

Accurate analysis method to detect rare events in VIP-2

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21 September 2022



X-Rays

Ejected K-shell electron

Incident radiation

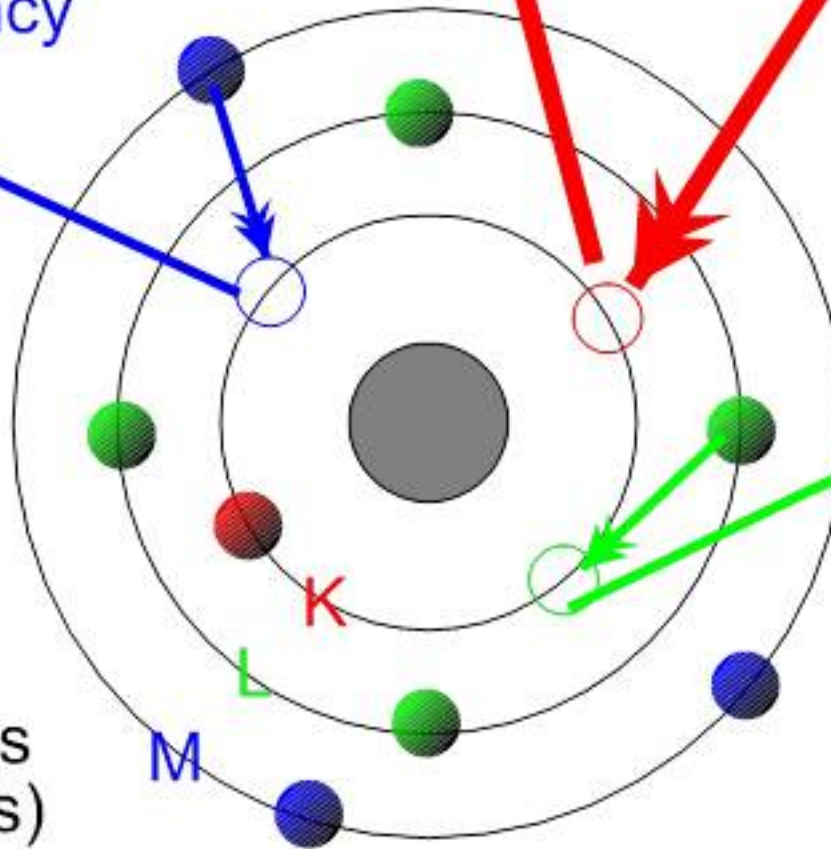
M-shell electron
fills vacancy

L-shell electron
fills vacancy

K_{β} x-ray emitted

K_{α} x-ray emitted

Shells
(orbits)



Searching for PEP violation

Why Fermi-Dirac and Bose-Einstein are distinct?

- ◆ **Green's general quantum field:** paronic particles
 - ◆ Order 1: fermionic/bosonic fields
 - ◆ Order >1: parafermionic/parabosonic fields
 - ◆ Messiah-Greenberg Super-Selection: no fermion/boson decays into parafermion/paraboson (and vice-versa)
 - ◆ **Paronic:** a mixture of fermionic/bosonic and parafermionic/parabosonic states
- ◆ **Non-Commutative Quantum Gravity**
 - ◆ **θ -Poincaré:** distortion of Lorentz symmetry (visible in a two identical particles system)

Both break the anti-/symmetric commutativity with an amplitude β .

In a system of two fermions (i.e., two electrons),

PEP is violated with a probability of $\beta^2/2$

[See Fabrizio's Talk for more details]

