





*Ab-initio* description of monopole resonances in light- and mediummass nuclei

Methods, uses and new preliminary results

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### Outline



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$$H \left| \Psi_{\nu} \right\rangle = E_{\nu} \left| \Psi_{\nu} \right\rangle$$



Input Hamiltonian

$$(H)\Psi_{\nu}\rangle = E_{\nu}(|\Psi_{\nu}\rangle)$$

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Many-body solution

# Global philosophy

The approximate solution must be systematically improvable and approach the exact solution in a well-defined limit.

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QCD

#### Nuclear forces



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- More elementary description
- Complexity in terms of elementary DOF
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- Collective picture
- Phenomena from effective description
- Energy Density Functional
- Collective models



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# X-EFT

- Structure-less Protons and Neutrons
- All nucleons are active
- Systematically improvable
- LEC from data (or simulations)
- Up to A-body forces



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#### Virtually exact methods

Early 2000's

- Factorial scaling
- Monte-Carlo methods and NCSM
- Explicit few-body solution



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**From 2005** Symmetry-conserving ref state MBPT, CC, SCGF, IMSRG

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#### Breakdown for open-shell systems !

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U(1) symmetry breaking Bogoliubov extensions

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Doubly-closed shell

#### From 2005 Symmetry-conserving ref state MBPT, CC, SCGF, IMSRG Singly-open shell From 2011 U(1) symmetry breaking Bogoliubov extensions Doubly-open shell From 2016

SU(2) symmetry breaking

Deformed calculations

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- Ground-state relies on previous calculations
- Excited states from the action of linear operators
- Similar to the Equation of Motion
- Linear system

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### State-specific expansion

Wave operator acting separately on excited states



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- Reduction to bound-state problem
- Numerical inversion issues
- The response should consist of 1/2 broad peaks

Coupled Cluster (CC-LIT)

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Application to <sup>4</sup>He



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- RPA, 2nd-RPA and QRPA (Darmstadt group) (CC-RPA, IMSRG-RPA, IMSRG-2nd-RPA) Limited to spherical systems
- SCGF, RPA with dressed propagators

For closed-shell systems

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- Early works in EDF essentially forgotten
- Large amplitude vibrations (possibly anharmonic)
- Present goal to revive it within ab-initio

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# **Collective excitations**

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### Role of three-body forces

- Systematic effect on the peaks' position
- Crucial aspect in ab-initio
- Different possible treatments



[From R. Trippel, PhD Thesis, Technischen Universität Darmstadt, 2016 ]

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### Chiral order dependence

- **Convergence** wrt the **chiral order** within given family
- Non-negligible dependence on the used fit
- Good agreement with exp for presently used family

[Y. Beaujeault-Taudière, M. Frosini, J.-P. Ebran, T. Duguet, R. Roth, V. Somà, arXiv:2203.13513]

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Quasi-Boson Approximation Pauli principle violation (But this is often asymptomatic)

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- 1. First step towards more sophisticated **Boson Expansions**
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GDR in <sup>16</sup>O



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### Valid motivation for ab-initio RPA !

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30  $\omega$  [MeV]

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shkanov et al.

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Well known issue in Quantum Chemistry

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ExcS correlation

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30  $\omega$  [MeV]

G<sub>0</sub>W vs GW approximation

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- Perturbation theory + PGCM (PGCM-PT) recently formulated
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- Consistent correction





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### Error cancellation for spectroscopy!



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PGCM

Triax. PGCM

PGCM-PT(2)

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Exp

[Frosini et al. , EPJA, 2022]



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#### PGCM promising ab-initio candidate for collective ecxs

[Frosini et al. , EPJA, 2022]



#### First ab-initio calculations of GMR for

- Closed- and open-shell nuclei
- Two complementary methods (PGCM and QRPA(QFAM))

# What to (and not to) expect?

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- Deformation mechanisms on GMR
- Pairing fluctuations on GMR

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### Nuclear structure study



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Shape coexistence Superfluidity **b** T = 0, Deformation

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- Treatment of anharmonicities
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[Blaizot 1995]

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#### Not right there yet

- Discussion about  $K_{\!\varpi}$
- Pairing/isospin effects «fluffiness»



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#### What to (and not to) expect? (From the present talk/study)

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Ni isotopes and more systematics coming soon !

