

The impact of radius data on nuclear density functional theory

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Outline

Acknowledgments:

J. Friedrich	formerly University Mainz / Germany
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exp. data	MAMI/Mainz, ISOLDE/CERN, NSCL/MSU, JLAB

1) Statistical aspects of empirical calibration of nuclear density functionals

Nuclear density functional theory (DFT)

nuclear energy-density functional: $E = E(\underbrace{\mathbf{p}}_{\text{params.}})[\underbrace{\rho(\mathbf{r}), \tau(\mathbf{r}), \mathbf{J}(\mathbf{r}), \xi(\mathbf{r})}_{\text{densities:local,kinetic,spin-orbit,pairing}}]$

eqs. of motion by variation of s.p. wavefunction & occupation amplitudes

typically a dozen free model parameters $\mathbf{p} \longleftrightarrow \chi^2$ fit to empirical data

\implies optimal \mathbf{p} , uncertainties of predictions, stat. correlations between observables

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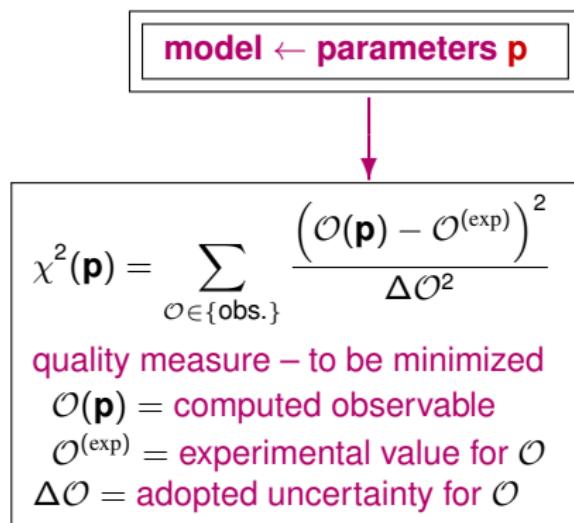
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2 functionals: **Skyrme:** density dependence, kinetic, surface, l^* 's, pairing

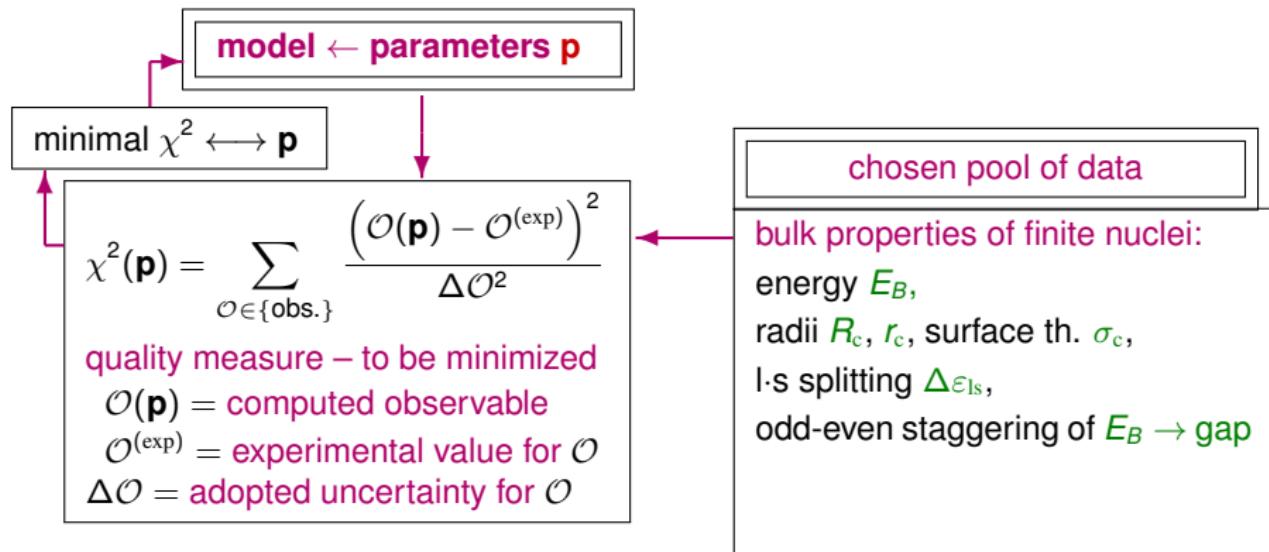
Fayans: " + gradient term in pairing & surface energy

Fayans pairing energy density $\mathcal{E} = \rho_{\text{pair}}^2(\mathbf{r}) \left(\underbrace{V_{\text{pair}} + V'_{\text{pair}} \rho^\gamma(\mathbf{r})}_{\text{volume&surface pairing}} + \underbrace{V_{\text{pair}}^{(\nabla)} (\nabla \rho(\mathbf{r}))^2}_{\text{gradient term}} \right)$

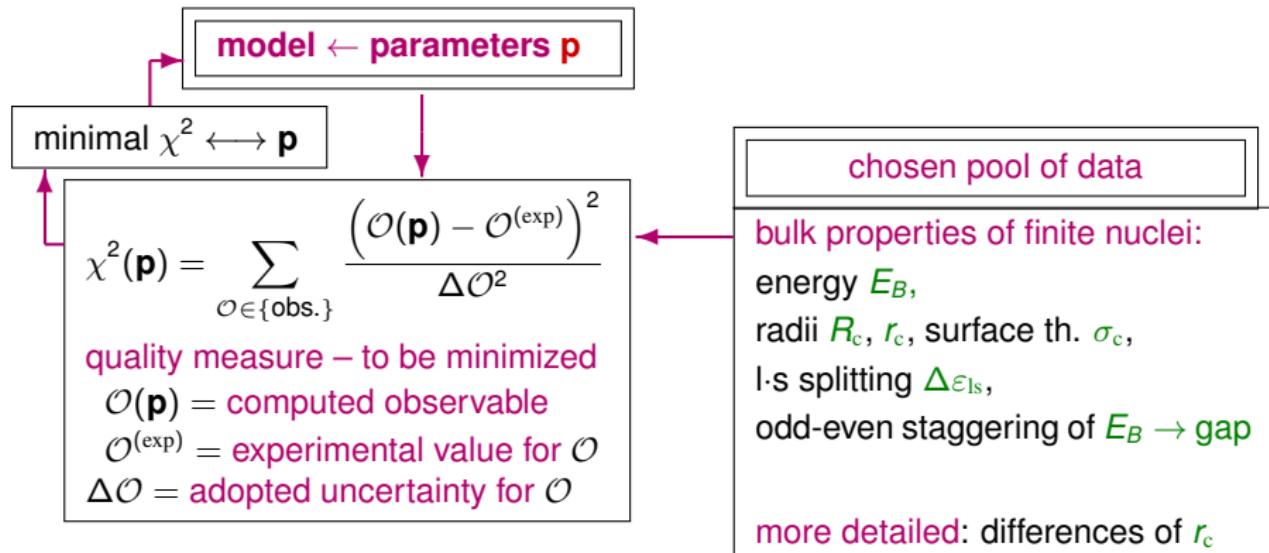
Strategy for adjusting the model parameters to phenomenological data