VIP-2 GOAL

searching **V**iolation of **P**auli Exclusion Principle



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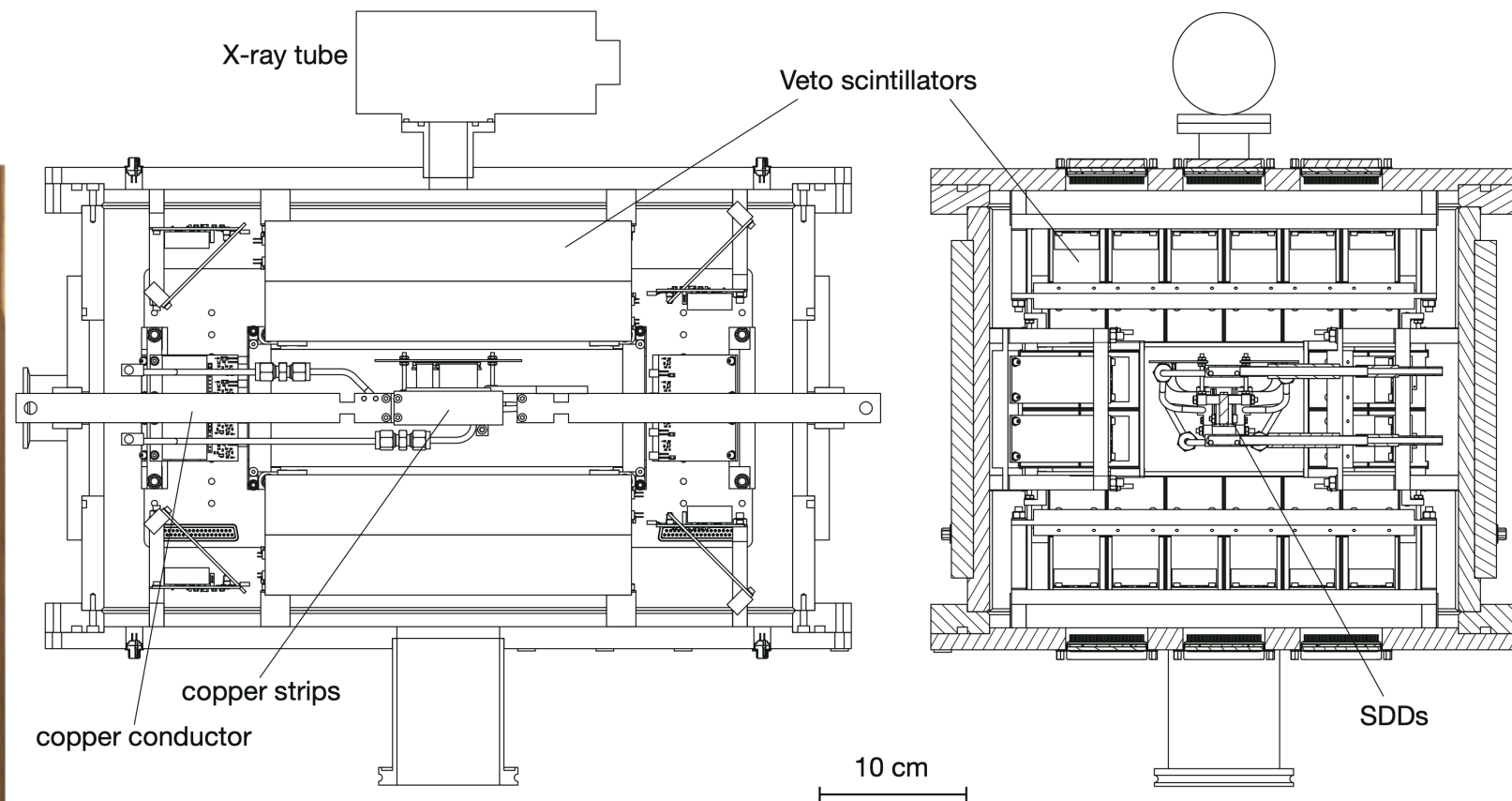
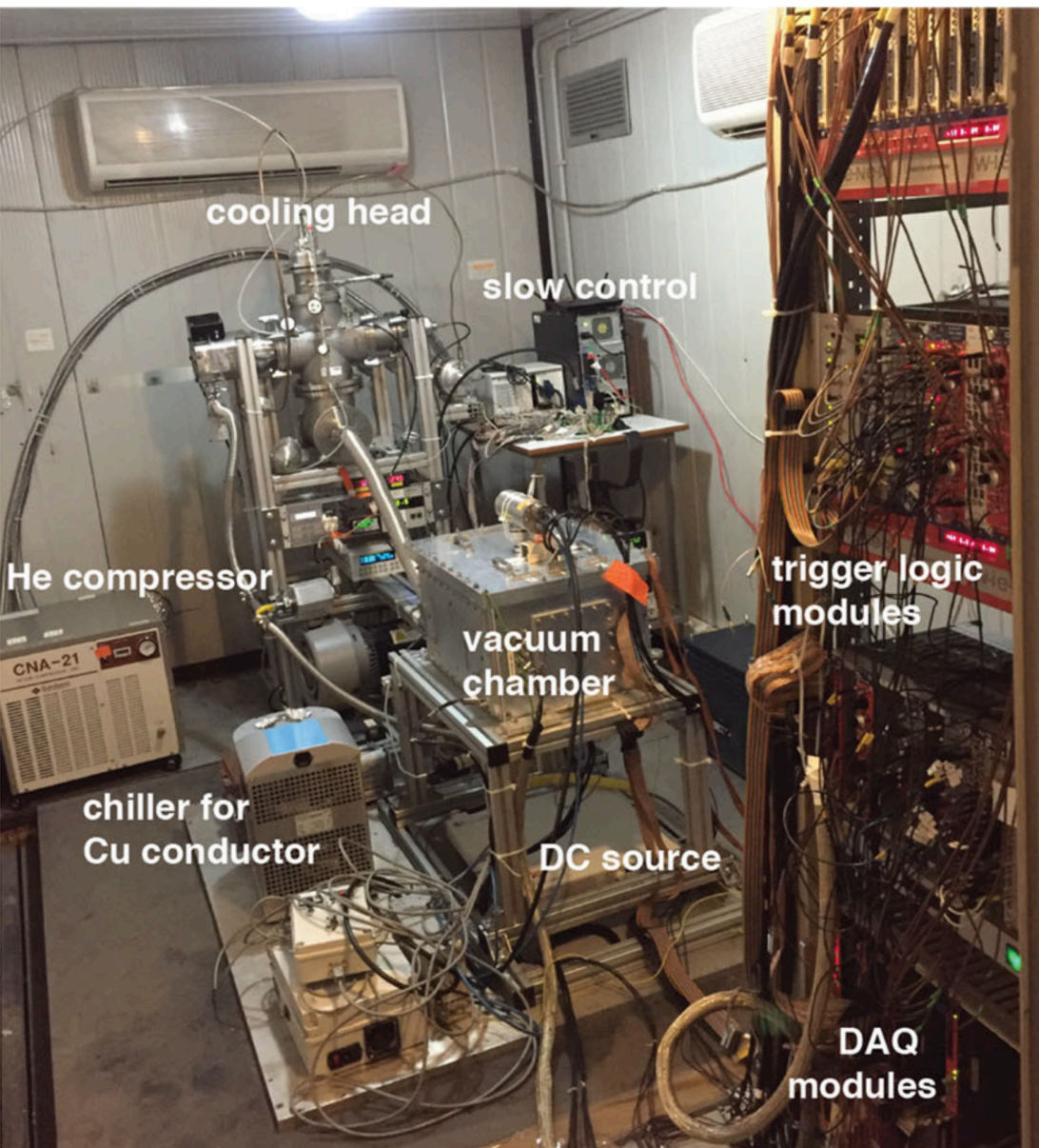
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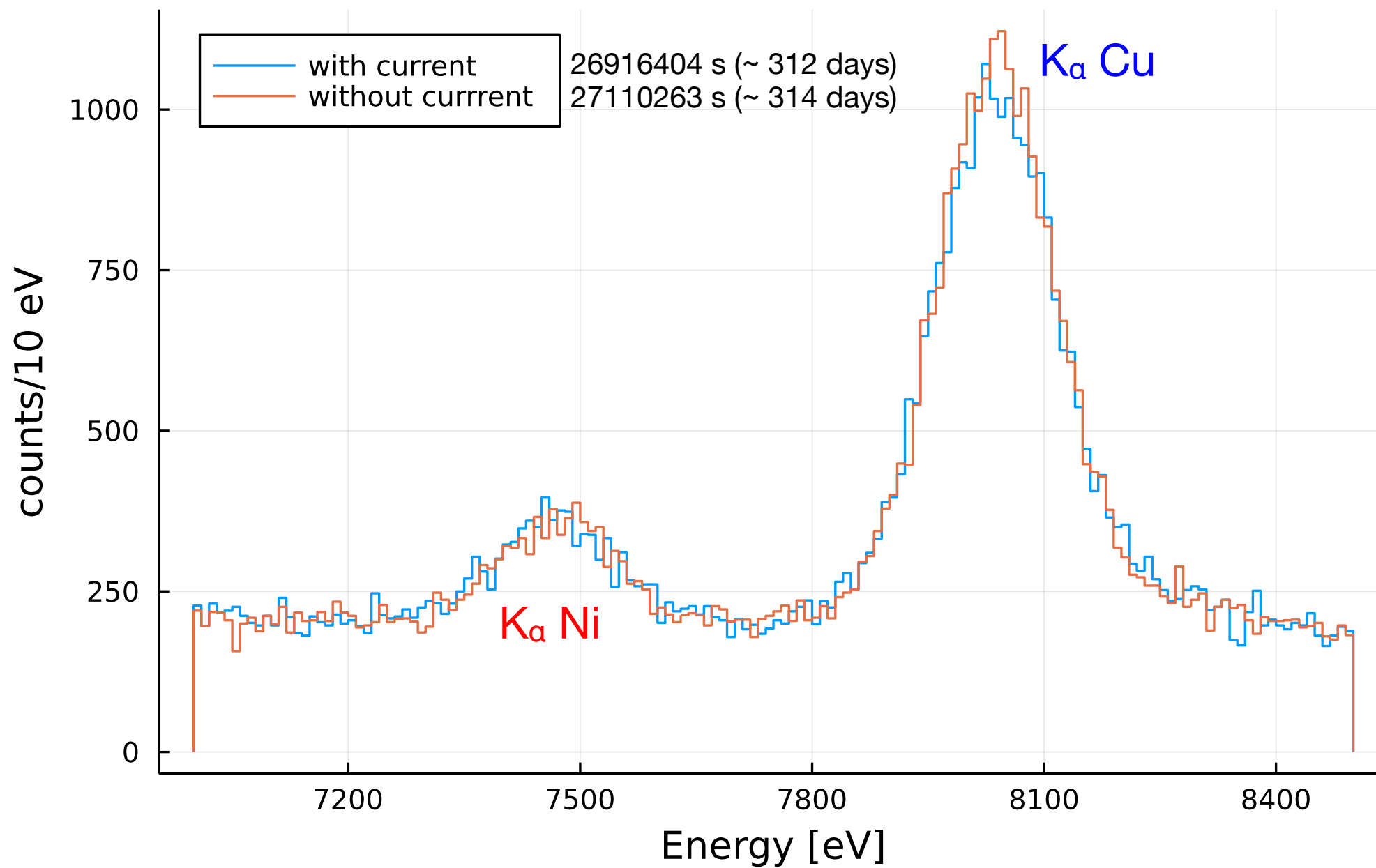
VIP-2



- ◆ **Target:** Copper strips
- ◆ **WITHOUT CURRENT** configuration: regime case (stable states: background)
- ◆ **WITH CURRENT** configuration (180 A): dynamic case (PEP violation through electron capture)
- ◆ **SDD:** 32 detectors by SDDs, stably kept @ -170_{-0}^{+1} °C even with the current in Cu
- ◆ **@LNGS Underground** (beneath Gran Sasso Mountain – IT): ~1400 m of rock shielding

