

CP Violation at LHCb

Discrete Symmetries in Particle, Nuclear and Atomic Physics
and implications for our Universe

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on behalf of the LHCb Collaboration

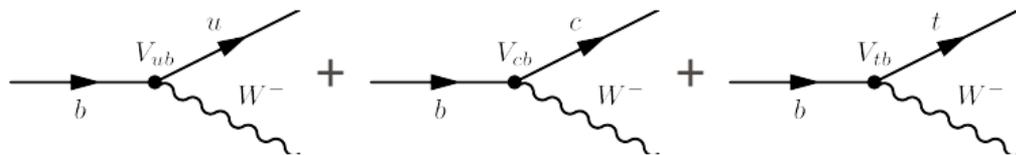
Trento, Oct 08-12, '18



Introduction

- Diagonalizing Yukawa matrices leads to **quark mass eigenstates** q , but skews **quark flavor eigenstates** q'

$$\mathcal{J}_{cc}^\mu \supset (\bar{u}', \bar{c}', \bar{t}') \gamma^\mu \frac{1 - \gamma^5}{2} \begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} \rightarrow (\bar{u}, \bar{c}, \bar{t}) \gamma^\mu \frac{1 - \gamma^5}{2} V_{\text{CKM}} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$



- This introduces the non-trivial transformation matrix V_{CKM} for quarks (CKM matrix)
- V_{CKM} is a **unitary** 3×3 matrix

The CKM Matrix

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \underbrace{\begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}}_{\text{cf. Wolfenstein: Phys. Rev. Lett. 51 (1945), 21}} + \mathcal{O}(\lambda^3)$$

$$\lambda \approx .23, A \approx .81, \rho \approx .14, \eta \approx .35$$

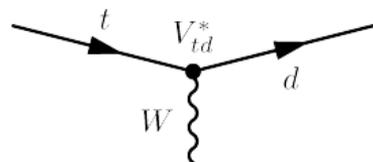
- One non-trivial complex phase, encoded in matrix elements
 - V_{ub} and V_{td} (up to $\mathcal{O}(\lambda^2)$)
 - V_{cd} , V_{cs} and V_{ts} (up to $\mathcal{O}(\lambda^6)$)
- CP violation (CPV) if and only if $\eta \neq 0$

CP violation (?)

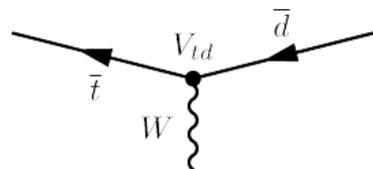
- If $\eta \neq 0$ some V_{ij} carry a complex phase (*weak phase*)
 - amplitude: $\mathcal{A}(t \rightarrow dW^+) \sim V_{td}^*$
 - amplitude: $\bar{\mathcal{A}}(\bar{t} \rightarrow \bar{d}W^-) \sim V_{td}$
 - CPV since $\mathcal{A} \neq \bar{\mathcal{A}}$ (?)

... not quite

- Amplitude \mathcal{A} is not observable ...
- ... but branching fraction $\mathcal{B} \sim |\mathcal{A}|^2$
 - CPV needs at least two interfering decay modes with ...
 - ... one **CP odd** and one **CP even** phase



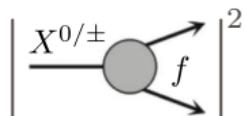
$t \rightarrow dW^+$



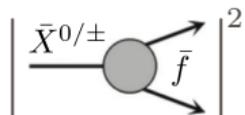
$\bar{t} \rightarrow \bar{d}W^-$

Types of CPV

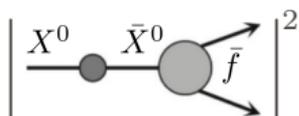
(1) Direct CPV



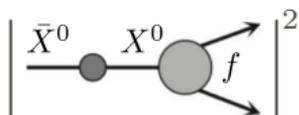
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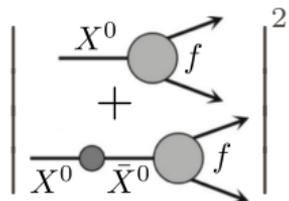
(2) CPV in mixing



