Neutrinoless double beta decay in effective field theory

with
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E. Mereghetti, M. Piarulli, S. Pastore,
B. van Kolck, A. Walker-Loud, B. Wiringa

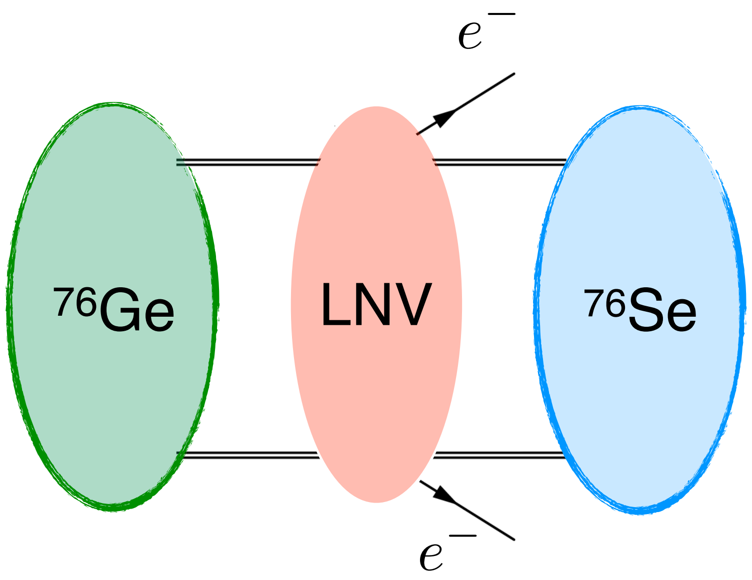
Based on:

arXiv:1806.02780, 1710.01729, 1802.10097,
1710.05026, 1708.09390

UC San Diego

Fluffy introduction

0νββ



= **BSM physics**

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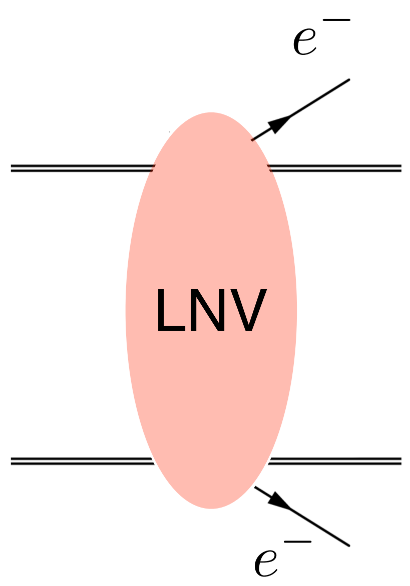

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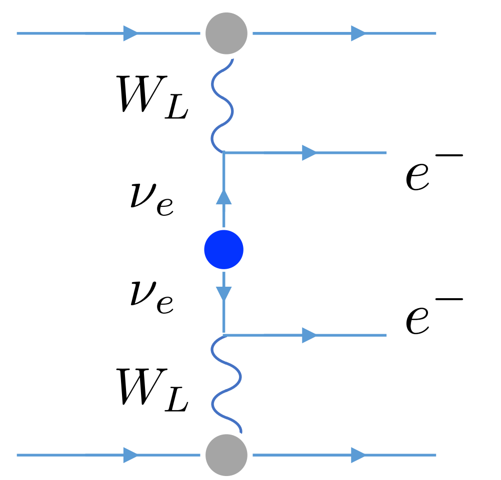
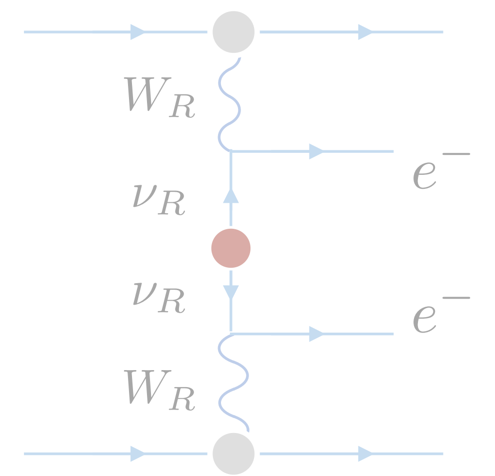
Light Majorana mass?

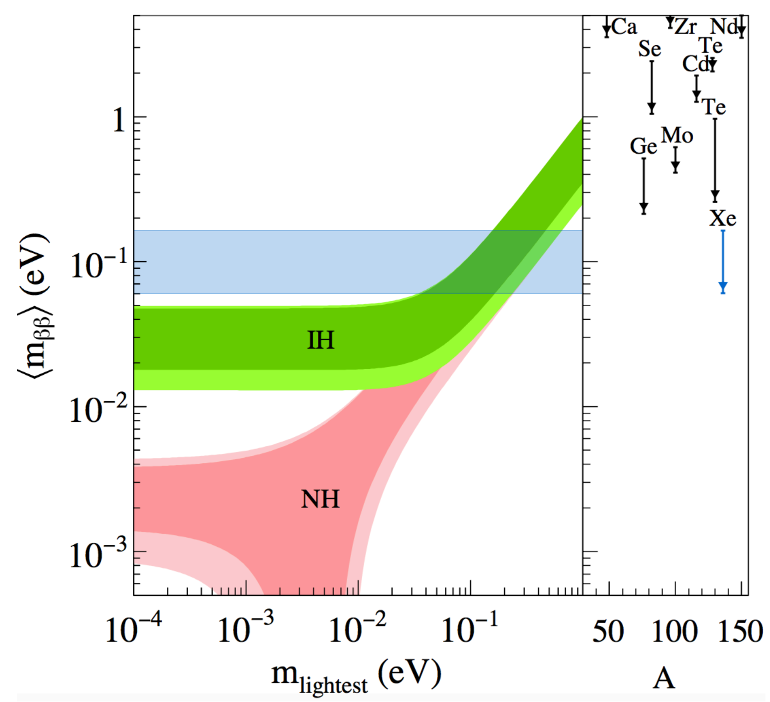


Left-right model?



+ ??

Well-known Majorana mass mechanism



- Implications for the mass hierarchy

Fluffy introduction

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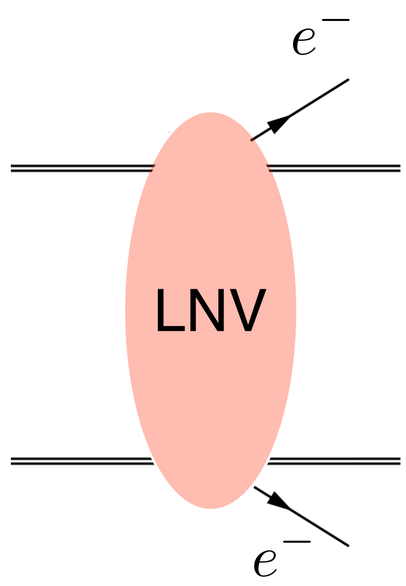

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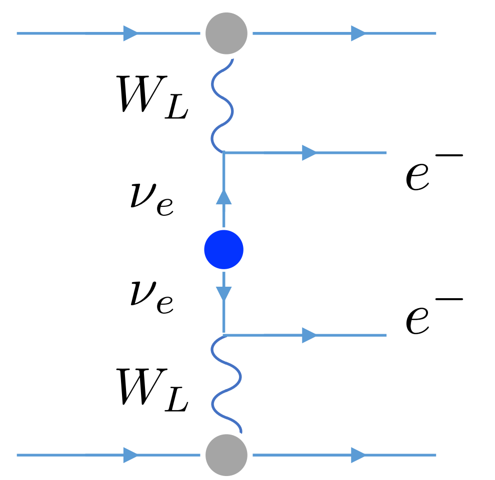

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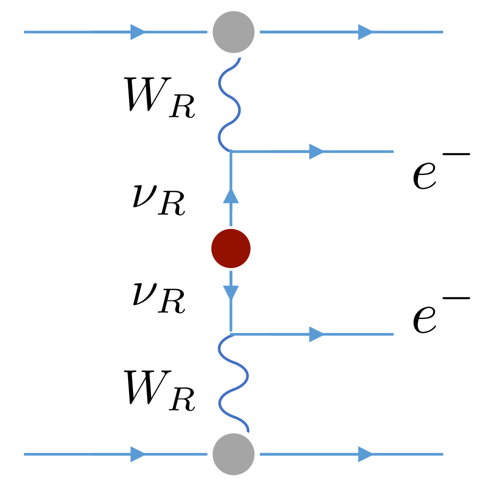

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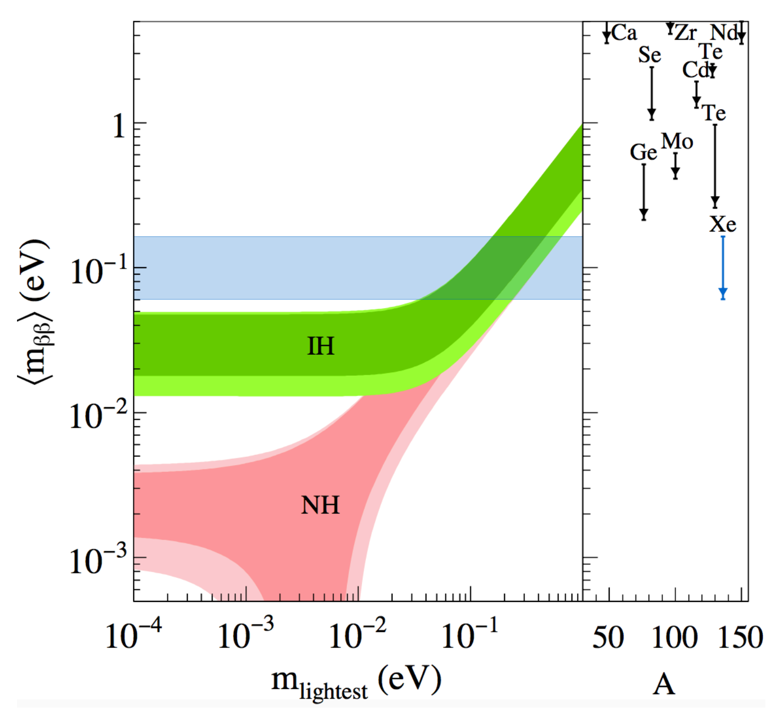

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Heavy BSM mechanisms

- Many possible scenarios
  - Left-right model,
  - R-parity violating SUSY
  - Leptoquarks…

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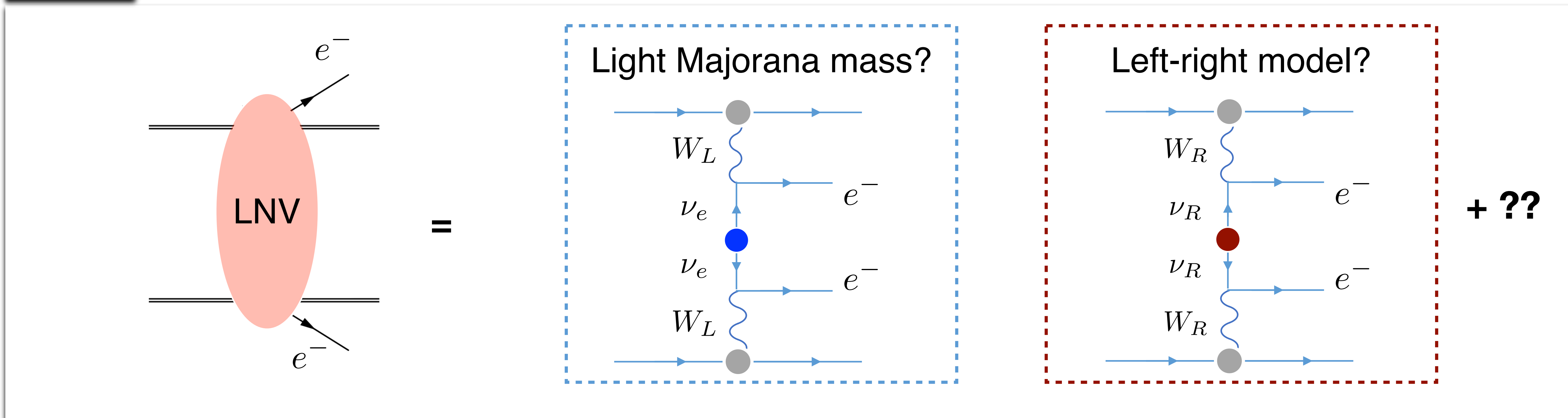

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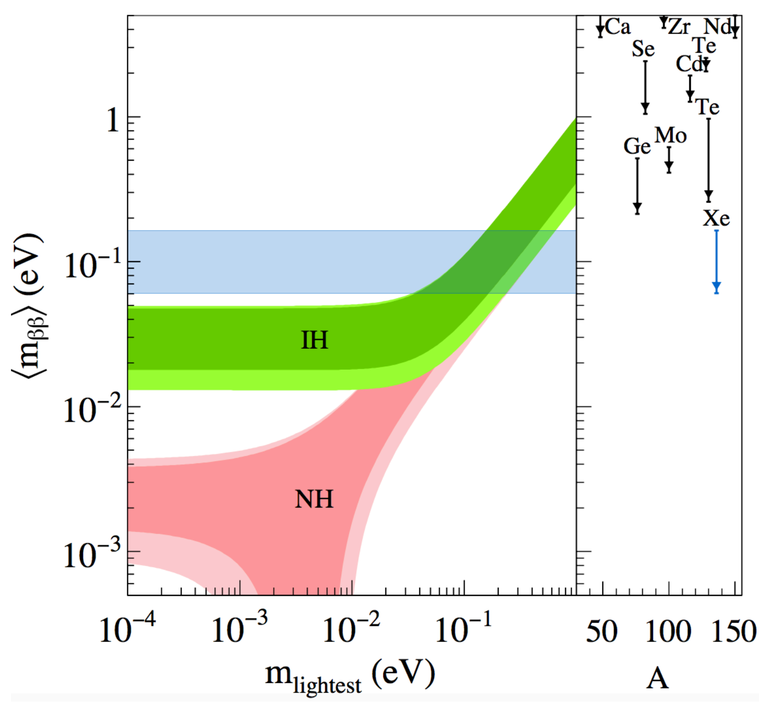

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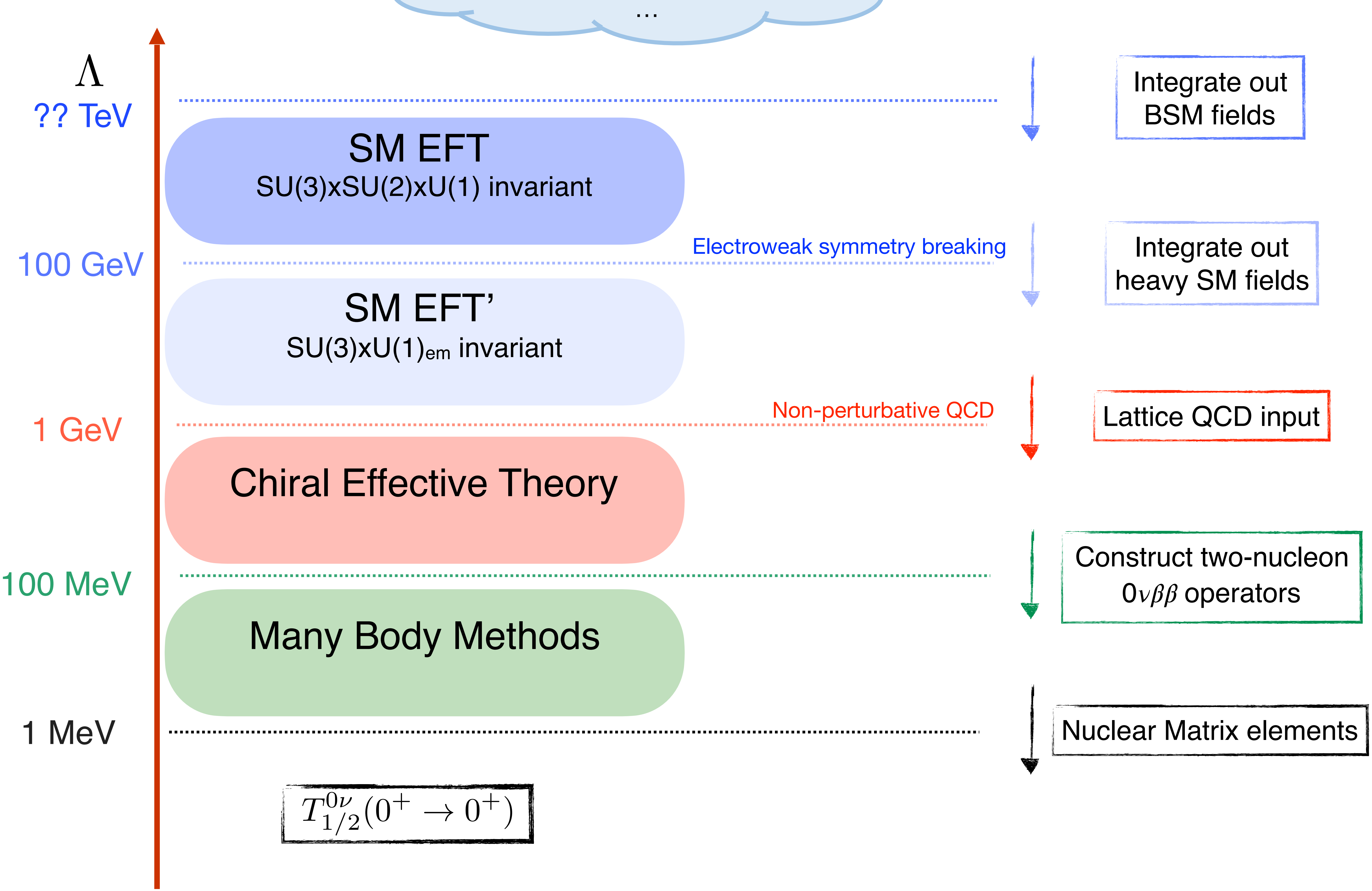
  - Leptoquarks…
- How to describe all LNV sources systematically?

Outline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Many Body Methods
$T^{0\nu}_{1/2}(0^+ \to 0^+)$
Integrate out BSM fields
Integrate out heavy SM fields
Lattice QCD input
Construct two-nucleon $0\nu\beta\beta$ operators
Nuclear Matrix elements

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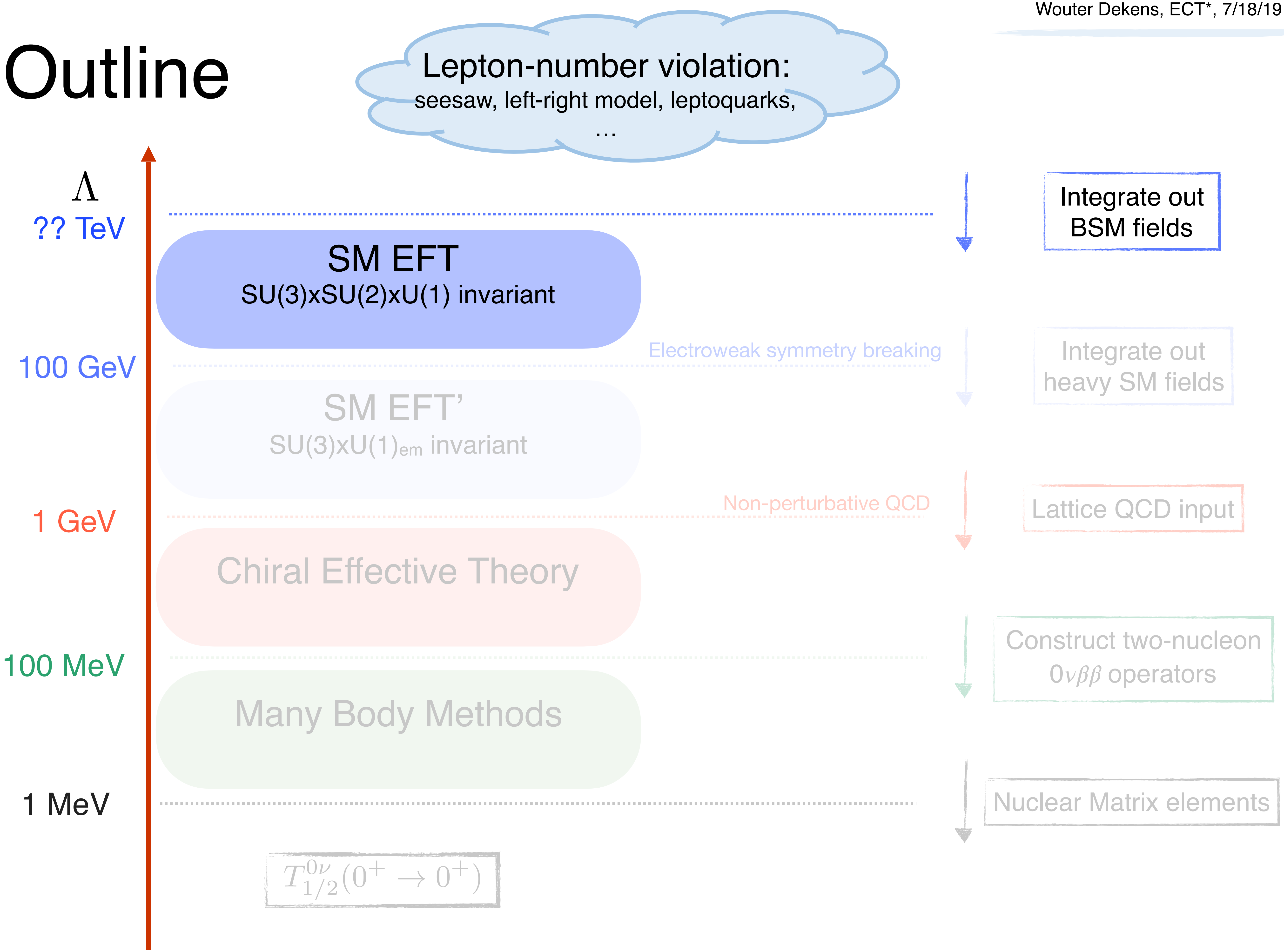
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$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \frac{c_i^{(5)}}{\Lambda} O_i^{(5)} + \frac{c_i^{(7)}}{\Lambda^3} O_i^{(7)} + \frac{c_i^{(9)}}{\Lambda^5} O_i^{(9)} + \ldots$$

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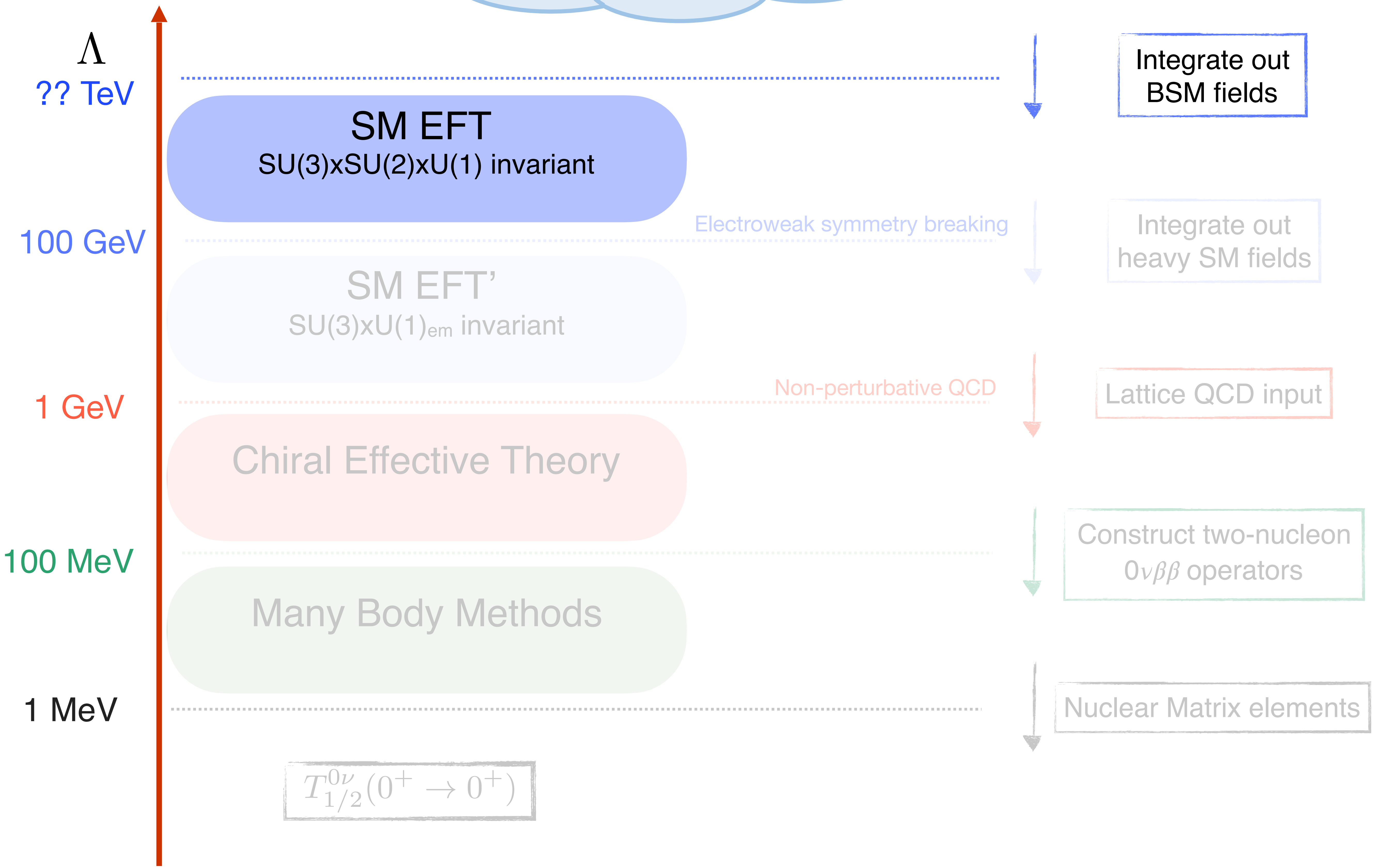

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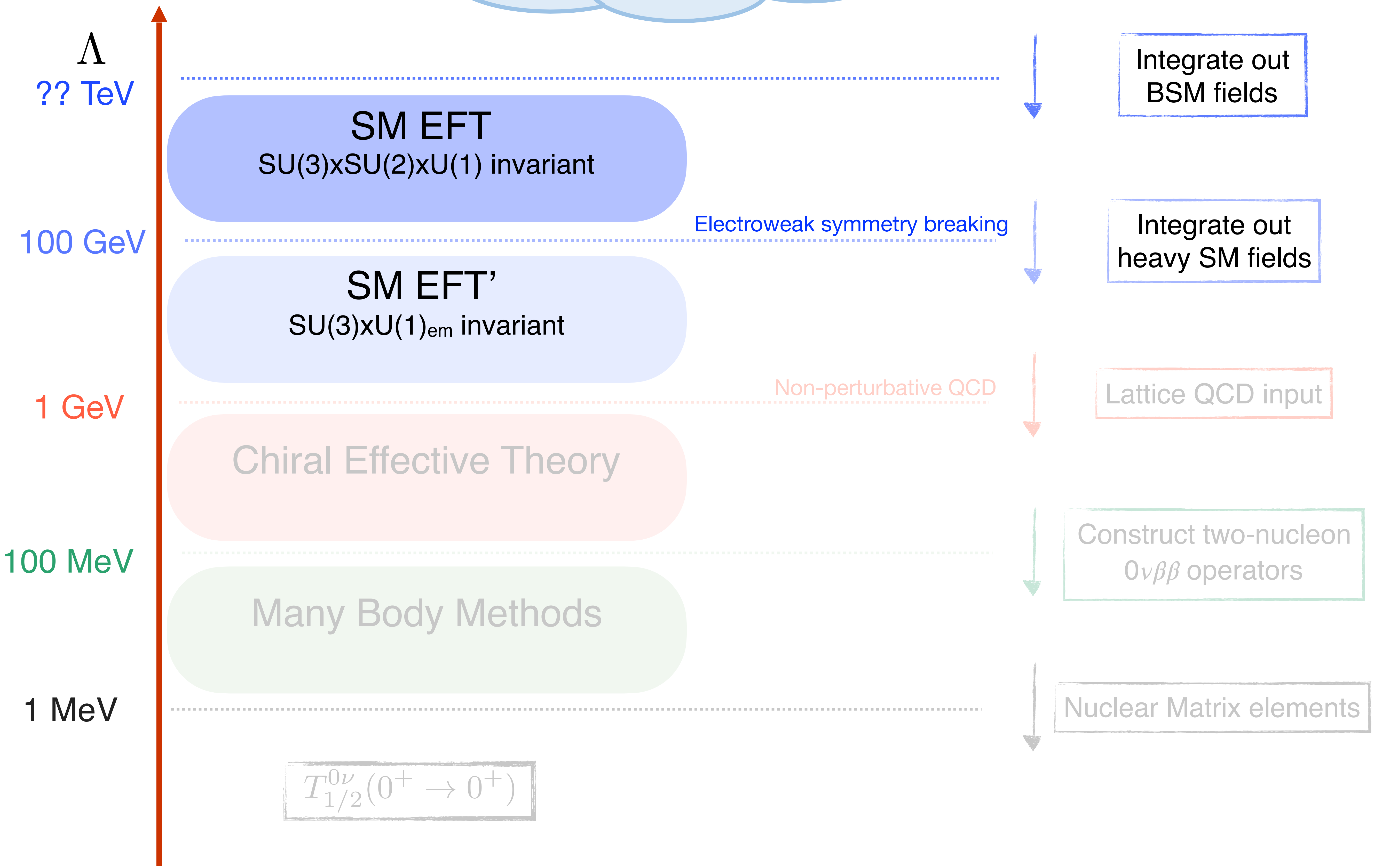


