Search for jet quenching in small systems

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QGP in small collision systems?

(d) CMS N \geq 110, 1.0GeV/c<p_<3.0GeV/c



- QGP-like signatures in high-multiplicity pp and pA
- How do QGP signatures that we see in large collision systems evolve when decreasing system size?
- Jet quenching is necessary consequence of a hot and dense fireball. Can we see evidence of it?

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Jet quenching observables





• Yield suppression relative to min. bias $pp \rightarrow energy$ transport out-of-cone

$$R_{\rm AA}^{h,j}(p_T,y) = \frac{1}{\langle T_{\rm AA} \rangle} \frac{(1/N_{\rm ev}) \ dN_{\rm AA}^{h,j}/dp_T dy}{d\sigma_{pp}^{h,j}/dp_T dy}$$

measurement of inclusive suppression R_{AA} requires Glauber scaling \rightarrow

- limited precision of $\langle T_{\scriptscriptstyle AA} \rangle$ for centrality biased events
- Glauber model does not account for conservation laws, geometry information smeared by fluctuations
- not defined in high-multiplicity pp collisions
- Jet substructure modification
- Jet deflection → dijet acoplanarity

J.P. Blaizot and L. McLerran, PRD 34, 2739 (1986)



Minimum bias R_{pPb} for π^0 in p-Pb @ $\sqrt{s_{NN}}$ = 8 TeV



ALICE, PLB 827 (2022) 136943

$$R_{\rm pPb} = \frac{1}{A_{\rm Pb}} \frac{{\rm d}^2 \sigma_{\rm pPb}}{{\rm d} p_{\rm T} {\rm d} y} \bigg/ \frac{{\rm d}^2 \sigma_{\rm pp}}{{\rm d} p_{\rm T} {\rm d} y}$$

Data disfavor more than 1% relative energy loss or an induced $p_{\rm T}$ shift larger than 100 MeV in the range 10-20 GeV/c

Suppression for $p_{\tau} < 10$ GeV/*c* described by

- NLO calculations using nPDFs EPPS16 [10] [Eskola et al. EPJC 77 (2017) 163] nCETQ15 [Kovarik et al. PRD 93 (2016) 085037]
- CGC-based calculations [Lappi et al. PRD 88 (2013) 114020]
- parton energy loss in cold nuclear matter with fully coherent energy loss (FCEL)

[Arleo et al. JHEP 09 (2020) 190]

Q_{pPb} for h⁺⁻ and heavy flavor hadrons in p-Pb



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$R_{_{XA}}$ of π^{0} in EA biased collisions by PHENIX



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Corrlation of hard processes and soft particle production in pp by ATLAS



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Prospects for OO run at LHC

Small system $\langle N_{\rm ch} \rangle_{\rm OO} \approx 2 \langle N_{\rm ch} \rangle_{\rm p-Pb}$ with AA geometry

$$R_{\rm AA}^{h,j}(p_T,y) = \frac{1}{\langle T_{\rm AA} \rangle} \frac{(1/N_{\rm ev}) \ dN_{\rm AA}^{h,j}/dp_T dy}{d\sigma_{pp}^{h,j}/dp_T dy}$$

 $\begin{array}{l} \langle {\cal T}_{_{AA}} \rangle \mbox{ nuclear overlap function depends on} \\ \mbox{ soft physics of tot. inel. pp Xsec. and } \langle {\sf N}_{_{\rm coll}} \rangle \\ \Rightarrow \mbox{ MB provides better precision} \end{array}$

$$R_{AA, \min bias}^{h,j}(p_T, y) = \frac{1}{A^2} \frac{d\sigma_{AA}^{h,j}/dp_T dy}{d\sigma_{pp}^{h,j}/dp_T dy}$$



Projection of hadron *R*_{AA} **for min bias OO**

Luminosities used in the projection : OO $\sqrt{s_{NN}} = 6.37 \text{ TeV}$ $L_{OO} = 1 \text{ nb}^{-1}$ pp $\sqrt{s} = 5.02 \text{ TeV}$ $L_{pp} = 3 \text{ pb}^{-1}$



