

# Effect of hadronic interactions on Lambda polarization

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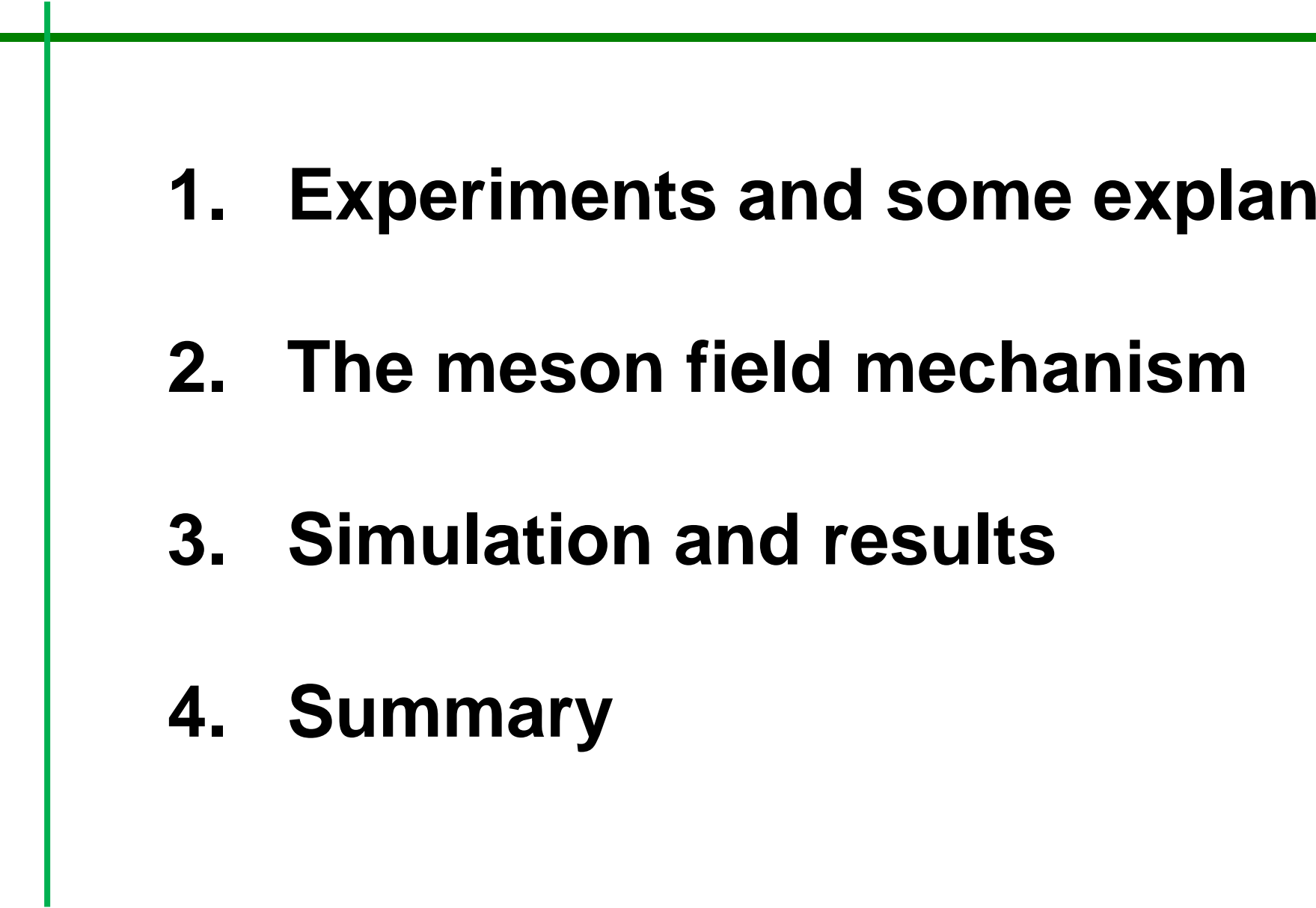
**Physics Department  
China University of Geosciences (Wuhan)**



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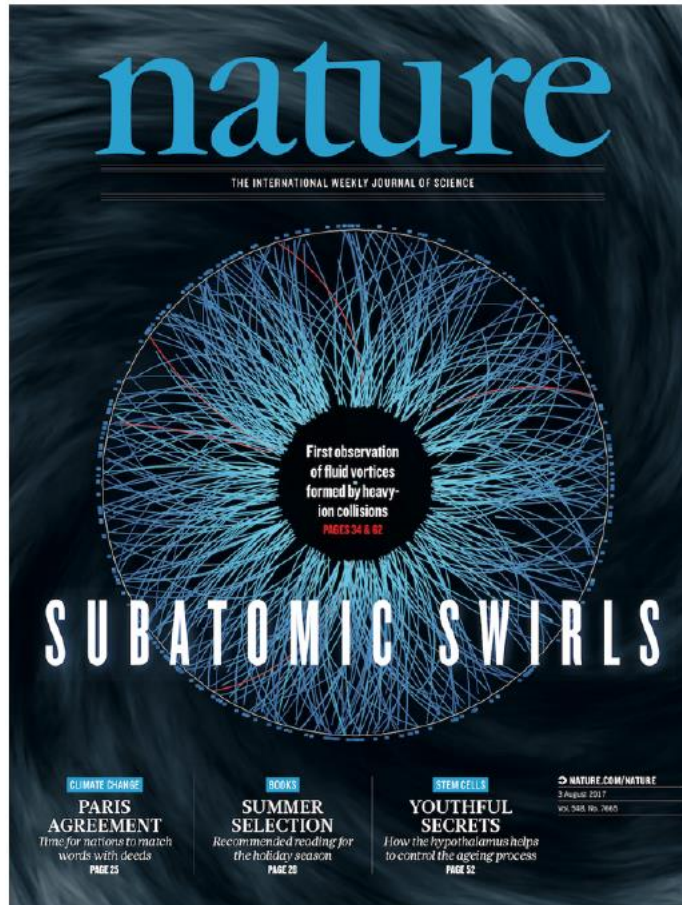
**L. P. Csernai, D.J. Wang, et. al.**

ECT\* Workshop on Spin and Hydrodynamics in Relativistic Nuclear Collisions

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- 1. Experiments and some explanations**
  - 2. The meson field mechanism**
  - 3. Simulation and results**
  - 4. Summary**

