

**Open Quantum Systems:
From atomic nuclei to ultracold atoms and quantum optics**

Trento, *September 30—October 4, 2019*

Nuclear astrophysics experiments with Trojan Horse Method

Giovanni Luca Guardo

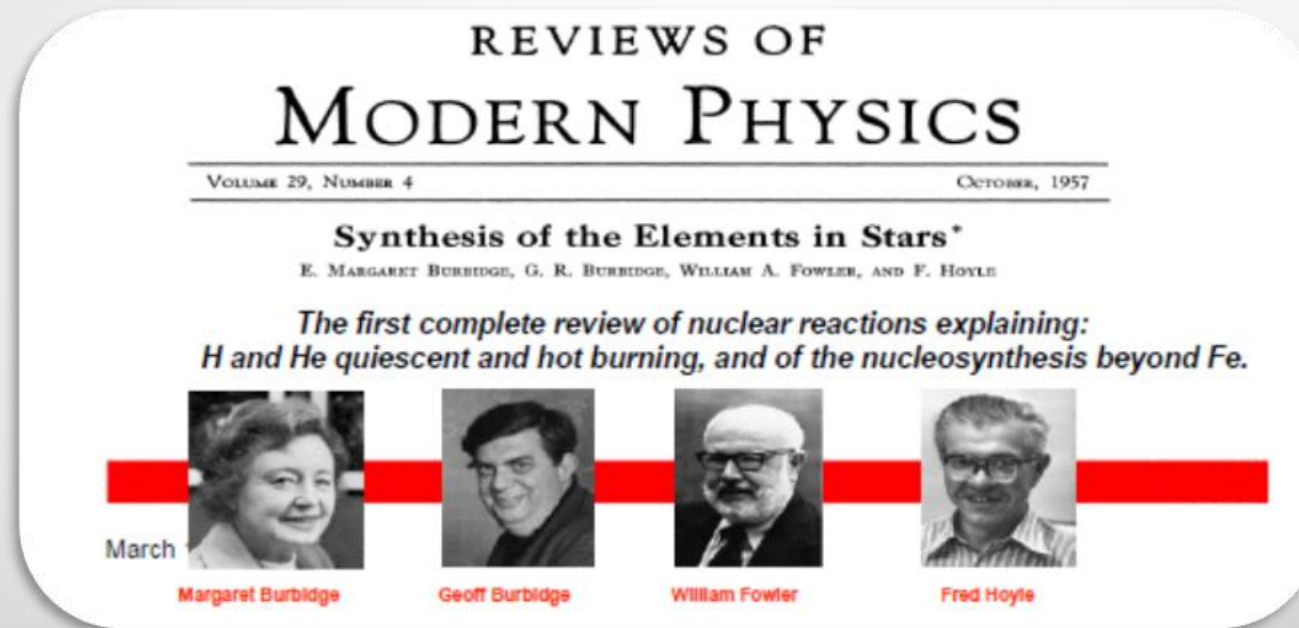
on behalf of the AsFiN group



Experimental Nuclear Astrophysics

... Everything starts from the **B²FH** review paper of 1957,
the basis of the modern nuclear astrophysics

this work has been considered as the greatest gift of astrophysics to modern civilization

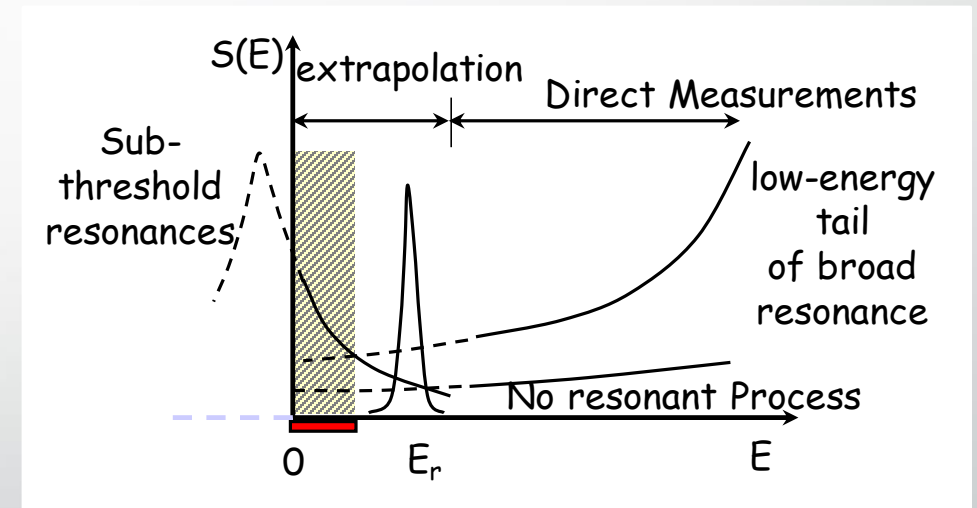
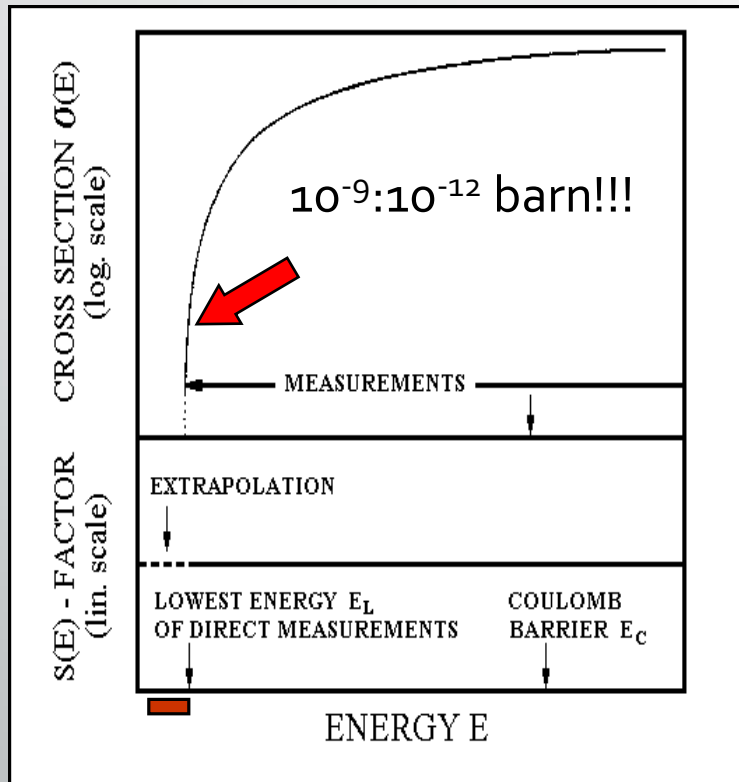
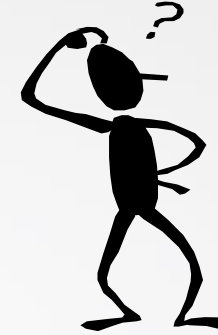


The elements composing everything from planets to life were forged inside earlier generations of stars!

Nuclear reactions responsible for both ENERGY PRODUCTION and SYNYHESIS OF ELEMENTS

Direct Measurements

- Very small cross section values reflect in a faint statistic;
- Very low signal-to-noise ratio makes hard the investigation at astrophysical energies;
- Instead of the cross section, the $S(E)$ -factor is introduced

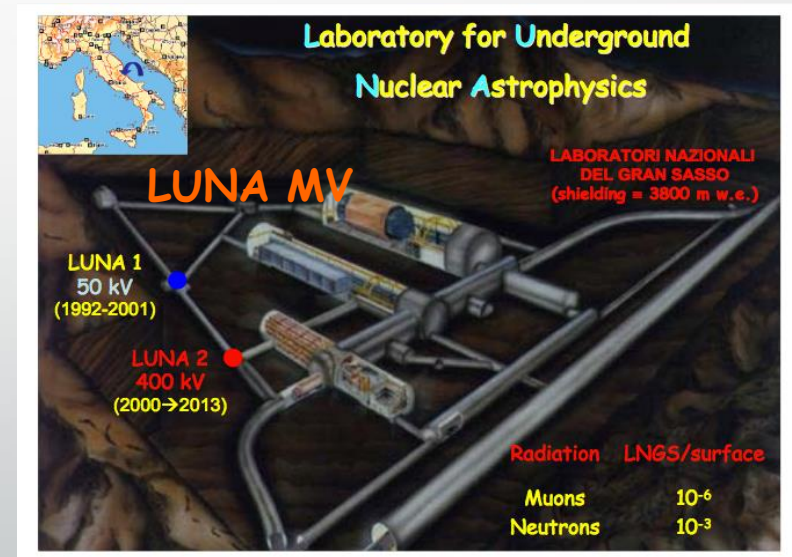


Direct Measurements

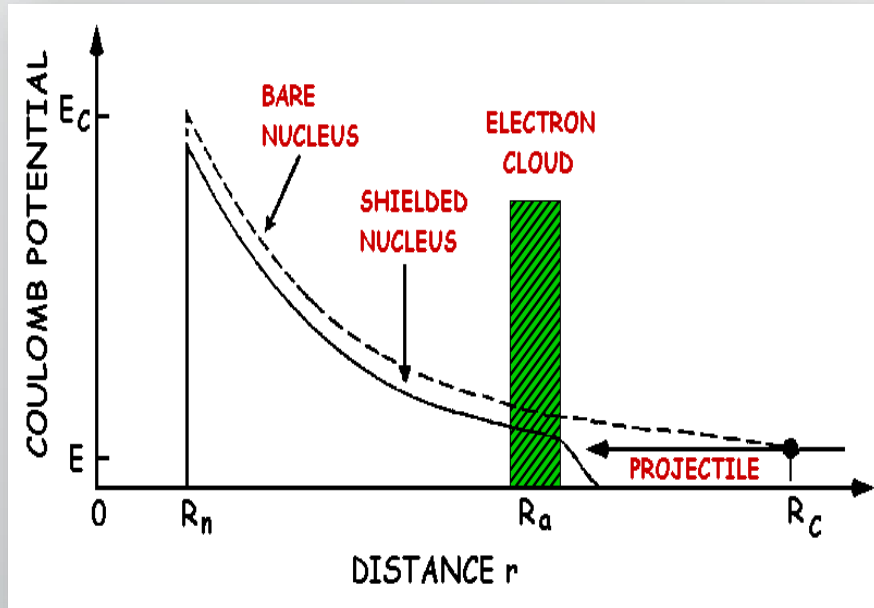


Several efforts have been made in the last years in order to **improve the signal-to-noise ratio** for low-energy cross section measurement.

