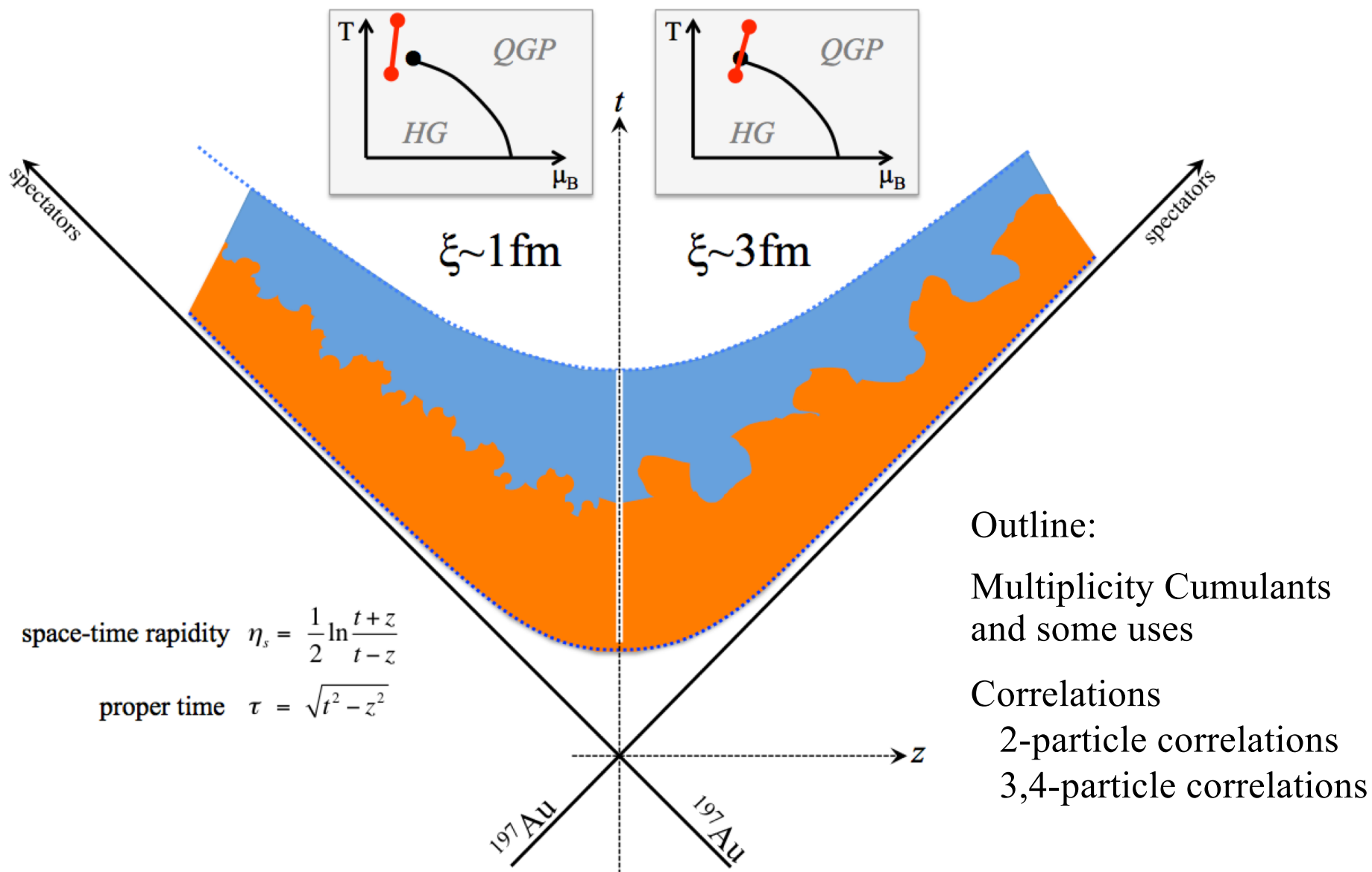
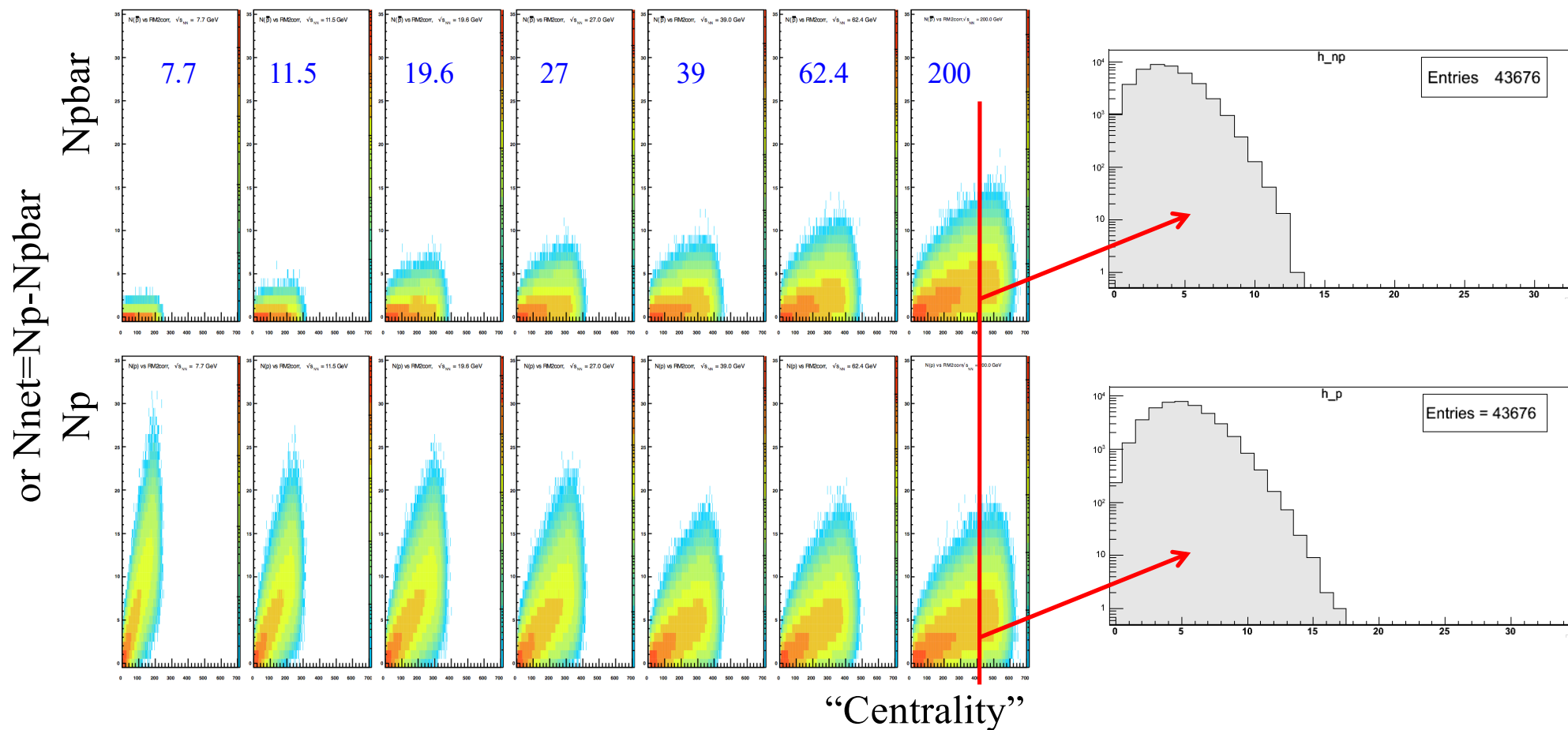


Fluctuations and Correlations in the STAR BES-I Data

W.J. Llope, Wayne State University





$$K_1 = \langle x \rangle \quad \delta x = x - \langle x \rangle$$

$$K_2 = \langle \delta x^2 \rangle$$

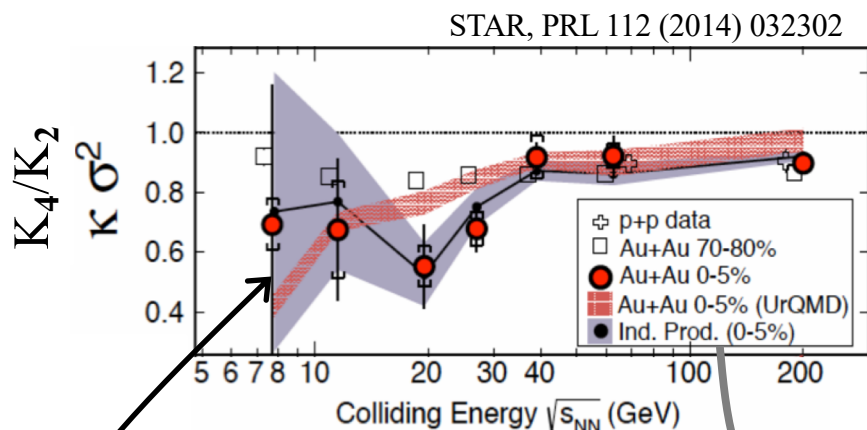
$$K_3 = \langle \delta x^3 \rangle$$

$$K_4 = \langle \delta x^4 \rangle - 3\langle \delta x^2 \rangle^2$$

Efficiency corrections assuming efficiencies
are binomially distributed...

STAR calculates weighted-averages of K_k and
K-ratios within 5%- or 10%-wide centrality bins

STAR net-p multiplicity cumulant ratios



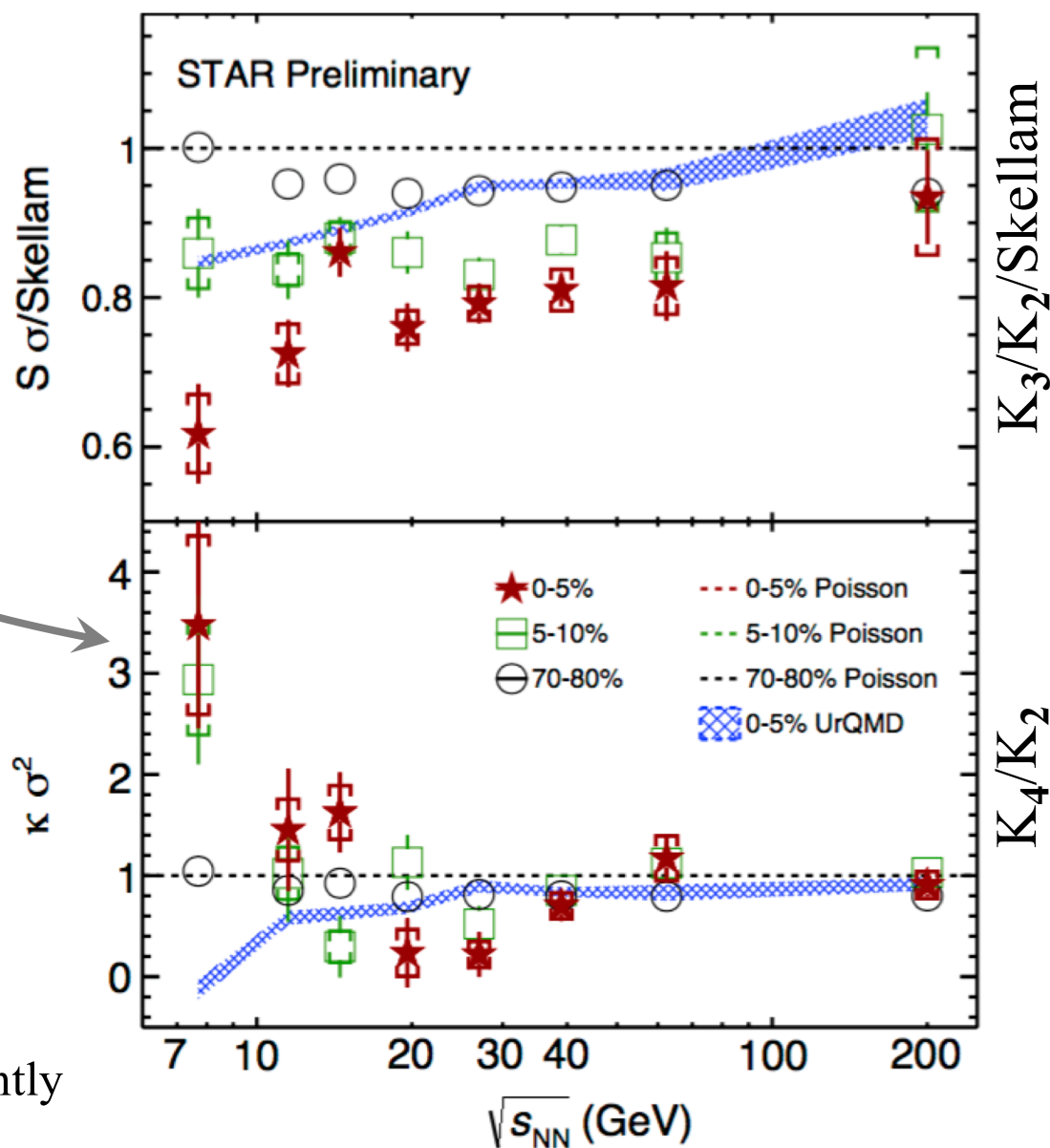
Widen the acceptance
from $0.4 < P_T < 0.8$ to $0.4 < P_T < 2.0$...

Gray band: “Independent Production”

$$K_{\text{net}} = K_{\text{pos}} + (-1)^k K_{\text{neg}}$$

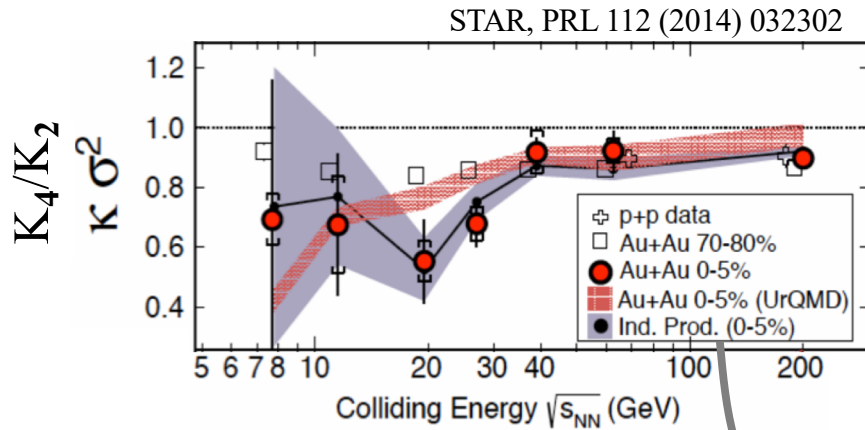
Also holds for $0.4 < P_T < 2.0$

→ numbers of protons are not significantly
correlated with the number of antiprotons.
(Does not rule out CP).



A wider acceptance increased the multiplicities
and made the deviations from Poisson larger

STAR net-p multiplicity cumulant ratios



Widen the acceptance
from $0.4 < P_T < 0.8$ to $0.4 < P_T < 2.0$...

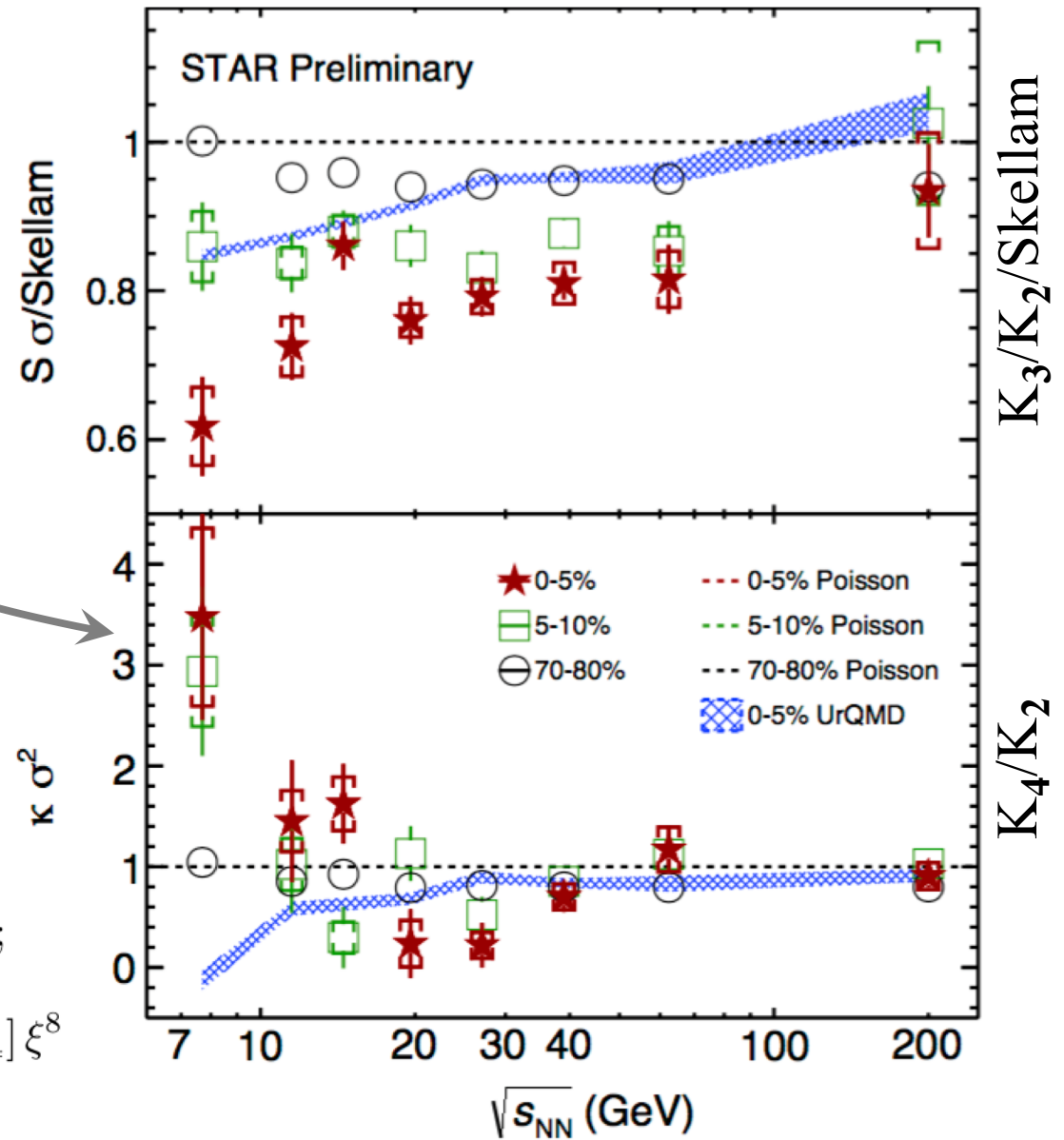
Critical point searches:

$$\kappa_2 = \langle \sigma_0^2 \rangle = \frac{T}{V} \xi^2; \quad \kappa_3 = \langle \sigma_0^3 \rangle = \frac{2\lambda_3 T}{V} \xi^6;$$

$$\kappa_4 = \langle \sigma_0^4 \rangle_c \equiv \langle \sigma_0^4 \rangle - \langle \sigma_0^2 \rangle^2 = \frac{6T}{V} [2(\lambda_3 \xi)^2 - \lambda_4] \xi^8$$

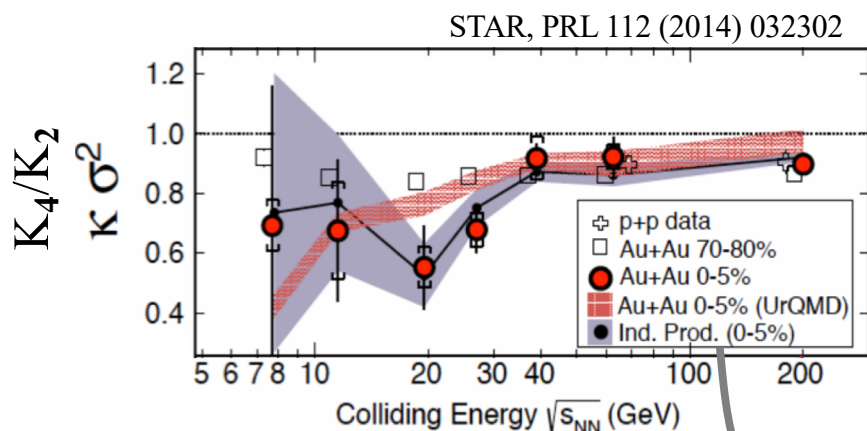
Sensitivity increases with k
in k-particle fluctuations.

