# How does kinetic theory remember about initial conditions? 

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## Introduction

- Initial conditions of kinetic theory transform into hydrodynamic information
- How is information from the pre-hydrodynamic stage kept?


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Answers in Bjorken flow described by BRSSS or Holography

Transients/Quasinormal modes/Non-hydrodynamic modes

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## Hydrodynamization in kinetic theory?

Similiarities between hydrodynamization at strong and weak coupling $\longrightarrow$ universal lessons

## Bjorken flow

## Highly symmetric flow

- Homogeneous and isotropic transverse to beam
- Boost-invariant in beam direction



## BRSSS in conformal Bjorken flow [Heller, Spalinski,1503.07514]

Pressure anisotropy $\mathcal{A}(w)=\frac{\mathcal{P}_{T}(\tau)-\mathcal{P}_{L}(\tau)}{\mathcal{P}(\tau)}$ as function of $w=\tau T$

- Governed by differential equation:

$$
C_{\tau_{\pi}} w\left(1+\frac{\mathcal{A}}{12}\right) \mathcal{A}^{\prime}+\left(\frac{C_{\tau_{\pi}}}{3}+\frac{C_{\lambda_{1}}}{8 C_{\eta}}\right) \mathcal{A}^{2}+\frac{3}{2} w \mathcal{A}-12 C_{\eta}=0
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Perturbations around the gradient expansion
$\mathcal{A}(w)+\delta \mathcal{A}(w) \rightarrow \delta \mathcal{A}(w) \sim e^{-\frac{3}{2 C_{T_{\pi}}} w} w^{\frac{C_{\eta}-2 C_{\lambda_{1}}}{C_{\tau_{\pi}}}}$

## The gradient expansion diverges and reveals transients



Figure 1: [Heller, Spalinski,1503.07514]

$$
f_{k} \sim \frac{\Gamma(k+\beta)}{(-S)^{k}} \xrightarrow{\text { Borel transform }} \text { poles of the form } \frac{1}{(S-x)^{\beta+1}}
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$\xrightarrow{\text { Resummation }}$ contributions of the form $e^{-S w} w^{\beta}$

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Figure 1: [Heller, Spalinski,1503.07514]


Figure 2: [Basar, Dunne, 1509.05046]

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## Similar story in holography

- Gradient expansion in Bjorken flow diverges [Heller, Janik, Witaszczyk,1302.06979]
- Initial conditions: Specify metric on a timeslice

Quasinormal modes in Bjorken flow

$$
\delta \mathcal{E}_{i}=\sigma_{i} \tau^{\alpha_{i}} e^{-S_{i} \tau^{2 / 3}}
$$

## Borel plane in holography



Figure 3: [1302.0697,1707.02282]

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## Initial conditions, gradient expansion and non-hydrodynamic modes in BRSSS and Holography

- Divergence of gradient expansion revealed non-hydrodynamic modes
- Each mode has a free parameter set by initial conditions
- In BRSSS: one free parameter $\rightarrow$ one transient
- Holography: infinite number of transients


## Kinetic theory

Distribution function
$f(x, p)=$ number of particles at position x with momentum p .

Evolution described by Boltzmann equation
Free streaming $\longrightarrow \quad p^{\mu} \partial_{\mu} f(x, p)=\mathcal{C}[f] \quad \longleftarrow$ Collision kernel

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Simplifies in Bjorken flow and Relaxation time approximation (RTA)

$$
\partial_{\tau} f(\tau, p)=\frac{f_{e q}(\tau, p)-f(\tau, p)}{\tau_{\mathrm{rel}}}
$$

$$
f_{e q}(\tau, p) \propto e^{-\frac{E(\rho)}{T(\tau)}}(\text { we consider classical, massless particles })
$$

$$
\tau_{\text {rel }} \propto T^{-\Delta}, \quad \Delta<3
$$

Energy density satisfies an integral equation [Florkowski,Ryblewski,Strickland,1305.7234]

$$
\mathcal{E}(\tau)=D\left(\tau, \tau_{0}\right) \mathcal{E}^{0}(\tau)+\int_{\tau_{0}}^{\tau} \frac{\mathrm{d} \tau^{\prime}}{2 \tau_{\mathrm{rel}}\left(\tau^{\prime}\right)} D\left(\tau, \tau^{\prime}\right) H\left(\frac{\tau^{\prime}}{\tau}\right) \mathcal{E}\left(\tau^{\prime}\right),
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Damping: $D\left(\tau, \tau^{\prime}\right)=e^{-\int_{\tau^{\prime}}^{\tau} \frac{d t}{\tau_{\text {rel }}}}$ Initial conditions: $\mathcal{E}^{0}(\tau)$
$H(x)=x^{2}+\frac{\arctan \sqrt{\frac{1}{x^{2}}-1}}{\sqrt{\frac{1}{x^{2}-1}}}$


## Gradient expansion in conformal RTA kinetic theory

$$
\mathcal{E}(\tau) \propto \frac{1}{\tau^{4 / 3}}\left(1+e_{1} \frac{\tau_{\mathrm{rel}}}{\tau}+e_{2}\left(\frac{\tau_{\mathrm{rel}}}{\tau}\right)^{2}+\ldots\right)
$$



Figure 4: [Heller,Kurkela,Spalinski,VS,1609.04803]

## Gradient expansion for different $\Delta$



## Perturbations around gradient expansion

## Use ansatz

$\mathcal{E}(\tau)=\mathcal{E}_{g e}(\tau)+\sigma e^{-S \frac{\tau}{\tau_{\mathrm{rel}}} \tau^{\beta} \mathcal{E}_{\beta}(\tau)}$
Exponential decay rate $S$

- Only one purely decaying solution: $S$ is real.
- No off-axis singularities.