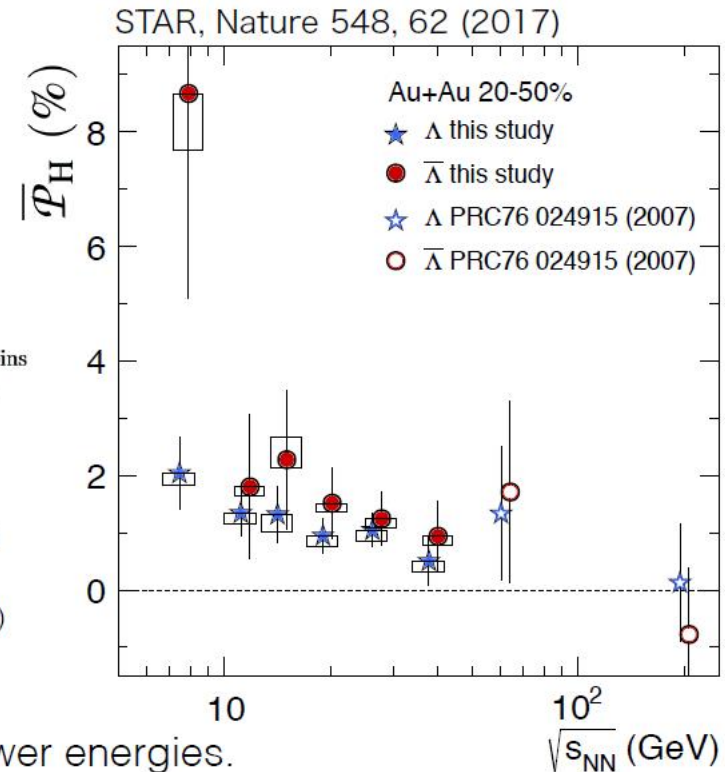
# Experiments: RHIC Au+Au collisions

- Global polarization measurements by RHIC BES II programme: Au-Au collisions at energies of 7.7, 11.5, 14.5, 19.6, 27.0, 39.0, 62.4, and 200 GeV



$$\omega = (P_{\Lambda} + P_{\bar{\Lambda}})k_B T / \hbar$$

$$\sim 10^{22} \text{ s}^{-1} \quad (T=160 \text{ MeV})$$

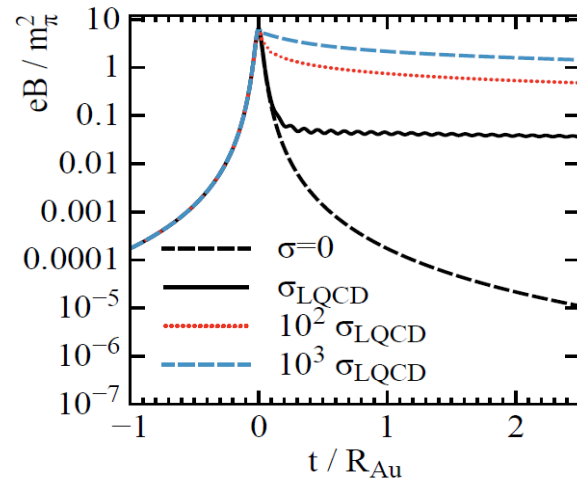


Positive signal at lower energies.  
The most vortical fluid ever observed!

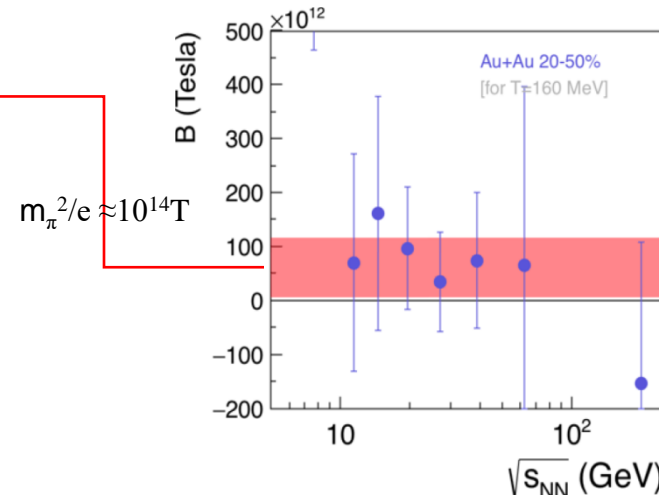
# Earlier explanations to Polarization Splitting

➤ Why is the anti- $\Lambda$ 's polarization larger than  $\Lambda$  polarization?

1. Polarization induced by magnetic field might split the vorticity induced polarization?

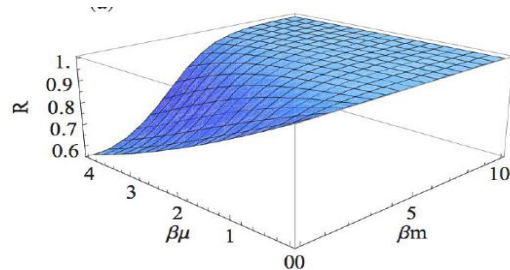


[ L. McLerran, V. Skokov, Nucl.Phys. A929 (2014) 184-190 ]



[ M.A. Lisa, invited talk in Zimany School, Amsterdam 2017 ]

2. Effect of baryon chemical potential : accounts for only 1%



$$\beta = \frac{1}{T}$$

$$\left. \begin{aligned} \beta\mu &\approx \frac{150 \text{ MeV}}{150 \text{ MeV}} = 1 \\ \beta m &\approx \frac{1100 \text{ MeV}}{150 \text{ MeV}} \approx 7 \end{aligned} \right\} 1 - \frac{1}{R} < 1\%$$

$$n_F(x, p) = \frac{1}{e^{\beta(p^\mu u_\mu \mp \mu)} + 1}$$

R.H. Fang, et al., PRC 94 (2016) 024904;  
M.A. Lisa, invited talk in WPCF, Budapest 2017

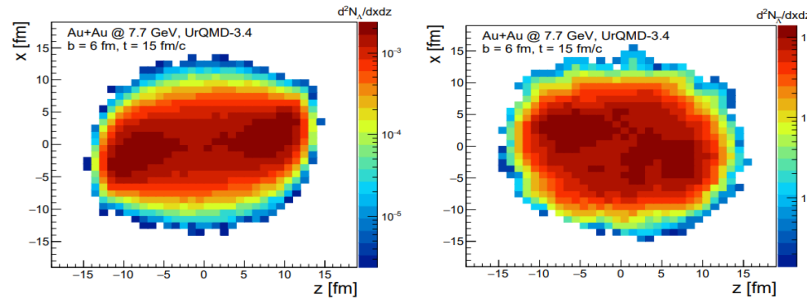
3. Axial Anomaly Charge: The same for  $\Lambda$ s and anti- $\Lambda$ s. But  $N_\Lambda > N_{\bar{\Lambda}}$

A. Sorin, O. Teryaev, PRC 94 011902(R) (2017)

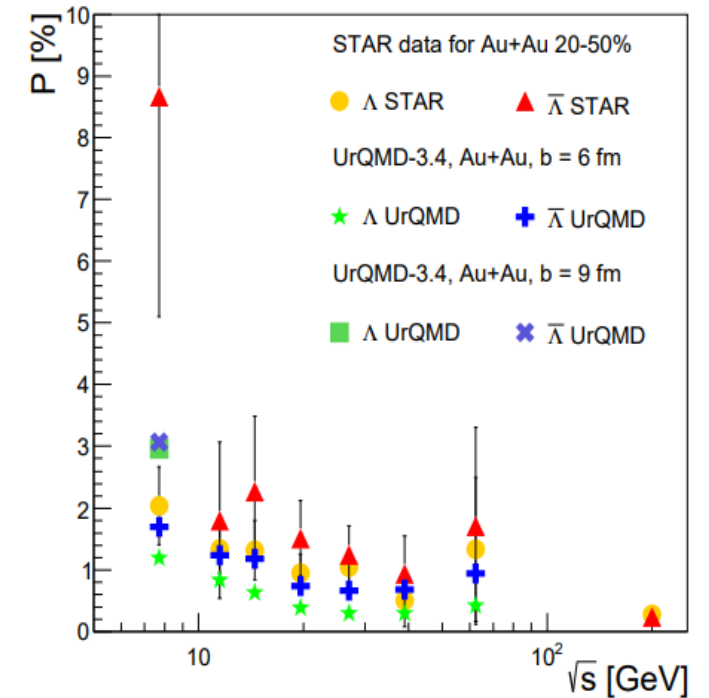


# Recent explanations

Different Space-time freeze-out of  $\Lambda$  and anti- $\Lambda$ , results into polarization difference?



