

Few-body reactions in neutron-rich systems

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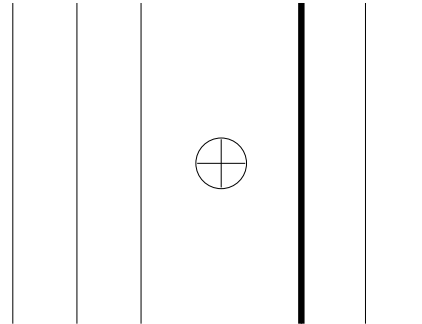
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Outline

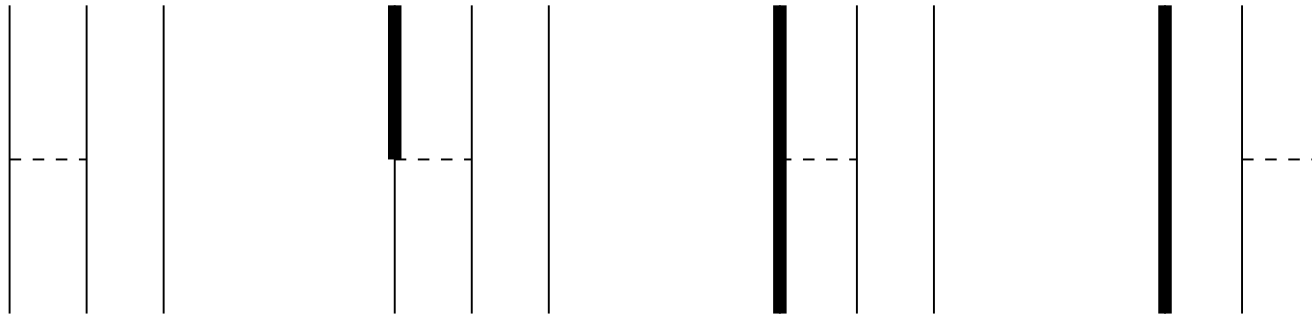
- core excitation in 3-body nuclear reactions:
 - extended Faddeev/AGS formalism
 - $^{20}\text{O}(d, p)$
- trineutron resonances
- tetraneutron resonances
- four-fermion universality

Core excitation (CX): extended Hilbert space

$$\mathcal{H} = \mathcal{H}_g \oplus \mathcal{H}_x$$

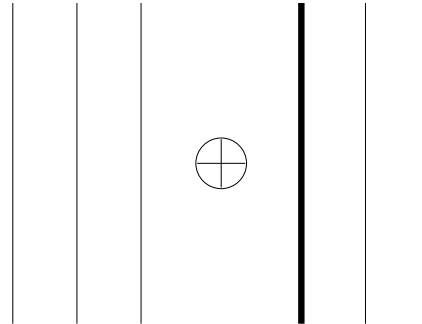


sector coupling by interaction

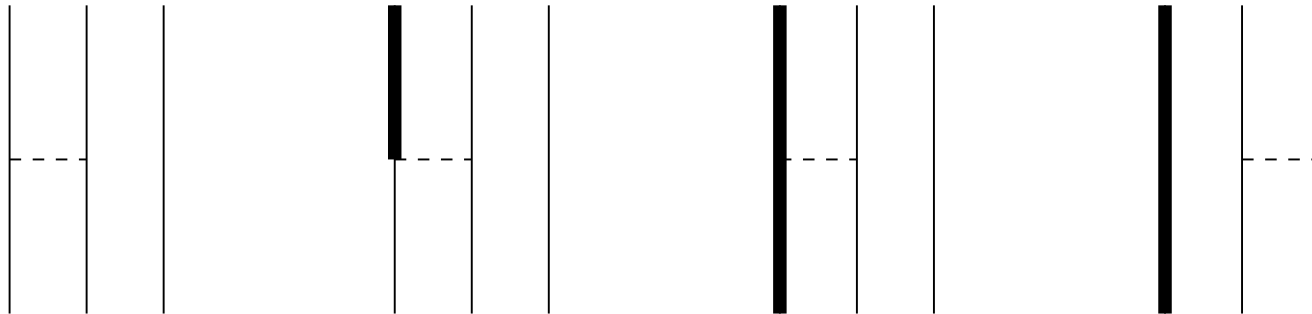


Core excitation (CX): extended Hilbert space

$$\mathcal{H} = \mathcal{H}_g \oplus \mathcal{H}_x$$



sector coupling by interaction



standard operator form of 3-body AGS equations

with $H_0 \rightarrow H_0 + h_A^{\text{int}}$

$$h_A^{\text{int}} |\mathcal{H}_a\rangle = (m_{A^*} - m_A) \delta_{ax} |\mathcal{H}_a\rangle$$

3-body AGS equations with core excitation

$$U_{\beta\alpha}^{ba} = \bar{\delta}_{\beta\alpha} \delta_{ba} G_0^{-1} + \sum_{\sigma} \sum_j \bar{\delta}_{\beta\sigma} T_{\sigma}^{bj} G_0 U_{\sigma\alpha}^{ja}$$

$$U_{0\alpha}^{ba} = \delta_{ba} G_0^{-1} + \sum_{\sigma} \sum_j T_{\sigma}^{bj} G_0 U_{\sigma\alpha}^{ja}$$

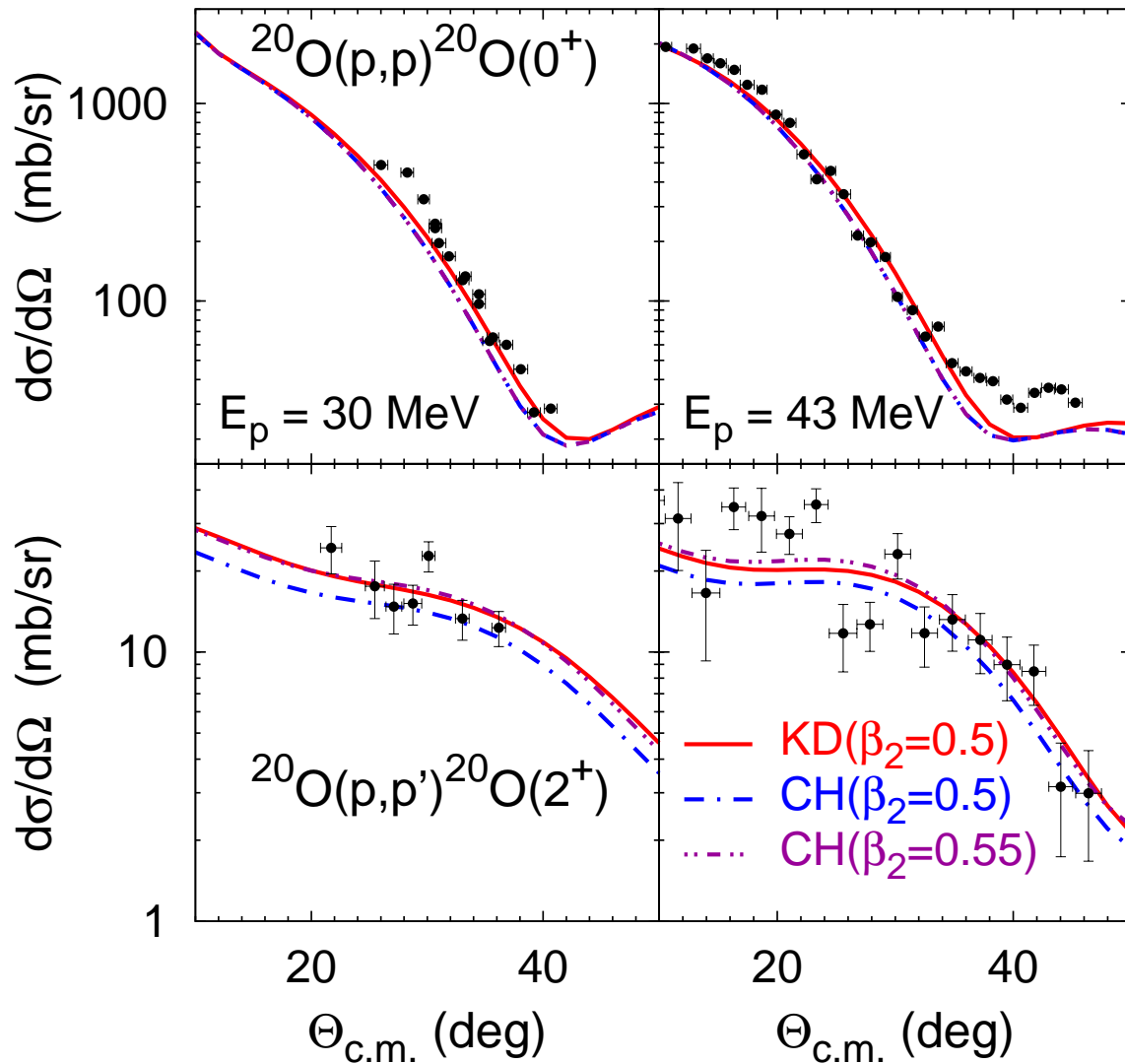
$$T_{\sigma}^{ba} = v_{\sigma}^{ba} + \sum_j v_{\sigma}^{bj} G_0 T_{\sigma}^{ja}$$

$$G_0 = (E + i0 - H_0)^{-1}$$

channel states $(E - H_0)|\phi_{\alpha}^a\rangle = \sum_j v_{\alpha}^{aj} |\phi_{\alpha}^j\rangle$

$$H_0 |\mathbf{p}_{\alpha} \mathbf{q}_{\alpha}\rangle^a = [p_{\alpha}^2/2\mu_{\alpha} + q_{\alpha}^2/2M_{\alpha} + (m_{A^*} - m_A)\delta_{ax}] |\mathbf{p}_{\alpha} \mathbf{q}_{\alpha}\rangle^a$$

