

From jet quenching to recombination and hydrodynamics: solving $R_{AA}-v_2$ puzzle

In collaboration with: Wenbin Zhao, Weiyao Ke, Wei Chen and Tan Luo

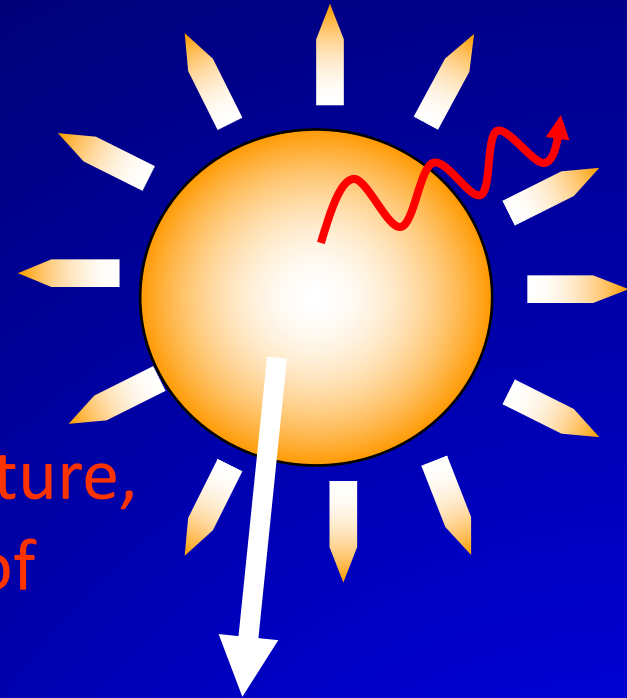
Xin-Nian Wang (LBNL)



Properties of QGP in A+A Collisions

Dynamical System:

- Soft probes: collective flow - bulk properties, EoS, transport properties
- EM Probes: EM emission – Temperature, EM response, medium modification of resonances
- Hard probes: Jet quenching – Jet transport coefficients



Collective flow of QGP

- Hydrodynamics: $\partial_\mu T^{\mu\nu} = 0$

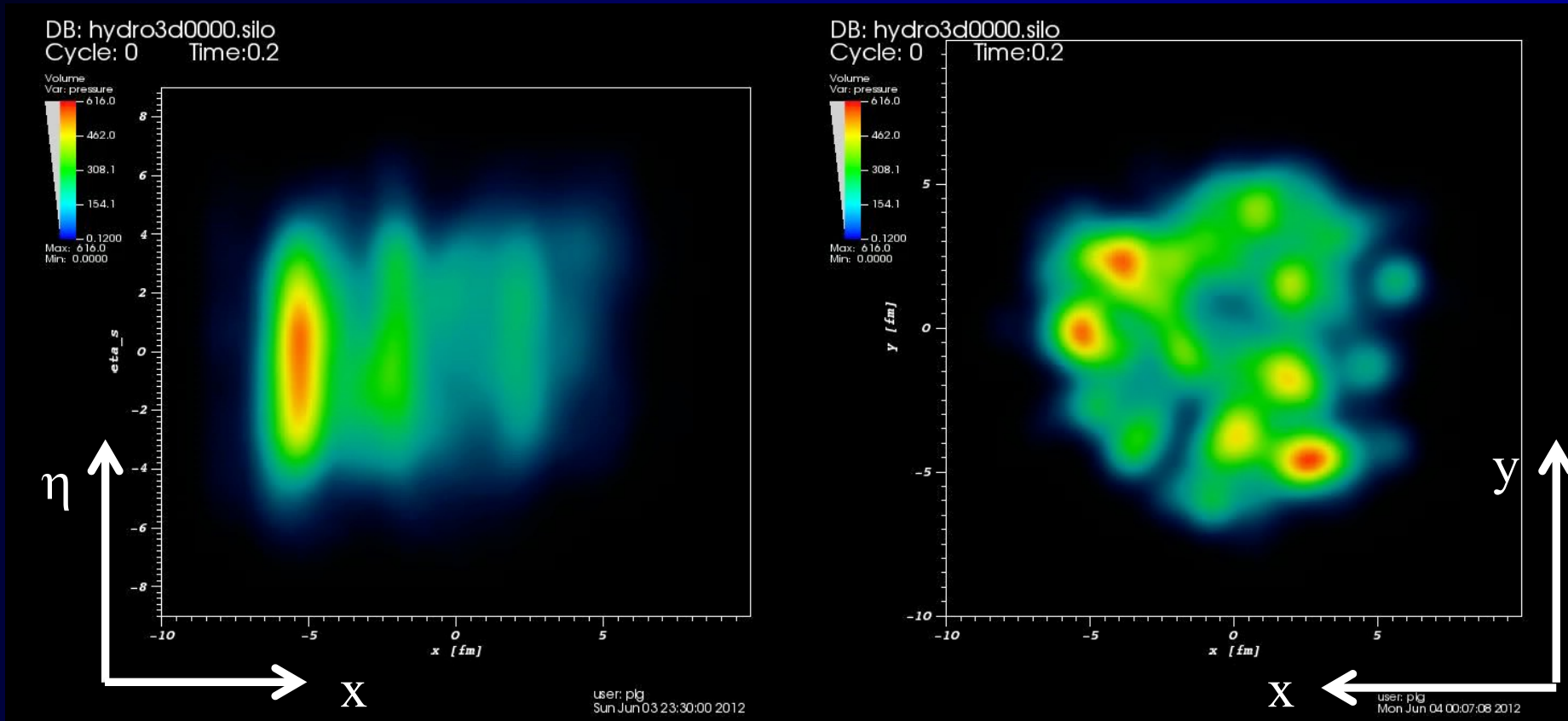
$$T^{\mu\nu} = (\epsilon + P)u^\mu u^\nu - P g^{\mu\nu} + \Delta T^{\mu\nu}$$

$$\Delta T^{\mu\nu} = \eta(\Delta^\mu u^\nu + \Delta^\nu u^\mu) + \left(\frac{2}{3}\eta - \zeta\right)H^{\mu\nu} \partial_\rho u^\rho$$

- a low-momentum effective theory
- Inputs from first principle QCD (lattice QCD)
EoS $p(\epsilon)$, transport coefficients $\xi(T)$, $\zeta(T)$
- Initial condition: parton prod. & thermalization

Anisotropic hydro expansion

with 3D fluctuating initial conditions



(3+1)D ideal hydro with AMPT initial condition (Pang & XNW'13)

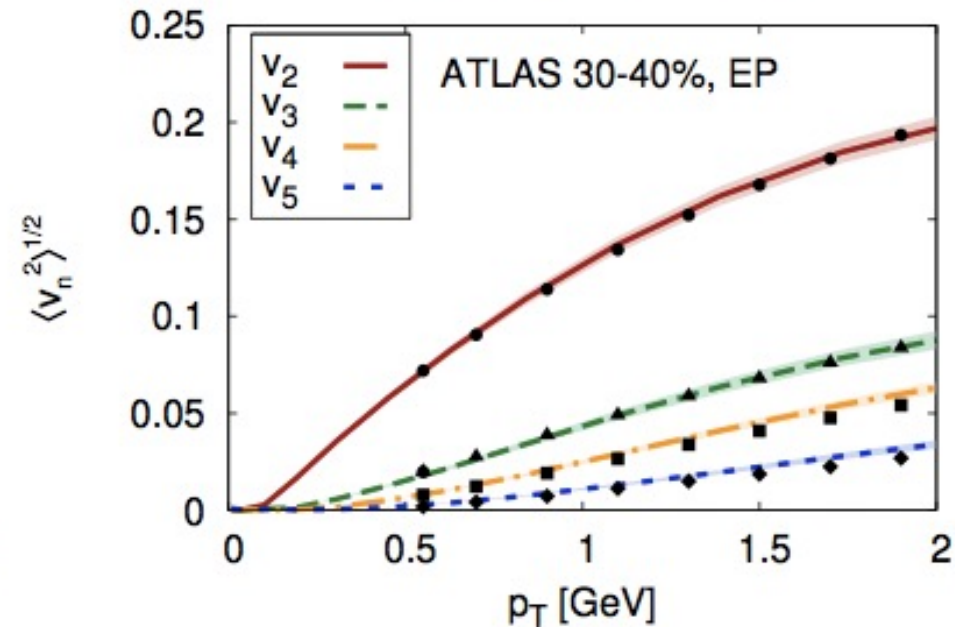
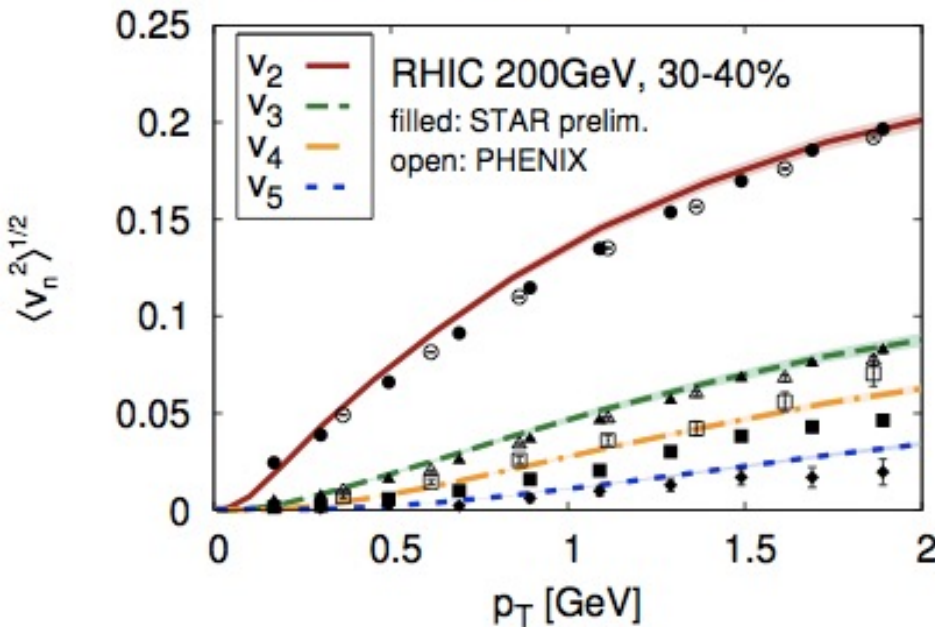
Viscosity of QGP in A+A collisions

Heinz & Song 2010

Gale, Jeon, Schenke, Tribedy & Venugopalan 2013

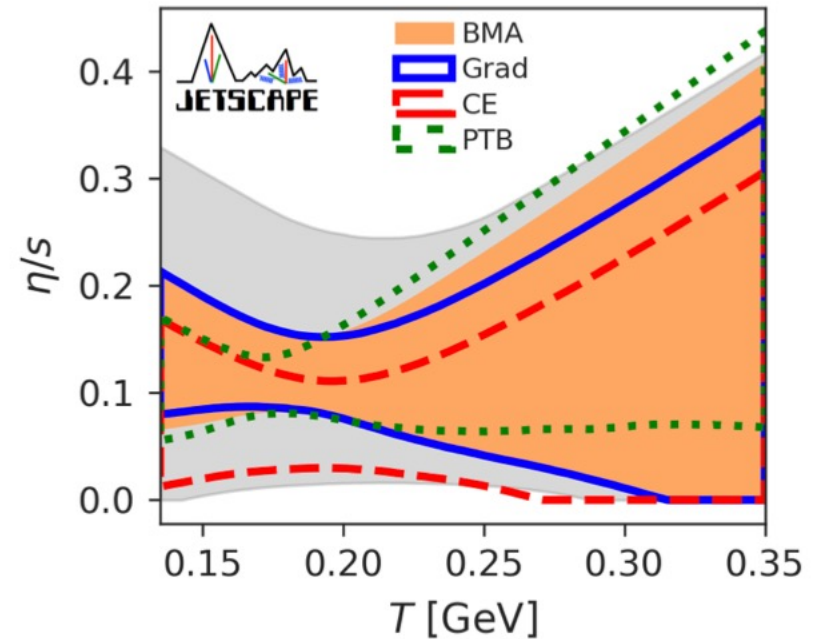
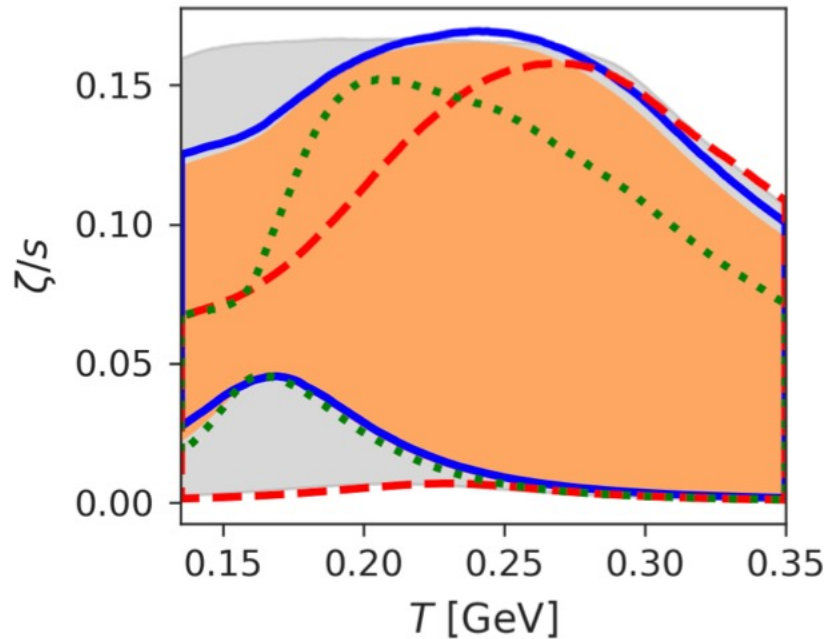
RHIC $\eta/s = 0.12$

LHC $\eta/s = 0.2$



Fluctuation + viscous hydro required to fit all v_n
Viscosity at LHC is larger than at RHIC

Extraction of bulk transport coefficients



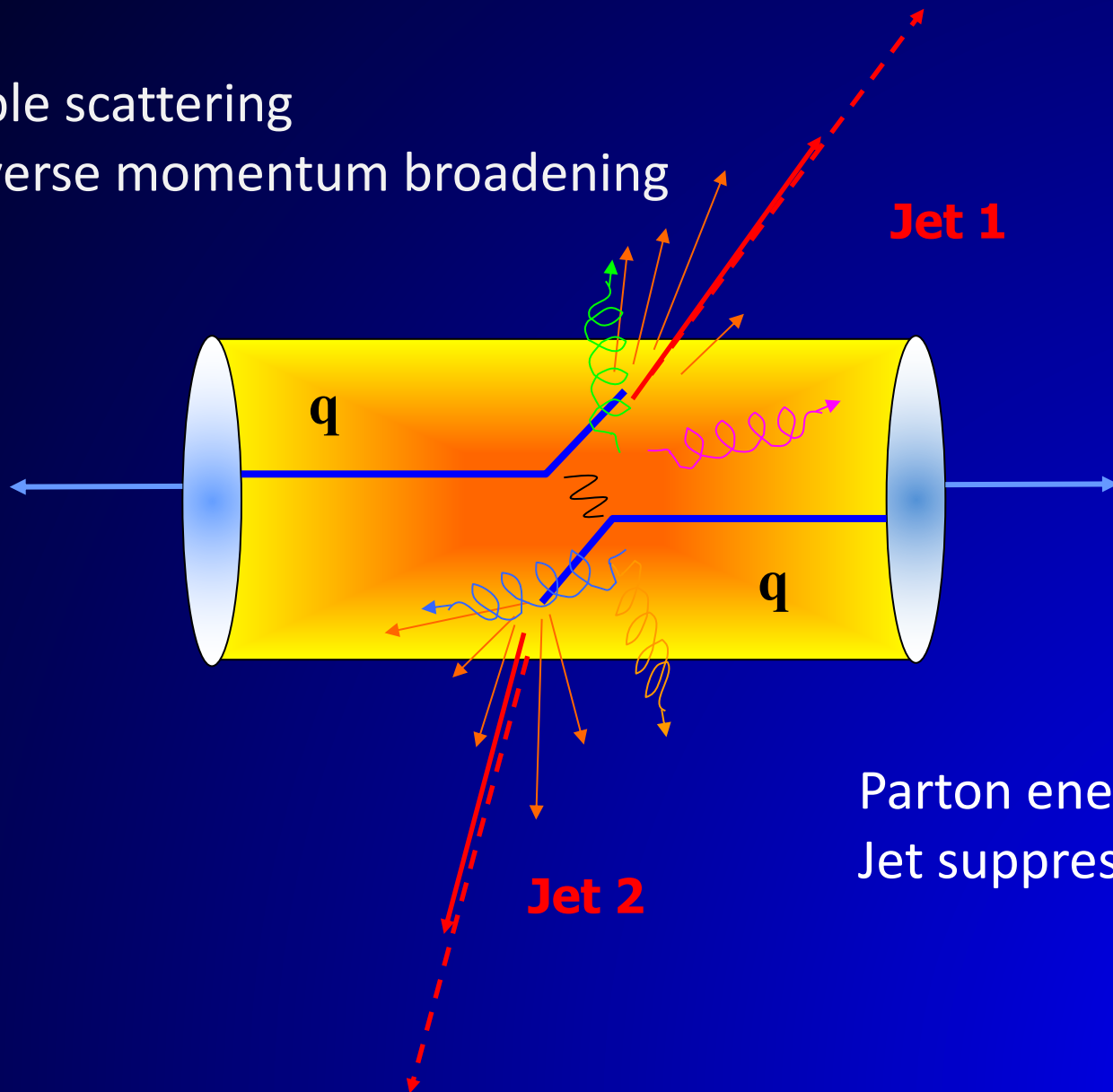
Constraints with Bayesian inference

JETSCAPE, Phys. Rev. Lett. 126, no.24, 242301 (2021)

