

# What have we learned from GW170817?

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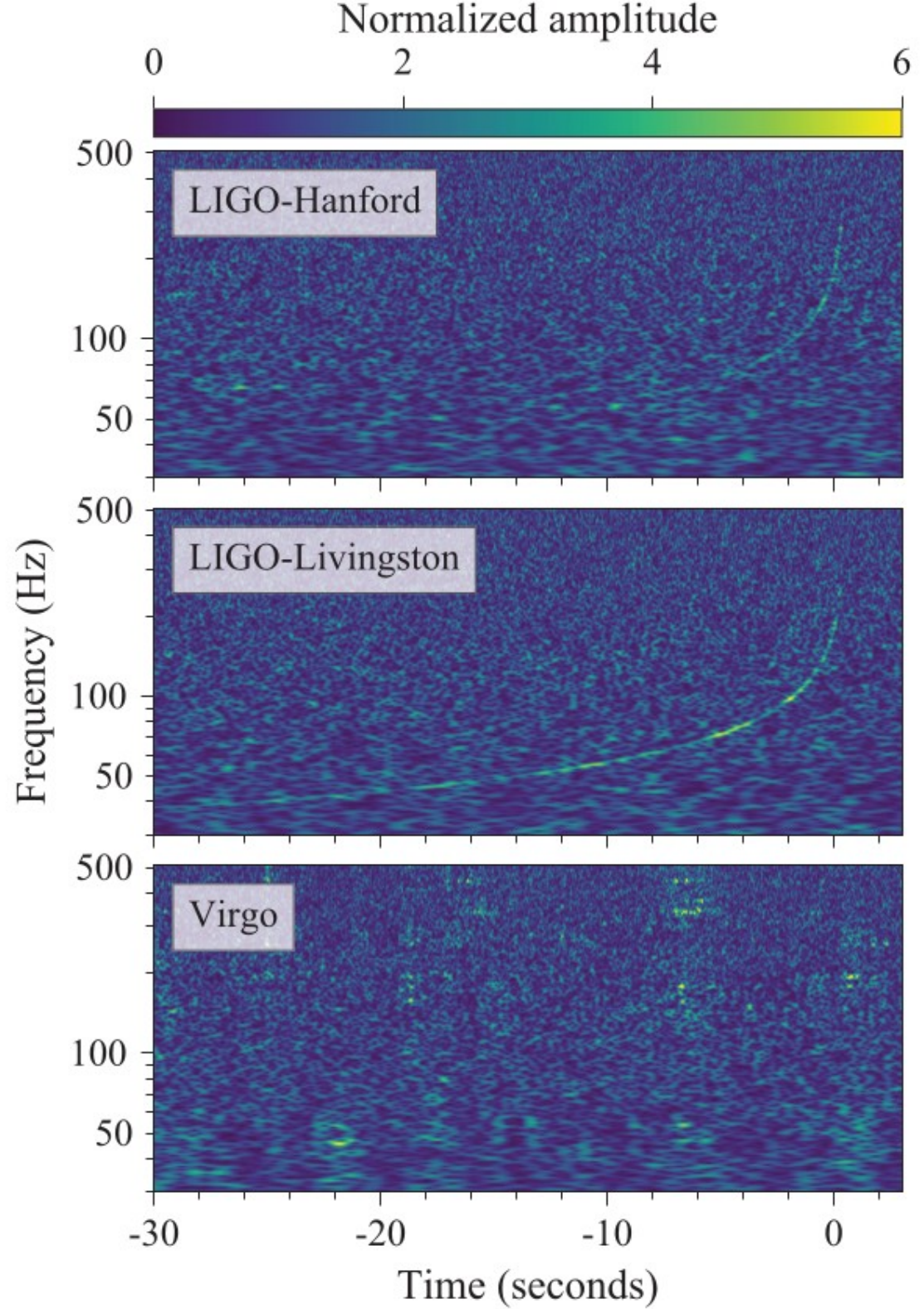
# Outline