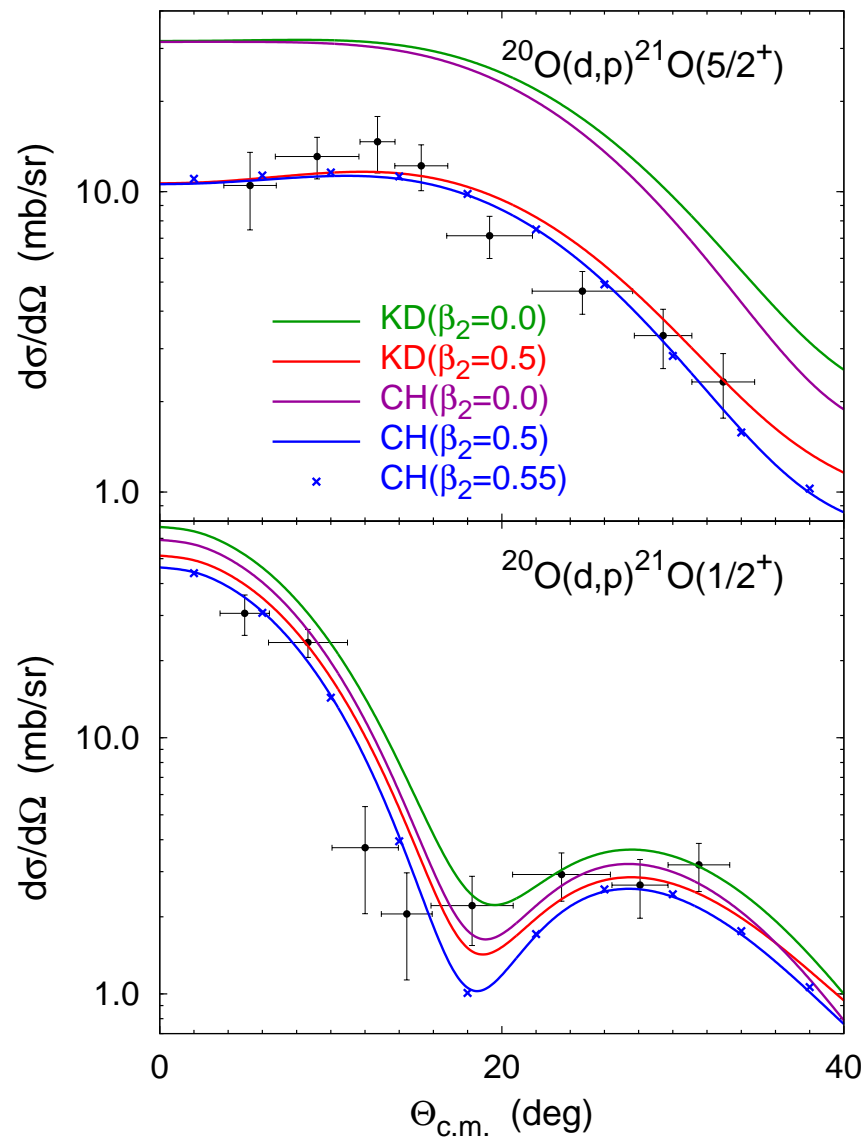
Potential test: N + ^{20}O



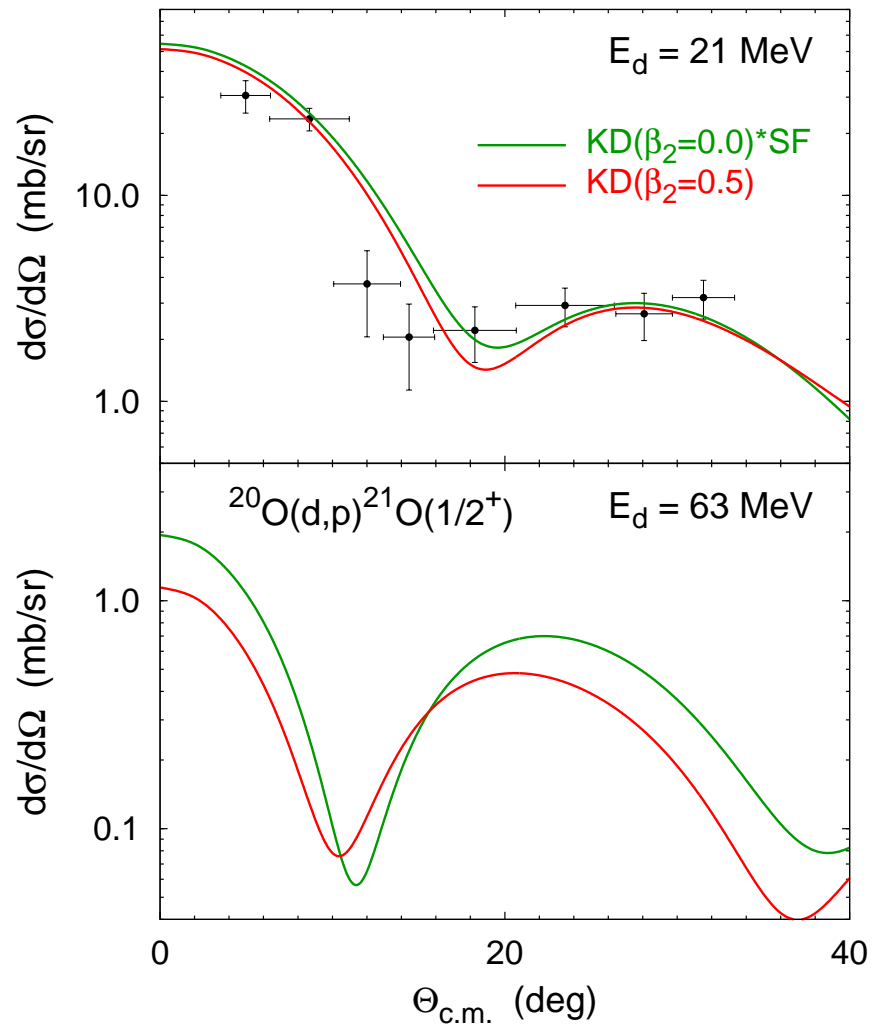
Vibrational model for V_{NA}

Shell-model SF for ^{21}O : $0.34(\frac{5}{2}^+)$, $0.82(\frac{1}{2}^+)$

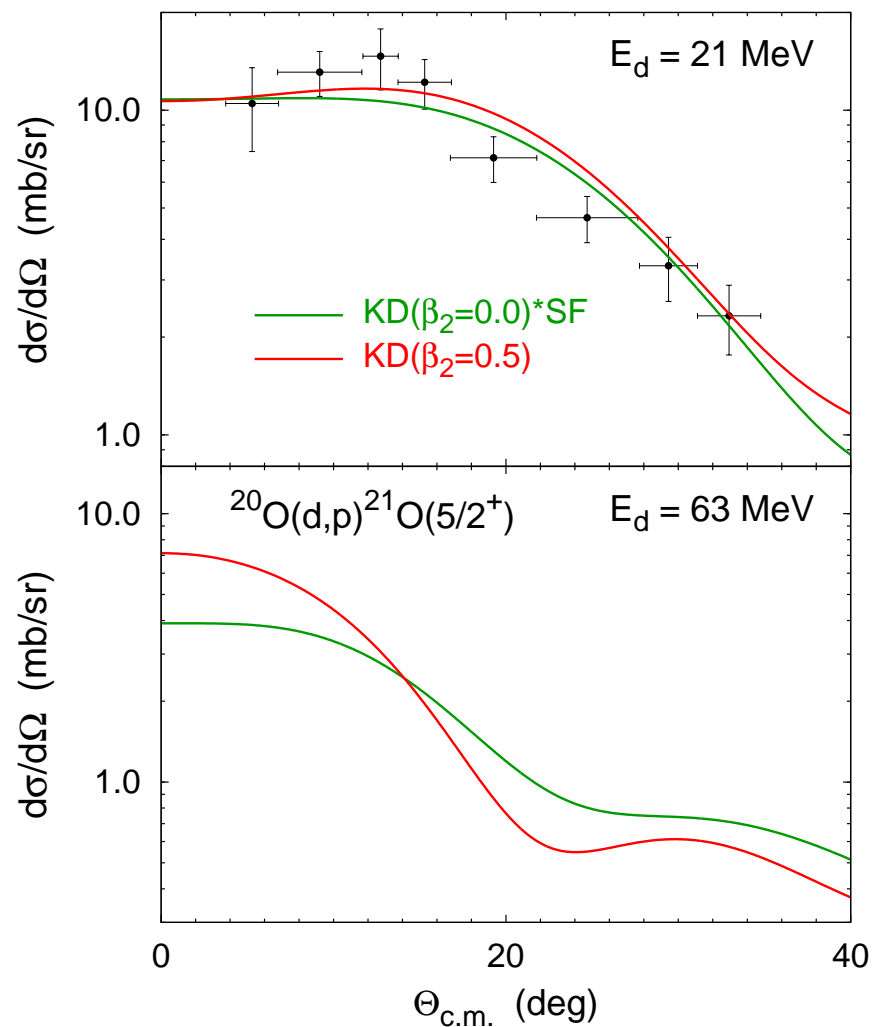
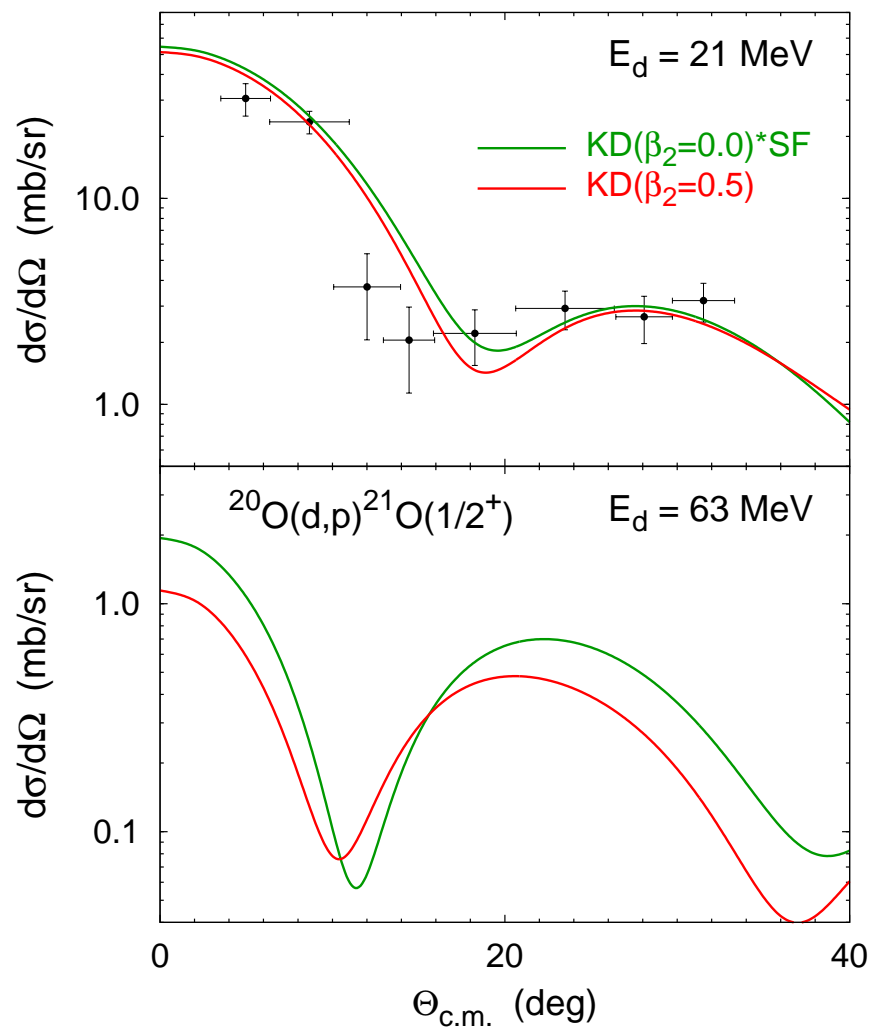
$^{20}\text{O}(d,p)^{21}\text{O}$ at 21 MeV



