

# DIFFERENT STRATEGIES FOR EFT PARAMETER ESTIMATION IN THE FEW-NUCLEON SECTOR

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- And many people in the *ab initio nuclear theory* community for enlightening discussions



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# INTRODUCTION

What information can be inferred<sup>1</sup> from available few-nucleon data<sup>2,3</sup> to state–of-the-art "models"<sup>4</sup> of the strong force between nucleons?

<sup>1</sup> using either frequentist or bayesian approaches.

- <sup>2</sup> Here:  $\pi N$  and NN scattering data
- <sup>3</sup>...plus A=2-3 bound-state observables
- <sup>4</sup> effective field theories => systematically improvable

# FIVE-BULLET OVERVIEW OF WHAT I WILL <u>NOT</u> TALK ABOUT

- What are the predictions of these few-body constrained Hamiltonians for many-body observables...
- Different power-counting schemes (but I will offer the Bayesian viewpoint on this)...
- Details of the computational tools that we are using: the nsopt code base, Automatic differentiation, MCMC algorithms, ...
- The facts that Sweden eliminated both the Netherlands and Italy in the World Cup qualifier (I'm actually a bit sorry about that), plus the observation that the German national team has not won a single game this year.
- Inference from the above data on to the probability that Sweden will beat Germany in the World Cup on June 23rd.



Note: Statistical analysis revisited following the *"Schiavilla-correction"* (relation between cD in the one-pion exchange plus contact NNN potential and the LEC multiplying the contact axial-vector current)

# A DATA-CENTRIC VIEW ON THE EFT DESCRIPTION OF NUCLEAR FORCES

## Inference

"the act of passing from one proposition, statement, or judgment considered as true to another whose truth is believed to follow from that of the former" (Webster)

Do premises  $A, B, \ldots \rightarrow$  hypothesis, H?

- Assume that **hypothesis**  $H_i$  is a model  $M_i$  with parameters  $\alpha_i$ .
- Inductive inference: Premises bear on truth/falsity of H, but don't allow its definite determination
- Statistical Inference: Quantify the strength of inductive inferences from data and other premises to hypotheses about the phenomena producing the data.
- Quantify via probabilities, or averages calculated using probabilities. Frequentists and Bayesians use probabilities very different for this.

# DATA AND ITS IMPORTANCE FOR EFFECTIVE FIELD THEORIES

- Short-range physics for an EFT is encoded in LECs.
- These LECs can be inferred by confrontation with (lowenergy) observables.
- Many issues for chiral EFT remain:
  - regulator artefacts,
  - convergence,
  - The choice of relevant degrees-of-freedom

such issues can only be adressed with proper uncertainty quantification (UQ).



Work with A. Johansson and A. Ekström

# **BAYESIAN POSTERIORS IN THE NUCLEON-NUCLEON SECTOR**



- Assume that hypothesis H<sub>i</sub> is a model M<sub>i</sub> with parameters p<sub>i</sub>.
- In Bayesian statistics we assess the hypotheses by calculating their probabilities p(H<sub>i</sub>|...) conditional on known and/or presumed information using the rules of probability theory.

### Parameter estimation:

Assume that the model  $M_i$  is true; Compute:  $p(\mathbf{p}_i | D_{obs}, M_i, I)$ 

### Model comparison:

Compute ratio:  $p(D_{obs} | M_i, I) / p(D_{obs} | M_j, I)$ 

**Bayes' theorem** (follows from probability product rule):

$$\begin{array}{lll} \textbf{posterior} & \textbf{likelihood} & \textbf{prior} \\ p(\pmb{p}|D,I) = \frac{p(D|\pmb{p},I)p(\pmb{p}|I)}{p(D|I)} \\ \textbf{Bayesian evidence} \end{array}$$

**Marginalization:**  $p(p_1|D, I) = \int dp_2 \dots dp_k p(\boldsymbol{p}|D, I)$ 

- For many lessons and suggestions on the use of Bayesian methods in Effective Field Theories, see work by the BUQEYE collaboration (and talk by Sarah).
- Here we report on progress in implementing Bayesian methods for parameter estimation in Chiral EFT (up to N3LO) using NN scattering data (phase shifts).

# **NN PHASE SHIFT OPTIMIZATION**

Let us consider the task of determining a set of low-energy constants (LECs) by fitting to phase shifts.



# $p(\tilde{C}_{1S0}, C_{1S0}|D, I) \propto p(D|\tilde{C}_{1S0}, C_{1S0}, I)p(\tilde{C}_{1S0}, C_{1S0}|I)$

posterior

likelihood

prior

1S0 channel





![](_page_15_Figure_1.jpeg)

![](_page_16_Figure_1.jpeg)

![](_page_17_Figure_1.jpeg)

# The 1SO(np) channel @ N3L0 with redundant parameter

![](_page_18_Figure_1.jpeg)

# The 1SO(np) channel @ N3LO

![](_page_19_Figure_1.jpeg)

![](_page_20_Figure_1.jpeg)

$$p(\pmb{\alpha}|D,I) = \frac{p(D|\pmb{\alpha},I)p(\pmb{\alpha}|I)}{p(D|I)}$$
Bayesian evidence (=Z)

## The deuteron channel

![](_page_21_Figure_1.jpeg)

# **Expectation integrals, error propagation**

Expectation integrals for observables can be performed using the posterior pdf

$$\langle O(\boldsymbol{\alpha}) \rangle = \int d\boldsymbol{\alpha} p(\boldsymbol{\alpha} | D, I) O(\boldsymbol{\alpha})$$
$$\approx \frac{1}{N} \sum_{j=1}^{N} O(\boldsymbol{\alpha}_j)$$
The MCMC elements

