Implementation and applications of 3N interactions and nuclear currents

Kai Hebeler Trento, June 6, 2018

ECT* workshop "New Ideas in Constraining Nuclear Forces"





Separation of long- and short-range physics



$$\mathbf{p} = (\mathbf{p}_1 - \mathbf{p}_2)/2$$

 $\mathbf{p}' = (\mathbf{p}'_1 - \mathbf{p}'_2)/2$
 $\mathbf{q} = (\mathbf{p}_1 - \mathbf{p}'_1)$

Separation of long- and short-range physics



nonlocal
$$V_{\rm NN}(\mathbf{p},\mathbf{p}') \to \exp\left[-\left((p^2+p'^2)/\Lambda^2\right)^n\right]V_{\rm NN}(\mathbf{p},\mathbf{p}')$$

Epelbaum, Glöckle, Meissner, NPA 747, 362 (2005) Entem, Machleidt, PRC 68, 041001 (2003)

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(momentum space)

$$V_{NN}(\mathbf{q}) \rightarrow \exp\left[-\left(q^2/\Lambda^2\right)^n\right] V_{NN}(\mathbf{q})$$

(f. Navratil, Few-body Systems 41, 117 (2007) Reinert et al., arXiv:1711.08821

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local (momentum space)

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(coordinate space)
$$V_{\rm NN}^{\pi}(\mathbf{r}) \rightarrow \left(1 - \exp\left[-\left(r^2/R^2\right)^n\right]\right) V_{\rm NN}^{\pi}(\mathbf{r})$$
$$\delta(\mathbf{r}) \rightarrow \alpha_n \exp\left[-\left(r^2/R^2\right)^n\right]$$

Gezerlis et. al, PRL, 111, 032501 (2013)

Separation of long- and short-range physics



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Epelbaum et. al, PRL, 115, 122301 (2015)

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$$V_{\rm NN}^{\pi}(\mathbf{r}) \to \left(1 - \exp\left[-\left(r^2/R^2\right)\right]\right)^n V_{\rm NN}^{\pi}(\mathbf{r})$$
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semi-local

Power counting in chiral 3N sector (nonlocal): Contributions of many-body forces at N³LO in neutron matter



n [fm⁻³]

 $n [fm^{-3}]$

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Power counting in chiral 3N sector (nonlocal): Contributions of many-body forces at N³LO in neutron matter



Representation of 3N interactions in momentum space

 $|pq\alpha\rangle_i \equiv |p_iq_i; [(LS)J(ls_i)j] \mathcal{J}\mathcal{J}_z(Tt_i)\mathcal{T}\mathcal{T}_z\rangle$



Due to the large number of matrix elements, the traditional way of computing matrix elements requires extreme amounts of computer resources.

$$N_p \simeq N_q \simeq 15$$

$$N_\alpha \simeq 30 - 180 \qquad \longrightarrow \quad \dim[\langle pq\alpha | V_{123} | p'q'\alpha' \rangle] \simeq 10^7 - 10^{10}$$

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A 'new' algorithm allows efficient calculation. KH, Krebs, Epelbaum, Golak, Skibinski, PRC 91, 044001(2015)

Calculation of 3N forces in momentum partial-wave representation

 $\langle pq\alpha | V_{123} | p'q'\alpha' \rangle \sim \sum_{m_i} \int d\hat{\mathbf{p}} \, d\hat{\mathbf{q}} \, d\hat{\mathbf{p}}' \, d\hat{\mathbf{q}}' Y_l^m(\hat{\mathbf{p}}) Y_{\bar{l}}^{\bar{m}}(\hat{\mathbf{q}}) \, \langle \mathbf{pq}ST | V_{123} | \mathbf{p'q'}S'T' \rangle \, Y_{l'}^{m'}(\hat{\mathbf{p}}') Y_{\bar{l}'}^{\bar{m}'}(\hat{\mathbf{q}}')$

traditional method:

- reduce dimension of angular integrals from 8 to 5 by using symmetry
- discretize angular integrals and perform all sums numerically

much more efficient method:

- use that all interaction contributions (except rel. corr.) are local: $\langle \mathbf{pq}|V_{123}|\mathbf{p'q'}\rangle = V_{123}(\mathbf{p} - \mathbf{p'}, \mathbf{q} - \mathbf{q'})$ $= V_{123}(p - p', q - q', \cos \theta)$
 - \rightarrow allows to perform all except for 3 integrals analytically
- only a few small discrete internal sums need to be performed for each external momentum and angular momentum

Hartree-Fock energy of infinite matter (unregularized 3NF)

8 6 4 **⊙** – –**⊚** ^C E_{3N} [MeV] \rightarrow N3LO 2π (no c_i) 2 - N3LO 2π - 1π → N3LO rings ▶ N3LO 2π -contact \triangleleft N3LO rel. corr. (C_s) - **N**3LO rel. corr. (C_T) -2 -4 -6 [1/2,+][3/2,-] [3/2,+][5/2,-] [5/2,+][7/2,-][7/2,+][9/2,-] [9/2,+]exact [1/2,-] $[\mathcal{J}, P]$ KH, Krebs, Epelbaum, Golak, Skibinski, PRC 91,044001 (2015)

neutron matter:

- in PNM only matrix elements with T = 3/2 contribute
- resummation up to $\mathcal{J} = 9/2$ leads to well converged results
- essentially perfect agreement with 'exact' results (cf. PRC88, 025802)

Hartree-Fock energy of infinite matter (unregularized 3NF)

symmetric nuclear matter:



- resummation up to $\mathcal{J} = 9/2$ leads to well converged results
- essentially perfect agreement with 'exact' results (cf. PRC88, 025802)

3NF power counting in 3H for different regulators



PRC 91,044001 (2015)

3NF power counting in 3H for different regulators



- size of N3LO contribution not suppressed for shown nonlocal interactions
- N3LO contributions suppressed for semilocal interactions
- technical challenges for semilocal interactions:
 - * forces non-perturbative, large basis spaces/RG evolution needed
 - * implementation of 3N forces hard, stability problems for scattering calculations
 - * derivation and implementation of nuclear currents hard

3NF power counting for different regulators







Computational strategy:

(1) calculate unregularized 3NF in sufficiently large partial-wave basis(2) fourier transform coordinate space regulator to momentum space



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(1) calculate unregularized 3NF in sufficiently large partial-wave basis (2) fourier transform coordinate space regulator to momentum space (3) decompose regulator f_{LR} in partial wave momentum basis (4) perform convolution integrals:

$$\langle pq\alpha | V_{123}^{\rm reg} | p'q'\alpha' \rangle = \int d\tilde{q} \,\tilde{q}^2 \int d\tilde{p} \,\tilde{p}^2 \sum_{\tilde{\alpha}} \langle pq\alpha | V_{123} | \tilde{p}\tilde{q}\tilde{\alpha} \rangle \,\langle \tilde{p}\tilde{q}\tilde{\alpha} | f_{LR} | p'q'\alpha' \rangle$$



Computational strategy:

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(5) regularize short-range parts in interactions with non-local regulator



Computational strategy:

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(5) regularize short-range parts in interactions with non-local regulator(6) antisymmetrize interactions (optional)

Naive calculation of convolution integrals: Problem!

$$V_{reg}(\mathbf{r}_{12}, \mathbf{r}_{13}, \mathbf{r}_{23}) = V(\mathbf{r}_{12}, \mathbf{r}_{13}, \mathbf{r}_{23})F(r_{12})F(r_{13})F(r_{23})$$
$$V_{reg}(\mathbf{r}_{12}, \mathbf{r}_{13}) = V(\mathbf{r}_{12}, \mathbf{r}_{13})F(r_{12})F(r_{13})$$
$$V_{reg}(\mathbf{r}_{12}) = V(\mathbf{r}_{12})F(r_{12})$$

with
$$F(r_{ij}) = (1 - \exp(-r_{ij}^2/R^2))^{n_{exp}}$$

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Problem:

for practical calculation of the convolution integrals we need to explicitly separate the delta function part

$$\langle \mathbf{p}\mathbf{q}|V_{reg}|\mathbf{p'q'}\rangle = \int d\tilde{\mathbf{p}}d\tilde{\mathbf{q}} \langle \mathbf{p}\mathbf{q}|V|\tilde{\mathbf{p}}\tilde{\mathbf{q}}\rangle \langle \tilde{\mathbf{p}}\tilde{\mathbf{q}}|R|\mathbf{p'q'}\rangle$$