Jets in heavy-ion collisions

Multiple scattering

Transverse momentum broadening



Jet 1

Jet 2

Parton energy loss
Jet suppression

Parton energy loss and jet transport

$$\frac{dE_{rad}}{dx} \approx \frac{\alpha_s N_c}{4} \hat{q} L$$

Radiative energy loss (BDMPS'96)

$$\frac{dE_{el}}{dx} = \int \frac{d^3k}{(2\pi)^3} dq_{\perp}^2 f(k) \frac{q_{\perp}^2}{2k} \frac{d\sigma}{dq_{\perp}^2} \approx \left\langle \frac{1}{2\omega} \right\rangle \hat{q}$$

Elastic energy loss

Jet transport coefficient:

$$\hat{q}(y) = \frac{4\pi^2 \alpha_s C_R}{N_c^2 - 1} \rho(y) x G(x) \Big|_{x \approx 0} = \frac{\langle q_{\perp}^2 \rangle}{\lambda}$$

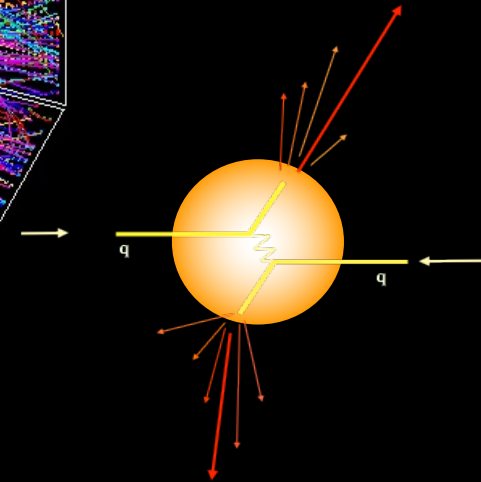
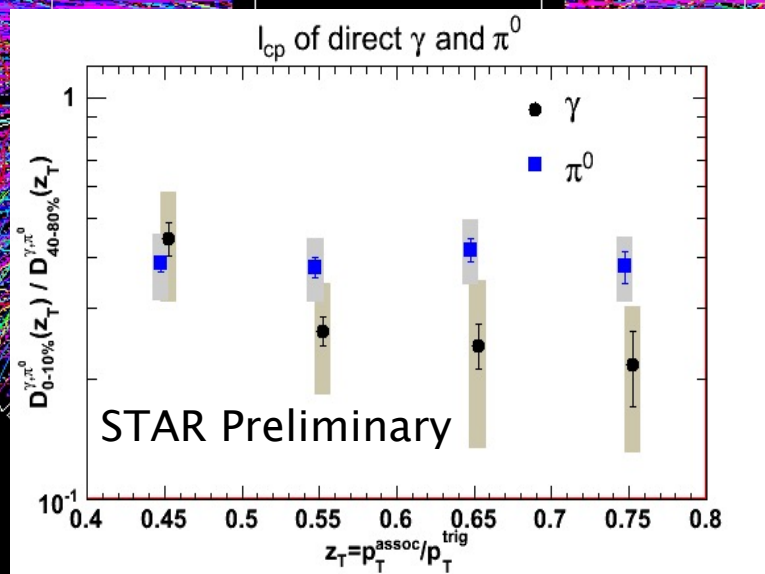
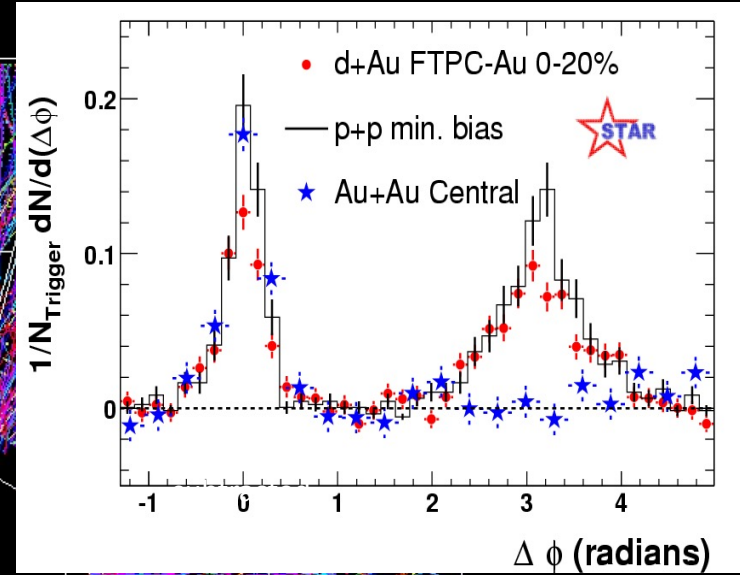
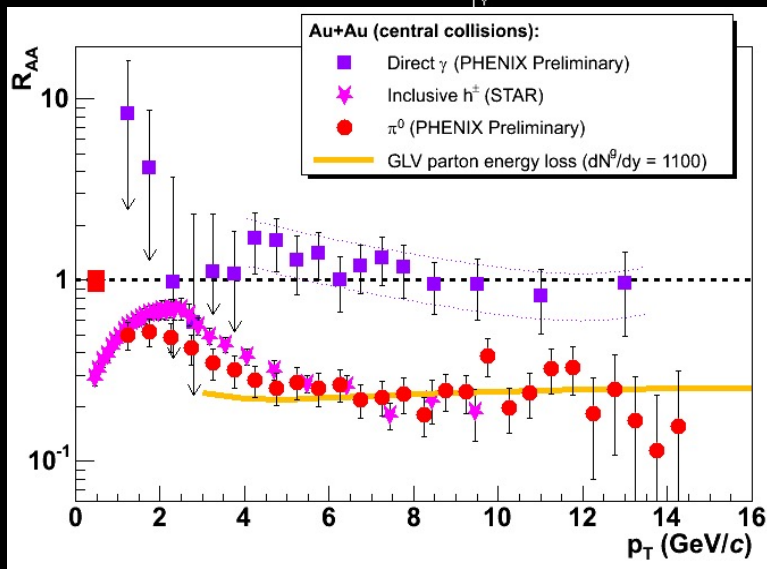
pQCD (BDMPS'96)

AdS/CFT (Liu, Rajagopal & Wideman'06)

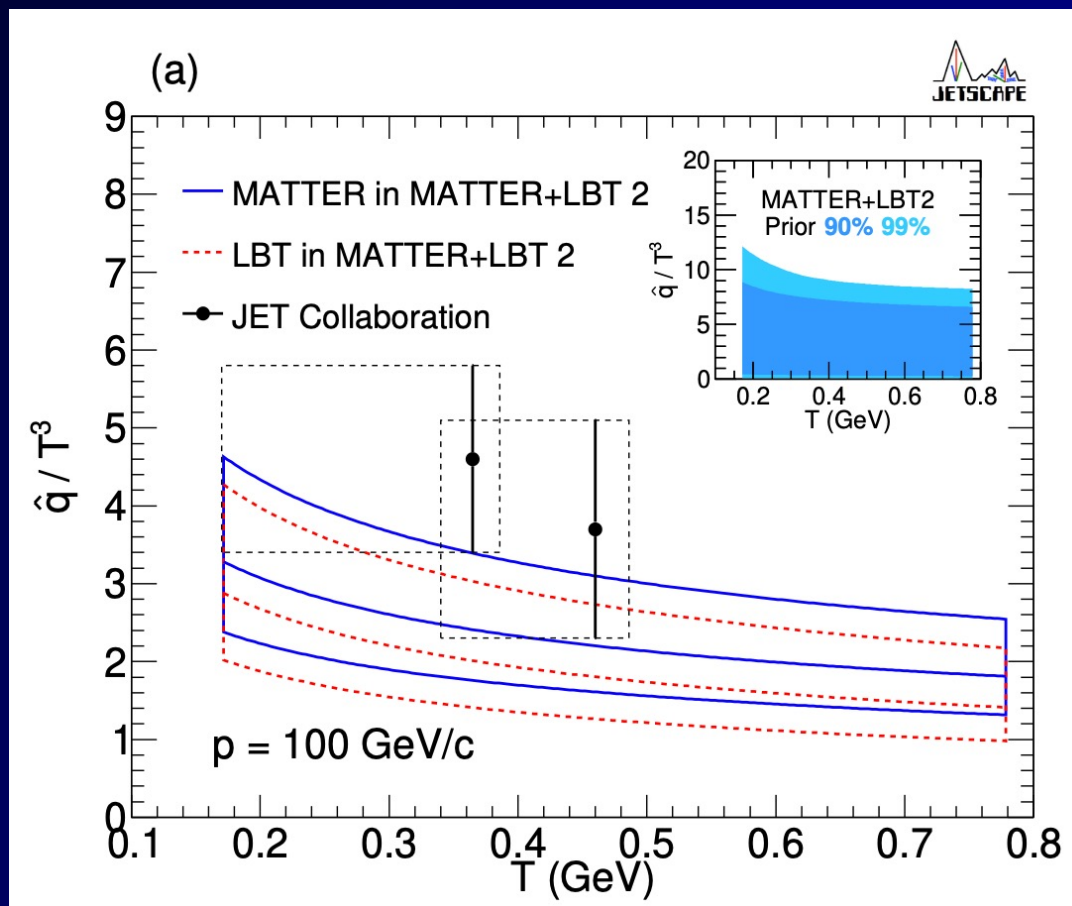
lattice QCD (Majumder'12)

Extract jet transport coefficient from parton energy loss

Jet Quenching phenomena at RHIC



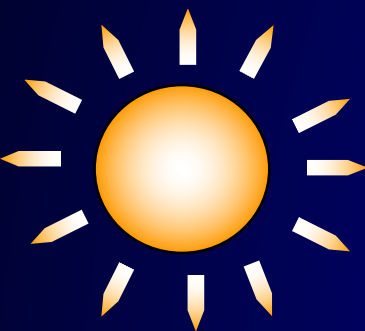
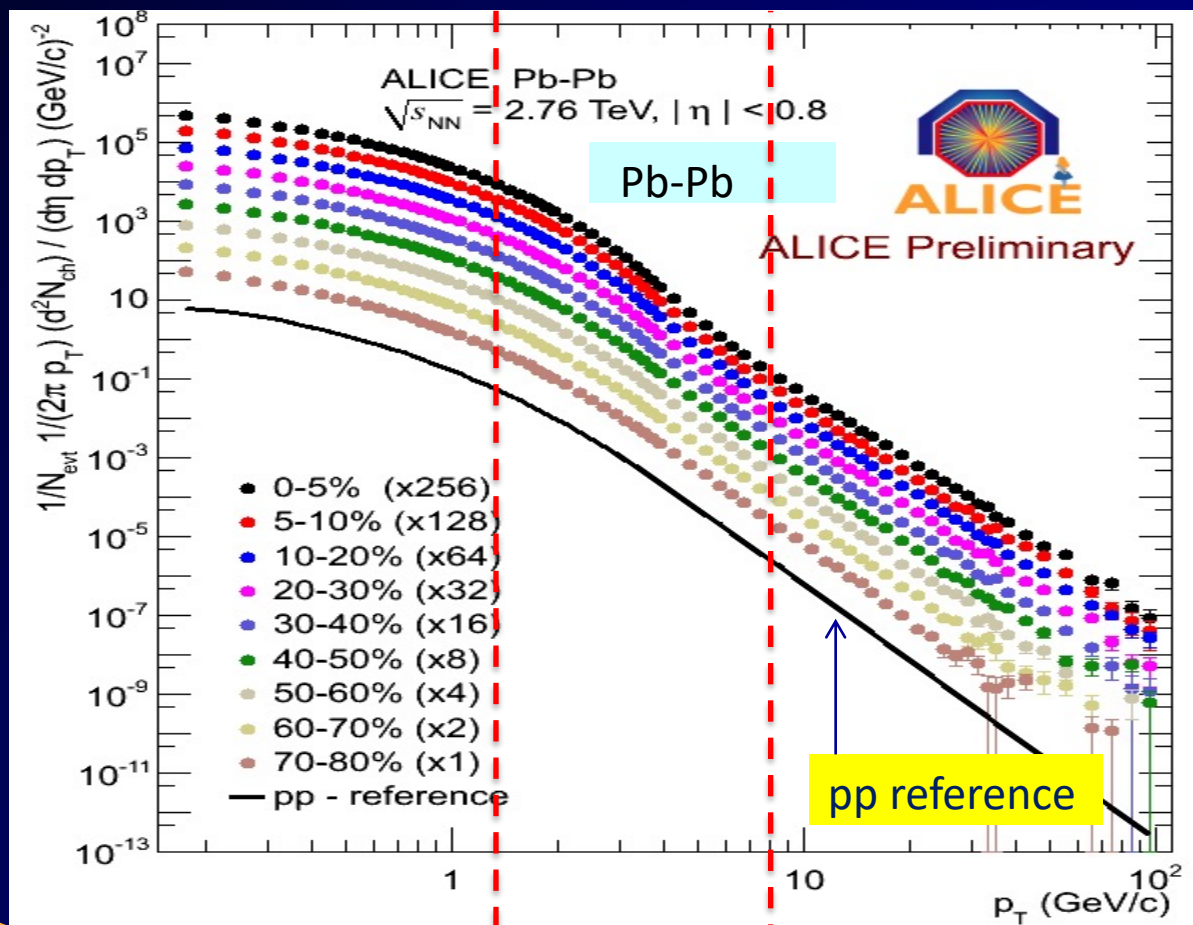
Extraction of jet transport coefficient



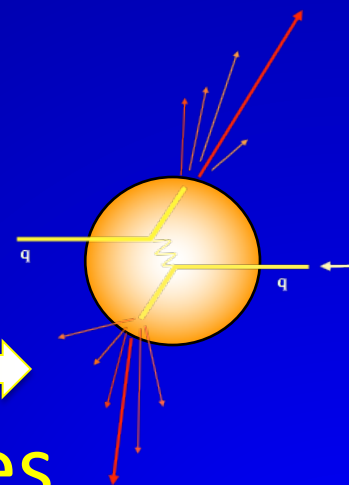
JET: *Phys.Rev.C* 90 (2014) 1, 014909

JETScape: [2102.11337](https://arxiv.org/abs/2102.11337)

Parton energy loss & medium response



soft response



hard probes

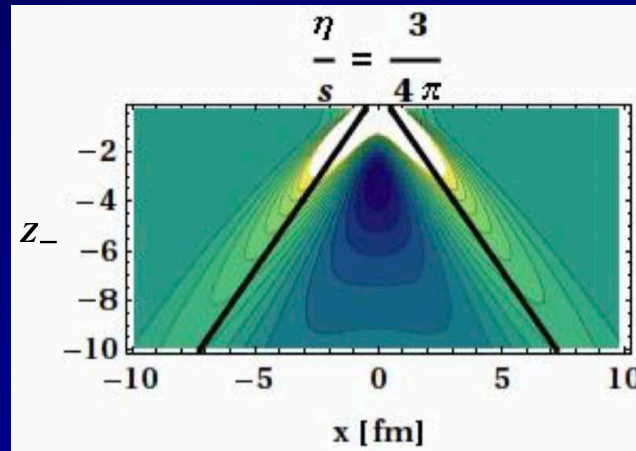
Physics at Intermediate pT

Jet-induced medium excitation

Casalderrey-Solana, Shuryak & Teaney (2005),
Stoecker (2005)

Jet induced Mach-cone in QGP

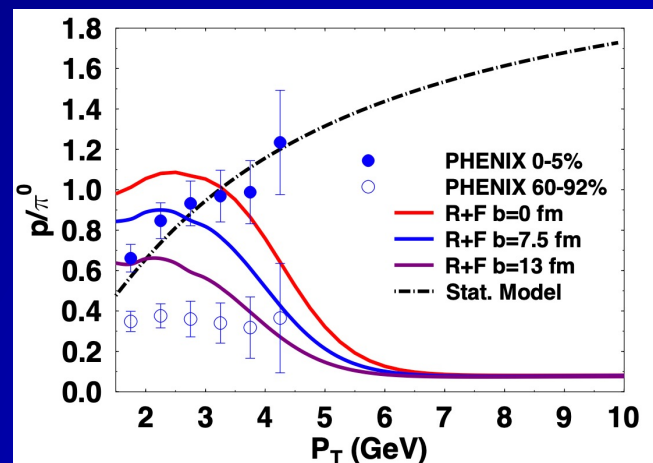
$$v = p/E > c_s$$



Ruppert & Muller (2005)

