

A spectral reconstruction test using the two-dimensional $O(3)$ -model

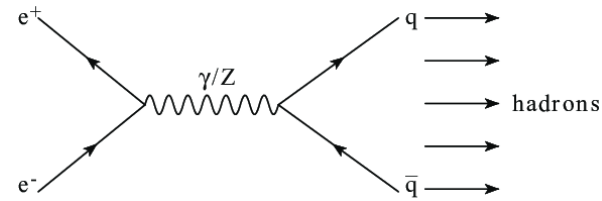
John Bulava

DESY-Zeuthen

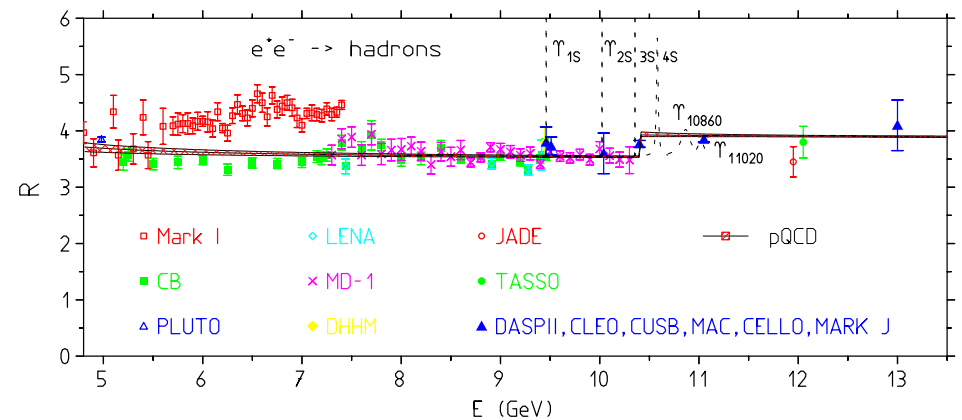
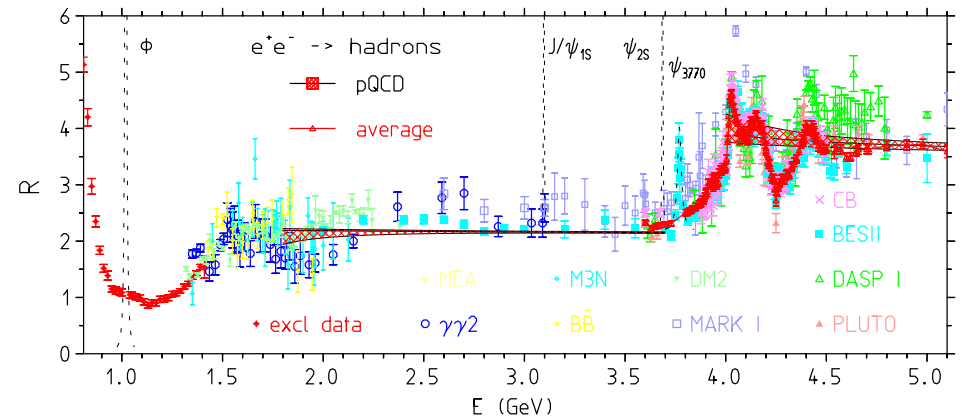
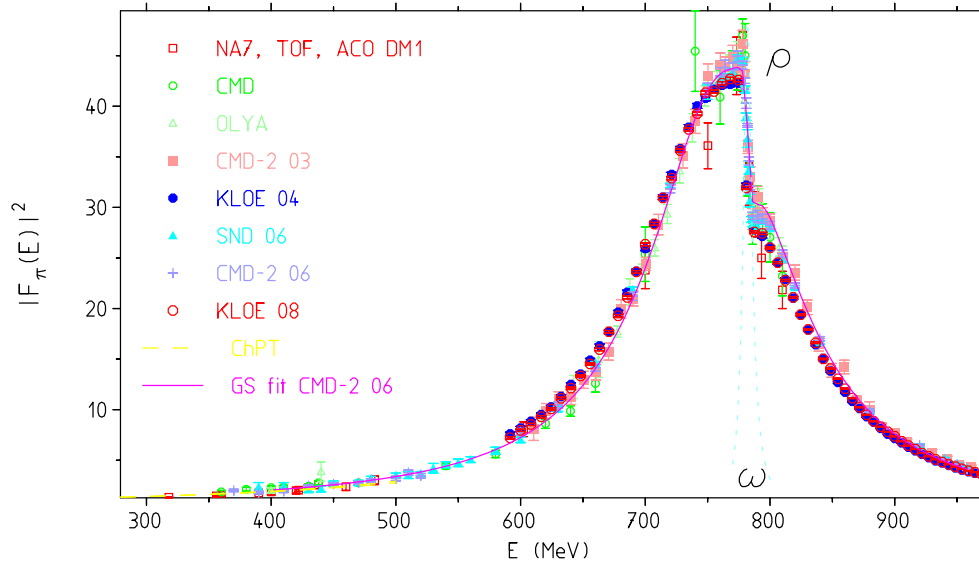


ECT* Workshop
Spectral Properties From Euclidean Path Integrals
Sep. 15th, 2021

A “simple” lattice QCD spectral reconstruction



$$R(s) = \frac{R(e^+e^- \rightarrow \text{hadrons})}{4\pi\alpha_{\text{em}}(s)/(3s)}, \quad \rho(s) = \frac{R(s)}{12\pi^2}$$

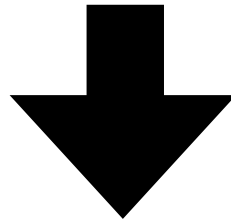


Higher energies => lower resolution.

