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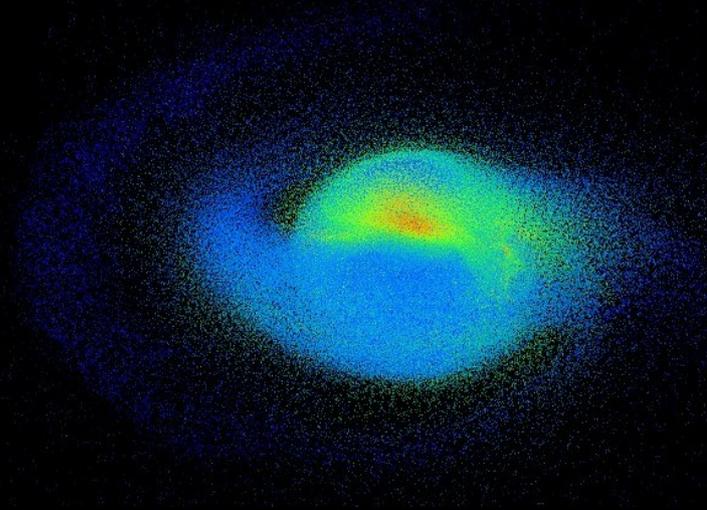
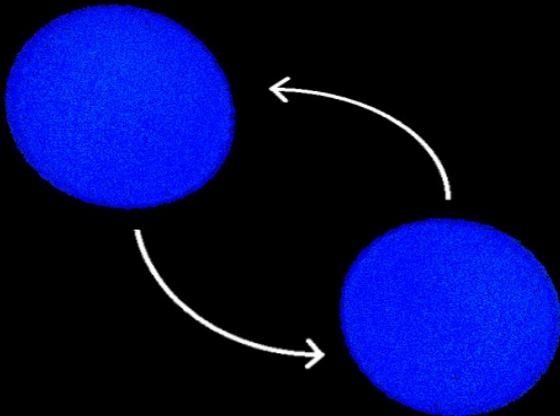
Neutron star mergers and matter under extreme conditions

The first compact star merger event – implications for nuclear and particle physics

ECT* Trento, 17/10/2019

Andreas Bauswein

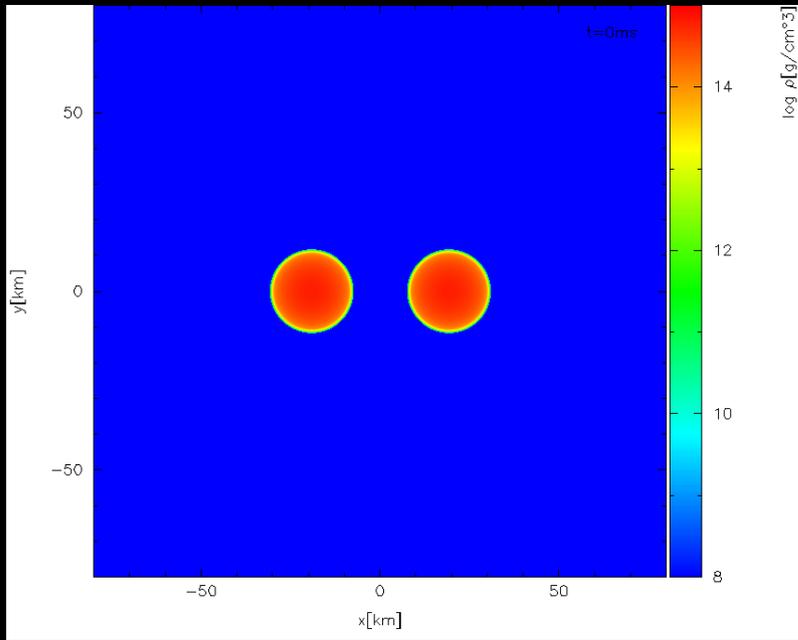
(GSI Darmstadt)



Outline

- ▶ R-process in NS mergers
- ▶ Multi-messenger interpretation of GW170817 → lower limit on NS radii
→ Collapse behavior (EoS dependence of BH formation)
- ▶ Postmerger GW emission
- ▶ Signatures of the QCD phase transition

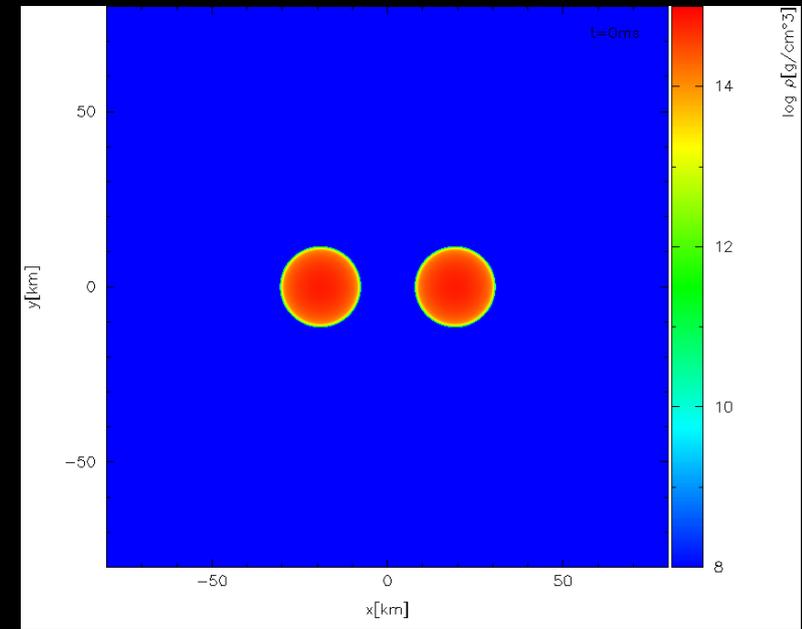
Collapse behavior and multi-messenger EoS constraints



$$M_{\text{tot}} = 3.4 M_{\odot}$$



$$M_{\text{tot}} = 3.5 M_{\odot}$$



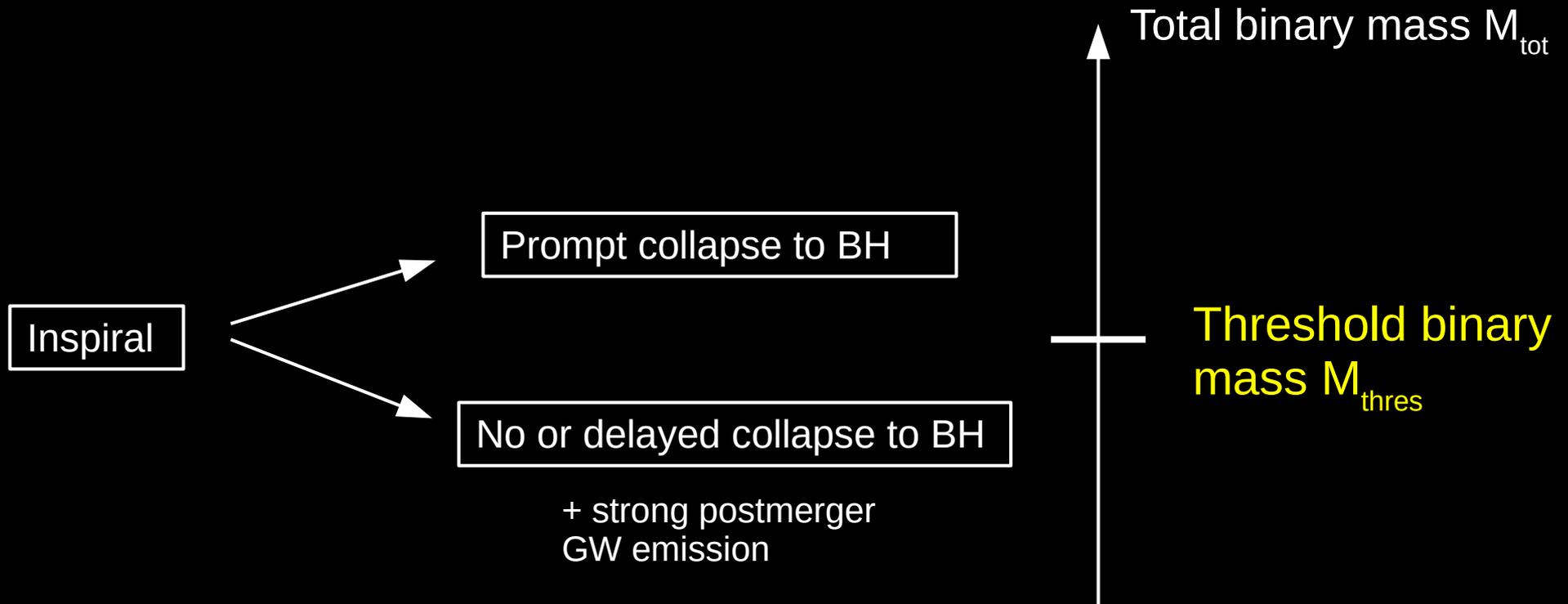
Shen EoS

$$\longrightarrow M_{\text{thres}} = (3.45 \pm 0.05) M_{\odot} \quad (\text{for this particular EoS})$$

Collapse behavior: Prompt vs. delayed (/no) BH formation

Relevant for: EoS constraints through M_{max} measurement, Conditions for short GRBs, Mass ejection, Electromagnetic counterparts powered by thermal emission, NS radius constraints !!!

Collapse behavior



EoS dependent - somehow M_{max} should play a role

Simulations reveal M_{thres}

TOV properties of nonrotating stars, i.e. EoS characteristics

Merger property from simulations

| EoS | M_{max} (M_{\odot}) | R_{max} (km) | C_{max} | $R_{1.6}$ (km) | M_{thres} (M_{\odot}) |
|-------------|-------------------------------------|--------------------------|------------------|-------------------|---------------------------------------|
| NL3 [37,38] | 2.79 | 13.43 | 0.307 | 14.81 | 3.85 |
| GS1 [39] | 2.75 | 13.27 | 0.306 | 14.79 | 3.85 |
| LS375 [40] | 2.71 | 12.34 | 0.325 | 13.71 | 3.65 |
| DD2 [38,41] | 2.42 | 11.90 | 0.300 | 13.26 | 3.35 |
| Shen [42] | 2.22 | 13.12 | 0.250 | 14.46 | 3.45 |
| TM1 [43,44] | 2.21 | 12.57 | 0.260 | 14.36 | 3.45 |
| SFHX [45] | 2.13 | 10.76 | 0.292 | 11.98 | 3.05 |
| GS2 [46] | 2.09 | 11.78 | 0.262 | 13.31 | 3.25 |
| SFHO [45] | 2.06 | 10.32 | 0.294 | 11.76 | 2.95 |
| LS220 [40] | 2.04 | 10.62 | 0.284 | 12.43 | 3.05 |
| TMA [44,47] | 2.02 | 12.09 | 0.247 | 13.73 | 3.25 |
| IUF [38,48] | 1.95 | 11.31 | 0.255 | 12.57 | 3.05 |

Bauswein et al. 2013

Smooth particle hydrodynamics + conformal flatness

Threshold binary mass

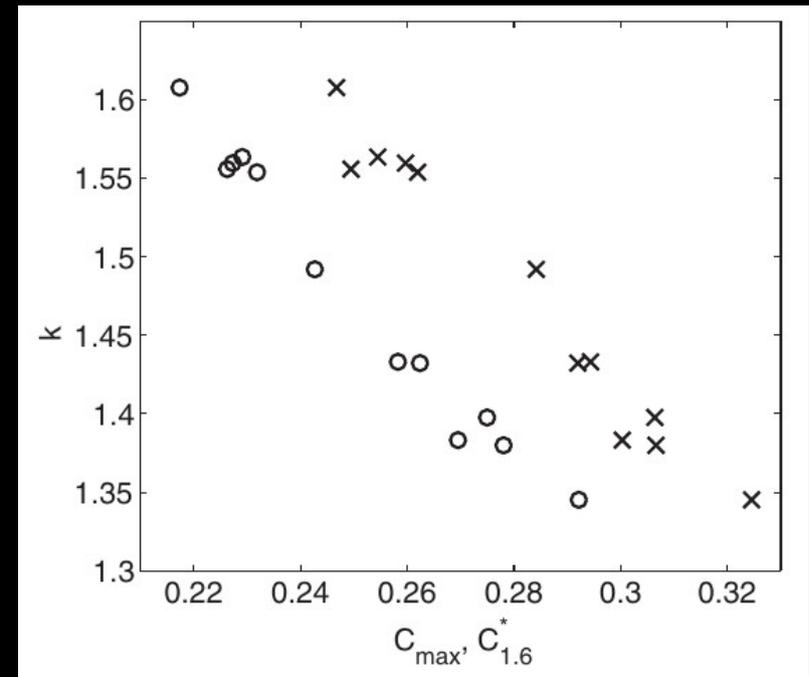
- ▶ Empirical relation from simulations with different M_{tot} and EoS
- ▶ Fits (to good accuracy):

$$M_{\text{thres}} = M_{\text{thres}}(M_{\text{max}}, R_{\text{max}}) = \left(-3.38 \frac{GM_{\text{max}}}{c^2 R_{\text{max}}} + 2.43 \right) M_{\text{max}}$$

$$M_{\text{thres}} = M_{\text{thres}}(M_{\text{max}}, R_{1.6}) = \left(-3.6 \frac{GM_{\text{max}}}{c^2 R_{1.6}} + 2.38 \right) M_{\text{max}}$$

- ▶ Both better than $0.06 M_{\text{sun}}$,

$$k = M_{\text{thres}}/M_{\text{max}}$$



EoS constraints from GW170817*

→ lower bound on NS radii

(recall: upper bound from tidal deformability)

* See also Margalit & Metzger 2017, Shibata et al. 2017, Radice et al. 2018, Rezzolla et al. 2018, Ruiz & Shapiro 2018, Capano et al 2019,... for other EoS constraints in the context of GW170817

A simple but robust NS radius constraint from GW170817

- ▶ High ejecta mass inferred from electromagnetic transient (high compared to simulations)
 - provides strong support for a delayed/no collapse in GW170817
 - even asymmetric mergers that directly collapse do not produce such massive ejecta

| Reference | $m_{\text{dyn}} [M_{\odot}]$ | $m_{\text{w}} [M_{\odot}]$ |
|-----------------------------|------------------------------|----------------------------|
| Abbott et al. (2017a) | 0.001 – 0.01 | – |
| Arcavi et al. (2017) | – | 0.02 – 0.025 |
| Cowperthwaite et al. (2017) | 0.04 | 0.01 |
| Chornock et al. (2017) | 0.035 | 0.02 |
| Evans et al. (2017) | 0.002 – 0.03 | 0.03 – 0.1 |
| Kasen et al. (2017) | 0.04 | 0.025 |
| Kasliwal et al. (2017b) | > 0.02 | > 0.03 |
| Nicholl et al. (2017) | 0.03 | – |
| Perego et al. (2017) | 0.005 – 0.01 | 10^{-5} – 0.024 |
| Rosswog et al. (2017) | 0.01 | 0.03 |
| Smartt et al. (2017) | 0.03 – 0.05 | 0.018 |
| Tanaka et al. (2017a) | 0.01 | 0.03 |
| Tanvir et al. (2017) | 0.002 – 0.01 | 0.015 |
| Troja et al. (2017) | 0.001 – 0.01 | 0.015 – 0.03 |

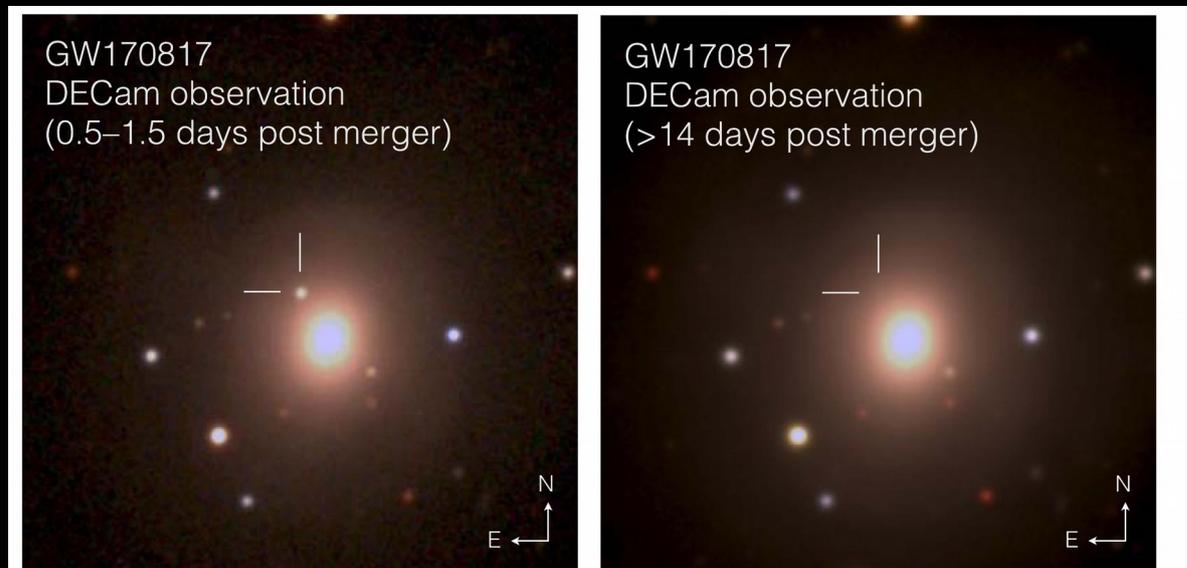


Figure 1. NGC4993 *grz* color composites ($1'5 \times 1'5$). Left: composite of detection images, including the discovery *z* image taken on 2017 August 18 00:05:23 UT and the *g* and *r* images taken 1 day later; the optical counterpart of GW170817 is at R.A., decl. = 197.450374, -23.381495 . Right: the same area two weeks later.