Parton recombination

Bayron/meson ratio,
NCQ scaling of v_2



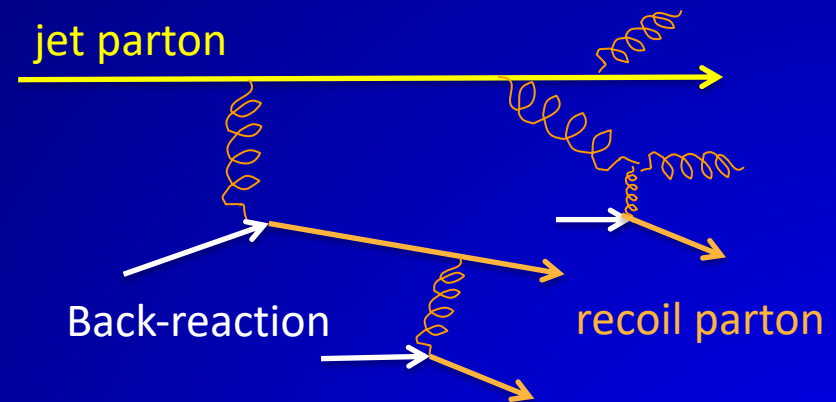
Fries, Muller, Nonaka and Bass, Phys.
Rev. C 68, 044902 (2003)

LBT: Linear Boltzmann Transport

$$p_1 \cdot \partial f_1 = - \int dp_2 dp_3 dp_4 (f_1 f_2 - f_3 f_4) |M_{12 \rightarrow 34}|^2 (2\pi)^4 \delta^4 \left(\sum_i p_i \right) + \text{inelastic}$$

Induced radiation $\frac{dN_g}{dz d^2 k_\perp dt} \approx \frac{2C_A \alpha_s}{\pi k_\perp^4} P(z) \hat{q} (\hat{p} \cdot u) \sin^2 \frac{k_\perp^2 (t - t_0)}{4z(1-z)E}$

- pQCD elastic and radiative processes (high-twist)
- Transport of medium recoil partons (and back-reaction)



CoLBT-hydro

(Coupled Linear Boltzmann Transport hydro)

Concurrent and coupled evolution of bulk medium and jet showers

$$p \cdot \partial f(p) = -C(p) \quad (p \cdot u > p_{cut}^0)$$

$$\partial_\mu T^{\mu\nu}(x) = j^\nu(x)$$

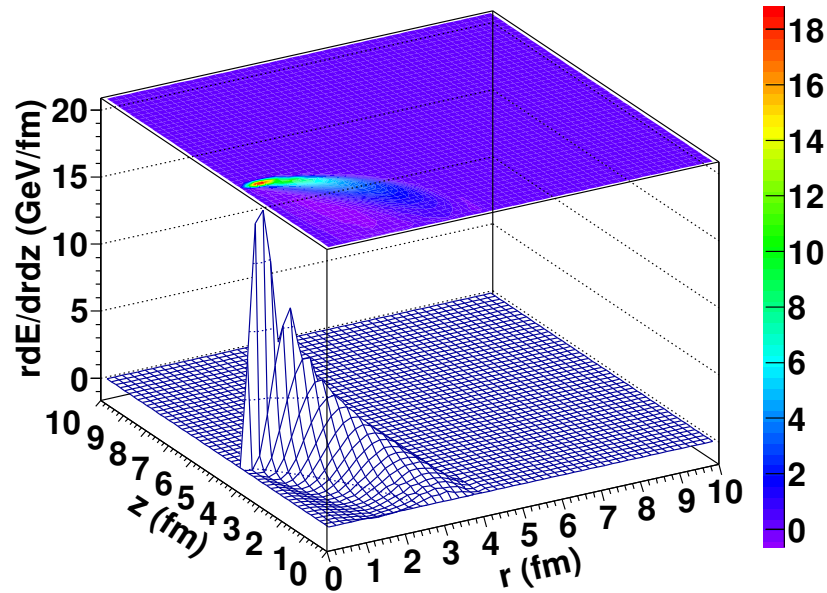
$$j^\nu(x) = \sum_i p_i^\nu \delta^{(4)}(x - x_i) \theta(p_{cut}^0 - p \cdot u)$$

- LBT for energetic partons (jet shower and recoil)
- Hydrodynamic model for bulk and soft partons: CLVisc
- Parton coalescence (thermal-shower)+ jet fragmentation
- Hadron cascade using UrQMD

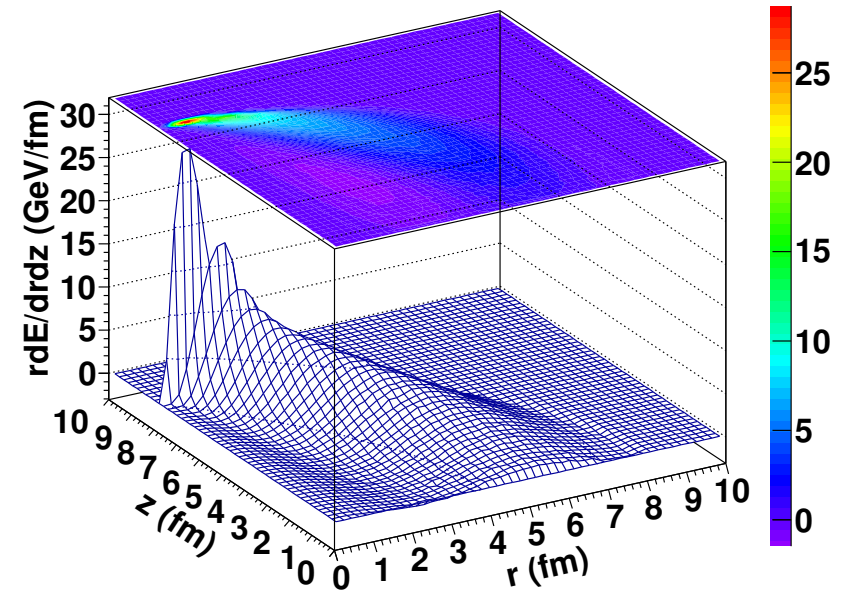
Chen, Cao, Luo, Pang & XNW, PLB777(2018)86

LBT: Jet-induced medium response

(a) $t=4$ fm/c



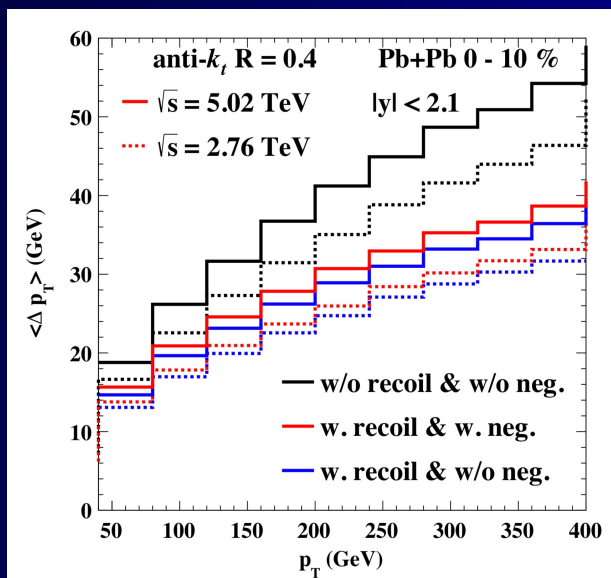
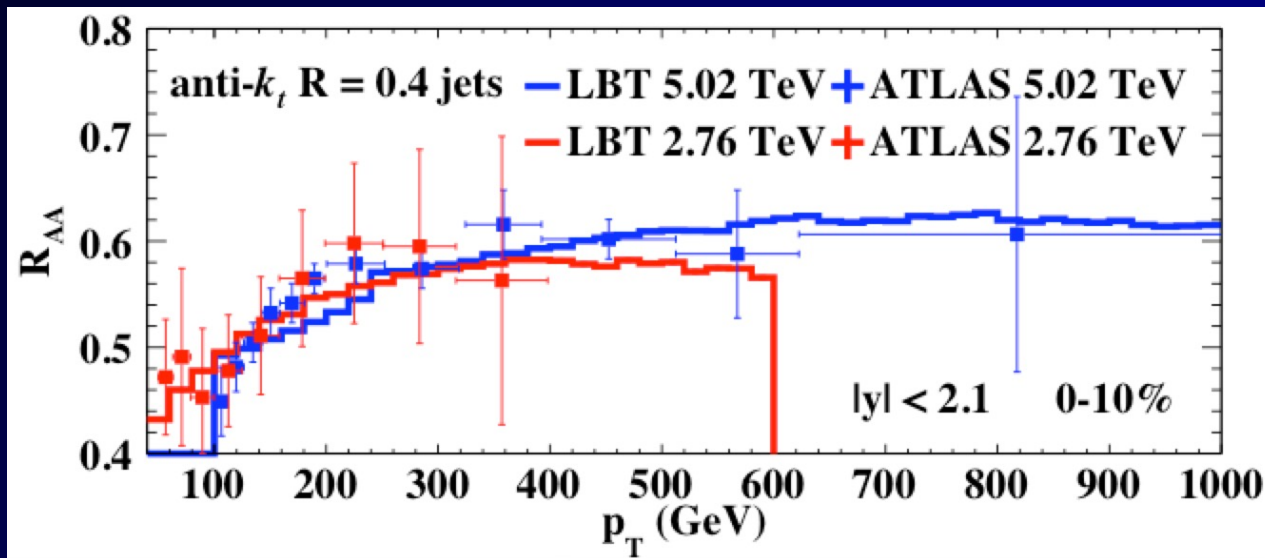
(b) $t=8$ fm/c



Energy distr. of medium response in a static medium

He, Luo, XNW & Zhu, PRC91 (2015) 054908

Jet suppression and energy loss

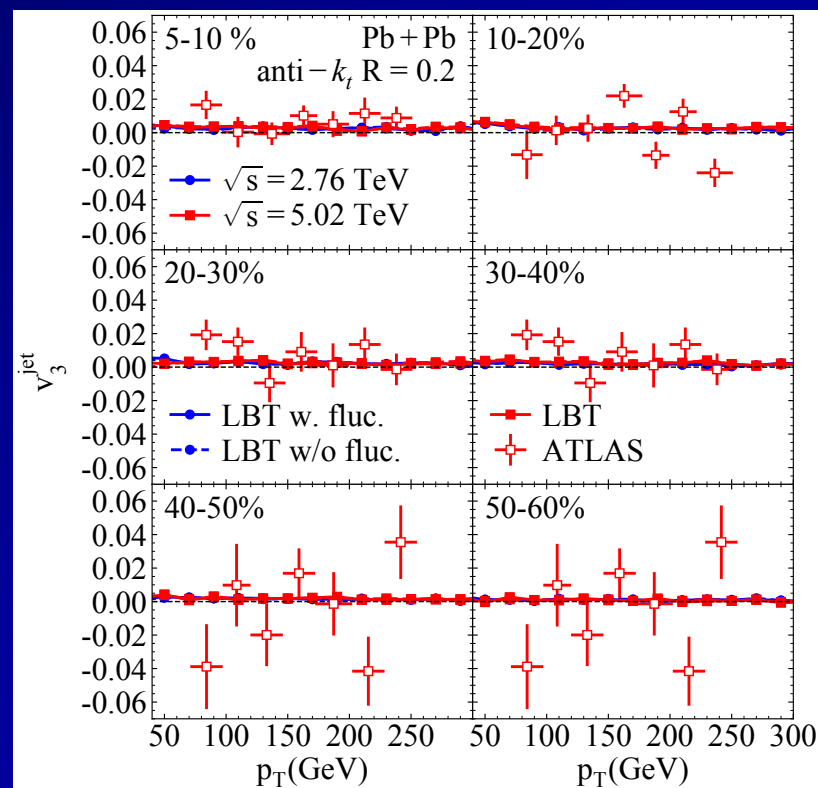
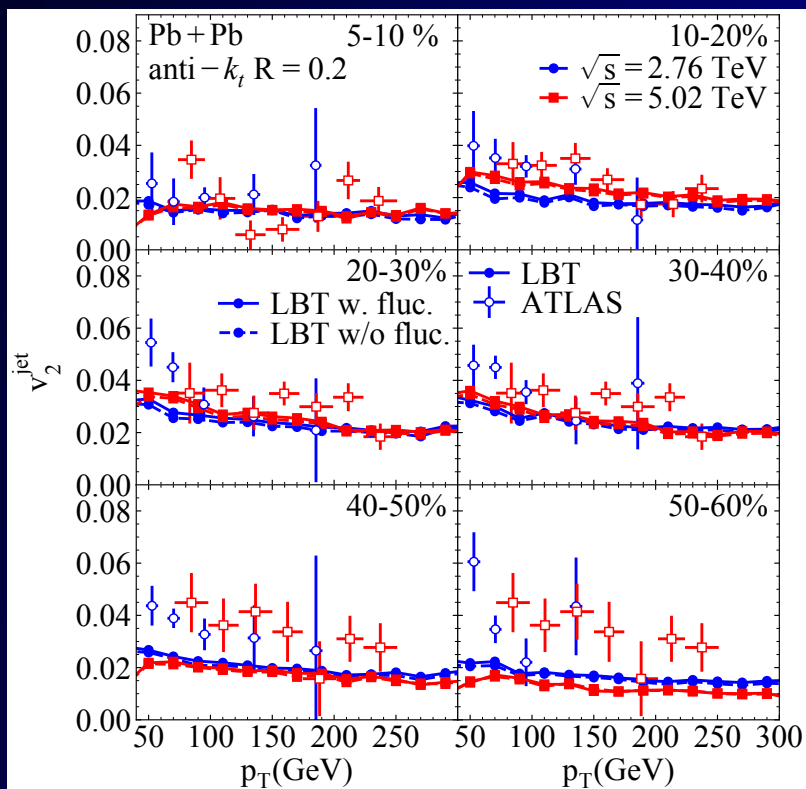


He, Cao, Chen, Luo, Pang & XNW 1809.02525

- Weak p_T dependence: initial jet spectra and p_T dependence of energy loss ΔE
- Weak energy dependence: increase of jet energy loss and the slope of initial spectra
- Medium response reduce jet net energy loss

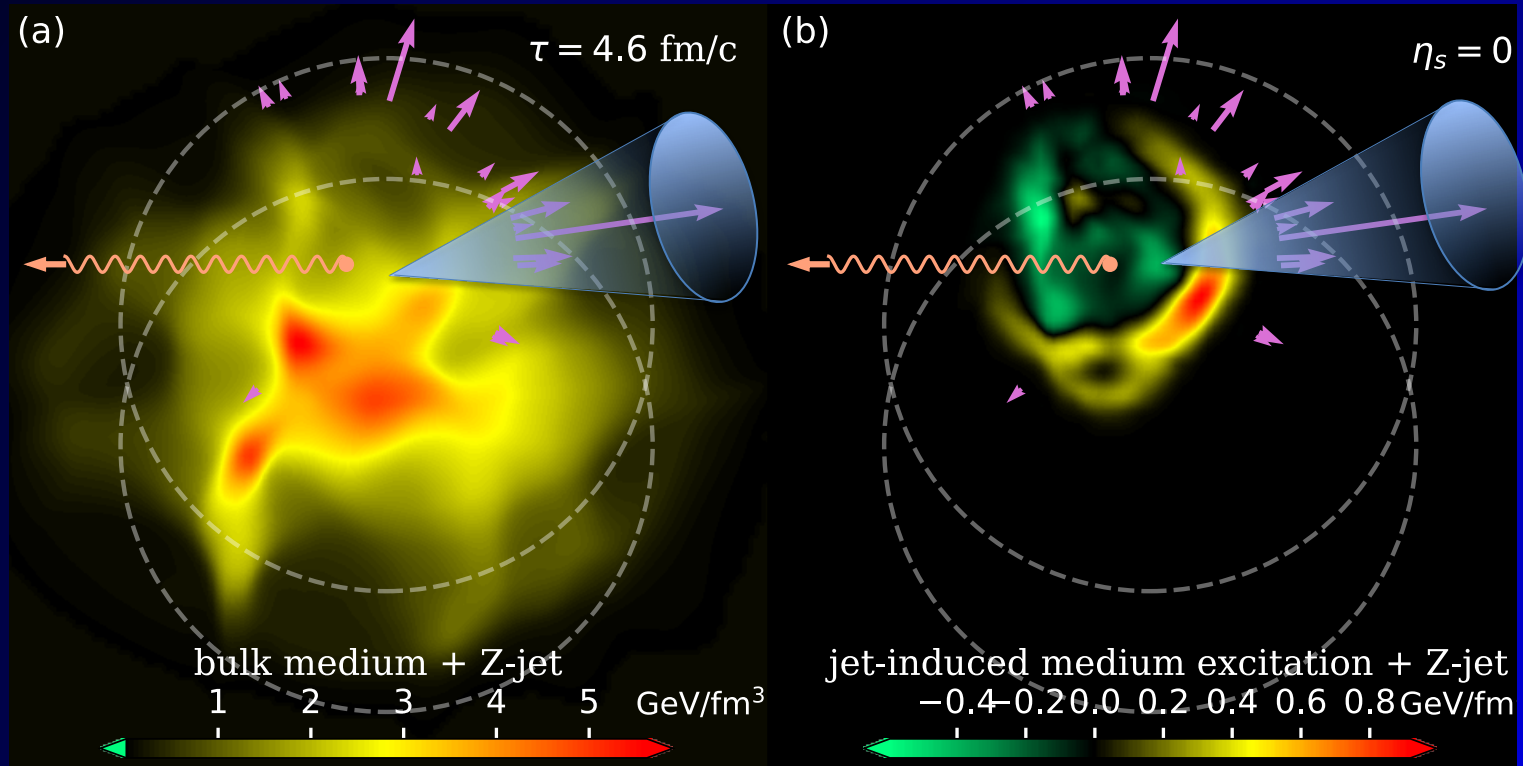
Single jet anisotropy

$$v_n^{\text{jet}} = \frac{\langle \langle v_n \cos[n(\phi^{\text{jet}} - \Psi_n)] \rangle \rangle}{\sqrt{\langle v_n^2 \rangle}}$$



