

The background of the slide is a large, dense cluster of small, semi-transparent spheres in various colors including white, light blue, light green, and light orange. The spheres are arranged in a way that suggests a hot, dense medium, possibly representing a quark-gluon plasma or a similar state of matter. The overall effect is a soft, glowing, and textured background.

# Direct Photon Flow at the LHC

ECT\* workshop:  
Electromagnetic Radiation from  
Hot and Dense Hadronic Matter

Klaus Reygers  
Heidelberg University  
28 November 2018

# Outline

- Direct-photon  $v_2$  at the LHC
- Statistical Methods (correlated uncertainties, Bayesian approach)
- Towards quantifying the statistical significance of the direct-photon puzzle
- Reducing systematic uncertainties
- Thoughts on the puzzle

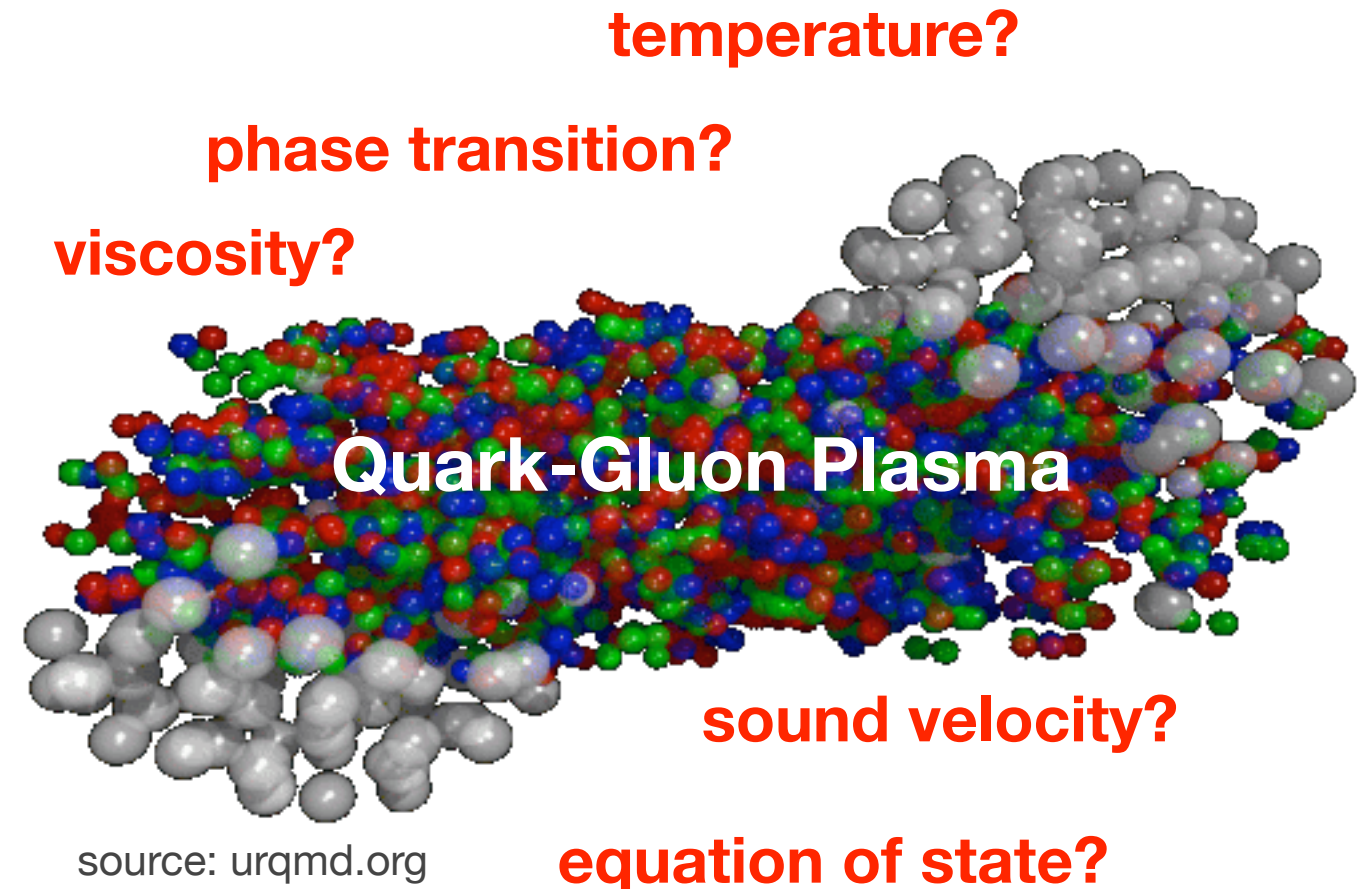
# Why Heavy-Ion Collisions?

- Particle physics: reductionism
- Heavy-ion physics:  
**Emergent** properties of QCD

More precisely:  
"Material properties" of the QGP?

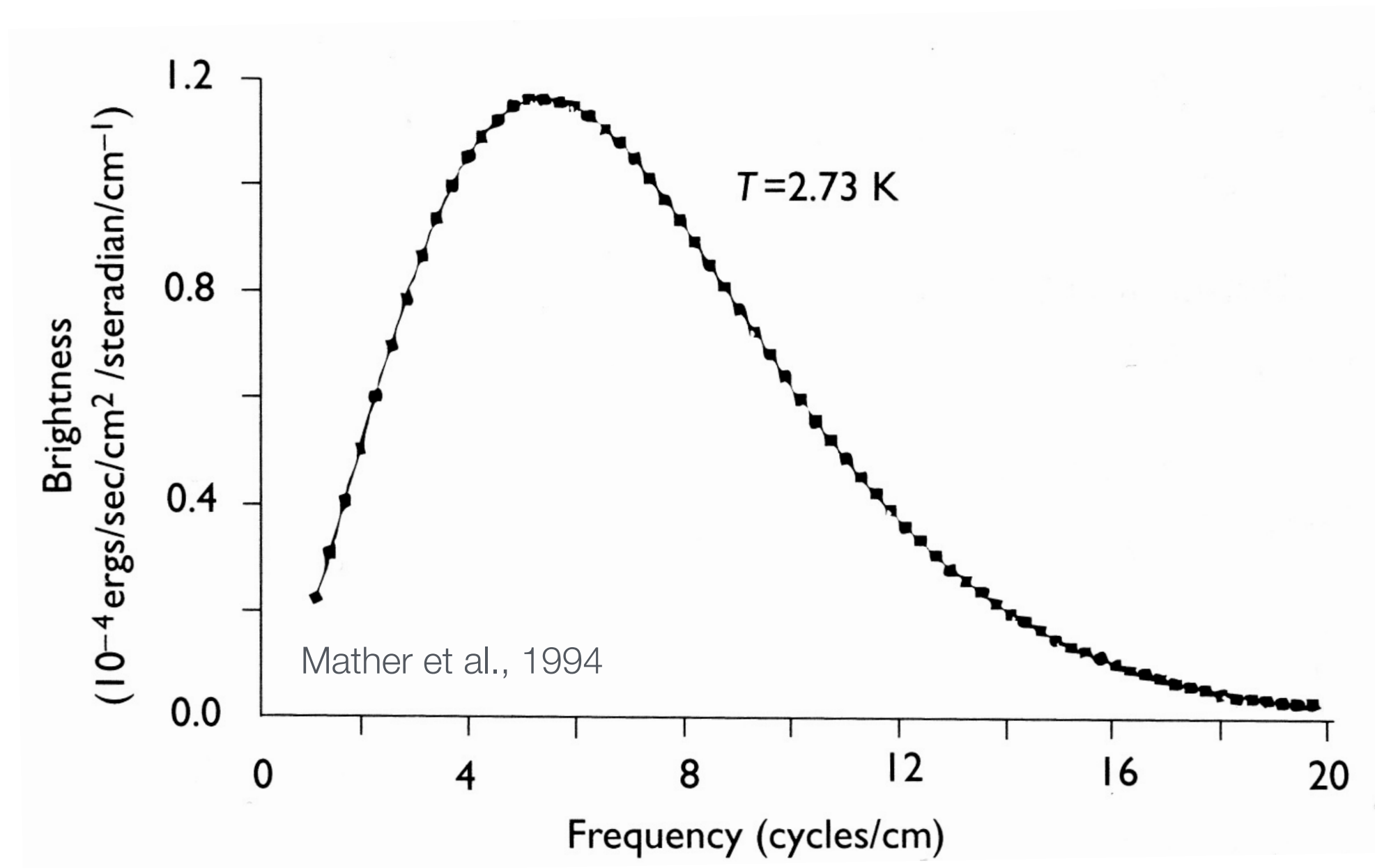
„More is different“

Philip W. Anderson,  
Science, 177, 1972, p. 393



# An Iconic Figure from Another Field

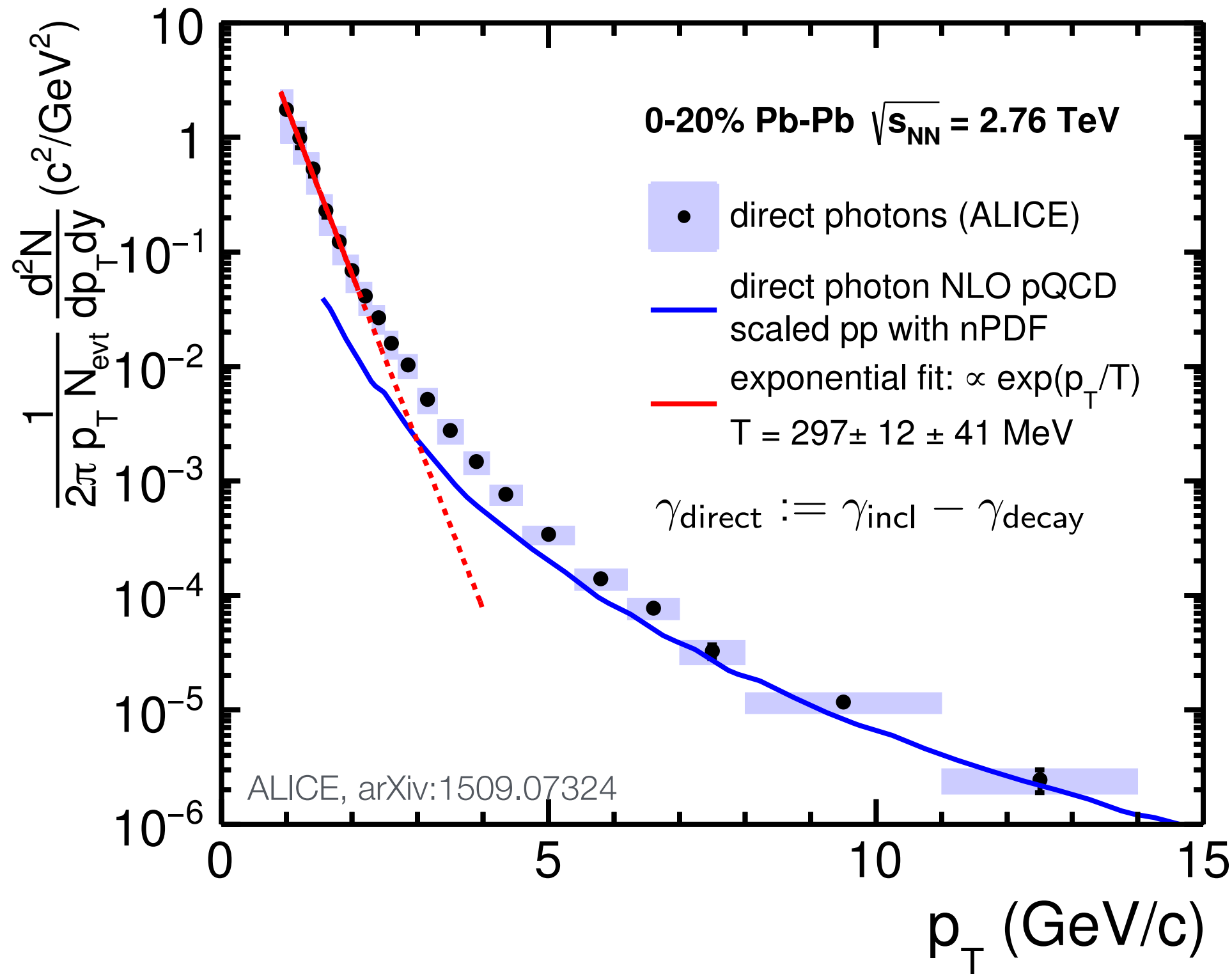
CMB black-body spectrum (COBE)



Recipe: Good data + well understood theory

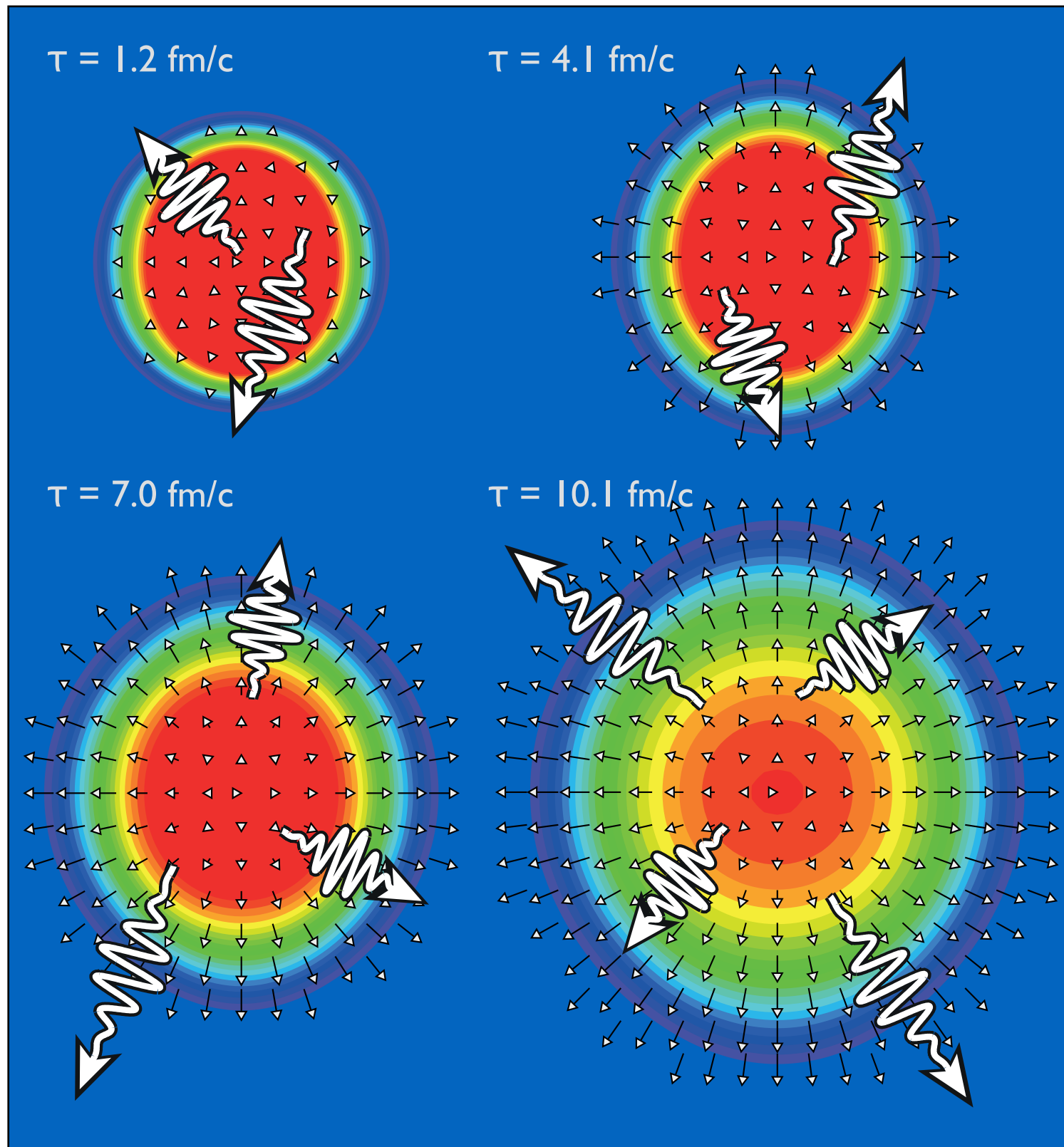


# A Candidate for an Iconic Figure from Heavy Ions: Planck-like Photon Spectrum



Current proxy (here from the LHC) looks already OK, but statistical significance needs to be improved

# Direct-Photon $v_n$ : An additional Handle to Check Our Understanding of Photon Production



- Photons produced over the entire duration of the collision
  - ▶ Test understanding of the space-time evolution and production mechanisms
  - ▶ Access to initial  $T_{\text{QGP}}$
  - ▶ Expect more photons per pion at low  $p_T$  than in pp
- But: Slope  $T_{\text{eff}} > T_{\text{QGP}}$  due to blue shift

QGP photon rate  $r_\gamma$ :

$$E_\gamma \frac{dr_\gamma}{d^3p} \propto \alpha \alpha_s T^2 e^{-E_\gamma/T} \log \frac{E_\gamma T}{k_c^2}$$

Total emission rate:

$$r_\gamma \propto T^4$$

Direct-photon  $v_2$  at the LHC

# The Master Formula

$v_2$  of all photons  
("inclusive photons")

calculated  $v_2$  of  
decay photons

$$v_2^{\gamma, \text{dir}} = \frac{R_\gamma v_2^{\gamma, \text{inc}} - v_2^{\gamma, \text{dec}}}{R_\gamma - 1}$$

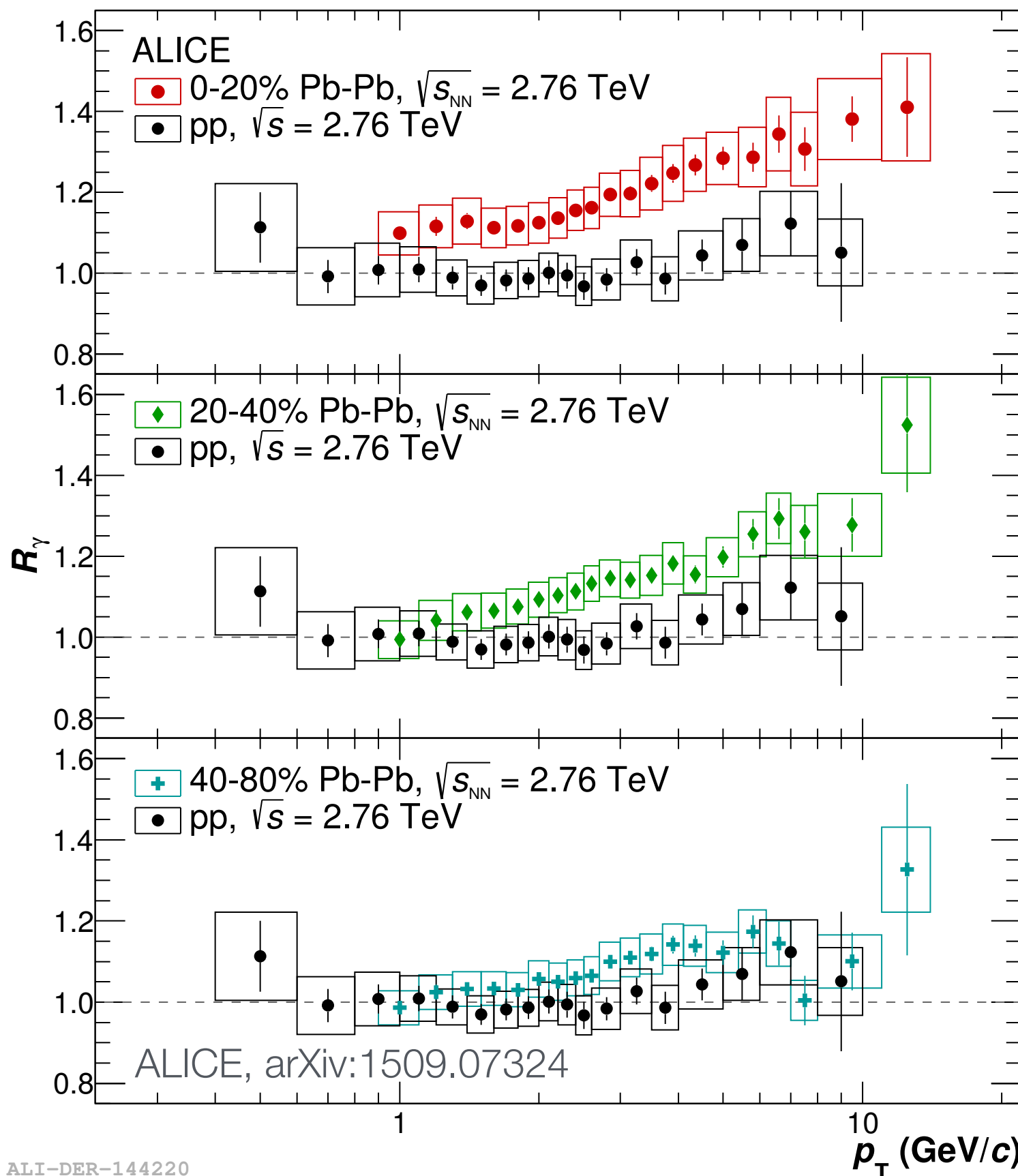
photon excess:

$$R_\gamma = \frac{\gamma_{\text{incl}}}{\gamma_{\text{decay}}} = 1 + \frac{\gamma_{\text{dir}}}{\gamma_{\text{decay}}}$$

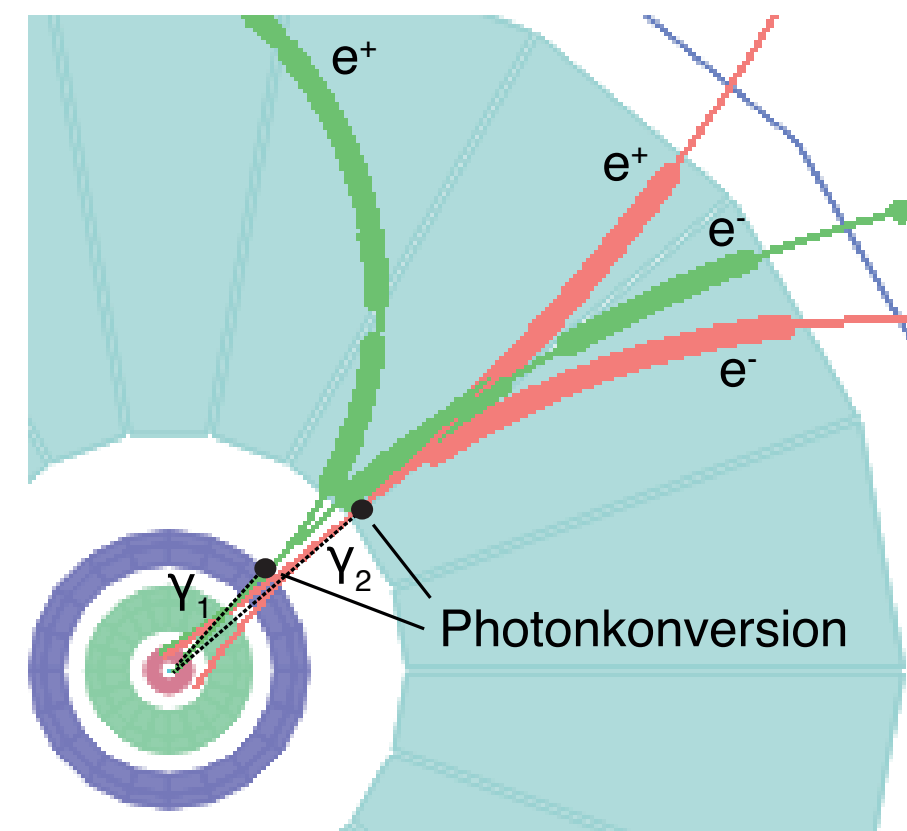


# $R_\gamma$ at the LHC (Pb-Pb at $\sqrt{s_{NN}} = 2.76$ TeV)

$$R_\gamma = \frac{\gamma_{incl}}{\gamma_{decay}}$$



- Low- $p_T$  excess in Pb-Pb
- No low- $p_T$  excess in pp
- But: Low significance of  $R_\gamma$  makes extraction  $v_{2,dir}$  challenging



+ PHOS calorimeter measurement

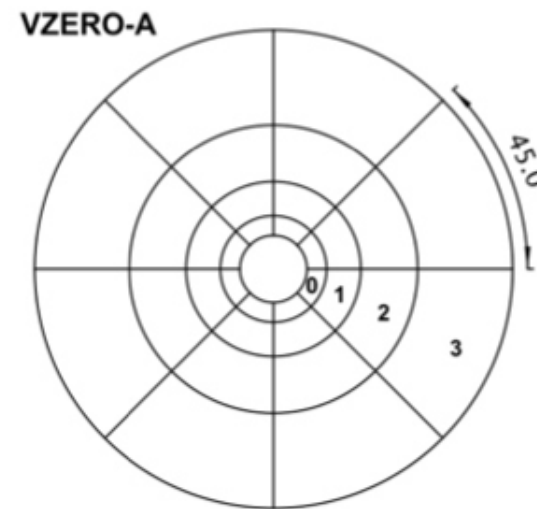
# Inclusive Photon $v_2$ : Scalar Product Method