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Probabilistic interpretation of the space of model parameters

probability distribution for model parameters \mathbf{p} :

$$P(\mathbf{p}) \propto \exp\left(-\chi^2(\mathbf{p})\right), \quad \chi^2 = \sum_{\mathcal{O} \in \{\text{obs.}\}} \frac{(\mathcal{O}(\mathbf{p}) - \mathcal{O}^{(\text{exp})})^2}{\Delta \mathcal{O}^2}$$

- ⇒ optimal parameter set \mathbf{p}_0 : $\chi^2(\mathbf{p}_0) = \text{minimal}$
range of “reasonable” parameters: \mathbf{p} such that $\chi^2(\mathbf{p}) \leq \chi^2(\mathbf{p}_0) + 1$

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statistical properties of the space of reasonable parameters:

average of an observable $A(\mathbf{p})$: $\bar{A} = \int d\mathbf{p} P(\mathbf{p}) A(\mathbf{p})$

variance (uncertainty of prediction): $\overline{\Delta^2 A} = \int d\mathbf{p} P(\mathbf{p}) (A(\mathbf{p}) - \bar{A})^2$.

covariance: $\overline{\Delta A \Delta B} = \int d\mathbf{p} P(\mathbf{p}) (A(\mathbf{p}) - \bar{A})(B(\mathbf{p}) - \bar{B})$.

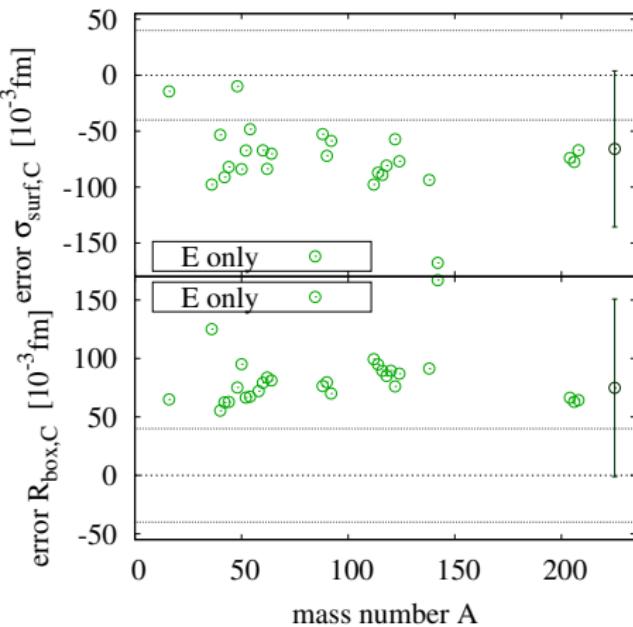
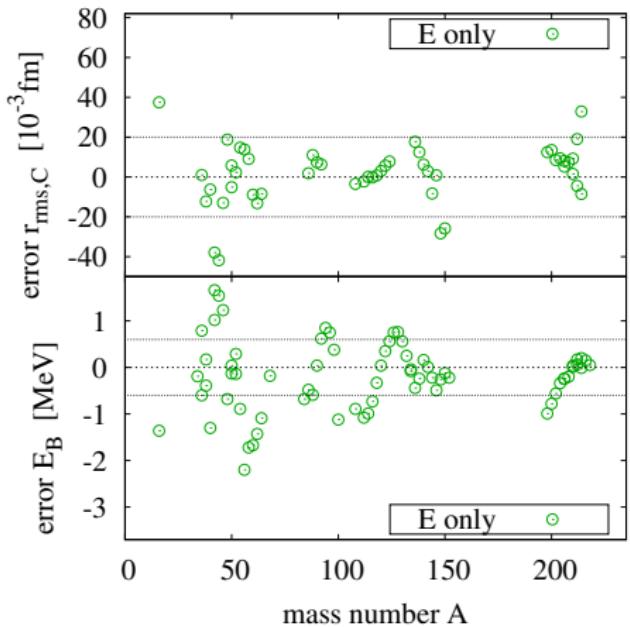
coefficient of determination (CoD): $r_{AB}^2 = \frac{|\overline{\Delta A \Delta B}|^2}{\overline{\Delta^2 A} \overline{\Delta^2 B}}$ (statistical correlation)

$$r_{AB}^2 = 1 \leftrightarrow \text{correlated}, \quad r_{AB}^2 = 0 \leftrightarrow \text{uncorrel.}$$

2) Application examples

Example 1: Impact of radii in bulk ground state properties

Distribution of errors on E_B , $r_{\text{rms,C}}$, $R_{\text{box,C}}$, $\sigma_{\text{surf,C}}$ – fit to E_B only

