

Using a nonlocal dispersive-optical-model to generate ingredients for ν -A cross sections

Mack C. Atkinson

Washington University in St. Louis

ECT* 2019

Using a nonlocal dispersive-optical-model to generate ingredients for ν -A cross sections

Willem Dickhoff
Bob Charity
Henk Blok
Louk Lapikás
Hossein Mahzoon
Cole Pruitt
Lee Sobotka

Mack C. Atkinson

Washington University in St. Louis

ECT* 2019

Introduction

- Goal is to help describe the nuclear aspects of ν -A interactions

Introduction

- Goal is to help describe the nuclear aspects of ν -A interactions
- The structure of large nuclei can be determined using a nonlocal dispersive optical model (DOM)

Introduction

- Goal is to help describe the nuclear aspects of ν -A interactions
- The structure of large nuclei can be determined using a nonlocal dispersive optical model (DOM)
- Experimental data is used to constrain the DOM

Introduction

- Goal is to help describe the nuclear aspects of ν -A interactions
- The structure of large nuclei can be determined using a nonlocal dispersive optical model (DOM)
- Experimental data is used to constrain the DOM
- The $(e, e' p)$ reaction can be described using the DOM

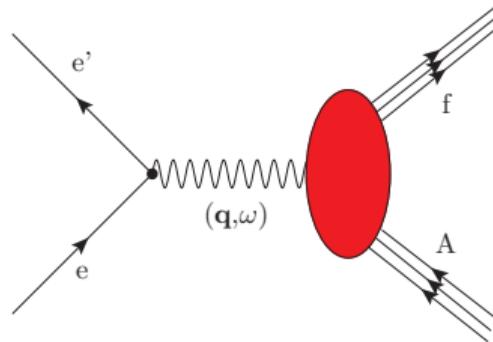
Introduction

- Goal is to help describe the nuclear aspects of ν -A interactions
- The structure of large nuclei can be determined using a nonlocal dispersive optical model (DOM)
- Experimental data is used to constrain the DOM
- The $(e, e' p)$ reaction can be described using the DOM
- This can be extended to different leptonic probes

Introduction

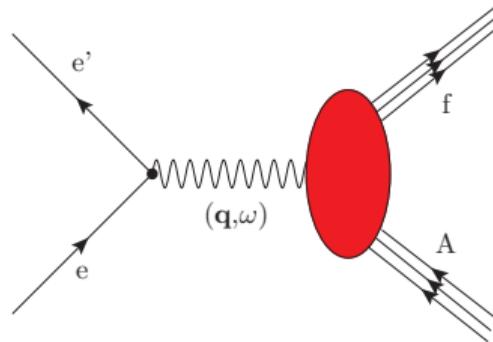
- Goal is to help describe the nuclear aspects of ν -A interactions
- The structure of large nuclei can be determined using a nonlocal dispersive optical model (DOM)
- Experimental data is used to constrain the DOM
- The $(e, e' p)$ reaction can be described using the DOM
- This can be extended to different leptonic probes
- In particular, a DOM analysis of ^{40}Ar (relevant for DUNE) is underway

Electron-Nucleus Scattering



Electron-Nucleus Scattering

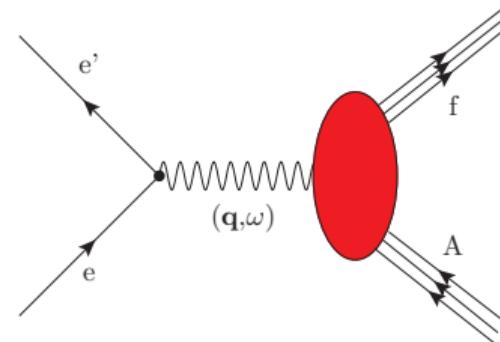
$$\frac{d^2\sigma}{d\Omega dE} = L_{\mu\nu} W^{\mu\nu}$$



Electron-Nucleus Scattering

$$\frac{d^2\sigma}{d\Omega dE} = L_{\mu\nu} W^{\mu\nu}$$

- Impulse Approximation (IA) \implies one-body current

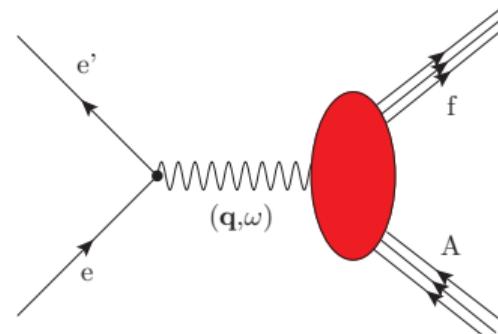


Electron-Nucleus Scattering

$$\frac{d^2\sigma}{d\Omega dE} = L_{\mu\nu} W^{\mu\nu}$$

- Impulse Approximation (IA) \implies one-body current

$$W^{\mu\nu} \propto \text{Im}\Pi(\mathbf{q}, \omega) \quad (\text{Polarization Propagator})$$



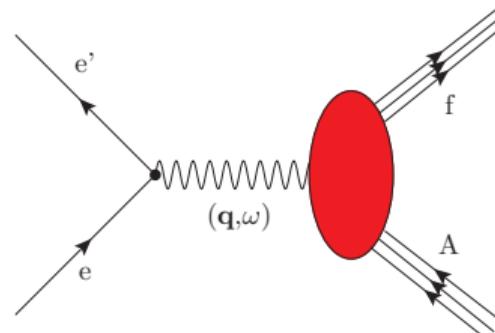
Electron-Nucleus Scattering

$$\frac{d^2\sigma}{d\Omega dE} = L_{\mu\nu} W^{\mu\nu}$$

- Impulse Approximation (IA) \implies one-body current

$$W^{\mu\nu} \propto \text{Im}\Pi(\mathbf{q}, \omega) \quad (\text{Polarization Propagator})$$

Fermi-Gas



Electron-Nucleus Scattering

$$\frac{d^2\sigma}{d\Omega dE} = L_{\mu\nu} W^{\mu\nu}$$

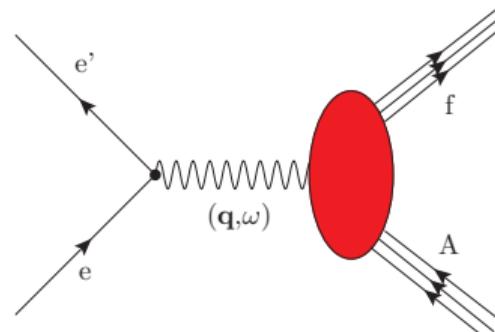
- Impulse Approximation (IA) \implies one-body current

$$W^{\mu\nu} \propto \text{Im}\Pi(\mathbf{q}, \omega) \quad (\text{Polarization Propagator})$$

Fermi-Gas



PWIA



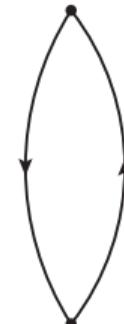
Electron-Nucleus Scattering

$$\frac{d^2\sigma}{d\Omega dE} = L_{\mu\nu} W^{\mu\nu}$$

- Impulse Approximation (IA) \implies one-body current

$W^{\mu\nu} \propto \text{Im}\Pi(\mathbf{q}, \omega)$ (Polarization Propagator)

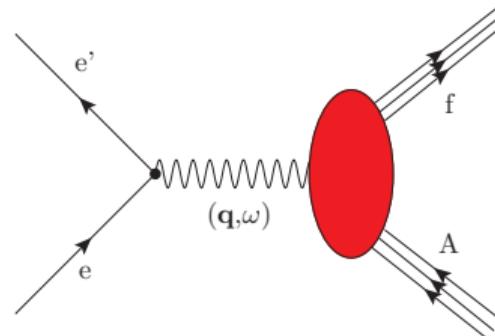
Fermi-Gas



PWIA



DWIA



Single-Particle Propagator and the Dyson Equation

$$G_{\ell j}(r, r'; E) = \sum_m \frac{\langle \Psi_0^A | a_{r\ell j} | \Psi_m^{A+1} \rangle \langle \Psi_m^{A+1} | a_{r'\ell j}^\dagger | \Psi_0^A \rangle}{E - (E_m^{A+1} - E_0^A) + i\eta} \\ + \sum_n \frac{\langle \Psi_0^A | a_{r'\ell j}^\dagger | \Psi_n^{A-1} \rangle \langle \Psi_n^{A-1} | a_{r\ell j} | \Psi_0^A \rangle}{E - (E_0^A - E_n^{A-1}) - i\eta}$$

Single-Particle Propagator and the Dyson Equation

$$G_{\ell j}(r, r'; E) = \sum_m \frac{\langle \Psi_0^A | a_{r\ell j} | \Psi_m^{A+1} \rangle \langle \Psi_m^{A+1} | a_{r'\ell j}^\dagger | \Psi_0^A \rangle}{E - (E_m^{A+1} - E_0^A) + i\eta} \\ + \sum_n \frac{\langle \Psi_0^A | a_{r'\ell j}^\dagger | \Psi_n^{A-1} \rangle \langle \Psi_n^{A-1} | a_{r\ell j} | \Psi_0^A \rangle}{E - (E_0^A - E_n^{A-1}) - i\eta}$$

- Poles correspond to excitation energies of $(A + 1)$ or $(A - 1)$ nucleus

Single-Particle Propagator and the Dyson Equation

$$G_{\ell j}(r, r'; E) = \sum_m \frac{\langle \Psi_0^A | a_{r\ell j} | \Psi_m^{A+1} \rangle \langle \Psi_m^{A+1} | a_{r'\ell j}^\dagger | \Psi_0^A \rangle}{E - (E_m^{A+1} - E_0^A) + i\eta} \\ + \sum_n \frac{\langle \Psi_0^A | a_{r'\ell j}^\dagger | \Psi_n^{A-1} \rangle \langle \Psi_n^{A-1} | a_{r\ell j} | \Psi_0^A \rangle}{E - (E_0^A - E_n^{A-1}) - i\eta}$$

- Poles correspond to excitation energies of $(A + 1)$ or $(A - 1)$ nucleus
- Numerator like a transition probability to given excitation

Single-Particle Propagator and the Dyson Equation

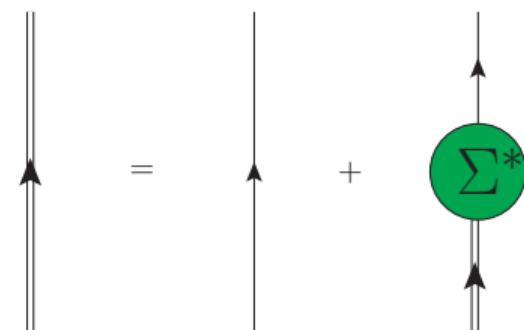
$$G_{\ell j}(r, r'; E) = \sum_m \frac{\langle \Psi_0^A | a_{r\ell j} | \Psi_m^{A+1} \rangle \langle \Psi_m^{A+1} | a_{r'\ell j}^\dagger | \Psi_0^A \rangle}{E - (E_m^{A+1} - E_0^A) + i\eta} \\ + \sum_n \frac{\langle \Psi_0^A | a_{r'\ell j}^\dagger | \Psi_n^{A-1} \rangle \langle \Psi_n^{A-1} | a_{r\ell j} | \Psi_0^A \rangle}{E - (E_0^A - E_n^{A-1}) - i\eta}$$

- Poles correspond to excitation energies of $(A + 1)$ or $(A - 1)$ nucleus
- Numerator like a transition probability to given excitation
- Close connection with experimental observables

Single-Particle Propagator and the Dyson Equation

$$G_{\ell j}(r, r'; E) = \sum_m \frac{\langle \Psi_0^A | a_{r\ell j} | \Psi_m^{A+1} \rangle \langle \Psi_m^{A+1} | a_{r'\ell j}^\dagger | \Psi_0^A \rangle}{E - (E_m^{A+1} - E_0^A) + i\eta} \\ + \sum_n \frac{\langle \Psi_0^A | a_{r'\ell j}^\dagger | \Psi_n^{A-1} \rangle \langle \Psi_n^{A-1} | a_{r\ell j} | \Psi_0^A \rangle}{E - (E_0^A - E_n^{A-1}) - i\eta}$$