Jegerlehner, Nyffeler '09

Useful ‘smeared’ spectral functions

$$G(\tau) = \int d^3\mathbf{x} \langle \Omega | \hat{j}_i^{\text{em}}(\mathbf{x}, \tau) \hat{j}_i^{\text{em}\dagger}(0) | \Omega \rangle = \int_0^\infty d\omega \omega^2 \rho(\omega^2) e^{-\omega\tau}$$



Adler function:

$$D(Q^2) = Q^2 \int_0^\infty ds \frac{\rho(s)}{(s + Q^2)^2}$$

Hadronic contribution to $(g - 2)_\mu$:

$$a_\mu^{\text{HVP}} = \int_0^\infty dQ^2 D(Q^2) g(Q^2, m_\mu)$$

Spectral Reconstruction

Backus, Gilbert '68, '70
F. Pijpers, M. Thompson '92
M. R. Hansen, A. Lupo, N. Tantalo, PRD99 (2019)

Linear ansatz:

$$\hat{\rho}_{\epsilon}(E) = \sum_{t=a}^{t_{\max}} q_t(\epsilon, E) C(t), \quad \hat{\delta}_{\epsilon}(E, \omega) = \sum_{t=a}^{t_{\max}} q_t(\epsilon, E) e^{-\omega t}$$

Two criteria when choosing $\{q_t(\epsilon, E)\}$

- Accuracy: $A[q] = \int_{E_0}^{\infty} d\omega \left\{ \delta_{\epsilon}(E - \omega) - \hat{\delta}_{\epsilon}(E, \omega) \right\}^2$
- Precision: $B[q] = \text{Var} \{ \rho_{\epsilon}(E) \}$

Optimal coeffs minimize:

$$W_{\lambda}[q] = (1 - \lambda)A[q]/A[0] + \lambda B[q]$$

Controlled Test

JB, M. W. Hansen, M. T. Hansen, (M. R. Hansen), A. Patella, N. Tantalo, *in prep.*

2d O(3)-model: ..., M. Lüscher, U. Wolff, Nucl. Phys. B339 (1990),...

$$S[\sigma] = -\beta \sum_{x, \mu} \sigma(x) \cdot \sigma(x + \hat{\mu}), \quad \sigma(x) \in \mathbb{R}^3, |\sigma(x)| = 1$$

Conserved (global) current:

$$j_{\mu}^a = \beta \epsilon^{abc} \sigma^b(x) \hat{\partial}_{\mu} \sigma^c(x)$$

Massive single-particle states. Target process: inclusive rate for $j \rightarrow X$

$$\rho(E) = \sum_{\alpha} \delta(\mathbf{P}_{\alpha}) \delta(E - E_{\alpha}) |_{\text{out}} \langle \alpha | \hat{j}(0) | 0 \rangle|^2$$

Controlled Test

Spectral function from current-current correlator:

$$C(\tau) = \int d\mathbf{x} \langle j_1^a(\tau, \mathbf{x}) j_1^a(0) \rangle = \int_0^\infty d\omega e^{-\omega\tau} \rho(\omega)$$

All (even) particle-number sectors contribute

$$\rho(E) = \sum_{n=1}^{\infty} \rho^{(2n)}(E), \quad \rho^{(2n)}(E) = 0 \text{ for } E < 2nm$$

Elastic region $E < 4m$ can be computed using finite-volume formalism

(e.g. timelike pion form factor in QCD)

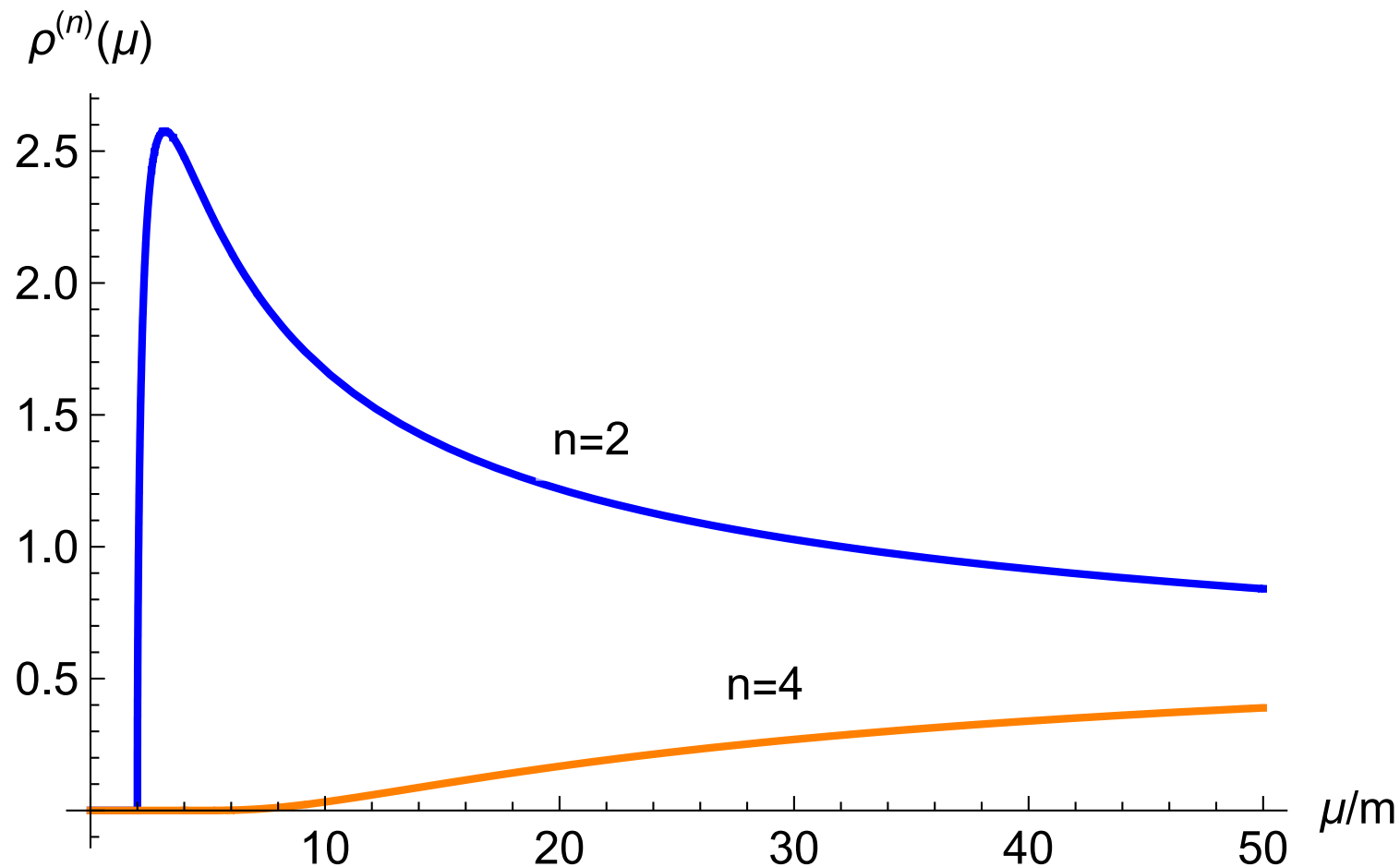
Controlled Test

Integrable model => spectral function known exactly:

M. Karowski, P. Weisz, Nucl. Phys. B139 (1978)

A. B. Zamolodchikov, A. B. Zamolodchikov, Nucl. Phys. B133 (1978)

J. Balog, M. Niedermaier, Nucl. Phys. B500 (1997)



Two-particle contribution dominant, four-particle ~2% near $E = 10m$

Controlled Test

ID	$(L/a) \times (T/a)$	am	mL	mT
A1	640×320	0.0453605(96)	29	15
A2	1280×640	0.0259551(49)	33	17
A3	1920×960	0.0176980(47)	34	17
A4	2880×1440	0.0112925(45)	33	16
B1	160×320	0.0623644(66)	10	20
B2	320×320		20	
B3	480×320		30	
B4	640×320		40	
B5	800×320		48	
B6	960×320		58	
C1	320×160	0.0623644(66)	20	10
C2	320×240			15
C3	320×320			20
C4	320×480			30

Simulations with Wolff two-cluster algorithm:

• A1-A4: continuum limit

• B1-B6: finite-L effects

• C1-C4: finite-T effects

Controlled Test

Five smearing kernels $\delta_\epsilon(E - \omega)$:

$$\delta_\epsilon^{\text{g}}(x) = \frac{1}{\sqrt{2\pi}\epsilon} \exp \left[-\frac{x^2}{2\epsilon^2} \right],$$

$$\delta_\epsilon^{\text{c0}}(x) = \frac{1}{\pi} \frac{\epsilon}{x^2 + \epsilon^2},$$

$$\delta_\epsilon^{\text{c1}}(x) = \frac{2}{\pi} \frac{\epsilon^3}{(x^2 + \epsilon^2)^2},$$

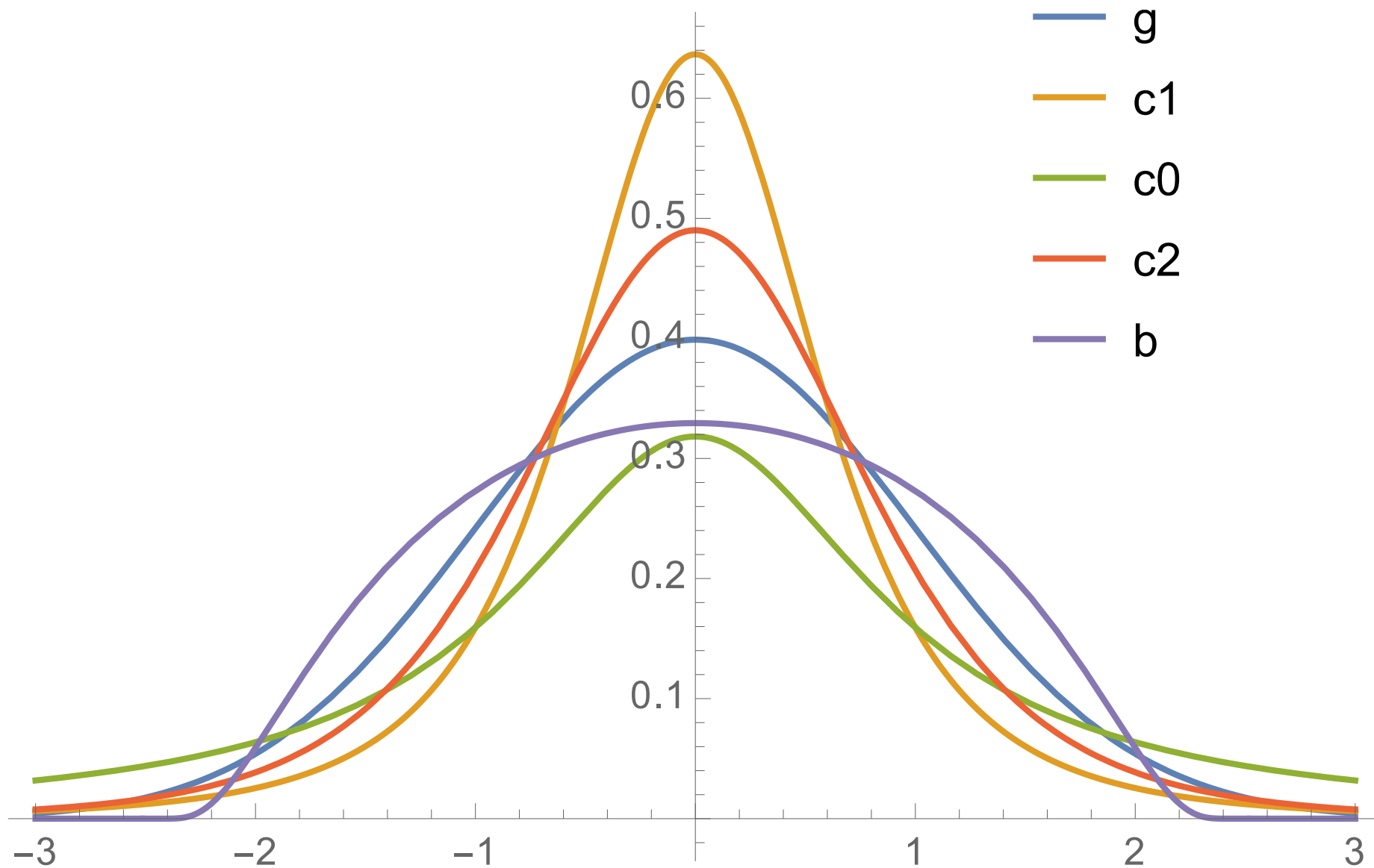
$$\delta_\epsilon^{\text{c2}}(x) = \frac{8}{3\pi} \frac{\epsilon^5}{(x^2 + \epsilon^2)^3},$$

$$\delta_\epsilon^{\text{b}}(x) = \frac{1}{N\epsilon} b(x/\epsilon), \quad b(y) = \begin{cases} \exp \left(-\frac{1}{1-y^2} \right), & -1 < y < 1, \\ 0, & \text{otherwise,} \end{cases}$$

Width of 'b' and 'c2' adjusted to coincide with Gaussian second moment.