- Observations
- NS star properties
- Fundamental physics
- Astrophysics

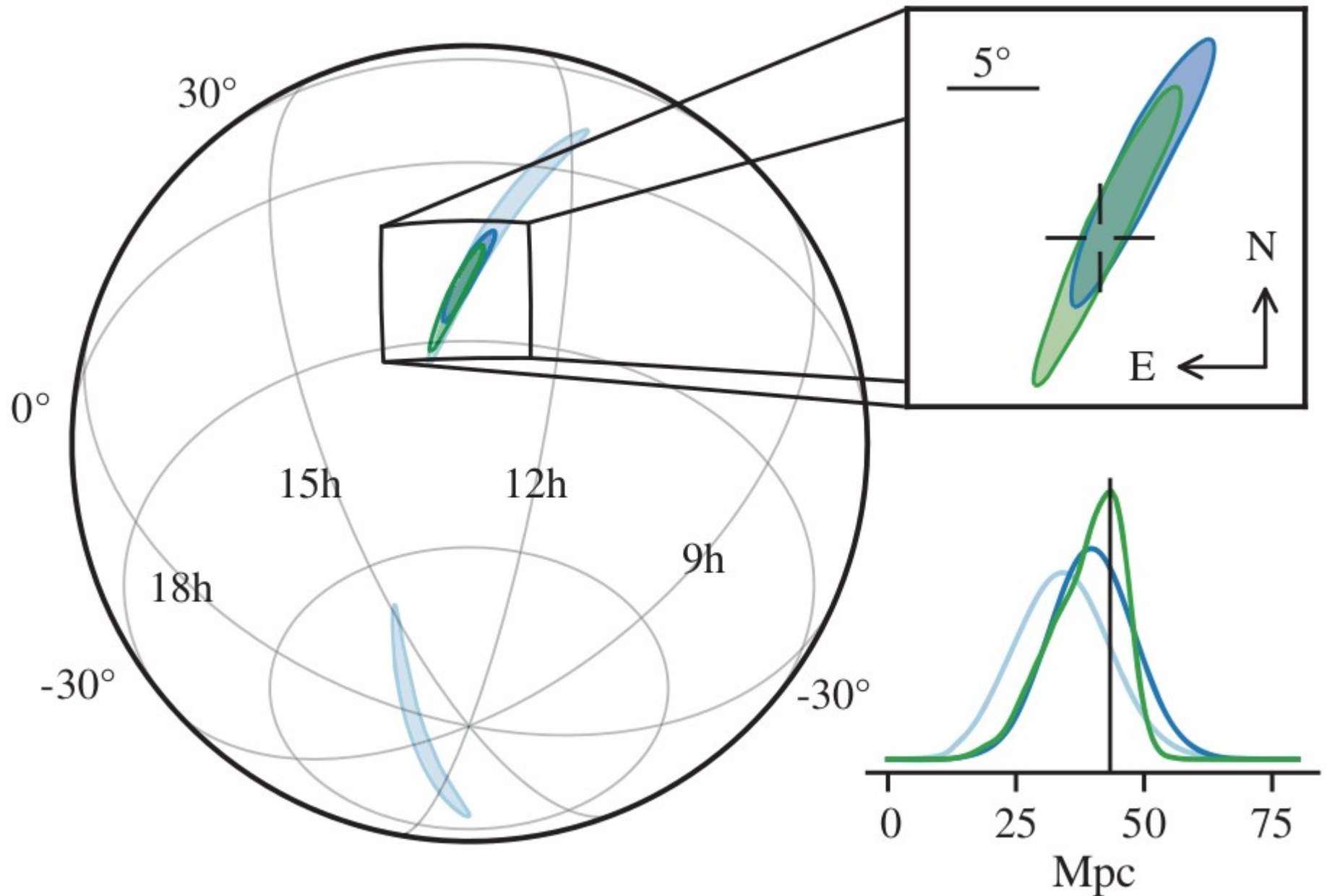
# Observations

Combined SNR = 32.6

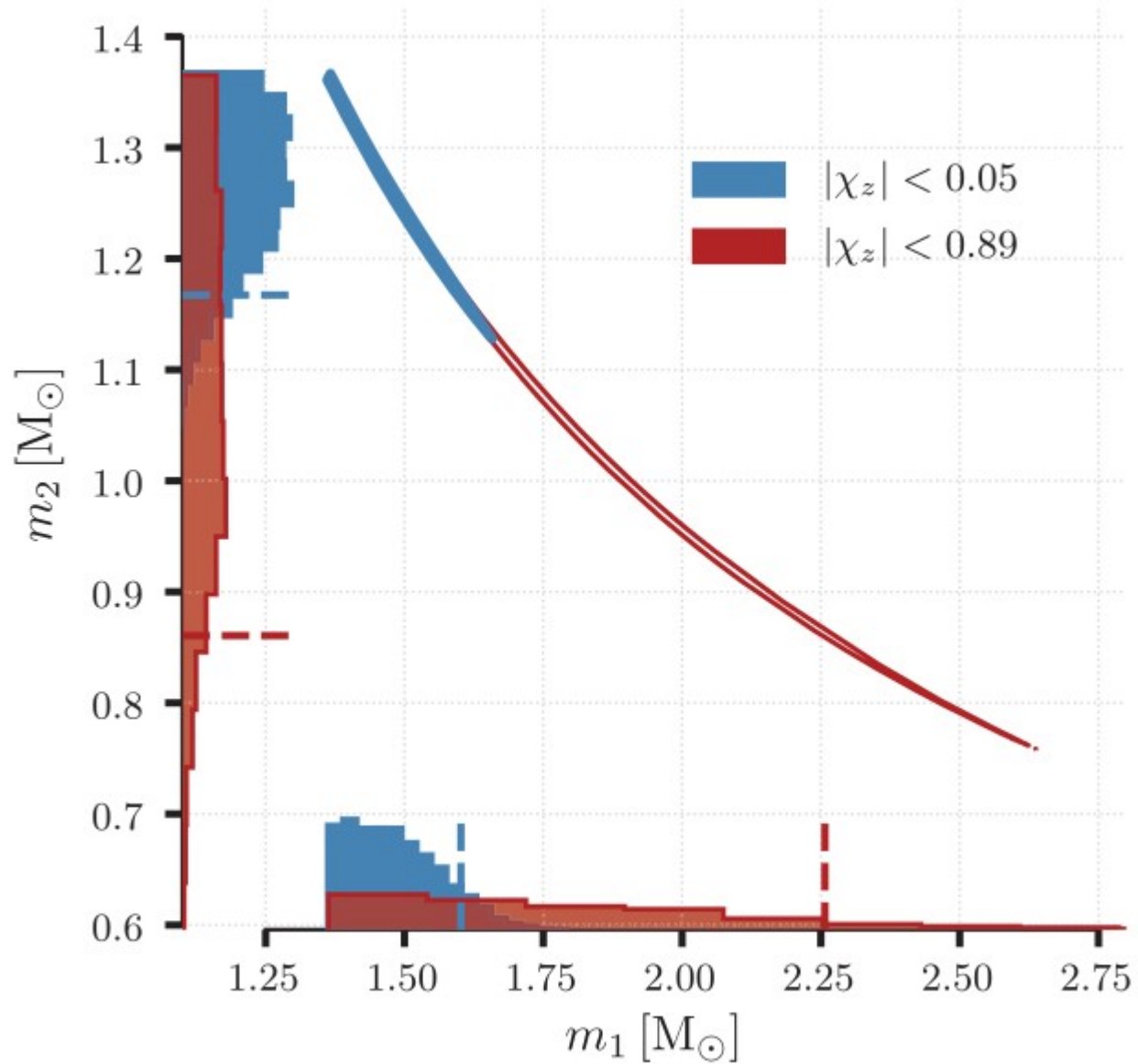
FAR < 1 in 80000 years



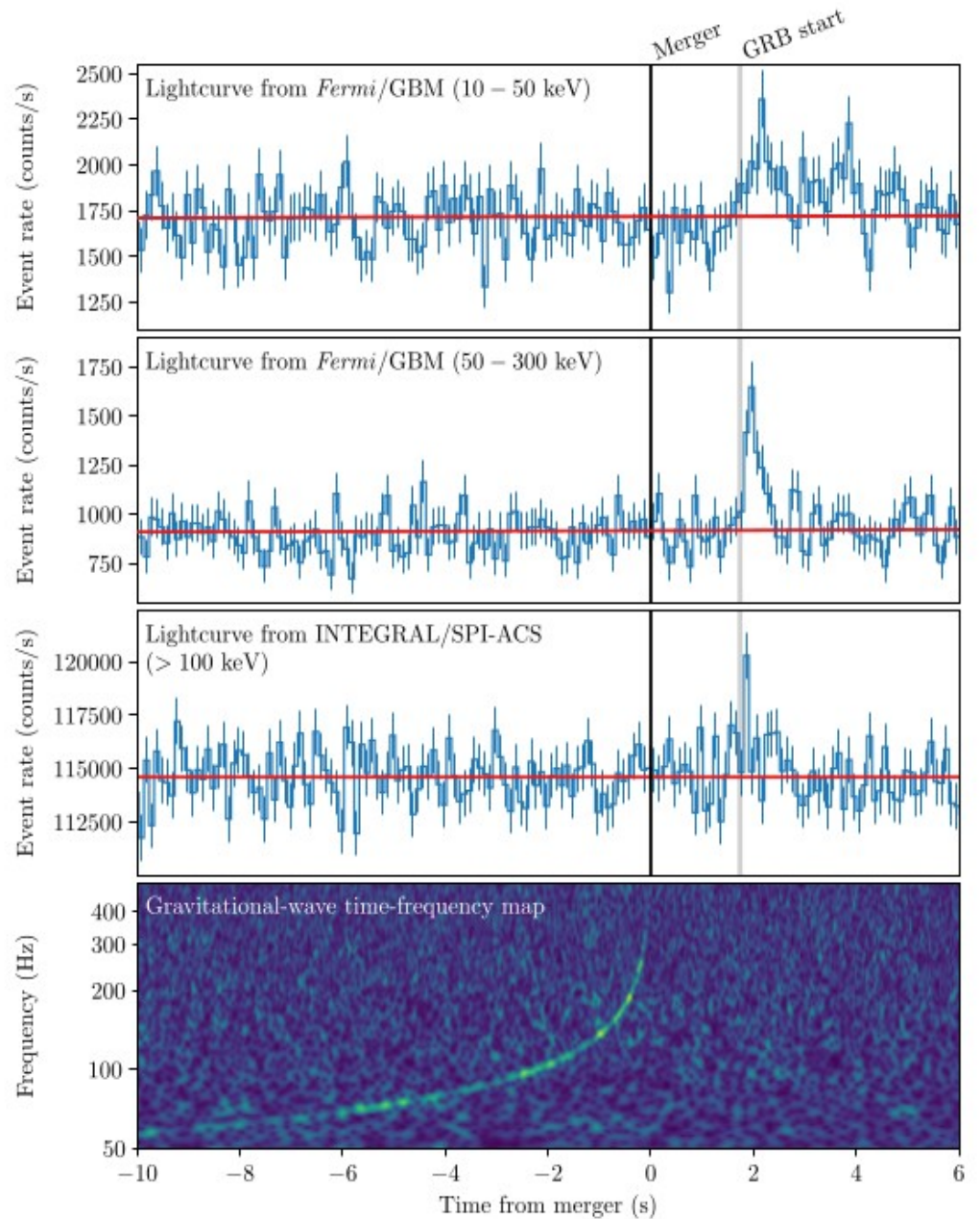
# Sky position and distance



# Masses



# GRB at the same time



# Optical counterpart

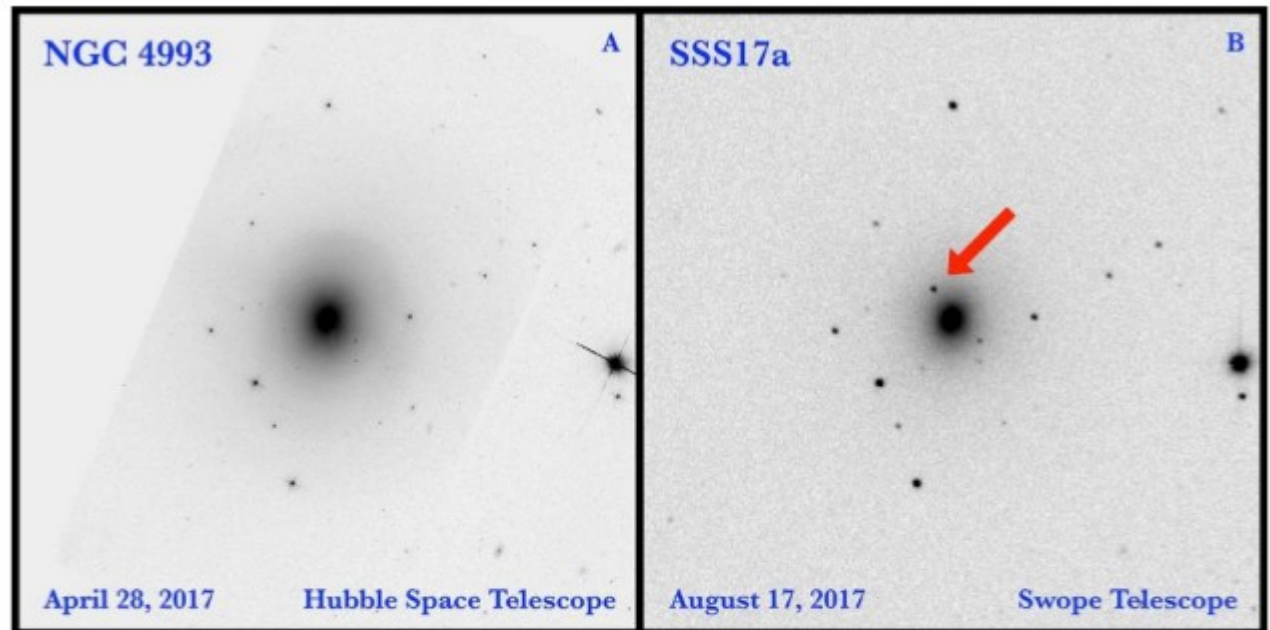
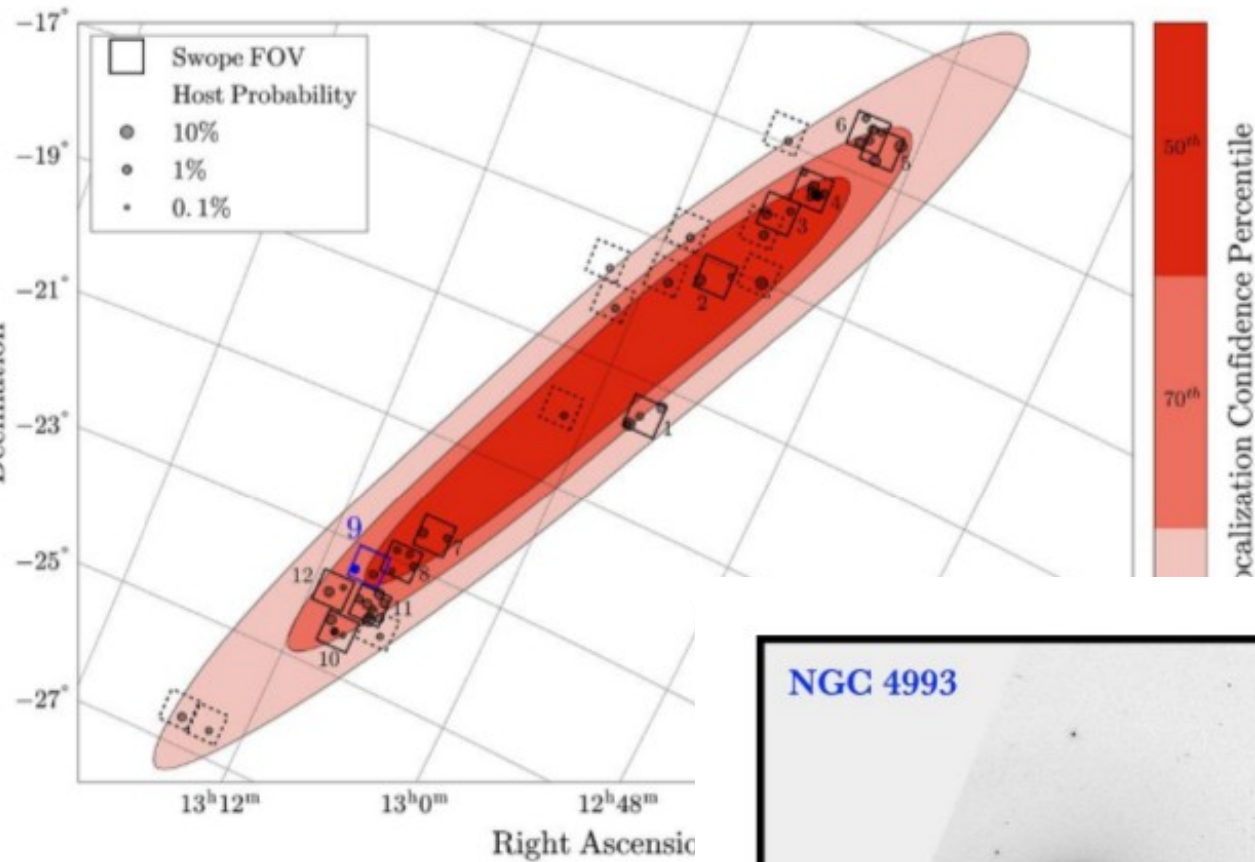


Fig. 4. 3×3 arcminute images centered on NGC 4993 with North up and East left. (A) *Hubble Space Telescope* F606W-band (broad V)

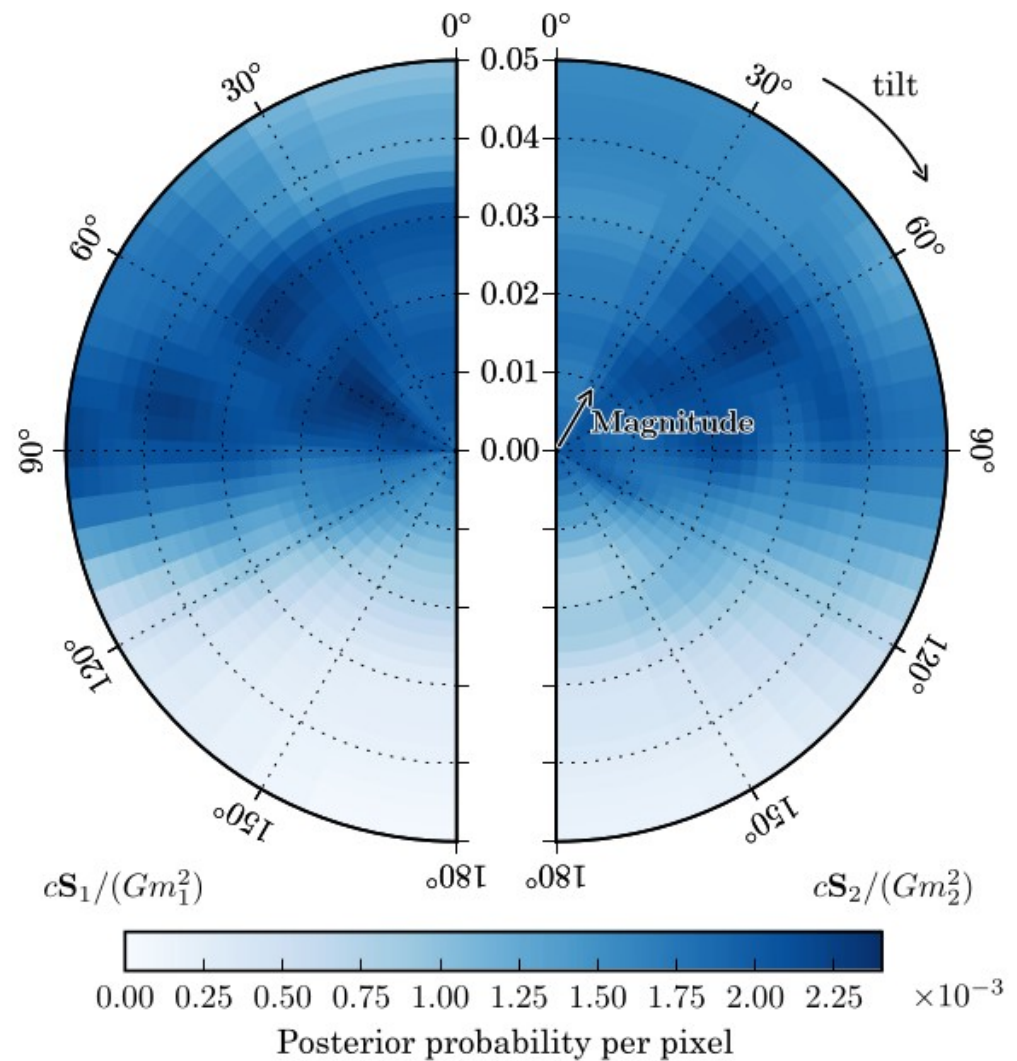
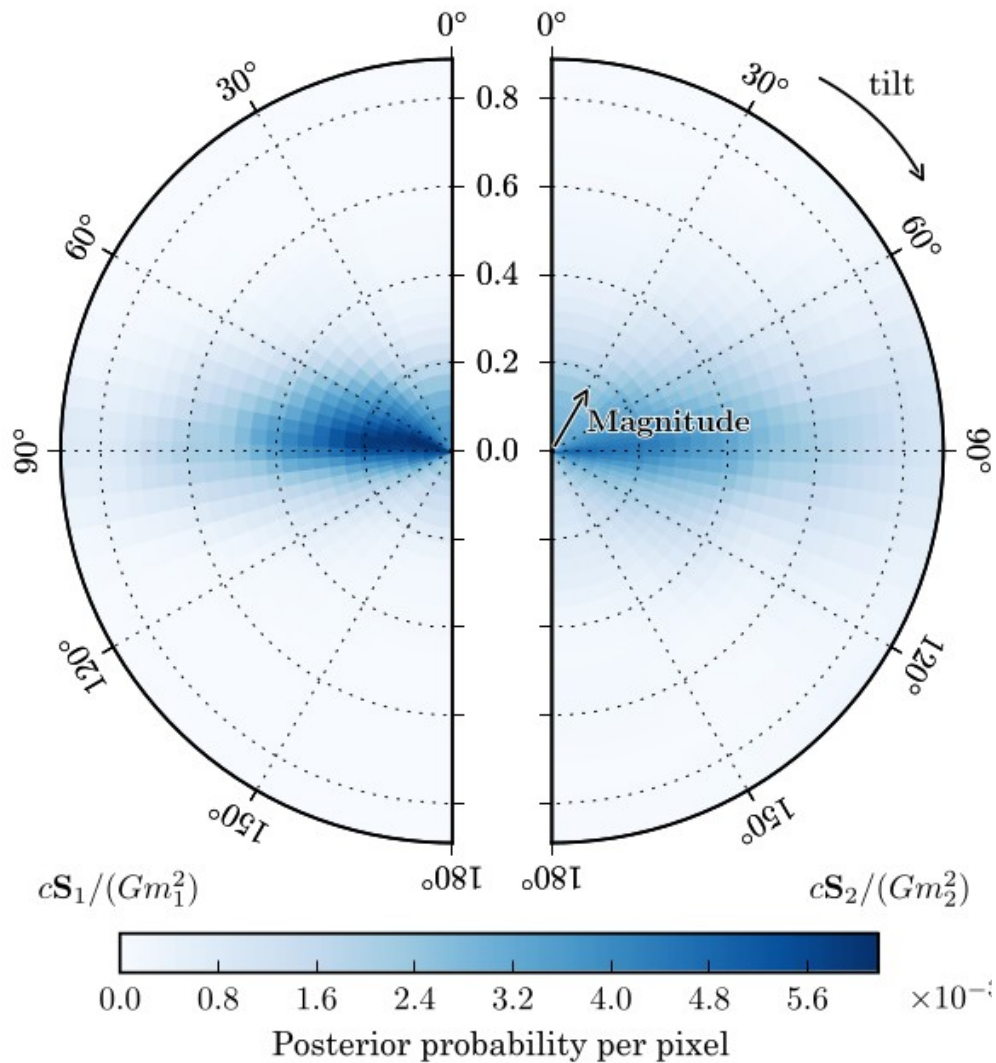
# Parameters

TABLE I. Source properties for GW170817: we give ranges encompassing the 90% credible intervals for different assumptions of the waveform model to bound systematic uncertainty. The mass values are quoted in the frame of the source, accounting for uncertainty in the source redshift.

