

# A lattice approach to hot neutrons & cold atoms

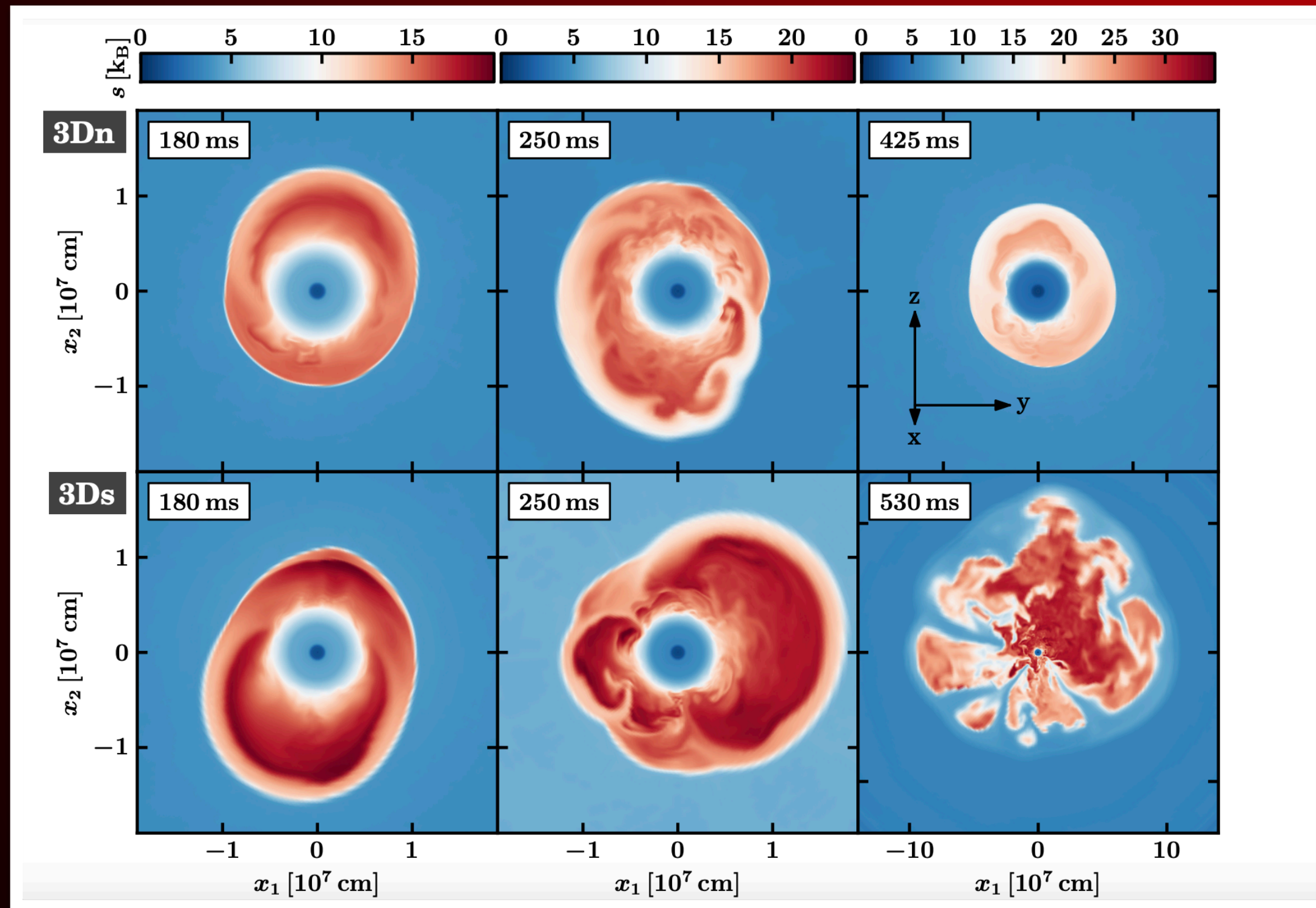
**Neill Warrington**

Institute for Nuclear Theory, University of Washington

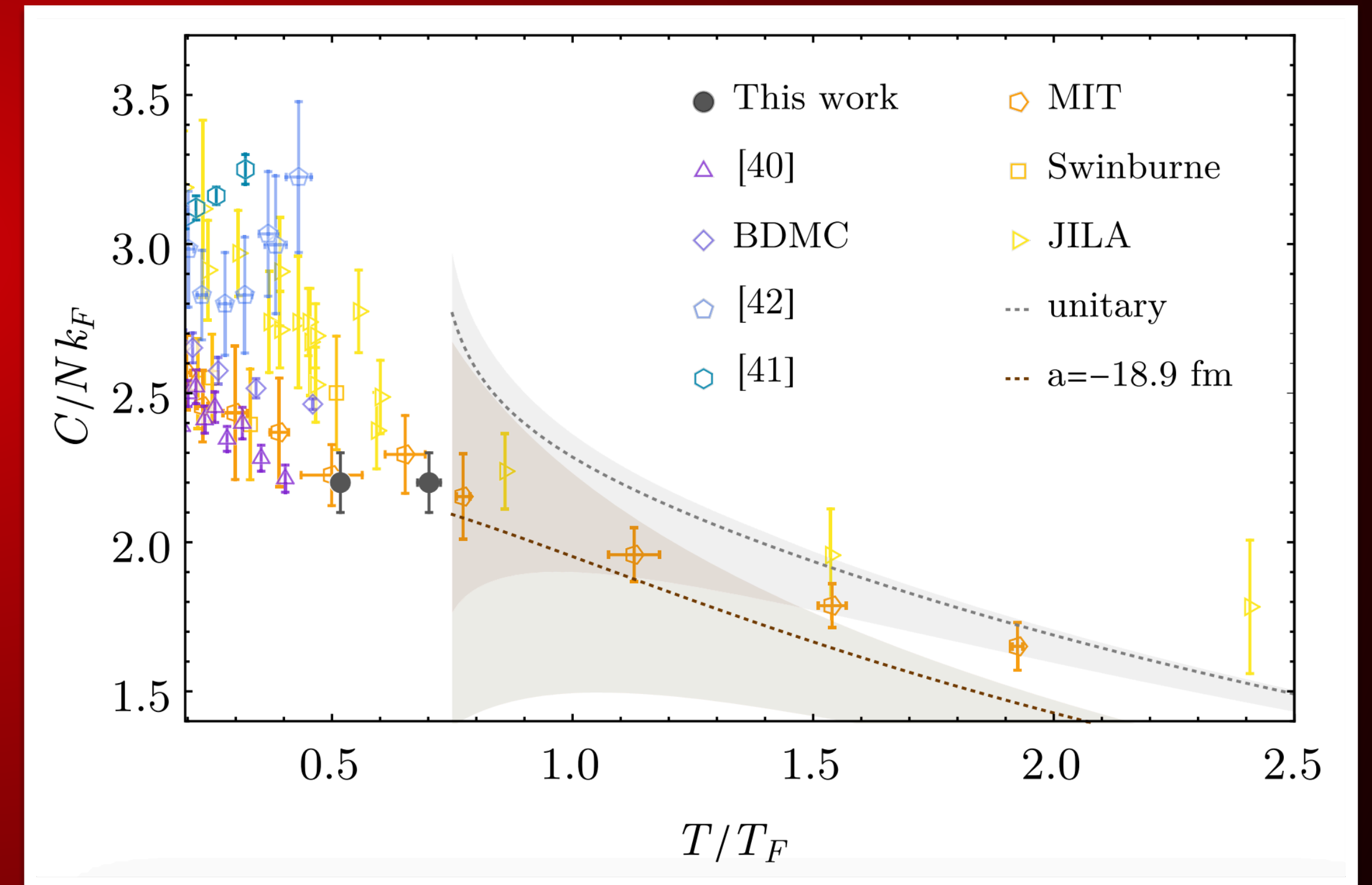


# Main Idea:

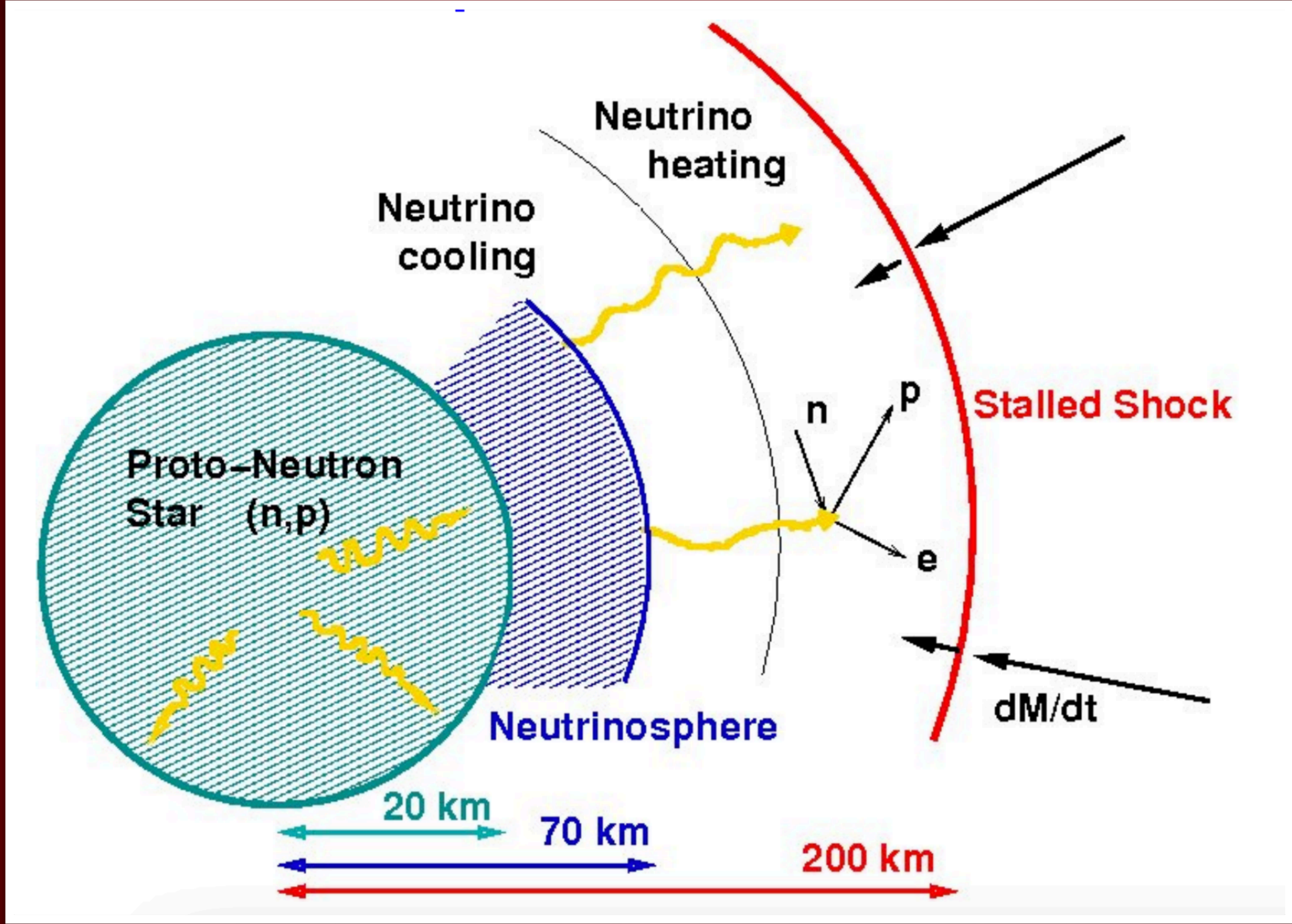
Obtain precision physics where it counts



Melson, Janka et. al., ApJL (2015)

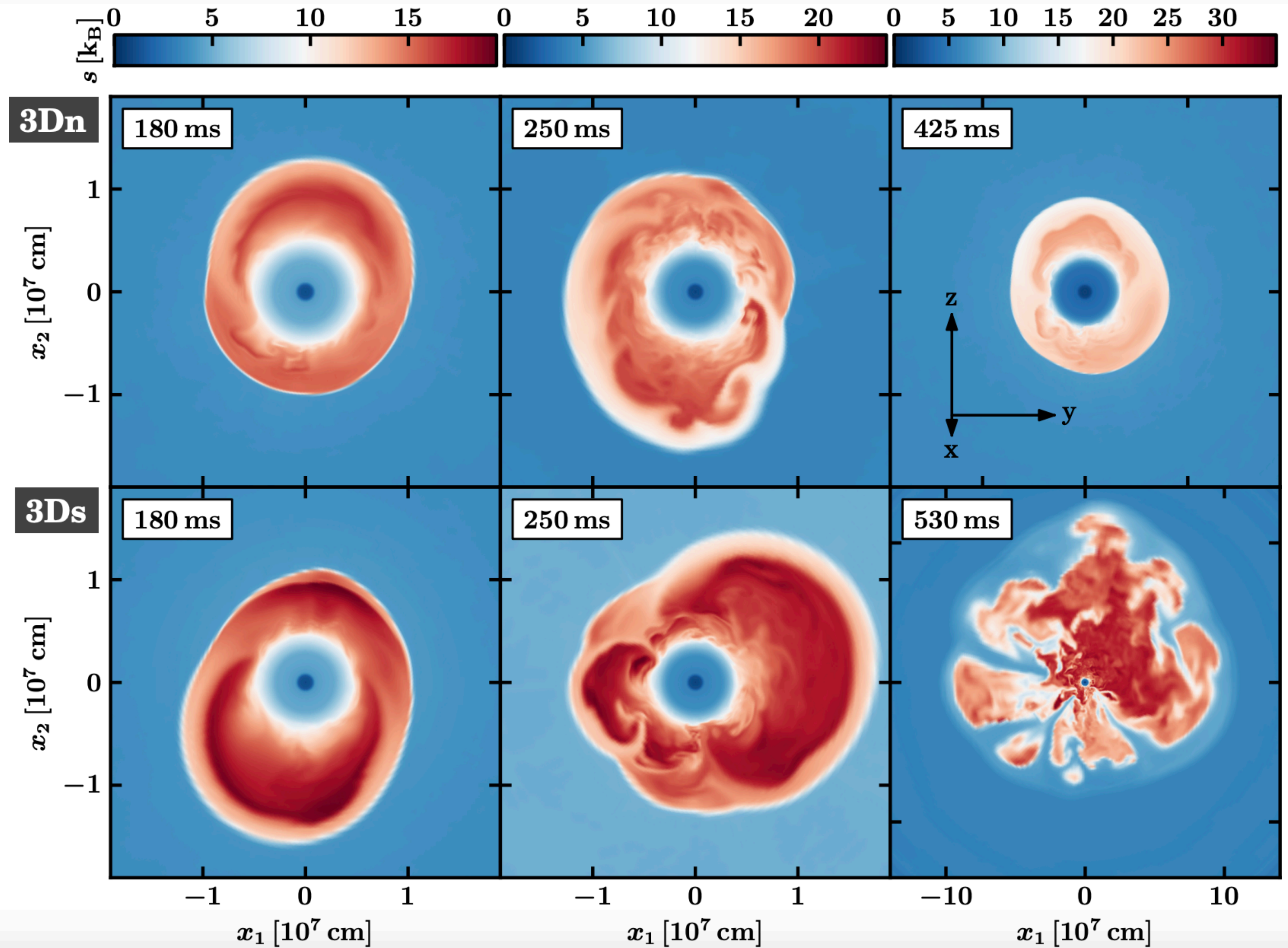


Berkowitz, NCW, et. al. PRL (2022)