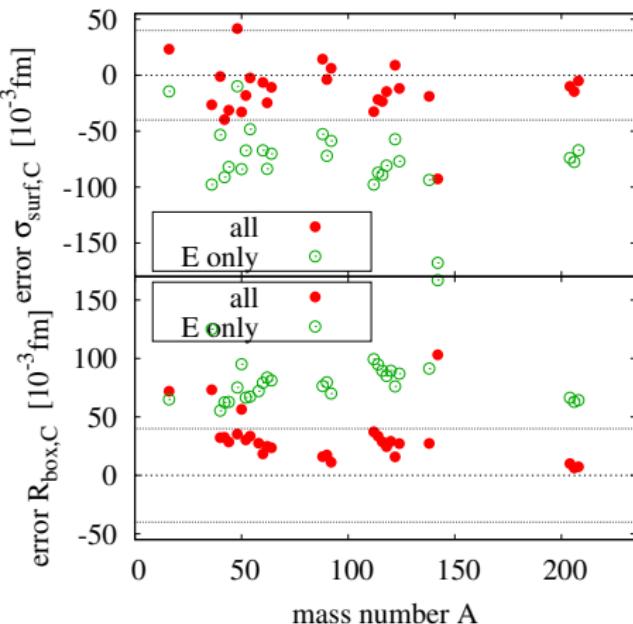
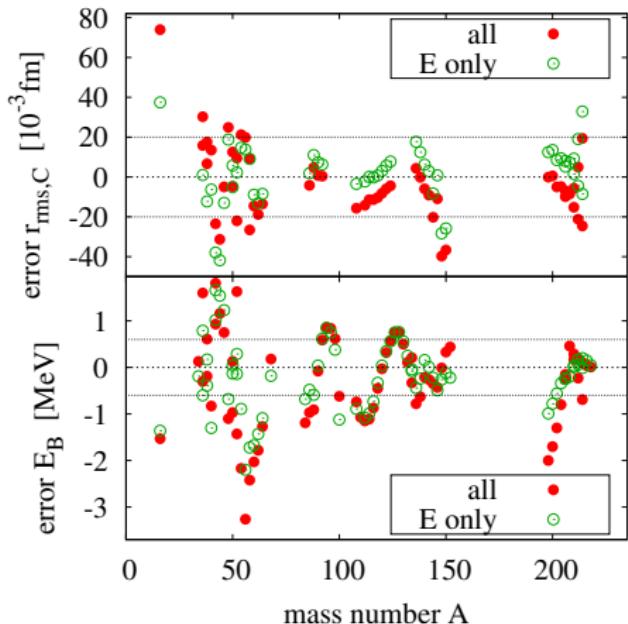


r.m.s. radius $r_{\text{rms,C}}$ well predicted by E_B fit

box radius $R_{\text{box,C}}$ and surface thickness $\sigma_{\text{surf,C}}$ systematically off the goal

large extrapolation uncertainty for $R_{\text{box,C}}$ and $\sigma_{\text{surf,C}}$ \Rightarrow can be adjusted within model

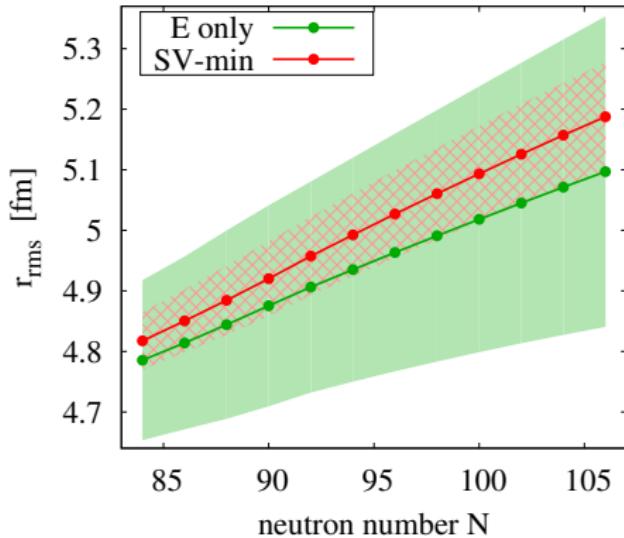
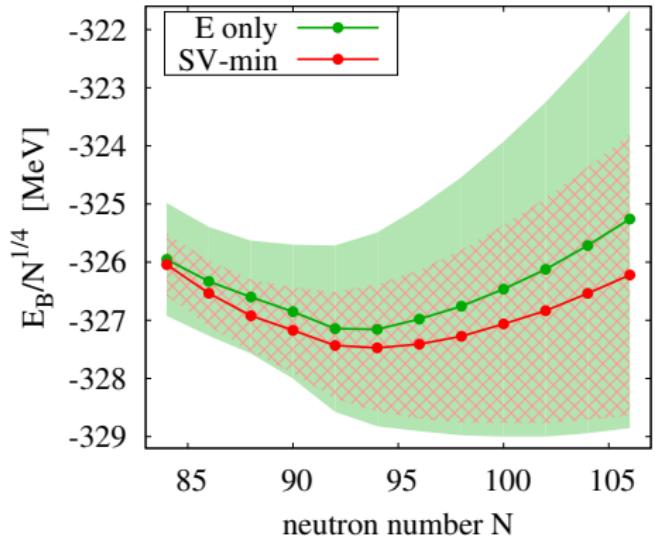
Distribution of errors on E_B , $r_{\text{rms,C}}$, $R_{\text{box,C}}$, $\sigma_{\text{surf,C}}$ – fit to all



include also $r_{\text{rms,C}}$, $R_{\text{box,C}}$, $\sigma_{\text{surf,C}}$ in the fit data

⇒ accomodate box radius $R_{\text{box,C}}$ and surf. thickn. $\sigma_{\text{surf,C}}$ at only small sacrifices in E_B