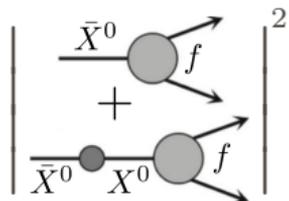
\neq



(3) CPV in interference of mixing and decay



\neq



- CP odd: from CKM matrix
- CP even:
 - (1): strong phase difference between both amplitudes (e.g. *tree* and *penguin*)
 - (2),(3): $\pi/2$ (**constant!**)

(Images: *CP Violation*, I. I. Bigi and A. I. Sanda)

The CKM Matrix

Another parametrization – Prog. Part. Nucl. Phys. 47 (2001)

$$V_{\text{CKM}} = \begin{pmatrix} |V_{ud}| & |V_{us}| & |V_{ub}|e^{-i\tilde{\gamma}} \\ -|V_{cd}|e^{+i\phi_4} & |V_{cs}|e^{-i\phi_6} & |V_{cb}| \\ |V_{td}|e^{-i\tilde{\beta}} & -|V_{ts}|e^{+i\phi_2} & |V_{tb}| \end{pmatrix}$$

$$\begin{aligned} \gamma &\equiv \tilde{\gamma} - \phi_4, & \phi_2 &\approx \eta\lambda^2 \\ \beta &\equiv \tilde{\beta} + \phi_4, & \phi_4 &\approx \eta A^2 \lambda^4 \\ \alpha &\equiv \pi - \beta - \gamma, & \phi_6 &\approx \eta A^4 \lambda^6 \\ \beta_s &\equiv \phi_s \equiv \phi_2 + \phi_6 \end{aligned}$$

→ From unitarity: 6 triangles

$$\rightarrow V_{ud}V_{us}^* + V_{cd}V_{cs}^* + V_{td}V_{ts}^* = 0$$

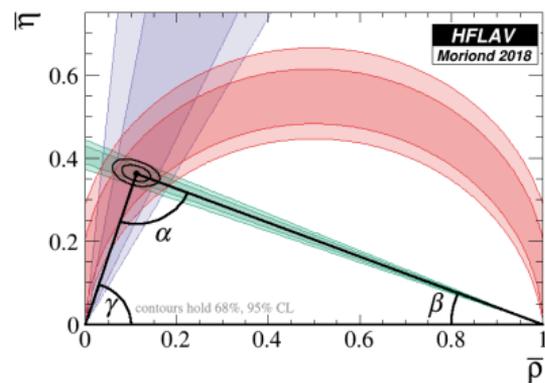
$$\rightarrow V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

→ ... 4 more

→ Angles $\tilde{\alpha}$, $\tilde{\beta}$, $\tilde{\gamma}$ and $\phi_{2,4,6}$ **depend on phase convention** (i.e. not observable)

→ Phases of products $V_{ij}V_{kl}V_{il}^*V_{kj}^*$ are **invariant** and **observable** (e.g. α , β , γ , ...)

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$



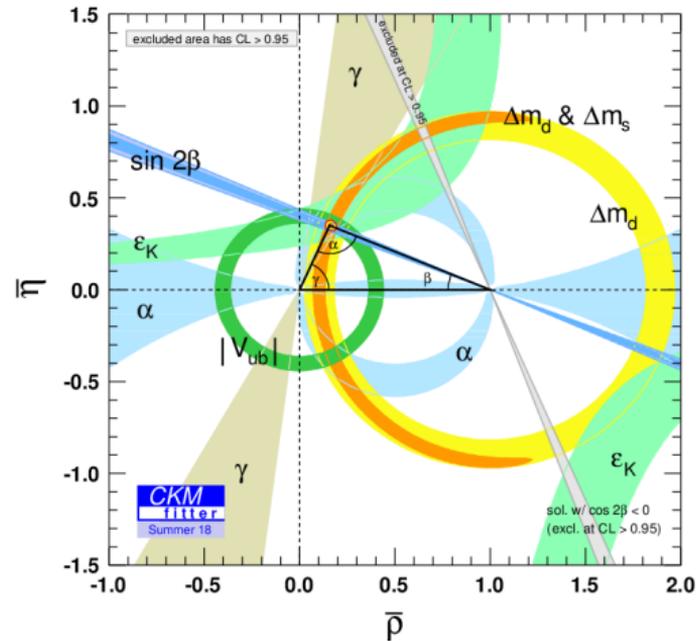
Probing the SM by Overconstraining

→ Motivation:

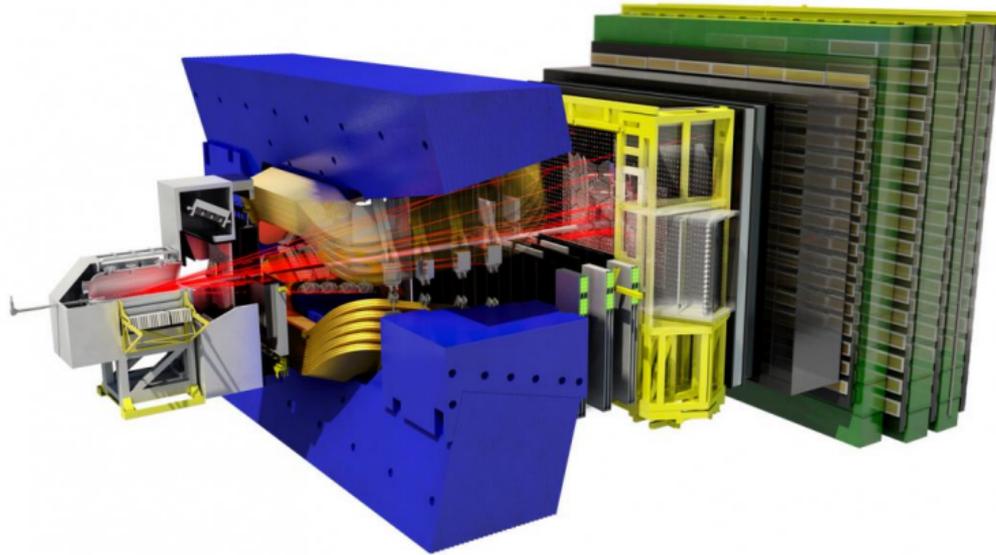
- We utterly fail explaining CPV on cosmological scale!
- Is the CKM matrix the only source of CPV?
- Is the SM incomplete / are there more particles?

→ Strategy:

- Overconstrain CKM triangles by measuring
 - sides - e.g. $|V_{td}V_{tb}^*|$
 - angles - e.g. $\gamma = \arg\left(-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}\right)$
- Is $\beta_s = \arg\left(-\frac{V_{ts}V_{tb}^*}{V_{cs}V_{cb}^*}\right)$ tiny? ($\mathcal{O}(\lambda^2)$)
- Any deviation would point towards **new physics**, e.g. 4th quark family?



The LHCb Detector

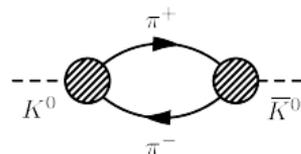


more information: *The LHCb Detector at the LHC*, JINST 3 (2008), S08005

or visit: <http://lhcb-public.web.cern.ch/>

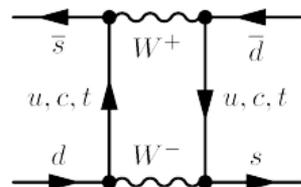
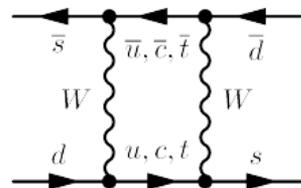
Neutral mesons mix through weak interactions

- CPV if $\mathcal{P}(X^0 \rightarrow \bar{X}^0) \neq \mathcal{P}(\bar{X}^0 \rightarrow X^0)$
- SM: CPV large for K^0/\bar{K}^0 (first discovery of CPV)
- SM: CPV small for $B_{(s)}^0/\bar{B}_{(s)}^0$



Example: CPV in mixing for $B_{(s)}^0/\bar{B}_{(s)}^0$ in semileptonic modes

- $B^0 \rightarrow D^{(*)-} \mu^+ \nu_\mu X^*$
- $B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu X^{**}$
- decays are theoretically clean
 - flavor specific
 - CP-conserving
 - tree-dominated