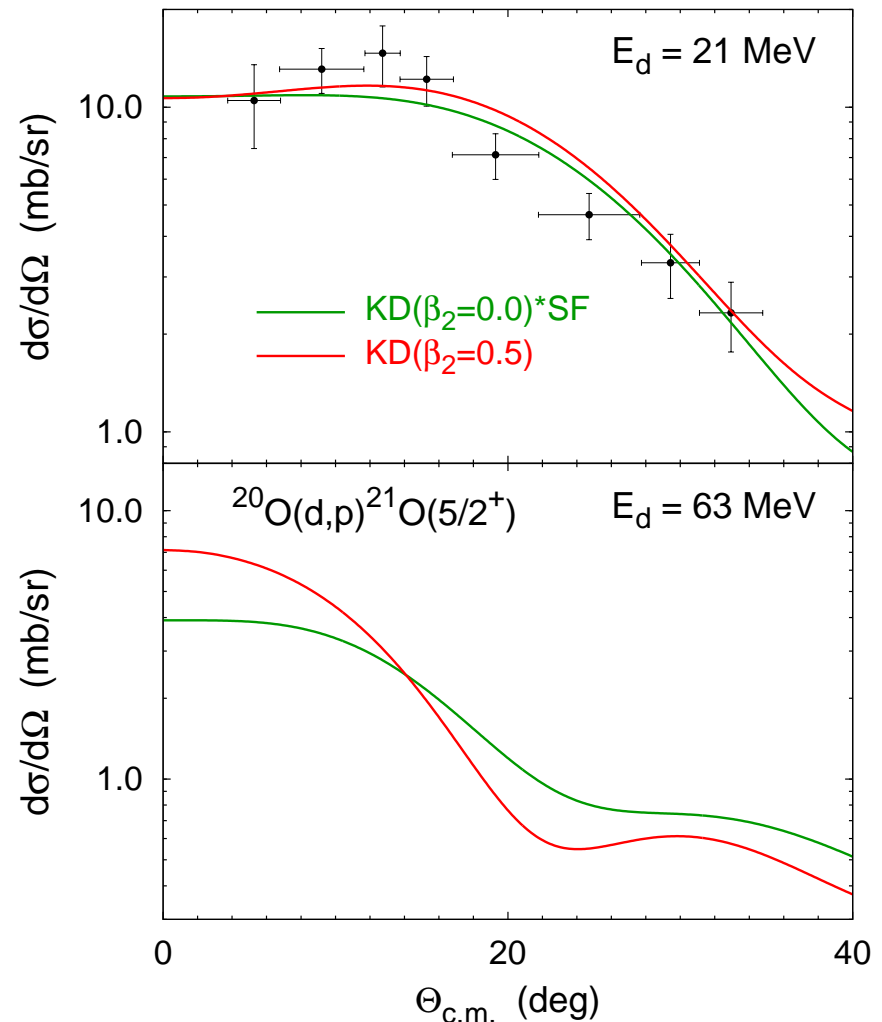
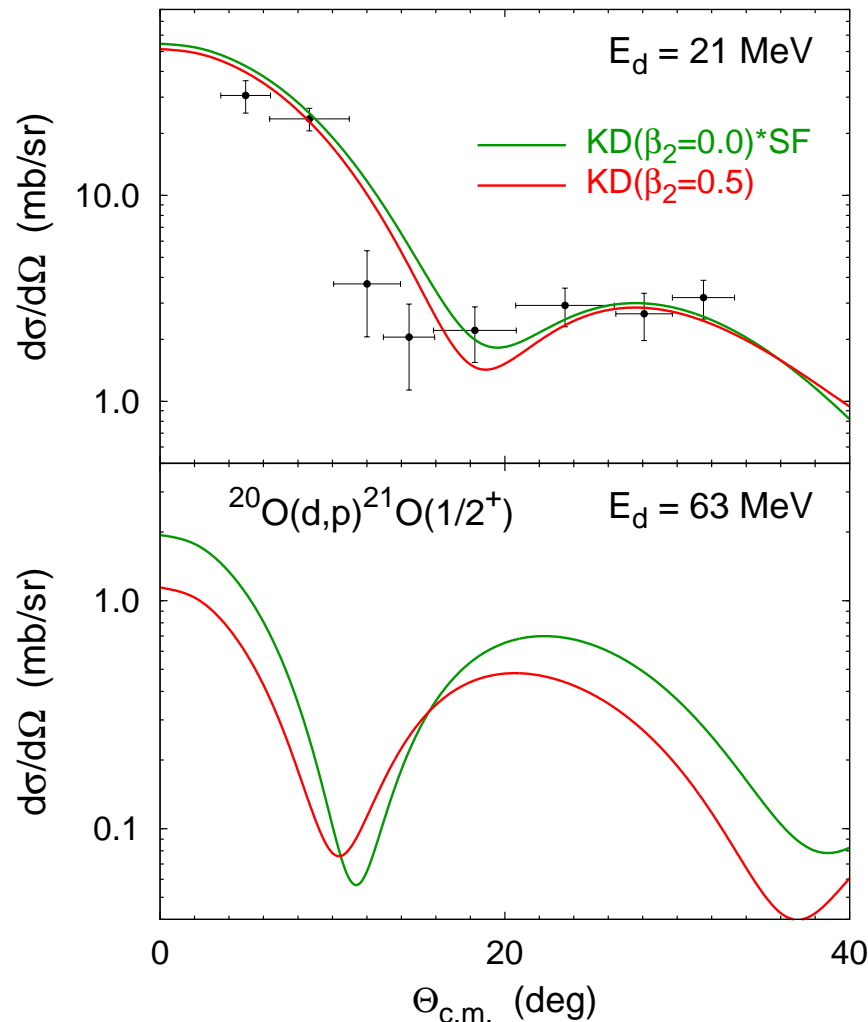
$^{20}\text{O}(d,p)^{21}\text{O}$: extracting SF?



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SF = $\sigma_{\text{exp}}/\sigma_{\text{SP}}$ in general unreliable !