OO run is planned in 2024

ALICE-PUBLIC-2021-004

Calculation which assumes no energy loss and which accounts just for nuclear PDFs

Calculations which assume energy loss models together with nuclear PDFs [Huss et al. arXiv 2007.13754]

ALICE projection:

data points follow a mean energy loss model

- In the range up to 50 GeV/c:
- statistical precision < 1.5%
- systematic precision 4–6%
 - interpolation error $\leq 3\%$
 - cross section normalization 3%

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- other systematics 2-4%

Measurement is potentially sensitive to the effect

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Flow of high- p_{T} **particles in p-Pb from ATLAS**

ATLAS, Eur. Phys. J. C 80 (2020) 73



Flow of jet fragments in p-Pb by ALICE





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Jet fragmentation in p-Pb from ATLAS



- Event activity measured with ZDC
- No modification in away-side
- Excess in near-side particle production
- Behavior of I_{pPb} is reproduced by PYTHIA ANGANTYR

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Measurement of jet particle yield in pp, p-Pb and Pb-Pb by ALICE

- Correlation of **8-15 GeV/c leading particle 4-6 GeV/c associated particle** both in $|\eta| < 0.8$
- Per trigger yield corrected for UE estimated in tranverse region relative to the leading particle

$$I_{pp,p-Pb,Pb-Pb} = \frac{Y^{pp,p-Pb,Pb-Pb} - Y^{pp,p-Pb,Pb-Pb}_{TS}}{Y^{pp \text{ min.bias}} - Y^{pp \text{ min.bias}}_{TS}}$$





Transverse

 $60^\circ < |\Delta \varphi| < 12$

 $-\Delta \phi$

Leading-particle

Toward $|\Delta \varphi| < 60^{\circ}$

Away $|\Delta \varphi| > 120^{\circ}$

 $\Delta \phi$

Transverse

 $60^\circ < |\Delta \phi| < 12$

Measurement of jet particle yield in pp, p-Pb and Pb-Pb by ALICE



Absence of away side yield suppression for pp and pPb \Rightarrow absence of jet quenching

Longitudinal fragmentation of jets in high-multiplicity pp events in ALICE



HM event activity selection:

5x larger multiplicity in V0 detector w.r.t min. bias 0-0.1% event activity percentile

HM event selection \rightarrow softer jet fragmentation

This is consistent with larger portion of jets comming from NLO processes

Search for jet quenching in p-Pb with h+jet correlations in ALICE

coil (GeV/c)



TT{X,Y} means X< p_{T,trig} < Y GeV/c

ALICE, PLB 783 (2018) 95

 $\Delta_{
m recoil}$

- Event activity measured by ZDC
- Jets recoiling from high- p_{T} trigger hadron (TT)
- Data-driven statistical approach to remove recoil-jet yield uncorrelated to TT including MPI

$$ZDC$$
rigger hadron (TT)
Deach to remove
to TT including MPI
$$\frac{et}{d\eta}\Big|_{p_{T,trig} \in TT\{12,50\}} = CRef \cdot \frac{1}{N_{trig}} \frac{d^2 N_{jet}}{dp_{T,jet}^{ch} d\eta}\Big|_{p_{T,trig} \in TT\{6,7\}}$$

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ALICE p-Pb $\sqrt{s_{NN}} = 5.02 \text{ TeV}$

Anti- k_{τ} charged jets, R = 0.4

 $-0.43 < y_{\tau\tau}^* < 1.36; -0.03 < y_{iet}^* < 0.96$

TT(12 50)