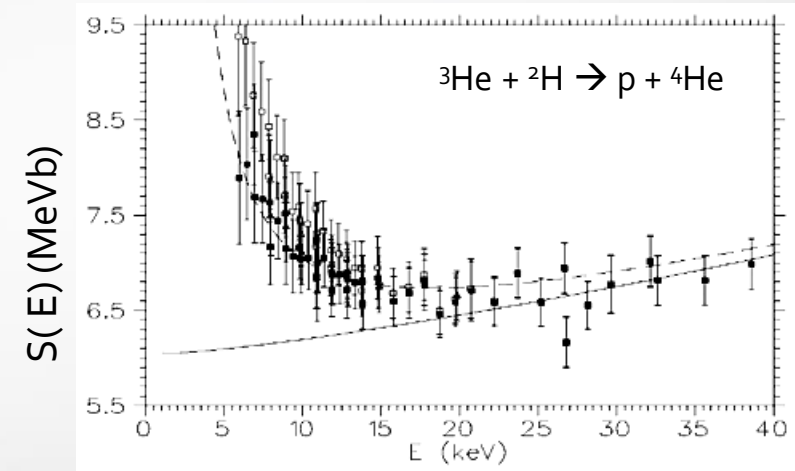
- Longer measurements
- Higher beam currents
- 4π detectors
- Pure targets
- Underground laboratories



Electron Screening



Due to the electron cloud surrounding the interacting ions the projectile feels a reduced barrier



Theory vs. Experiment
 → Far to be understood...
Stellar Plasma

Reaction	U_{ad} (eV)	U_{exp} (eV)	Reference
${}^6\text{Li}(p,\alpha){}^3\text{He}$	186	440 ± 150	[Engstler et al.(1992)]
${}^6\text{Li}(d,\alpha){}^4\text{He}$	186	330 ± 120	[Engstler et al.(1992)]
$\text{H}({}^7\text{Li},\alpha){}^4\text{He}$	186	300 ± 160	[Engstler et al.(1992)]
${}^2\text{H}({}^3\text{He},p){}^4\text{He}$	65	109 ± 9	[Aliotta et al.(2004)]
${}^3\text{He}({}^2\text{H},p){}^4\text{He}$	120	219 ± 7	[Aliotta et al.(2004)]
$\text{H}({}^9\text{Be},\alpha){}^6\text{Li}$	240	900 ± 50	[Zahnaw et al.(1997)]
$\text{H}({}^{11}\text{B},\alpha){}^8\text{Be}$	340	430 ± 80	[Angulo et al. (1993)]

Indirect Methods

❖ Coulomb dissociation

G. Baur et al. Annu. Rev. Nucl. Part. Sci. 46,321,(1996)

to determine the absolute $S(E)$ factor of a radiative capture reaction $A+x \rightarrow B+\gamma$
studying the reversing photodisintegration process $B+\gamma \rightarrow A+x$

❖ Asymptotic Normalization Coefficients (ANC)

A.M. Mukhamedzhanov et al.: PRC 56,1302,(1997)

to determine the $S(0)$ factor of the radiative capture reaction, $A+x \rightarrow B+\gamma$ studying a peripheral transfer reaction into a bound state of the **B** nucleus

❖ Trojan Horse Method (THM)

C. Spitaleri, *Problems of Fundamental Modern Physics, II*, (World Sci.,1991), p. 21.

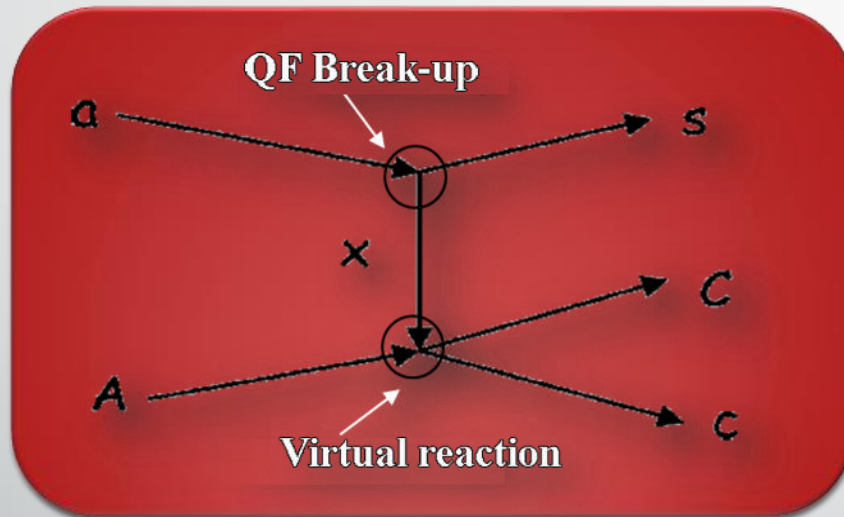
C. Spitaleri et al., Phys. of Atomic Nuclei, 74 (2011) 1725

to determine the $S(E)$ factor of a charged particle reaction $A+x \rightarrow c+C$

The Trojan Horse Method



The idea of the *THM* is to extract the cross section of an astrophysically relevant two-body reaction $A+x \rightarrow c+C$ at low energies from a suitable three-body reaction $a+A \rightarrow c+C+s$



Quasi free kinematics is selected

✓ only $x - A$ interaction

✓ $s = \text{spectator}$ ($p_s \sim 0$)

$$E_A > E_{\text{Coul}} \rightarrow$$

- NO coulomb suppression
- NO electron screening
- NO centrifugal barrier

- *THM Review papers* → [Spitaleri C. et al., PAN, 2011](#)
[Tribble R. et al., Rep. Prog. Phys. 2014](#)
[Spitaleri C. et al., EpJ A, 2019](#)

Theoretical Approach

The TH-nucleus is chosen because of:

- its large amplitude in the $a=x\oplus s$ cluster configuration;
- its relatively low-binding energy;
- Its known x - s momentum distribution $|\Phi(p_s)|^2$ in a .

$$E_{Ax} = \frac{m_x}{m_x + m_A} E_A - B_{xs}$$

B_{x-s} plays a key role in compensating for the beam energy thanks to the x - s *intercluster motion* inside a , it is possible to span an energy range of several hundreds of keV with only one beam energy

In the Plane Wave Impulse Approximation (PWIA) the cross section of the three body reaction can be factored as:

$$\frac{d^3\sigma}{d\Omega_c d\Omega_C dE_c} \propto KF \cdot |\Phi(p_s)|^2 \cdot \frac{d\sigma_{Ax}}{d\Omega}$$

Three body
measured
cross section

Calculated
kinematical
factor

Fourier transform
for the x - s
intercluster
motion

Astrophysically
relevant two
body cross
section

THM Cross Section

Virtual nature of x particle \rightarrow $A+x$ interaction is off-energy shell

Cross section of the bare nucleus but **NO absolute value** \rightarrow
normalization to direct data available at higher energies

Standard R-Matrix approach cannot be applied to extract the resonance parameters \rightarrow Modified R-Matrix is introduced instead

$$\frac{d^2\sigma}{dE_{xA}d\Omega_s} = NF \sum_i (2J_i + 1) \times \left| \frac{\sqrt{k_f(E_{xA})} \sqrt{2P_l_i(k_{cC}R_{cC})} M_i(p_{xA}R_{xA}) \gamma_{cC}^i \gamma_{xA}^i}{\mu_{cC} D_i(E_{xA})} \right|^2$$

where:

- $M_i(p_{xA}R_{xA})$ describes the transfer amplitude for the QF-process;
- γ_{xA} and γ_{cC} represents the reduced partial widths for the resonant excited states that are the same of the direct measurements

La Cognata et al., ApJ, 777, 143, 2013

**Study of the $^{17}\text{O}(n,\alpha)^{14}\text{C}$
reaction: extension of the
Trojan Horse Method to
neutron induced reactions**

Astrophysical Scenario

- **Inhomogeneous Big Bang Nucleosynthesis (IBBN)**

The reaction $^{17}\text{O}(n,\alpha)^{14}\text{C}$ represents one of the main channels for ^{14}C production, a key element for the ^{22}Ne production via $^{14}\text{C}(\alpha,\gamma)^{18}\text{O}(n,\gamma)^{19}\text{O}(\beta)^{19}\text{F}(n,\gamma)^{20}\text{F}(\beta)^{20}\text{Ne}(n,\gamma)^{21}\text{Ne}(n,\gamma)^{22}\text{Ne}$

- **Weak component s-process**

$^{17}\text{O}(n,\alpha)^{14}\text{C}$ and $^{17}\text{O}(\alpha,n)^{20}\text{Ne}$ since they act as a neutron poison and a recycle channel during s-process nucleosynthesis in massive stars ($M > 8M_{\text{SUN}}$)

Astrophysical Scenario

• Inhomogeneous Big Bang Nucleosynthesis (IBBN)

The reaction $^{17}\text{O}(n, \alpha)^{14}\text{C}$ represents the main channel for ^{14}C



Temperature $\rightarrow 0.8 < T_8 < 11$ K
Energy range $\rightarrow \sim 0-100$ keV

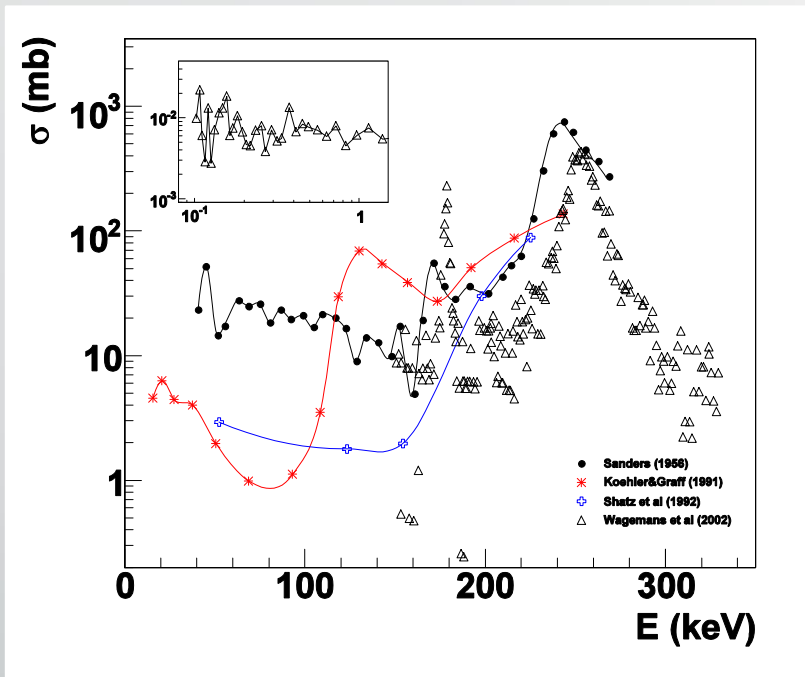


recycle channel

$M > 8 M_{\text{SUN}}$

poison and a
synthesis in massive stars

Status of the Art



• R. M. Sanders, Phys. Rev., 104, 1434 (1956)
 INVERSE REACTION $^{14}\text{C}(\alpha, n)^{17}\text{O}$