	Low-spin priors ( $ \chi  \leq 0.05$ )	High-spin priors ( $ \chi  \leq 0.89$ )
Primary mass $m_1$	1.36–1.60 $M_\odot$	1.36–2.26 $M_\odot$
Secondary mass $m_2$	1.17–1.36 $M_\odot$	0.86–1.36 $M_\odot$
Chirp mass $\mathcal{M}$	$1.188^{+0.004}_{-0.002} M_\odot$	$1.188^{+0.004}_{-0.002} M_\odot$
Mass ratio $m_2/m_1$	0.7–1.0	0.4–1.0
Total mass $m_{\text{tot}}$	$2.74^{+0.04}_{-0.01} M_\odot$	$2.82^{+0.47}_{-0.09} M_\odot$
Radiated energy $E_{\text{rad}}$	$> 0.025 M_\odot c^2$	$> 0.025 M_\odot c^2$
Luminosity distance $D_L$	$40^{+8}_{-14}$ Mpc	$40^{+8}_{-14}$ Mpc
Viewing angle $\Theta$	$\leq 55^\circ$	$\leq 56^\circ$
Using NGC 4993 location	$\leq 28^\circ$	$\leq 28^\circ$
Combined dimensionless tidal deformability $\tilde{\Lambda}$	$\leq 800$	$\leq 700$
Dimensionless tidal deformability $\Lambda(1.4M_\odot)$	$\leq 800$	$\leq 1400$



# Spins – two cases

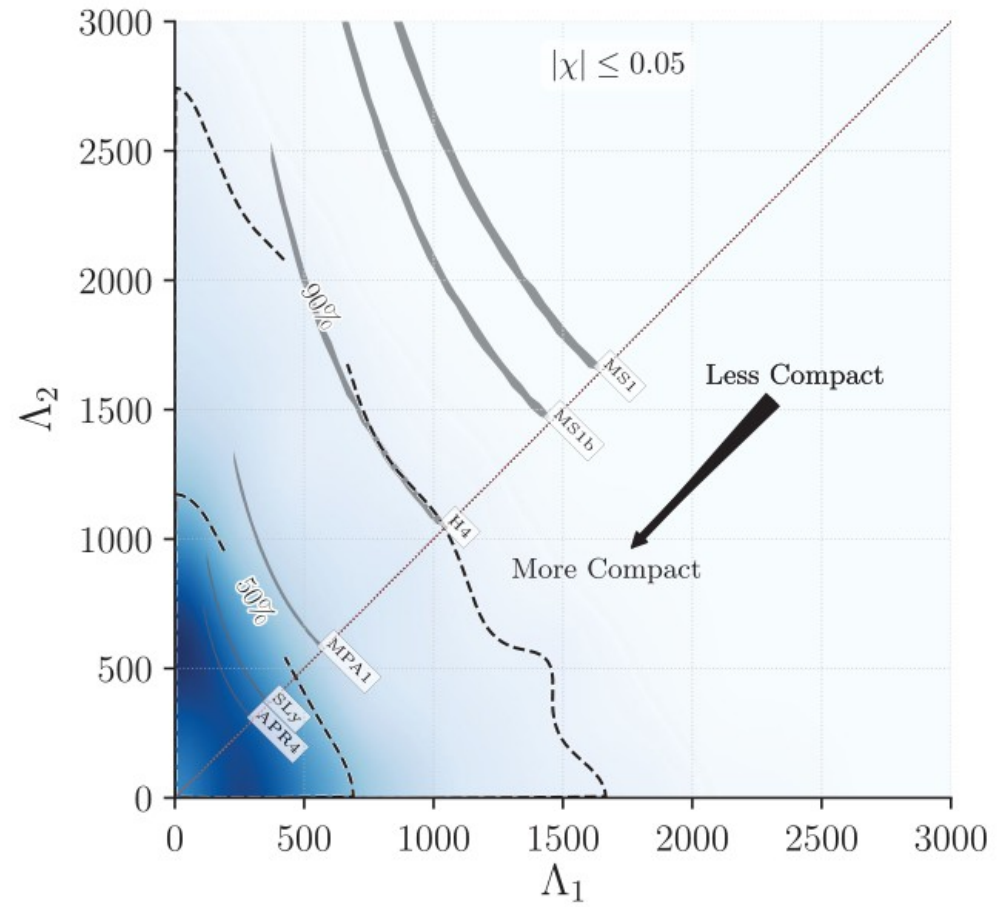
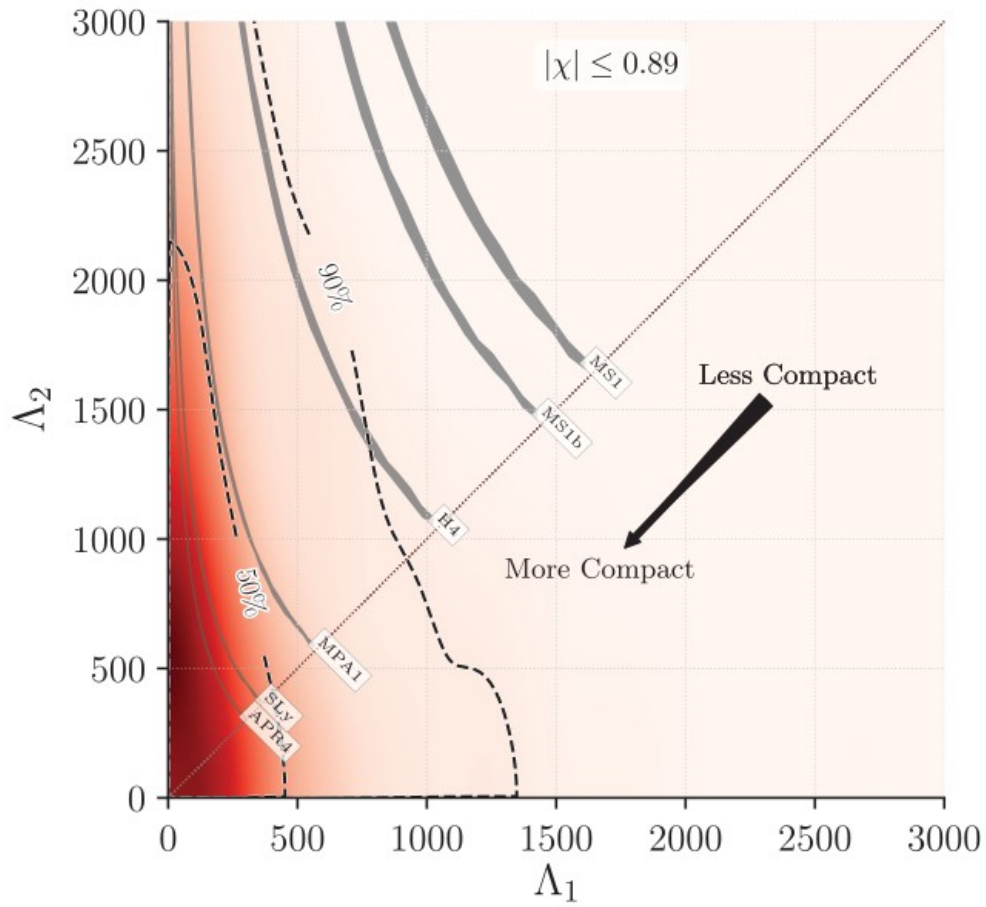


# Observations summary

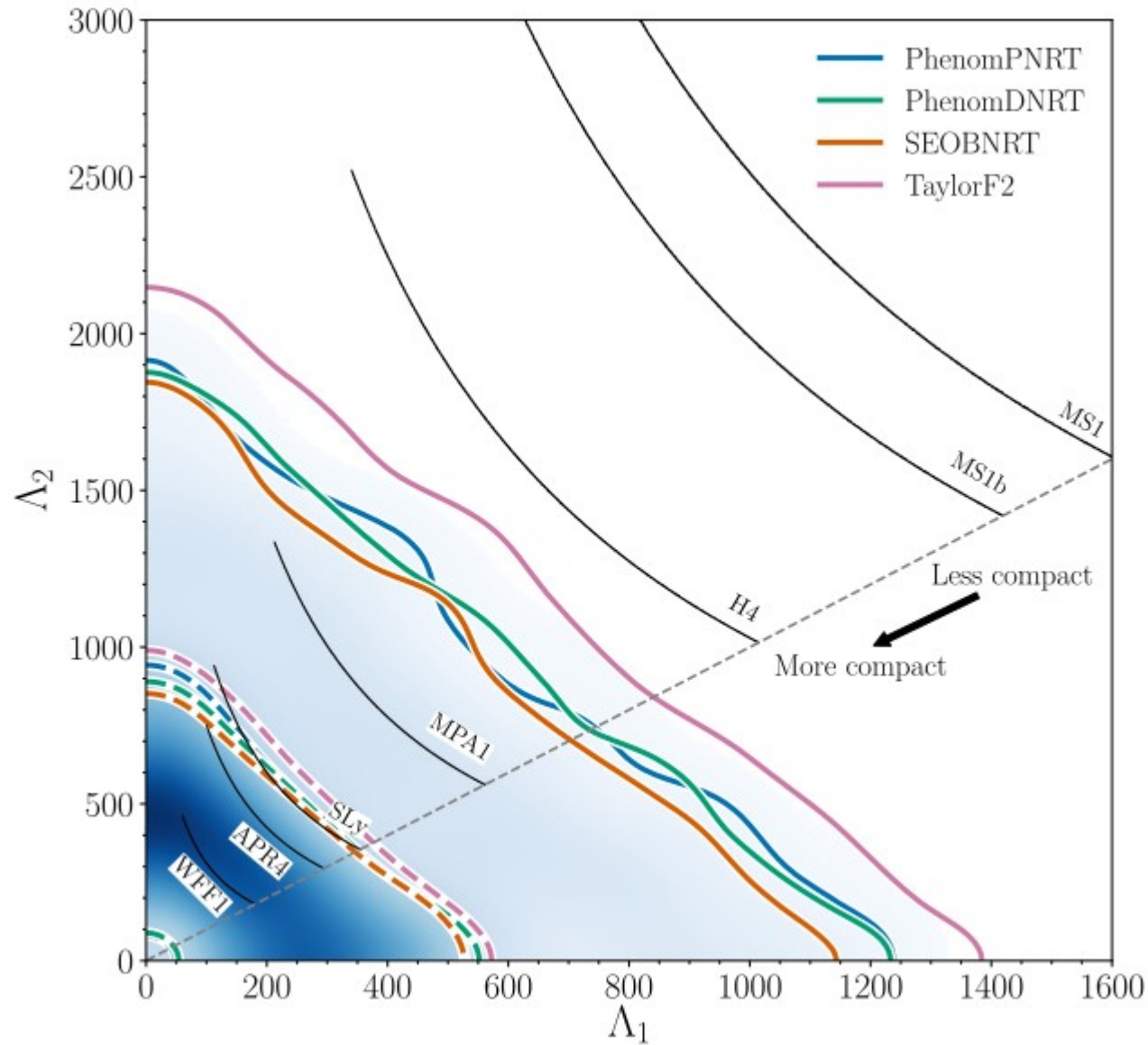
- GW observations for 30sec
- GRB at the same time
- Multiwavelength afterglow
- No VHE photons detected
- No neutrinos detected
- Masses in the NS range
- Low spins likely