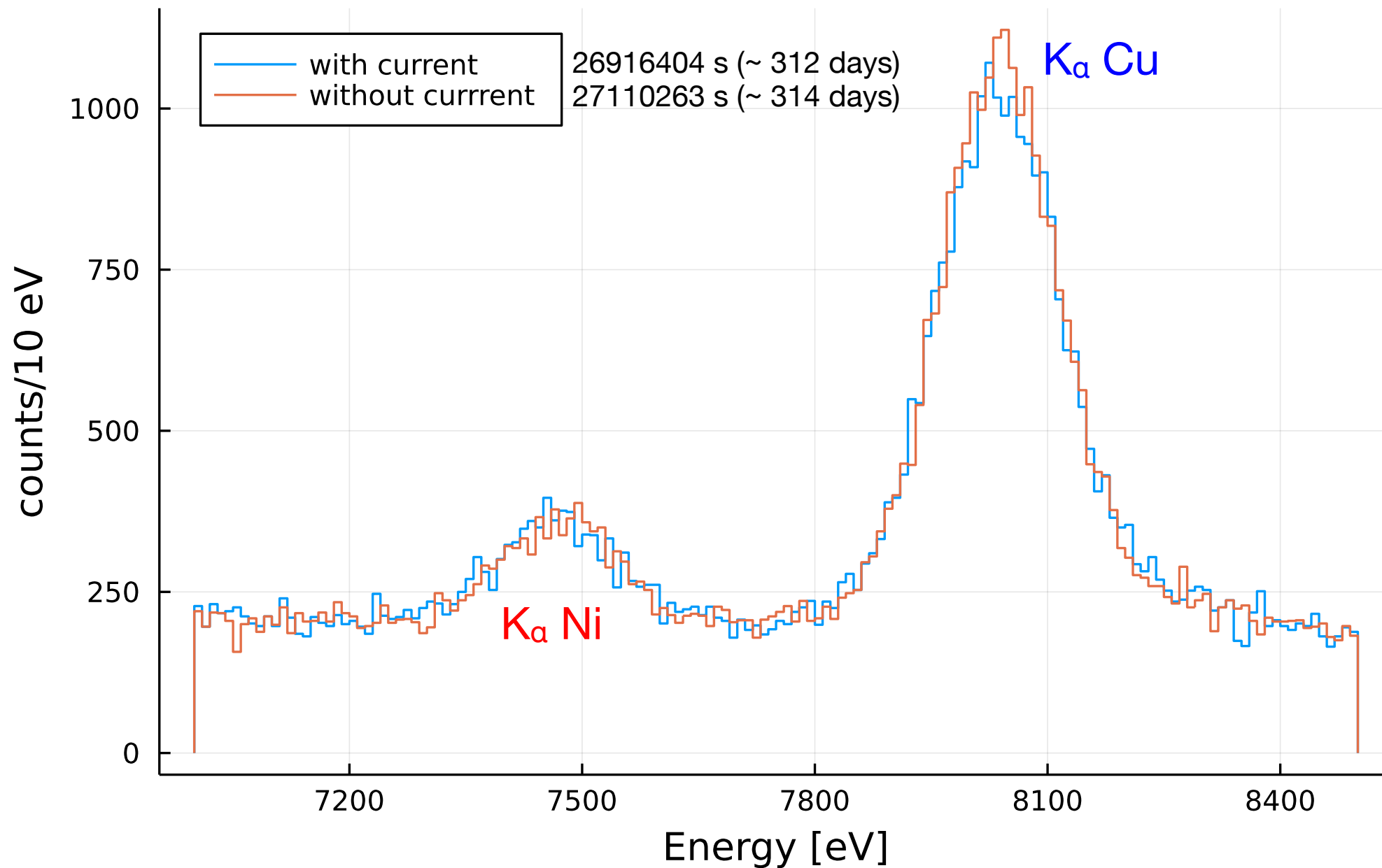


Data model



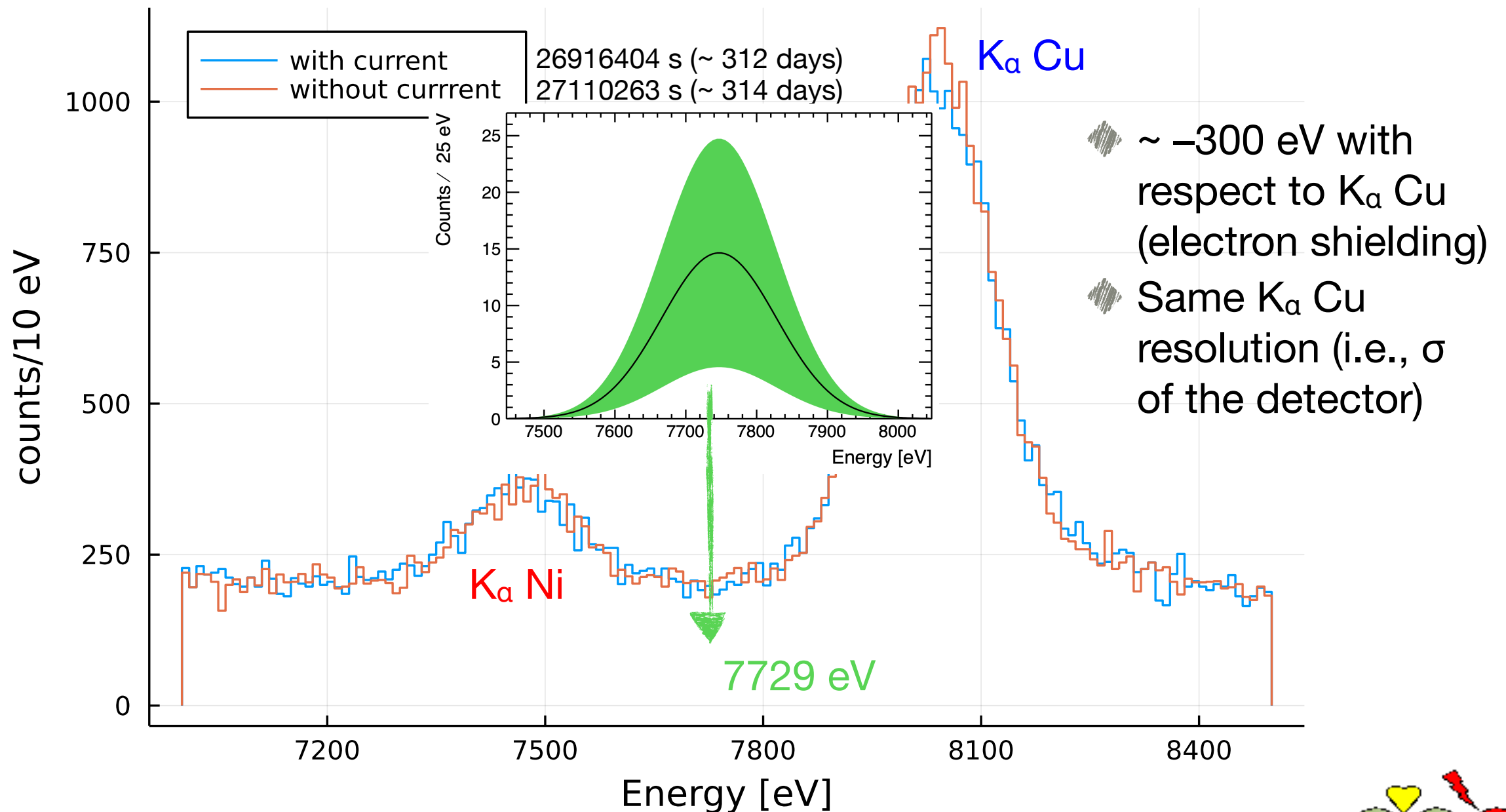
Data model

$$\mathcal{F}^{woc}(\theta, y) = y_1 \times Ni(\theta_1, \theta_2) + y_2 \times Cu(\theta_3, \theta_4) + y_3 \times pol_1(\theta_5)$$