* P.E.Koehler & S.M.Graff, Phys. Rev., C44(6),
 2788 (1991)

○ H. Schatz et al., Astroph. J., 413, 750
 (1993)

△ J. Wagemans et al., Phys. Rev., C65(3), 34614
 (2002)

Subthreshold Level

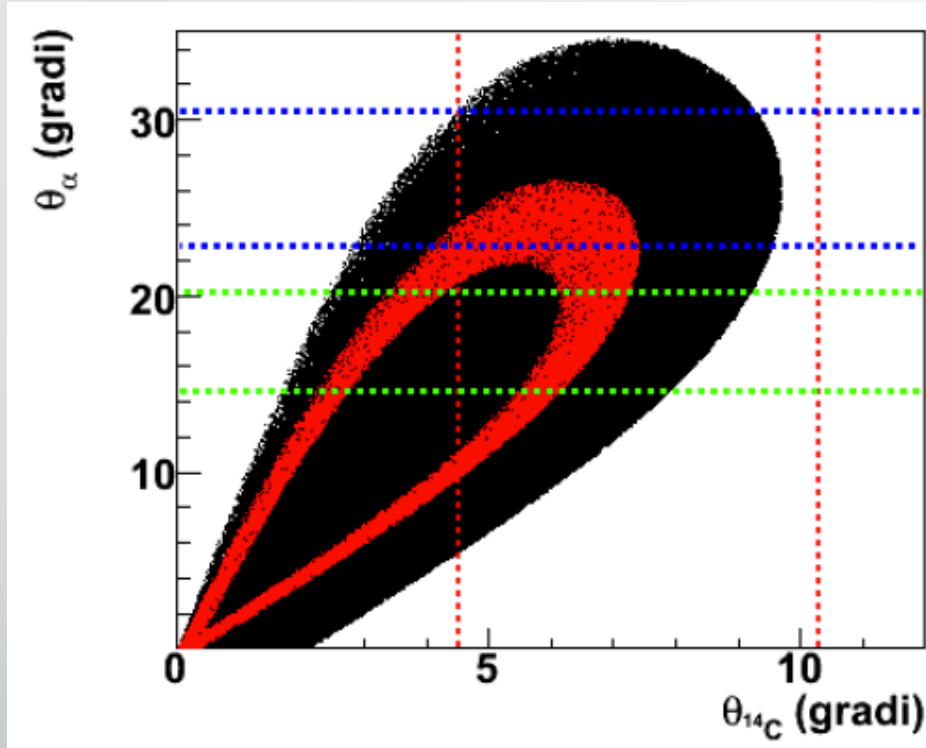
Suppressed due to the centrifugal barrier

Available in literature

$E_{c.m.}$ (keV)	$^{18}\text{O}^*$ (MeV)	J^π
-7	8.039	1^-
75	8.125	5^-
166	8.213	2^+
236	8.282	3^-

F. Ajzenberg-Selove, Nucl. Phys., A475, 1 (1987)

Preliminary Study



Black points: kinematic calculations
Red points: kinematic calculations + $|p_s| < 5 \text{ MeV}/c$

$$E_{cm} + B_{xs} = E_{ax} \approx 2.45 \text{ MeV} = E_{fascio} \frac{m_N}{m_N + m^{17}\text{O}}$$

$^{17}\text{O} + ^2\text{H}$, $E_{\text{beam}} = 43.3 \text{ MeV}$

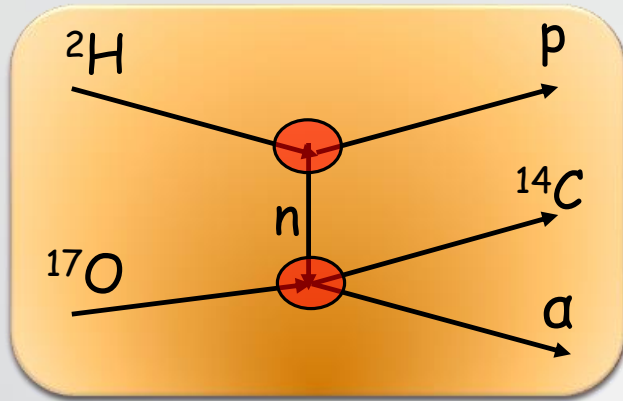
Deuteron as source of virtual
neutrons!!

This allows then to determine the
experimental apparatus:

^{14}C detection $\rightarrow 6^\circ < \theta_c < 10^\circ$

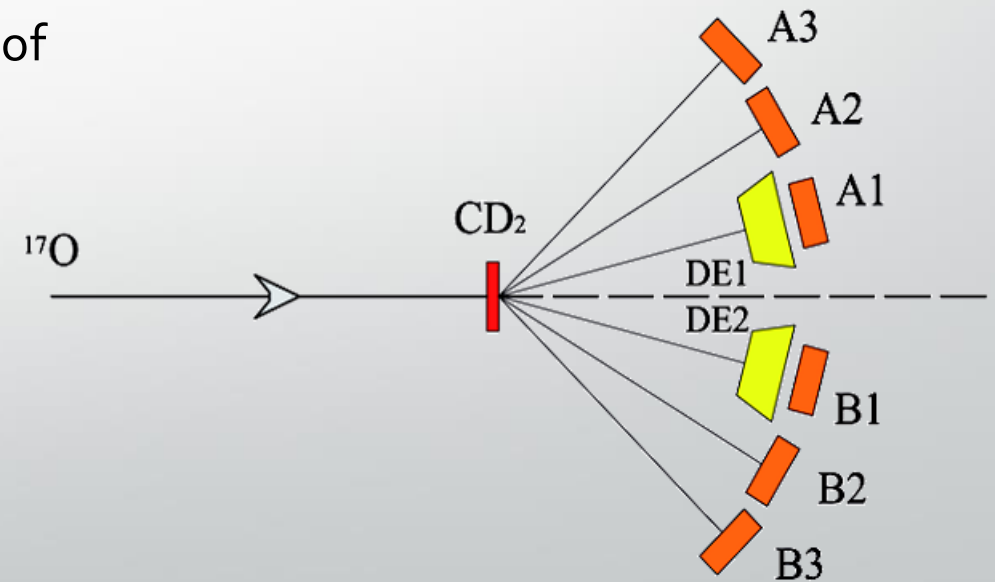
α detection $\rightarrow 15^\circ < \theta_\alpha < 20^\circ$
 $\rightarrow 23^\circ < \theta_\alpha < 31^\circ$

Experimental Setup



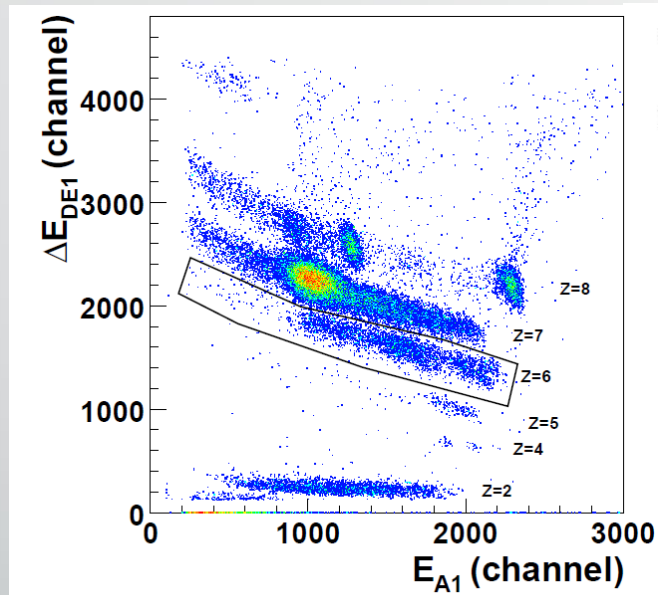
- The reaction $^{17}\text{O}(n, \alpha)^{14}\text{C}$ was studied via the $^2\text{H}(^{17}\text{O}, \alpha^{14}\text{C})p$, $V_{\text{coul}}=2.3$ MeV;
- The deuteron is the TH nucleus. Strong cluster $n+p$; $B=2.2$ MeV, $|p_s|=0$ MeV/c .

- ✓ Experiments performed at ISNAP at the University of Notre Dame (USA) and LNS of Catania;
- ✓ $E_{\text{beam}}(^{17}\text{O})= 43.5$ MeV;
- ✓ Target thickness $\text{CD}_2 \sim 150 \mu\text{g}/\text{cm}^2$;
- ✓ IC filled with ~ 50 mbar isobutane gas;
- ✓ Angular position to cover the QF angular region
- ✓ Symmetric set-up in order to increase the statistic.