He, Cao, Luo, Pang & XNW, in preparation

Z/ γ -jet: a better probe

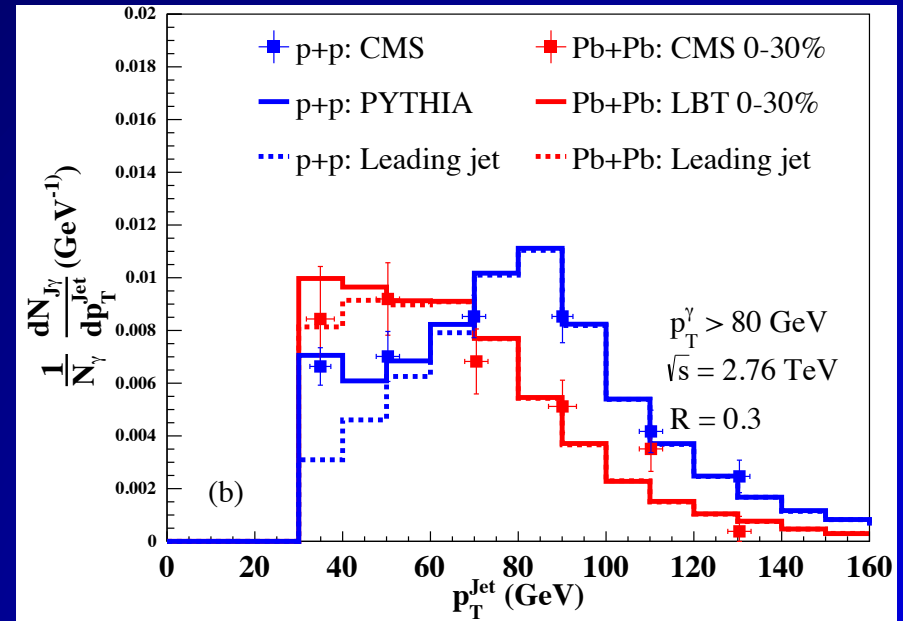
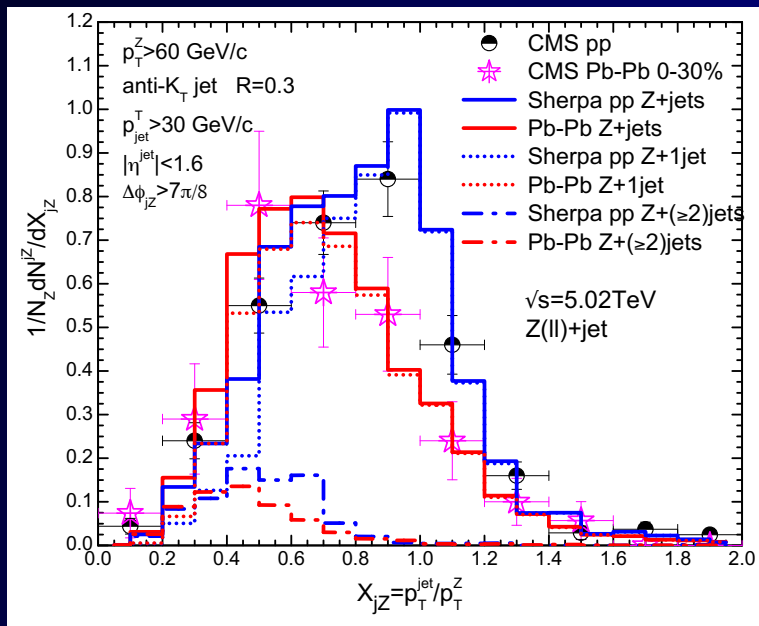


Chen, Cao, Luo, Pang & XNW, PLB777(2018)86

Chen, Yang, He, Ke, Pang and XNW, 2101.05422

Energy loss in γ/Z -jet at LHC

Suppression of leading and multiple jets

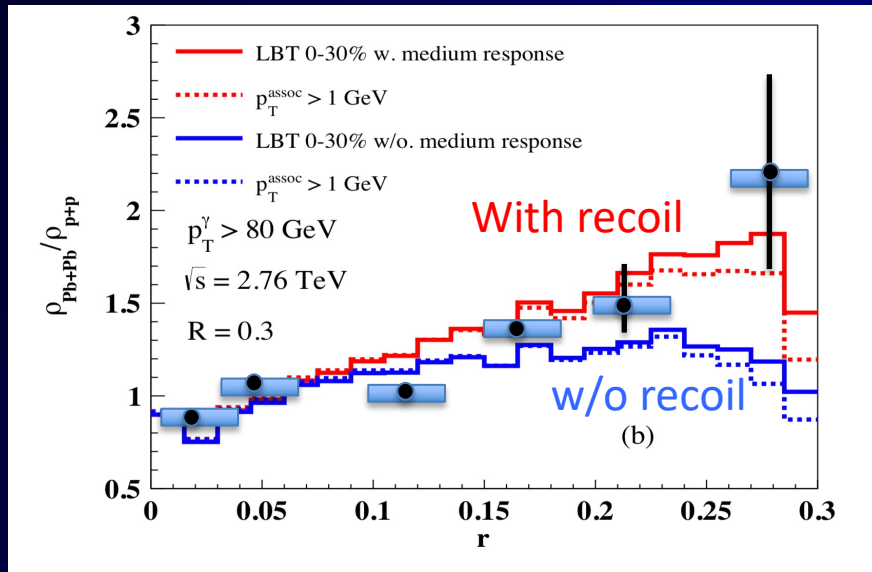


Zhang, Luo, XNW, Zhang, arXiv:1804.11041

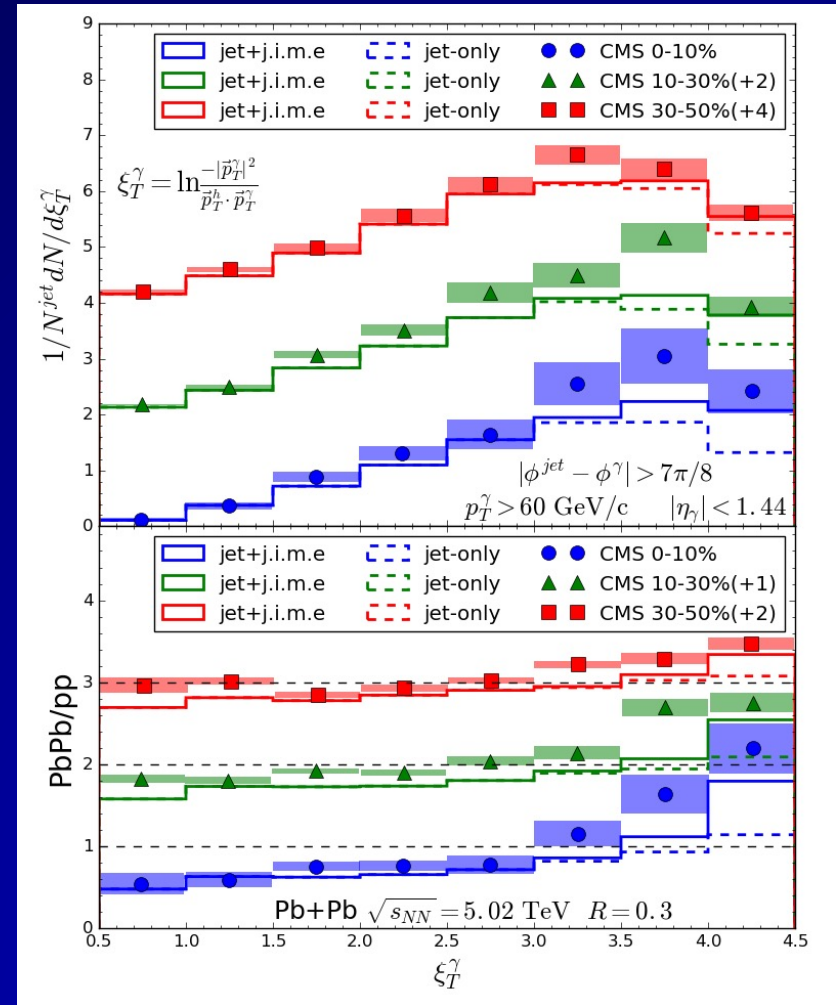
Luo, Cao, He & XNW, arXiv:1803.06785

Medium modification of γ -jets

Enhancement of soft hadrons
in large angles

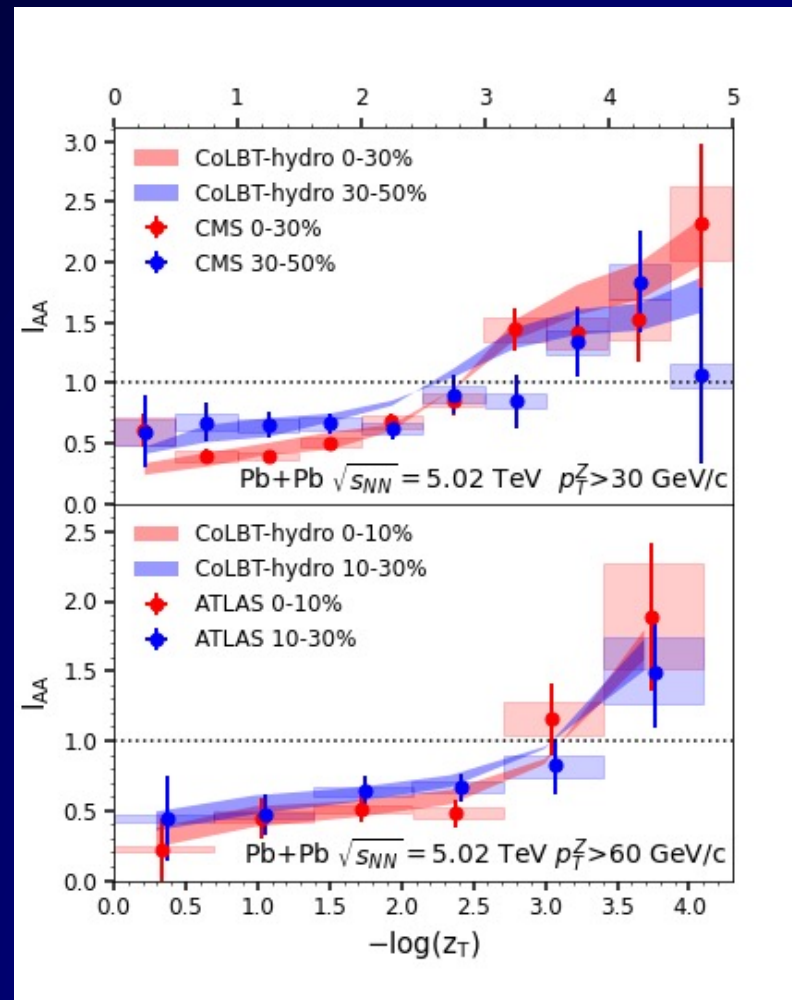


Luo, Cao, He & XNW, arXiv:1803.06785



Chen, Cao, Luo, Pang & XNW, 2005.09678

Z-hadron correlation at LHC



Chen, Yang, He, Ke, Pang and XNW, 2101.05422

Medium response & soft gluon radiation

Medium response: $\delta f(p) \sim e^{-p \cdot u/T}$

Medium-induced gluon radiation:

Formation time: $\tau_f = \frac{2\omega}{k_T^2} \quad k_T^2 \approx \tau_f \hat{q}$

$\hookrightarrow \tau_f \approx \sqrt{2\omega/\hat{q}}$

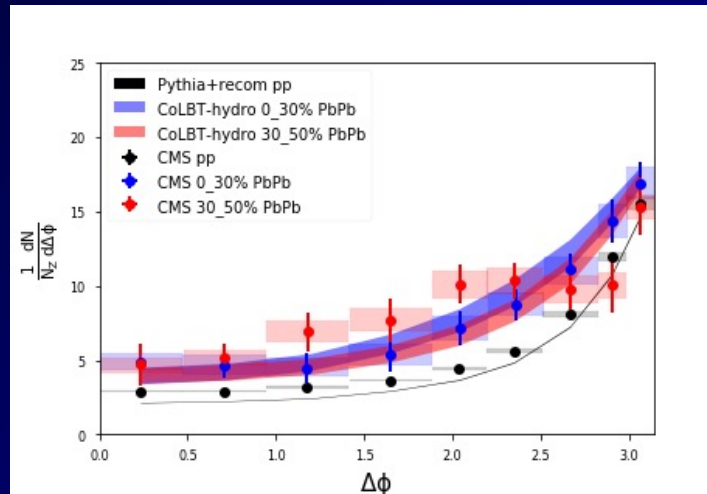
Mean-free-path
limits the formation time

$$\tau_f \leq \lambda \sim 1/T \quad \hat{q} \sim T^3$$

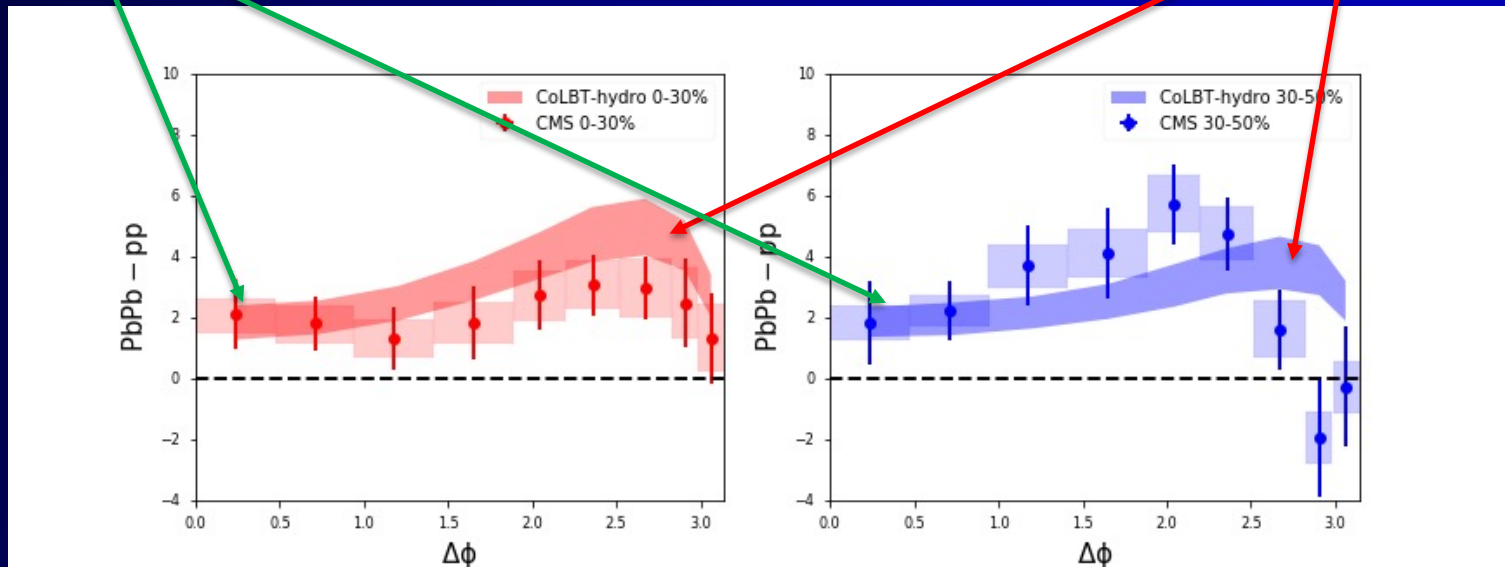
$$\omega \approx \lambda^2 \hat{q}/2 \sim T$$

