

Neutrinoless Double Beta Decay and Baryogenesis

Frank Deppisch

`f.deppisch@ucl.ac.uk`

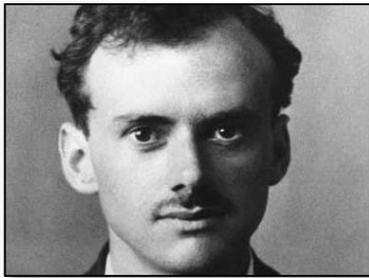
University College London

Fluff Ahead

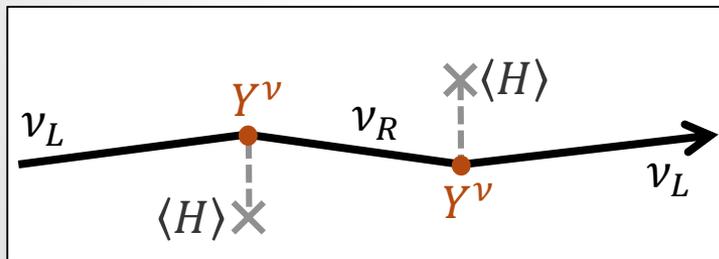


Dirac vs Majorana

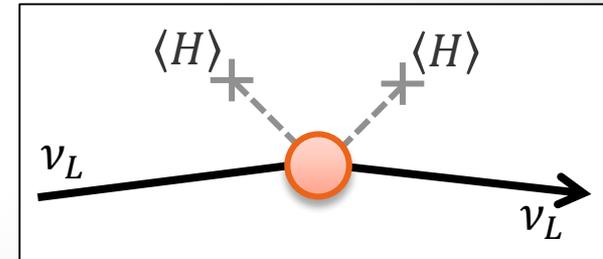
- ▶ Origin of neutrino masses beyond the Standard Model
- ▶ Two possibilities to define neutrino mass



Dirac mass analogous to other fermions but with $m_\nu / \Lambda_{EW} \approx 10^{-12}$ couplings to Higgs



Majorana mass, using only a left-handed neutrino
→ Lepton Number Violation

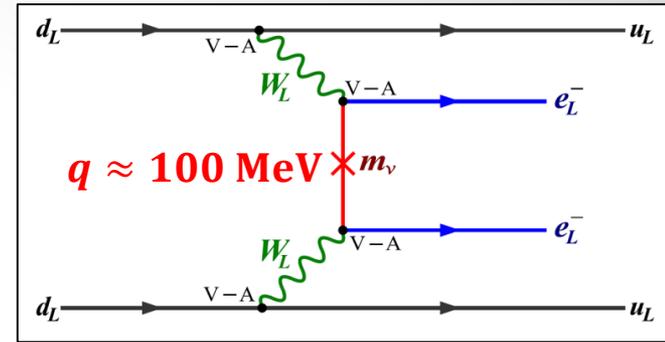


$0\nu\beta\beta$

▶ Half-life

$$T_{1/2}^{-1} = |m_{\beta\beta}|^2 G^{0\nu} |M^{0\nu}|^2$$

▶ Particle Physics



$$\mathcal{A}_{\mu\nu}^{lep} = \frac{1}{4} \sum_{i=1}^3 U_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{\not{q} + m_{\nu_i}}{q^2 - m_{\nu_i}^2} \gamma_\nu (1 - \gamma_5) \approx \frac{\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4q^2} \sum_{i=1}^3 U_{ei}^2 m_{\nu_i} \rightarrow m_{\beta\beta}$$

▶ Atomic Physics

- Leptonic phase space $G^{0\nu}$

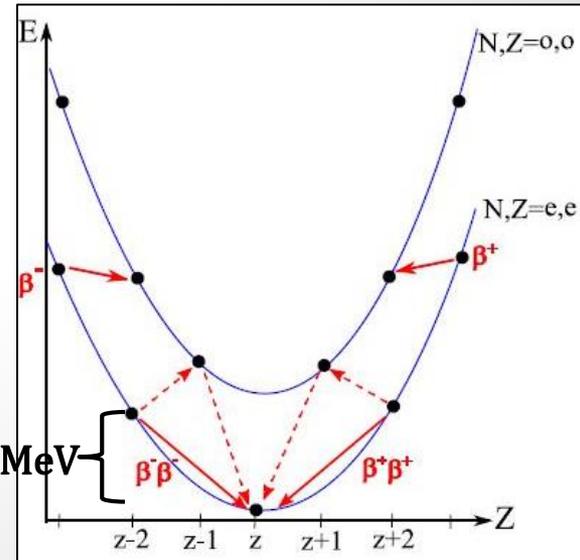
▶ Nuclear Physics

- Nuclear transition matrix element $M^{0\nu}$

$$T_{1/2}^{-1} \propto \frac{|m_{\beta\beta}|^2}{q^4} G_F^4 Q^5$$

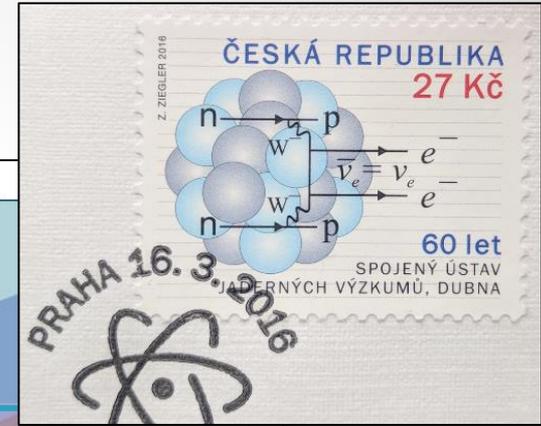
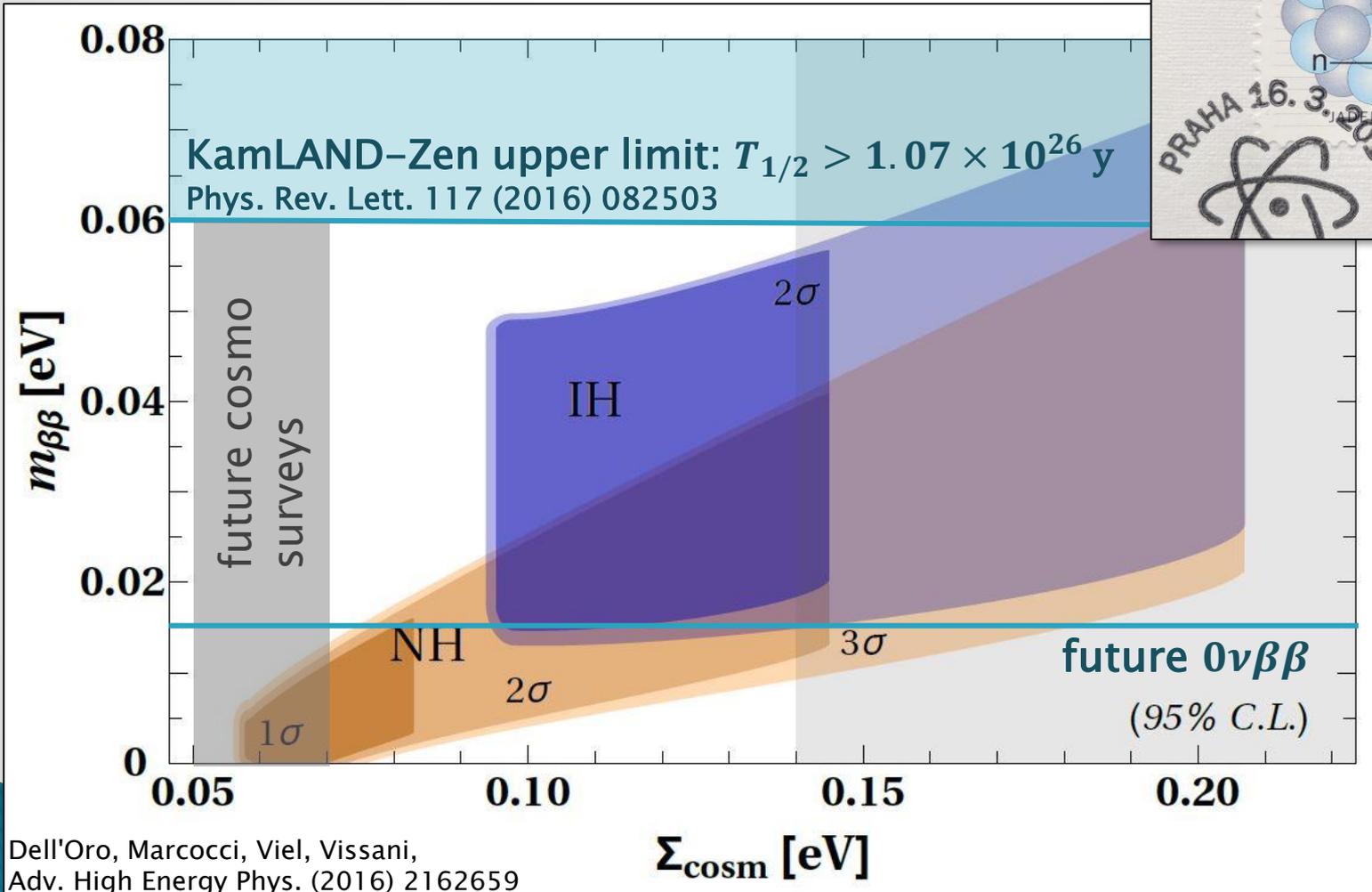
$$\frac{10^{25} \text{yr}}{T_{1/2}} \approx \left(\frac{|m_{\beta\beta}|}{\text{eV}} \right)^2$$

$Q \approx 2-4 \text{ MeV}$



Three Active Neutrinos

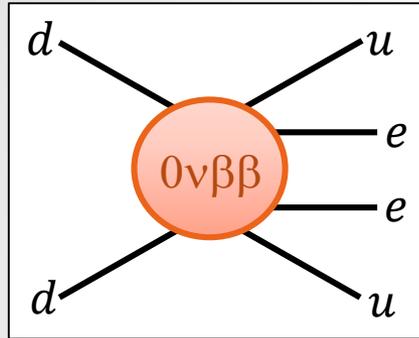
▶ Effective $0\nu\beta\beta$ Mass



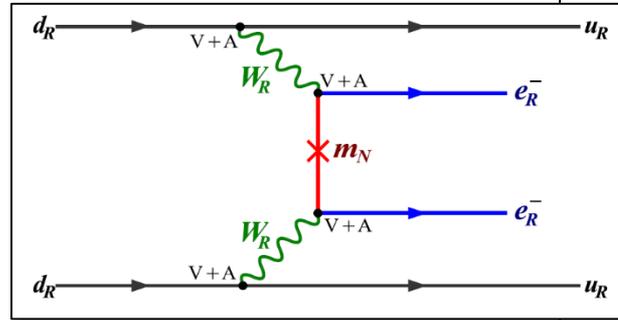
Dell'Oro, Marcocci, Viel, Vissani,
Adv. High Energy Phys. (2016) 2162659

New Physics and $0\nu\beta\beta$

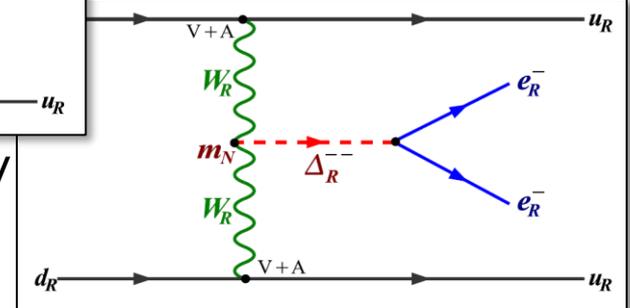
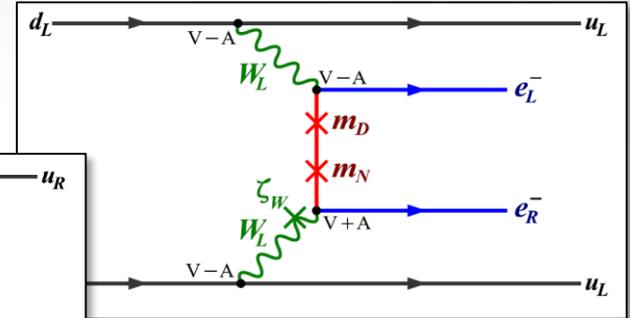
► Plethora of New Physics scenarios