M. Stephanov

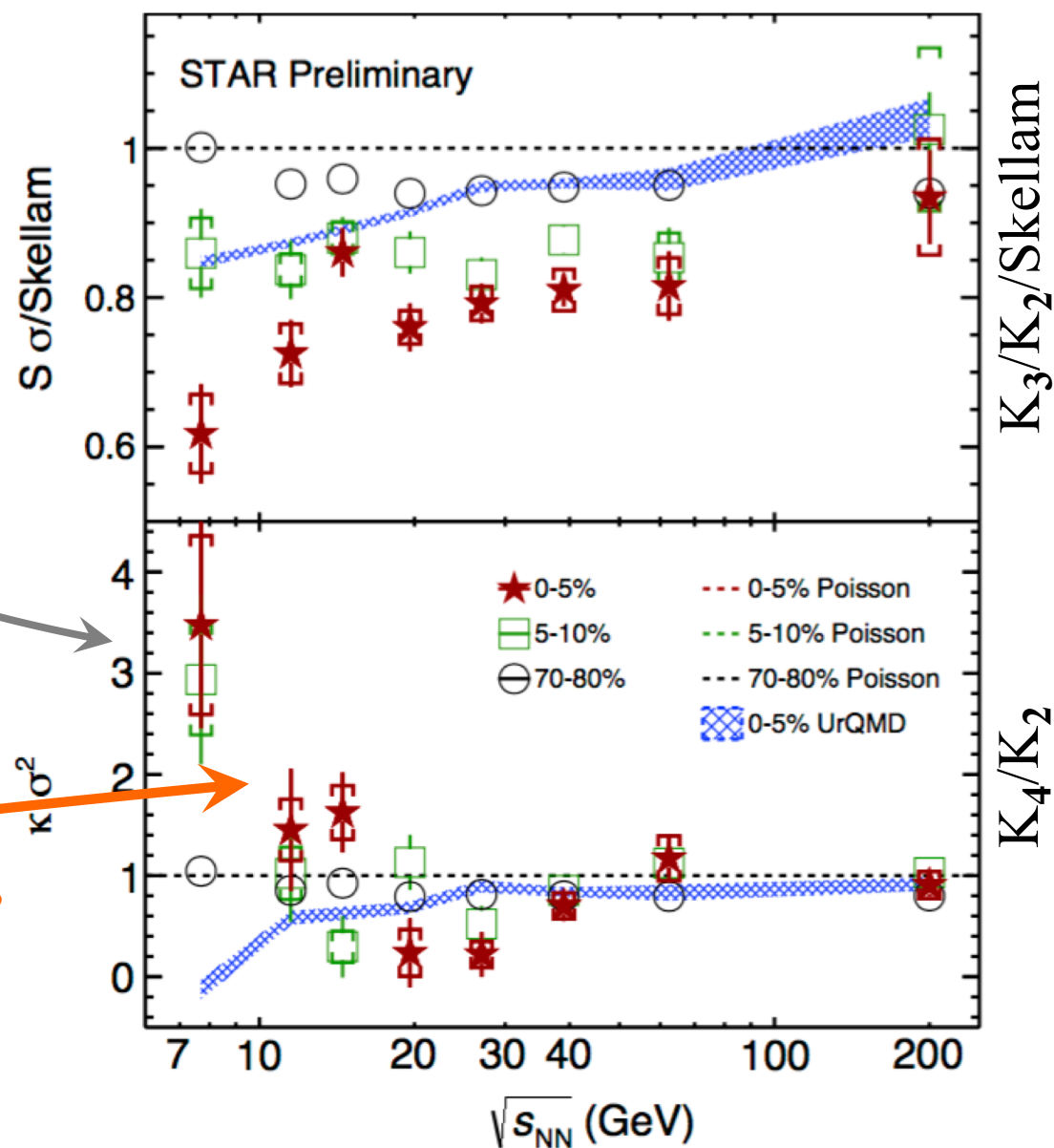
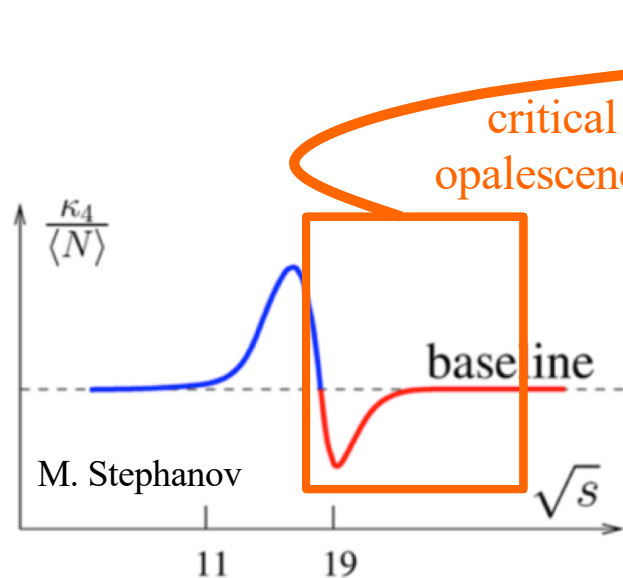


A wider acceptance increased the multiplicities
and made the deviations from Poisson larger

STAR net-p multiplicity cumulant ratios

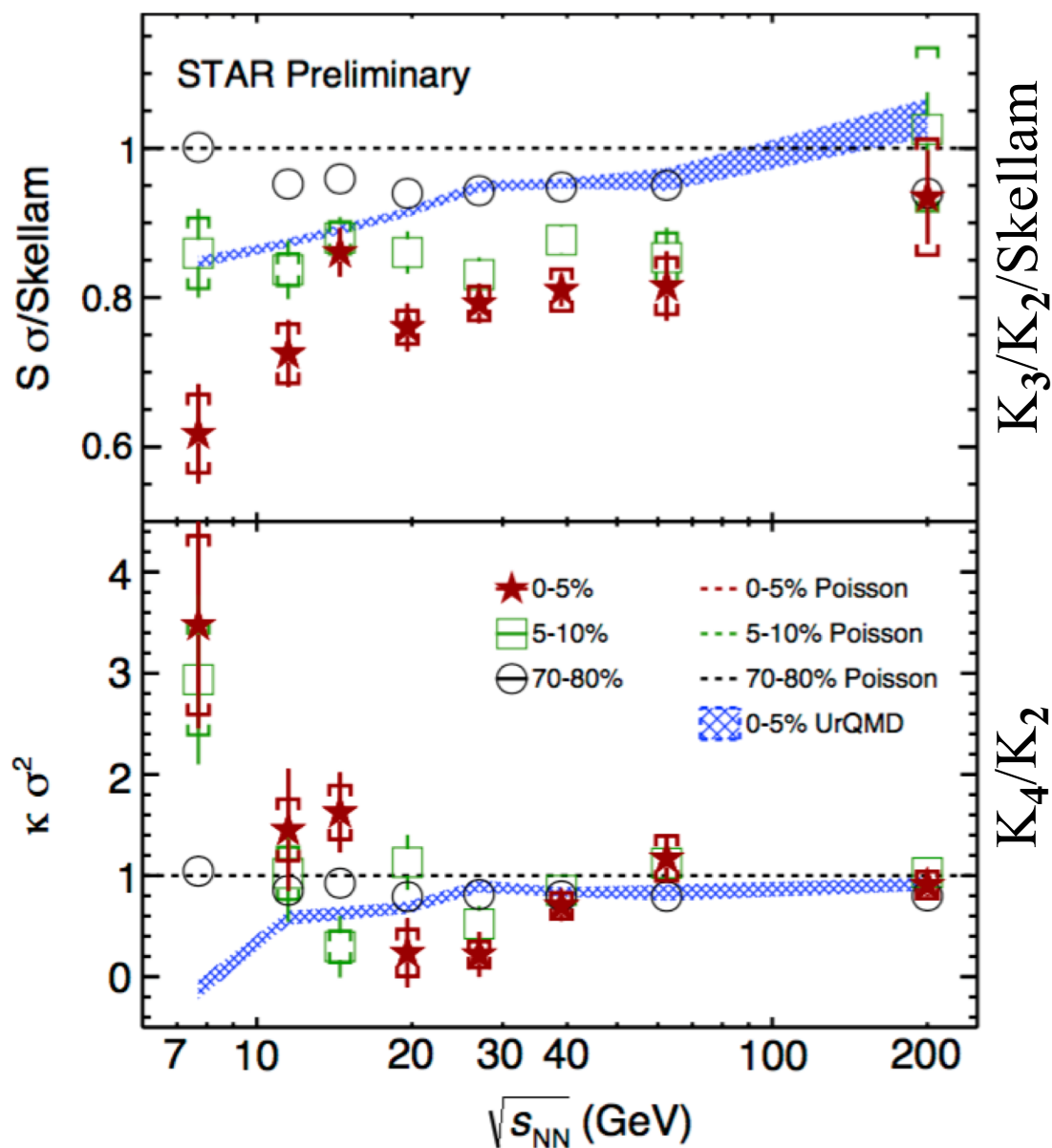
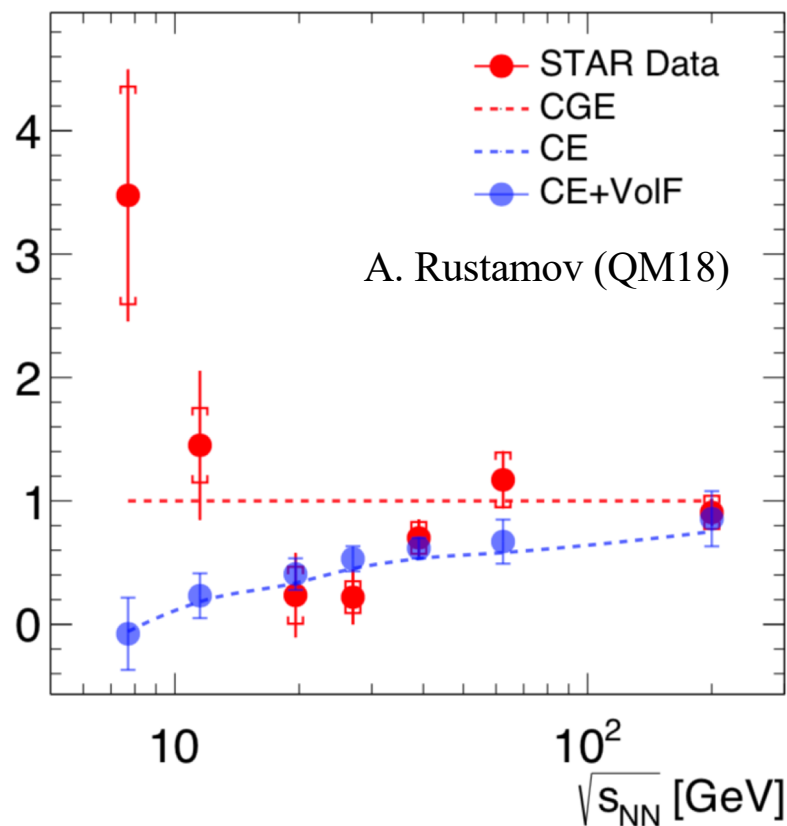
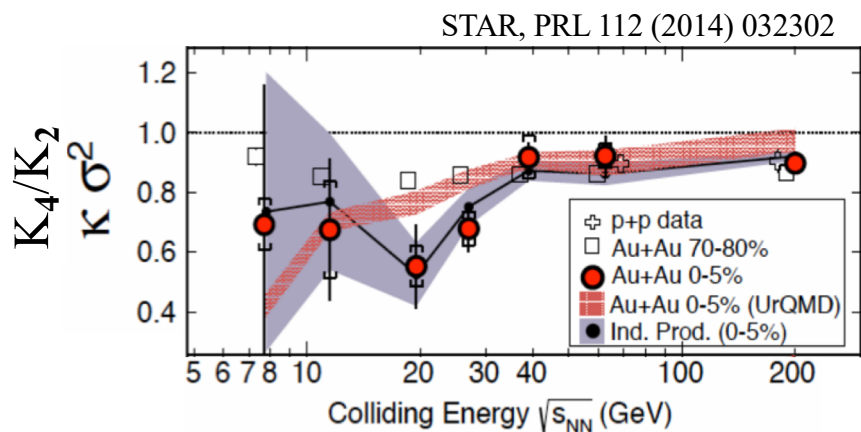


Widen the acceptance
from $0.4 < P_T < 0.8$ to $0.4 < P_T < 2.0$...



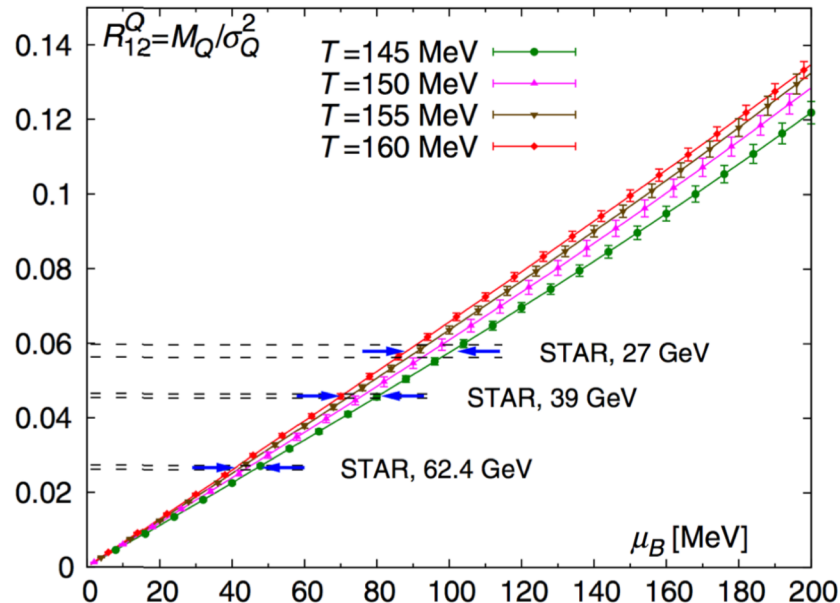
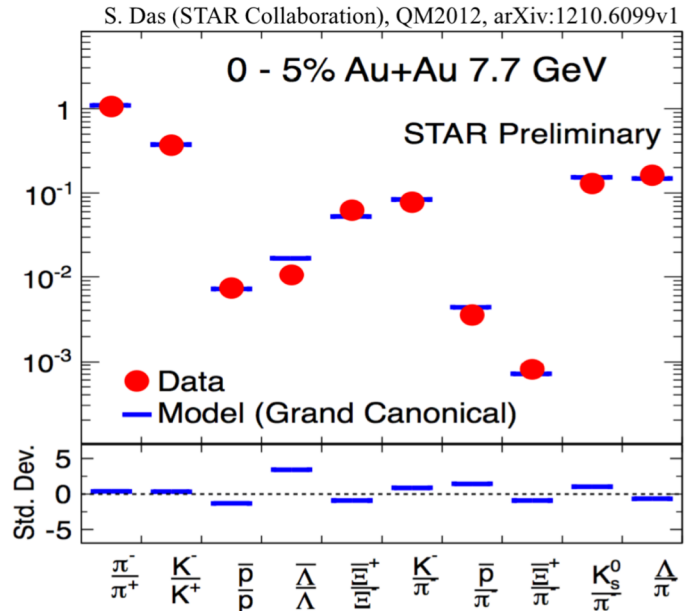
A wider acceptance increased the multiplicities
and made the deviations from Poisson larger

STAR net-p multiplicity cumulant ratios



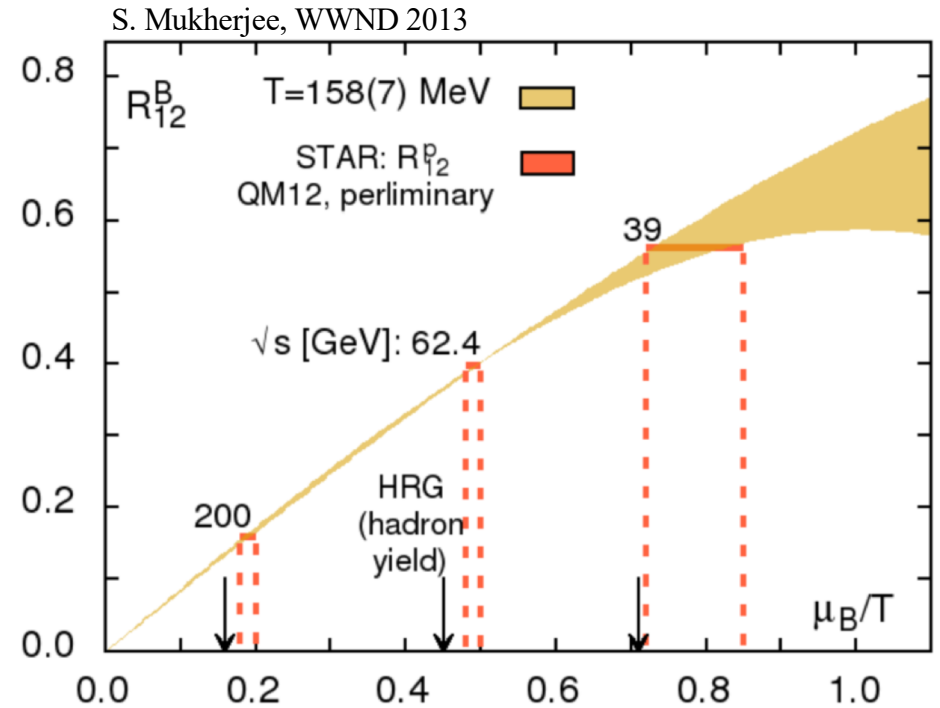
A wider acceptance increased the multiplicities
and made the deviations from Poisson larger

...with statistical hadronization models (THERMUS, FIST, etc.)