Data model

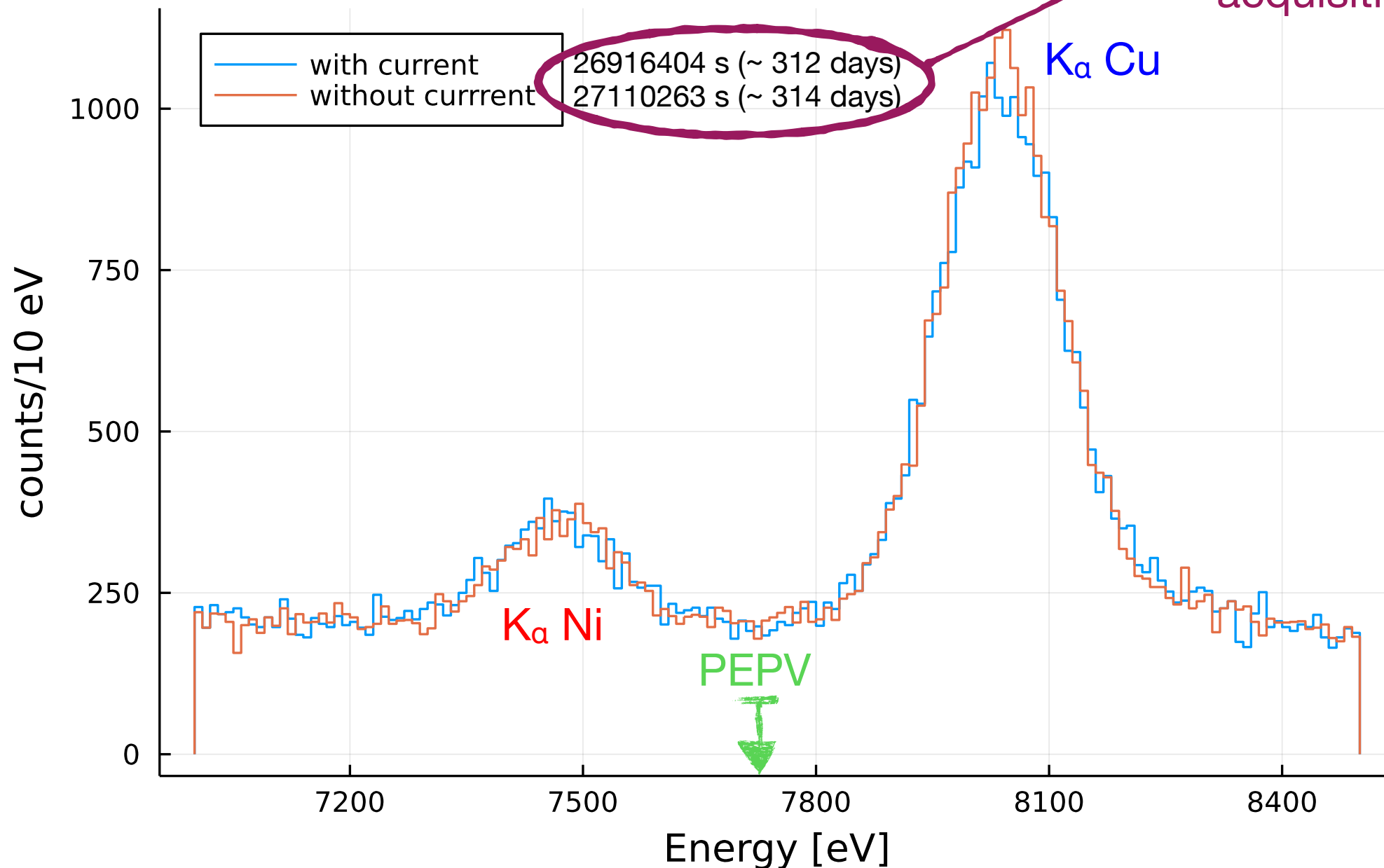
$$\mathcal{F}^{wc}(\theta, y, \mathcal{S}) = y_1 \times Ni(\theta_1, \theta_2) + y_2 \times Cu(\theta_3, \theta_4) + y_3 \times pol_1(\theta_5) + \mathcal{S} \times PEPV(\theta_4)$$



Data Likelihood

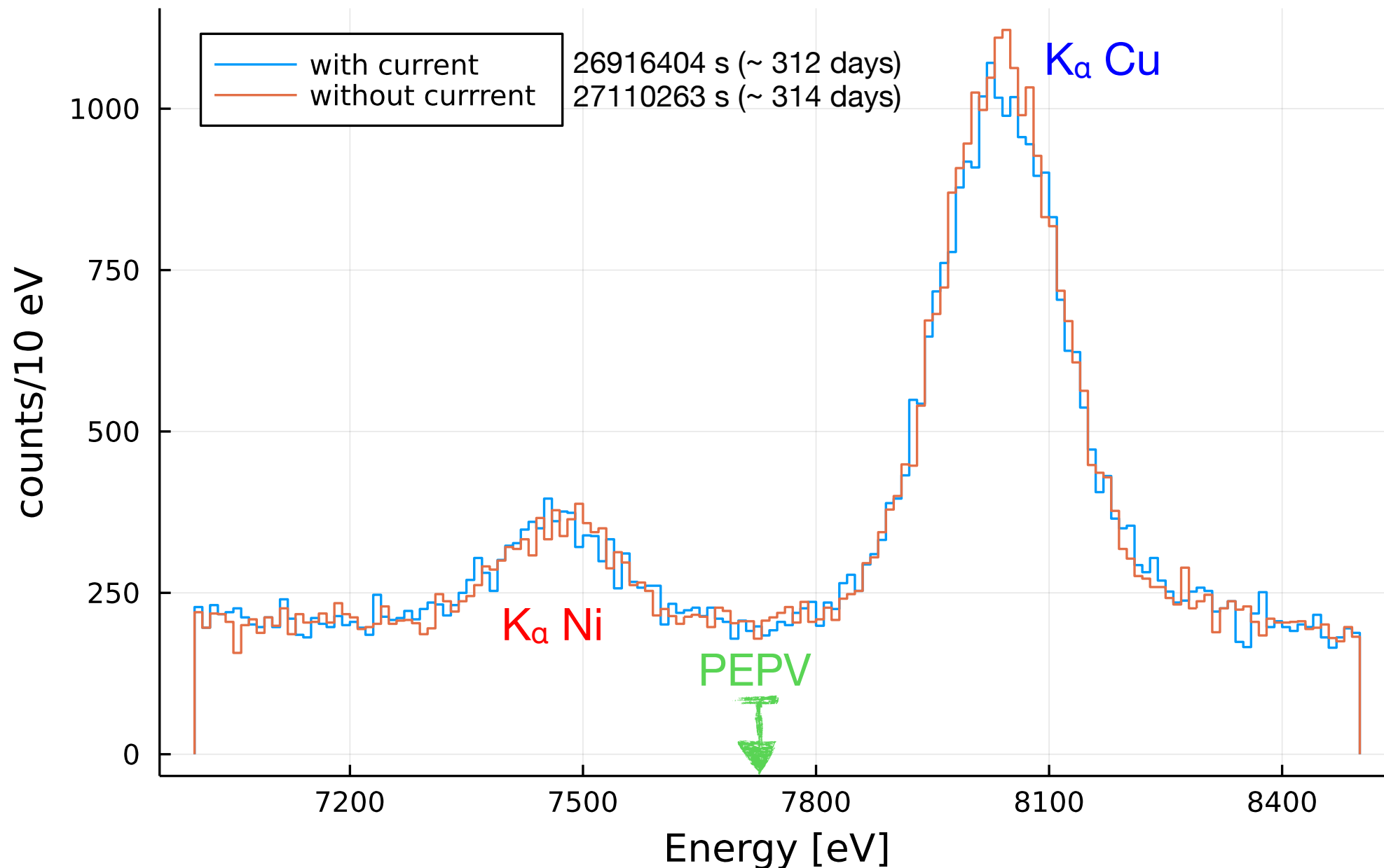
$$\mathcal{L}(\mathcal{D}^{wc}, \mathcal{D}^{woc} | \theta, y, \mathcal{S}) = \text{Poiss}(\mathcal{D}^{wc} | \mathcal{F}^{wc}(\theta, y, \mathcal{S})) \times \text{Poiss}(\mathcal{D}^{woc} | \mathcal{F}^{woc}(\theta, y \times \mathcal{R}))$$