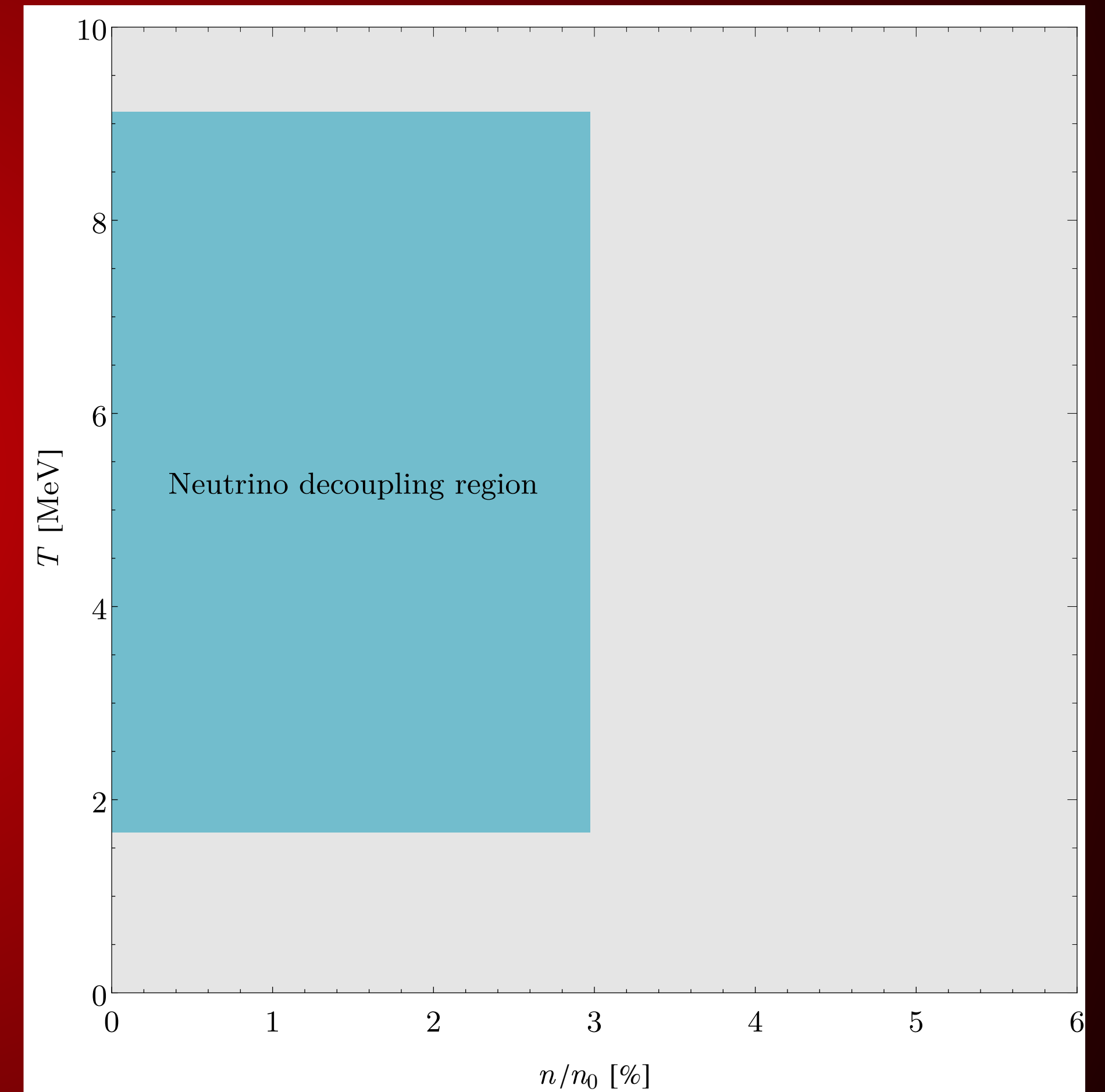
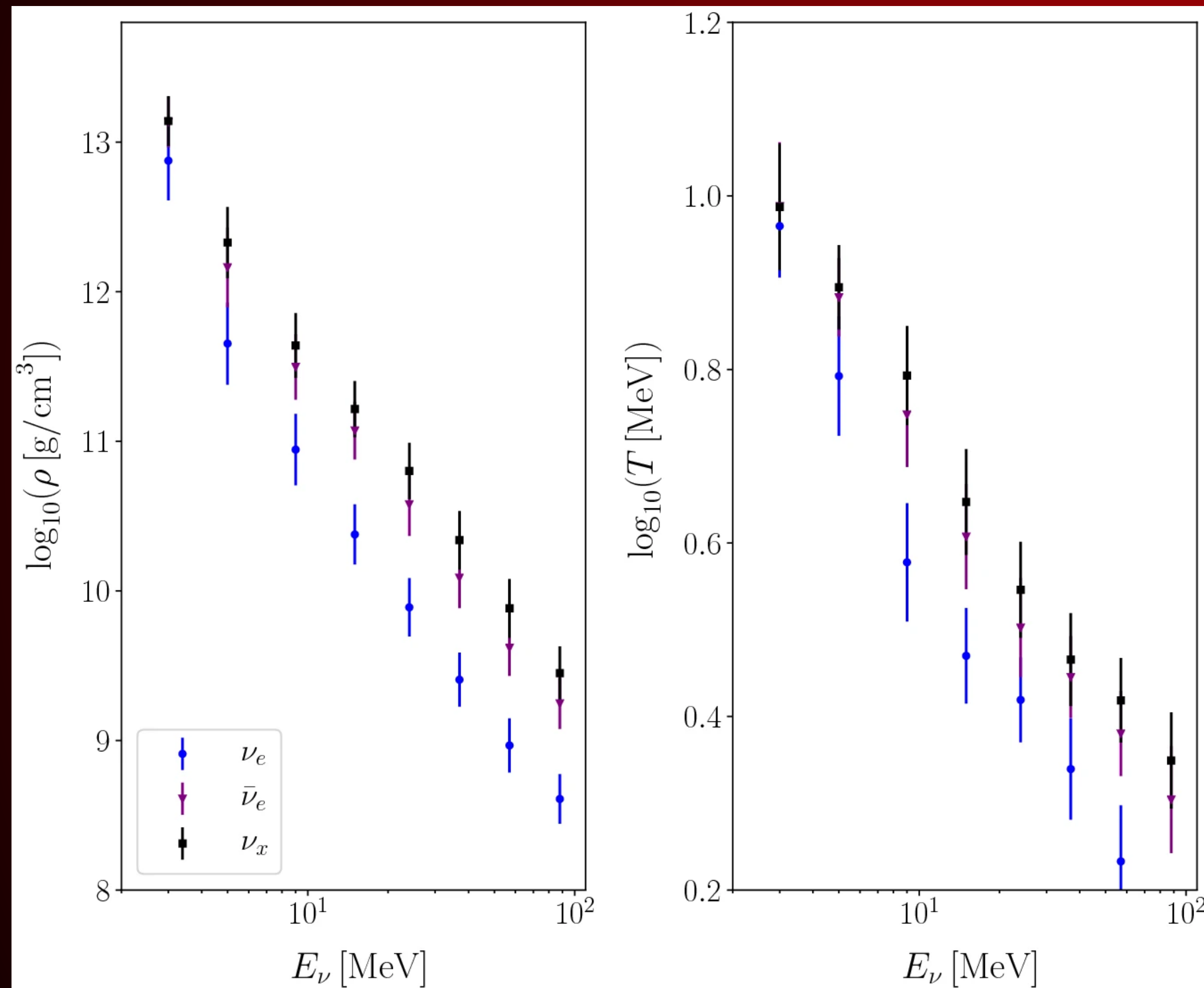


Neglect  
many-body

Include  
many-body



# The Neutrinosphere



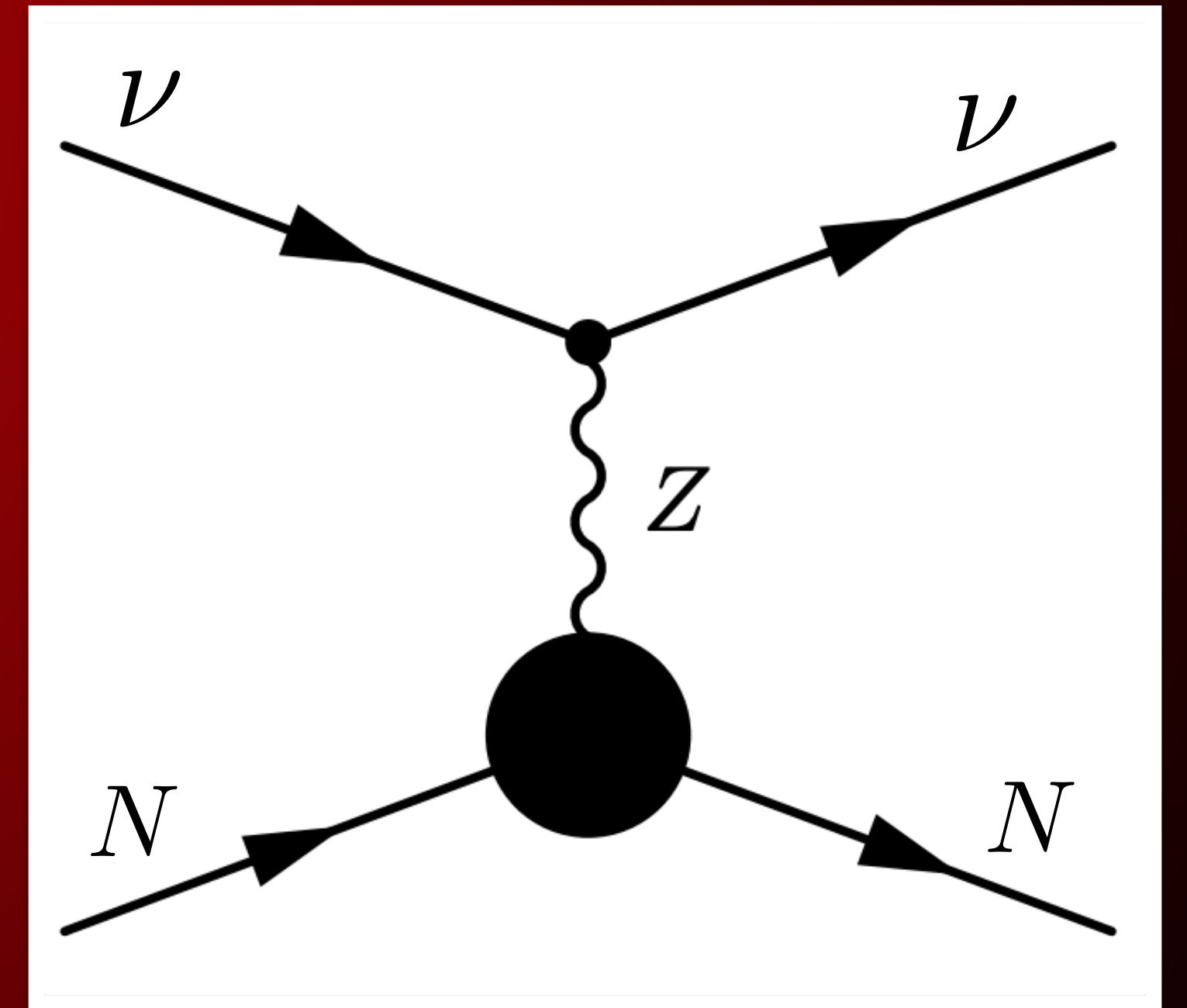
# Neutrino Scattering & Structure Factors

Pethick & Iwamoto PRD (1982) , Horowitz & Schwenk PLB (2006)

$$\frac{d\Gamma(E_\nu)}{d\cos\theta} = \frac{G_F^2 E_\nu^2}{4\pi^2} \left[ c_V^2 (1 + \cos\theta) S_V(\mathbf{q}) + c_A^2 (3 - \cos\theta) S_A(\mathbf{q}) \right]$$

$$S_V(\mathbf{q}) = \int d^3x e^{-i\mathbf{q}\cdot\mathbf{x}} \langle \delta n(\mathbf{x}, 0) \delta n(\mathbf{0}, 0) \rangle$$

$$S_A(\mathbf{q}) = \int d^3x e^{-i\mathbf{q}\cdot\mathbf{x}} \langle \delta s(\mathbf{x}, 0) \delta s(\mathbf{0}, 0) \rangle$$



# Pionless EFT:

Bedaque & van Kolck Annu. Rev. Nucl. Part. Sci. (2002)