Controlled Test

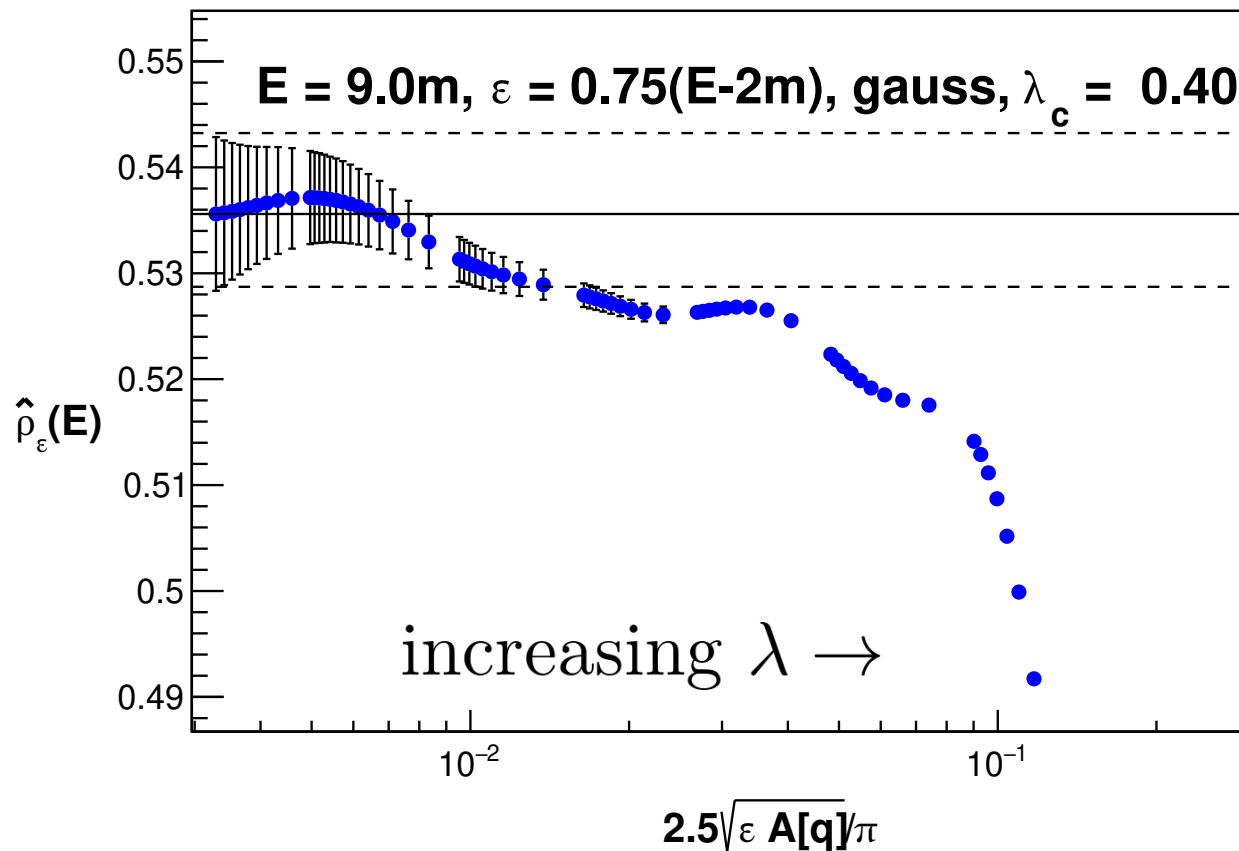
Five smearing kernels $\delta_\epsilon(E - \omega)$:



Spectral Reconstruction

$$W_\lambda[q] = (1 - \lambda)A[q]/A[0] + \lambda B[q]$$

Tradeoff parameter (λ) balances systematic (A) and statistical (B) error

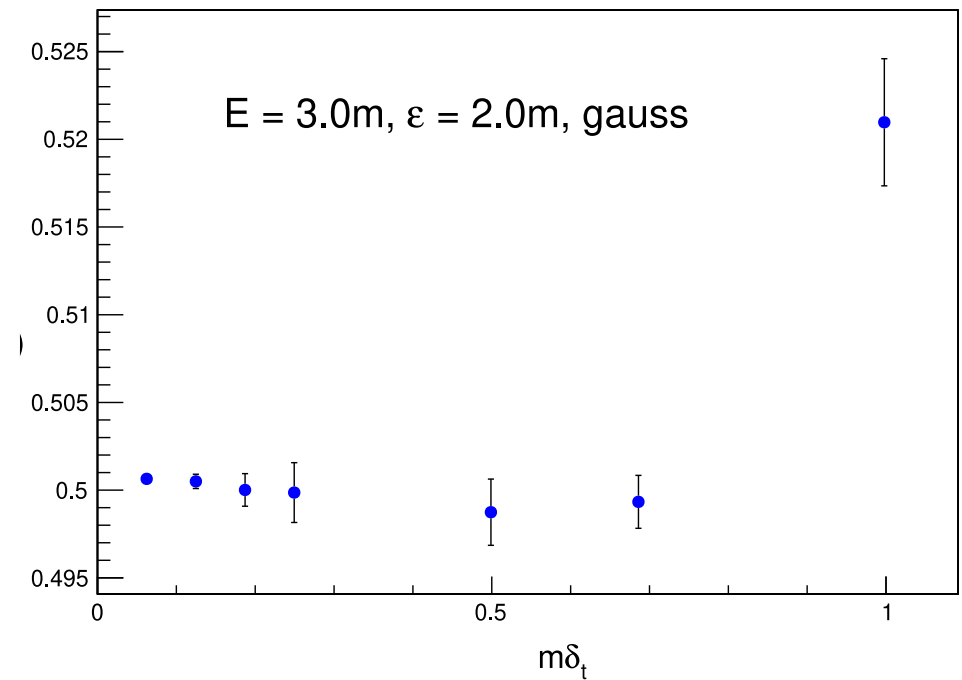
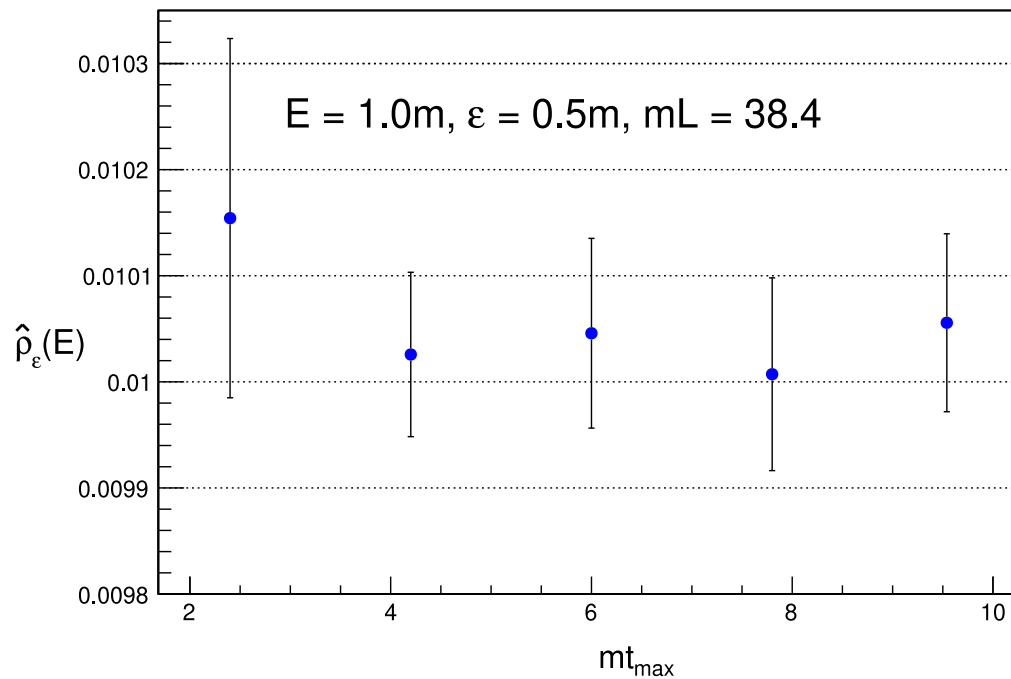