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$$
\begin{gathered}
\int_{0}^{1} \mathrm{~d} x H(x) x^{\beta}=0 \\
\frac{{ }_{3} F_{2}\left(1, \frac{\beta}{2}+2, \frac{\beta}{2}+2 ; \frac{\beta}{2}+\frac{5}{2}, \frac{\beta}{2}+3 ; 1\right)}{2 \beta^{2}+14 \beta+24}+\frac{1}{2(\beta+4)}=0
\end{gathered}
$$

## Allowed power laws $\beta$



Leading power law: $\beta=-3.43 \ldots$
Logarithmic oscillations: $\Re\left(\sigma \tau^{\beta}\right) \propto \tau^{\Re(\beta)} \cos (\theta+\Im(\beta) \log (\tau))$

Which contour to use in the integral equation?


## Which contour to use in the integral equation?



## Finding transients in numerical solutions

$$
\mathcal{E}(\tau)=D\left(\tau, \tau_{0}\right) \mathcal{E}^{0}(\tau)+\int_{\tau_{0}}^{\tau} \frac{\mathrm{d} \tau^{\prime}}{2 \tau_{\mathrm{rel}}} D\left(\tau, \tau^{\prime}\right) H\left(\frac{\tau^{\prime}}{\tau}\right) \mathcal{E}\left(\tau^{\prime}\right)
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Numerically solvable by iteration, starting with $\mathcal{E}^{0}$
[Phys.Lett. B224, 16 (1989), Banerjee, Bhalerao, Ravishankar]
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Can we study transients numerically?
Need $w \equiv \frac{\tau}{\tau_{\text {rel }}} \sim 100 \Rightarrow e^{-w} \sim 10^{-50}$

## Precision numerics with spectral methods


error $\propto \quad(\# \text { grid points) })^{-m}$

$e^{-\# \text { grid points }}$

## Subtracting solutions to see transients

Gradient expansion of pressure anisotropy $\mathcal{A}(w)=\frac{\mathcal{P}_{T}(\tau)-\mathcal{P}_{L}(\tau)}{\mathcal{P}(\tau)}$ is universal

Gradient expansion: $\mathcal{A}_{0}=\mathcal{A}$
Leading transient: $\mathcal{A}_{1}=\frac{\mathrm{d}}{\mathrm{d} w} \log \left(\mathcal{A}_{0}-\mathcal{A}_{0}^{\prime}\right)$
Subleading transient: $\mathcal{A}_{2}=\frac{\mathrm{d}}{\mathrm{d} w} \log \left(\mathcal{A}_{1}-\mathcal{A}_{1}^{\prime}\right)$

## Exponential decay of leading transient

$$
\mathcal{A}_{1}(w)=-1+\frac{\beta_{1}+7 / 3}{w}+\ldots
$$



## Power law of leading transient

$$
\beta_{1}+\ldots=w\left(\mathcal{A}_{1}(w)+1\right)-7 / 3
$$



## Subleading transient

$$
w \mathcal{A}_{2}(w) \approx-3.0271-0.5614 \tan (\theta+0.5614 \log (w))
$$

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## Signatures of off-axis modes at $\Delta=2.5$

## Pressure anisotropy



## Signatures of off-axis modes at $\Delta=2.5$

Log[Pressure anisotropy - truncated gradient expansion]


Two types of modes are extractable from the divergent behaviour of the gradient expansion


- Each $\sigma_{i}$ is a free parameter
- Off axis modes are fixed

$$
\begin{aligned}
\mathcal{E}(\tau)=\mathcal{E}_{g e}(\tau) & +e^{-S_{\text {real axis }} \frac{\tau}{\tau_{\text {rel }}}} \sum_{i} \sigma_{i} \tau^{\beta_{i}} \mathcal{E}_{\beta_{i}}(\tau) \\
& +e^{-S_{\text {off axis }} \frac{\tau}{\tau_{\text {rel }}}} \mathcal{E}_{\text {off axis }}(\tau)+\ldots
\end{aligned}
$$

## Outlook

Three ways to study transients

- Borel transform of gradient expansion
- Perturbations around gradient expansion
- Subtracting numerical solutions


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## Non-hydrodynamic sector in other models?

- More general relaxation time, e.g. momentum dependent...
- Other flows, e.g. Gubser...
- Beyond RTA: more realistic collision kernels, EKT for QCD...

