J/ψ SDMEs with GlueX at 17 GeV

Keigo Mizutani Opportunities with JLab Energy and Luminosity Upgrade 27 September 2022

Why measure J/ψ SDMEs?

- Threshold J/ ψ photoproduction \rightarrow properties of proton mass?
- High E_{γ} region is well explained by Pomeron model.
- Understanding of production mechanism near threshold is desired.



In addition to σ_{total} and d σ/dt measurements, Spin observables (SDME) place stringent constraints on models!

Polarization measurements: Naturality

Decay angle wrt polarization plane is sensitive to naturality $P(-1)^J$ of the t-channel exchanged particle.



Naturality $\frac{\rho_{1-1}^1 - \text{Im}\rho_{1-1}^2}{2}$ is +0.5 when Pomeron exchange is dominant. Deviation from +0.5 implies unnatural parity exchanges.

Angular distributions of $J/\psi \rightarrow e^+e^-$



Decay angles (θ, ϕ) in helicity frame is sensitive to helicity conservation.

When helicity of the photon is fully transferred to J/ ψ , $W(\cos \theta) \sim 1 + \cos^2 \theta \quad (\rho_{00}^0 = 0)$ $W(\phi) = \text{flat} \quad (\text{Re}\rho_{1-1}^0 = 0)$

 Φ . angle btw polarization & production plane

$$\begin{array}{l} \mbox{Helicity} \\ \mbox{conservation} \end{array} \begin{cases} W(\cos\theta) = \frac{3}{8} \left(1 + \rho_{00}^0 + (1 - 3\rho_{00}^0)\cos^2\theta \right), \\ W(\phi) = \frac{1}{2\pi} \left(1 + \operatorname{Re}\rho_{1-1}^0\cos 2\phi \right), \\ W(\phi - \Phi) = \frac{1}{2\pi} \left(1 - P_\gamma \frac{\rho_{1-1}^1 - \operatorname{Im}\rho_{1-1}^2}{2}\cos\left[2\left(\phi - \Phi\right)\right]\right), \\ \mbox{Helicity} \\ \mbox{conservation} \end{cases} \\ W(\phi + \Phi) = \frac{1}{2\pi} \left(1 - P_\gamma \frac{\rho_{1-1}^1 + \operatorname{Im}\rho_{1-1}^2}{2}\cos\left[2\left(\phi + \Phi\right)\right]\right), \\ \mbox{Beam Asym.} (\Sigma) \qquad W(\Phi) = 1 - P_\gamma \left(2\rho_{11}^1 + \rho_{00}^1\right)\cos 2\Phi. \end{cases}$$

What we expect about J/ψ SDMEs?

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- No SDME measurements for threshold J/ ψ photoproduction so far.
- Helicity conservation is known to be broken for light vector mesons. The same for J/ψ or not?
- In high E_{γ} region, naturality close to +0.5 is observed for light vector mesons, corresponding to t-channel Pomeron exchange.
 - How about threshold J/ψ production?
 - Close to +0.5 because of 2 gluon exchange?
 - Or close to -0.5 because of 3 gluon exchange?
- Beam asymmetry $\sum \sim 0$ is observed for light vector mesons. $\sum \sim 0$ for J/ ψ as well?

Basically, we have no knowledge about J/ψ SDMEs near thr. That's why we measure them.

GlueX can do unique measurements of naturality.

$\gamma p \to J/\psi (\to e^+ e^- p)$ analysis

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GlueX-I + 30% GlueX-II data are used for this analysis. J/ ψ is identified using $M(e^+e^-)$ distribution (10-20% background). Calorimeter response (p/E) is used to subtract π misidentification background.



Polarized measurements





 $n \cos 2(\phi - \phi) = P (o^{1} - Imo^{2})/2$

 $-n \cos^2(\phi + \phi) = P (o^1 + lmo^2)/2$



 9_{9} 10₁₀ 11₁₁ 12₁₂ 13₃

0

The same assumption as Lubomir's talk.

14 14

15



0,1 0.1

00

9⁹

Polarization FOM increases ~30 times in the coherent peak region. Precise measurement of naturality is possible with 17 GeV data.

 10^{10} 11^{11} 12^{12} 13^{13} 14^{14} 15^{15} 16^{16} 17 17 E_{γ} (GeV) eV)

Polarized SDMEs at 17 GeV



Also, other SDMEs provide strong constraints on the production mechanism.