$$= \int d\tilde{\mathbf{p}}d\tilde{\mathbf{q}} \langle \mathbf{p}\mathbf{q}|V|\tilde{\mathbf{p}}\tilde{\mathbf{q}}\rangle \langle \tilde{\mathbf{p}}\tilde{\mathbf{q}}|R-1|\mathbf{p'q'}\rangle + \langle \mathbf{p}\mathbf{q}|V|\mathbf{p'q'}\rangle$$

$$delicate$$

$$cancellation!$$

consider a N2LO long-range topology:

$$V(\mathbf{r}_{13}, \mathbf{r}_{23}) = \int \frac{d\mathbf{q}_1}{(2\pi)^3} \int \frac{d\mathbf{q}_2}{(2\pi)^3} e^{i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{i\mathbf{q}_3 \cdot \mathbf{r}_{23}} V(\mathbf{q}_2, \mathbf{q}_3)$$

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for the calculation of the regularized interaction we insert an identity

$$V_{\rm reg}(\mathbf{r}_{13}, \mathbf{r}_{23}) = V(\mathbf{r}_{13}, \mathbf{r}_{23}) \frac{Q(r_{13}^2)}{Q(r_{13}^2)} \frac{Q(r_{23}^2)}{Q(r_{23}^2)} \left(1 - e^{-r_{13}^2/R^2}\right)^6 \left(1 - e^{-r_{23}^2/R^2}\right)^6$$

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and define a *preregularized* interaction:

$$V_{\text{prereg}}(\mathbf{q}_2, \mathbf{q}_3) = \int d\mathbf{r}_{13} \int d\mathbf{r}_{23} e^{-i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{-i\mathbf{q}_3 \cdot \mathbf{r}_{23}} Q(r_{13}^2) Q(r_{23}^2) V(\mathbf{r}_{13}, \mathbf{r}_{23}) = Q(-\Delta_{q_2}) Q(-\Delta_{q_3}) V(\mathbf{q}_2, \mathbf{q}_3)$$

consider a N2LO long-range topology:

$$V(\mathbf{r}_{13}, \mathbf{r}_{23}) = \int \frac{d\mathbf{q}_1}{(2\pi)^3} \int \frac{d\mathbf{q}_2}{(2\pi)^3} e^{i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{i\mathbf{q}_3 \cdot \mathbf{r}_{23}} V(\mathbf{q}_2, \mathbf{q}_3)$$

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the preregularized regulator reads accordingly:

$$R_{\text{prereg}}(\mathbf{q}_2, \mathbf{q}_3) = \int \frac{d\mathbf{r}_{13}}{(2\pi)^3} \int \frac{d\mathbf{r}_{23}}{(2\pi)^3} e^{-i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{-i\mathbf{q}_3 \cdot \mathbf{r}_{23}} \frac{\left(1 - e^{-r_{13}^2/R^2}\right)^6 \left(1 - e^{-r_{23}^2/R^2}\right)^6}{Q(r_{13}^2)Q(r_{23}^2)}$$

consider a N2LO long-range topology:

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For $Q(r^2) = r^2$ all integrals are finite and can be calculated without subtraction!



Fits of N2LO 3NF (semi-local momentum space) to 3H binding energy





 currently exploring behaviour of NN+3N interactions in many-body systems

• Which set of observables for fixing 3NF couplings cD and cE?

$$(1+\delta_R)t = \frac{K/G_V^2}{f_V \langle \mathbf{F} \rangle^2 + f_A g_A^2 \langle \mathbf{GT} \rangle^2}$$

 $\langle \mathbf{F} \rangle = \langle^{3} \mathbf{H} \mathbf{e} \| \sum_{i=1}^{3} \tau_{i}^{+} \|^{3} \mathbf{H} \rangle \qquad \langle \mathbf{G} \mathbf{T} \rangle = \frac{1}{g_{A}} \langle^{3} \mathbf{H} \mathbf{e} \| \sum_{i=1}^{3} \mathbf{J}_{i,1\mathbf{b}}^{+} + \sum_{i < j} \mathbf{J}_{ij,2\mathbf{b}}^{+} \|^{3} \mathbf{H} \rangle$