# Neutron star physics

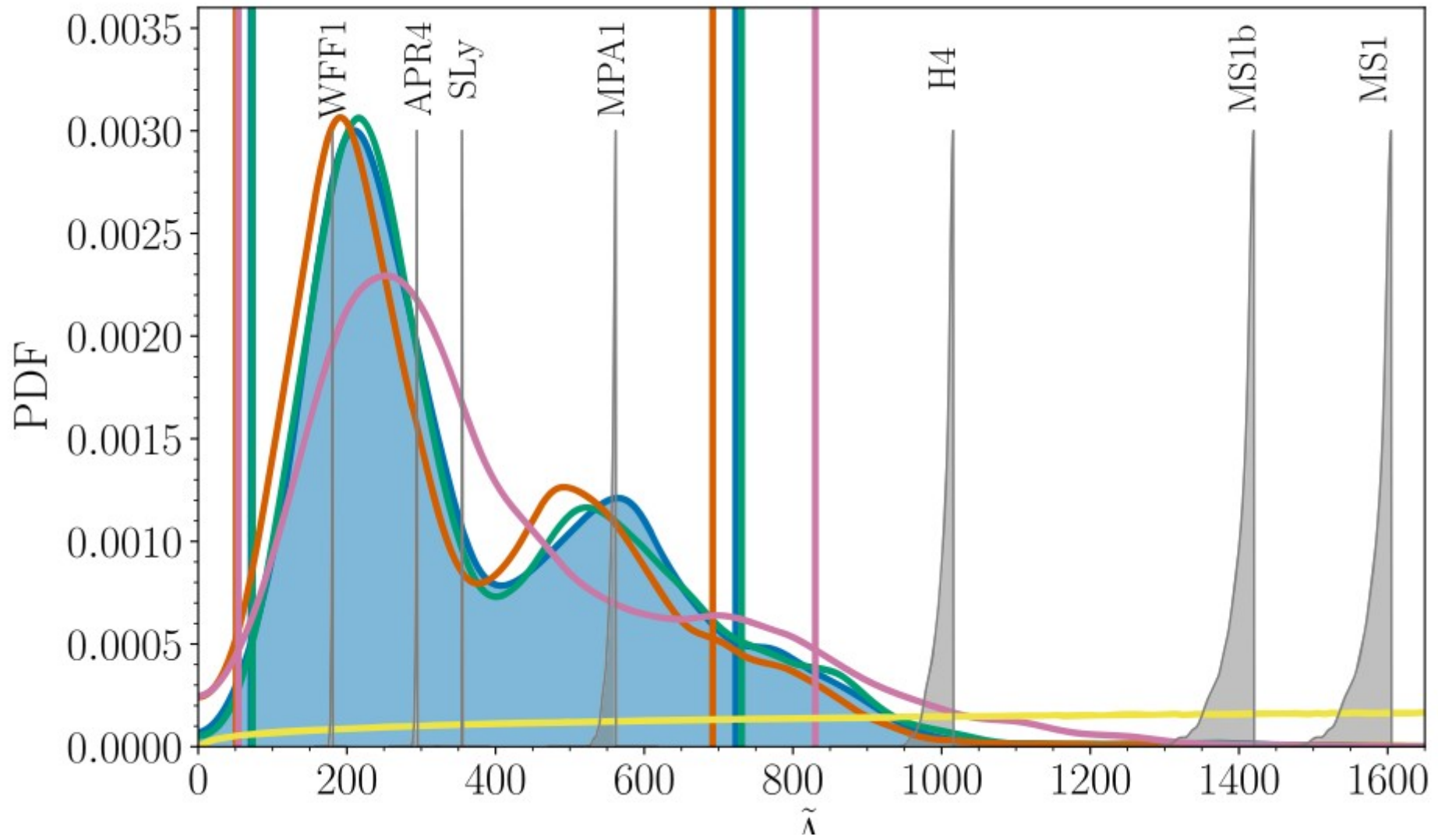
Tidal deformability - first look upper limits



# Tidal deformability - detailed analysis



# Tidal deformability



$$\tilde{\Lambda} = \frac{16(m_1 + 12m_2)m_1^4\Lambda_1 + (m_2 + 12m_1)m_2^4\Lambda_2}{(m_1 + m_2)^5}.$$

# Tidal deformability summary

- Hi spin case – upper limits
- Lo spin case – constraints

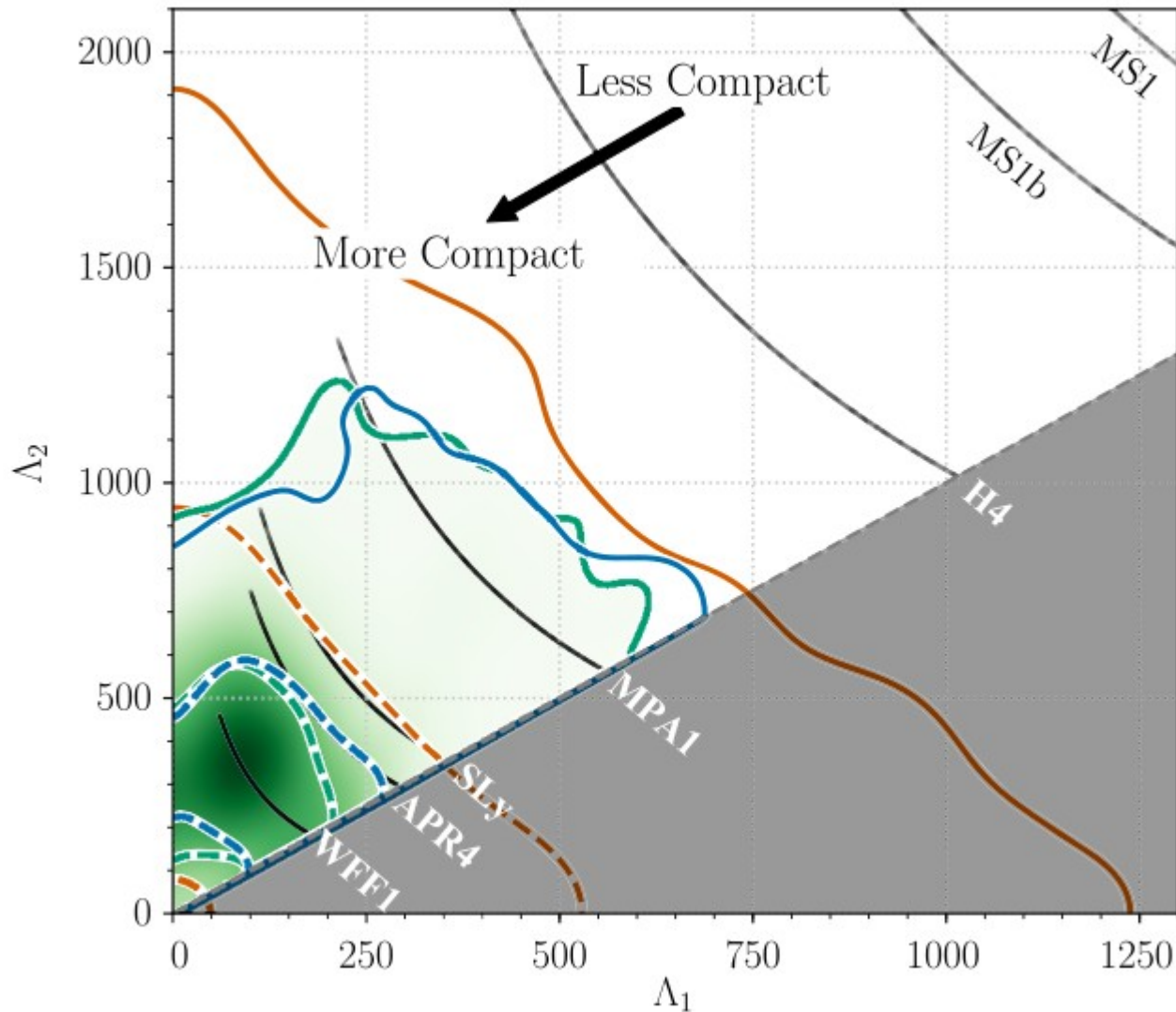
$$\Lambda = \begin{array}{ccc} 340^{+580}_{-240} / 340^{+490}_{-290} & 280^{+490}_{-190} / 280^{+410}_{-230} & 300^{+520}_{-190} / 300^{+430}_{-230} \\ 280^{+430}_{-220} / 280^{+280}_{-280} & 340^{+520}_{-260} / 340^{+350}_{-330} & 310^{+510}_{-240} / 310^{+380}_{-290} \end{array}$$

- Favors low values – compact NS

# NS radii and EOS

- Assume common EOS for two NS
- Parametrize EOS as polynomial
- Stitch to Sly Eos for low densities
- Impose causality
- Must support 1.97 $M_{\text{sun}}$  NS, or the highest mass in fit

# Constrain tidal deformability again

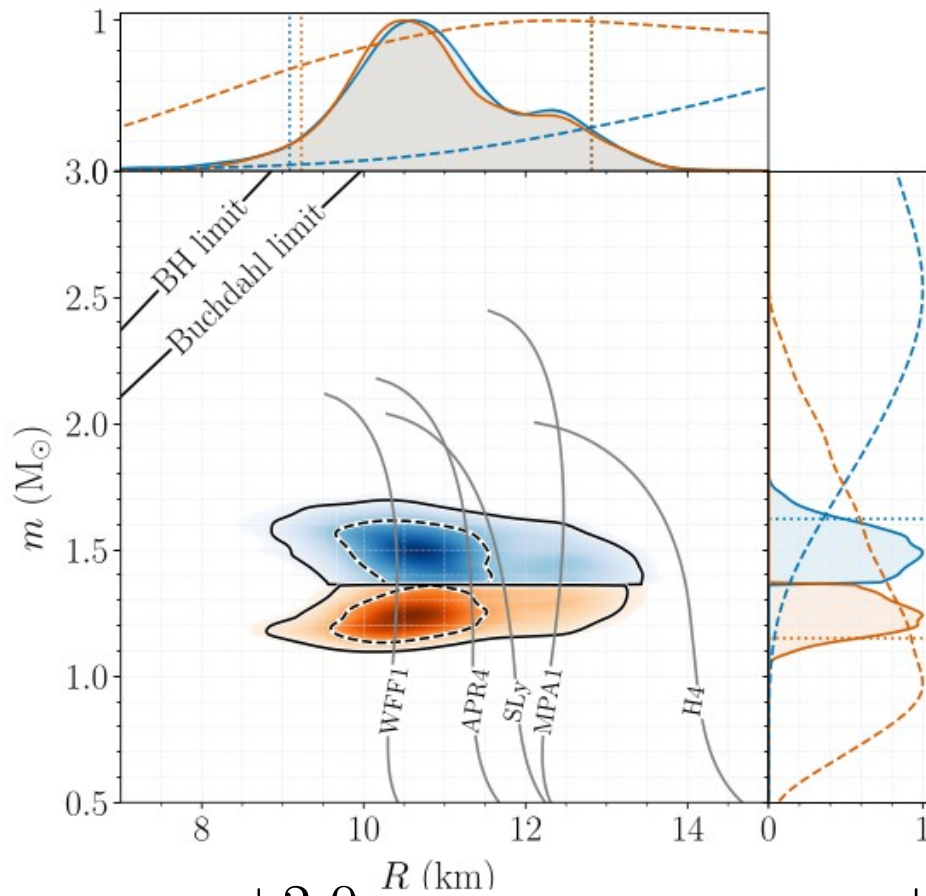


Then use the Lambda-Compactness relation to constrain the radii.