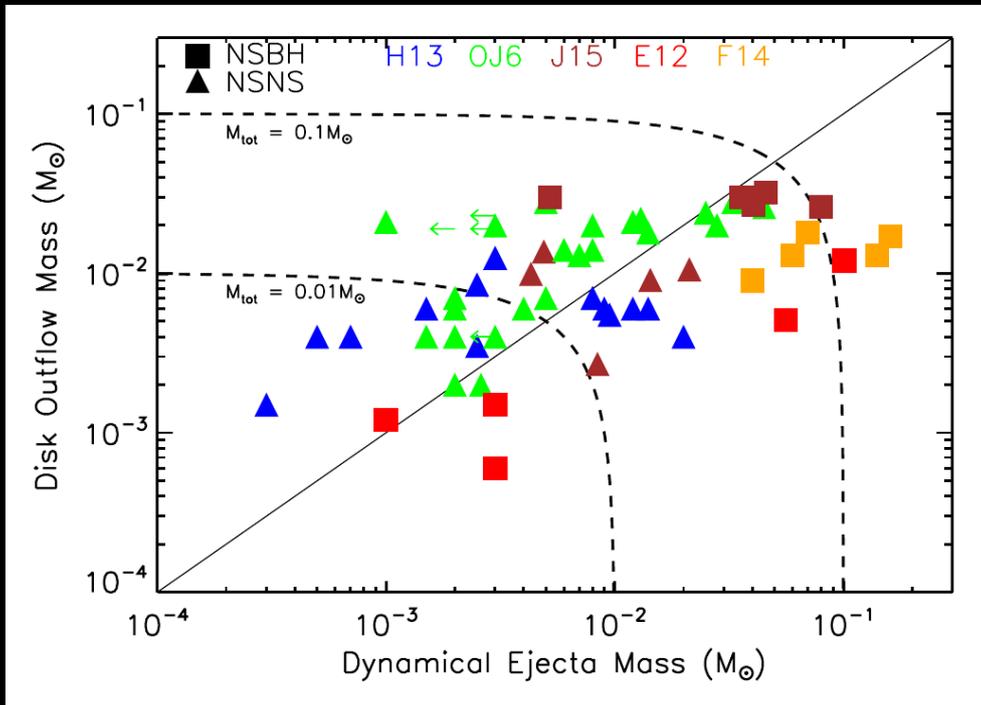
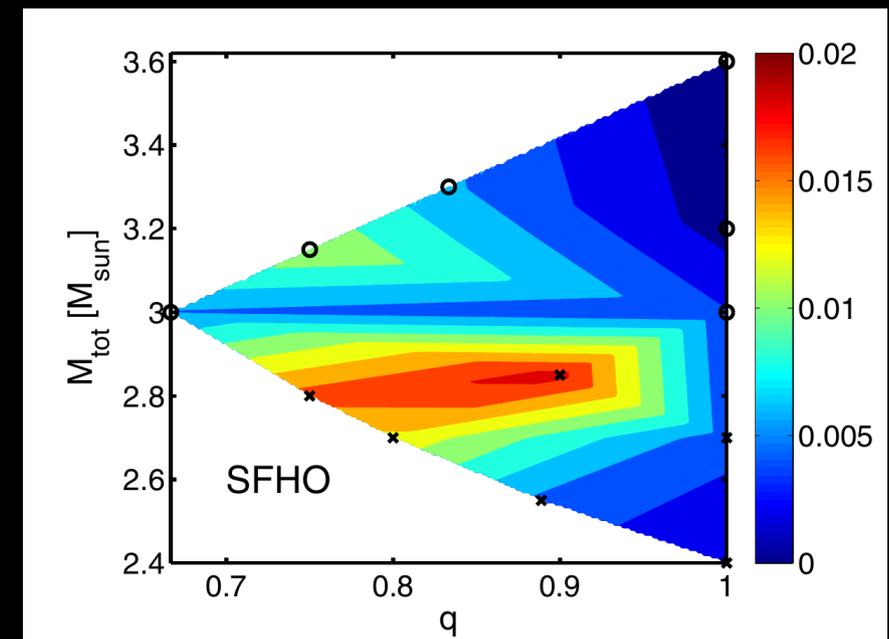
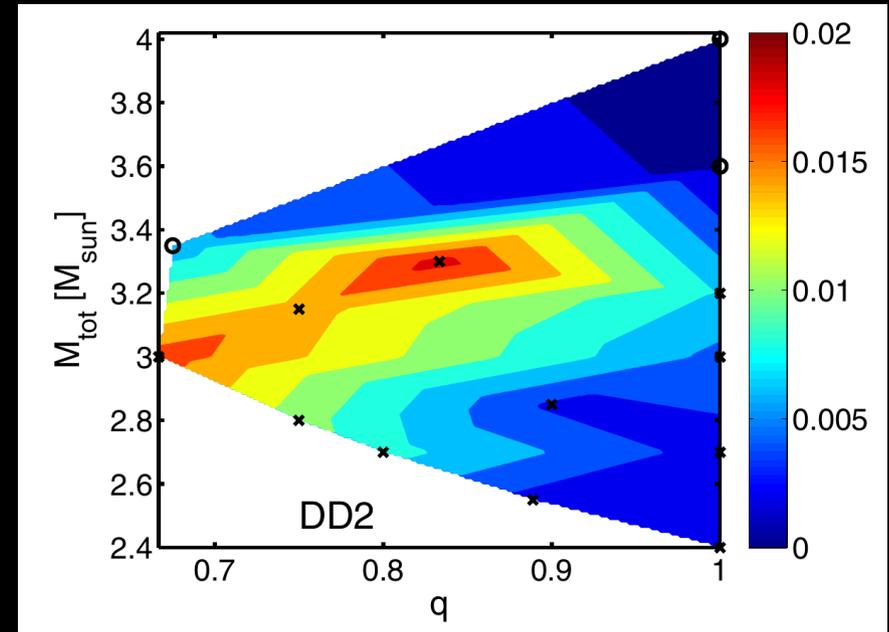
Soares-Santos et al 2017

Compilation in Cote et al 2018

- ▶ Ejecta masses depend on EoS and binary masses
- ▶ Note: high mass points already to soft EoS (tentatively/qualitatively)
- ▶ Prompt collapse leads to reduced ejecta mass
- ▶ Light curve depends on ejecta mass:
→ 0.02 - 0.05 M_{sun} point to delayed collapse

EoS and binary-mass dependence:

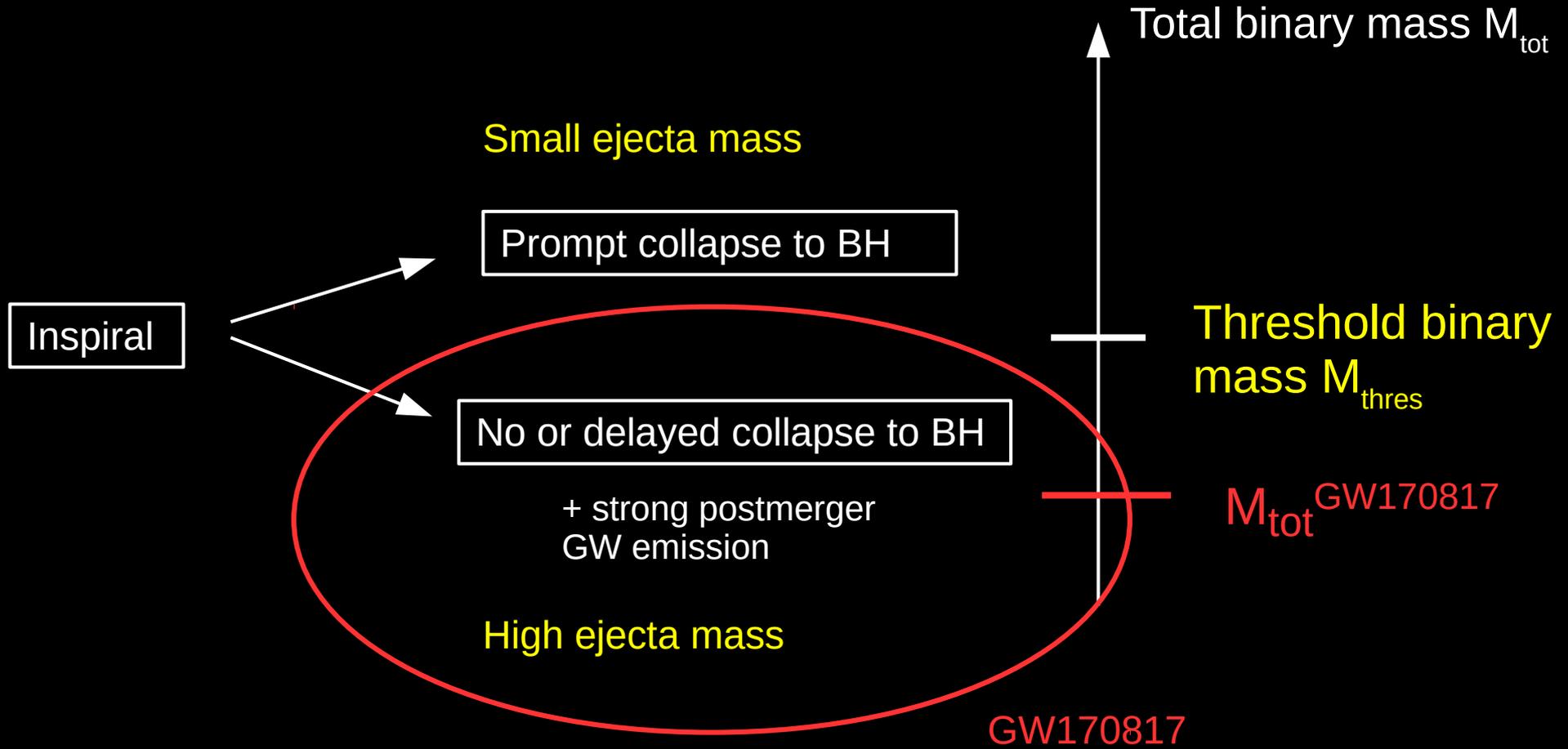
Bauswein et al. 2013



Compilation Wu et al 2016: dynamical and secular ejecta comparable

Only dynamical ejecta

Collapse behavior



(1) If GW170817 was a delayed (/no) collapse:

$$M_{\text{thres}} > M_{\text{tot}}^{\text{GW170817}}$$

(2) Recall: empirical relation for threshold binary mass for prompt collapse:

$$M_{\text{thres}} = \left(-3.38 \frac{G M_{\text{max}}}{c^2 R_{\text{max}}} + 2.43 \right) M_{\text{max}} > 2.74 M_{\odot}$$

(with $M_{\text{max}}, R_{\text{max}}$ unknown)

(3) Causality: speed of sound $v_s \leq c$

$$\Rightarrow M_{\text{max}} \leq \frac{1}{2.82} \frac{c^2 R_{\text{max}}}{G}$$

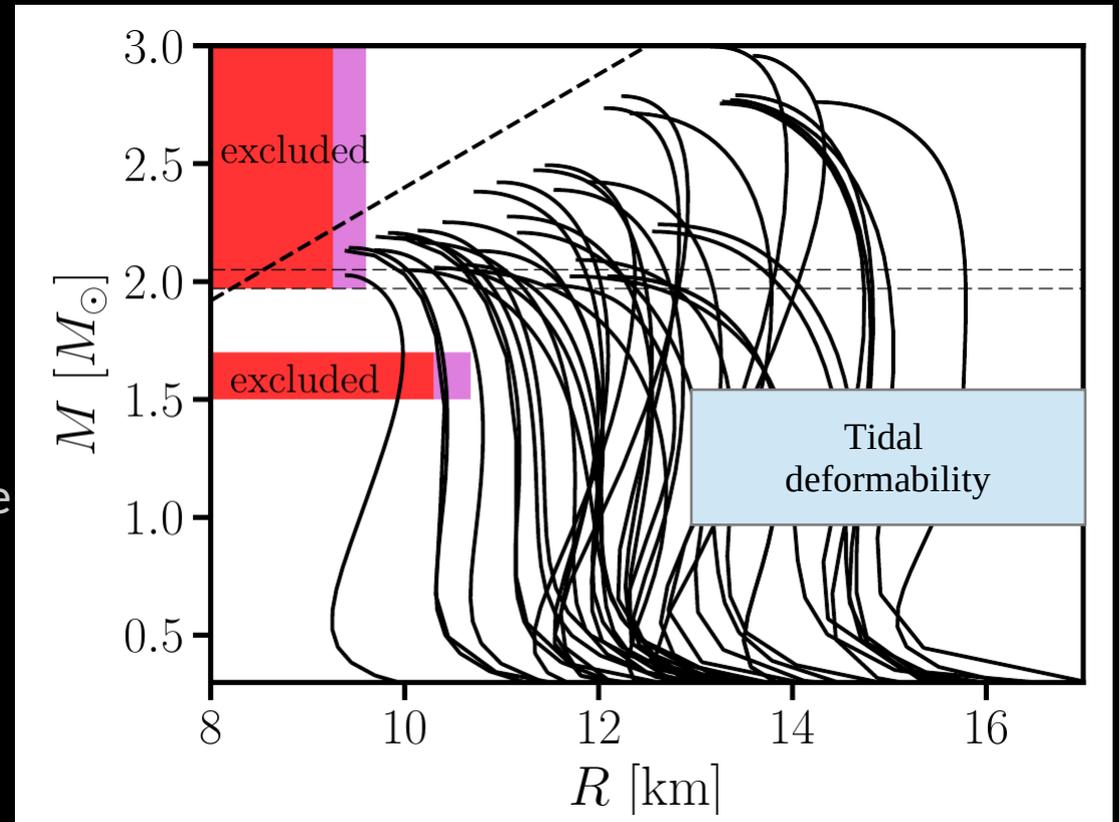
► Putting things together:

$$M_{\text{tot}}^{\text{GW170817}} \leq \left(-3.38 \frac{G M_{\text{max}}}{c^2 R_{\text{max}}} + 2.43 \right) M_{\text{max}} \leq \left(-\frac{3.38}{2.82} + 2.43 \right) \frac{1}{2.82} \frac{c^2 R_{\text{max}}}{G}$$