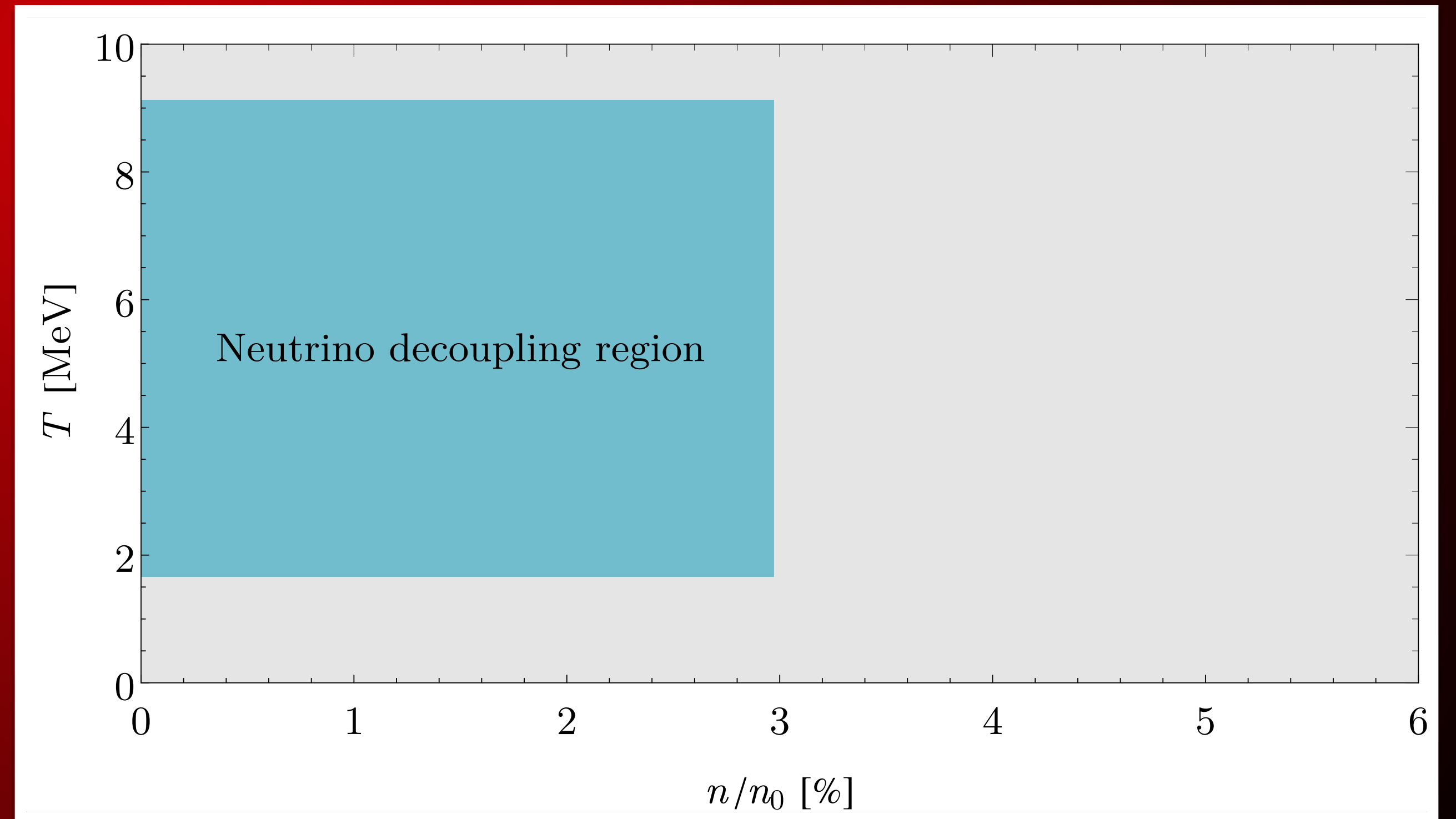
$$H = \int d^3x \frac{|\nabla\psi|^2}{2M} + C_0 (\psi^\dagger\psi)^2 + C_2 (\psi^\dagger\psi) \nabla^2 (\psi^\dagger\psi) + \dots$$

Converges for  $k < m_\pi$



$$T \lesssim 10 \text{ MeV}$$

$$n \lesssim 6\% n_0$$



# Pionless EFT:

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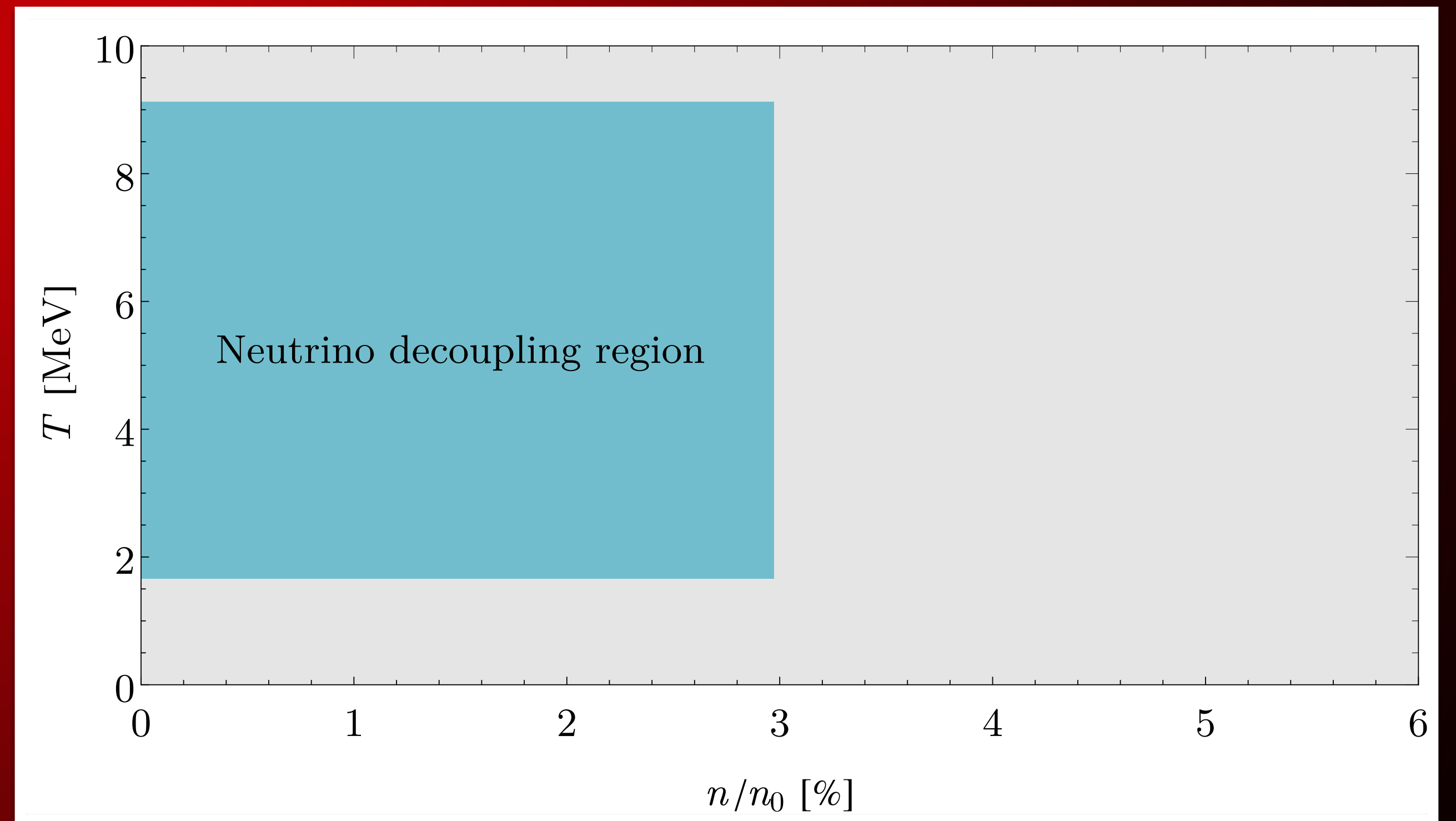
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# Lattice Methods

**Goal:**  $\text{tr}(\mathcal{O} \rho)$  , where  $\rho = \frac{e^{-\beta(H - \mu N)}}{\text{tr}(e^{-\beta(H - \mu N)})}$

**Step 1:** Discretize to finite cubic lattice

$$H = \sum_x \Delta x^3 \left\{ \frac{|\nabla \psi(x)|^2}{2M} + C_0 (\psi^\dagger \psi)^2 \right\}$$

# Lattice Methods

**Goal:**  $\text{tr}(\mathcal{O} \rho)$ , where  $\rho = \frac{e^{-\beta(H - \mu N)}}{\text{tr}(e^{-\beta(H - \mu N)})}$

**Step 2:** Discretize in “time”

$$\text{tr}\left(e^{-\beta(H - \mu N)}\right) = \text{tr}\left(e^{-\Delta t(K - \mu N)} e^{-\Delta t V} \dots e^{-\Delta t(K - \mu N)} e^{-\Delta t V}\right) + \mathcal{O}(\Delta t^2)$$

where  $\beta = \Delta t N_t$

# Lattice Methods

**Goal:**  $\text{tr}(\mathcal{O} \rho)$  , where  $\rho = \frac{e^{-\beta(H - \mu N)}}{\text{tr}(e^{-\beta(H - \mu N)})}$

## Step 3: Hubbard-Stratanovich

$$\text{tr}\left(e^{-\Delta t(K - \mu N)} e^{-\Delta t V} \dots e^{-\Delta t(K - \mu N)} e^{-\Delta t V}\right) = \int \prod_{xt} dA_{xt} e^{-\sum_{xt} A_{xt} M_{xy} A_{yt}} \det D(A)$$

# Lattice Methods

**Goal:**  $\text{tr}(\mathcal{O} \rho)$ , where  $\rho = \frac{e^{-\beta(H - \mu N)}}{\text{tr}(e^{-\beta(H - \mu N)})}$

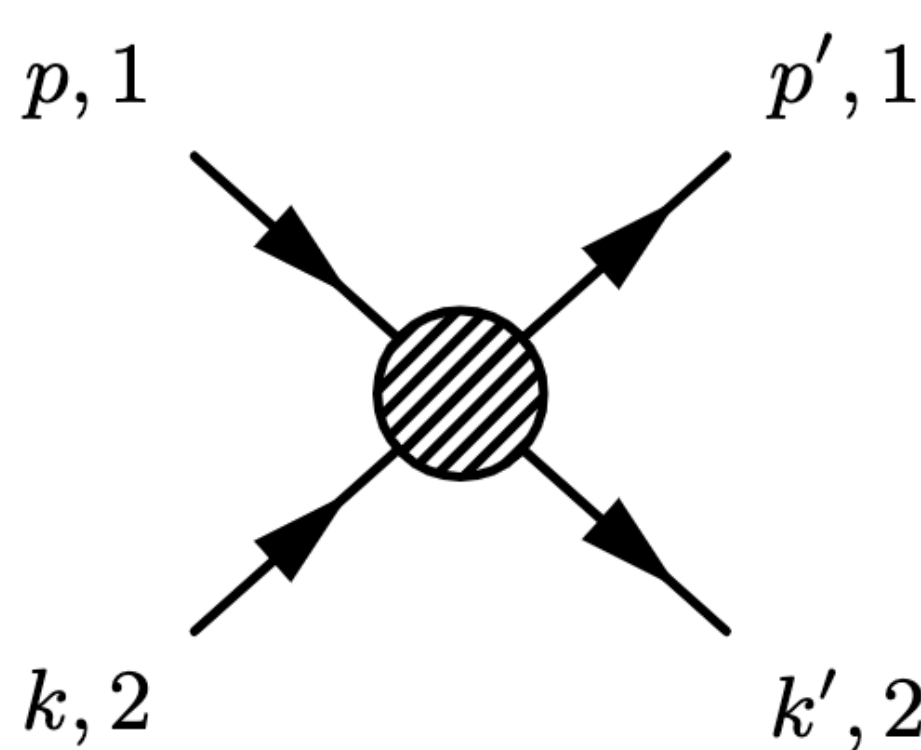
## Step 4: Monte Carlo

$$\frac{\text{tr}(\mathcal{O} \prod_t e^{-\Delta t V} e^{-\Delta t(K - \mu N)})}{\text{tr}(\prod_t e^{-\Delta t V} e^{-\Delta t(K - \mu N)})} = \int DA p(A) \mathcal{O}(A)$$

where  $p(A) = \frac{e^{-\sum_{xt} A_{xt} M_{xy} A_{yt}} \det D(A)}{\int DA e^{-\sum_{xt} A_{xt} M_{xy} A_{yt}} \det D(A)}$

# Renormalization

Kaplan, Savage & Wise (1998), Bedaque & van Kolck (2002)



A Feynman diagram showing a four-point interaction. A central shaded circle is connected to four external lines. The top-left line is labeled  $p, 1$  with an arrow pointing towards the circle. The top-right line is labeled  $p', 1$  with an arrow pointing away from the circle. The bottom-left line is labeled  $k, 2$  with an arrow pointing towards the circle. The bottom-right line is labeled  $k', 2$  with an arrow pointing away from the circle.