$\sqrt{s}$ (GeV)	7.7	11.5	14.5	19.6
$\langle t_{\Lambda}^{FO} \rangle$ (fm/c)	21.3009	21.9568	23.066	24.3462
$\langle t_{\bar{\Lambda}}^{FO} \rangle$ (fm/c)	19.7806	21.0302	21.959	23.1288

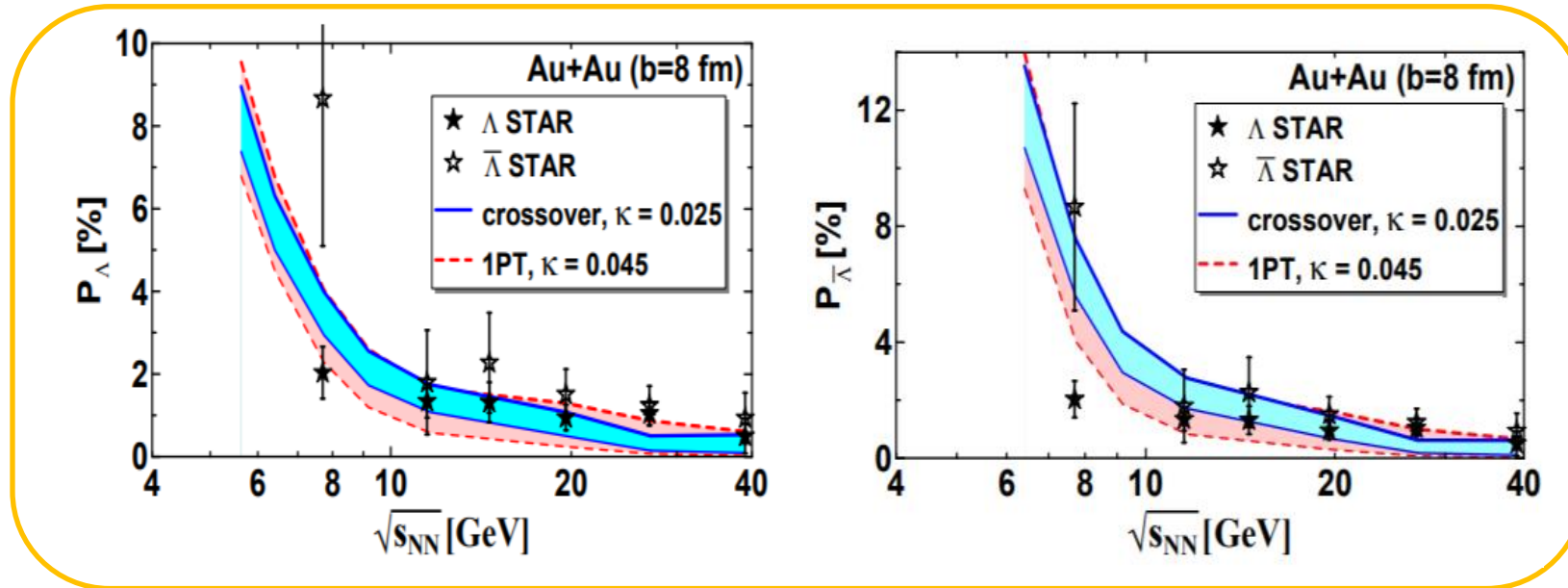


[O. Vitiuk, L.V. Bravina, E. E. Zabrodin, PLB 803, 135298 (2020).]



# Recent explanations

Axial Anomaly Charge: The same for  $\Lambda$ s and anti- $\Lambda$ s. But  $N_\Lambda > N_{\bar{\Lambda}}$



$$P_\Lambda = \int d^3x (J_{5s}^0/u_y)/(N_\Lambda + N_{\bar{K}^*}),$$

$$P_{\bar{\Lambda}} = \int d^3x (J_{5s}^0/u_y)/(N_{\bar{\Lambda}} + N_{K^*}),$$

[Yu. B. Ivanov, arXiv:2006.14328, accepted by PRC]



# Meson field in rotating system

Considering the system during hadron rescattering, the strong interactions between  $\Lambda$  and baryons are mediated by the scalar meson  $\sigma$  and vector meson  $V^\mu$

$$\mathcal{L}_{\text{eff}} = \sum_i \bar{\psi}_i (i \not{\partial} - m_i + g_{\sigma i} \sigma - g_{V i} \not{V}) \psi_i + \frac{1}{2} (\partial_\mu \sigma \partial^\mu \sigma - m_\sigma^2 \sigma^2) - \frac{1}{4} V^{\mu\nu} V_{\mu\nu} + \frac{1}{2} m_V^2 V_\mu V^\mu.$$



$$H_{\text{spin}}^V = -\frac{g_{VH}}{M_H} \beta \boxed{S \cdot B_V} - i \frac{g_{VH}}{4M_H^2} S \cdot \nabla \times E_V - \frac{g_{VH}}{2M_H^2} S \cdot E_V \times p$$

Dirac  $4 \times 4$  matrix:  $\beta = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$$\boxed{B_V = \frac{\bar{g}_V}{m_V^2} (\nabla \times J_B)}$$

----- Foldy-Wouthuysen Hamiltonian

[L.P. Csernai, J. Kapusta, et al, PRC 99, 021901(R) (2019)]