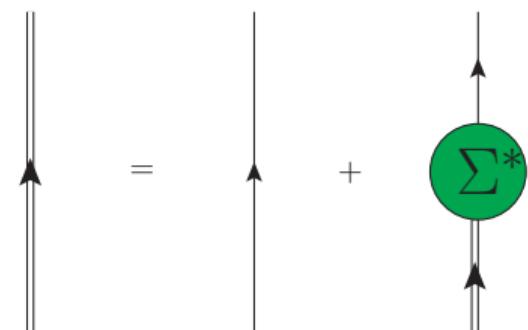
- Poles correspond to excitation energies of $(A + 1)$ or $(A - 1)$ nucleus
- Numerator like a transition probability to given excitation
- Close connection with experimental observables
- Perturbation expansion of G leads to the Dyson equation



Single-Particle Propagator and the Dyson Equation

$$G_{\ell j}(r, r'; E) = \sum_m \frac{\langle \Psi_0^A | a_{r\ell j} | \Psi_m^{A+1} \rangle \langle \Psi_m^{A+1} | a_{r'\ell j}^\dagger | \Psi_0^A \rangle}{E - (E_m^{A+1} - E_0^A) + i\eta} \\ + \sum_n \frac{\langle \Psi_0^A | a_{r'\ell j}^\dagger | \Psi_n^{A-1} \rangle \langle \Psi_n^{A-1} | a_{r\ell j} | \Psi_0^A \rangle}{E - (E_0^A - E_n^{A-1}) - i\eta}$$

- Poles correspond to excitation energies of $(A + 1)$ or $(A - 1)$ nucleus
- Numerator like a transition probability to given excitation
- Close connection with experimental observables
- Perturbation expansion of G leads to the Dyson equation
- If the irreducible self-energy (Σ^*) is known, then so is G



Spectral function

$$S_{\ell j}^h(r; E) = \frac{1}{\pi} \text{Im} G_{\ell j}(r, r; E) \theta(E - (E_0^A - E_0^{A-1}))$$

Spectral function

$$S_{\ell j}^h(r; E) = \frac{1}{\pi} \text{Im} G_{\ell j}(r, r; E) \theta(E - (E_0^A - E_0^{A-1}))$$

- Reveals effects of many-body correlations beyond the independent particle model

Spectral function

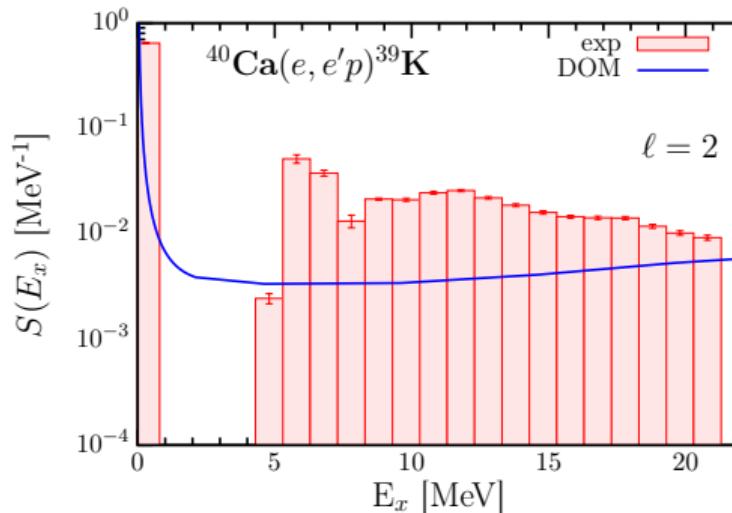
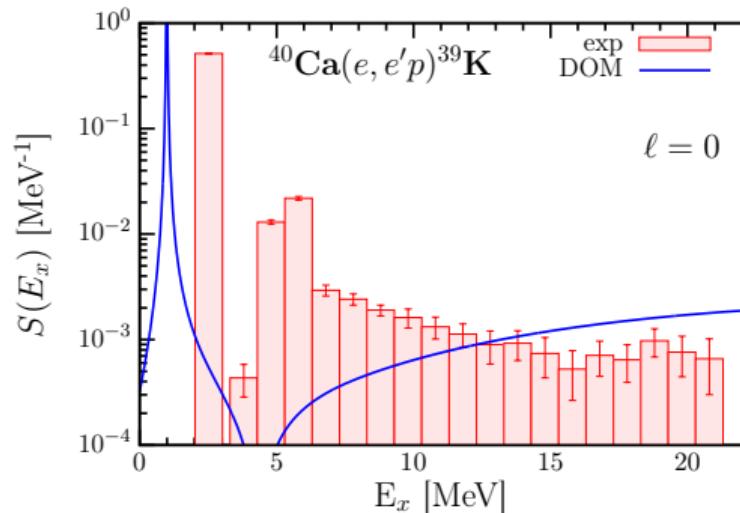
$$S_{\ell j}^h(r; E) = \frac{1}{\pi} \text{Im} G_{\ell j}(r, r; E) \theta(E - (E_0^A - E_0^{A-1}))$$

- Reveals effects of many-body correlations beyond the independent particle model
- Can be observed with excitation spectrum from knockout reactions

Spectral function

$$S_{\ell j}^h(r; E) = \frac{1}{\pi} \text{Im} G_{\ell j}(r, r; E) \theta(E - (E_0^A - E_0^{A-1}))$$

- Reveals effects of many-body correlations beyond the independent particle model
- Can be observed with excitation spectrum from knockout reactions



The Dispersive Optical Model (DOM)

- Irreducible self-energy at positive energies corresponds to an optical potential

The Dispersive Optical Model (DOM)

- Irreducible self-energy at positive energies corresponds to an optical potential
- Use same functional form as standard optical potentials to parametrize self-energy

The Dispersive Optical Model (DOM)

- Irreducible self-energy at positive energies corresponds to an optical potential
- Use same functional form as standard optical potentials to parametrize self-energy
- $\Sigma^*(\mathbf{r}, \mathbf{r}'; E)$ is explicitly **nonlocal**

The Dispersive Optical Model (DOM)

- Irreducible self-energy at positive energies corresponds to an optical potential
- Use same functional form as standard optical potentials to parametrize self-energy
- $\Sigma^*(\mathbf{r}, \mathbf{r}'; E)$ is explicitly **nonlocal**
- Dispersion relation connects to negative energies

The Dispersive Optical Model (DOM)

- Irreducible self-energy at positive energies corresponds to an optical potential
- Use same functional form as standard optical potentials to parametrize self-energy
- $\Sigma^*(r, r'; E)$ is explicitly **nonlocal**
- Dispersion relation connects to negative energies

Dispersive Correction

$$\begin{aligned} Re\Sigma_{\ell j}(r, r'; E) = & Re\Sigma_{\ell j}(r, r'; \epsilon_F) - \frac{1}{\pi}(\epsilon_F - E)\mathcal{P} \int_{\epsilon_T^+}^{\infty} dE' Im\Sigma_{\ell j}(r, r'; E')[\frac{1}{E - E'} - \frac{1}{\epsilon_F - E'}] \\ & + \frac{1}{\pi}(\epsilon_F - E)\mathcal{P} \int_{-\infty}^{\epsilon_T^-} dE' Im\Sigma_{\ell j}(r, r'; E')[\frac{1}{E - E'} - \frac{1}{\epsilon_F - E'}] \end{aligned}$$

The Dispersive Optical Model (DOM)

- Irreducible self-energy at positive energies corresponds to an optical potential
- Use same functional form as standard optical potentials to parametrize self-energy
- $\Sigma^*(r, r'; E)$ is explicitly **nonlocal**
- Dispersion relation connects to negative energies

Dispersive Correction

$$\begin{aligned} \text{Re}\Sigma_{\ell j}(r, r'; E) = & \text{Re}\Sigma_{\ell j}(r, r'; \epsilon_F) - \frac{1}{\pi}(\epsilon_F - E)\mathcal{P} \int_{\epsilon_T^+}^{\infty} dE' \text{Im}\Sigma_{\ell j}(r, r'; E') \left[\frac{1}{E - E'} - \frac{1}{\epsilon_F - E'} \right] \\ & + \frac{1}{\pi}(\epsilon_F - E)\mathcal{P} \int_{-\infty}^{\epsilon_T^-} dE' \text{Im}\Sigma_{\ell j}(r, r'; E') \left[\frac{1}{E - E'} - \frac{1}{\epsilon_F - E'} \right] \end{aligned}$$

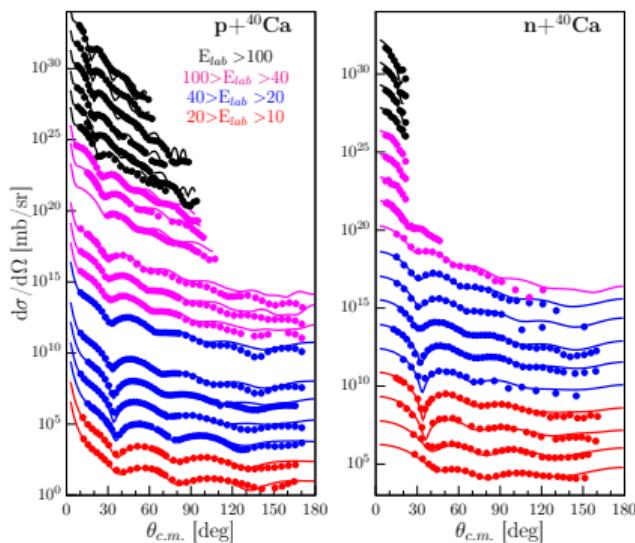
- This constraint ensures bound and scattering quantities are simultaneously described

Fitting the Self-energy (^{40}Ca)

- Parameters of self-energy varied to minimize χ^2

Fitting the Self-energy (^{40}Ca)

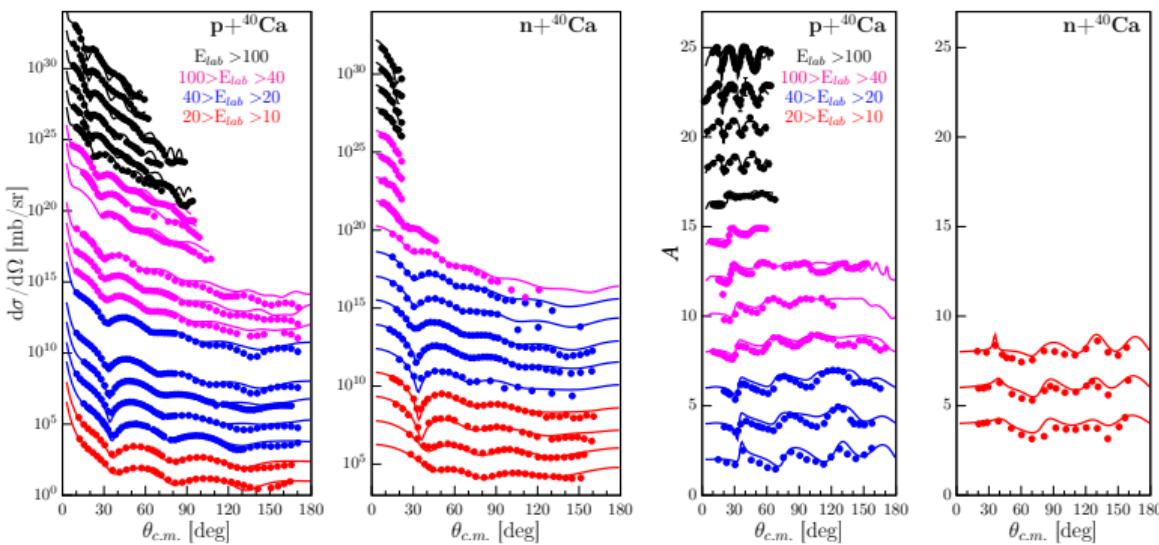
- Parameters of self-energy varied to minimize χ^2