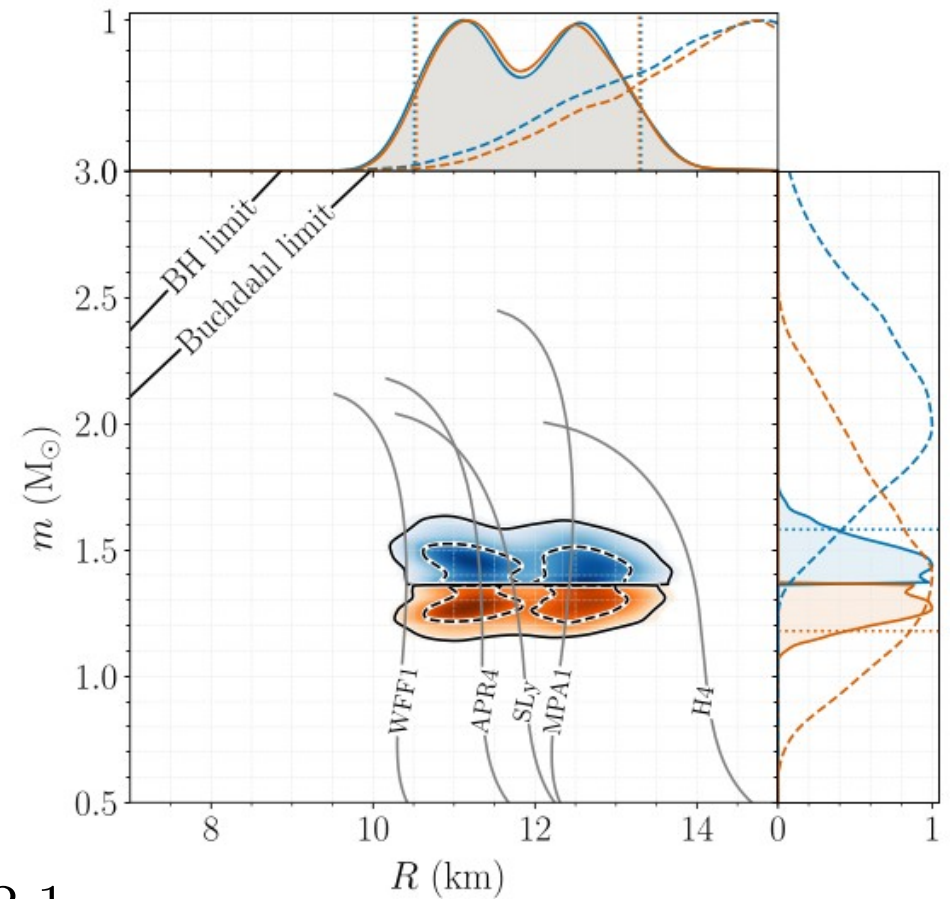
# Mass radius plot

EOS insensitive



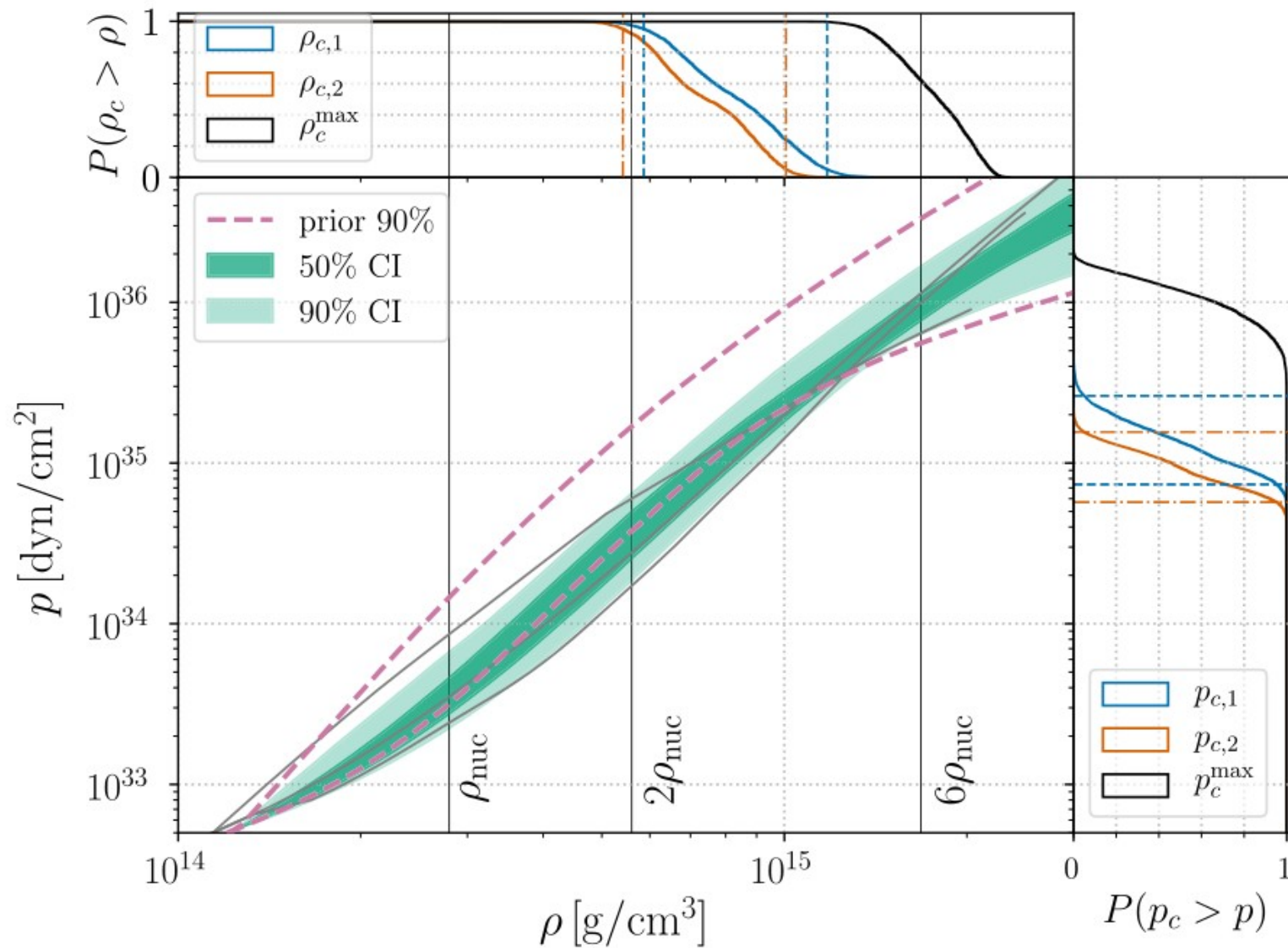
$$R_1 = 10.8^{+2.0}_{-1.7} \text{ km}, \quad R_2 = 10.7^{+2.1}_{-1.5} \text{ km}$$

Parameterized EOS



$$R_1 = 11.9^{+1.4}_{-1.4} \text{ km}, \quad R_2 = 11.9^{+1.4}_{-1.4} \text{ km}$$

# Constraints on EOS

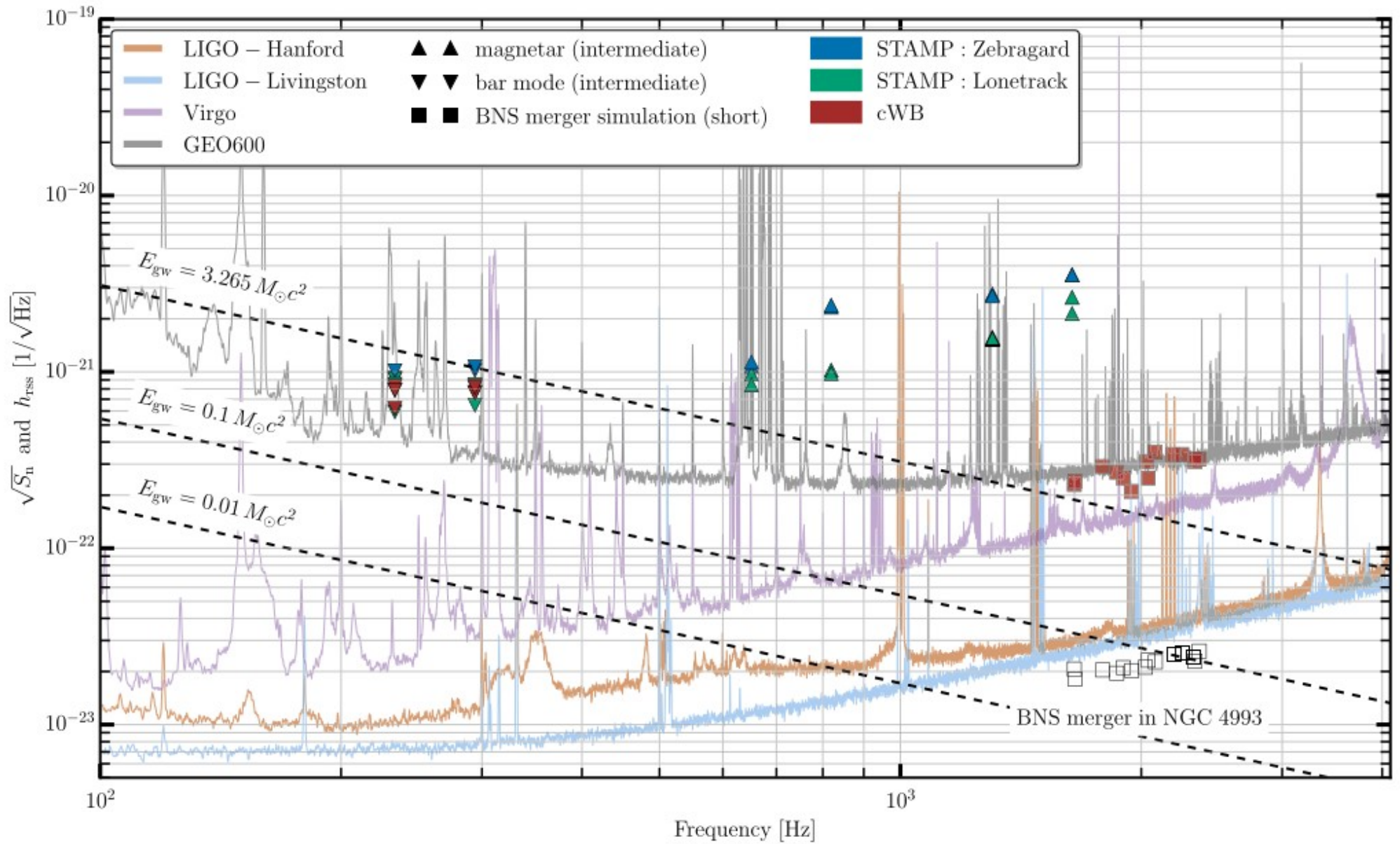


# Summary of NS -EOS properties

- NS are rather small at 1.3-1.5 Msun
- Tidal defomability is on the lower end of scale
- Yet must sustain maximal mass up to 2.14 Msun
- One can start measuring the EOS!

# Post-merger and remnant search

# Short duration search



# Short duration search

**Table 1**  
Sensitivity of the cWB Pipeline to Waveforms Generated by Binary NS Simulations

Equation of State	$m_1$ ( $M_\odot$ )	$m_2$ ( $M_\odot$ )	$\bar{f}$ (Hz)	Simulation	$h_{\text{rss}}$ (expected) ( $10^{-22}/\sqrt{\text{Hz}}$ )	$h_{\text{rss}}^{50\%}$ ( $10^{-22}/\sqrt{\text{Hz}}$ )
H4 (Glendenning & Moszkowski 1991)	1.25	1.25	1946	Takami et al. (2015)	0.21	2.1
H4 (Glendenning & Moszkowski 1991)	1.3	1.3	2083	Takami et al. (2015)	0.23	3.5
H4 (Glendenning & Moszkowski 1991)	1.35	1.35	2247	Ciolfi et al. (2017)	0.26	3.4
H4 (Glendenning & Moszkowski 1991)	1.42	1.29	2192	Ciolfi et al. (2017)	0.26	3.4
H4 (Glendenning & Moszkowski 1991)	1.54	1.26	2030	Kawamura et al. (2016)	0.22	3.1
LS220 (Lattimer & Swesty 1991)	1.20	1.50	1900	Bauswein et al. (2013a)	0.22	2.5
SHT (Shen et al. 2010)	1.40	1.40	1788	Kastaun et al. (2017)	0.21	2.9
SFHx (Steiner et al. 2013)	1.2	1.5	1650	Bauswein et al. (2013a)	0.21	2.3
SFHx (Steiner et al. 2013)	1.35	1.35	2040	Bauswein et al. (2013a)	0.24	2.5
SLy (Douchin & Haensel 2001)	1.25	1.25	2333	Takami et al. (2015)	0.23	3.2
SLy (Douchin & Haensel 2001)	1.3	1.3	2325	Takami et al. (2015)	0.25	3.1
SLy (Douchin & Haensel 2001)	1.35	1.25	2363	Takami et al. (2015)	0.27	3.2
TMA (Toki et al. 1995)	1.20	1.50	1864	Bauswein et al. (2013a)	0.19	3.2
TMA (Toki et al. 1995)	1.35	1.35	1653	Bauswein et al. (2013a)	0.20	2.4

**Note.** Waveforms were selected to represent a variety of equations of state (first column) and progenitor mass configurations (second and third columns). The fourth column is the mean frequency for each waveform  $\bar{f}$ , and the fifth column is the reference for the BNS simulation. The sixth column is the root-sum-squared strain  $h_{\text{rss}}$  predicted by that simulation for a post-merger signal from a BNS with a distance and inclination consistent with estimates from the inspiral analysis (Abbott et al. 2017a). The seventh column shows the  $h_{\text{rss}}$  required for 50% detection efficiency with a false-alarm probability of  $10^{-4}$ ,  $h_{\text{rss}}^{50\%}$ .

# Long duration search

**Table 5**

FreqHough Search Sensitivities Estimated from Simulated Signals (Injections) Following the Power-law Spin-down Model with Braking Index  $n = 5$

$T_{\text{SFT}}$ (s)	$f_{\text{start}}$ (Hz)	$\tau$ (s)	$\epsilon$	$d^{90\%}$ (Mpc)	$E_{\text{gw}}^{90\%} (M_{\odot}c^2)$
2	390	$4.15 \times 10^2$	$5.91 \times 10^{-2}$	$0.78_{-0.04}^{+0.05}$	$4.76_{-0.50}^{+0.59} \times 10^0$
2	440	$4.65 \times 10^2$	$4.39 \times 10^{-2}$	$0.83_{-0.05}^{+0.06}$	$5.42_{-0.62}^{+0.75} \times 10^0$
2	490	$5.15 \times 10^2$	$3.36 \times 10^{-2}$	$0.79_{-0.05}^{+0.06}$	$7.39_{-0.89}^{+1.08} \times 10^0$
2	540	$5.65 \times 10^2$	$2.64 \times 10^{-2}$	$0.72_{-0.04}^{+0.05}$	$1.06_{-0.12}^{+0.15} \times 10^1$
2	590	$6.15 \times 10^2$	$2.12 \times 10^{-2}$	$0.75_{-0.05}^{+0.06}$	$1.19_{-0.16}^{+0.20} \times 10^1$
...	...	...	...	...	...

**Note.** Each row corresponds to injections marginalized over a 50 Hz band in  $f_{\text{start}}$  (where  $f_{\text{start}} = f_{\text{gw}}(t_{\text{start}} = t_c + \Delta t)$ ) and random  $\cos \iota$ ; sensitivities are at 90% confidence.