Unpolarized SDME (cosθ_{hel})





Helicity conservation at 17 GeV

Non-zero values mean the helicity is not conserved.



Luminosity near threshold region ($E_{\gamma} < 12.5$ GeV) is 5.5 times larger with 17 GeV beam.

Precise check of helicity conservation is possible with 17 GeV beam.

E_v dep. of Unpolarized SDMEs



Measurements of E_{γ} dep. of helicity conservation will be improved with 17 GeV beam as seen above.



- J/ψ SDMEs are measured to help determine the production mechanism near threshold.
- We have no SDME data so far. No reliable predictions either.
 - GlueX can provide unique polarized SDME measurements.
- 17 GeV energy upgrade gives 30 times larger polarization FOM, and significantly increase the precision of polarized measurements.
 - Precise measurements of naturality.
- For unpolarized measurements, 17 GeV upgrade gives 5.5 times larger yields near threshold region. Helicity conservation can be checked precisely.

Backup

GJ frame to check helicity conservation 17



Decay asymmetry measurements in the helicity frame are used to check "helicity conservation" where incoming photon helicity is fully transferred to J/ψ .

GJ frame allows us to check "helicity conservation in $\gamma J/\psi \rightarrow \bar{p}p$ ".

For light vector mesons (ρ , ω , ϕ), it is known that helicity conservation in $\gamma V \rightarrow \bar{p}p$ is badly broken.



E_v dep. of Unpolarized SDMEs (GJ)



 E_{γ} dep. of helicity conservation in $\gamma J/\psi \rightarrow \bar{p}p$ will be more precisely measured with 17 GeV beam, which place strict constraints on models.

GJ frame analysis at 17 GeV



5.5 times larger luminosity (17 GeV beam) allows us more precise study about helicity conservation in $\gamma J/\psi \to \bar{p}p$.

JPAC Pomeron model by Daniel Winney 21

Helicity frame Solid: vector Pomeron Dashed: Pomeron with Helicity conserved

GJ frame Solid: vector Pomeron Dashed: Pomeron with Helicity conserved



Ref.) A. I. Titov et al., PRC60,035205 (1999)

In the case of $\phi \to K^+K^-$ ($a = K^+$), we obtain:

$$W(\cos\theta,\phi) = \frac{3}{4\pi} \left(\frac{\rho_{11} + \rho_{-1-1}}{2} \sin^2\theta + \rho_{00} \cos^2\theta - \frac{\text{Re}\rho_{10} - \text{Re}\rho_{-10}}{\sqrt{2}} \sin 2\theta \cos\phi + \frac{\text{Im}\rho_{10} + \text{Im}\rho_{-10}}{\sqrt{2}} \sin 2\theta \sin\phi - \text{Re}\rho_{1-1} \sin^2\theta \cos 2\phi + \text{Im}\rho_{1-1} \sin^2\theta \sin 2\phi \right), \quad (2)$$

In the case of $J/\psi \rightarrow e^-e^+$ ($a = e^-$), we obtain:

$$W(\cos\theta,\phi) = \frac{3}{8\pi} \left(\frac{\rho_{11} + \rho_{-1-1}}{2} \left(1 + \cos^2\theta \right) + \rho_{00} \sin^2\theta + \frac{\text{Re}\rho_{10} - \text{Re}\rho_{-10}}{\sqrt{2}} \sin 2\theta \cos\phi - \frac{\text{Im}\rho_{10} + \text{Im}\rho_{-10}}{\sqrt{2}} \sin 2\theta \sin\phi + \text{Re}\rho_{1-1} \sin^2\theta \cos 2\phi - \text{Im}\rho_{1-1} \sin^2\theta \sin 2\phi \right), \quad (3)$$

Linearly polarized distributions



For K^+K^- case,

$$W^{0}(\cos\theta,\phi) = \frac{3}{4\pi} \left(\frac{1-\rho_{00}^{0}}{2} + \frac{3\rho_{00}^{0}-1}{2}\cos^{2}\theta - \sqrt{2}\operatorname{Re}\rho_{10}^{0}\sin2\theta\cos\phi - \operatorname{Re}\rho_{1-1}^{0}\sin^{2}\theta\cos2\phi \right),$$
(4)

$$W^{1}(\cos\theta,\phi) = \frac{3}{4\pi} \left(\rho_{11}^{1} \sin^{2}\theta + \rho_{00}^{1} \cos^{2}\theta - \sqrt{2} \operatorname{Re} \rho_{10}^{1} \sin 2\theta \cos\phi - \operatorname{Re} \rho_{1-1}^{1} \sin^{2}\theta \cos 2\phi \right),$$
(5)

$$W^{\alpha}(\cos\theta,\phi) = \frac{3}{4\pi} \left(\sqrt{2}\mathrm{Im}\rho_{10}^{\alpha}\sin2\theta\sin\phi + \mathrm{Im}\rho_{1-1}^{\alpha}\sin^2\theta\sin2\phi\right) \quad (\alpha = 2,3).$$
(6)