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$$\mathbf{J}_{i,1b}^{+} = g_{A} \tau_{i}^{+} \sigma_{i}$$

$$\mathbf{J}_{12,2b}^{+} = -\frac{g_{A}}{2F_{\pi}^{2}} \frac{1}{\mathbf{k}^{2} + m_{\pi}^{2}} \Big[4c_{3} \mathbf{k} \, \mathbf{k} \cdot (\tau_{1}^{+} \sigma_{1} + \tau_{2}^{+} \sigma_{2})$$

$$+ \Big(c_{4} + \frac{1}{4m_{N}} \Big) (\tau_{1} \times \tau_{2})^{+} \mathbf{k} \times [(\sigma_{1} \times \sigma_{2}) \times \mathbf{k}]$$

$$- \frac{i}{8m_{N}} (\tau_{1} \times \tau_{2})^{+} (\mathbf{p}_{1} + \mathbf{p}_{1}' - \mathbf{p}_{2} - \mathbf{p}_{2}') (\sigma_{1} - \sigma_{2}) \cdot \mathbf{k} \Big]$$

$$- 2id_{1} (\tau_{1}^{+} \sigma_{1} + \tau_{2}^{+} \sigma_{2}) - id_{2} (\tau_{1} \times \tau_{2})^{+} (\sigma_{1} \times \sigma_{2})$$

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Question:

How should the current be regularized when computing <F> and <GT>? What are the uncertainties related to this choice?

Using for the 2b-currents: $f_{\Lambda}^{\text{loc}}(\mathbf{p},\mathbf{p}') = \exp\left[-(\mathbf{p}-\mathbf{p}')^4/\Lambda^4\right]$ Gazit et al., PRL103, 102502 (2009)



Klos et al., EPJA A53, 168 (2018), EPJA A54, 76 (2018) [erratum]

- varying Lambda in current leads to a significant range in cD
- how to choose the cutoffs consistently in currents and interaction (continuity equation?)

Implementation of nuclear current operators for applications in few- and many-body frameworks

 \mathbf{k}'

 \mathbf{k}

Warmup problem: single-nucleon point charge operator

$$\langle \mathbf{k}'_i | \rho | \mathbf{k}_i \rangle = \begin{bmatrix} G_E^p \frac{(1 + \tau_z)}{2} + G_E^n \frac{(1 - \tau_z)}{2} \end{bmatrix} \delta(\mathbf{k}'_i - \mathbf{k}_i - \mathbf{Q})$$
with $G_E^p = 1$ and $G_E^n = 0$

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$$\mathbf{Q}$$
with $G_E^p = 1$ and $G_E^n = 0$

charge form factors of 2H and 3H:



Implementation of nuclear current operators for applications in few- and many-body frameworks

 \mathbf{k}'

k

Q

Spin vector current

$$\langle \mathbf{k}'_i | \mathbf{j} | \mathbf{k}_i \rangle = \frac{1}{2M} G_M \left[i \boldsymbol{\sigma} \times \mathbf{Q} \right] \, \delta(\mathbf{k}'_i - \mathbf{k}_i - \mathbf{Q})$$

Spin form factors 3H:

$$\langle \psi, M_J = 1/2 | j_+ | \psi, M_J = -1/2 \rangle = \langle \psi, M_J = -1/2 | j_- | \psi, M_J = 1/2 \rangle$$



Summary

• power counting in 3NF sector sensitively depends on regularization scheme

 developed efficient framework to compute general 3NF for basically arbitrary regulators in coordinate or momentum space

• adaption of framework to calculation of **general one- and two-body nuclear currents** (benchmarks finished for first simple test cases)

Summary

• power counting in 3NF sector sensitively depends on regularization scheme

 developed efficient framework to compute general 3NF for basically arbitrary regulators in coordinate or momentum space

• adaption of framework to calculation of **general one- and two-body nuclear currents** (benchmarks finished for first simple test cases)

Open Questions

- more systematic study of power counting uncertainty estimates in fits?
- how to avoid fine tuning/cancellations in fits at N3LO?
- inclusion of new observables in fitting frameworks
 - \star identify relevant nuclear currents
 - * regularization of currents?