Z-hadron correlation

enhancement
of the away-side
background

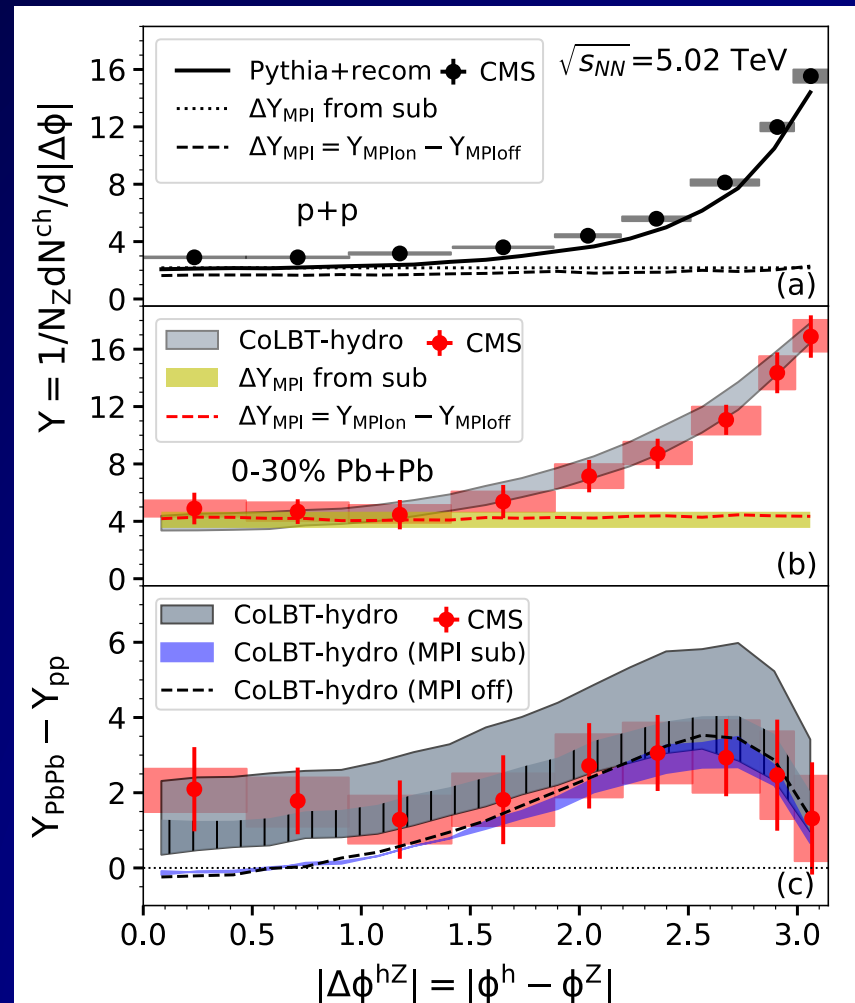


enhancement
and broadening
of the jet peak



Chen, Yang, He, Ke, Pang and XNW, 2101.05422

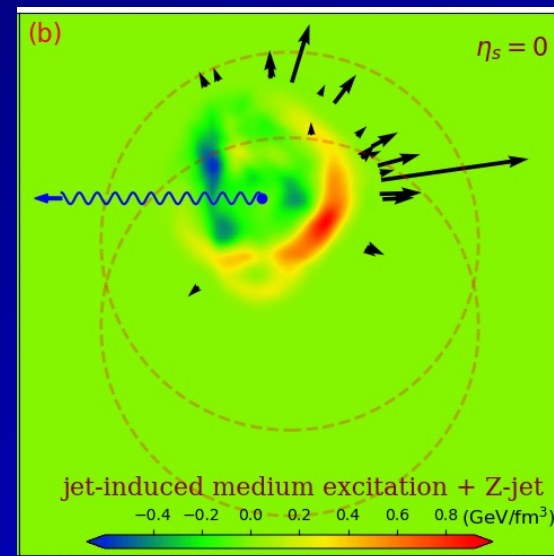
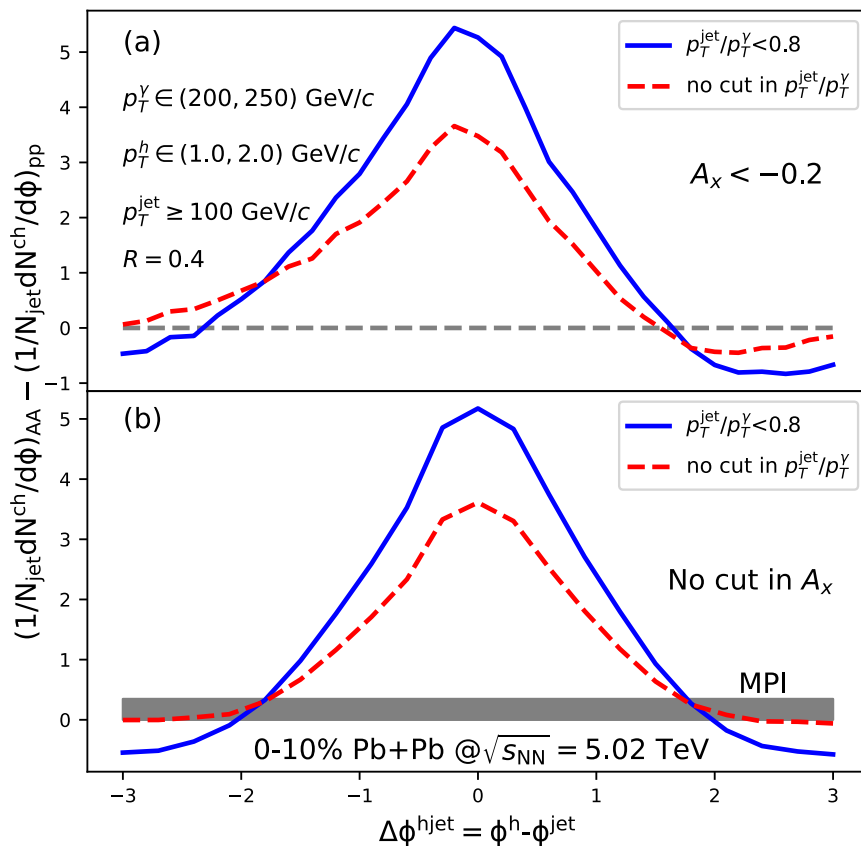
MPI subtraction in Z-hadron correlation



Mixed event subtraction

$$\frac{dN_{MPI}^{hZ}}{d\phi} = \frac{dN_{mix}^{hZ}}{d\phi} - \int_1^\pi \frac{d\phi}{\pi} \left(\frac{dN^{hZ}}{d\phi} - \frac{dN^{hZ}}{d\phi} \Big|_{\phi=1} \right)$$

Enhancing the diffusion wake

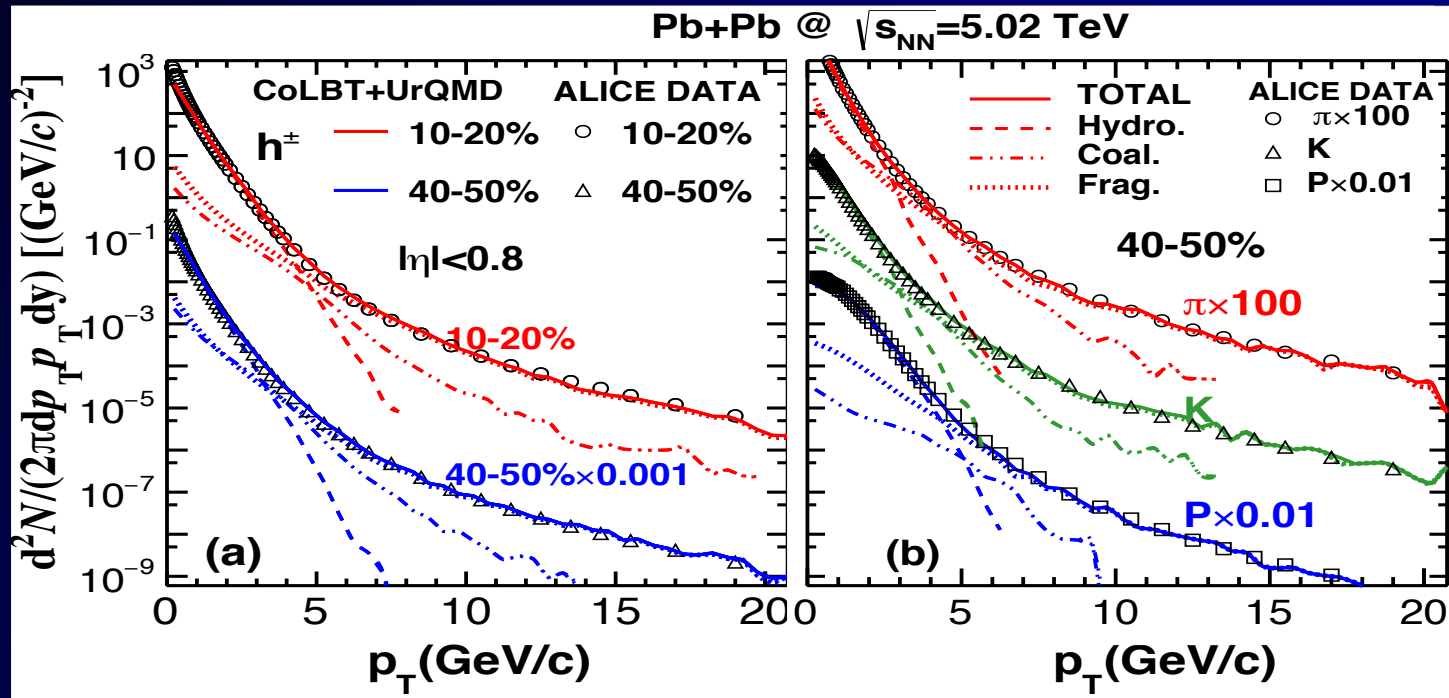


Chen, Yang, He, Ke, Pang and XNW, 2101.05422

Hydro, coalescence, fragmentation and hadron cascade

- Thermal hadrons (CLVisc)
 - Hydro with Cooper-Frye ($p_T < p_{T1}$)
- Coalescence hadrons
 - Include thermal-thermal, thermal-hard and hard-hard coalescence
- Fragmentation hadrons
 - Lund fragmentation
- UrQMD hadronic afterburner
 - All hadrons are fed into UrQMD for hadronic evolution and decay

Hadron spectra from low to high p_T



Hydro : $p_T < 2$ GeV/c

radial flow

Coal.: $2 < p_T < 6$ GeV/c

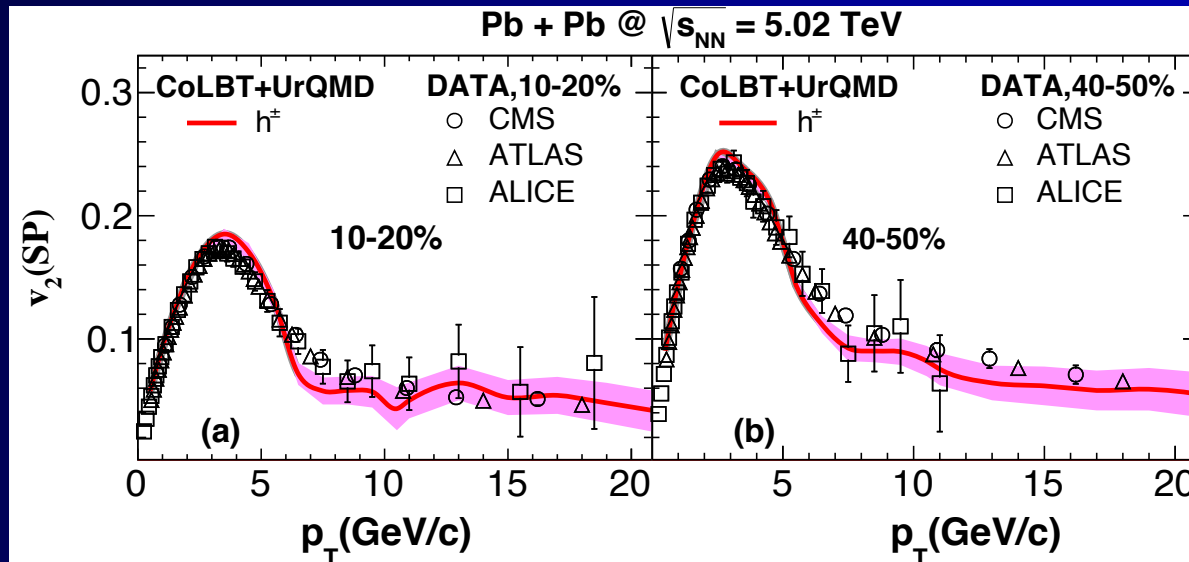
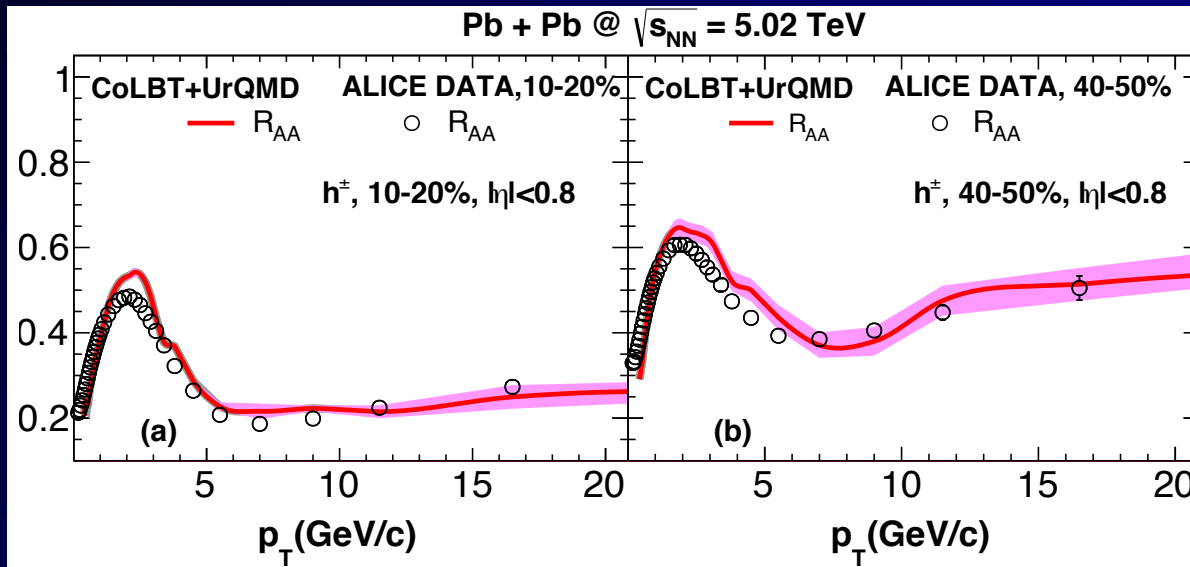
Coalescence