Faddeev/AGS: (V_{NA} - SF - data) compatibility check

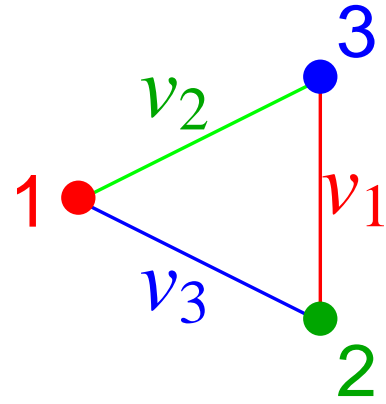
Experiment and theory: 3n

- no experimental evidence for 3n resonance
- Hemmdan, Kamada, Glöckle, PRC 66, 054001,
Lazauskas, Carbonell, Hiyama, Kamimura,
PRC 71, 044004, PRC 93, 044004:
no observable 3n resonance
- Gandolfi, Hammer, Klos, Lynn, Schwenk,
PRL 118, 232501:
3n resonance $E_r = 1.1$ MeV

Experiment and theory: 4n

- Kisamori et al, PRL 116, 052501:
few events in ${}^4\text{He}({}^8\text{He}, {}^8\text{Be}) \rightarrow 4n$ state at
 $E_r = 0.83 \pm 0.65 \pm 1.25$ MeV, $\Gamma < 2.6$ MeV
- Lazauskas, Carbonell, Hiyama, Kamimura,
PRC 72, 034003, PRC 93, 044004:
no observable 4n resonance
- Shirokov, Papadimitriou, Mazur, Mazur, Roth, Vary,
PRL 117, 182502:
4n resonance $(E_r, \Gamma) = (0.8, 1.4)$ MeV
- Gandolfi, Hammer, Klos, Lynn, Schwenk, PRL 118,
232501: 4n resonance $E_r = 2.1$ MeV
- Fosse, Rotureau, Michel, Ploszajczak, PRL 119,
032501: 4n resonance $(E_r, \Gamma) = (7.3, 3.8)$ MeV

Three-particle system



Hamiltonian $H_0 + \sum_{\alpha} v_{\alpha}$

Alt, Grassberger, and Sandhas equations

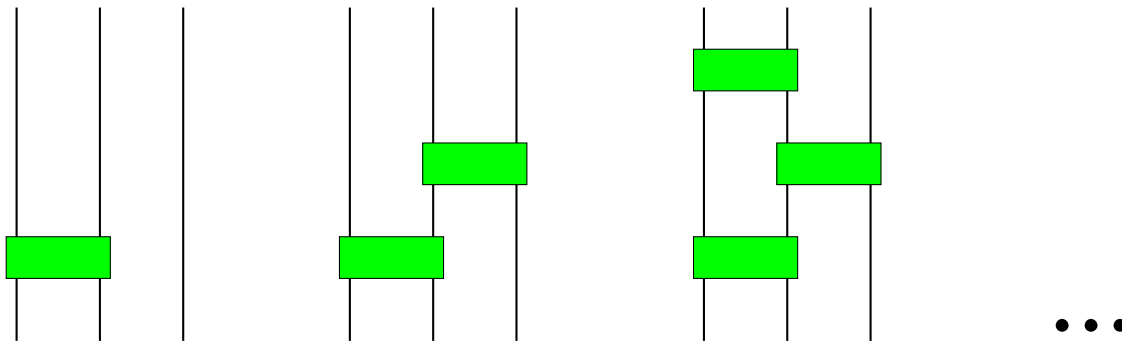
$$U_{\beta\alpha} = \bar{\delta}_{\beta\alpha} G_0^{-1} + \sum_{\sigma} \bar{\delta}_{\beta\sigma} t_{\sigma} G_0 U_{\sigma\alpha}$$

$$U_{00} = \sum_{\alpha} t_{\alpha} + \sum_{\beta\alpha} t_{\beta} G_0 U_{\beta\alpha} G_0 t_{\alpha}$$

$$t_{\sigma} = v_{\sigma} + v_{\sigma} G_0 t_{\sigma}$$

$$G_0 = (E + i0 - H_0)^{-1}$$

channel states $(E - H_0 - v_{\alpha})|\phi_{\alpha}\rangle = 0$



Symmetrized AGS equations

$$U = PG_0^{-1} + PtG_0U$$

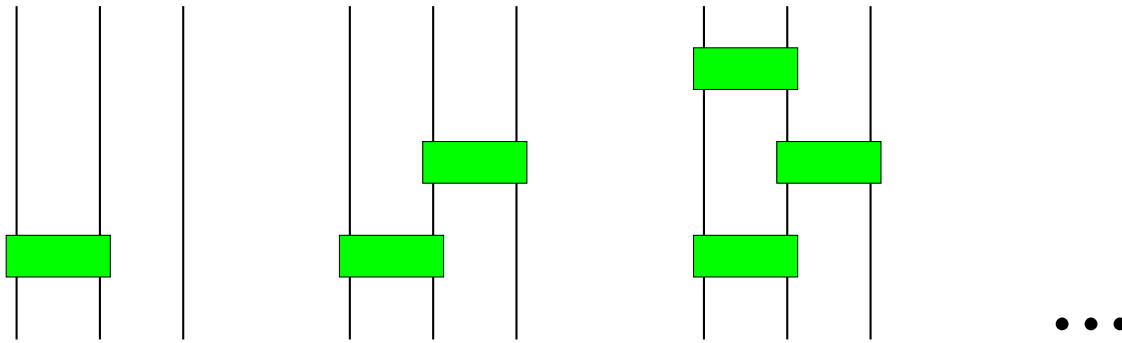
$$T = tG_0UG_0t$$

$$T = tG_0Pt + tG_0PT$$

$$U_{00} = (1 + P)t(1 + P) + (1 + P)T(1 + P)$$

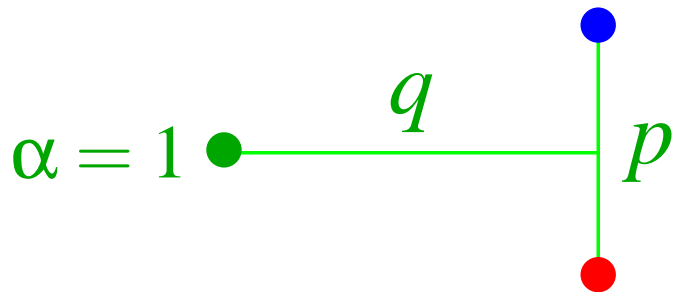
$$t = v + vG_0t$$

$$P = P_{12}P_{23} + P_{13}P_{23}$$



Resonance: pole of T

$$T_{J\Pi} = \sum_{n=-1}^{\infty} \tilde{T}_{J\Pi}^{(n)} (E - E_r + i\Gamma/2)^n$$



Basis states $|pq(l\{[L(s_2s_3)s]js_1\}S)JM\rangle$

Trial states:

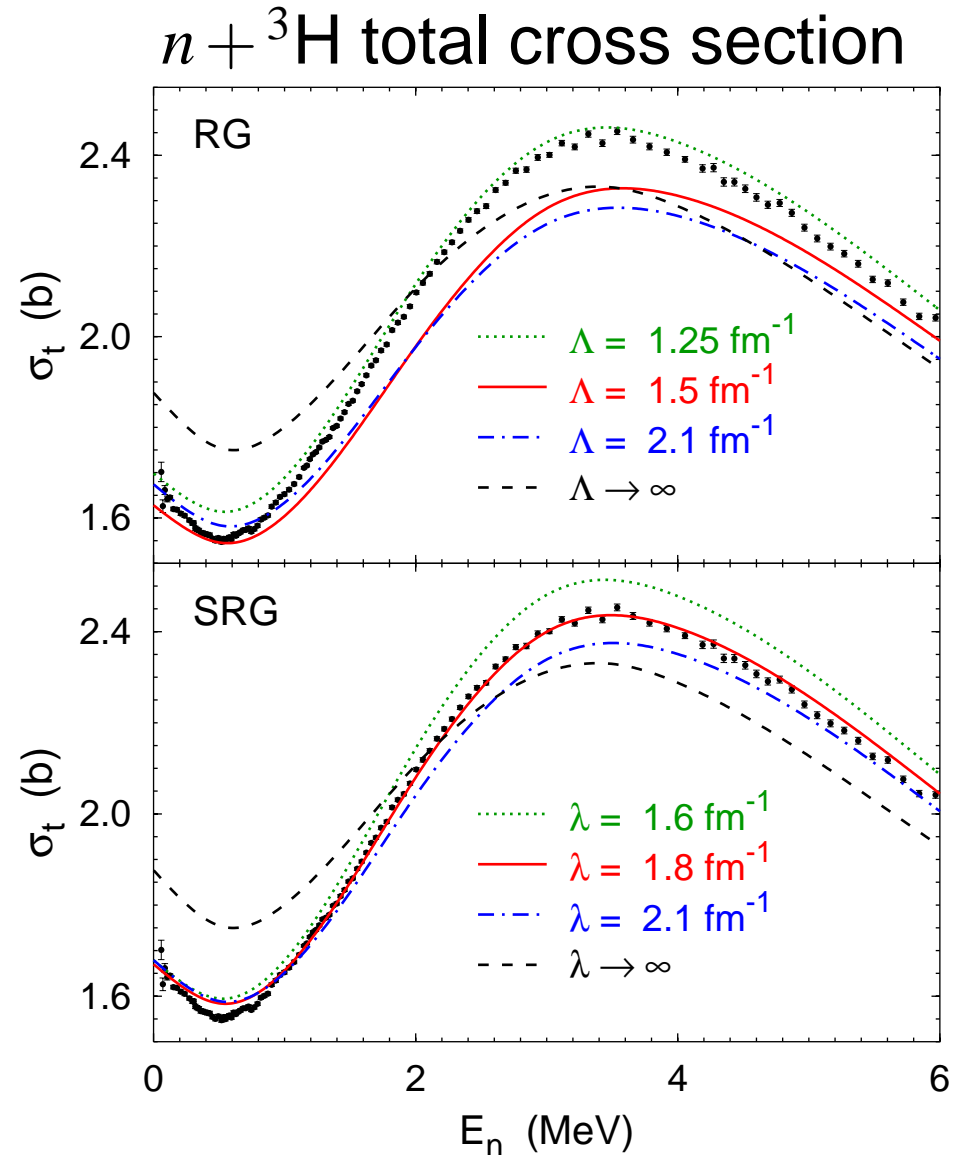
- p-state ${}^{2s+1}L_j^p$: $p = \sqrt{mE}$, $q = 0$
- q-state ${}^1S_0^q$: $p = 0$, $q = \sqrt{4mE/3}$
- off-shell state ${}^1S_0^{\text{off}}$: Gaussian for p , $q = \sqrt{4m(E + \epsilon_{\text{off}})/3}$