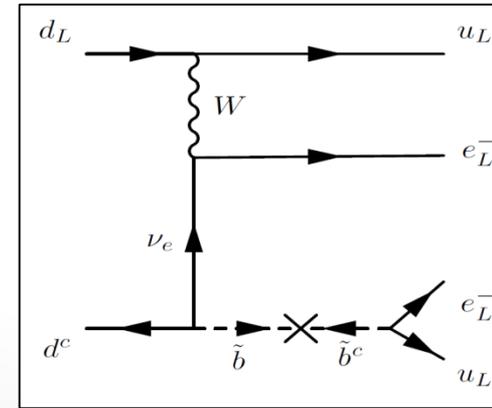
=



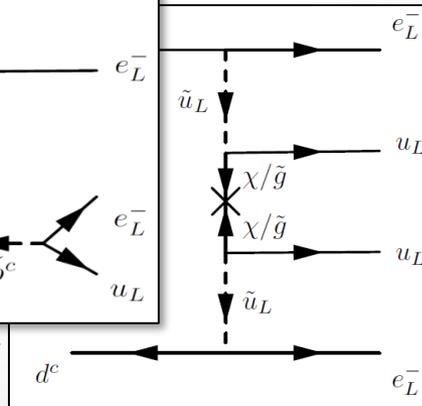
Left-Right Symmetry



$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$



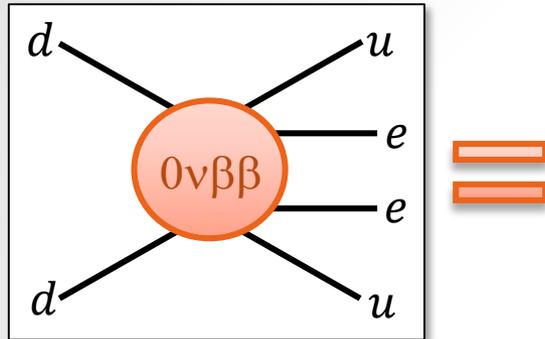
R-Parity Violating SUSY



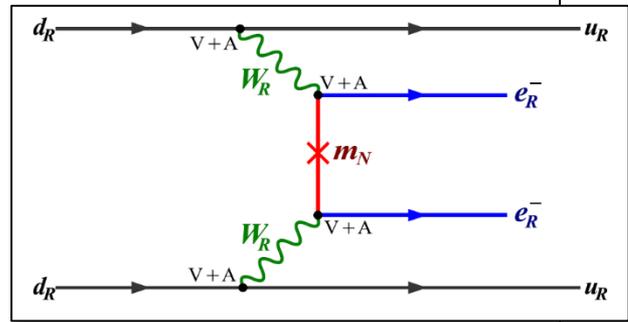
- Extra Dimensions
- Majorons
- Leptoquarks
- ...

New Physics and $0\nu\beta\beta$

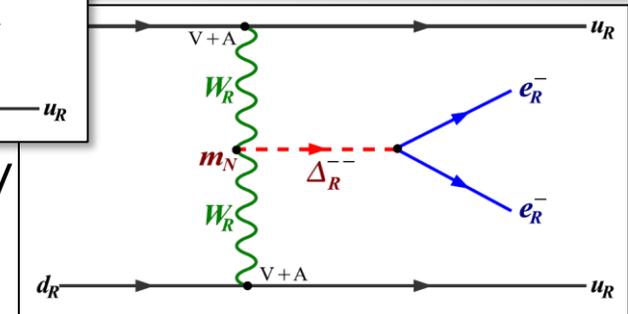
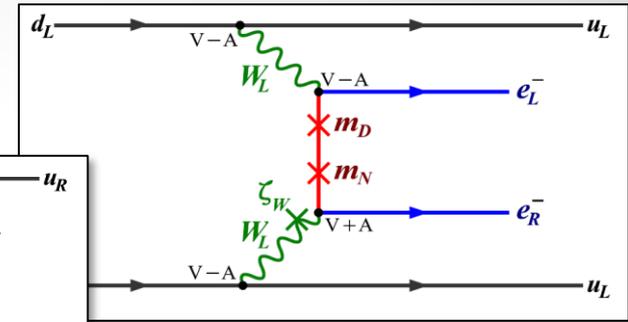
- ▶ Plethora of New Physics scenarios



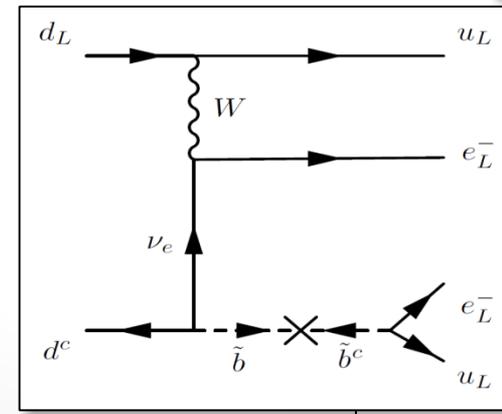
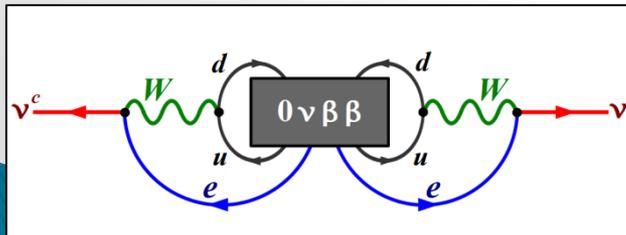
$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$



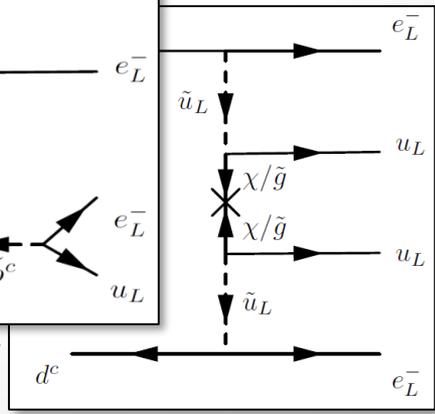
Left-Right Symmetry



- ▶ Neutrinos still Majorana



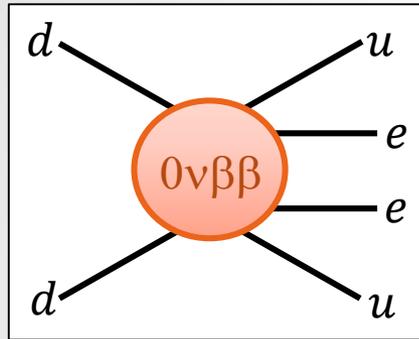
R-Parity Violating SUSY



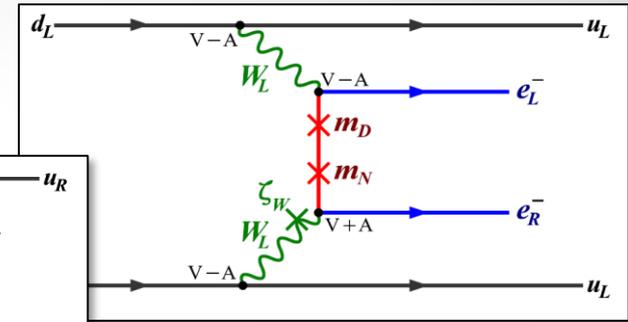
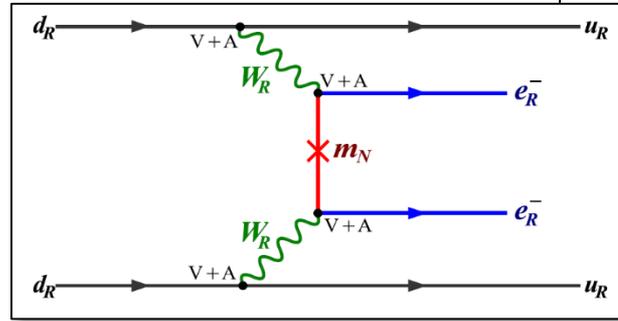
- Extra Dimensions
- Majorons
- Leptoquarks
- ...

New Physics and $0\nu\beta\beta$

▶ Examples in Left-Right Symmetry



=

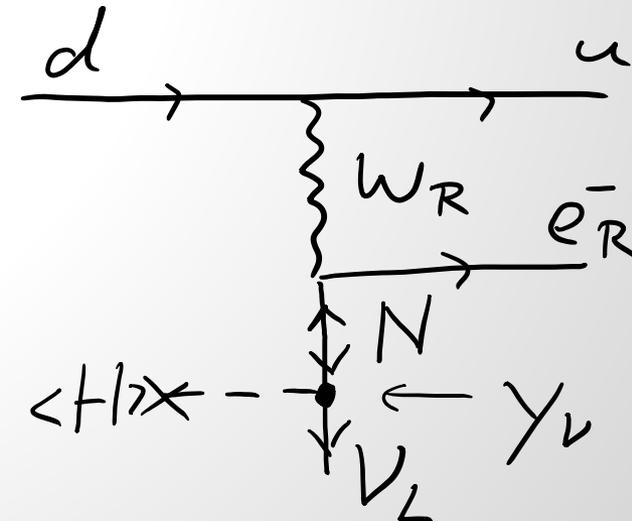


$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$

$$\epsilon_3^{RRZ} = \sum_{i=1}^3 V_{ei}^2 \frac{m_p}{m_N} \frac{m_W^4}{m_{WR}^4} \approx \frac{10^{-8}}{(\Lambda/1 \text{ TeV})^5}$$

$$\epsilon_{V-A}^{V+A} = \sum_{i=1}^3 U_{ei} W_{ei} \tan \zeta_W \approx \frac{10^{-9}}{(\Lambda/100 \text{ TeV})^3}$$

▶ $0\nu\beta\beta$ probes the TeV scale



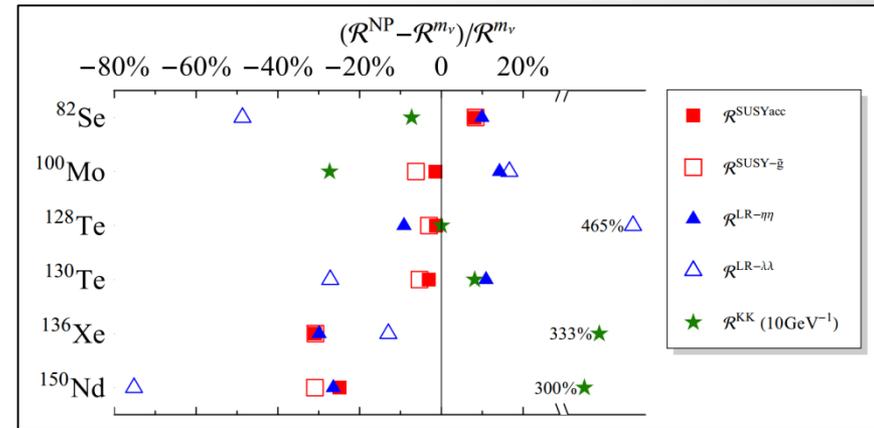
Disentangling New Physics

Comparison of $0\nu\beta\beta$ in multiple isotopes

(FFD, Päs PRL 2007, Meroni et al. 2013)

- Depends on $0\nu\beta\beta$ mechanism
- Independent of details of new physics (if one mechanism dominates)

$$\frac{T_{1/2}(X)}{T_{1/2}(Y)} = \frac{G(Y)|M(Y)|^2}{G(X)|M(X)|^2}$$

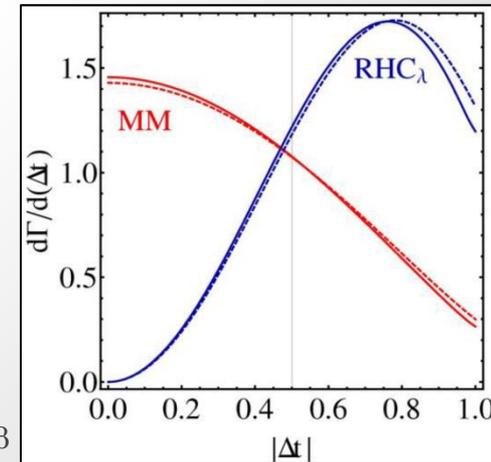


Angular and energy distribution of emitted electrons

(Doi et al. '83; Ali et al. '06; Arnold et al. '10; FFD, Jackson, Nasteva, Söldner-Rembold '10)

$$\frac{d\Gamma}{dE_{e_1} dE_{e_2} d\cos\theta} = \frac{\Gamma}{2} (1 - k(E_{e_1}, E_{e_2}) \cos\theta), \quad -1 < k < 1$$