...with Lattice QCD

Basically, assume K ratios are χ ratios



A. Bazavov *et al.*, PRL 109, 192302 (2012)

S. Borsányi *et al.*, PRL 111, 062005 (2013)

Comparisons to HRG can include
acceptance/cuts.

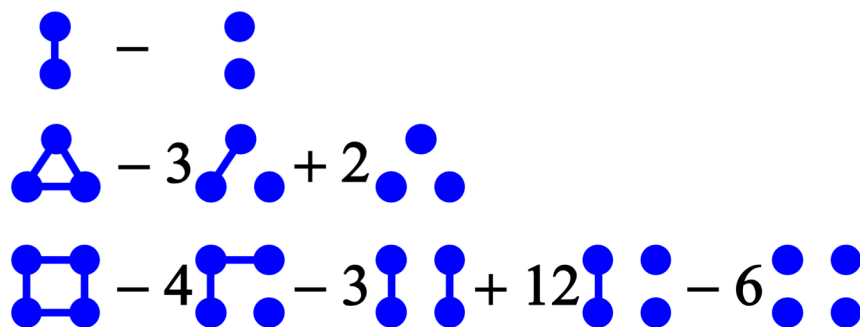
(C. Ratti *et al.*)

“Correlation Functions”

$$C_2 = \rho_2 - \rho_1 \rho_1$$

$$C_3 = \rho_3 - 3\rho_2\rho_1 + 2\rho_1\rho_1\rho_1$$

$$C_4 = \rho_4 - 4\rho_3\rho_1 - 3\rho_2\rho_2 + 12\rho_2\rho_1\rho_1 - 6\rho_1\rho_1\rho_1\rho_1$$



$$C_2 = C_2(\Delta y, \Delta \varphi)$$

$$C_2 = C_2(y_1, y_2)$$

$$C_3 = C_3(y_1, y_2, y_3)$$

$$C_4 = C_4(y_1, y_2, y_3, y_4)$$

Explicit subtraction of lower-order correlations...

“Multiplicity Cumulants”

$$K_1 = \langle N \rangle$$

$$K_2 = \langle N \rangle + c_2$$

$$K_3 = \langle N \rangle + 3c_2 + c_3$$

$$K_4 = \langle N \rangle + 7c_2 + 6c_3 + c_4$$

E.L. Berger, NPB **85**, 61 (1975)

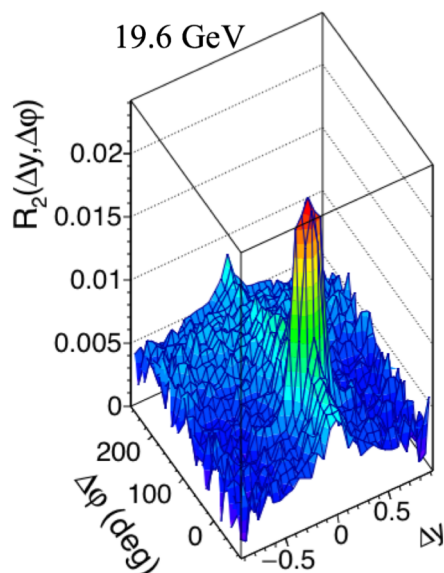
P. Carruthers *et al.*, PRL **63**, 1562 (1989)

P. Carruthers, PRA **43**, 2632 (1991)

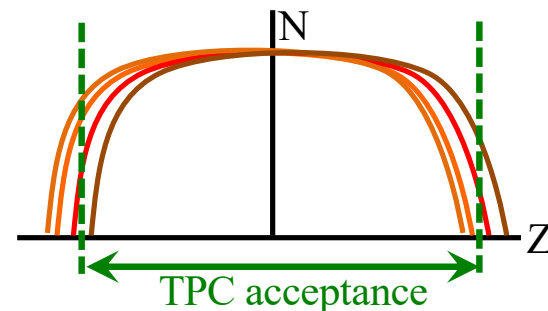
A. Bzdak *et al.*, PRC **95**, 054906 (2017)

$$c_k = \int C_k(y_1, \dots, y_k) dy_1 \dots dy_k$$

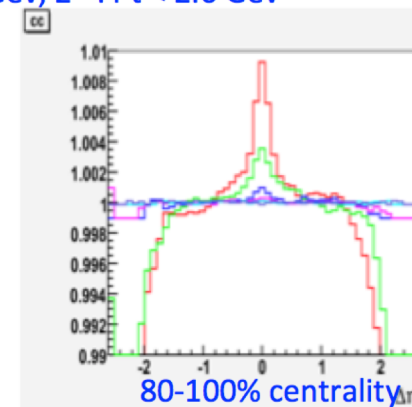
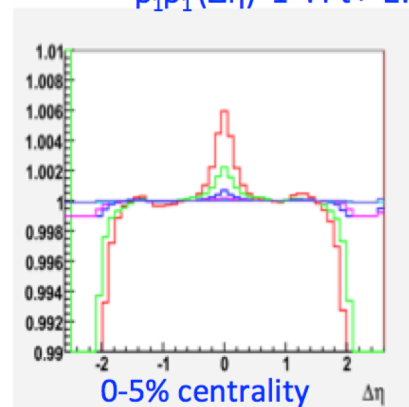
A few possible advantages from going through the correlations C_k to get the fluctuations K_k



If K_k deviate from Poisson, one can locate the kinematical location of these deviations by looking at C_k

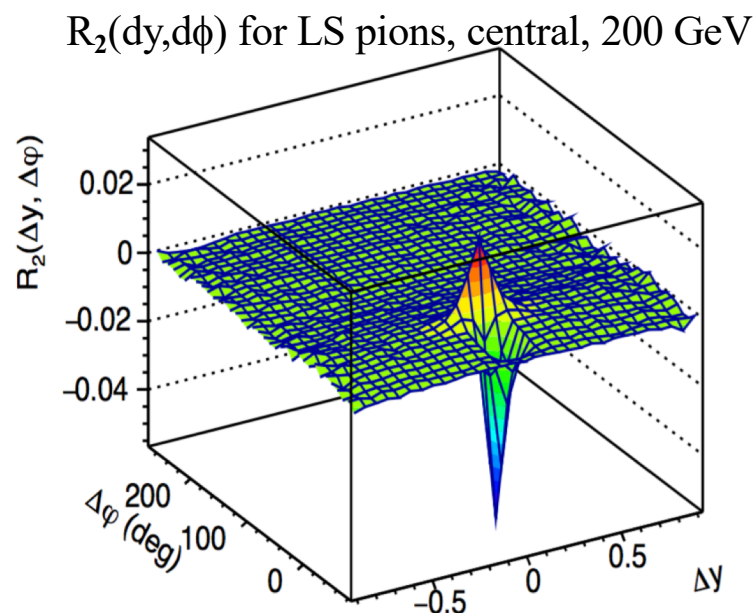


$\rho_1 \rho_1(\Delta\eta)$ 1st: $P_t > 2.0$ GeV, 2nd: $P_t < 2.0$ GeV



Zvtx-smearing pseudo-correlations:
 C_k **require** “Zvtx-averaging” over narrow (1-2cm) bins. Not done for available STAR K_k results.

(note STAR net-p results use 1m-wide Zvtx bins at 7.7 GeV)



“track crossing” is a *pair inefficiency*
which may not be entirely treated with some
single-particle efficiency correction techniques

Common to use “reduced” form of C_k to remove trivial pair multiplicity dependence

$$R_k = C_k(y_1, \dots, y_k) / \prod_{i=1}^k \rho_1(y_i) \qquad R_2 = \frac{C_2}{\rho_1 \rho_1} = \frac{\rho_2 - \rho_1 \rho_1}{\rho_1 \rho_1} = \frac{\rho_2}{\rho_1 \rho_1} - 1$$

With the integral:

$$r_k = \frac{\int \rho_1(y_1) \dots \rho_1(y_k) R_k(y_1, \dots, y_k) dy_1 \dots dy_k}{\int \rho_1(y_1) \dots \rho_1(y_k) dy_1 \dots dy_k} = c_k / \langle N \rangle^k$$

Then the multiplicity cumulants can equivalently be obtained via:

$$K_1 = \langle N \rangle$$

$$K_2 = \langle N \rangle + \langle N \rangle^2 r_2$$

$$K_3 = \langle N \rangle + 3 \langle N \rangle^2 r_2 + \langle N \rangle^3 r_3$$

$$K_4 = \langle N \rangle + 7 \langle N \rangle^2 r_2 + 6 \langle N \rangle^3 r_3 + \langle N \rangle^4 r_4$$

Alternate approach to measuring multiplicity cumulants that retains all the knowledge of full correlation functions underneath...

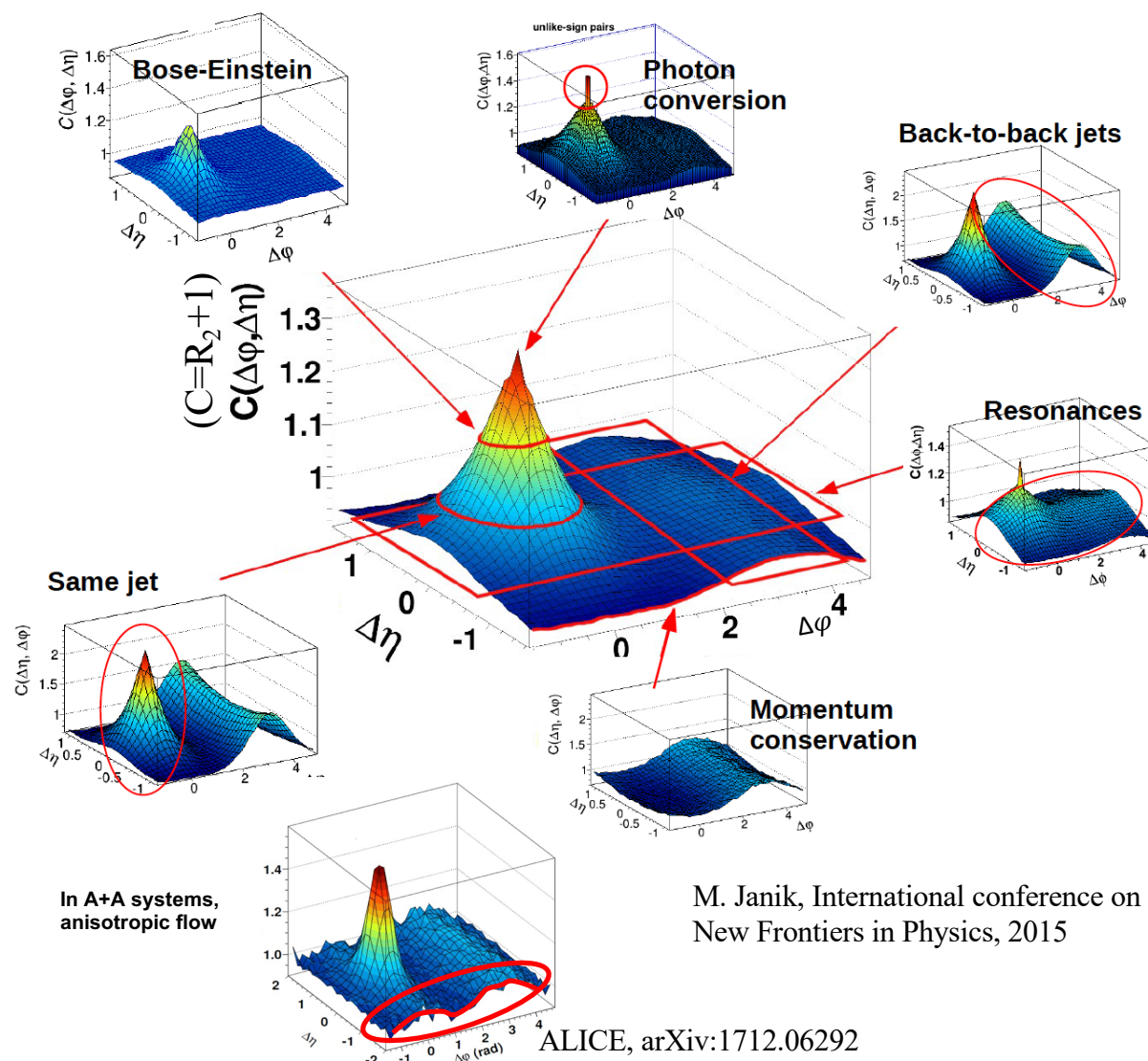
$$R_2(\Delta y, \Delta \varphi) = \frac{\rho_2(\Delta y, \Delta \varphi)}{\rho_1(y_1, \varphi_1) \otimes \rho_1(y_2, \varphi_2)} - 1$$

Same correlator normalized to
number of uncorrelated pairs

- uncorrelated: $R_2 = 0$
- correlations: $R_2 > 0$
- anticorrelations: $R_2 < 0$

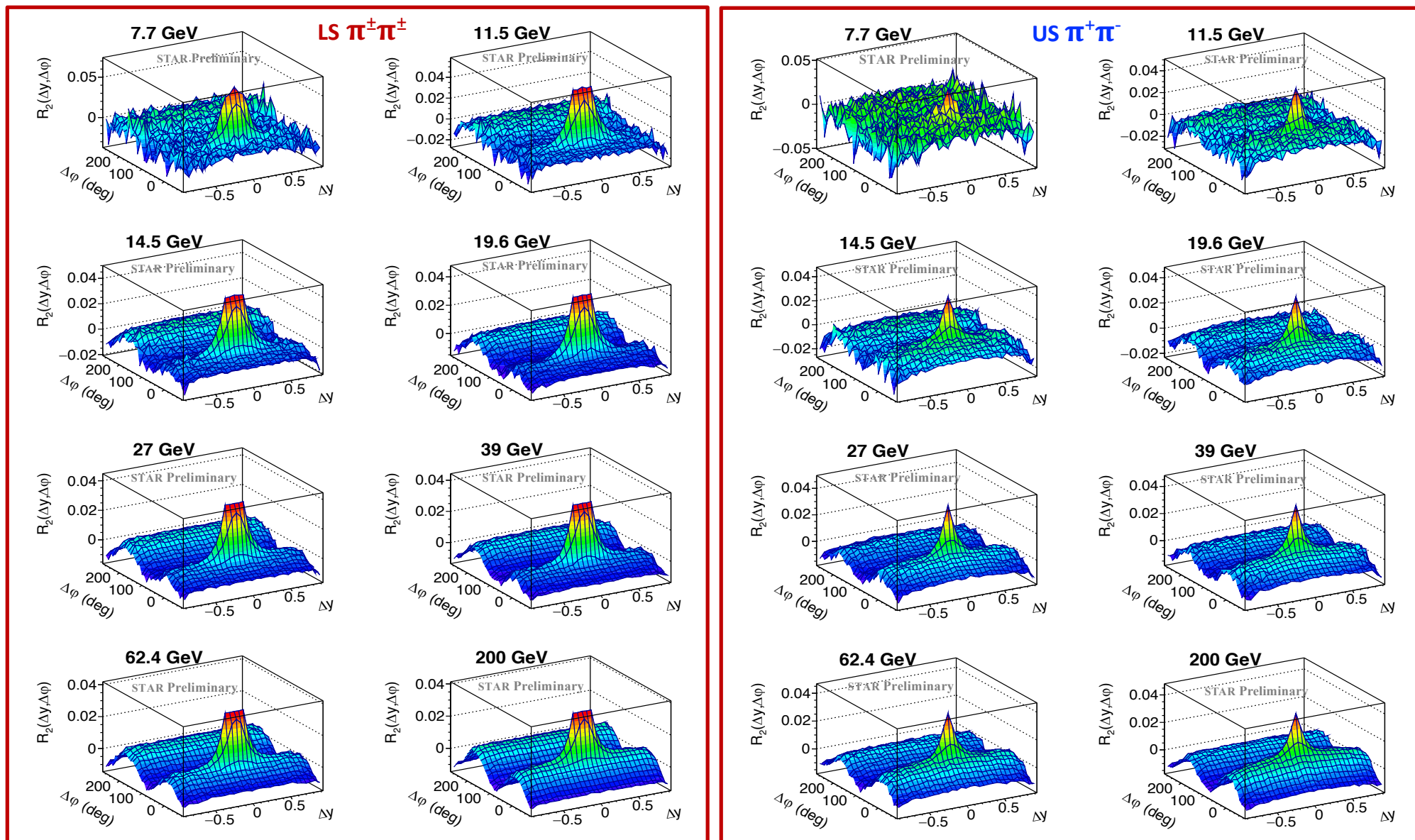
- Different physical phenomena exhibit different “shapes” in the $R_2(\Delta y, \Delta \phi)$ distributions