Running/matching at the weak scale

SM EFT

SU(3)xSU(2)xU(1) invariant

$$\mathcal{L} = \frac{c_i^{(5)}}{\Lambda} O_i^{(5)} + \frac{c_i^{(7)}}{\Lambda^3} O_i^{(7)} + \frac{c_i^{(9)}}{\Lambda^5} O_i^{(9)}$$

M_{EW}

100 GeV

Electroweak symmetry breaking

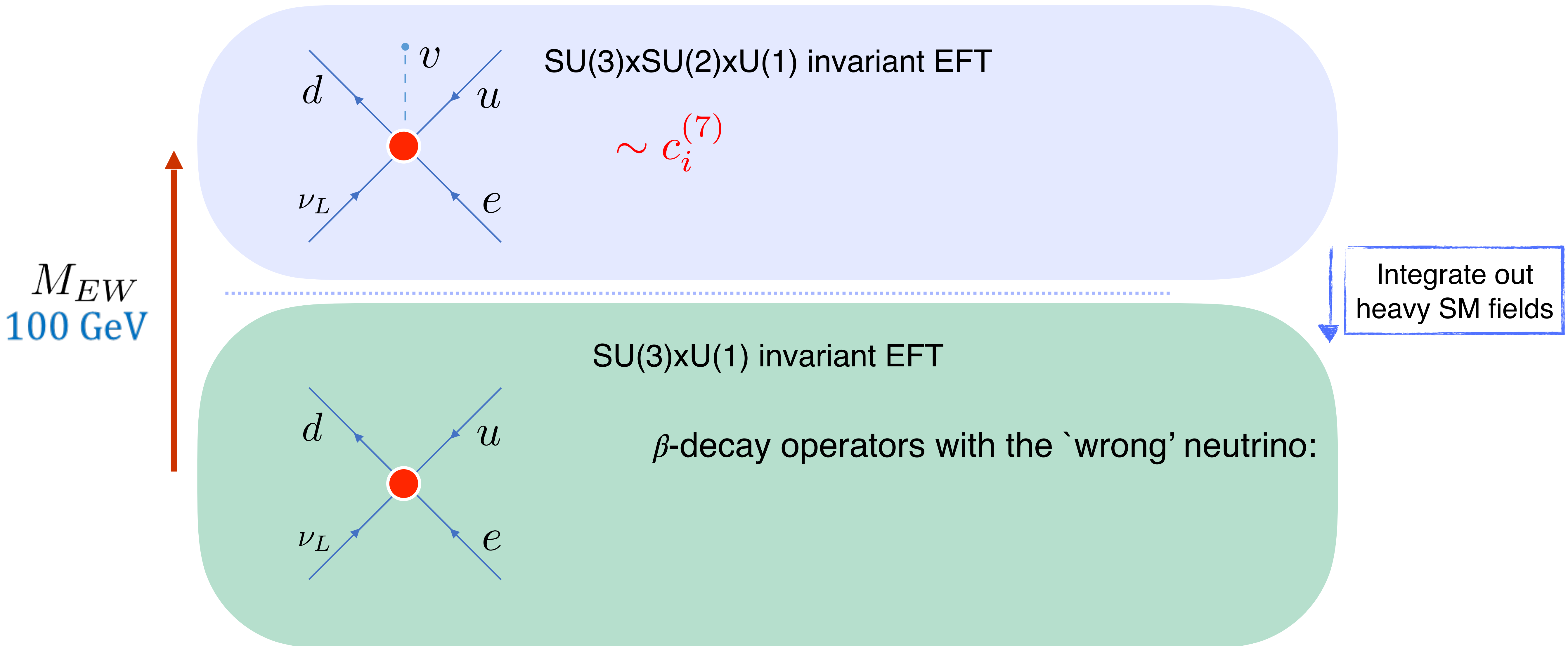
SM EFT'

SU(3)xU(1) invariant

$$\mathcal{L} = \frac{c_i^{(3)}}{\Lambda} O_i^{(3)} + \frac{c_i^{(6)}}{\Lambda^3} O_i^{(6)} + \frac{c_i^{(7)}}{\Lambda^3} O_i^{(7)} + \frac{c_i^{(9)}}{\Lambda^5} O_i^{(9)}$$

- Mismatch in dimensions due to insertions of the Higgs vacuum expectation value

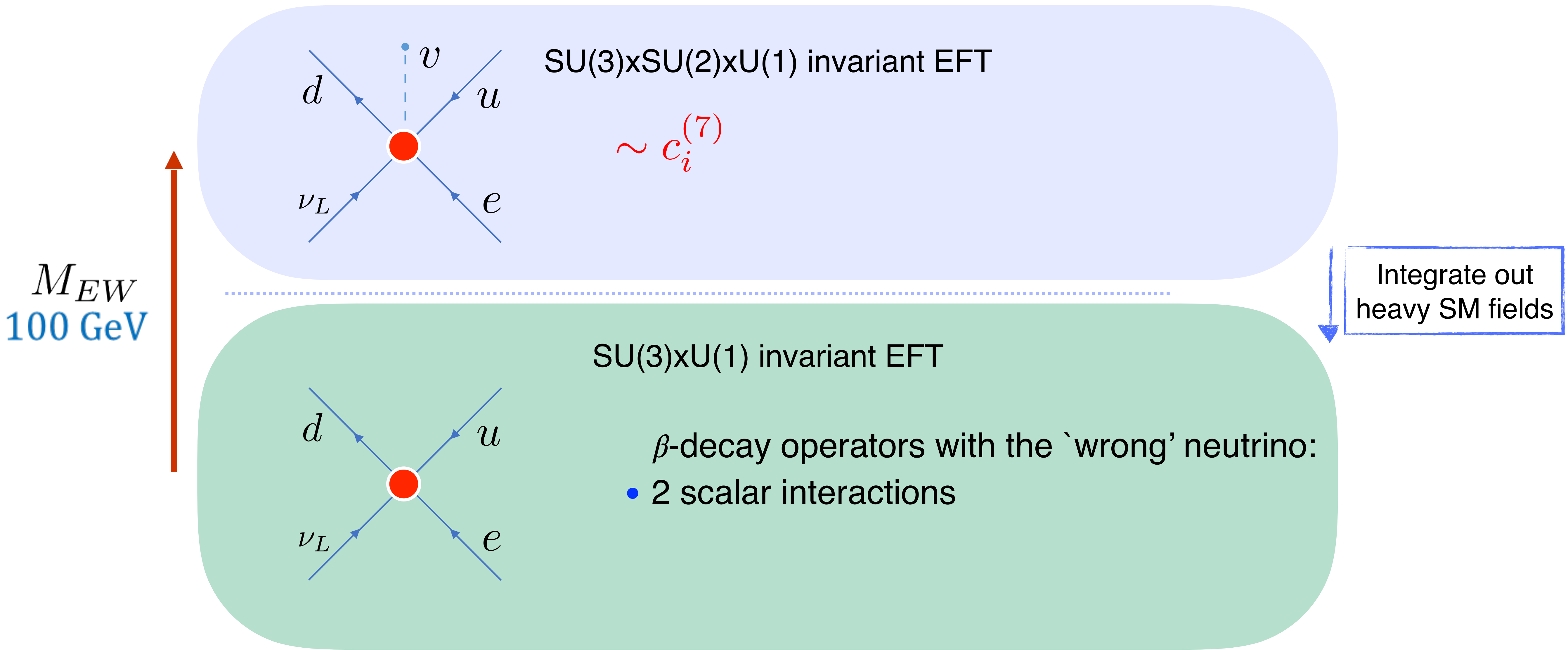
Low-energy operators
Dimension-6



$$\mathcal{L}^{(6)}_{\Delta L=2} = \frac{2G_F}{\sqrt{2}}\Big\{ C^{(6)}_{\mathrm{VL},ij}\,\bar u_L\gamma^\mu d_L\,\bar e_{R,i}\,\gamma_\mu\,C\bar\nu^T_{L,j} + C^{(6)}_{\mathrm{VR},ij}\,\bar u_R\gamma^\mu d_R\,\bar e_{R,i}\,\gamma_\mu\,C\bar\nu^T_{L,j}$$

$$+C^{(6)}_{\mathrm{SR},ij}\,\bar u_L d_R\,\bar e_{L,i}\,C\bar\nu^T_{L,j} + C^{(6)}_{\mathrm{SL},ij}\,\bar u_R d_L\,\bar e_{L,i}\,C\bar\nu^T_{L,j} + C^{(6)}_{\mathrm{T},ij}\,\bar u_L\sigma^{\mu\nu}d_R\,\bar e_{L,i}\sigma_{\mu\nu}\,C\bar\nu^T_{L,j}\Big\}$$

Low-energy operators
Dimension-6



$$\mathcal{L}^{(6)}_{\Delta L=2} = \frac{2G_F}{\sqrt{2}}\Big\{ C^{(6)}_{\mathrm{VL},ij}\,\bar u_L\gamma^\mu d_L\,\bar e_{R,i}\,\gamma_\mu\,C\bar\nu^T_{L,j} + C^{(6)}_{\mathrm{VR},ij}\,\bar u_R\gamma^\mu d_R\,\bar e_{R,i}\,\gamma_\mu\,C\bar\nu^T_{L,j}$$

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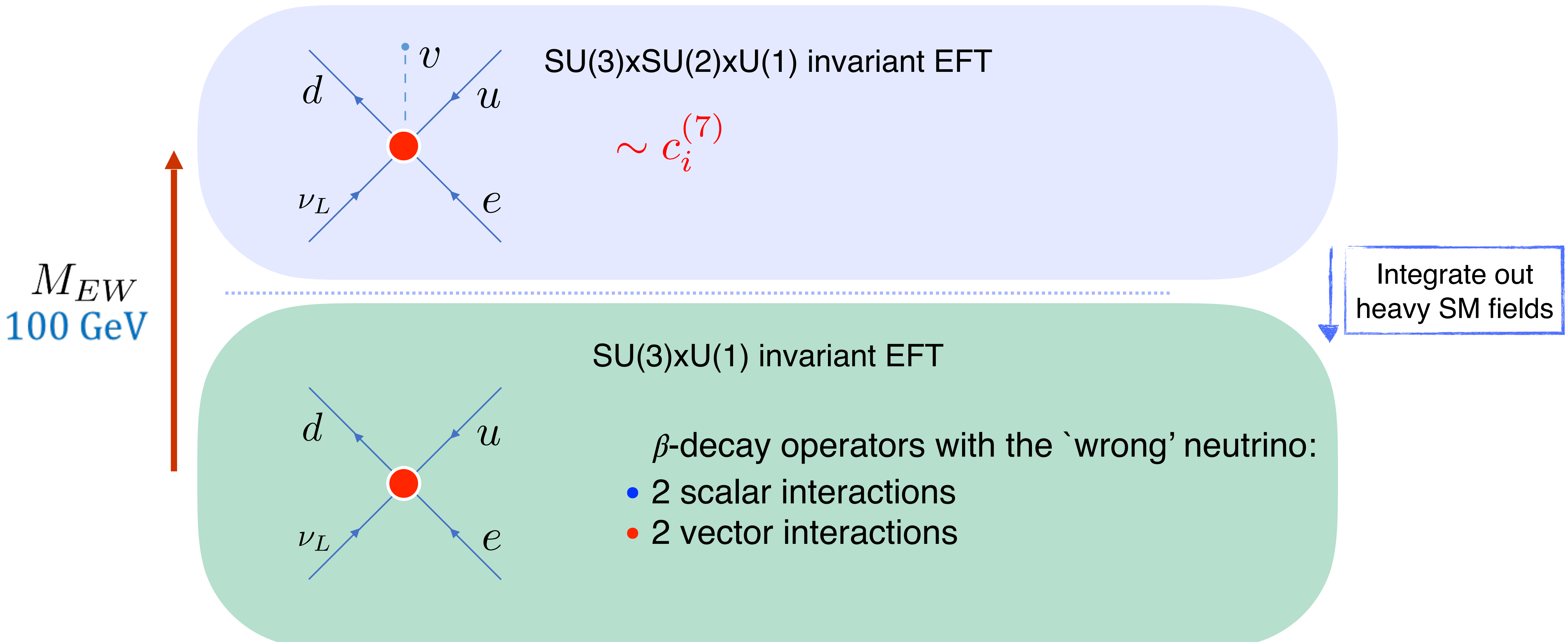
Low-energy operators

Dimension-6



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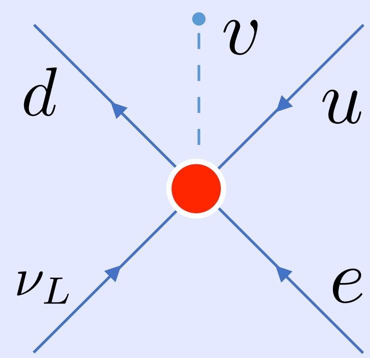
$$+ C^{(6)}_{\mathrm{SR},ij}\,\bar u_L d_R\,\bar e_{L,i}\, C\bar\nu^T_{L,j} + C^{(6)}_{\mathrm{SL},ij}\,\bar u_R d_L\,\bar e_{L,i}\, C\bar\nu^T_{L,j} + C^{(6)}_{\mathrm{T},ij}\,\bar u_L\sigma^{\mu\nu} d_R\,\bar e_{L,i}\sigma_{\mu\nu}\, C\bar\nu^T_{L,j}\Big\}$$

Low-energy operators
Dimension-6

M_{EW}
100 GeV

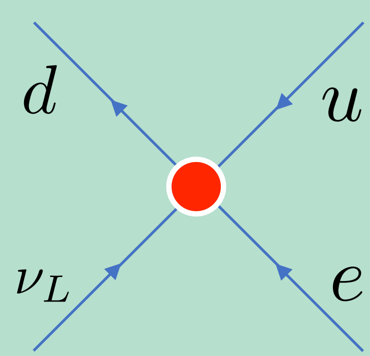


SU(3)xSU(2)xU(1) invariant EFT

$\sim c_i^{(7)}$

Integrate out heavy SM fields

SU(3)xU(1) invariant EFT



β-decay operators with the `wrong' neutrino:

- 2 scalar interactions
- 2 vector interactions
- one tensor interaction

$$\mathcal{L}^{(6)}_{\Delta L=2} = \frac{2G_F}{\sqrt{2}}\Big\{ C^{(6)}_{\mathrm{VL},ij}\,\bar{u}_L\gamma^\mu d_L\,\bar{e}_{R,i}\,\gamma_\mu\,C\bar{\nu}^T_{L,j} + C^{(6)}_{\mathrm{VR},ij}\,\bar{u}_R\gamma^\mu d_R\,\bar{e}_{R,i}\,\gamma_\mu\,C\bar{\nu}^T_{L,j}$$

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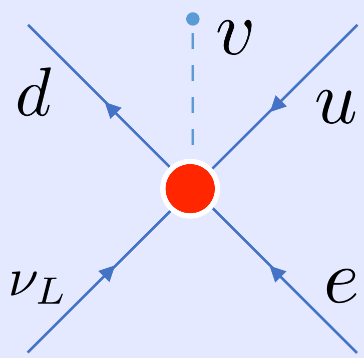
Low-energy operators
Dimension-6

M_{EW}
100 GeV



SU(3)xSU(2)xU(1) invariant EFT

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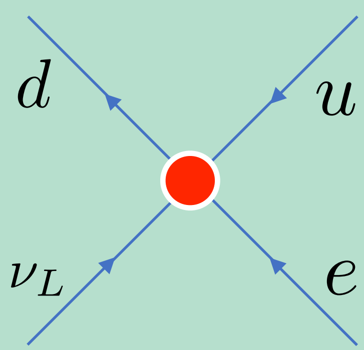
Integrate out heavy SM fields



SU(3)xU(1) invariant EFT

β-decay operators with the `wrong' neutrino:

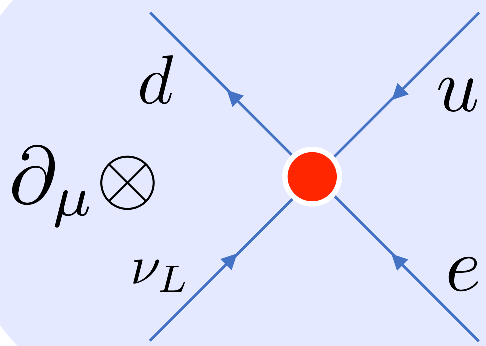
- 2 scalar interactions
- 2 vector interactions
- one tensor interaction

Induced by dimension-7 SU(2)-invariant operators

$$C_i^{(6)} \sim v^3/\Lambda^3$$

Low-energy operators
Dimension-7

SU(3)xSU(2)xU(1) invariant EFT



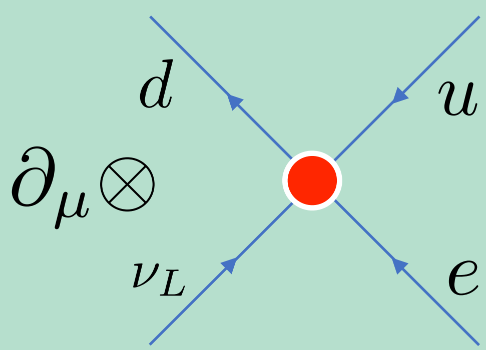
$\sim c_i^{(7)}$

M_{EW}

100 GeV

Integrate out heavy SM fields

SU(3)xU(1) invariant EFT



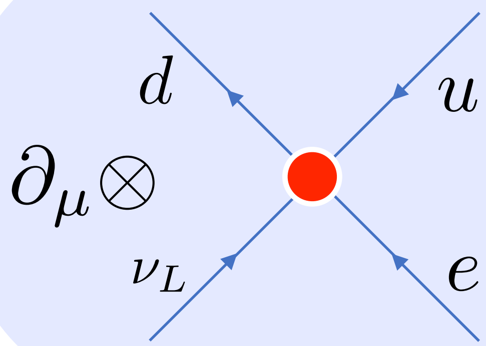
β-decay operators with the `wrong' neutrino and a derivative

- Two vector-like operators

$$\mathcal{L}^{(7)}_{\Delta L=2} = \frac{2G_F}{\sqrt{2}v}\left\{C^{(7)}_{\mathrm{VL},ij}\,\bar{u}_L\gamma^\mu d_L\,\bar{e}_{L,i}\,C\,i\overleftrightarrow{\partial}_\mu\bar{\nu}^T_{L,j} + C^{(7)}_{\mathrm{VR},ij}\,\bar{u}_R\gamma^\mu d_R\,\bar{e}_{L,i}\,Ci\overleftrightarrow{\partial}_\mu\bar{\nu}^T_{L,j}\right\} + \mathrm{h.c.}$$

Low-energy operators
Dimension-7

SU(3)xSU(2)xU(1) invariant EFT



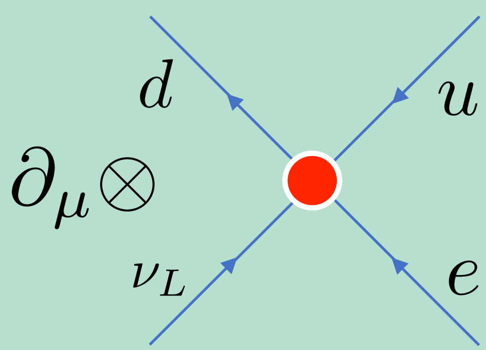
$\sim c_i^{(7)}$

M_{EW}

100 GeV

Integrate out heavy SM fields

SU(3)xU(1) invariant EFT



β-decay operators with the `wrong' neutrino and a derivative

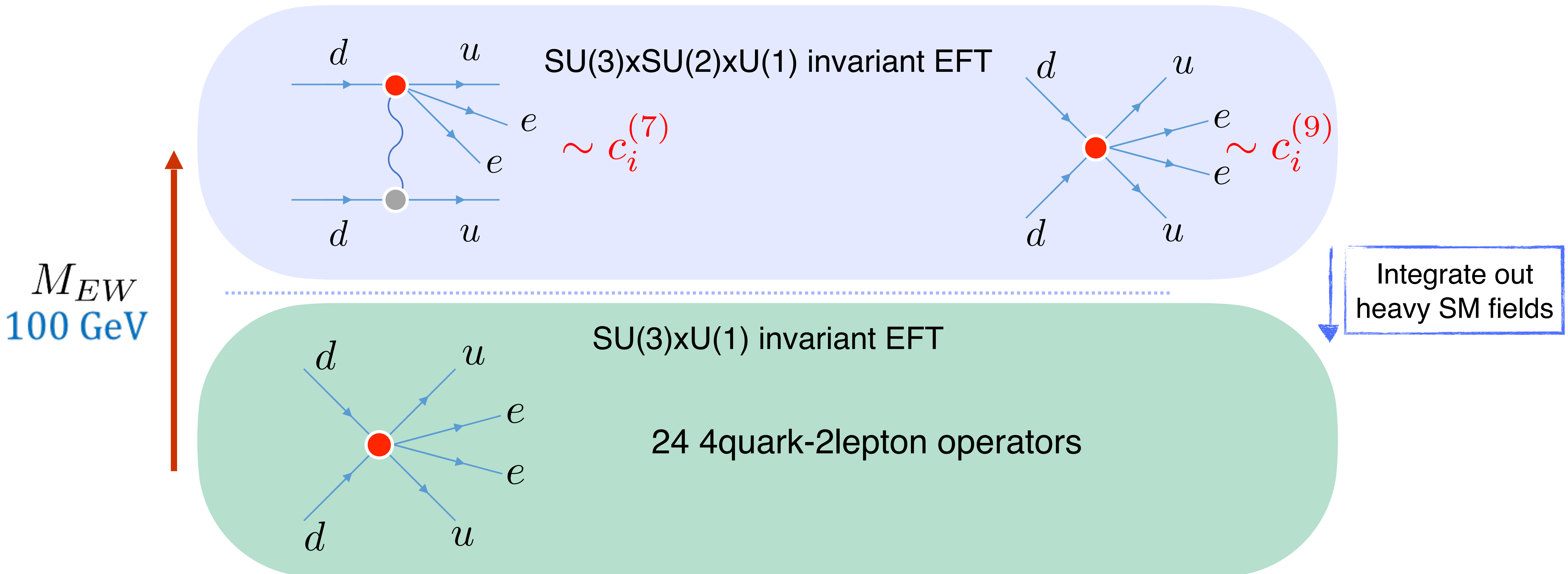
- Two vector-like operators

$\mathcal{L}^{(7)}_{\Delta L=2} =$ … $\overleftrightarrow{\partial}_\mu \bar{\nu}^T_{L,j} \Big\} + \text{h.c.}$

Induced by dimension-7 SU(2)-invariant operators

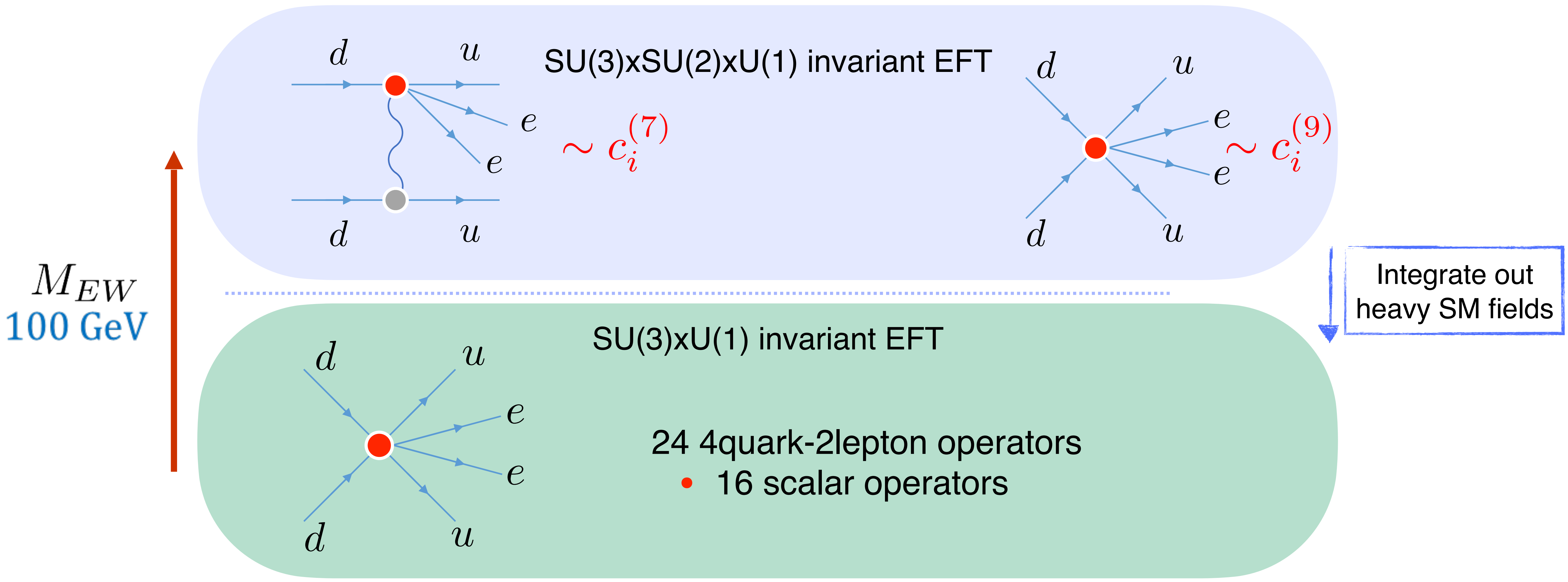
$$C_i^{(7)} \sim v^3/\Lambda^3$$

Low-energy operators
Dimension-9



$$\mathcal{L}^{(9)}_{\Delta L=2} = \frac{1}{v^5}\sum_i \left[\left(C^{(9)}_{i\,\mathrm{R}}\, \bar e_R C \bar e_R^T + C^{(9)}_{i\,\mathrm{L}}\, \bar e_L C \bar e_L^T \right) O_i + C^{(9)}_i \bar e \gamma_\mu\gamma_5 C \bar e^T O_i^\mu \right],$$

Low-energy operators
Dimension-9



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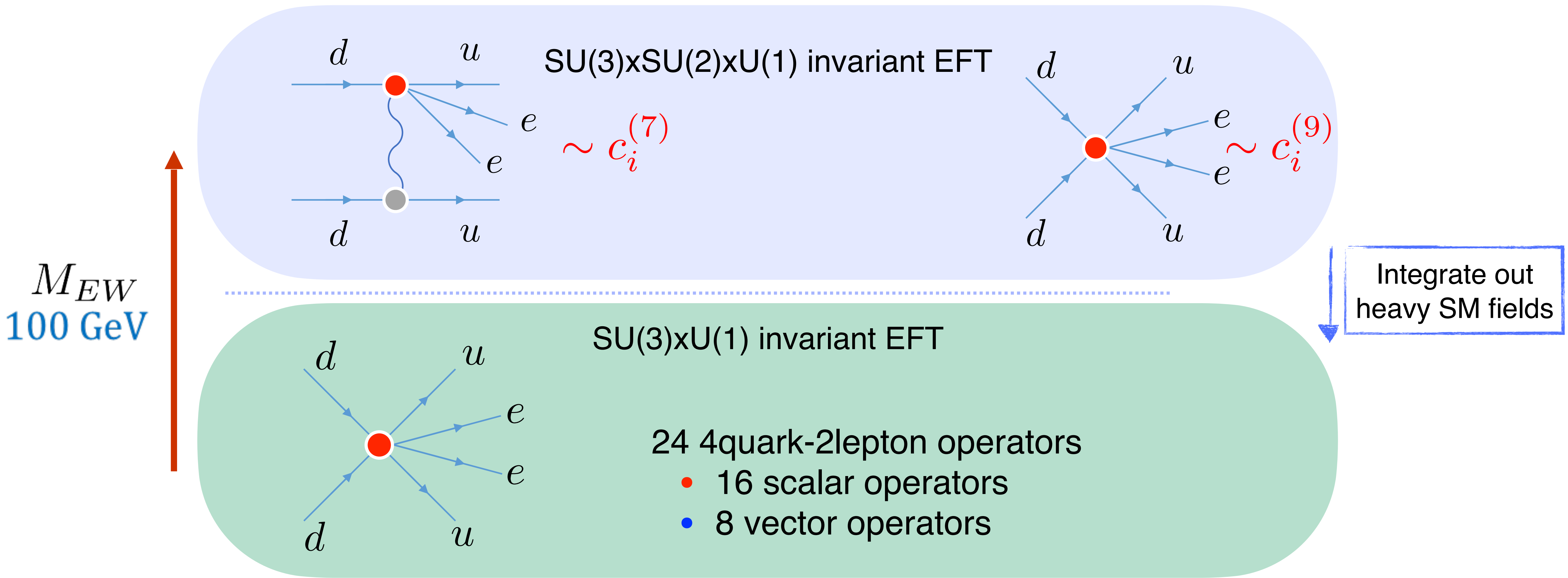
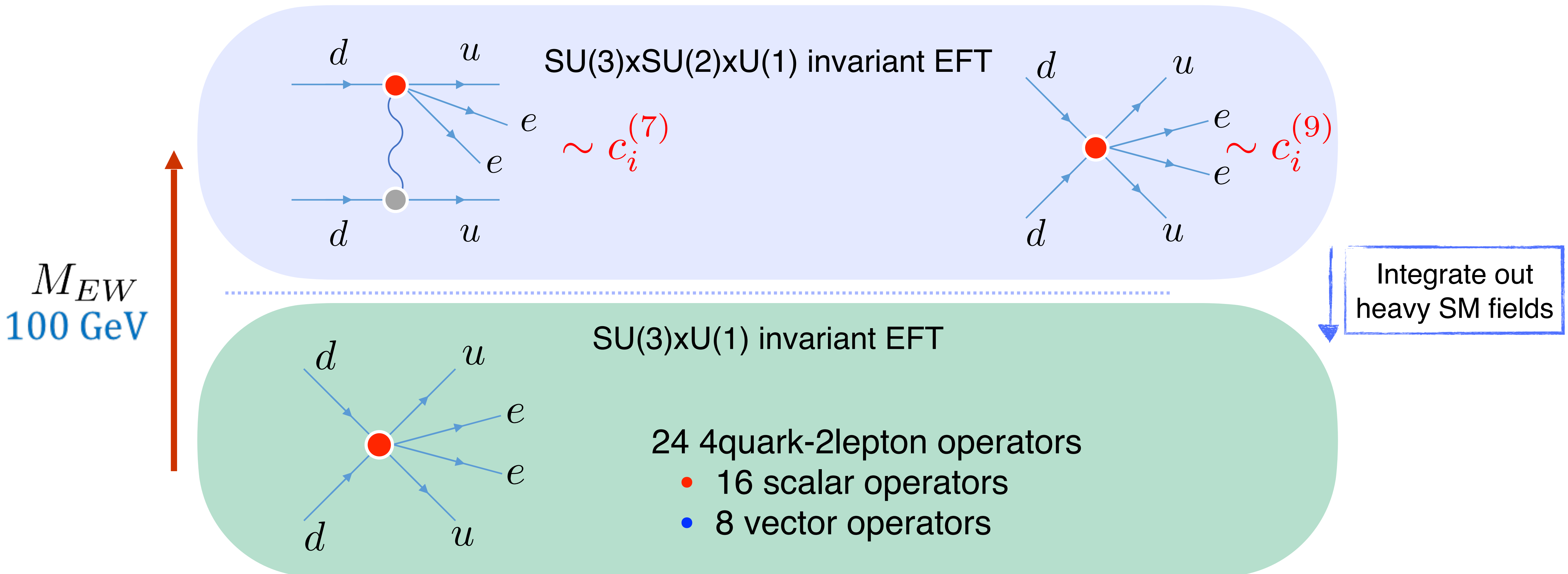
Low-energy operators