- Linear in $\cos\theta$
- $k(E_{e_1}, E_{e_2})$ depends on $0\nu\beta\beta$ mechanism



Disentangling New Physics

▶ Comparison of $0\nu\beta\beta$ in multiple isotopes

(FFD, Päs PRL 2007, Meroni et al. 2013)

- Depends on $0\nu\beta\beta$ mechanism
- Independent of details of new physics (if one mechanism dominates)

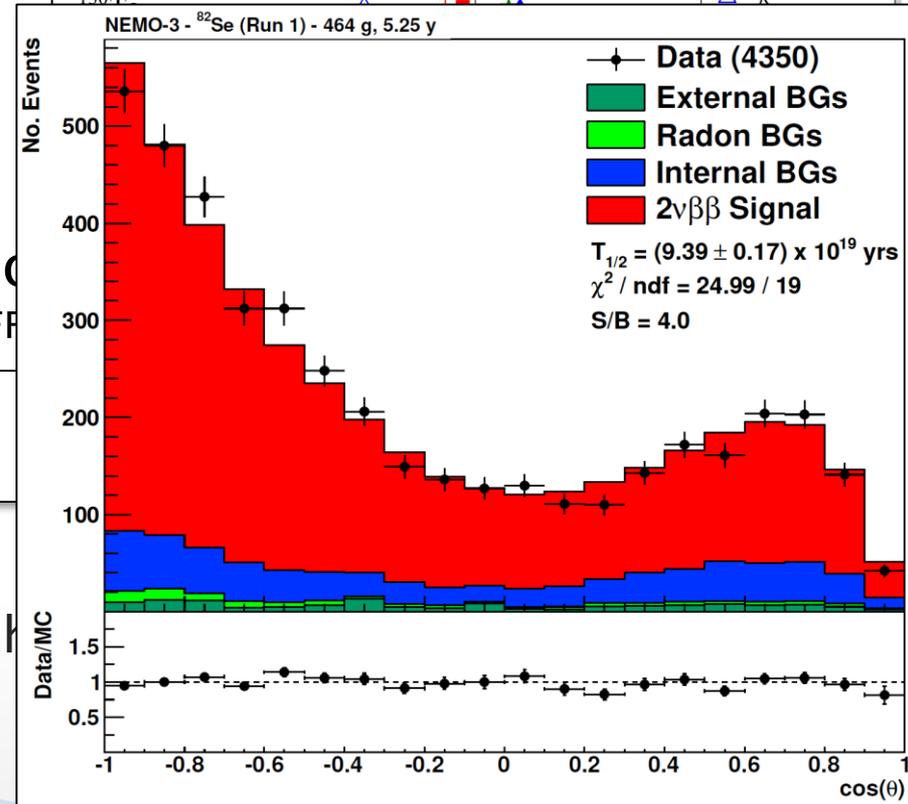
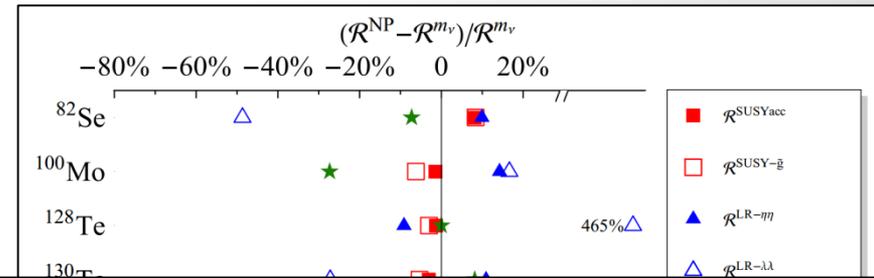
$$\frac{T_{1/2}(X)}{T_{1/2}(Y)} = \frac{G(Y)|M(Y)|^2}{G(X)|M(X)|^2}$$

▶ Angular and energy distribution

(Doi et al. '83; Ali et al. '06; Arnold et al. '10; FF)

$$\frac{d\Gamma}{dE_{e_1} dE_{e_2} d\cos\theta} = \frac{\Gamma}{2} (1 - k(E_{e_1}, E_{e_2}) \cos\theta),$$

- Linear in $\cos\theta$
- $k(E_{e_1}, E_{e_2})$ depends on $0\nu\beta\beta$ mechanism



Baryon Asymmetry

- ▶ The Universe is not matter–antimatter symmetric
 - CMB Anisotropy
 - Primordial Nucleosynthesis
 - No matter–antimatter annihilation
- ▶ Observed asymmetry

$$\eta_B \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} = (6.20 \pm 0.15) \times 10^{-10}$$

- Very small... Universe may have begun symmetric
- Still too large... to be compatible with the Standard Model

Baryon Asymmetry

- ▶ Dynamic generation of baryon asymmetry requires (Sakharov '66)
 - Baryon number violation
 - C and CP Violation
 - Out-of-equilibrium dynamics
- ▶ Standard Model
 - Baryon number violated at quantum level (Sphalerons)
 - C and CP violated but effect too small

$$\frac{\text{Im det}(m_u m_u^+ m_d m_d^+)}{v^{12}} = J \frac{m_t^4 m_c^2 m_b^4 m_s^2}{v^{12}} \approx 10^{-19}$$

- Electroweak phase transition out-of-equilibrium if first order but requires

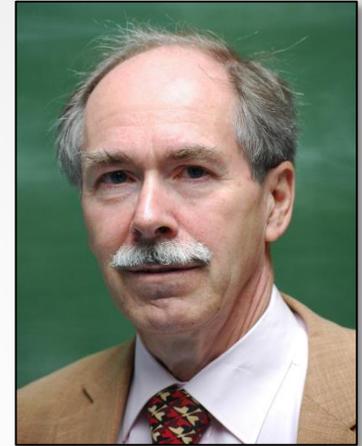
$$m_h < 60 - 80 \text{ GeV}$$



Sphalerons

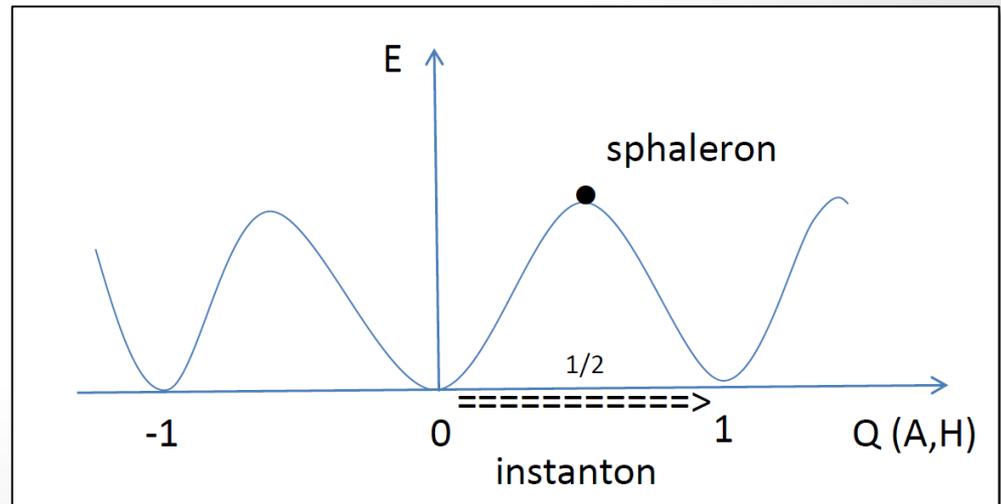
- ▶ Baryon and Lepton numbers accidental, classical symmetries in the Standard Model
- ▶ Violated at the quantum level (t' Hooft '76)

$$\partial_\mu J_B^\mu = \partial_\mu J_L^\mu = \frac{g^2}{32\pi^2} F_{\mu\nu} \tilde{F}^{\mu\nu}$$



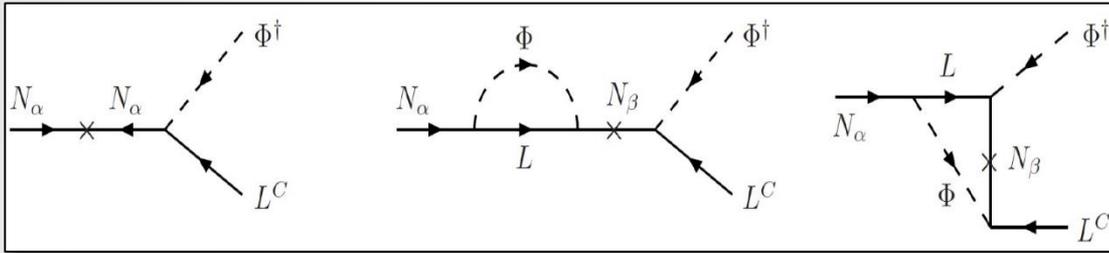
- $B + L$ violated
- $B - L$ remains conserved
- ▶ Sphaleron transitions in equilibrium $\frac{\Gamma_{Sph}}{H} > 1$ for

$$\Lambda_{EW} \approx 10^2 \text{ GeV} < T < 10^{12} \text{ GeV}$$



Leptogenesis – Vanilla

- ▶ Decays of heavy Majorana neutrinos violating L and CP (Fukugita, Yanagida '86)

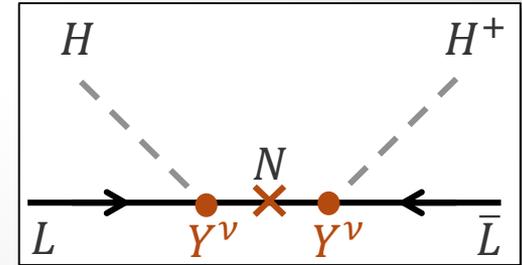


- CP asymmetry

$$\epsilon_1 = \frac{\Gamma(N_1 \rightarrow LH^+) - \Gamma(N_1 \rightarrow \bar{L}H)}{\Gamma(N_1 \rightarrow LH^+) + \Gamma(N_1 \rightarrow \bar{L}H)} \approx \frac{3}{8\pi} \frac{\text{Im}[(Y_\nu Y_\nu^+)_{1k}^2] M_1}{(Y_\nu Y_\nu^+)_{11} M_k}$$

- ▶ Competition with washout processes eradicating L asymmetry

$$M_N \gtrsim 10^8 \left(\frac{\eta_B}{5 \times 10^{-11}} \right) \left(\frac{0.06\text{eV}}{m_3} \right) \text{GeV}$$

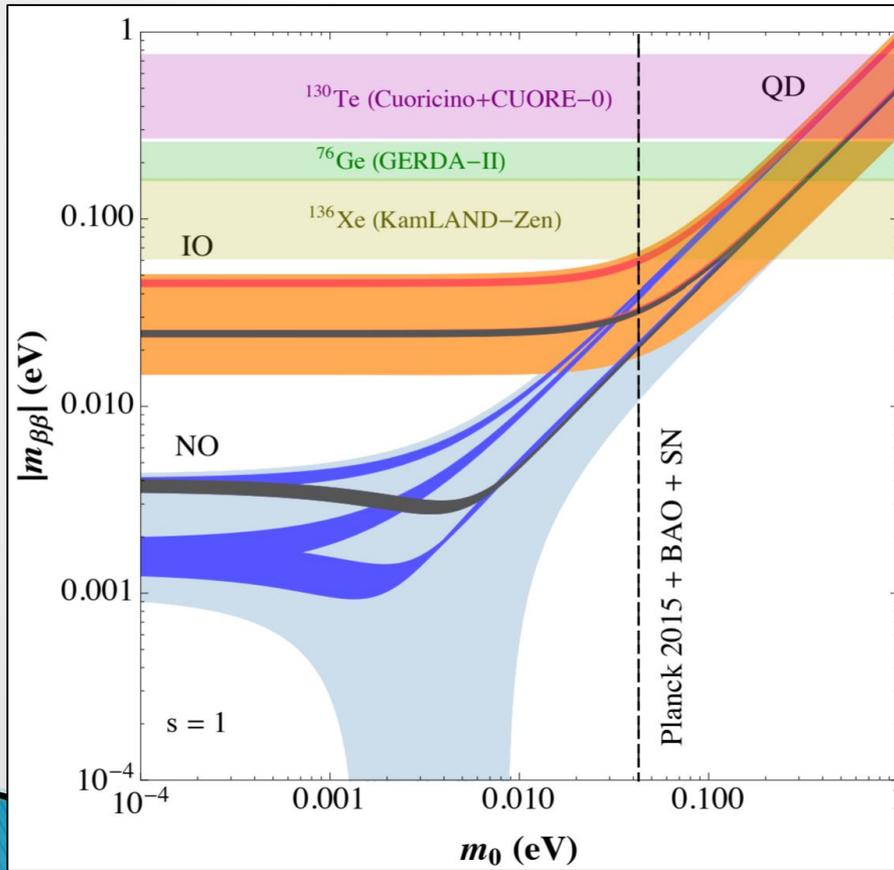


Conversion to baryon asymmetry via sphaleron processes

$$\eta_B \approx \eta_L$$

Leptogenesis – Vanilla

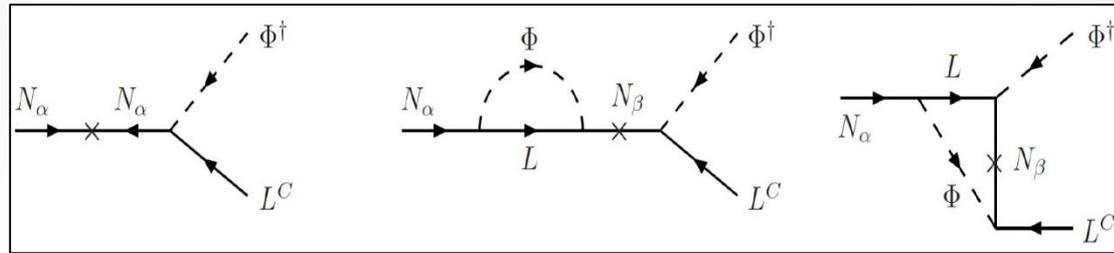
- ▶ Consequence for $0\nu\beta\beta$
 - Only standard light neutrino exchange
 - But CP Majorana phases expected?