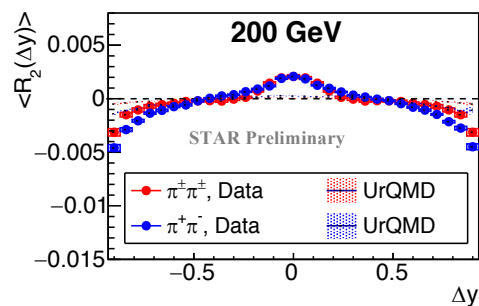
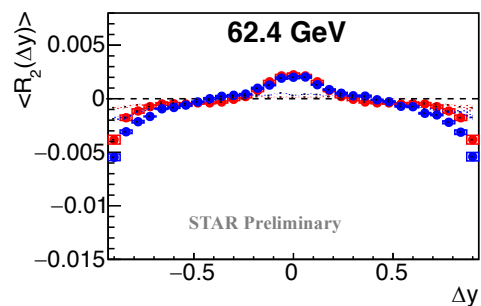
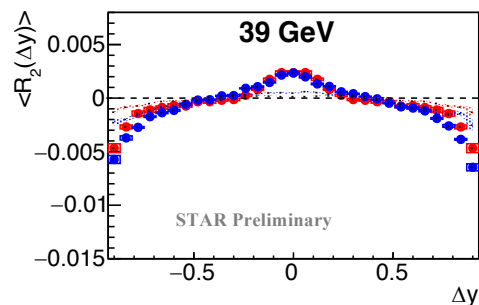
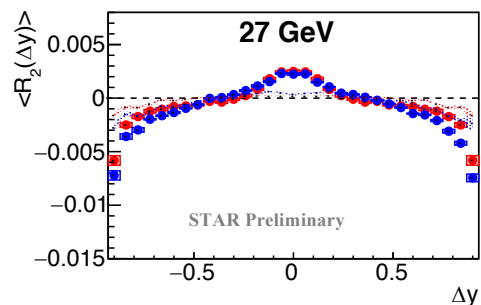
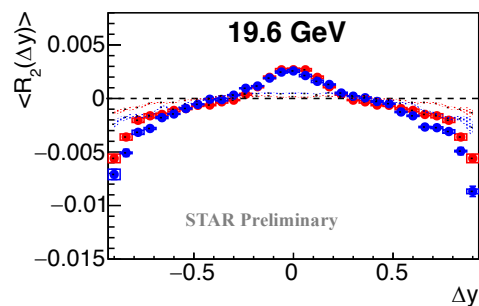
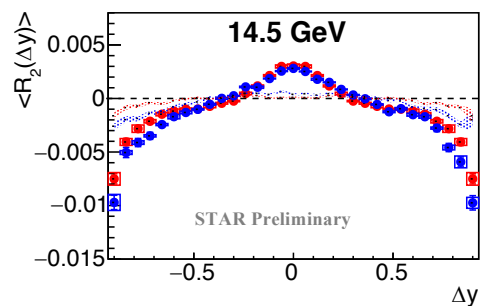
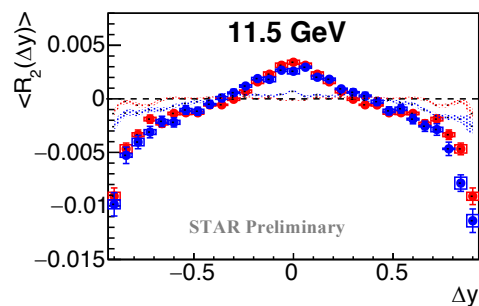
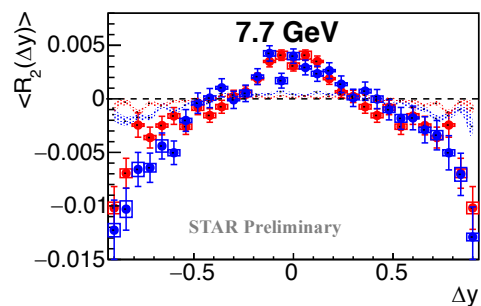
- Short-range correlations (resonance decays, jet fragmentation, Bose-Einstein correlations)
- Long-range correlation (di-jets, the “ridge,” conservation laws)
- QCD Critical Point?



M. Janik, International conference on
New Frontiers in Physics, 2015

ALICE, arXiv:1712.06292



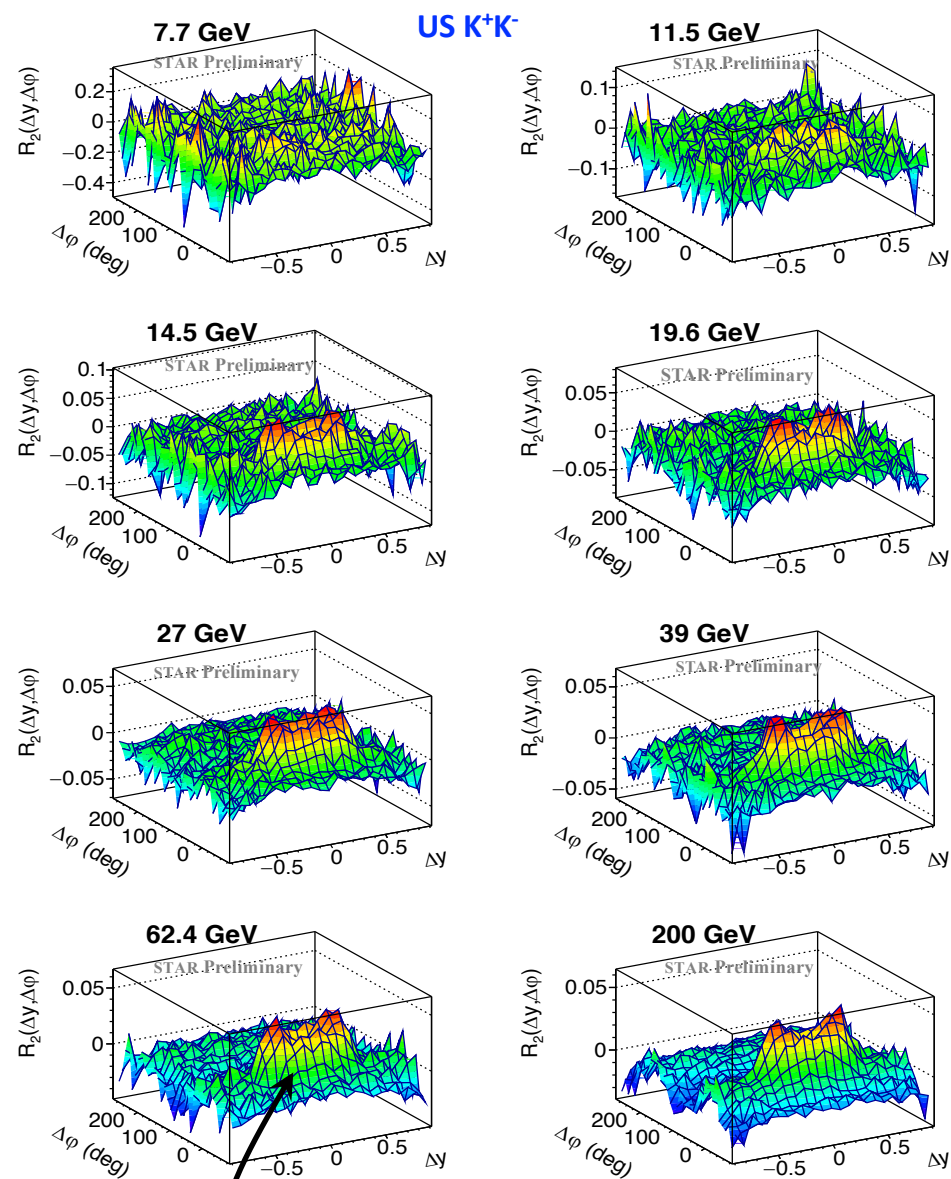
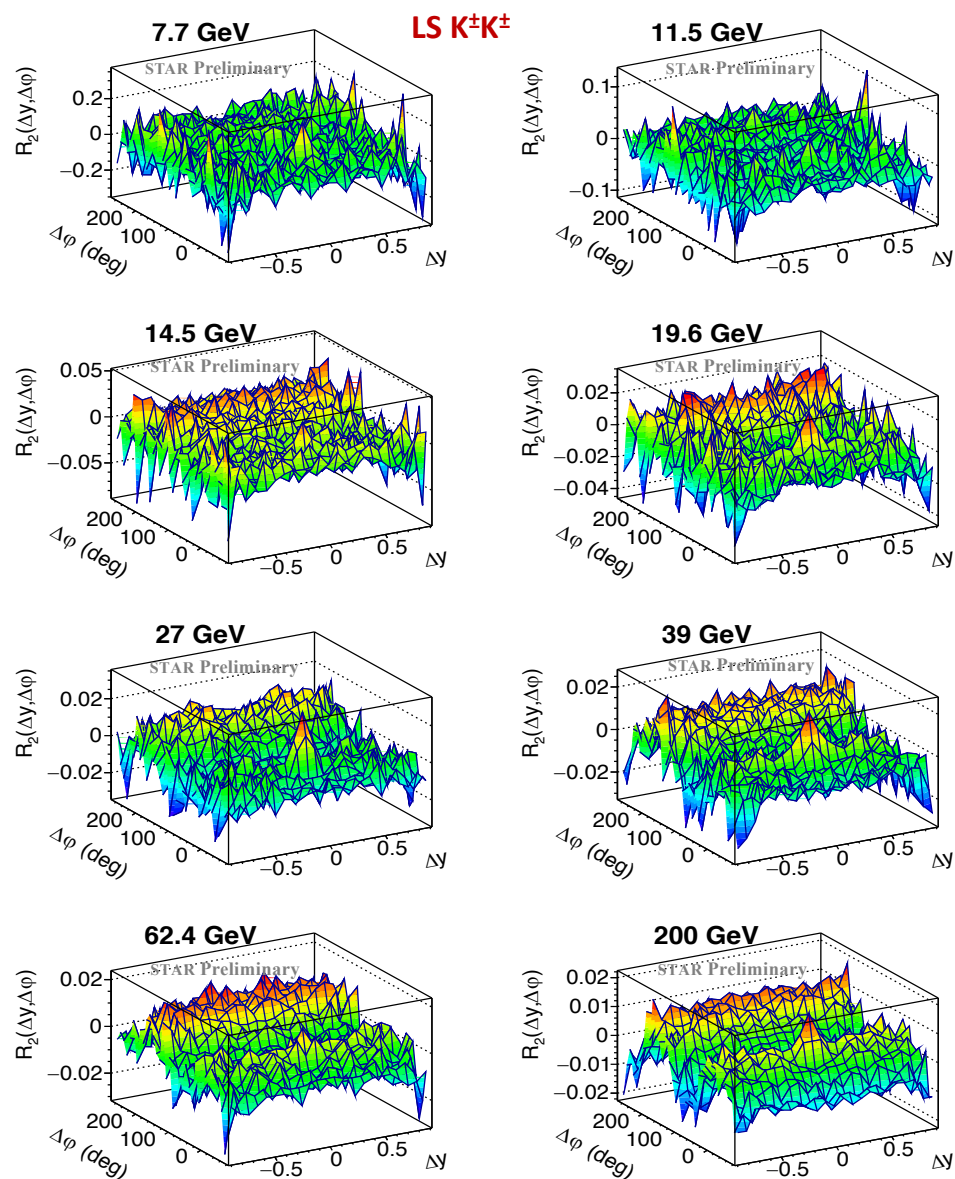


Near-side enhancement in $\pi^+\pi^+$ & $\pi^+\pi^-$
SRC from jet fragmentation and
resonance decays

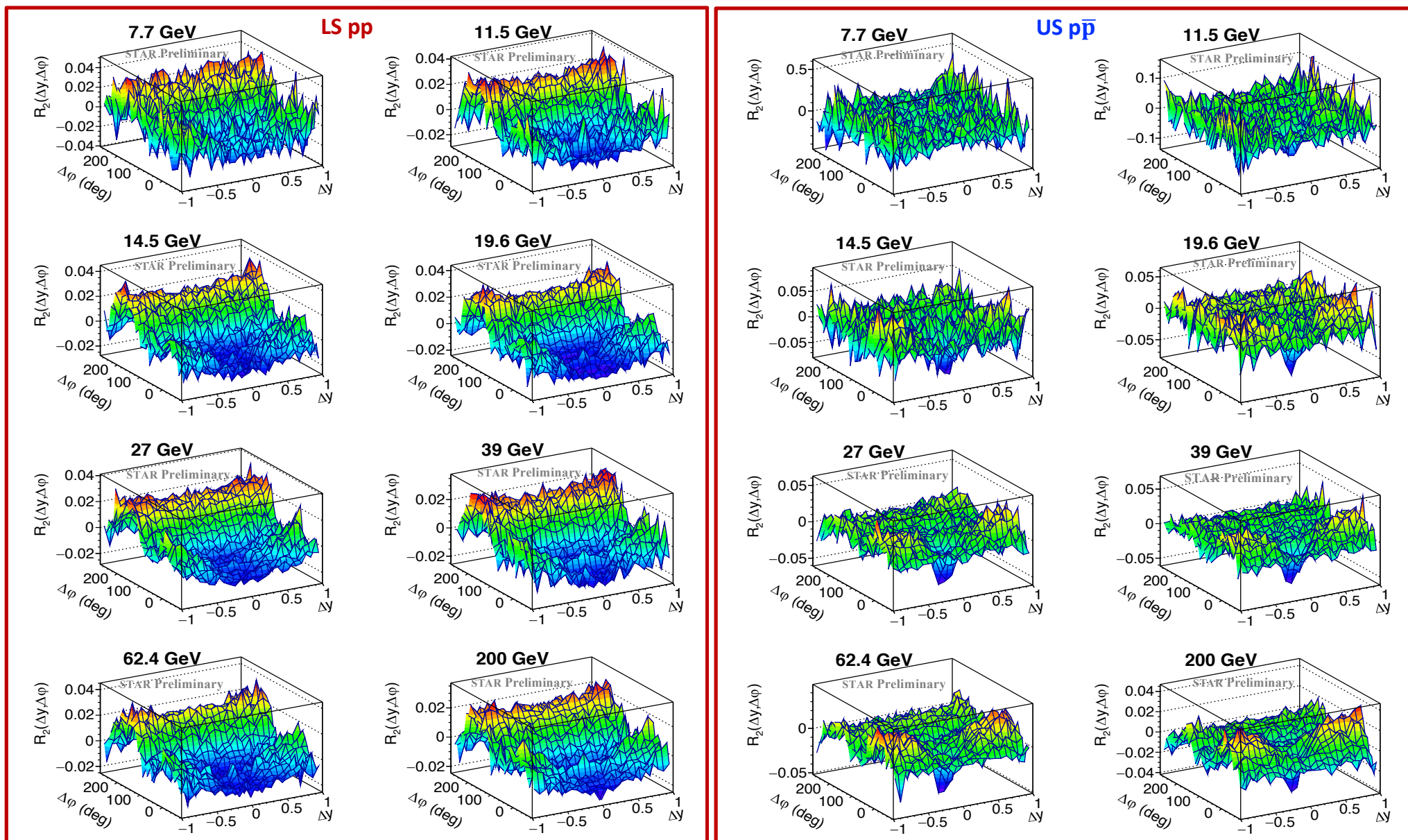
Small differences between like-sign &
unlike-sign charge combinations

Correlations decrease slightly with
increasing beam energy

UrQMD does not reproduce the
observed SRC



$\phi(1020) \rightarrow K^+ + K^-$



Strong, nearly beam energy independent, near-side anticorrelation in LS protons.
US results suggest proton annihilation, but LS proton results much longer range...