Dimension-9



$$\mathcal{L}^{(9)}_{\Delta L=2} = \frac{1}{v^5} \sum_i \left[\left(C^{(9)}_{i\,\mathrm{R}}\, \bar{e}_R C \bar{e}_R^T + C^{(9)}_{i\,\mathrm{L}}\, \bar{e}_L C \bar{e}_L^T \right) O_i + C^{(9)}_i \bar{e} \gamma_\mu \gamma_5 C \bar{e}^T O_i^\mu \right],$$

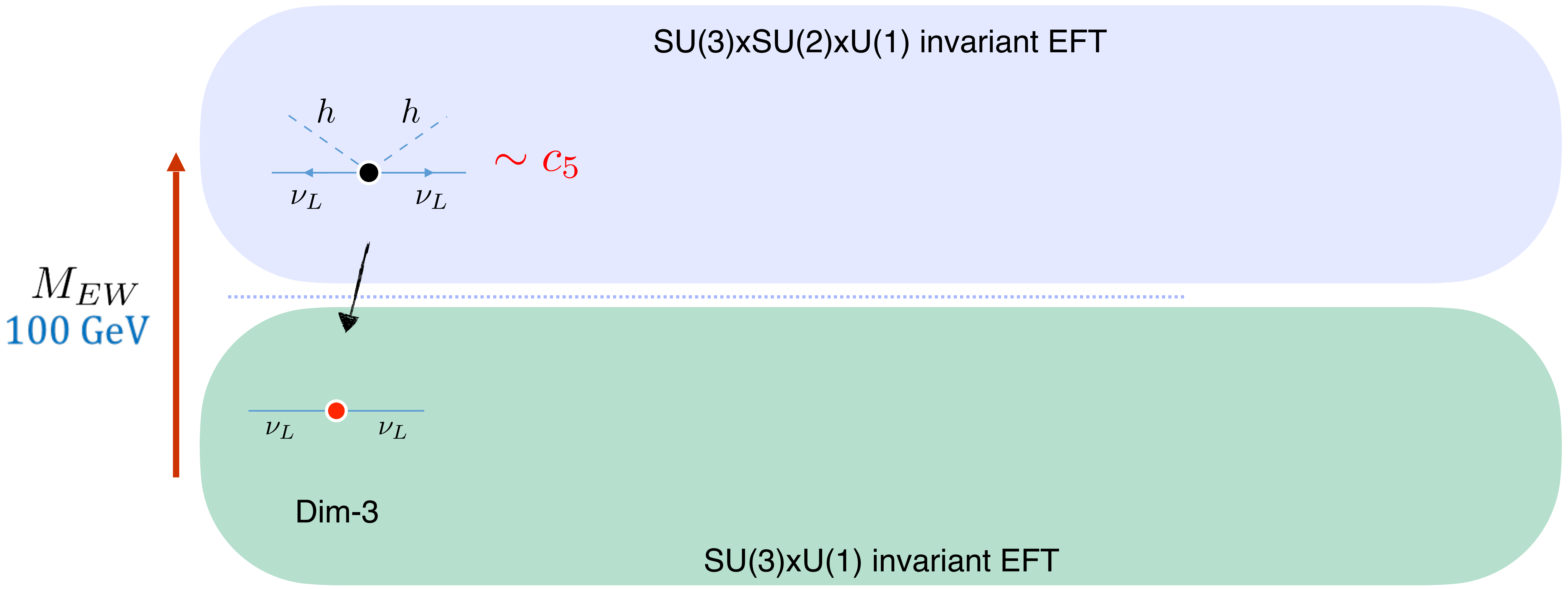
Low-energy operators
Dimension-9



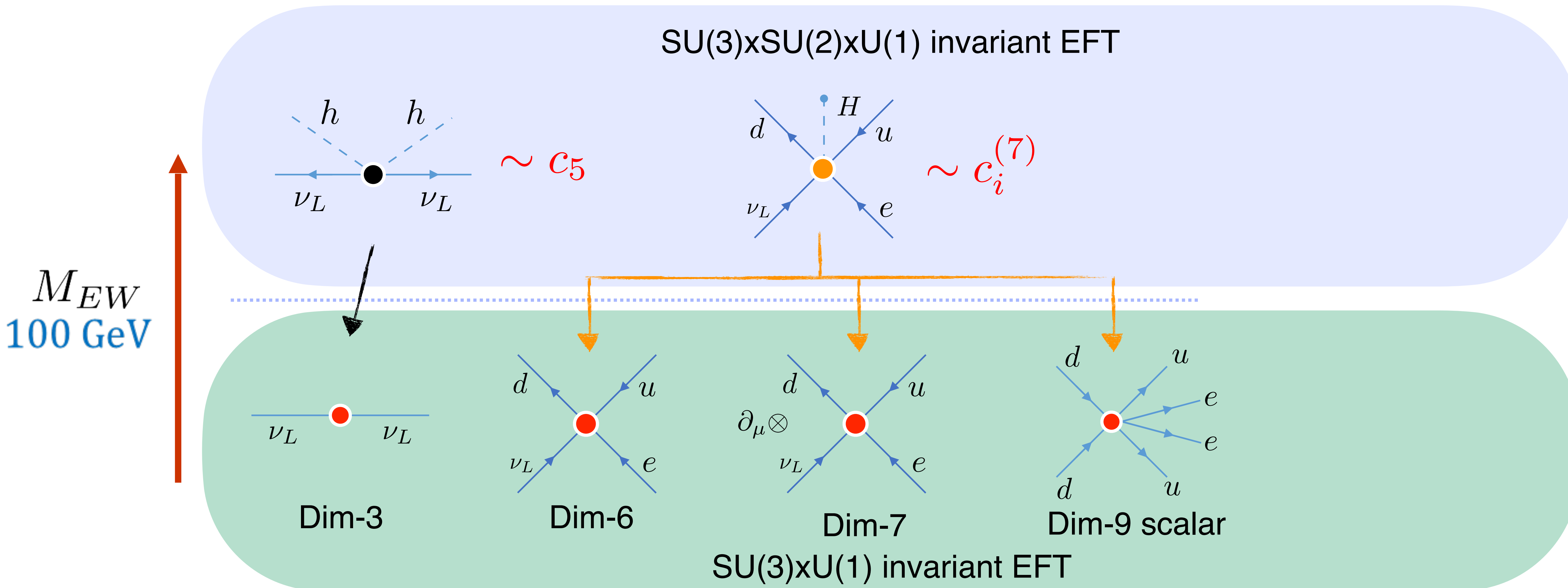
- 3 can be induced by dimension-7 operators $C_i^{(9)} \sim v^3/\Lambda^3$
- 19 can be induced by dimension-9 operators $C_i^{(9)} \sim v^5/\Lambda^5$

Low-energy operators

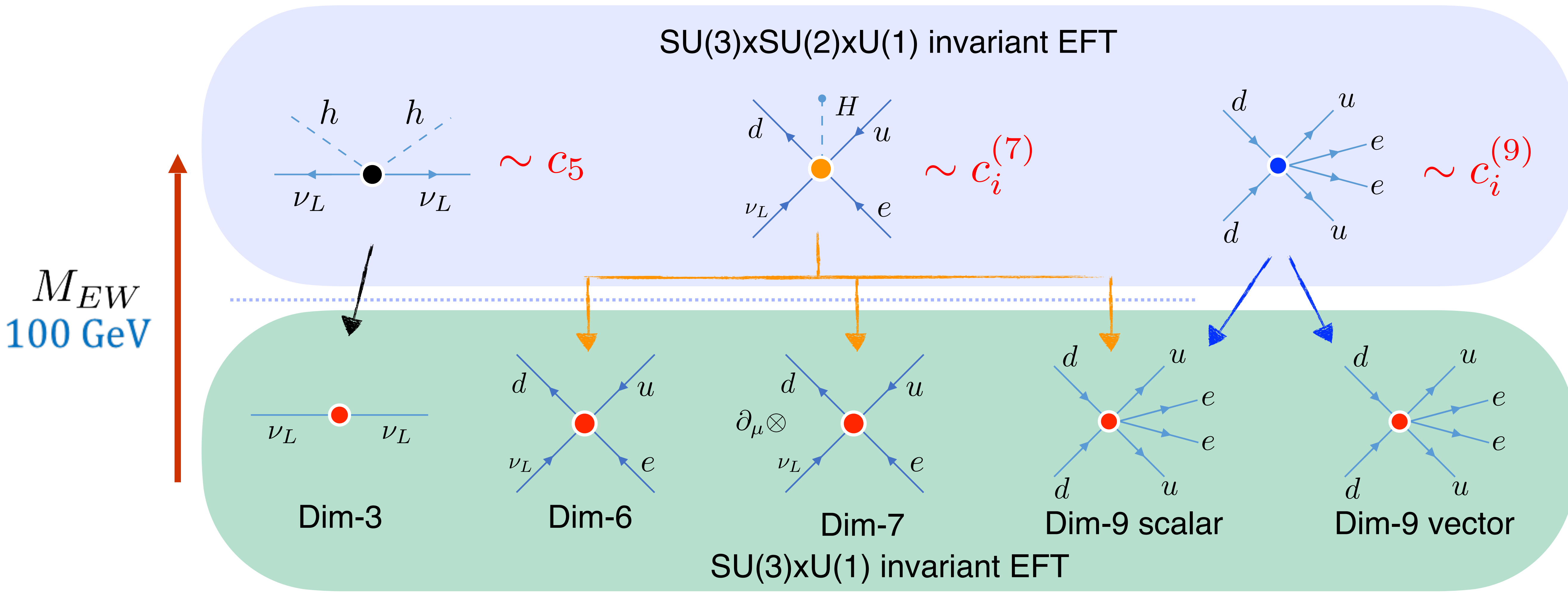
Summary



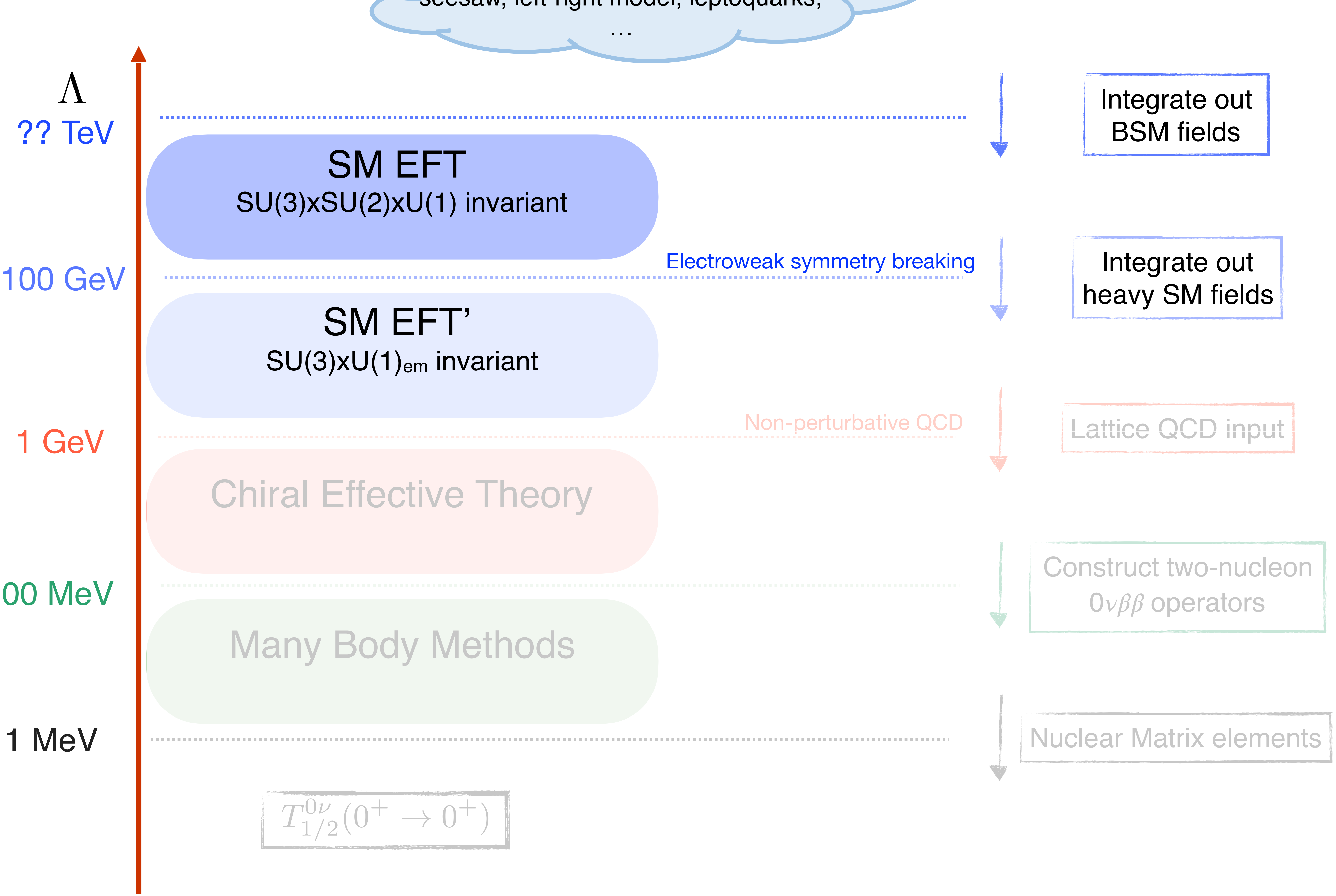
Low-energy operators
Summary



Low-energy operators
Summary

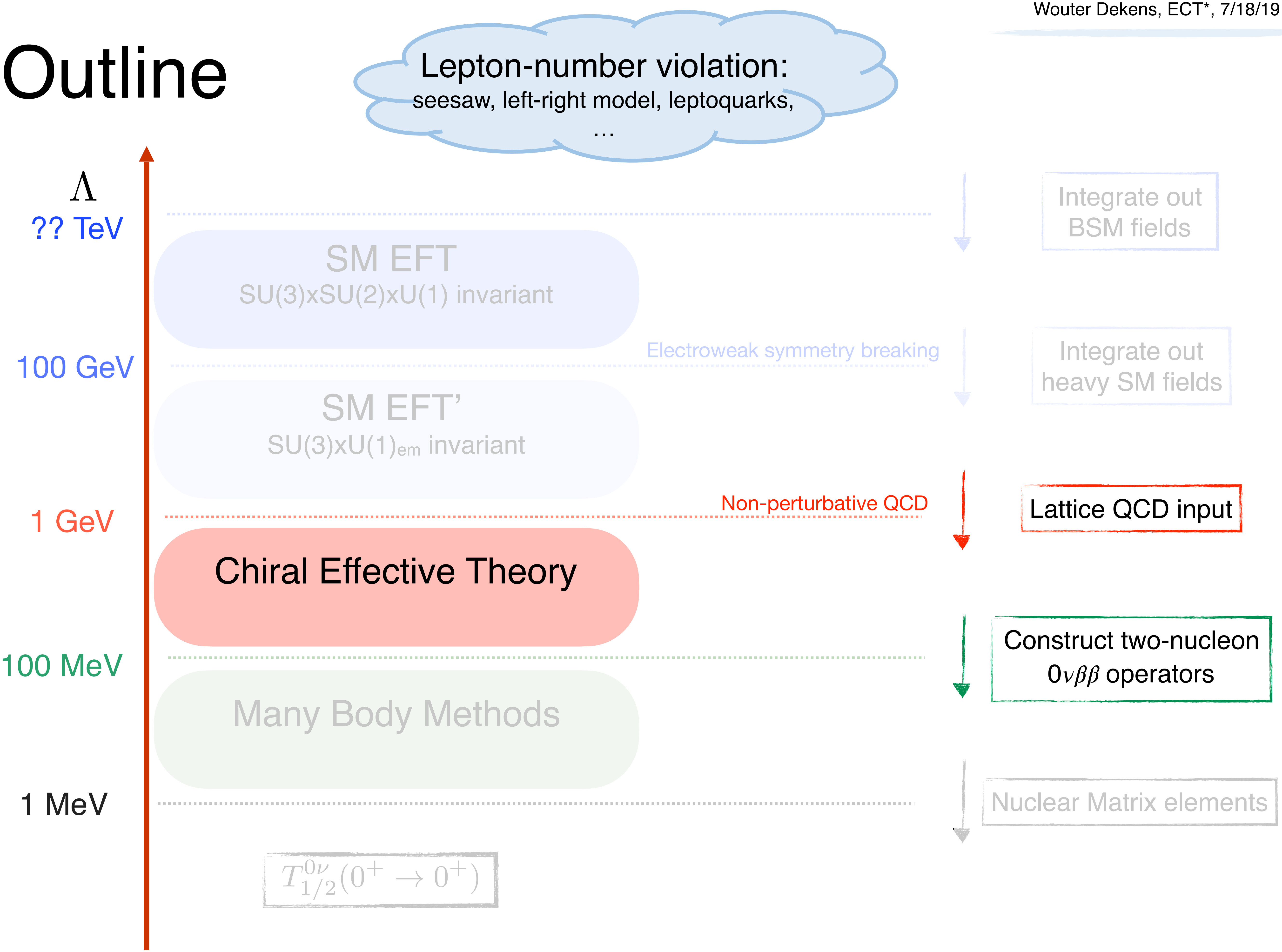


Outline



Lepton-number violation:
seesaw, left-right model, leptoquarks, …
Λ
?? TeV
100 GeV
1 GeV
100 MeV
1 MeV
SM EFT
SU(3)xSU(2)xU(1) invariant
Electroweak symmetry breaking
SM EFT’
SU(3)xU(1)$_{em}$ invariant
Non-perturbative QCD
Chiral Effective Theory
Many Body Methods
$T^{0\nu}_{1/2}(0^+ \to 0^+)$
Integrate out BSM fields
Integrate out heavy SM fields
Lattice QCD input
Construct two-nucleon $0\nu\beta\beta$ operators
Nuclear Matrix elements

Outline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Lattice QCD input
Construct two-nucleon $0\nu\beta\beta$ operators
Nuclear Matrix elements

Matching to Chiral EFT

M_{QCD}
1 GeV

SM EFT'
SU(3)xU(1)$_{em}$ invariant

In terms of quarks, gluons, leptons $\mathcal{L}(u, d, g, \nu, e)$

Chiral Effective Theory

In terms of nucleons, pions, leptons $\mathcal{L}_{\chi PT}(N, \pi, \nu, e)$

Form of operators determined by chiral symmetry

The operators come with unknown constants (LECs)

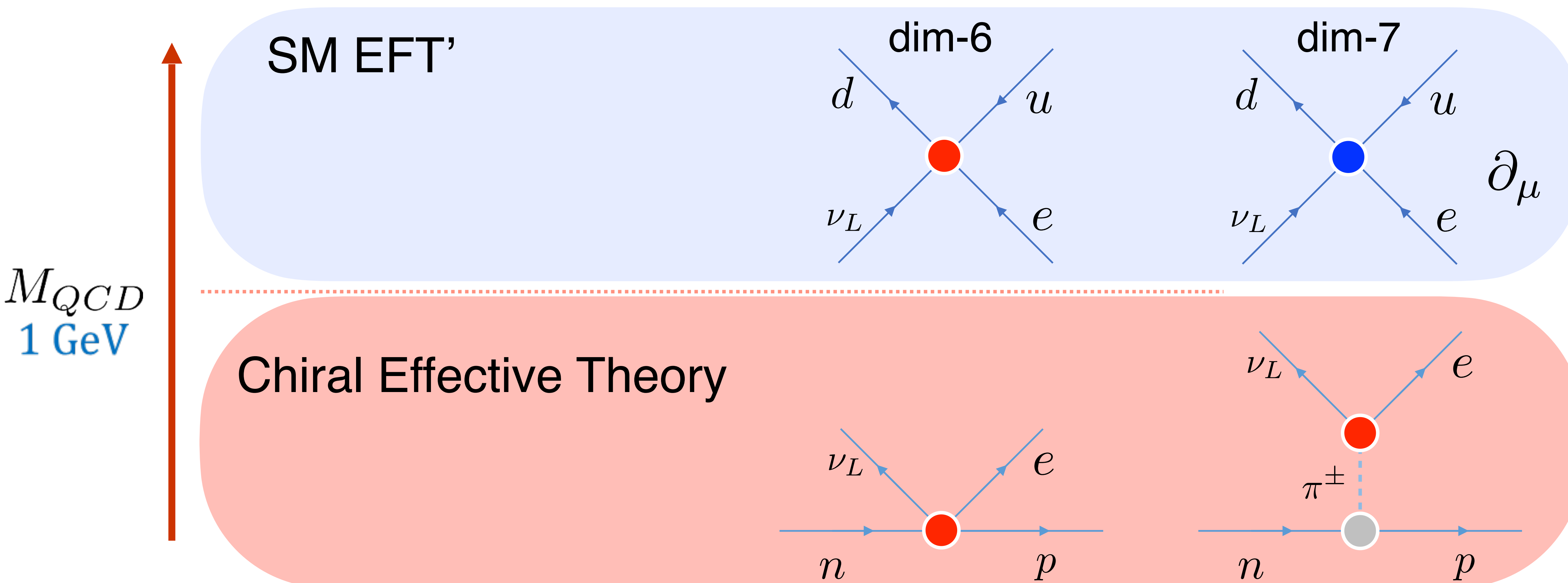
Need a power-counting scheme
- Start by assuming Naive dimensional analysis (NDA)
- Will come back to whether it breaks down

Manohar, Georgi, `84; Weinberg, `90, `91

Matching to Chiral EFT

Dimension-6 and -7: vector & scalar

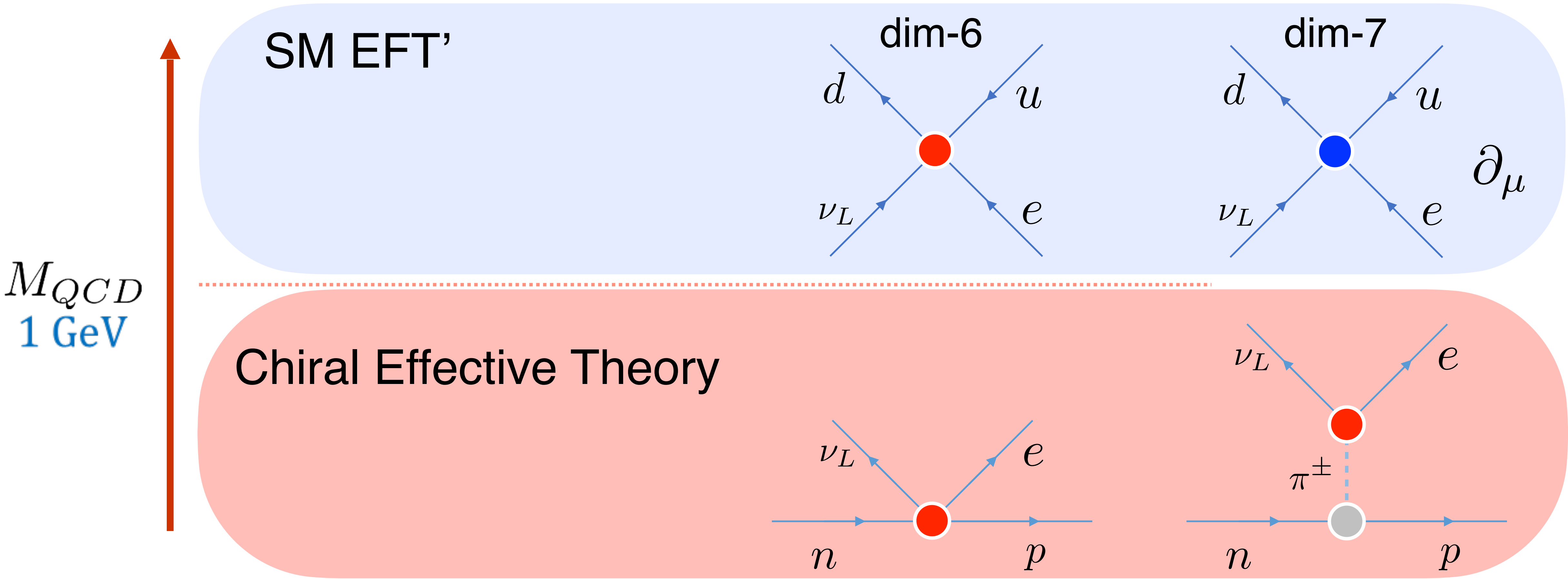
Warning: Based on NDA



Matching to Chiral EFT

Dimension-6 and -7: vector & scalar

Warning: Based on NDA

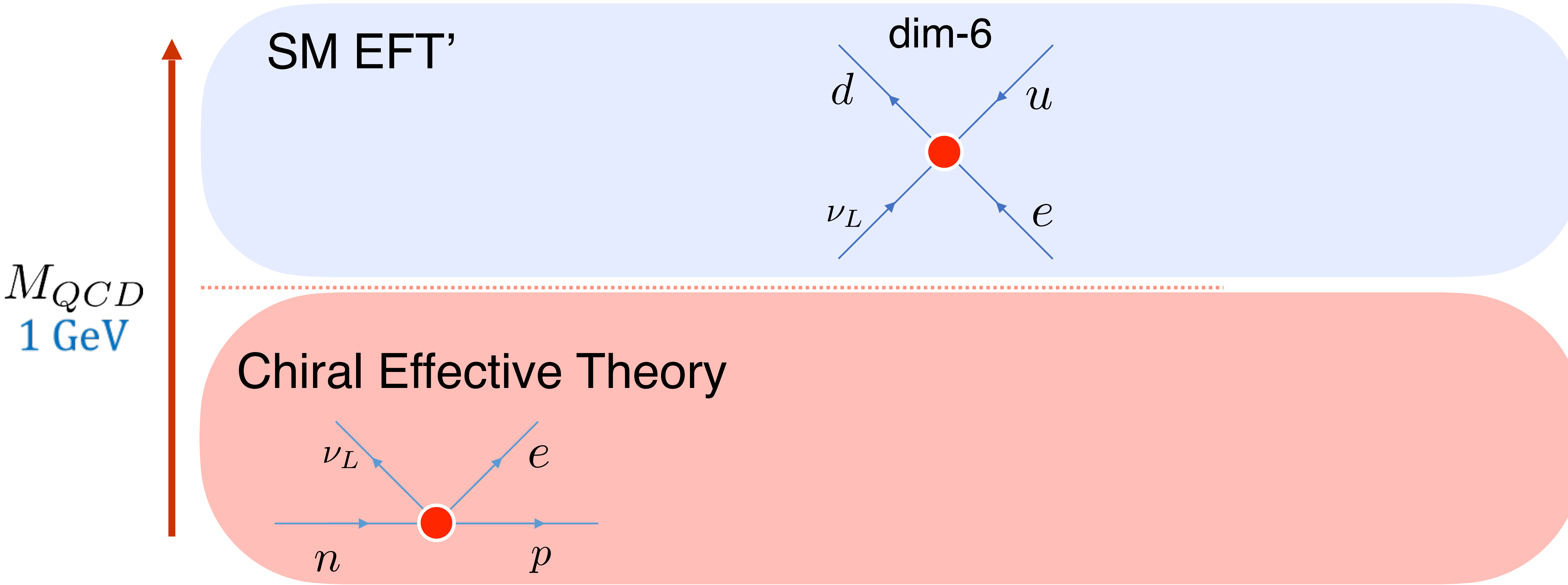


- Needed low-energy constants are the (scalar, vector) nucleon charges
  - gV, gA, gS, gM
  - Known from experiment and/or Lattice QCD