→ Lower limit on NS radius

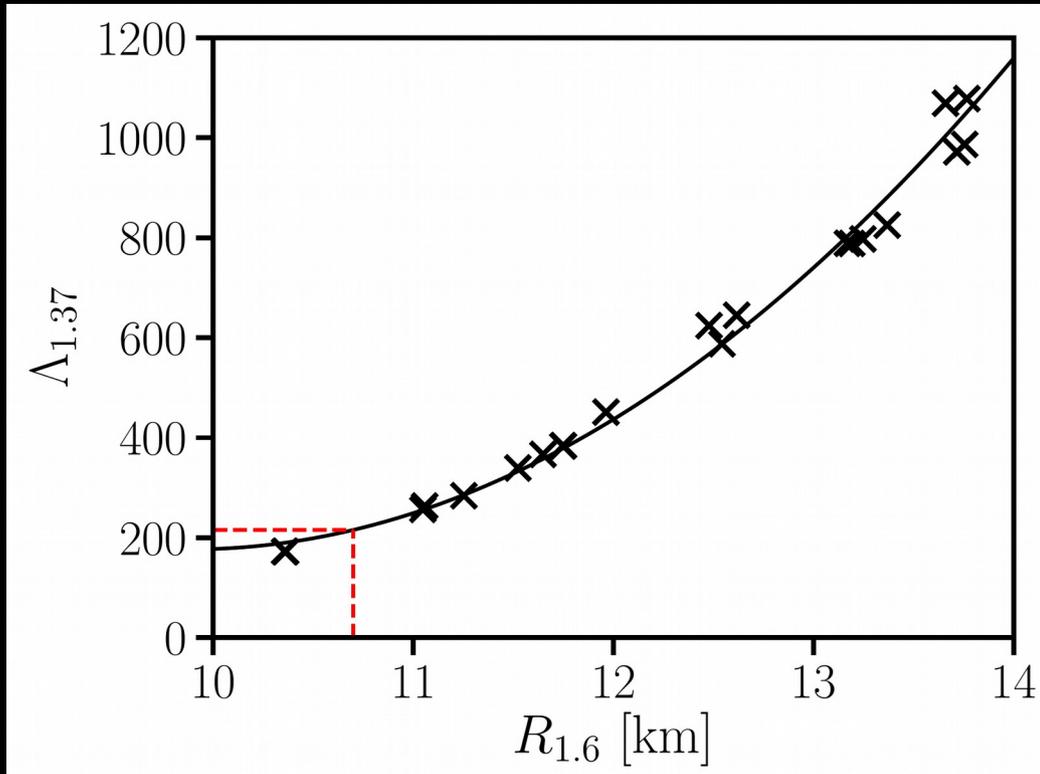
NS radius constraint from GW170817

- ▶ $R_{\text{max}} > 9.6$ km
- ▶ $R_{1.6} > 10.7$ km
- ▶ Excludes very soft nuclear matter
- ▶ Similar argument for Lambda in Radice et. al 2018
- ▶ follow-up Koepfel et al 2019 (same idea) arriving at similar constraints of 10.7 km
- ▶ See Capano et al. 2019 for an application within Bayesian statistics framework



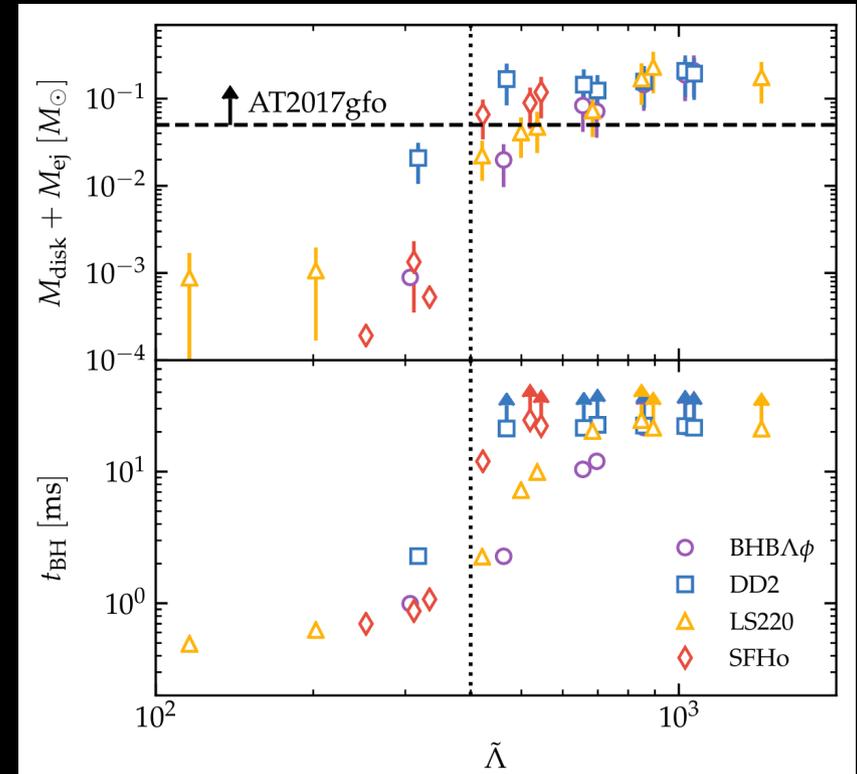
Bauswein et al. 2017

Radius vs. tidal deformability



Bauswein et al. 2019

Radice et al 2018



- ▶ Radius and tidal deformability scale tightly → $\Lambda > 210$
- ▶ Limit cannot be much larger otherwise we could get no direct collapse / dim counterpart (unless one weakens some of the conservative assumptions)
- ▶ Radice et al. 2018 followed a very similar argument claiming $\Lambda > 400$ (300 in Dai 2019)
 - only 4 EoS considered – no complete coverage existing simulation data/parameter space
 - no argument why fifth EoS shouldn't lie at $\Lambda < 400$ (see also Tews et al. 2018, Kiuchi 2019)
 - full EoS dependence (including M_{max}) has to be investigated via M_{thres}

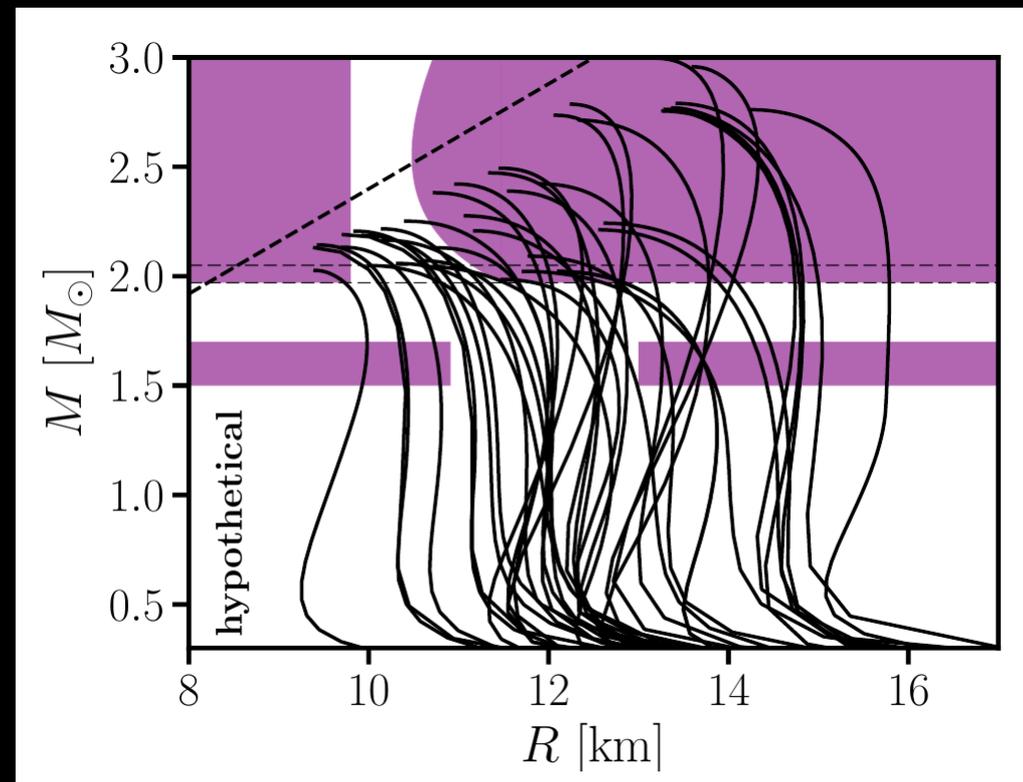
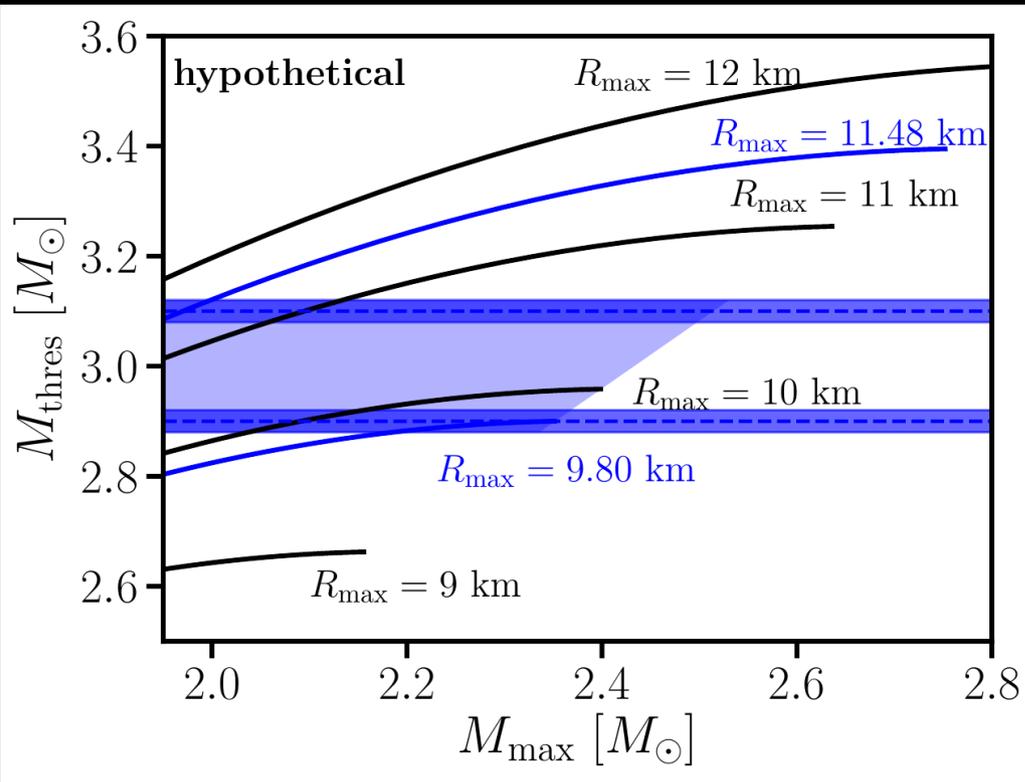
Discussion - robustness

- ▶ Binary masses well measured with high confidence error bar
- ▶ Clearly defined working hypothesis: delayed collapse
 - testable by refined emission models
 - as more events are observed more robust distinction
- ▶ Very conservative estimate, errors can be quantified
- ▶ Empirical relation can be tested by more elaborated simulations (but unlikely that MHD or neutrinos can have strong impact on M_{thres})
- ▶ Confirmed by semi-analytic collapse model
- ▶ Low-SNR constraint !!!

Future

- ▶ Any new detection can be employed if it allows distinction between prompt/delayed collapse
- ▶ With more events in the future our comprehension of em counterparts will grow → more robust discrimination of prompt/delayed collapse events
- ▶ Low-SNR detections sufficient !!! → that's the potential for the future
 - we don't need louder events, but more
 - complimentary to existing ideas for EoS constraints

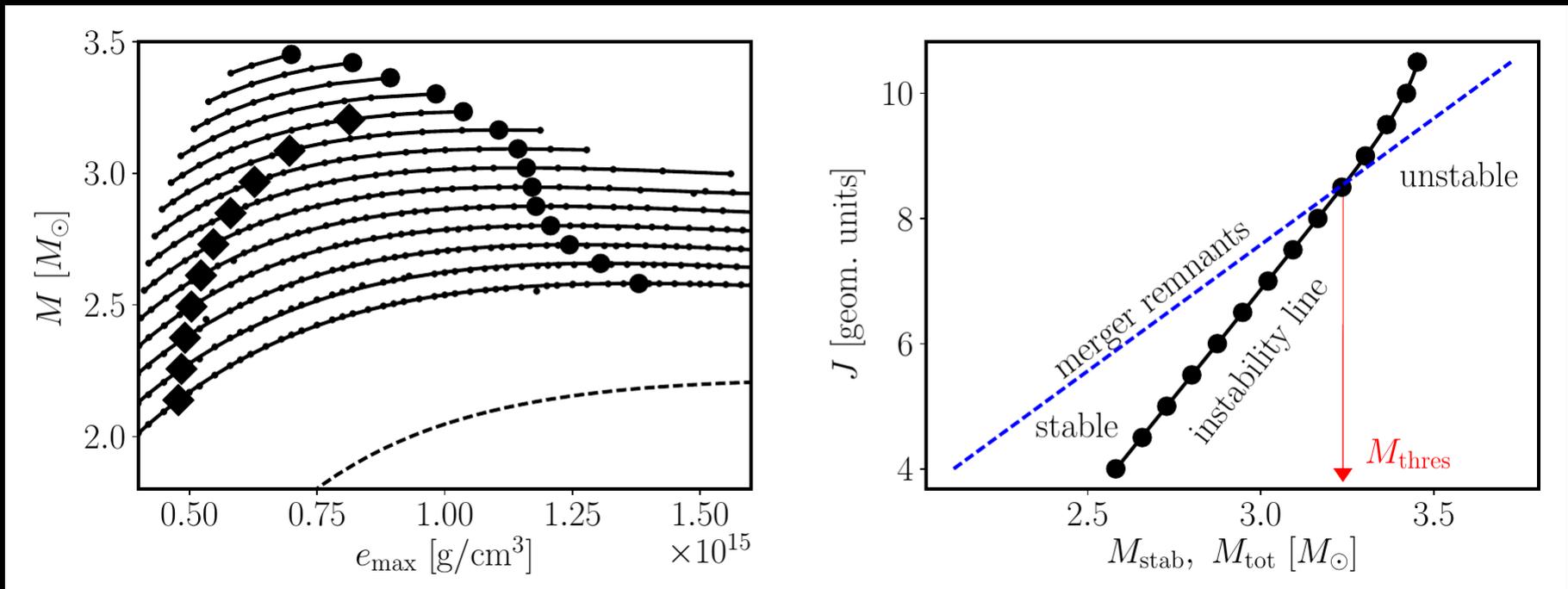
Future detections (hypothetical discussion)