A near-side anticorrelation is observed in
2-proton correlations

Observed for the first time in an A+A system

Anti-correlation in like-sign baryons has been
seen in small systems:

e^+e^- collisions at 29 GeV by TPC/Two-
Gamma Collaboration (SLAC)

H. Aihara et al. PRL 57 (1986) 3140

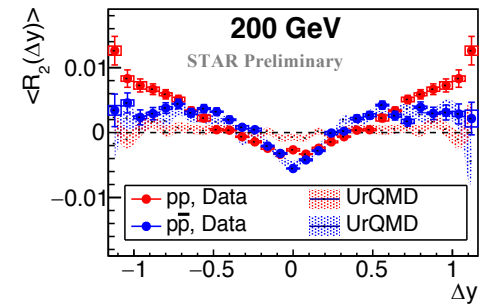
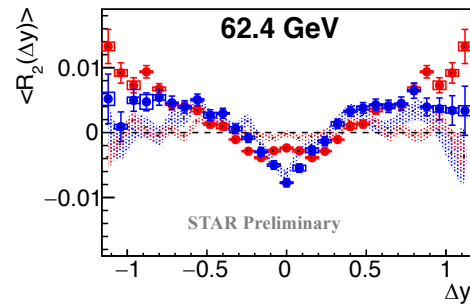
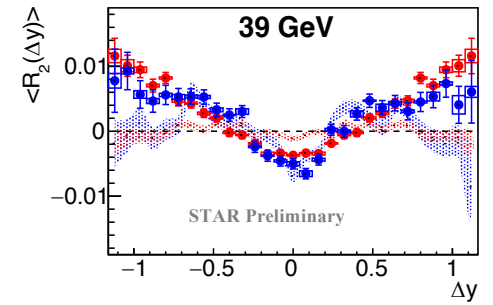
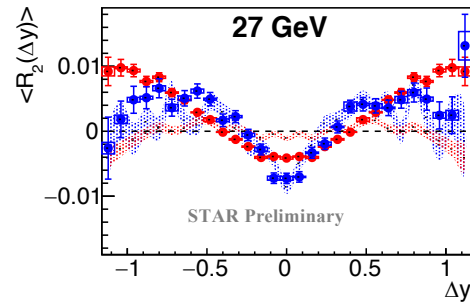
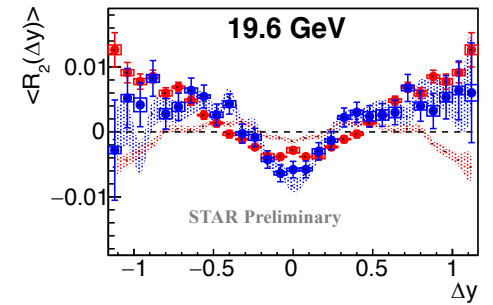
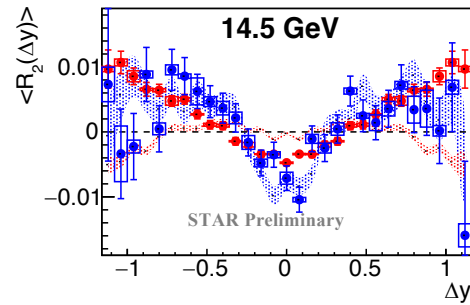
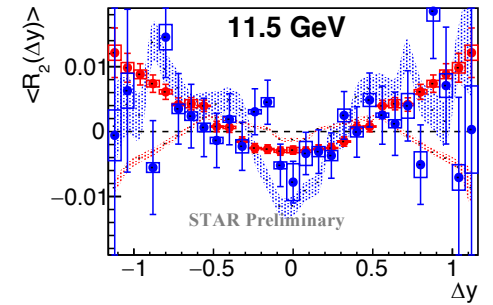
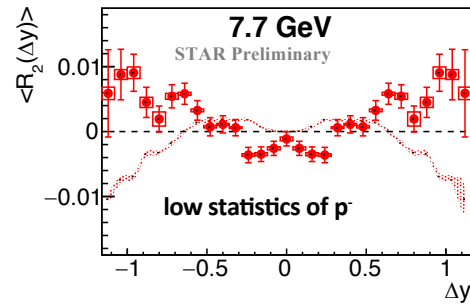
$p+p$ at $\sqrt{s_{NN}}=7$ TeV by ALICE (LHC)

ALICE Coll., EPJ. C (2017) 77:569

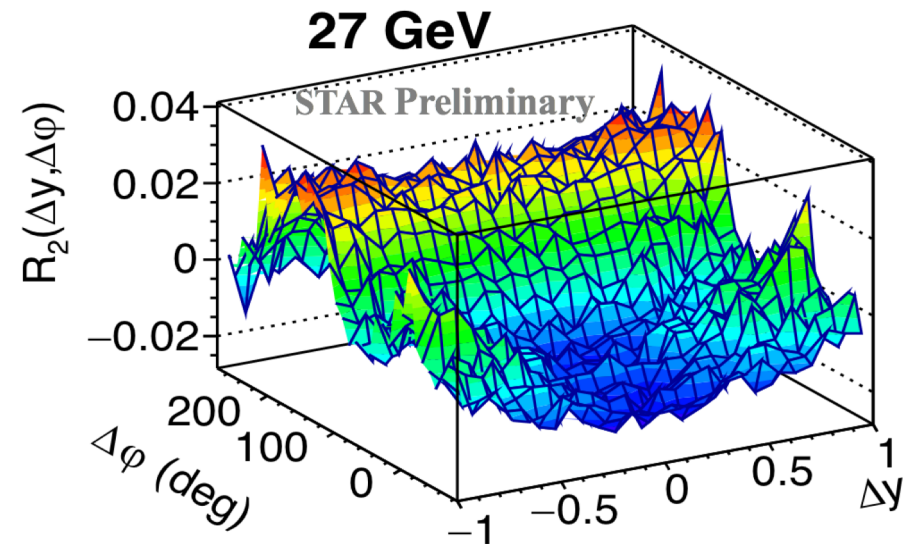
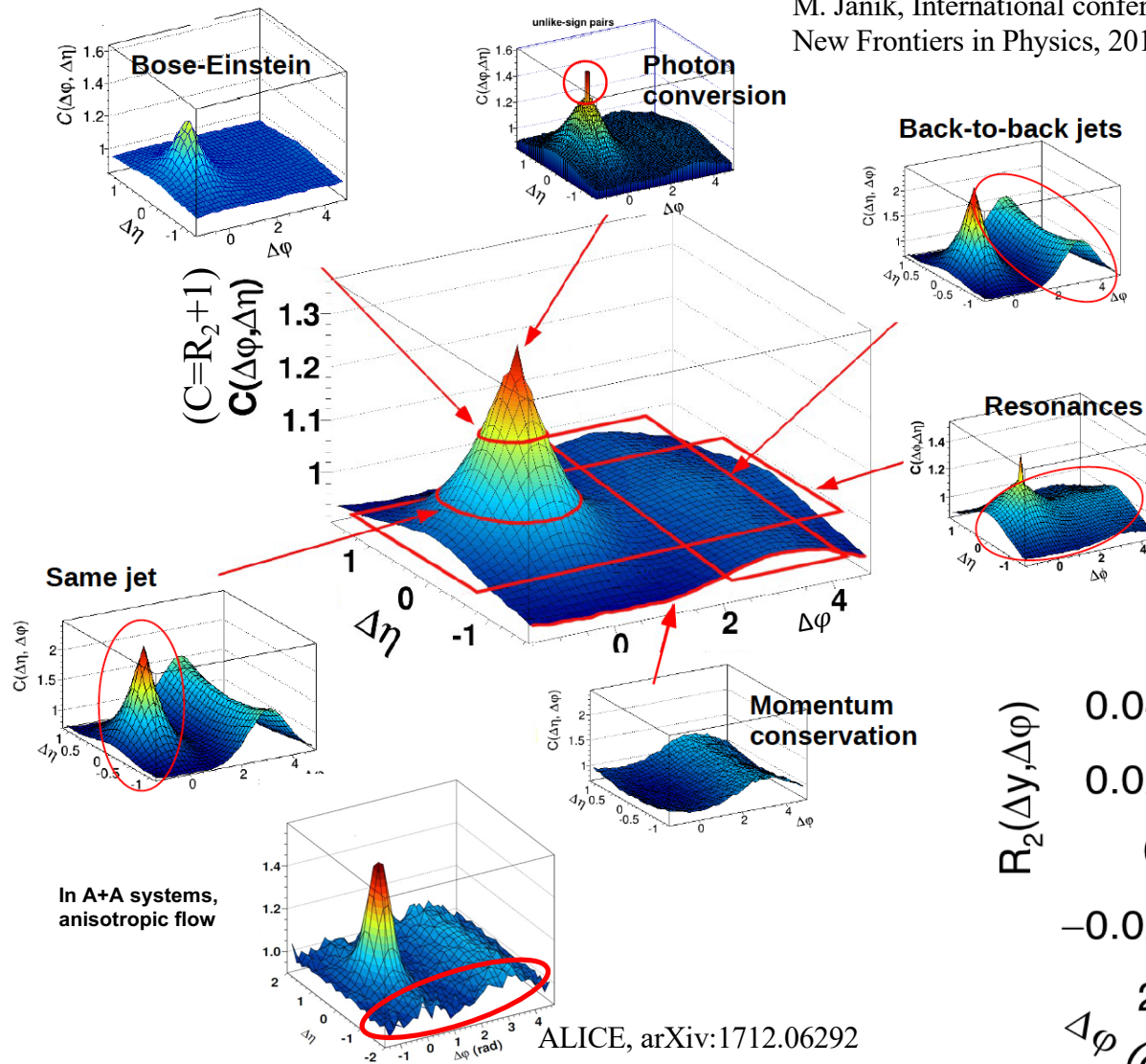
The anticorrelation is observed in both like-sign
and unlike-sign charge combinations

It is narrower in Δy for $p\bar{p}$ w.r.t. pp

UrQMD reproduces the observed
anti-correlation in unlike-sign pairs,
but not like-sign pairs...



M. Janik, International conference on
New Frontiers in Physics, 2015



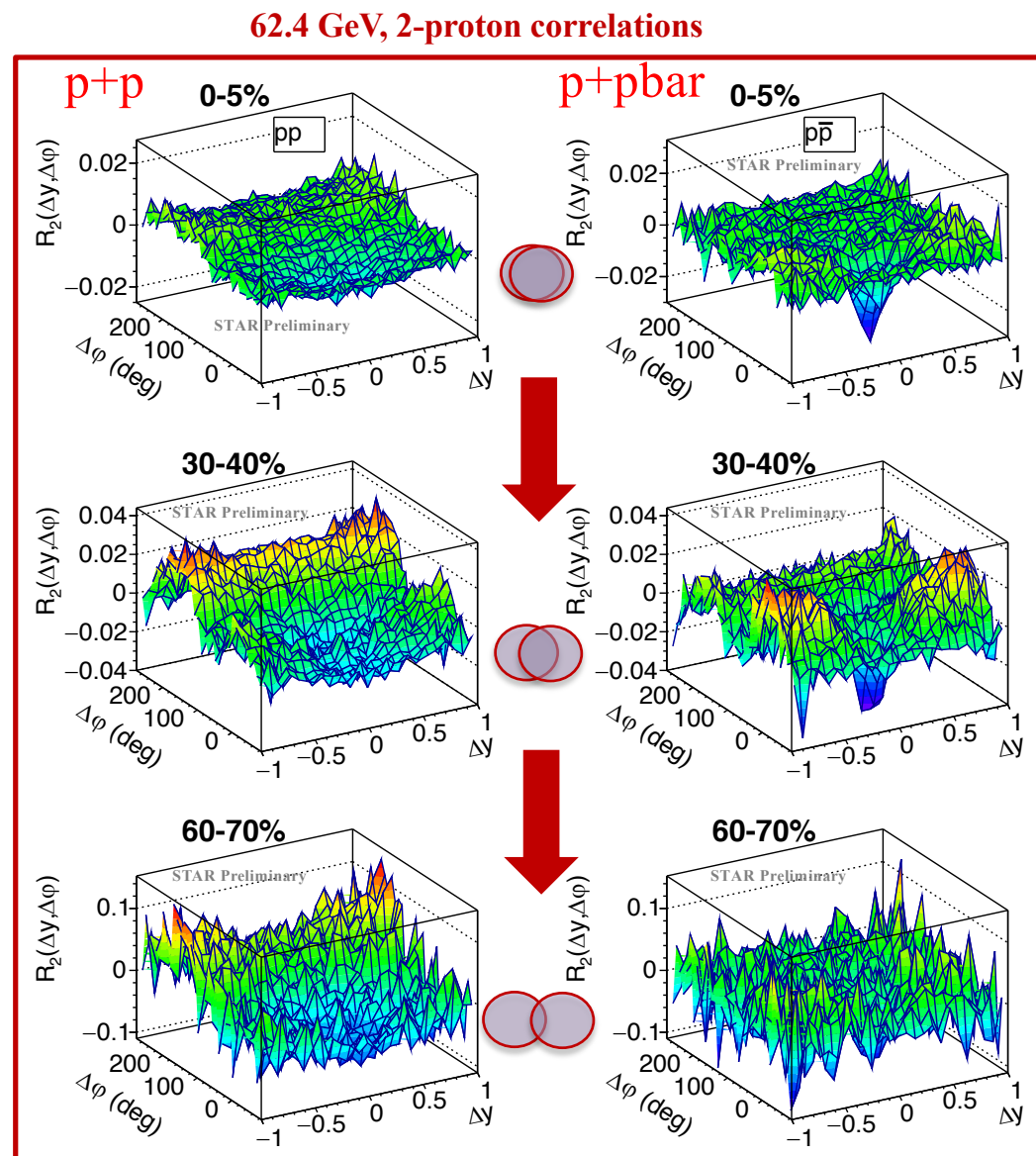
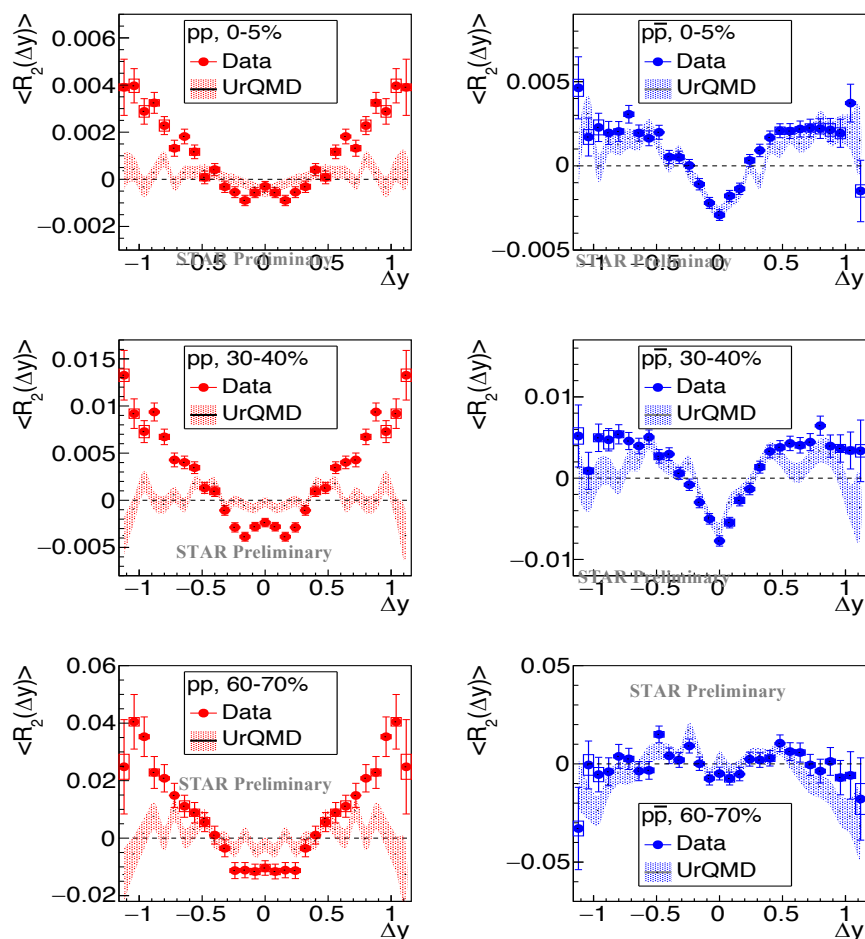
see Malgorzata Janik's excellent CERN seminar on the ALICE results:

https://mediastream.cern.ch/MediaArchive/Video/Public/WebLectures/2017/632396/632396_mobile_480p_1000.mp4

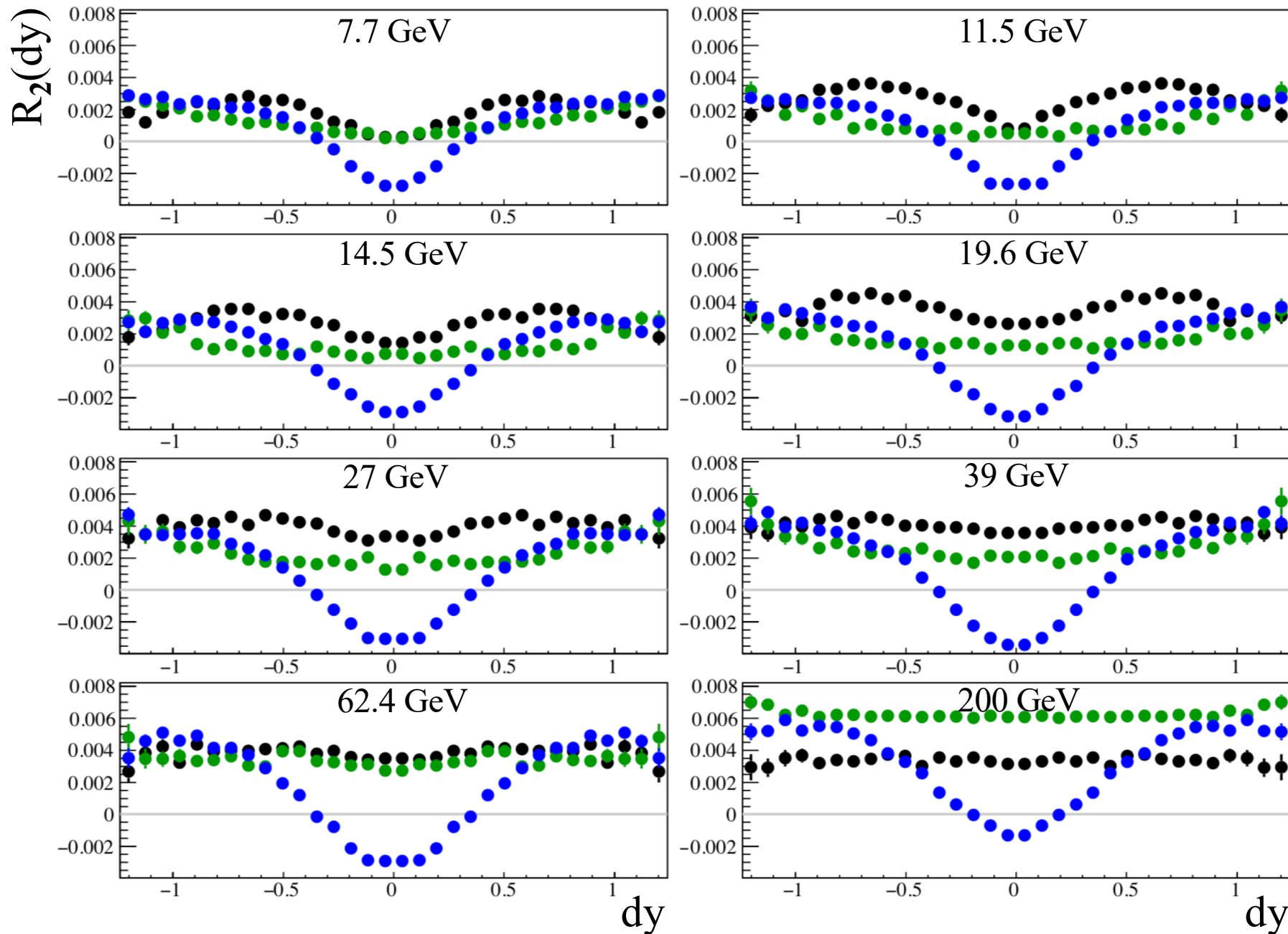
Anticorrelations increase with decreasing centrality. Insignificant dependence on P_T