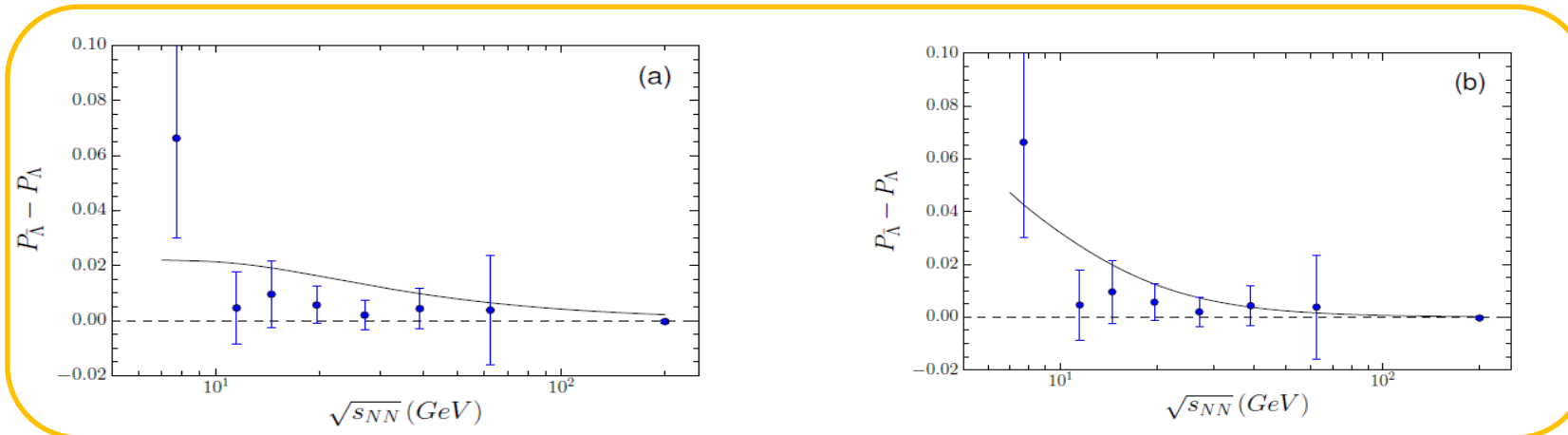


# Meson field in rotating system

In conclusion, the baryons (fermions) with vorticity will induce vector meson (bosons) magnetic field, this magnetic&electric field would interact with the  $\Lambda$ 's spin, resulting into the polarization and polarization difference between  $\Lambda$  and anti-  $\Lambda$ .

$$P_y = \beta \frac{g_{V\Lambda}}{m_\Lambda} \frac{|\mathbf{B}_V|}{2T} = g_{V\Lambda} \bar{g}_V \frac{n_B(t_{\text{ch}})}{m_\Lambda m_V^2} \frac{\Delta c \beta}{2t_{\text{ch}} T(t_{\text{ch}})}.$$

$$P_{\bar{\Lambda}} - P_{\Lambda} = C \left( \frac{n_B(t_{\text{ch}})}{0.15/\text{fm}^3} \right) \left( \frac{140 \text{ MeV}}{T(t_{\text{ch}})} \right).$$



[L.P. Csernai, J. Kapusta, et al, PRC 99, 021901(R) (2019)]





## Meson field in rotating system

We modify the polarization splitting formula therein, by removing the free parameter  $C$  and explicitly bringing out the vorticity, which is essential in polarization study.

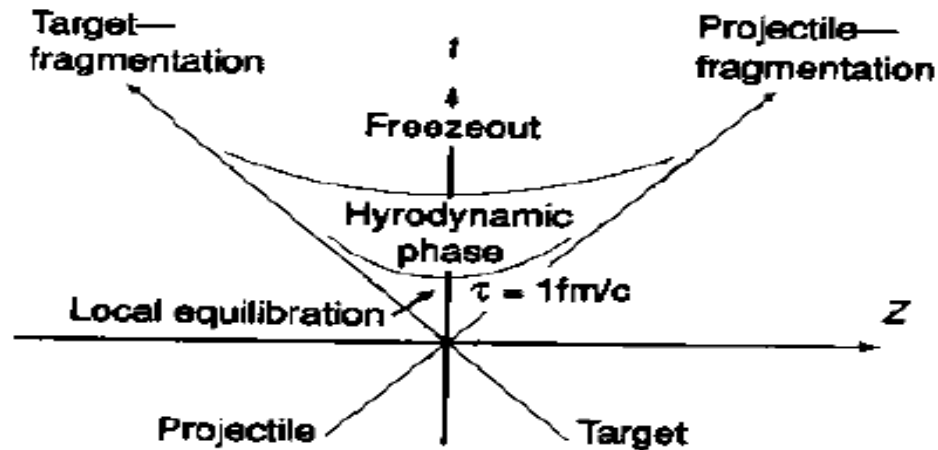
$$P = 2S = \tanh\left(\frac{\Omega}{2T}\right) \hat{\Omega} \simeq \frac{\Omega}{2T} = -\beta \frac{g_{VH}}{M_H} \frac{B_V}{T}$$

$$\begin{aligned} \Delta P_J &= \left\langle C \frac{\nabla \times \mathbf{J}_B}{T} \right\rangle = C \left\langle \frac{\rho_B \boldsymbol{\omega}}{T} \right\rangle + C \left\langle \frac{\nabla \rho_B \times \mathbf{v}}{T} \right\rangle \\ &= \Delta P_\omega + \Delta P_\rho. \end{aligned}$$

where  $C = 2(g_{VH} \bar{g}_V)/(M_H m_V^2)$  is a coefficient determined by the strong coupling constants, hyperon mass and meson mass.

Assumption: the post-freeze-out system is near Boltzmann limit and  $\Lambda$  particles are non-relativistic

# Hydrodynamic Simulation

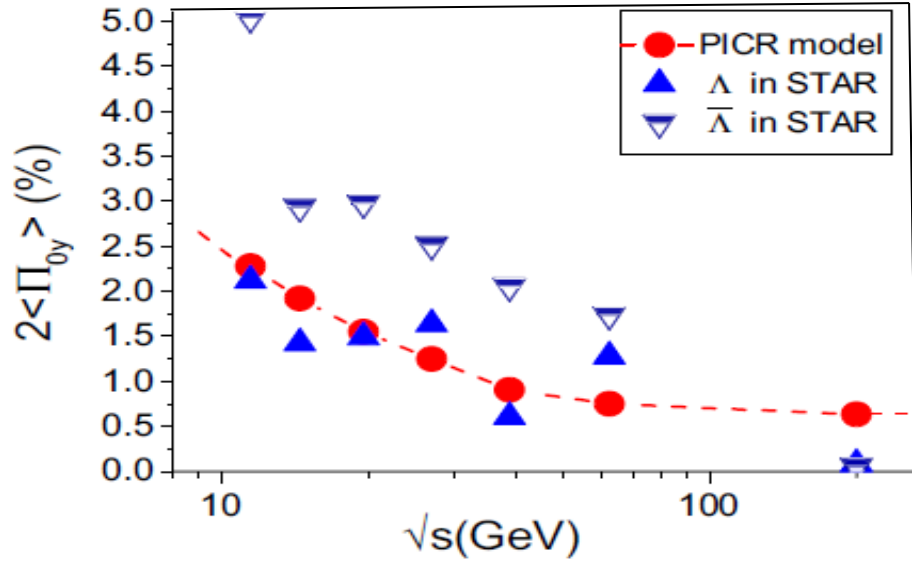


Space-time evolution in Bjorken model

- ❑ An initial state based on Yang-Mills fields (flux tube) is formed after Lorentz contracted nuclei penetrate each other, then system evolves with a (3+1)D fluid dynamics: the high resolution Particle-In-Cell Relativistic (PICR) hydrodynamic model.
- ❑ The major part of freeze out hypersurface is assumed to be **time like** here. We use the ideal-gas post-FO distribution. The precise FO prescription in Hydro fluid dynamics was discussed in Ref. [Yu Cheng, 2010].

