How does kinetic theory remember about initial conditions?

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Introduction

- Initial conditions of kinetic theory transform into hydrodynamic information
- How is information from the pre-hydrodynamic stage kept?

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Answers in Bjorken flow described by BRSSS or Holography Transients/Quasinormal modes/Non-hydrodynamic modes

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Hydrodynamization in kinetic theory?

Similiarities between hydrodynamization at strong and weak coupling \longrightarrow universal lessons

Bjorken flow

Highly symmetric flow

- Homogeneous and isotropic transverse to beam
- Boost-invariant in beam direction



BRSSS in conformal Bjorken flow [Heller, Spalinski,1503.07514]

Pressure anisotropy $\mathcal{A}(w) = \frac{\mathcal{P}_T(\tau) - \mathcal{P}_L(\tau)}{\mathcal{P}(\tau)}$ as function of $w = \tau T$

• Governed by differential equation: $C_{\tau_{\pi}}w(1+\frac{A}{12})\mathcal{A}' + \left(\frac{C_{\tau_{\pi}}}{3} + \frac{C_{\lambda_{1}}}{8C_{\eta}}\right)\mathcal{A}^{2} + \frac{3}{2}w\mathcal{A} - 12C_{\eta} = 0$

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• Formal solution in gradient expansion: $\mathcal{A}(w) = f_0 + \frac{f_1}{w} + \frac{f_2}{w^2} + \dots$

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Perturbations around the gradient expansion

$$\mathcal{A}(w) + \delta \mathcal{A}(w) \rightarrow \delta \mathcal{A}(w) \sim e^{-rac{3}{2C_{ au\pi}}w} w^{rac{C_{\eta}-2C_{\lambda_1}}{C_{ au\pi}}}$$

The gradient expansion diverges and reveals transients



Figure 1: [Heller, Spalinski,1503.07514]

$$f_k \sim rac{\Gamma(k+eta)}{(-S)^k} \xrightarrow[\text{Resummation}]{Borel transform}$$
 poles of the form $rac{1}{(S-x)^{eta+1}}$
 $\xrightarrow{ ext{Resummation}}$ contributions of the form $e^{-Sw}w^{eta}$

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- Gradient expansion in Bjorken flow diverges [Heller, Janik, Witaszczyk, 1302.06979]
- Initial conditions: Specify metric on a timeslice

Quasinormal modes in Bjorken flow

$$\delta \mathcal{E}_i = \sigma_i \, \tau^{\alpha_i} \, e^{-S_i \tau^{2/3}}$$

Borel plane in holography



Figure 3: [1302.0697,1707.02282]

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Initial conditions, gradient expansion and non-hydrodynamic modes in BRSSS and Holography

- Divergence of gradient expansion revealed non-hydrodynamic modes
- Each mode has a free parameter set by initial conditions
- In BRSSS: one free parameter \rightarrow one transient
- Holography: infinite number of transients

Kinetic theory

Distribution function f(x, p) = number of particles at position x with momentum p. Evolution described by Boltzmann equation

Free streaming $\longrightarrow p^{\mu}\partial_{\mu}f(x,p) = \mathcal{C}[f] \quad \leftarrow \text{Collision kernel}$

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Simplifies in Bjorken flow and Relaxation time approximation (RTA)

$$\partial_{\tau} f(\tau, p) = rac{f_{eq}(\tau, p) - f(\tau, p)}{ au_{
m rel}}$$

 $f_{eq}(au, p) \propto e^{-rac{E(p)}{T(au)}}$ (we consider classical, massless particles) $au_{
m rel} \propto \mathcal{T}^{-\Delta}, \qquad \Delta < 3$

$$\mathcal{E}(\tau) = D(\tau, \tau_0) \mathcal{E}^0(\tau) + \int_{\tau_0}^{\tau} \frac{\mathrm{d}\tau'}{2\tau_{\mathrm{rel}}(\tau')} D(\tau, \tau') H\left(\frac{\tau'}{\tau}\right) \mathcal{E}(\tau'),$$

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 $H(x) = x^{2} + \frac{\arctan\sqrt{\frac{1}{x^{2}}-1}}{\sqrt{\frac{1}{x^{2}}-1}}$