Matching to Chiral EFT

Dimension-6: tensor, left-handed vector

Warning: Based on NDA

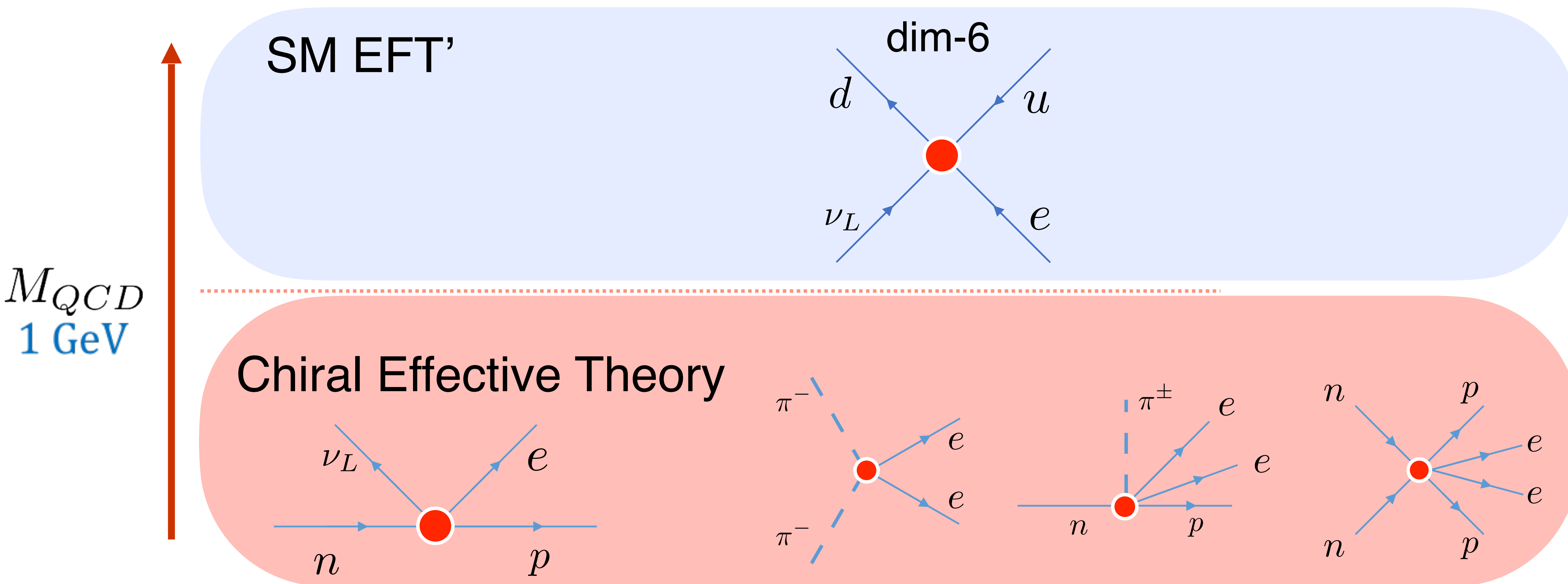


- Generate beta-decay like operators

Matching to Chiral EFT

Dimension-6: tensor, left-handed vector

Warning: Based on NDA

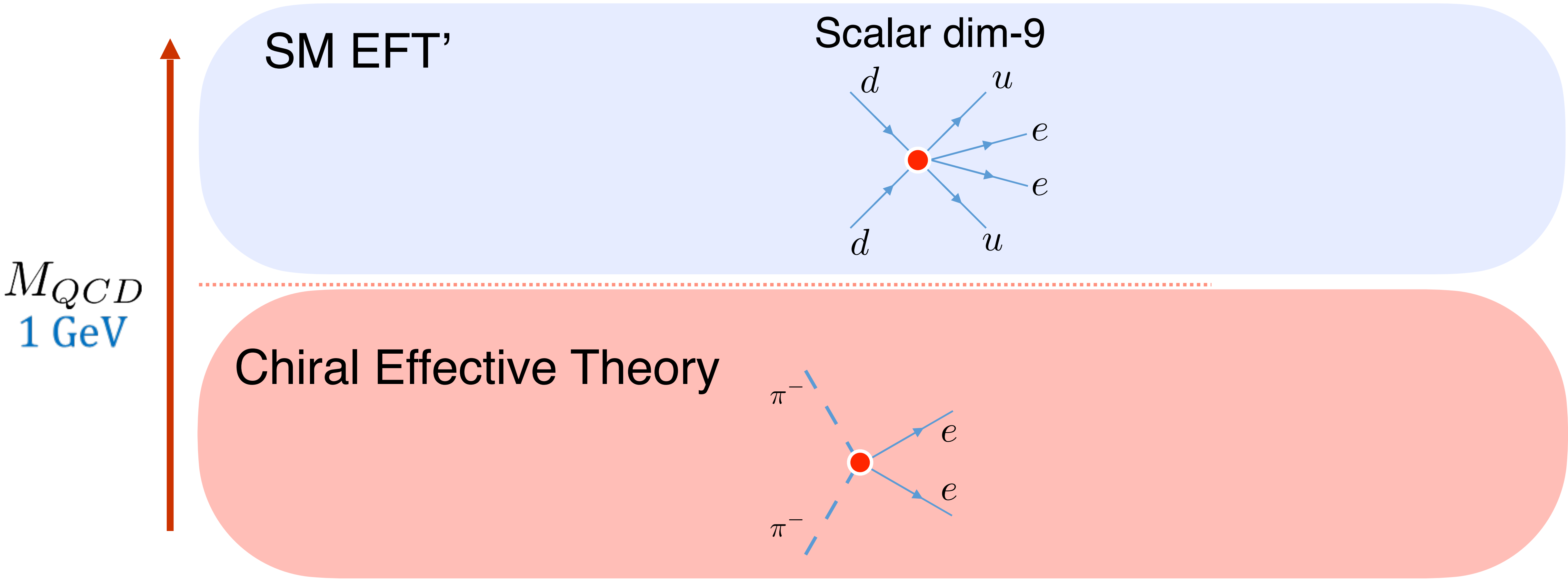


- Generate beta-decay like operators
- Also induce $\pi\pi$, πN, and NN interactions
  - Come with unknown LECs

Matching to Chiral EFT
Dimension-9

Warning: Based on NDA

$$\mathcal{L}^{(9)}_{\Delta L=2} = \frac{1}{v^5}\sum_i \left[\left(C^{(9)}_{i\,\mathrm{R}}\, \bar{e}_R C \bar{e}^T_R + C^{(9)}_{i\,\mathrm{L}}\, \bar{e}_L C \bar{e}^T_L \right) O_i + C^{(9)}_i \bar{e}\gamma_\mu\gamma_5 C \bar{e}^T\, O^\mu_i \right],$$



Matching to Chiral EFT

Dimension-9

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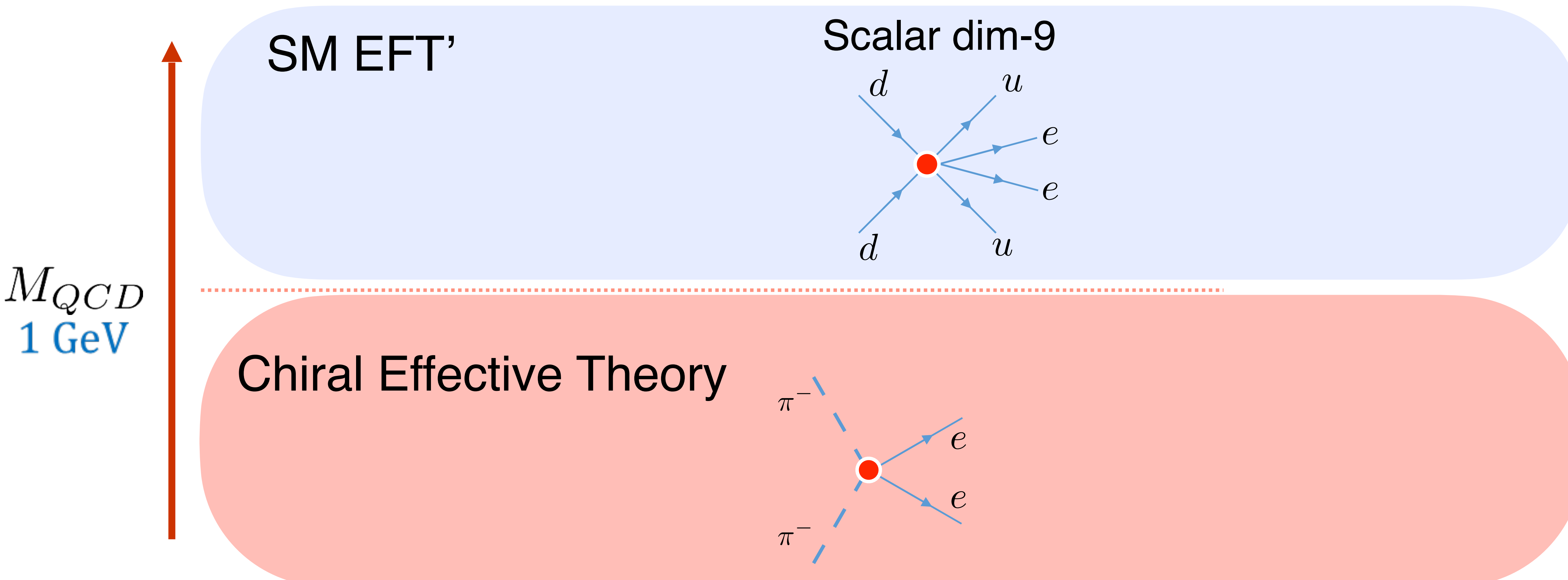

- Most scalar operators only induce $\pi\pi$ interactions
  - Known from Lattice QCD / SU(3) chiral symmetry

Nicholson et al.’18; Cirigliano et al. ‘17

Matching to Chiral EFT

Dimension-9

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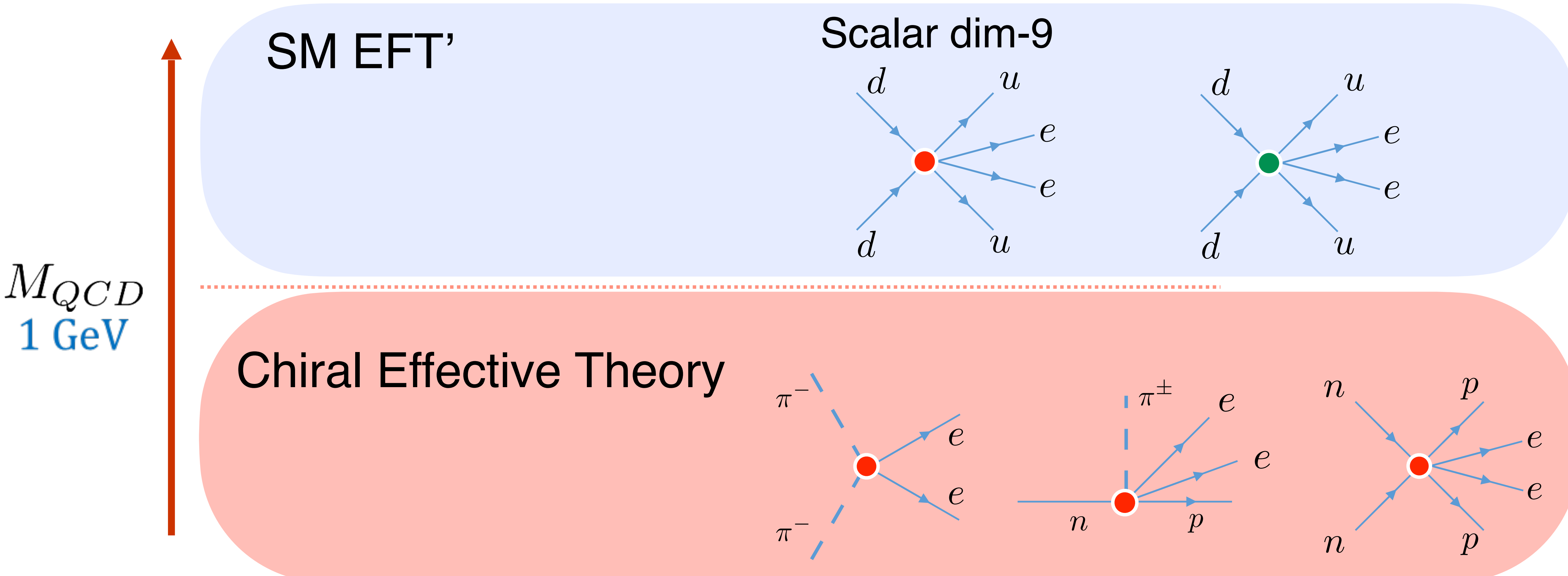
$$\mathcal{L}^{(9)}_{\Delta L=2} = \frac{1}{v^5}\sum_i \left[\left(C^{(9)}_{i\,\mathrm{R}}\,\bar{e}_R C \bar{e}_R^T + C^{(9)}_{i\,\mathrm{L}}\,\bar{e}_L C \bar{e}_L^T\right) O_i + C_i^{(9)}\bar{e}\gamma_\mu\gamma_5 C \bar{e}^T O_i^\mu\right],$$



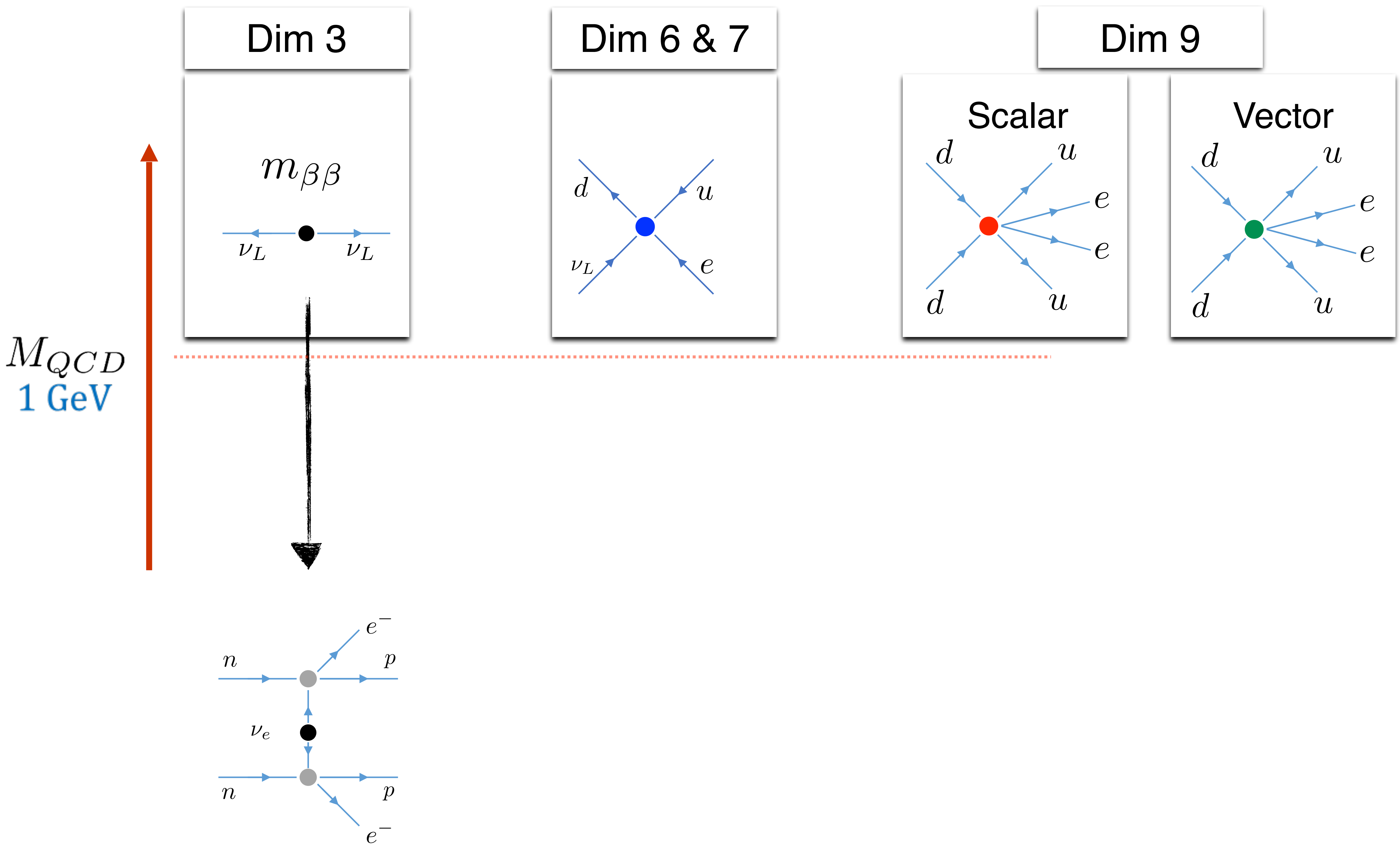
- Most scalar operators only induce $\pi\pi$ interactions
  - Known from Lattice QCD / SU(3) chiral symmetry

- One scalar structure + vector operators induce πN & NN terms
  - The low-energy constants for the πN and NN interactions are unknown