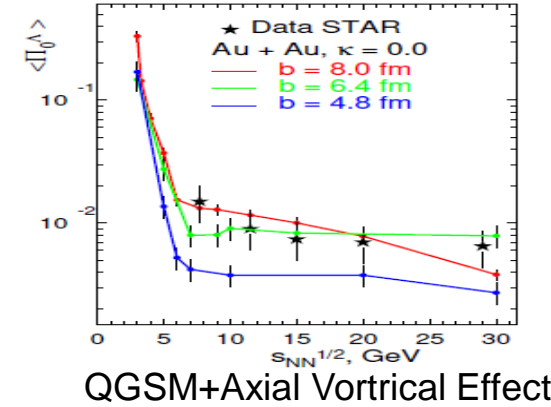
# Previous simulation

[ Y.L. Xie et al., PRC 93, 031901(R) (2017)]

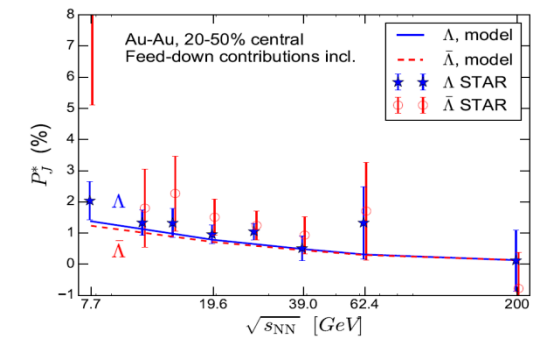
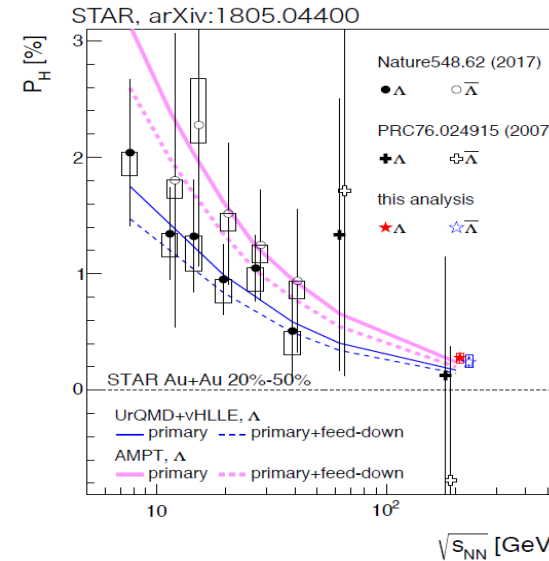
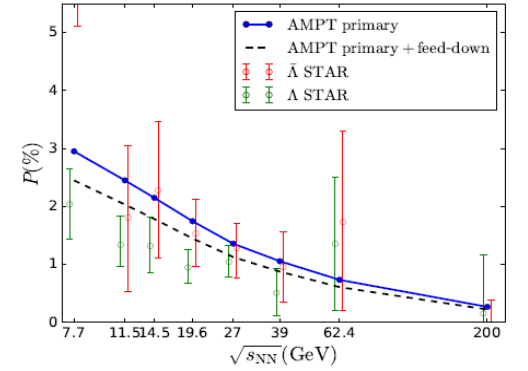


The global polarization,  $2\langle\Pi_{0y}\rangle$ , in our PICR model (red circle) and STAR BES experiments (green triangle), at energies  $\sqrt{s}$  of 11.5, 14.5, 19.6, 27.0, 39.0, 62.4, and 200 GeV.

$$\langle\Pi_{0y}\rangle_p = \frac{\int dp dx \Pi_{0y}(p, x) n_F(x, p)}{\int dp dx n_F(x, p)} = \frac{\int dp \Pi_{0y}(p) n_F(p)}{\int dp n_F(p)}$$



M. Baznat, et al., arXiv:1701.00923. H. Li, X.N. Wang et al., PRC 96,054908 (2017).



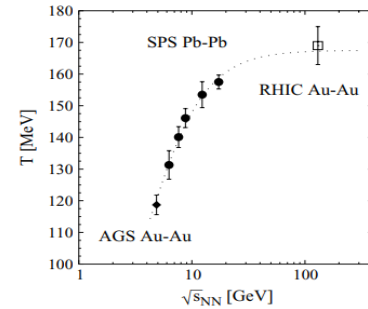
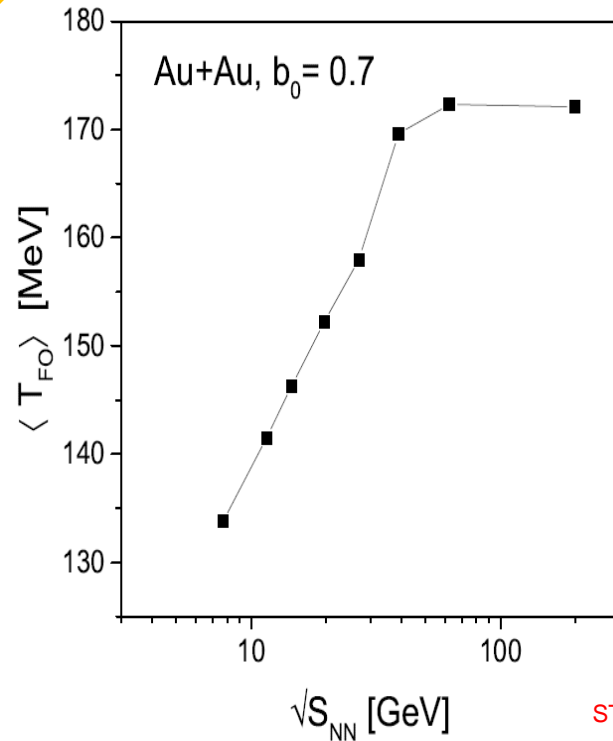
UrQMD-vHLLLE hybrid Model

I. Karpenko, F. Becattini, et al., Eur. Phys. J. C 77, 213 (2017).

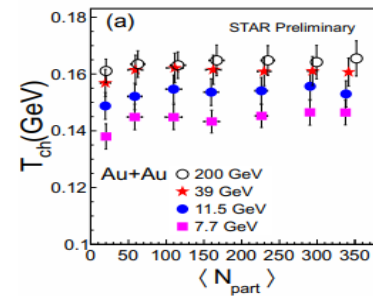
We use the same simulation data, but vary the freeze-out time for different collision energies.



