

Virtual photon polarization and dilepton anisotropy in relativistic nucleus-nucleus collisions

Enrico Speranza

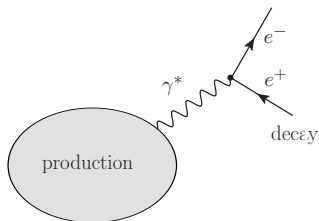
E.S., A. Jaiswal, B. Friman, PLB **782**, 395 (2018)
B. Friman, R. Rapp, E.S., J. Wambach (in preparation)



Electromagnetic Radiation from Hot and Dense Hadronic Matter

ECT*, Trento, November 26, 2018

The goal



Angular distribution of the dilepton



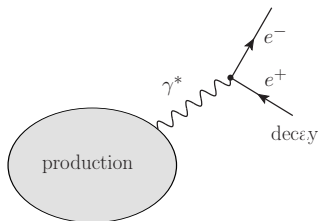
Information on the polarization states of the virtual photon



Information on the production mechanism

- ▶ Early stages and onset of thermalization in HIC (P. Hoyer, PLB **187**, 162; E. Shuryak, arXiv:1203.1012)
- ▶ Parton anisotropic momentum distributions in HIC (G. Baym, T. Hatsuda, and M. Strickland, PRC **95**, 044907)
- ▶ Help disentangle different production mechanisms in elementary reactions (E.S., M. Zétényi, and B. Friman, PLB **764**, 282)

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Spin-density matrix

- ▶ **Pure state:** $|\psi\rangle = \sum_{\lambda} c_{\lambda} |\lambda\rangle$
Expectation value of an operator $\langle O \rangle = \langle \psi | O | \psi \rangle$
- ▶ **Mixed state:** incoherent mixture of $|\psi_i\rangle$ with statistical weight a_i

$$\rho = \sum_i a_i |\psi_i\rangle \langle \psi_i| = \sum_{\lambda, \lambda'} \rho_{\lambda\lambda'} |\lambda\rangle \langle \lambda'|$$

$$\rho_{\lambda\lambda'} = \sum_i a_i c_{\lambda}^{(i)} c_{\lambda'}^{(i)*}. \text{ Expectation value: } \langle O \rangle = \text{Tr}(\rho O)$$

Example: Spin-1/2 particle (2×2 hermitian matrix):

$$\rho = \frac{1}{2}(1 + \vec{P} \cdot \vec{\sigma})$$

- ▶ Spin polarization vector: $\vec{P} = \langle \vec{\sigma} \rangle = \text{Tr}(\rho \vec{\sigma})$

$$|\vec{P}| = 1 \quad \text{Pure state}$$

$$0 < |\vec{P}| < 1 \quad \text{Mixed state}$$

$$|\vec{P}| = 0 \quad \text{Completely unpolarized mixed state}$$

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Spin-density matrix for spin-1 particles

- ▶ Three polarization states (in rest frame)

$$\text{Transverse to } \vec{q}: \quad \epsilon(\pm 1) = \mp \frac{1}{\sqrt{2}} (0, 1, \pm i, 0)$$

$$\text{Longitudinal to } \vec{q}: \quad \epsilon(0) = (0, 0, 0, 1)$$

- ▶ **Spin-density matrix:** hermitian 3×3 matrix

$$\rho = \frac{1}{3} \left[1 + \frac{3}{2} \vec{P} \cdot \vec{S} + \sqrt{\frac{3}{2}} \sum_{i,j} T_{ij} (S_i S_j + S_j S_i) \right]$$

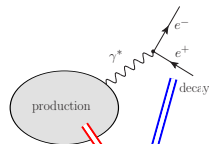
$\vec{S}$ are the spin-1 operators

- ▶ $\text{Tr} \rho = 1$ (8 parameters)
- ▶ Vector polarization: $\vec{P} = \langle \vec{S} \rangle$ (3 parameters)
- ▶ Tensor polarization: $T_{ij} = \frac{1}{2} \sqrt{\frac{3}{2}} (\langle S_i S_j + S_j S_i \rangle - \frac{4}{3} \delta_{ij})$, $\sum_i T_{ii} = 0$ (5 parameters)