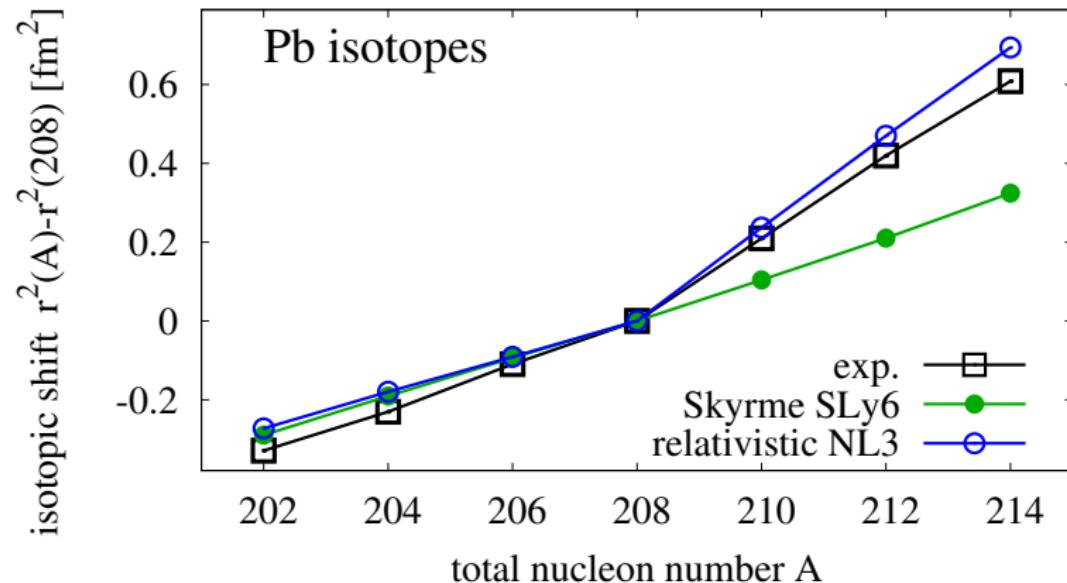
Extrapolation r -process nuclei – neutron rich Sn chain



fit to E_B only fixed model parameters too weakly \Rightarrow lousy extrapolation errors
large fit pool E_B , $r_{\text{rms,C}}$, $R_{\text{box,C}}$, $\sigma_{\text{surf,C}}$ \Rightarrow more reliable extrapolations
still demand for improvement \longleftrightarrow can more detailed radius data help?

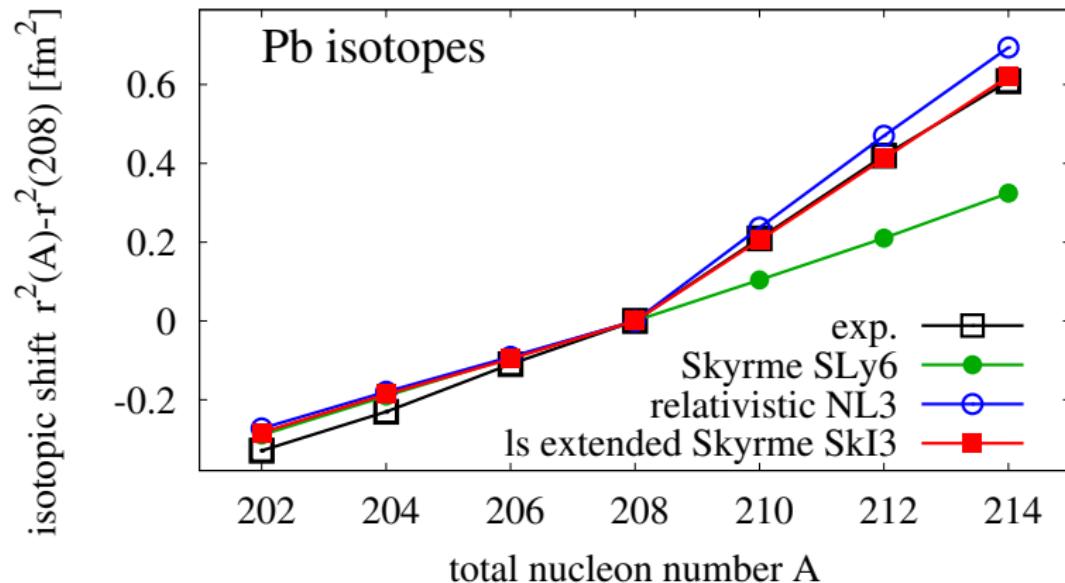
Example 2: Kink of r.m.s. radii at ^{208}Pb and spin-orbit term

Isotopic shifts $r_{\text{rms},C}^2(^A\text{Pb}) - r_{\text{rms},C}^2(^{208}\text{Pb})$ – problem with kink at ^{208}Pb



all Skyrme models up to year 1995 failed at the kink – RMF reproduced it at once
statistical analysis indicated sensitivity to pairing and spin-orbit term

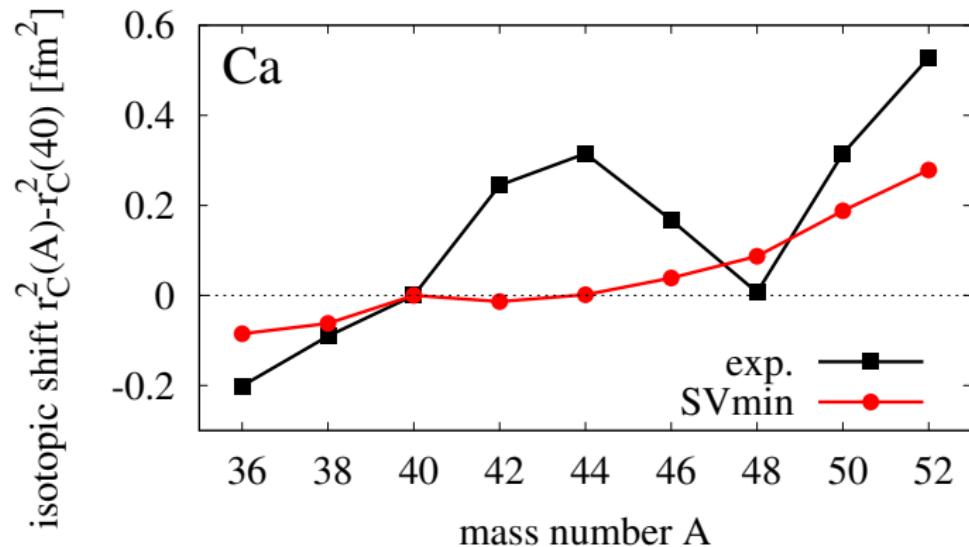
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non-relativistic limit of RMF revealed different isospin structure of spin-orbit term
extend Skyrme functional by isovector spin-orbit term \Rightarrow kink also reproduced