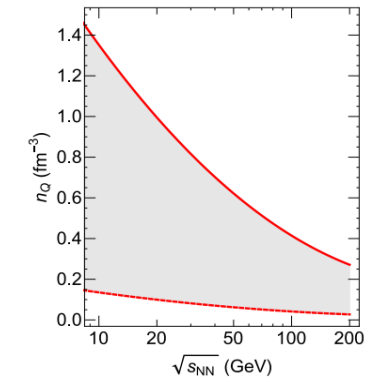
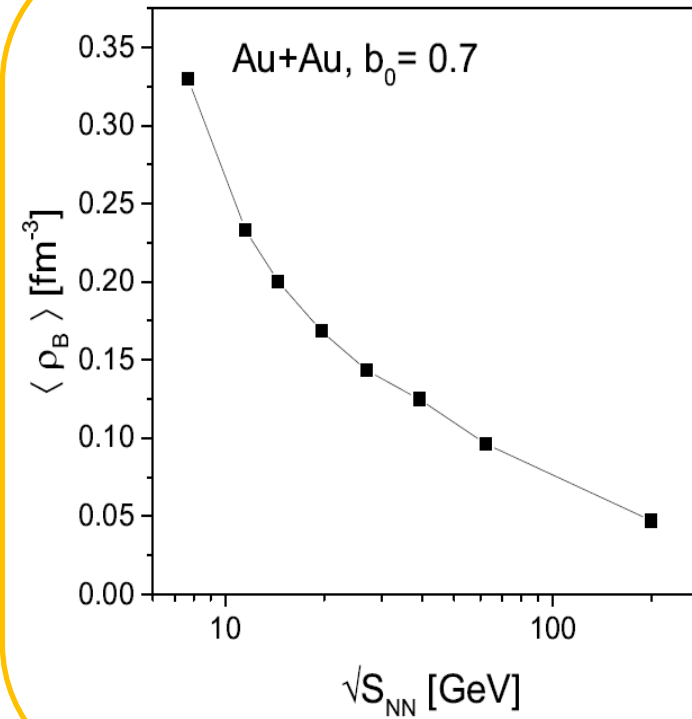
# Simulation and results



F. Becattini et al., PRC 73, 044905 (2006)



STAR Collaboration, NPA 904-905, 891c (2013).

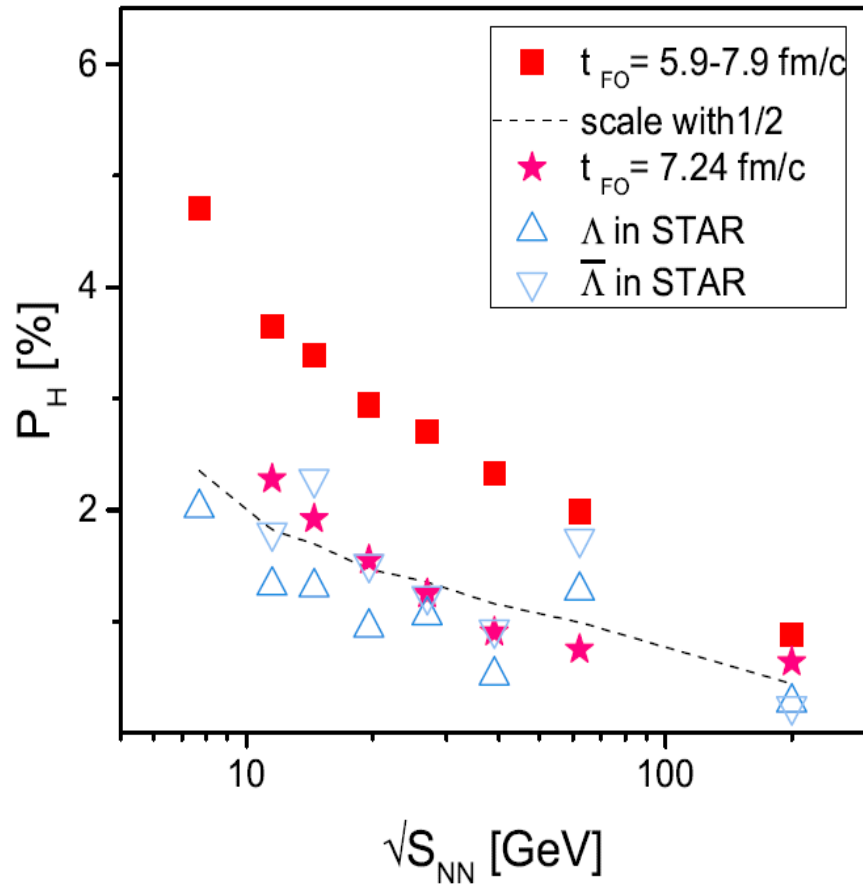


X. Y. Guo et al., arXiv:1904.04704.

The freeze-out time is varied from 5.9 -7.9 fm/c for  $\sqrt{s} = 7.7$ -200GeV, so that the freeze-out  $T$  and  $\rho$  are consistent with theoretical expectations and experiments.



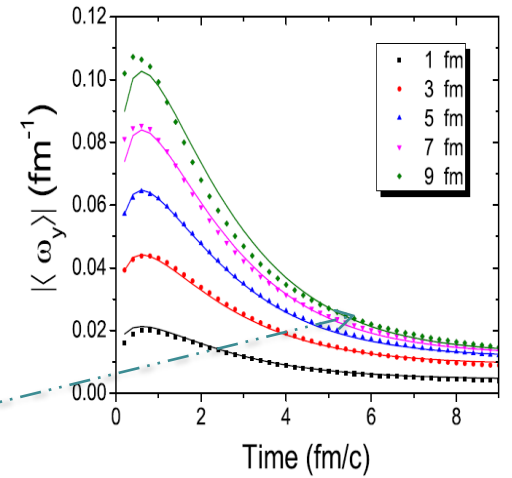
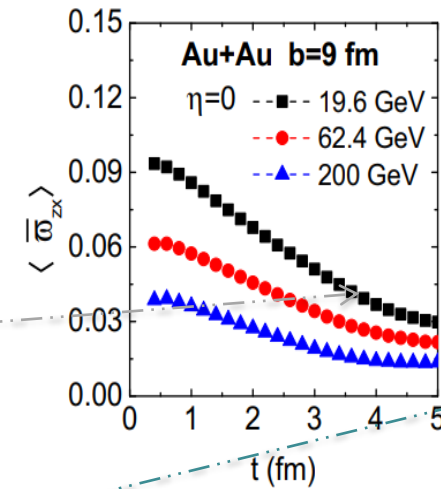
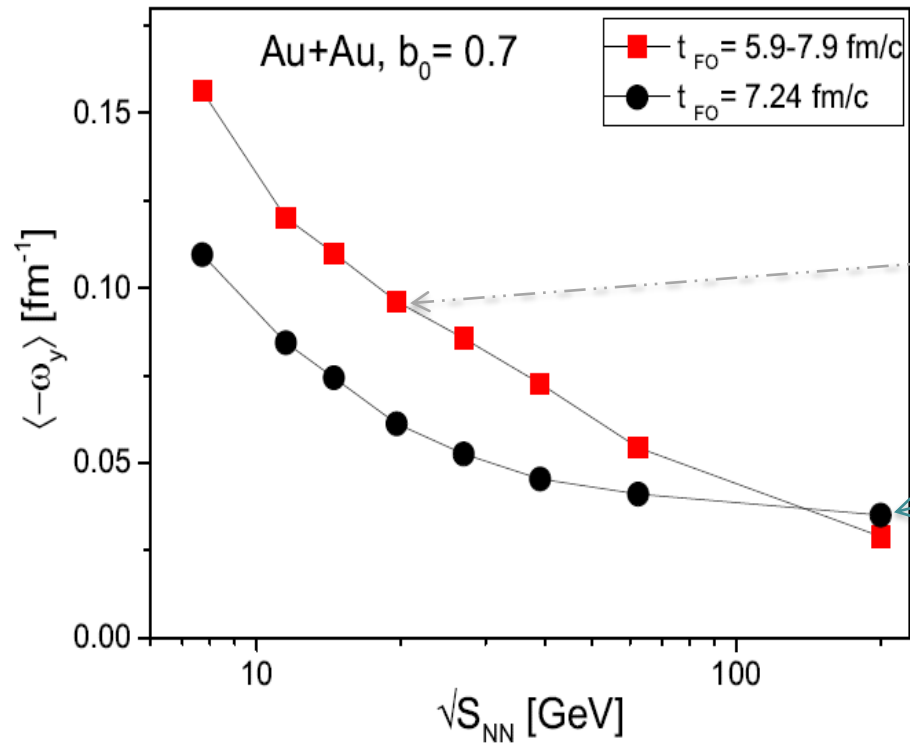
# Simulation and results