Hagedorn, Molinaro '17

Leptogenesis – Resonant

- ▶ Dominance of self-energy loop for small mass difference between heavy neutrinos (Pilaftsis '97)



- CP asymmetry can be $O(1)$ for $\Delta M_N \approx \Gamma_N$
- Viable leptogenesis for neutrino masses as light as $M_N \approx 100$ GeV

Leptogenesis – Resonant

▶ Seesaw I mechanism with TeV scale heavy neutrinos

- Standard Seesaw with small Yukawa couplings

$$Y_\nu \approx 10^{-6} \sqrt{M_N/\text{TeV}}$$

- “Bent” Seesaw I mechanisms (e.g. Inverse Seesaw)

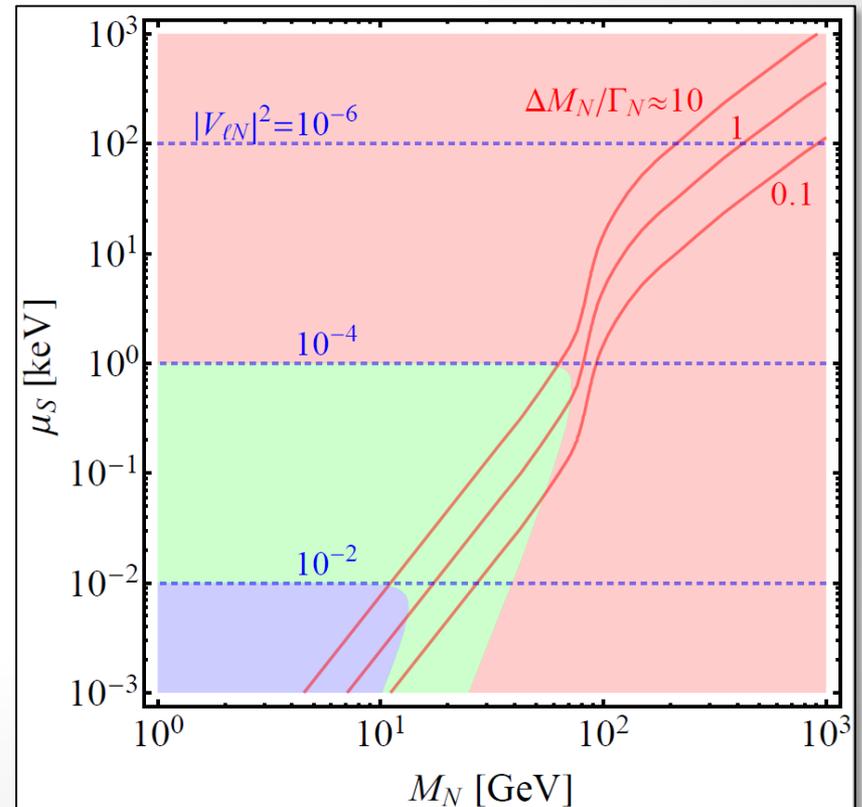
- Decouple Λ_{LNV} from heavy neutrino mass

- Example

$$\begin{pmatrix} 0 & Y_\nu \langle H \rangle & 0 \\ Y_\nu \langle H \rangle & \mu & M \\ 0 & M & \mu \end{pmatrix}$$

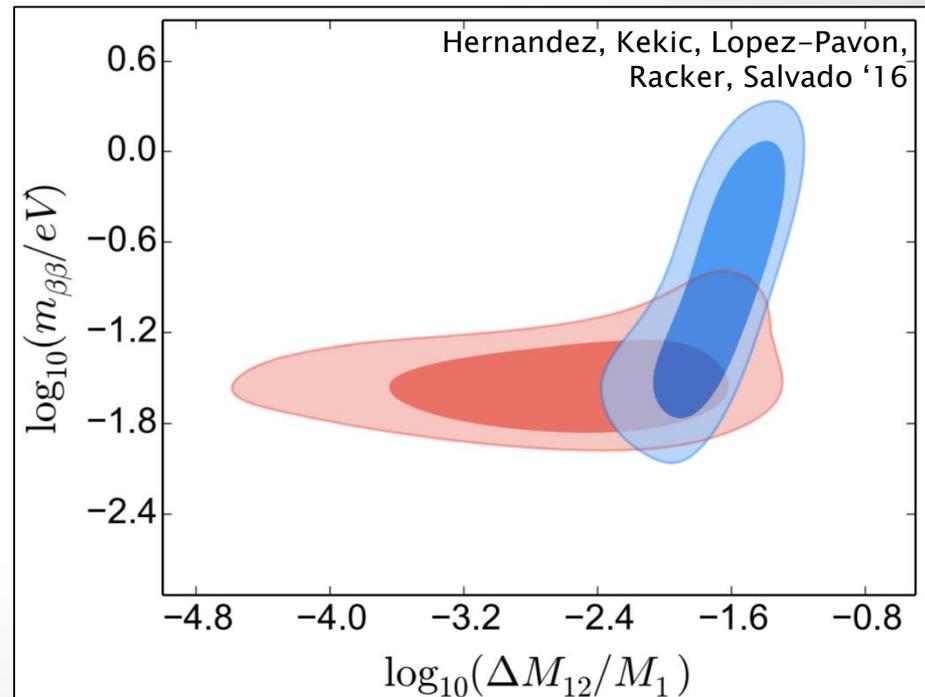
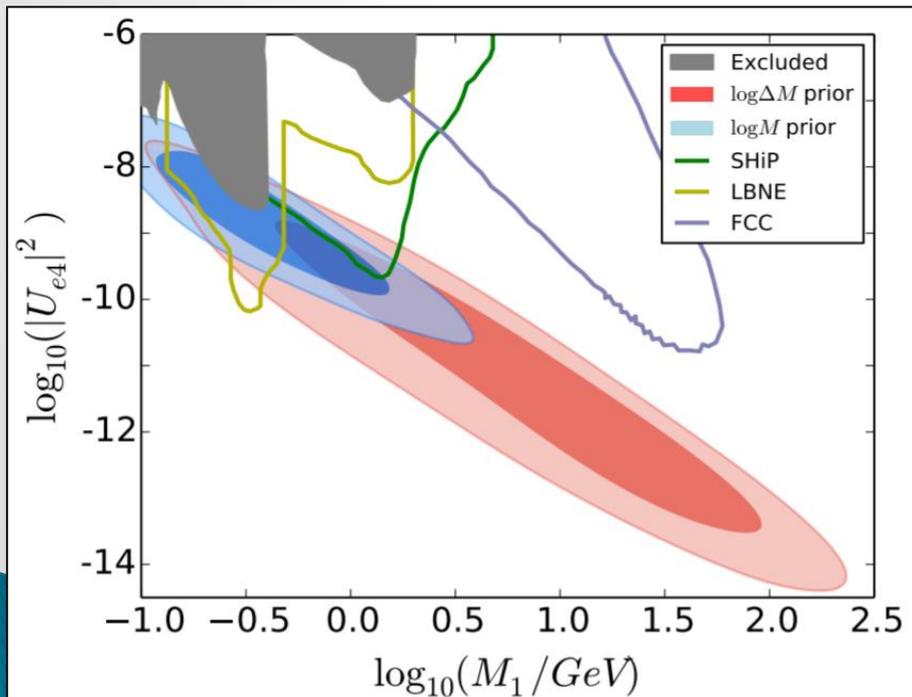
- LNV in $0\nu\beta\beta$ suppressed by $\frac{\Delta m_N}{m_N}$

- LNV in resonant N production suppressed by $\frac{\Delta m_N}{\Gamma_N} \approx \frac{\mu}{\Gamma_N}$



Leptogenesis – Oscillations

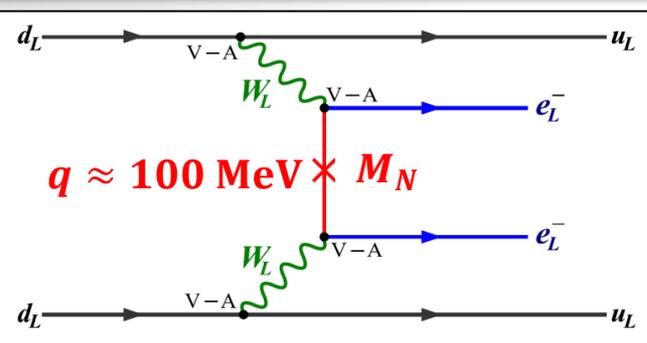
- ▶ Sterile neutrinos with small hierarchical Yukawa couplings (Akhmedov, Rubakov, Smirnov '98)
 - One neutrino not in equilibrium before critical sphaleron T
 - CP and flavor violating oscillations between sterile neutrinos generate asymmetry
 - Viable mechanism for $m_N \approx 1 - 100 \text{ GeV}$



Leptogenesis – Oscillations

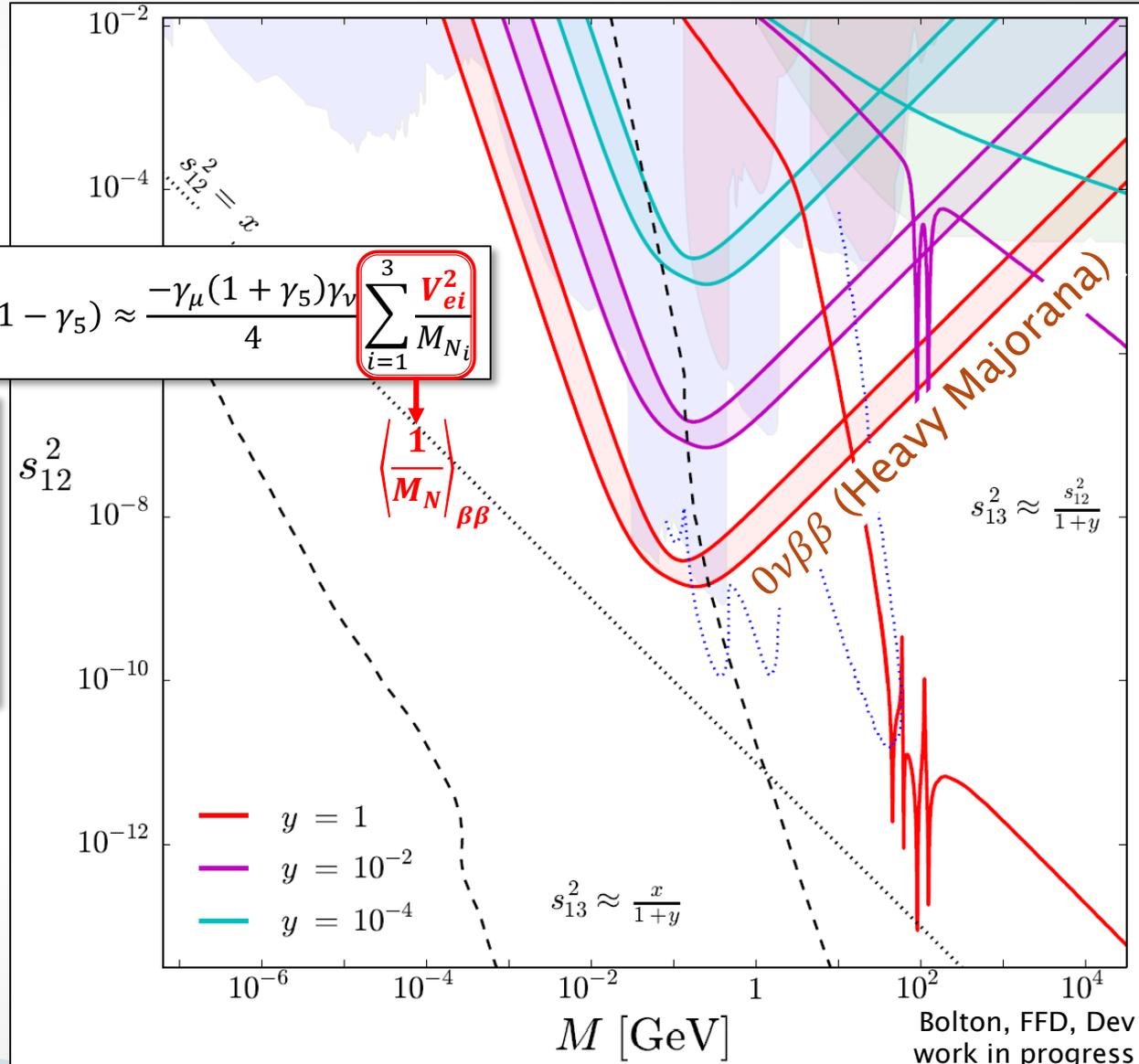
- ▶ Heavy neutrinos with masses larger than ≈ 100 MeV