One can have tensor polarization without vector polarization

Lepton angular distribution

$$\text{spin-1} \rightarrow \text{spin-}\frac{1}{2} + \text{spin-}\frac{1}{2}$$

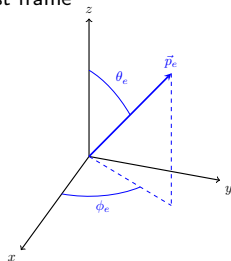


$$\frac{d\sigma}{d\Omega_e} \propto \text{Tr}(\rho O^{\text{dec}})$$

$$= \mathcal{N} \left(1 + \lambda_\theta \cos^2 \theta_e + \lambda_\phi \sin^2 \theta_e \cos 2\phi_e + \lambda_{\theta\phi} \sin 2\theta_e \cos \phi_e \right. \\ \left. + \lambda_\phi^\perp \sin^2 \theta_e \sin 2\phi_e + \lambda_{\theta\phi}^\perp \sin 2\theta_e \sin \phi_e + \text{parity violating terms} \right)$$

$$\lambda_\theta = \frac{\rho_T - \rho_L}{\rho_T + \rho_L}$$

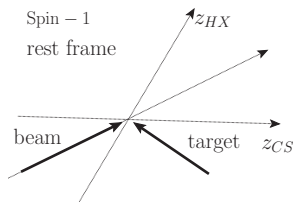
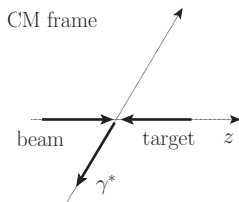
Photon rest frame



- ▶ Transverse: $\rho_T = \rho_{-1-1} + \rho_{+1+1}$
Longitudinal: $\rho_L = 2\rho_{00}$
($\rho = \sum_{\lambda, \lambda'} \rho_{\lambda\lambda'} |\lambda\rangle \langle \lambda'|$)
- ▶ Completely **transverse** polarized: $\lambda_\theta = +1$
Completely **longitudinal** polarized: $\lambda_\theta = -1$
- ▶ Photon polarization reflected in angular distribution

Reference frames

Anisotropy coefficients depend on the reference frame

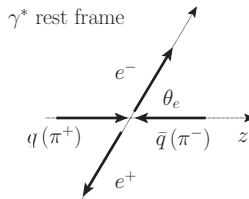


- ▶ Helicity (HX): z-axis along photon momentum
- ▶ Collins-Soper (CS): z-axis along bisector between beam and target
- ▶ Different frames are related by rotation

Examples

- Drell-Yan process: $q\bar{q} \rightarrow \gamma^* \rightarrow e^+e^-$

$$\frac{d\sigma}{d\Omega_e} \sim 1 + \cos^2 \theta_e$$



$\lambda_\theta = +1$. Virtual photon is completely transverse polarized along beam axis

- Pion annihilation process: $\pi^+\pi^- \rightarrow \gamma^* \rightarrow e^+e^-$

$$\frac{d\sigma}{d\Omega_e} \sim 1 - \cos^2 \theta_e$$

$\lambda_\theta = -1$. Virtual photon is completely longitudinal polarized along beam axis

Virtual photon emission from a thermal medium

$$q\bar{q} \rightarrow \gamma^* \rightarrow e^+e^-$$
$$\pi^+\pi^- \rightarrow \gamma^* \rightarrow e^+e^-$$

- ▶ Thermal average of initial particles momenta p through Fermi or Bose distribution

$$f(p) = \frac{1}{e^{(u \cdot p)/T} \pm 1}$$

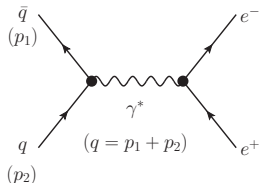
- ▶ Fluid rest frame $u^\mu = (1, 0, 0, 0) \Rightarrow$ Distribution is spherical symmetric