Force models

- CD Bonn
- Reid93
- NLO [Epelbaum et al, PRL 115, 122301]
- **SRG** of AV18
with $\lambda = 1.8 \text{ fm}^{-1}$

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Enhancing nn interaction

$v = v_{nn}$ for 1S_0 (bound if enhanced by ~ 1.1)

$v = f v_{nn}$ for higher waves $L > 0$

Bound dineutron
at threshold:

$^{2s+1}L_j$	f
	SRG
$^3P_2 - ^3F_2$	7.24
3P_0	7.97
	Reid93
$^3P_2 - ^3F_2$	4.00
3P_0	5.95

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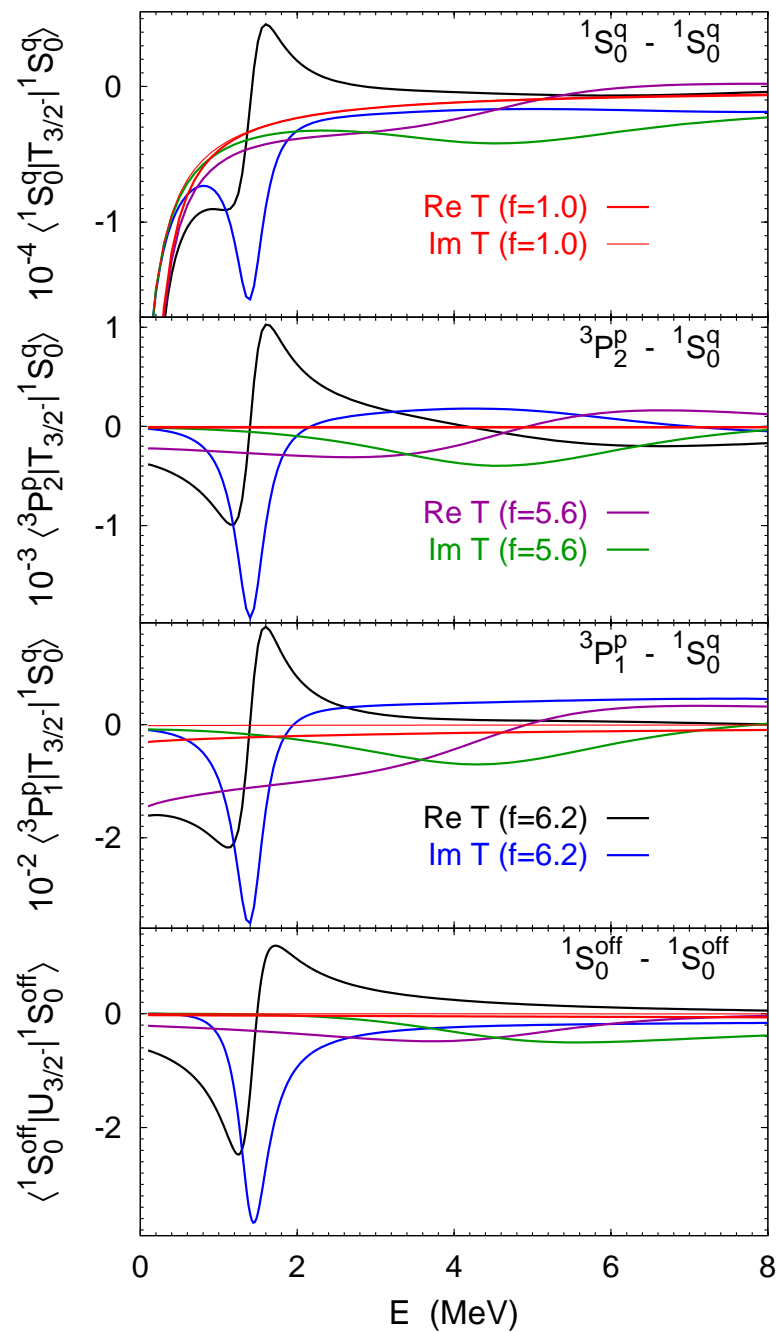
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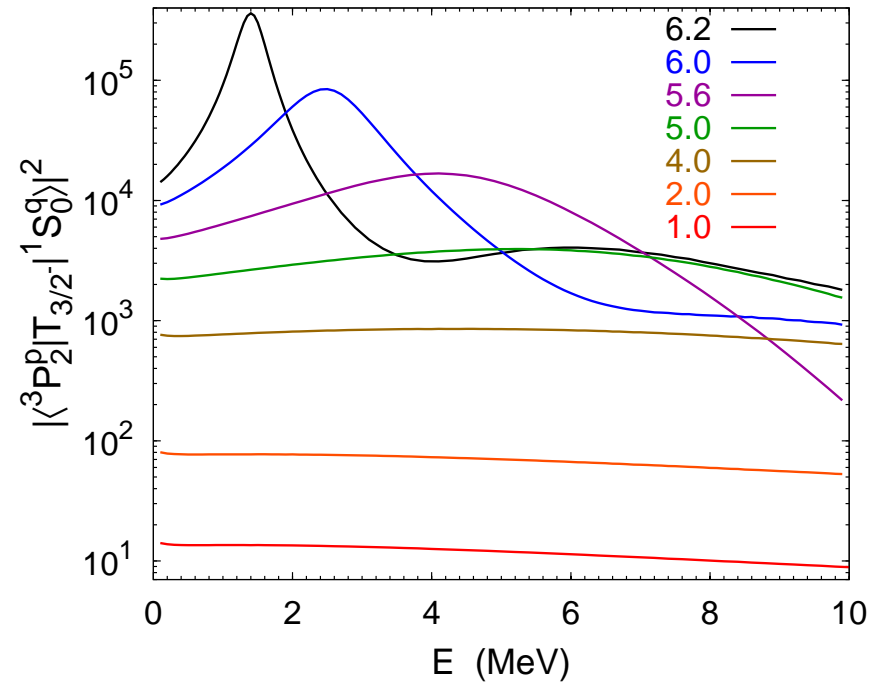
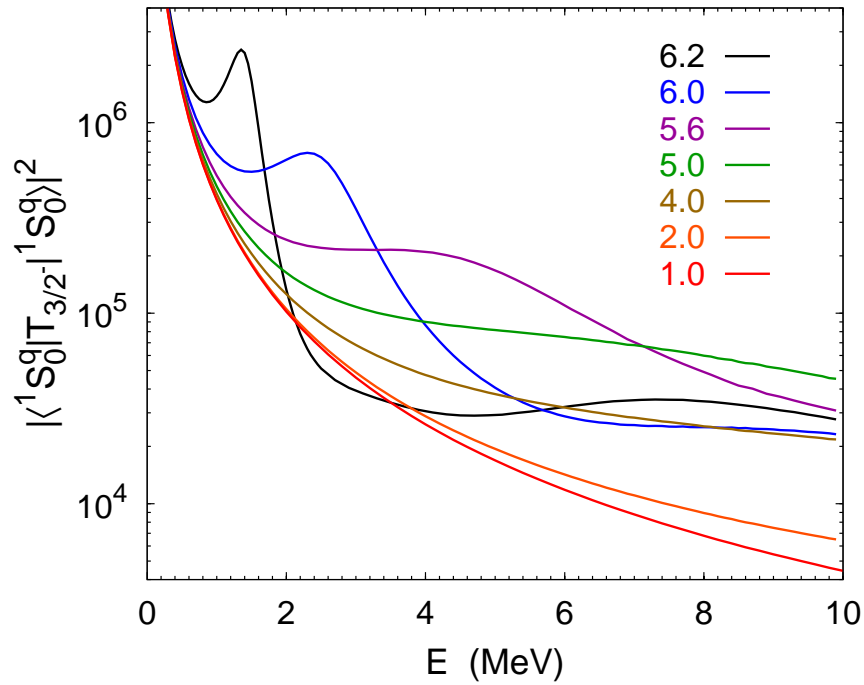
Bound trineutron
at threshold:

J^Π	$f(\text{SRG})$
$1/2^+$	> 7.24
$1/2^-$	> 7.24
$3/2^+$	6.71
$3/2^-$	6.42
$5/2^+$	6.02
$5/2^-$	6.94