$$\mathcal{A}_{\mu\nu}^{lep} = \frac{1}{4} \sum_{i=1}^3 V_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{\not{q} + M_{N_i}}{q^2 - M_{N_i}^2} \gamma_\nu (1 - \gamma_5) \approx \frac{-\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4} \sum_{i=1}^3 \frac{V_{ei}^2}{M_{N_i}}$$



Different nuclear matrix elements

- Short-range operator



Baryon Asymmetry Generation and Washout

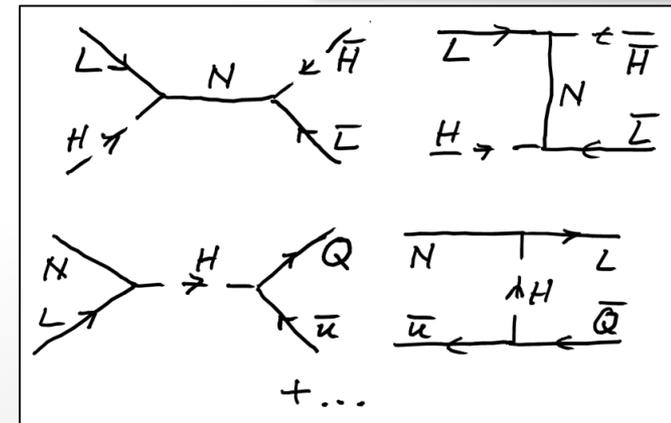
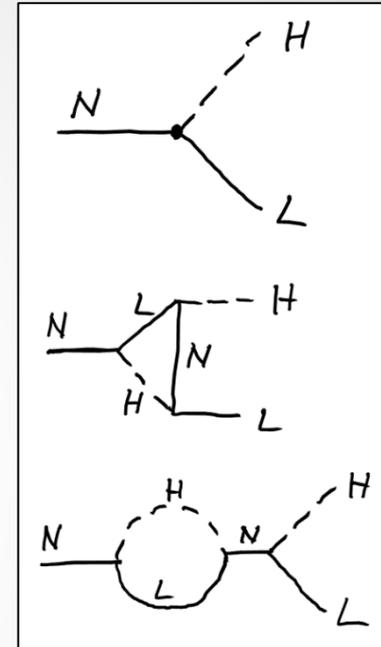
▶ Classic Example: High-Scale Leptogenesis

- Generation via heavy neutrino decays
- Competition with LNV washout processes
- Conversion to baryon asymmetry
 - EW sphaleron processes at $T \approx 100$ GeV
 - Observed asymmetry

$$\eta_B \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} = (6.20 \pm 0.15) \times 10^{-10}$$

▶ Other possible scenarios

- For us only important:
($B - L$) asymmetry generated
above LHC scale



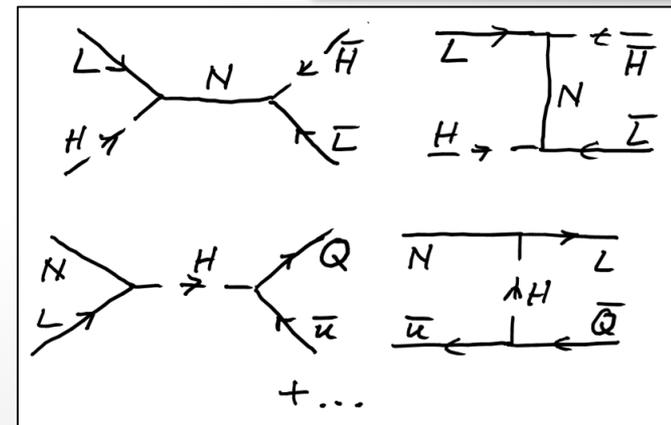
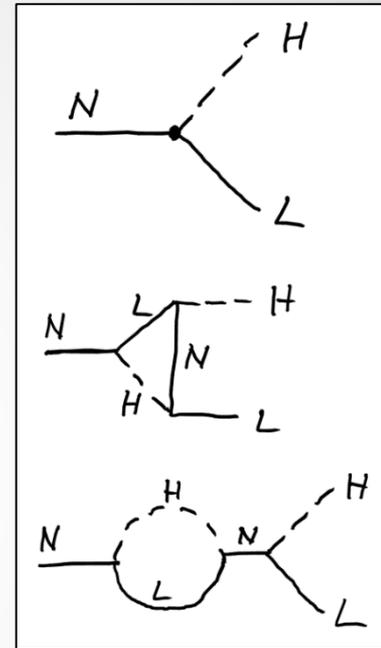
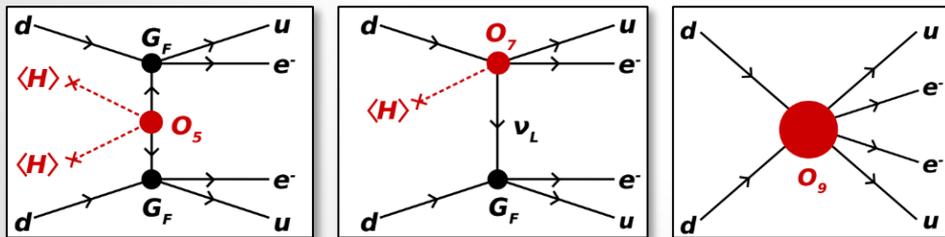
Baryon Asymmetry Generation and Washout

▶ Classic Example: High-Scale Leptogenesis

- Generation via heavy neutrino decays
- Competition with LNV washout processes
- Conversion to baryon asymmetry
 - EW sphaleron processes at $T \approx 100$ GeV
 - Observed asymmetry

$$\eta_B \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} = (6.20 \pm 0.15) \times 10^{-10}$$

▶ What if we observe lepton number violating processes in $0\nu\beta\beta$?



Washout via $0\nu\beta\beta$ operators

▶ Analogous analysis using LNV effective operators of mass dimensions 5, 7, 9, 11

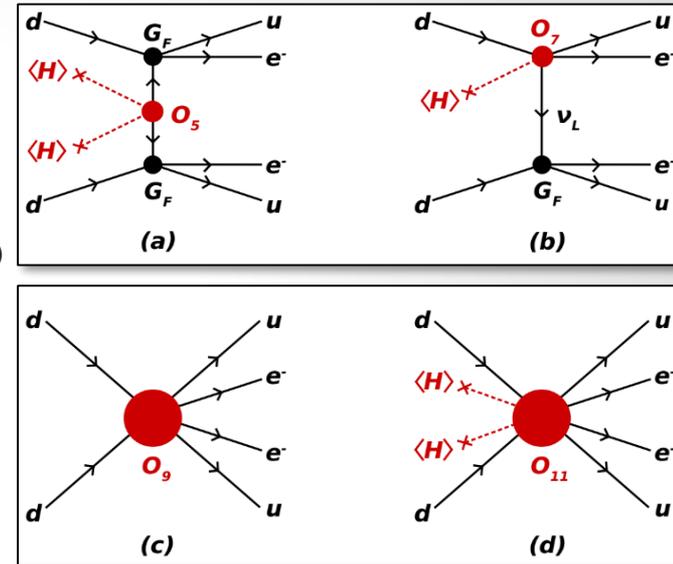
- 129 Operators (Babu, Leung '01, de Gouvea, Jenkins '08)
- Examples

$$\begin{aligned} \mathcal{O}_5 &= (L^i L^j) H^k H^l \epsilon_{ik} \epsilon_{jl}, \\ \mathcal{O}_7 &= (L^i d^c) (\bar{e}^c \bar{u}^c) H^j \epsilon_{ij}, \\ \mathcal{O}_9 &= (L^i L^j) (\bar{Q}_i \bar{u}^c) (\bar{Q}_j \bar{u}^c), \\ \mathcal{O}_{11} &= (L^i L^j) (Q_k d^c) (Q_l d^c) H_m \bar{H}_i \epsilon_{jk} \epsilon_{lm}, \end{aligned}$$

- Matching to $0\nu\beta\beta$ operators

$$m_e \epsilon_5 = \frac{g^2 v^2}{\Lambda_5}, \quad \frac{G_F \epsilon_7}{\sqrt{2}} = \frac{g^3 v}{2\Lambda_7^3}, \quad \frac{G_F^2 \epsilon_{\{9,11\}}}{2m_p} = \left\{ \frac{g^4}{\Lambda_9^5}, \frac{g^6 v^2}{\Lambda_{11}^7} \right\}.$$

$$T_{1/2} = 2.1 \times 10^{25} \text{ y} \cdot \left(\Lambda_D / \Lambda_D^0 \right)^{2d-8}$$



\mathcal{O}_D	λ_D^0 [GeV]	Λ_D^0 [GeV]
\mathcal{O}_5	9.2×10^{10}	9.1×10^{13}
\mathcal{O}_7	1.2×10^2	2.6×10^4
\mathcal{O}_9	4.3×10^1	2.1×10^3
\mathcal{O}_{11}	7.8×10^1	1.0×10^3

Washout via $0\nu\beta\beta$ operators

- ▶ Boltzmann equation including washout of D -dim effective operator

$$n_\gamma H T \frac{d\eta_L}{dT} = c_D \frac{T^{2D-4}}{\Lambda_D^{2D-8}} \eta_L$$

- $c_{\{5,7,9,11\}} = \left\{ \frac{8}{\pi^5}, \frac{27}{2\pi^7}, \frac{3.2 \times 10^4}{\pi^9}, \frac{3.9 \times 10^5}{\pi^{13}} \right\}$

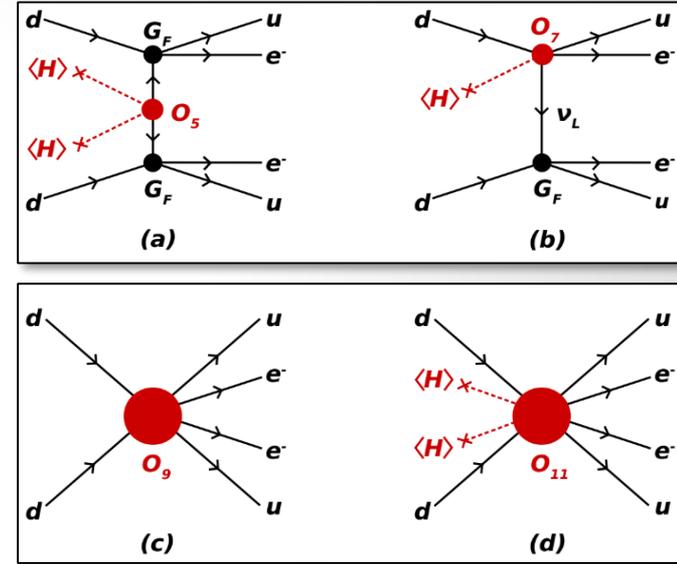
- ▶ Effective washout if

$$\frac{\Gamma_W}{H} \equiv \frac{c_D}{n_\gamma H} \frac{T^{2D-4}}{\Lambda_D^{2D-8}} = c'_D \frac{\Lambda_{\text{Pl}}}{\Lambda_D} \left(\frac{T}{\Lambda_D} \right)^{2D-9} \gtrsim 1$$

$$\Lambda_D \left(\frac{\Lambda_D}{c'_D \Lambda_{\text{Pl}}} \right)^{\frac{1}{2D-9}} \equiv \lambda_D \lesssim T \lesssim \Lambda_D$$