The MCMC algorithm generates N samples  $\{\alpha_j\}$  according to the posterior pdf

## **Deuteron observables**

![](_page_23_Figure_1.jpeg)

## **Deuteron observables**

![](_page_24_Figure_1.jpeg)

![](_page_25_Picture_0.jpeg)

Note: Statistical analysis revisited following the *"Schiavilla-correction"* (relation between cD in the one-pion exchange plus contact NNN potential and the LEC multiplying the contact axial-vector current)

# A DATA-CENTRIC VIEW ON NUCLEAR Forces: CHI-Squared Minimization

Low-energy constants (LECs) need to be fitted to experimental data.

$$\chi^{2}(\vec{p}) \equiv \sum_{i} r_{i}^{2}(\vec{p}) = \sum_{j \in NN} r_{j}^{2}(\vec{p}) + \sum_{k \in \pi N} r_{k}^{2}(\vec{p}) + \sum_{l \in 3N} r_{l}^{2}(\vec{p})$$

#### # parameters that are allowed to vary:

![](_page_26_Figure_4.jpeg)

# **INPUT AND TECHNOLOGY**

### $\pi N$ scattering

- WI08 database
- T<sub>lab</sub> between 10-70 MeV
- N<sub>data</sub> = 1347
- R. Workman et al. (2012)

### **NN scattering**

- Granada '13 database
- T<sub>lab</sub> between 0-290 MeV
- N<sub>data</sub> = 4753 (np + pp)
- R. Navarro Pérez et al. (2013)

### All 6000 residuals computed on 1 node in ~90 sec.

### A=2,3 bound states

 <sup>2</sup>H,<sup>3</sup>H,<sup>3</sup>He [binding energy, radius, Q(<sup>2</sup>H), <sup>3</sup>H half life] On 1 node in ~10 sec

+ derivatives! (×2-20 cost)

### Alternatively... theoretical analysis of data

#### $\pi N$ scattering

Roy-Steiner analysis
 M. Hoferichter et al. (2015)

### **NN scattering**

• Phase shifts from partial wave analysis

# **OPTIMIZATION STRATEGY**

#### Low-energy constants (LECs) need to be fitted to experimental data.

$$\chi^2(\vec{p}) \equiv \sum_i \left(\frac{O_i^{\text{theo}}(\vec{p}) - O_i^{\text{expr}}}{\sigma_{\text{tot},i}}\right)^2 \equiv \sum_i r_i^2(\vec{p})$$

- Derivative-free optimization using POUNDerS was used in our earliest works
- More efficient minimization algorithms (Levenberg-Marquardt, Newton), and statistical error analysis require derivatives

$$\frac{\partial r_i}{\partial p_j}$$
 and  $\frac{\partial^2 r_i}{\partial p_j \partial p_k}$ 

- Numerical derivation using finite differences is plagued by low numerical precision and is computationally costly.
- Instead, we use Automatic Differentiation (AD)

# TOTAL ERROR BUDGET

- The total error budget is  $\sigma_{\text{tot}}^2 = \sigma_{\text{exp}}^2 + \sigma_{\text{method}}^2 + \sigma_{\text{numerical}}^2 + \sigma_{\text{model}}^2$
- At a given chiral order v, the omitted diagrams should be of order

 $\mathcal{O}\left((Q/\Lambda_{\chi})^{\nu+1}\right)$ 

- Still needs to be converted to actual numbers  $\sigma_{model}$
- We translate this EFT knowledge into an error in the scattering amplitudes

$$\sigma_{\text{model},x}^{(\text{amp})} = C_x \left(\frac{Q}{\Lambda_{\chi}}\right)^{\nu+1}, \quad x \in \{NN, \pi N\}$$

which is then propagated to an error in the observable.

# **TOTAL NP CROSS SECTION**

![](_page_30_Figure_1.jpeg)

# **Quadratic error propagation vs Brute force sampling**

![](_page_31_Figure_1.jpeg)

# Systematic uncertainties: input data, regulator cutoff

![](_page_32_Figure_1.jpeg)

 ▶ 6 different NN-scattering datasets  $T_{lab} \in [0, T_{lab,max}]$ , with  $T_{lab,max}$ =125, ..., 290 MeV

# Systematic uncertainties: input data, regulator cutoff

**Previous version** 

![](_page_33_Figure_1.jpeg)

- 7 different regulator cutoffs:
  Λ=450, 475, ..., 575, 600 MeV
- 6 different NN-scattering datasets  $T_{lab} \in [0, T_{lab,max}]$ , with  $T_{lab,max}$ =125, ..., 290 MeV

# CONCLUSION

# OUTLOOK

- The inclusion of more data in the objective function requires other approaches to the optimization problem needed. (See also Andreas' presentation.)
- The frequentist approach does not offer an easy and transparent method for handling systematic uncertainties or imposing prior knowledge.
- Bayesian parameter estimation is advantageous, but costly.
  - avoiding the need to 'judge', a priori, what data can be included in order to safely avoid overfitting.
  - offers a viable approach to include prior knowledge of certain parameters from Roy-Steiner analysis.
- Investigate other chiral EFT power-counting schemes via Bayesian evidence.

# QUESTIONS

- How to best combine different kinds of experimental data (such as NN scattering and A=3 bound-state observables)?
- And how to combine this with (prior) information from a theoretical analysis (RS eqs., sub-threshold parameters)?
- Feasible ways to evaluate multi-dimensional posterior distributions, and perform sub-sequent error propagation?
- What are the most relevant tests of various power counting schemes given the Bayesian framework?