Photon momentum $\vec{q}$ breaks spherical symmetry,
but not azimuthal symmetry

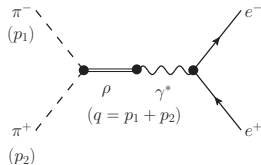
- ▶ Photons are only **tensor** polarized
- ▶ $|\vec{q}| \rightarrow 0 \Rightarrow$ No anisotropy $\Rightarrow$ **No photon polarization**

Boltzmann limit

$$q + \bar{q} \rightarrow \gamma^* \rightarrow e^+ e^-$$



$$\pi^+ \pi^- \rightarrow \gamma^* \rightarrow e^+ e^-$$

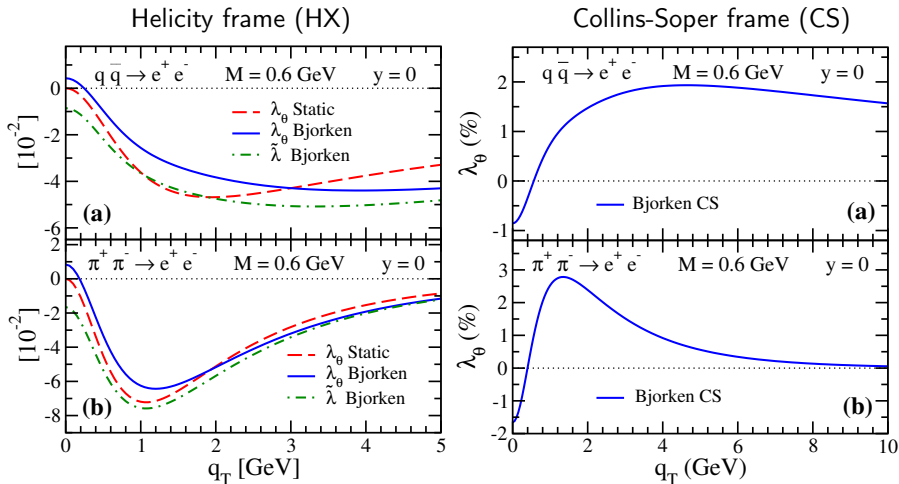


$$\int \frac{d^3 p_1}{E_1} \frac{d^3 p_2}{E_2} \frac{1}{e^{(u \cdot p_1)/T} \pm 1} \frac{1}{e^{(u \cdot p_2)/T} \pm 1} \sim e^{-(u \cdot q)/T} \int \frac{d^3 p_1}{E_1} \frac{d^3 p_2}{E_2}$$

- No photon polarization independently of photon momentum

Photon polarization is due to quantum statistics!

Results (static and Bjorken expansion)

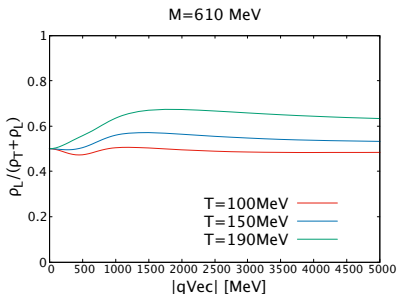
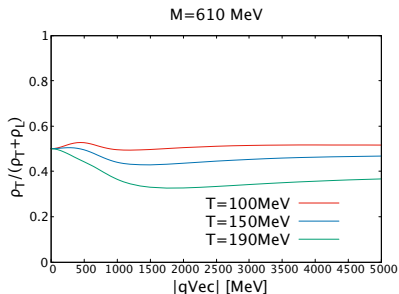


- ▶ Static case: $\lambda_\theta \rightarrow 0$ for $q_T \rightarrow 0$, and for $q_T \rightarrow \infty$ (Boltzmann limit)
- ▶ λ_θ changes sign from the Helicity to Collins-Soper frame
- ▶ Frame invariant combination: $\tilde{\lambda}_\theta = \frac{\lambda_\theta + 3\lambda_\phi}{1 - \lambda_\phi}$
- ▶ Experiments: sum over q_T

Realistic models (Preliminary)