Reference particles in VZERO-A and VZERO-C:

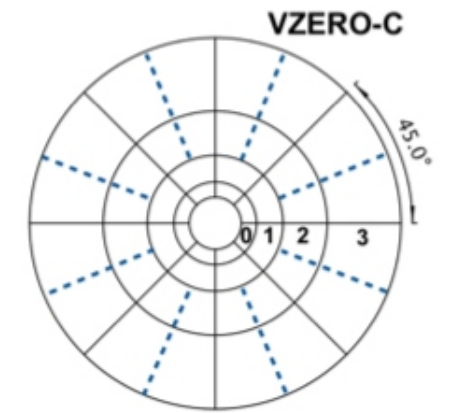
$$\vec{Q}_n = \sum_{i \in \text{RFP}} w_i e^{in\varphi_i}$$

For each photon:

$$\vec{u}_2 = e^{i2\varphi}$$



$$2.8 < \eta < 5.1$$



$$-3.7 < \eta < -1.7$$

$$v_2 = \sqrt{\frac{\langle \langle \vec{u}_2 \cdot \frac{\vec{Q}_2^{A*}}{M_A} \rangle \rangle \langle \langle \vec{u}_2 \cdot \frac{\vec{Q}_2^{C*}}{M_C} \rangle \rangle}{\langle \frac{\vec{Q}_2^A}{M_A} \cdot \frac{\vec{Q}_2^{C*}}{M_C} \rangle}}$$

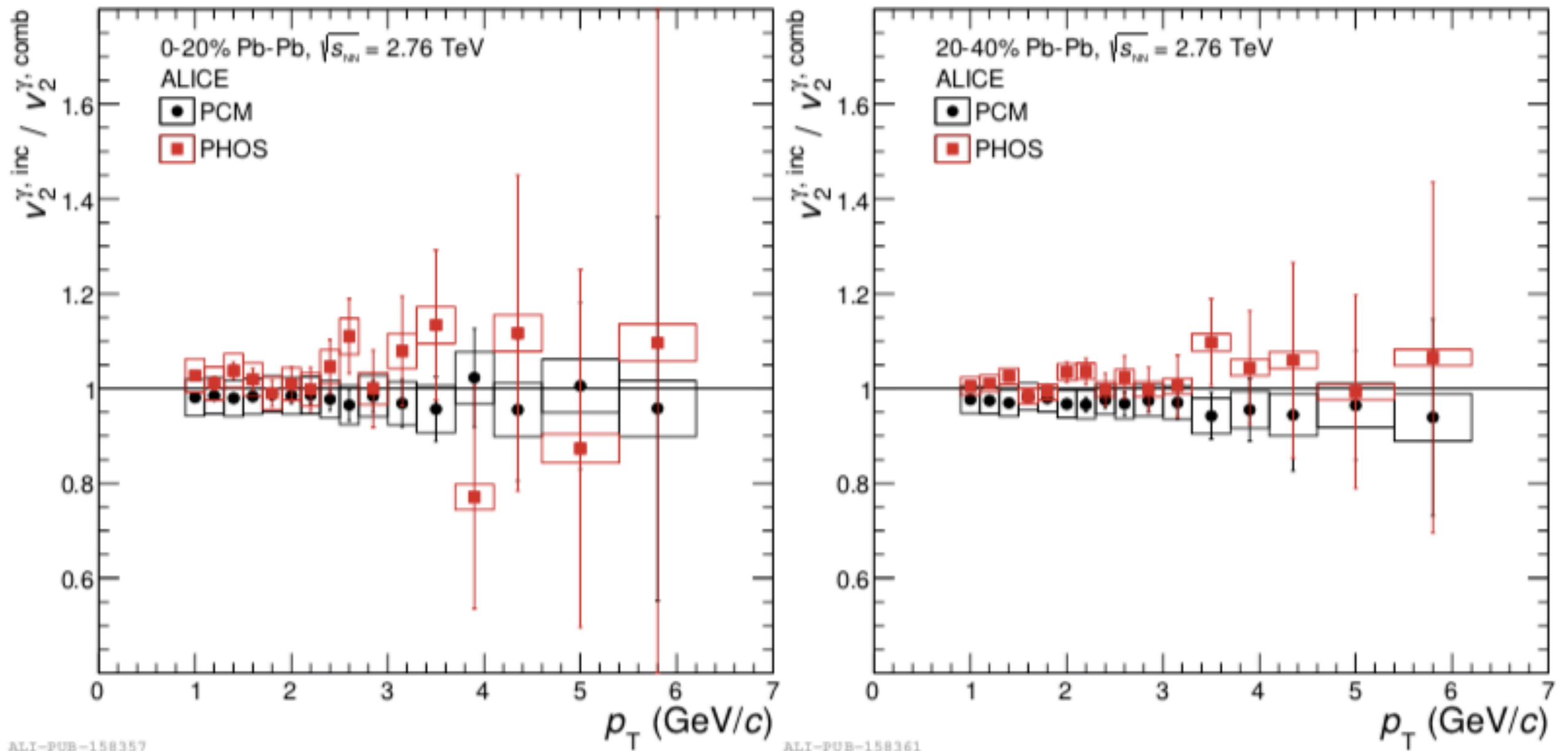
Advantage over event plane method: result not sensitive to event plane resolution

# $v_{2,\text{inc}}$ and $v_{2,\text{dec}}$ Systematic Uncertainties

Centrality	0–20%		20–40%	
$p_T$ (GeV/c)	2.0	5.0	2.0	5.0
<b>PCM</b>				
Photon selection	2.4	4.2	2.1	4.0
Energy resolution	1.0	1.0	1.0	1.0
Efficiency	3	3	1.9	1.9
Total	4.0	5.3	3.0	4.5
<b>PHOS</b>				
Efficiency & contamination	3.0	3.0	0.7	0.7
Event plane flatness	2.0	2.0	1.4	1.4
Total	3.5	3.5	1.6	1.6
<b>Decay photon calculation</b>				
Parameterization of $v_2^\pi$	1.3	3.6	0.8	2.2
$\eta/\pi^0$ normalization	1.7	3.2	1.7	2.4
Total	2.2	4.8	1.9	3.3

# Agreement between $v_{2,inc}$ from PCM and PHOS

ALICE arXiv:1805.04403

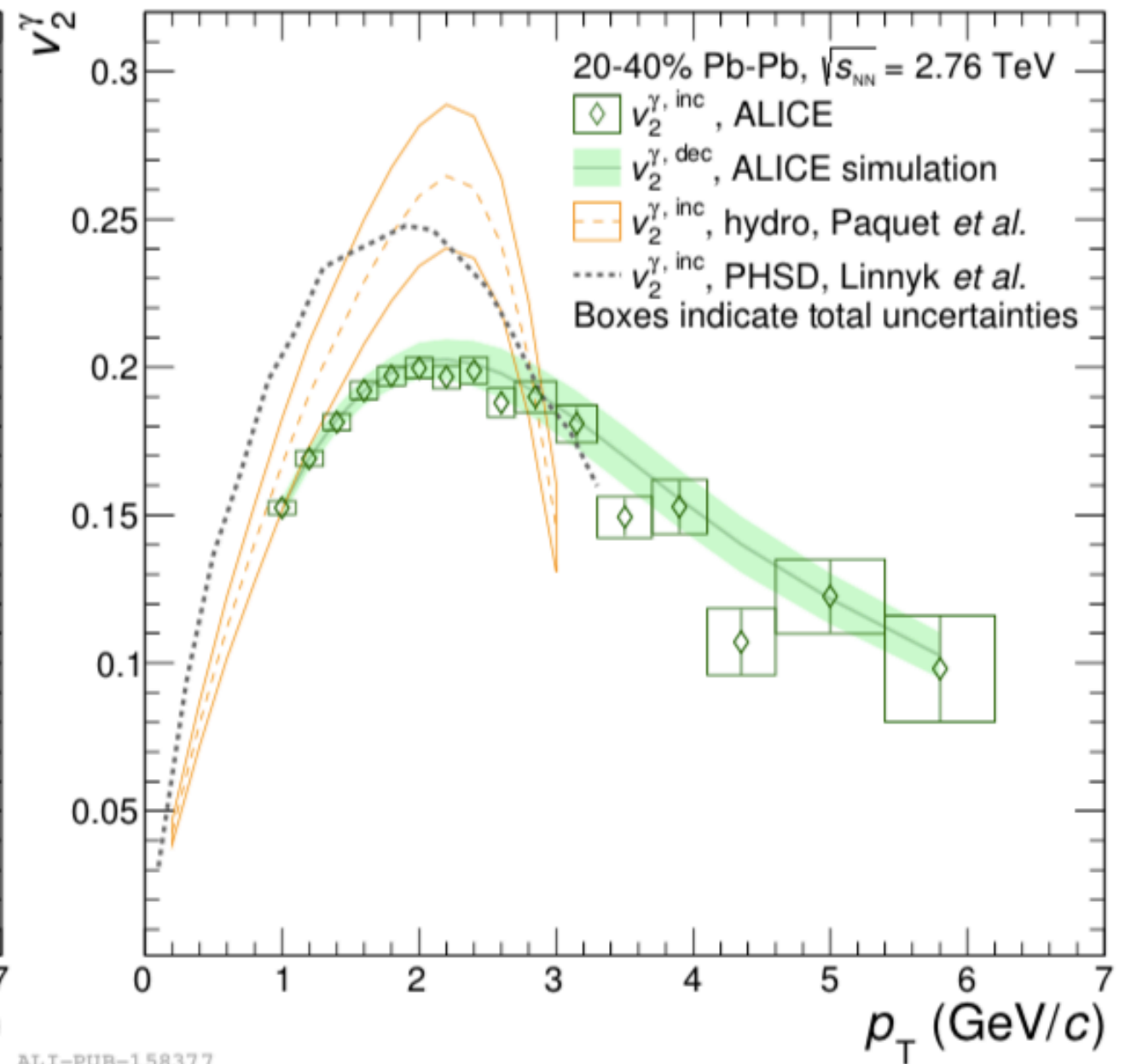
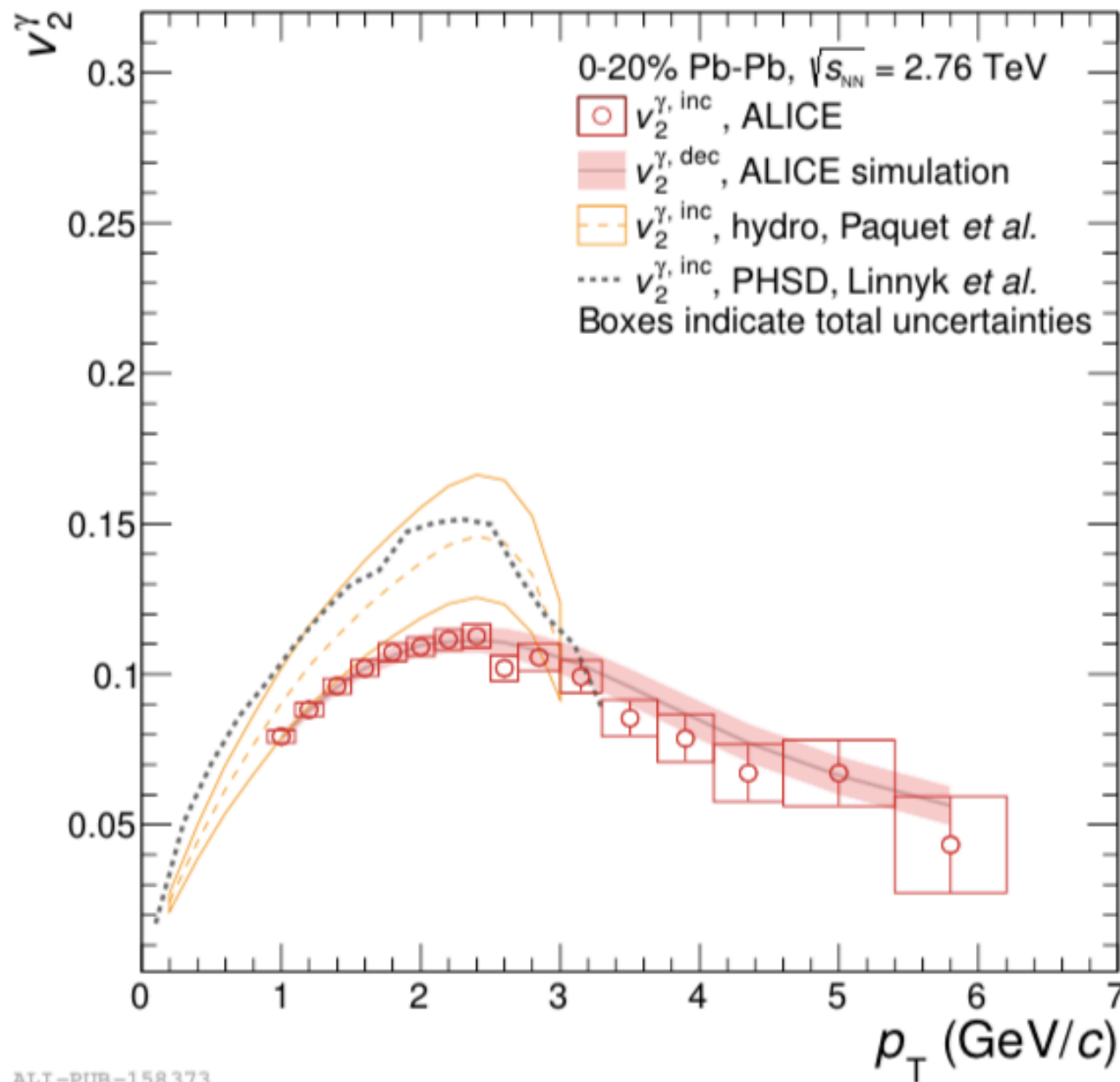


Statistical uncertainty starts to be dominant for  $p_T > 2$  GeV/c



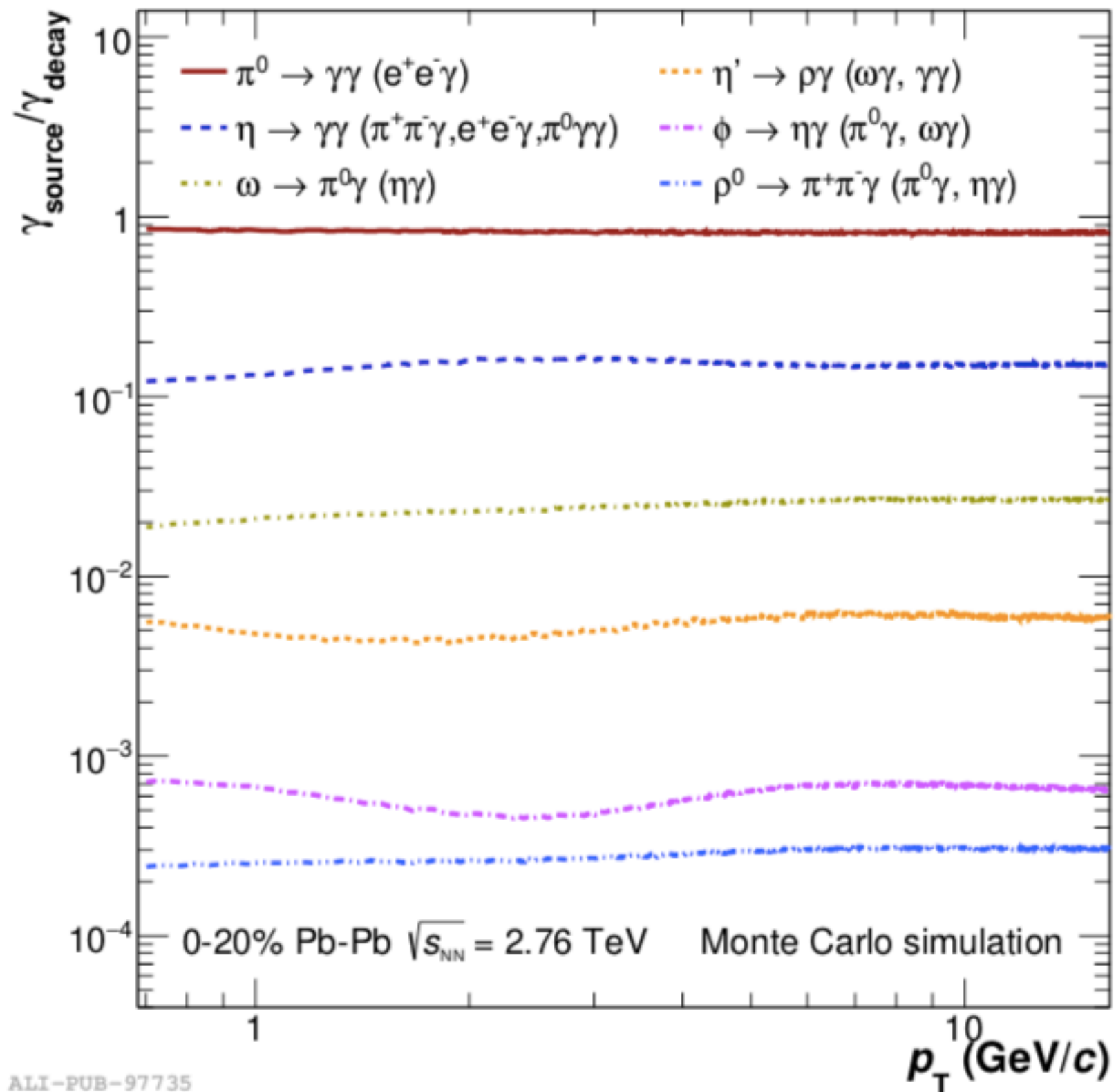
# Inclusive Photons: $v_{2,inc} \approx v_{2,dec}$

ALICE arXiv:1805.04403



- Either  $R_\gamma$  small or  $v_{2,dir} \approx v_{2,inc} \approx v_{2,dec}$
- Model comparison: important to compare inclusive *and* direct photon  $v_2$

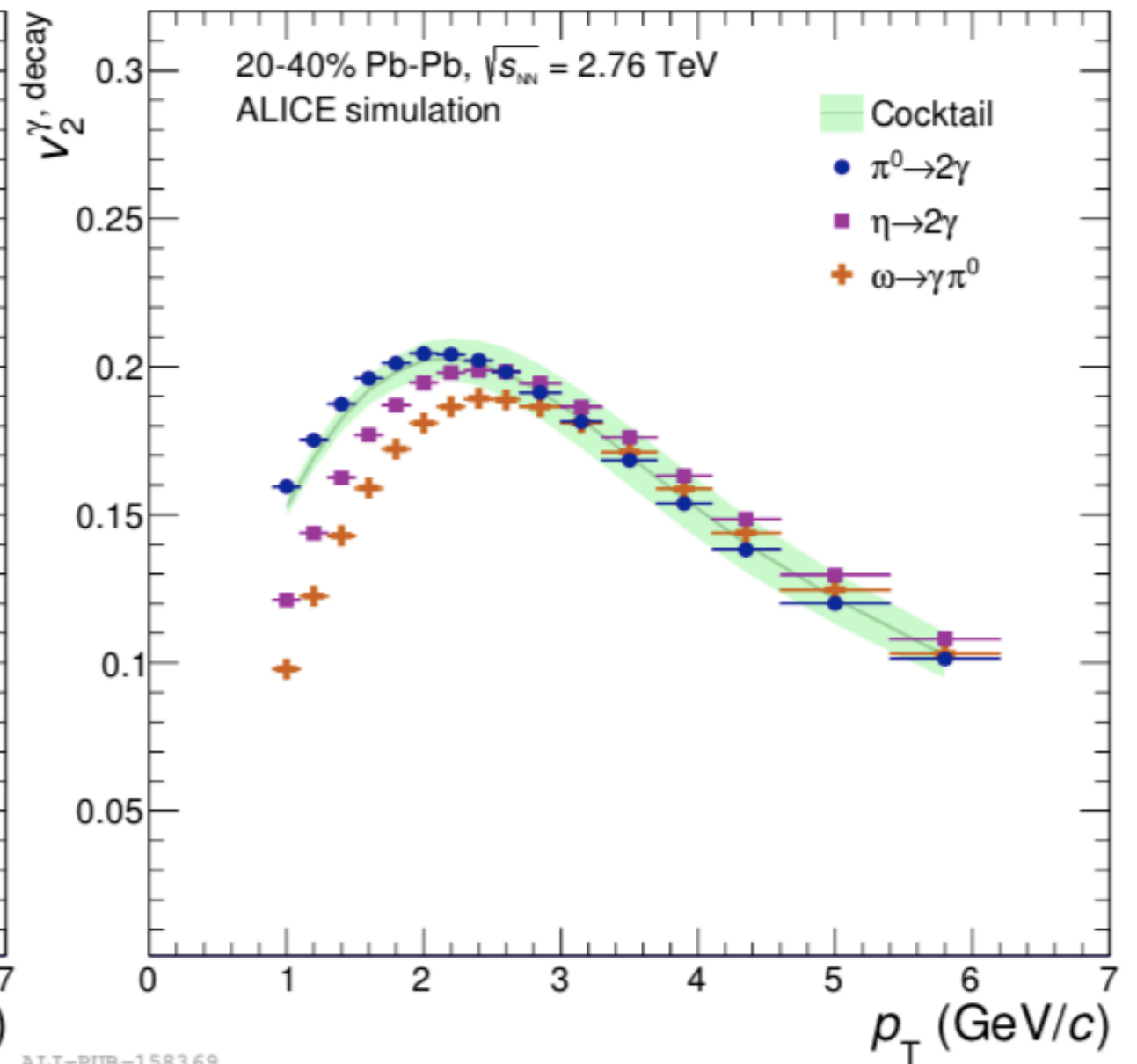
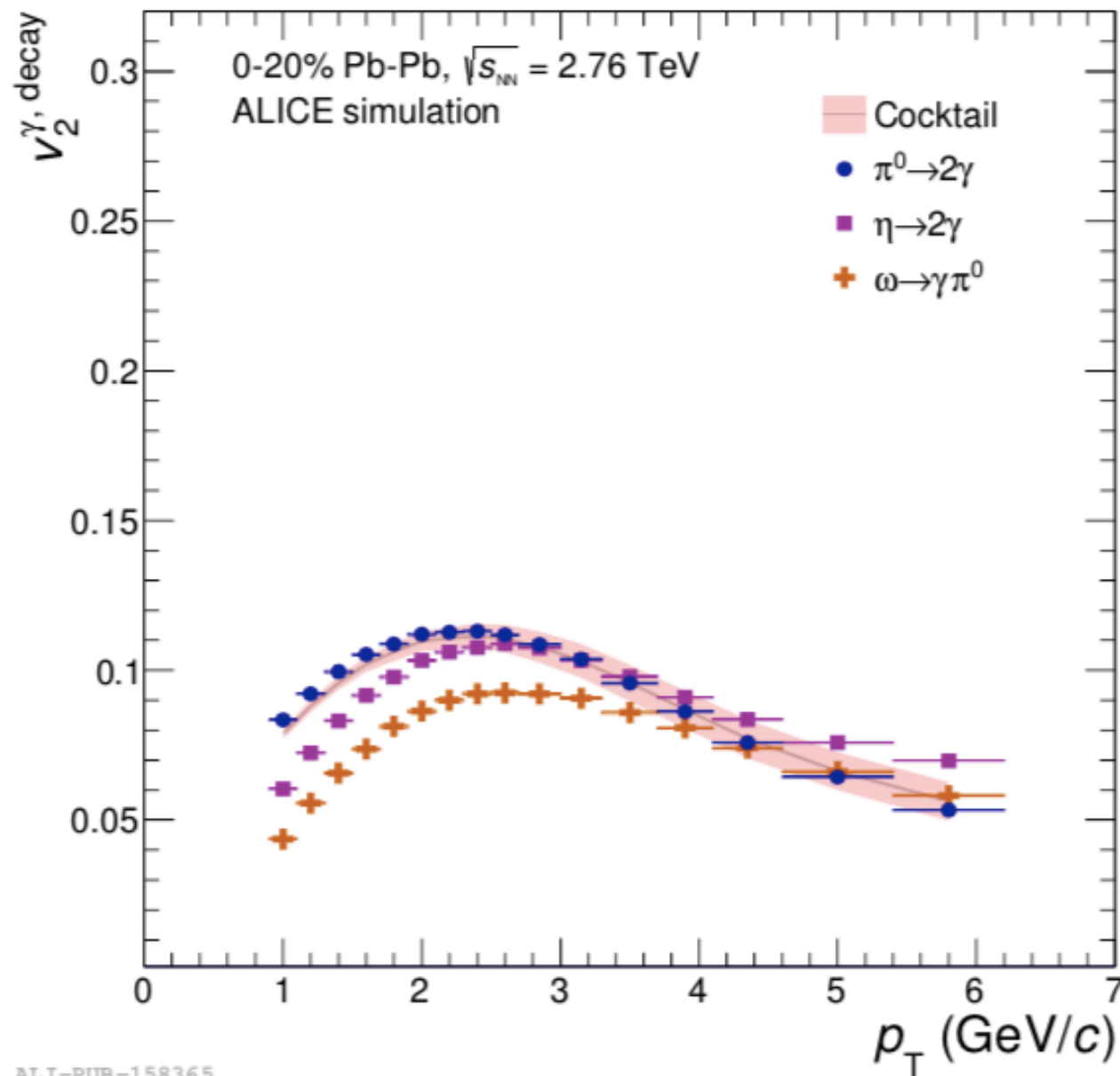
# Decay Photon Cocktail: $\pi^0$ , $\eta$ , $\omega$ relevant