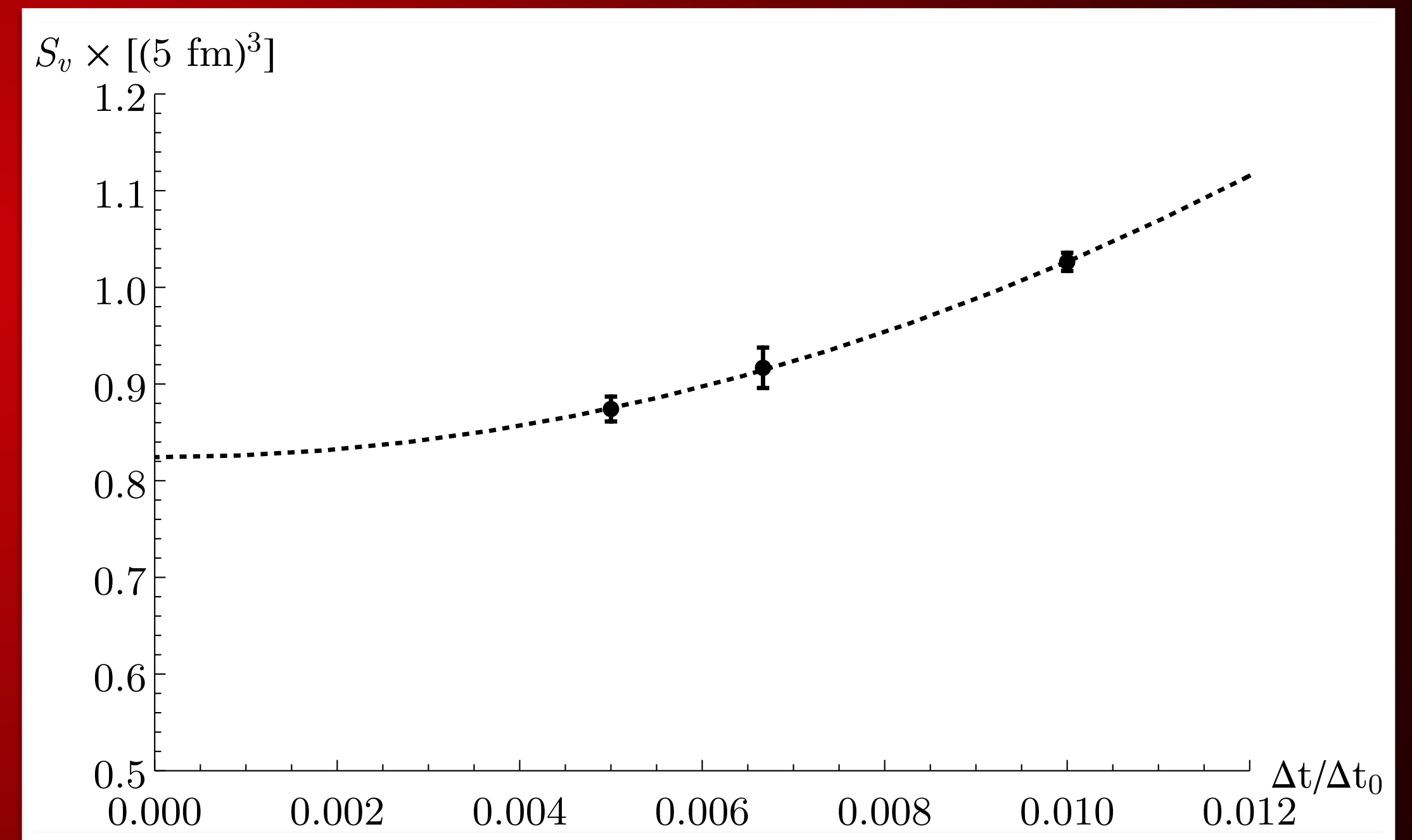
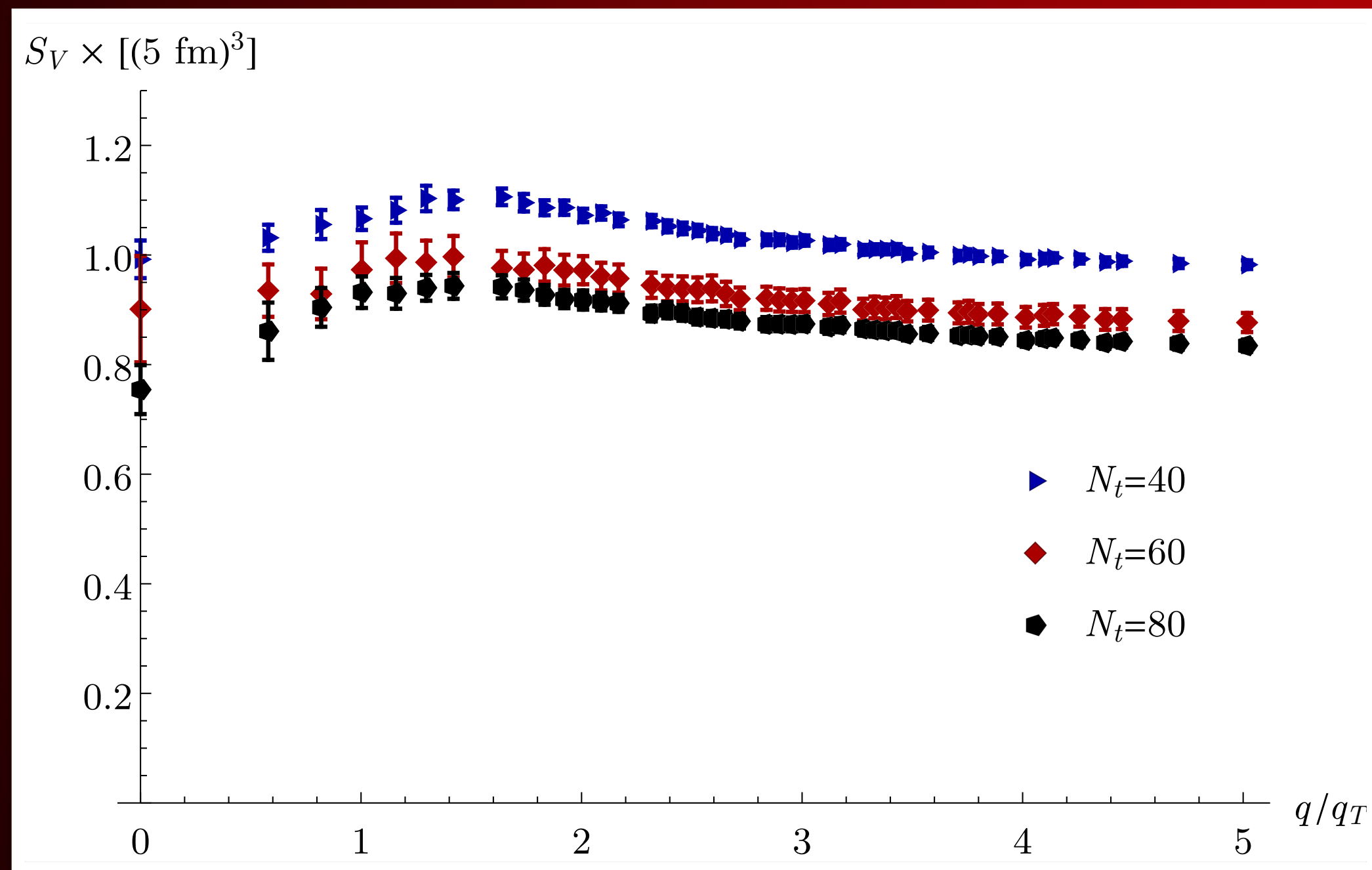
$$= \left( \frac{\Delta x}{MC_0} + \mathcal{P} \int \frac{d^3 q}{(2\pi)^3} \frac{1}{q^2 - EM\Delta x^2} + i \frac{\sqrt{EM\Delta x^2}}{4\pi} \right)^{-1}$$

$$iA = \frac{4\pi/M}{k \cot \delta(k) - ik} \implies C_0(\Delta x) = -\frac{\Delta x}{M} \left( a - \frac{\Delta x}{4\pi a} \right)^{-1}$$

$$= \frac{4\pi/M}{(-1/a + rk^2/2 + \dots) - ik}$$

# Strategy

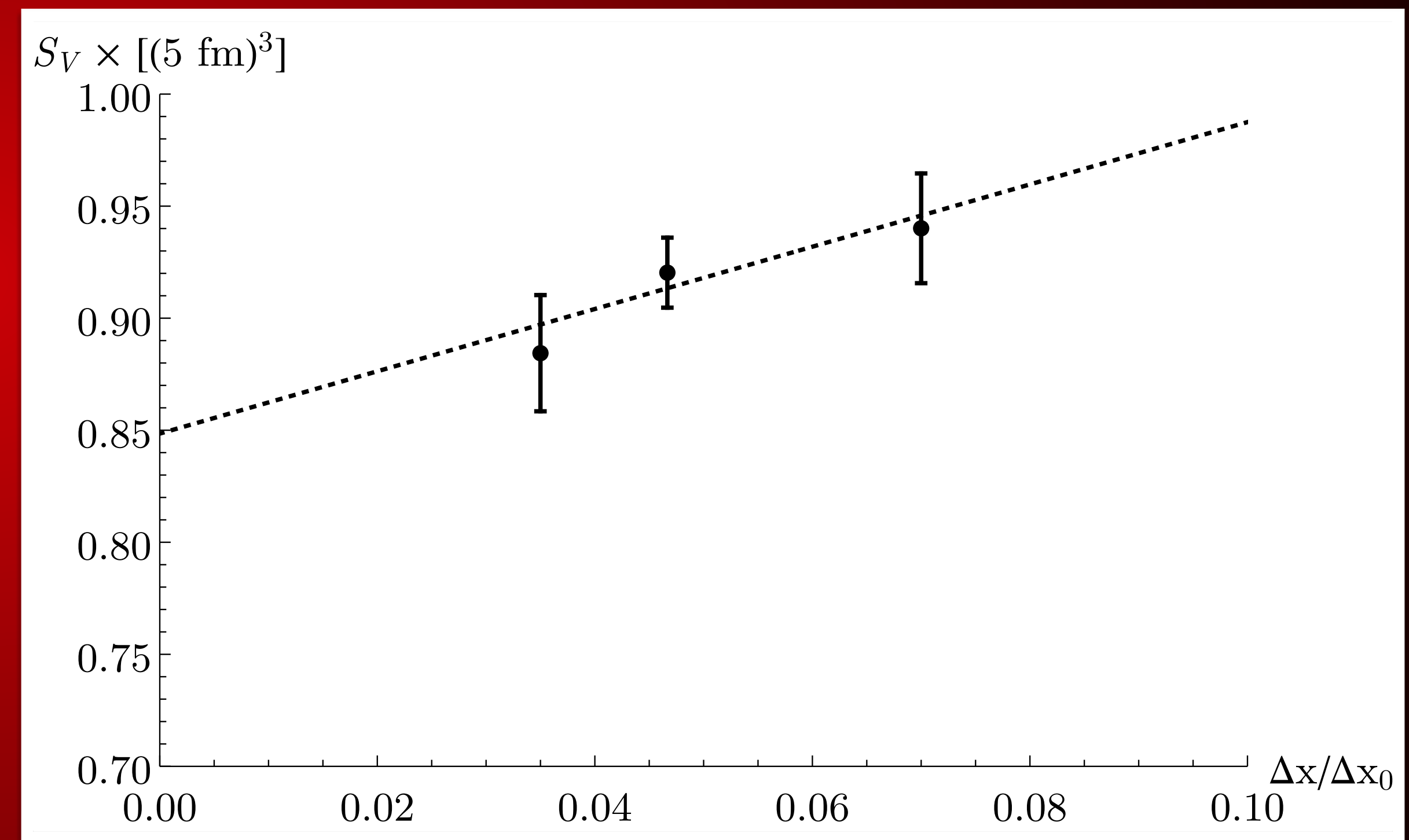
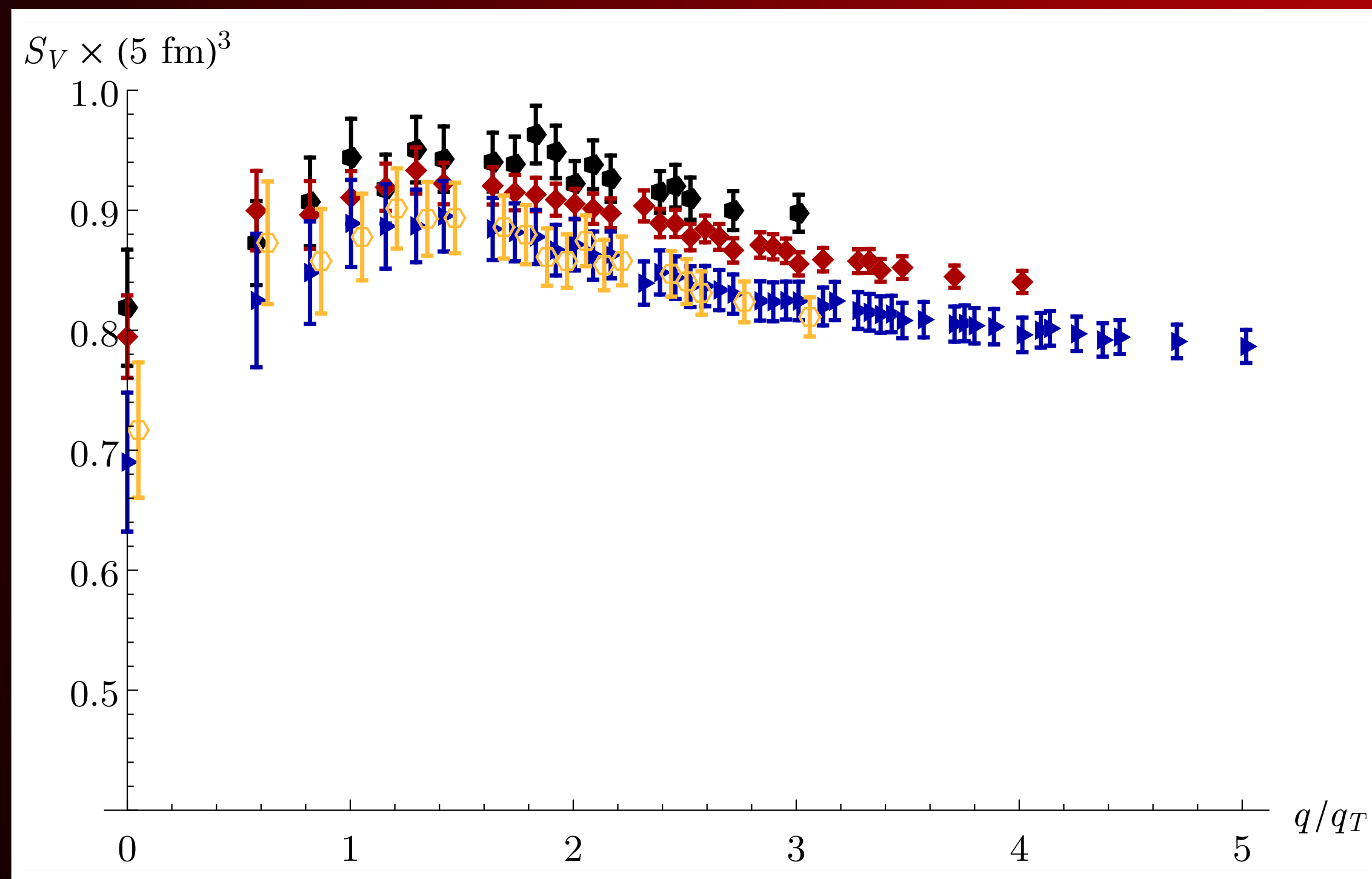
1.  $\Delta t \rightarrow 0$  first