Nicholson et al.'18; Cirigliano et al. '17

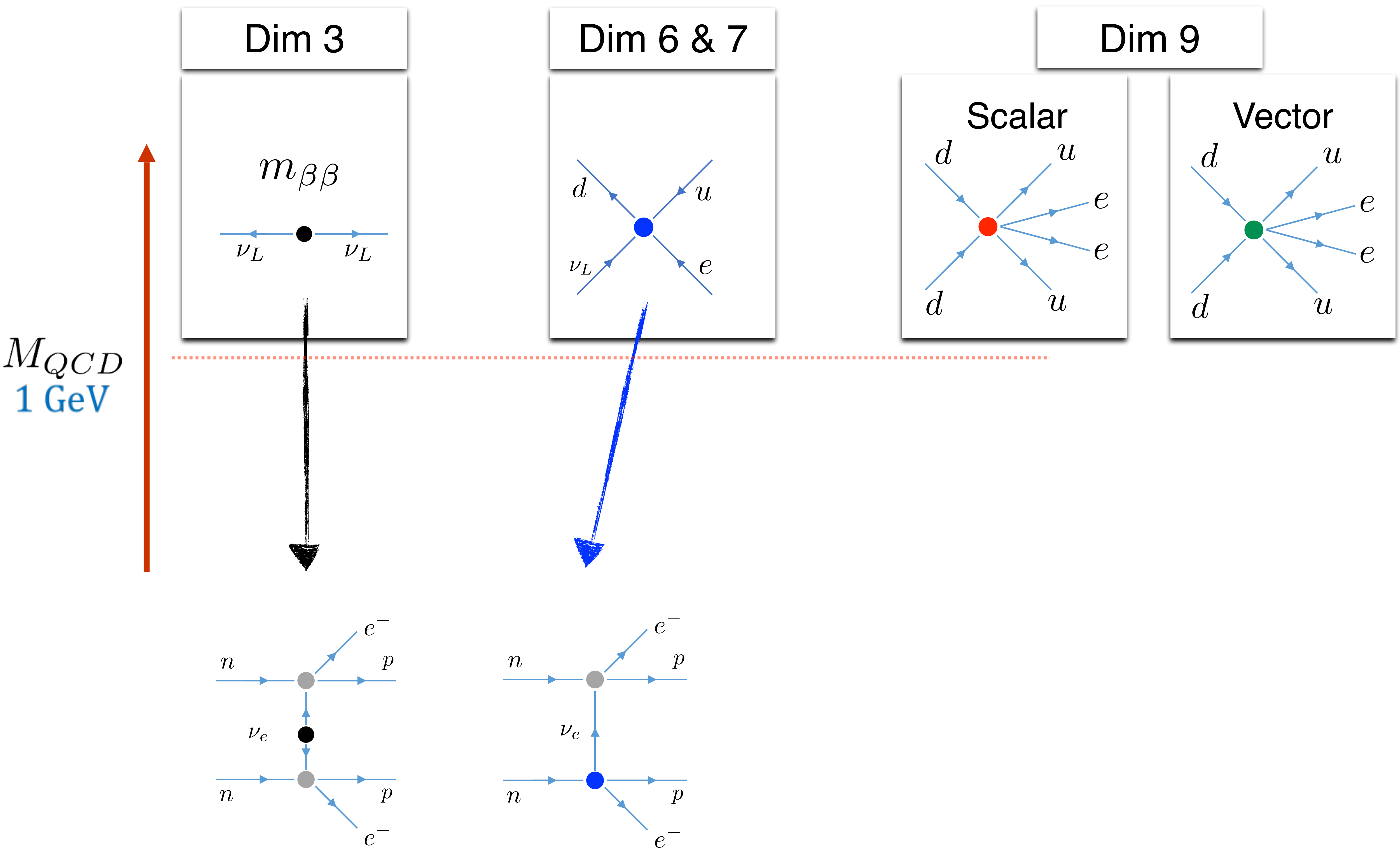
Chiral EFT

Summary



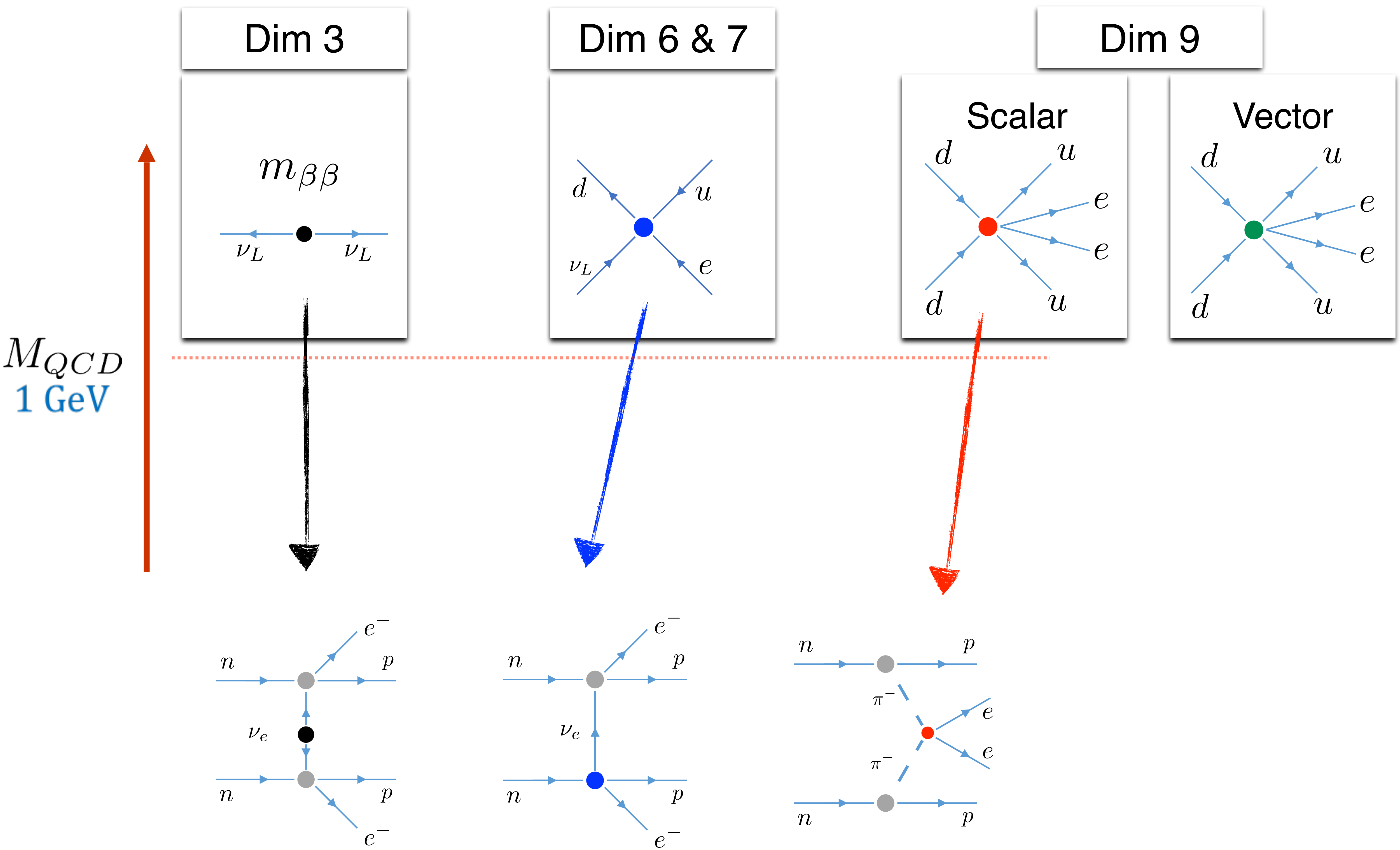
Chiral EFT

Summary



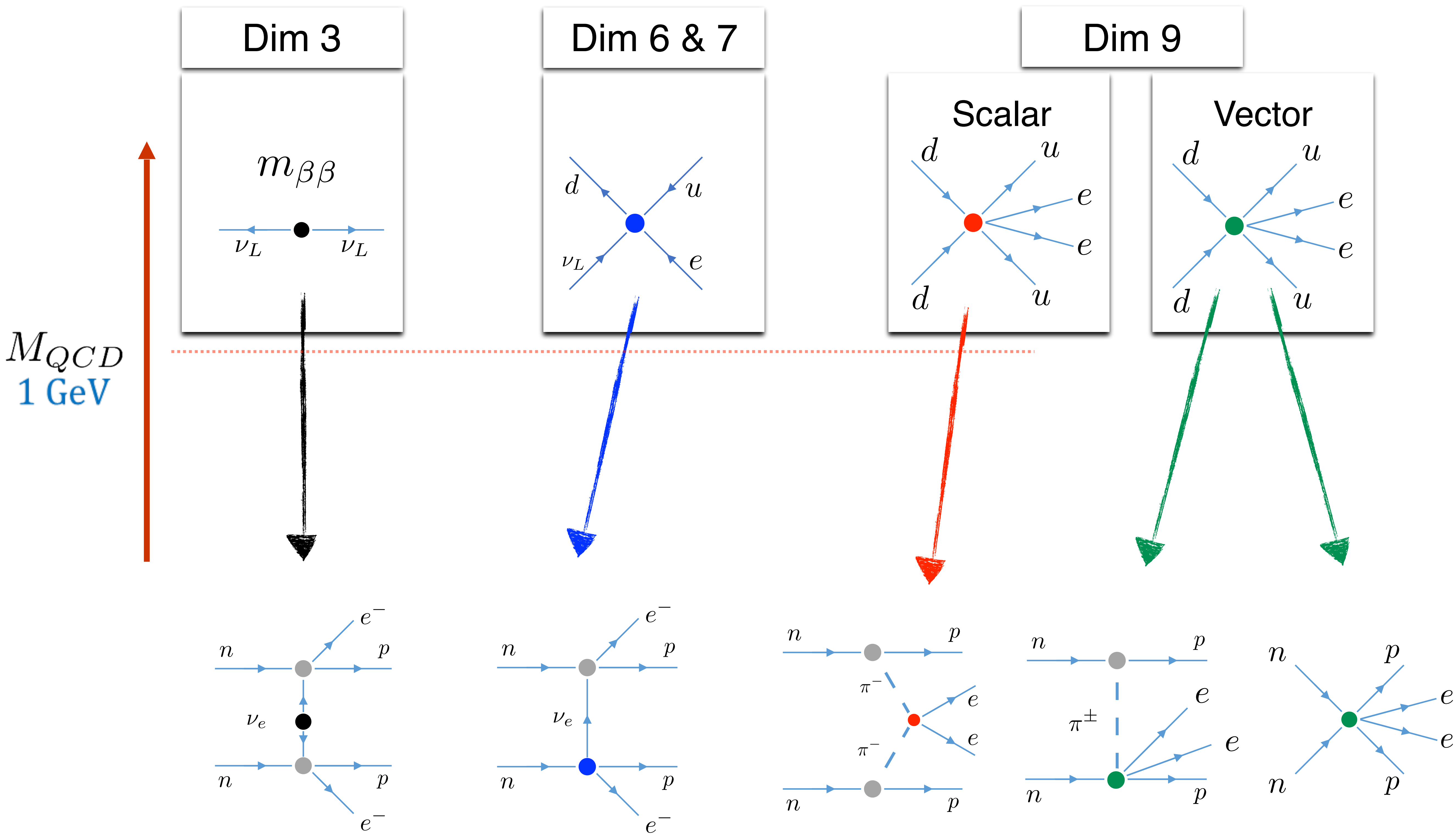
Chiral EFT

Summary

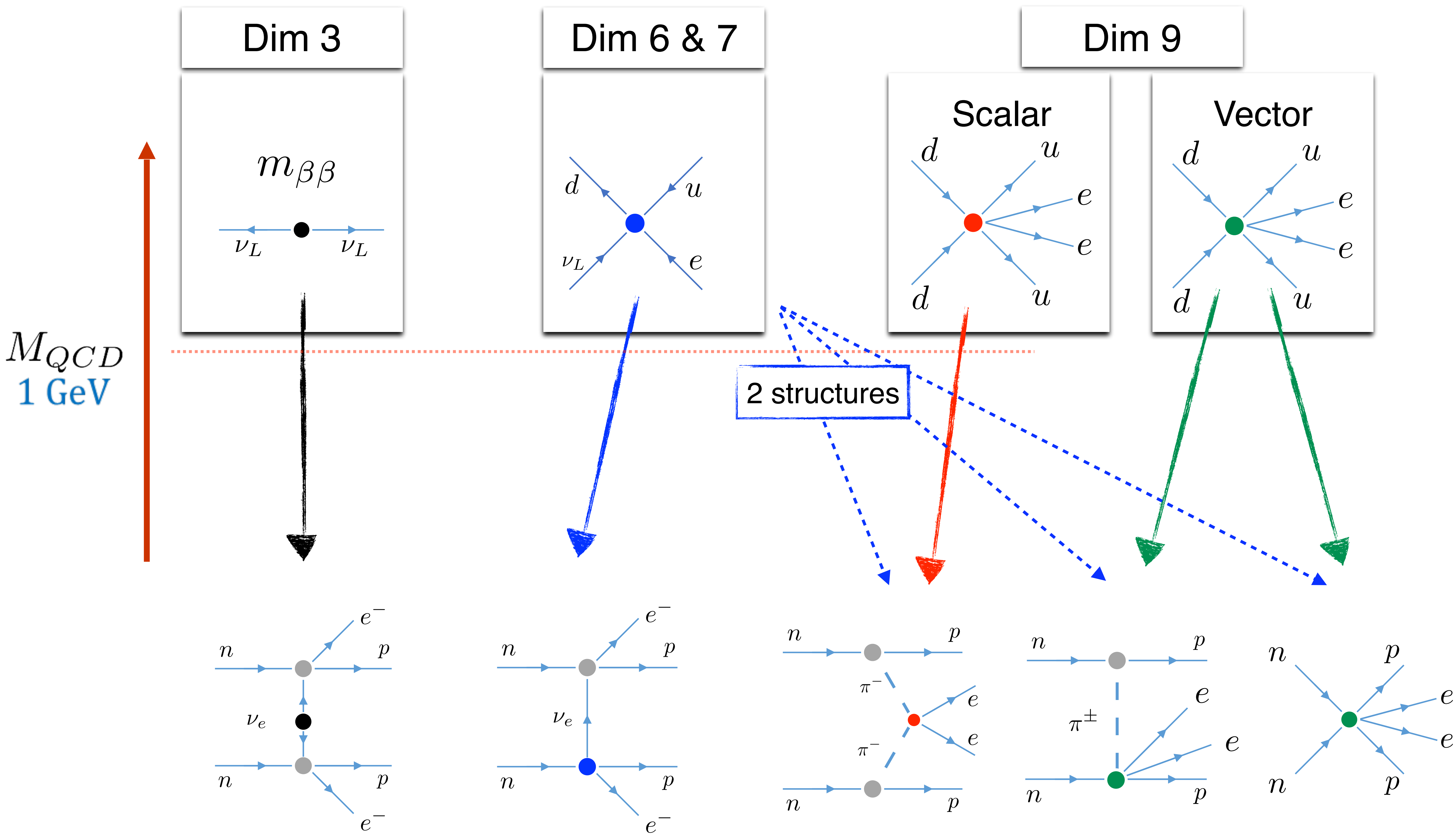


Chiral EFT

Summary



Chiral EFT
Summary



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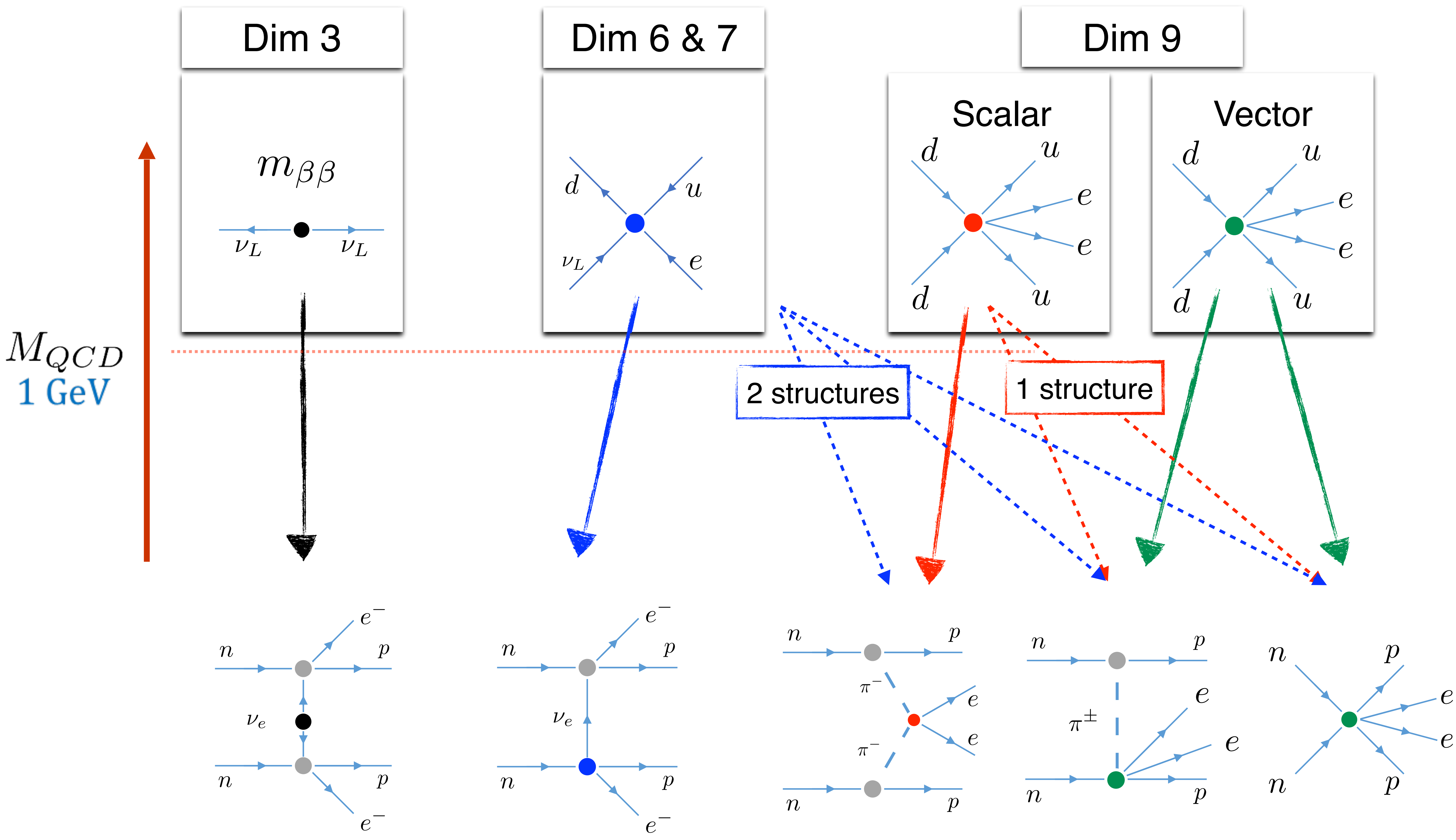
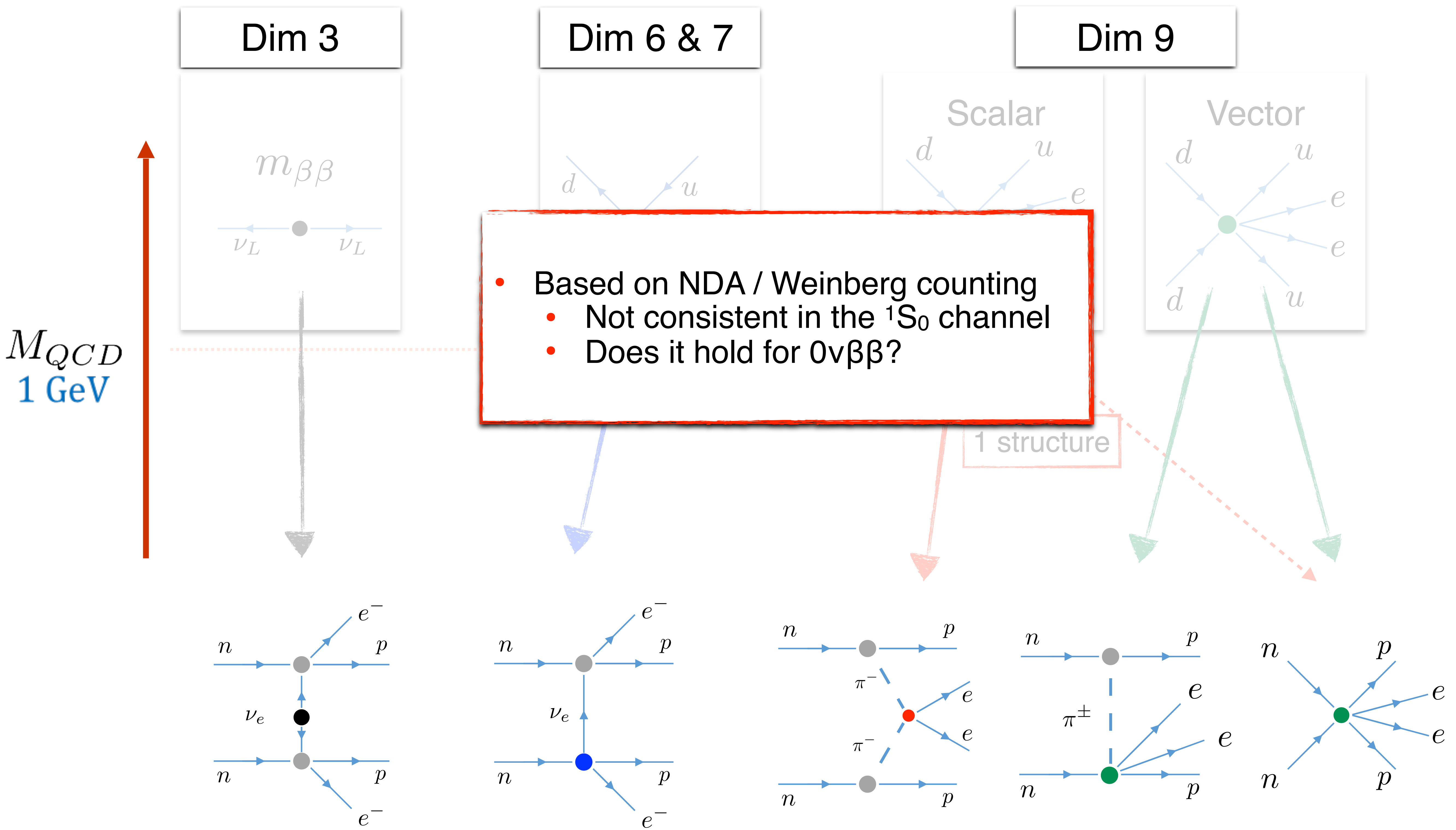


Chiral EFT
Summary



- Based on NDA / Weinberg counting
  - Not consistent in the $^1S_0$ channel
  - Does it hold for 0νββ?

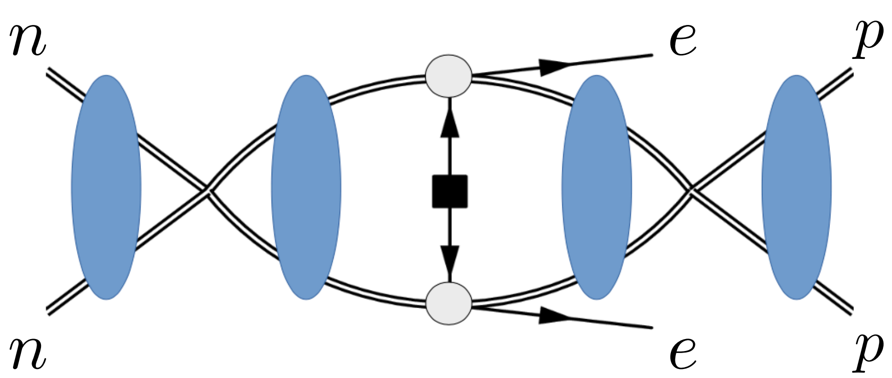
Kaplan, Savage, Wise, '96; Beane, Bedaque, Savage, van Kolck, '03, Nogga, Timmermans, van Kolck, '05, Long, Yang, '12;

Checking the Weinberg counting

Any effect for the dim-6,7,9 terms?

- In the Majorana-mass case, the LNV potential leads to a divergence
  - Due to the potential at large momenta $V_{\Delta L=2} \sim 1/\vec{q}^2$

See Jordy's talk



$$\sim m_N^2 \int d^3q d^3k \frac{1}{m_N E - \vec{q}^2} \frac{1}{(\vec{q}-\vec{k})^2} \frac{1}{m_N E' - \vec{k}^2}$$

Checking the Weinberg counting

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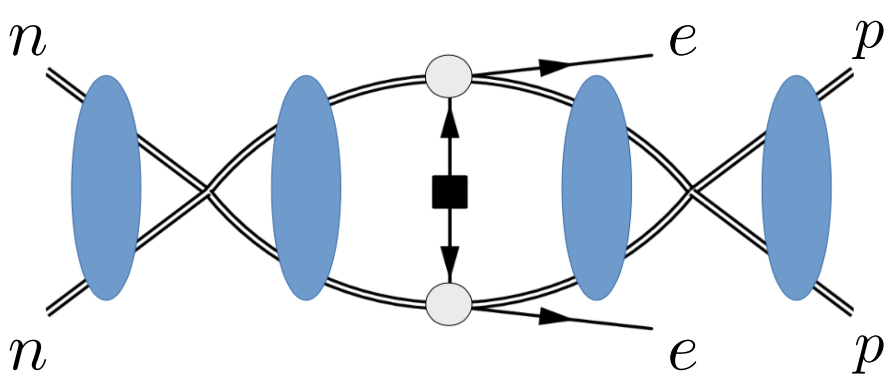

$V_{\Delta L=2}$

$$\sim m_N^2 \int d^3q d^3k \frac{1}{m_N E - \vec{q}^2} \frac{1}{(\vec{q}-\vec{k})^2} \frac{1}{m_N E' - \vec{k}^2}$$

Checking the Weinberg counting

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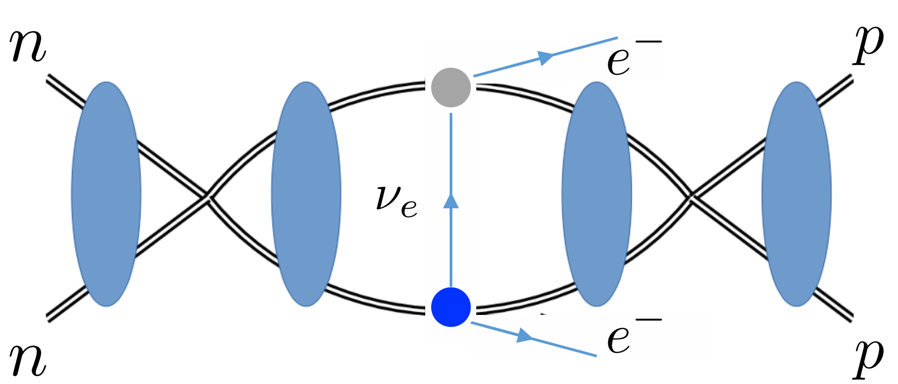

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Dimension-6,7,9

- Several potentials have the same behavior
  - The case for the vector operators

$$C^{(6)}_{VL,VR}: \quad V_{\Delta L=2} \sim 1/\vec{q}^2$$

$$C^{(9)}_{1-9}: \quad V_{\Delta L=2} \sim \frac{1}{\vec{q}^2 + m_\pi^2}$$

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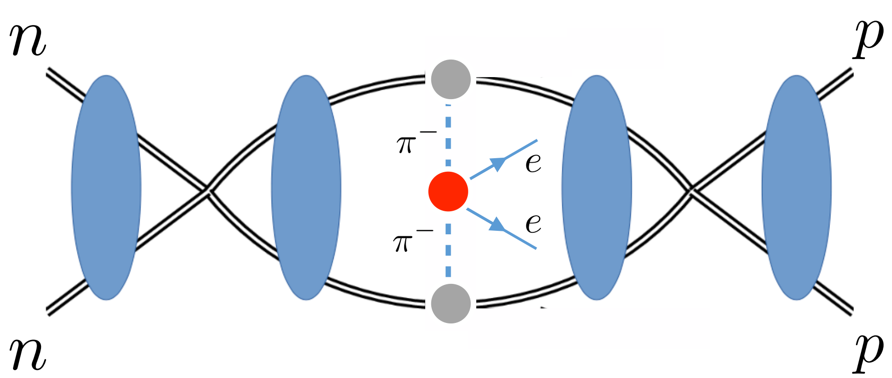
Any effect for the dim-6,7,9 terms?

- In the Majorana-mass case, the LNV potential leads to a divergence
  - Due to the potential at large momenta $V_{\Delta L=2} \sim 1/\vec{q}^2$

See Jordy's talk



$$\sim m_N^2 \int d^3q d^3k \frac{1}{m_N E - \vec{q}^2} \frac{1}{(\vec{q}-\vec{k})^2} \frac{1}{m_N E' - \vec{k}^2}$$

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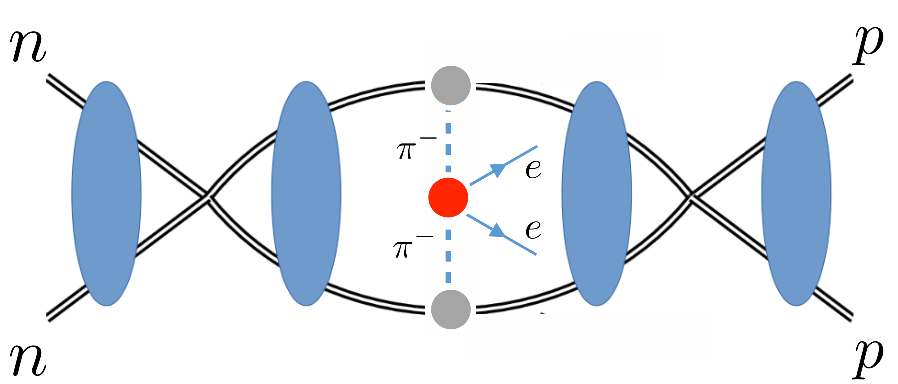

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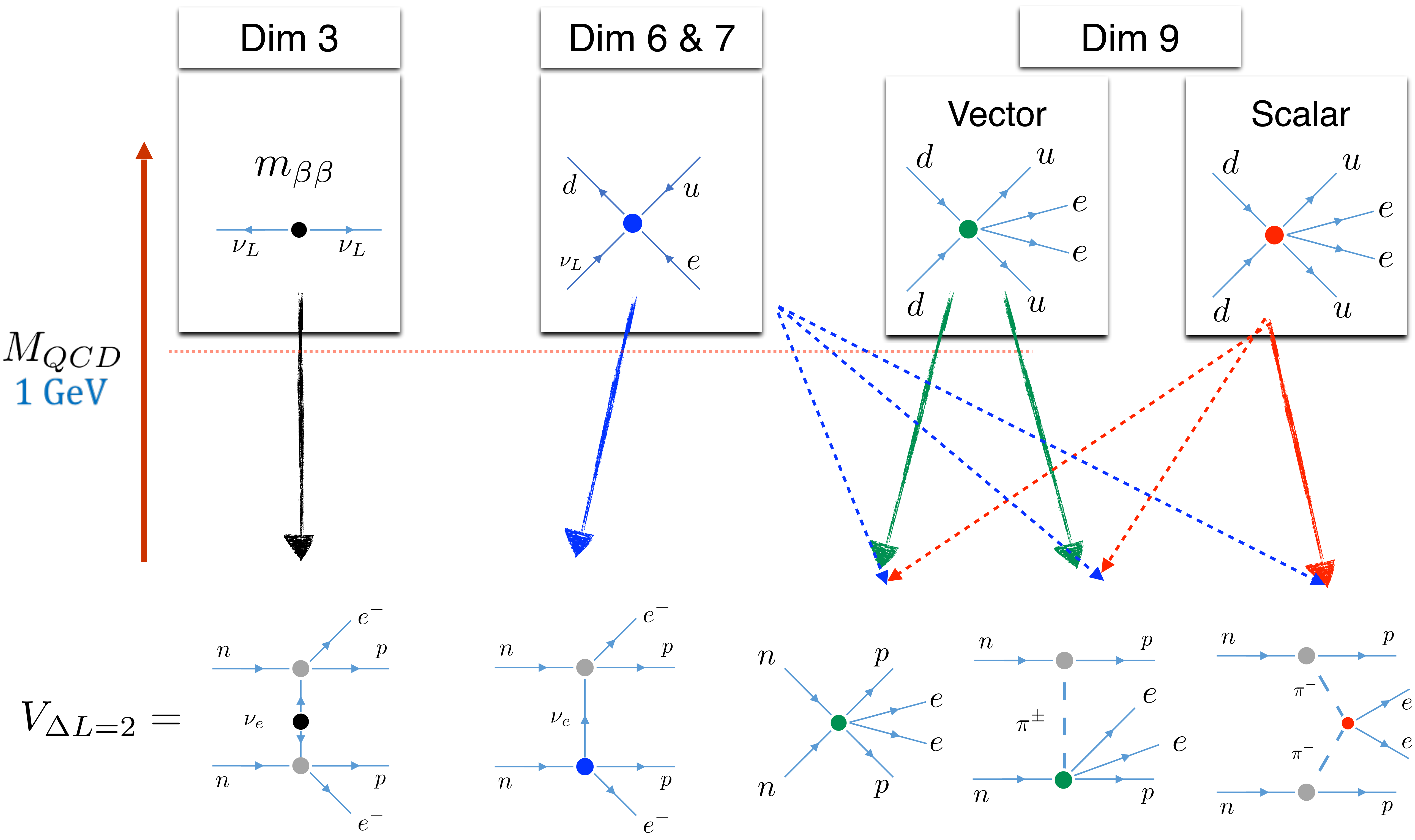
- Need to include contact interactions at LO in these cases
  - Often disagrees with the Weinberg / NDA counting

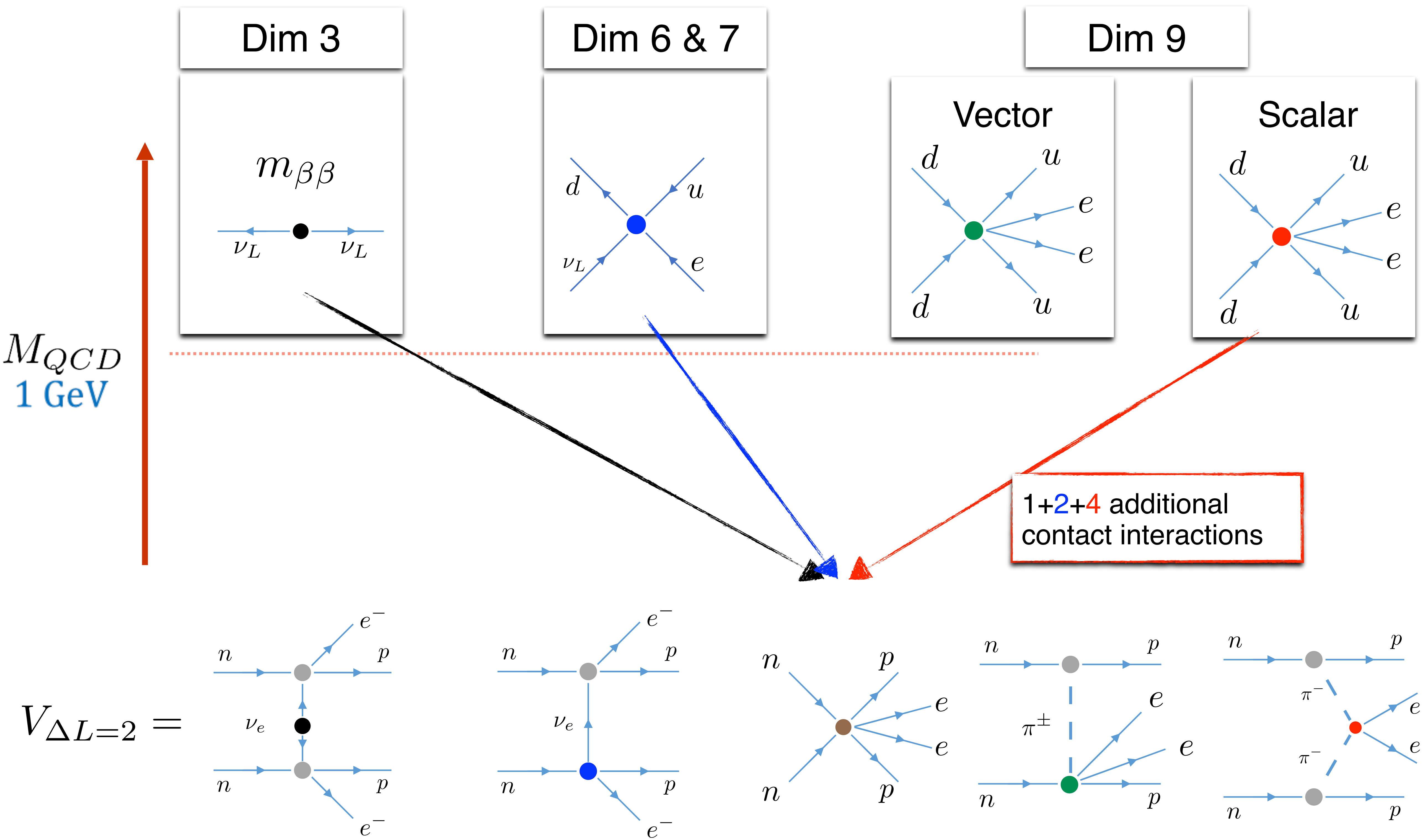
Chiral EFT

NDA / Weinberg



Chiral EFT

Beyond NDA / Weinberg



Chiral EFT
The potential

Can finally derive the potential and amplitude

$$\mathcal{A} = \langle 0^+ | \sum_{m,n} \int \frac{d^3\mathbf{q}}{(2\pi)^3} e^{i\mathbf{q}\cdot\mathbf{r}} V(\mathbf{q}^2) | 0^+ \rangle \qquad V(\mathbf{q}^2) = V_3(\mathbf{q}^2) + V_6(\mathbf{q}^2) + V_7(\mathbf{q}^2) + V_9(\mathbf{q}^2)$$

The contributions scale as

	$d{=}3$	$C^{(6)}_{\rm SL,\,SR}$	$C^{(6)}_{\rm T}$	$C^{(6)}_{\rm VL}$	$C^{(6)}_{\rm VR}$	$C^{(7)}_{\rm VL,\,VR}$	$C^{(9)\,(\prime)}_{\rm 1R}$	$C^{(9)(\prime)}_{\rm 1L}$	$C^{(9)(\prime)}_{\rm 2R-5R}$	$C^{(9)(\prime)}_{\rm 2L-5L}$	$C^{(9)}_{\rm vector}$
$m_e\mathcal{A}$	$m_{\beta\beta}$	Λ_χ	$\Lambda_\chi\epsilon_\chi^2$	$\Lambda_\chi\epsilon_\chi^2$	$\Lambda_\chi\epsilon_\chi^3$	$\frac{\Lambda_\chi^2}{v}\,\epsilon_\chi^2$	$\frac{\Lambda_\chi^2}{v}\,\epsilon_\chi^2$	$\frac{\Lambda_\chi^2}{v}\,\epsilon_\chi^2$	$\frac{\Lambda_\chi^2}{v}$	$\frac{\Lambda_\chi^2}{v}$	$\frac{\Lambda_\chi^2}{v}\,\epsilon_\chi^2$

$$\Lambda_\chi = 1\,\mathrm{GeV}, \qquad \epsilon_\chi = m_\pi/\Lambda_\chi$$

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- Several operators are suppressed by two or three powers of $\epsilon_\chi = m_\pi/\Lambda_\chi$

- Scaling of Wilson coefficients needed to see which are important
  - To be determined in explicit models of new physics