Data: J.M. Mueller et al. *Phys. Rev. C*, **83** 064605, 2011

Fitting the Self-energy (^{40}Ca)

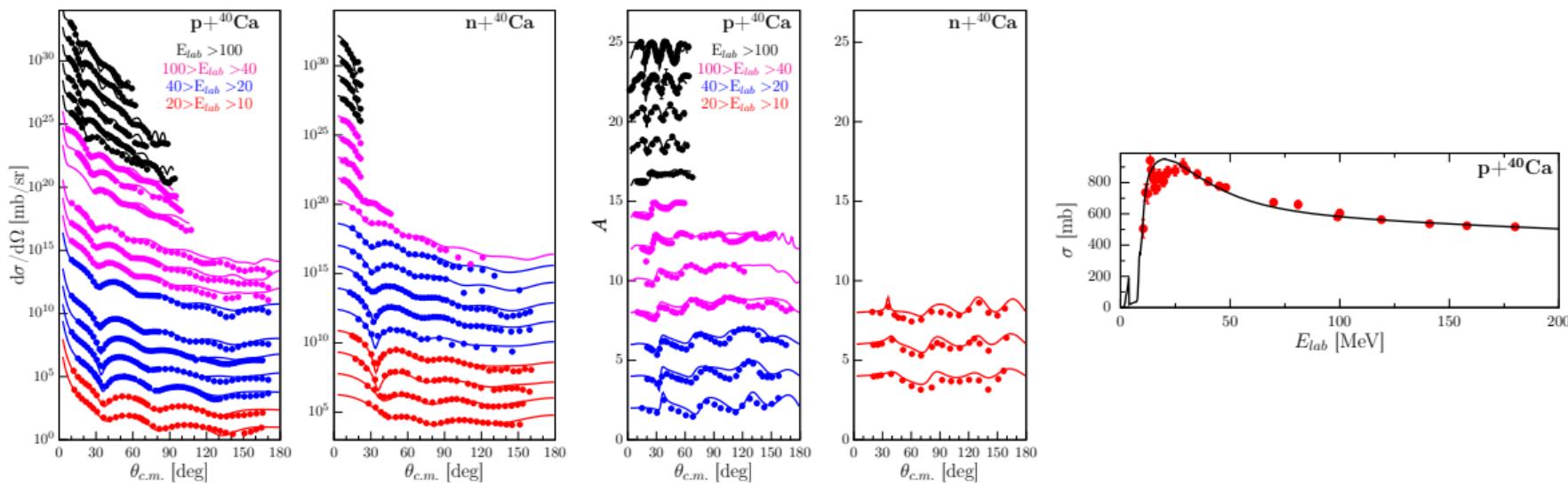
- Parameters of self-energy varied to minimize χ^2



Data: J.M. Mueller et al. *Phys. Rev. C*, **83** 064605, 2011

Fitting the Self-energy (^{40}Ca)

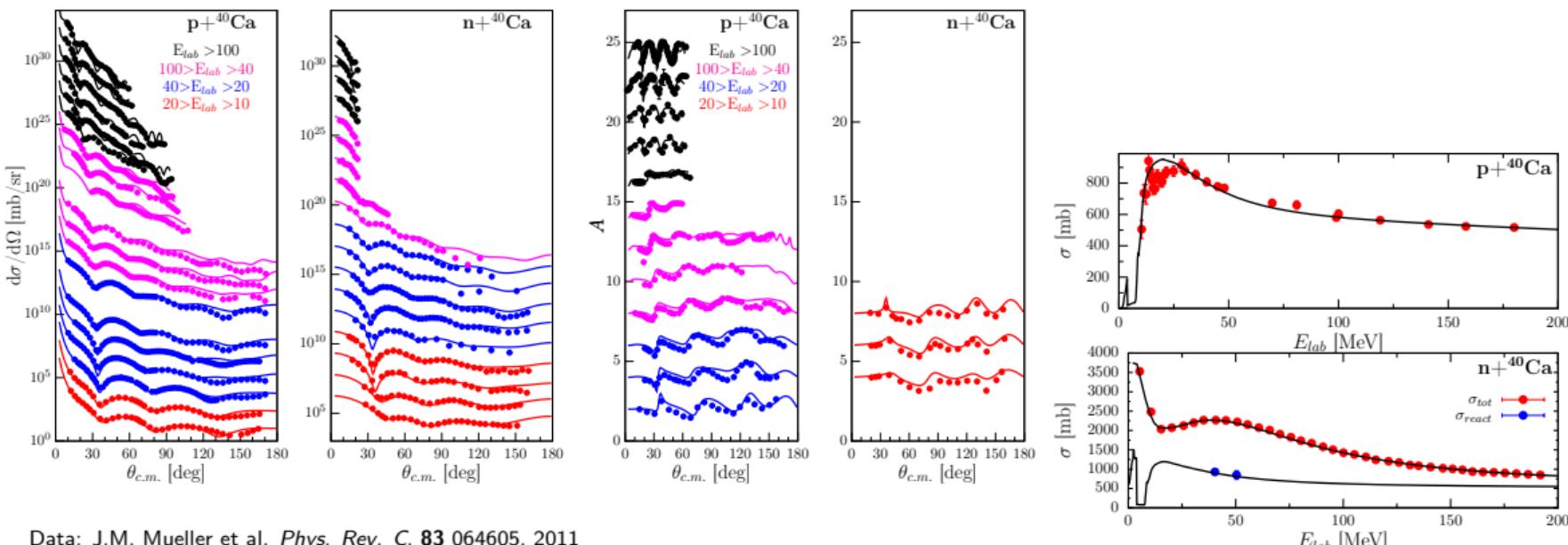
- Parameters of self-energy varied to minimize χ^2



Data: J.M. Mueller et al. *Phys. Rev. C*, **83** 064605, 2011

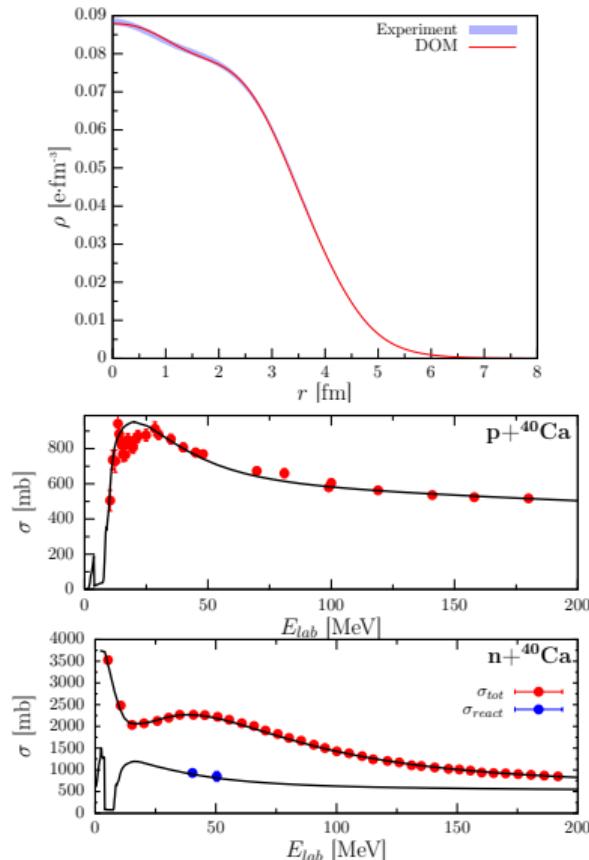
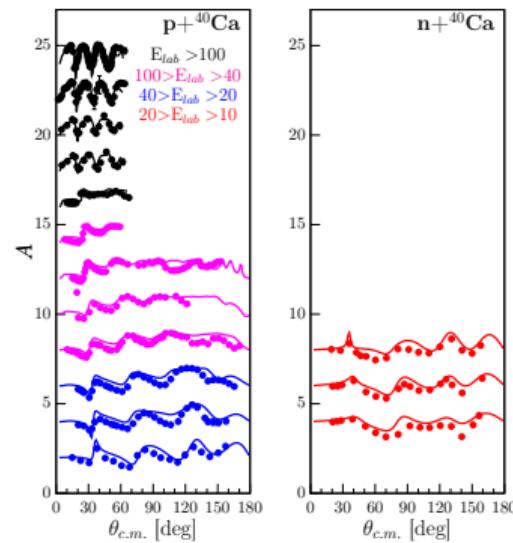
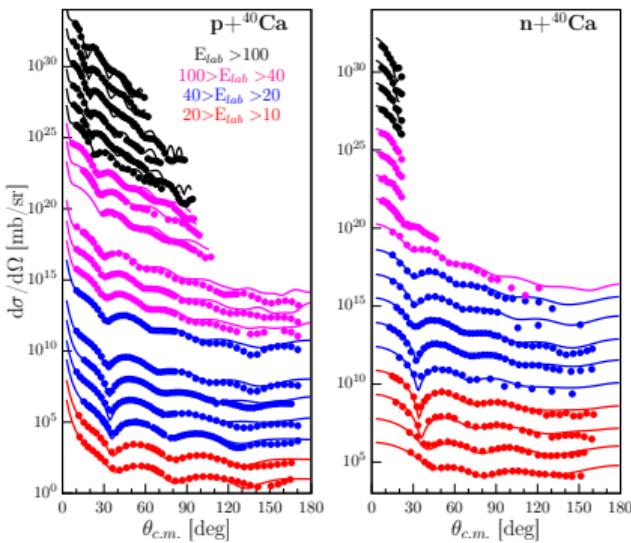
Fitting the Self-energy (^{40}Ca)

- Parameters of self-energy varied to minimize χ^2



Fitting the Self-energy (^{40}Ca)

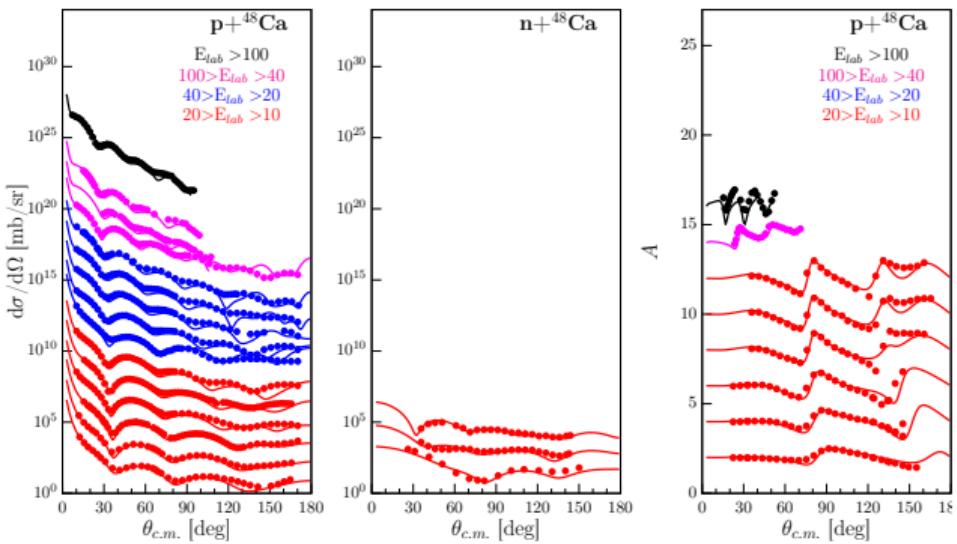
- Parameters of self-energy varied to minimize χ^2
- Reproducing the data means self-energy is found