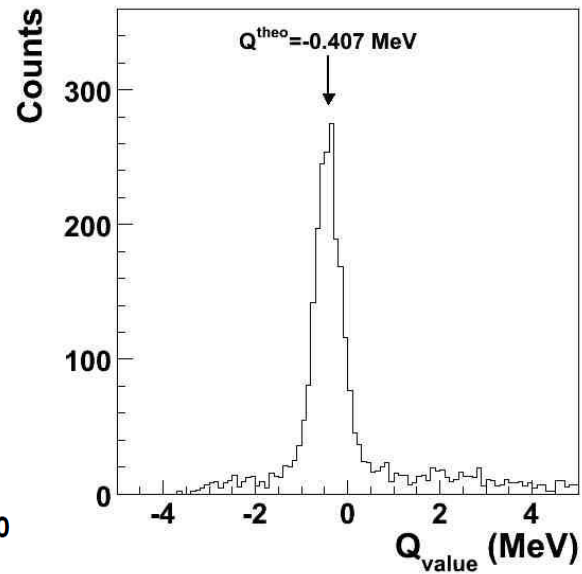


Data Analysis: Channel Selection

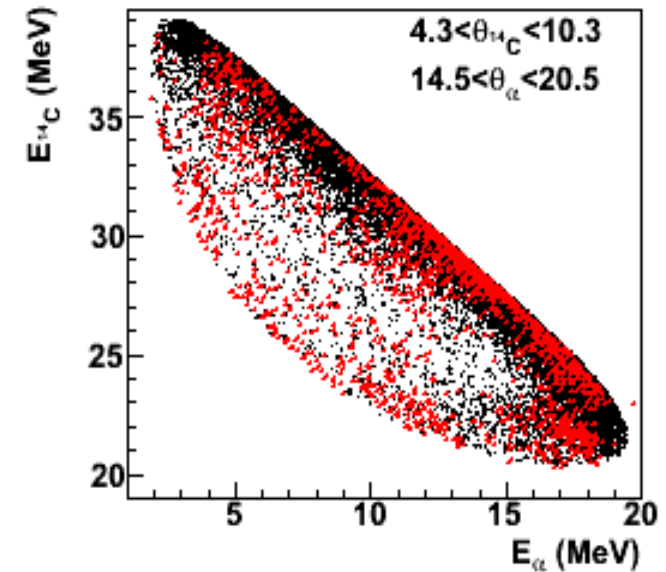
$${}^2\text{H}({}^{17}\text{O}, \alpha {}^{14}\text{C})\text{p} \rightarrow Q_{\text{value}} = -0.407 \text{ MeV}$$



» Selection of the events with carbon in the final state

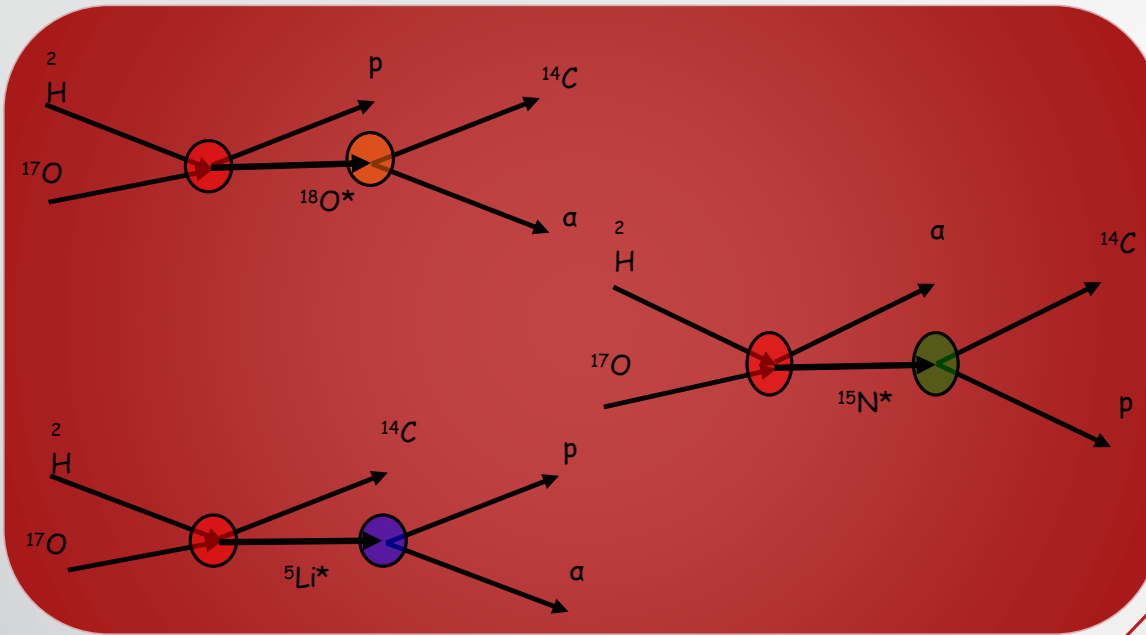


» Presence of a peak centered at the expected value

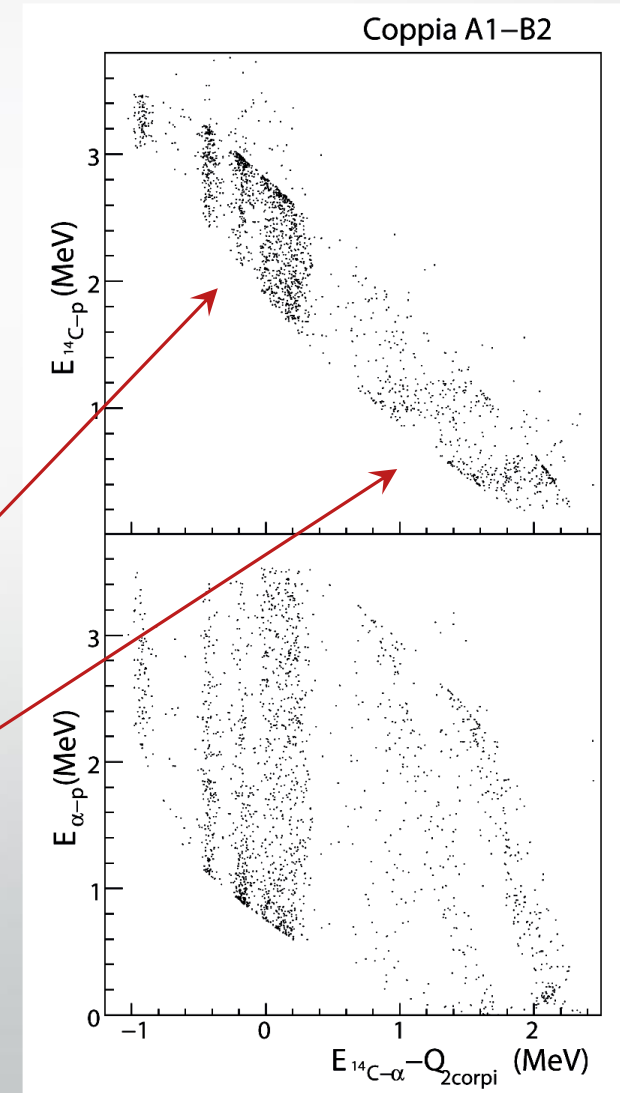


» Comparison between experimental and simulated kinematical locus

Data Analysis: Sequential Mechanisms

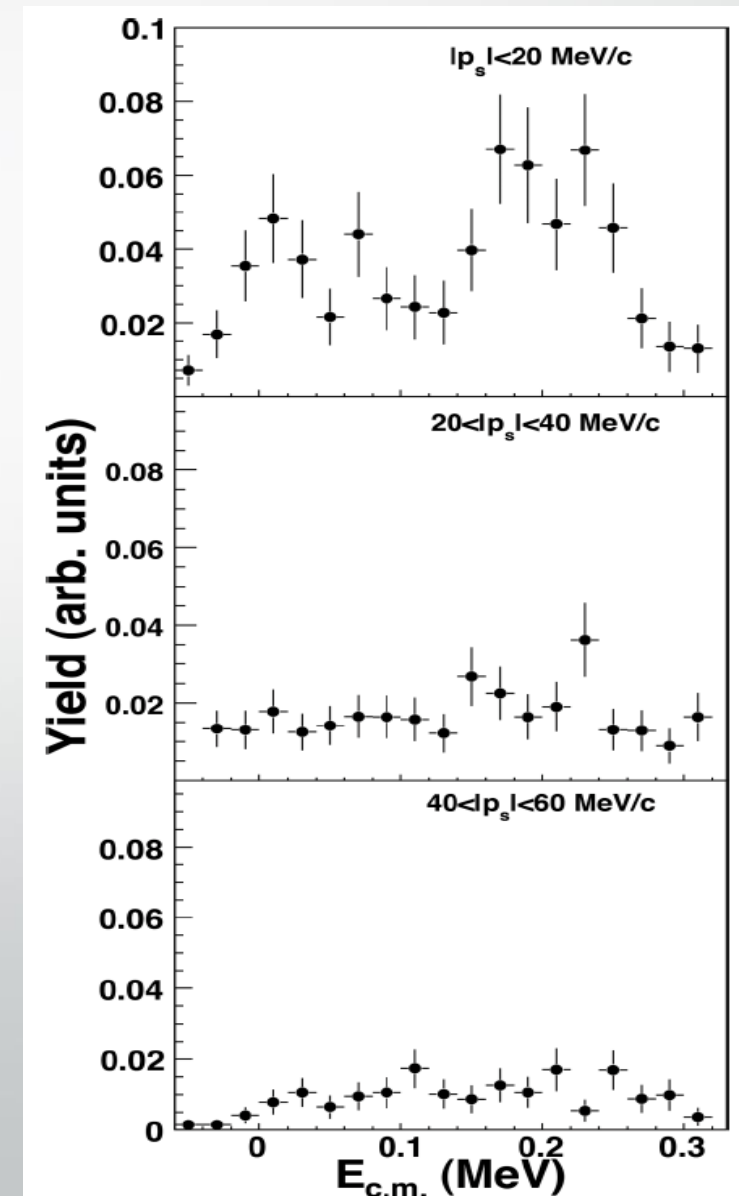


- » Presence of vertical loci corresponding to ^{18}O ;
- » Presence of ^{15}N levels but far from the zone of interest;
- » Absence of horizontal loci from ^5Li .