‘Plateau’ indicates statistics-limited regime. Choose point where:

$$q^* = \max_\lambda W_\lambda[q] \quad \Rightarrow \quad A[q^*]/A[0] = B[q^*]$$

Spectral Reconstruction

Variation of timeslices used in the reconstruction:

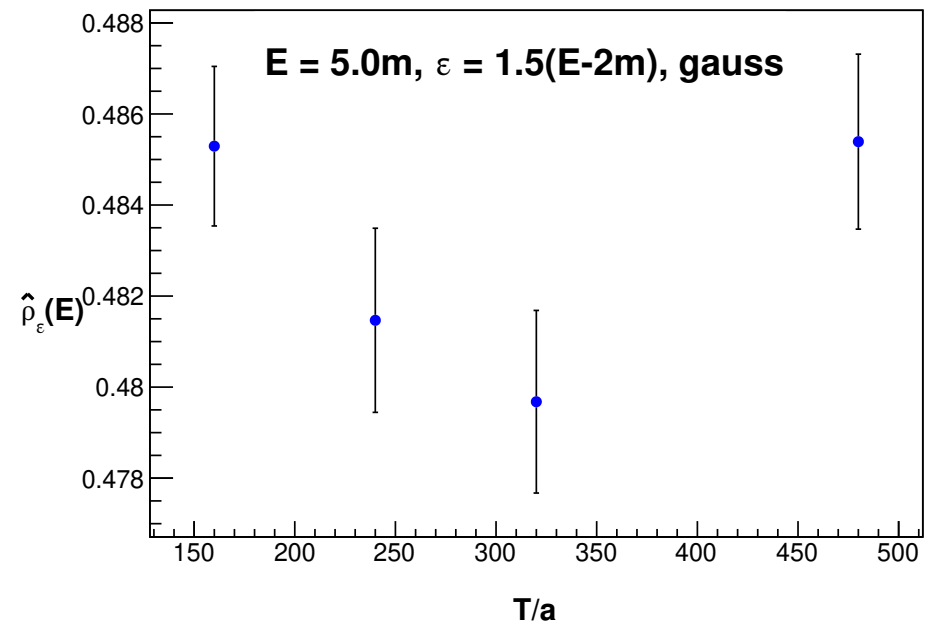
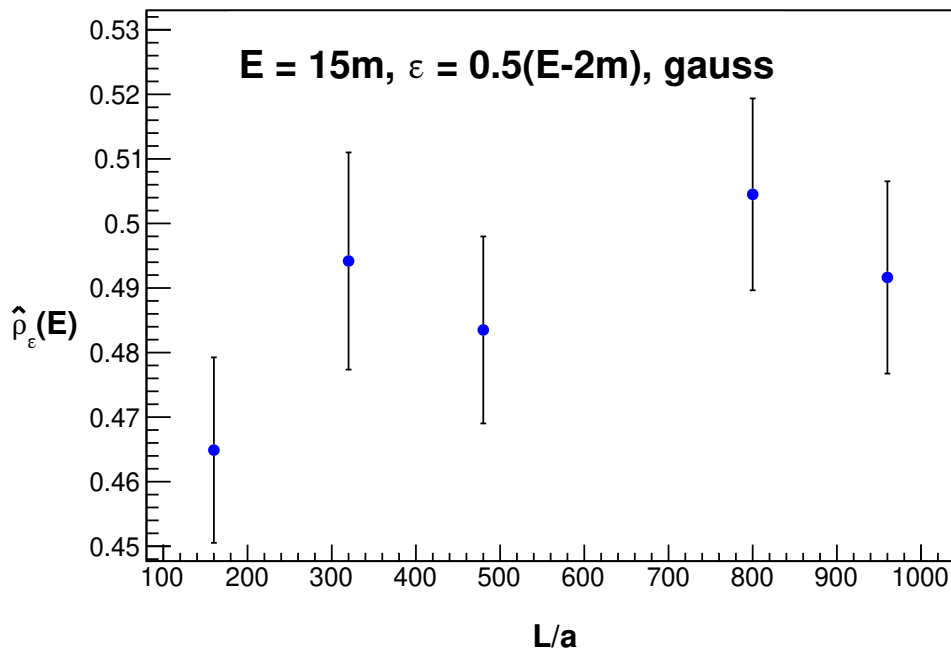


We take $t_{\max} \approx 160a$ and spacing $\delta_t = a$

Finite L and T effects

Data-driven estimate of sys. errors for each (E, ϵ) and kernel.

- Typically at or below few-percent level
- Can be large for large (E, ϵ) => investigation ongoing.