ALI-PUB-97735

# Calculated Decay Photon $v_2$

ALICE arXiv:1805.04403

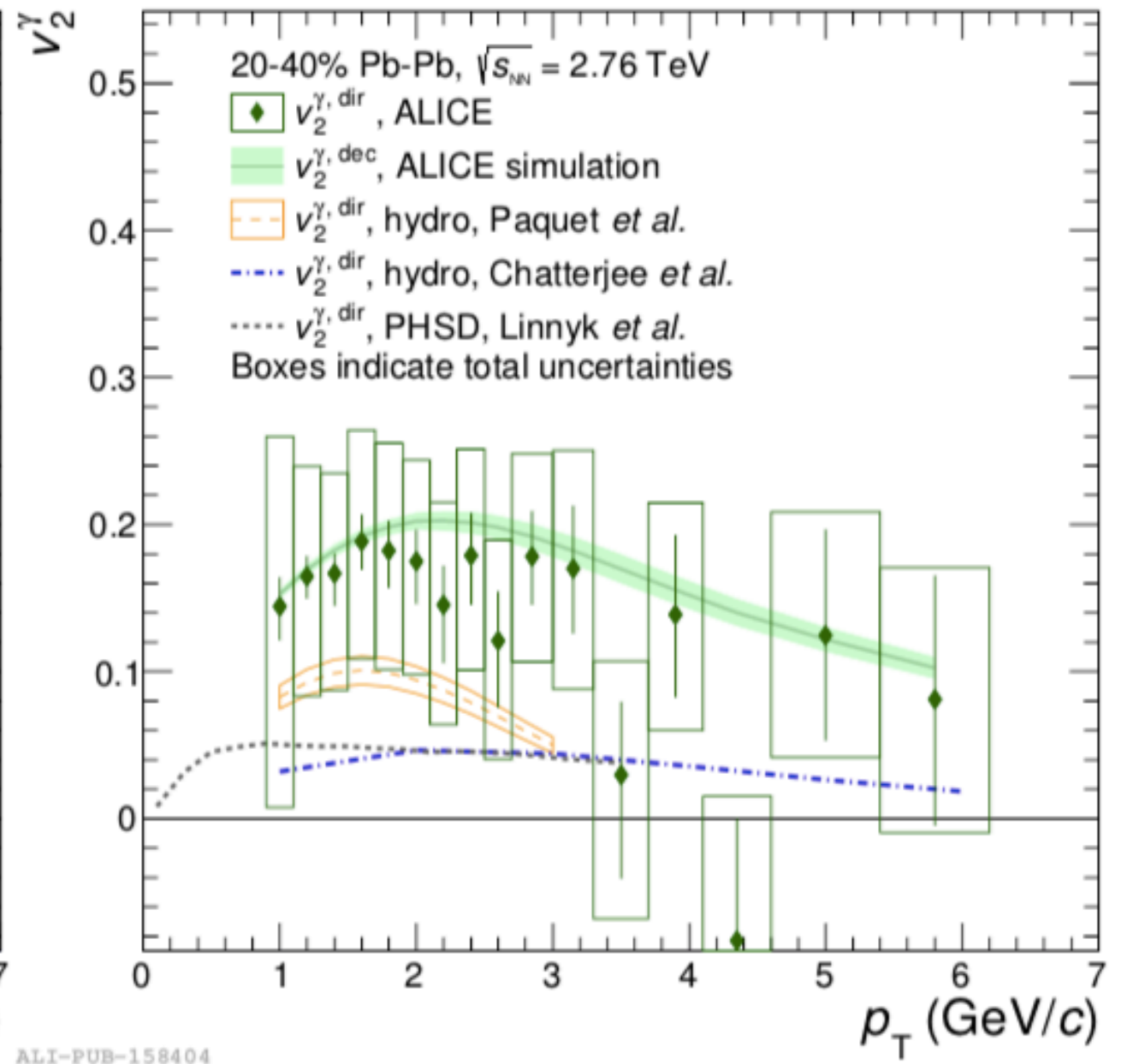
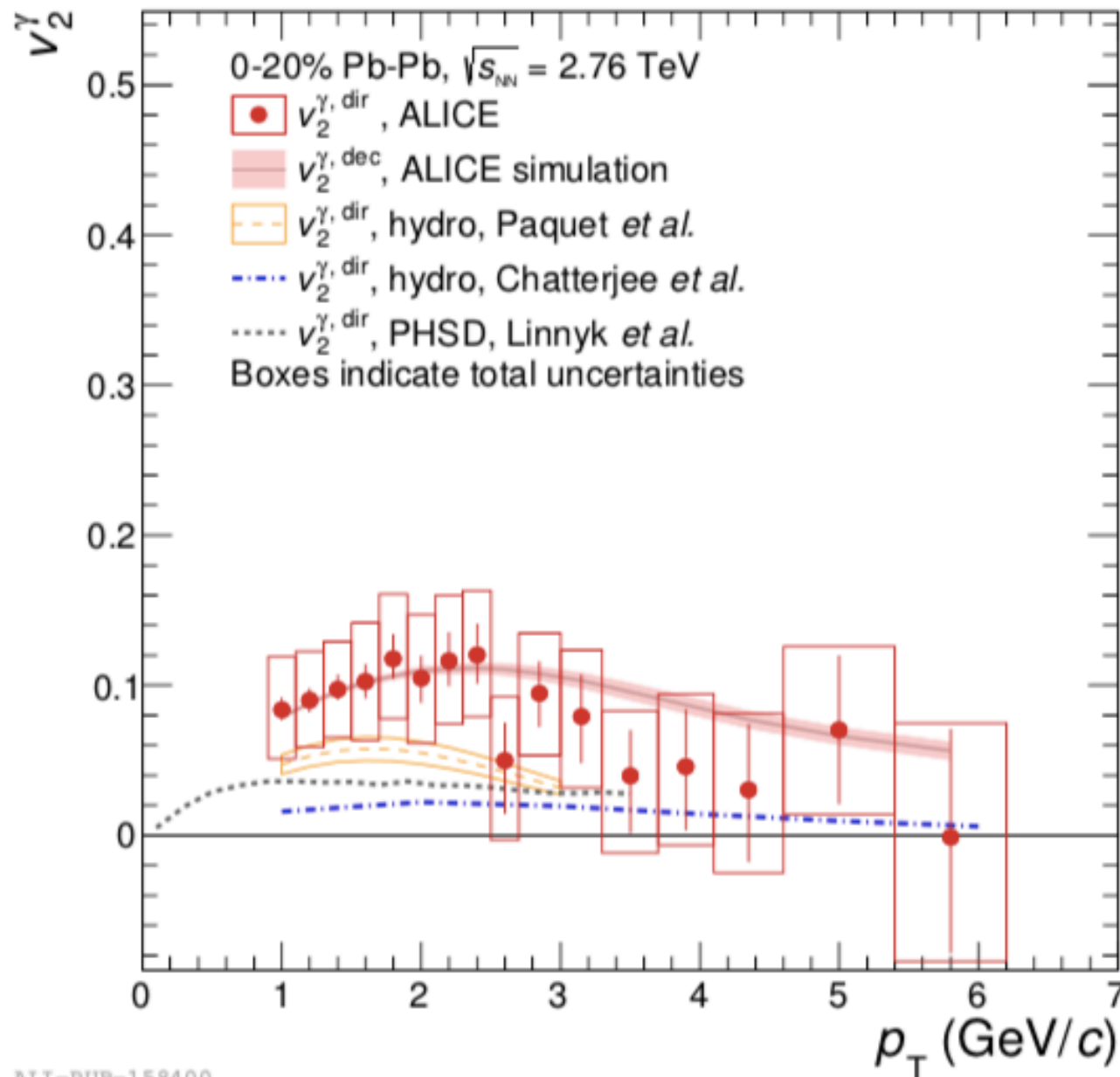


- Input: Measured  $\pi^\pm, K^{\pm,0} v_2$
- Scaling of  $v_2$  in transverse kinetic energy  $m_T - m$  for  $\eta$  and  $\omega$

# Large Direct Photon $v_2$ (but take error bars seriously)

ALICE arXiv:1805.04403

error bars = statistical uncertainties, boxes = total uncertainties

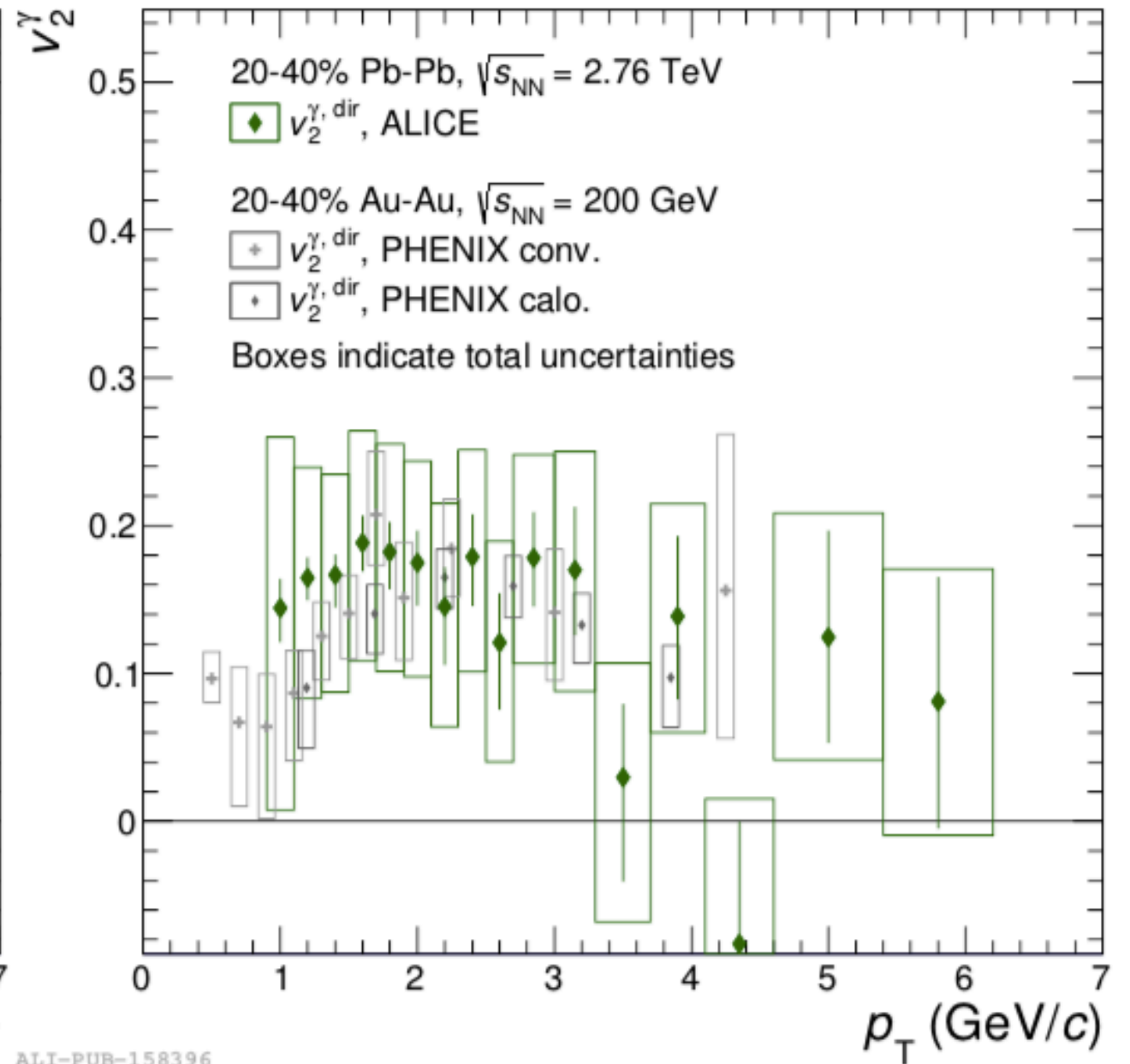
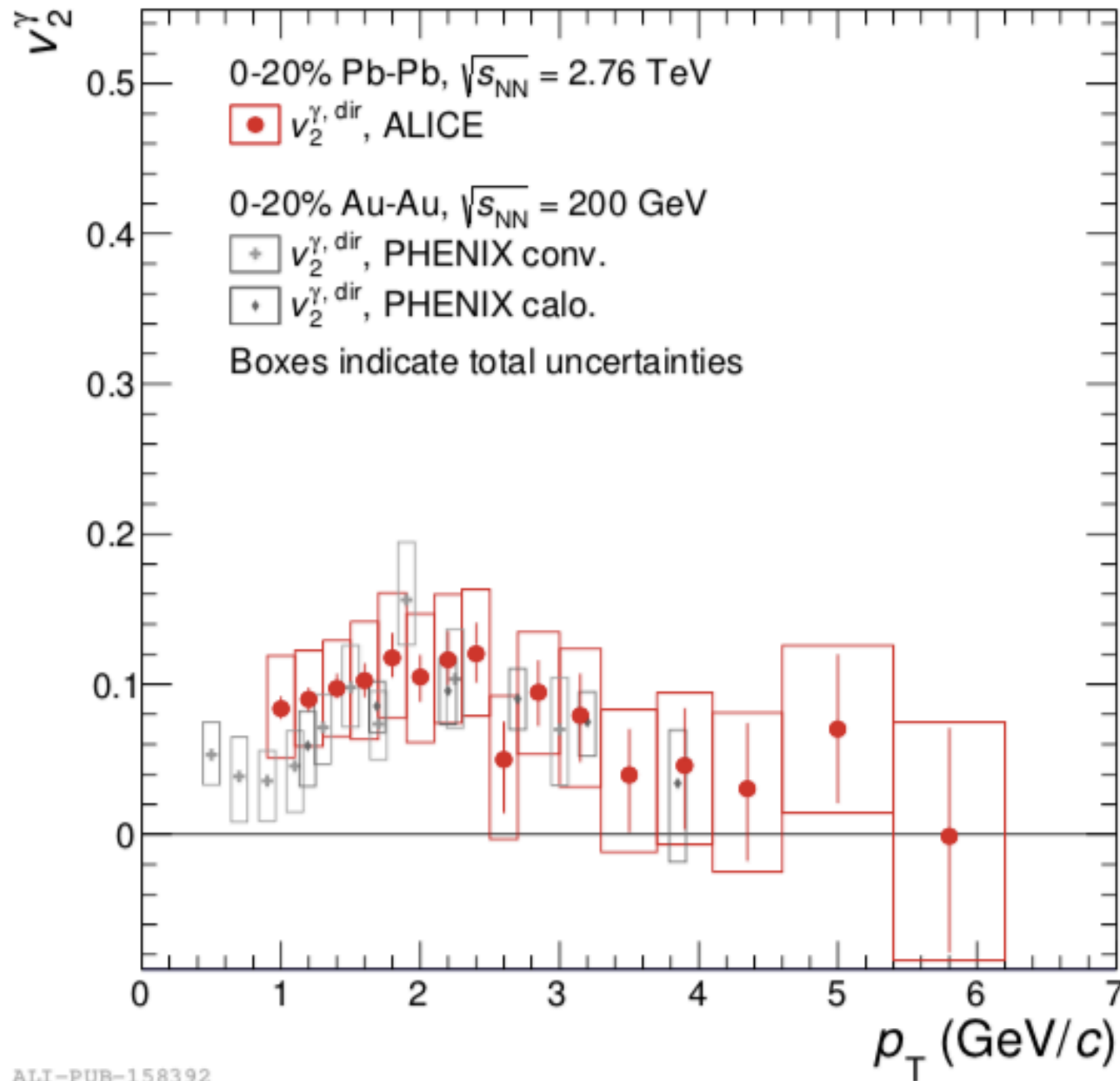


- $v_{2,dir}$  larger than models predictions (in qualitative agreement with PHENIX)
- But: null hypothesis  $v_{2,dir} = 0$  not inconsistent with the data



$$V_{2,\text{dir}}(\text{LHC}) \approx V_{2,\text{dir}}(\text{RHIC})$$

ALICE arXiv:1805.04403



# Statistical Methods

# Correlated Systematic Uncertainties

"PHENIX system":

Uncertainties categorized as

$\sigma_{i,\text{stat}}, \sigma_{i,A}, \sigma_{i,B}, \sigma_{C,\text{rel}},$

A - point-by-point uncorrelated

B - correlated, size of relative error varies point-by-point

C - constant fractional error

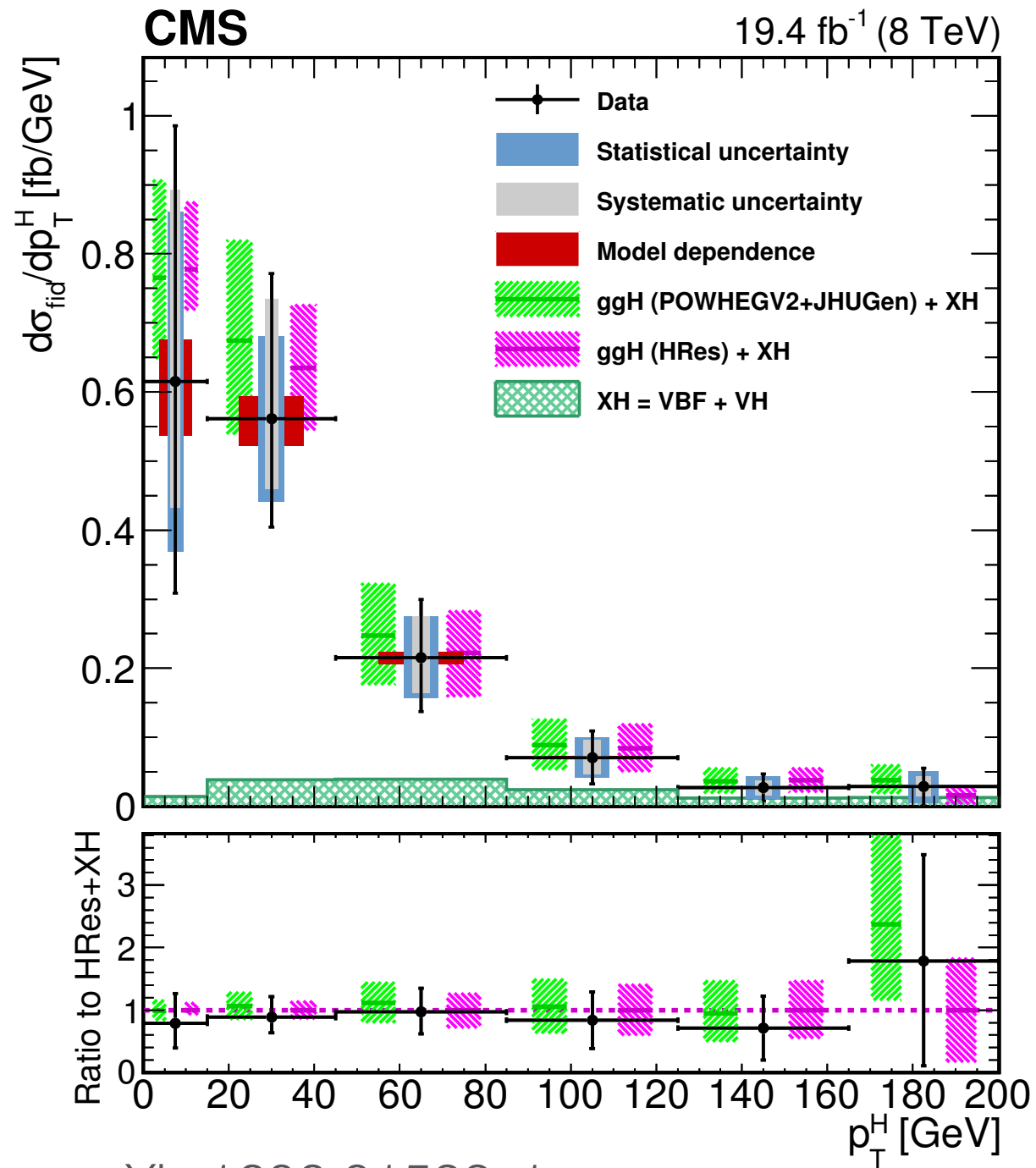
$$\chi^2 = \sum_{i=1}^n \frac{(\tilde{y}_i - \mu_i)^2}{\tilde{\sigma}_i^2} + \varepsilon_B^2 + \varepsilon_C^2$$

$$\tilde{y}_i = y_i + \varepsilon_B \sigma_{B,i} + \varepsilon_C \sigma_{C,\text{rel}} y_i$$

$$\tilde{\sigma}_i^2 = \frac{\sigma_{i,\text{stat}}^2 + \sigma_{i,A}^2}{y_i^2} \cdot \tilde{y}_i^2$$

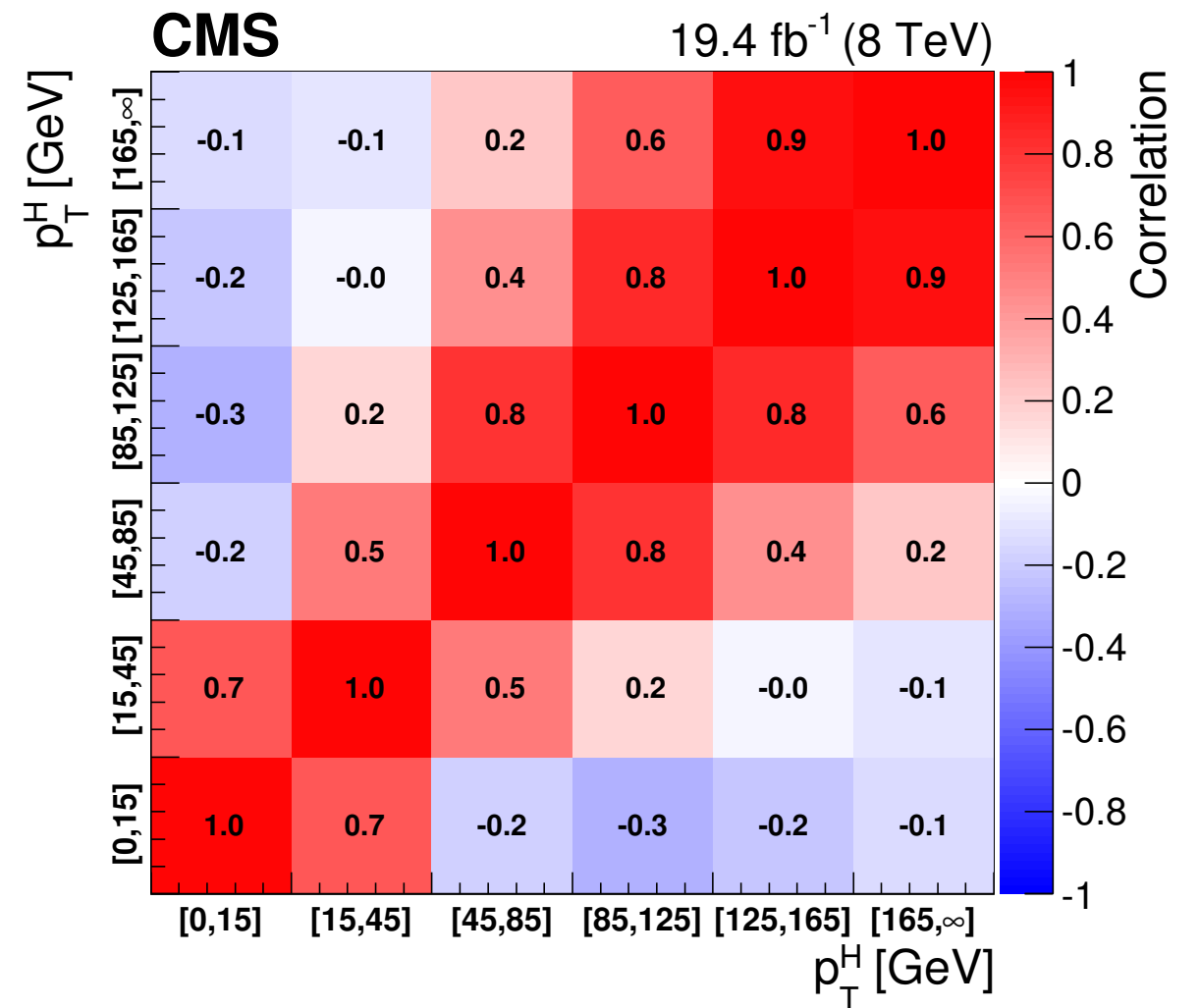
Quite useful, but often more flexibility needed

# Correlated Uncertainties: Covariance Matrix



arXiv:1606.01522v1

Correlation matrix  $\rho_{ij}$  of the  $p_T$  bins:



Covariance matrix:

$$V_{ij} = \begin{cases} \sigma_i^2 & i = j \\ \rho_{ij}\sigma_i\sigma_j & i \neq j \end{cases}$$



# Significances from Covariance Matrix

My preference: full covariance matrix  $V$

$$\chi^2 = (\underbrace{\vec{x}}_{\text{data}} - \underbrace{\vec{\mu}}_{\text{model prediction}})^T V^{-1} (\underbrace{\vec{x}}_{\text{data}} - \underbrace{\vec{\mu}}_{\text{model prediction}})$$