Example 3: Isotopic shifts and odd-even staggering in Ca isotopes

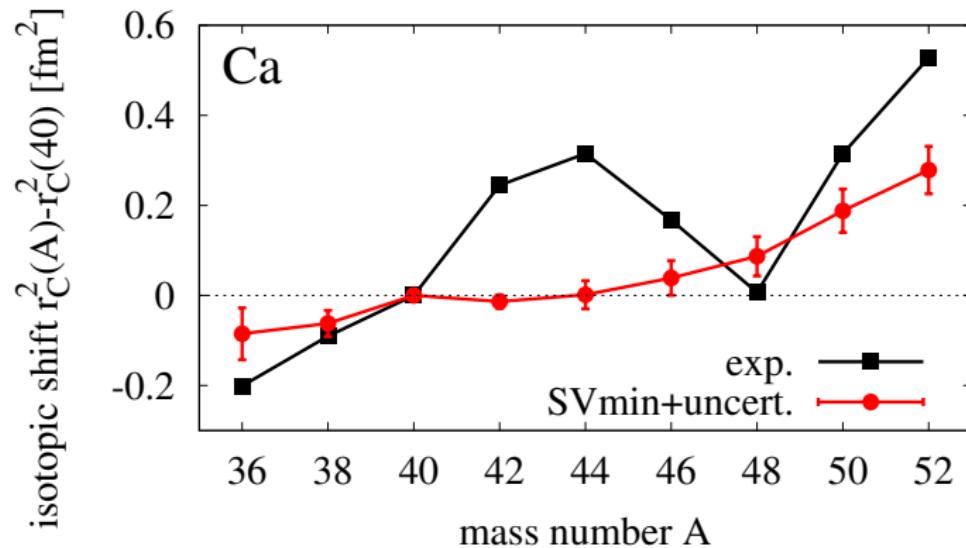
The problem: trend of charge r.m.s. radii in Ca chain



SVmin = fit of Skyrme functional with “traditional pairing” (contact force & density dep.)
⇒ theory averages nicely, but fails to reproduce the trend

Statistical uncertainty \leftrightarrow explore leeway of the Skyrme functional

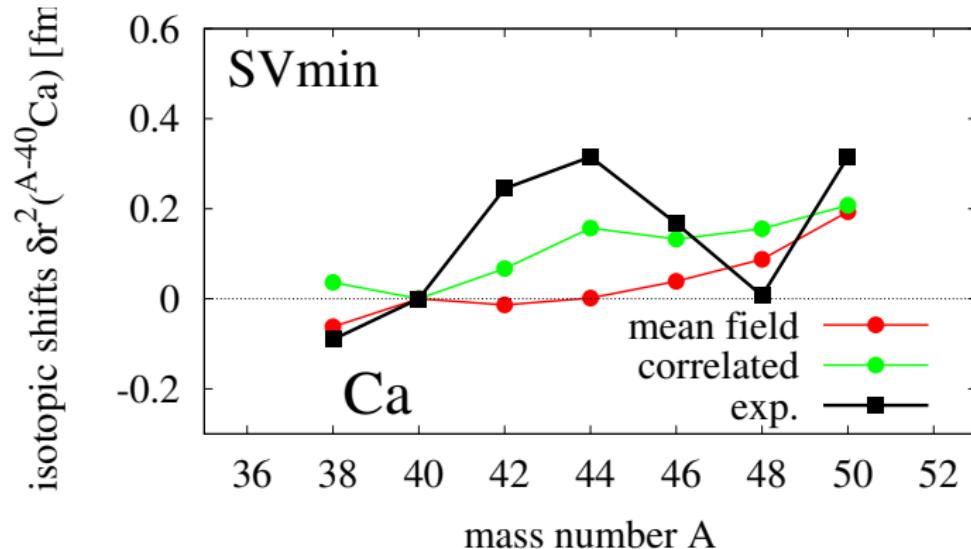
look at error bands to check chances for fitting the trend



extrapolation error too small \implies enforcing isotopic shifts will destroy overall quality

Impact of collective ground-state correlations

g.s. vibrations from low-lying 2^+ states contribute to radii

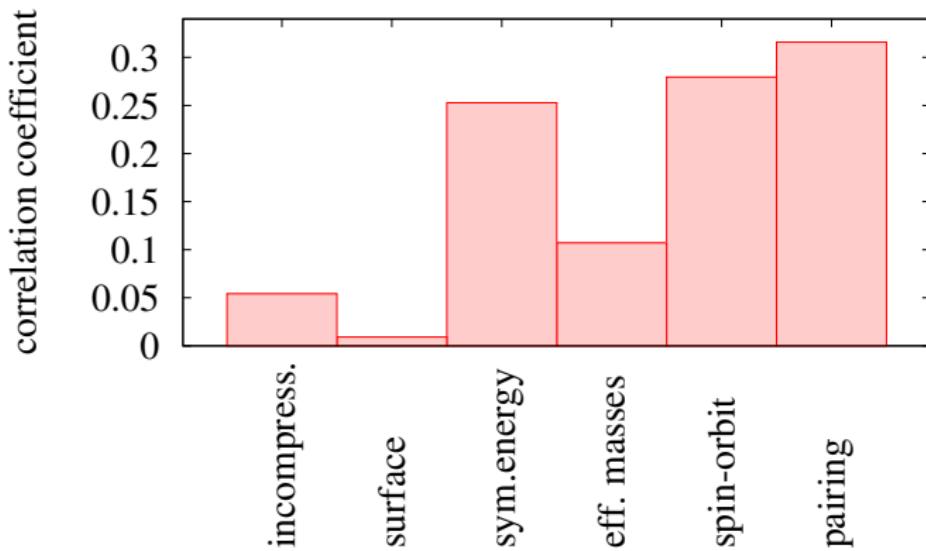


the effect is visible, but too small \Rightarrow the problem is the functional

Find most promising feature of functional

look at statistical correlations to find strongest lever

correlation with iso.-shift $\delta r^2(^{44-40}\text{Ca})$ for SV-min



symmetry energy:

fixed by polarizability

⇒ no!

spin-orbit splittings:

conceivable but limited changeability

⇒ unlikely!

pairing:

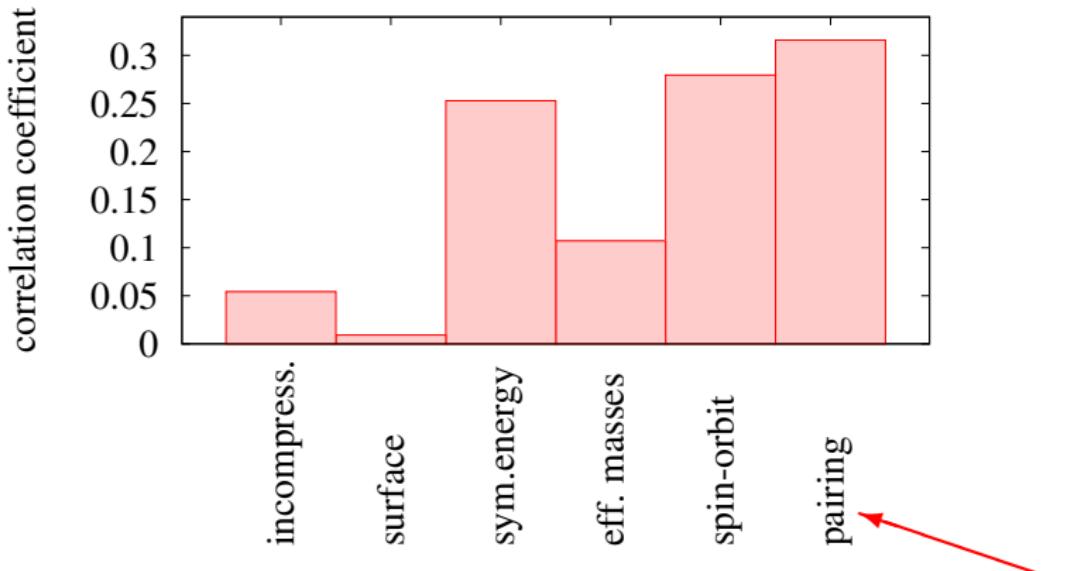
least well known, least well fixed

⇒ try that!

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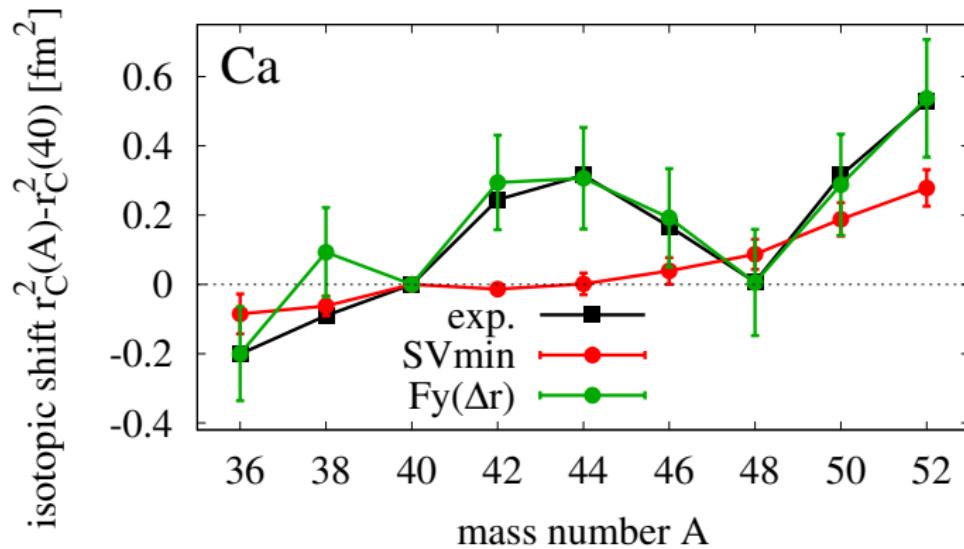
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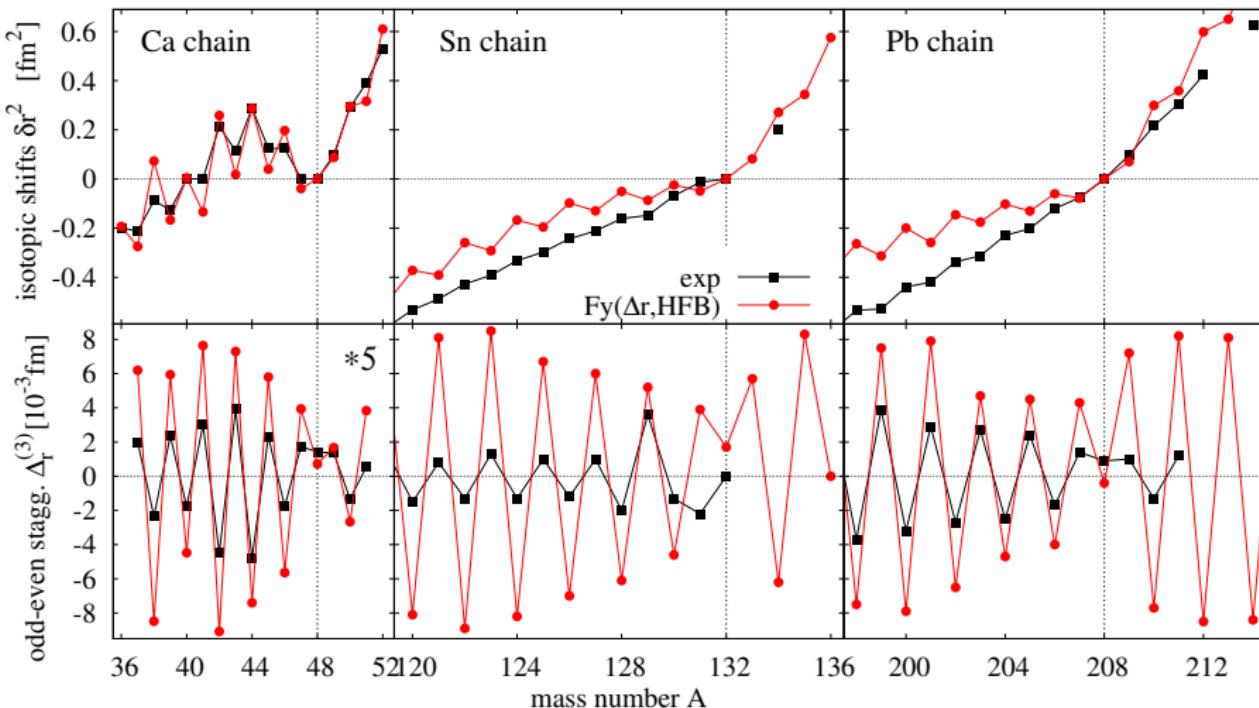
Fit isotopic trend with Fayans functional: $Fy(\Delta r, HFB)$ the solution