0-20% ZNA

 $\pi - \Delta \omega < 06$

Hadron-jet observables and T_{AA}

$$\frac{1}{N_{\rm trig}^{\rm AA}} \frac{{\rm d}^2 N_{\rm jet}^{\rm AA}}{{\rm d} p_{\rm T,jet}^{\rm ch} {\rm d} \eta_{\rm jet}} \Big|_{p_{\rm T,trig} \in {\rm TT}} = \left(\frac{1}{\sigma^{\rm AA \to h+X}} \cdot \frac{{\rm d}^2 \sigma^{\rm AA \to h+jet+X}}{{\rm d} p_{\rm T,jet}^{\rm ch} {\rm d} \eta_{\rm jet}} \right) \Big|_{p_{\rm T,h} \in {\rm TT}}$$

In case of no nuclear effects

$$\frac{1}{N_{\text{trig}}^{\text{AA}}} \frac{d^2 N_{\text{jet}}^{\text{AA}}}{d p_{\text{T,jet}}^{\text{ch}} d \eta_{\text{jet}}} \Big|_{p_{\text{T,trig}} \in \text{TT}} = \left(\frac{1}{\sigma^{\text{pp} \to \text{h} + X}} \cdot \frac{d^2 \sigma^{\text{pp} \to \text{h} + jet + X}}{d p_{\text{T,jet}}^{\text{ch}} d \eta_{\text{jet}}} \right) \Big|_{p_{\text{T,h}} \in \text{TT}} \times \frac{T_{\text{AA}}}{T_{\text{AA}}}$$

- This coincidence observable is self-normalized, no requirement of T_{AA} scaling
- No requirement to assume correlation between Event Activity and collision geometry

Limit on energy transport out of *R* = 0.4 in p-Pb



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Search for jet quenching in high multiplicity pp collisions using hadron-jet acoplanarity

- pp minimum bias (MB)
- pp high-multiplicity (HM) : 5x larger multiplicity in V0 detector w.r.t. MB

$$\Delta_{\text{recoil}} (\Delta \varphi) = \frac{1}{N_{\text{trig}}} \frac{dN_{\text{jet}}}{d\Delta \varphi} \Big|_{\text{TT}\{20,30\}} \bigotimes_{p_{\text{t,jet}}^{\text{ch}}} - c_{\text{ref}} \cdot \frac{1}{N_{\text{trig}}} \frac{dN_{\text{jet}}}{d\Delta \varphi} \Big|_{\text{TT}\{6,7\}} \bigotimes_{p_{\text{t,jet}}^{\text{ch}}} + \frac{1}{N_{\text{trig}}} \sum_{q_{\text{ch}}^{\text{ch}}} \frac{dN_{\text{jet}}}{d\Delta \varphi} \Big|_{\text{TT}\{6,7\}} \bigotimes_{p_{\text{t,jet}}^{\text{ch}}} + \frac{1}{N_{\text{trig}}} \sum_{q_{\text{ch}}^{\text{ch}}} \frac{dN_{\text{jet}}}{d\Delta \varphi} \Big|_{\text{TT}\{6,7\}} \bigotimes_{p_{\text{t,jet}}^{\text{ch}}} + \frac{1}{N_{\text{trig}}} \sum_{q_{\text{ch}}^{\text{ch}}} \frac{1}{N_{\text{trig}}} \sum_{q_{\text{ch}}^{\text{ch}}} \frac{dN_{\text{jet}}}{d\Delta \varphi} \Big|_{\text{TT}\{6,7\}} \bigotimes_{p_{\text{t,jet}}^{\text{ch}}} + \frac{1}{N_{\text{trig}}} \sum_{q_{\text{ch}}^{\text{ch}}} \frac{1}{N_{trig}} \sum_{q_{\text{ch}}^{\text{ch}}} \frac$$



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Comparison of hadron-jet acoplanarity with PYTHIA



Quantitative comparison to PYTHIA 8 Monash shows similar suppression pattern

The effect is not due to jet quenching

Use PYTHIA to explore the origin of the effect

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PYTHIA : recoil jet η_{jet} **versus** $p_{T,jet}$





HM events:

- significant bias in distribution of high- p_{τ} recoil jets
- enhancement in forward trigger detector acceptance
- V0A and V0C have asymmetric coverage