Translating statistical and type A, B, C systematic uncertainties into a covariance matrix:

$$V_{i,j} = V_{i,j}^{\text{stat}} + V_{i,j}^{\text{A}} + V_{i,j}^{\text{B}} + V_{i,j}^{\text{C}} = \begin{cases} \sigma_{i,\text{stat}}^2 + \sigma_{i,\text{A}}^2 + \sigma_{i,\text{B}}^2 + y_i^2 \sigma_{\text{C,rel}}^2, & \text{if } i = j. \\ \sigma_{i,\text{B}} \sigma_{j,\text{B}} + y_i y_j \sigma_{\text{C,rel}}^2, & \text{if } i \neq j. \end{cases}$$

The other way round (e.g. for plotting purposes) is not straightforward

→ That's why ALICE error boxes = total errors

# Combining PCM and PHOS $v_{2,inc}$

Correlated uncertainties in  $p_T$  for PCM and PHOS, but no correlation between PCM and PHOS;

$$\vec{v}_2^{\gamma,inc} = (V_{v_2,PCM}^{-1} + V_{v_2,PHOS}^{-1})^{-1} (V_{v_2,PCM}^{-1} \vec{v}_2^{\gamma,inc,PCM} + V_{v_2,PHOS}^{-1} \vec{v}_2^{\gamma,inc,PHOS})$$

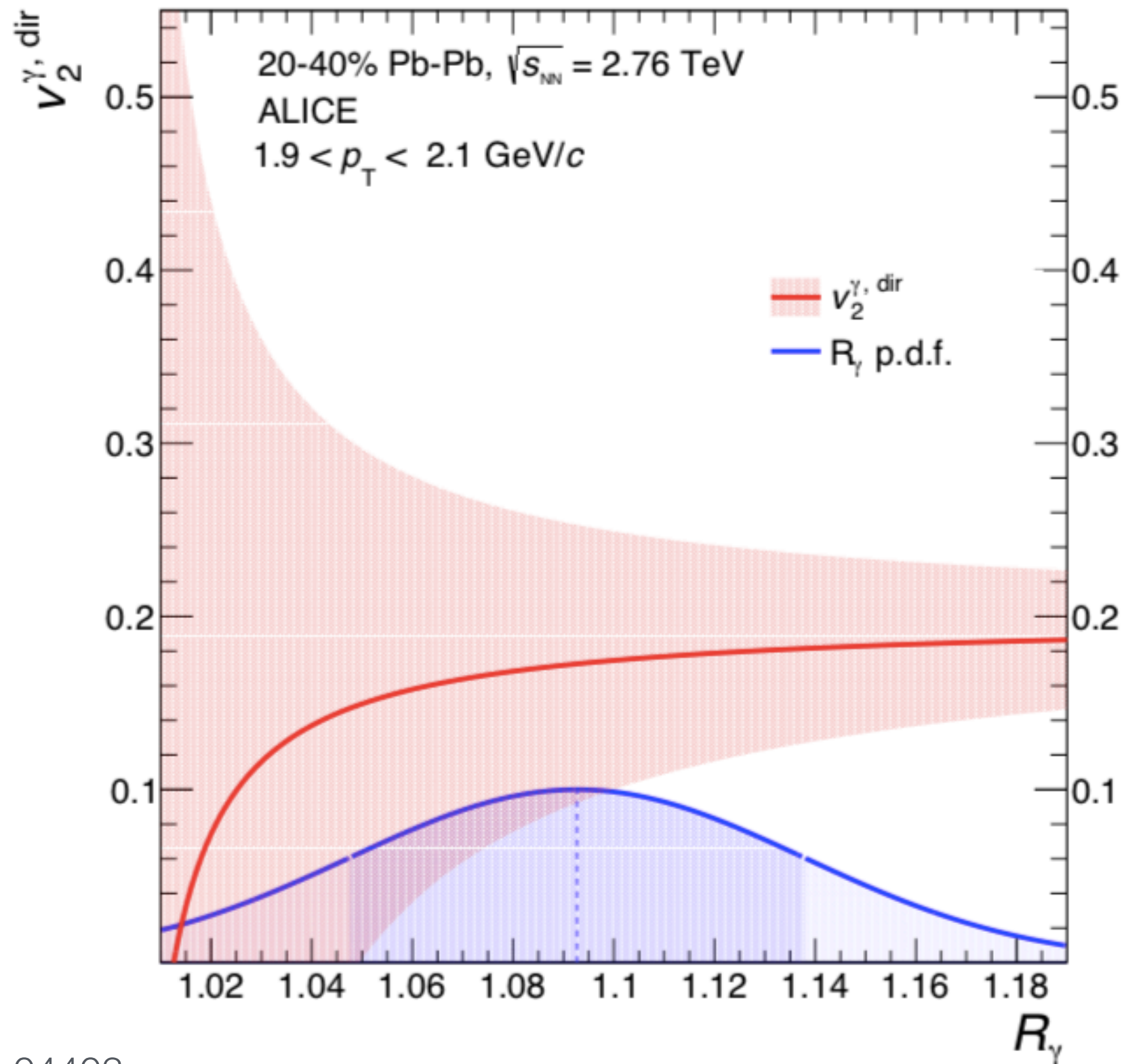
$$V_{ij} = V_{stat,ij} + V_{syst,ij}, \text{ where } V_{syst,ij} = \rho \sigma_{syst,i} \sigma_{syst,j}$$

(linear combination with the smallest  $\chi^2$ )

Covariance matrices from estimating the correlations coefficient.

Better, but more work: covariance matrices from toy Monte Carlo studies

# Error Propagation: Small measured $R_\gamma$ Requires Special Consideration

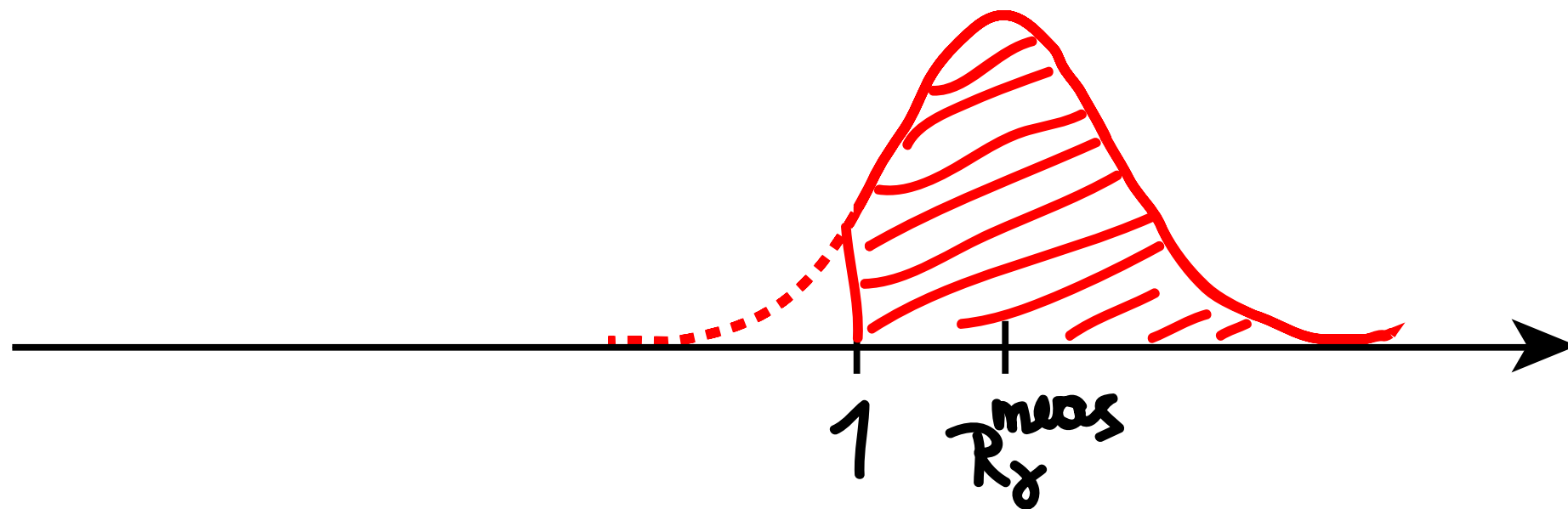


# Bayesian Approach

Bayes:

$$P(R_{\gamma,t}|R_{\gamma,m}) \propto L(R_{\gamma,m}|R_{\gamma,t})\pi(R_{\gamma,t}), \quad \pi(R_{\gamma,t}) = \Theta(R_{\gamma,t} - 1)$$

Posterior distribution of  $R_{\gamma}$ :



## $v_{2,\text{dir}}$ Calculation: One $p_T$ bin

- The formula for  $v_{2,\text{dir}}$  is only defined for  $R_\gamma > 1$ : 
$$v_{2,\text{dir}} = \frac{R_\gamma v_{2,\text{incl}} - v_{2,\text{decay}}}{R_\gamma - 1}$$
- The measured value of  $R_\gamma$  can be below unity. Same problem in MC error propagation for  $R_\gamma$  slightly larger than 1. How to handle this case?

- Bayes: 
$$P(\vec{\theta}|\vec{m}) \propto L(\vec{m}|\vec{\theta}) \pi(\vec{\theta})$$

- In our case: 
$$\vec{\theta} = (v_{2,\text{inc,t}}, v_{2,\text{dec,t}}, R_{\gamma,t}) \quad \vec{m} = (v_{2,\text{inc,m}}, v_{2,\text{dec,m}}, R_{\gamma,m})$$

$$\pi(R_{\gamma,t}) = \Theta(R_{\gamma,t} - 1)$$

$$L(\vec{m}|\vec{\theta}) = L(v_{2,\text{inc,m}}|v_{2,\text{inc,t}}) L(v_{2,\text{dec,m}}|v_{2,\text{dec,t}}) L(R_{\gamma,m}|R_{\gamma,t})$$

Likelihoods modeled as Gaussians, e.g.:

$$L(R_{\gamma,m}|R_{\gamma,t}) = G(R_{\gamma,m}; R_{\gamma,t}, \sigma(R_{\gamma,t}))$$

- We obtain a posterior distribution for  $v_{2,\text{inc,t}}$ ,  $v_{2,\text{dec,t}}$ , and  $R_{\gamma,t}$  from which we obtain a posterior distribution for  $v_{2,\text{dir,t}}$

# $v_{2,\text{dir}}$ Calculation: Several $p_T$ Bins

$$\vec{m} = (\vec{R}_{\gamma,m}, \vec{v}_2^{\gamma,\text{inc},m}, \vec{v}_2^{\gamma,\text{dec},m}) \quad \vec{\vartheta} = (\vec{R}_{\gamma,t}, \vec{v}_2^{\gamma,\text{inc},t}, \vec{v}_2^{\gamma,\text{dec},t})$$

$$P(\vec{\vartheta}|\vec{m}) \propto P(\vec{m}|\vec{\vartheta})\pi(\vec{\vartheta})$$

$$\pi(\vec{\vartheta}) \equiv \pi(\vec{R}_{\gamma,t}) = \Theta(R_{\gamma,t,1} - 1, \dots, R_{\gamma,t,n} - 1)$$

$$P(\vec{m}|\vec{\vartheta}) = \prod_{x=R_{\gamma}, v_2^{\gamma,\text{inc}}, v_2^{\gamma,\text{dec}}} G(\vec{x}_m; \vec{x}_t, V_x)$$

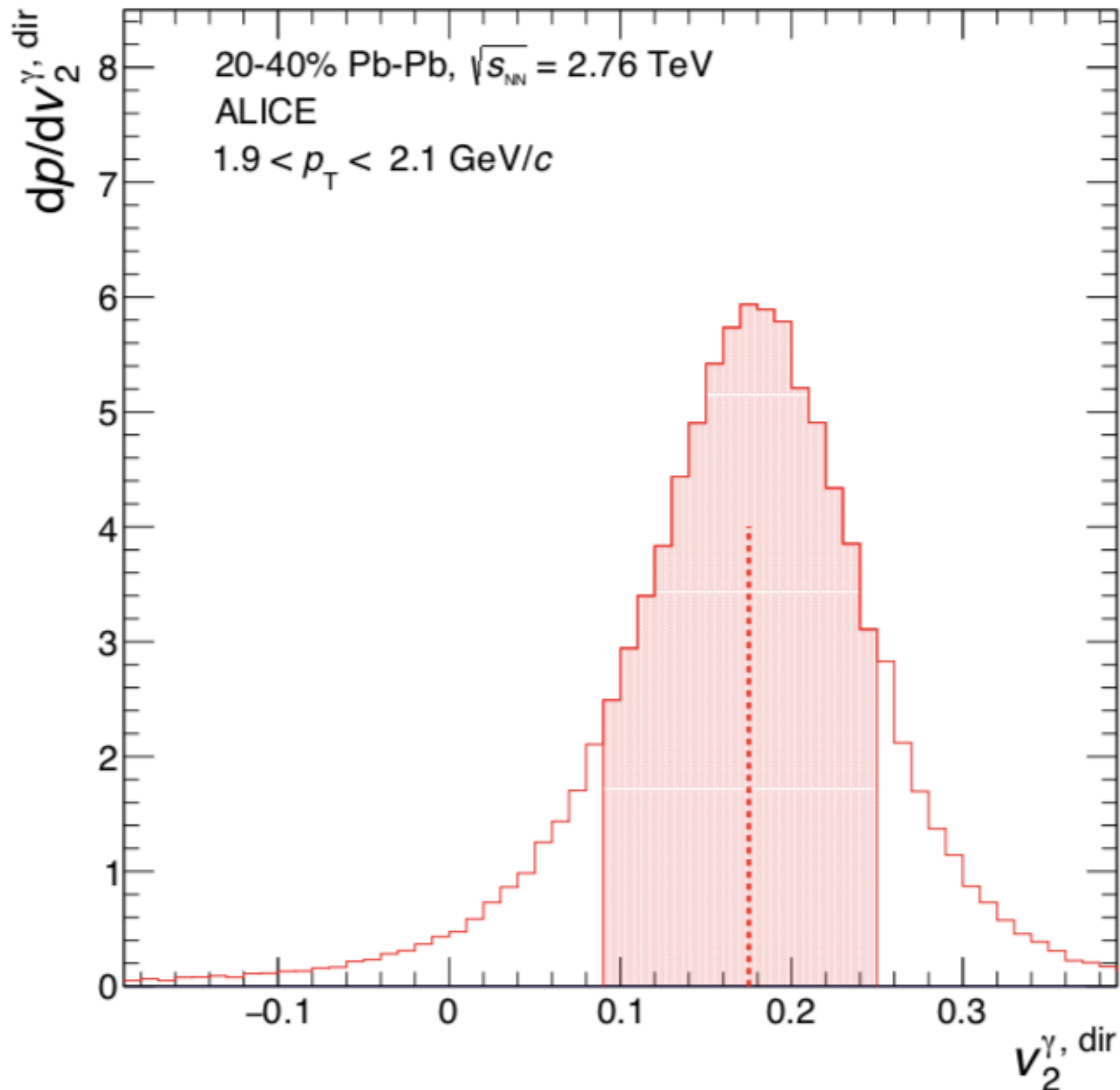
likelihood

multivariate Gaussian

covariance matrix



# $V_{2,dir}$ Posterior Distribution for One $p_T$ Bin



Towards quantifying the statistical significance of the  
direct-photon puzzle

# Hypothesis Testing / Significance

Q: Significance of the difference between data and a certain  $v_{2,\text{dir}}$  hypothesis?

Need a test statistic  $t$  that quantifies deviation/difference between a hypothesis and the data.

Often used test statistics: likelihood  $L(x; \theta)$  and  $\chi^2$

$$p\text{-value} = \int_{t_{\text{obs}}}^{\infty} p(t|H) dt \quad (\text{frequentist concept})$$

One way to handle nuisance parameters: marginal likelihoods

integrate over nuisance  
parameter  $v$ :

prior knowledge about  $v$ :

$$t := L_m(x|s) = \int L(x|\nu, s) \pi'_\nu(\nu) d\nu$$

Further information: lecture G. Cowan (especially p. 88ff.):

- <http://www.ippp.dur.ac.uk/~ross/invisibles13/school/talks/GlenCowanStatisticalandDataAnalysis.pdf>
- 1307.2487 (section 5.2)

# Likelihood Function (One $p_T$ bin)

Likelihood function (note that we use  $v_{2,\text{dir}}^t$  as parameter):

$$L(v_{2,\text{inc}}^m, v_{2,\text{dec}}^m, R_\gamma^m; v_{2,\text{dir}}^t, v_{2,\text{dec}}^t, R_\gamma^t) = G(v_{2,\text{inc}}^m; v_{2,\text{inc}}^t, \sigma_{v_{2,\text{inc}}}) \cdot G(v_{2,\text{dec}}^m; v_{2,\text{dec}}^t, \sigma_{v_{2,\text{dec}}}) \cdot G(R_\gamma^m; R_\gamma^t, \sigma_{R_\gamma})$$

$G(x; \mu, \sigma)$ : 1d Gaussians

Here  $v_{2,\text{inc}}^t$  is a function of the other three parameters:

$$v_{2,\text{inc}}^t \equiv v_{2,\text{inc}}^t(v_{2,\text{dir}}^t, v_{2,\text{dec}}^t, R_\gamma^t) = \frac{(R_\gamma^t - 1)v_{2,\text{dir}}^t + v_{2,\text{dec}}^t}{R_\gamma^t}$$

# Marginalized Likelihood (I)

Treat  $v_{2,\text{dec}}^t$  and  $R_\gamma^t$  as nuisance parameters:

$$L_m(v_{2,\text{inc}}^m | v_{2,\text{dir}}^t) = \int dv_{2,\text{dec}}^t dR_\gamma^t L(v_{2,\text{inc}}^m | v_{2,\text{dir}}^t, v_{2,\text{dec}}^t, R_\gamma^t) \pi(v_{2,\text{dec}}^t) \pi(R_\gamma^t)$$

$$\pi(v_{2,\text{dec}}^t) \propto G(v_{2,\text{dec}}^t; v_{2,\text{dec}}^m, \sigma_{v_{2,\text{dec}}}), \quad \pi(R_\gamma^t) \propto G(R_\gamma^t; R_\gamma^m, \sigma_{R_\gamma}) \theta(R_\gamma^t - 1)$$

This gives:

$$v_{2,\text{inc}}^t \equiv v_{2,\text{inc}}^t(v_{2,\text{dir}}^t, v_{2,\text{dec}}^t, R_\gamma^t) = \frac{(R_\gamma^t - 1)v_{2,\text{dir}}^t + v_{2,\text{dec}}^t}{R_\gamma^t}$$

$$L(v_{2,\text{inc}}^m | v_{2,\text{dir}}^t) = \int_{-0.5}^{0.5} dv_{2,\text{dec}}^t \int_1^\infty dR_\gamma^t G(v_{2,\text{inc}}^m; v_{2,\text{inc}}^t, \sigma_{v_{2,\text{inc}}}) \cdot G(v_{2,\text{dec}}^t; v_{2,\text{dec}}^m, \sigma_{v_{2,\text{dec}}}) \cdot G(R_\gamma^t; R_\gamma^m, \sigma_{R_\gamma})$$

# Marginalized Likelihood (II)

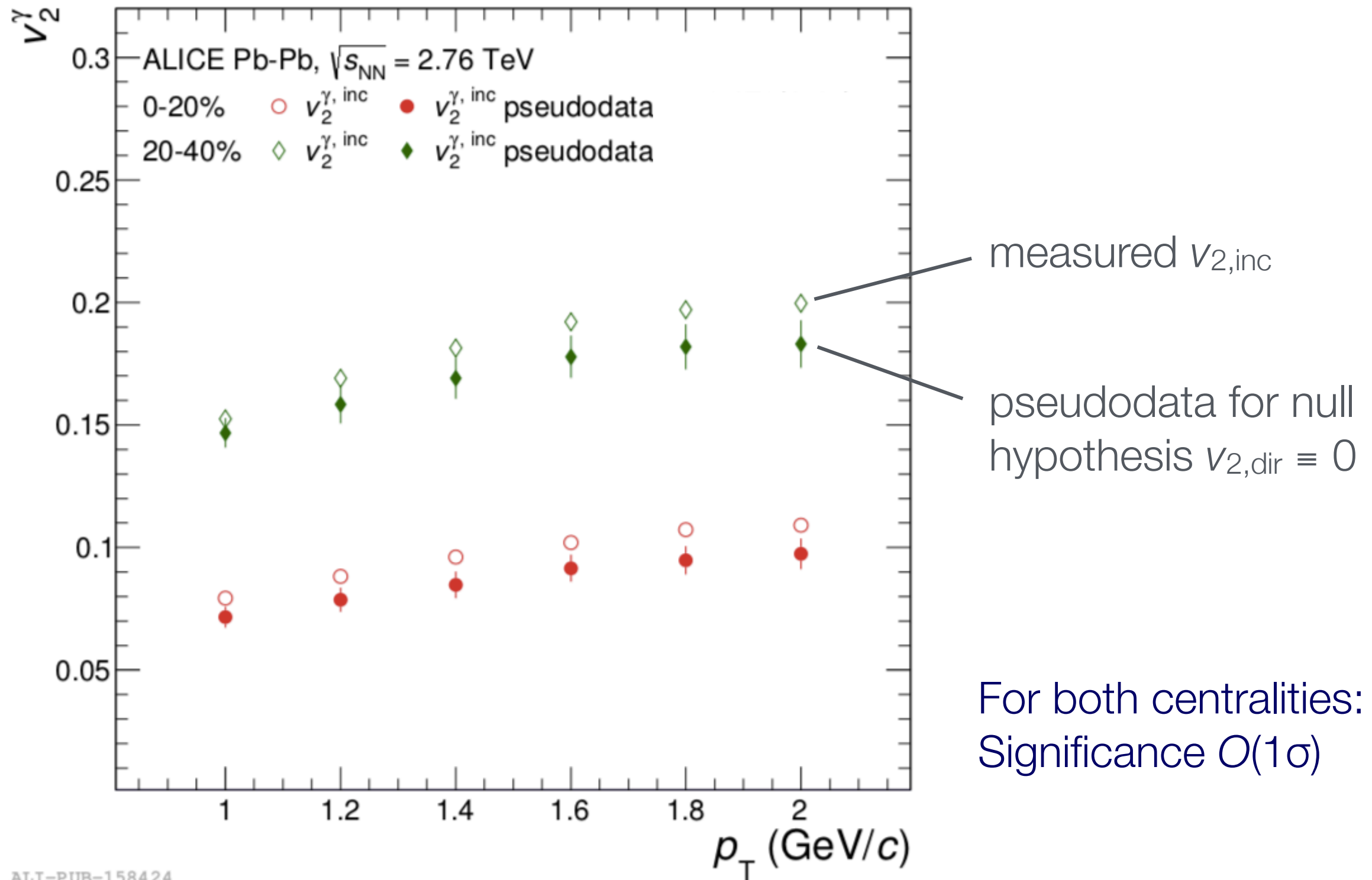
From previous slide:

$$L(v_{2,\text{inc}}^m | v_{2,\text{dir}}^t) = \int_{-0.5}^{0.5} dv_{2,\text{dec}}^t \int_1^{\infty} dR_{\gamma}^t G(v_{2,\text{inc}}^m; v_{2,\text{inc}}^t, \sigma_{v_{2,\text{inc}}}) \cdot G(v_{2,\text{dec}}^t; v_{2,\text{dec}}^m, \sigma_{v_{2,\text{dec}}}) \cdot G(R_{\gamma}^t; R_{\gamma}^m, \sigma_{R_{\gamma}})$$

MC sampling instead of solving the integrals  $\rightsquigarrow$  histogram of  $v_{2,\text{inc}}^m$  values:

- Draw  $v_{2,\text{dec}}^t$  and  $R_{\gamma}^t$  from the corresponding Gaussians
- Calculate  $v_{2,\text{inc}}^t$  from these values and the given  $v_{2,\text{dir}}^t$  hypothesis
- Generate  $v_{2,\text{inc}}^m$  pseudo-data

# Significance of the Deviation from the Null Hypothesis $v_{2,\text{dir}} \equiv 0$





# Significance of Puzzles

Another example:

<https://home.cern/news/news/experiments/lhcb-finds-new-hints-possible-standard-model-deviations>

## LHCb finds new hints of possible Standard Model deviations

The LHCb experiment finds intriguing anomalies in the way some particles decay

18 APRIL, 2017 | By [Stefania Pandolfi](#)

"While potentially exciting, the discrepancy with the Standard Model occurs at the level of **2.2 to 2.5 sigma**, which is not yet sufficient to draw a firm conclusion."