Frag.: $p_T > 5$ GeV

energy loss

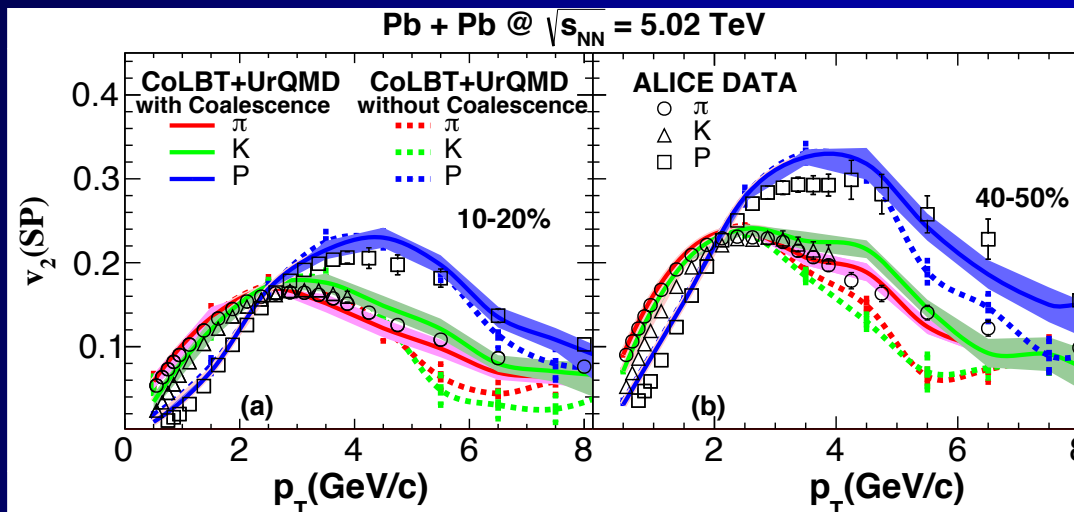
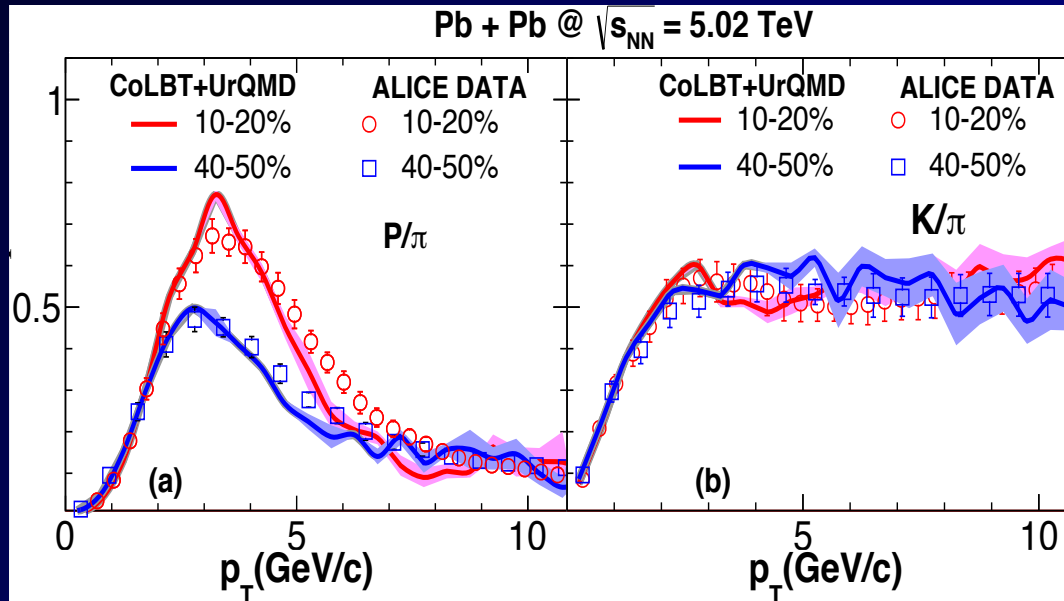
Zhao, Ke, Chen, Luo and XNW,
 [arXiv:2103.14657 [hep-ph]]

Solving $R_{AA}-v_2$ puzzle



Zhao, Ke, Chen, Luo and XNW, arXiv:2103.14657

Flavor dependence

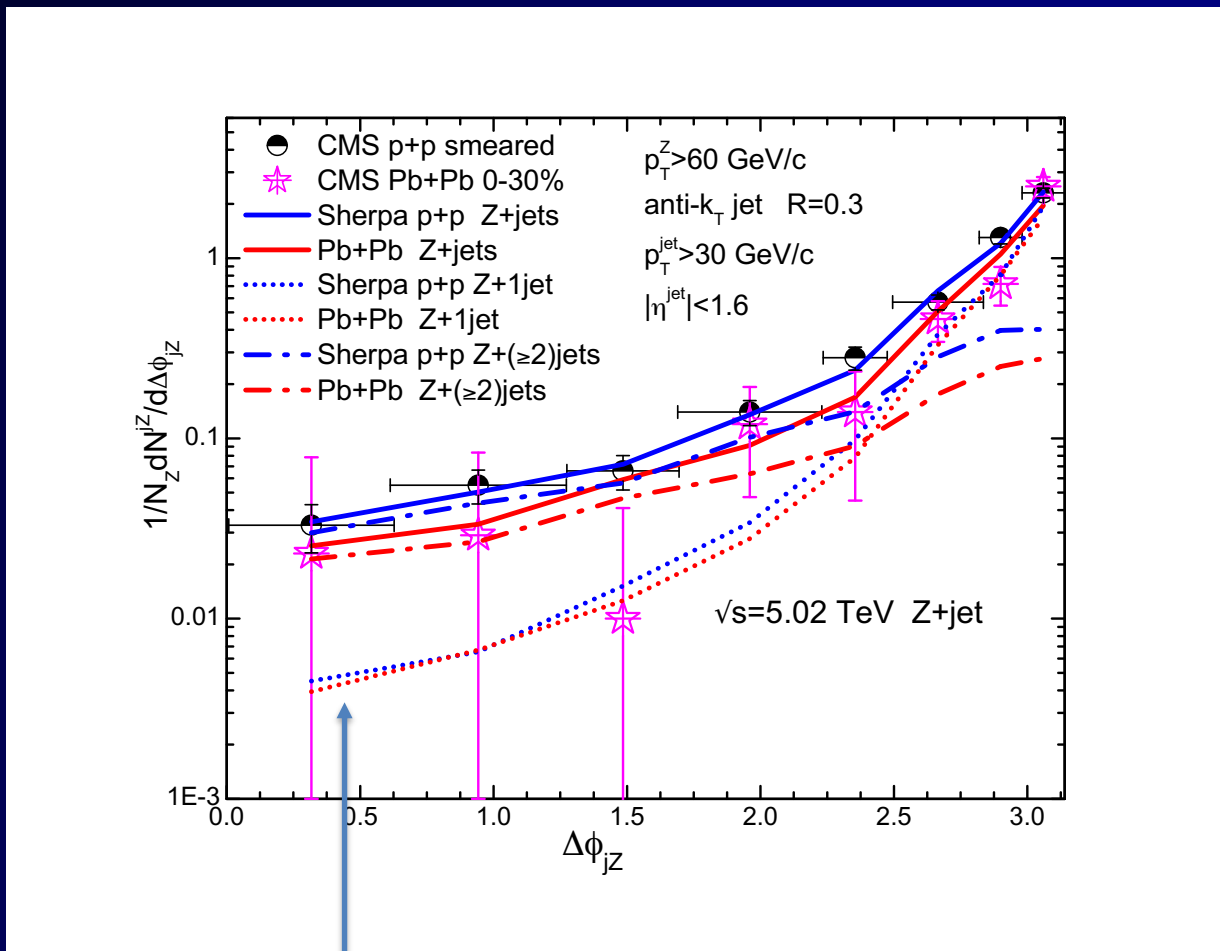


Zhao, Ke, Chen, Luo and XNW, arXiv:2103.14657

Conclusion & Discussion

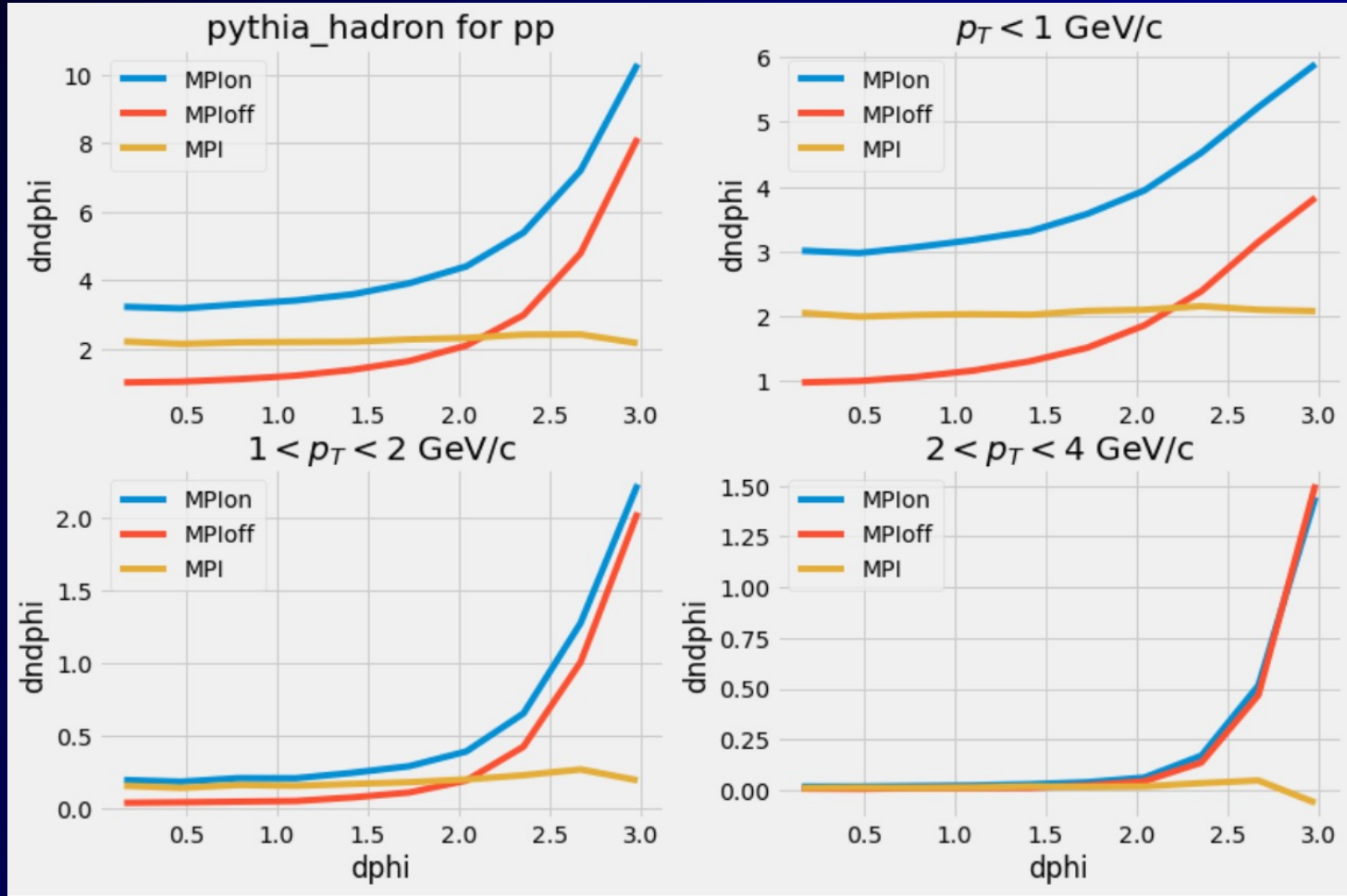
- CoLBT-hydro: describe h spectra from low to high p_T
- Coalescence at intermediate p_T solves R_{AA} v_2 puzzle
- Flavor dependence of spectra and v_2
- Medium response leads to
 - enhancement of soft hadrons in jet direction
 - depletion of soft hadron on the away side
- Use 2D jet tomography to reveal the angular structure of Mach-cone excitation

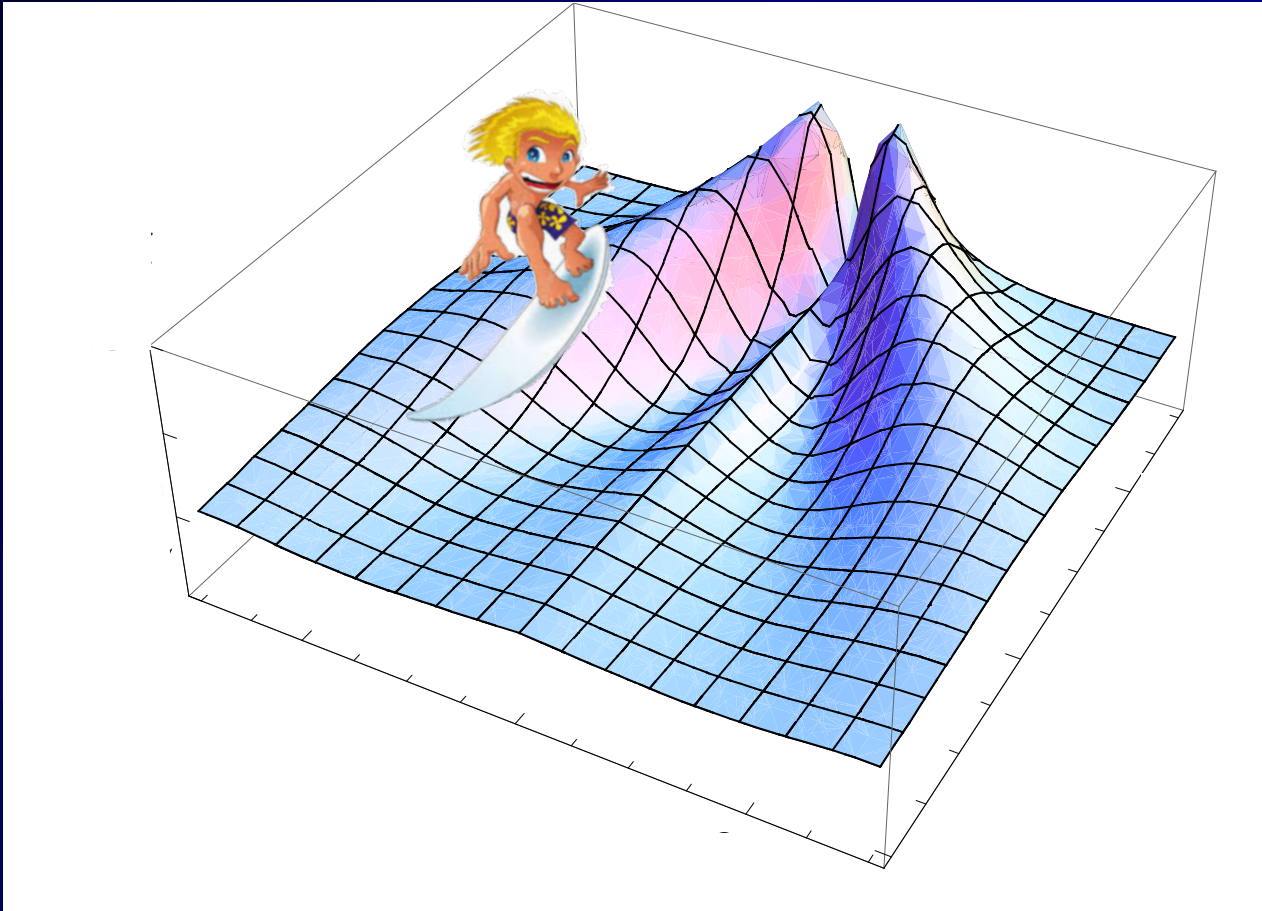
Multiple jets in Z-jet events



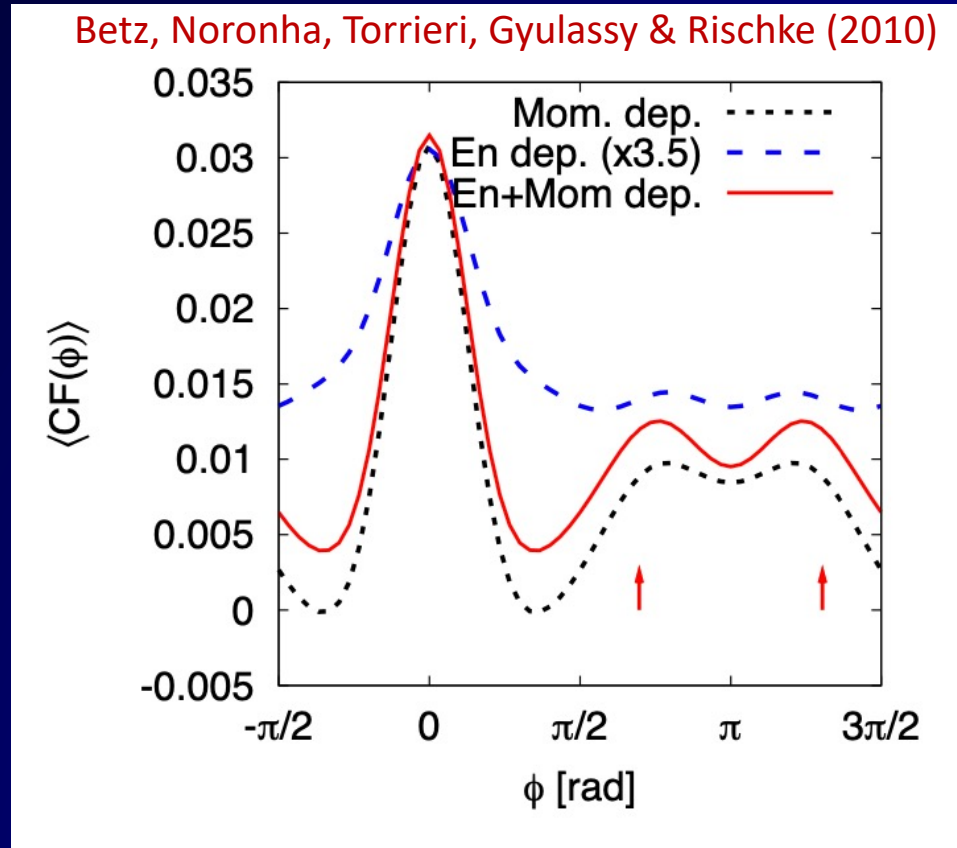
Multiple jets contribution is negligible in Z direction