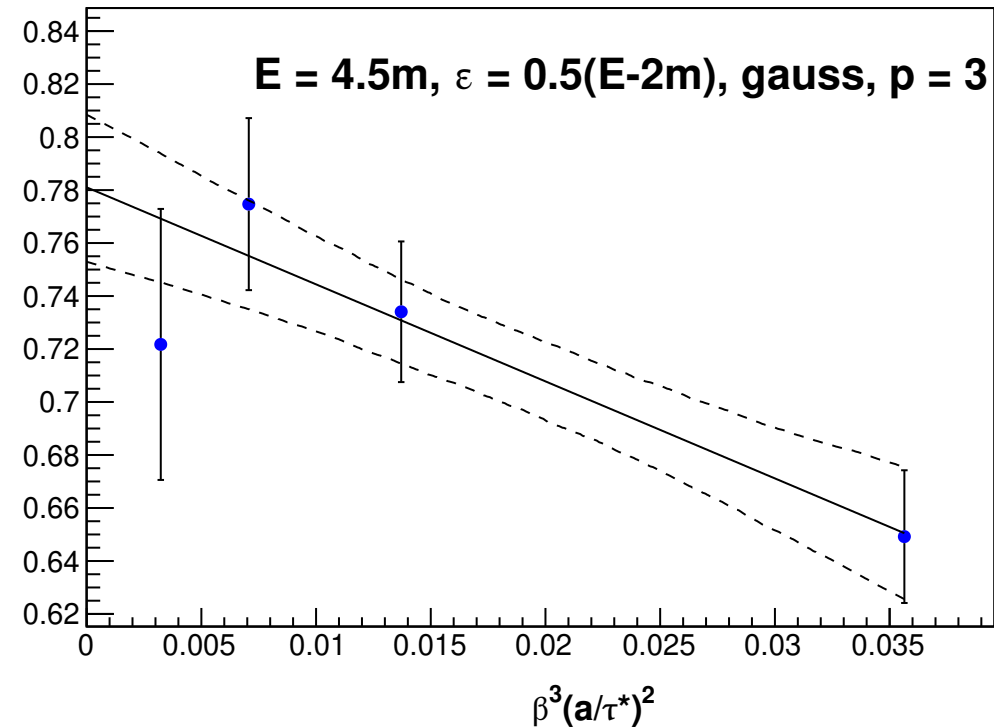
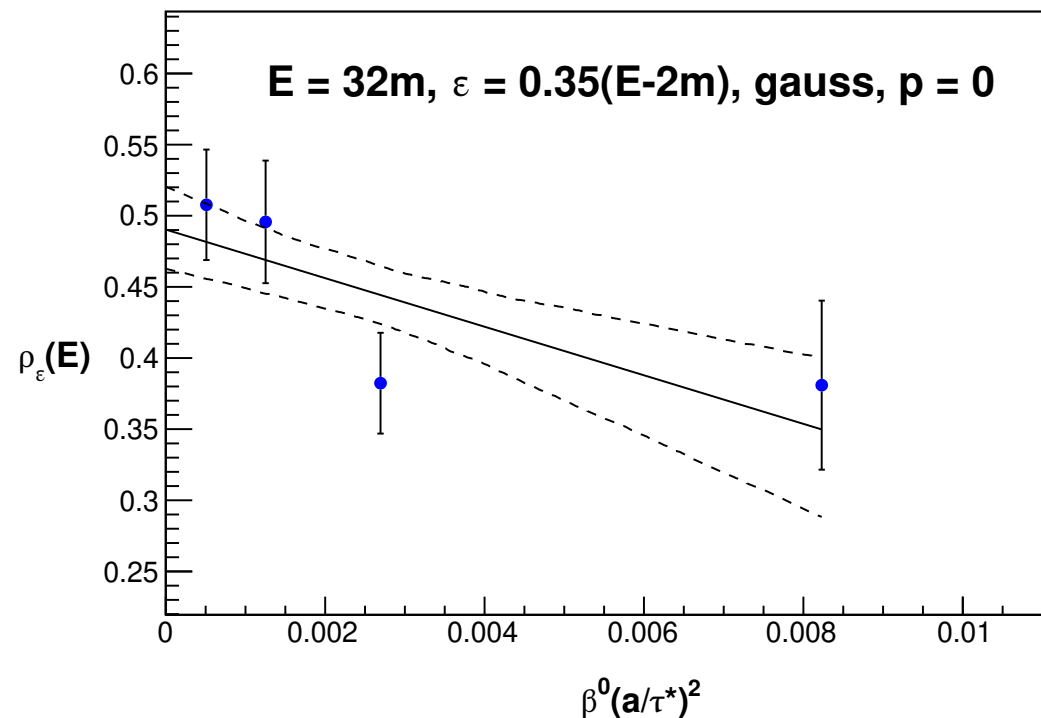


Theoretical investigations (cf. Max's talk tomorrow), higher statistics ongoing...

Continuum Limit

Long History! For spectral quantities: ... , J. Balog, F. Niedermayer, P. Weisz, Nucl. Phys. B824 (2010)

$$\lim_{a \rightarrow 0} E(a) = E^{\text{phys}} + A\beta^3 a^2 \left(1 + \frac{r}{\beta} + \frac{c}{\beta^2} + \dots\right)$$