#### Outline





#### Schrödinger equation $H |\Psi_{\nu}\rangle = E_{\nu} |\Psi_{\nu}\rangle$

Schrödinger equation  $H |\Psi_{\nu}\rangle = E_{\nu} |\Psi_{\nu}\rangle$ 

1 Constrained HFB solutions  $|\Phi(r^2, \beta_2)\rangle$ 















Schrödinger-like equation





















Schrödinger equation

$$H |\Psi_{\nu}\rangle = E_{\nu} |\Psi_{\nu}\rangle$$



Schrödinger equation

PGCM

 $H |\Psi_{\nu}\rangle = E_{\nu} |\Psi_{\nu}\rangle$ 

$$|\Psi_{\nu}\rangle \equiv \sum_{r^2,q} f_{\nu}(r^2,q) |\Phi(r^2,q)\rangle$$

r<sup>2</sup> to study GMR q to couple to other modes Symmetry breaking and restoration Variational method



Schrödinger equation PGCM  $|\Psi_{\nu}\rangle \equiv \sum_{r^2,q} f_{\nu}(r^2,q) |\Phi(r^2,q)\rangle$ 

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$$H |\Psi_{\nu}\rangle = E_{\nu} |\Psi_{\nu}\rangle$$
QRPA
$$|\Psi_{\nu}\rangle \equiv Q_{\nu}^{\dagger} |\Psi_{0}\rangle$$
Boson-like excitation operator
QRPA matrix diagonalization

 $_{\rm rs}Q_{\nu}^{\dagger}$ **QFAM** formulation frequencies  $\mathbb C$ 



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Boson-like excitation operators  $Q_{\nu}^{\dagger}$ QRPA matrix diagonalization QFAM formulation frequencies  $\mathbb{C}$ 

#### Pros and Cons



Handle anharmonicities and shape coexistance Select on few collective coordinates Symmetries are restored Computationally expensive Harmonic limit of GCM All coordinates are explored Symmetries are not restored Low computational cost

Schrödinger equation **PGCM**  $|\Psi_{\nu}\rangle \equiv \sum_{r^{2},q} f_{\nu}(r^{2},q) |\Phi(r^{2},q)\rangle$ 

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Total  $H |\Psi_{\nu}\rangle = E_{\nu} |\Psi_{\nu}\rangle$ Schrödinger equation con-field Energy PGCM QRPA 0  $|\Psi_{\nu}\rangle \equiv \sum_{r^{2},q} f_{\nu}(r^{2},q) |\Phi(r^{2},q)\rangle$  $|\Psi_{\nu}\rangle \equiv Q_{\nu}^{\dagger}|\Psi_{0}\rangle$ Boson-like excitation operators  $Q_{\nu}^{\dagger}$ r<sup>2</sup> to study GMR **QRPA** matrix diagonalization **q** to couple to other modes Symmetry breaking and restoration QFAM formulation frequencies C Variational method Pros and Cons



Square

redius

Harmonic limit PGCM points

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Select on few collective coordinates

Symmetries are restored

Computationally expensive

All coordinates are explored

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Low computational cost



Computationally expensiv

First **ab-initio** realization very recently developed

- 1) PGCM (M. Frosini, CEA Saclay)
- 2) QFAM (Y. Beaujeault-Taudière, CEA DAM)



• Studied quantity: monopole strength

$$S_{00}(\omega) \equiv \sum_{\nu} |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \delta(E_{\nu} - E_0 - \omega)$$

- Transition amplitudes: height of peaks
- Energy difference: position of peaks



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- Related moments  $m_k \equiv \int_0^\infty S_{00}(\omega) \, \omega^k \, d\omega$ =  $\sum_{\nu} (E_{\nu} - E_0)^k |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2$ =  $\langle \Psi_0 | \check{M}_k(i, j) | \Psi_0 \rangle$



[Bohigas et al., 1979]

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- Energy difference: position of peaks
- Related moments  $m_k \equiv \int_0^\infty S_{00}(\omega) \,\omega^k \, d\omega$   $= \sum_{\nu} (E_{\nu} - E_0)^k |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \longrightarrow \text{Must know excited states}$   $\equiv \langle \Psi_0 | \check{M}_k(i,j) | \Psi_0 \rangle$ [Bohigas et al., 1979]

200

150

100

50

 $S_{00} \ [fm^4MeV^{-1}]$ 

• Studied quantity: monopole strength

$$S_{00}(\omega) \equiv \sum_{\nu} |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \delta(E_{\nu} - E_0 - \omega)$$

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 $S_{00}$  [fm<sup>4</sup>MeV<sup>-1</sup>]