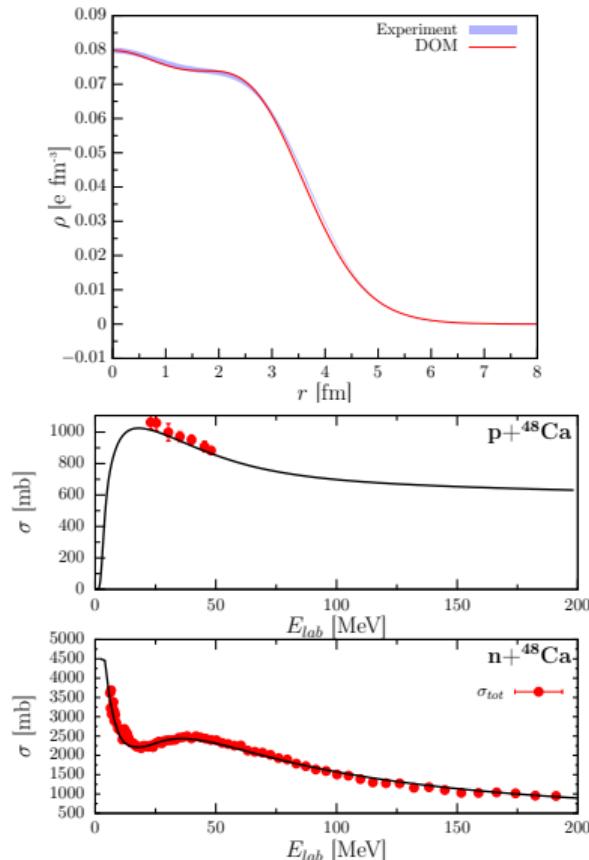
Data: J.M. Mueller et al. *Phys. Rev. C*, **83** 064605, 2011

Fitting the Self-energy (^{48}Ca)

- Parameters of self-energy varied to minimize χ^2
- Reproducing the data means self-energy is found



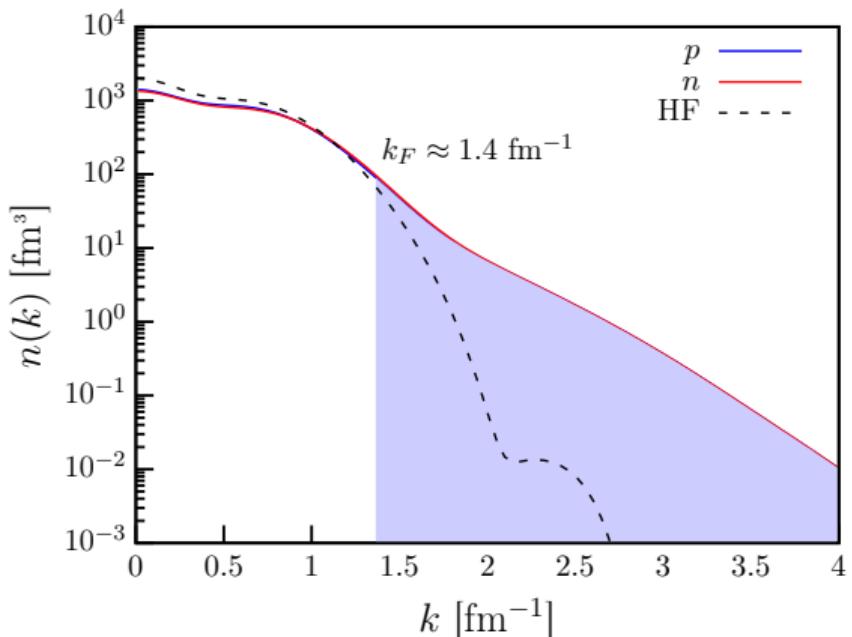
Data: J.M. Mueller et al. *Phys. Rev. C*, **83** 064605, 2011



Momentum Distributions

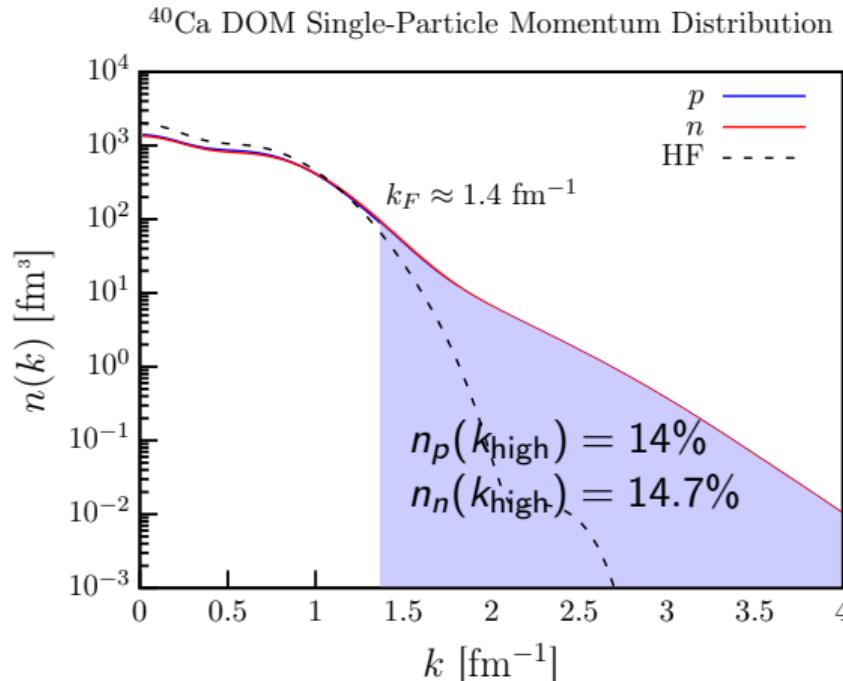
$$n(\mathbf{k}) = \int d^3r \int d^3r' e^{i\mathbf{k}\cdot(\mathbf{r}-\mathbf{r}')} \rho(\mathbf{r}, \mathbf{r}')$$

^{40}Ca DOM Single-Particle Momentum Distribution



Momentum Distributions

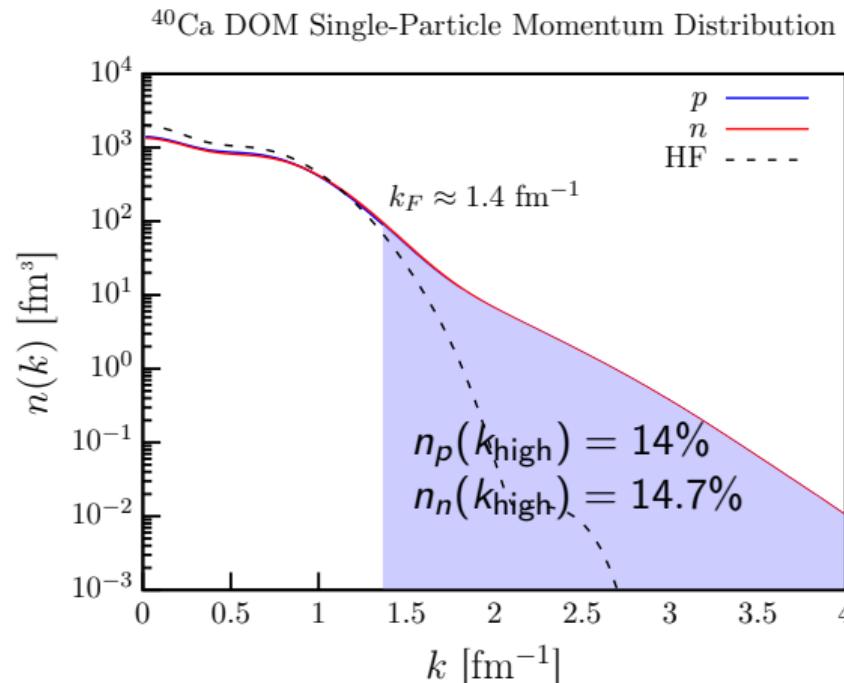
$$n(\mathbf{k}) = \int d^3r \int d^3r' e^{i\mathbf{k}\cdot(\mathbf{r}-\mathbf{r}')} \rho(\mathbf{r}, \mathbf{r}')$$



Momentum Distributions

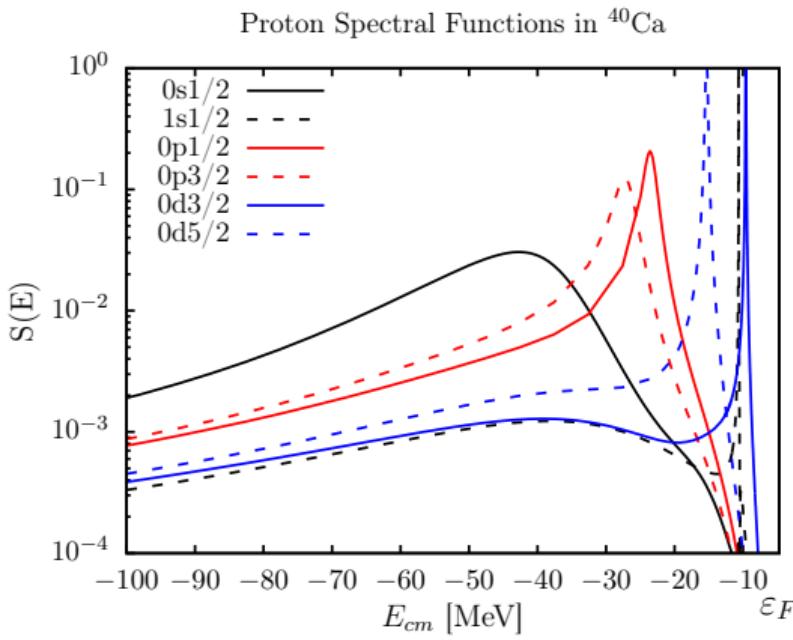
- Short-range correlations (SRC) responsible for this high-momentum content

$$n(\mathbf{k}) = \int d^3r \int d^3r' e^{i\mathbf{k}\cdot(\mathbf{r}-\mathbf{r}')} \rho(\mathbf{r}, \mathbf{r}')$$

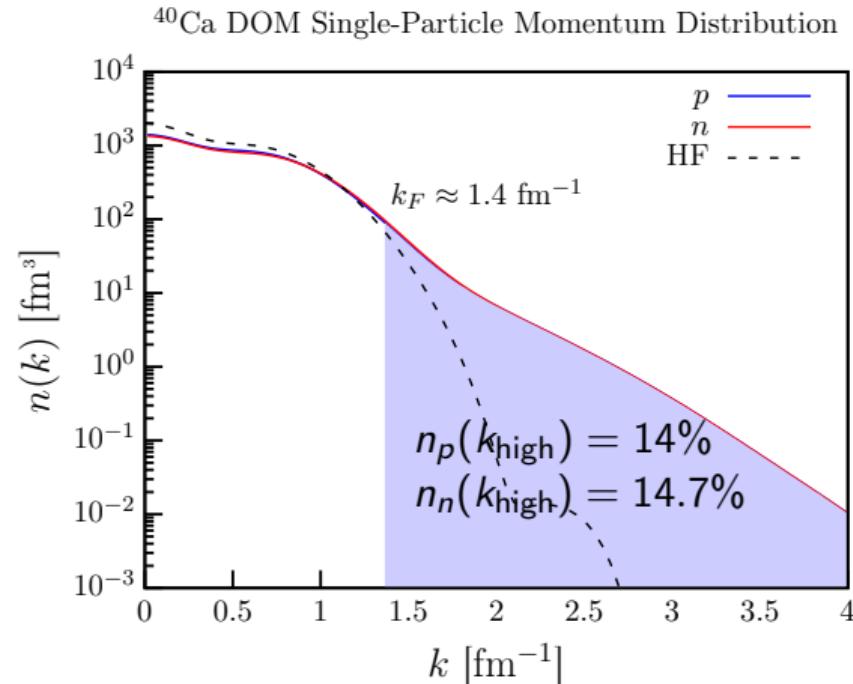


Momentum Distributions

- Short-range correlations (SRC) responsible for this high-momentum content

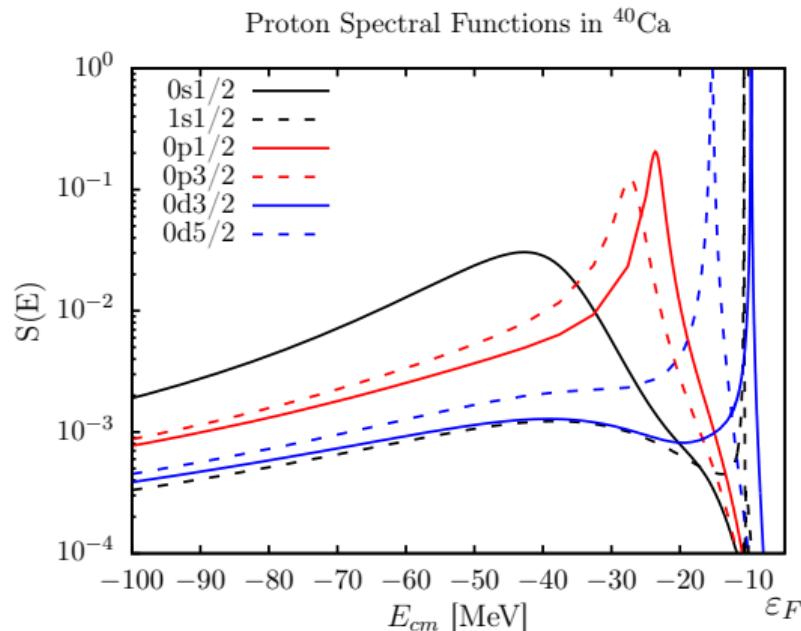


$$n(\mathbf{k}) = \int d^3r \int d^3r' e^{i\mathbf{k} \cdot (\mathbf{r} - \mathbf{r}')} \rho(\mathbf{r}, \mathbf{r}')$$