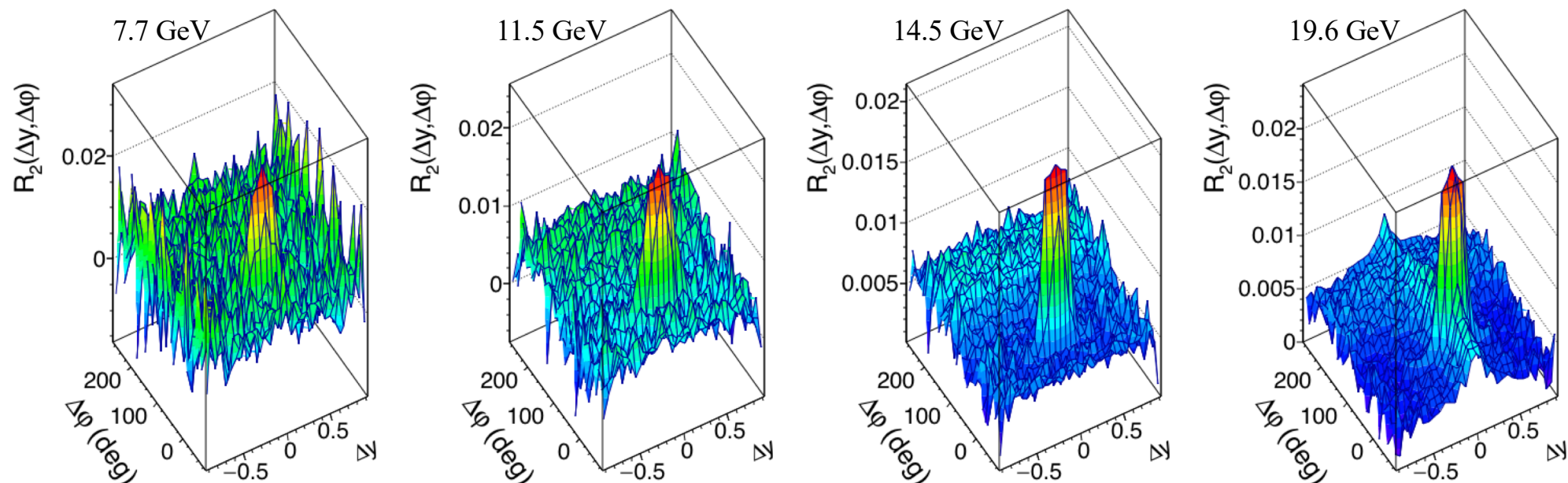
For all centralities, the anticorrelation in unlike-sign proton pairs $p\bar{p}$ is narrower in Δy w.r.t. like-sign pairs pp

UrQMD simulation well reproduces the anti-correlation for unlike-sign proton pairs at the different centralities



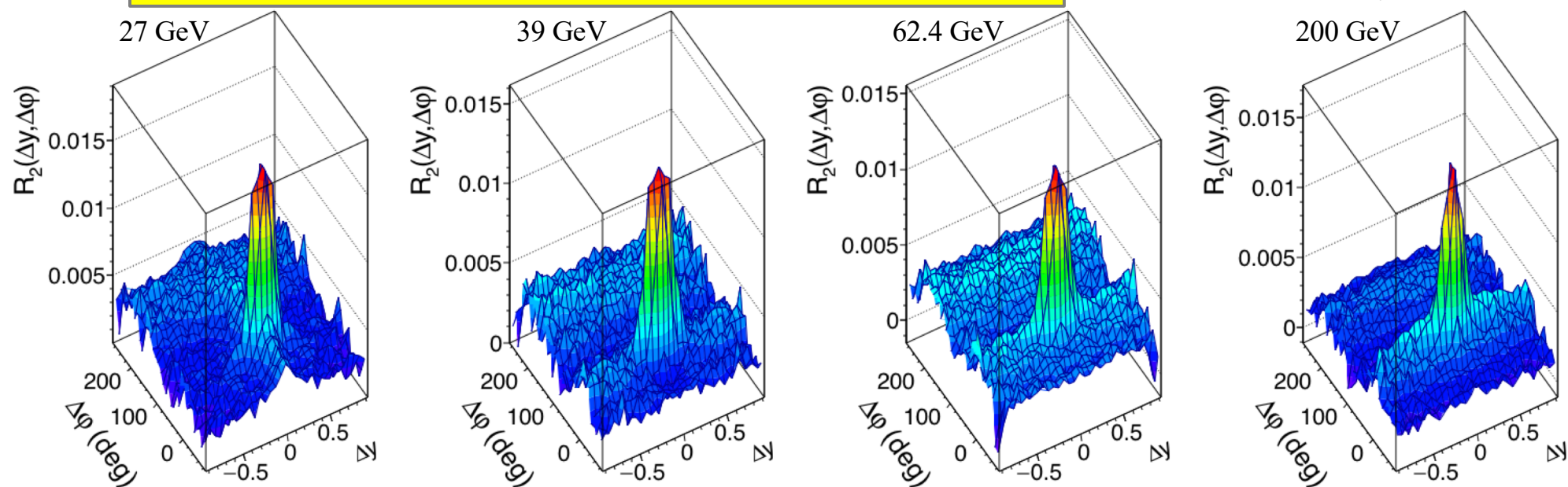
UrQMD 3.4 Hijing 1.411 AMPT 2.26t7b

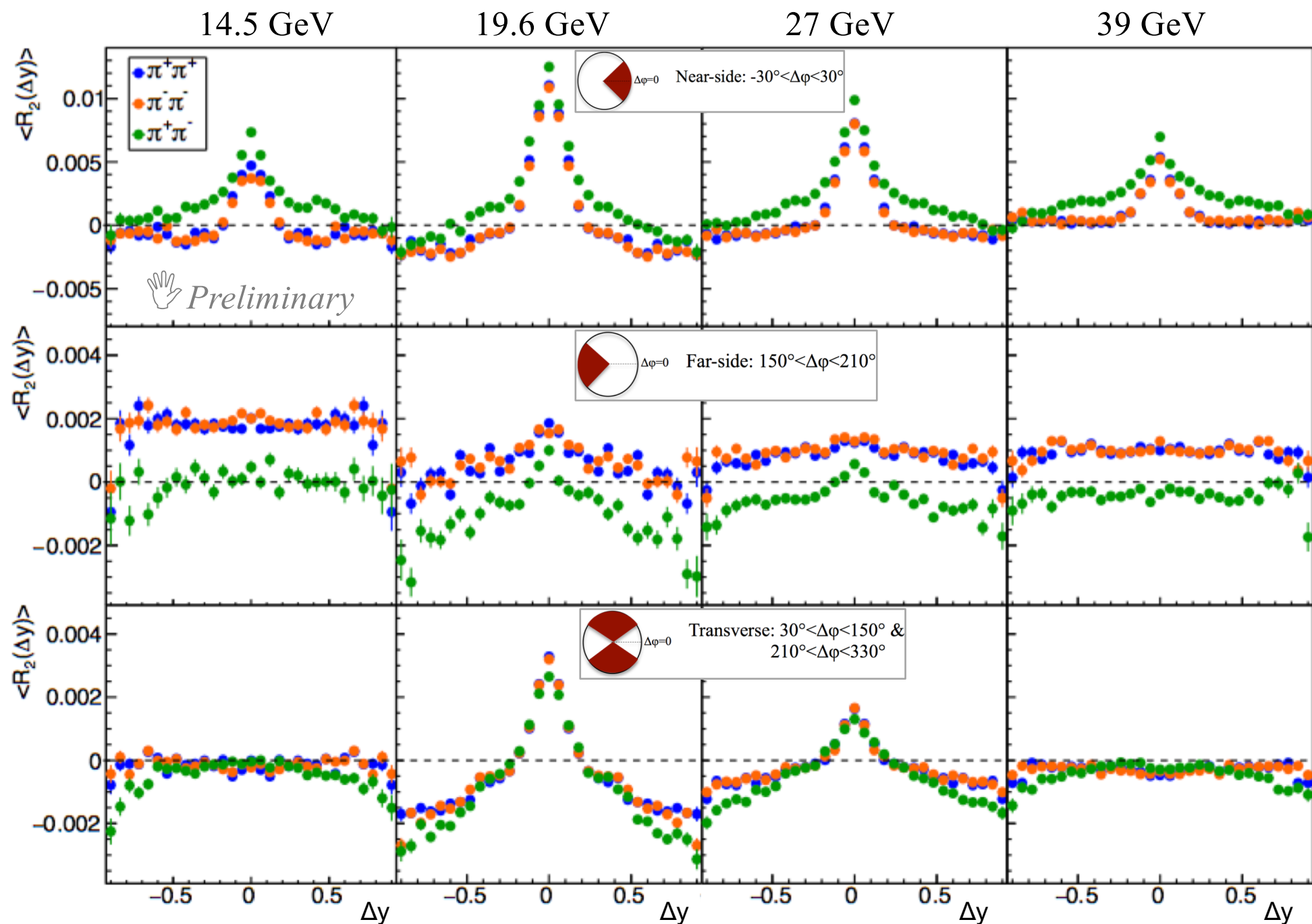




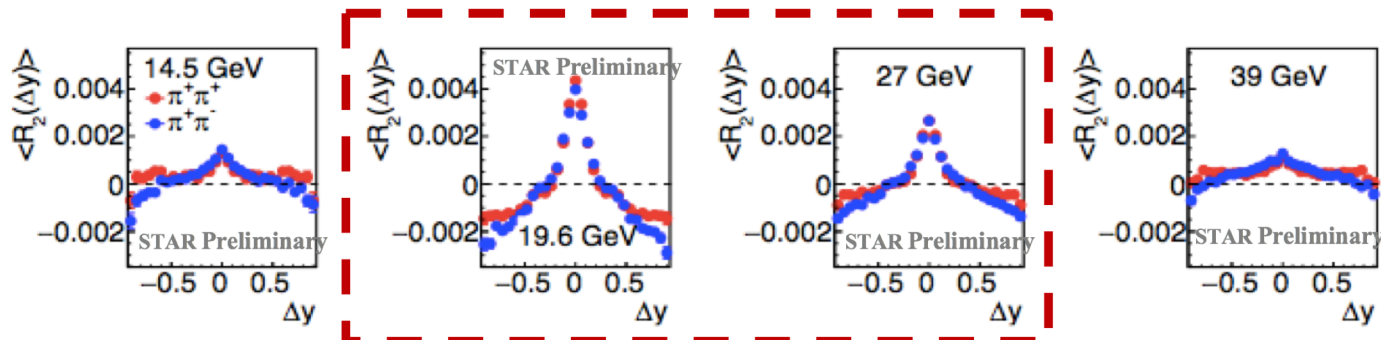
Older 2-pion CF results shown at QM17 and CPOD17

STAR Preliminary



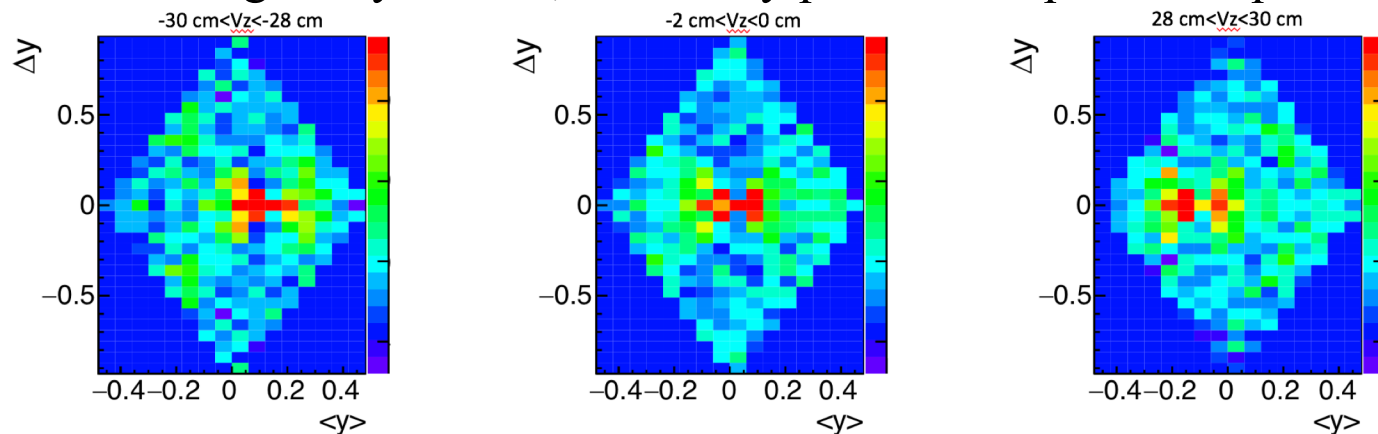


QM/CPOD 2017 results based on PID: TPC+TOF in all p range



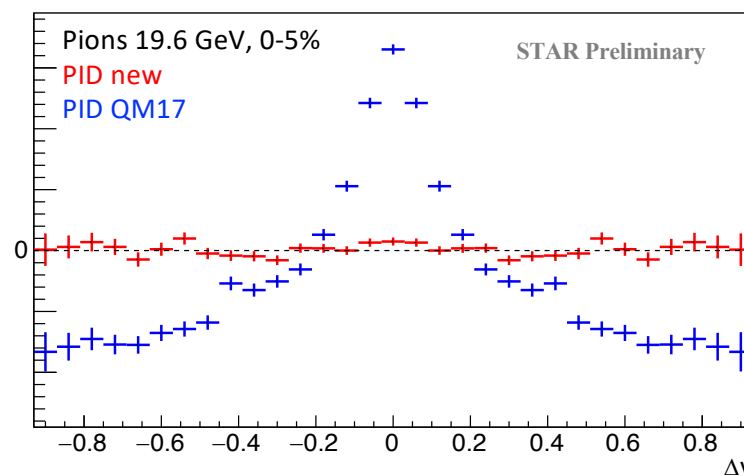
root(s_{NN})-localized..
charge-independent..
but in *pions*...

After doing many checks, we finally proved this peak is experimental, not physical...



Peak stays near $Z=0$
when Z_{vtx} moved around.

Can be removed either with a Z_{vtx} -dependent
rapidity cut, or by a more efficient PID method...



A paper on our identified 2-particle correlations versus the beam energy and centrality in the Au+Au BES-I data is in the STAR internal review process.

My main interest at the moment is measuring R₃ & R₄ (and C₃ & C₄)

$$R_3 = \frac{C_3}{\rho_1 \rho_1 \rho_1} = \frac{\rho_3}{\rho_1 \rho_1 \rho_1} - 3 \frac{\rho_2 \rho_1}{\rho_1 \rho_1 \rho_1} + 2$$

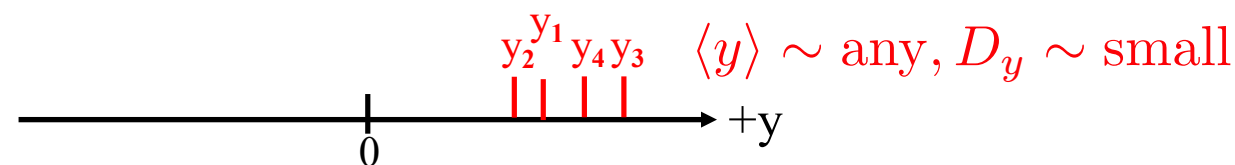
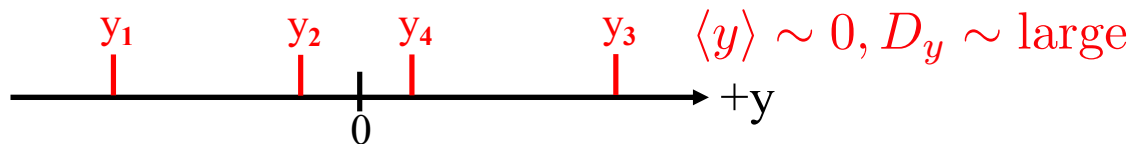
$$R_4 = \frac{C_4}{\rho_1 \rho_1 \rho_1 \rho_1} = \frac{\rho_4}{\rho_1 \rho_1 \rho_1 \rho_1} - 4 \frac{\rho_3 \rho_1}{\rho_1 \rho_1 \rho_1 \rho_1} - 3 \frac{\rho_2 \rho_2}{\rho_1 \rho_1 \rho_1 \rho_1} + 12 \frac{\rho_2 \rho_1 \rho_1}{\rho_1 \rho_1 \rho_1 \rho_1} - 6$$

Direct measure of “true” 3- and 4-particle correlations...

Insensitive to single-particle inefficiencies, explicitly subtract any pair inefficiencies...

As interest is focused on beam energies of ~20 GeV and lower, consider only LS protons.