- ▶ Better: Solve Boltzmann such that initial asymmetry is washed out at the EW scale

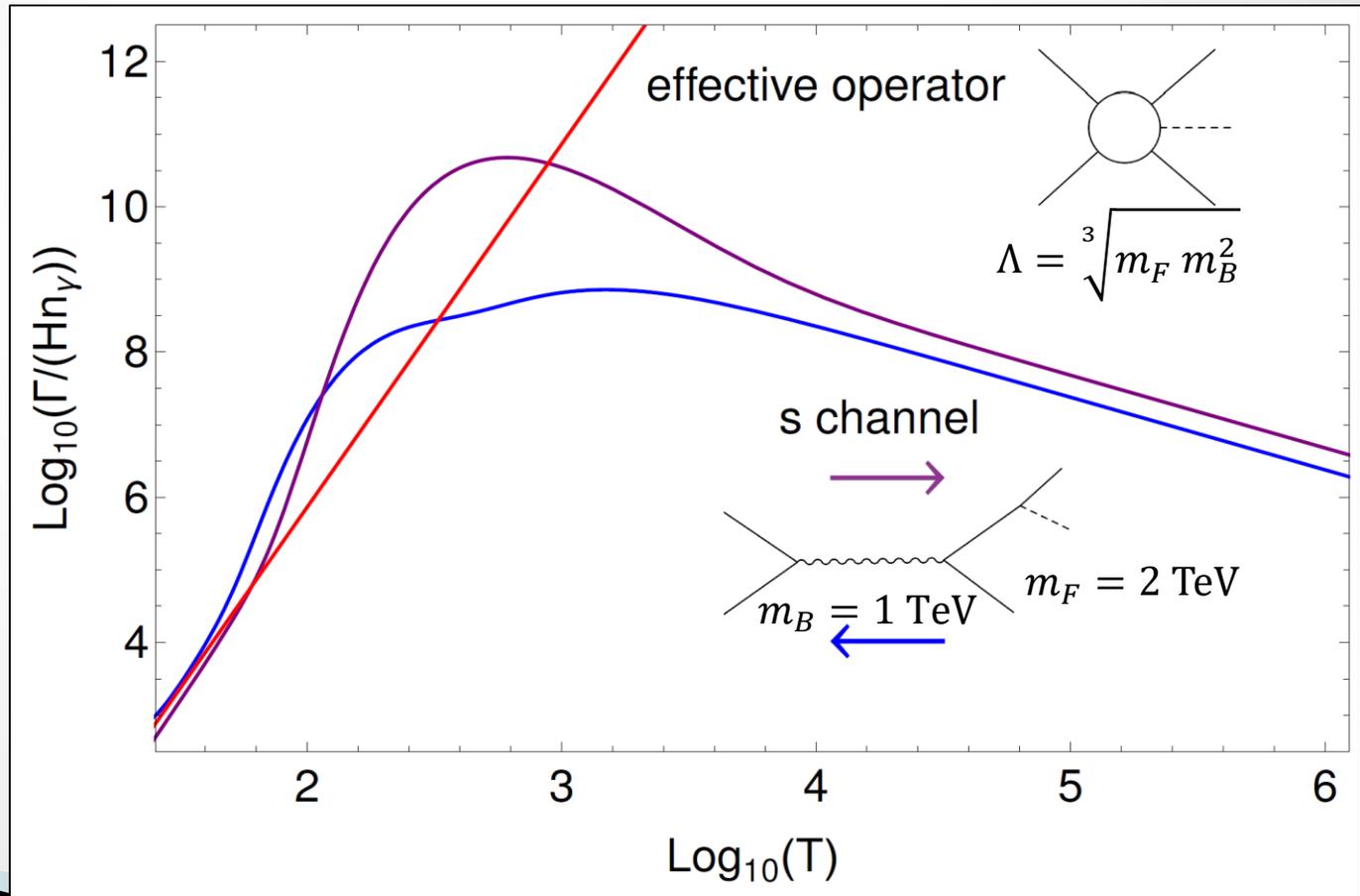
$$\hat{\lambda}_D \approx \left[(2D-9) \ln \left(\frac{10^{-2}}{\eta_B^{\text{obs}}} \right) \lambda_D^{2D-9} + v^{2D-9} \right]^{\frac{1}{2D-9}},$$



\mathcal{O}_D	λ_D^0 [GeV]	Λ_D^0 [GeV]
\mathcal{O}_5	9.2×10^{10}	9.1×10^{13}
\mathcal{O}_7	1.2×10^2	2.6×10^4
\mathcal{O}_9	4.3×10^1	2.1×10^3
\mathcal{O}_{11}	7.8×10^1	1.0×10^3

Washout via $0\nu\beta\beta$ operators

- ▶ Even better:
UV-completed operators for behaviour around Λ



Effect of LFV operators

- ▶ Analogous analysis for eff. 6-dim LFV operators

$$\mathcal{O}_6(ll\gamma H) = \bar{L}_i \sigma^{\mu\nu} e_j^c H^+ F_{\mu\nu}$$

$$\mathcal{O}_6(llll) = (\bar{L}_i \gamma^\mu L_j) (\bar{L}_k \gamma^\mu L_l)$$

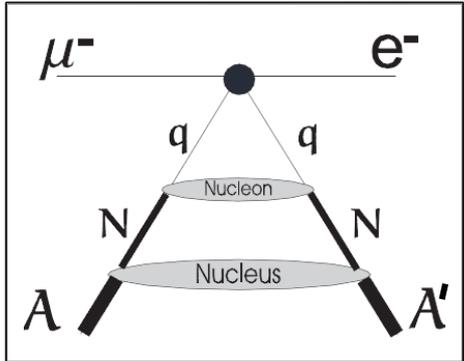
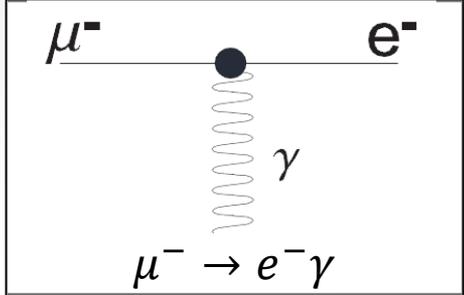
$$\mathcal{O}_6(llqq) = (\bar{L}_i \gamma^\mu L_j) (\bar{Q}_k \gamma^\mu Q_l)$$

- Do not washout total lepton number asymmetry but equilibrate lepton flavours

- ▶ Matching to LFV process rate

$$\mathcal{C}_{ll\gamma} = \frac{eg^3}{16\pi^2 \Lambda_{ll\gamma}^2}, \quad \mathcal{C}_{llqq} = \frac{g^2}{\Lambda_{llqq}^2}$$

$$\text{Br}_{\mu \rightarrow e\gamma} = 5.7 \times 10^{-13} \cdot (\Lambda_{\mu e\gamma}^0 / \Lambda_{\mu e\gamma})^4$$



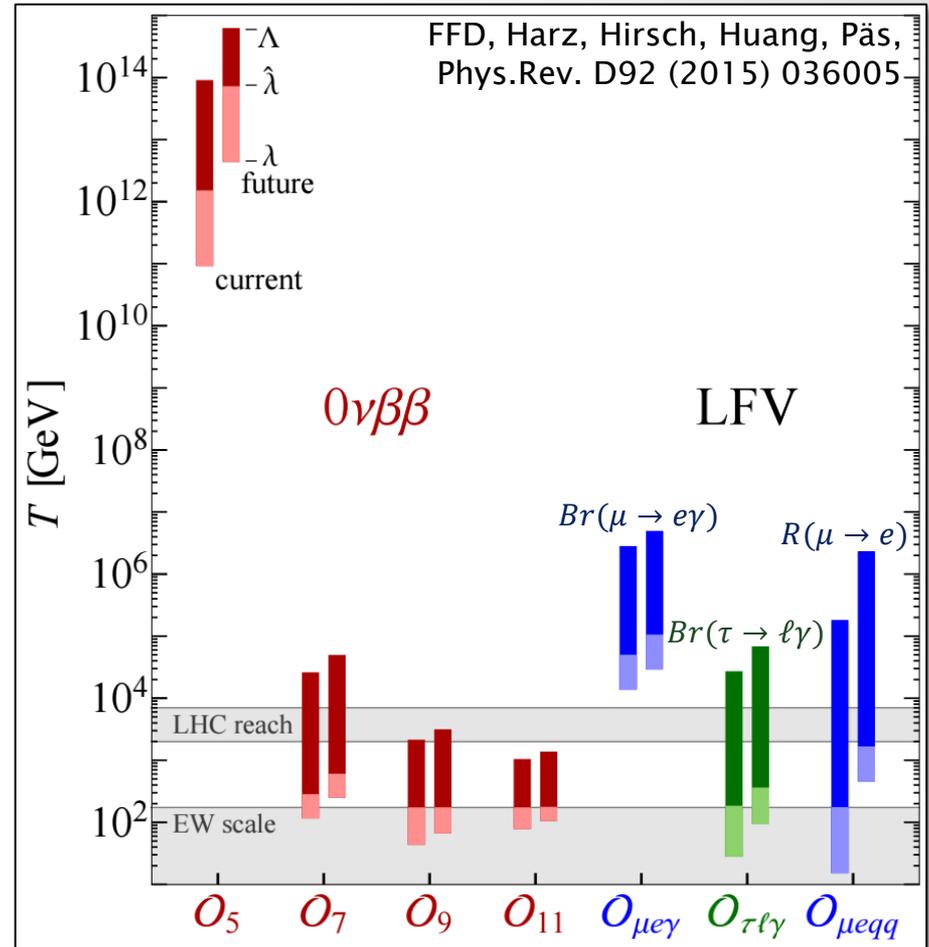
$\mu^- \rightarrow e^-$ conversion in nuclei

\mathcal{O}_i	λ_i^0 [GeV]	Λ_i^0 [GeV]
$\mathcal{O}_{\mu e\gamma}$	1.4×10^4	2.8×10^6
$\mathcal{O}_{\tau l\gamma}$	2.8×10^1	2.7×10^4
$\mathcal{O}_{\mu eqq}$	1.5×10^1	1.8×10^5