Constraints for Spectral Function

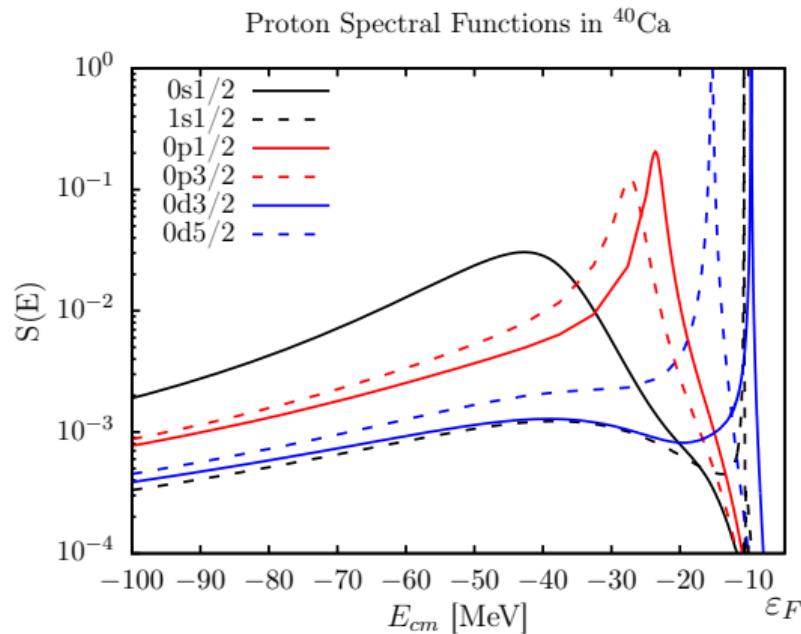
$$S^h(\alpha, \beta; E) = \frac{1}{\pi} \text{Im}\{G(\alpha, \beta; E)\}$$
$$S^h(E) = \sum_{\alpha} S(\alpha, \alpha; E)$$



Constraints for Spectral Function

$$S^h(\alpha, \beta; E) = \frac{1}{\pi} \text{Im}\{G(\alpha, \beta; E)\}$$
$$S^h(E) = \sum_{\alpha} S(\alpha, \alpha; E)$$

$$\rho_{\alpha, \beta} = \int_{-\infty}^{\varepsilon_F} dE S(\alpha, \beta; E)$$

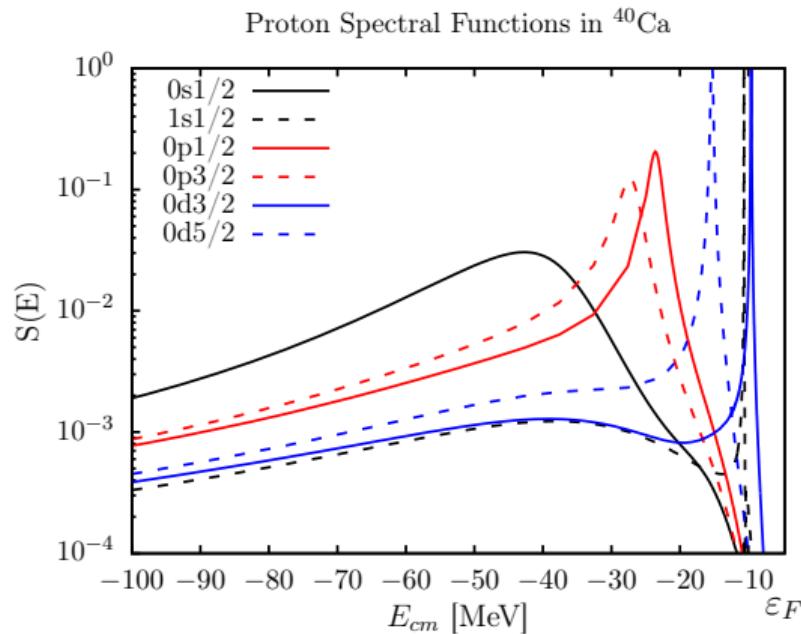


Constraints for Spectral Function

$$S^h(\alpha, \beta; E) = \frac{1}{\pi} \text{Im}\{G(\alpha, \beta; E)\}$$

$$S^h(E) = \sum_{\alpha} S(\alpha, \alpha; E)$$

$$\rho_{\alpha, \beta} = \int_{-\infty}^{\varepsilon_F} dE S(\alpha, \beta; E) \quad N, Z = \sum_{\alpha} \rho_{\alpha, \alpha}^{N, Z}$$



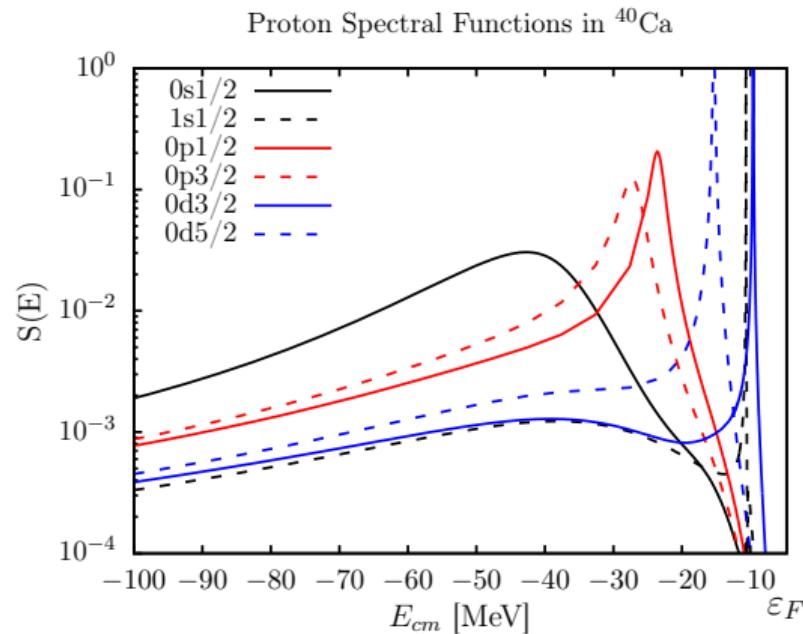
Constraints for Spectral Function

$$S^h(\alpha, \beta; E) = \frac{1}{\pi} \text{Im}\{G(\alpha, \beta; E)\}$$

$$S^h(E) = \sum_{\alpha} S(\alpha, \alpha; E)$$

$$\rho_{\alpha, \beta} = \int_{-\infty}^{\varepsilon_F} dE S(\alpha, \beta; E) \quad N, Z = \sum_{\alpha} \rho_{\alpha, \alpha}^{N, Z}$$

$$E_0^A = \frac{1}{2} \sum_{\alpha \beta} \left[T_{\beta \alpha} \rho_{\alpha \beta} + \delta_{\alpha \beta} \int_{-\infty}^{\varepsilon_F^-} dE E S_h(\alpha; E) \right]$$



Constraints for Spectral Function

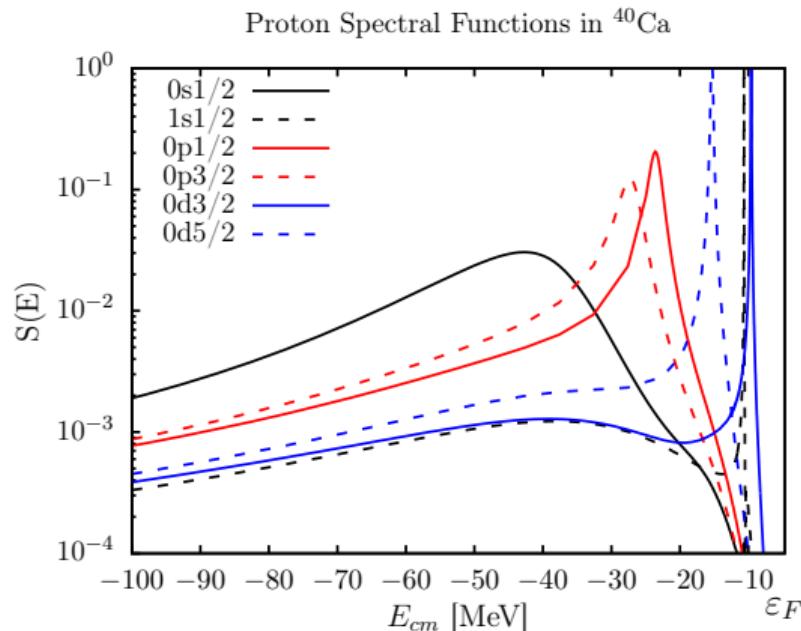
$$S^h(\alpha, \beta; E) = \frac{1}{\pi} \text{Im}\{G(\alpha, \beta; E)\}$$

$$S^h(E) = \sum_{\alpha} S(\alpha, \alpha; E)$$

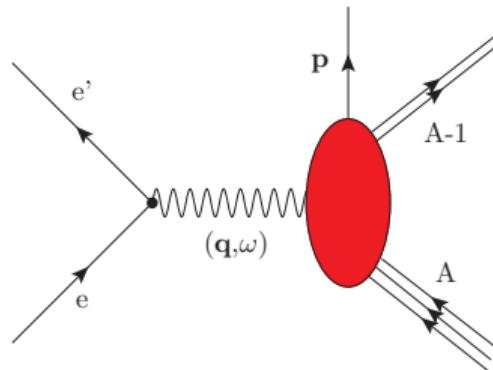
$$\rho_{\alpha, \beta} = \int_{-\infty}^{\varepsilon_F} dE S(\alpha, \beta; E) \quad N, Z = \sum_{\alpha} \rho_{\alpha, \alpha}^{N, Z}$$

$$E_0^A = \frac{1}{2} \sum_{\alpha\beta} \left[T_{\beta\alpha} \rho_{\alpha\beta} + \delta_{\alpha\beta} \int_{-\infty}^{\varepsilon_F^-} dE E S_h(\alpha; E) \right]$$

	N	Z	DOM E_0^A/A	Exp. E_0^A/A
^{40}Ca	19.9	19.8	-8.49	-8.55
^{48}Ca	27.9	19.9	-8.7	-8.66