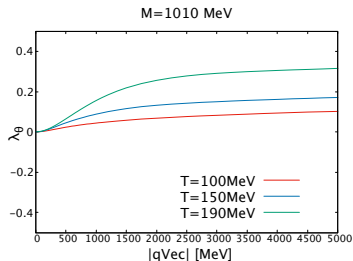
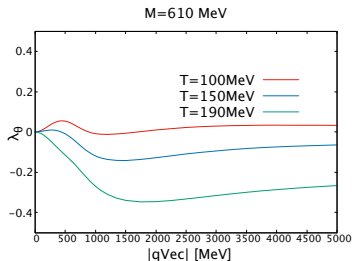
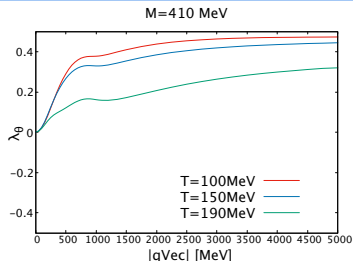
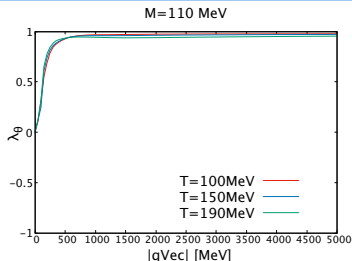
$$\frac{d\sigma}{d\Omega_e} \propto \rho^{\mu\nu} L_{\mu\nu}$$

$$\rho^{\mu\nu} = \rho_T P_T^{\mu\nu} + \rho_L P_L^{\mu\nu}$$



- ρ_T , ρ_L taken from: Rapp, Chanfray, Wambach, NPA **617**, 472;
Rapp, Wambach EPJA **6**, 415; Urban, Buballa, Rapp, Wambach, NPA **673**, 357

Polarization with realistic models (Preliminary)

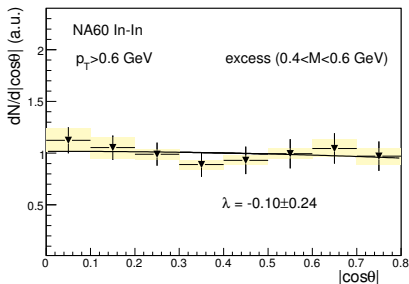


$$\lambda_\theta = \frac{\rho_T - \rho_L}{\rho_T + \rho_L}$$

Large polarization!

Experimental results (NA60 and HADES)

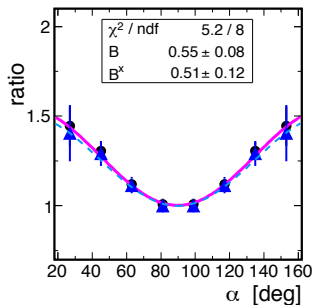
Collins-Soper frame



In-In at 158A GeV

(NA60 Collaboration), PRL **96**, 222301 (2009)

Helicity frame



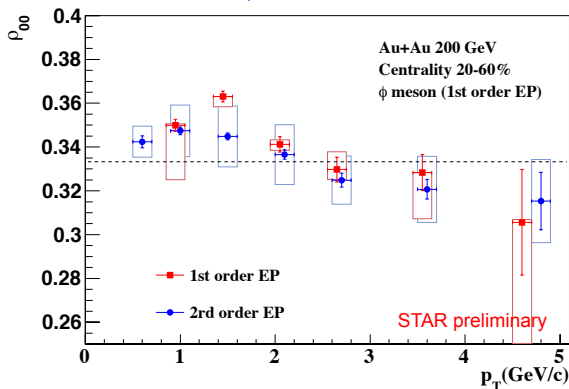
Ar-KCl at 1.76A GeV

(HADES Collaboration), PRC **84**, 014902 (2011)

- ▶ NA60: $\lambda_\theta \simeq 0$, but large error bars
- ▶ HADES: large polarization $\lambda_\theta \simeq 0.5$

ϕ meson polarization (STAR)

$$\lambda_\theta = \frac{3\rho_{00} - 1}{1 - \rho_{00}}, \quad \rho_{00} = \rho_L$$



A. Tang, Chirality Workshop, Florence 2018; C. Zhou QM2018

- ▶ Noncentral collisions: large global angular momentum
⇒ Vorticity ⇒ Particle polarization
- ▶ Vorticity or just thermalized medium?