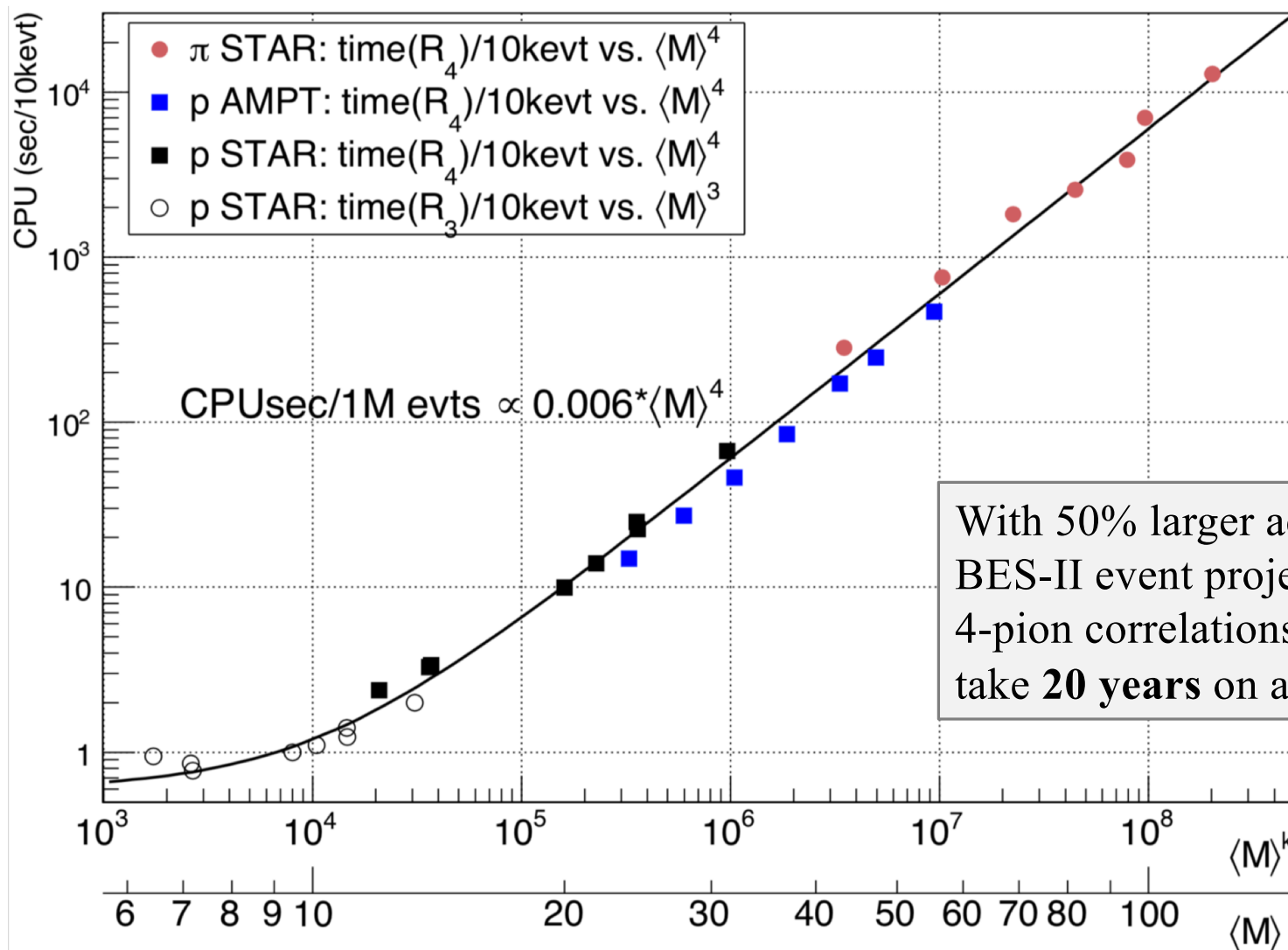
Projection approach used for now: $R_k(y_1, \dots, y_k) \rightarrow \langle R_k(\langle y \rangle, D_y) \rangle$



$$\langle y \rangle = \frac{1}{k} \sum_{i=1}^k y_i$$

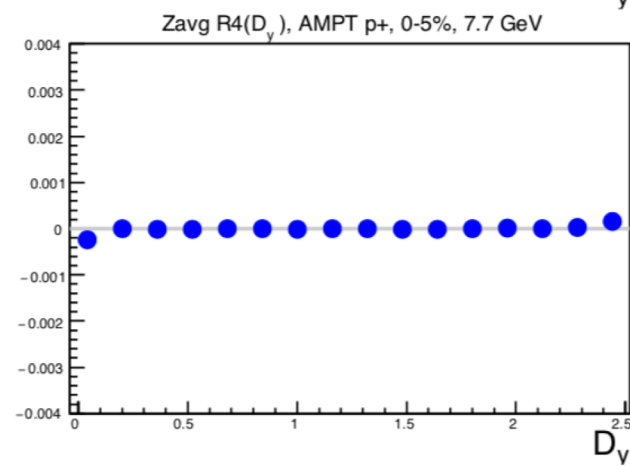
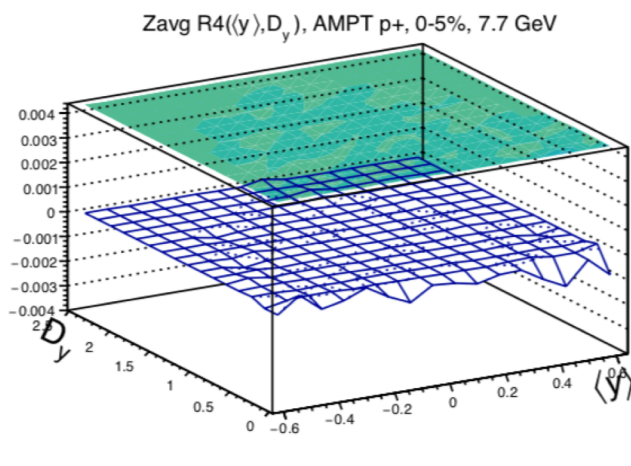
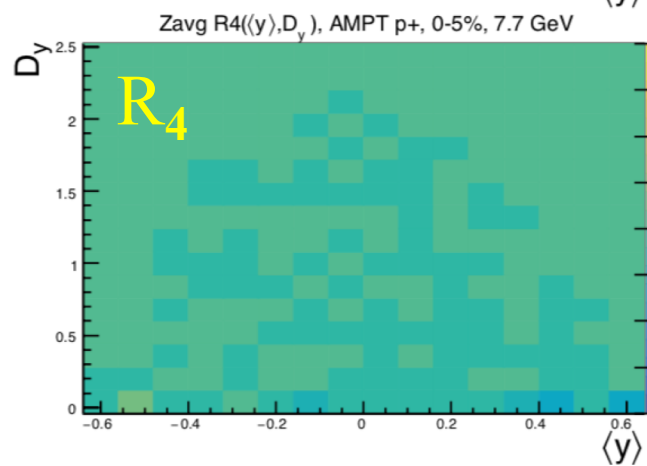
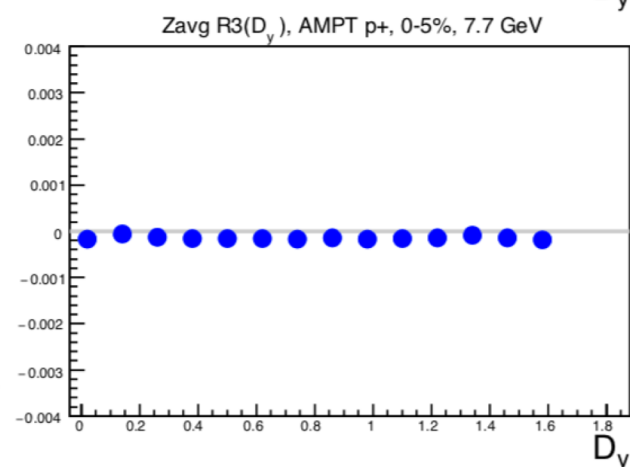
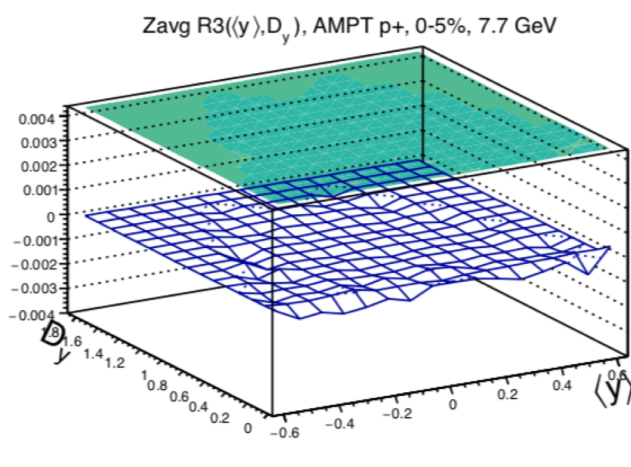
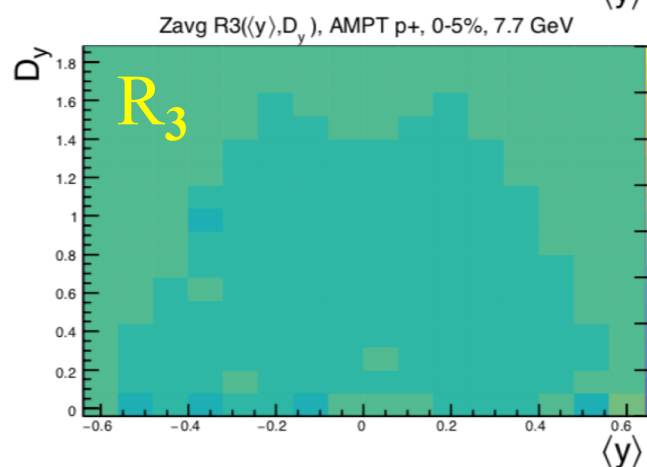
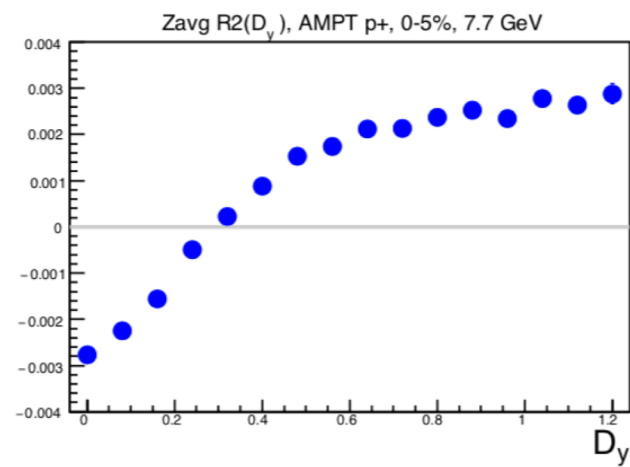
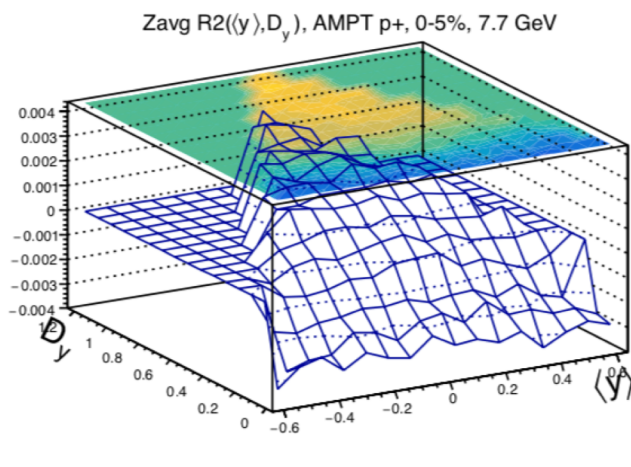
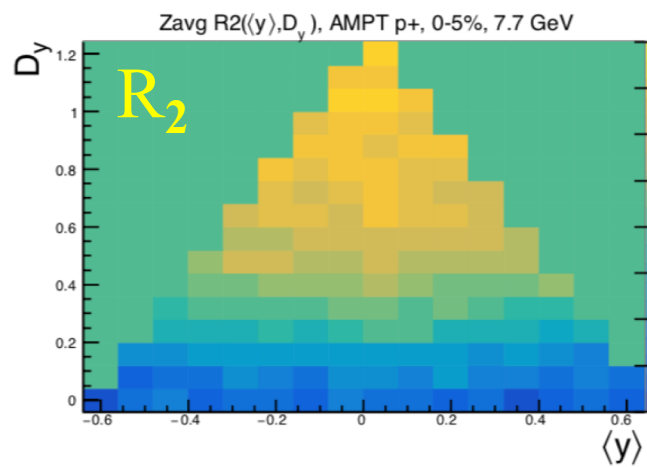
$$D_y = \sum_{i=1}^k |y_i - \langle y \rangle|$$

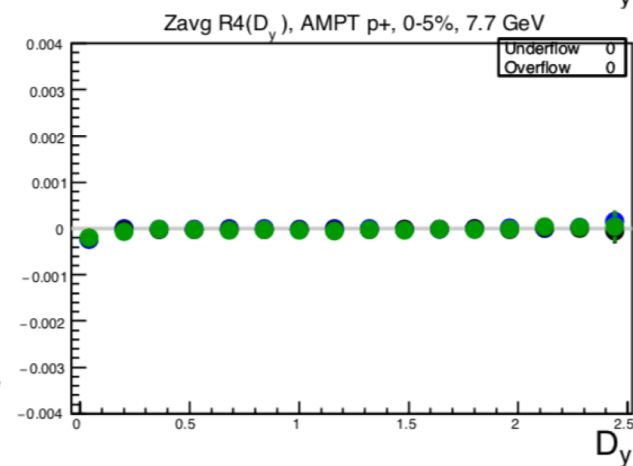
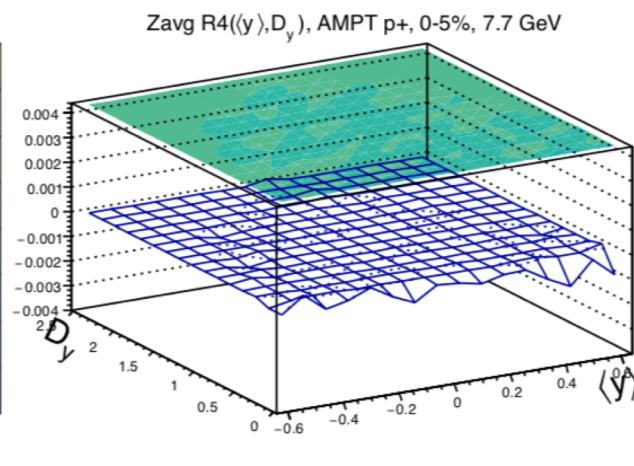
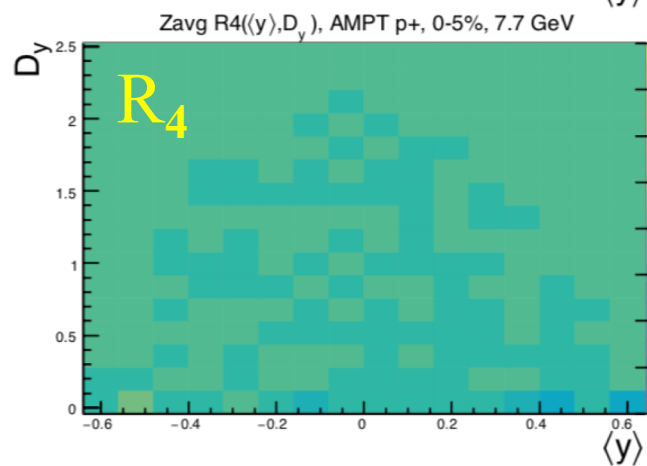
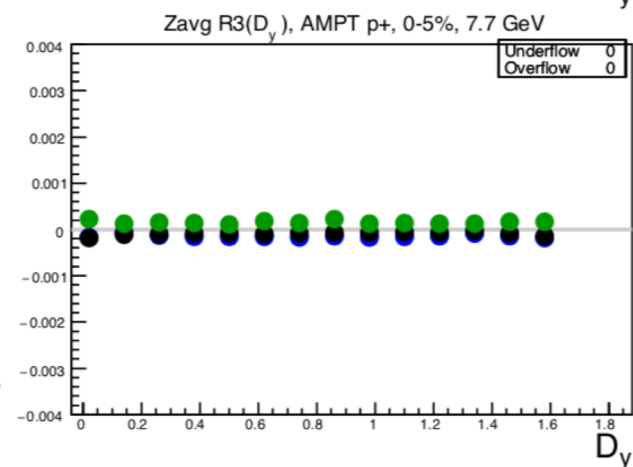
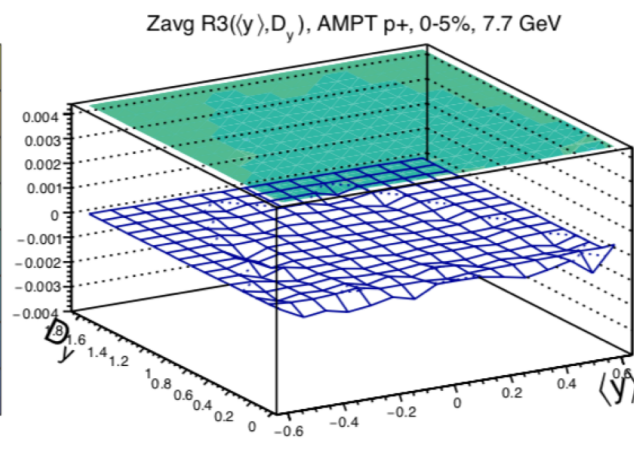
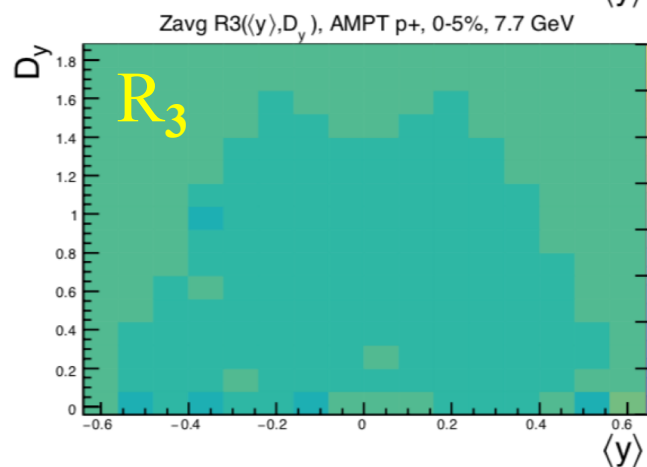
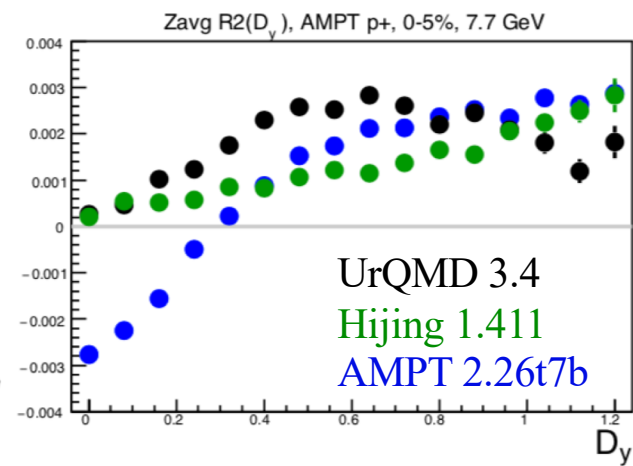
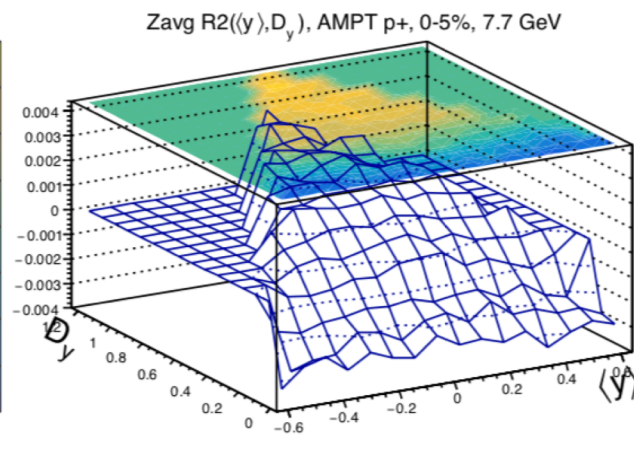
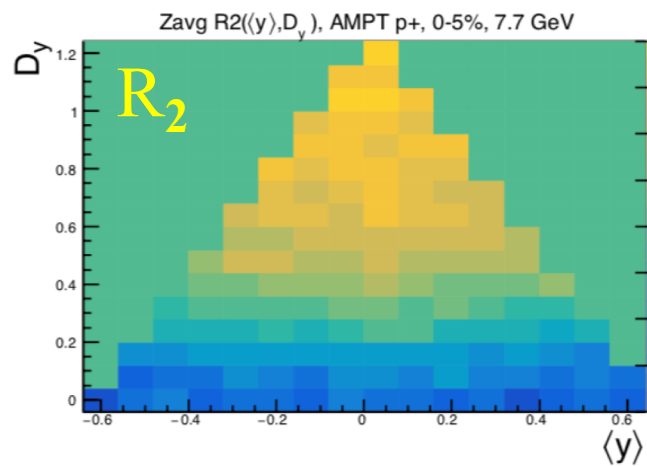
R_4 requires 4 nested loops over all M identified particles in each event