Baryon Asymmetry

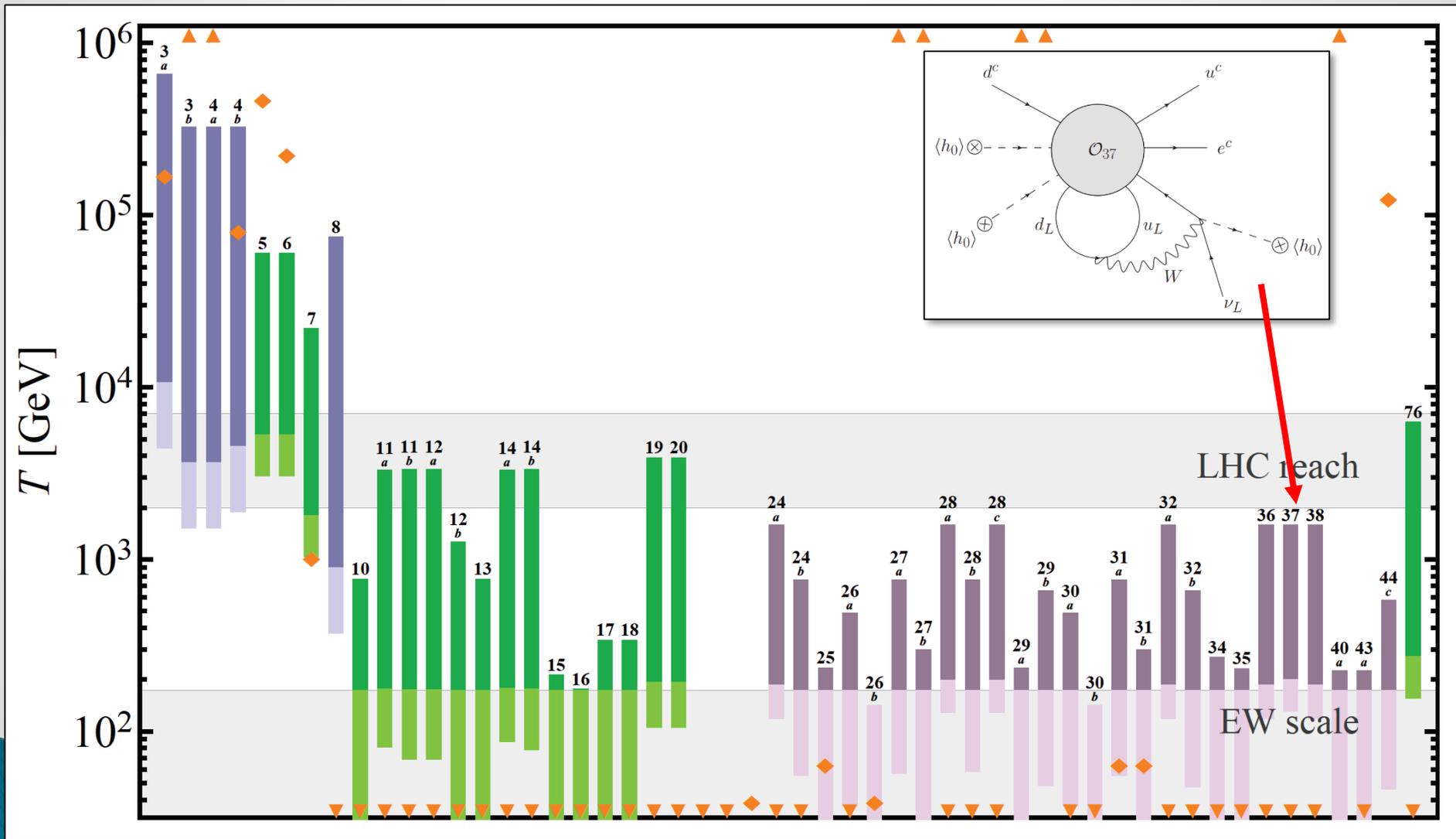
Lepton Asymmetry Washout

- ▶ Temperature ranges of strong equilibration
 - Assumes observation of corresponding process!
- ▶ Observation of LN(F)V
 - gives information at what temperatures operators are in equilibrium
 - **can falsify high-scale baryogenesis scenarios**



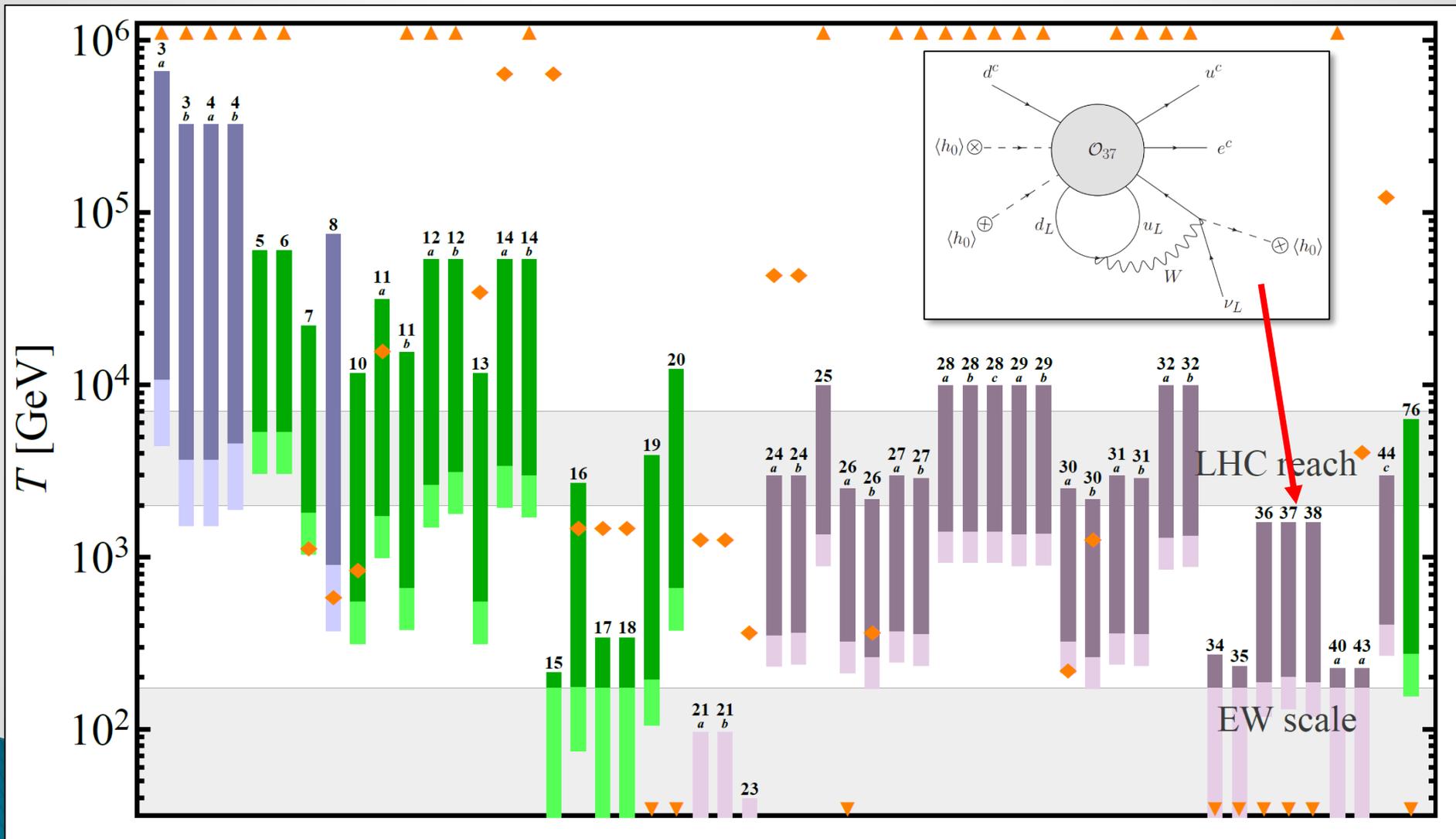
Baryon Asymmetry

Lepton Asymmetry Washout



Baryon Asymmetry

Lepton Asymmetry Washout



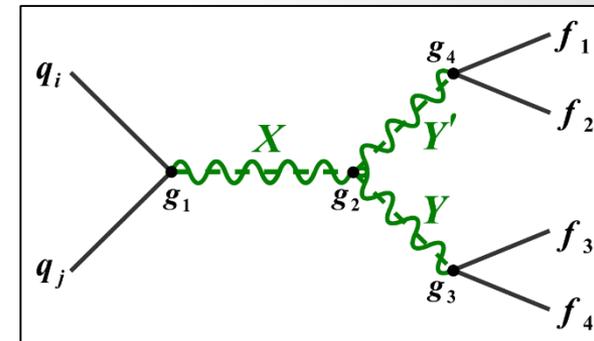
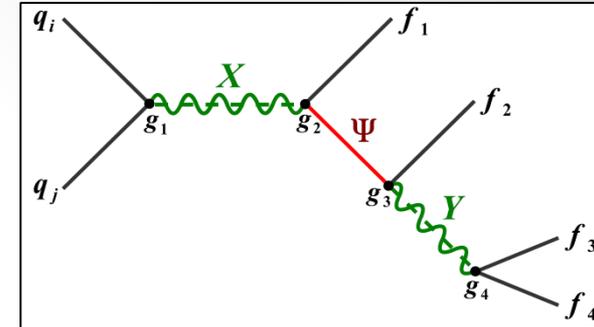
- ▶ Compare LHC cross section with lepton number asymmetry washout

$$\frac{\Gamma_W}{H} > 3 \times 10^{-3} \frac{M_P M_X^3}{T^4} \frac{K_1(M_X/T)}{f_{q_1 q_2}(M_X/\sqrt{s})} \times (s \sigma_{\text{LHC}})$$

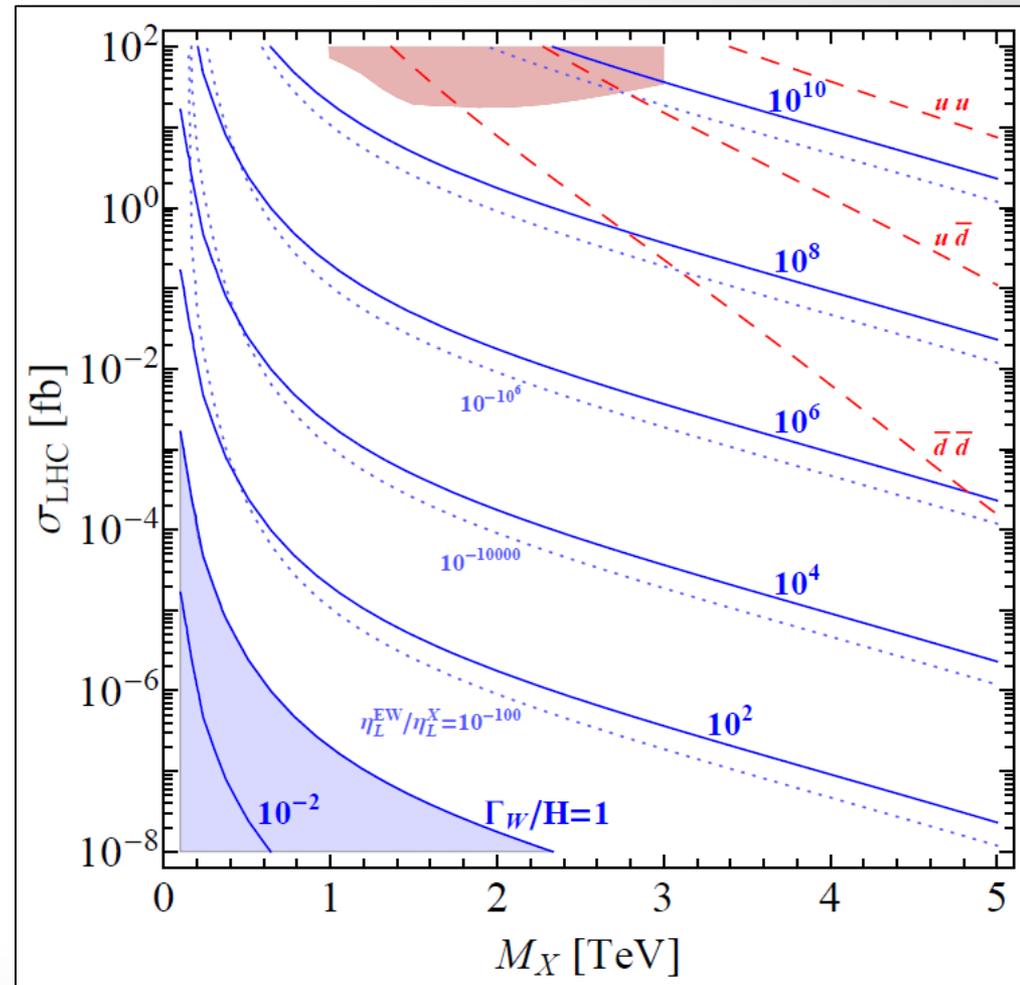
- Lower limit on total washout rate
 - Neglecting other washout processes

$$\log_{10} \frac{\Gamma_W}{H} > 7 + 0.6 \left(\frac{M_X}{\text{TeV}} - 1 \right) + \log_{10} \frac{\sigma_{\text{LHC}}}{\text{fb}}$$

- Observation of LNV @ LHC corresponds to highly effective washout $\Gamma_W/H \gg 1$
 - Excludes Leptogenesis models that generate asymmetry above M_X