[mind: $\text{Poiss}(\mathcal{D} | \mathcal{F}) = \frac{\mathcal{F}^{\mathcal{D}}}{\mathcal{D}!} e^{-\mathcal{F}}$, \mathcal{D} are data, \mathcal{F} is the model]



Bayesian approach

$$p(\theta, y, \mathcal{S} | \mathcal{D}^{wc}, \mathcal{D}^{woc}) = \frac{\mathcal{L}(\mathcal{D}^{wc}, \mathcal{D}^{woc} | \theta, y, \mathcal{S})p(\theta, y, \mathcal{S})}{\int d\theta dy \mathcal{L}(\mathcal{D}^{wc}, \mathcal{D}^{woc} | \theta, y, \mathcal{S})p(\theta, y, \mathcal{S})}$$



Bayesian approach

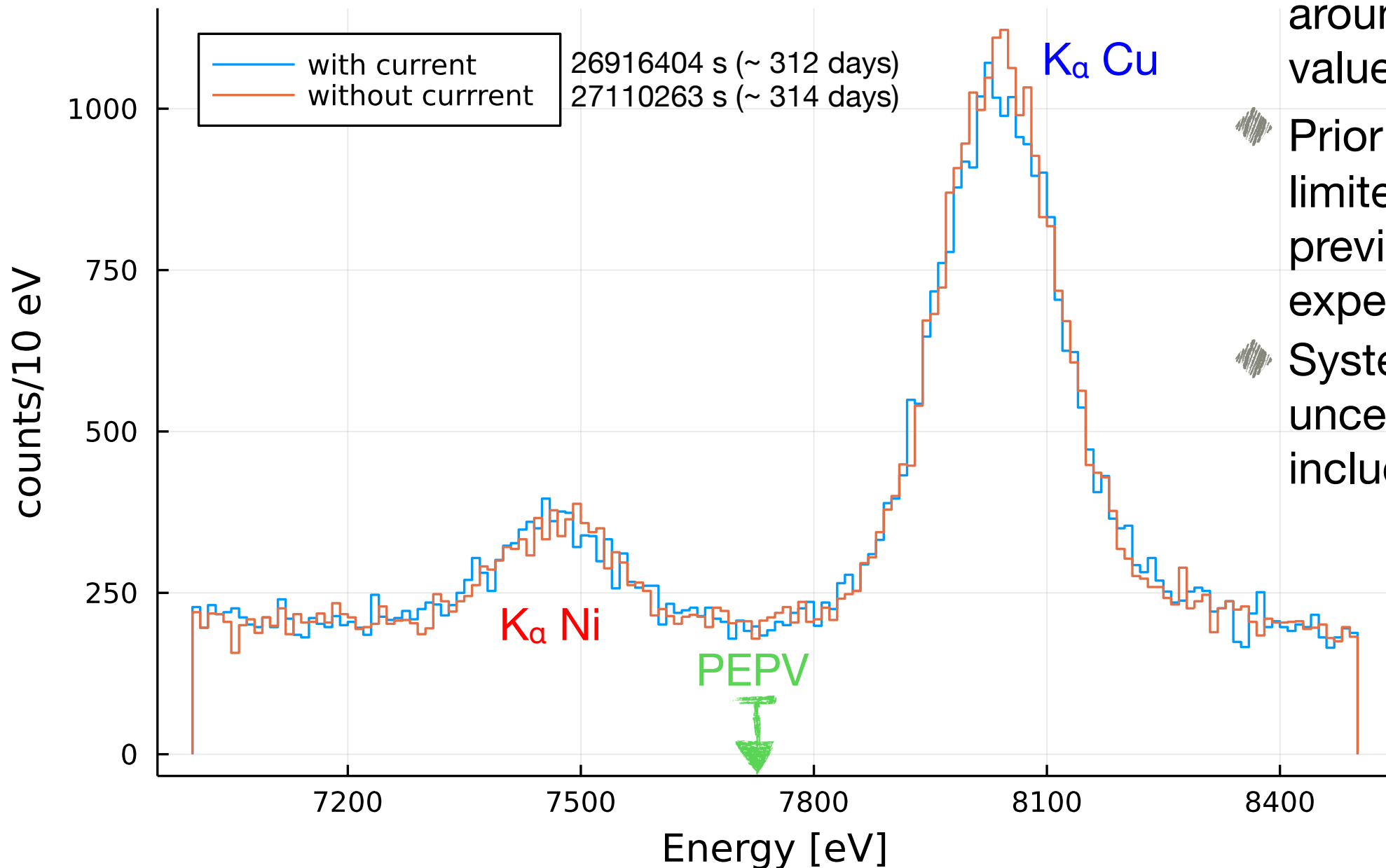
Posterior

$$p(\theta, y, \mathcal{S} | \mathcal{D}^{wc}, \mathcal{D}^{woc}) = \frac{\mathcal{L}(\mathcal{D}^{wc}, \mathcal{D}^{woc} | \theta, y, \mathcal{S}) p(\theta, y, \mathcal{S})}{\int d\theta dy \mathcal{L}(\mathcal{D}^{wc}, \mathcal{D}^{woc} | \theta, y, \mathcal{S}) p(\theta, y, \mathcal{S})}$$

Priors of θ and y are Gaussians: statistical fluctuations around known values

Prior of \mathcal{S} is flat, limited from previous experiments

Systematic uncertainties included

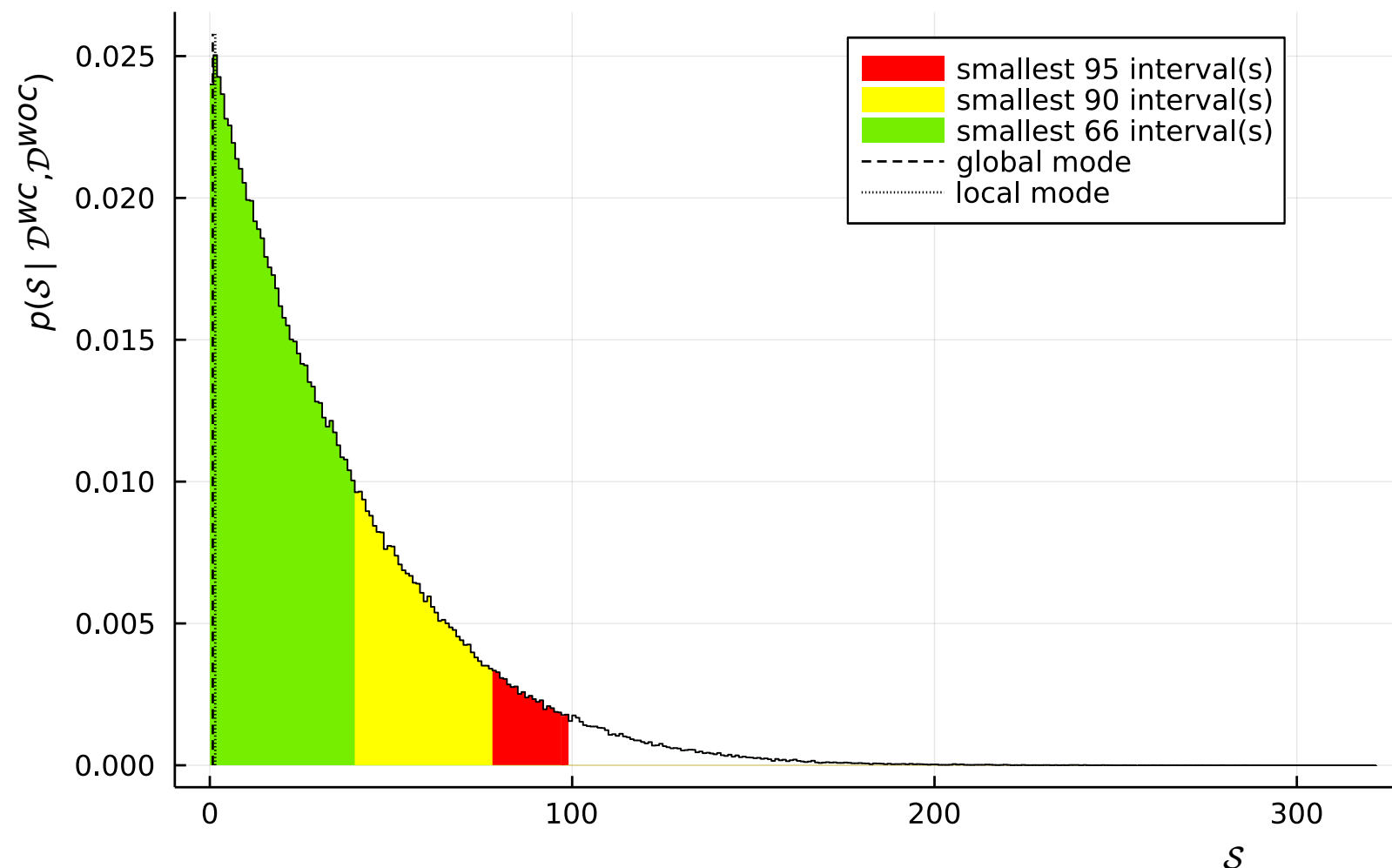


Bayesian result

(marginalized Posterior)