So, I break up the dataset into separate jobs of $\sim 100\text{k}$ events each and run in parallel...

Natural then to calculate uncertainties via “subgroup” method...





The correlations C_k (and the pair-normalized R_k) provide a more differential view of the particle distributions than does the multiplicity cumulants.

Allow direct treatment of Zvtx smearing, track crossing, ...

Able to localize any enhanced (anti)correlations kinematically

2 proton correlations do not look like 2 pion or 2 kaon correlations

mid-range anti-correlation for US protons

deepens with centrality and beam energy – likely annihilation

long-range anti-correlation for LS protons

cause unknown, not reproduced by some popular event generators

Moving toward 3- and 4- particle correlations in the STAR BES-I data

Presently projecting in a way that should “see” mid-range clustering

Any suggestions appreciated!

No observable 3- and 4-proton rapidity correlations in the models...

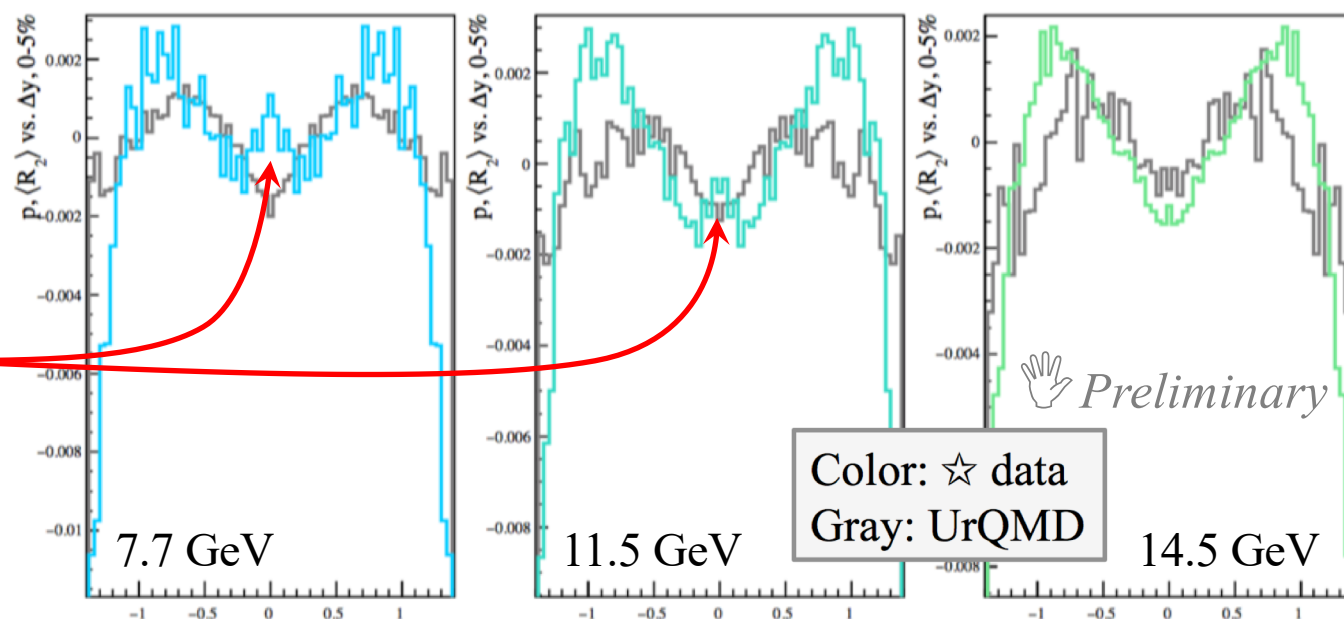
BACKUP SLIDES

Pseudocorrelations

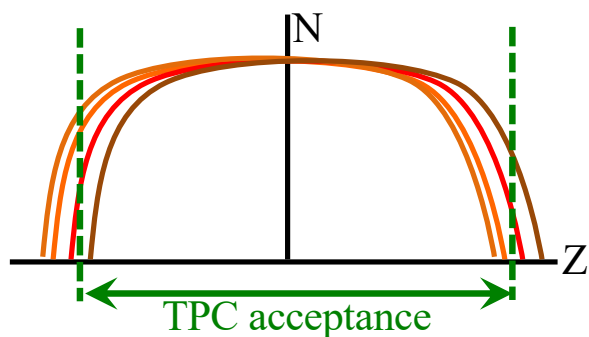
$\langle R_2 \rangle$ vs Δy

low Δy enhancement...

not seen in UrQMD evts...

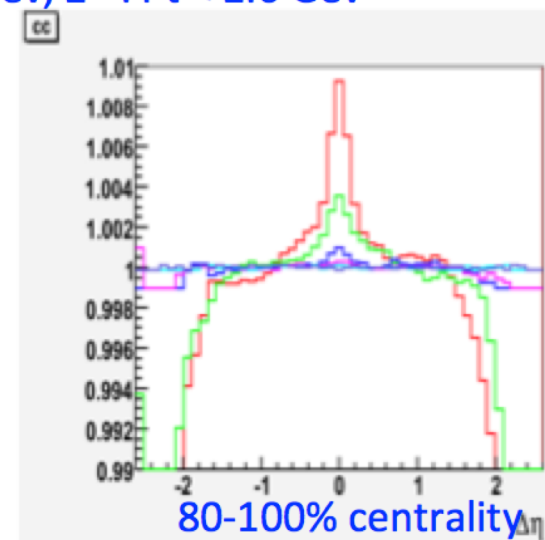
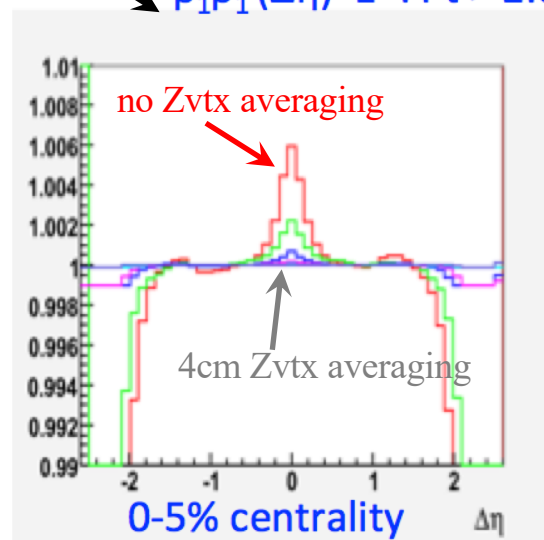


Caused by rapidity dependence of experimental efficiency coupled with Zvtx smearing...



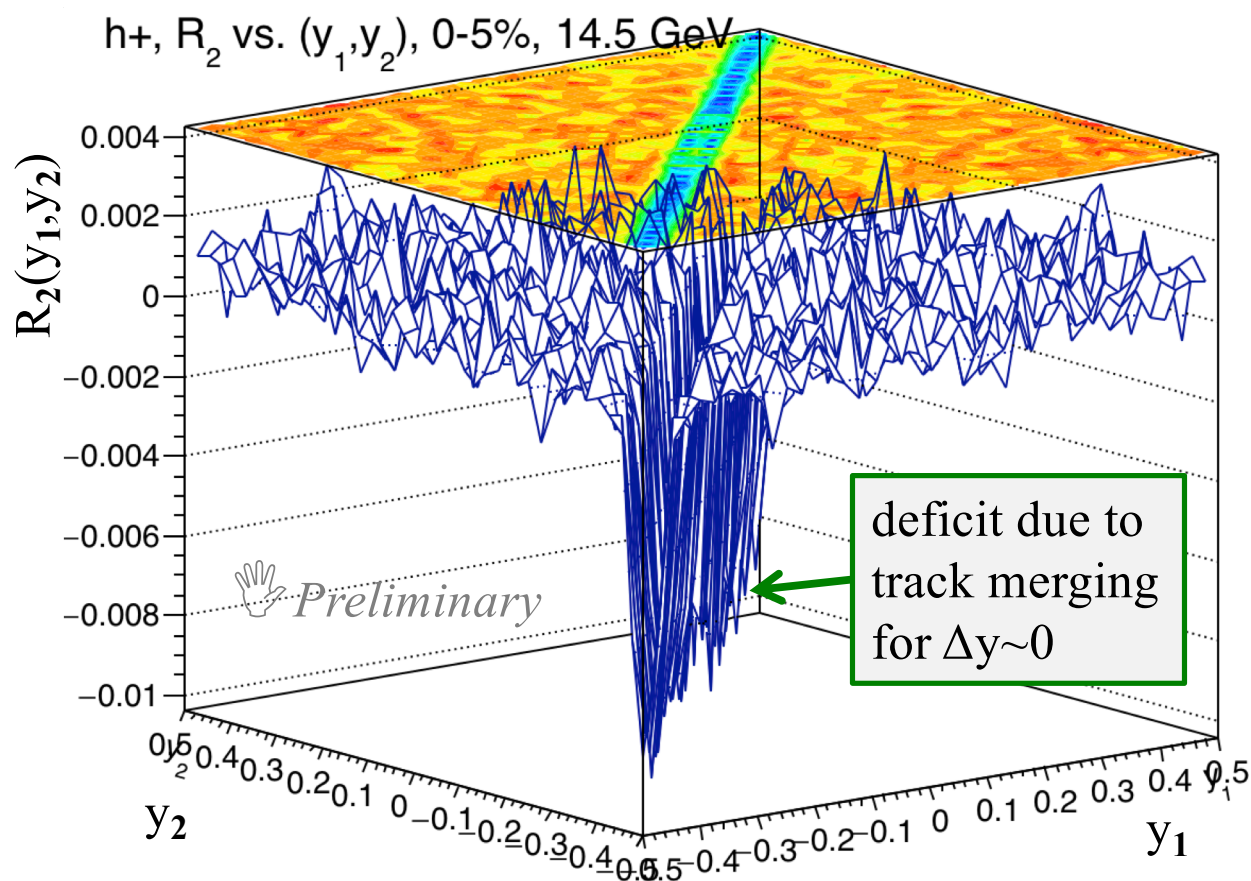
✓ Analyze in 2cm-wide Zvtx bins
then weight-average the results...

$p_1 p_1(\Delta\eta)$ 1st: $P_t > 2.0$ GeV, 2nd: $P_t < 2.0$ GeV



Very strong trench in R_2 when particle multiplicities/event of POI get large:
 $h^\pm$ for all centralities and $\sqrt{s_{NN}}$, and only most central for $K^\pm$

Numerator and denominator of R_2 & C_N uses only measured tracks...
but there is a slight **2-particle** inefficiency when two tracks are close ($\Delta y \sim 0$)

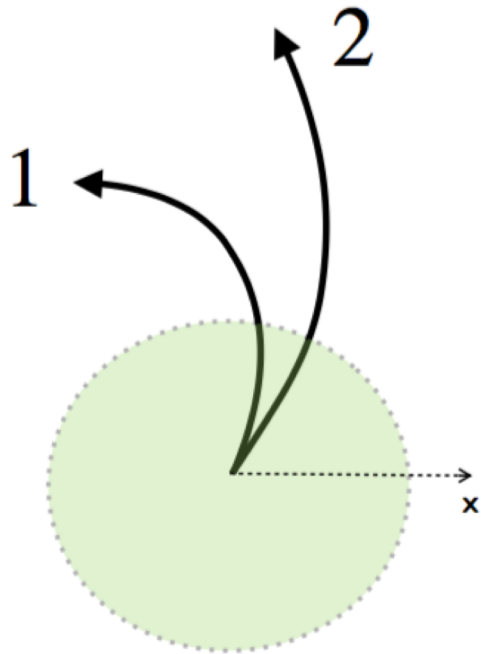


The STAR track-finder "sti"
does not share spacepoints!

a new one does "stiCA" (10%)

like-sign

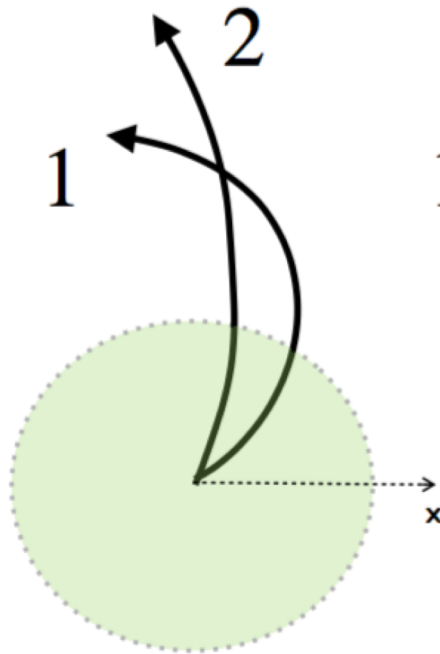
unlike-sign



No merging/crossing
No losses

$$p_{t,1} < p_{t,2}$$

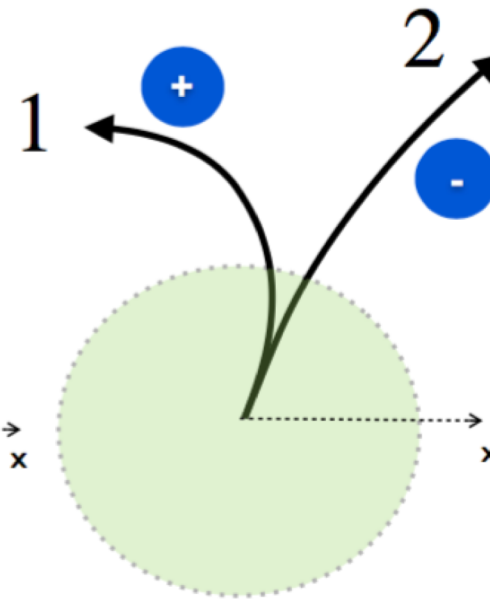
$$\Delta\varphi_{12} > 0$$



Merging/crossing
Hit losses, Pair Loss

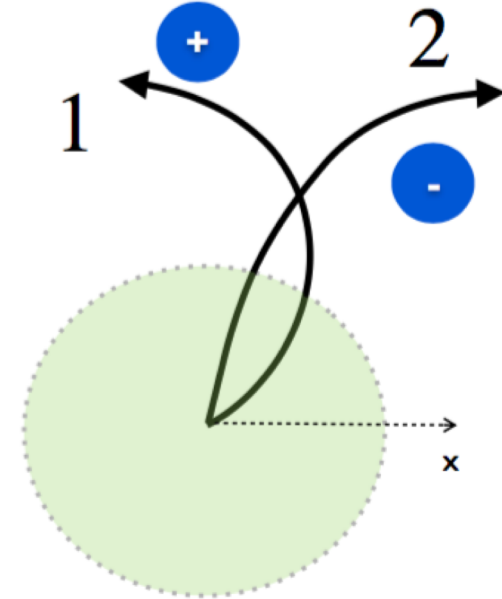
$$p_{t,1} < p_{t,2}$$

$$\Delta\varphi_{12} < 0$$



No merging/crossing
No losses

$$\Delta\varphi_{12} > 0$$



Merging/crossing
Hit losses, Pair Loss

$$\Delta\varphi_{12} < 0$$

Image from P. Pujahari

LS & US: reflect clean area in $\Delta\varphi$ to replace problem area

US: nothing special in fill method

LS: p_T order the tracks, fill numerator for upper triangle only, then symmetrize