MPI at RHIC



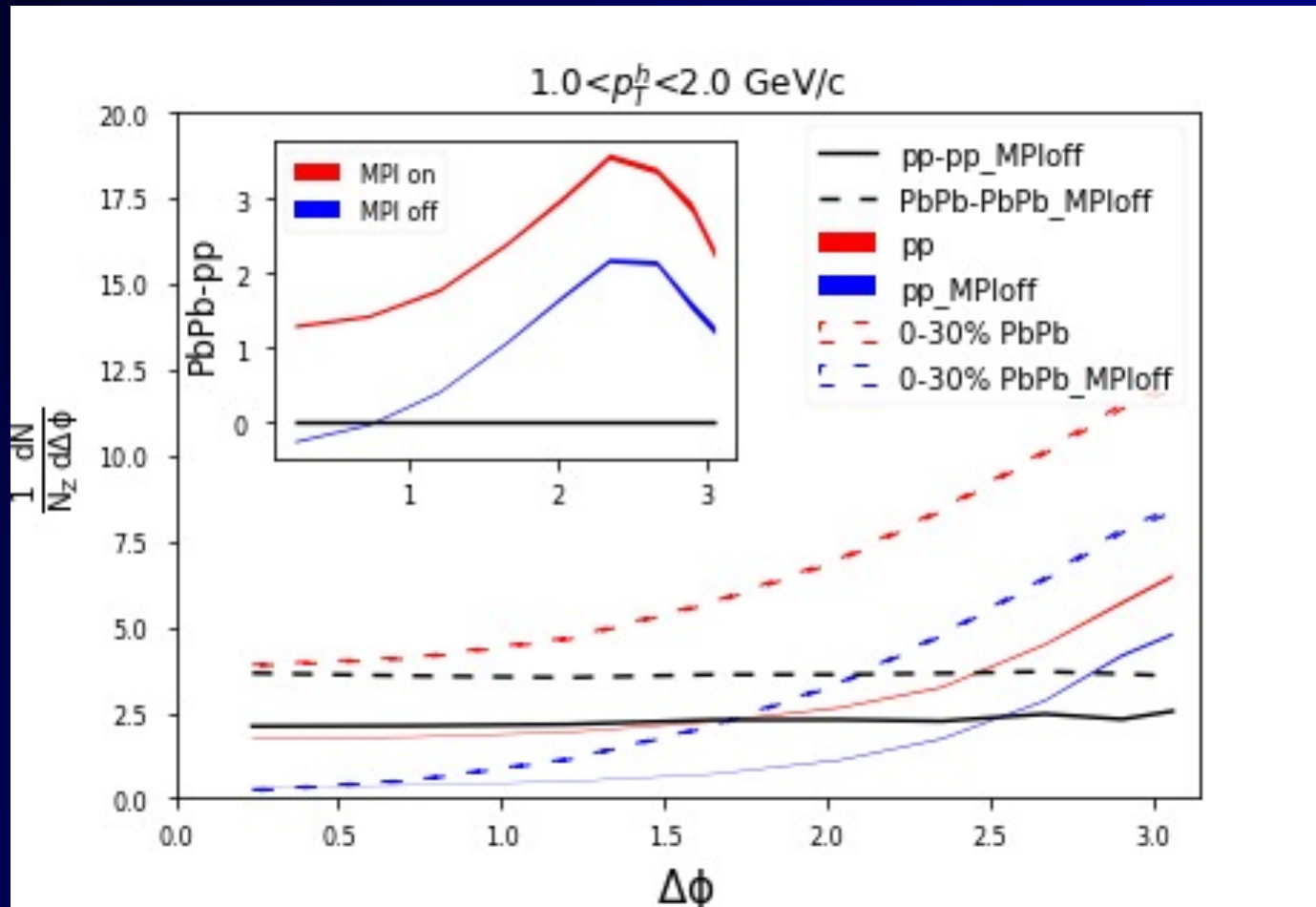


Signals of Mach-cone?



- Angular structure disappears after average over initial production points
- Medium-induced gluon radiation can overwhelm in some phase-space
- Complication by anisotropic flow v_n

MPI contribution to Z-hadron correlation



MPI negligible at RHIC

MPI: Multiple parton interaction

XNW & Gyulassy (1991)

$$g_j(b, p_T) = \frac{[\Delta\sigma(p_T)T(b)]^j}{j!} e^{-\Delta\sigma(p_T)T(b)}$$

Multiple jet production in pp:

$$g_j(b) = \frac{[\sigma(p_0)T(b)]^j}{j!} e^{-\sigma(p_0)T(b)}$$

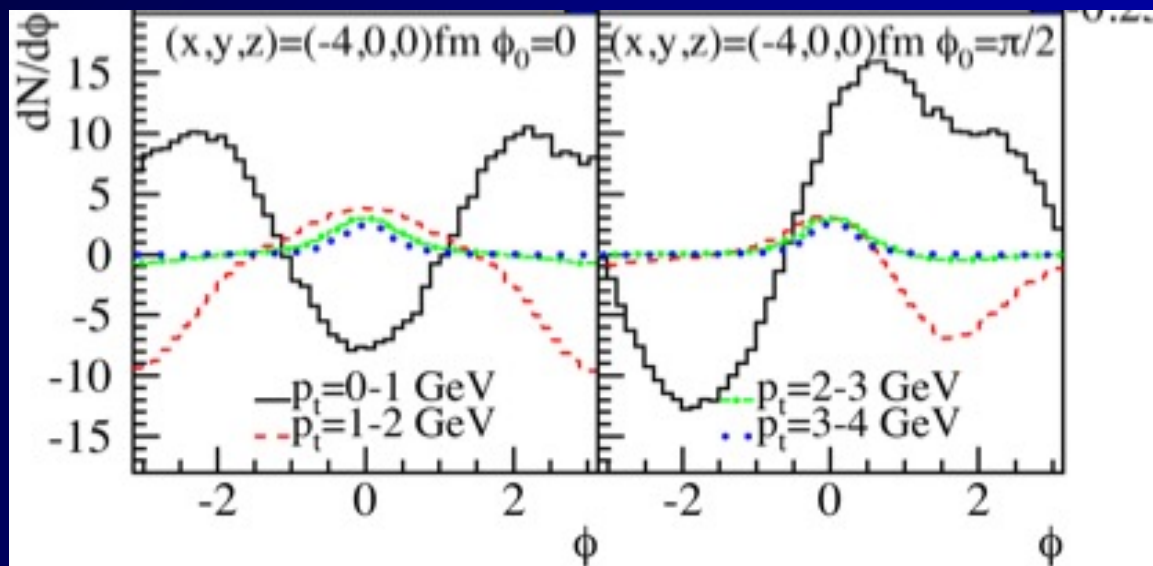
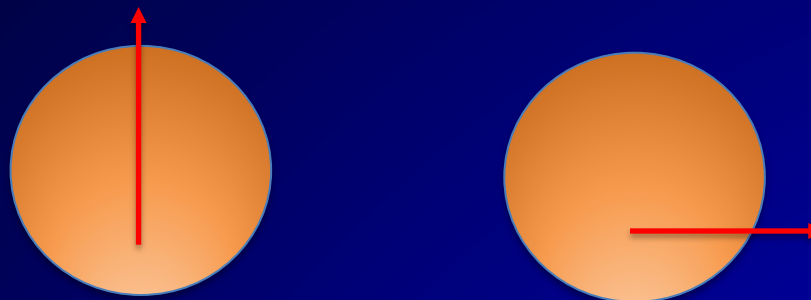
Probability of multiple jets ($p_T > p_0$) with at least one jet with $p_T > p_T^{\text{trig}}$

$$g_j^{\text{trig}}(b) = \frac{[\sigma(p_0)T(b)]^j}{j!} \left\{ 1 - \frac{[(\sigma(p_0) - \sigma(p_T^{\text{trig}}))]^j}{\sigma(p_0)^j} \right\} e^{-\sigma(p_0)T(b)}$$

$$\approx j \frac{\sigma(p_T^{\text{trig}})}{\sigma(p_0)} g_j(b)$$

**Enhanced multiple minijet
Production in triggered jet events**

Distorted Mach-cone-like excitation



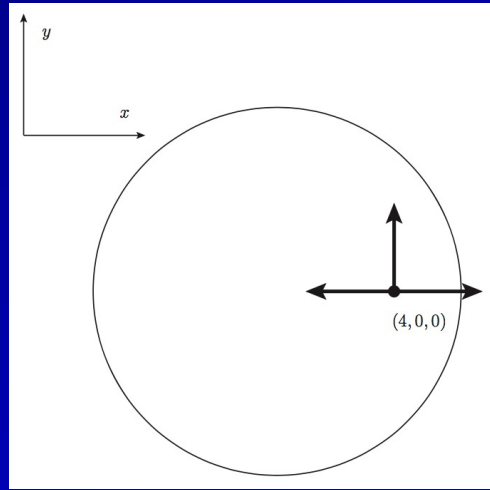
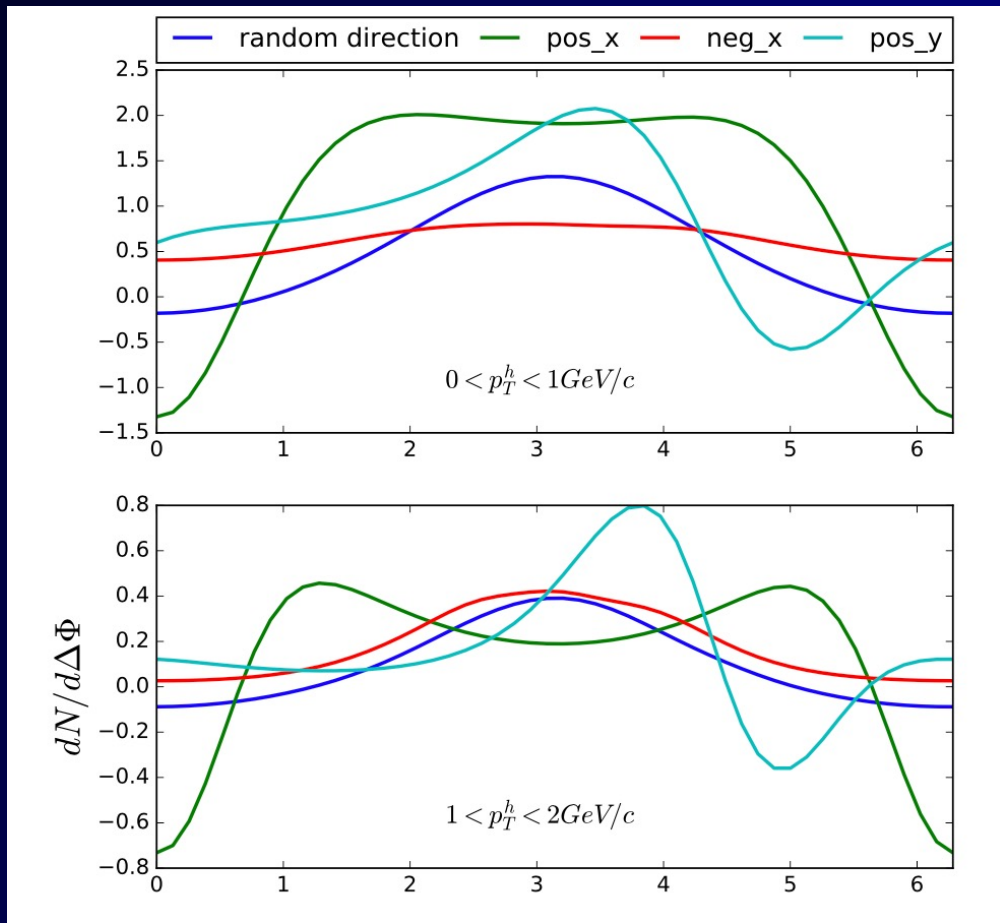
Li, Liu, Ma, XNW and Zhu, Phys. Rev. Lett. 106, 012301 (2011)

Tachibana, Shen & Majumder [2001.08321](#) (2020)

Initial position & azimuthal correlation

W Chen & XNW (2018)

γ -hadron correlation

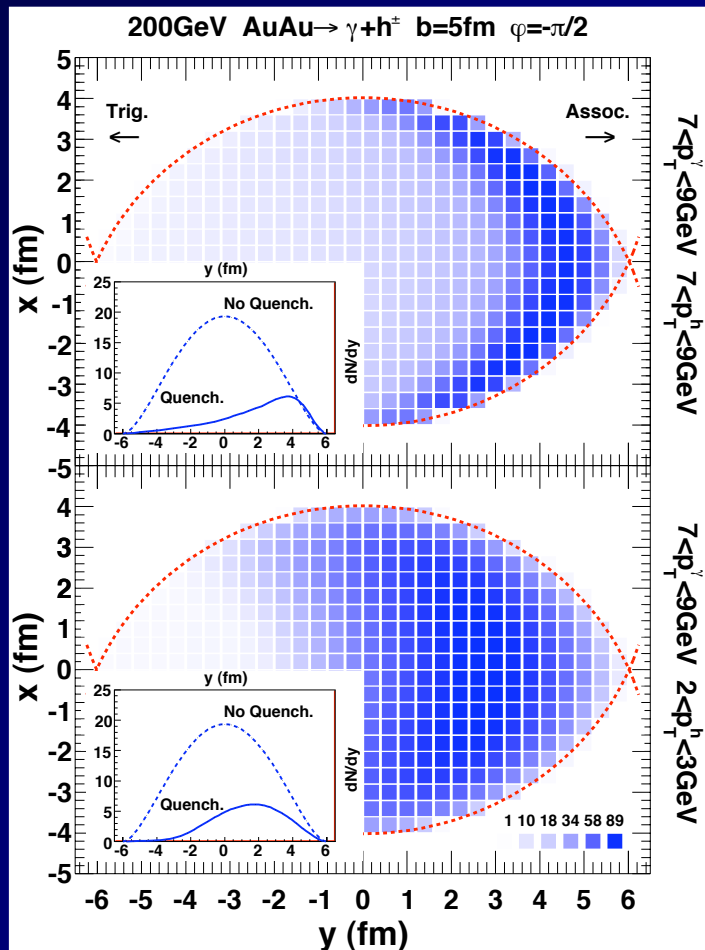


$\Delta\Phi$

Longitudinal jet tomography

Zhang, Owens, Wang and XNW, Phys. Rev. Lett. 103, 032302 (2009)