Data Analysis: QF Selection

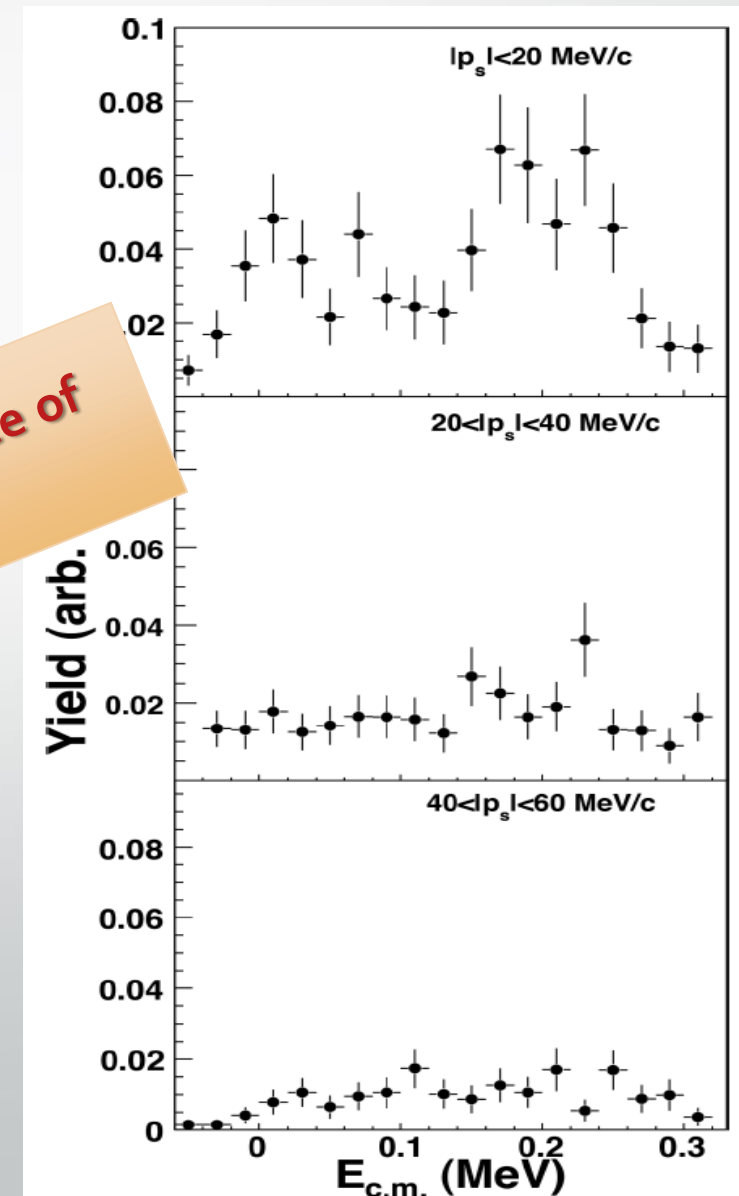
- 1) The experimental yield of the $E_{c.m.}(\text{MeV})=E_{\alpha^{14}\text{C}}-1.818$ has been studied;
- 2) Spectra are obtained with different condition on the *undetected* proton momentum;
- 3) The yield is enhanced around low-neutron momenta and decrease for $40 < |p_p| < 60$ MeV/c;
- 4) The coincidence yield appears strongly influenced by the p-n momentum distribution in the deuteron (having its maximum at 0 MeV/c).



Data Analysis: QF Selection

- 1) The experimental yield of the $E_{c.m.}(\text{MeV})=E_{\alpha^{14}\text{C}}-1.818$ has been studied;
- 2) Spectra are obtained with different condition on the *undetected* proton momentum;
- 3) The yield is enhanced around low-neutron momenta and decrease for $40 < p_s < 60$ MeV/c;
- 4) The coincidence yield is influenced by the distribution in the deuteron (maximum at 0 MeV/c).

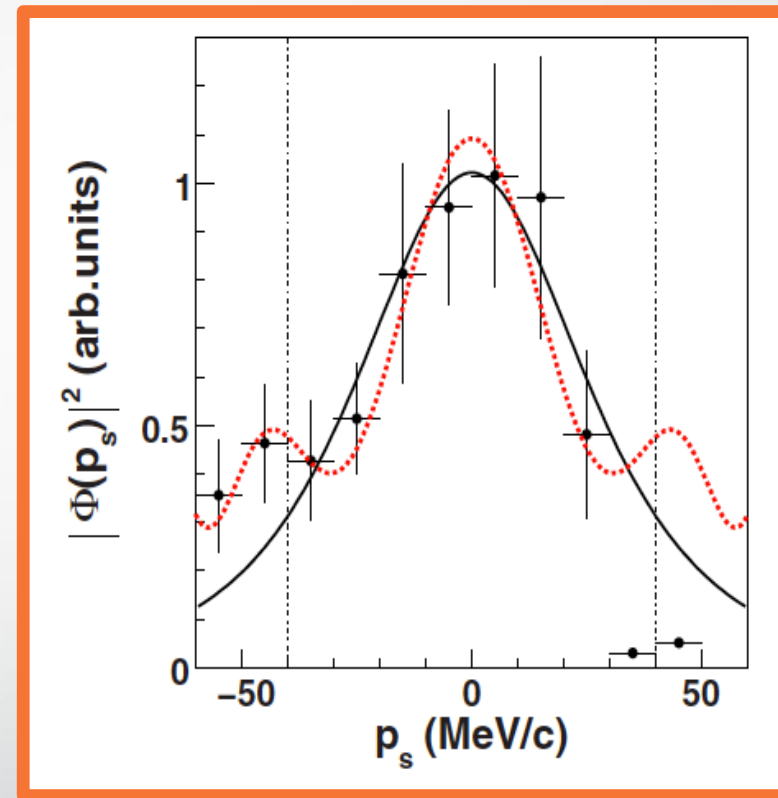
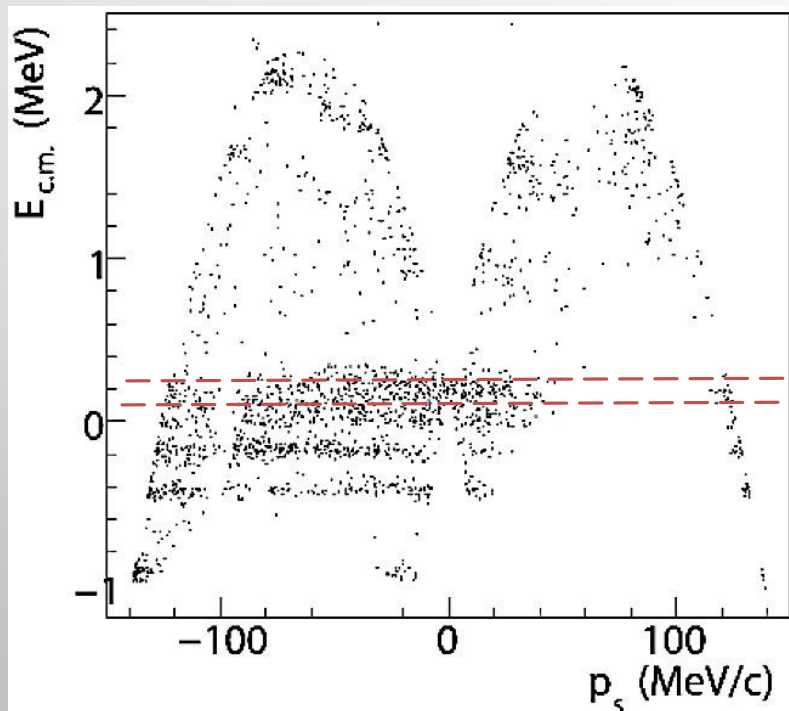
Necessary condition for the presence of the QF mechanism



Data Analysis: QF Selection

By following the PWIA approach it is possible to extract the experimental momentum distribution:

$$|\Phi(p_s)|^2 \propto \frac{d^3\sigma}{dE_c d\Omega_c d\Omega_c} / KF$$

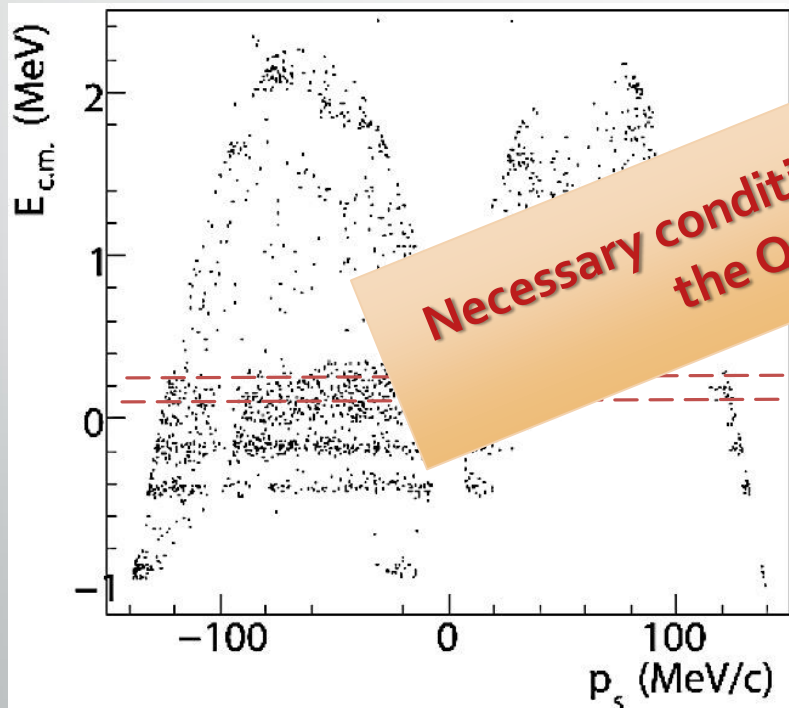


- Theoretical Hulthen Function
- - - DWBA calculation (FRESCO code)

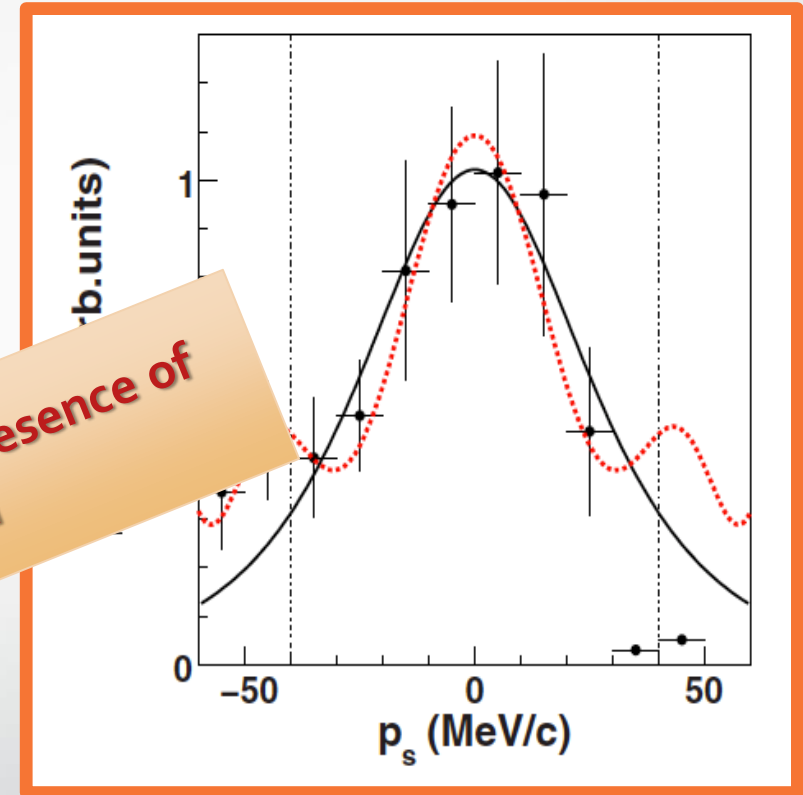
Data Analysis: QF Selection

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$$|\Phi(p_s)|^2 \propto \frac{d^3\sigma}{dE_c d\Omega_c d\Omega_c} / KF$$