Outline

Lepton-number violation:
seesaw, left-right model, leptoquarks, …

Λ

?? TeV

SM EFT
SU(3)xSU(2)xU(1) invariant

Integrate out BSM fields

Electroweak symmetry breaking

100 GeV

SM EFT'
SU(3)xU(1)$_{em}$ invariant

Integrate out heavy SM fields

Non-perturbative QCD

1 GeV

Lattice QCD input

Chiral Effective Theory

Construct two-nucleon $0\nu\beta\beta$ operators

100 MeV

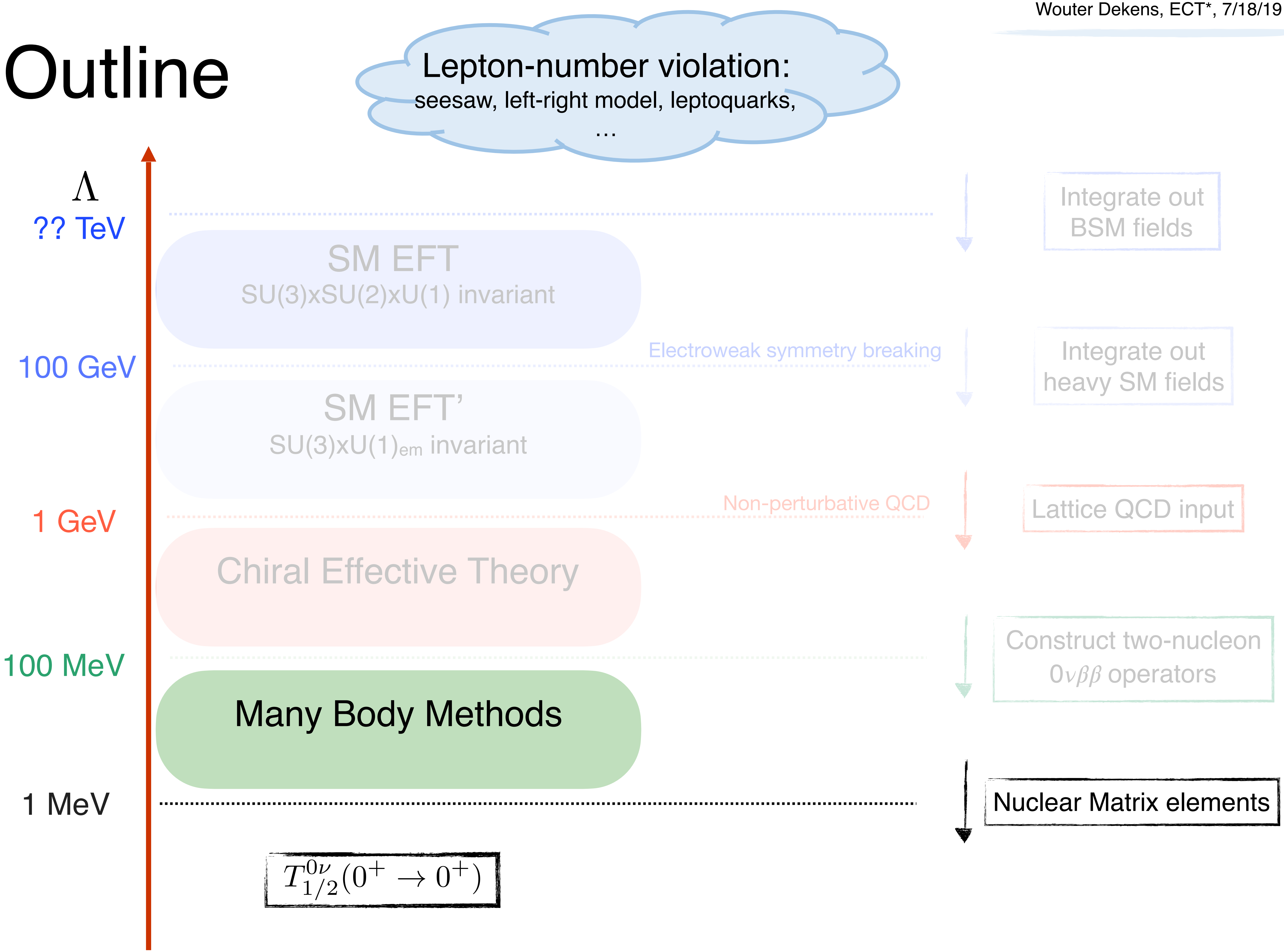
Many Body Methods

1 MeV

Nuclear Matrix elements

$T^{0\nu}_{1/2}(0^+ \to 0^+)$

Outline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The $0\nu\beta\beta$ half-life

$$\Gamma^{0\nu}(0^+ \to 0^+) \sim \left| \langle 0^+ | \sum_{\text{nucleons}} \int \frac{d^3\boldsymbol{q}}{(2\pi)^3} e^{i\boldsymbol{q}\cdot\boldsymbol{r}}\, V(\boldsymbol{q}^2)\, |0^+\rangle \right|^2 = \sum_{i,j} G_{i,j} M_i M_j g_i g_j C_i C_j^*$$

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Nuclear matrix elements

- All NMEs can be obtained from those of light/heavy neutrino exchange
  - 9 long-distance & 6 short-distance
  - Have been determined in literature

- Follow ChiPT expectations fairly well
  - E.g. all *O*(1) and

$$M_{GT,sd}^{PP} = -\frac{1}{2}M_{GT,sd}^{AP} - M_{GT}^{PP}\,, \qquad M_{T,sd}^{PP} = -\frac{1}{2}M_{T,sd}^{AP} - M_{T}^{PP}\,,$$

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NMEs	^{76}Ge			
	[74]	[31]	[81]	[82, 83]
M_F	-1.74	-0.67	-0.59	-0.68
M_{GT}^{AA}	5.48	3.50	3.15	5.06
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M_{GT}^{PP}	0.66	0.33	0.30	
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M_{T}^{AA}	–	–	–	
M_{T}^{AP}	-0.35	0.01	-0.01	
M_{T}^{PP}	0.10	0.00	0.00	
M_{T}^{MM}	-0.04	0.00	0.00	

NMEs	^{76}Ge			
$M_{F,\,sd}$	-3.46	-1.55	-1.46	-1.1
$M_{GT,\,sd}^{AA}$	11.1	4.03	4.87	3.62
$M_{GT,\,sd}^{AP}$	-5.35	-2.37	-2.26	-1.37
$M_{GT,\,sd}^{PP}$	1.99	0.85	0.82	0.42
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Barea et al. '15; Hyvarinen et al, '15; Horoi et al. '17, Menendez et al, '18

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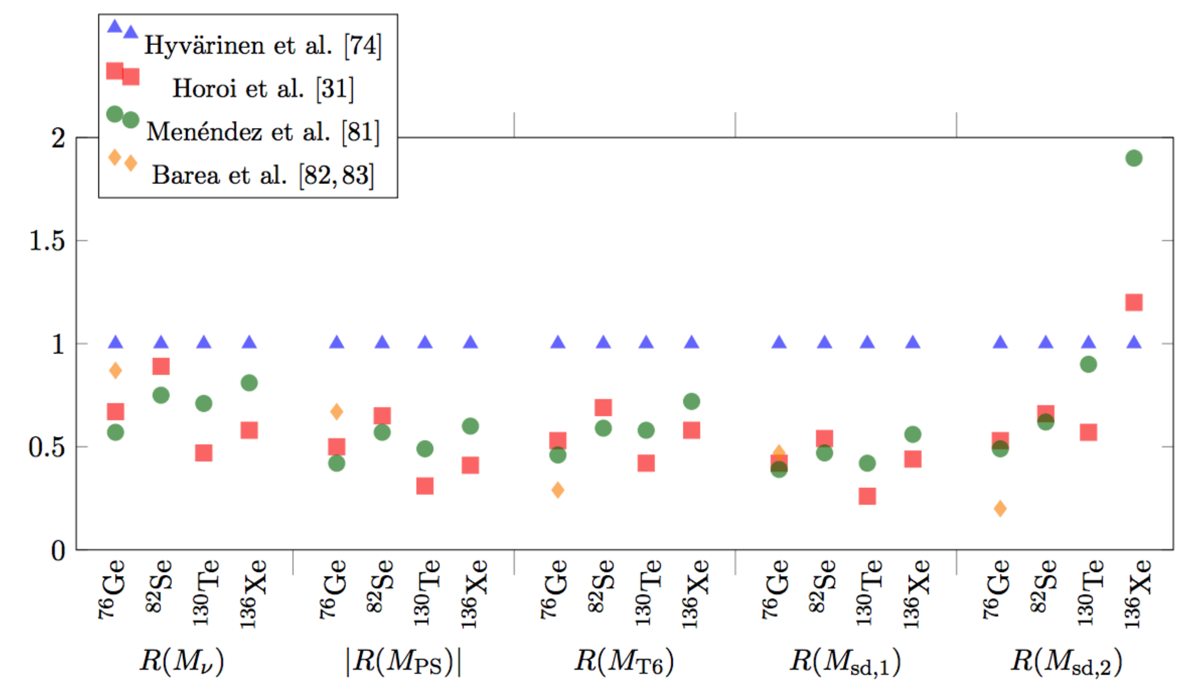
- The NMEs differ by a factor 2-3 between methods
  - For Majorana-mass term & other LNV sources



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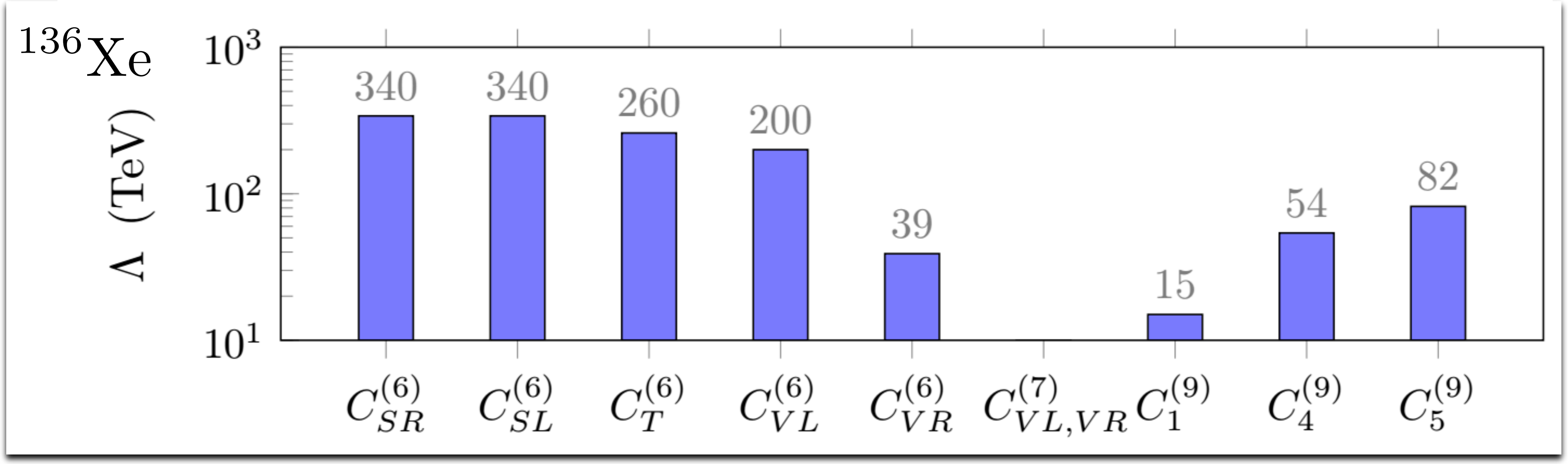
Phenomenology

Current limits

- Assumes $C_i = v^3/\Lambda^3$



- Uncertainties:
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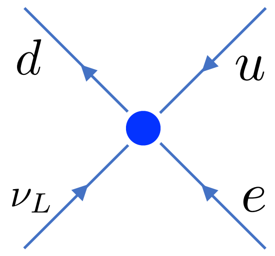
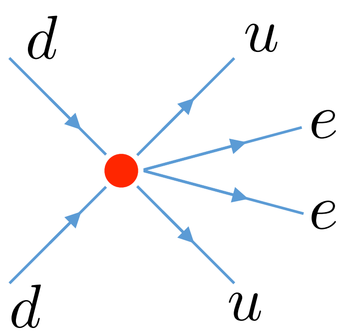
  - Nuclear Matrix elements



Dim 6 & 7



Dim 9

Current limits

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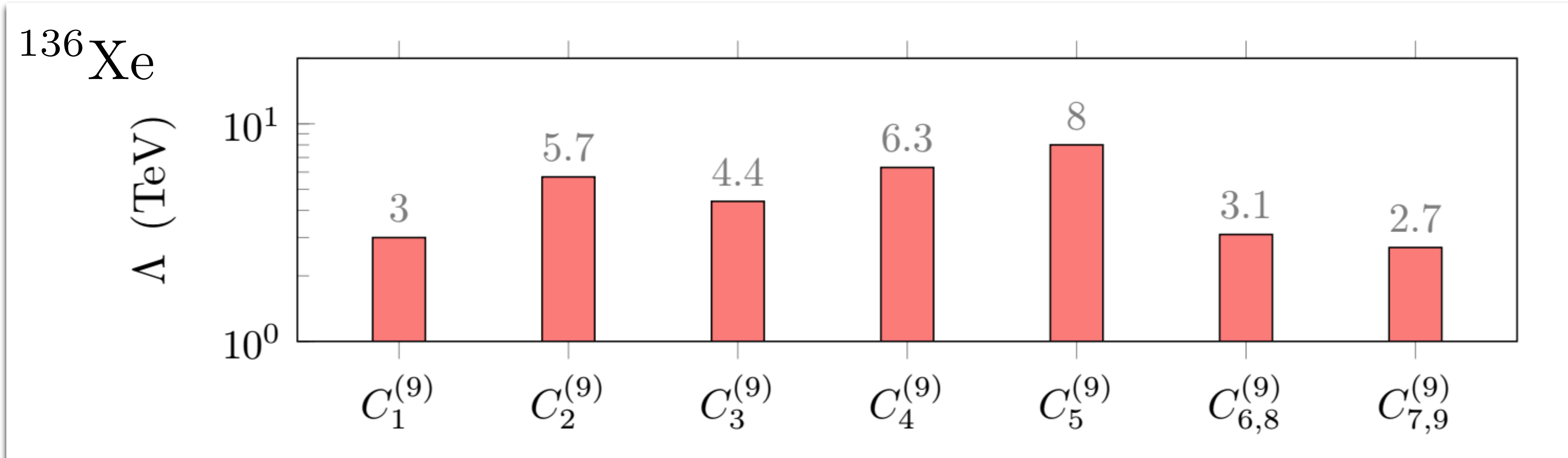

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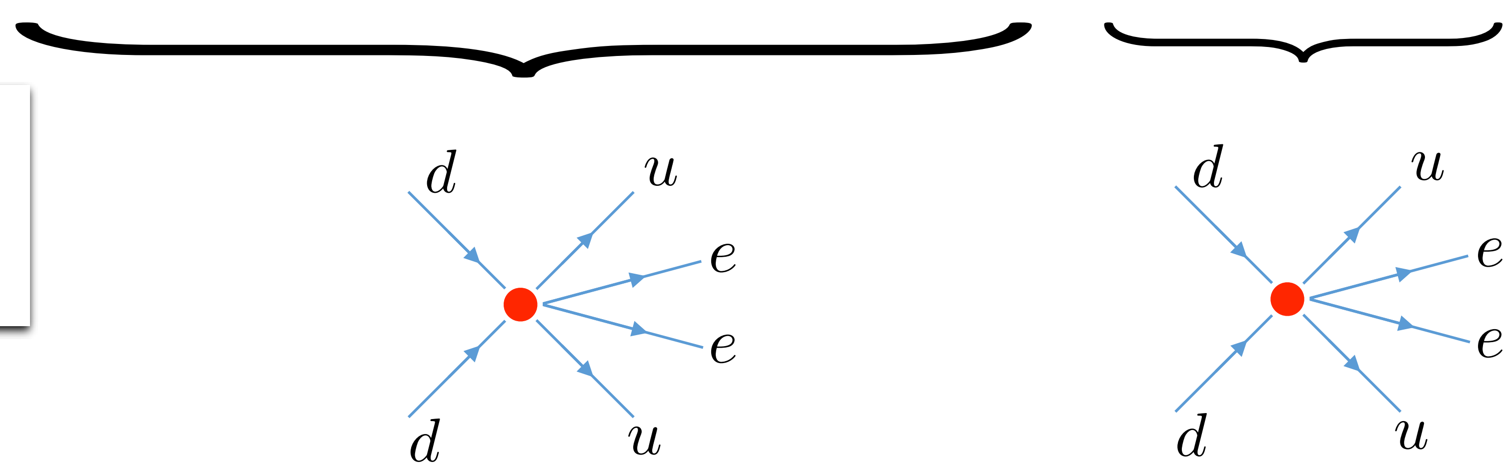
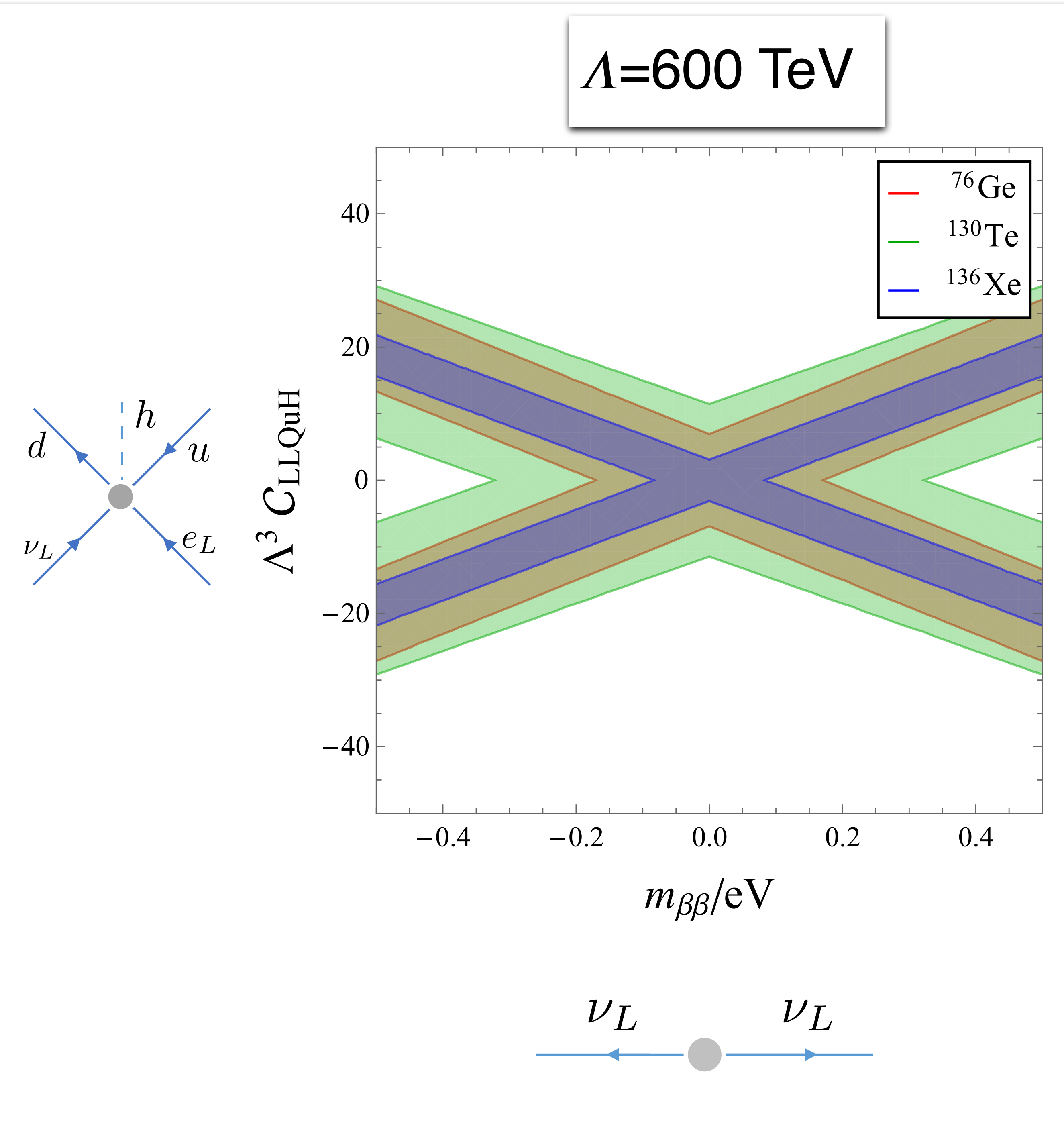


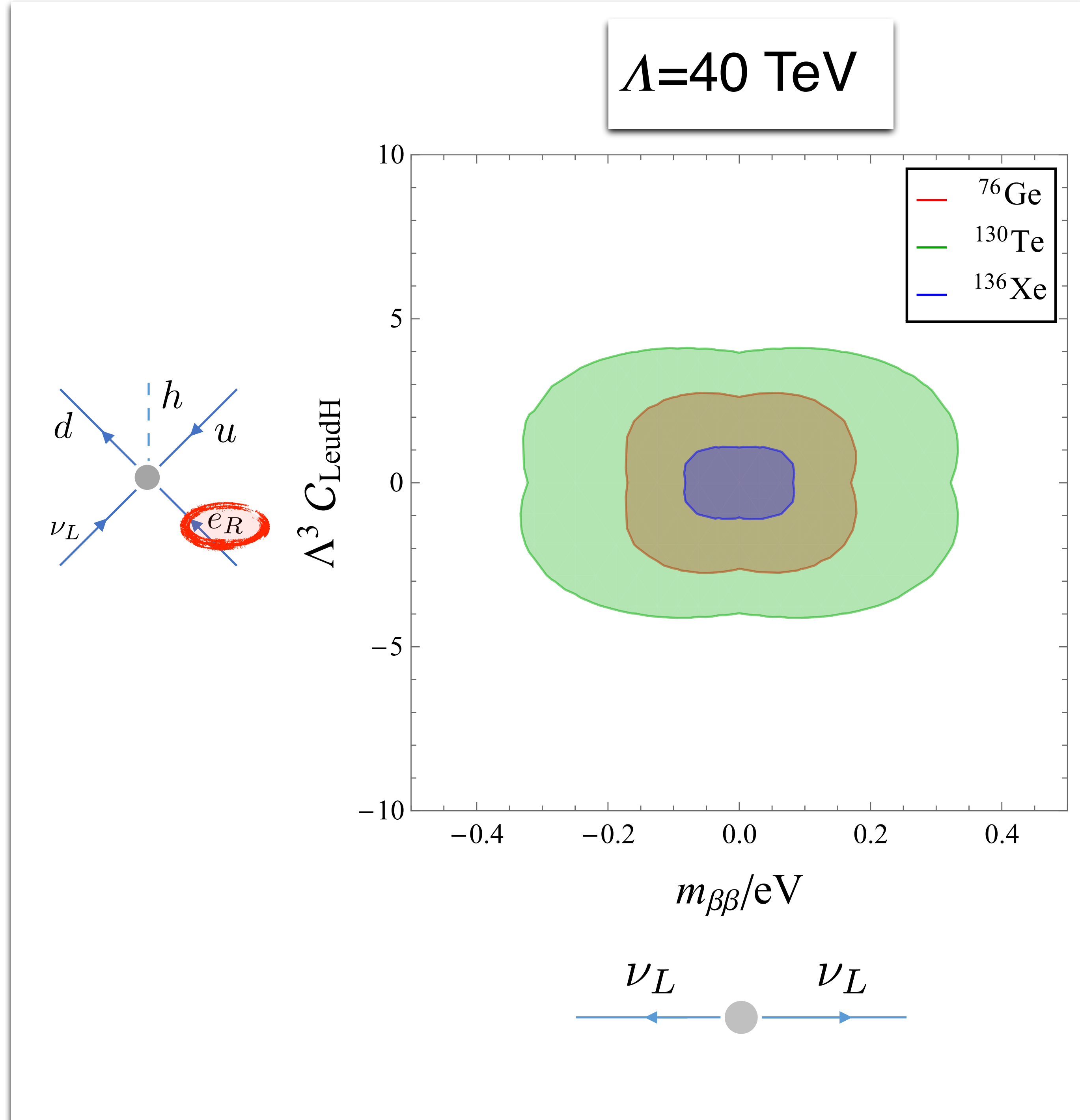
Dim 9 Scalar

Dim 9 Vector

Current limits

Two-coupling analysis





An example: LR model

In Left-Right models:

- SM gauge symmetry is extended to $SU(2)_R \times SU(2)_L \times U(1)_{B-L}$
- Allows for parity or charge-conjugation to be restored at high energies
- Explains neutrino masses through the see-saw mechanism (Type-I & Type-II)

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Violates lepton number

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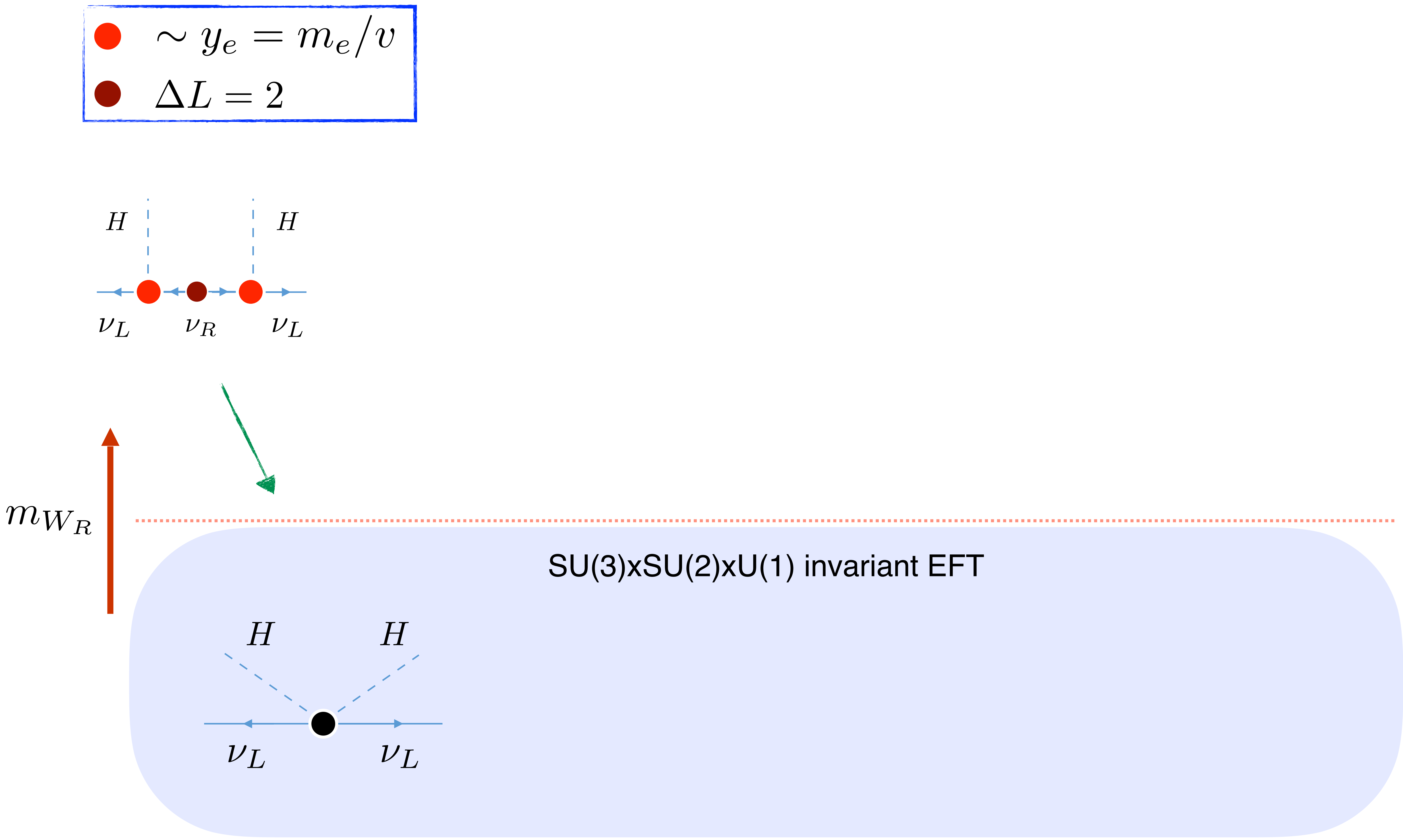