⇒ $Fy(\Delta r, HFB)$ reproduces trend for Ca isotopes almost perfectly
other isotopic chains? → to be checked

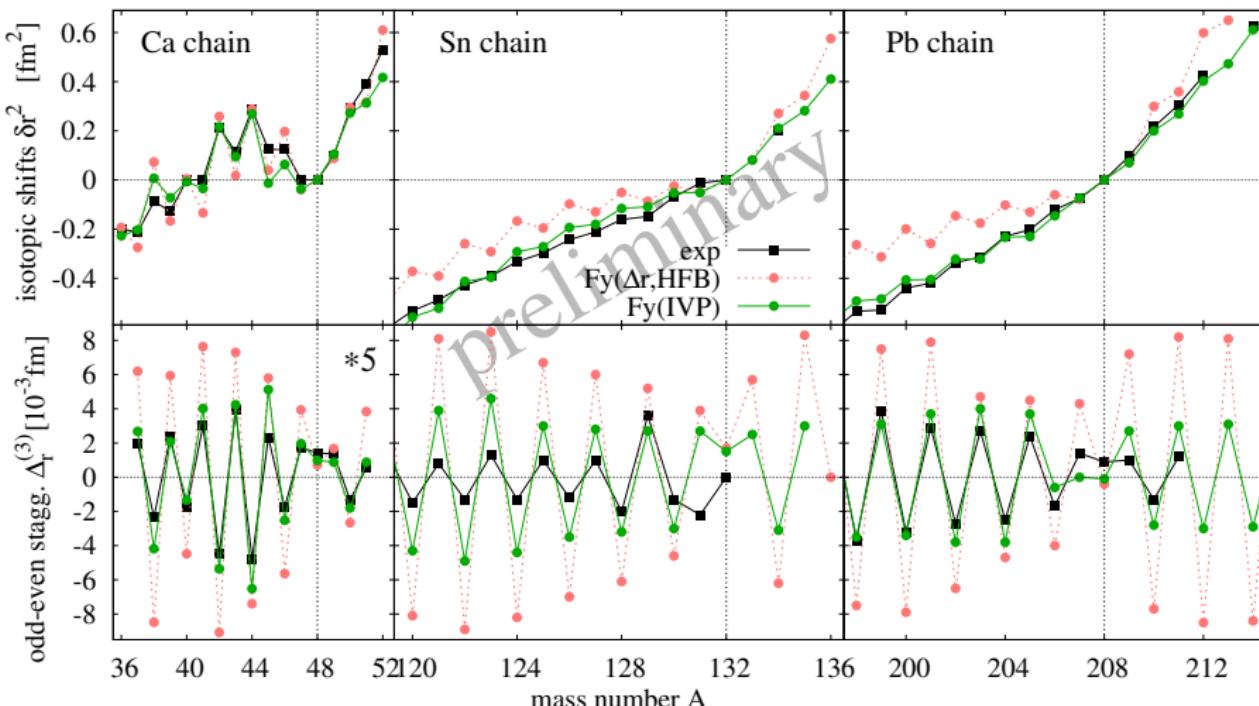
Example 4: Isotopic shifts and odd-even staggering in Sn & Pb isotopes

The problem: trends in Sn & Pb isotopes, odd-even staggerings



Fy(Δr , HFB) fails in: staggering $\Delta_r^{(3)}$, kinks at ^{132}Sn & ^{208}Pb , isot. trends Ca & Pb
 statistical analysis: all these observables \longleftrightarrow pairing
 yet missing in Fayans functional: isovector pairing (IVP) \implies try!

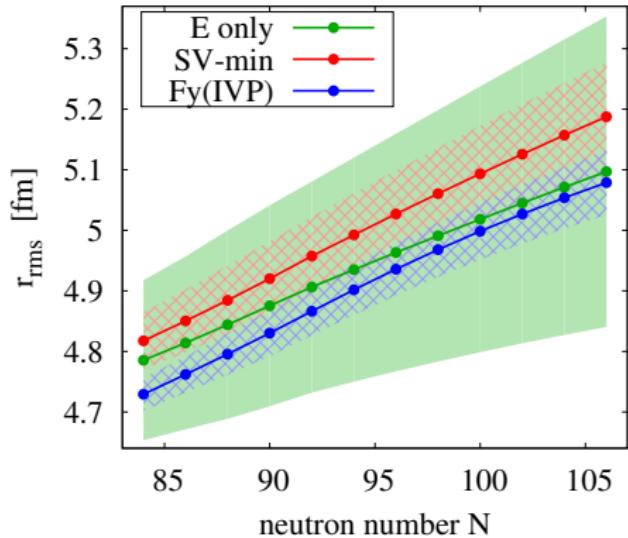
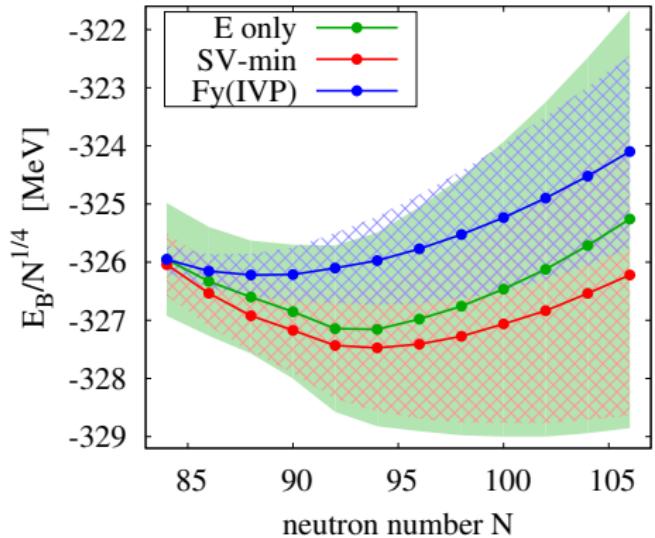
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fit functional with IVP adding data on $\Delta_r^{(3)}(\text{Ca})$ and isot.trends Ca&Pb \Rightarrow Fy(IPV)

Fy(IPV): solves nearly all problems, still slightly to large $\Delta_r^{(3)}(\text{Sn})$

Extrapolation r -process nuclei – neutron rich Sn chain



data on isotopic shifts & odd-even staggerings + extended model (Fayans)
⇒ somewhat more reduction of errors

r -process ≡ neutron rich nuclei ⇒ better information from neutron radius r_n ?