1. New values of polarization are larger than old ones, showing the sensitivity to freeze-out time, while the energy dependence behavior is still kept.
2. Estimates of the global polarization at  $c=20-50\%$  by scaling of 0.5, is very close to the experimental results.



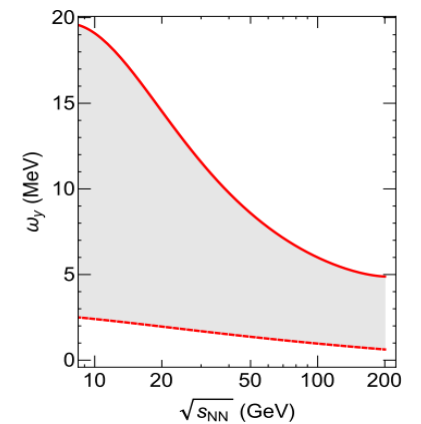
# Simulation and results



Y. Jiang et al., Phys. Rev. C 94, 044910 (2016).

D.-X. Wei et al., PRC 99, 014905 (2019).

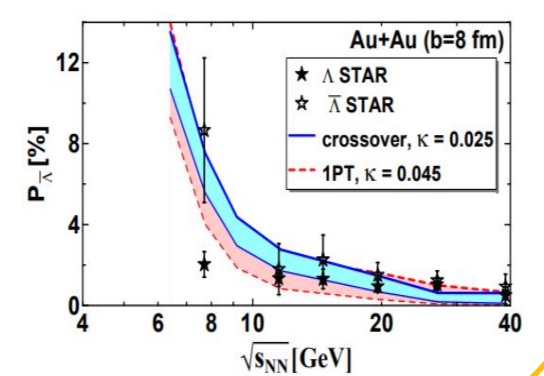
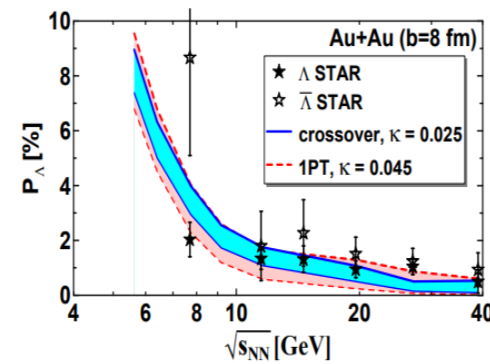
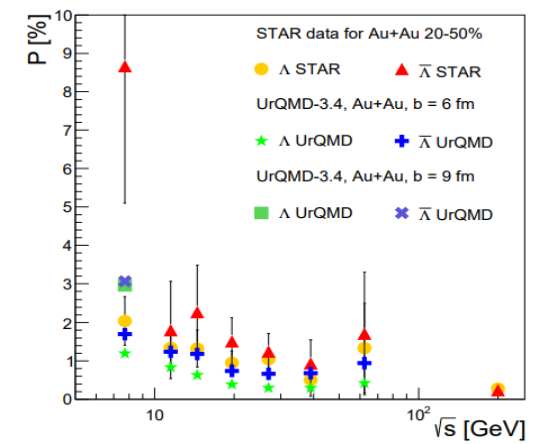
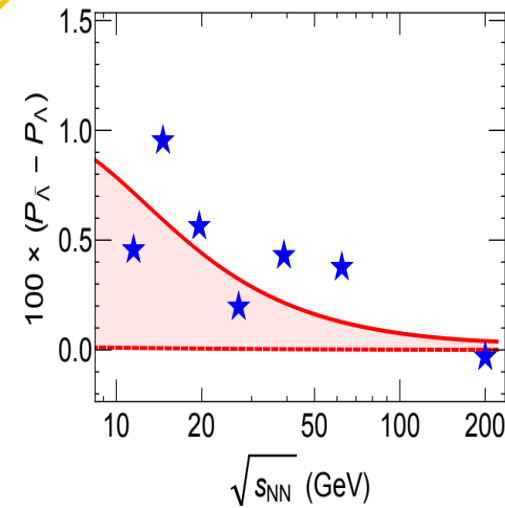
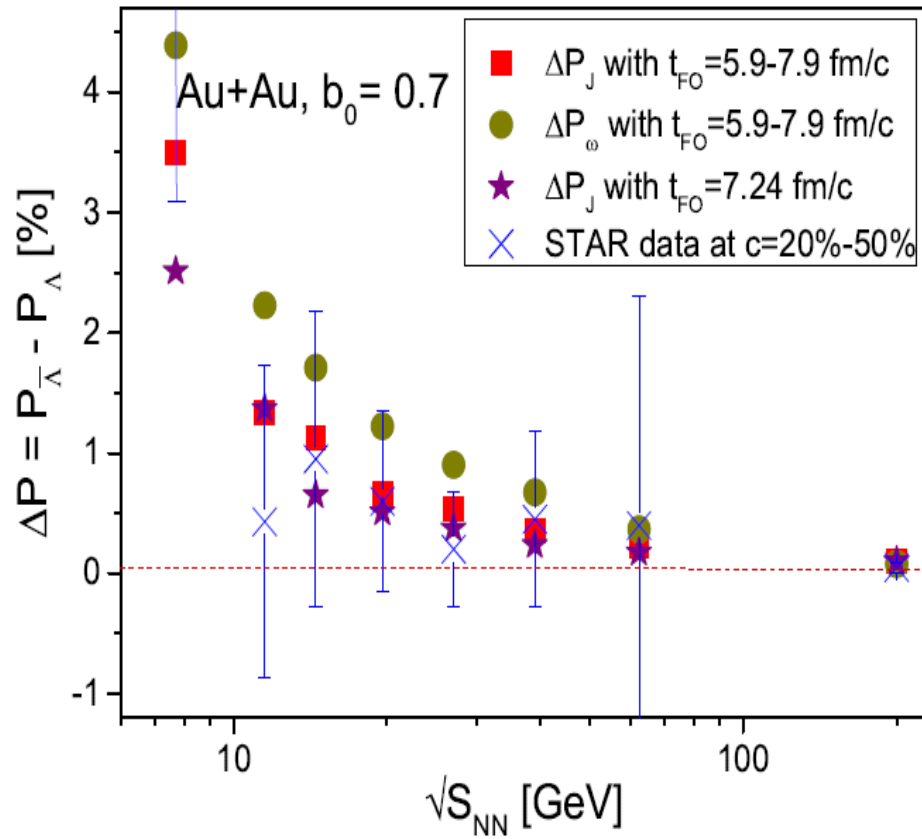
PICR model has similar magnitude of vorticity as the AMPT model