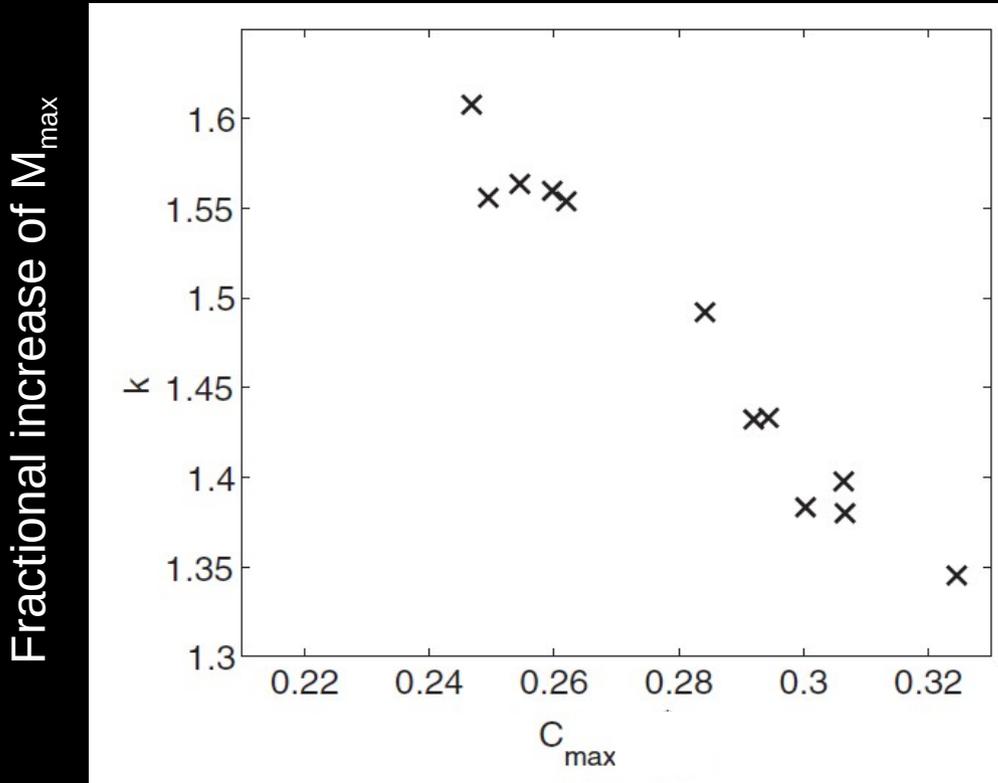
- as more events are observed, bands converge to true M_{thres}
- prompt collapse constrains M_{\max} from above

Semi-analytic model: details

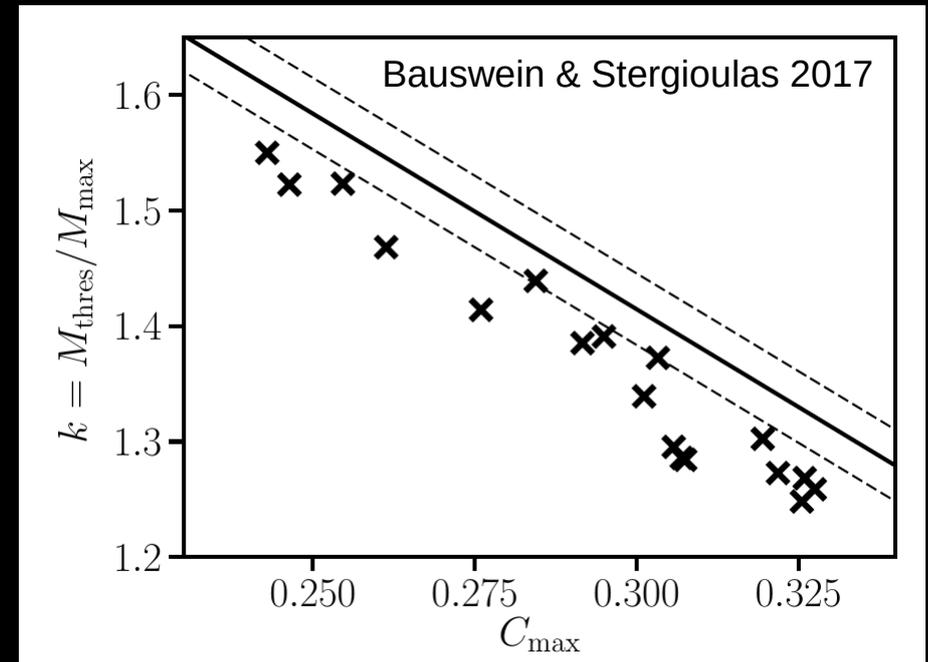
- ▶ Stellar equilibrium models computed with RNS code (diff. Rotation, $T=0$, many different microphysical EoS) \Rightarrow turning points $\Rightarrow M_{\text{stab}}(J)$
- ▶ Compared to $J(M_{\text{tot}})$ of merger remnants from simulations (very robust result) \rightarrow practically independent from simulations



Semi-analytic model reproducing collapse behavior



Bauswein et al 2013: numerical determination of collapse threshold through hydrodynamical simulations



Solid line fit to numerical data

Crosses stellar **equilibrium models**:

- prescribed (simplistic) diff. rotation
- many EoSs at $T=0$
- detailed angular momentum budget !
=> equilibrium models qualitatively reproduce collapse behavior
- even quantitatively good considering the adopted approximations

Future: Maximum mass

- ▶ Empirical relation

$$M_{\text{thres}} = \left(-3.6 \frac{G M_{\text{max}}}{c^2 R_{1.6}} + 2.38 \right) M_{\text{max}}$$

- ▶ Sooner or later we'll know $R_{1.6}$ (e.g. from postmerger) and M_{thres} (from several events – through presence/absence of postmerger GW emission or em counterpart)

=> direct inversion to get precise estimate of M_{max}

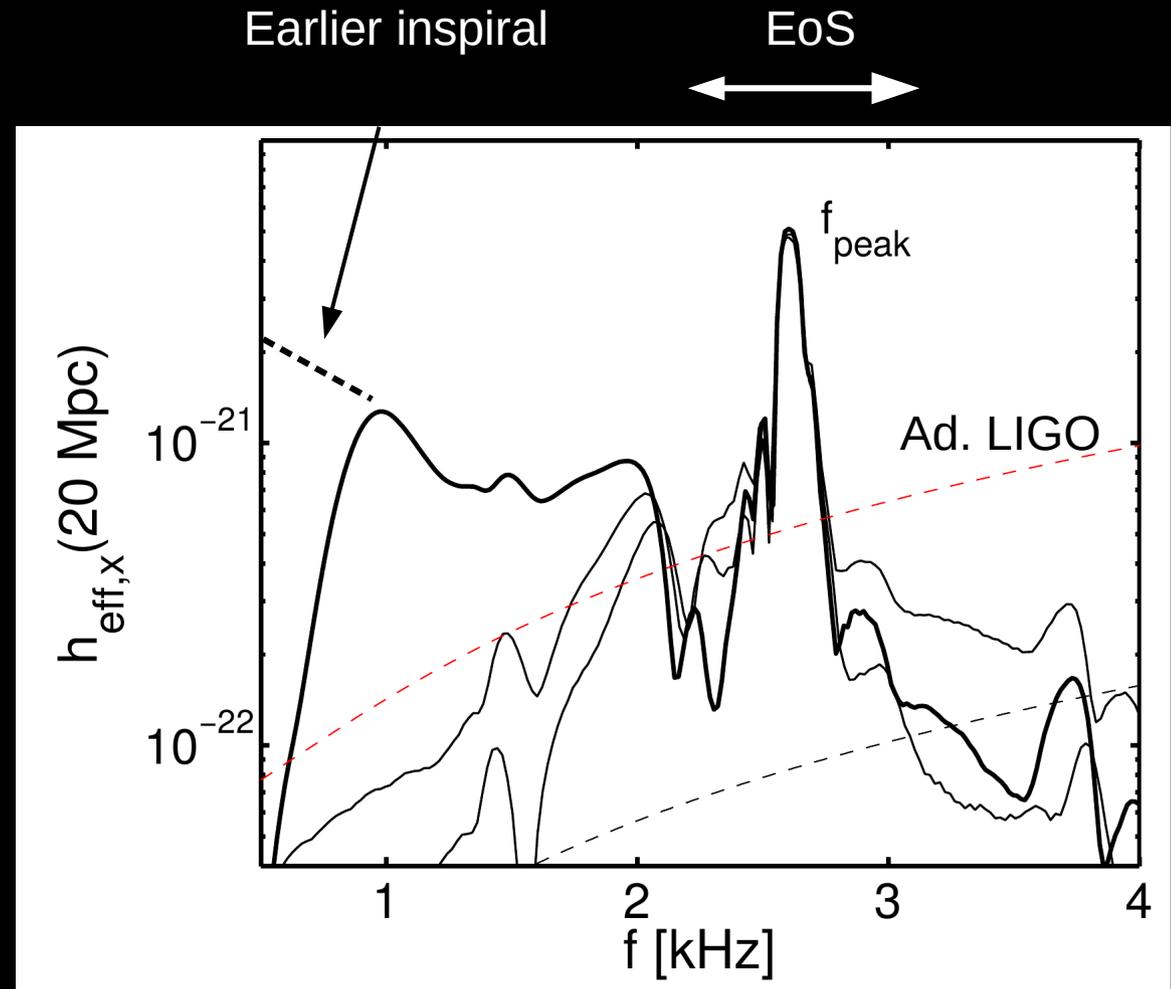
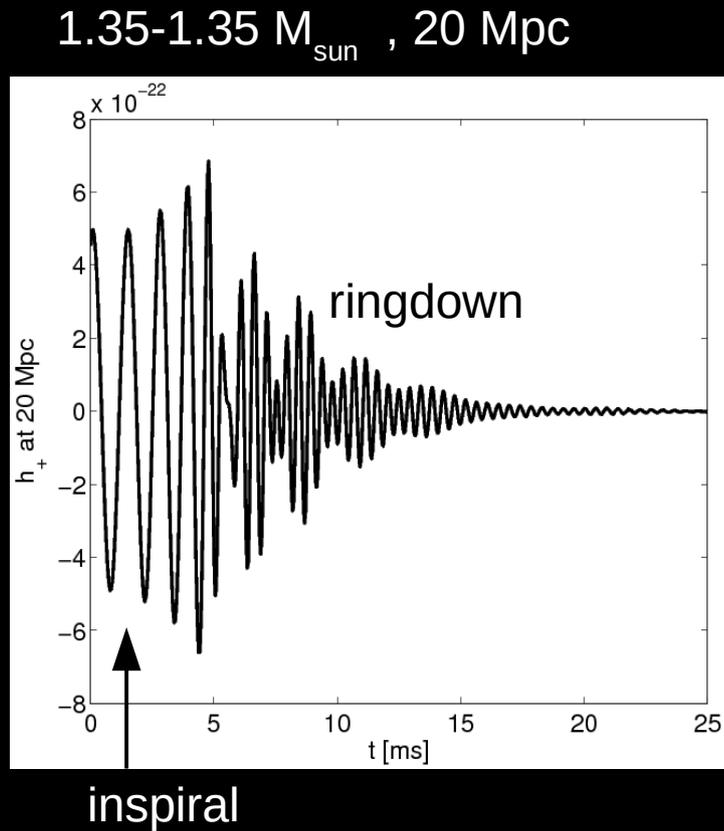
Postmerger GW emission*

(dominant frequency of postmerger phase)

- determine properties of EoS/NSs
- complementary to inspiral

* not detected for GW170817 – expected for current sensitivity and $d=40$ Mpc
(Abbott et al. 2017)

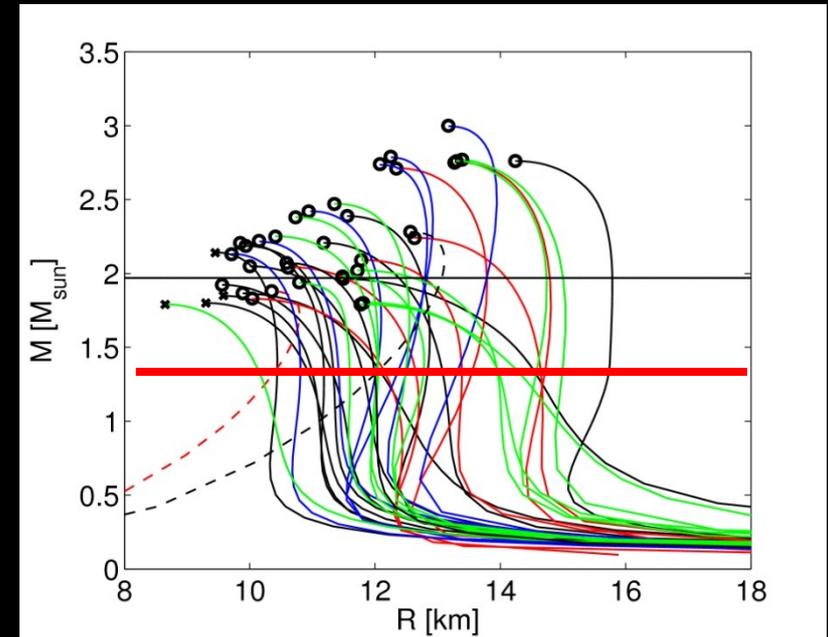
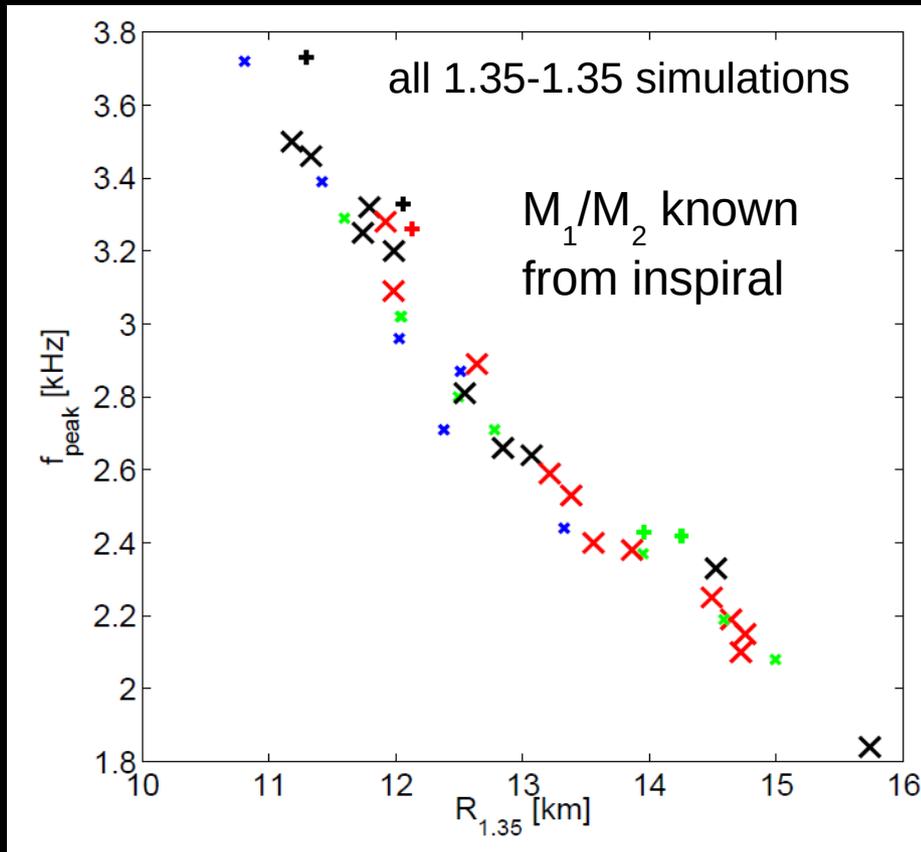
Postmerger