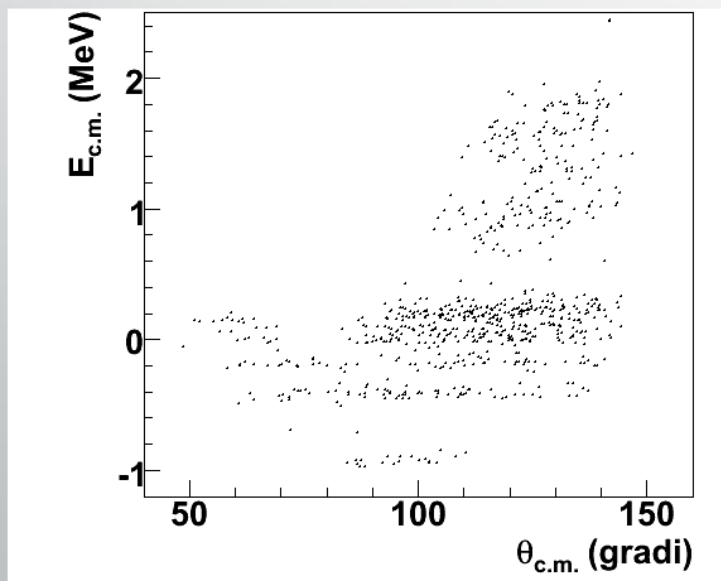
Necessary condition for the presence of the QF mechanism



- Theoretical Hulthen Function
- - - DWBA calculation (FRESCO code)

Data Results: Angular Distributions

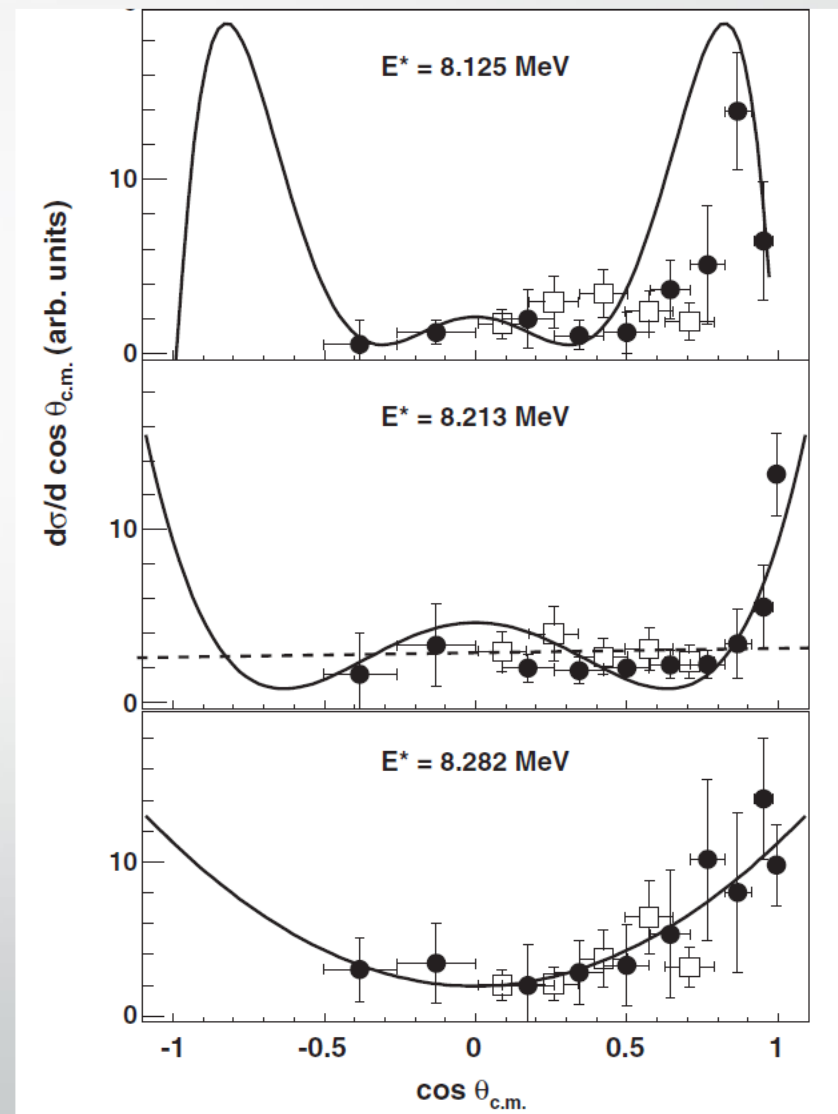
After the study of the 3-body channel and the QF selection, it is important to study the 2-body one. The angular range covered in the experiments in the c.m. system allows one to study the angular distributions



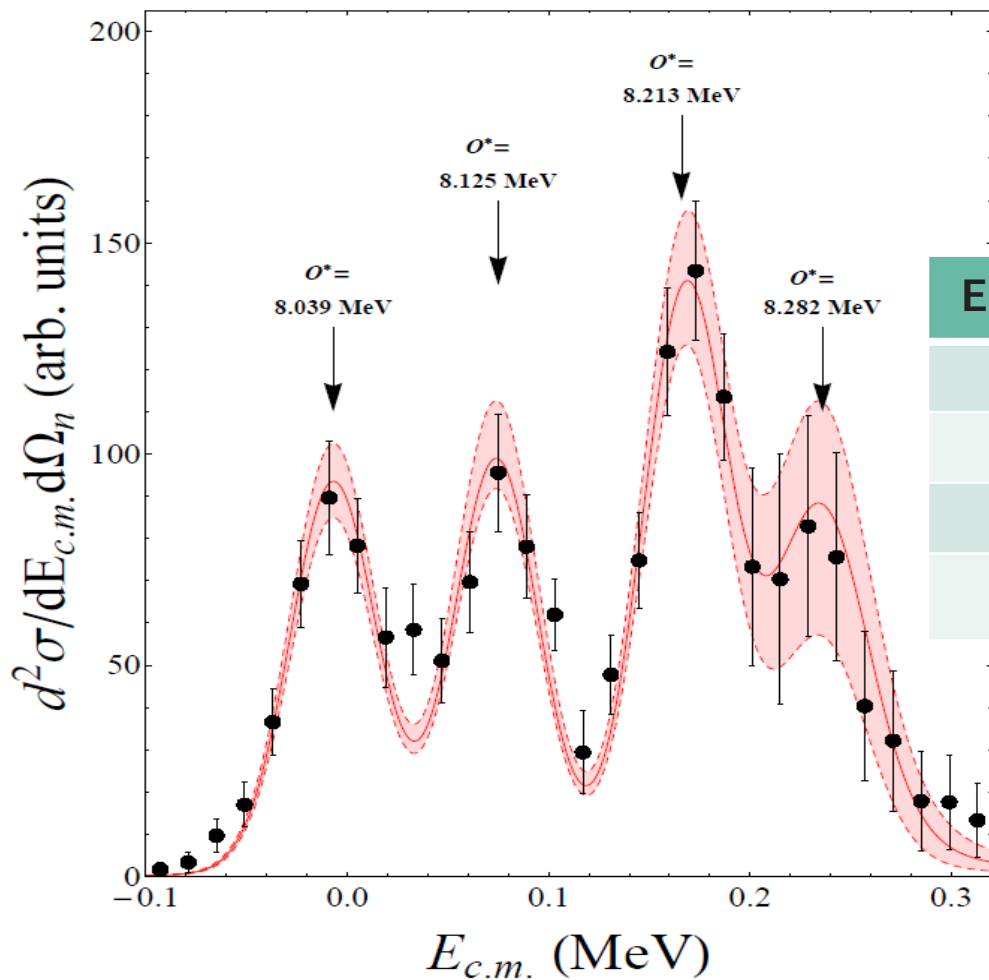
$\ell=3$ distribution: no data present in literature (suppressed in direct measurements)

$\ell=2$ distribution: no data present in literature

$\ell=1$ distribution: will be compared with the available data



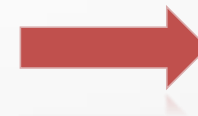
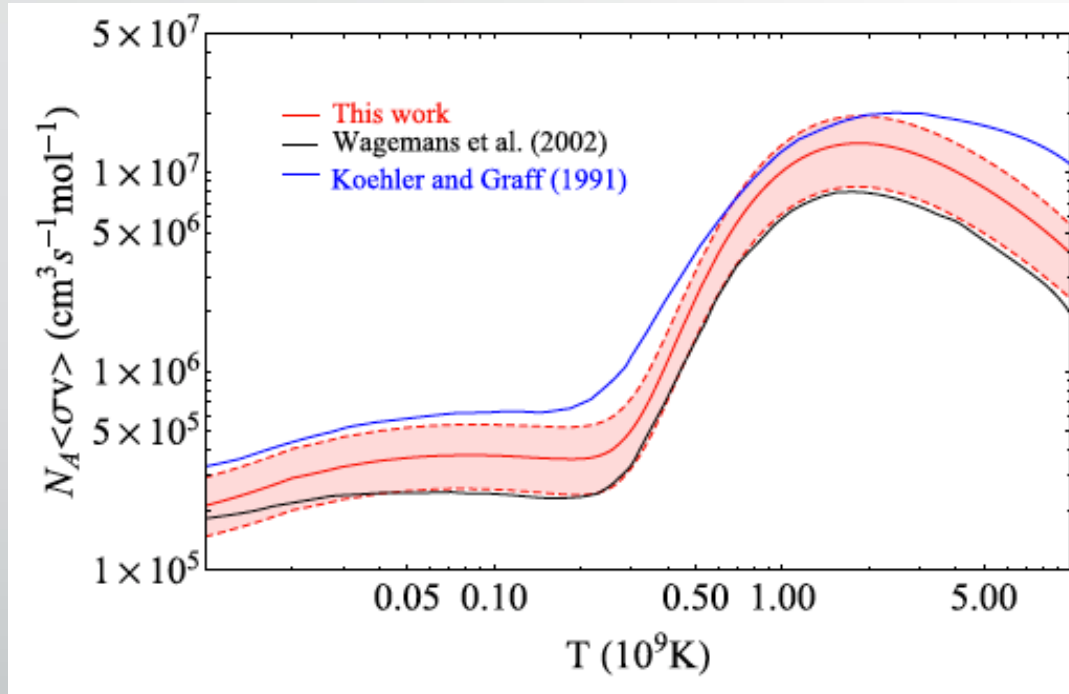
Data Results: R-Matrix Fit



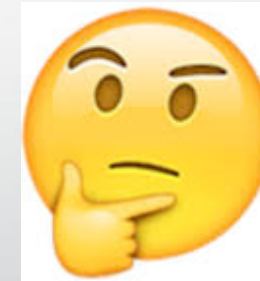
E_{cm} (keV)	Γ_n (eV)	Γ_α (eV)	Γ_{TOT} (eV)	$\Gamma_{wag.}$ (eV)
-7	$0,01 \pm 0,001$	2362 ± 307	2362 ± 307	2400
75	$0,05 \pm 0,006$	36 ± 5	36 ± 5	-
166	86 ± 11	2171 ± 282	2257 ± 293	2258 ± 135
236	1714 ± 446	13021 ± 3386	14735 ± 3832	14739 ± 590

Guardo et al., Phys. Rev. C, 95, 025807, 2017

Data Results: Astrophysical Rate



Possible impact
on stellar
nucleosynthesis?