# Post merger summary

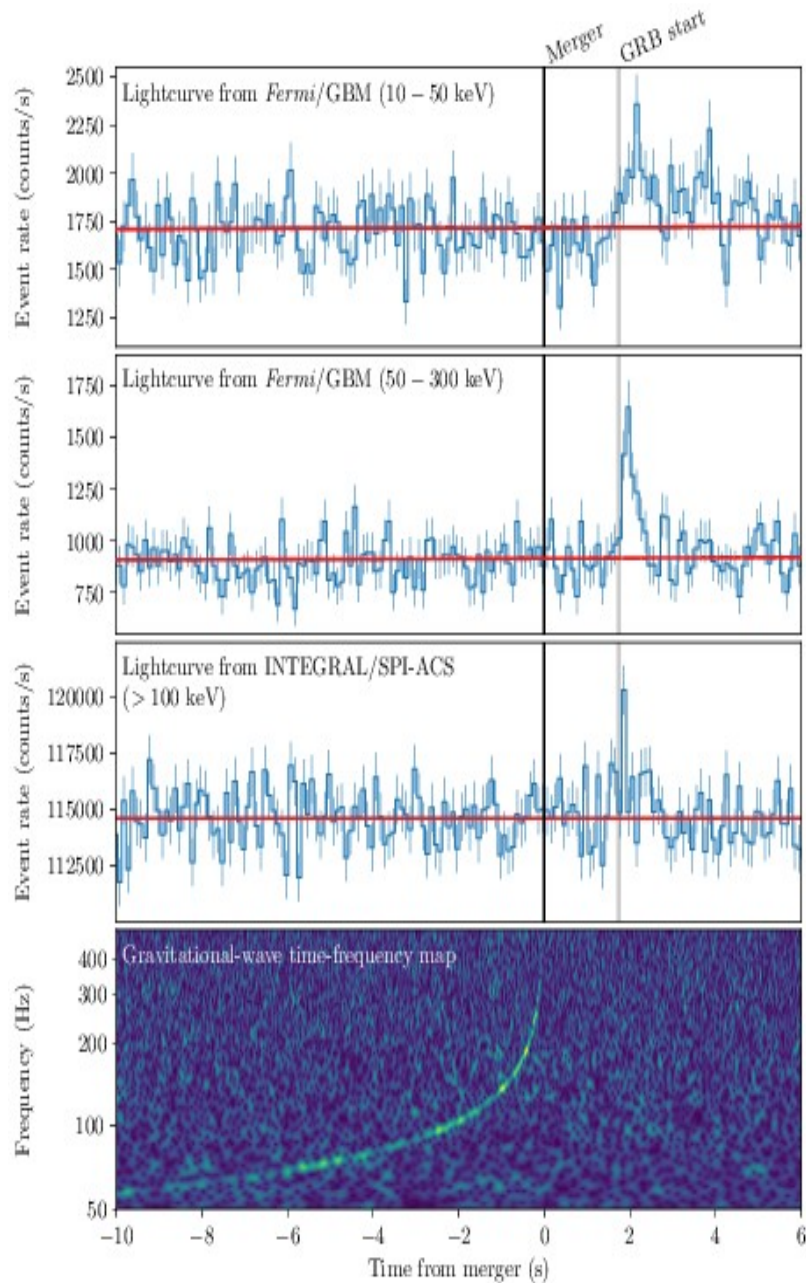
- Search for short duration -yields no results
- Sensitivity – 10 times too low
- ET, Cosmic Explorer will be sensitive enough



# Fundamental physics

- Speed of gravity
- Dispersion relation
- Dimension of space
- GR validity
  - Motion of two bodies, PN corrections
  - GW polarizations

# Speed of gravity



Time delay - 1.7 s, let us assume it is less than 10s

Distance 40 Mpc =  $4.10 \times 10^{15}$  light s, let us assume a lower limit of 26Mpc

Relative difference of speed

$$-3 \times 10^{-15} < \frac{\delta c_g}{c} < 7 \times 10^{-15}$$

# Dispersion relation

- Does speed of GW depend on frequency?

$$E^2 = p^2 c^2 + A p^\alpha c^\alpha$$

$$m_g < 9.51 \times 10^{-22} \text{eV}$$

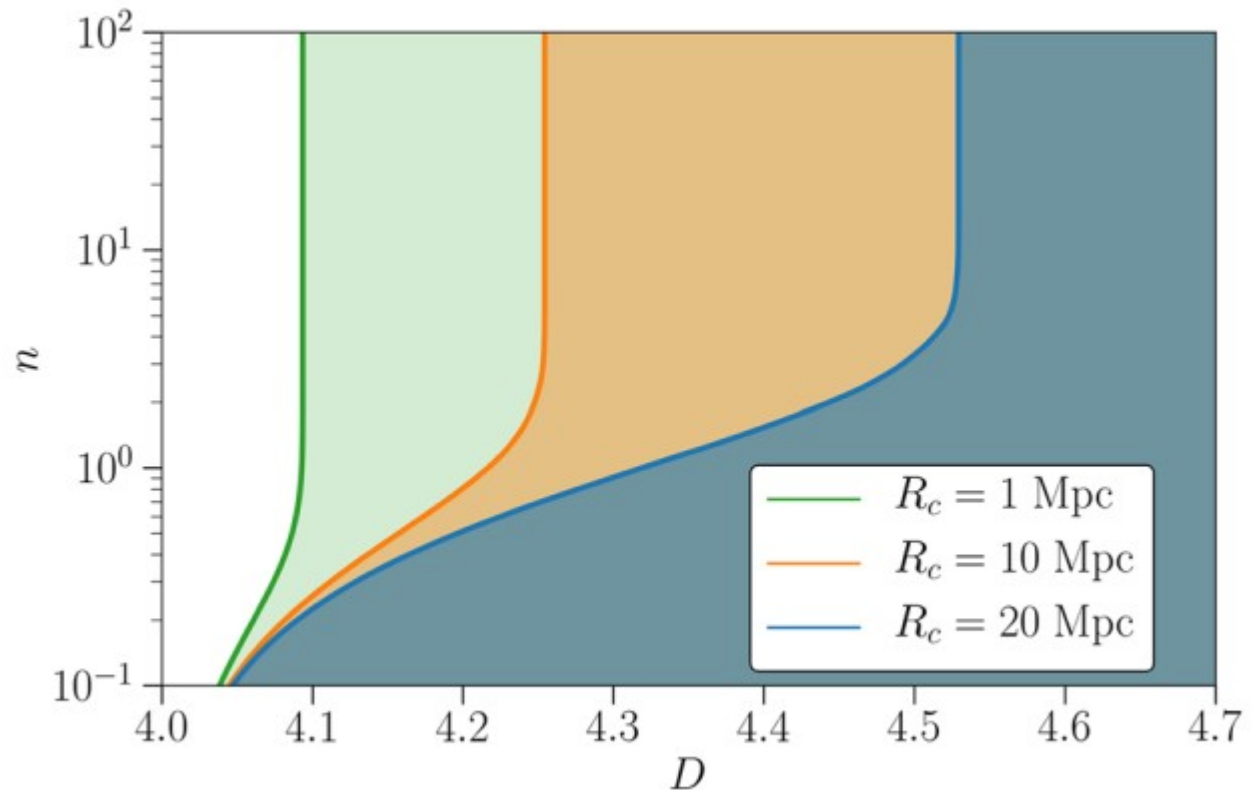
- Limits less stringent than for BBH sources.

# Space dimensions

- Leakage of GW energy to other dimensions

$$h \propto \frac{1}{d_L^{GW}} = \frac{1}{d_L^{EM}} \left[ 1 + \left( \frac{d_L^{EM}}{R_c} \right)^n \right]^{-(D-4)/2n}$$

- Difference in estimate of the luminosity distance



# GR validity

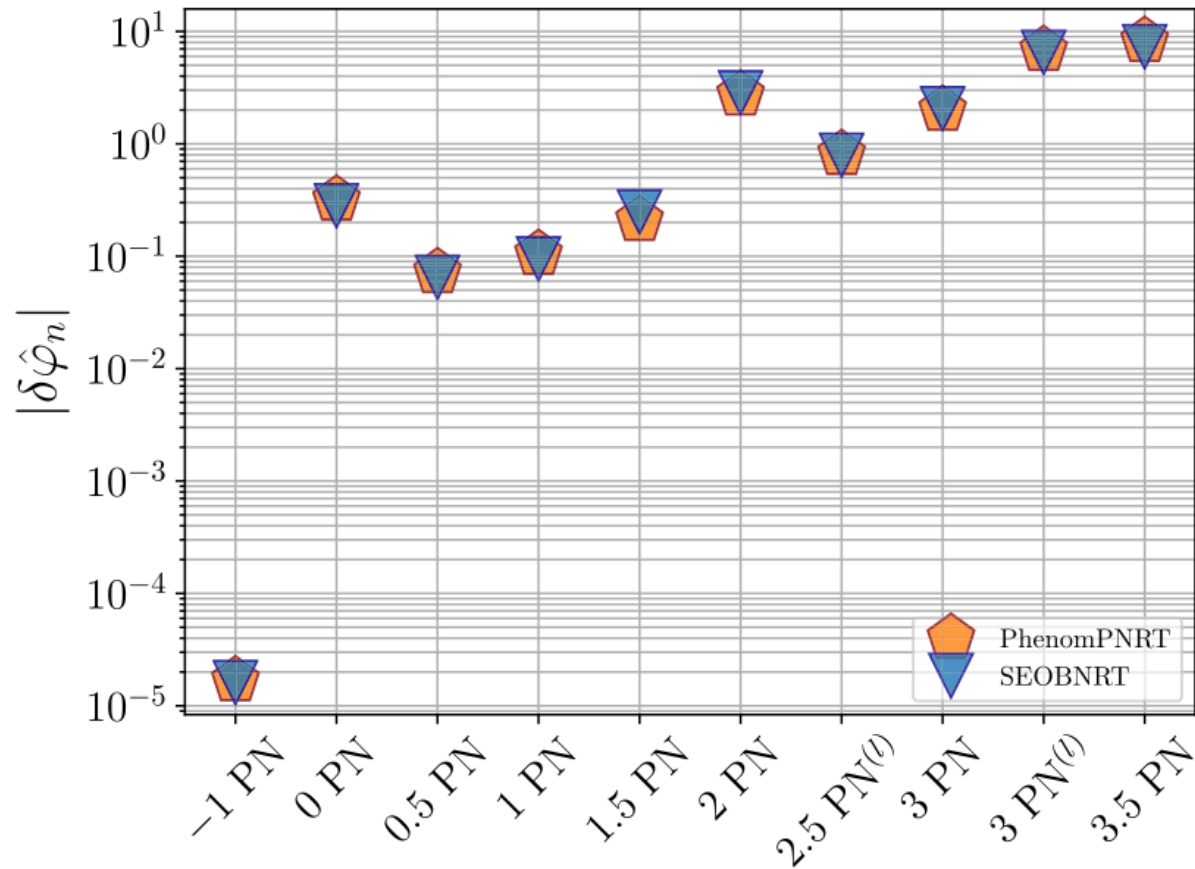


FIG. 2. 90% upper bounds on deviations  $|\delta\hat{\varphi}_n|$  in the PN coefficients following from the posterior density functions shown in Fig. 1.

# GW polarizations

- GR predicts only two polarizations
- Search for other modes
- Very strong Bayesian preference for tensor modes  $\rightarrow$  GR is valid.

# Summary of fundamental physics

- GR works
- Very stringent constraints on deviations from GR
- Future observations may improve these by a factor of ten

# Astrophysics

- Rates
- Progenitors
- GRB science
- Heavy elements production
- Cosmology



# The merger rate densities

- BBH estimate  $R = 9.7 - 101 \text{Gpc}^{-3} \text{yr}^{-1}$
- BNS estimate  $R = 110 - 3840 \text{Gpc}^{-3} \text{yr}^{-1}$
- The local supernova rate  $\sim 10^5 \text{Gpc}^{-3} \text{yr}^{-1}$
- The BH formation rate is  $\sim 10^4 \text{Gpc}^{-3} \text{yr}^{-1}$
- About 1 black hole in a 100-1000 ends up in a merging binary
- Similarly NS: 1 in 100-1000 is in a merging binary!