$$p(\mathcal{S} | \mathcal{D}^{wc}, \mathcal{D}^{woc}) = \int p(\theta, y, \mathcal{S} | \mathcal{D}^{wc}, \mathcal{D}^{woc}) d\theta dy$$

Integrals with Markov Chain Monte Carlo method



Modified frequentist CLs

one-sided test statistic

$$t_{\mathcal{S}} = -2 \ln \Lambda(\mathcal{S}) = -2 \ln \frac{\mathcal{L}(\hat{\theta}, \hat{y}, \mathcal{S})}{\mathcal{L}(\hat{\theta}, \hat{y}, \hat{\mathcal{S}})}$$



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Profile Likelihood;

\mathcal{L} now includes multiplicative penalties given by experimental uncertainties: i.e., the priors in the Bayesian

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$t_{\mathcal{S}}$ distribution, given \mathcal{S}

$$p_{\mathcal{S}} = \int_{t_{obs}}^{\infty} f(t_{\mathcal{S}} | \mathcal{S}) dt_{\mathcal{S}}$$

$t_{\mathcal{S}}$ of observed \mathcal{S}

$$CLs = \frac{p_{\mathcal{S}}}{1 - p_0} < 1 - \text{C.L.} \quad (\text{i.e., } 90\% \text{ C.L.} \Rightarrow CLs < 0.1)$$

background case (i.e., $\mathcal{S} = 0$)