Recent Results of THM

THE ASTROPHYSICAL JOURNAL, 845:19 (13pp), 2017 August 10

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<https://doi.org/10.3847/1538-4357/aa7de7>



New Improved Indirect Measurement of the $^{19}\text{F}(p, \alpha)^{16}\text{O}$ Reaction at Energies of Astrophysical Relevance

I. Indelicato¹, M. La Cognata¹, C. Spitaleri^{1,2}, V. Burjan³, S. Cherubini^{1,2}, M. Gulino^{1,4}, S. Hayakawa⁵, Z. Hons³, V. Kroha³, L. Lamia¹, M. Mazzocco^{6,7}, J. Mrazek³, R. G. Pizzone¹, S. Romano^{1,2}, E. Strano^{6,7}, D. Torresi^{6,7}, and A. Tumino^{1,4}

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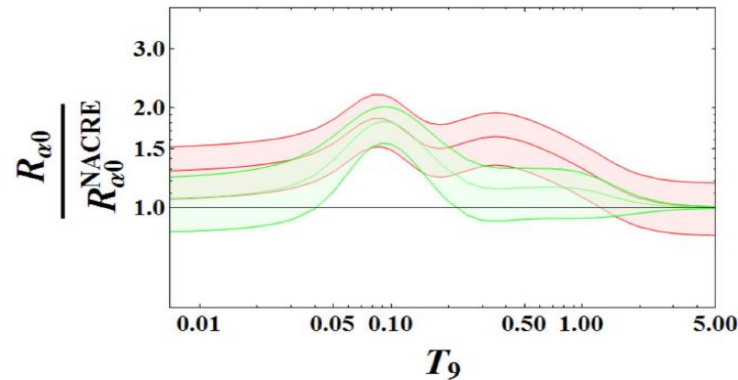
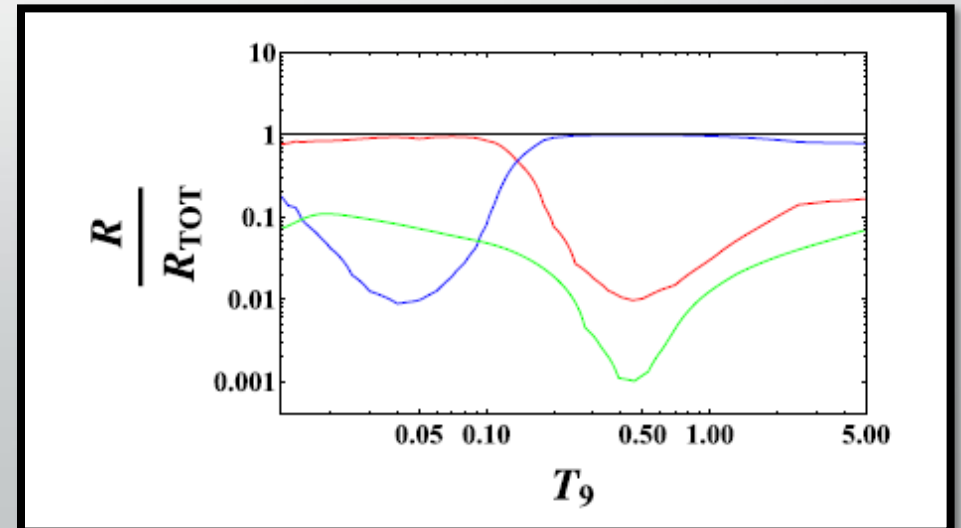


Figure 15. Ratio of the reaction rate calculation obtained from the THM astrophysical factor (red band) to the rate recommended in NACRE (Angulo et al. 1999). T_9 is the temperature in GK ($T_9 = T/10^9$ K). The black line corresponds to $R/R_{\alpha_0}^{\text{nacre}} = 1$. For comparison, the $R/R_{\alpha_0}^{\text{nacre}}$ ratio given in La Cognata et al. (2015) is shown as a green band.

Three open channels:

- $^{19}\text{F}(p, \alpha_0)^{16}\text{O}$ —
- $^{19}\text{F}(p, \alpha_\pi)^{16}\text{O}$ —
- $^{19}\text{F}(p, \alpha_\gamma)^{16}\text{O}$ —



Recent Results of THM

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CrossMark

The $^{19}\text{F}(\alpha, p)^{22}\text{Ne}$ Reaction at Energies of Astrophysical Relevance by Means of the Trojan Horse Method and Its Implications in AGB Stars

G. D'Agata^{1,2}, R. G. Pizzone¹, M. La Cognata¹, I. Indelicato¹, C. Spitaleri^{1,2}, S. Palmerini^{3,4}, O. Trippella^{3,4}, D. Vescovi^{3,4}, S. Blagus⁵, S. Cherubini^{1,2}, P. Figuera¹, L. Grassi⁵, G. L. Guardo¹, M. Gulino^{1,6}, S. Hayakawa^{1,7}, R. Kshetri^{1,8}, L. Lamia^{1,2}, M. Lattuada^{1,2}, T. Mijatović⁵, M. Milin⁹, Đ. Miljanić^{5,10}, L. Prepolec⁵, G. G. Rapisarda¹, S. Romano^{1,2}, M. L. Sergi^{1,2}, N. Skukan⁵, N. Soić⁵, V. Tokić⁵, A. Tumino^{1,6}, and M. Uroić⁵

THE ASTROPHYSICAL JOURNAL, 836:57 (6pp), 2017 February 10

<https://doi.org/10.3847/1538-4357/836/1/57>

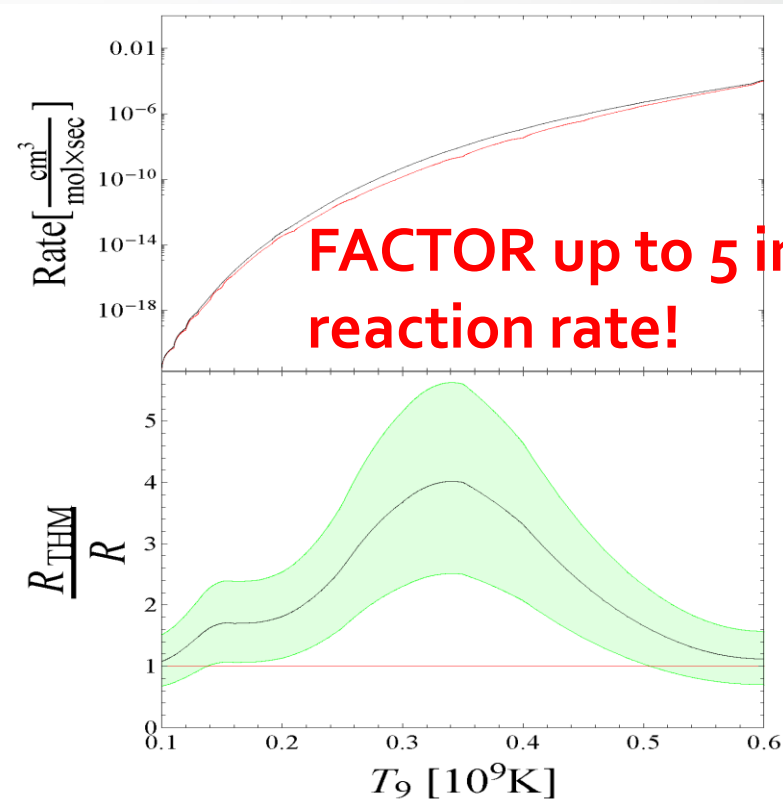
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CrossMark

First Measurement of the $^{19}\text{F}(\alpha, p)^{22}\text{Ne}$ Reaction at Energies of Astrophysical Relevance

R. G. Pizzone¹, G. D'Agata^{1,2}, M. La Cognata¹, I. Indelicato¹, C. Spitaleri^{1,2}, S. Blagus³, S. Cherubini^{1,2}, P. Figuera¹, L. Grassi³, G. L. Guardo¹, M. Gulino^{1,4}, S. Hayakawa^{1,5}, R. Kshetri^{1,6}, L. Lamia^{1,2}, M. Lattuada^{1,2}, T. Mijatović³, M. Milin⁷, Đ. Miljanić^{3,8}, L. Prepolec³, G. G. Rapisarda¹, S. Romano^{1,2}, M. L. Sergi¹, N. Skukan³, N. Soić³, V. Tokić³, A. Tumino^{1,4}, and M. Uroić³