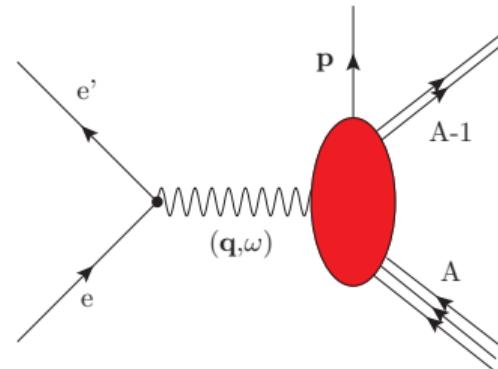


DWIA for exclusive $^{40}\text{Ca}(e, e' p)^{39}\text{K}$ reaction



DWIA for exclusive $^{40}\text{Ca}(e, e' p)^{39}\text{K}$ reaction

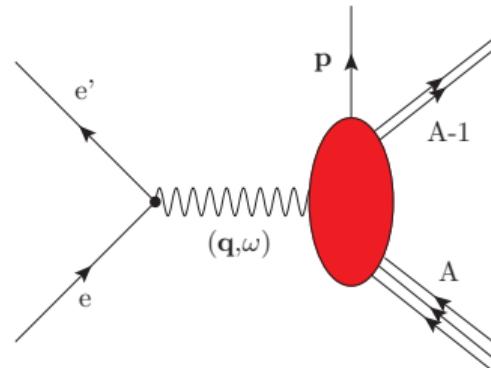
- DWIA for exclusive reaction (C. Giusti's DWEEPY code)



DWIA for exclusive $^{40}\text{Ca}(e, e' p)^{39}\text{K}$ reaction

- DWIA for exclusive reaction (C. Giusti's DWEEPY code)

$$J^\mu(\mathbf{q}) = \int \chi_{E\alpha}^{(-)*}(\mathbf{r}) j^\mu(\mathbf{r}) \phi_{E\alpha}(\mathbf{r}) [\mathcal{Z}_\alpha(E)]^{1/2} e^{i\mathbf{q}\cdot\mathbf{r}} d^3 r$$

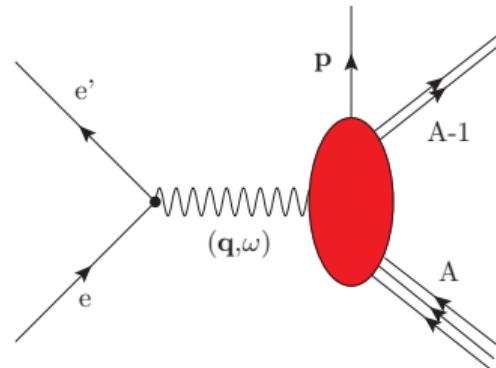


DWIA for exclusive $^{40}\text{Ca}(e, e' p)^{39}\text{K}$ reaction

- DWIA for exclusive reaction (C. Giusti's DWEEPY code)

$$J^\mu(\mathbf{q}) = \int \chi_{E\alpha}^{(-)*}(\mathbf{r}) j^\mu(\mathbf{r}) \phi_{E\alpha}(\mathbf{r}) [\mathcal{Z}_\alpha(E)]^{1/2} e^{i\mathbf{q}\cdot\mathbf{r}} d^3 r$$

- Spectroscopic factor, \mathcal{Z} , quantifies correlations

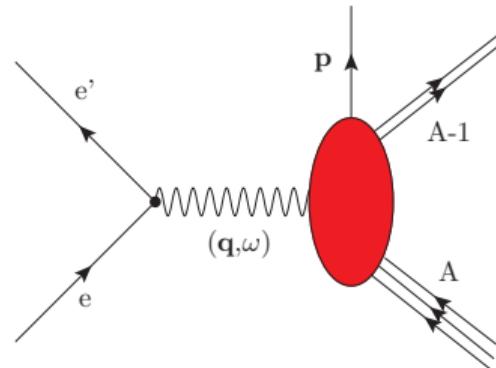


DWIA for exclusive $^{40}\text{Ca}(e, e' p)^{39}\text{K}$ reaction

- DWIA for exclusive reaction (C. Giusti's DWEEPY code)

$$J^\mu(\mathbf{q}) = \int \chi_{E\alpha}^{(-)*}(\mathbf{r}) j^\mu(\mathbf{r}) \phi_{E\alpha}(\mathbf{r}) [\mathcal{Z}_\alpha(E)]^{1/2} e^{i\mathbf{q}\cdot\mathbf{r}} d^3 r$$

- Spectroscopic factor, \mathcal{Z} , quantifies correlations
- DOM provides all ingredients

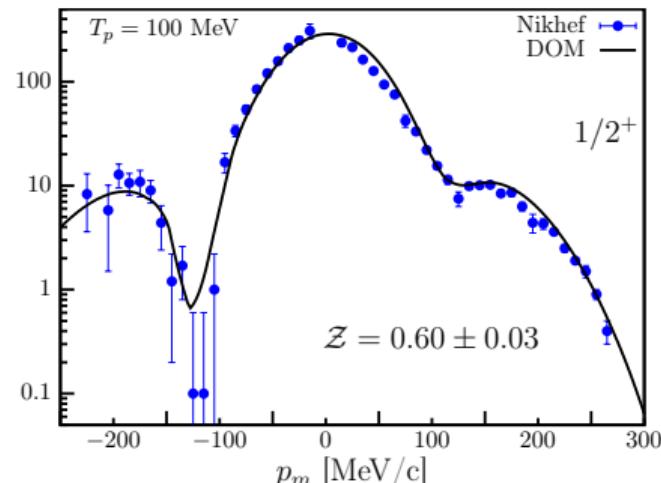
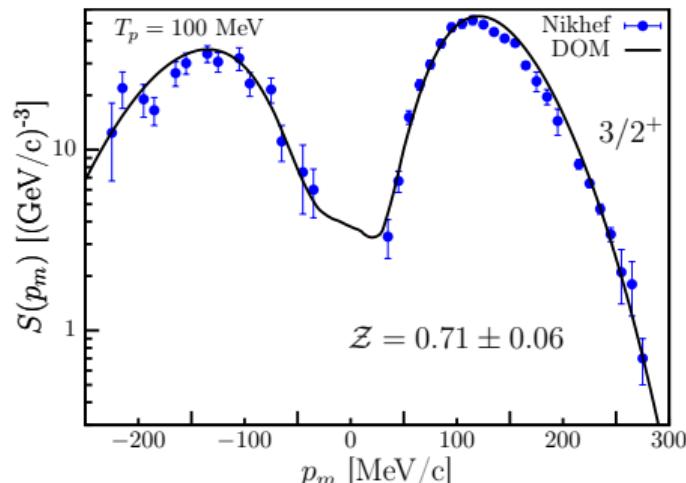
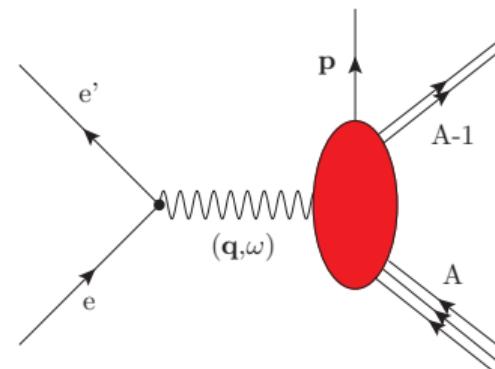


DWIA for exclusive $^{40}\text{Ca}(e, e' p)^{39}\text{K}$ reaction

- DWIA for exclusive reaction (C. Giusti's DWEEPY code)

$$J^\mu(\mathbf{q}) = \int \chi_{E\alpha}^{(-)*}(\mathbf{r}) j^\mu(\mathbf{r}) \phi_{E\alpha}(\mathbf{r}) [\mathcal{Z}_\alpha(E)]^{1/2} e^{i\mathbf{q}\cdot\mathbf{r}} d^3 r$$

- Spectroscopic factor, \mathcal{Z} , quantifies correlations
- DOM provides all ingredients

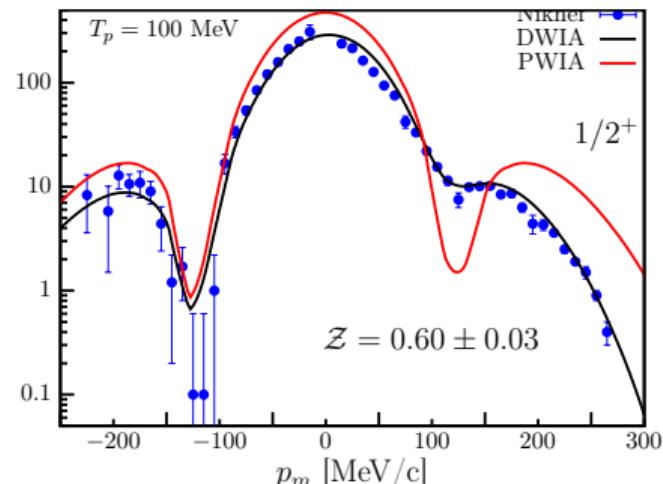
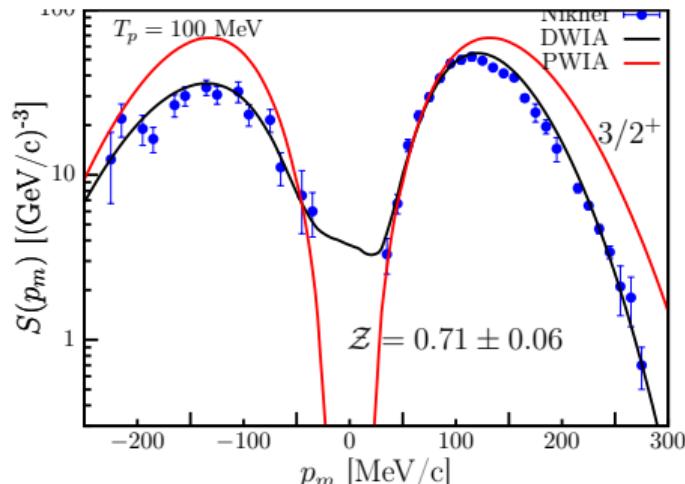
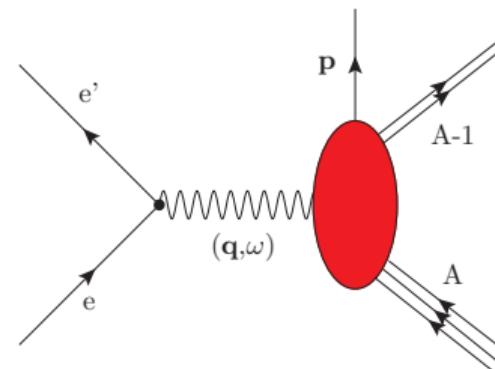


DWIA for exclusive $^{40}\text{Ca}(e, e' p)^{39}\text{K}$ reaction

- DWIA for exclusive reaction (C. Giusti's DWEEPY code)

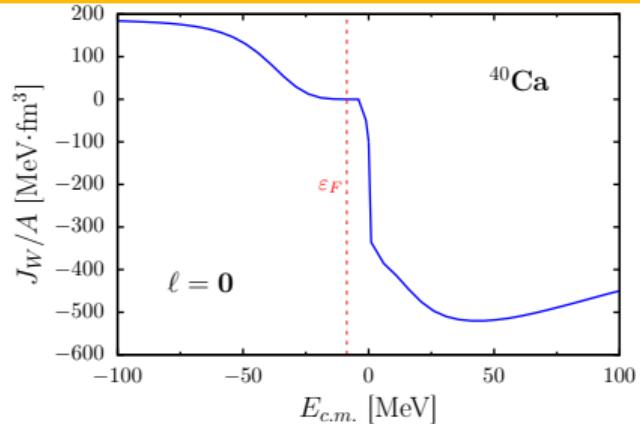
$$J^\mu(\mathbf{q}) = \int \chi_{E\alpha}^{(-)*}(\mathbf{r}) j^\mu(\mathbf{r}) \phi_{E\alpha}(\mathbf{r}) [\mathcal{Z}_\alpha(E)]^{1/2} e^{i\mathbf{q}\cdot\mathbf{r}} d^3 r$$

- Spectroscopic factor, \mathcal{Z} , quantifies correlations
- DOM provides all ingredients