Identical scaling demonstrated by Alhassid et. al., PRL (2020)

# Strategy

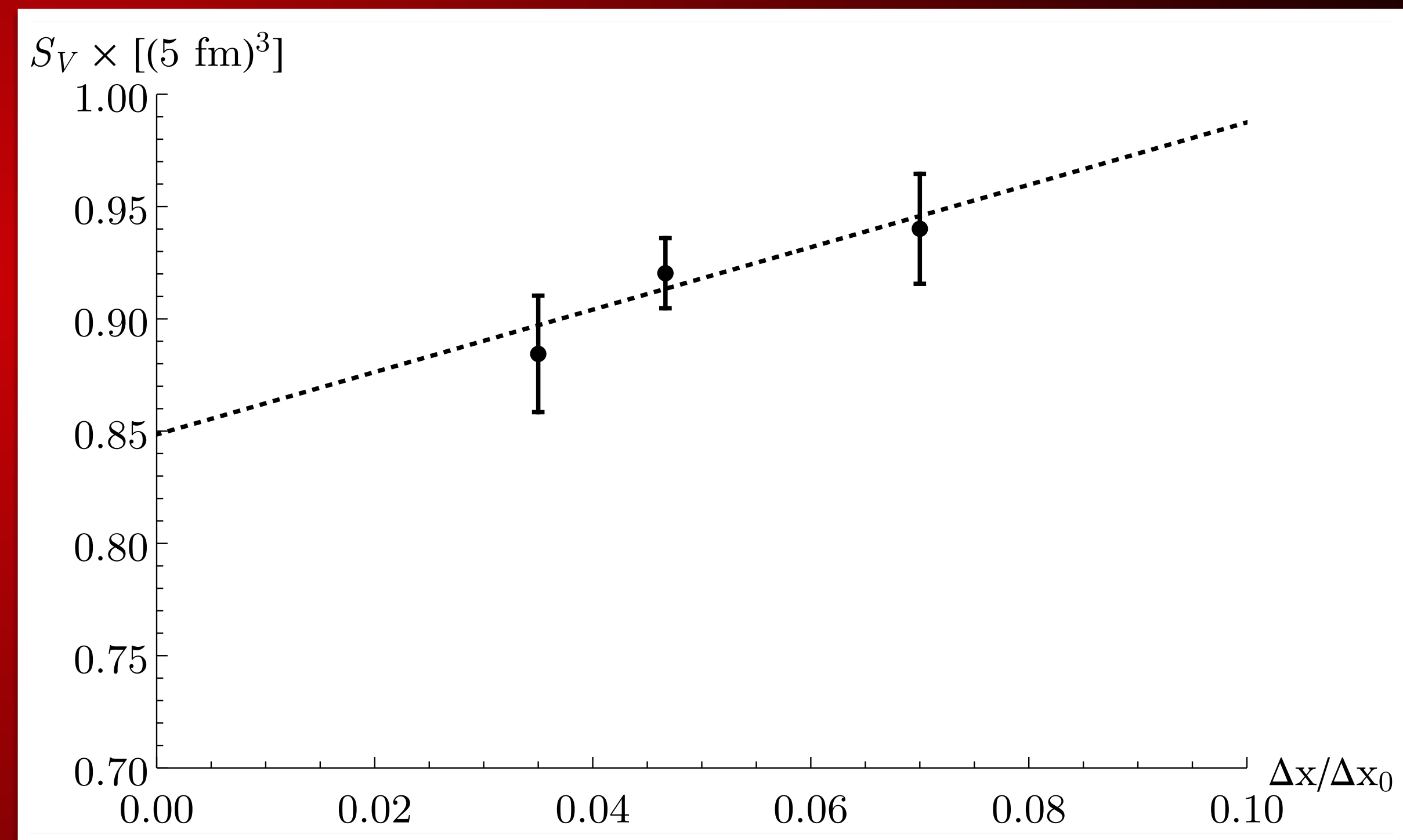
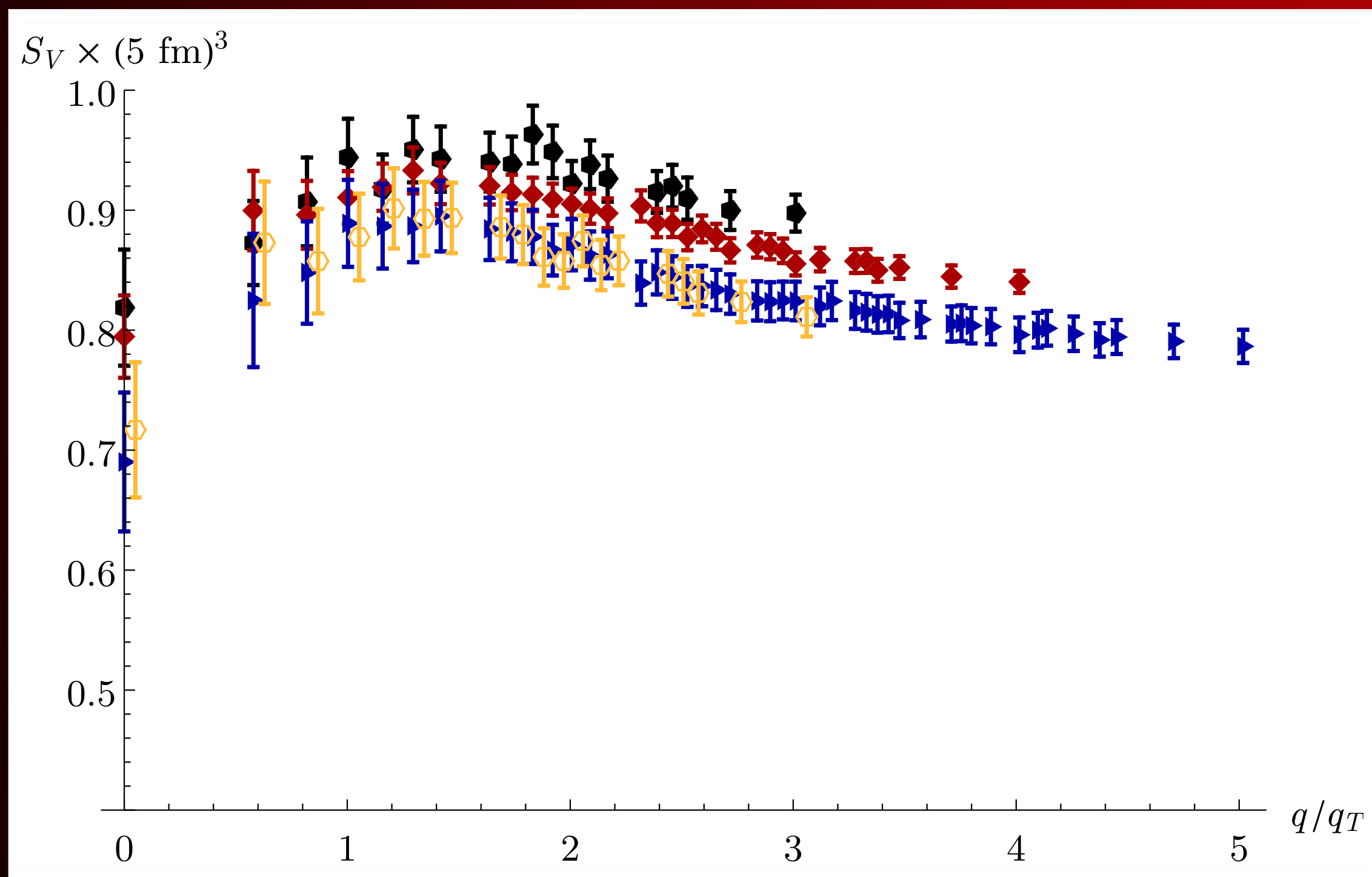
## 2. $\Delta x \rightarrow 0$ second



Identical scaling demonstrated by Alhassid et. al., PRL (2020)

# Strategy

## 3. $V \rightarrow \infty$ third





# Structure Factors

Berkowitz, NCW, et. al. PRL (2022)

NCW, et. al. PRC (2020)

## OPE Prediction:

Braaten & Platter PRL (2008)

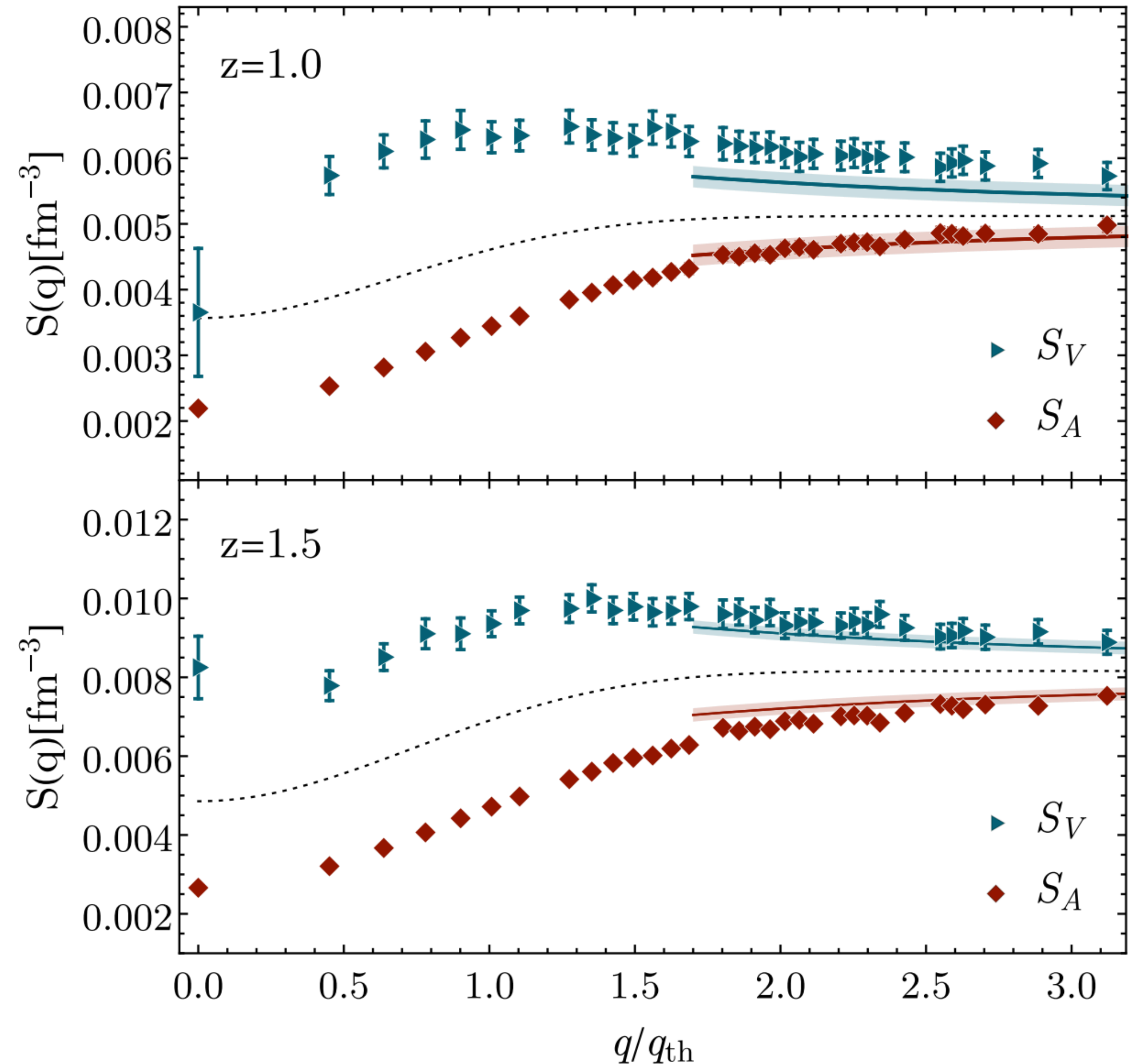
$$S_V(q) = \langle n \rangle + \frac{C}{8Vq} + \mathcal{O}(q^{-2})$$

$$S_A(q) = \langle n \rangle - \frac{C}{8Vq} + \mathcal{O}(q^{-2})$$

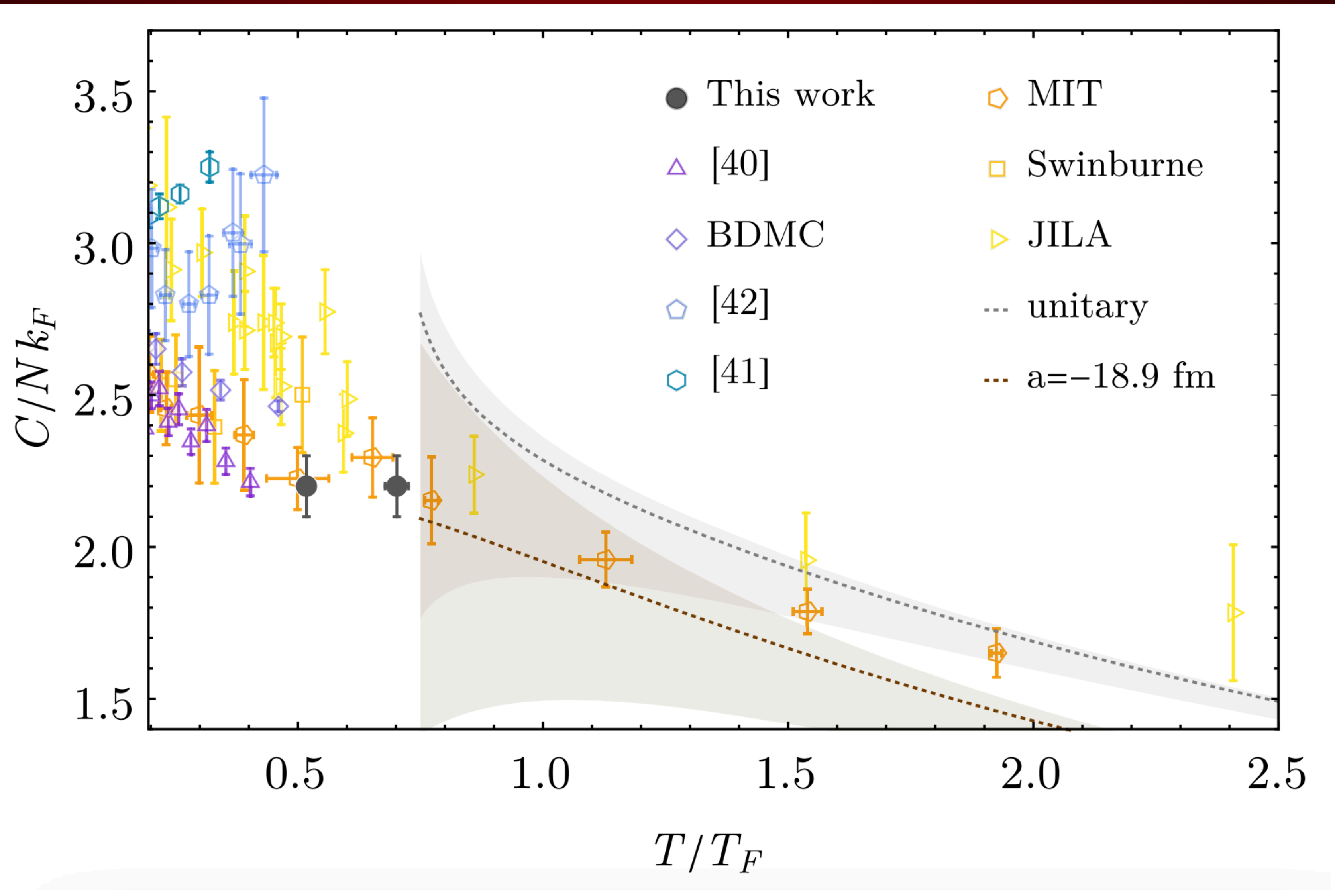
## Fluctuations:

$$S_V(0) = T \partial n / \partial \mu$$

$$S_A(0) = T \partial s / \partial h$$



# **Symanzik Improvement for the unitary gas**

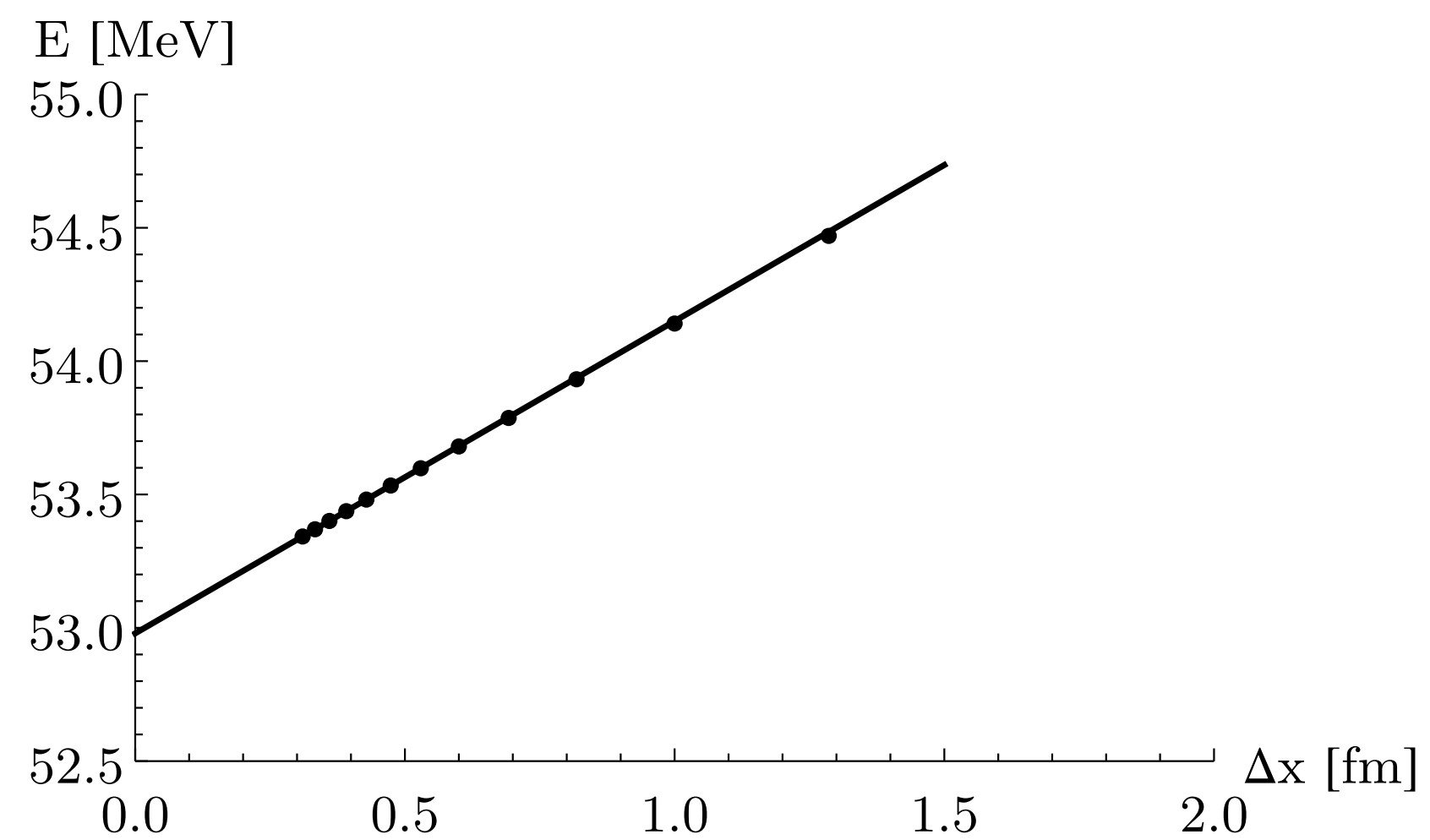
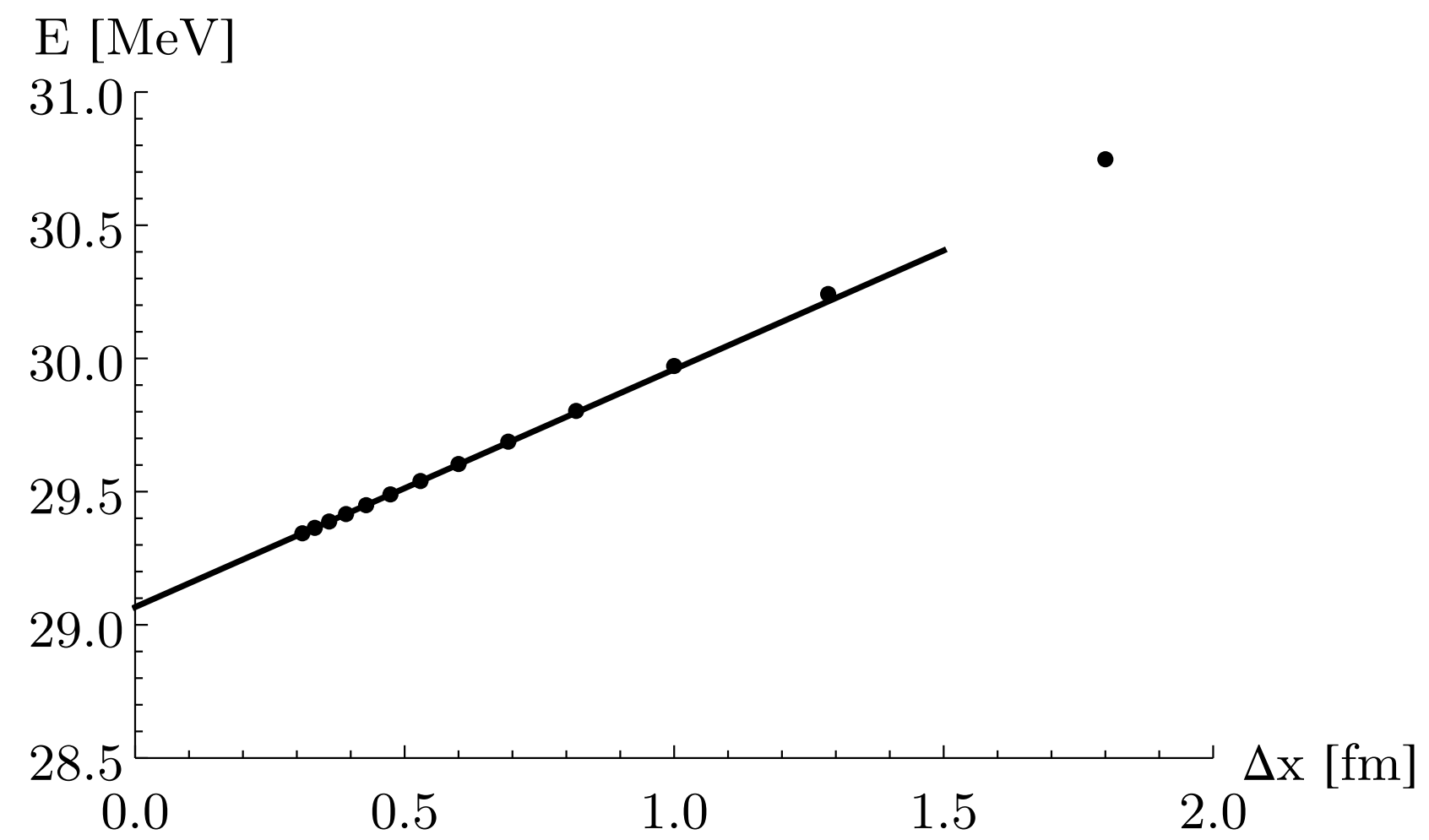


# Symanzik Improvement speeds-up convergence

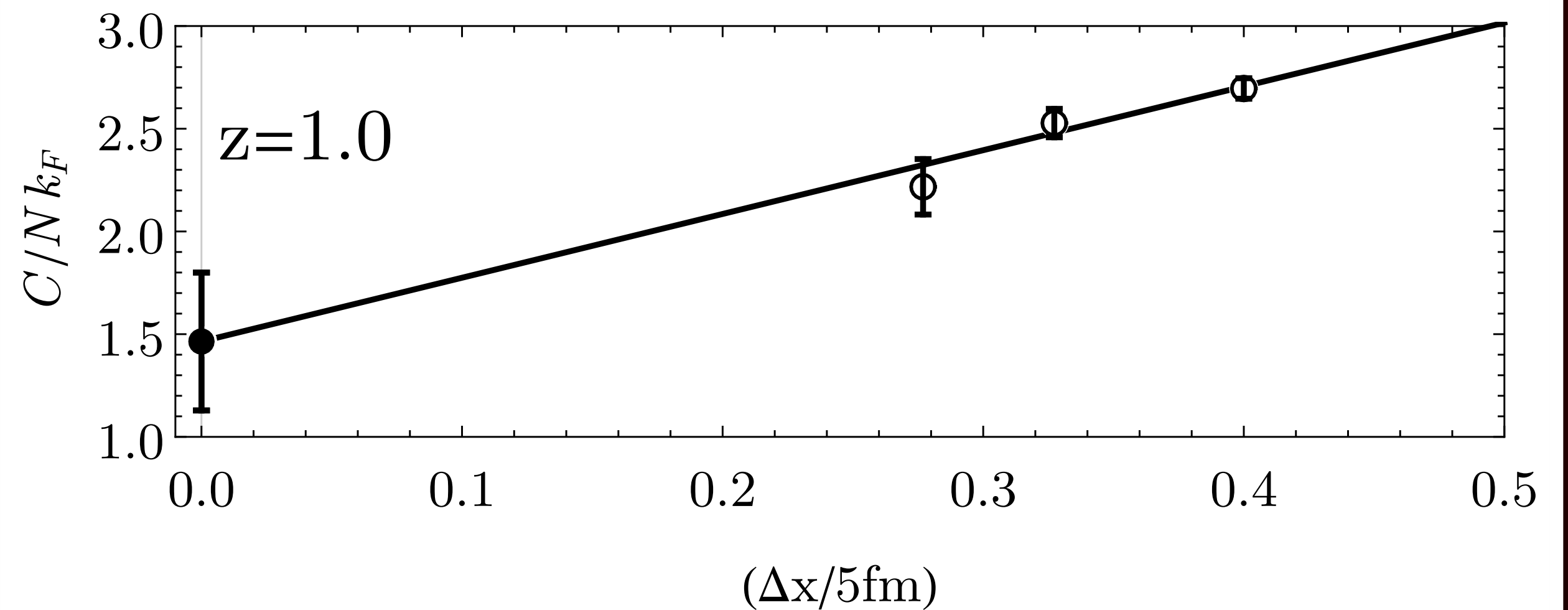
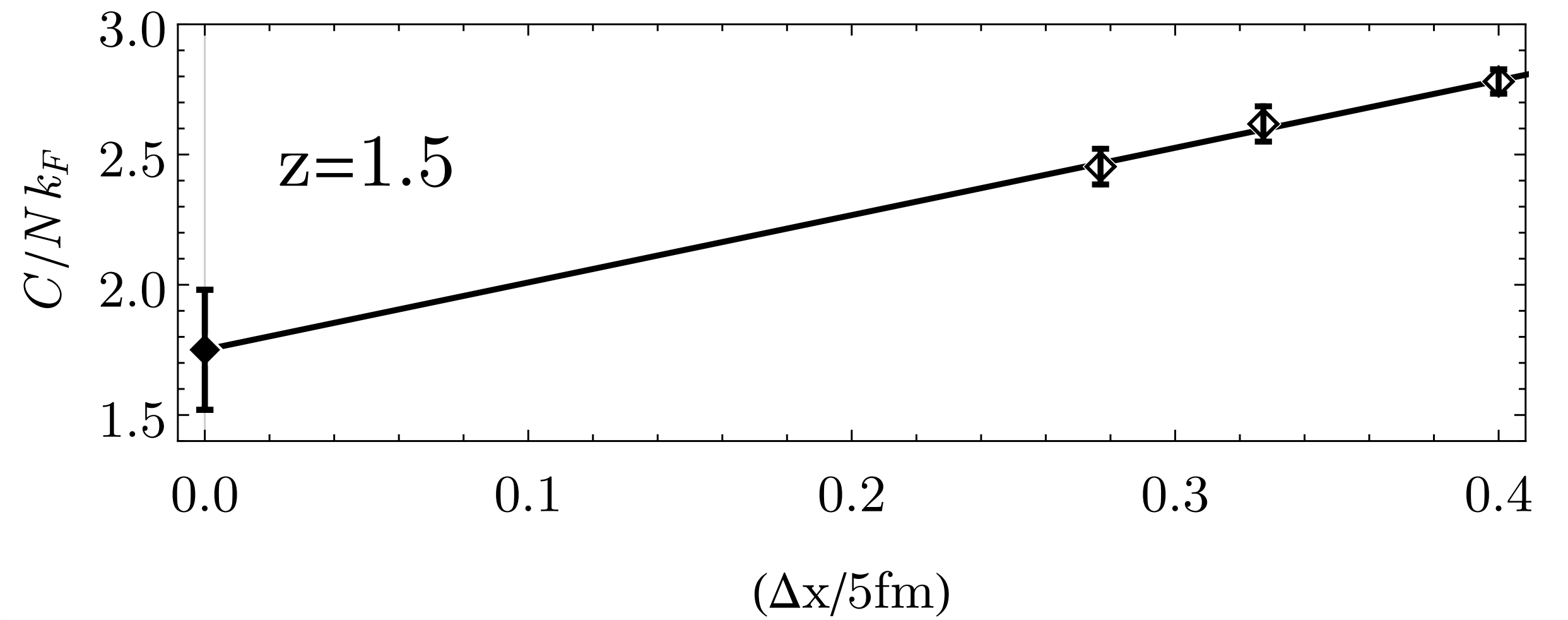
K. Symanzik, NPB 226, 187 (1983)