Recent Results of THM

nature

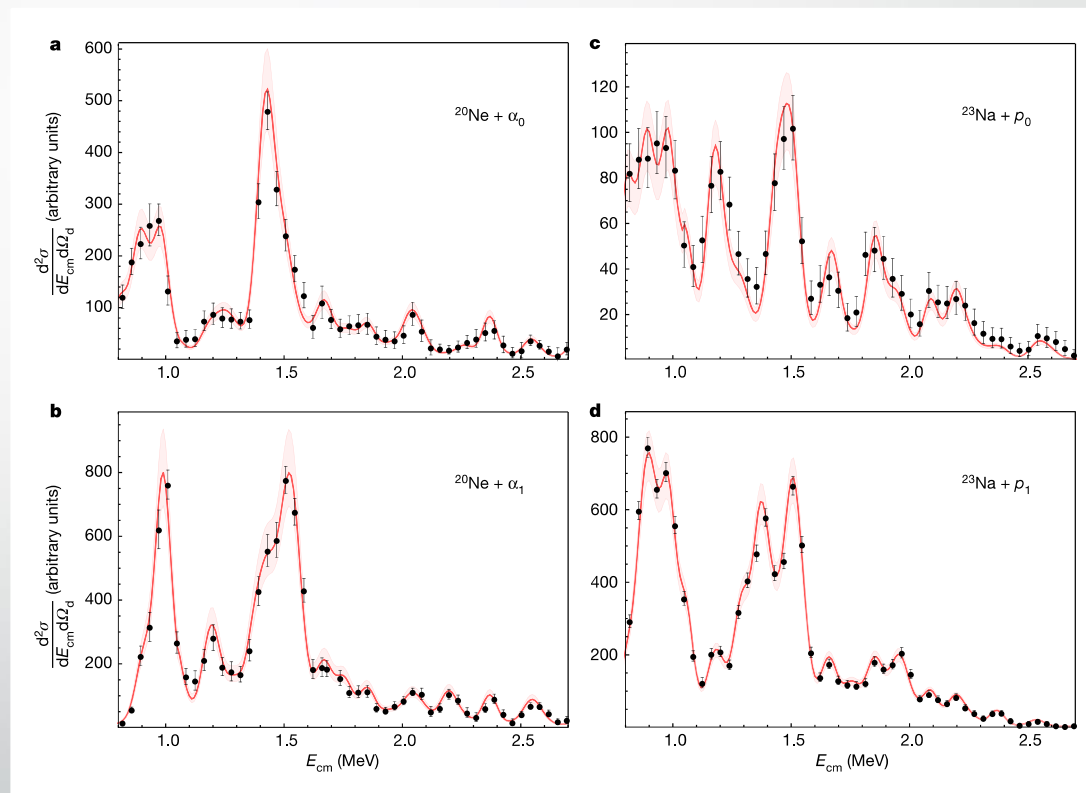
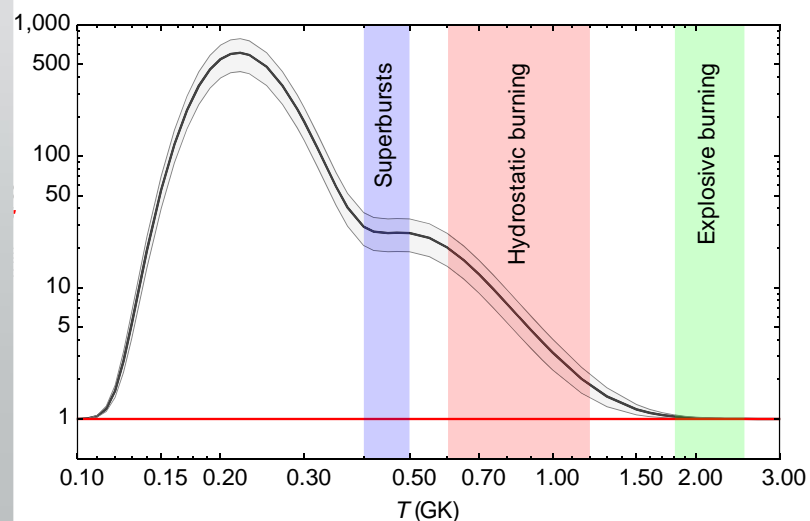
International journal of science

Letter | Published: 23 May 2018

An increase in the $^{12}\text{C} + ^{12}\text{C}$ fusion rate from resonances at astrophysical energies

A. Tumino , C. Spitaleri, M. La Cognata, S. Cherubini, G. L. Guardo, M. Gulino, S. Hayakawa, I. Indelicato, L. Lamia, H. Petrascu, R. G. Pizzone, S. M. R. Puglia, G. G. Rapisarda, S. Romano, M. L. Sergi, R. Spartá & L. Trache

Nature **557**, 687–690 (2018) | [Download Citation](#) ↓






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A Trojan Horse Approach to the Production of ^{18}F in Novae

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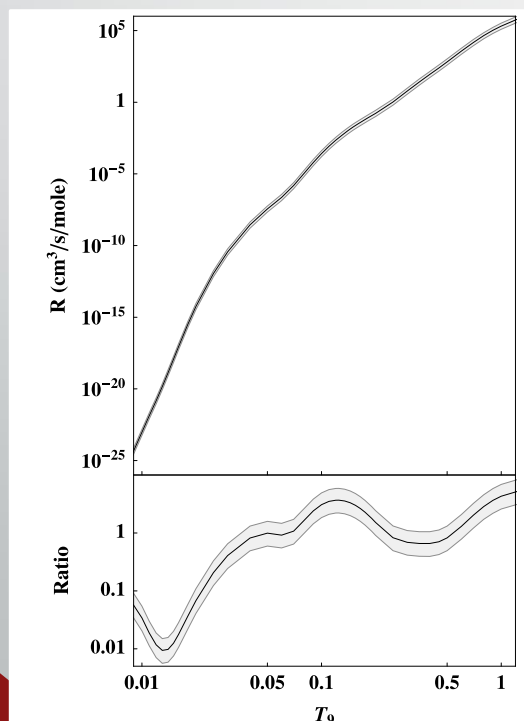
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THM successfully applied to RIBs

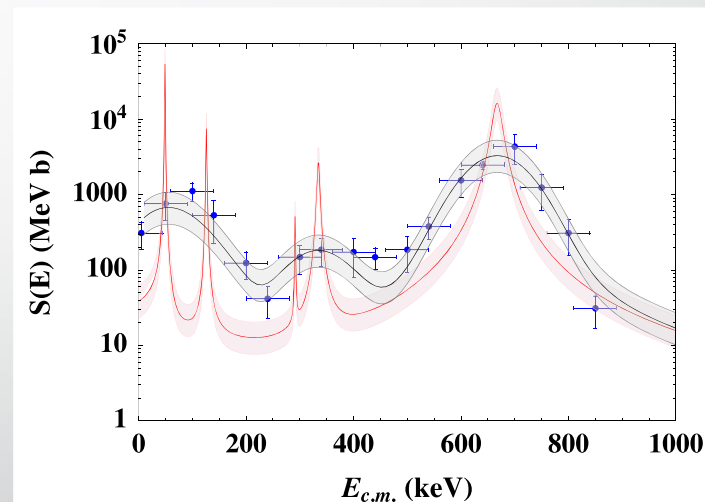


Figure 2. R-matrix analysis of the THM astrophysical factor (blue points) as in Figure 1. The evaluated uncertainty in the R-matrix fit is reported as a shadowed gray area and as a red band for the corresponding deconvoluted S(E)-factor.

Experiment in CNS RIKEN and Texas A&M

Recent Results of THM

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On the Determination of the ${}^7\text{Be}(n, \alpha){}^4\text{He}$ Reaction Cross Section at BBN Energies

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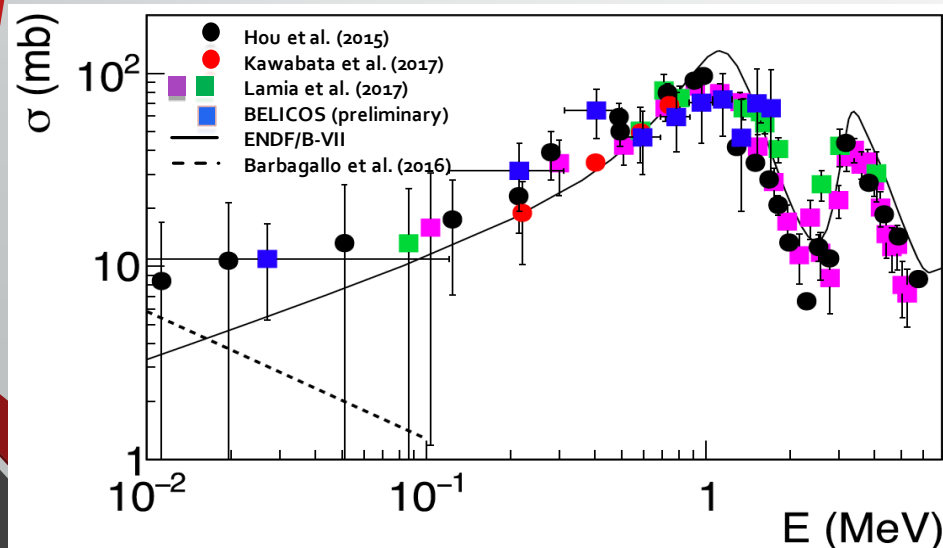
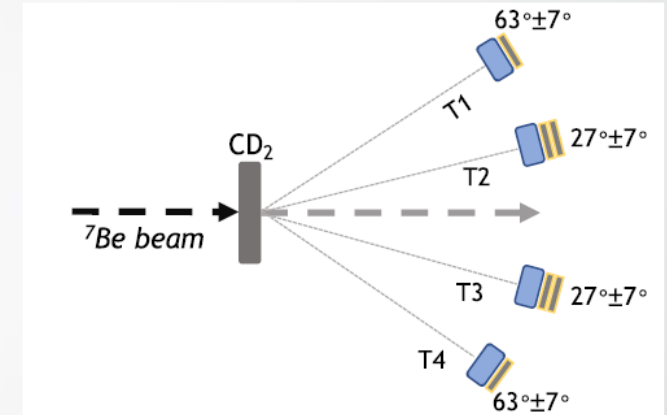
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Application of THM with RIBs and neutron induced reactions



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Cross-section Measurement of the Cosmologically Relevant ${}^7\text{Be}(n, \alpha){}^4\text{He}$ Reaction over a Broad Energy Range in a Single Experiment

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Advantages of THM

A - It is possible to measure the bare nucleus cross section s_b (or the bare nucleus Astrophysical Factor $S_b(E)$) at Gamow energy for reactions involving charged particles and neutron.

No extrapolation

B - It is possible to measure excitation function in a “ relatively” short time because typical order of magnitude for a three- body cross- section is mb;

C - One of the few ways to measure the electron screening effect; comparison with direct data;

D - Application to the radioactive beam measurements;

Main limitations of THM

- A- Preliminary study of quasi-free mechanism and tests of validity are necessary.
 - Presence of different 3-body reaction mechanisms (Sequential Decay – Quasi-Free)
- B- No absolute cross section is measurable:
 - The excitation functions at energies above/below Coulomb barrier must be known from direct measurements;
- C- Measurements with high angular and energy resolutions are needed;
- D- Theoretical analysis is needed:
 - PWIA, MPWBA

TH Method is complementary to direct measurements as well as other indirect methods.

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*Thank you for your
attention*