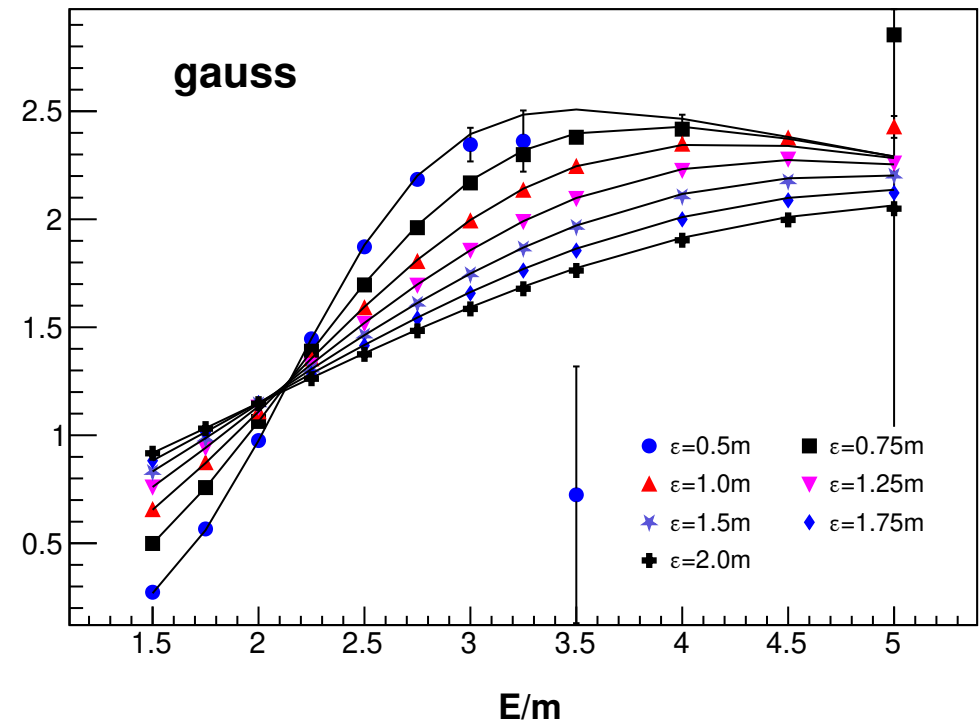
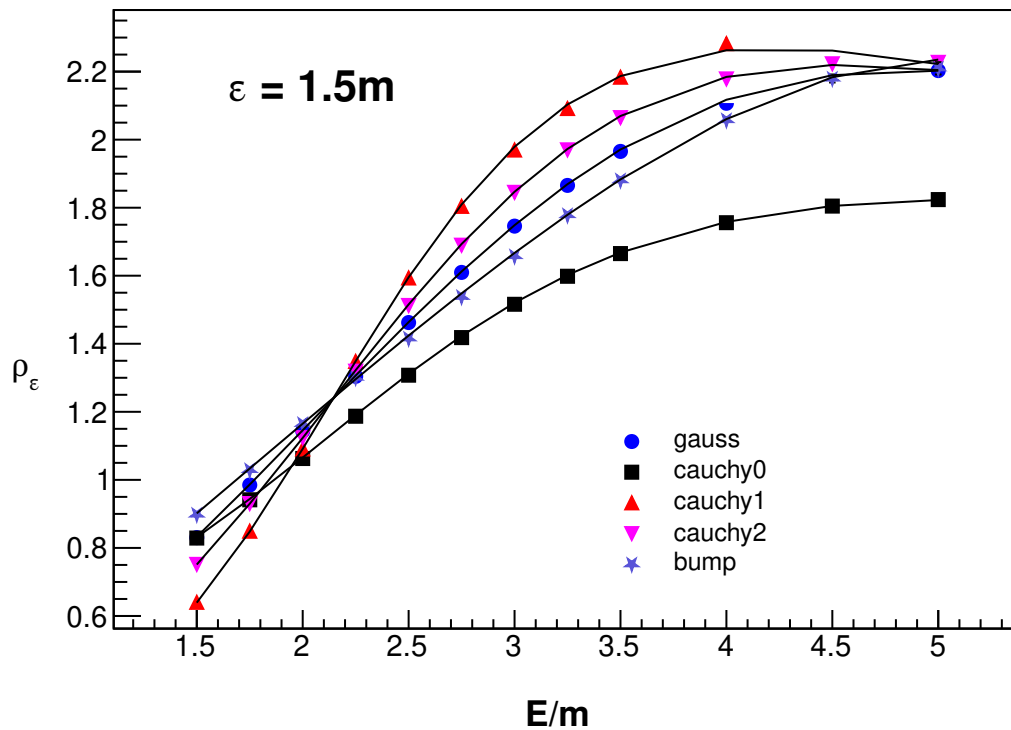


Not applicable here.

Sys. error estimate from three (arbitrary) fit forms: $\beta^p (a/m)^2$, $p = 0, 3, 6$

Results: fixed smearing width

PRELIMINARY: systematic error estimates not finalized.



Solid lines: exact smeared spectral function, using $N=2$ and 4 particle contributions.

Results: extrapolation to zero width

All kernels have the same $O(\epsilon^2)$ coefficient (up to a sign):

$$\rho_\epsilon^x(E) = \rho(E) + \sum_{k=1}^{\infty} w_k^{(x)} a_k(E) \epsilon^k ,$$

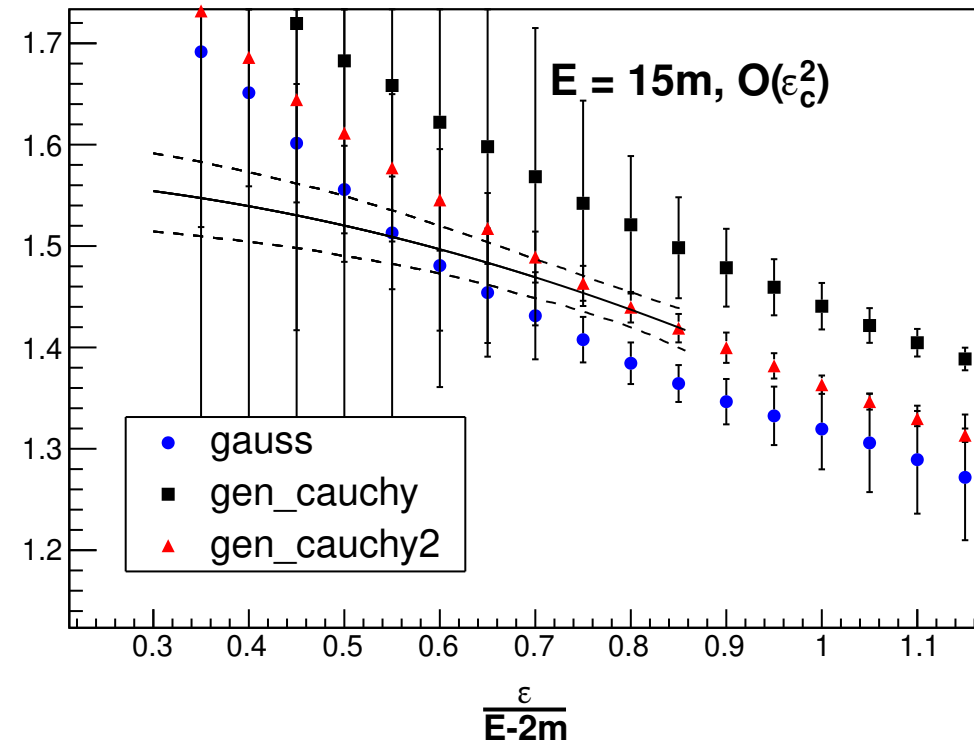
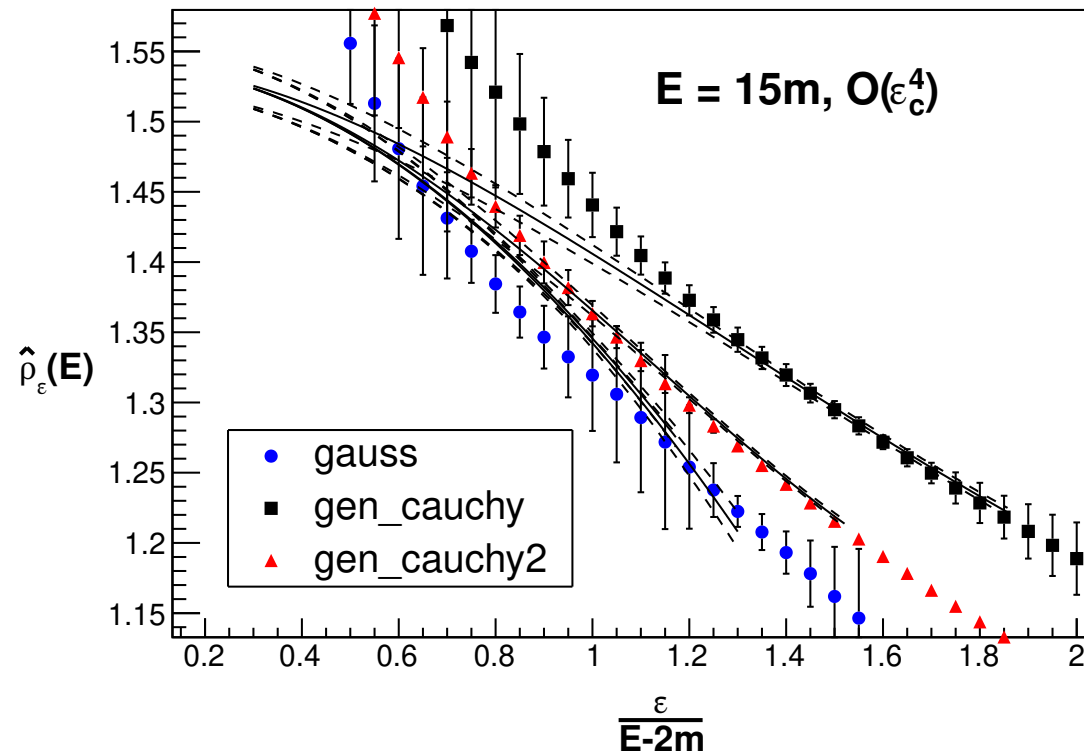
x	$w_1^{(x)}$	$w_2^{(x)}$	$w_3^{(x)}$	$w_4^{(x)}$
g	0	-1	0	1
c0	1	1	1	1
c1	0	-1	-2	-3
c2	0	-1	0	9
b	0	-1	0	2.119