Dominant postmerger oscillation frequency f_{peak}

Very characteristic (robust feature in all models)

Gravitational waves – EoS survey



characterize EoS by radius of nonrotating NS with $1.35 M_{\text{sun}}$

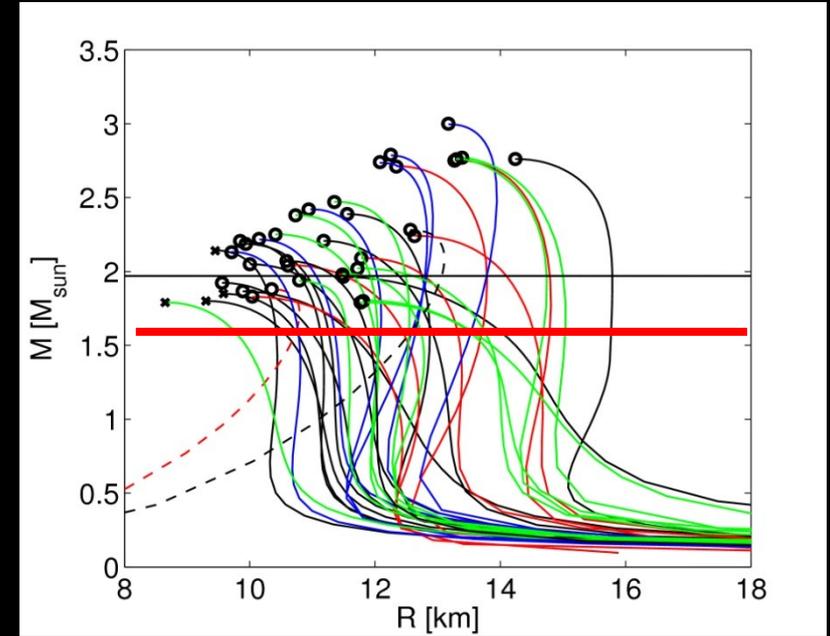
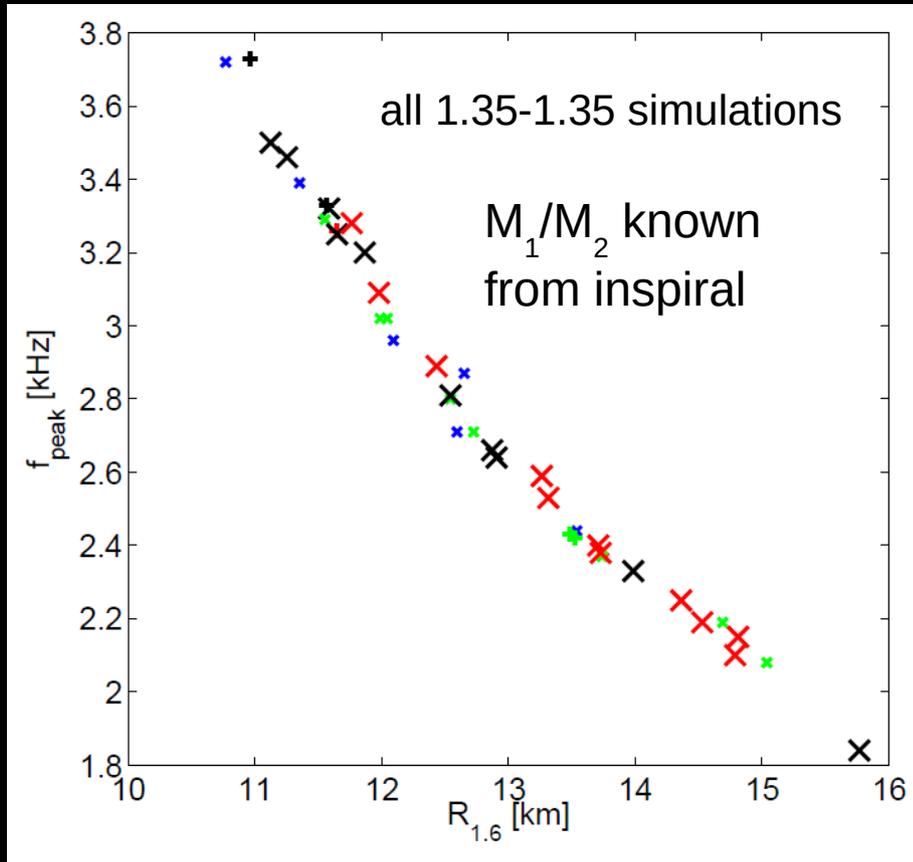
Bauswein et al. 2012

Pure TOV/EoS property \Rightarrow **Radius measurement** via f_{peak}

Here only 1.35-1.35 M_{sun} mergers (binary masses measurable) – similar relations exist for other fixed binary setups !!!

~ 40 different NS EoSs

Gravitational waves – EoS survey



characterize EoS by radius of nonrotating NS with $1.6 M_{\text{sun}}$

Bauswein et al. 2012

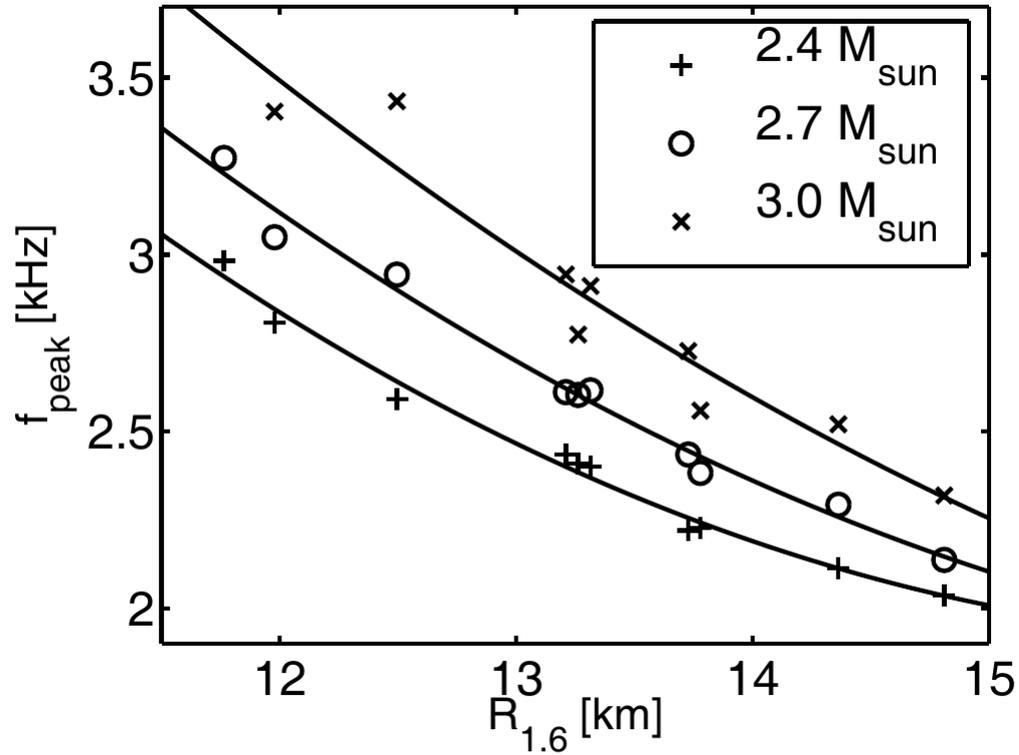
Pure TOV/EoS property \Rightarrow **Radius measurement** via f_{peak}

Smaller scatter in empirical relation (< 200 m) \rightarrow smaller error in radius measurement

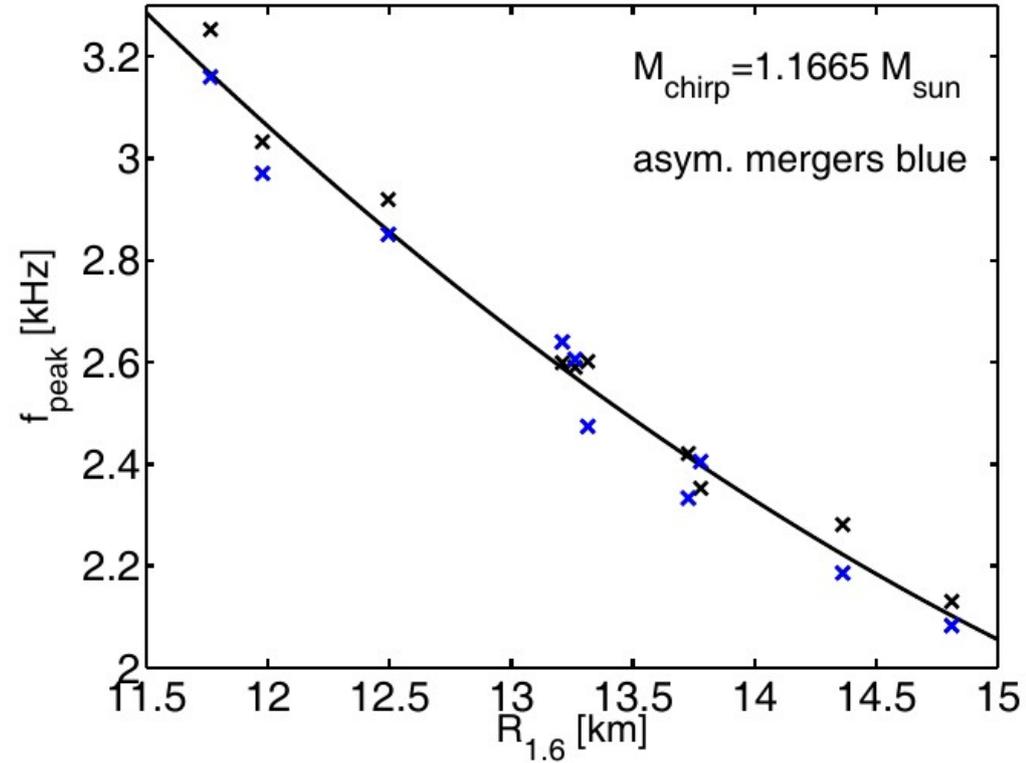
Note: R of $1.6 M_{\text{sun}}$ NS scales with f_{peak} from 1.35 - $1.35 M_{\text{sun}}$ mergers (density regimes comparable)

GW data analysis: Clark et al 2014, Clark et al 2016, Chatziioannou et al 2017, Bose et al. 2018, Torres-Riva et al 2019, Breschi et al 2019, ... \rightarrow detectable at a few 10 Mpc

Binary mass variations

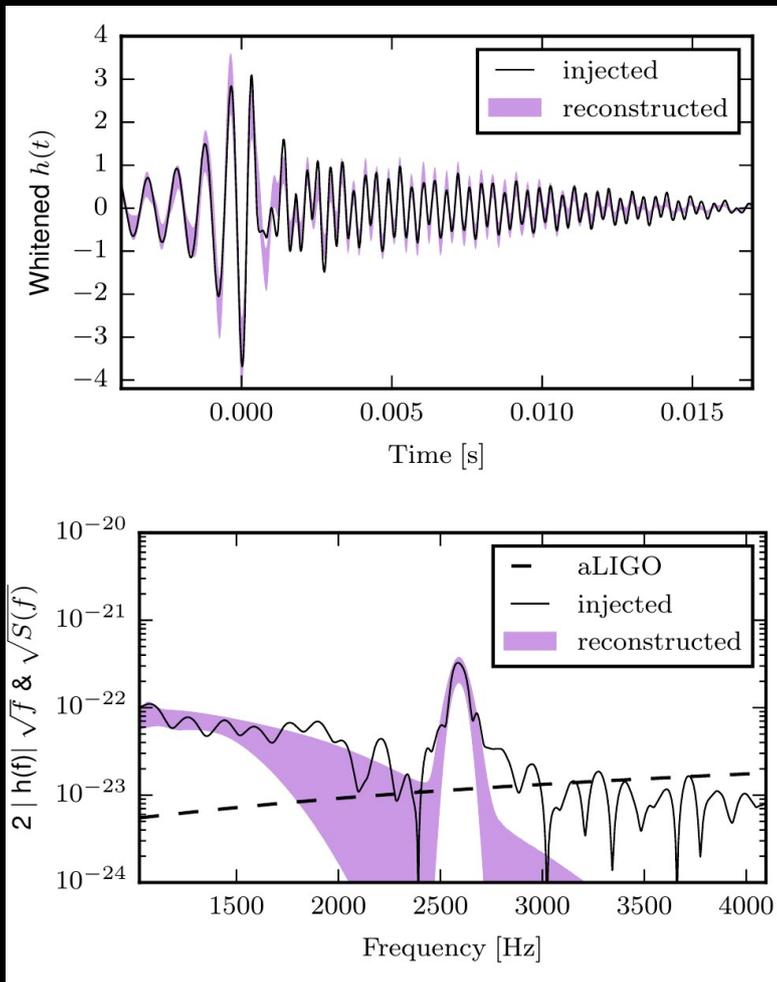


Different total binary masses
(symmetric)

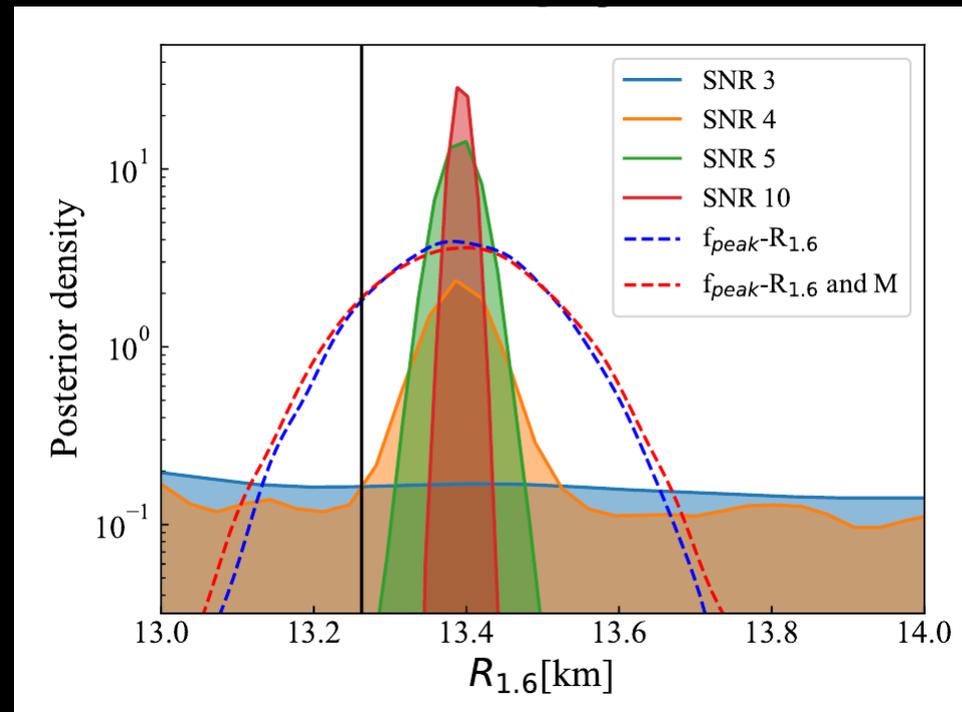


Fixed chirp mass (asymmetric 1.2-1.5
 M_{sun} binaries and symmetric 1.34-
1.34 M_{sun} binaries)