Summary

- ▶ Anisotropy coefficients as a tool to understand heavy-ion collisions
- ▶ Virtual photons from (unpolarized) thermal sources **are polarized**
- ▶ Collective flow affects shape of anisotropy coefficients
- ▶ Realistic models give large polarization

Outlook

- ▶ Analyze different elementary reactions
- ▶ Anisotropic momentum distributions $\Rightarrow$ nonequilibrium
- ▶ Effect of vorticity and magnetic field (polarized medium)

BACKUP

Diagonal form of the spin-density matrix

$$\rho = \frac{1}{3} \begin{pmatrix} 1 + \frac{3}{2}P_z + \sqrt{\frac{3}{2}}T_{zz} & 0 & 0 \\ 0 & 1 - \sqrt{6}T_{zz} & 0 \\ 0 & 0 & 1 - \frac{3}{2}P_z + \sqrt{\frac{3}{2}}T_{zz} \end{pmatrix}$$

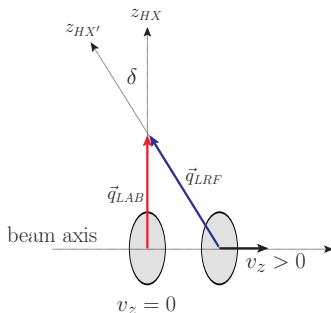
- ▶ In unpolarized system $P_z = 0$, but often $T_{zz} \neq 0$,
i.e., no vector but tensor polarization!
- ▶ In general vector and tensor polarization axes can be different

Medium and flow

Static uniform medium

- ▶ Photon rest frame: fluid velocity $\vec{v}$ opposite to photon "direction"
- ▶ Only $\lambda_\theta \neq 0$

Longitudinal Bjorken expansion



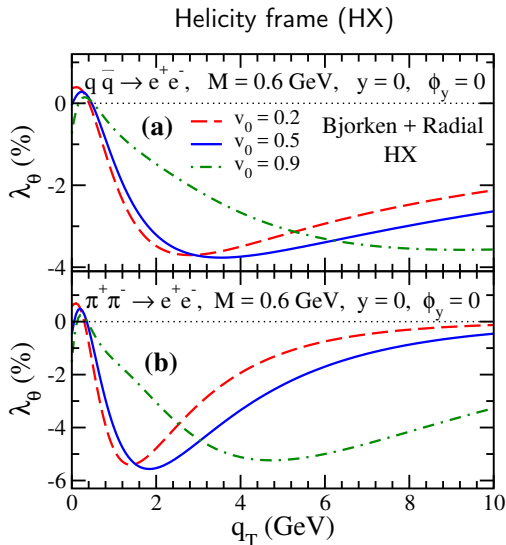
- ▶ $v_z = z/t$ along beam axis
- ▶ Photon polarized along $z_{HX'}$, defined by its momentum in local rest frame
- ▶ Rotation δ between $z_{HX'}$ and z_{HX} (Wick helicity rotation)
 $\Rightarrow \lambda_\theta, \lambda_\phi, \lambda_{\theta\phi} \neq 0$
- ▶ Frame invariant combination:

$$\tilde{\lambda}_\theta = \frac{\lambda_\theta + 3\lambda_\phi}{1 - \lambda_\phi}$$

Longitudinal Bjorken + Radial expansion

- ▶ All coefficients $\lambda_\theta, \lambda_\phi, \lambda_{\theta\phi}, \lambda_\phi^\perp, \lambda_{\theta\phi}^\perp \neq 0$

Results (Bjorken + radial expansion)



- ▶ $v_\perp = v_0 r_\perp / R_0$ (Transverse to beam axis)
- ▶ The position of the minimum shifts towards higher q_T as v_0 increases