$$H \rightarrow H + \underbrace{\sum_i C_i \mathcal{O}_i}$$

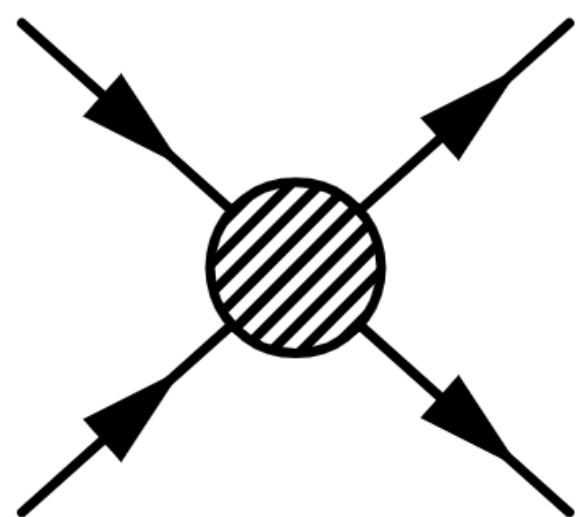
Tuned to cancel leading order  $\Delta x$  dependence



**Two-body**



**Many-body**



$$= \left( \frac{\Delta x}{MC_0} + \mathcal{P} \int \frac{d^3 q}{(2\pi)^3} \frac{1}{q^2 - EM\Delta x^2} + i \frac{\sqrt{EM\Delta x^2}}{4\pi} \right)^{-1}$$

$$= \left( \frac{\Delta x}{MC_0} + \alpha + \beta(EM\Delta x^2) + i \frac{\sqrt{EM\Delta x^2}}{4\pi} + \mathcal{O}\left((EM\Delta x^2)^2\right) \right)^{-1}$$

$$iA = \frac{4\pi/M}{k \cot \delta(k) - ik} \quad \Longrightarrow \quad C_0(\Delta x) = -\frac{\Delta x}{M} \left( \alpha - \frac{\Delta x}{4\pi a} \right)^{-1}$$

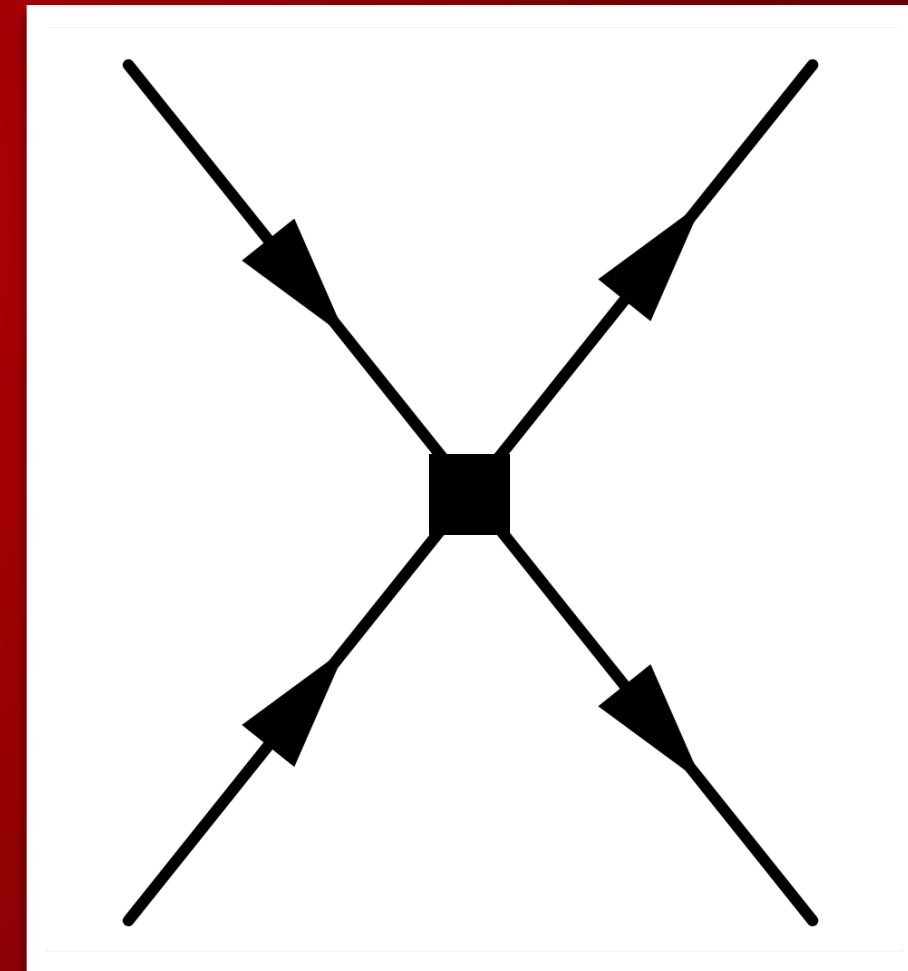
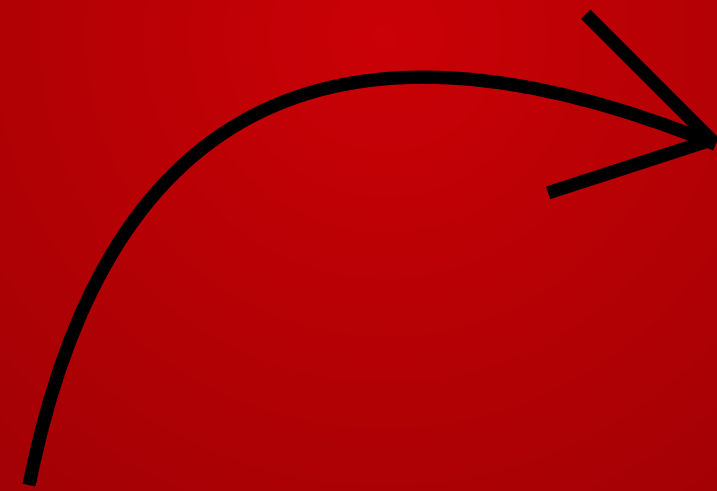
$$= \frac{4\pi/M}{(-1/a + rk^2/2 + \dots) - ik} \quad r(\Delta x) = 0.337\Delta x$$

**A solution:**  $H \rightarrow H + C_2 \sum_{xi} n(x)n(x + \delta x_i)$

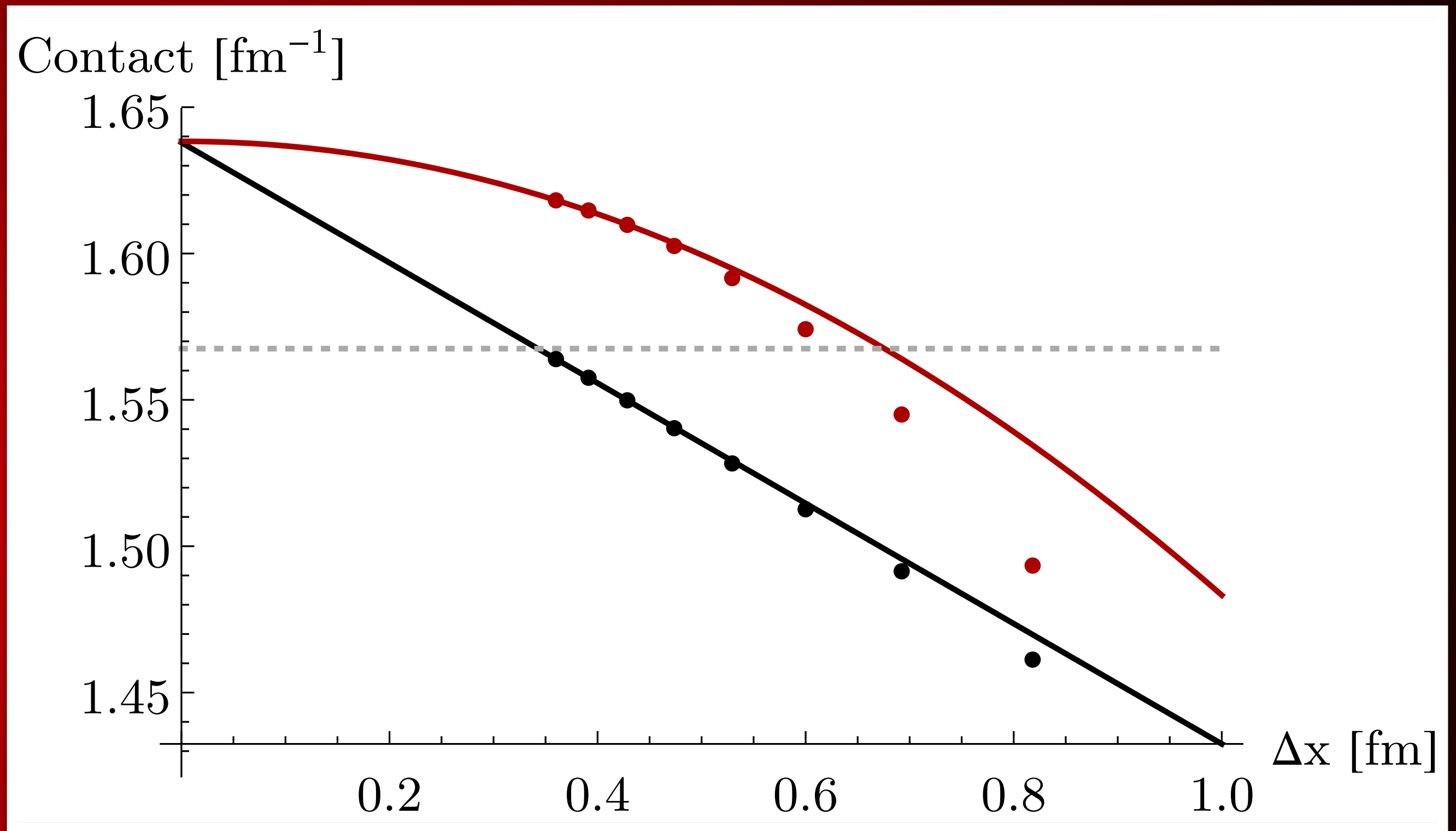
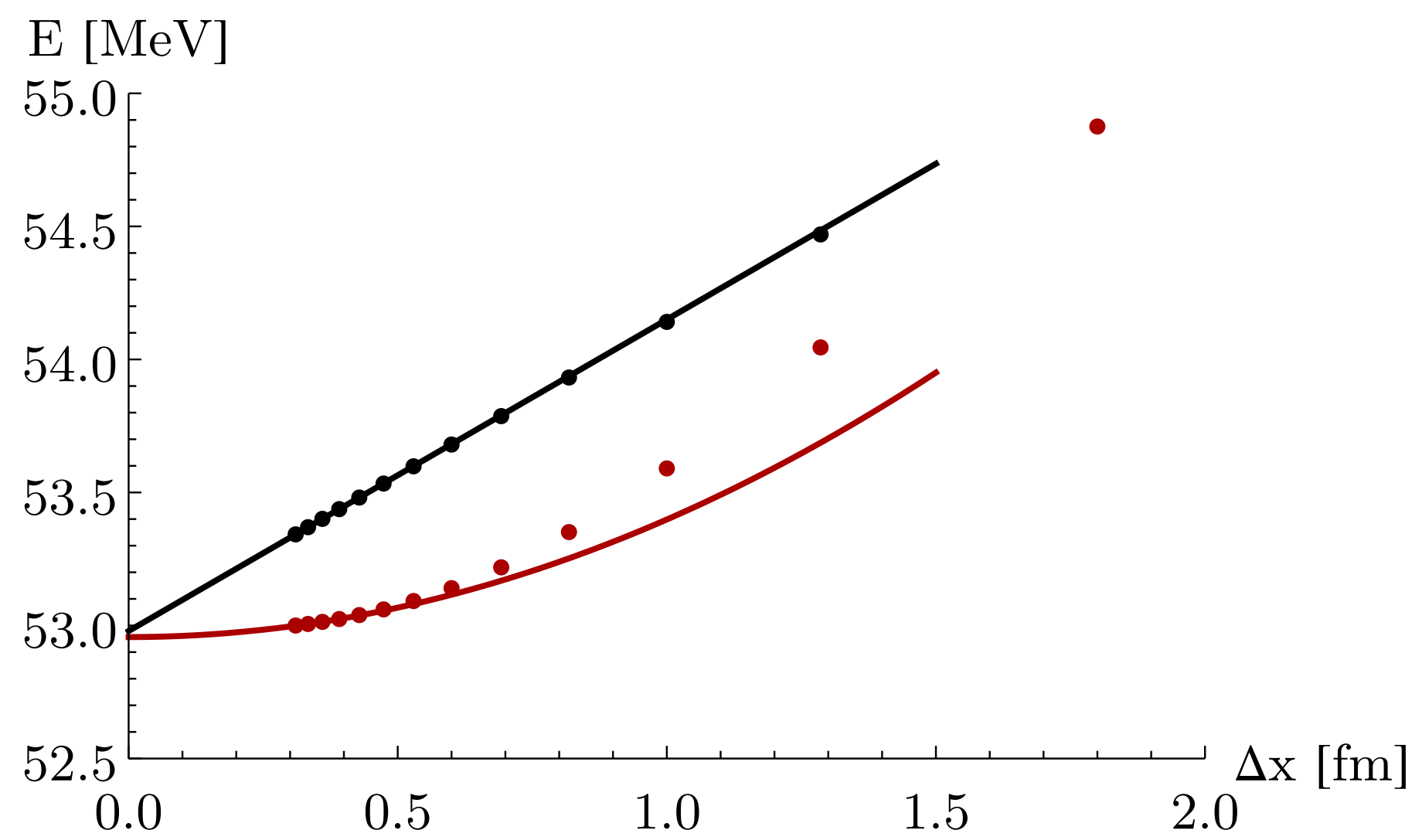
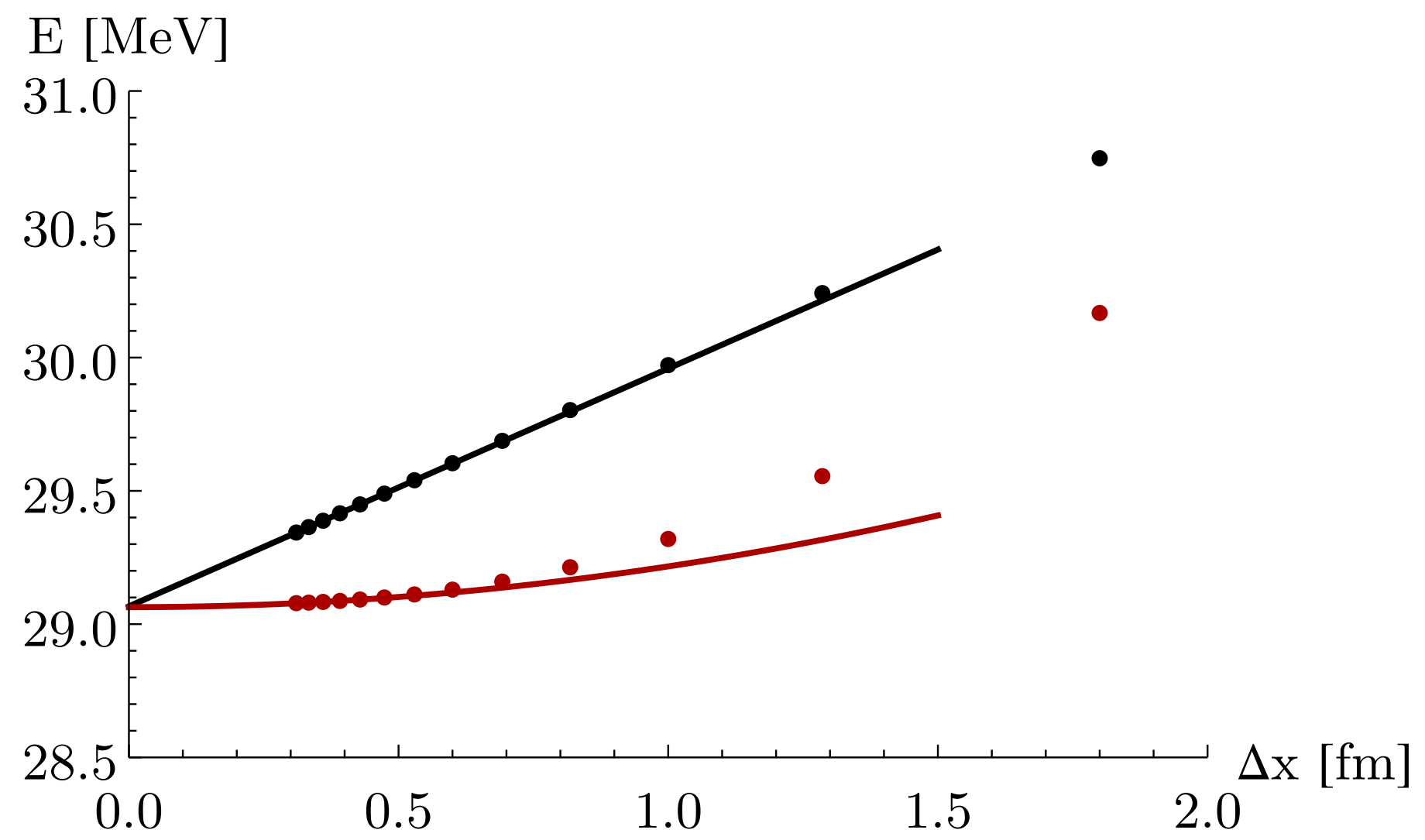
**Reason:**