Model-agnostic data analysis



Based on wavelets

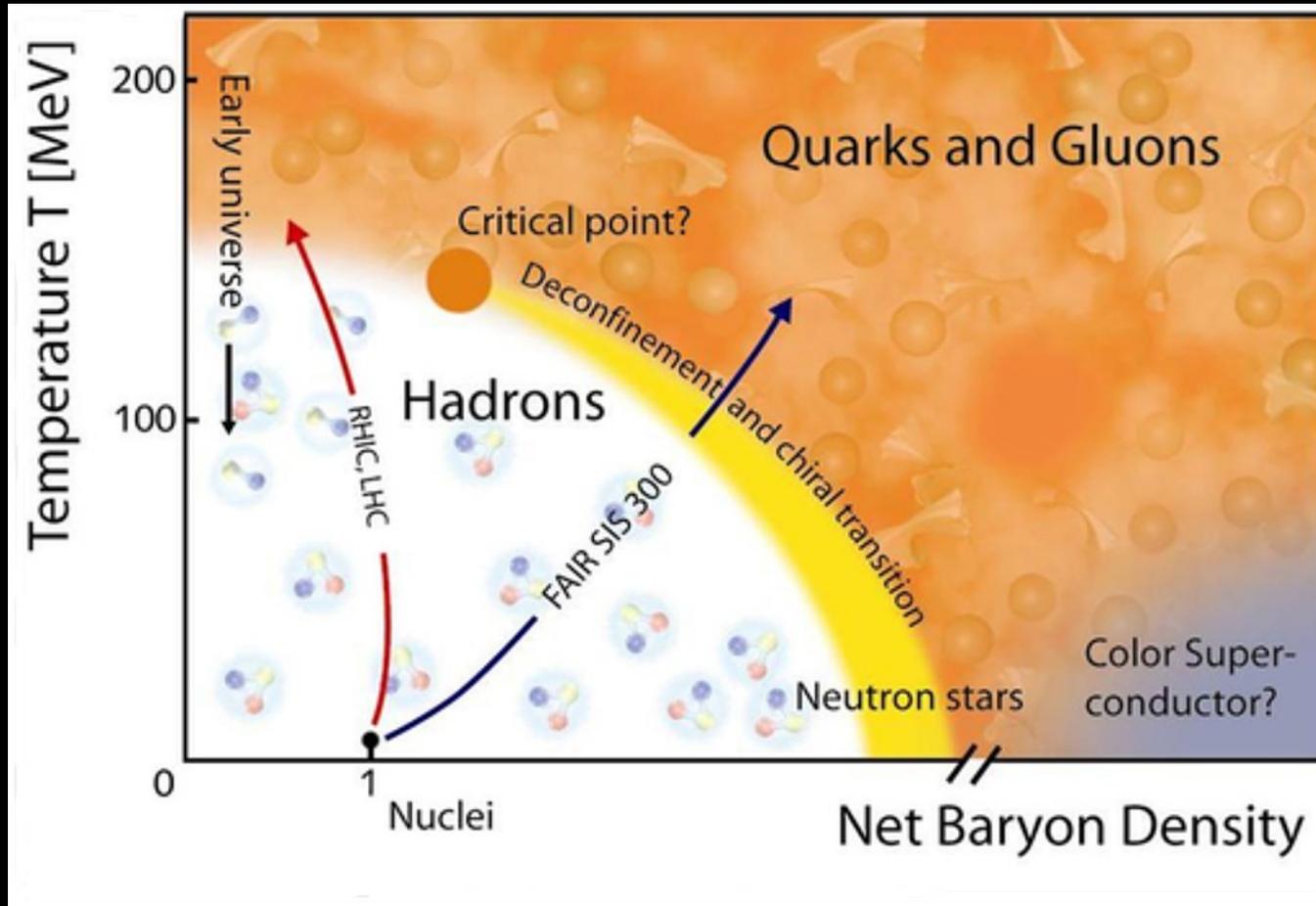


Chatziioannou et al. (2017) → detectable at a few 10 Mpc

See also Bauswein et al 2012, Clark et al 2014, Clark et al 2016, Chatziioannou et al 2017, Bose et al. 2018, Yang et al 2019, Torres-Riva et al 2019, Breschi et al 2019, Martynov et al 2019, ...

Observable signature of (QCD) phase transition

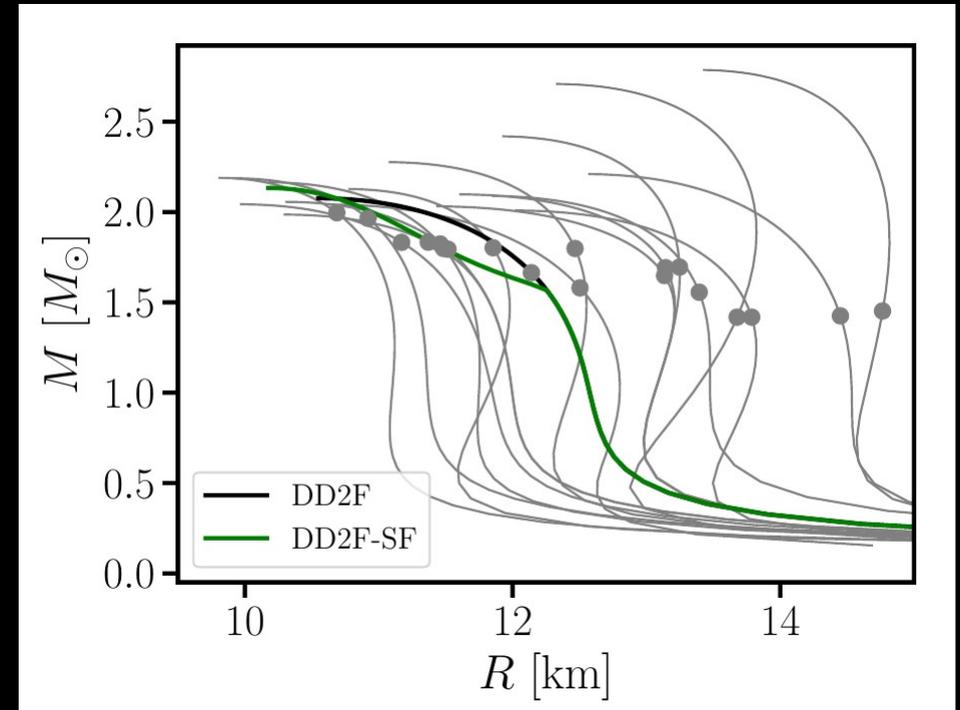
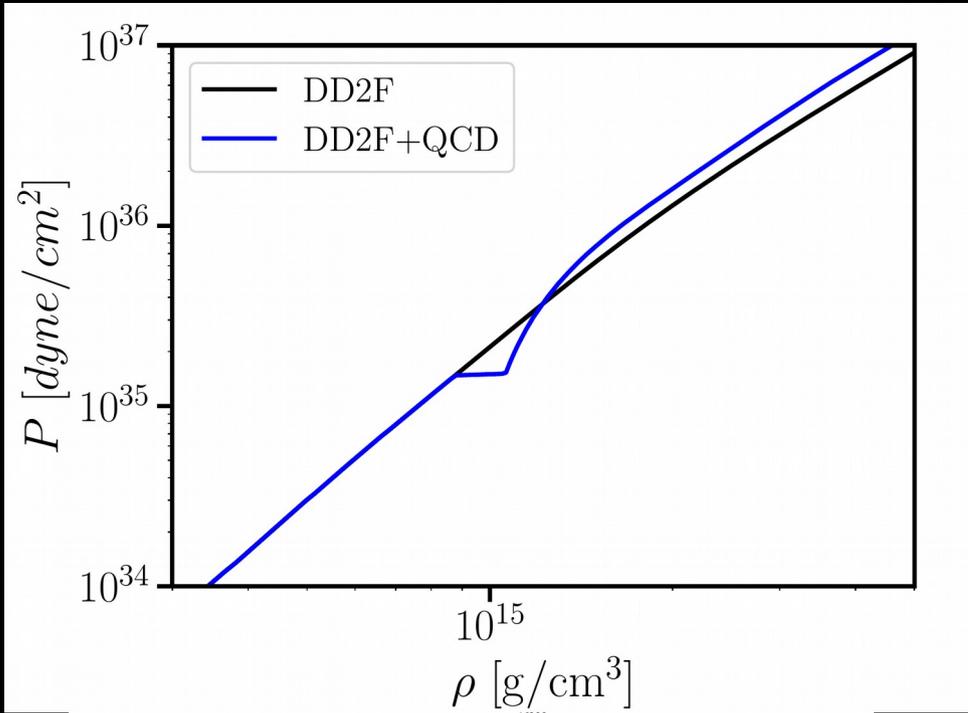
Phase diagram of matter



Does the phase transition to quark-gluon plasma occur (already) in neutron stars or only at higher densities ?

EoS with 1st-order phase transition to quark matter

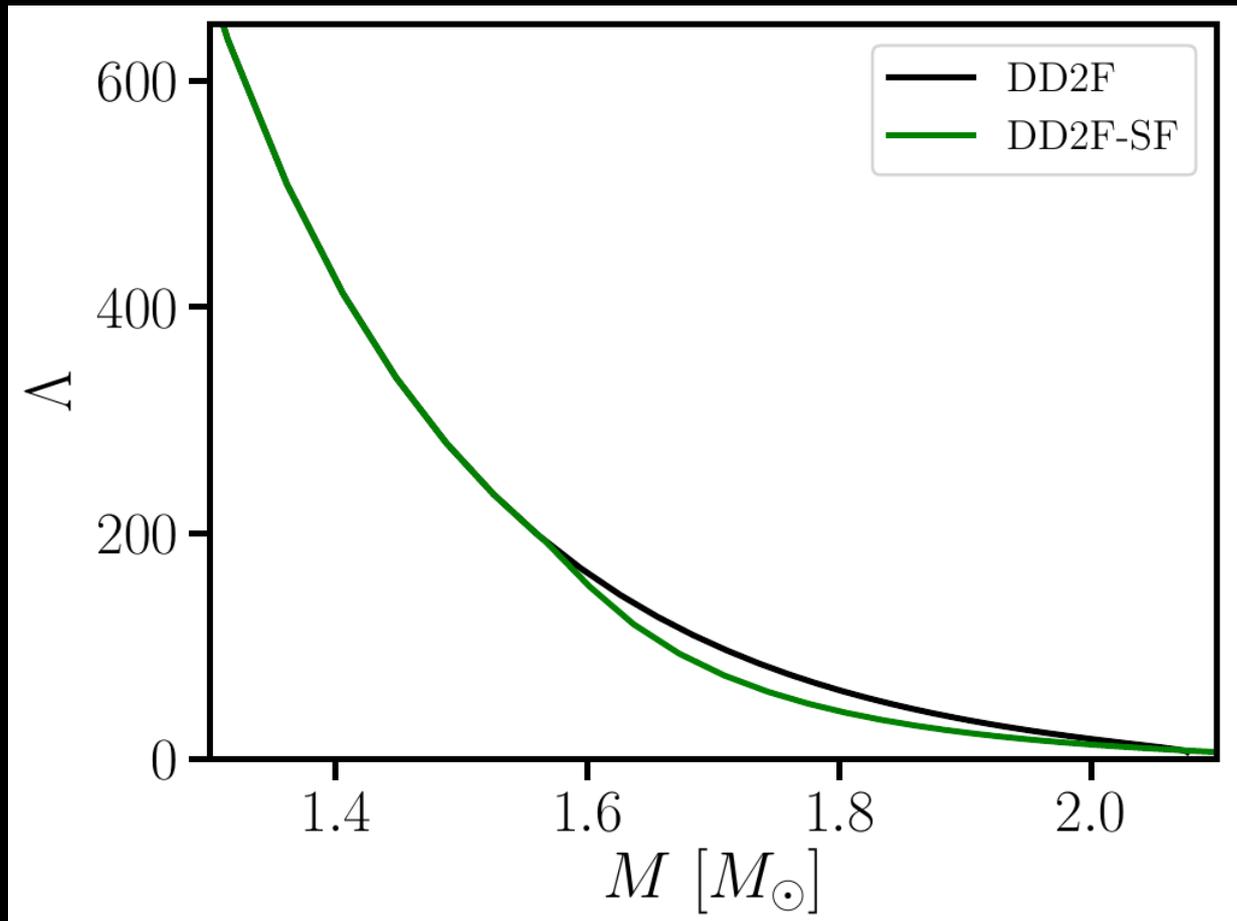
Bauswein et al. 2019



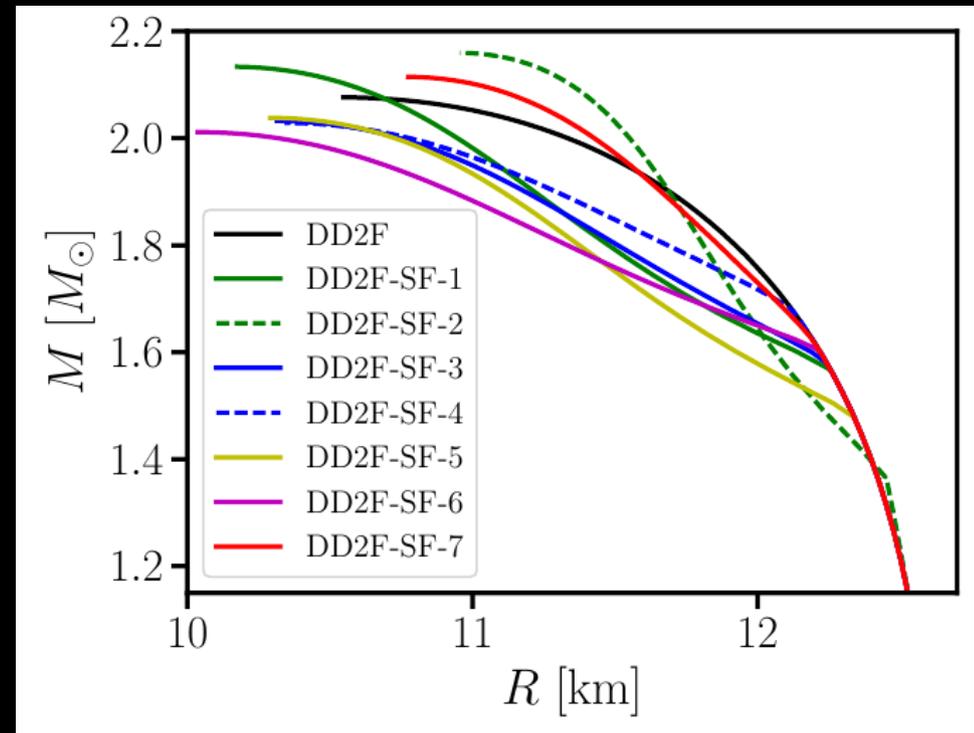
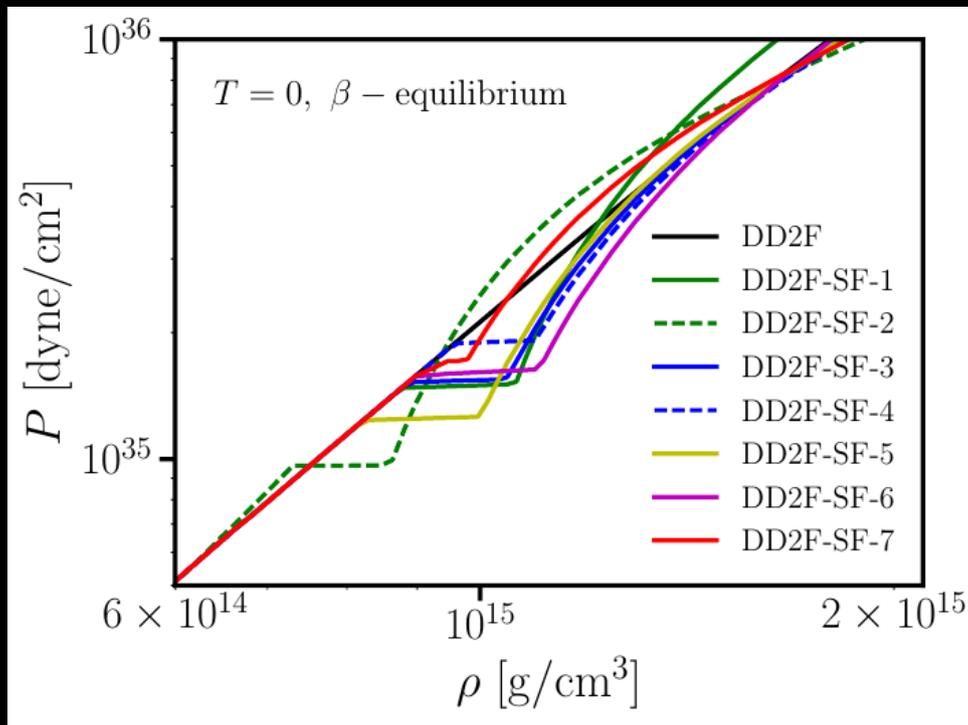
- ▶ EoS from Wroclaw group (Fischer, Bastian, Blaschke; Fischer et al. 2018) – as one example for an EoS with a strong 1st-order phase transition to deconfined quarks
- ▶ Difficult to measure transition in mergers through inspiral: Lambda very small, high mass star probably less frequent