Complexity is shifted to the operator structure

$$\begin{split} \breve{M}_k(i,j) &\equiv (-1)^i C_i C_j \quad \forall \ k \ge 0 \\ M_k(i,j) &\equiv \frac{1}{2} (-1)^i [C_i, C_j] \quad \text{if} \ k = 2n+1, \ n \in \mathbb{N} \end{split} \qquad \begin{array}{c} C_l &\equiv \begin{bmatrix} H, [H, \dots [H, [H, r^2]] \dots ] \\ I \text{ times} \end{bmatrix} \end{split}$$

Studied quantity: monopole strength

$$S_{00}(\omega) \equiv \sum_{\nu} |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \delta(E_{\nu} - E_0 - \omega)$$

- Transition amplitudes: height of peaks
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- Related moments  $m_k \equiv \int_0^\infty S_{00}(\omega) \, \omega^k \, d\omega$ 0 10 30 20 40 50 0  $\omega$  [MeV]  $= \sum_{\nu} (E_{\nu} - E_{0})^{k} |\langle \Psi_{\nu} | r^{2} | \Psi_{0} \rangle|^{2} \longrightarrow \text{Must know excited states}$  $= \langle \Psi_{0} | \check{M}_{k}(i, j) | \Psi_{0} \rangle \longrightarrow \text{Ground state only} \text{IBobia}$  $\rightarrow$  Ground state only [Bohigas et al., 1979]

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 $m_0$ 

 $m_0$ 

 $S_{00}$  [fm<sup>4</sup>MeV<sup>-1</sup>]

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Encode the main physical features of the strength 
$$\begin{split} \bar{E}_{1} &= \frac{m_{1}}{m_{0}} \qquad \sigma^{2} = \frac{m_{2}}{m_{0}} - \left(\frac{m_{1}}{m_{0}}\right)^{2} \geq 0 \end{split}$$

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 $S_{00}$  [fm<sup>4</sup>MeV<sup>-1</sup>]

Complexity is shifted to the operator structure

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#### First comparison ever of the two approaches !

Derived and implemented in an ab-initio PGCM code

• Studied quantity: monopole strength

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$$C_{l} \equiv \begin{bmatrix}H, [H, \dots [H, [H, r^{2}]] \dots] \\ I \text{ times} \end{bmatrix}$$

$$I \text{ times}$$

|          | m0    | m1   | m1/m0 |
|----------|-------|------|-------|
| QRPA     | 358,2 | 8532 | 23,82 |
| QFAM     | 358,2 | 8532 | 23,82 |
| PGCM sum | 356,4 | 8105 | 22,74 |
| PGCM gs  | 380,6 | 8543 | 22,45 |

• Studied quantity: monopole strength

$$S_{00}(\omega) \equiv \sum_{\nu} |\langle \Psi_{\nu} | r^2 | \Psi_0 \rangle|^2 \delta(E_{\nu} - E_0 - \omega)$$

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Complexity is shifted to the operator structure

16O

$$C_l \equiv [H, [H, \dots [H, [H, r^2]] \dots]]$$

$$l \text{ times}$$

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 $S_{00}$  [fm<sup>4</sup>MeV<sup>-1</sup>]

| 24 | M | g |
|----|---|---|
|----|---|---|

|          | m0    | m1     | m1/m0 |
|----------|-------|--------|-------|
| QFAM     | 852,4 | 17.441 | 20,46 |
| PGCM sum | 880,0 | 17.049 | 19,37 |
| PGCM gs  | 960,1 | 17.760 | 18,50 |
|          |       |        |       |

Sum rules are important for the extraction of experimental data (MDA)

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Usually computed within **EDF** theory

Standard assumption :  $H(r) \equiv H[\rho(r)] = T + V[\rho(r)]$ 

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Momentum-independent interactions

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Standard assumption :

$$H(r) \equiv H[\rho(r)] = T + V[\rho(r)]$$

Analytic expression

$$m_1 = \frac{1}{2} \langle \Psi | [r^2, [H(r), r^2]] | \Psi \rangle$$

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# Analytic expression

Momentum-independent interactions

$$m_1 = \frac{1}{2} \langle \Psi | [r^2, [H(r), r^2]] | \Psi \rangle$$

$$= \frac{1}{2} \langle \Psi | [r^2, [T, r^2]] | \Psi \rangle = \frac{2\hbar^2}{m} A \langle \Psi | r^2 | \Psi \rangle$$

Sum rules are important for the extraction of experimental data (MDA)