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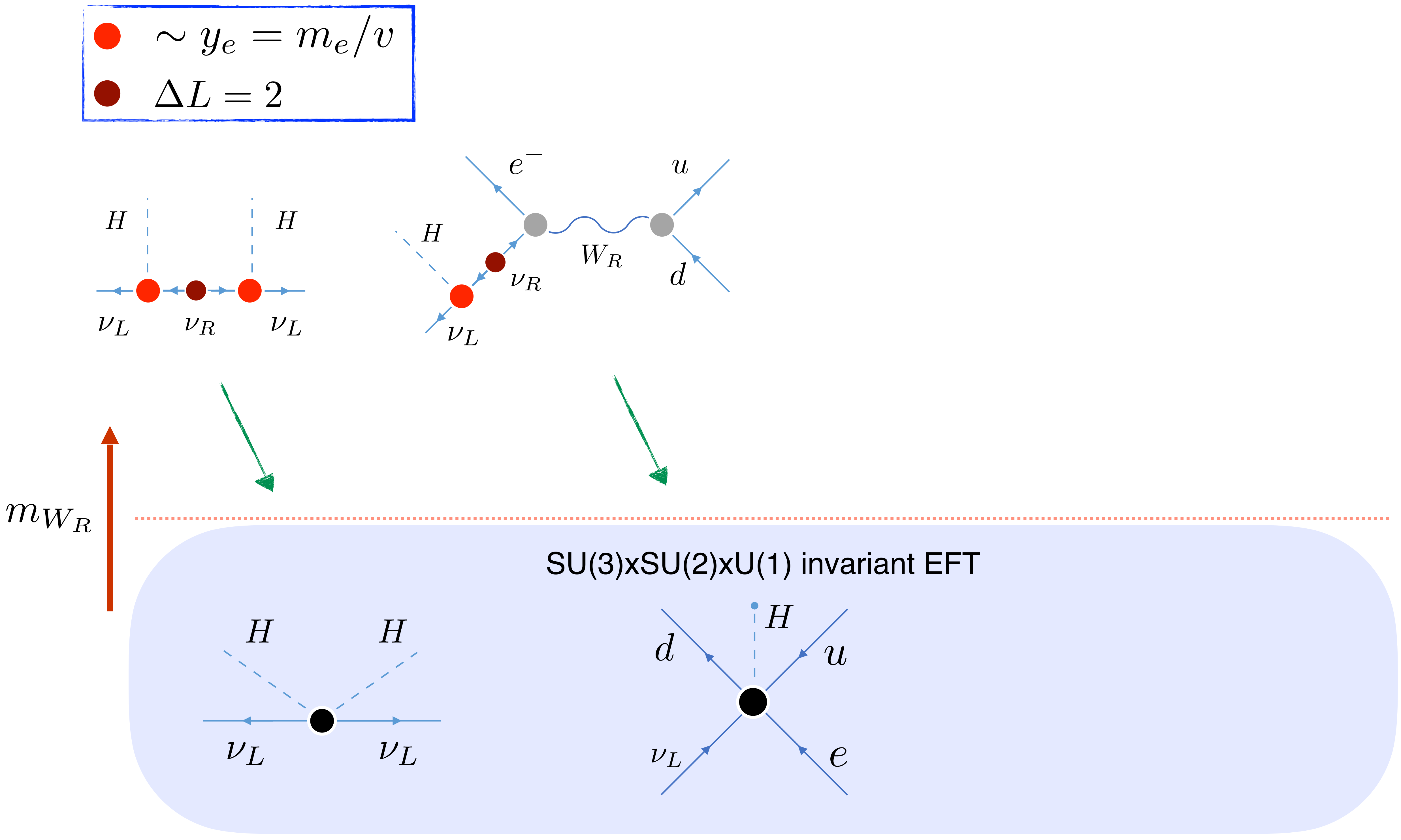

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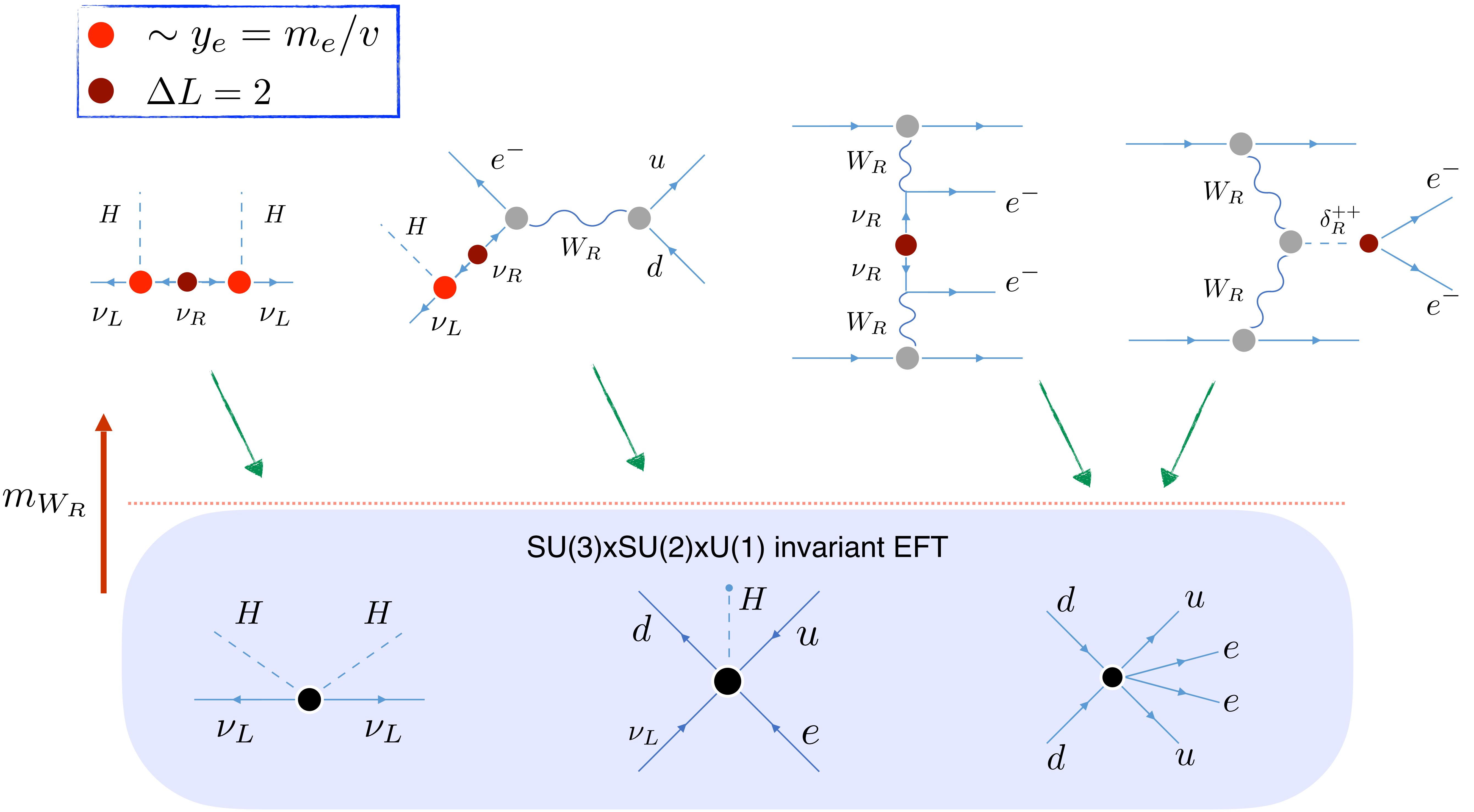
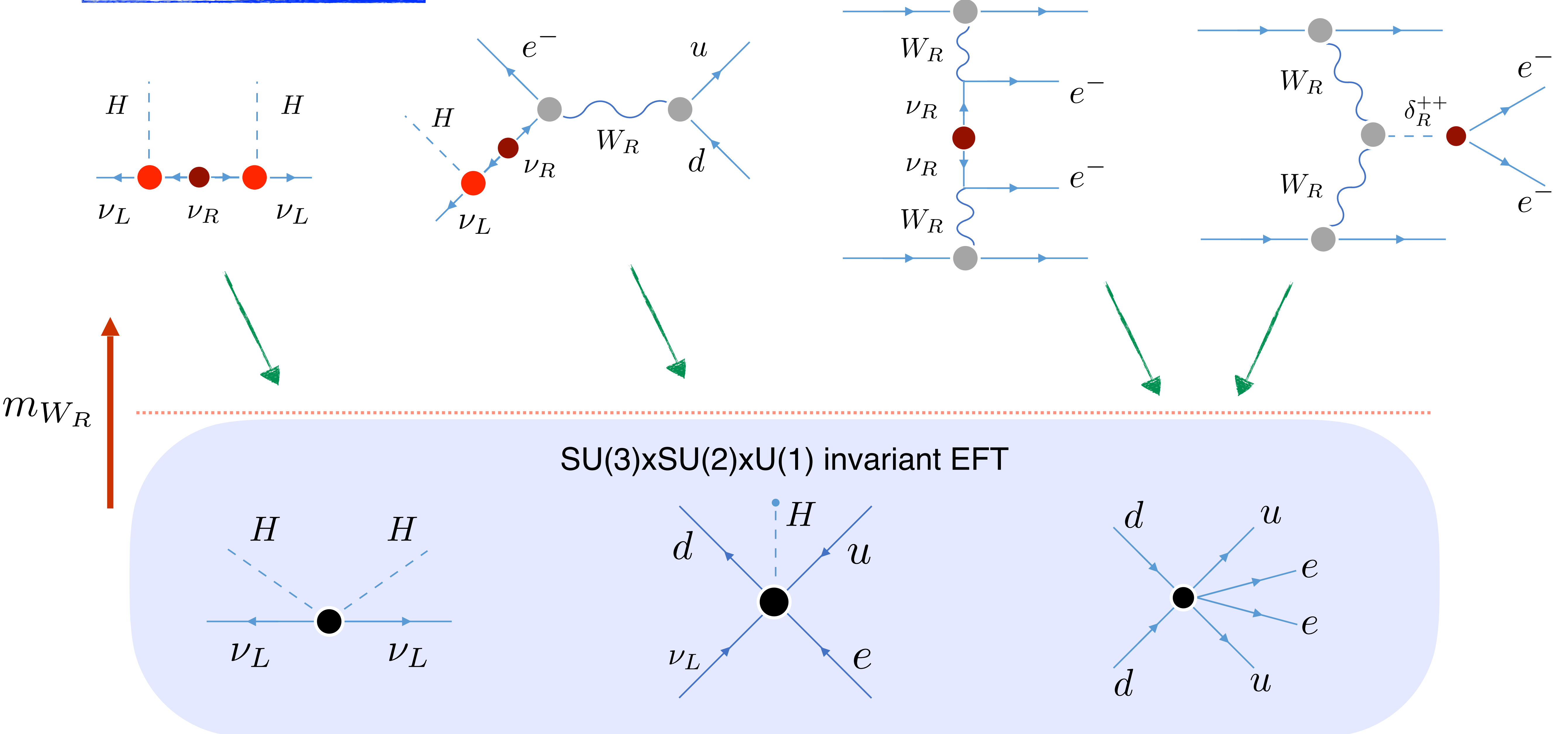


An example: LR model

- (red dot) $\sim y_e = m_e/v$
- (dark red dot) $\Delta L = 2$



dim-5 $\sim y_e^2 \left(\frac{v}{\Lambda}\right)$ dim-7 $\sim y_e \left(\frac{v}{\Lambda}\right)^3$ Dim-9 $\sim \left(\frac{v}{\Lambda}\right)^5$

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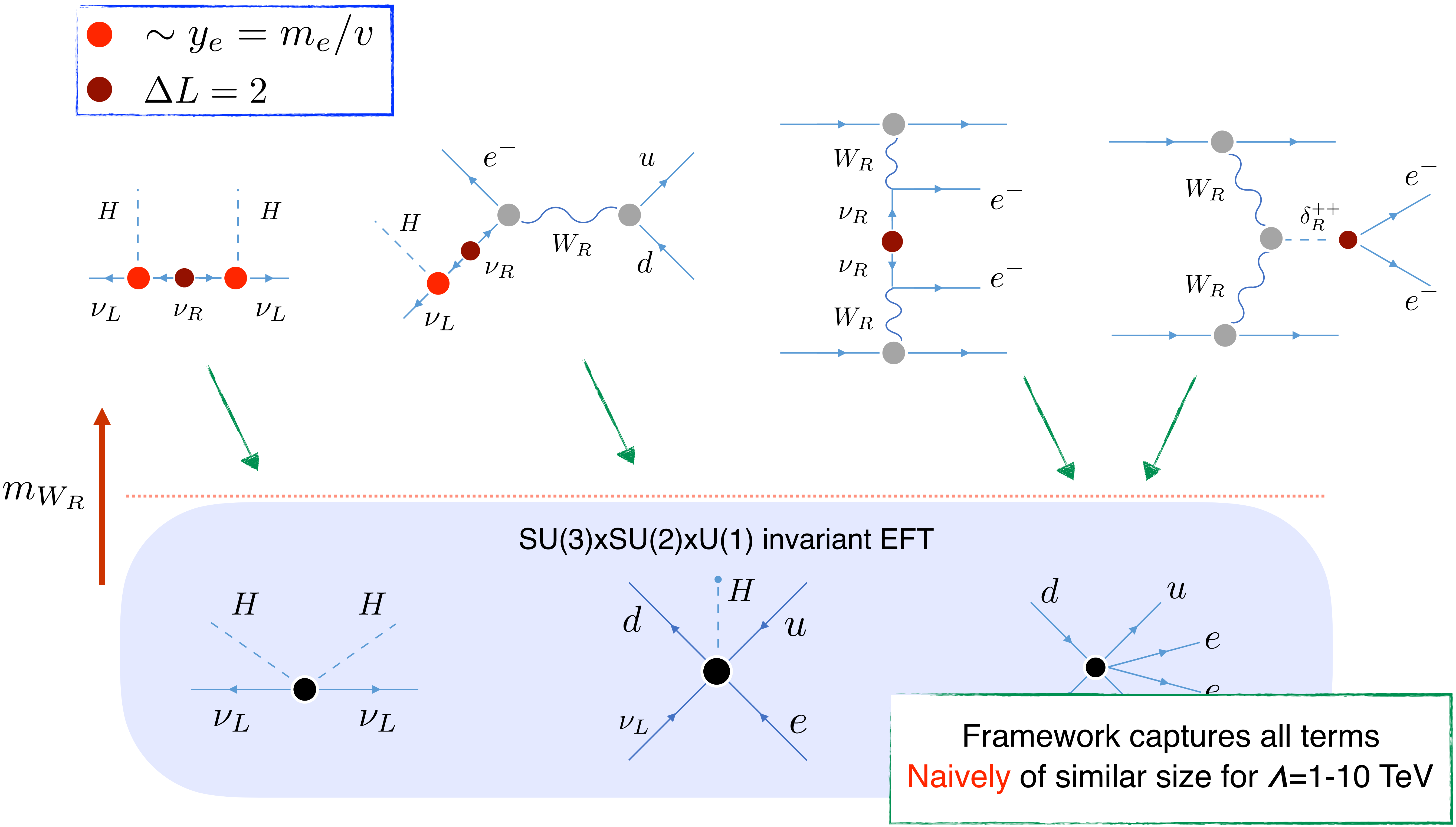

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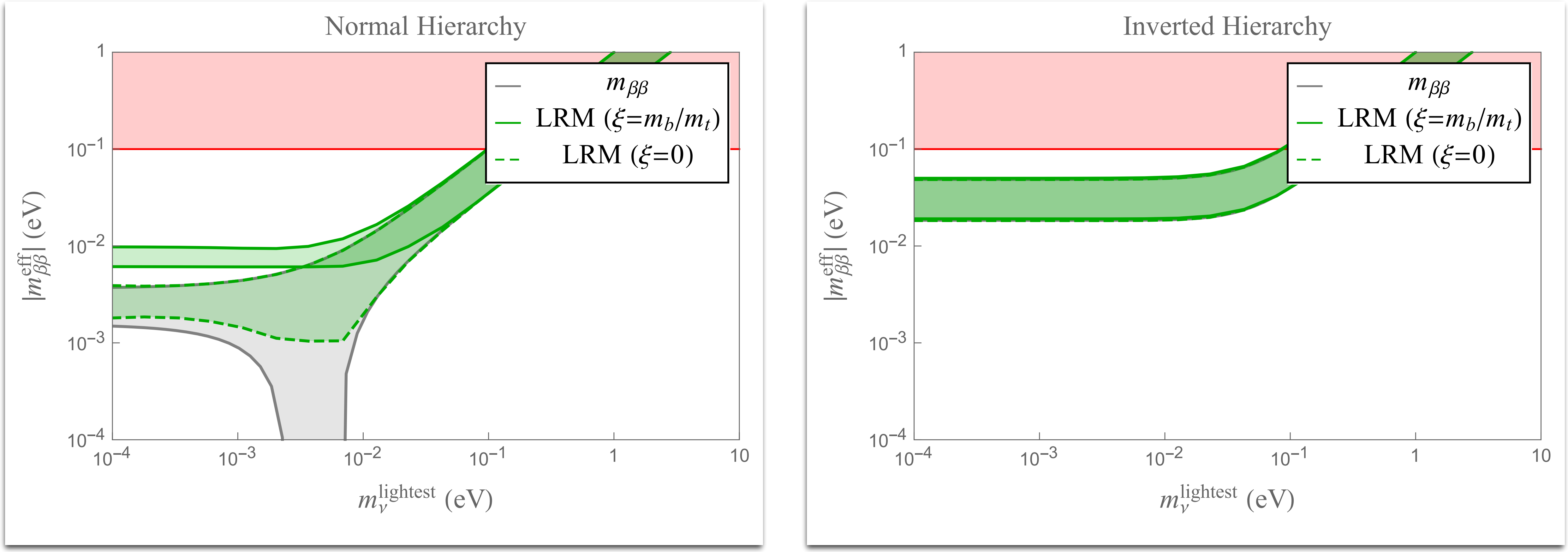
$$m_{W_R} = 4.5\,\text{TeV}, \qquad m_{\nu_R} = 10\,\text{TeV}, \qquad m_{\delta_R^{++}} = 4\,\text{TeV}$$

- Assume right-handed neutrino mixing follows the PMNS matrix



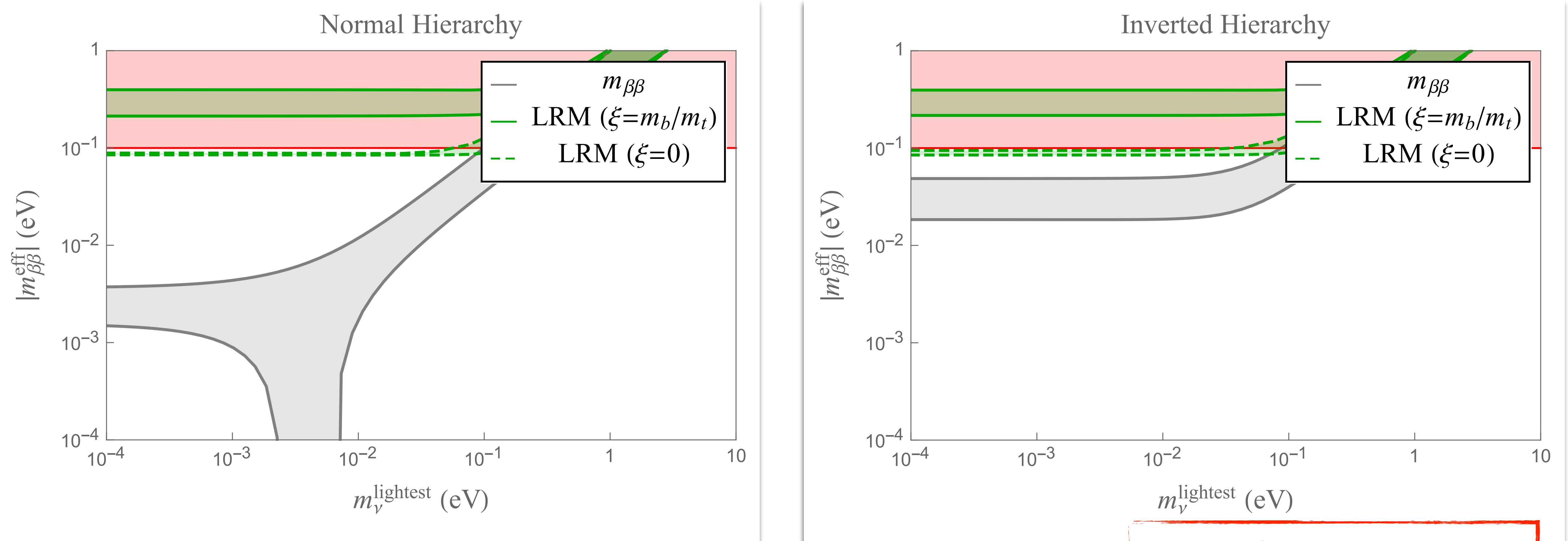
- Mild effect on NH (due to dim-9)
- Negligible effect in IH case, dim-5 terms dominate
  - Due to chiral suppression of the induced dim-6,7,9 operators

An example: LR model

Not excluded by collider

$$m_{W_R} = 4.5\,\text{TeV}, \qquad m_{\nu_R} = 10\,\text{GeV}, \qquad m_{\delta_R^{++}} = 4\,\text{TeV}$$

- Assume right-handed neutrino mixing follows the PMNS matrix

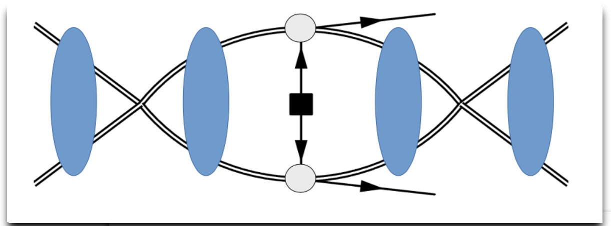


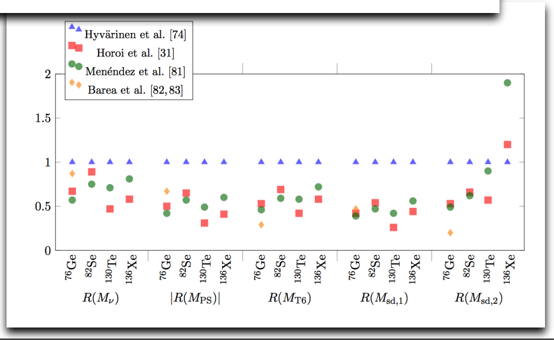
Subject to NME / LEC uncertainties

- Large effect in both NH & IH
- Now dominated by dim-9 terms

Summary

- EFTs allow one to systematically describe ΔL=2 sources
  - Matching to chiral EFT involves unknown LECs
  - Several more required by renormalization
  - Can in principle be determined from LQCD
  - Needed Nuclear Matrix Elements determined in literature





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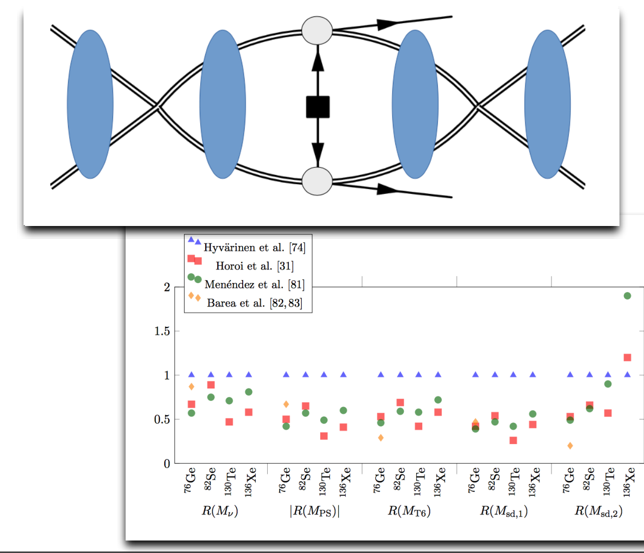
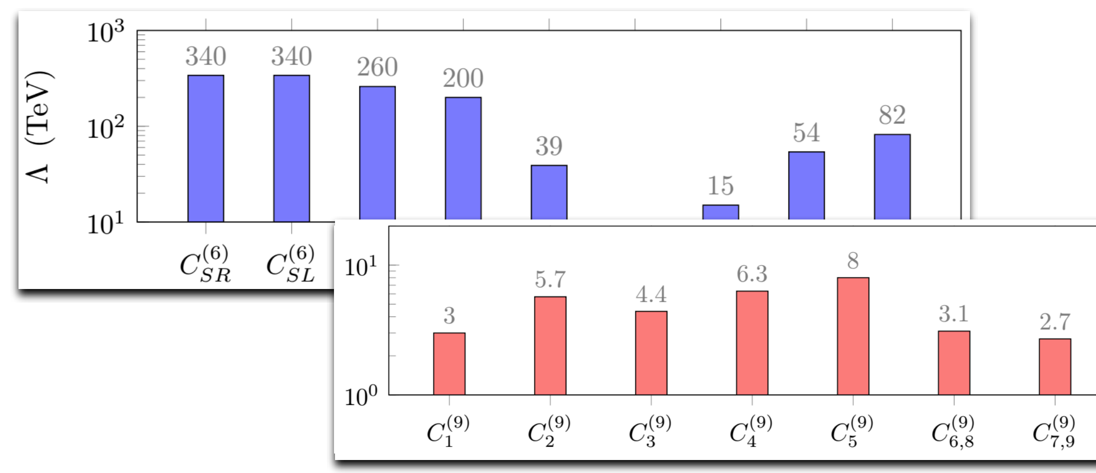


- Limits on higher-dimensional operators probe
  - O(1-10) TeV scales for dim-9
  - O(100) TeV scales for dim-7
- Order 1 uncertainties
  - Unknown LECs + NMEs



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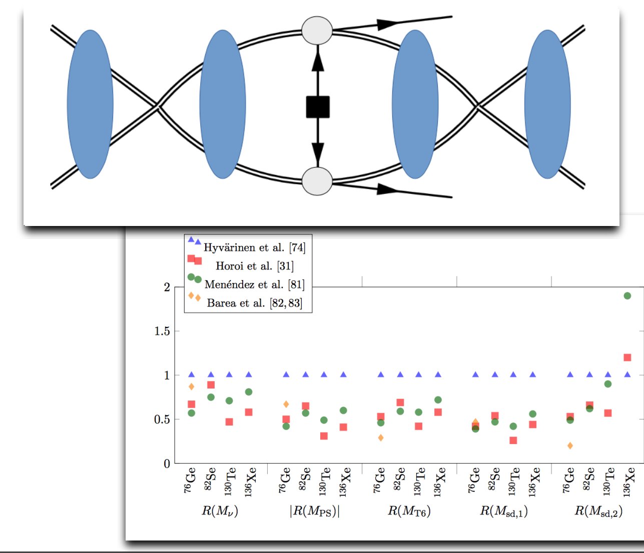

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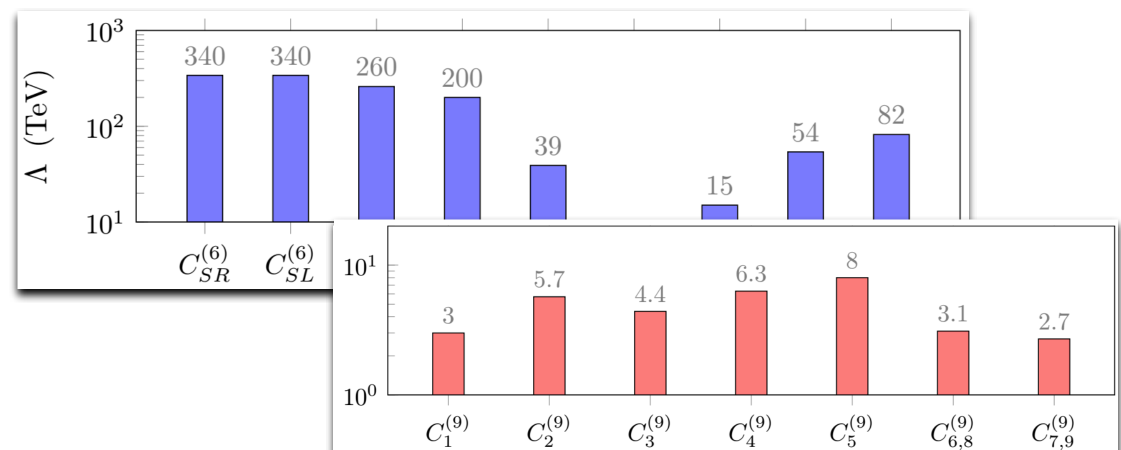
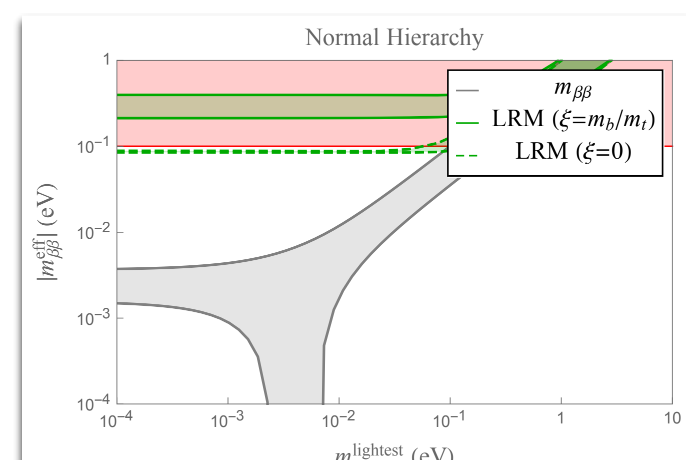


- Explicit example: Left-right model
- Induces dimension-5, -7, and -9
  - Captured by the EFT
  - Higher-dim. Operators can be important