## The rate implications

- Total GW luminosity density in the sky from NSNS mergers

$$\mathcal{L}_{GW} = 1500 \frac{0.025 M_{\odot} c^2}{3.1 \times 10^7 \text{s}} \approx 2.5 \times 10^{48} \text{ergs}^{-1} \text{Gpc}^{-3}$$

- The luminosity density of BHBH mergers is about 10 times larger

$$\mathcal{L}_{GW} = 50 \frac{2.0 M_{\odot} c^2}{3.1 \times 10^7 \text{s}} \approx 10^{49} \text{ergs}^{-1} \text{Gpc}^{-3}$$

- EM luminosity density of all galaxies:

$$\mathcal{L}_{EM} \approx 10^{50} \text{erg s}^{-1} \text{Gpc}^{-3}$$

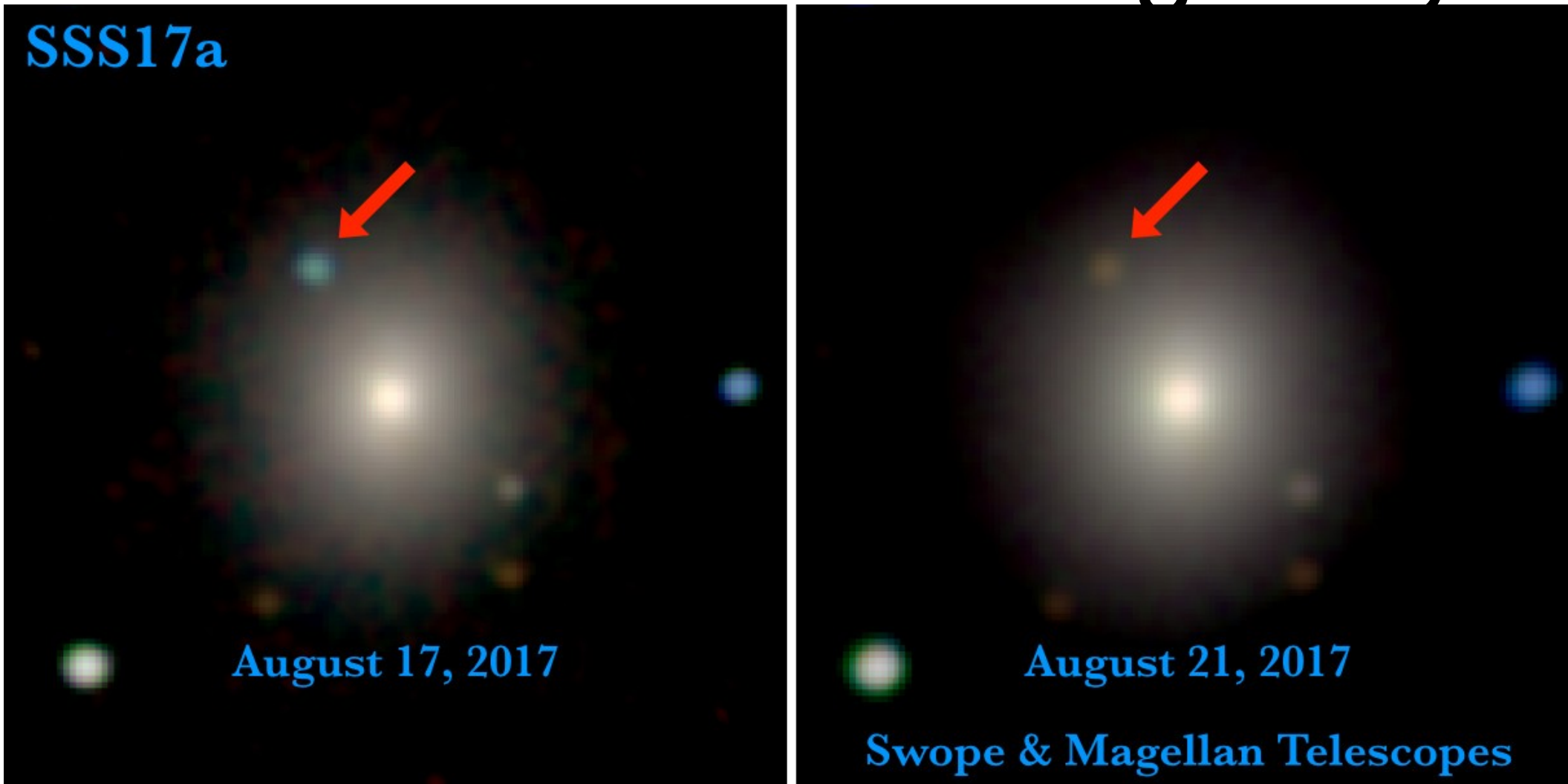
# Basic rate arguments

- Formation scenario must be generic
- Exceptional environments must produce BBH and BNS with extremely high efficiency
- Globular clusters are not favoured, but can contribute
- I am sceptical about exotic models

# What options do we have?

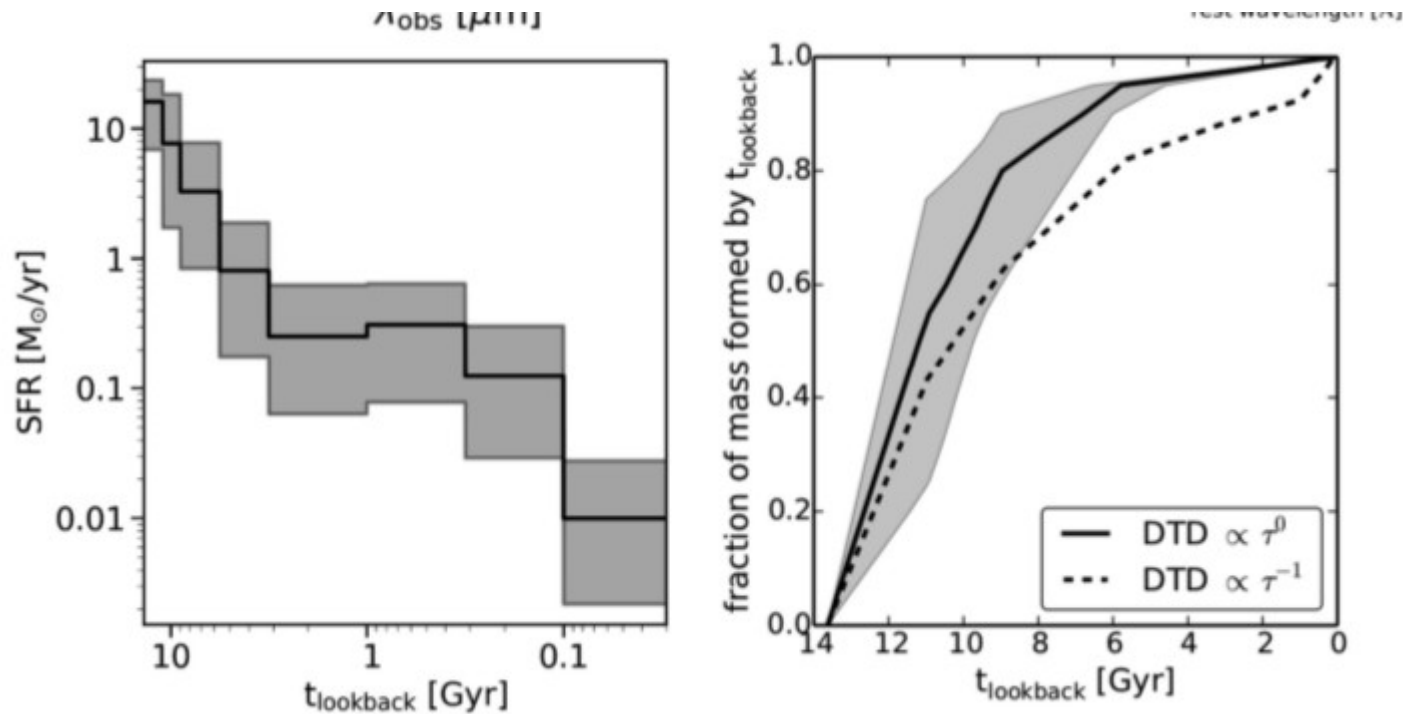
- Binary evolution
  - Standard
  - Chemically homogenous case
- Evolution in the clusters
- Exotica
  - Pop III stars
  - Exceptional environments.

# BNS: all that + host galaxy



NGC 4993 – old elliptical with no traces of str formation for th last 1-2Gyrs,  
merger on the ourskirts of the galaxy.

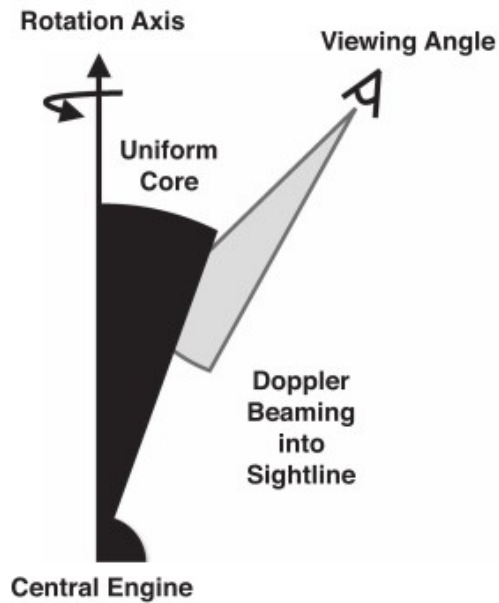
# Star formation history estimate



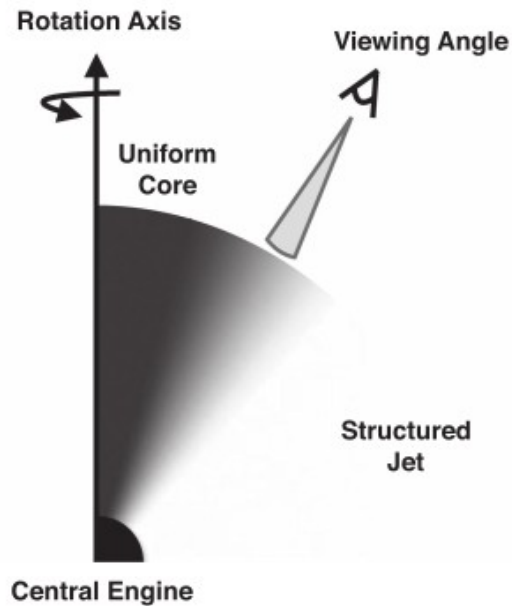
Is there something missing about formation of such systems?

# GRB science

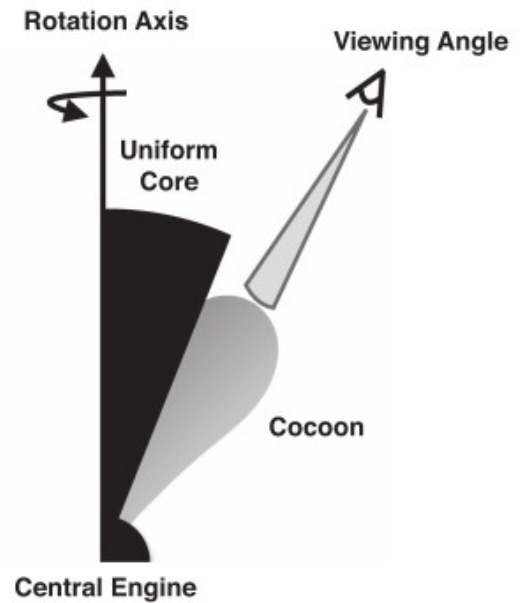
Scenario i: Uniform Top-hat Jet



Scenario ii: Structured Jet

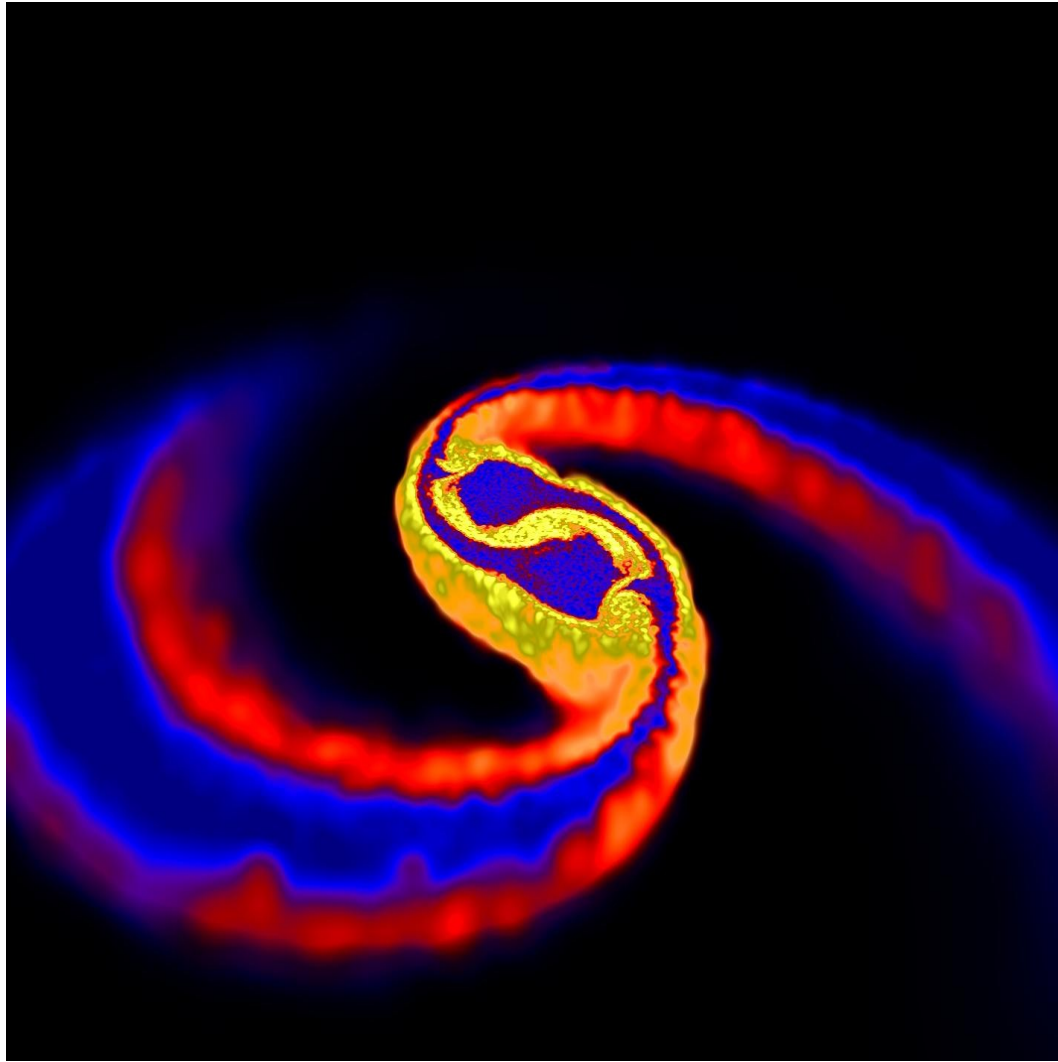


Scenario iii: Uniform Jet + Cocoon



**Figure 5.** Three potential jet viewing geometries and jet profiles that could explain the observed properties of GRB 170817A, as described by scenarios (i)–(iii) in Section 6.2.