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Gradient expansion in conformal RTA kinetic theory

$$\mathcal{E}(au) \propto rac{1}{ au^{4/3}} \left(1 + e_1 rac{ au_{\mathrm{rel}}}{ au} + e_2 \left(rac{ au_{\mathrm{rel}}}{ au}
ight)^2 + \ldots
ight)$$



Figure 4: [Heller,Kurkela,Spalinski,VS,1609.04803]

Gradient expansion for different Δ



[Heller, VS, 1802.08225]

Perturbations around gradient expansion

Use ansatz

$$\mathcal{E}(\tau) = \mathcal{E}_{ge}(\tau) + \sigma e^{-Srac{ au}{ au_{
m rel}}} au^{eta} \mathcal{E}_{eta}(au)$$

Exponential decay rate S

- Only one purely decaying solution: S is real.
- No off-axis singularities.

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$$\int_0^1 \mathrm{d}x H(x) \, x^\beta = 0$$

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$$\frac{{}_{3}F_{2}\left(1, \frac{\beta}{2} + 2, \frac{\beta}{2} + 2; \frac{\beta}{2} + \frac{5}{2}, \frac{\beta}{2} + 3; 1\right)}{2\beta^{2} + 14\beta + 24} + \frac{1}{2(\beta + 4)} = 0$$

Allowed power laws β



Leading power law: $\beta = -3.43...$ Logarithmic oscillations: $\Re(\sigma\tau^{\beta}) \propto \tau^{\Re(\beta)} \cos(\theta + \Im(\beta) \log(\tau))$

Which contour to use in the integral equation?



Which contour to use in the integral equation?



Finding transients in numerical solutions

$$\mathcal{E}(\tau) = D(\tau, \tau_0) \mathcal{E}^0(\tau) + \int_{\tau_0}^{\tau} \frac{\mathrm{d}\tau'}{2\tau_{\mathrm{rel}}} D(\tau, \tau') H\left(\frac{\tau'}{\tau}\right) \mathcal{E}(\tau')$$

Numerically solvable by iteration, starting with \mathcal{E}^0 [Phys.Lett. B224, 16 (1989), Banerjee, Bhalerao, Ravishankar] [1305.7234, Florkowski, Ryblewski, Strickland]

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> Can we study transients numerically? Need $w \equiv rac{ au}{ au_{
> m rel}} \sim 100 \Rightarrow e^{-w} \sim 10^{-50}$

Precision numerics with spectral methods



Gradient expansion of pressure anisotropy $\mathcal{A}(w) = \frac{\mathcal{P}_T(\tau) - \mathcal{P}_L(\tau)}{\mathcal{P}(\tau)}$ is universal

Gradient expansion:
$$A_0 = A$$

Leading transient: $A_1 = \frac{d}{dw} \log (A_0 - A'_0)$
Subleading transient: $A_2 = \frac{d}{dw} \log (A_1 - A'_1)$

Exponential decay of leading transient

$$\mathcal{A}_1(w) = -1 + \frac{\beta_1 + 7/3}{w} + \dots$$



Power law of leading transient



$$\beta_1+\ldots=w(\mathcal{A}_1(w)+1)-7/3$$

W

Subleading transient

 $w\mathcal{A}_2(w) \approx -3.0271 - 0.5614 \tan{(\theta + 0.5614 \log(w))}$

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Signatures of off-axis modes at $\Delta=2.5$



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Two types of modes are extractable from the divergent behaviour of the gradient expansion



$$\begin{split} \mathcal{E}(\tau) &= \mathcal{E}_{ge}(\tau) + e^{-S_{\text{real axis}}\frac{\tau}{\tau_{\text{rel}}}} \sum_{i} \sigma_{i} \tau^{\beta_{i}} \mathcal{E}_{\beta_{i}}(\tau) \\ &+ e^{-S_{\text{off axis}}\frac{\tau}{\tau_{\text{rel}}}} \mathcal{E}_{\text{off axis}}(\tau) + \dots \end{split}$$

Outlook

Three ways to study transients

- Borel transform of gradient expansion
- Perturbations around gradient expansion
- Subtracting numerical solutions

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Non-hydrodynamic sector in other models?

- More general relaxation time, e.g. momentum dependent...
- Other flows, e.g. Gubser...
- Beyond RTA: more realistic collision kernels, EKT for QCD...