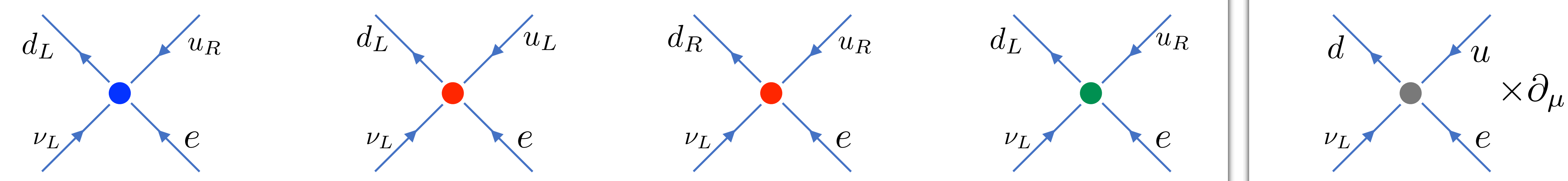


Back up slides

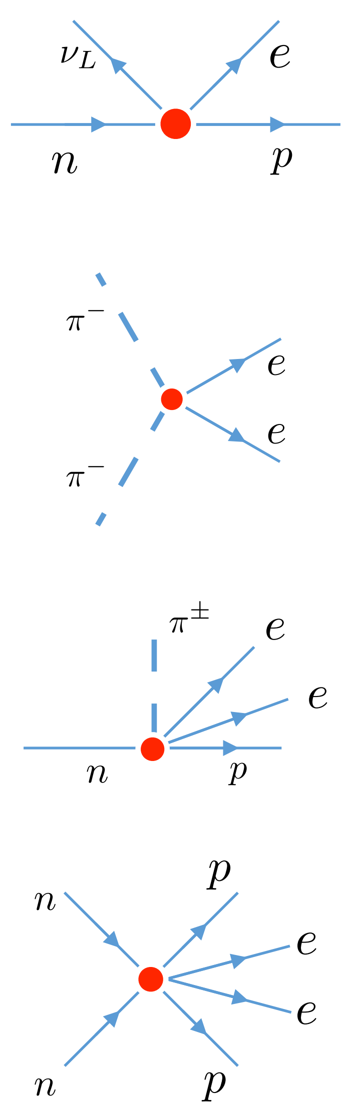
Low energy constants

LECs

	Dimension 6				Dimension 7
	$C^{(6)}_{SL,SR}$	$C^{(6)}_{VL}$	$C^{(6)}_{VR}$	$C^{(6)}_{T}$	$C^{(7)}_{VL,VR}$

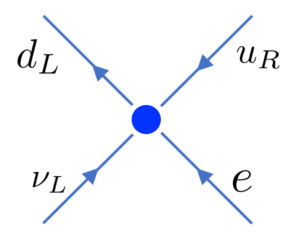
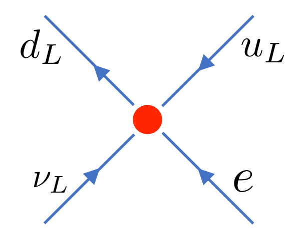
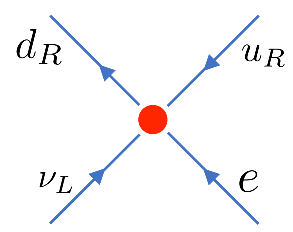
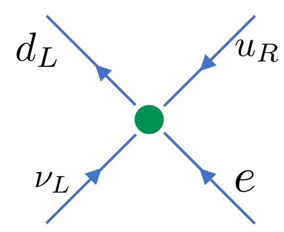
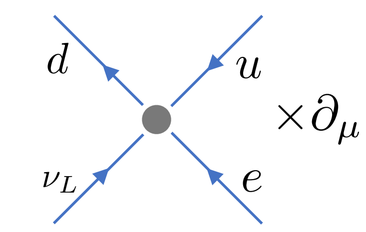
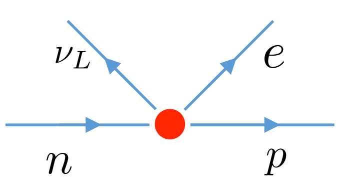
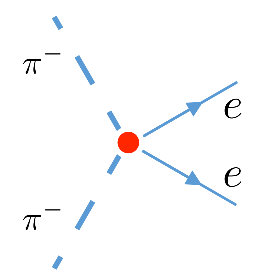
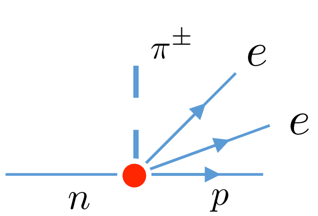
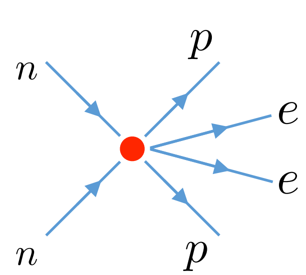


Low energy constants



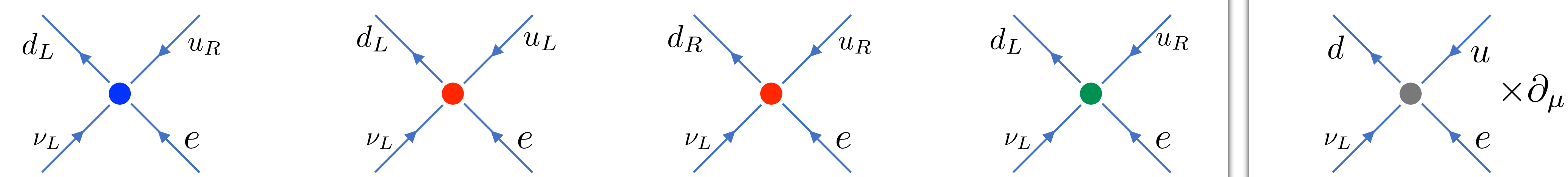
	$C^{(6)}_{SL,SR}$	$C^{(6)}_{VL}$	$C^{(6)}_{VR}$	$C^{(6)}_{T}$	$C^{(7)}_{VL,VR}$
$n \to p\, e\, \nu_L$	Quark condensate	Nucleon charges	Nucleon charges	Tensor charge NLO LEC	Quark condensate
$\pi^- \pi^- \to e\, e$				x1	
$n \to p\, \pi^\pm\, e\, e$		x1		x1	
$n\, n \to p\, p\, e\, e$		x1	x2	x1	

LECs

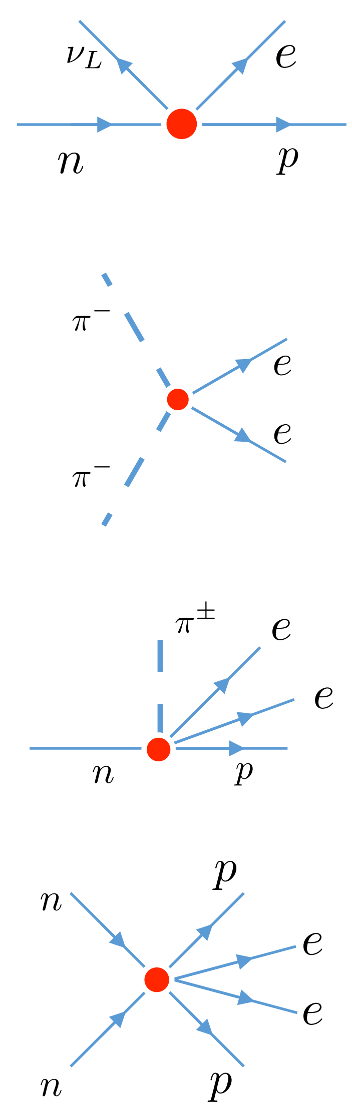
Low energy constants	Dimension 6: $C^{(6)}_{SL,SR}$	Dimension 6: $C^{(6)}_{VL}$	Dimension 6: $C^{(6)}_{VR}$	Dimension 6: $C^{(6)}_{T}$	Dimension 7: $C^{(7)}_{VL,VR}$
					
	Quark condensate ✓ LQCD	Nucleon charges ✓ expt	Nucleon charges ✓ expt	Tensor charge ✓ LQCD NLO LEC	Quark condensate ✓ LQCD
				x1	
		x1		x1	
		x1	x2	x1	

LECs

Dimension 6				Dimension 7
$C^{(6)}_{SL,SR}$	$C^{(6)}_{VL}$	$C^{(6)}_{VR}$	$C^{(6)}_{T}$	$C^{(7)}_{VL,VR}$



Low energy constants



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$\nu_L\, n \to p\, e$	Quark condensate ✓ LQCD	Nucleon charges ✓ expt	Nucleon charges ✓ expt	Tensor charge ✓ LQCD; NLO LEC ✗	Quark condensate ✓ LQCD
$\pi^-\pi^- \to e\, e$				x1 ✗	
$\pi^\pm\, n \to p\, e\, e$		x1 ✗		x1 ✗	
$n\, n \to p\, p\, e\, e$		x1 ✗	x2 ✗	x1 ✗	

LECs

Dimension 6				Dimension 7
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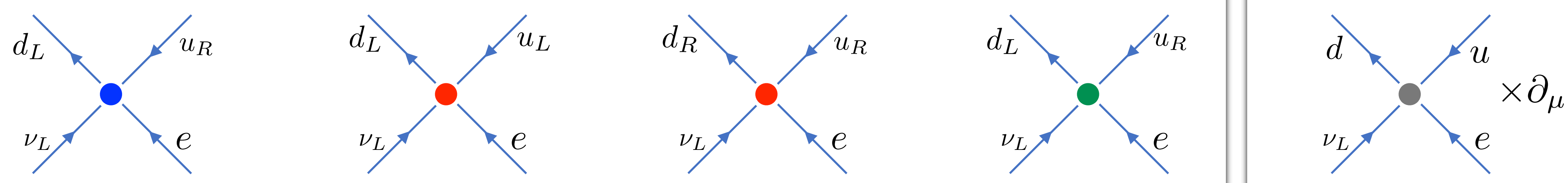
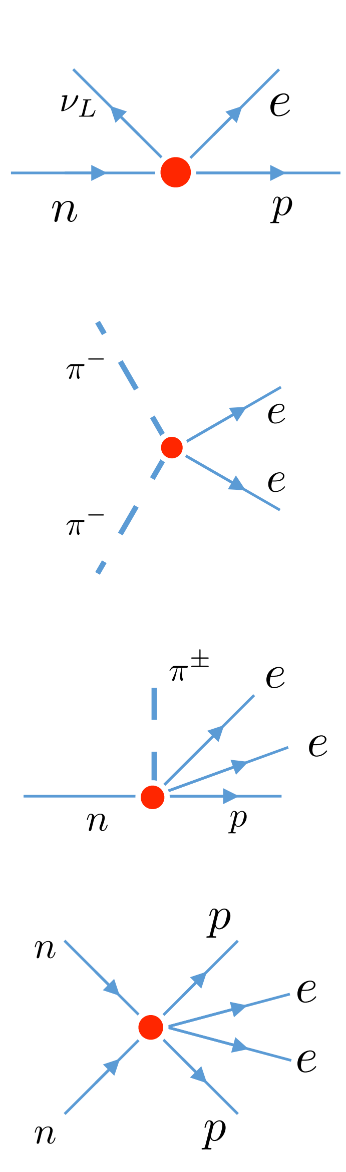


Low energy constants



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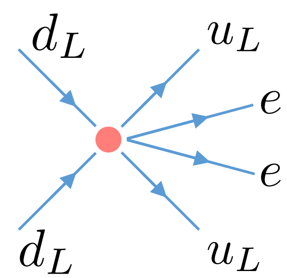
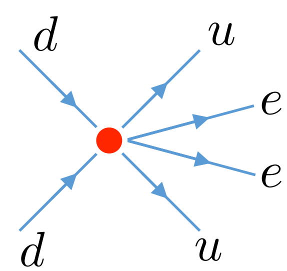
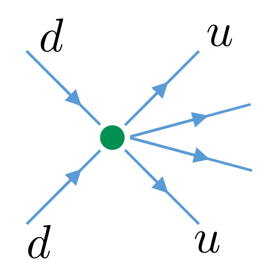
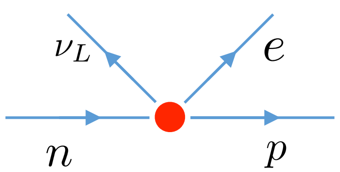
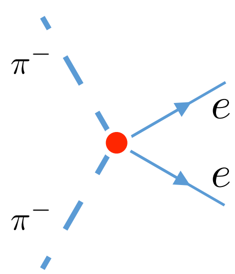
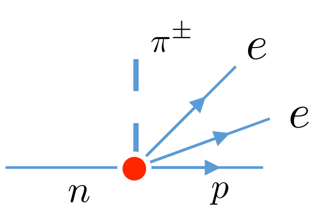
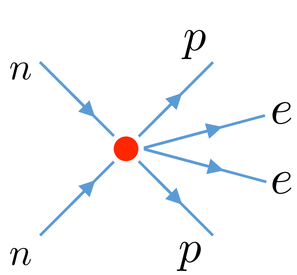
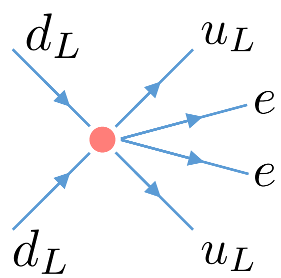
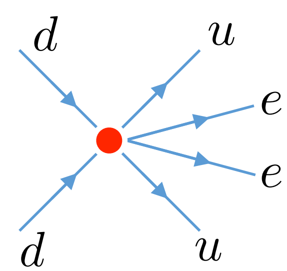
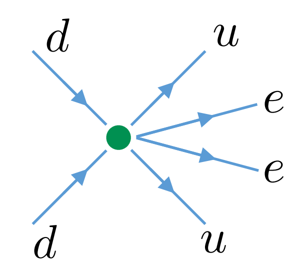
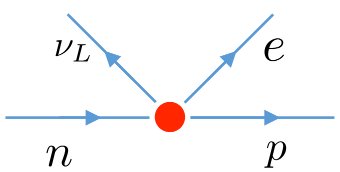
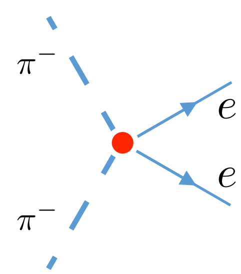
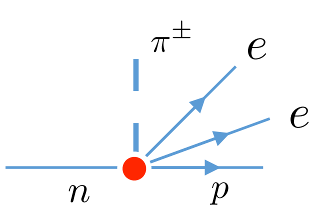
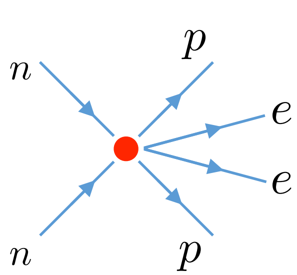
$\pi^- \pi^- \to e\, e$				x1 ✗	
$\pi^\pm\, n \to p\, e\, e$		x1 ✗		x1 ✗	
$n\, n \to p\, p\, e\, e$		x1 ✗	Non-NDA x2 ✗	x1 ✗	

LECs

Low energy constants

	Dimension 9 - scalar		Dimension 9 - vector
	$C^{(9)}_{1L,1R}$ 	$C^{(9)}_{2,3,4,5}$ 	$C^{(9)}_{6,7,8,9}$ 
			
	x1	x4	
	x1		x2
	x1	x4	x2

LECs

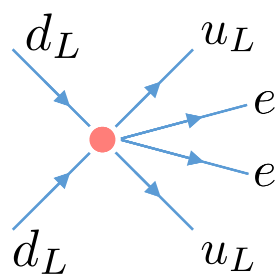
Low energy constants	Dimension 9 - scalar: $C^{(9)}_{1L,1R}$ 	Dimension 9 - scalar: $C^{(9)}_{2,3,4,5}$ 	Dimension 9 - vector: $C^{(9)}_{6,7,8,9}$ 
			
	x1 ✓ LQCD	x4 ✓ LQCD	
	x1		x2
	x1	x4	x2

Nicholson et al.'18

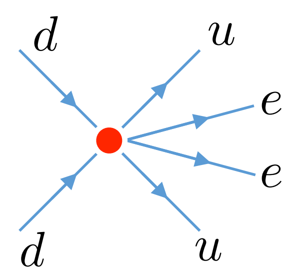
LECs

Dimension 9 - scalar

$C^{(9)}_{1L,1R}$

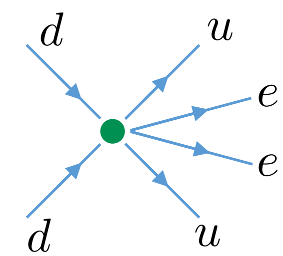


$C^{(9)}_{2,3,4,5}$

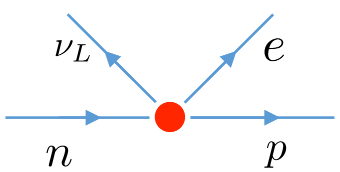
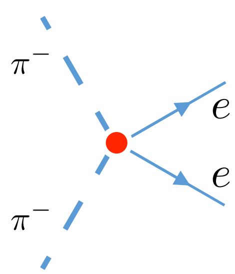
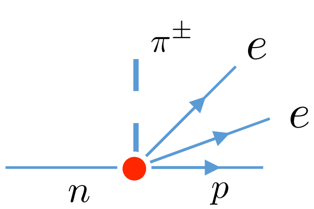
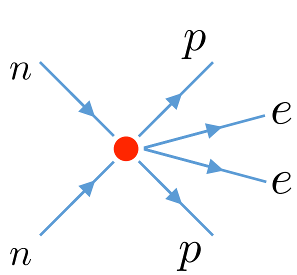


Dimension 9 - vector

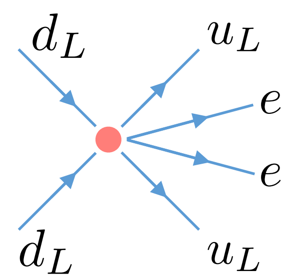
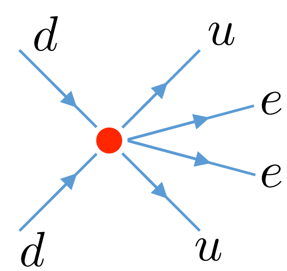
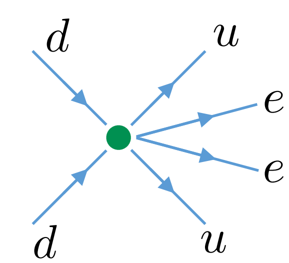
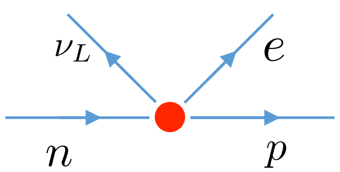
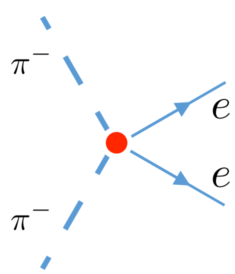
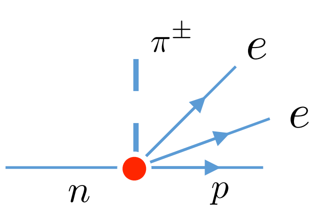
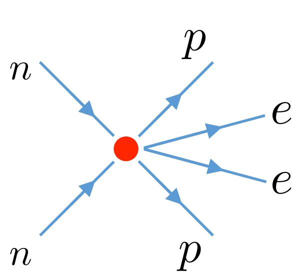
$C^{(9)}_{6,7,8,9}$



Low energy constants

Low energy constants	$C^{(9)}_{1L,1R}$	$C^{(9)}_{2,3,4,5}$	$C^{(9)}_{6,7,8,9}$
			
	x1 ✓ LQCD	x4 ✓ LQCD	
	x1 ✗		x2 ✗
	x1 ✗	x4 ✗	x2 ✗

LECs

Low energy constants	Dimension 9 - scalar: $C^{(9)}_{1L,1R}$ 	Dimension 9 - scalar: $C^{(9)}_{2,3,4,5}$ 	Dimension 9 - vector: $C^{(9)}_{6,7,8,9}$ 
			
	x1 ✓ LQCD	x4 ✓ LQCD	
	x1 ✗		x2 ✗
	x1 ✗	x4 ✗ Non-NDA	x2 ✗

Disentangling operators

Disentangling operators

What if a $0\nu\beta\beta$ signal is measured?

- Measurement in a single isotope could be due to any operator
- Could measure the rate in several nuclei, however
  - Different isotopes have similar sensitivity to LNV

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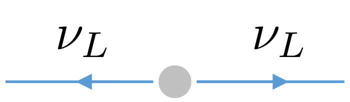
- Instead look at angular & energy distributions of the

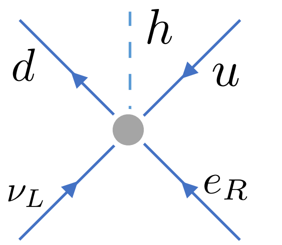
Disentangling operators

What if a $0\nu\beta\beta$ signal is measured?

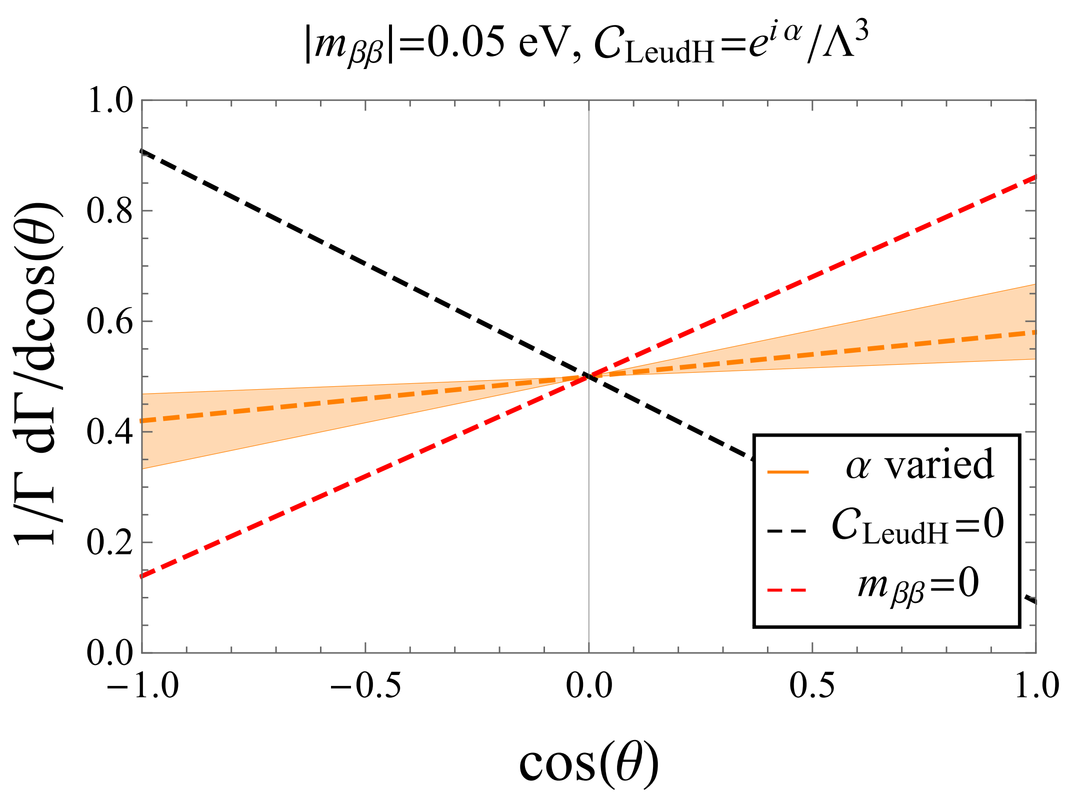
- Picking the allowed values



$m_{\beta\beta} = 0.05\,\mathrm{eV}$



$\mathcal{C}_{\mathrm{LeudH}} = e^{i\alpha}/\Lambda^3$ $\Lambda = 40\,\mathrm{TeV}$



Disentangling operators

What if a $0\nu\beta\beta$ signal is measured?

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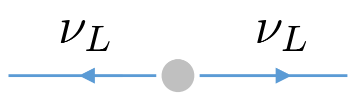

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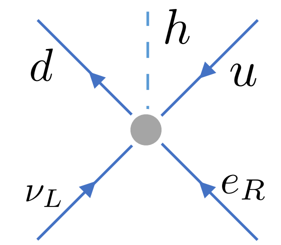
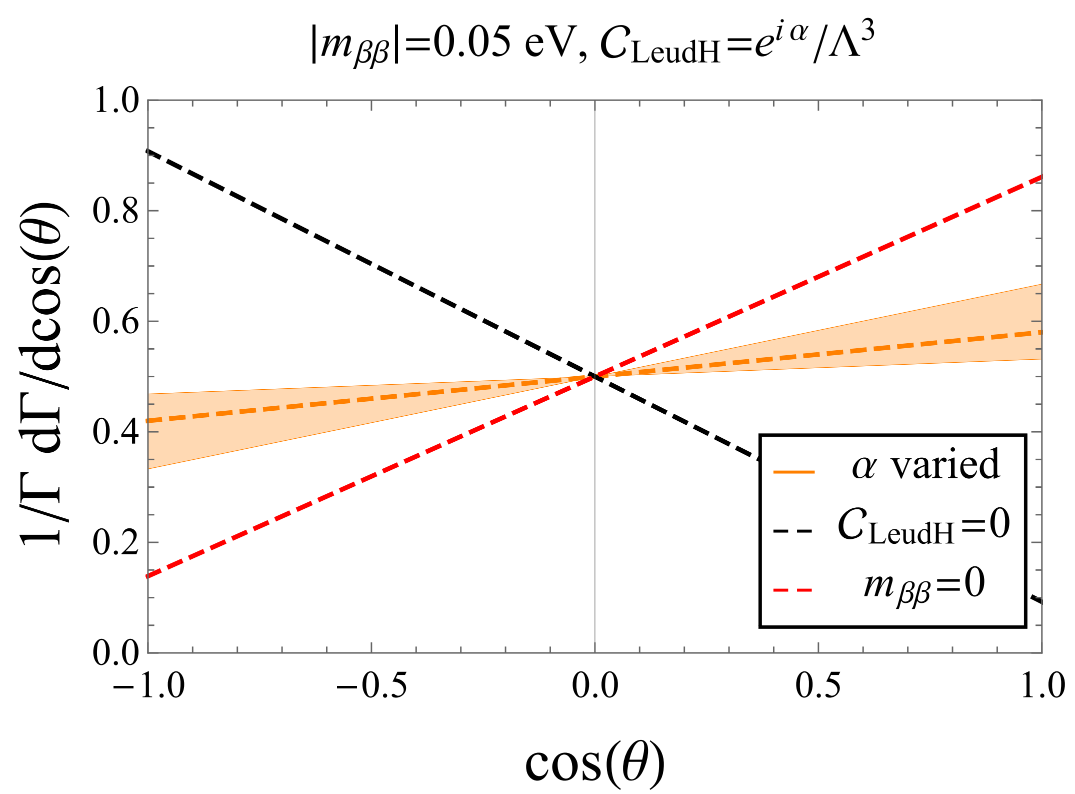
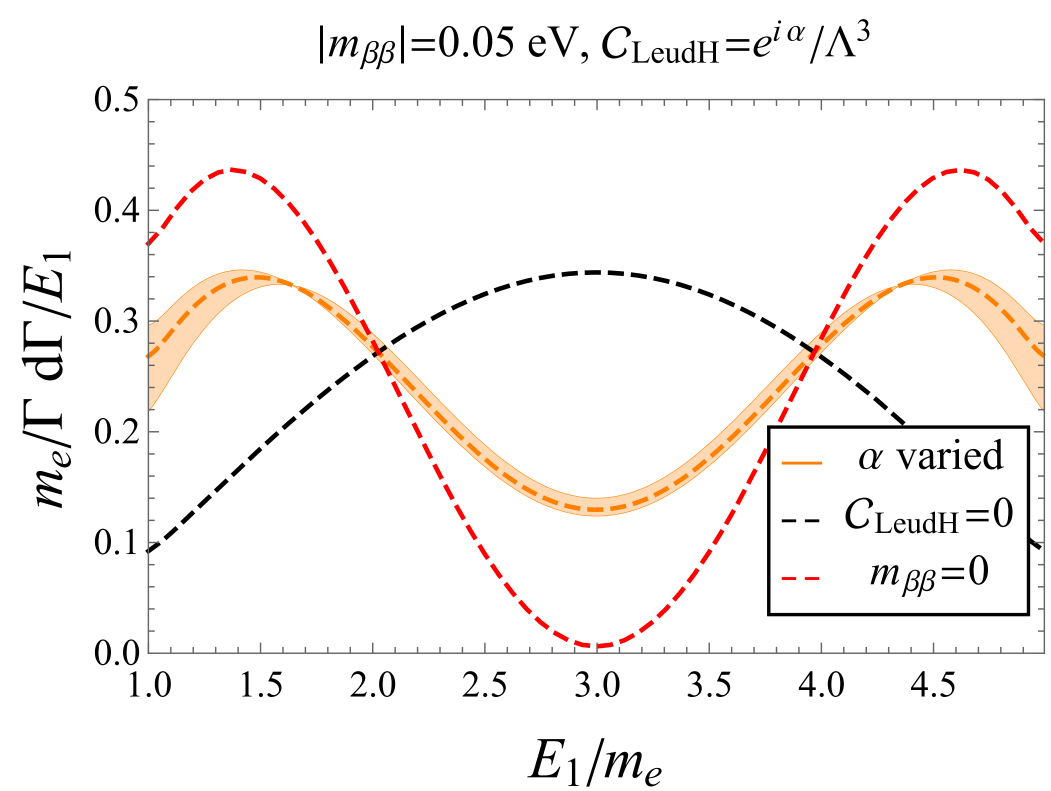


$C_{\mathrm{LeudH}} = e^{i\alpha}/\Lambda^3$ $\quad\Lambda = 40\,\mathrm{TeV}$





Why keep Dimension 7 and 9?

Effective Field Theory

Naive scaling of Dimension 5, 7, 9 operators

$$\mathcal{A}_{0\nu\beta\beta} \sim \frac{c_5}{\Lambda}\left[1 + \left(\frac{v}{\Lambda}\right)^2 \frac{c_7}{c_5} + \left(\frac{v}{\Lambda}\right)^4 \frac{c_9}{c_5}\right]$$

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- This happens in well-known BSM models
  - For example the Left right model gives

$$c_9 = \mathcal{O}(1), \qquad c_7 = \mathcal{O}(y_e), \qquad c_5 = \mathcal{O}(y_e^2)$$
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