# Heavy element production

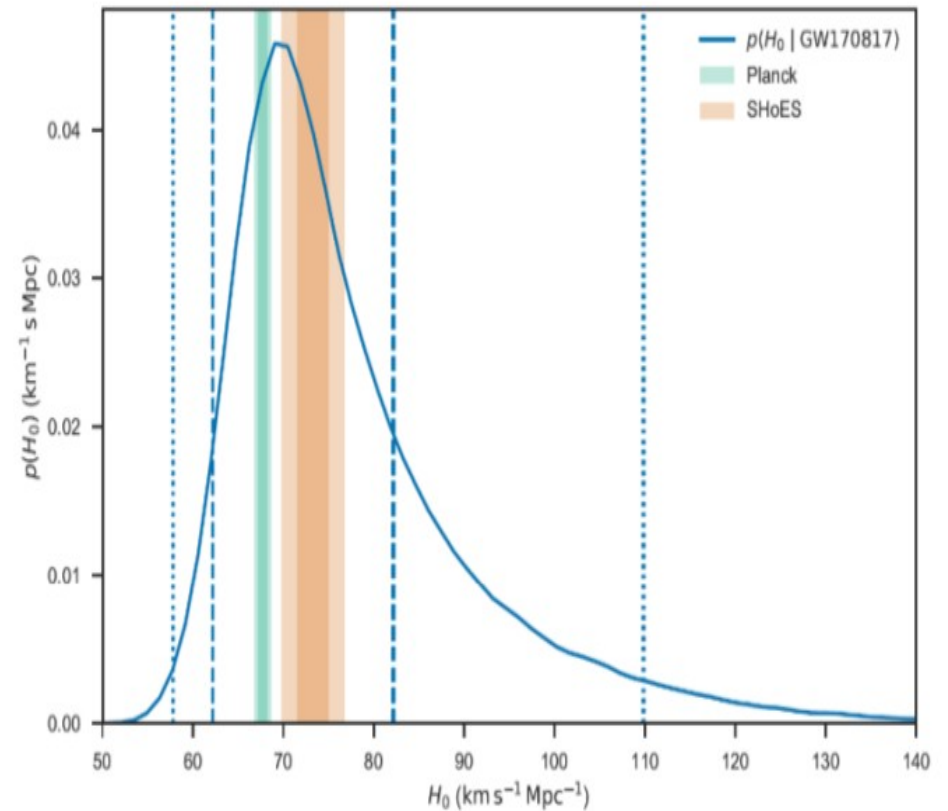




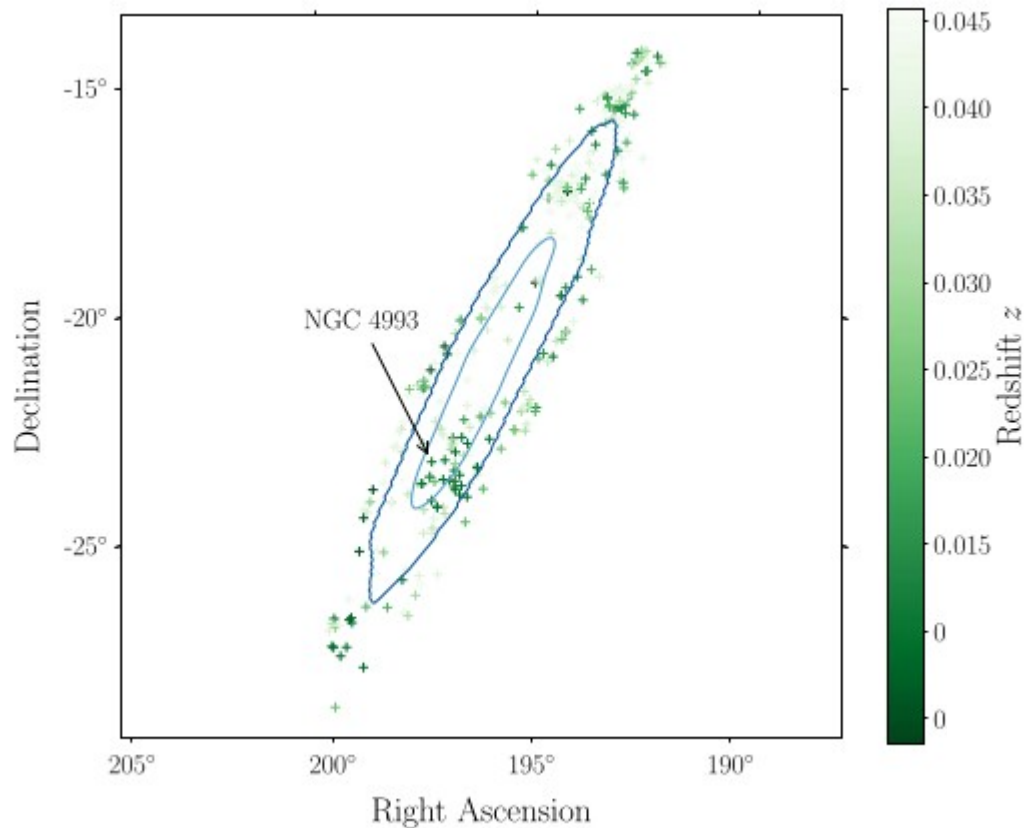
# Cosmology

Binary as a standard siren:  
Location and orientation needed  
LHV: detection, location  
From GW: distance  
Optical detection: redshift  
Together – Hubble constant !!!

$$H_0 = 70_{-8}^{+12} \text{ km s}^{-1} \text{ Mpc}^{-1}$$

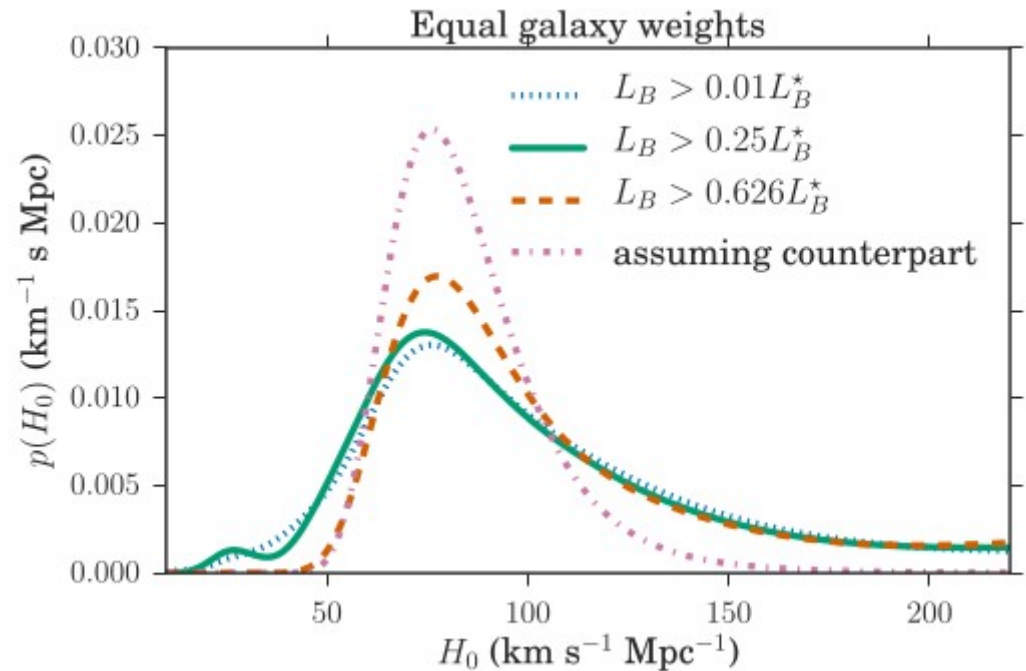


# Hubble without EM

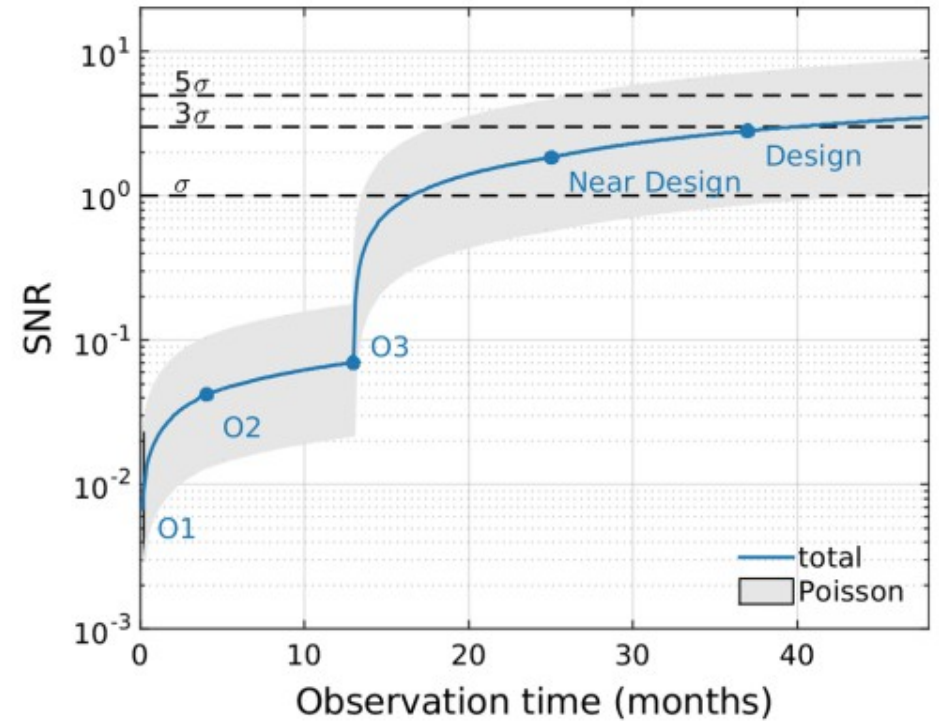
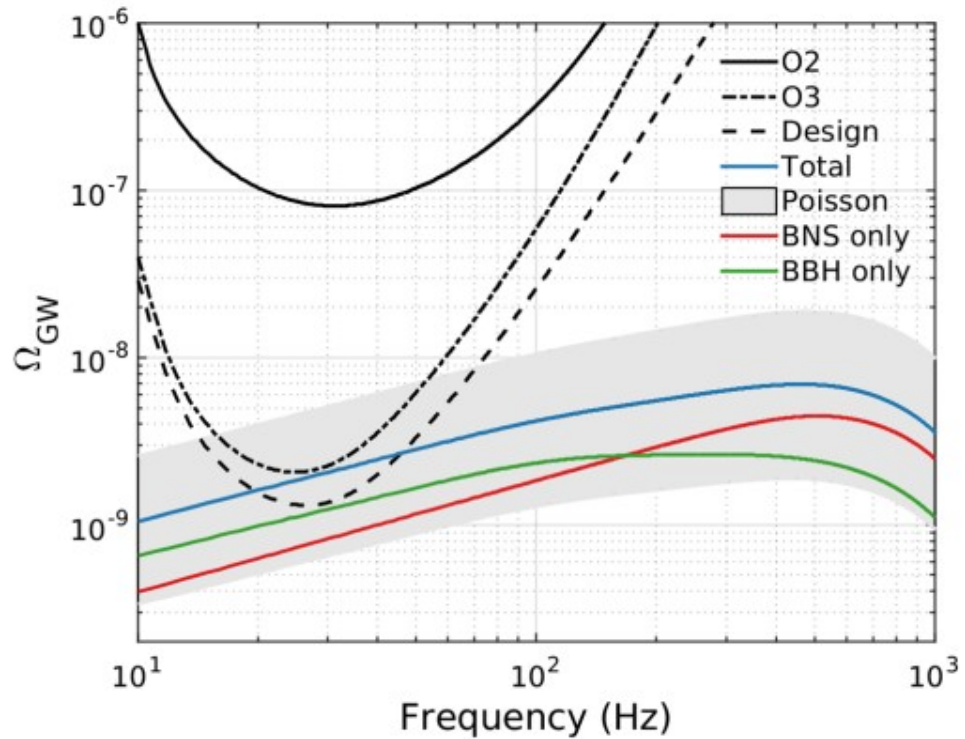


Great potential for future observations.

May solve the Hubble constant problem,



# GW background



# Astrophysics summary

- A number of astrophysical uses
- Origin – binary evolution but other options may work too
- Cosmology – need hundreds of sources
- High redshift mergers – through GW background
- Chemical evolution
- Connection with GRBs finally settled for good

# When can we expect a next GW10817 event?

Assuming the rate is  $500 \text{ Gpc}^{-1} \text{ yr}^{-1}$

We get one event per 4 years within 50 Mpc

And 2 events per year within 100 Mpc

Detectability depends on uptime. Assuming 3 detector time of 50% for O4 (two years) , we should see 2 events within 100Mpc, and  $\frac{1}{4}$  of event within 50Mpc.

GW170817 was very unique!

# Summary

- Neutron stars are small
- EOS is .....
- GW works
- GW binaries are great tools for astrophysics