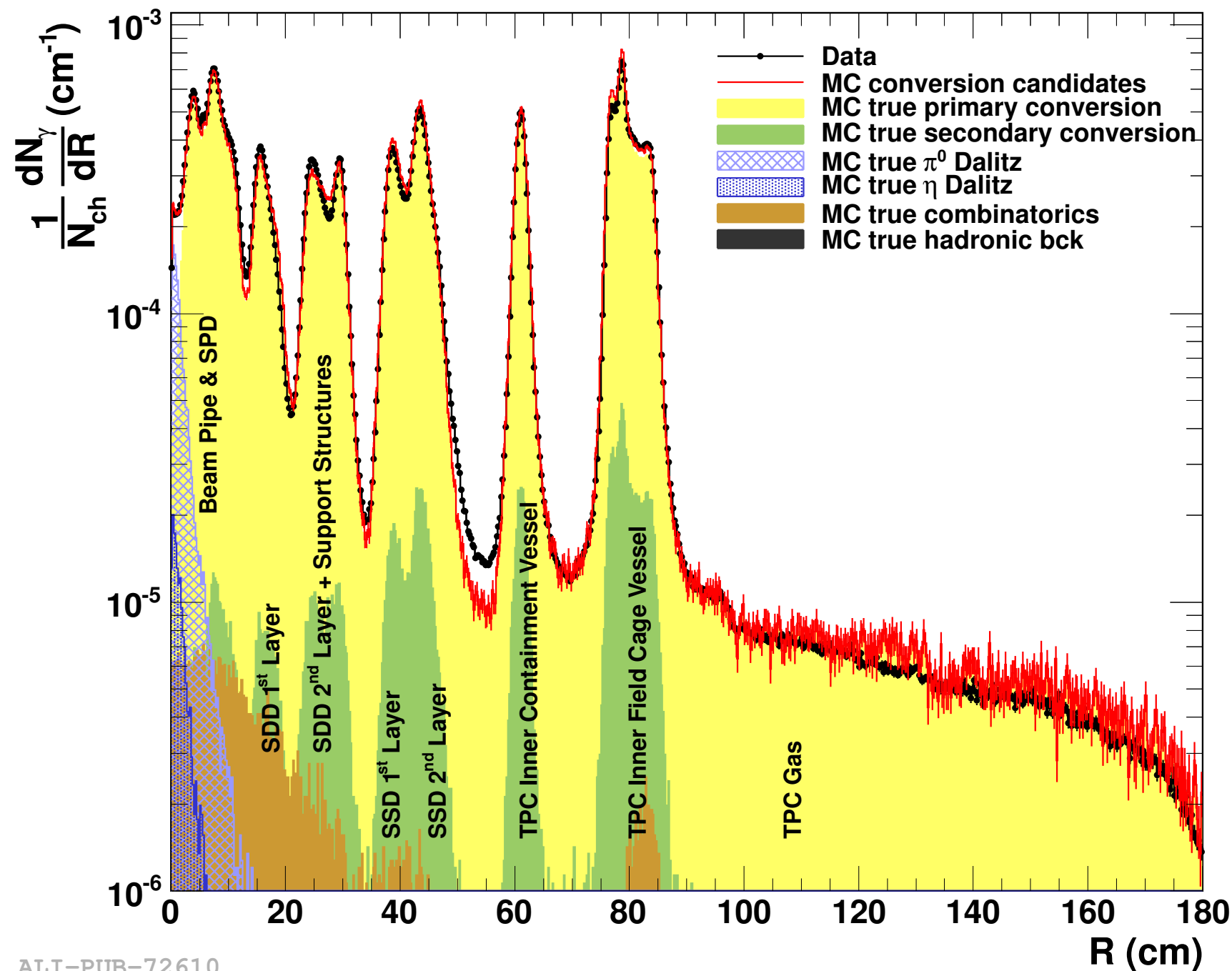
What is the statistical significance of the direct-photon puzzle?

In a joint effort, we could come up with a statement.

[reasonable timing: as spectra at  $v_2$  at RHIC and the LHC are now published]

Reducing systematic uncertainties

# Reducing the Material Budget Uncertainty in the Photon Conversion Method: Weights



$$w_i = \frac{(N_\gamma^{\text{rec}}(r_i)/N_\gamma^{\text{rec}}(r_{\text{ref}}))_{\text{data}}}{(N_\gamma^{\text{rec}}(r_i)/N_\gamma^{\text{rec}}(r_{\text{ref}}))_{\text{MC}}}$$

$$\varepsilon_\gamma^{\text{rec}}(p_T) = \frac{\sum_i w_i N_\gamma^{\text{rec}}(r_i, p_T)}{N_\gamma^{\text{prod}}(r_i, p_T)}$$

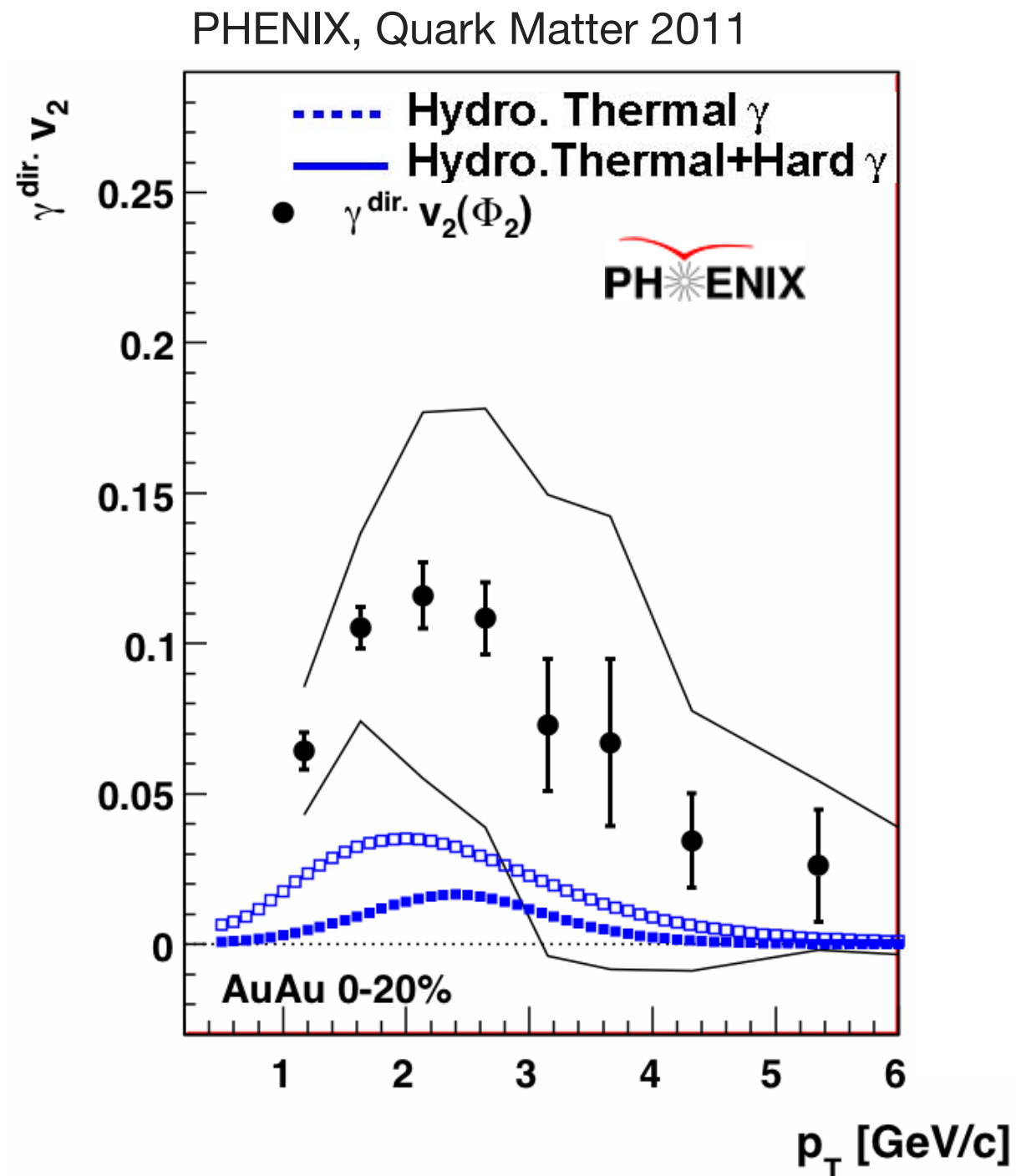
Reference material:

- ▶ TPC gas (need to keep track of pressure)
- ▶ tungsten wires (Run 3)

Might reduce current material budget uncertainty of 4.5% by more than a factor 2

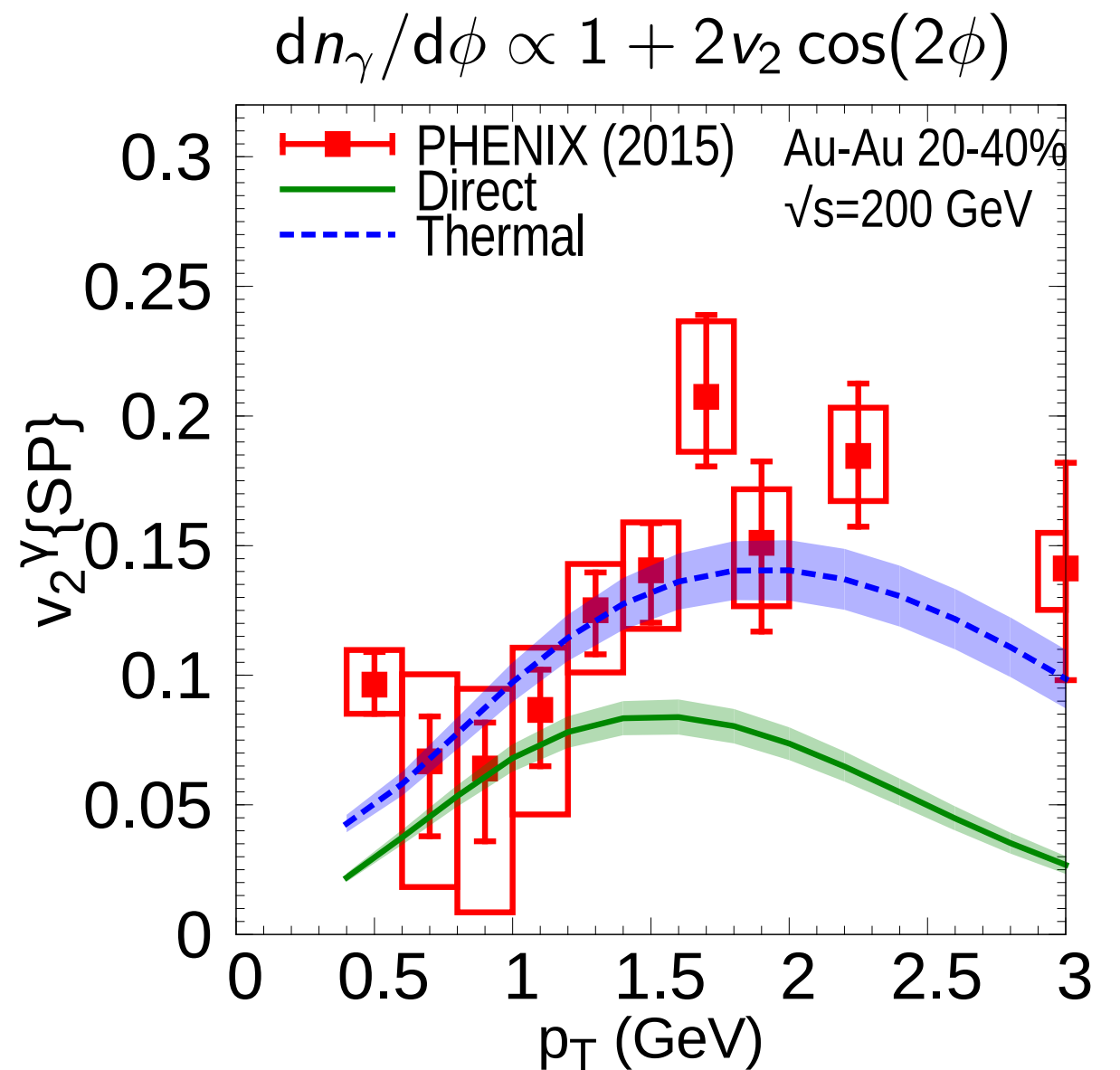
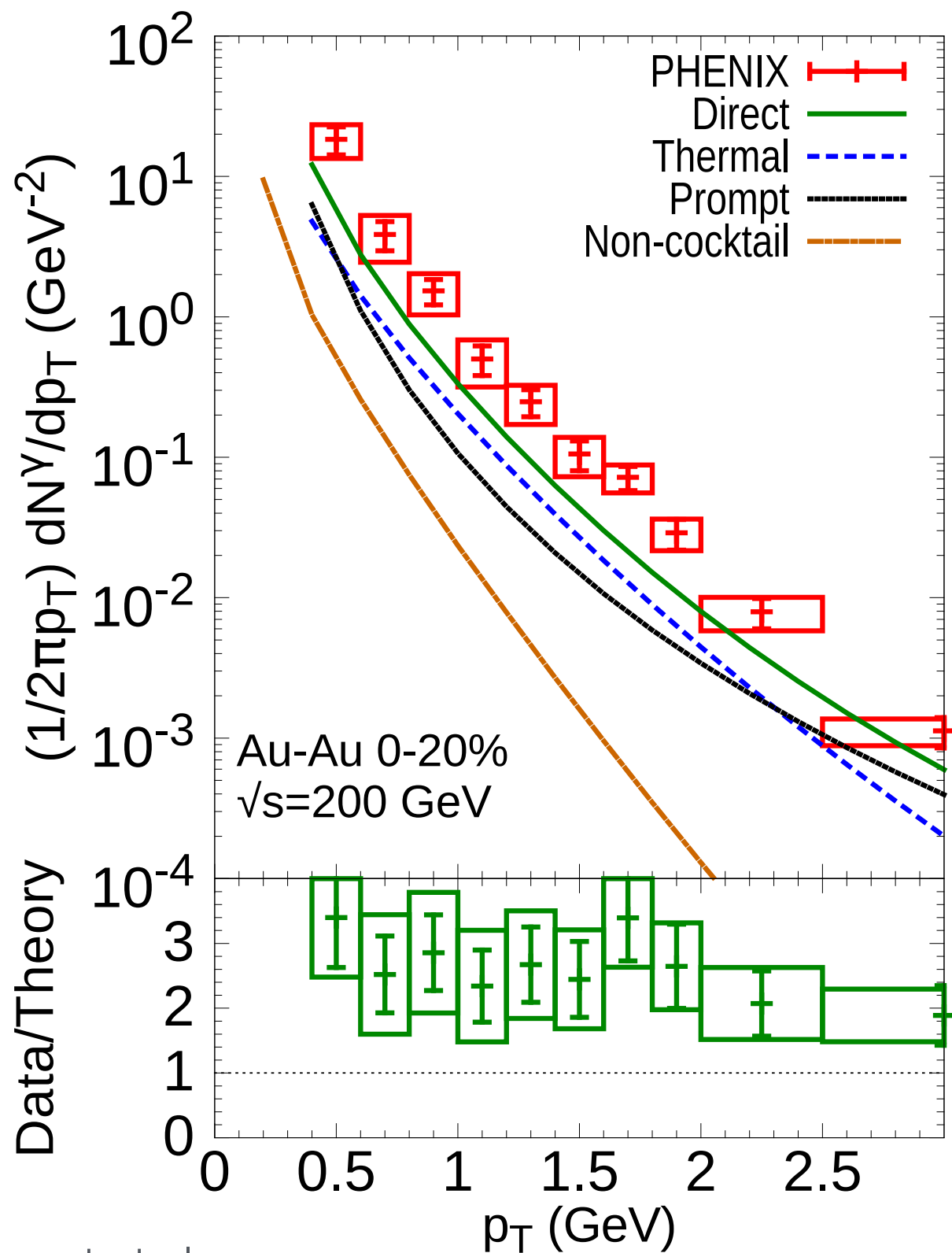
# Direct-Photon Puzzle

# Birth of the Direct Photon Puzzle



- Quark Matter 2011
- "Data a challenge to theory" or
- "Theory a challenge to the data"

# Direct-Photon Puzzle: Status



- Challenging for hydro models to describe  $v_2$  and yield
- ALICE  $\gamma_{\text{dir}}$  and  $v_2$ :  
"No puzzle with current errors"

# EMMI Rapid Reaction Task Force on the Direct Photon Flow Puzzle

- Feb. 2014, 25 participants (theory + experiment)
- Open Symposium:  
<https://indico.gsi.de/conferenceDisplay.py?confId=2662>
- Detailed discussions on
  - ▶ Averaging of  $v_n$  over large centrality bins, definition of  $v_n$  in models
  - ▶ Definition of decay photon cocktail in experiment and models, contribution from short-lived resonances
  - ▶ Comparison of the space-time evolution (hydro models, PHSD, parameterized fireball evolution)
  - ▶ pQCD contribution in various models
  - ▶ Initial flow, near  $T_c$  enhancement of photon rates, bremsstrahlung photons in the hadrons gas, Glasma photons, role of fragmentation photons, ...
- Puzzle remains after checking various aspects of the data/theory comparison

Helmholtz Alliance  
Extremes of Density and Temperature: Cosmic Matter in the Laboratory

## ExtreMe Matter Institute EMMI

EMMI Rapid Reaction Task Force  
**Direct-Photon Flow Puzzle**  
February 24-28, 2014, GSI, Darmstadt, Germany

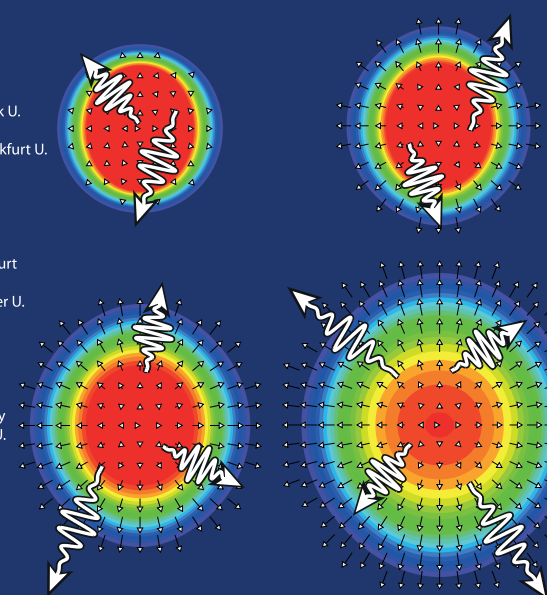
**Open Symposium**  
Monday, Feb. 24, 09:30 - 13:00h  
KBW Lecture Hall

**Participants**  
Benjamin Bannier, Stony Brook U.  
Friederike Bock, LBNL  
Elena Bratkovskaya, FIAS, Frankfurt U.  
Wolfgang Cassing, Gießen U.  
Gabor David, BNL  
Gabriel Denicol, McGill U.  
Axel Drees, Stony Brook U.  
Charles Gale, McGill U.  
Hendrik van Hees, FIAS Frankfurt  
Ulrich Heinz, Ohio State U.  
Christian Klein-Börsing, Münster U.  
Volker Koch, LBNL  
Olena Linnyk, Gießen U.  
Daniel Lohner, Heidelberg U.  
Constantin Loizides, LBNL  
Larry McLerran, BNL  
Jean-Yves Ollitrault, CEA Saclay  
Jean-François Paquet, McGill U.  
Ralf Rapp, Texas A&M U.  
Ilya Selyuzhenkov, EMMI  
Chun Shen, Ohio State U.  
Martin Wilde, Münster U.  
Li Yan, CEA Saclay

**Organizers**  
Johanna Stachel, Heidelberg U.  
Klaus Reygers, Heidelberg U.

**Further Information**  
[www.gsi.de/emmi/rtrf](http://www.gsi.de/emmi/rtrf)

**More about EMMI**  
[www.gsi.de/emmi](http://www.gsi.de/emmi)





# Resolution of the direct photon puzzle

## experiment

- decay photon cocktail?
- ...

## theory

### early stage

- Glasma?
- transport models (e.g. BAMPS) [2]
- Initial  $B$  field?

### late stage ( $T \approx T_c$ )

- ( $\rightarrow$  large  $T_{\text{eff}}$  due to blue shift)
- $\pi+\pi \rightarrow \pi+\pi+\gamma$  (e.g. PHSD model [1])
- "radiative hadronization"?
- ...

?

Possible paradigm shift concerning role of photons as QGP messengers?

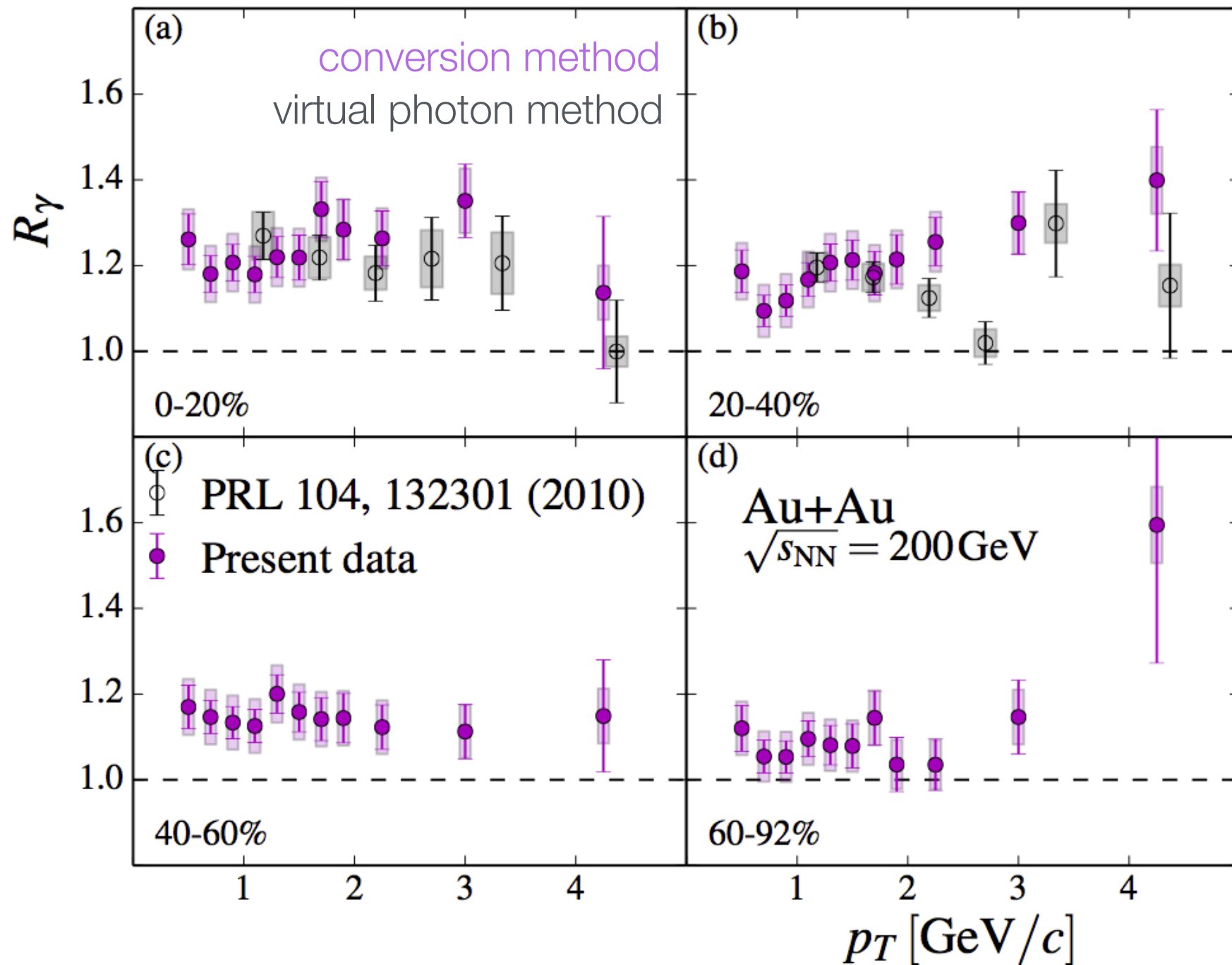
[1]: O. Linnyk et al, 1512.08126

[2]: M. Greif et al, 1612.05811

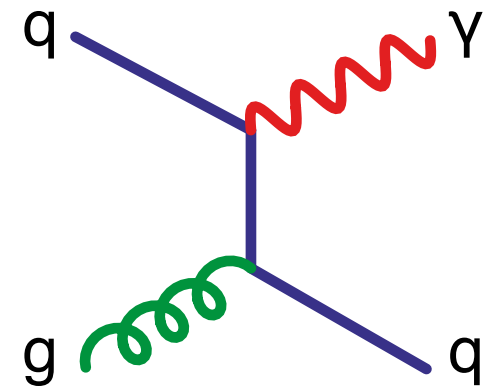
# RHIC: Two Methods, Same Answer

$$R_\gamma = \frac{\gamma_{\text{incl}}}{\gamma_{\text{decay}}} = 1 + \frac{\gamma_{\text{dir}}}{\gamma_{\text{decay}}}$$

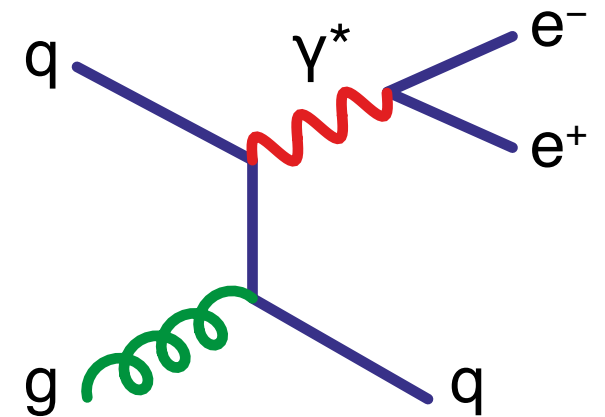
PHENIX, arXiv:1405.3940



Real photons (through conversion):