X. Y. Guo et al., arXiv:1904.04704.



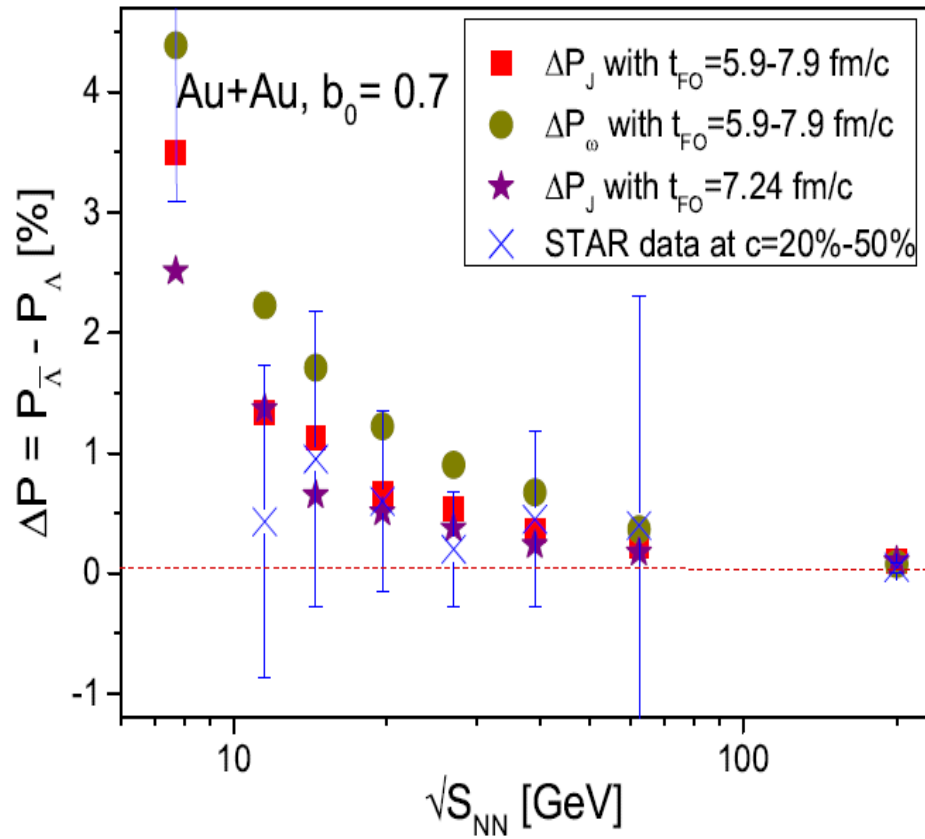
# Simulation and results







# Simulation and results



1. The polarization splitting based on meson field mechanism is larger than many other approaches.
2.  $\Delta P_J = \Delta P_\omega + \Delta P_\rho < \Delta P_\omega$ , and it means  $\Delta P_\rho$  is negative, so it decreases the final splitting effect by about 1/3~1/4.



# Simulation and results

$b_0$ (c)	0.45 (20%)	0.6 (36%)	0.7 (49%)
$\langle T_{FO} \rangle$ (MeV)	134.3	134.8	133.8
$t_{FO}$ (fm/c)	4.2 <	5.1 <	5.9
$\langle \rho_B \rangle$ (fm <sup>-3</sup> )	0.36	0.345	0.33
$\langle -\omega_y \rangle$ (fm <sup>-1</sup> )	0.140	0.163	0.156
$\Delta P_\omega$	4.49% $\approx$	4.77% $\approx$	4.39%
$\Delta P_J$	4.28% >	4.12% >	3.49%

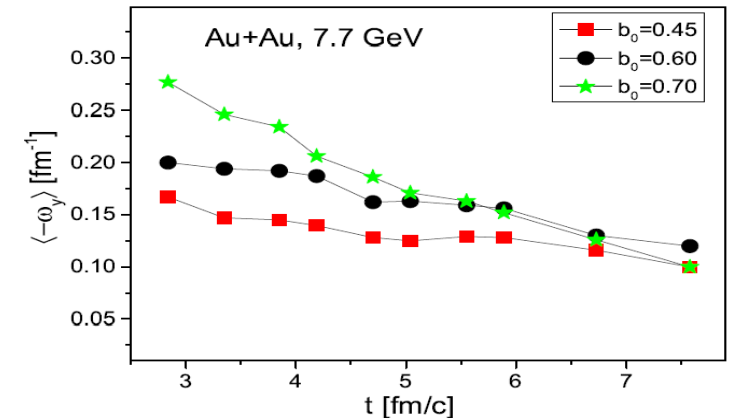
**Table 1.** The average freeze-out temperature  $T_{FO}$ , freeze-out time  $t_{FO}$ , average baryon density  $\langle \rho_B \rangle$ , average vorticity  $\langle -\omega_y \rangle$ , and  $\Delta P_\omega$ ,  $\Delta P_J$  defined in eq. (15), for Au+Au 7.7 GeV collisions at different centralities  $c = 20\%$ ,  $36\%$ ,  $49\%$ .

$b_0$ (c)	0.45 (20%)	0.7 (49%)
$\langle T_{FO} \rangle$ (MeV)	142.2	141.5
$t_{FO}$ (fm/c)	7.9 >	5.9
$\langle -\omega_y \rangle$ (fm <sup>-1</sup> )	0.140	0.156
$\Delta P_\omega$	1.43% <	2.23%
$\Delta P_J$	0.83% <	1.33%

**Table 2.** The average freeze-out temperature  $T_{FO}$ , freeze-out time  $t_{FO}$ , average vorticity  $\langle -\omega_y \rangle$ , and  $\Delta P_\omega$ ,  $\Delta P_J$  defined in eq. (15), for Au+Au 11.5 GeV collisions at different centralities  $c = 20\%$ ,  $49\%$ .

Things are different for 7.7 GeV case:

1. Freeze-out time is larger in peripheral collisions and the decreasing tendency of vorticity vs time is very mild. Thus similar  $\Delta P_\omega$
2. The larger fluctuations of baryon density in peripheral collisions lead to larger  $|\Delta P_\rho|$ . Thus  $\Delta P_J = \Delta P_\omega + \Delta P_\rho$  is smaller for peripheral collisions.





# Summary

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We have presented the polarization splitting based on the meson field mechanism.

Thank  
You



$$\Pi_0(p) = \Pi(p) - \frac{p}{p^0(p^0 + m)} \Pi(p) \cdot p .$$