length
dependence
of parton
Energy loss

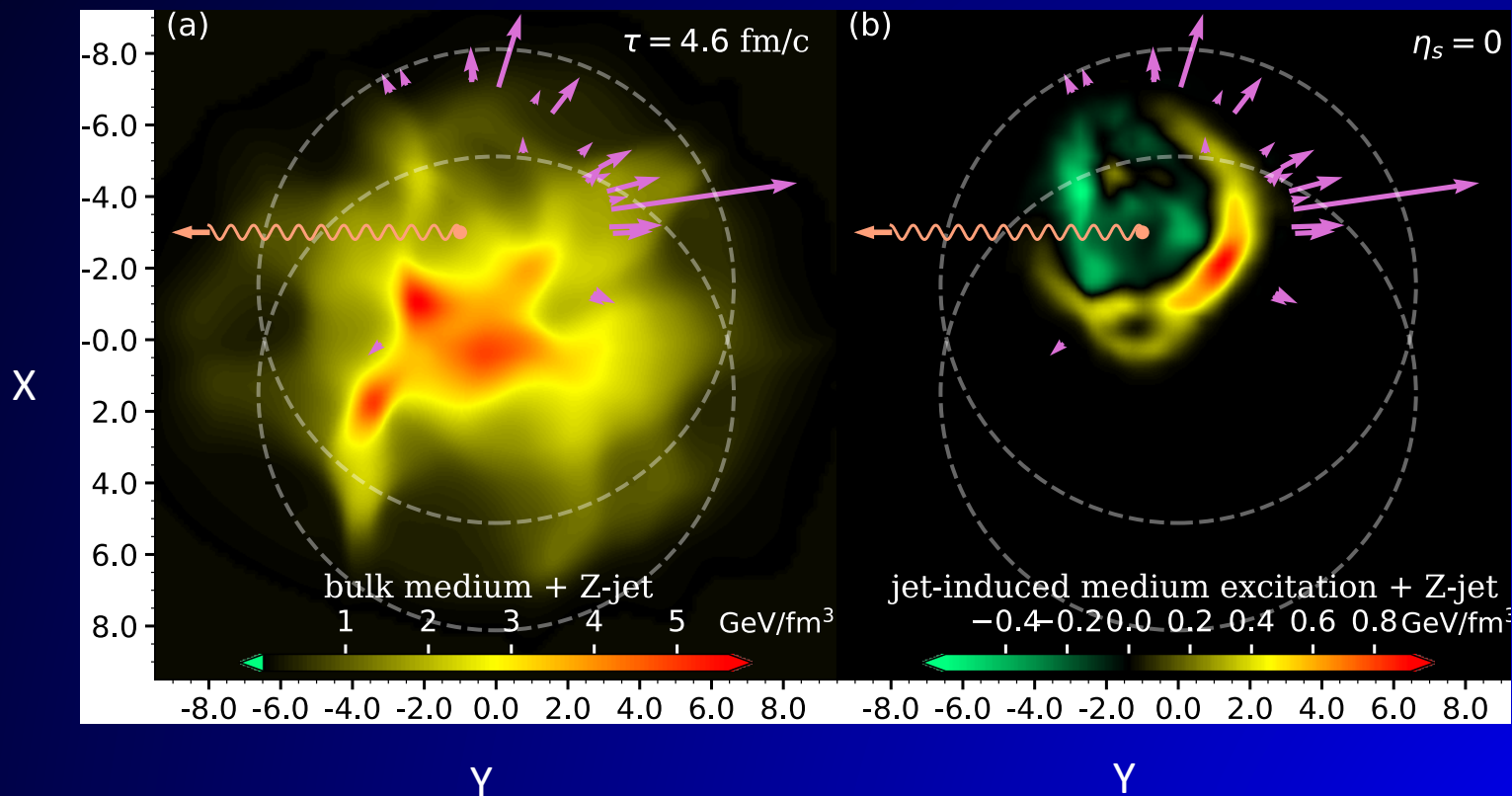


$$p_T^h / p_T^\gamma \sim 1$$

$$p_T^h / p_T^\gamma \sim 0.3$$

Transverse gradient tomography

gradient dependence of p_T broadening

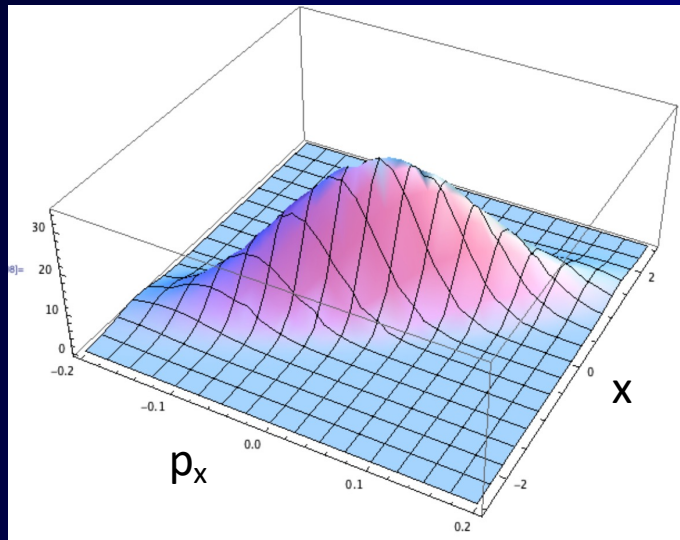


Drift-diffusion equation: uniform medium

Boltzmann equation under approximation of small angle elastic scattering, no drag:

$$\frac{\partial f}{\partial t} + \frac{\vec{p}_\perp}{E} \cdot \frac{\partial f}{\partial \vec{r}_\perp} = \frac{\hat{q}}{4} \nabla_{p_\perp}^2 f(\vec{p}, \vec{r})$$

Initial distr. $f(\vec{p}, \vec{r})_{t=0} = (2\pi)^2 \delta^2(\vec{r}_\perp) \delta^2(\vec{p}_\perp)$



$$f = 3 \left(\frac{4E}{\hat{q}t^2} \right)^2 \exp \left[-(\vec{r}_\perp - \frac{\vec{p}_\perp}{2E}t)^2 \frac{12E^2}{\hat{q}t^3} - \frac{p_\perp^2}{\hat{q}t} \right]$$

$$\begin{array}{ccc} \int d^2 p_\perp & & \int \frac{d^2 r}{(2\pi)^2} \\ \downarrow & & \downarrow \\ 4\pi \frac{3E^2}{\hat{q}t^3} \exp \left(-r_\perp^2 \frac{3E^2}{\hat{q}t^3} \right) & & \frac{1}{\pi \hat{q}t} \exp \left(-\frac{p_\perp^2}{\hat{q}t} \right) \end{array}$$

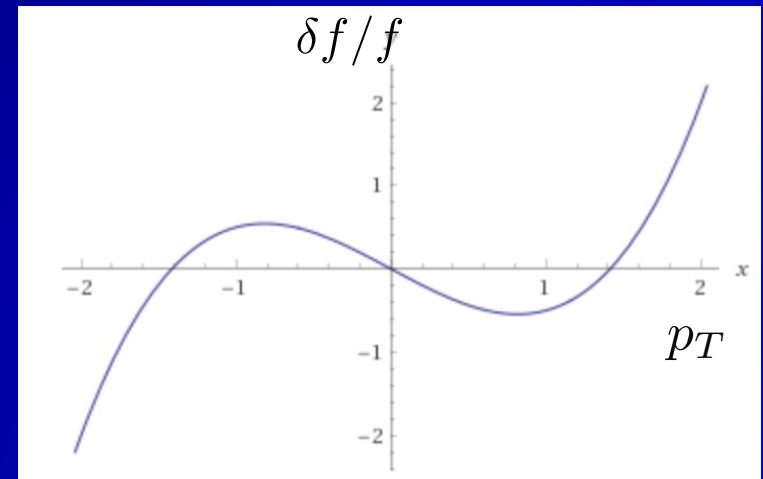
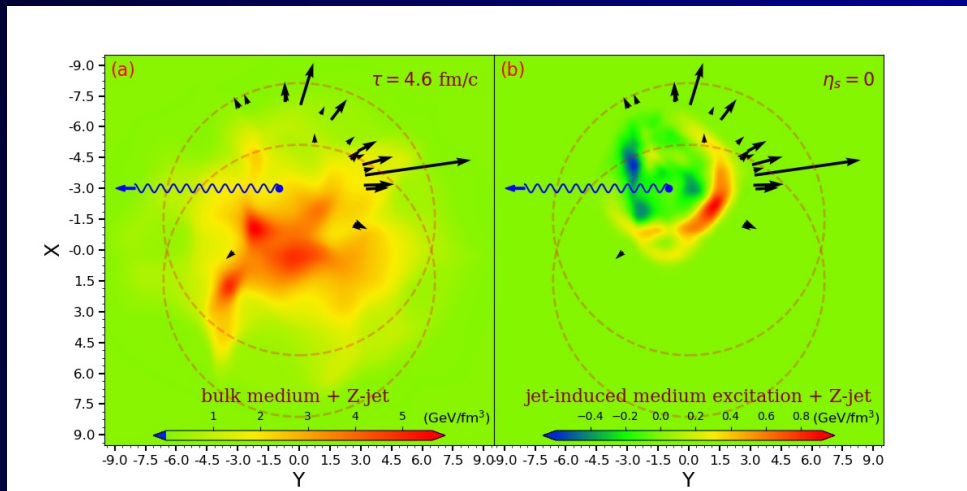
He, Pang & XNW, *PRL* 125 (2020) 12, 122301

Drift-diffusion equation: non-uniform medium

Linear spatial dependence $\hat{q} = \hat{q}_0 + \vec{x}_\perp \cdot \vec{a}$

Momentum asymmetry:

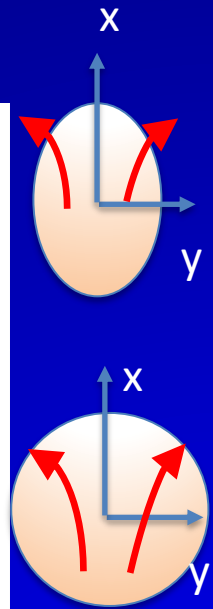
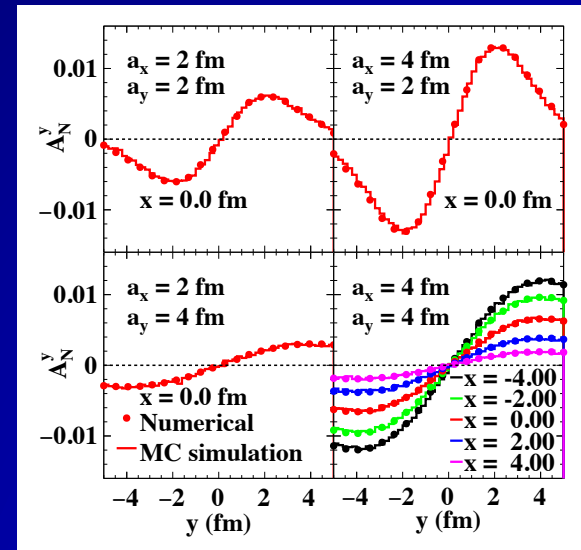
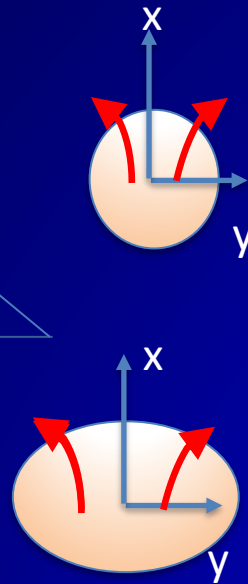
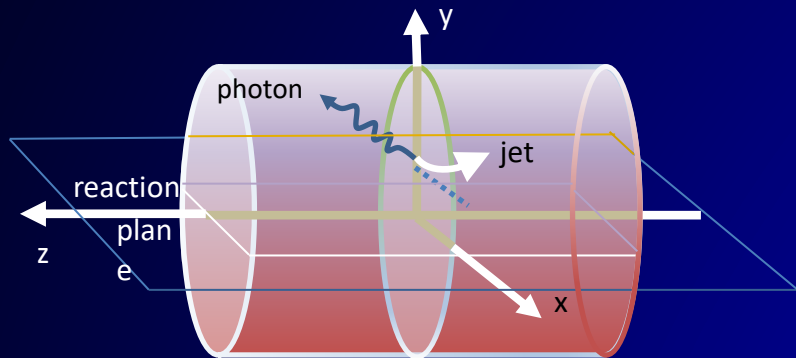
$$\delta f(\vec{p}_\perp) = -\frac{t}{3\omega\hat{q}_0} \vec{a} \cdot \vec{p}_\perp \left(1 - \frac{p_\perp^2}{2\hat{q}_0 t} \right) f_s(\vec{p}_\perp, t) + \mathcal{O}(a^2)$$



Diffusion in a non-uniform medium

Momentum asymmetry

$$A_N^{\vec{n}} = \frac{\int d^3r d^3k f_a(\vec{k}, \vec{r}) \text{Sign}(\vec{k} \cdot \vec{n})}{\int d^3r d^3k f_a(\vec{k}, \vec{r})}$$



$$\hat{q}(\vec{r}_\perp, t) = \frac{\hat{q}_0 t_0}{t_0 + t} e^{-x^2/a_x^2 - y^2/a_y^2}$$

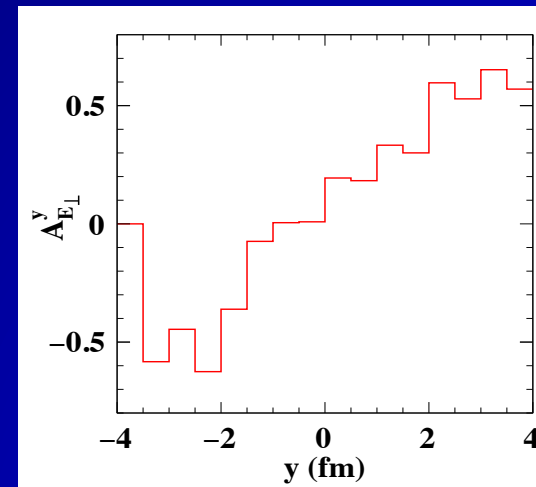
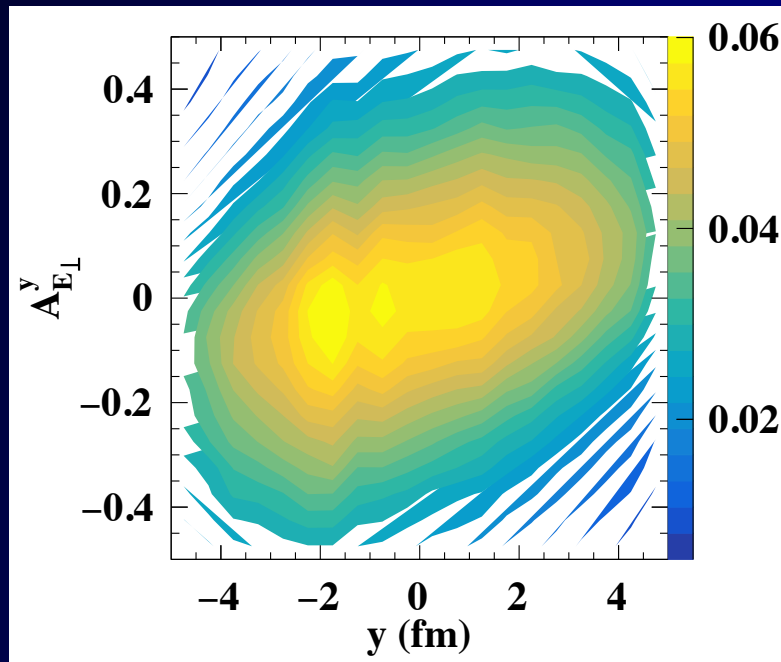
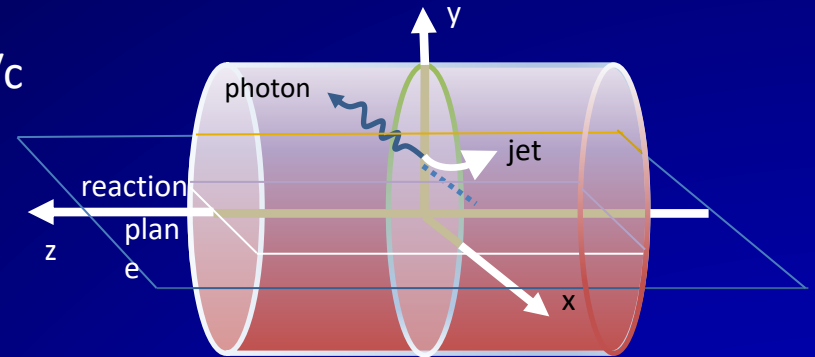
$$\hat{q}_0 = 5 \text{ GeV}^2/\text{fm}, t_0 = 0.5 \text{ fm}/c, E = 20 \text{ GeV}$$

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Momentum asymmetry wrt event plane

$$A_{E_{\perp}}^{\vec{n}} = \frac{\int d^3r d^3p f_a(\vec{p}, \vec{r}) \vec{p}_T \cdot \vec{n}}{\int d^3r d^3p f_a(\vec{p}, \vec{r})}$$

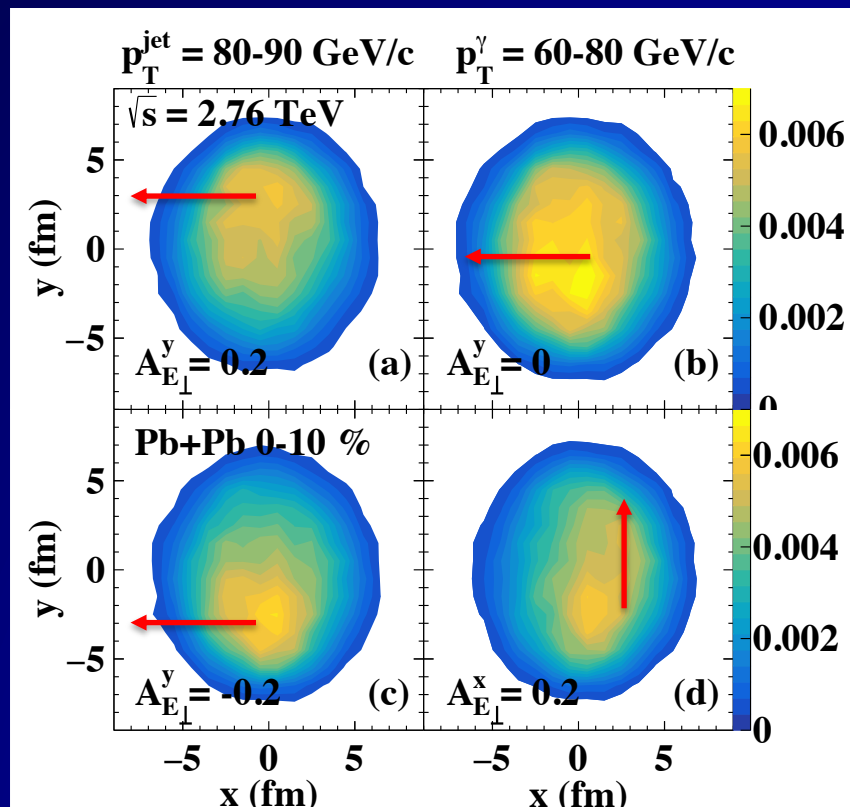
$p_T > 3 \text{ GeV}/c$



He, Pang & XNW, *PRL* 125 (2020) 12, 122301

Gradient tomography

With trigger on the transverse asymmetry energetic hadrons one can localize the initial production point of gamma-jet in the transverse plane



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