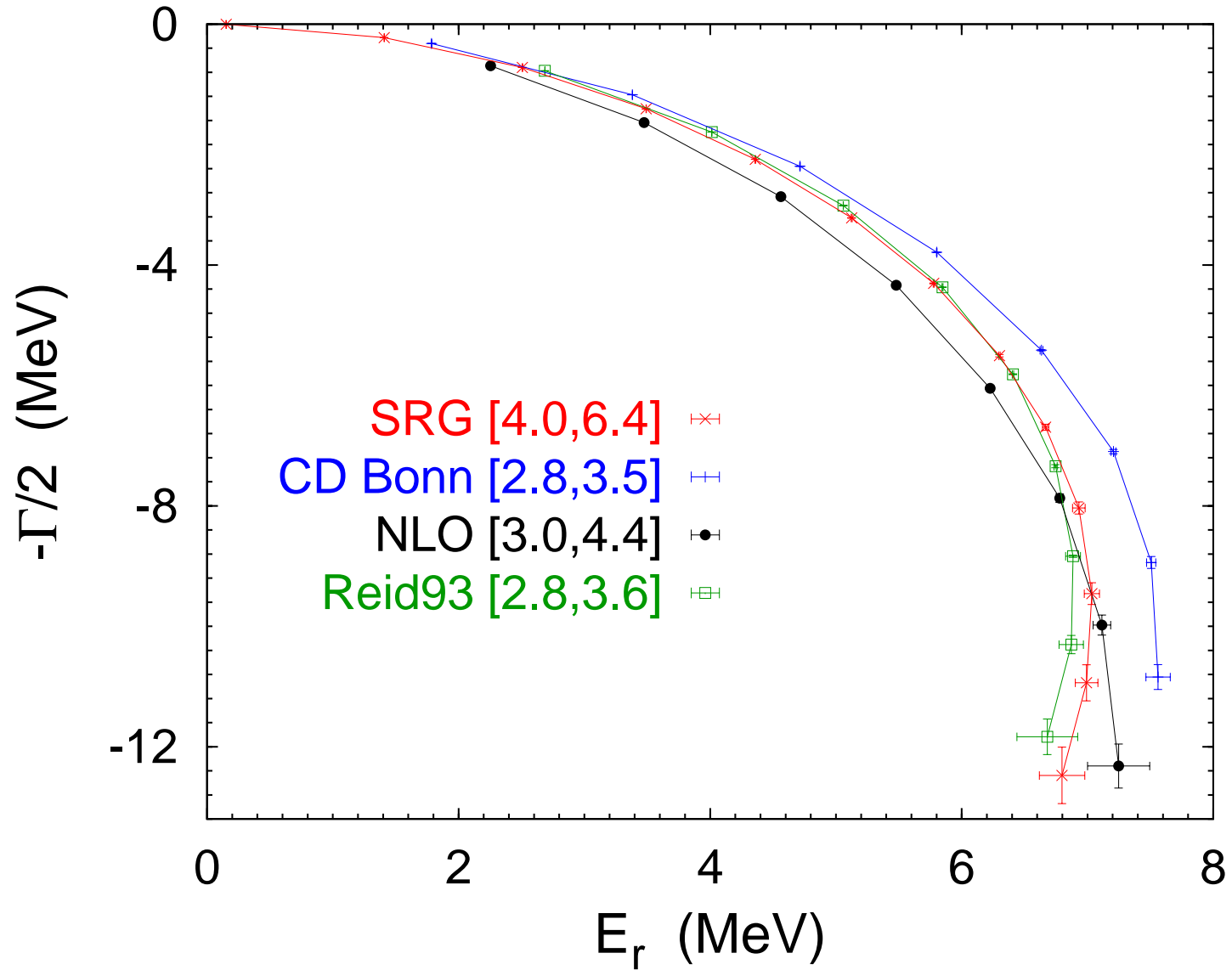
Resonance in all $J^\Pi = 3/2^-$ channels



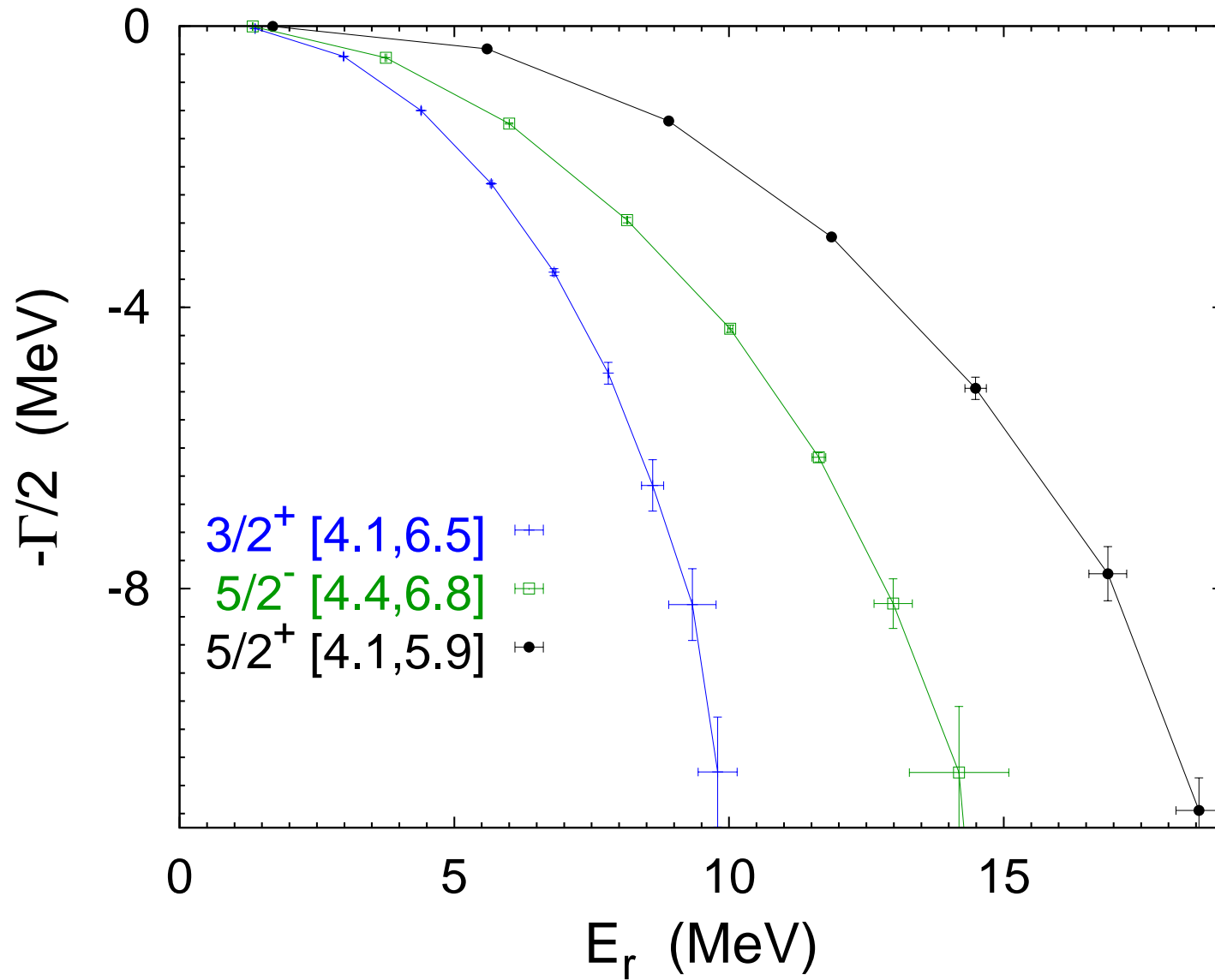
No $3/2^-$ resonance at physical strength



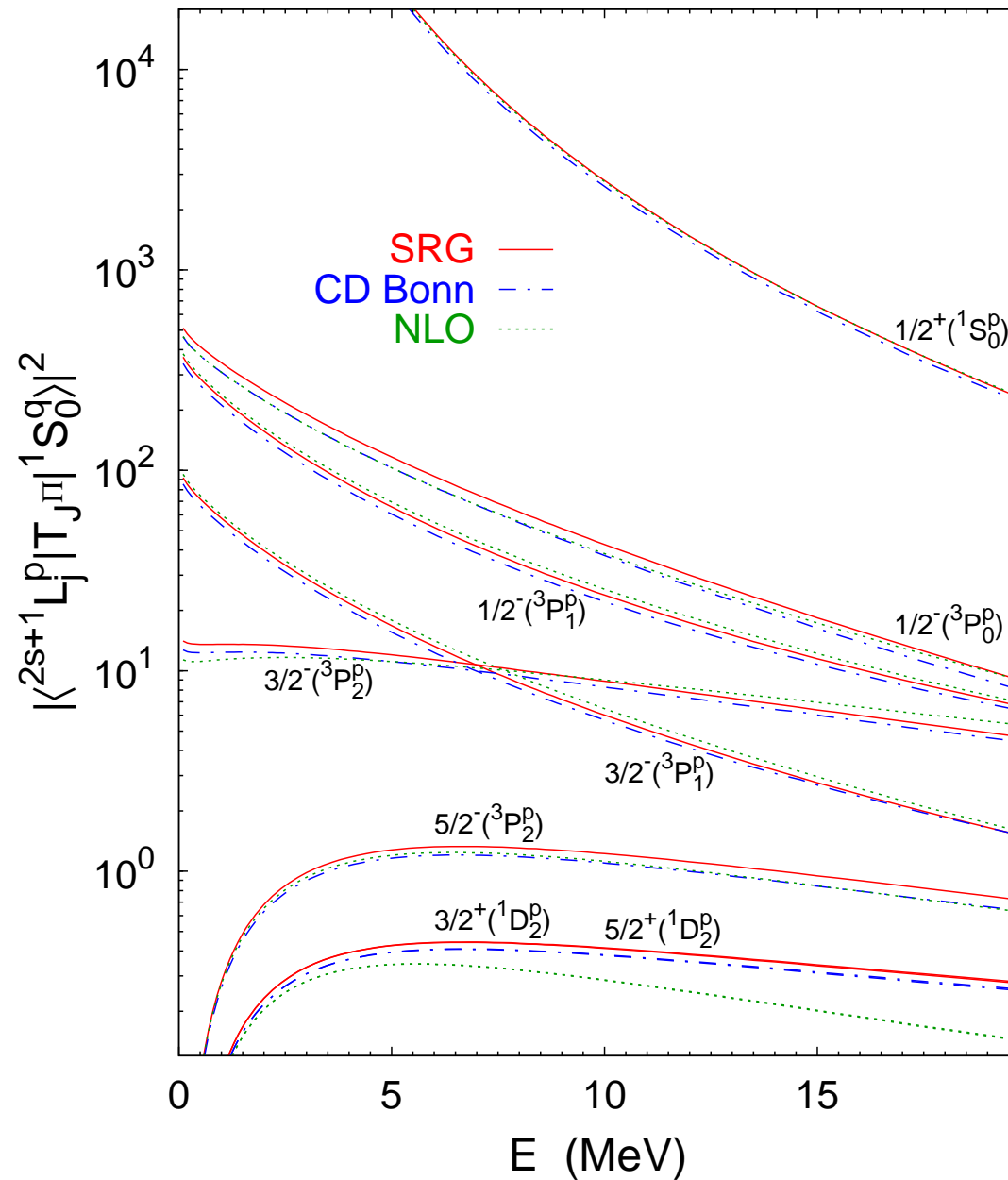
$3/2^-$ resonance trajectory



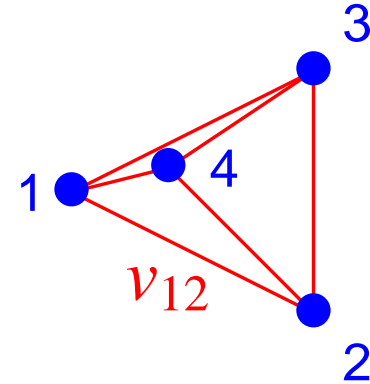
Resonance trajectories



$3n \rightarrow 3n$ transition strengths: resonance?