Phase transition

- ▶ Even strong phase transitions leave relatively weak impact on tidal deformability

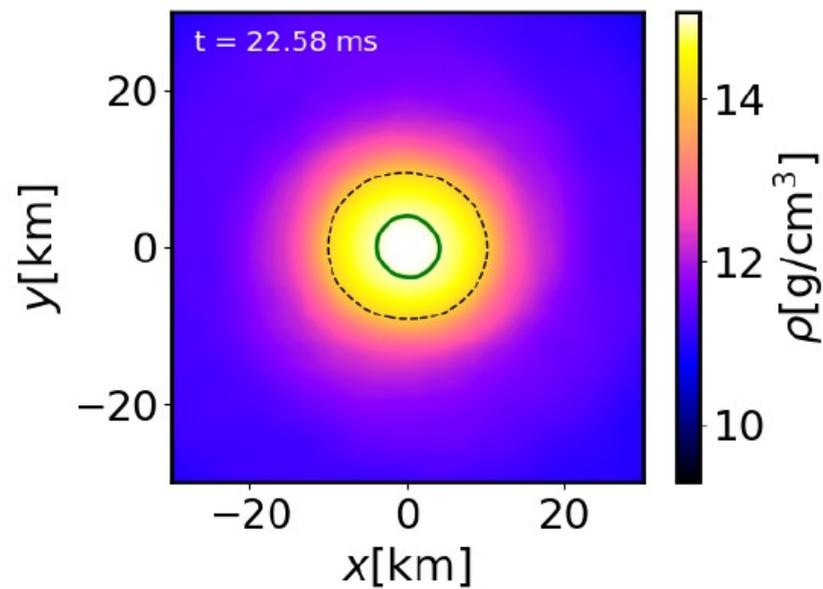
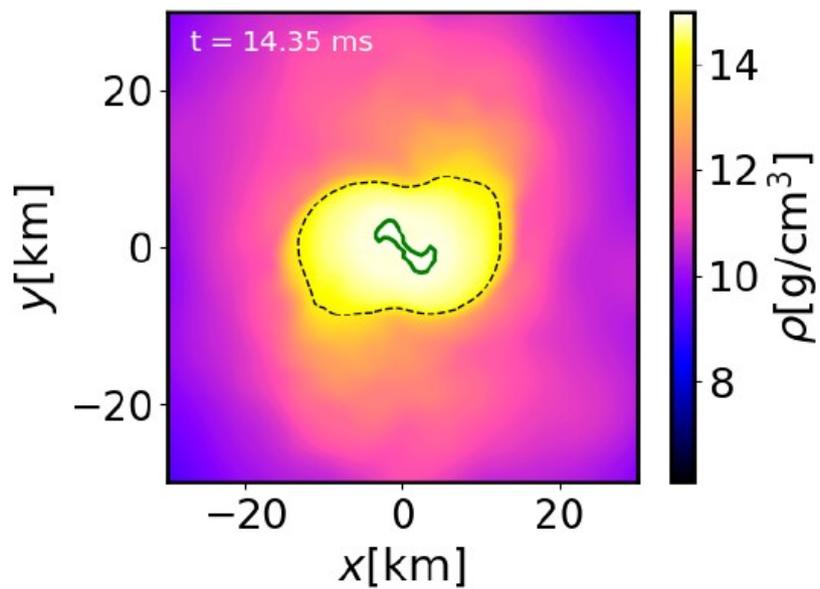
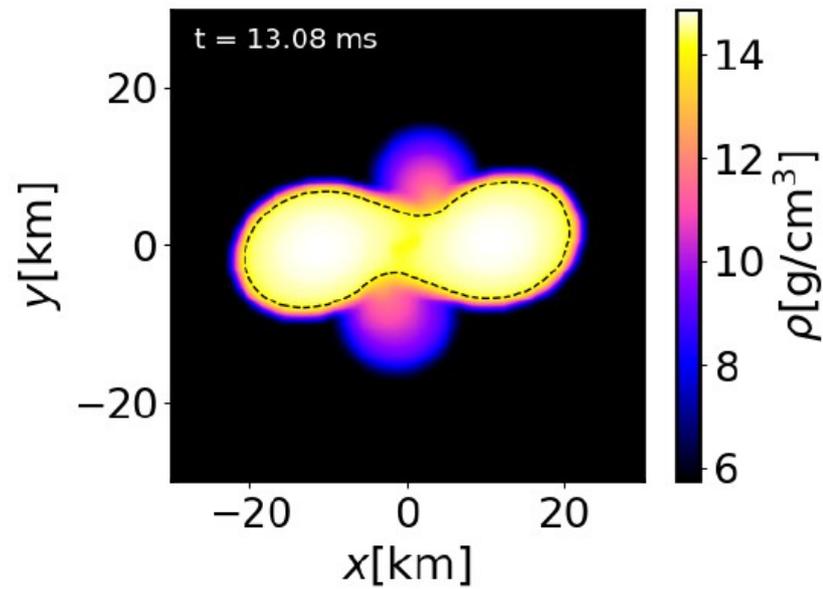
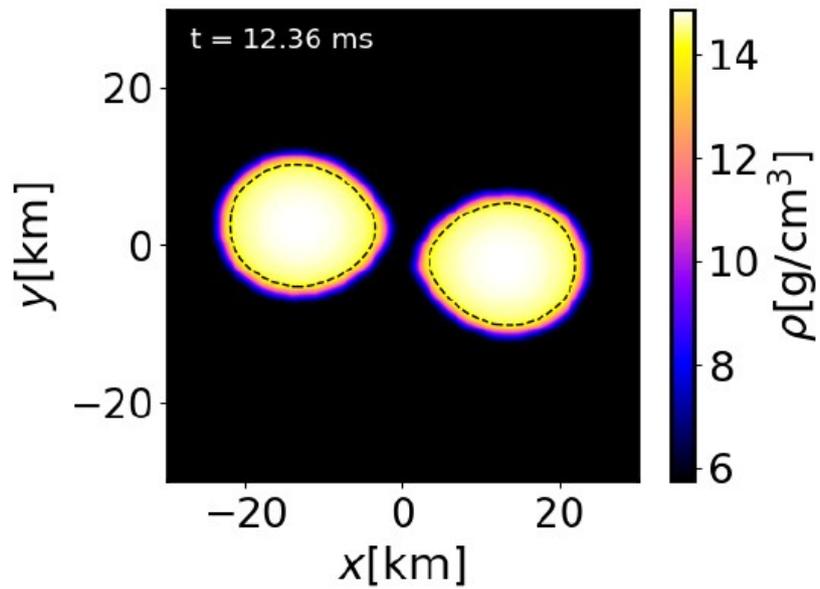


- 7 different models for quark matter: different onset density, different density jump, different stiffness of quark matter phase

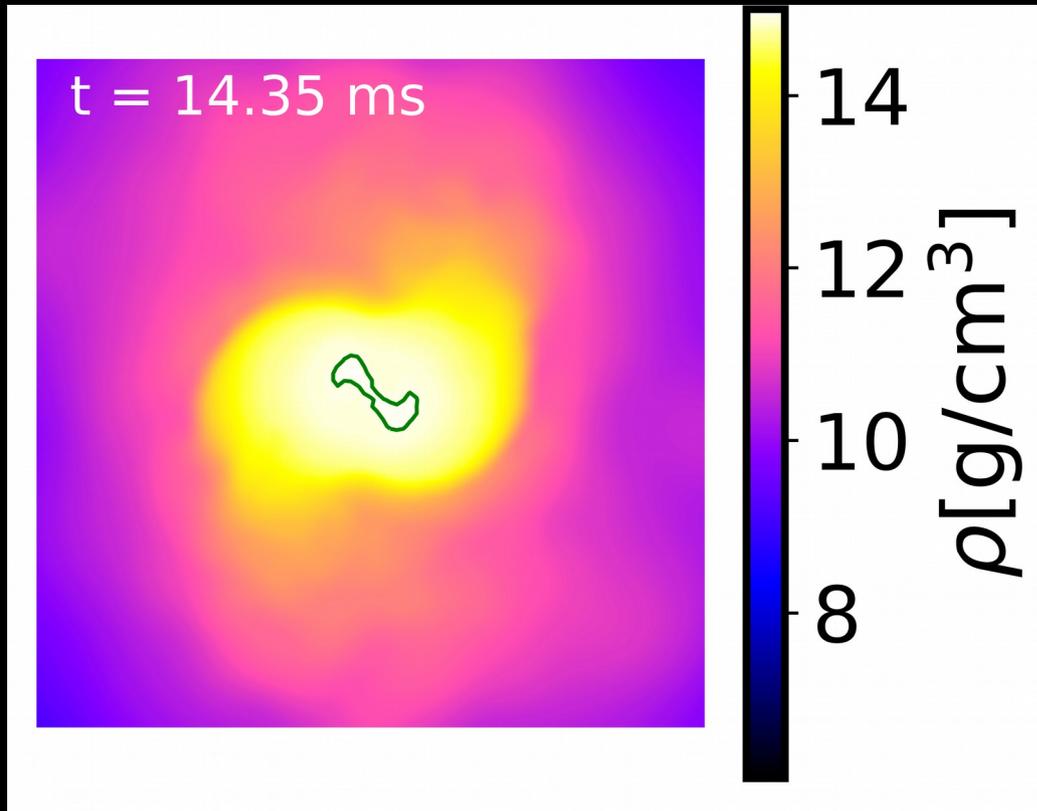


Bauswein et al. 2019

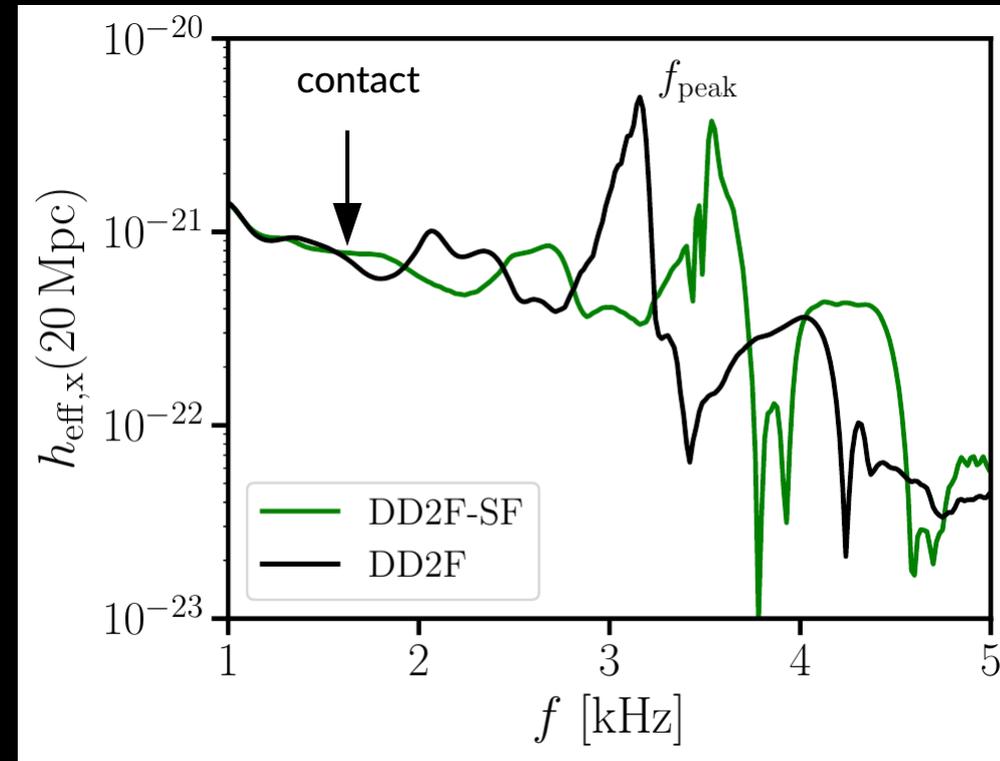
EoSs from Wroclaw group



Merger simulations



► GW spectrum 1.35-1.35 Msun



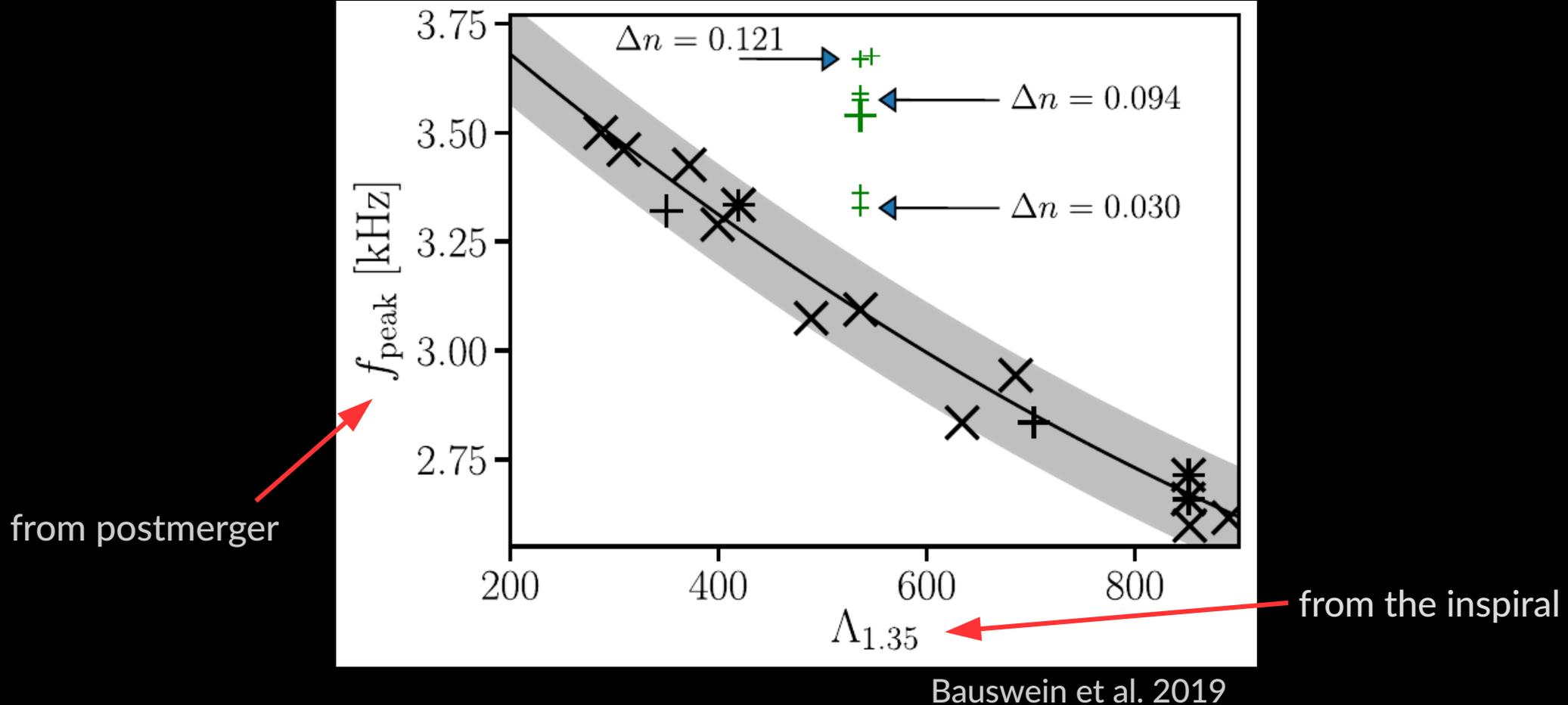
Bauswein et al. 2019

But: a high frequency on its own may not yet be characteristic for a phase transition