Virtual photons (at  $m_{ee} \gtrsim m_\pi$  extrapolated to  $m_{ee} = 0$ ):

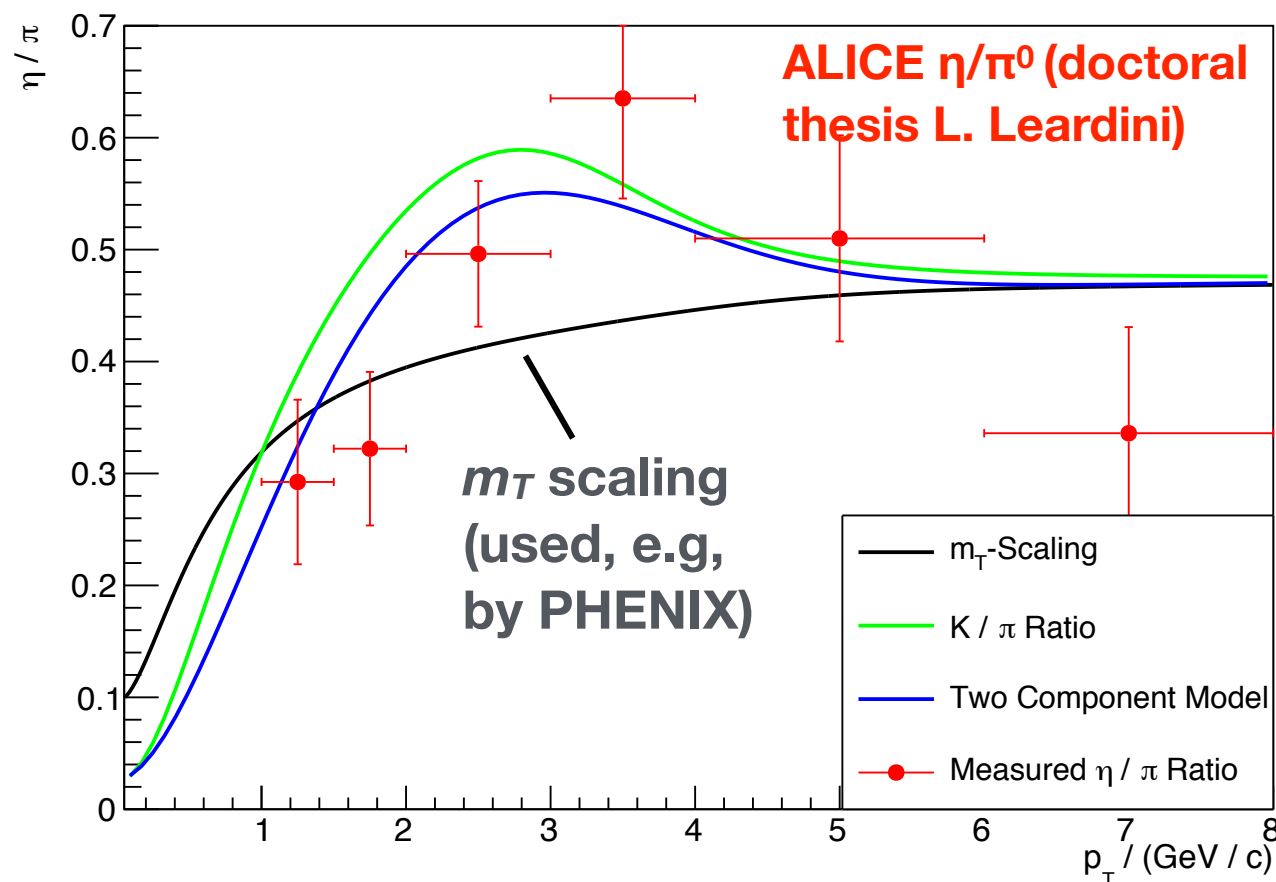


# Decay Photon Cocktail: Beyond $m_T$ scaling for $\eta$ , $\omega$ , $\eta'$ , ...

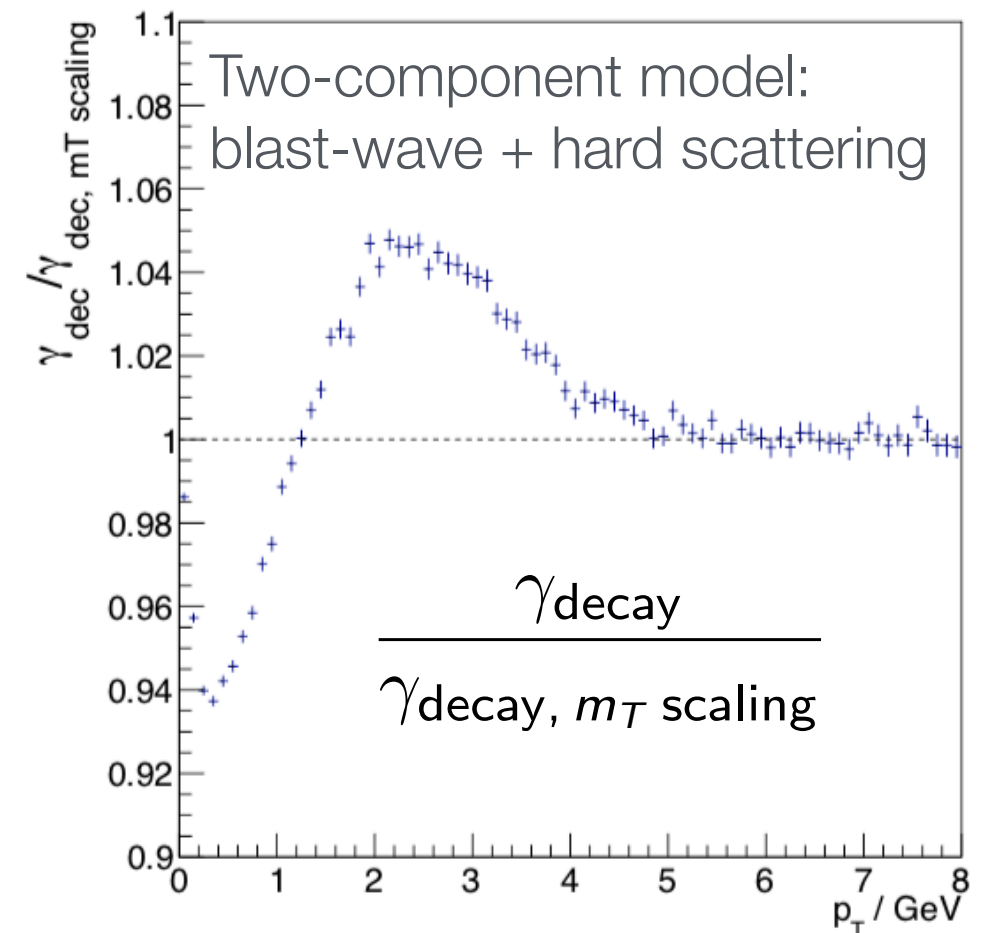
$m_T$  scaling often used to model spectra of  $\eta$ ,  $\omega$ , ...:

$$\frac{1}{p_T} \frac{dn}{dp_T} \propto f(m_T), \quad m_T = \sqrt{m^2 + p_T^2}$$

→ Include effect of radial flow (which breaks  $m_T$  scaling)



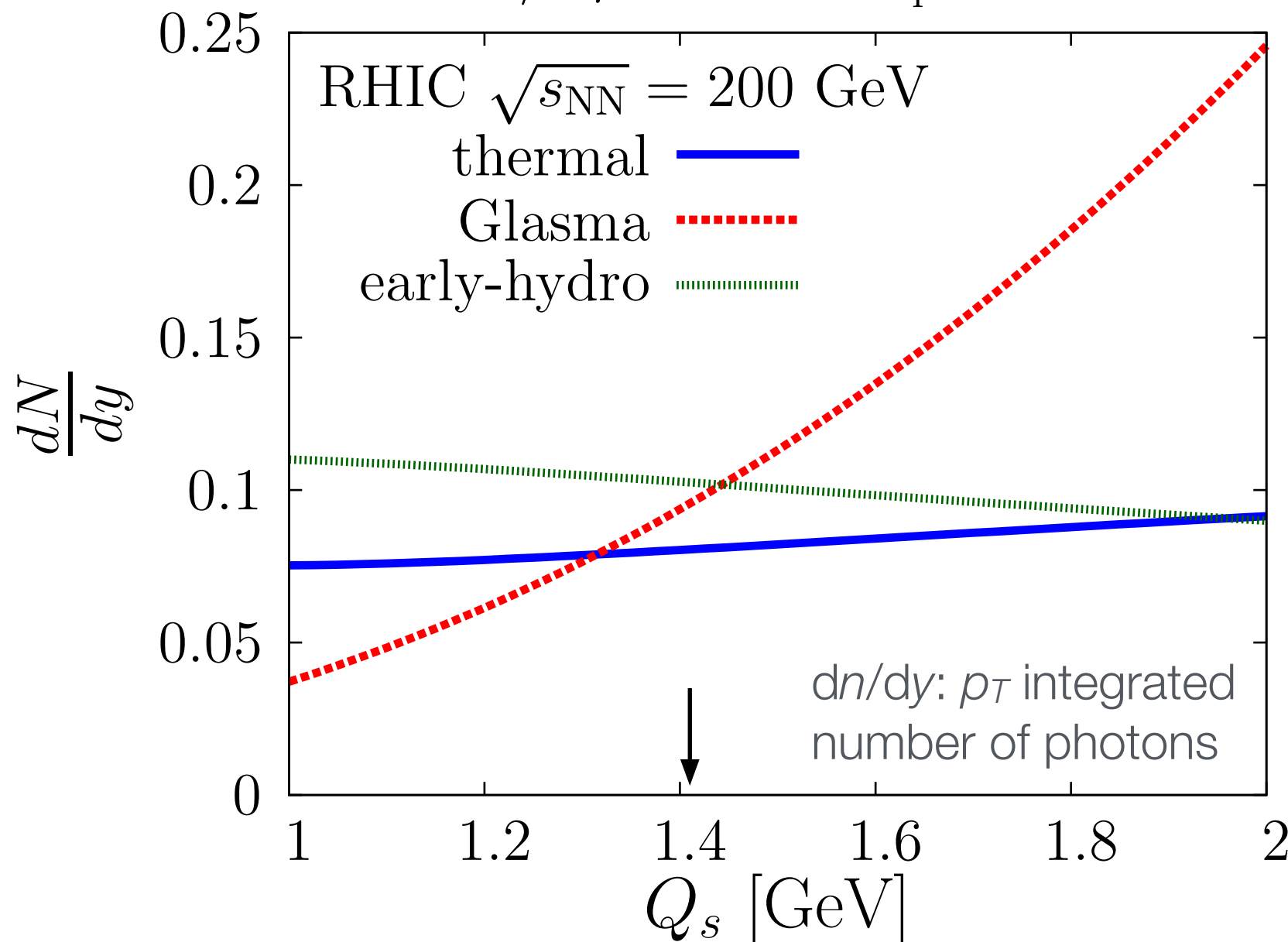
bachelor's thesis Ilya Fokin



Know your baseline!

# Early Stage: Glasma Contribution to Total Photon Yield Might be Sizable

$$dN_{\text{ch}}/d\eta = 687 \text{ at } N_{\text{part}} = 353$$



Parametric estimate

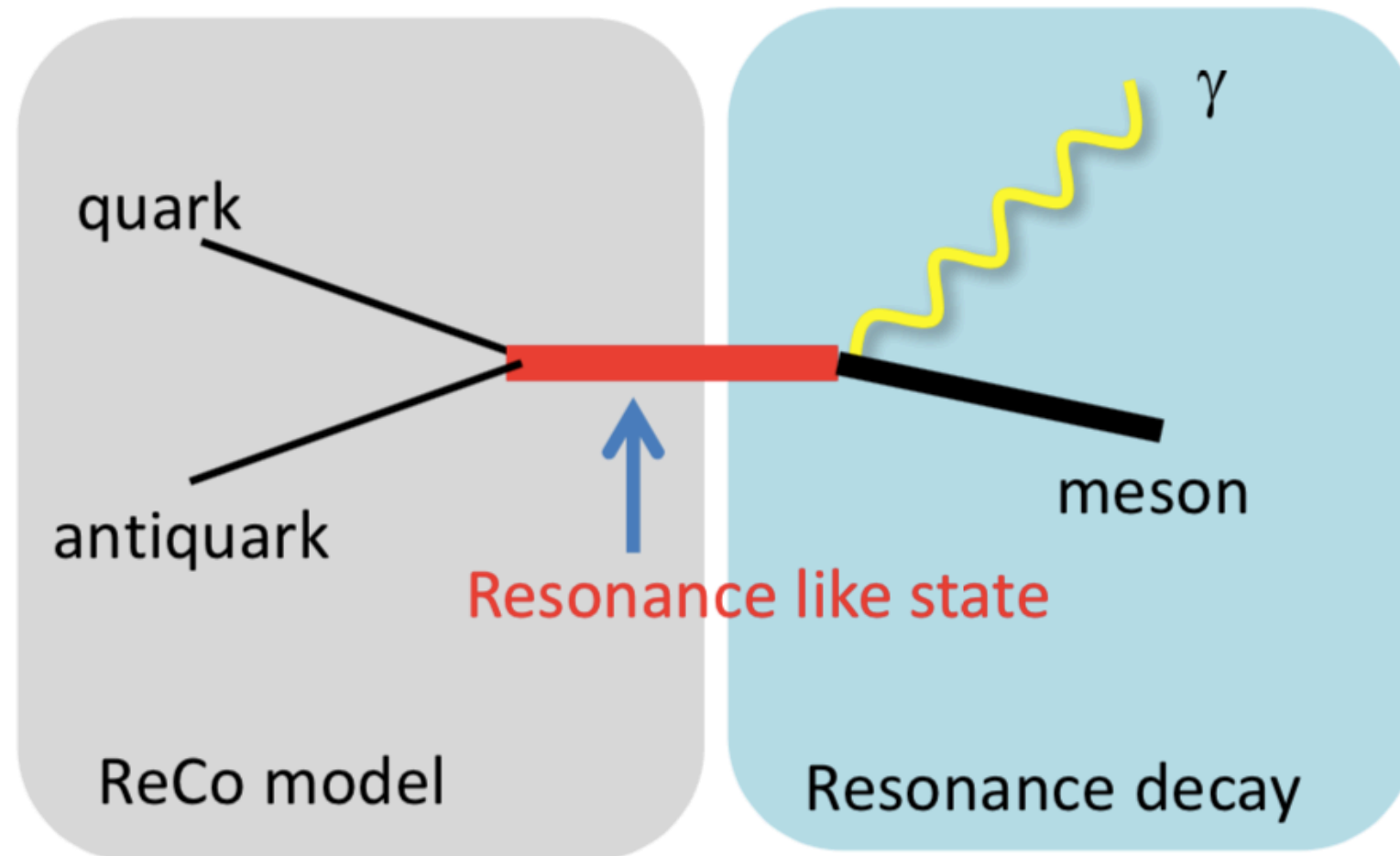
"bottom-up thermalization"  
[R.Baier et al.]

current hydro models:  
early-hydro + thermal

consistent weak coupling  
approach:  
Glasma + thermal

arXiv:1701.05064 (J. Berges,  
KR, N. Tanji, R. Venugopalan)

# Late Stage: Radiative Recombination?



- Naturally:

$$v_2(\gamma) \approx v_2(\text{hadron})$$

- Large  $T_{\text{eff}}$  due to blue shift

$$T_{\text{eff}} \approx \sqrt{\frac{1 + \beta}{1 - \beta}} T$$

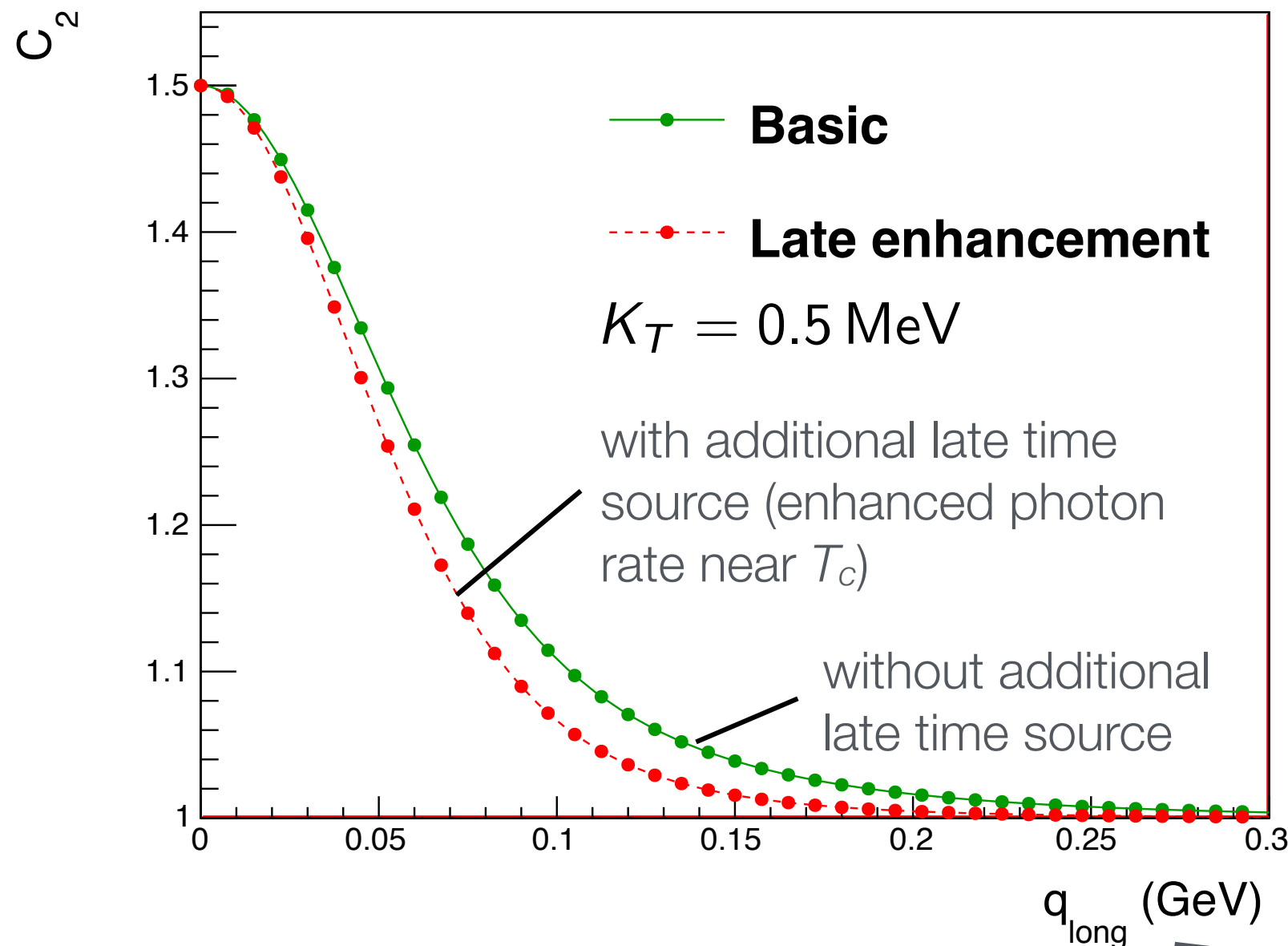
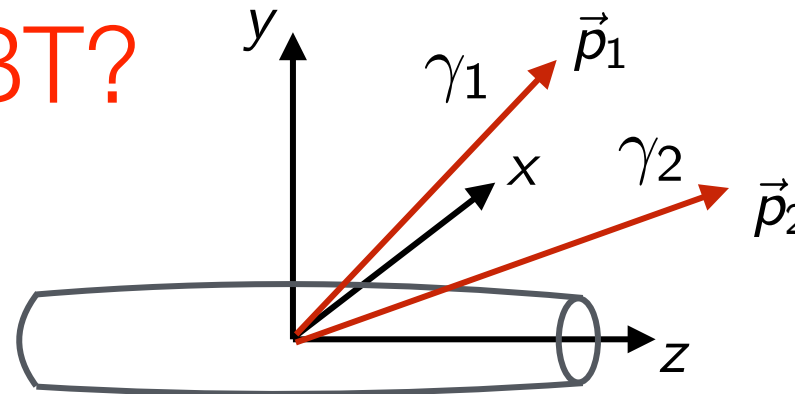
- "Saves" energy conservation in recombination models

Fujii, Itakura, Nonaka, Nucl.Phys. A967 (2017) 704-707

Young, Pratt, 1511.03147

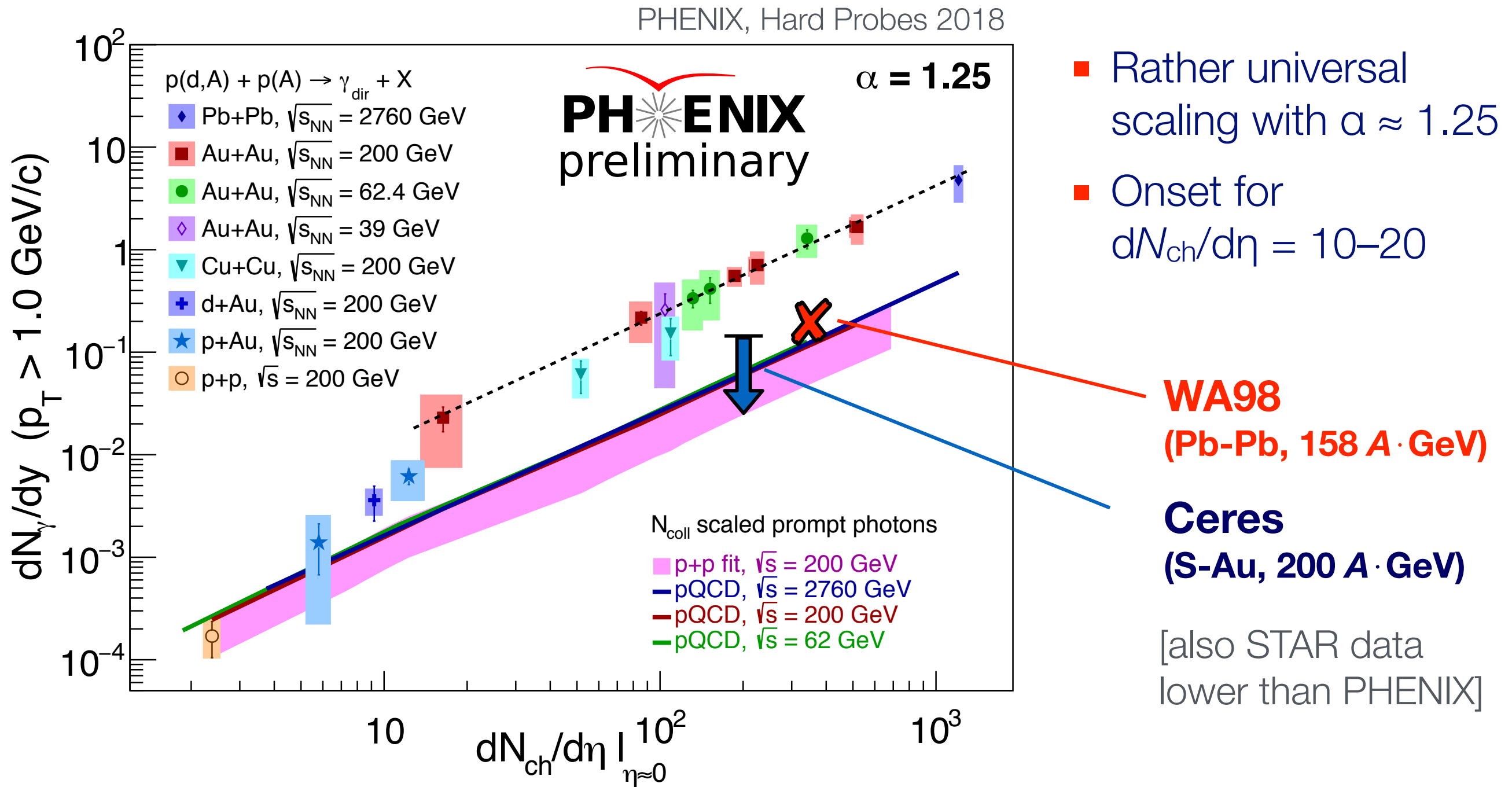
# Early or Late Stage Production: Constraints from Photon HBT?

$$C_2 = \frac{f(\vec{p}_1, \vec{p}_2)}{f(\vec{p}_1)f(\vec{p}_2)}$$



- Photon rate around  $T_c$  increased to describe data
- Narrower correlation for scenario with enhanced photon rate near  $T_c$
- Will be hard to measure, even in future high-statistics runs at the LHC

# Do the Data Speak for Themselves? Universal Scaling?



$\alpha = 1.25 \triangleq N_{\text{coll}} \text{ scaling} \Rightarrow$  photons related to initial parton scattering?