Example 5: Neutron radii and Pb Radius EXperiment (PREX)

Measuring the neutron radius

PREX = Pb Radius EXperiment:

scattering of high-energy polarized electrons (beam energy $E_{\text{in}} = 953 \text{ MeV}$)

\Rightarrow Parity-Violating Asymmetry $A_{\text{PV}}(q) \propto (\sigma_{\uparrow} - \sigma_{\downarrow})/\sigma_{\text{total}}$
at transferred momentum $q = 0.39/\text{fm}$

isovector dipole polarizability α_D :

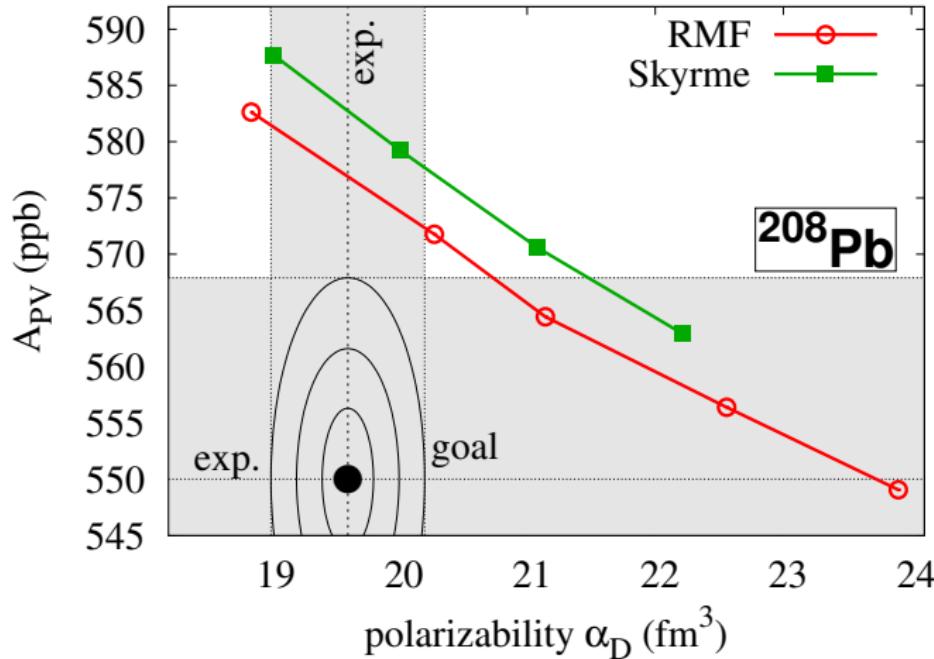
from photo-absorption strength σ_{γ} as $\alpha_D = \int_0^{\infty(E_{\text{max}})} dE E^{-2} \sigma_{\gamma}(E)$

correlations α_D , A_{PV} , **neutron radius** r_{neut} :

	A_{PV}	α_D	r_{neut}
A_{PV}	1	0.99	0.99
α_D	0.99	1	0.98
r_{neut}	0.99	0.98	1

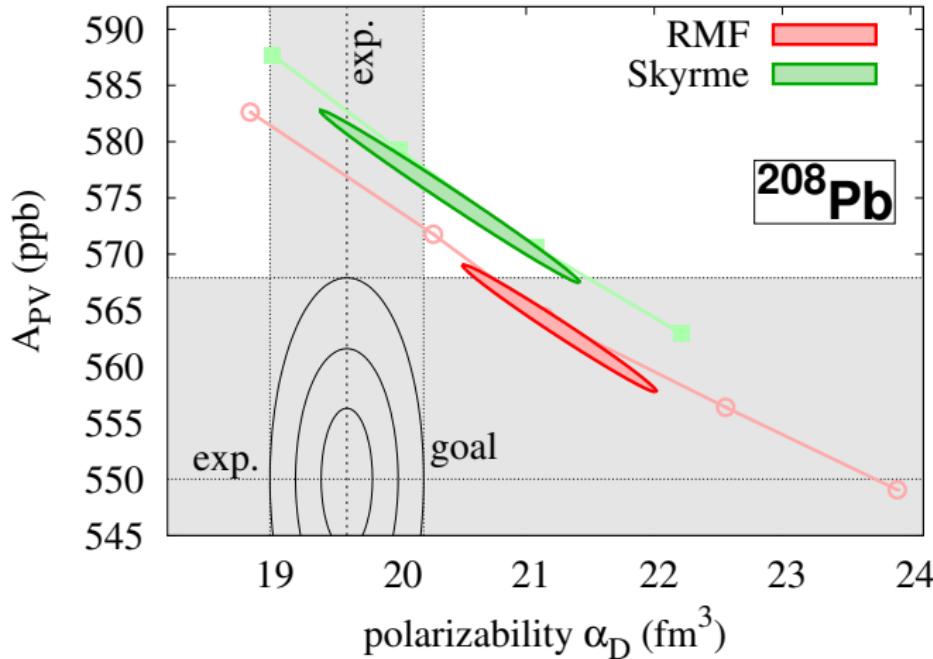
$\Rightarrow A_{\text{PV}}$ & α_D equivalent, model assisted measurement of r_{neut}

Fy(IPV) and extrapolation uncertainty to neutron matter



collection of Skyrme and RMF parametrizations line up to a linear trend
the trend avoids the matching point in plane of α_D and A_{PV}

Fy(IPV) and extrapolation uncertainty to neutron matter



the (1σ) uncertainty ellipsoids also follow the linear trends

incompatible data: either A_{PV} or α_D can be tuned in the given models

⇒ put the data point on hold, work on resolving the dilemma

Conclusions

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1. deliver valuable data for better confining models
2. challenge models by revealing missing features
3. more info from exotic nuclei needed to improve extrapolations
4. info on neutron radius highly desirable, still an enigma

