  $(\kappa, b)$ ? Two-point? Linear regression? Not conclusive!
- Is the von Neumann entropy the only relevant/useful QI quantity? Renyi entropy? Other Schul convex functions?
- 2RDM instead of 1RDM as optimization variable?

$$A(\mathbf{n}) = - \sum_i n_i^{0.8} \log(n_i^{0.8}) + (1 - n_i)^{0.8} \log((1 - n_i)^{0.8})$$

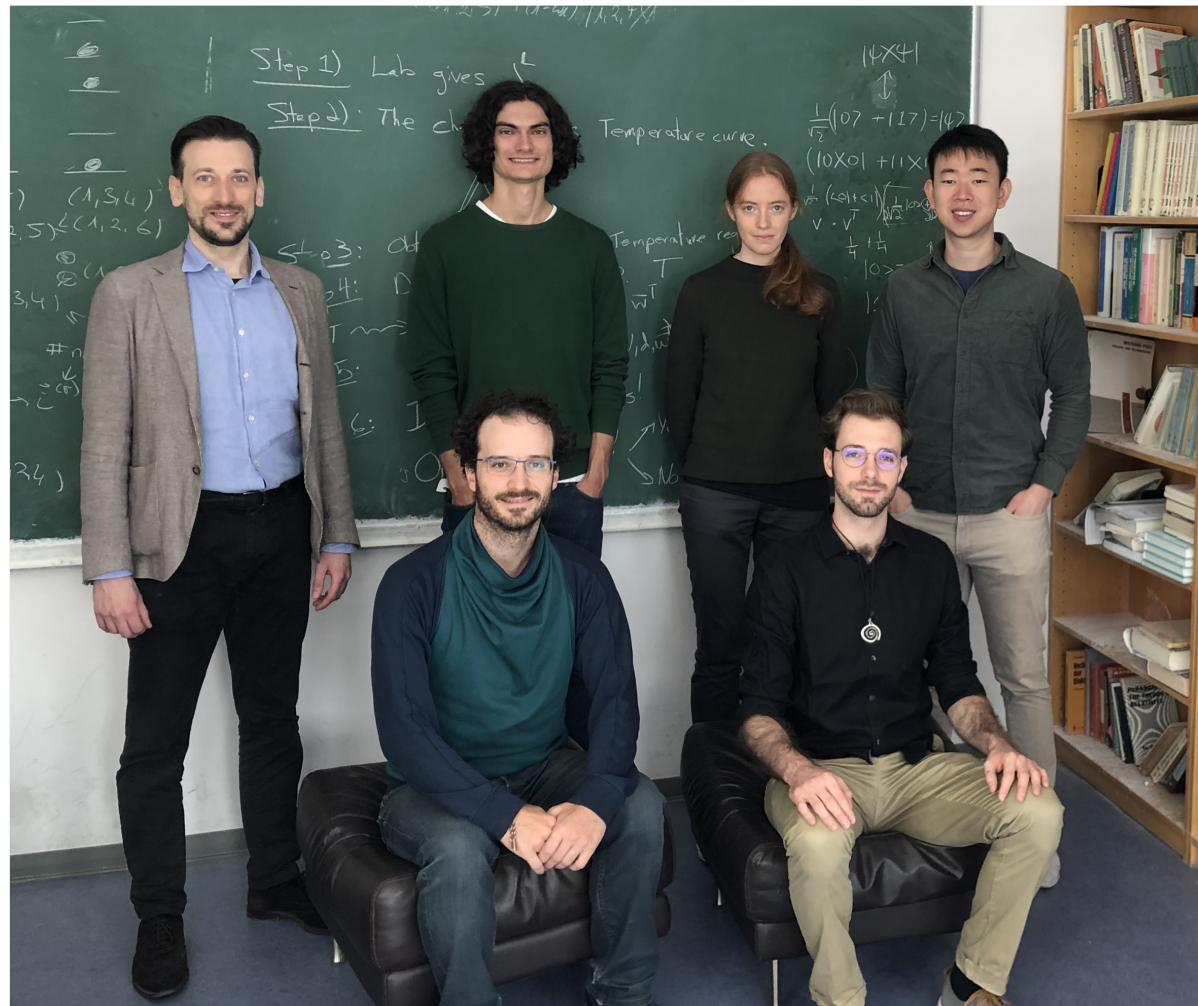


Ding, Liebert, Schilling, in preparation

# What people look for in a good RDMFT?

- correct static correlation in dissociation limit
- correct scaling at boundary, exchange force  
J. Liebert, C. Schilling Phys. Rev. Research 3, 013282 (2021)  
C. Schilling, R. Schilling Phys. Rev. Lett. 122, 013001 (2019)
- size versus volume extensivity  
J. Cioslowsky J. Chem. Phys. 120, 10364 (2004)
- accuracy of other 1-pt quantities other than energy, dipole moments, polarisability, kinetic energies...  
C. A. Coulson, Rev. Mod. Phys. 32, 170, (1960)
- ...

# Schilling group: Theoretical Quantum Physics



# Thank you!