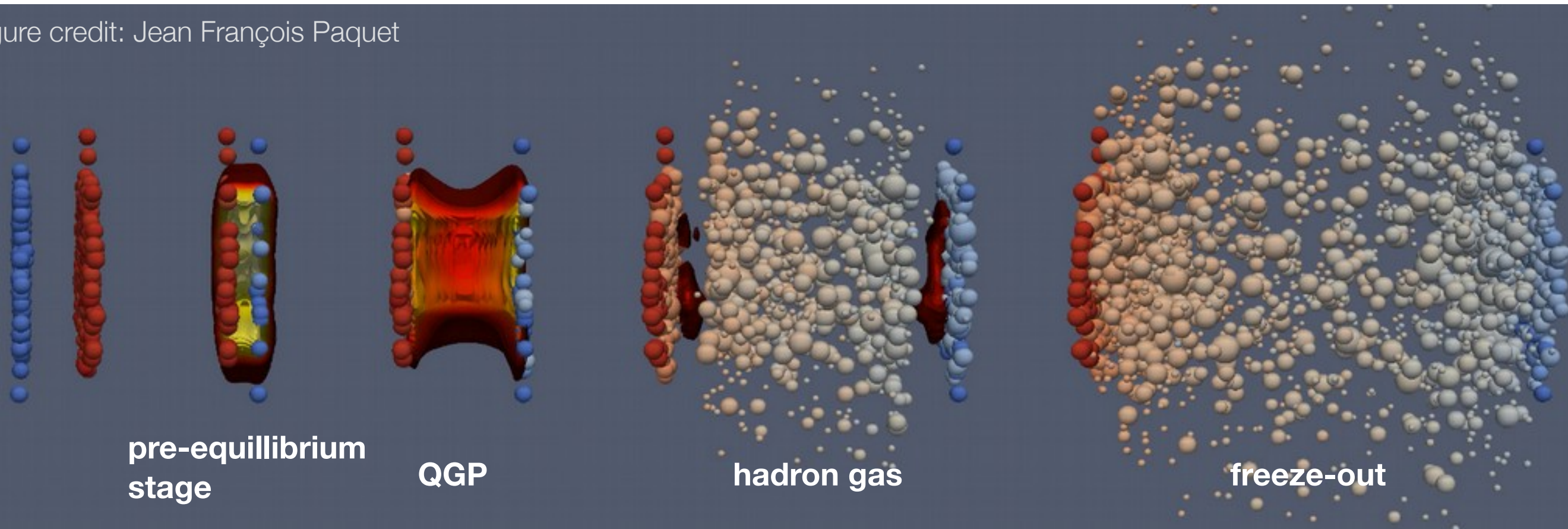


# Conclusions

- ALICE: Large  $v_{2,\text{dir}}$  ( $\approx v_{2,\text{dec}}$ ), but also large uncertainties
- Direct photon puzzle mostly at RHIC
- Quantifying the statistical significance of the puzzle would be a nice joint project
- Possible paradigm shift:  
Photon production dominated by late stage around  $T_c$ ?

# Extra Slides

figure credit: Jean François Paquet



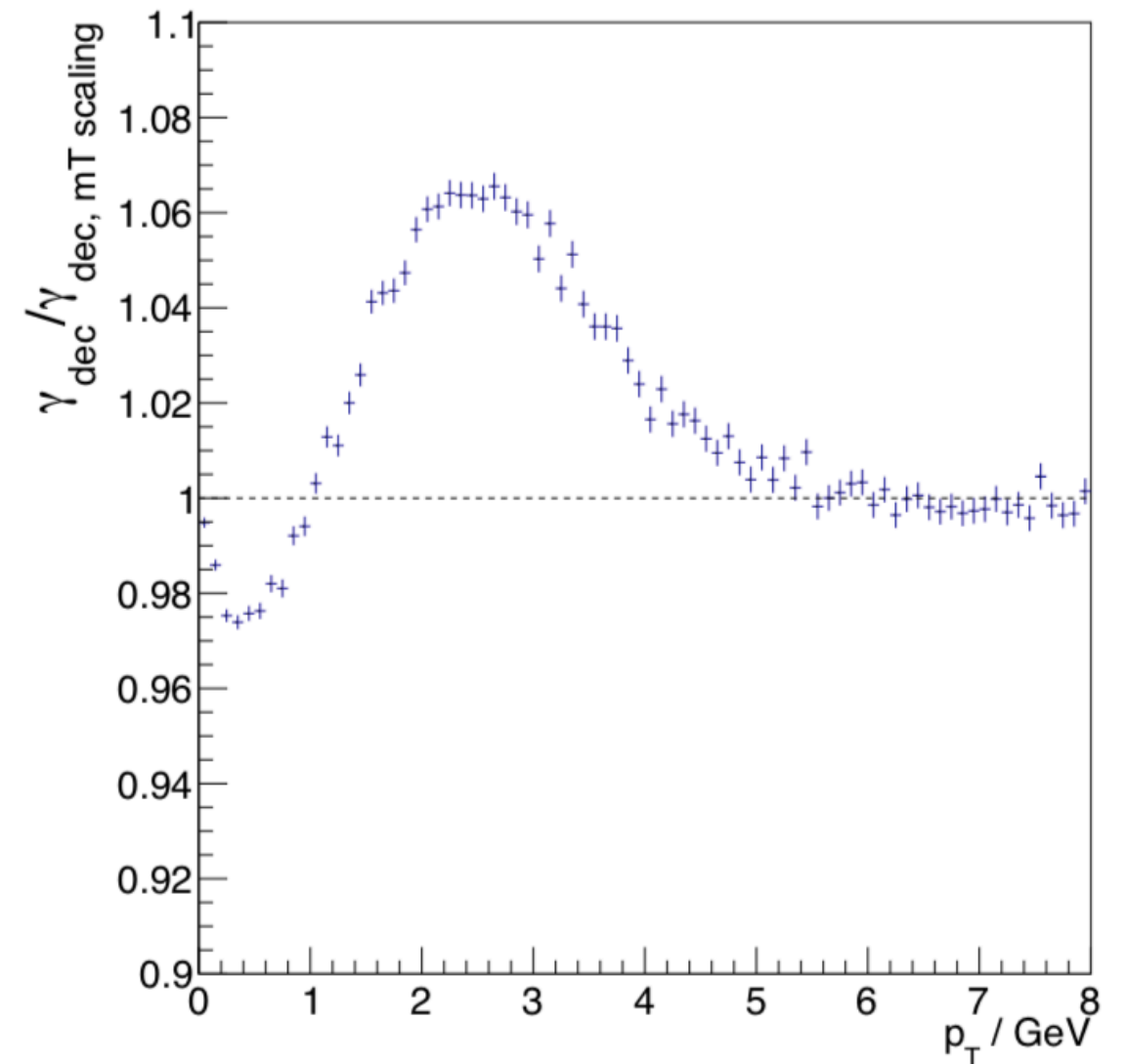
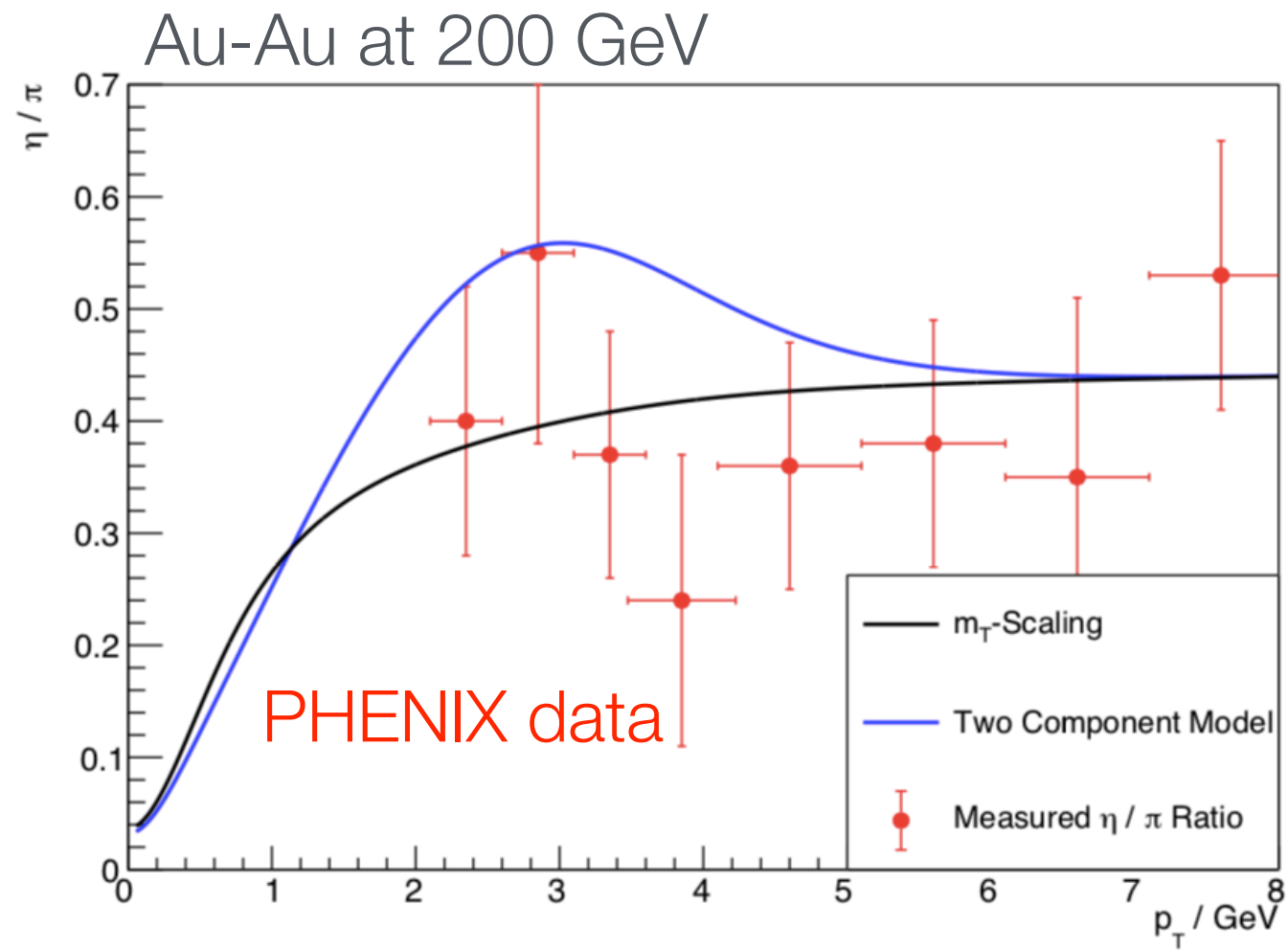
thermalization:  $\tau_{\text{th}} = 1-2 \text{ fm}/c$

hadronization:  $\tau_c \approx 10 \text{ fm}/c$  (LHC)

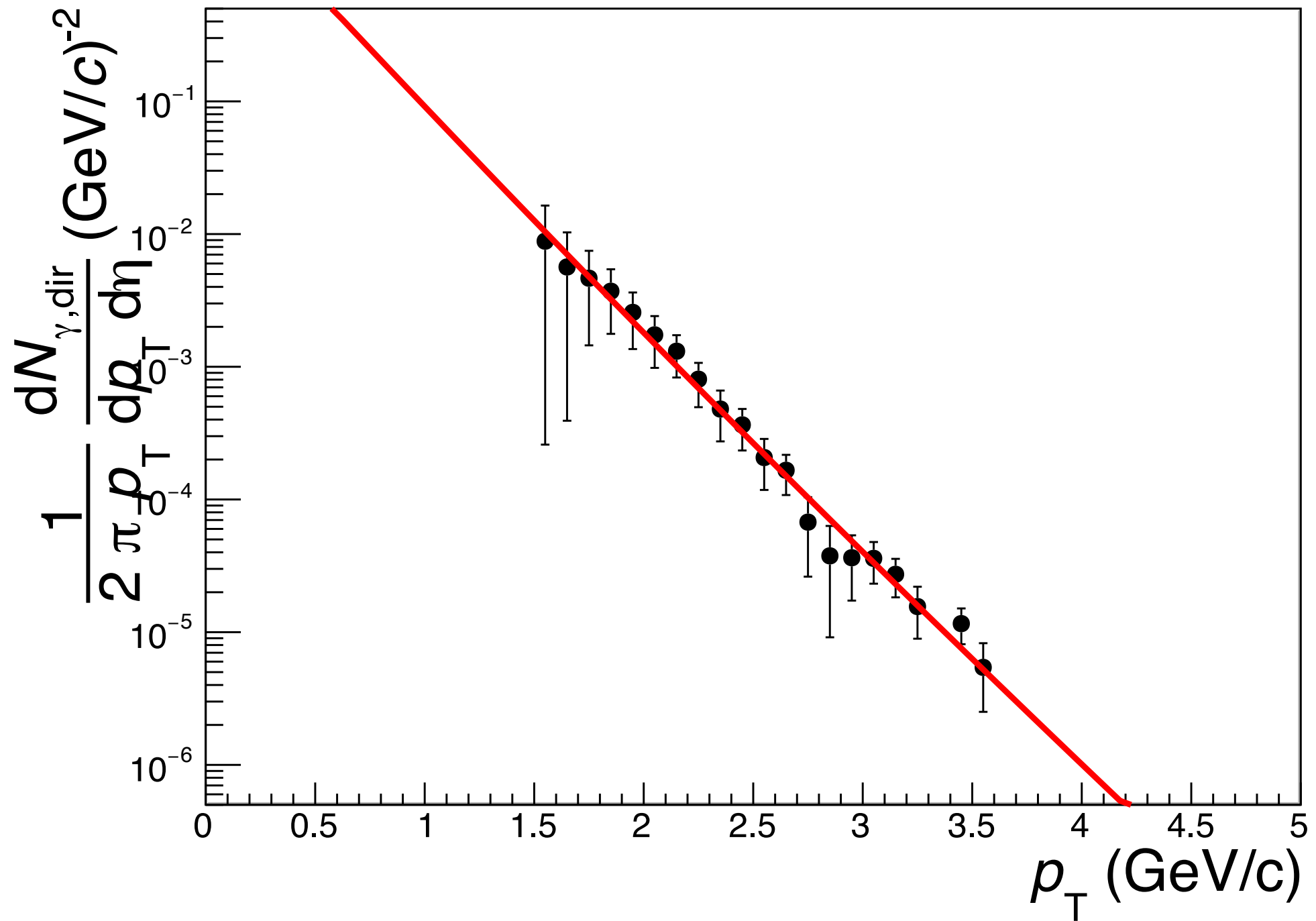
QGP  $\rightarrow$  hadron gas at

$T \approx 155 \text{ MeV}$

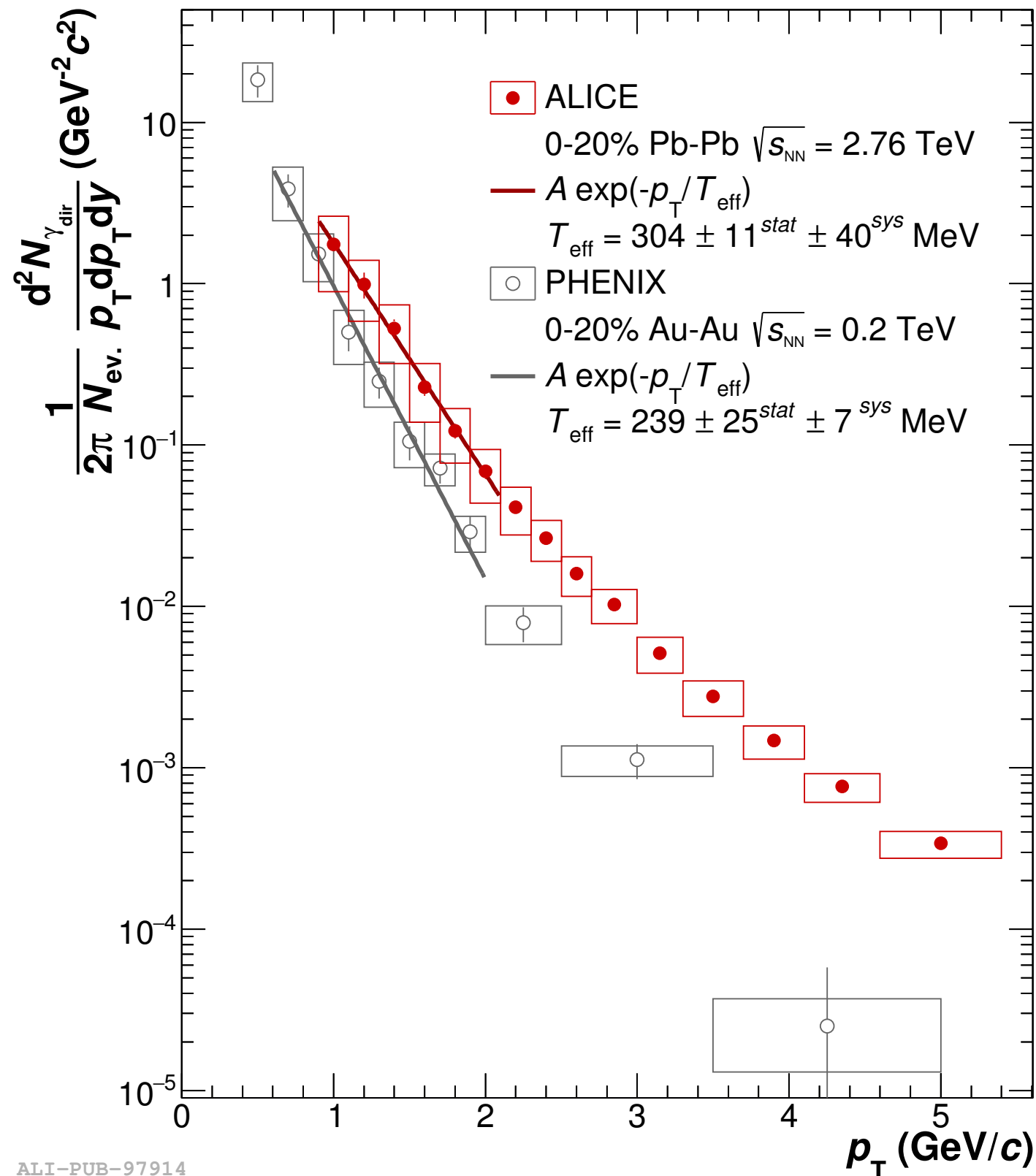
# Improved Cocktail: Au-Au at 200 GeV



# WA98 Data



# Larger $T_{\text{eff}}$ at the LHC



## ■ $T_{\text{eff}}$ LHC

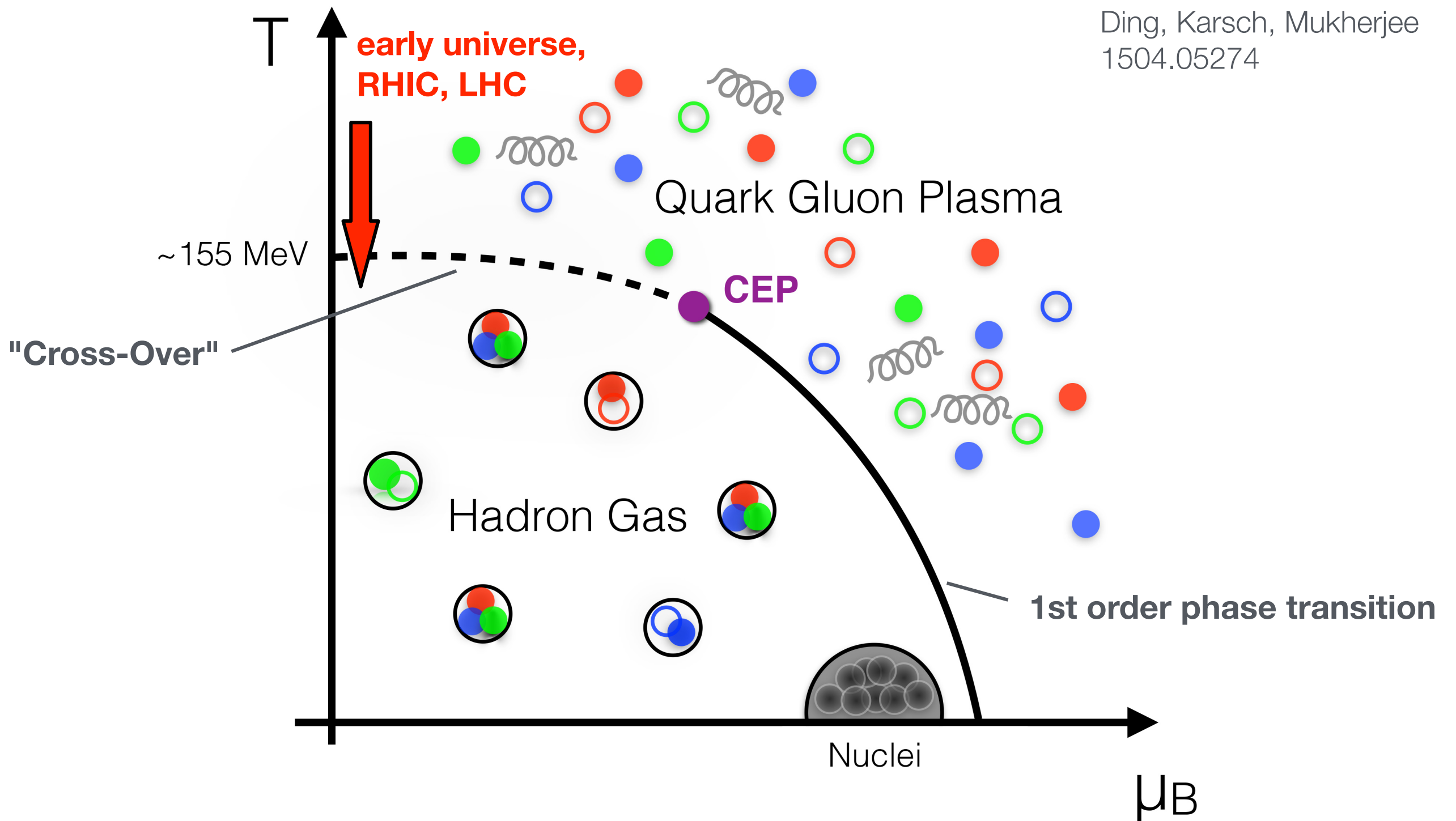
- ▶ 0-20% Pb-Pb@2.76 TeV
- ▶ without pQCD subtraction:  
 $T_{\text{eff}} = 304 \pm 11^{\text{stat}} \pm 40^{\text{sys}}$  MeV
- ▶ with pQCD subtraction:  
 $T_{\text{eff}} = 297 \pm 12^{\text{stat}} \pm 41^{\text{sys}}$  MeV

## ■ $T_{\text{eff}}$ RHIC

- ▶ 0-20% Au-Au@0.2 TeV
- ▶  $T_{\text{eff}} = 239 \pm 25^{\text{stat}} \pm 7^{\text{sys}}$  MeV  
(pp parameterization subtracted)