$$C_2 \sum_{xi} n(x)n(x + \delta x_i)$$

$$= C_2 \sum_{xi} n(x) \left[ n(x) + \Delta x^2 \nabla^2 n(x) + \dots \right]$$



$$= C_2 \Delta x^2 k^2$$





# Modification of Tan Contact

Tan Energy Relation: 
$$-\frac{C}{4\pi M} = \frac{dE}{da^{-1}}$$

Tan, Annals of Physics (2008)

Braaten showed: 
$$-\frac{C}{4\pi M} = \left\langle \frac{\partial H}{\partial a^{-1}} \right\rangle = \left\langle \frac{\partial H}{\partial C_0} \frac{\partial C_0}{\partial a^{-1}} \right\rangle$$

Braaten & Platter PRL (2008)

$$= C_0^2 M^2 \left\langle \int d^3x n_1(x) n_2(x) \right\rangle$$

# Modification of Tan Contact

Contact in improved theory:

$$-\frac{C}{4\pi M} = \left\langle \frac{\partial H}{\partial a^{-1}} \right\rangle = \left\langle \frac{\partial H}{\partial C_0} \frac{\partial C_0}{\partial a^{-1}} + \frac{\partial H}{\partial C_2} \frac{\partial C_2}{\partial a^{-1}} \right\rangle$$

Braaten operator



new operator



# The Team



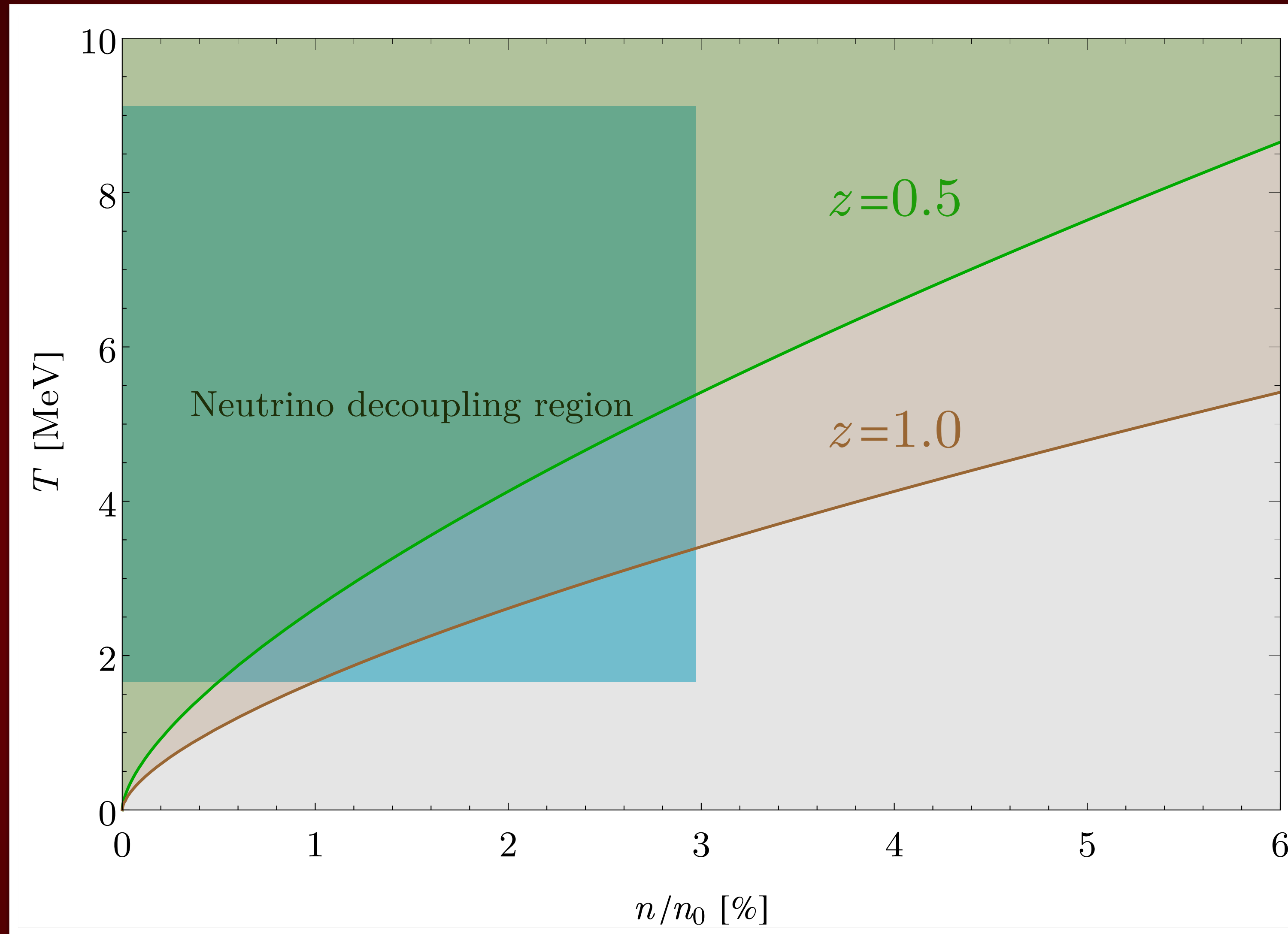
**Paulo Bedqaue**  
University of Maryland



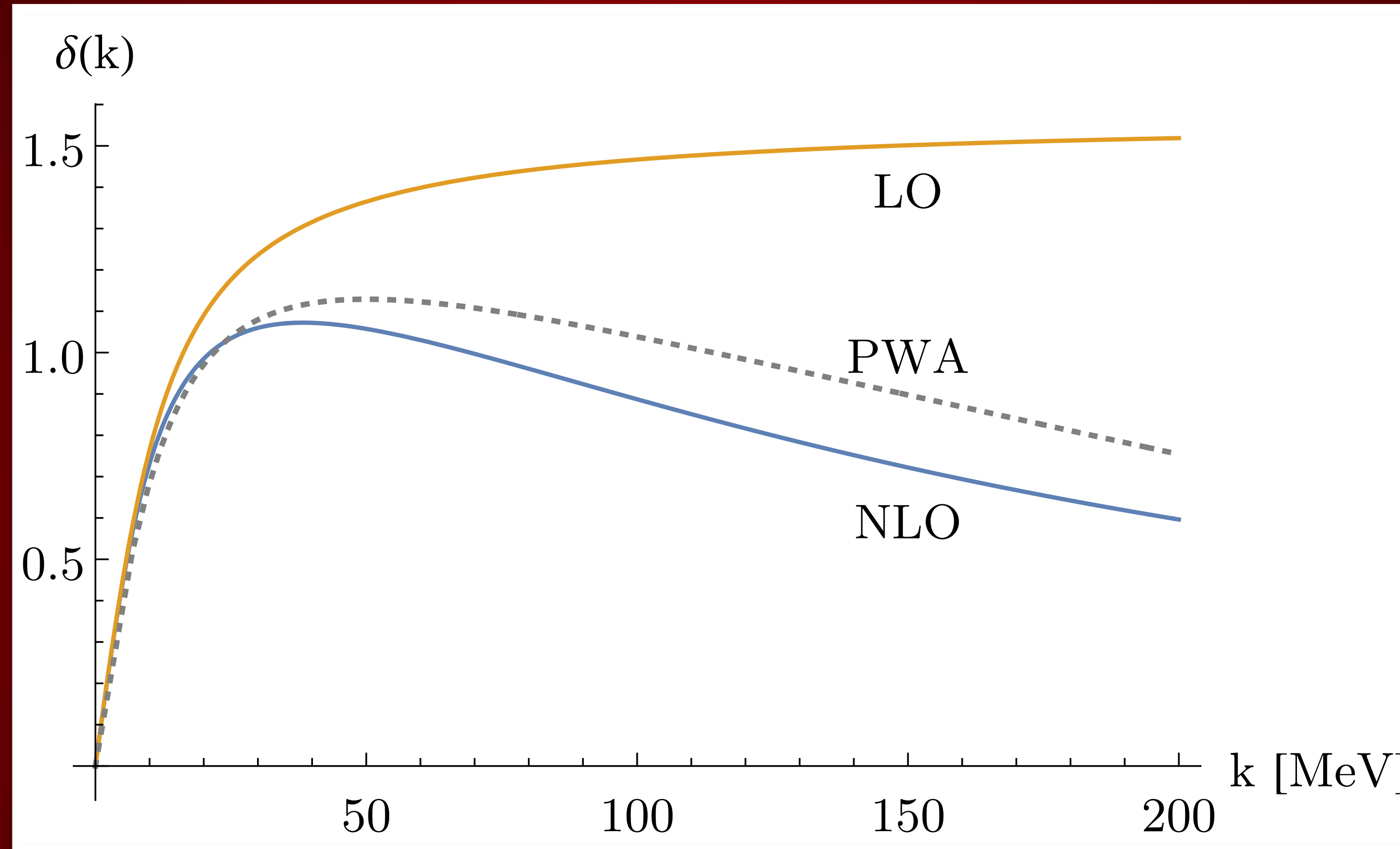
**Evan Berkowitz**  
Forschungszentrum Jülich



**Andrei Alexandru**  
GW University



**Virial comparison**



**1S0 nn phase shifts**