Spectroscopic factor, Occupation, and Depletion

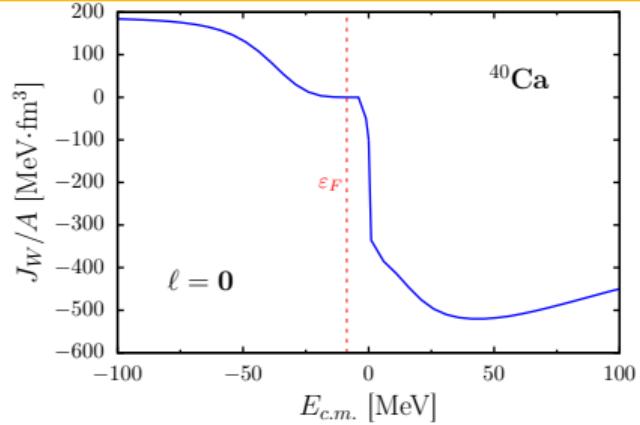
- No imaginary component of Σ^* around ϵ_F



Spectroscopic factor, Occupation, and Depletion

- No imaginary component of Σ^* around ϵ_F
- Spectroscopic factor for states near ϵ_F is well defined from Σ^*

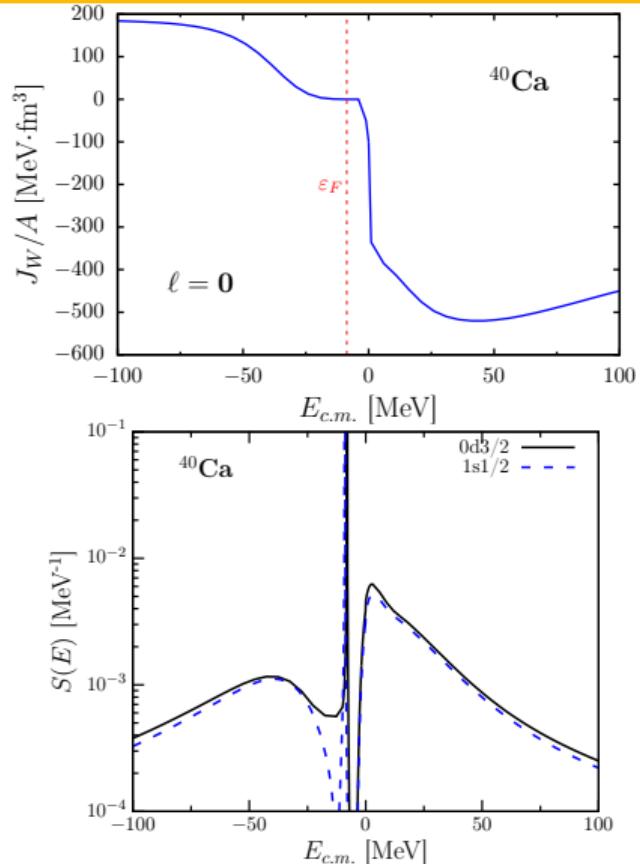
$$\mathcal{Z} = \left(1 - \frac{\partial \Sigma^*(\alpha_{qh}, \alpha_{qh}; E)}{\partial E} \Big|_{\epsilon} \right)^{-1}$$



Spectroscopic factor, Occupation, and Depletion

- No imaginary component of Σ^* around ϵ_F
- Spectroscopic factor for states near ϵ_F is well defined from Σ^*

$$\mathcal{Z} = \left(1 - \frac{\partial \Sigma^*(\alpha_{qh}, \alpha_{qh}; E)}{\partial E} \Big|_{\epsilon} \right)^{-1}$$

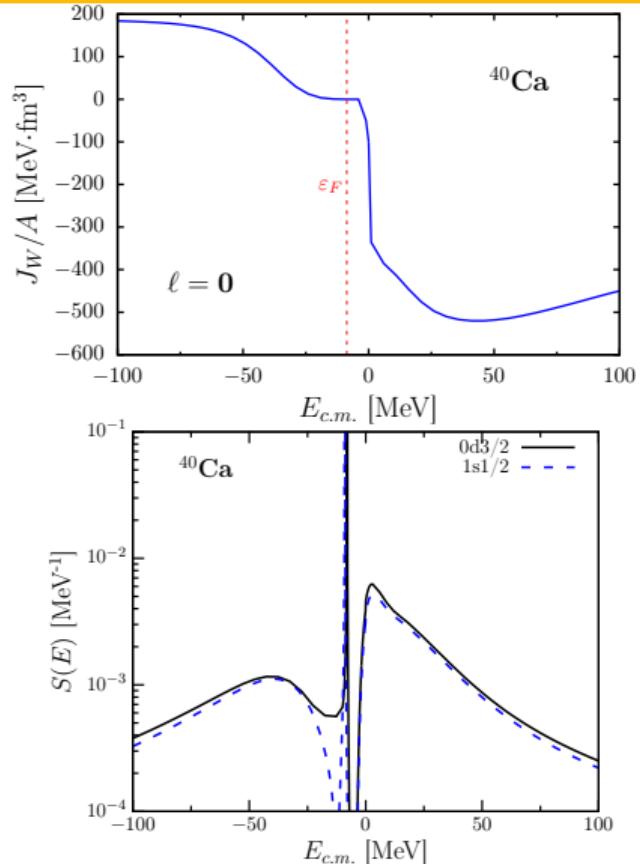


Spectroscopic factor, Occupation, and Depletion

- No imaginary component of Σ^* around ϵ_F
- Spectroscopic factor for states near ϵ_F is well defined from Σ^*

$$\mathcal{Z} = \left(1 - \frac{\partial \Sigma^*(\alpha_{qh}, \alpha_{qh}; E)}{\partial E} \Big|_{\epsilon} \right)^{-1}$$

$$n_{n\ell j} = \int_{-\infty}^{\epsilon_f} dE S_{n\ell j}^h(E) \quad d_{n\ell j} = \int_{\epsilon_f}^{\infty} dE S_{n\ell j}^p(E)$$



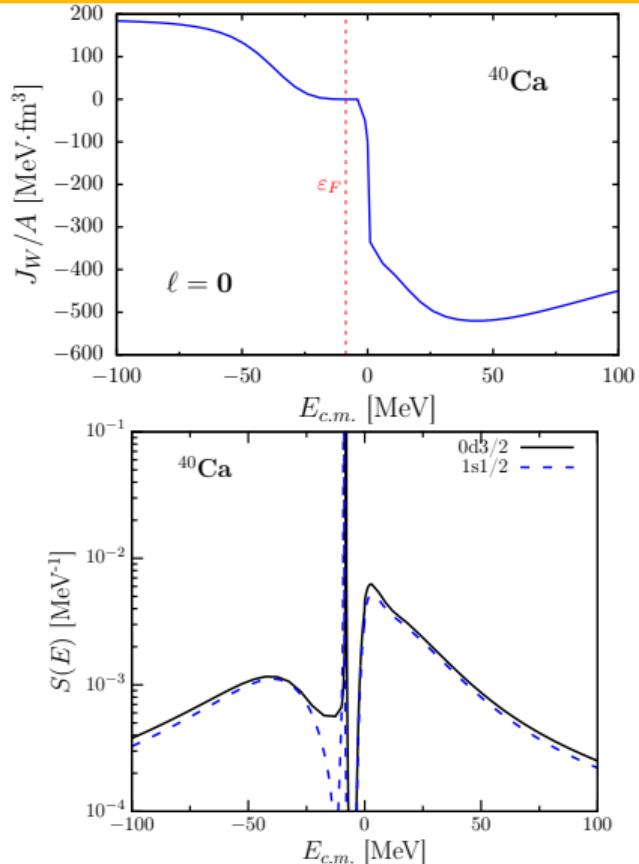
Spectroscopic factor, Occupation, and Depletion

- No imaginary component of Σ^* around ϵ_F
- Spectroscopic factor for states near ϵ_F is well defined from Σ^*

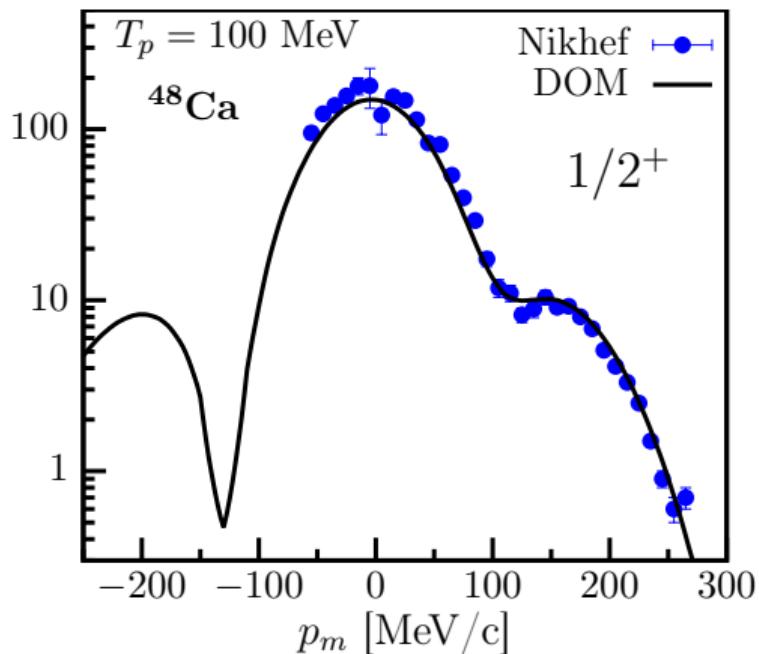
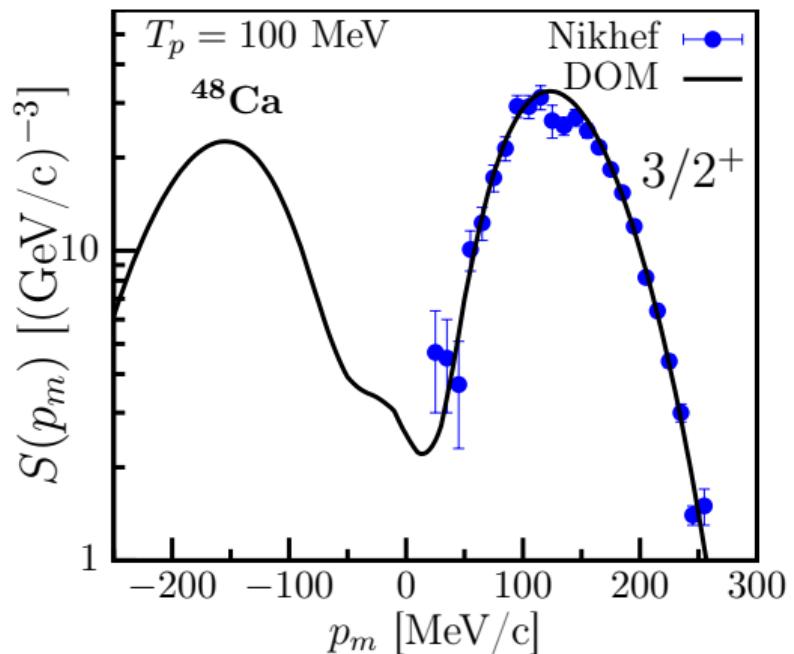
$$\mathcal{Z} = \left(1 - \frac{\partial \Sigma^*(\alpha_{qh}, \alpha_{qh}; E)}{\partial E} \Big|_{\epsilon} \right)^{-1}$$

$$n_{nlj} = \int_{-\infty}^{\epsilon_f} dE S_{nlj}^h(E) \quad d_{nlj} = \int_{\epsilon_f}^{\infty} dE S_{nlj}^p(E)$$

Orbital	\mathcal{Z}	n_{nlj}	d_{nlj}
$0d\frac{3}{2}$	0.71	0.80	0.17
$1s\frac{1}{2}$	0.60	0.82	0.15



$^{48}\text{Ca}(e,e'p)^{47}\text{K}$ Momentum Distribution



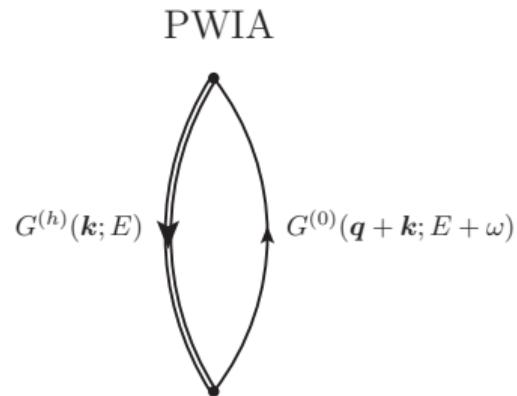
Data: G. J. Kramer *et. al.*, Nucl. Phys. A, **679**, 267 (2001)