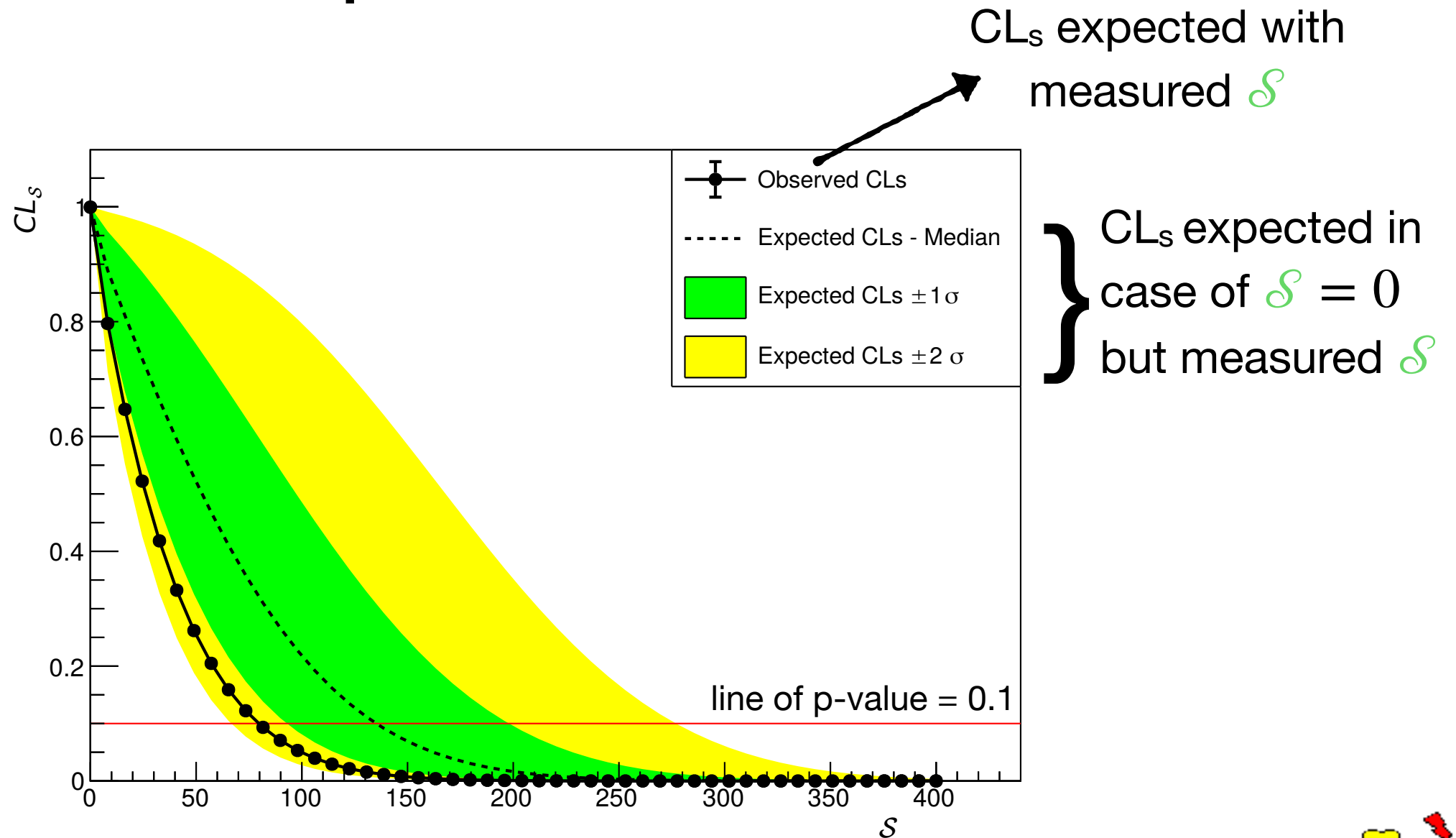


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Computation with RooFit

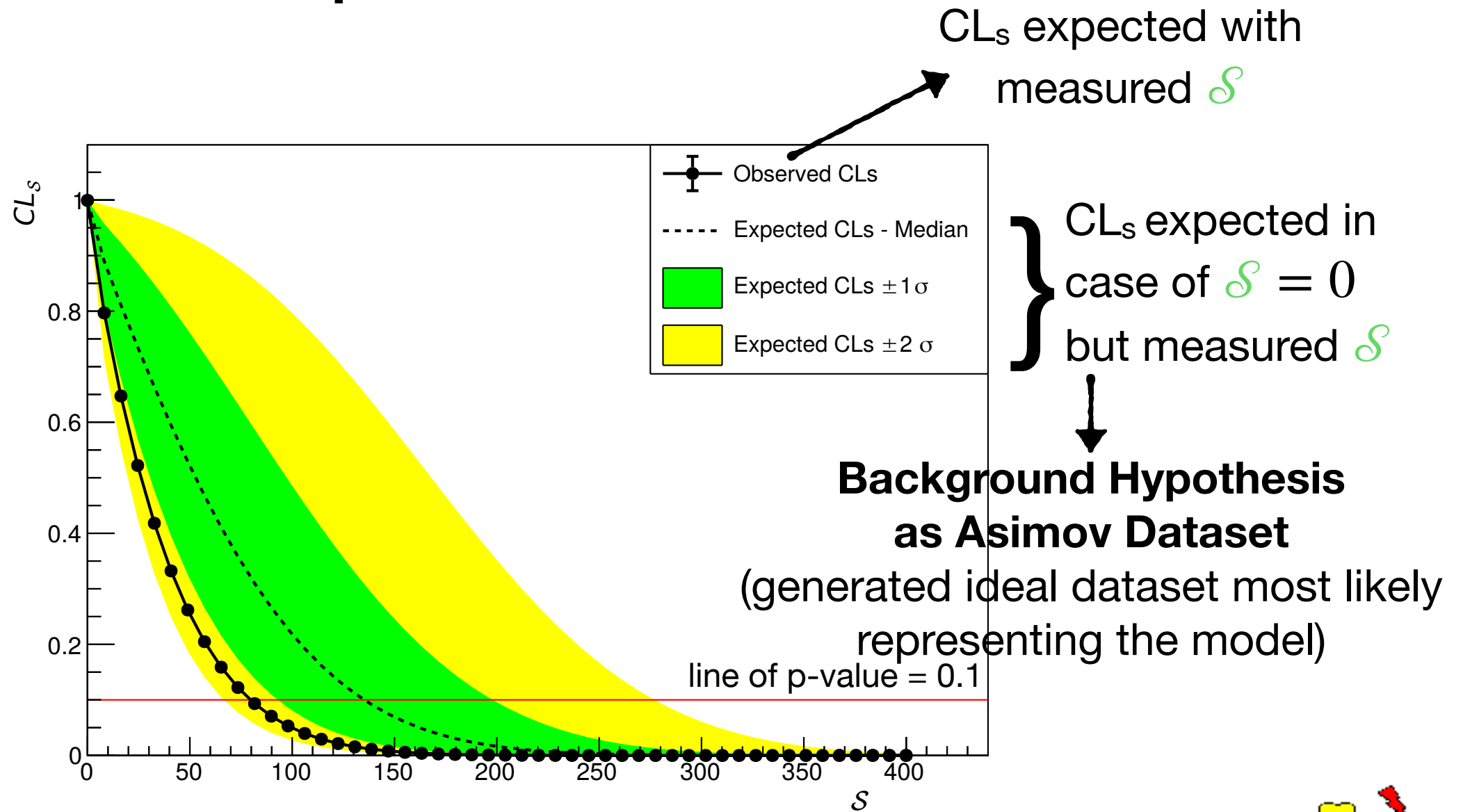


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Computation with RooFit



From \mathcal{S} to $\beta^2/2$

$$N_x \simeq \frac{\beta^2}{2} \cdot N_{\text{new}} \cdot \frac{N_{\text{int}}}{10} \cdot 7.25 \times 10^{-2}$$



From \mathcal{S} to $\beta^2/2$

This is our \mathcal{S} !

$$\rightarrow N_x \approx \frac{\beta^2}{2} \cdot N_{\text{new}} \cdot \frac{N_{\text{int}}}{10} \cdot 7.25 \times 10^{-2}$$



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Newly injected electrons!

$$\sum_i^{\text{runs}} I_i \Delta t_i / e \quad (= I \Delta t / e \text{ for simplicity})$$



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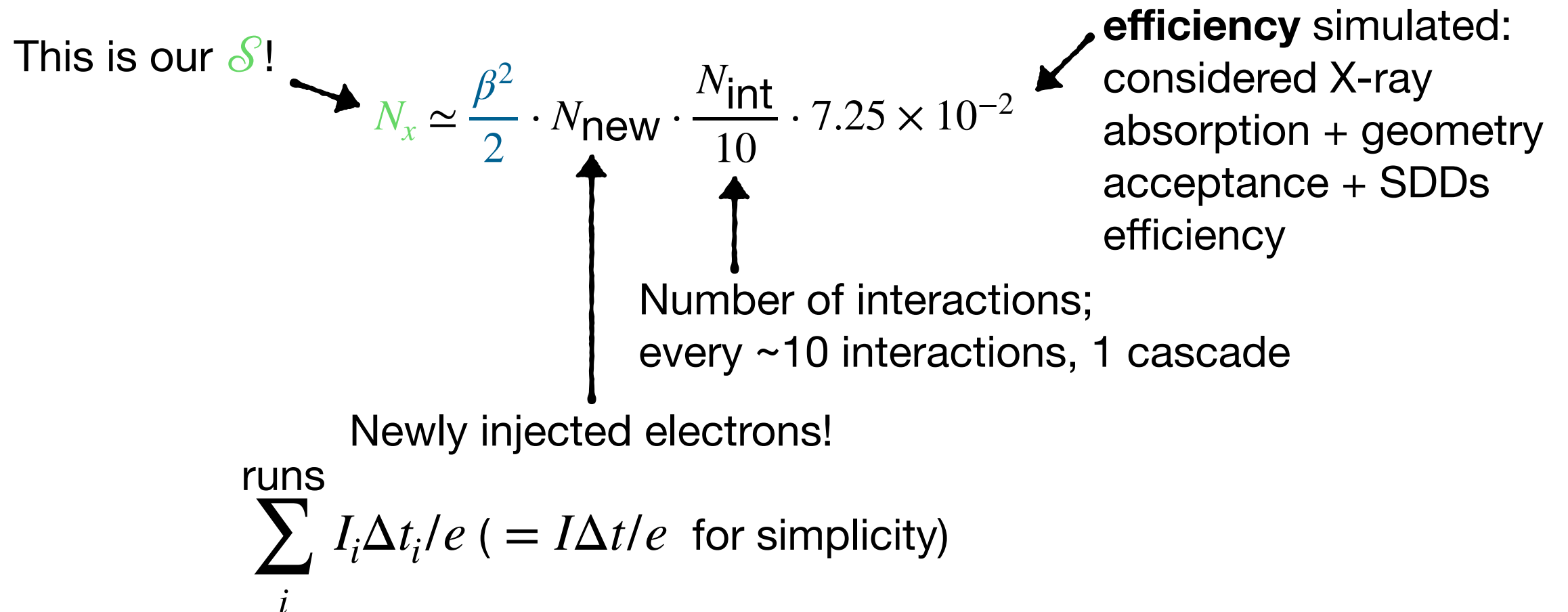
Number of interactions;
every ~10 interactions, 1 cascade

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efficiency simulated:
considered X-ray
absorption + geometry
acceptance + SDDs
efficiency

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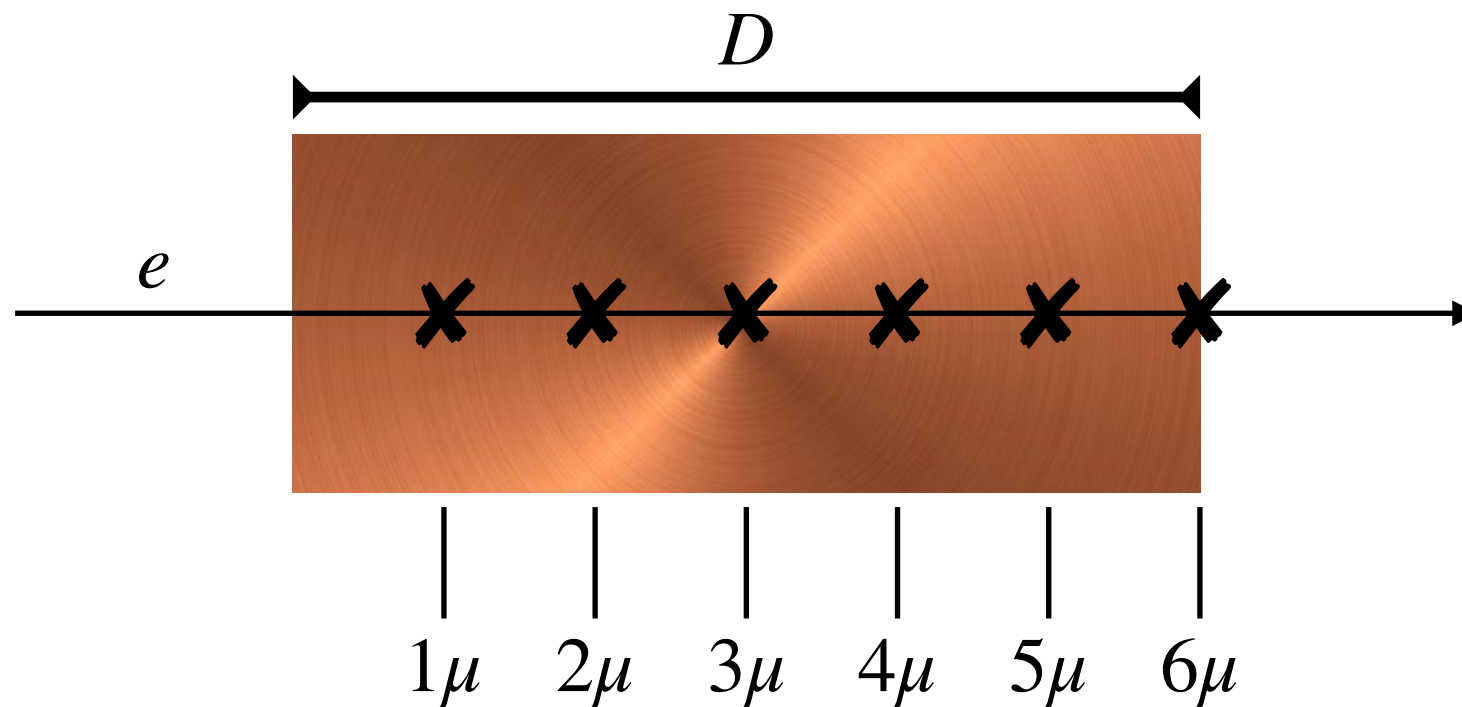
$$\frac{\beta^2}{2} \simeq \mathcal{S} \cdot \frac{10}{N_{\text{int}}} \cdot \frac{e}{I \Delta t} \cdot \frac{1}{7.25 \times 10^{-2}}$$