PYTHIA: Number of recoil jets versus event activity in ALICE acceptance

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MB, TT{20,30}

HM, TT{20,30} $4 < V0M/\langle V0M \rangle < 9$

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Distrib. of the number of recoil jets above p_{τ} threshold:

- HM trigger suppresses events with 1 hard recoil jet in the ALICE central barrel
- HM trigger enhances multi-jet events in small system



Summary

- Jet quenching in small collision systems is likely to be a small effect
- Jet quenching signatures can be created by event selection biases
 - picking up fluctuations in particle wavefunction when imposing event activity bias
 - NLO processes with multi jet topology in final state, soft fragmentation
 - cold nuclear matter effects, CGC
- Need to understand positive v_2 of high- p_T particles in p-Pb. Would we see the same effect in HM pp?
- New systems comming soon OO, pO



π⁰ production in min. bias p-Pb and pp @ $\sqrt{s_{NN}}$ = 8 TeV



ALICE, PLB 827 (2022) 136943

- Reach up to 200 GeV/c
- pp reference corrected to 8.16 TeV using PYTHIA 8 Monash
- NLO with NNFF1.0 frag. functions [NNPDF EPJC 77 (2017) 516]
- NLO with DSS14 frag. functions [de Florian et al. PRD 91 (2015) 014035]

Two particle correlations to measure yield assocated to jet fragments

- trig. and assoc. particles have same charge, measured in $|\eta| < 0.8$
- signal/background as a function $\Delta \phi, \Delta \eta$



Correlate trigger particle with a hit in Forward Multiplicity Detector $1.7 < \eta_{FMD} < 5.1$ - Non flow removal : subtraction of correlations for high mult. and low mult. events



Factorization: $\Delta V_2(\Delta \varphi, \Delta \eta) = v_2(\Delta \varphi, \Delta \eta) v_2^{\text{FMD}}$



Sum of 5 harmonics



Pb-Pb:

- Jet particle v_2 is compatible with v_2 of charged-particle jets for $p_T^{trig} > 7 \text{ GeV}$
- Both interpreted by pathlength dependnece of energy loss

p-Pb:

- Inclusive v_2 in Pb-Pb is about 1.7 greater w.r.t. p-Pb.
- Nonzero v_2 for jet fragments Initial-state effects (CGC) or final-state scatterings

$\textit{\textbf{R}}_{_{xA}}$ of $\pi^{_{0}}$ in min. bias collisions by PHENIX

PHENIX, arXiv 2111.05756



- $R_{xA} \approx 0.9$ but still compatible with unity within uncertainties
- No system size dependence at high $p_{\tau} \rightarrow$ little or no modification of hard scattering

Event activity selection in pp at $\sqrt{s} = 13$ TeV

- Trigger:
 - Minimum bias (MB) $L_{int} \approx 32 \text{ nb}^{-1}$
 - High multiplicity (HM) $L_{int} \approx 10 \text{ pb}^{-1}$
- Event activity (EA) selection: VOM = VOA + VOC
- HM is 0.1% of MB cross section





Charged-particle jets in MB and event activity biased pp collisions at √s = 13 TeV

arXiv:2202.01548



How does the imposed event activity selection bias the spectrum shape?

Ratios of EA-biased jet p_{T} spectra to MB



- Event activity (EA) bias affects the shape mostly for $p_{\rm T.ch\,iet}$ < 20 GeV/c
- Bias on high-EA causes increase of jet yield per event
 May arise from increase in average number of hard scatterings per event

Self-normalized jet yield versus self-normalized multiplicity

- Jets with p^{ch}_{T, jet}>9 GeV/c follow non-linear trend similar to J/ψ in midrapidity Phys. Lett. B 810 (2020) 135758 and heavy-flavor electrons and prompt D
- Electrons from W decay follow linear trend
- Overshoot of the trend by PYTHIA at high charged-particle multiplicities



Ratios of jet p_{T} spectra with different *R*



MB ratio of p_{T} -differential cross section spectra: independent of \sqrt{s} EA-selected ratio of spectra:

- small R : independent of EA
- large R : hint of EA dependence