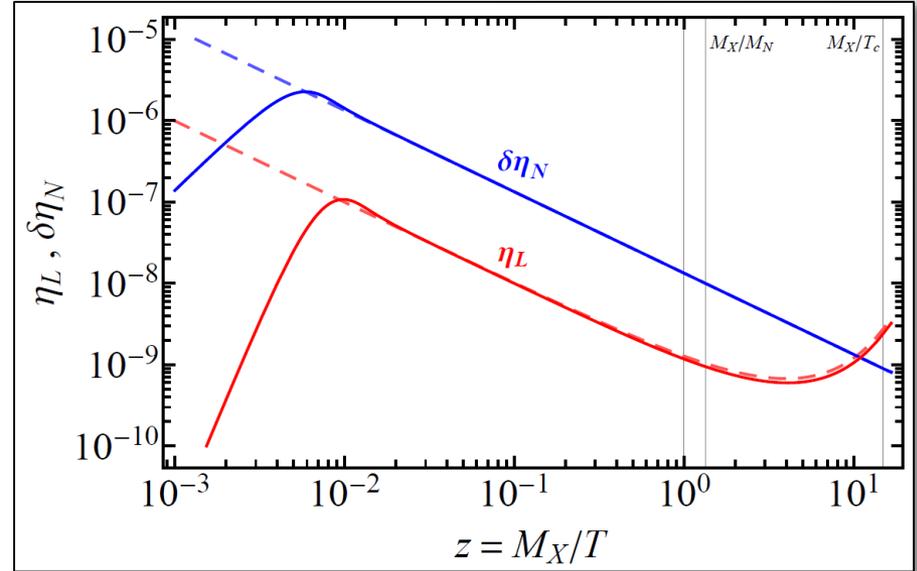


- ▶ Compare LHC cross section with lepton number asymmetry washout
 - Lower limit on total washout rate
 - Observation of LNV @ LHC corresponds to highly effective washout $\Gamma_W/H \gg 1$
 - Excludes Leptogenesis models that generate asymmetry above M_X



Caveats

- ▶ Cannot exclude scenarios that generate a lepton number asymmetry below observed scale M_X
 - But strong limits still apply
- ▶ Asymmetry can be present in one lepton generation only
 - Unambiguous falsification requires observation of LNV in all flavours (or observation of low energy LFV such as $\tau \rightarrow e\gamma$)
- ▶ Sphalerons only affect LH leptons...
 What if LNV is observed for RH leptons only?
 - Not an issue as all LH and RH charged fermions are in thermal equilibrium $\approx M_{EW}$
- ▶ Symmetry in new sector coupled via hypercharge induces $(B - L)$ chemical potential (Antaramian, Hall, Rašin '93)

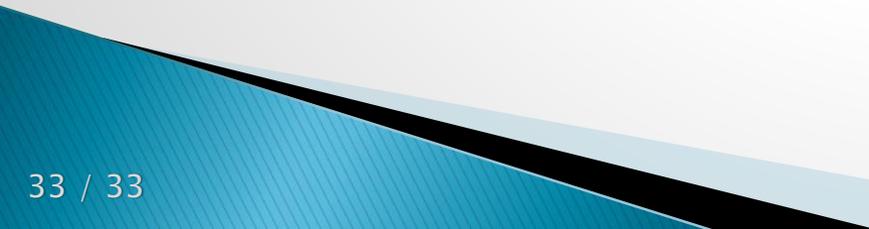


Conclusion

- ▶ **LNV a crucial BSM signature**
 - Majorana neutrino mass models
 - Baryogenesis via Leptogenesis
 - ▶ **Observations of LNV (and LFV) processes**
 - Tell us the temperature regime where leptons–antileptons (and different flavours) are equilibrated
 - Can falsify high scale baryogenesis scenarios
 - ▶ **Bottom–up approach**
 - Experimental data → Constrained model–landscape
 - ▶ **Important information for model selection, e.g.**
 - Observation of $0\nu\beta\beta$
 - Observation of LNV @ LHC
- } LNV @ TeV Scale
} Disfavours high–scale seesaw

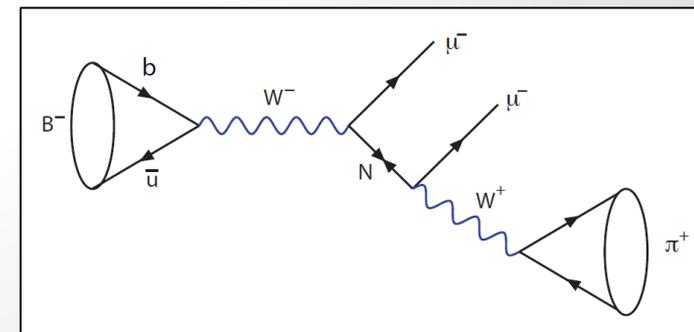
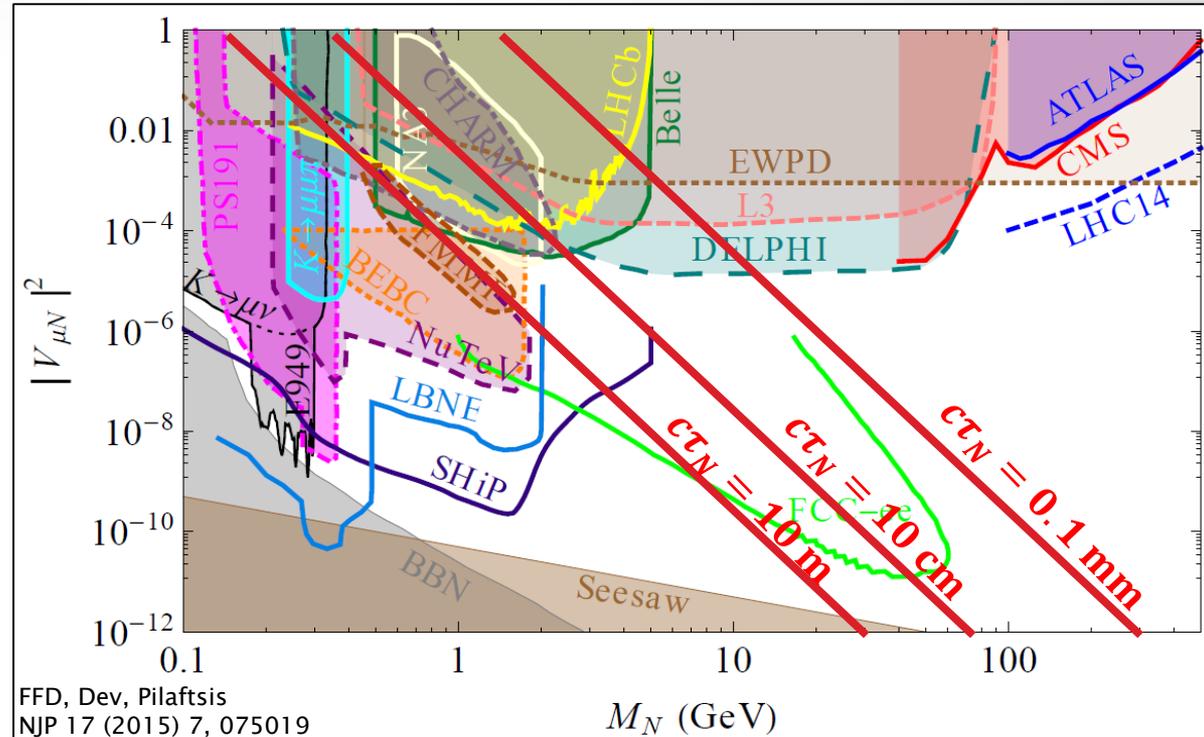
Conclusion

- ▶ **LNV a crucial BSM signature**
 - Majorana neutrino mass models
 - Baryogenesis via Leptogenesis
 - ▶ **Observations of LNV (and LFV) processes**
 - Tell us the temperature regime where leptons–antileptons (and different flavours) are equilibrated
 - Can falsify high scale baryogenesis scenarios
 - ▶ **Bottom–up approach**
 - Experimental data → Constrained model–landscape
 - ▶ **Important information for model selection, e.g.**
 - Observation of $0\nu\beta\beta$
 - No observation of LNV @ LHC
- } Improved confidence in standard $0\nu\beta\beta$ mechanism



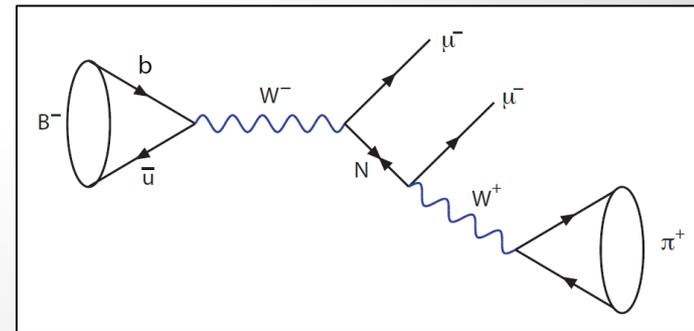
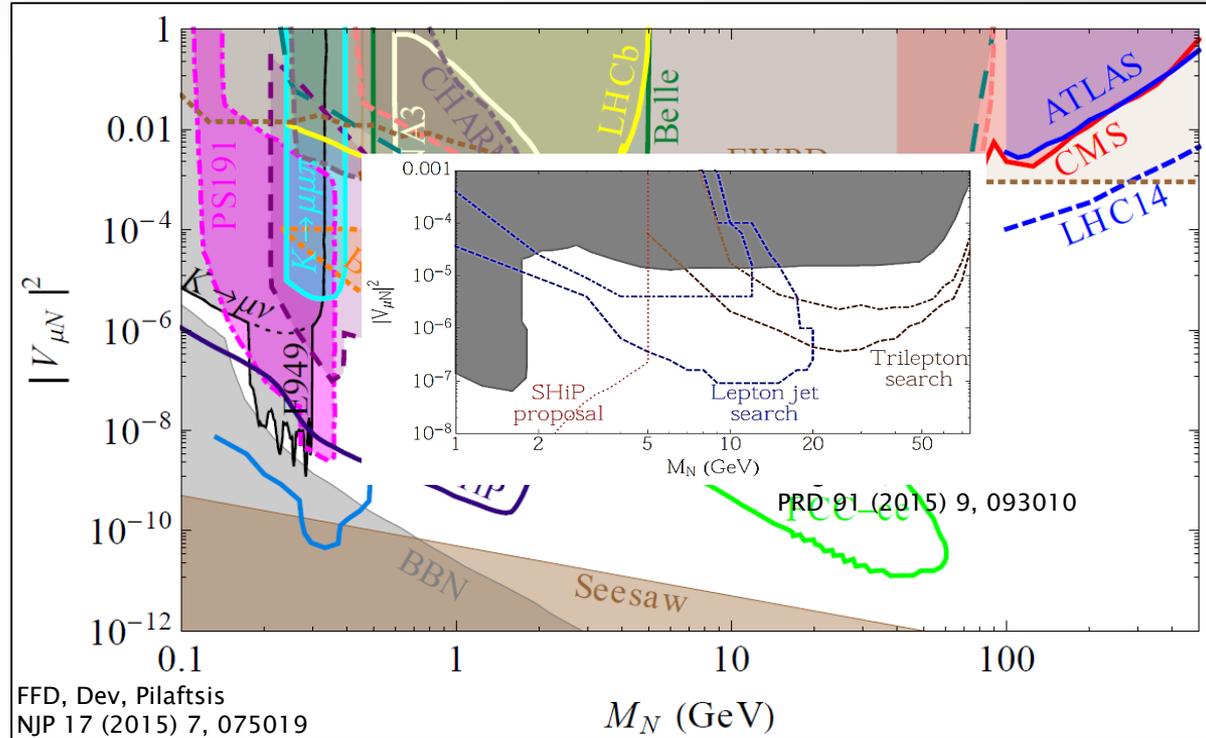
Heavy Sterile Neutrinos

- ▶ Constraints on coupling to leptons $|V_{lN}|$
- ▶ Neutrinoless Double Beta Decay
 - GERDA
 - stringent for pure Majorana N
- ▶ Peak Searches in Meson Decays
 - $\pi, K \rightarrow e\nu$
 - Belle
- ▶ Beam Dump Experiments
 - e.g. PS191, CHARM
 - LBNE
- ▶ LNV Meson Decays
 - $K \rightarrow e e \pi$
 - SHiP
- ▶ Z Decays
 - LEP: L3, Delphi
 - FCC-ee
- ▶ Electroweak Precision Tests
 - EWPd: Fit of electroweak precision observables, lepton universality observables



Heavy Sterile Neutrinos

- ▶ Constraints on coupling to leptons $|V_{lN}|$
- ▶ LEP2, ILC
 $e^+e^- \rightarrow N\nu$, $N \rightarrow eW, \nu Z, \nu H$
- ▶ LHC (ATLAS, CMS, LHC14)
 - Drell-Yan Production
 - Majorana N
 - Same-sign dilepton signal
 - (Quasi-)Dirac N
 - Trilepton signal
 - Modified searches for
 - Lighter neutrinos
 - Long-lived neutrinos



Not so Heavy Sterile Neutrinos