N_{int} is the normalization that decides the order of magnitude of $\beta^2/2$

Let's discuss e -atoms interaction Models!



N_{int} by Linear Scattering



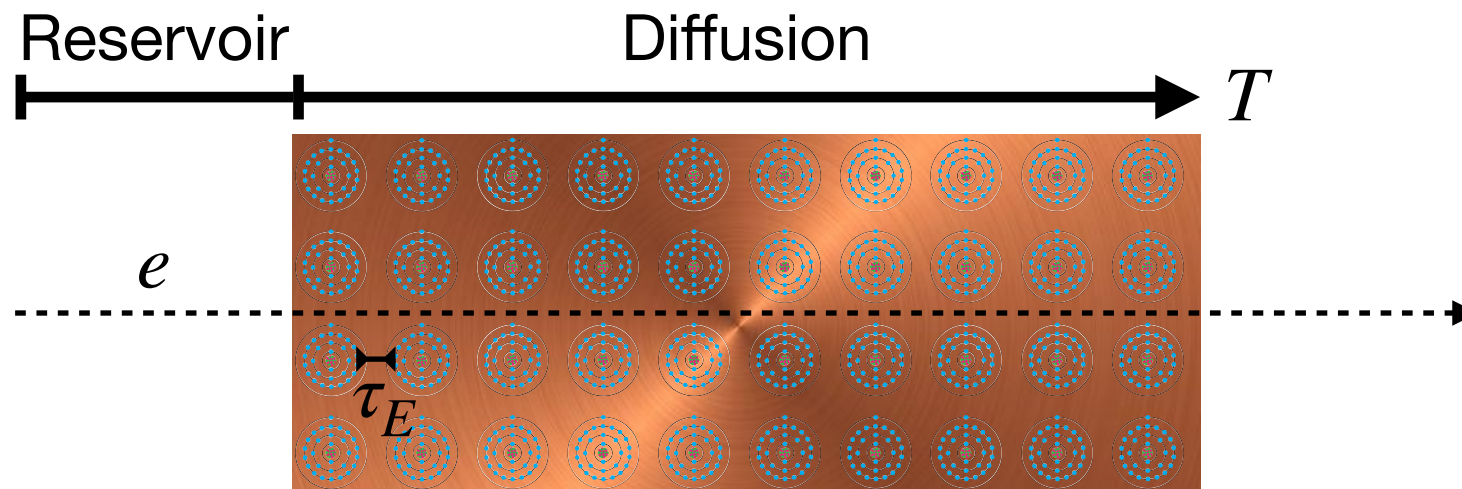
Through Copper Resistance,
we know the average interaction length μ

$$N_{\text{int}} = D/\mu \simeq 1.95 \times 10^6$$

$$\Rightarrow \frac{\beta^2}{2} \lesssim 10^{-31}$$



N_{int} by Close Encounters



Through Diffusion-Transport theory and Copper atomic density, we know:

- the average time τ_E on atomic encounter for a diffused electron
- the average time T of target crossing by an electron

$$N_{\text{int}} = T/\tau_E \simeq 4.29 \times 10^{17}$$

$$\Rightarrow \frac{\beta^2}{2} \approx 10^{-43}$$



Outlook

Bayesian

- ▶ Well established: excellent for low statistical signals
- ▶ Systematic uncertainty is the combination of different priors for the various factors

CL_s

- ▶ Models with little or no sensitivity to the null hypothesis, e.g., if the data fluctuate very low relative to the expectation of the background-only hypothesis: the lower/upper limit might be anomalously low; more robust compared to the classic p-value
- ▶ Sensible to small parameter fluctuations

N_{int}

- ◆ **Linear Scattering:** due to phonons and lattice irregularities
 - ✓ Safest hypothesis
 - ✗ Largely underestimation of how many interactions an electron does
- ◆ **Close Encounters:** a more realistic model of e -atom encounters, but still approximated
 - ✓ 12 order of magnitudes larger than Linear Scattering!
- ▶ **This is the key element to improve the measurement!**



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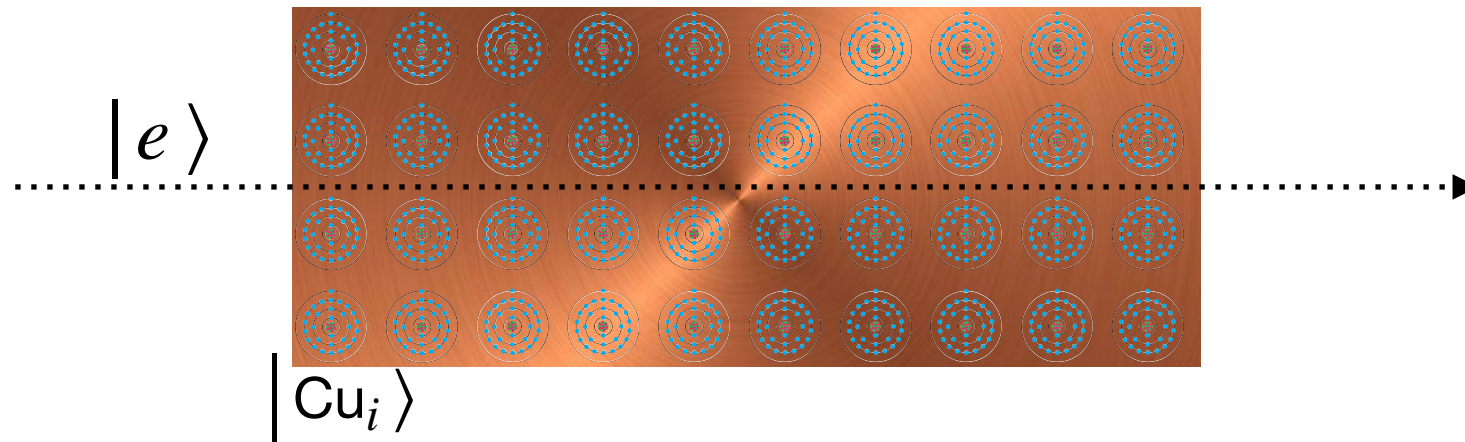
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THANK YOU



BACKUPS

TO DO: a quantum N_{int} ?



How many interactions between Cu atomic and electron fields occur?