 $H(r) \equiv H[\rho(r)] = T + V[\rho(r)]$ 

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# Analytic expression

$$m_{1} = \frac{1}{2} \langle \Psi | [r^{2}, [H(r), r^{2}]] | \Psi \rangle$$
$$= \frac{1}{2} \langle \Psi | [r^{2}, [T, r^{2}]] | \Psi \rangle = \frac{2\hbar^{2}}{m} A \langle \Psi | r^{2} | \Psi \rangle$$

#### Has this relevant consequences ? Ab-initio evaluation of commutators

Momentum-independent interactions

Sum rules are important for the extraction of experimental data (MDA)



# Has this relevant consequences ?

Ab-initio evaluation of commutators

Sum rules are important for the extraction of experimental data (MDA)



### Outline



# **Common features**

#### PGCM and QFAM have **consistent numerical settings**

- One-body spherical harmonic oscillator basis
  - e<sub>max</sub> = 10
  - ħω = 20 MeV
- Chiral two-plus-three-nucleon in-medium interaction
  - T. Hüther, K. Vobig, K. Hebeler, R. Machleidt and R. Roth, "Family of chiral twoplus three-nucleon interactions for accurate nuclear structure studies", *Phys. Lett. B*, 808, 2020
  - M. Frosini, T. Duguet, B. Bally, Y. Beaujeault-Taudière, J.-P. Ebran and V. Somà, "In-medium k-body reduction of n-body operators", *The European Physical Journal A*, *57*(4), 2021
- Only monopole strength is addressed
- The PGCM wavefunction explores the  $\beta_2$  and  $r^2$  collective coordinates (quadrupolar coupling)





Benchmark on existing spherical QRPA code



Difficulty



Benchmark on existing spherical QRPA code

Monopole Strength



 $\begin{bmatrix} & & & & & \\$ 

 $\omega$  [MeV]

Keruikr\_\_\_

- Single spherical harmonic energy minimum
- Exact QRPA/QFAM superposition

Difficulty

Benchmark on existing spherical QRPA code



Results

- Single spherical harmonic energy minimum
- Exact QRPA/QFAM superposition
- Excellent QFAM/PGCM agreement
- o Harmonic approximation clearly valid



Difficulty

Benchmark on existing spherical QRPA code

![](_page_131_Figure_4.jpeg)

Results

- Single spherical harmonic energy minimum
- Exact QRPA/QFAM superposition
- Excellent QFAM/PGCM agreement
- Harmonic approximation clearly valid
- No coupling with quadrupolar vibrations

![](_page_131_Figure_11.jpeg)

![](_page_132_Figure_1.jpeg)

![](_page_133_Figure_1.jpeg)

![](_page_134_Figure_1.jpeg)

![](_page_135_Figure_1.jpeg)

![](_page_136_Figure_1.jpeg)

(1) [Dowie et al., 2020]

![](_page_137_Figure_1.jpeg)

(1) [Dowie et al., 2020]

![](_page_138_Figure_1.jpeg)

![](_page_139_Figure_1.jpeg)

![](_page_140_Figure_1.jpeg)

![](_page_140_Figure_2.jpeg)

![](_page_141_Figure_1.jpeg)

![](_page_142_Figure_1.jpeg)

![](_page_143_Figure_1.jpeg)






























iThemba, Bahini 2021

- 1. PGCM superior to QRPA
- 2. Experiments useful and promising
- 3. Data are **not unambiguous**





















# Comparison to experiment



1. PGCM superior to QRPA, i.e. coupling to quadrupole deformation/fluctuations captured

- 2. Experimental data in doubly open-shell nuclei very useful and promising
- 3. Data are not unambiguous, i.e. better data would be beneficial

#### Outline



#### **Conclusions and Perspectives**

First **ab-initio** systematic description of GMR

Choose physics according to selected coordinates

No limitation on the nucleus choice

#### Plan of the complete study

- 🗹 🛛 Static quadrupolar deformation
- Z Coupling to quadrupolar vibrations
- 🗹 Shape isomers
- Theoretical comparison of moment computation
- Hamiltonian uncertainty through different chiral EFT orders
- Pairing: isospin dependence and coupling to pairing vibration
- Bubble structure ( <sup>34</sup>Si and <sup>36</sup>S )
  - Nuclei of current experimental interest ( <sup>68</sup>Ni and <sup>70</sup>Ni )

[ACTAR TPC]

# Thanks for the attention



#### Pairing effects in <sup>20</sup>O



In QRPA another mode seems to be important !