→ unambiguous signature

(→ show that all purely baryonic EoS behave differently)

Signature of 1st order phase transition

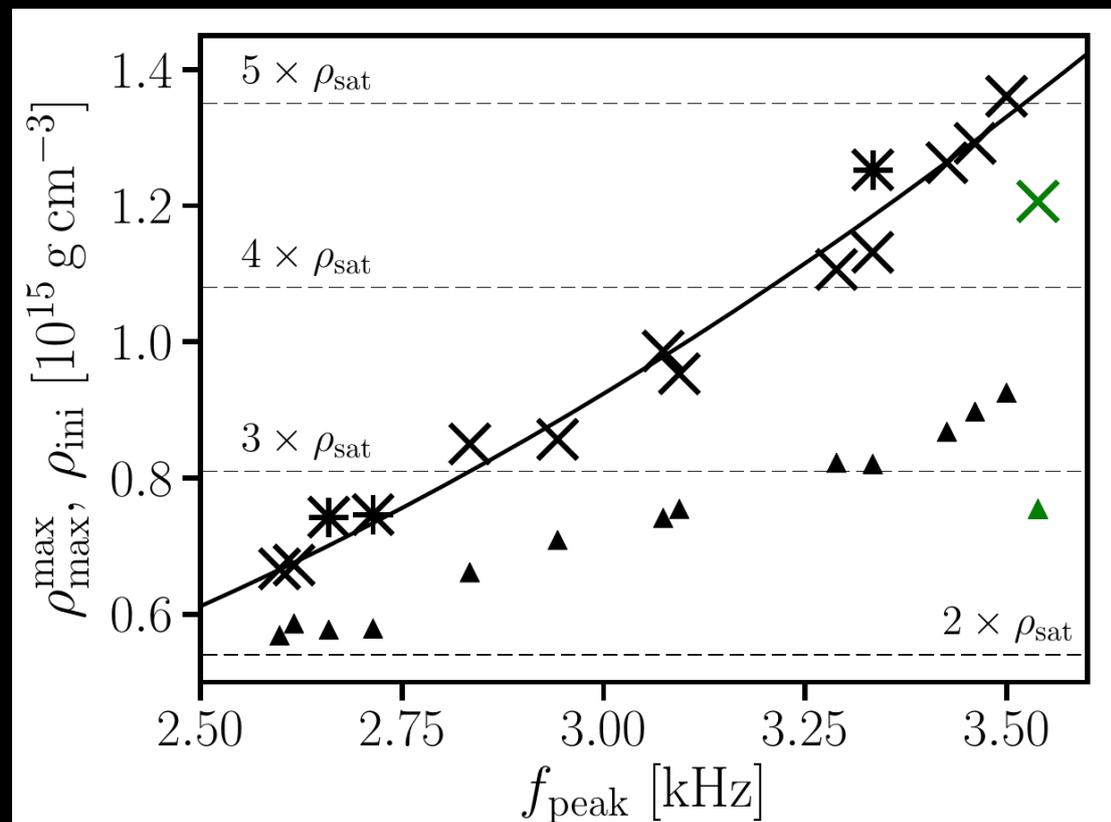


- ▶ Tidal deformability measurable from inspiral to within 100-200 (Adv. Ligo design)
- ▶ Postmerger frequency measurable to within a few 10 Hz @ a few 10 Mpc (either Adv. Ligo or upgrade: e.g Clark et al. 2016, Chatzioannou et al 2017, Bose et al 2018, Torres-Rivas et al 2019)
- ▶ Important: “all” purely hadronic EoS (including hyperonic EoS) follow f_{peak} - Λ relation → deviation characteristic for strong 1st order phase transition

Discussion

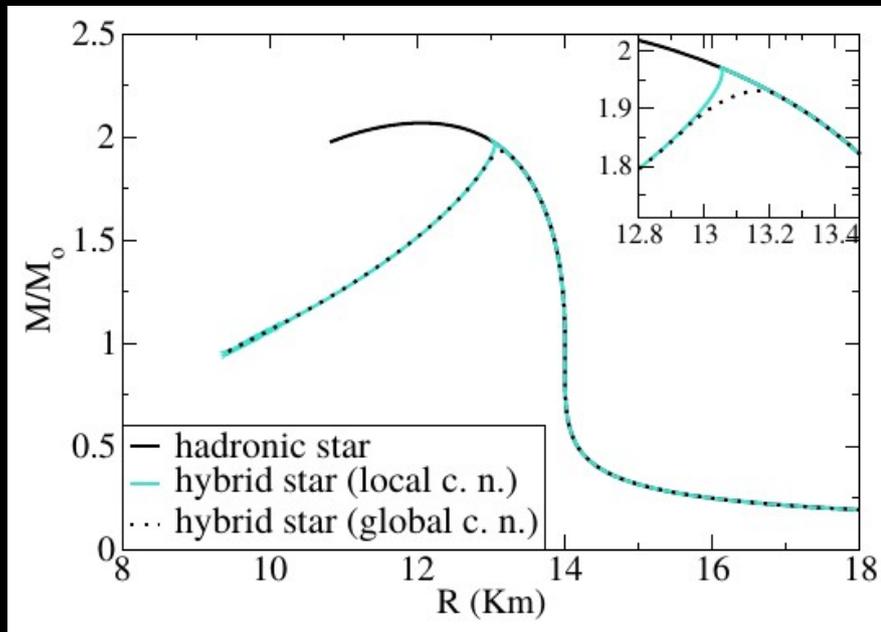
- ▶ Consistency with f_{peak} -Lambda relation points to
 - purely baryonic EoS
 - (or an at most weak phase transition \rightarrow no strong compactification)in the tested (!) density regime
- ▶ f_{peak} also determines maximum density in postmerger remnant
- ▶ postmerger GW emission provides complimentary information to inspiral \rightarrow probes higher density regime

Bauswein et al. 2019

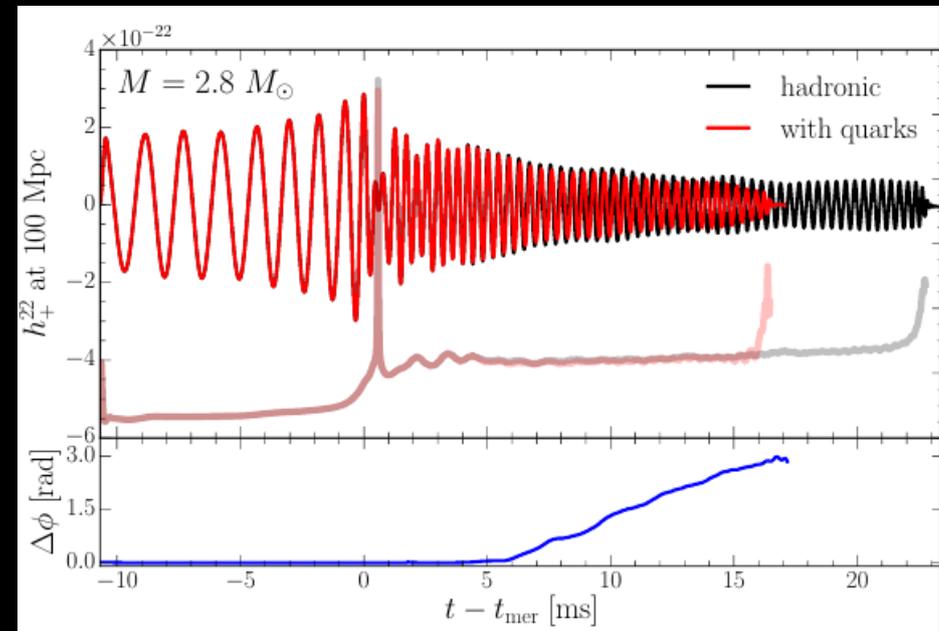


Different transitions

- ▶ Absolutely stable strange quark matter leads to similar shift (if still considered viable → requires re-hadronization to be compatible with GW170817)
- ▶ Phase transition without extended stable hybrid star branch → earlier collapse (not necessarily characteristic) → slight dephasing (small quark matter fraction) – similar frequencies, i.e. no strong and unambiguous signature of quark matter as hadronic EoS can lead to similar effects



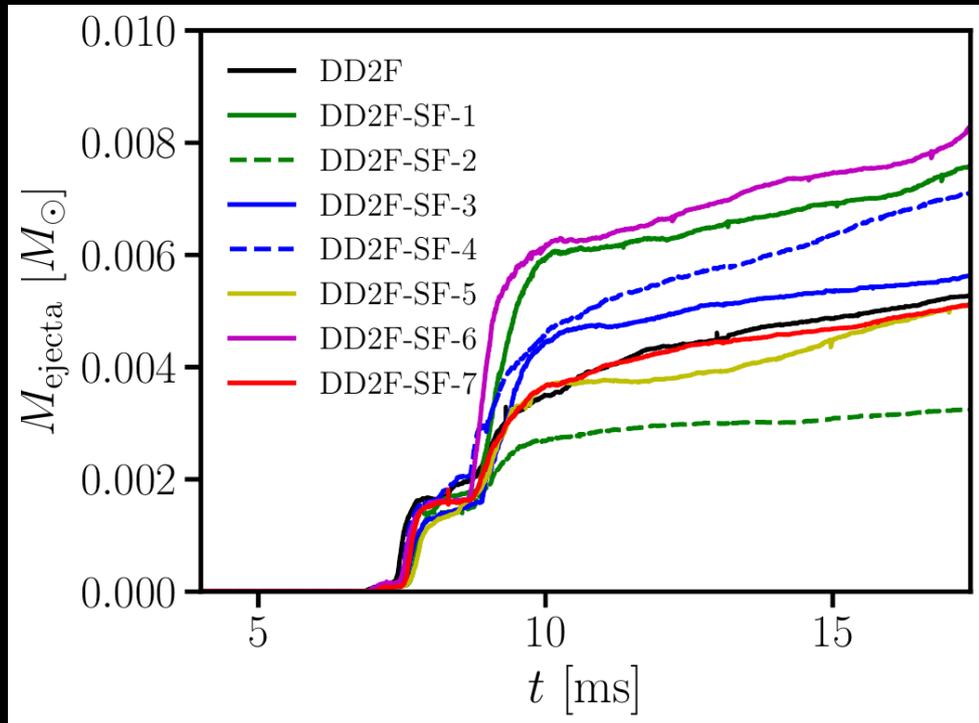
Dexheimer & Schramm 2010



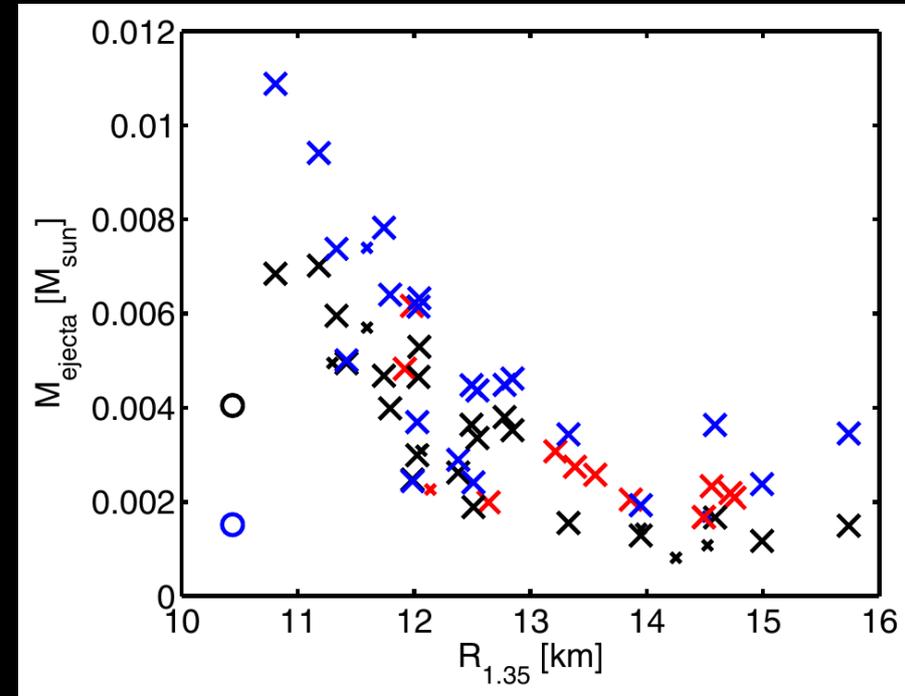
Most et. al 2019

Em counterpart / nucleosynthesi

- ▶ Electromagnetic transient powered by radioactive decays (during / after r-process)
 - quasi-thermal emission in UV, optical, infrared
- ▶ Different ejecta components: dynamical, disk ejecta
- ▶ No obvious qualitative differences – quantitative differences within expected “hadronic” scatter (simplistic considerations)
- ▶ More subtle impact possible, but unlikely (simple model wo neutrinos, network, disk evolution ...) - also other characteristic similar: outflow velocity, disk mass, ...



Bauswein et al 2019 – only dynamical ejecta



Bauswein et al 2013

Conclusions

- ▶ NS radius must be larger than 10.7 km (very robust and conservative), corresponds to $\Lambda > 200$
- ▶ More stringent constraints from future detections
- ▶ Limit on M_{\max} !
- ▶ NS radius measurable from dominant postmerger frequency
- ▶ Explicitly shown by GW data analysis
- ▶ Threshold binary mass for prompt collapse \rightarrow maximum mass M_{\max}
- ▶ Strong 1st order phase transitions leave characteristic imprint on GW (postmerger frequency higher than expected from inspiral)
- ▶ Complementarity of inspiral and postmerger phase \rightarrow postmerger probes higher density regime