4N scattering



Hamiltonian $H_0 + \sum_{i>j} v_{ij}$

- Wave function:
Schrödinger equation (HH + Kohn VP, r -space)
[M. Viviani, A. Kievsky, L. E. Marcucci, S. Rosati, L. Girlanda]
- Wave function components:
Faddeev-Yakubovsky equations (r -space)
[R. Lazauskas, J. Carbonell]
- Transition operators:
Alt-Grassberger-Sandhas equations (p -space)
[AD, A. C. Fonseca]

4-body scattering: AGS equations

4-body transition operators

$$t_i = v_i + v_i G_0 t_i$$

$$G_0 = (E + i0 - H_0)^{-1}$$

$$U_\gamma^{jk} = G_0^{-1} \bar{\delta}_{jk} + \sum_i \bar{\delta}_{ji} t_i G_0 U_\gamma^{ik}$$

$$\mathcal{U}_{\beta\alpha}^{ji} = (G_0 t_i G_0)^{-1} \bar{\delta}_{\beta\alpha} \delta_{ji} + \sum_{\gamma k} \bar{\delta}_{\beta\gamma} U_\gamma^{jk} G_0 t_k G_0 \mathcal{U}_{\gamma\alpha}^{ki}$$

i, j, k : pairs (\equiv three-cluster (2+1+1) partitions)

α, β, γ : two-cluster (1+3 or 2+2) partitions

4-body scattering: AGS equations

4-body transition operators

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$$\mathcal{U}_{\beta\alpha}^{ji} = (G_0 t_i G_0)^{-1} \bar{\delta}_{\beta\alpha} \delta_{ji} + \sum_{\gamma k} \bar{\delta}_{\beta\gamma} U_\gamma^{jk} G_0 t_k G_0 \mathcal{U}_{\gamma\alpha}^{ki}$$

i, j, k : pairs (\equiv three-cluster (2+1+1) partitions)

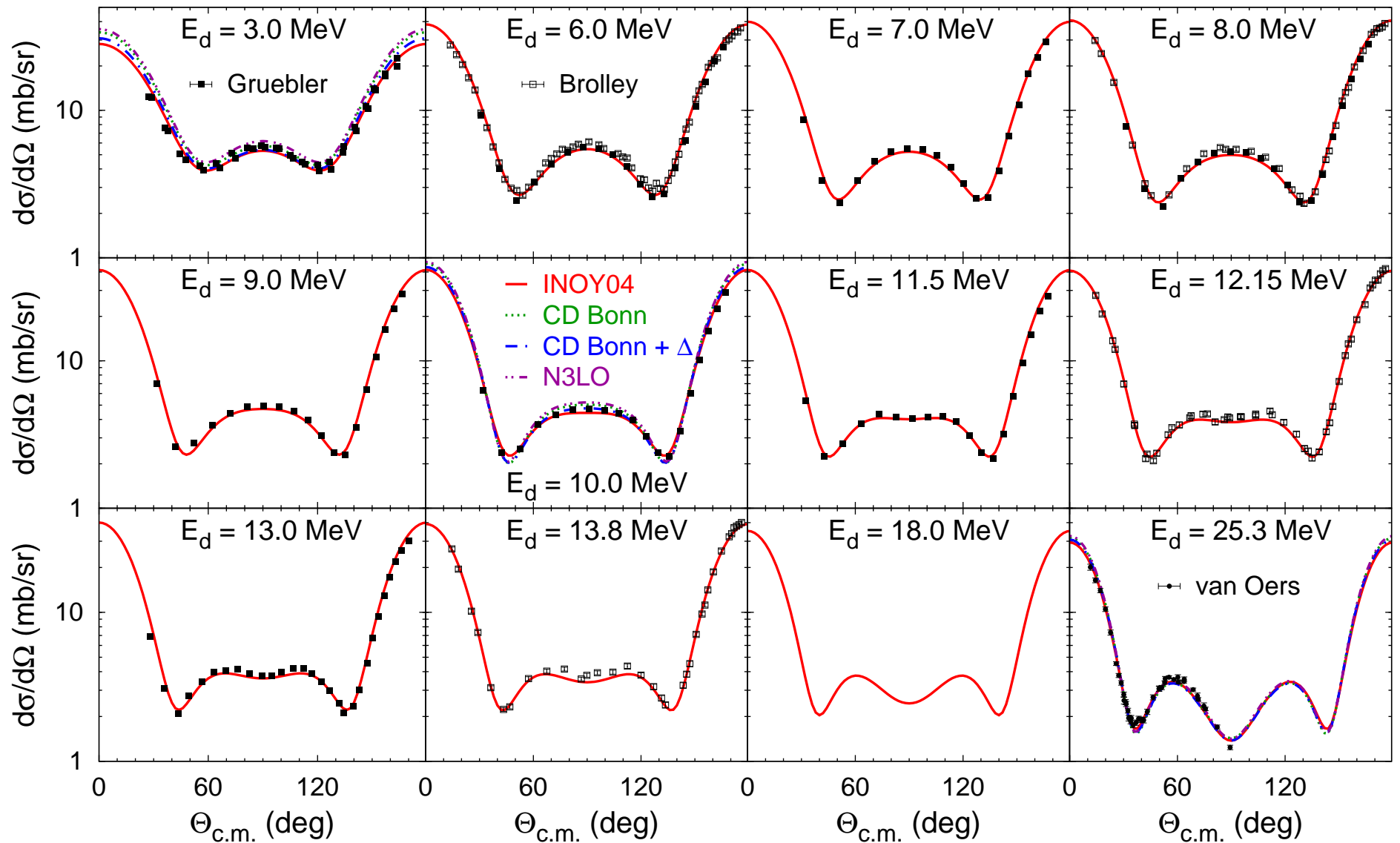
α, β, γ : two-cluster (1+3 or 2+2) partitions

wave function

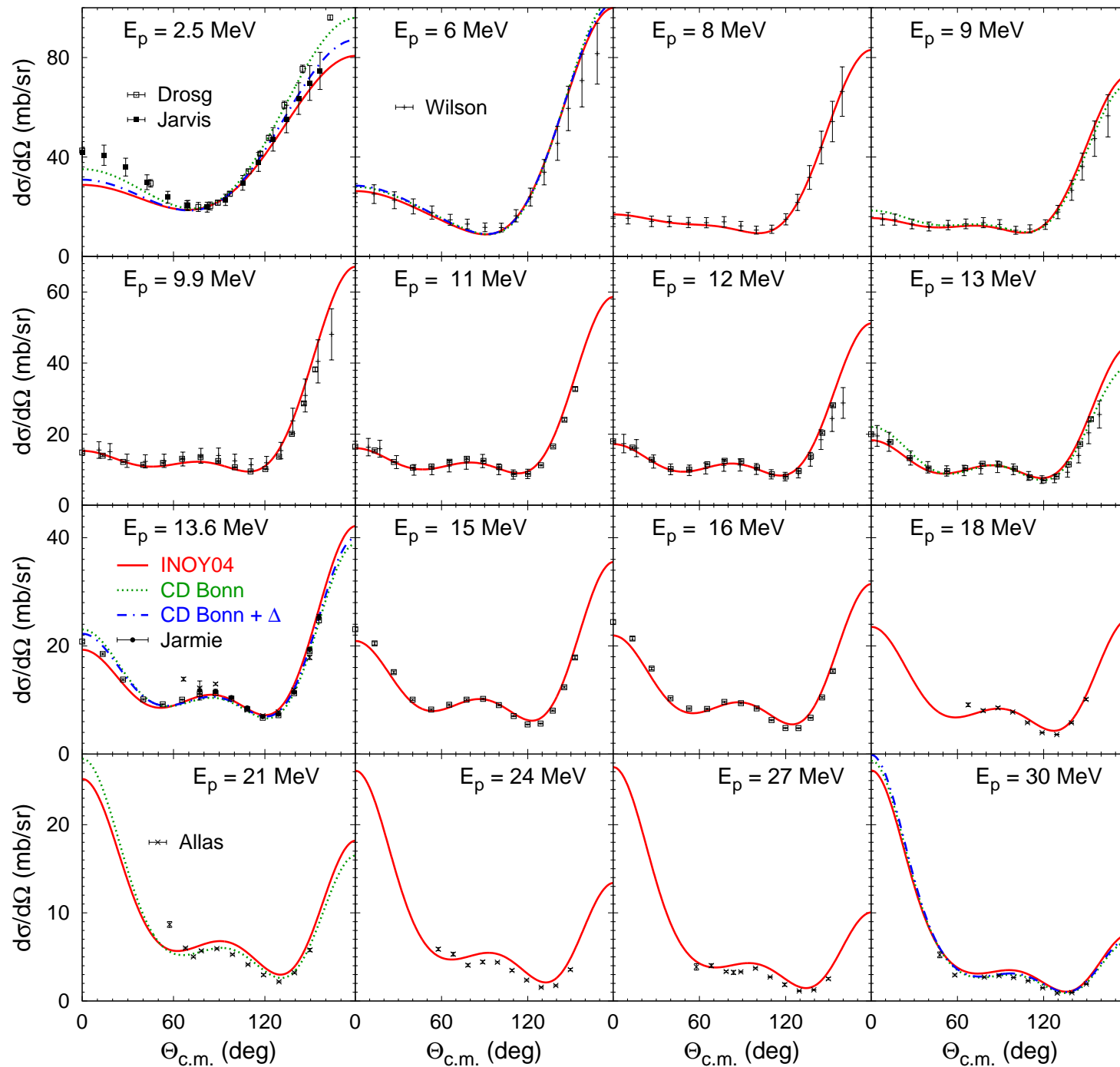
$$|\Psi_\alpha\rangle = |\Phi_\alpha\rangle + \sum_{\gamma j k i} G_0 t_j G_0 U_\gamma^{jk} G_0 t_k G_0 \mathcal{U}_{\gamma\alpha}^{ki} |\Phi_\alpha^i\rangle$$