# (Conjectured) QCD Phase Diagram

Ding, Karsch, Mukherjee  
1504.05274

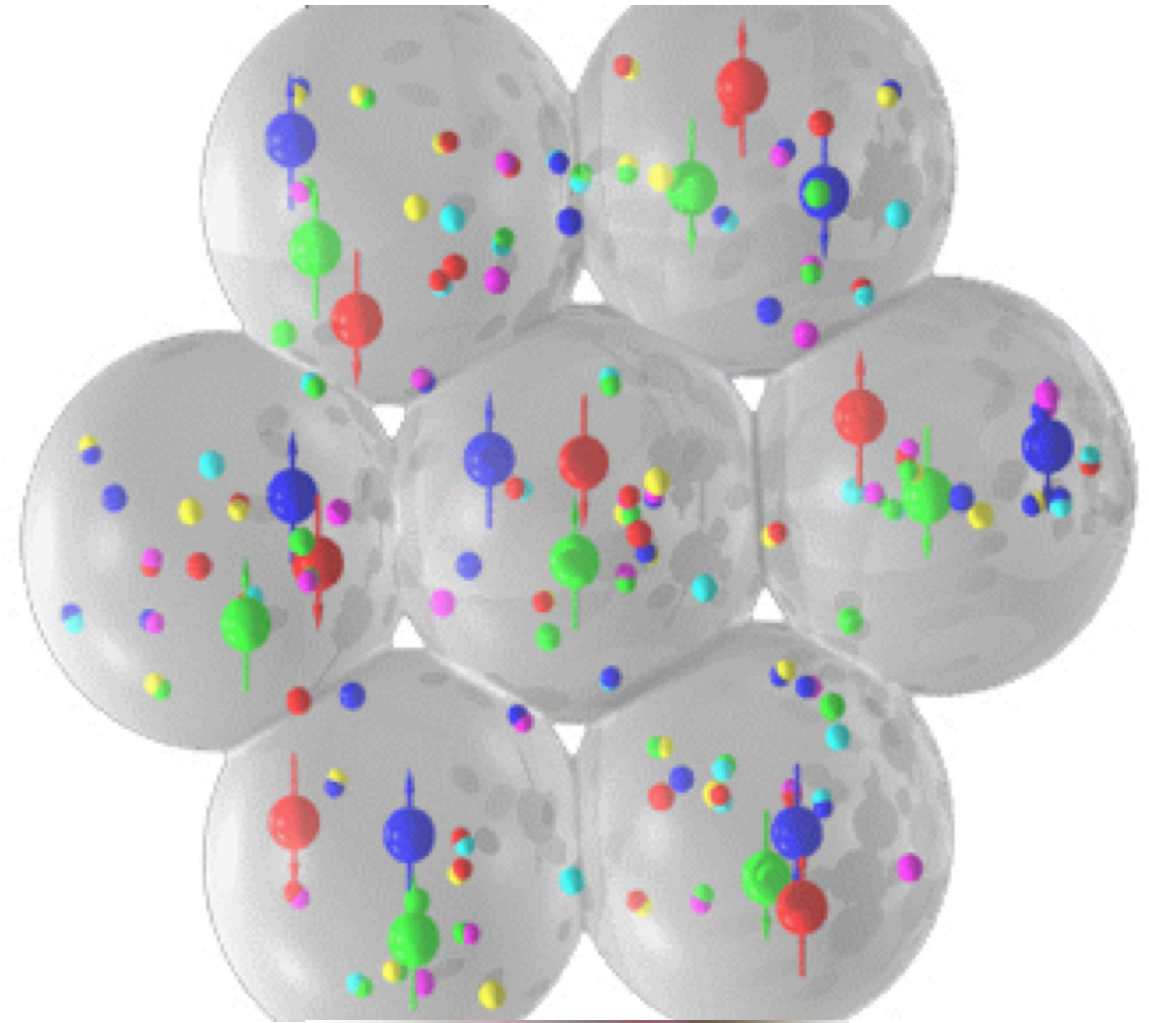




# What is the question?

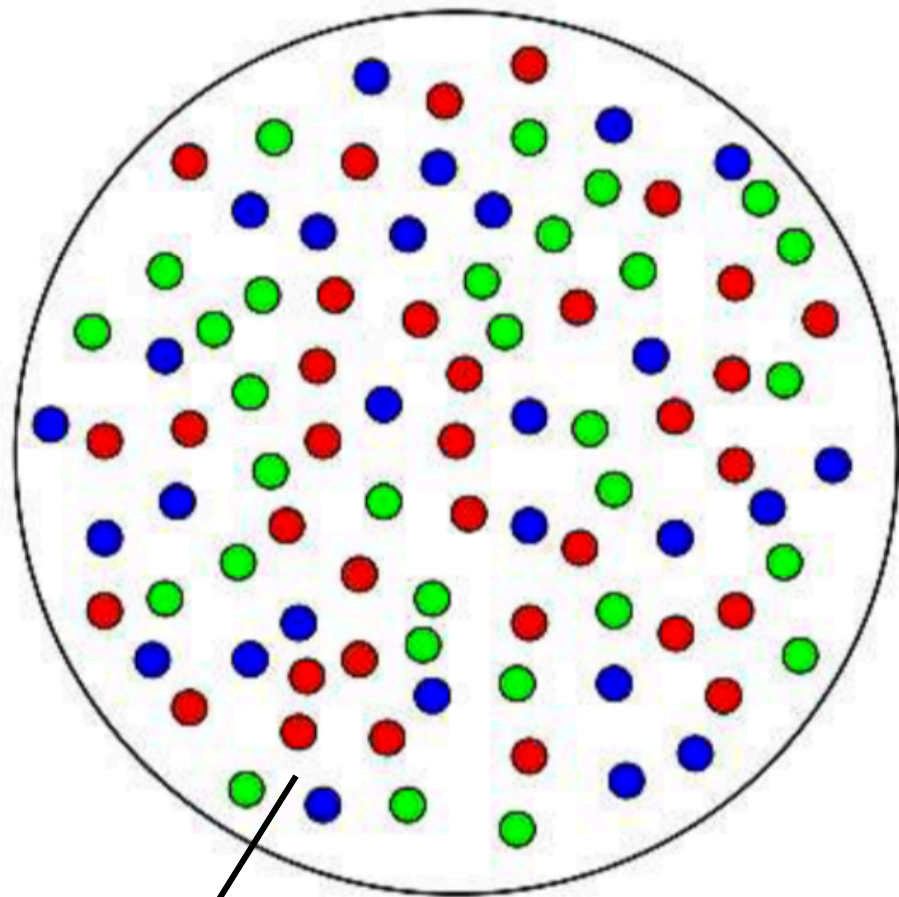
What happens if make nuclear matter

- hotter and hotter?
- denser and denser?



**solid → liquid → gas → plasma → hadron gas → QGP**

# Let the data Speak: Empirical Scaling Law for $n_\gamma$ vs $n_{\text{hadron}}$ ?



QGP at fixed temperature  $T$

In recombination models:

$$n_\gamma \propto n_h$$

Parameterization:

$$n_\gamma \propto n_h^\alpha$$

Bjorken expansion (only QGP):

$$\alpha \approx 2$$

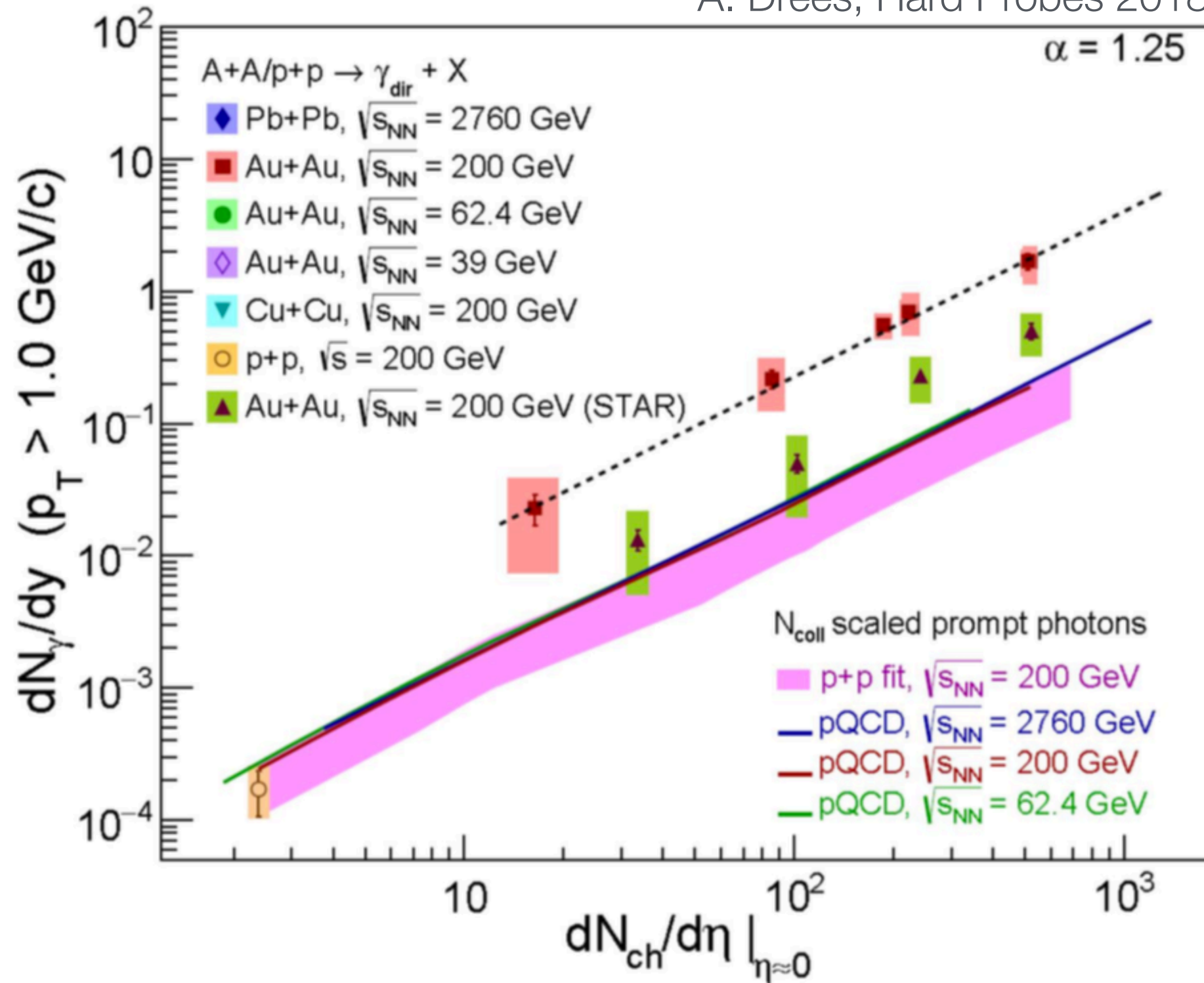
Realistic hydro model:  
( $p_{T,\gamma} > 1 \text{ GeV}/c$ )

$$\alpha \approx 1.6\text{--}1.7$$

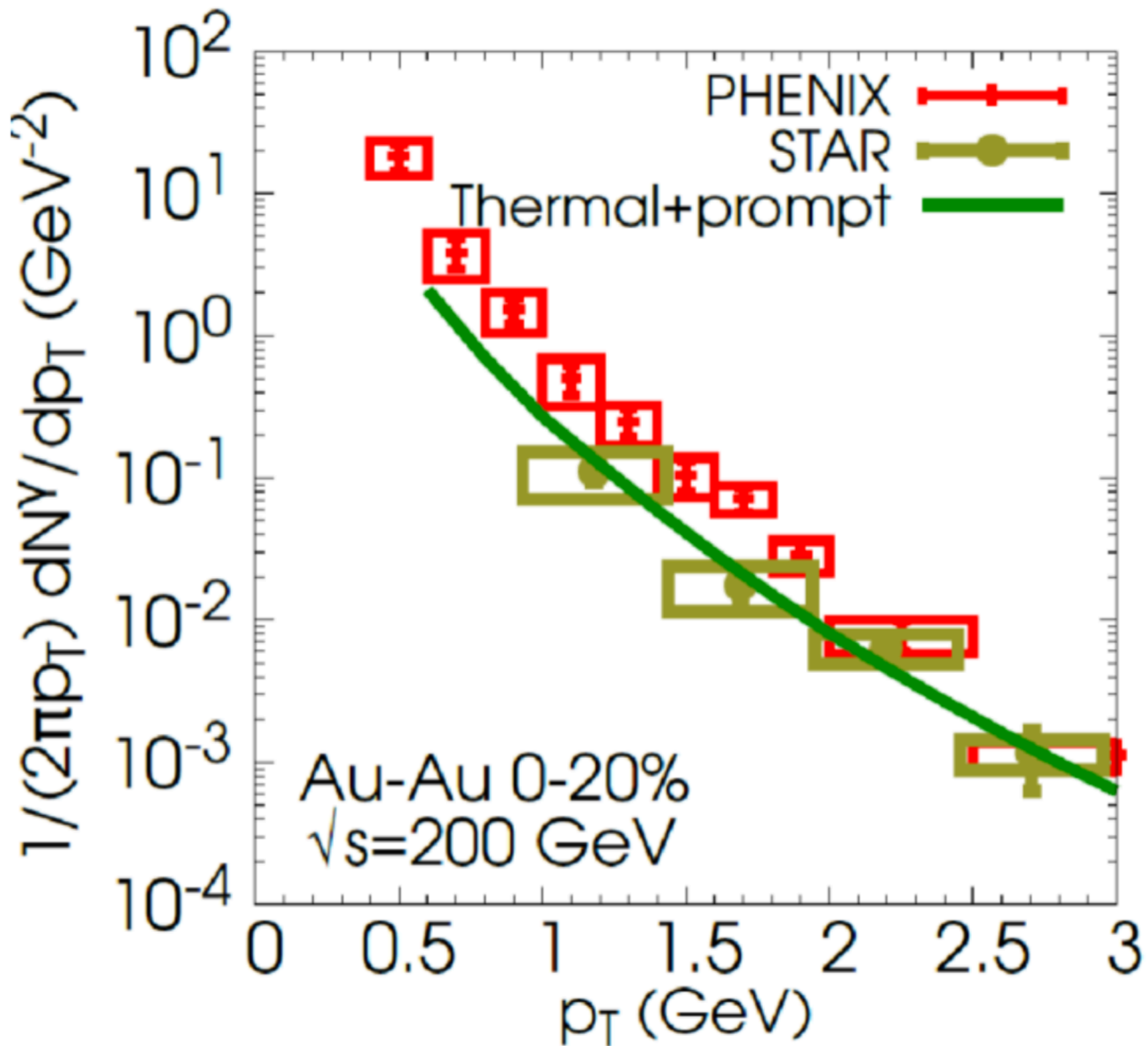
Jean-François Paquet,  
Hard Probes 2018

# Direct Photons: PHENIX vs. STAR

A. Drees, Hard Probes 2018



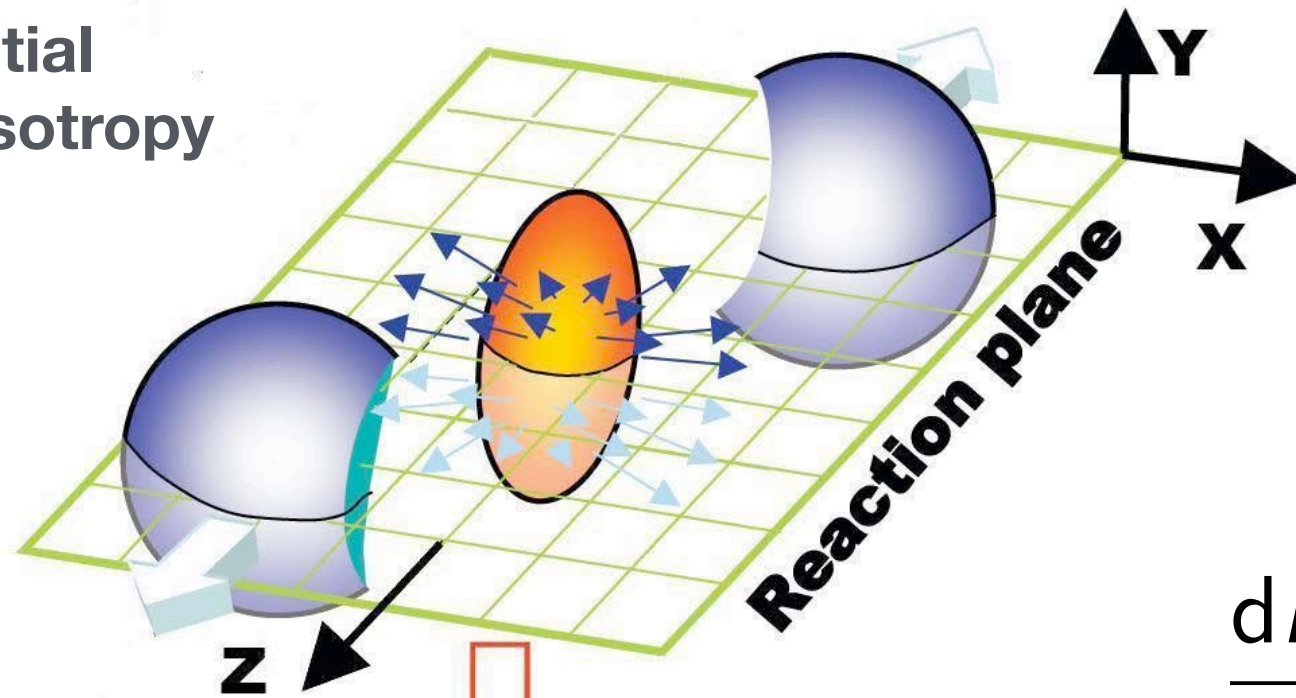
# Direct Photons: PHENIX vs. STAR



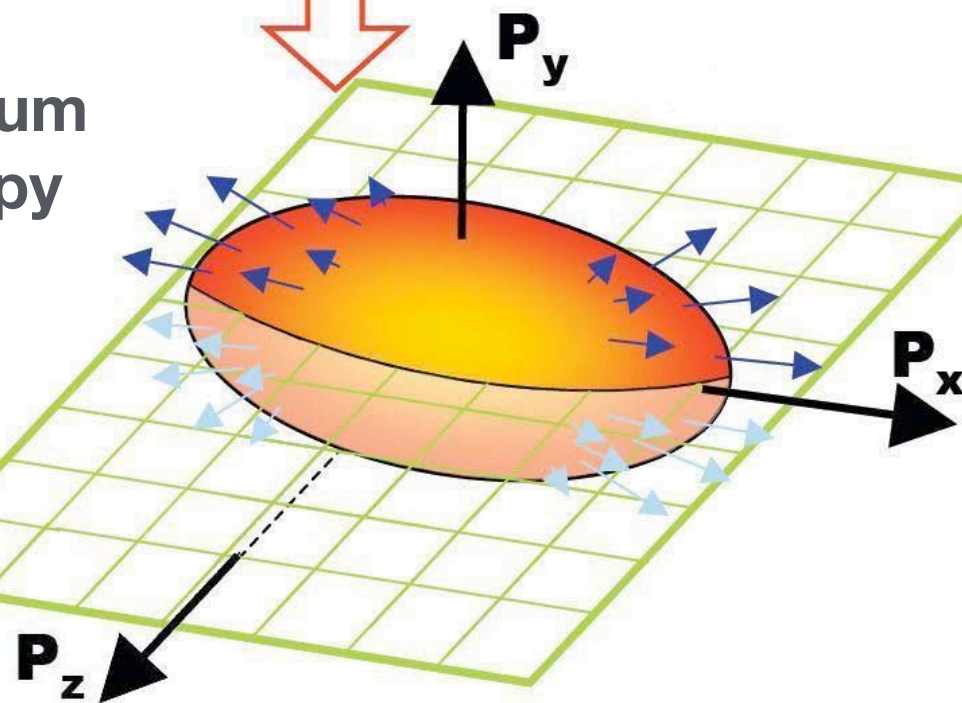


# Elliptic Flow

spatial  
anisotropy



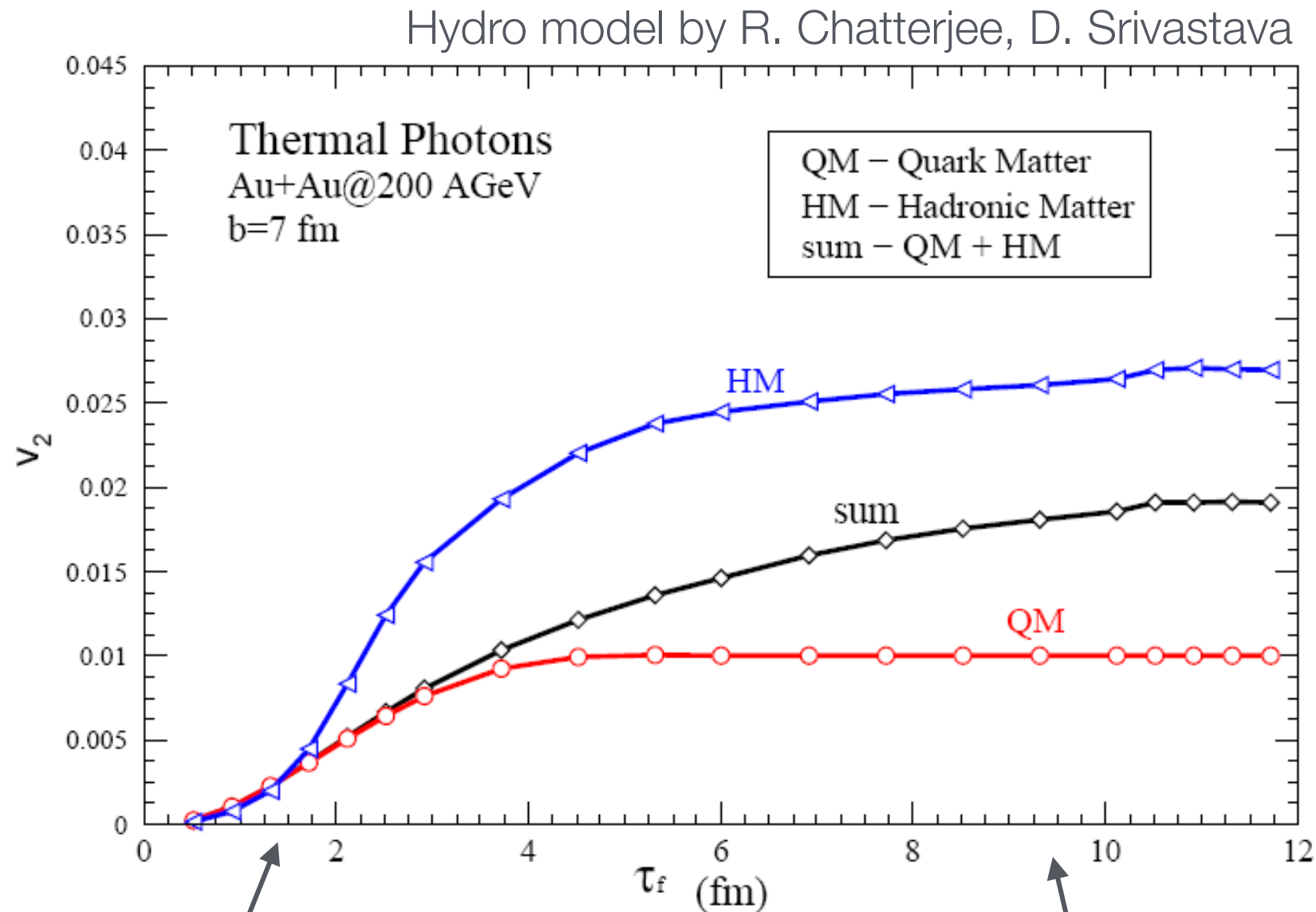
momentum  
anisotropy



$$\frac{dN}{d\varphi} = N_0 (1 + 2v_2 \cos(2\varphi))$$

# What's actually so puzzling?

Elliptic flow builds up gradually with time in hydro models:



Expect large fraction of thermal photons from early times

Expect bulk of hadrons to be produced at late times

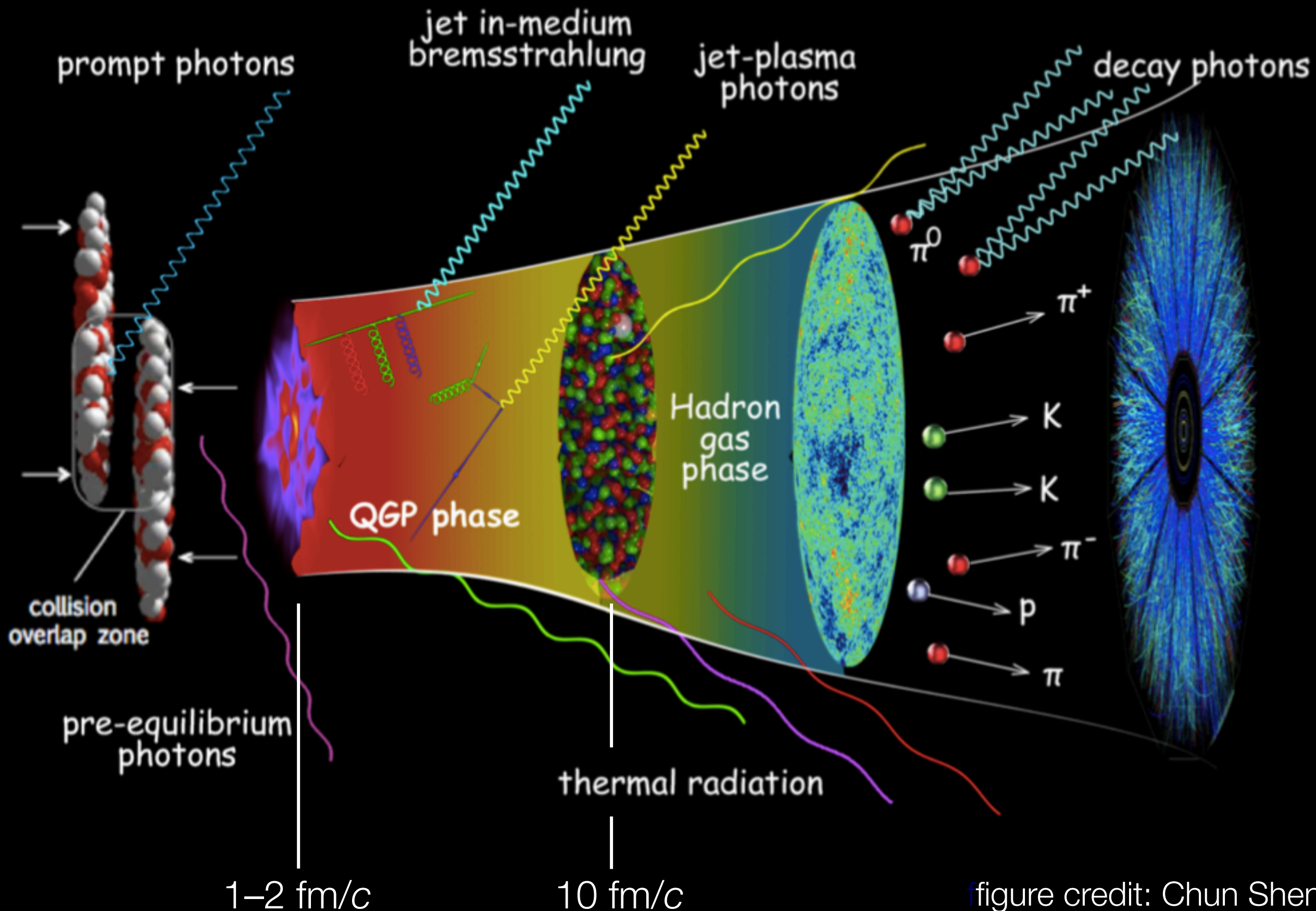


figure credit: Chun Shen