Extending to inclusive cross section

- First implement PWIA
- Using DWIA for inclusive cross section will involve energies that require a relativistic treatment

Extending to inclusive cross section

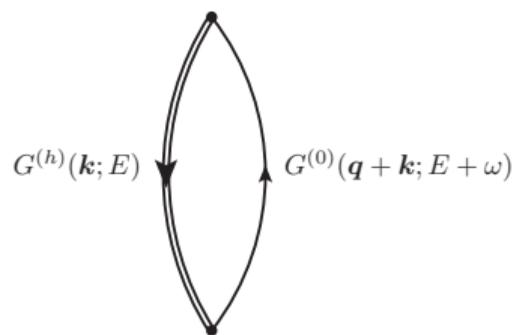
- First implement PWIA
- Using DWIA for inclusive cross section will involve energies that require a relativistic treatment



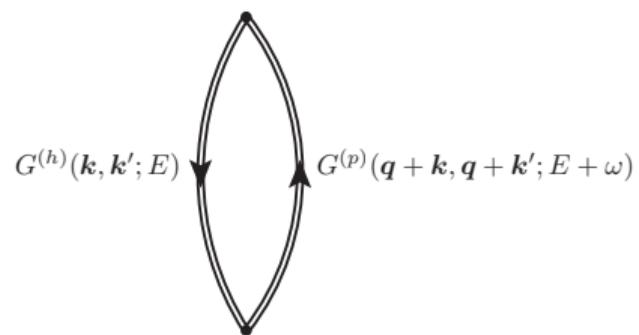
Extending to inclusive cross section

- First implement PWIA
- Using DWIA for inclusive cross section will involve energies that require a relativistic treatment

PWIA

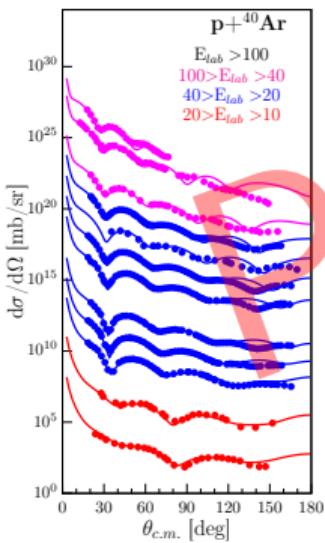


DWIA

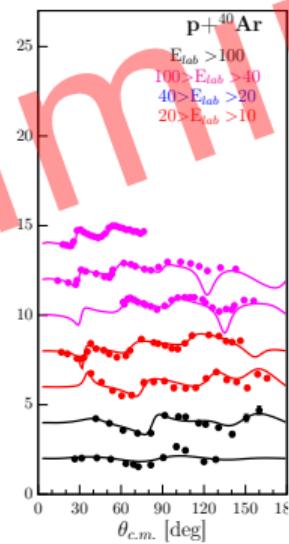
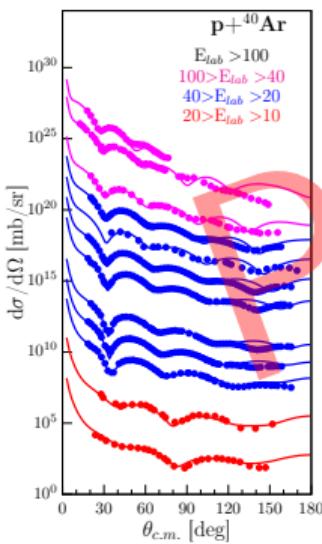


Preliminary Fit of ^{40}Ar

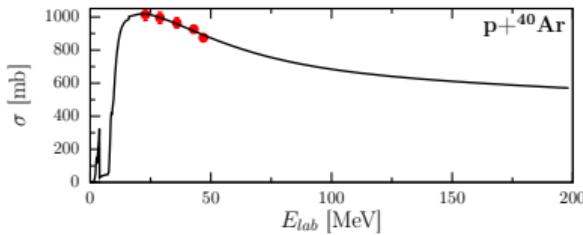
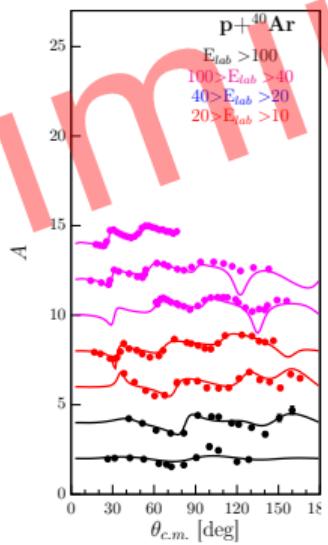
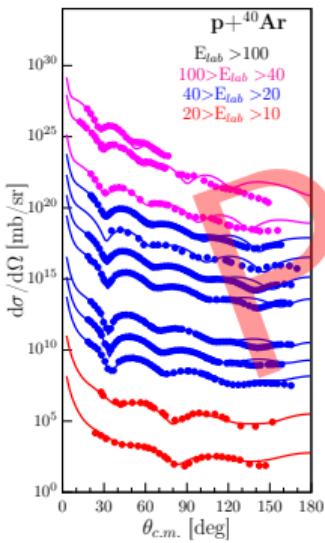
preliminary



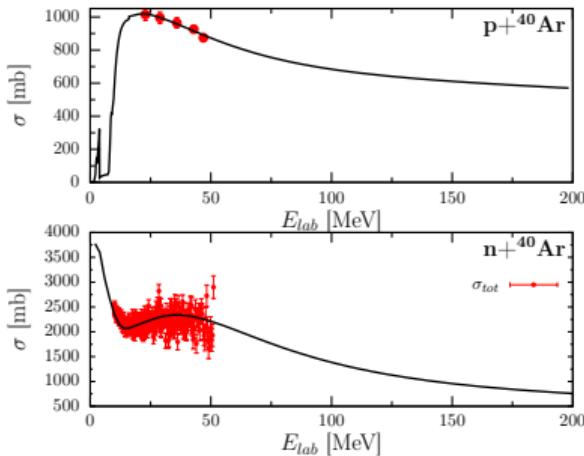
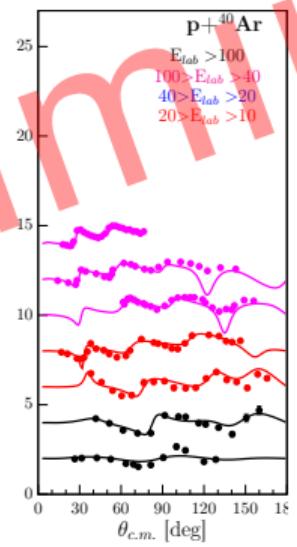
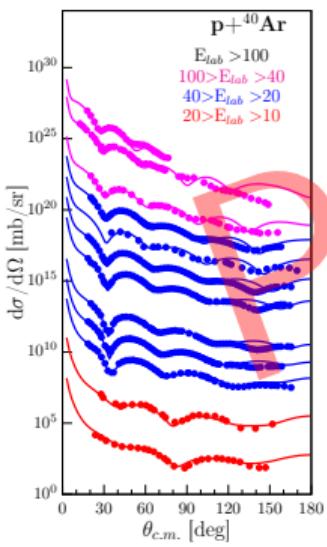
Preliminary Fit of ^{40}Ar



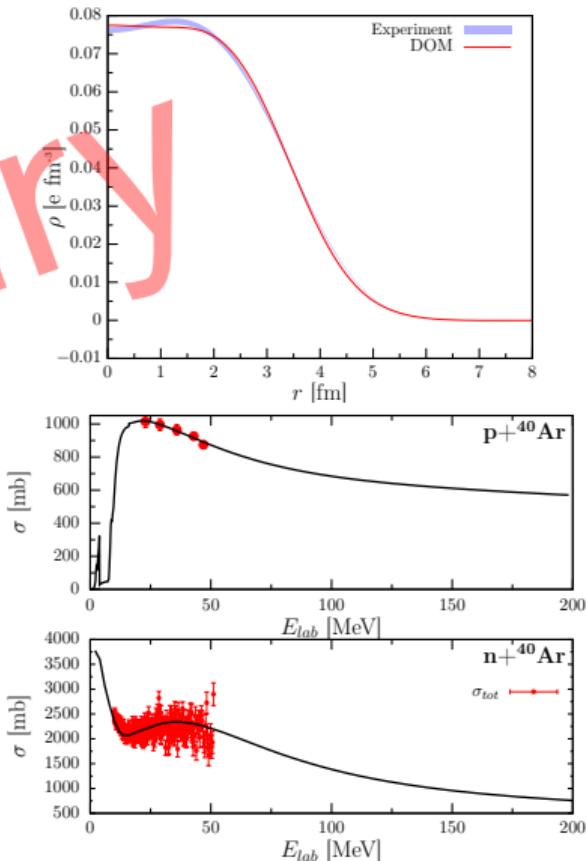
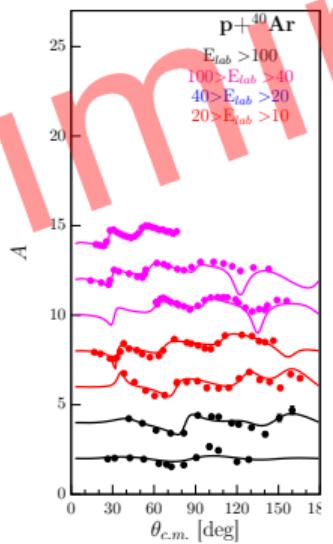
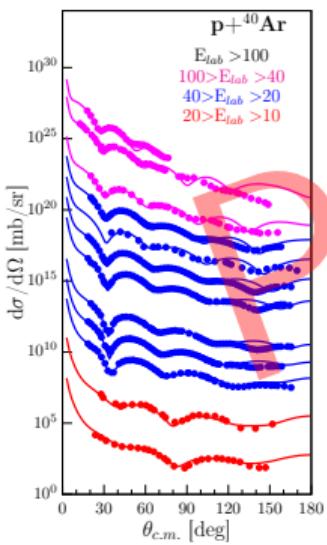
Preliminary Fit of ^{40}Ar



Preliminary Fit of ${}^{40}\text{Ar}$



Preliminary Fit of ^{40}Ar



Conclusions and Outlook

- The DOM is a robust model that can describe both positive and negative energy data

Conclusions and Outlook

- The DOM is a robust model that can describe both positive and negative energy data
- The DOM provides a consistent description of ${}^{40}\text{Ca}(\text{e},\text{e}'\text{p}){}^{39}\text{K}$ data

Conclusions and Outlook

- The DOM is a robust model that can describe both positive and negative energy data
- The DOM provides a consistent description of $^{40}\text{Ca}(\text{e},\text{e}'\text{p})^{39}\text{K}$ data
- Since this works for electrons, it should work for neutrinos

Conclusions and Outlook

- The DOM is a robust model that can describe both positive and negative energy data
- The DOM provides a consistent description of $^{40}\text{Ca}(\text{e},\text{e}'\text{p})^{39}\text{K}$ data
- Since this works for electrons, it should work for neutrinos
- DOM analysis of ^{40}Ar is underway

Thanks

- Willem Dickhoff - Advisor
- Robert Charity - DOM and data for DOM
- Henk Blok - $(e, e'p)$ data at Nikhef
- Louk Lapikás - $(e, e'p)$ data at Nikhef
- Carlotta Giusti - DWEEPY Code
- Hossein Mahzoon - DOM
- Lee Sobotka - Data for DOM



Backup

- “Smearing” of self-energy poles inflates \mathcal{Z}
- Renormalize with experimental excitation energy spectrum

$$\frac{\mathcal{Z}_F^{\text{DOM}}}{\int dE S^{\text{DOM}}(E)} = \frac{\mathcal{Z}_F^{\text{exp}}}{\int dE S^{\text{exp}}(E)}$$