$$|\Phi_\alpha\rangle = \sum_i |\phi_\alpha^i\rangle, \quad |\phi_\alpha^i\rangle = G_0 \sum_j \bar{\delta}_{ij} t_j |\phi_\alpha^j\rangle$$

Example: ${}^2\text{H}(d, p){}^3\text{H}$

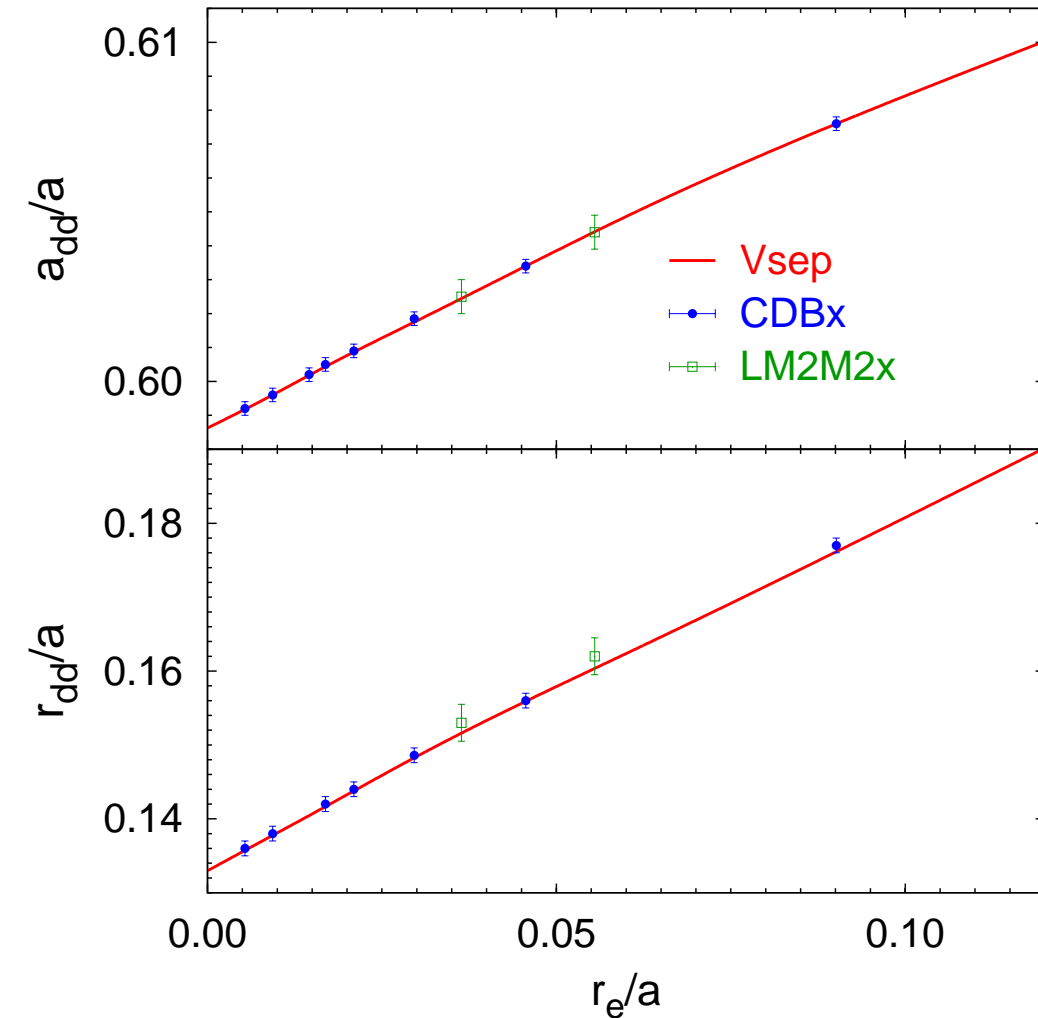


Example: ${}^3\text{H}(p,n){}^3\text{He}$



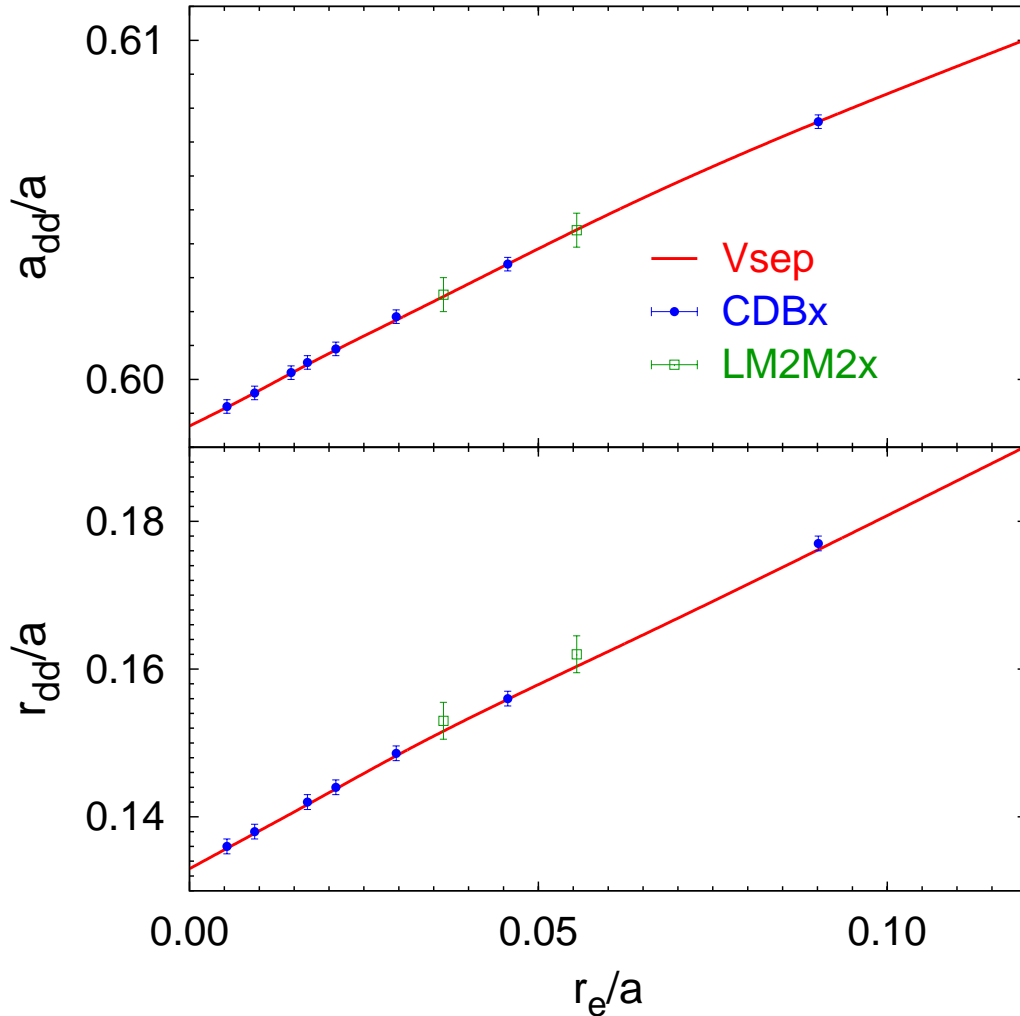
Fermionic dimer-dimer scattering: range effects

$$\text{ERE : } ap_{dd} \cot \delta_0 \approx -\frac{a}{a_{dd}} + \frac{1}{2} \frac{r_{dd}}{a} (ap_{dd})^2 - \frac{1}{4} c_{dd} (ap_{dd})^4$$



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$$\frac{a_{dd}}{a} = 0.5986 + 0.105 \frac{r_e}{a}$$

$$\frac{r_{dd}}{a} = 0.133 + 0.51 \frac{r_e}{a}$$

$$c_{dd} = 0.026 - 0.1 \frac{r_e}{a}$$

Greene et al - **OK**

Elhatisari et al - **wrong**

Conclusion

- L- and energy-dependent CX effects in transfer reactions, no simple relation to SF