For e^-e^+ case,

$$W^{0}(\cos\theta,\phi) = \frac{3}{8\pi} \left(\frac{1+\rho_{00}^{0}}{2} - \frac{3\rho_{00}^{0}-1}{2}\cos^{2}\theta + \sqrt{2}\operatorname{Re}\rho_{10}^{0}\sin2\theta\cos\phi + \operatorname{Re}\rho_{1-1}^{0}\sin^{2}\theta\cos2\phi \right),$$
(7)

$$W^{1}(\cos\theta,\phi) = \frac{3}{8\pi} \left(\rho_{11}^{1} \left(1 + \cos^{2}\theta \right) + \rho_{00}^{1} \sin^{2}\theta + \sqrt{2} \operatorname{Re} \rho_{10}^{1} \sin 2\theta \cos\phi + \operatorname{Re} \rho_{1-1}^{1} \sin^{2}\theta \cos 2\phi \right),$$
(8)

$$W^{\alpha}(\cos\theta,\phi) = \frac{3}{8\pi} \left(-\sqrt{2} \mathrm{Im}\rho_{10}^{\alpha} \sin 2\theta \sin \phi - \mathrm{Im}\rho_{1-1}^{\alpha} \sin^2\theta \sin 2\phi \right) \quad (\alpha = 2,3).$$
(9)

Averaged one-dimensional distributions 24

For $V \rightarrow 2$ spinless particles,

$$W(\cos\theta) = \frac{3}{2} \left(\frac{1 - \rho_{00}^0}{2} \sin^2\theta + \rho_{00}^0 \cos^2\theta \right),$$
(10)

$$W(\phi) = \frac{1}{2\pi} \left(1 - 2\text{Re}\rho_{1-1}^0 \cos 2\phi \right),$$
(11)

$$W(\phi - \Phi) = \frac{1}{2\pi} \left(1 + 2P_{\gamma} \frac{\rho_{1-1}^1 - \mathrm{Im}\rho_{1-1}^2}{2} \cos\left[2\left(\phi - \Phi\right)\right] \right),\tag{12}$$

$$W(\phi + \Phi) = \frac{1}{2\pi} \left(1 + 2P_{\gamma} \frac{\rho_{1-1}^{1} + \operatorname{Im} \rho_{1-1}^{2}}{2} \cos\left[2\left(\phi + \Phi\right)\right] \right), \tag{13}$$

$$W(\Phi) = 1 - P_{\gamma} \left(2\rho_{11}^1 + \rho_{00}^1 \right) \cos 2\Phi.$$
(14)

For $J/\psi \rightarrow e^-e^+$,

$$W(\cos\theta) = \frac{3}{8} \left(1 + \rho_{00}^{0} + (1 - 3\rho_{00}^{0})\cos^{2}\theta \right),$$
(15)

$$W(\phi) = \frac{1}{2\pi} \left(1 + \operatorname{Re}\rho_{1-1}^{0}\cos 2\phi \right),$$
(16)

$$W(\phi - \Phi) = \frac{1}{2\pi} \left(1 - P_{\gamma} \frac{\rho_{1-1}^{1} - \operatorname{Im}\rho_{1-1}^{2}}{2}\cos \left[2\left(\phi - \Phi\right)\right] \right) \underbrace{\operatorname{Naturality}}_{Naturality} (17)$$

$$W(\phi + \Phi) = \frac{1}{2\pi} \left(1 - P_{\gamma} \frac{\rho_{1-1}^{1} + \operatorname{Im}\rho_{1-1}^{2}}{2}\cos \left[2\left(\phi + \Phi\right)\right] \right),$$
(18)

$$W(\Phi) = 1 - P_{\gamma} \left(2\rho_{11}^{1} + \rho_{00}^{1}\right)\cos 2\Phi. \qquad \operatorname{Beam Asym.} \left(\sum\right) (19)$$