Expansion coefficients worked out to arbitrary orders. (Numerically for ‘bump’)

Results: extrapolation to zero width

Fits to determine $\{a_k(E)\}$:

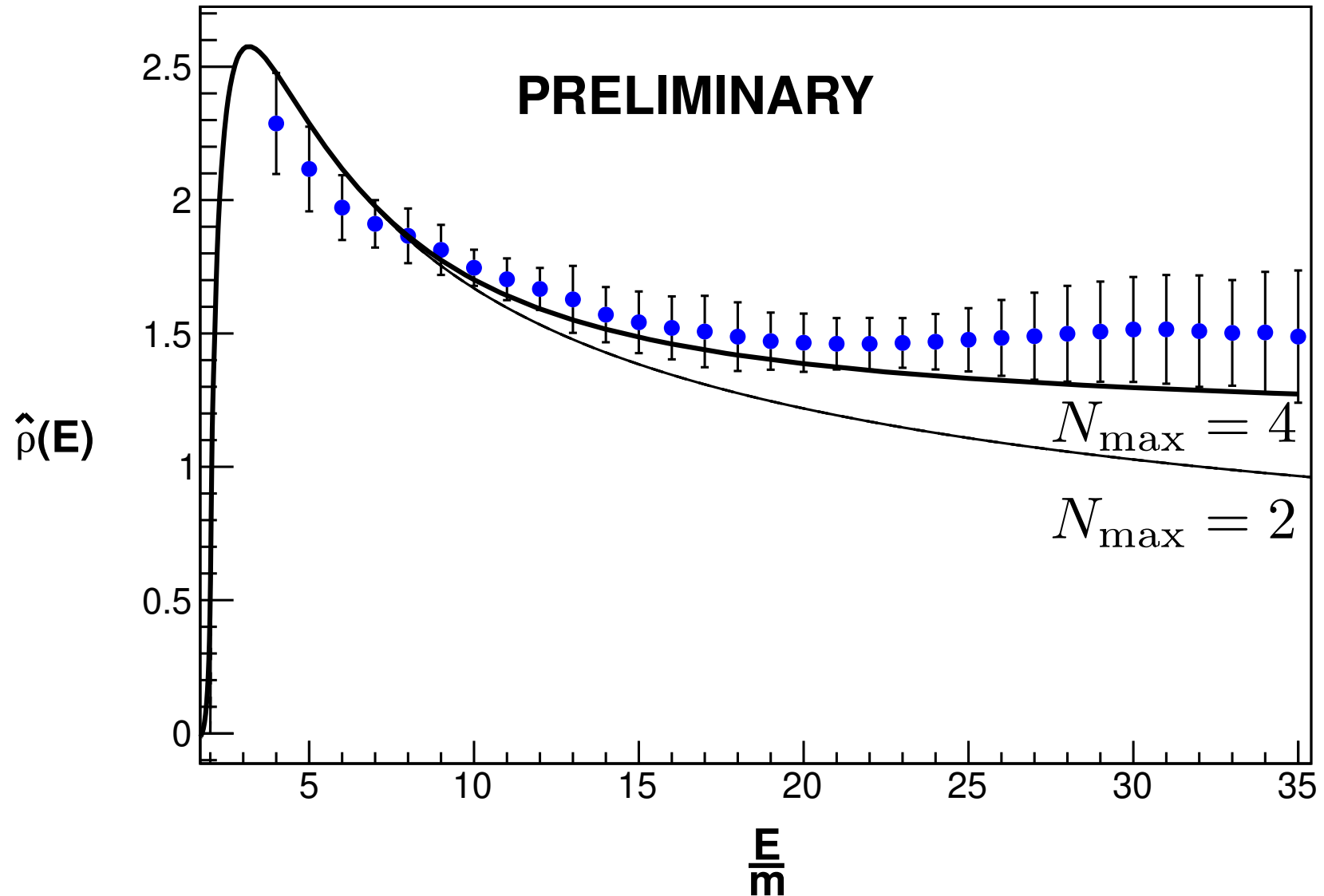
systematics estimated from $O(\epsilon^2)$, $O(\epsilon^3)$, $O(\epsilon^4)$



‘Cauchy0’ kernel not useful due to linear approach. ‘Bump’ noisy.

Results: extrapolation to zero width

Inelastic region only:



Extrapolation systematics not yet controlled below $E=4m$.

Conclusions

- Inelastic region accessible in the 2d $O(3)$ model, using the $\epsilon \rightarrow 0$ extrapolation.
- Our approach to the inverse problem is controlled and flexible, works with arbitrary smearing kernels and basis functions.
- Stay tuned for scattering amplitudes from spectral functions (cf. Max's talk tomorrow)
- Application to QCD:
 - Signal-to-noise and precision similar to vector-vector correlator
 - Masterfield paradigm ideal: see Marco Ce' and Patrick Fritzsch talks at Lattice21