* Phys. Rev. Lett. 114 (2015), 041601, ** Phys. Rev. Lett. 117 (2016), 061803

Example: $B_s^0 \rightarrow D_s^- \mu^+ \nu_\mu X$

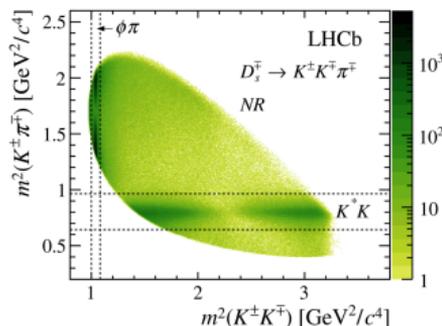
- CPV parametrization:
 - Semileptonic CP asymmetry (**const. in time**)

$$a_{sl}^s = \frac{\Gamma(\bar{B}_s^0 \rightarrow D_s^- \mu^+ |t) - \Gamma(B_s^0 \rightarrow D_s^+ \mu^- |t)}{\Gamma(\bar{B}_s^0 \rightarrow D_s^- \mu^+ |t) + \Gamma(B_s^0 \rightarrow D_s^+ \mu^- |t)}$$

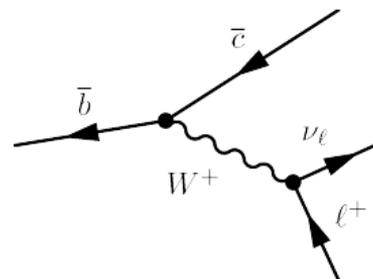
- Observables:

- $\Gamma[D_s^- \mu^+, t] \equiv \Gamma(\bar{B}_s^0 \rightarrow D_s^- \mu^+ |t) + \Gamma(B_s^0 \rightarrow D_s^- \mu^+ |t)$
- $\Gamma[D_s^+ \mu^-, t] \equiv \Gamma(B_s^0 \rightarrow D_s^+ \mu^- |t) + \Gamma(\bar{B}_s^0 \rightarrow D_s^+ \mu^- |t)$

$$\frac{\Gamma[D_s^- \mu^+, t] - \Gamma[D_s^+ \mu^-, t]}{\Gamma[D_s^- \mu^+, t] + \Gamma[D_s^+ \mu^-, t]} = \frac{a_{sl}^s}{2} \underbrace{\left(1 - \frac{\cos \Delta m_s t}{\cosh \Delta \Gamma t / 2}\right)}_{\approx 1 \text{ for fast oscillation}}$$

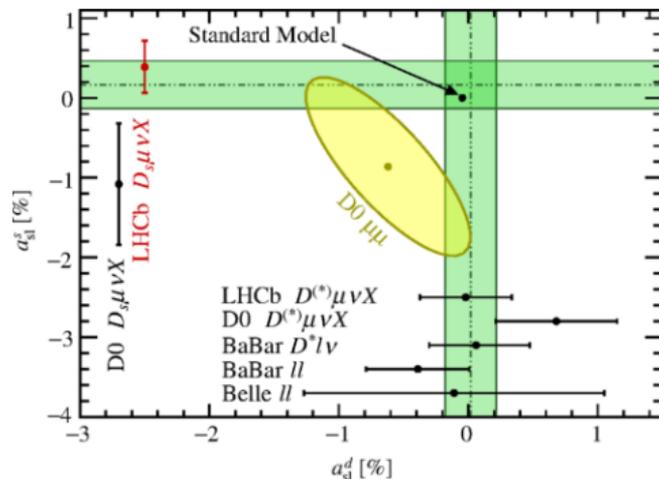


Phys. Rev. Lett. 117 (2016), 061803



$$A_{\text{raw}} \equiv \frac{\Gamma[D_s^- \mu^+, t] - \Gamma[D_s^+ \mu^-, t]}{\Gamma[D_s^- \mu^+, t] + \Gamma[D_s^+ \mu^-, t]} \approx \frac{a_{\text{sl}}^s}{2}$$

- Fast oscillation of B_s :
 - no dependency on production asymmetry between B_s and \bar{B}_s
 - measure A_{raw} time-integrated
 - include detection and background asymmetries
- Reconstruct $D_s^\mp \rightarrow K^+ K^- \pi^\mp$ in various subsamples of
 - magnet polarity
 - Dalitz regions
 - data taking periods



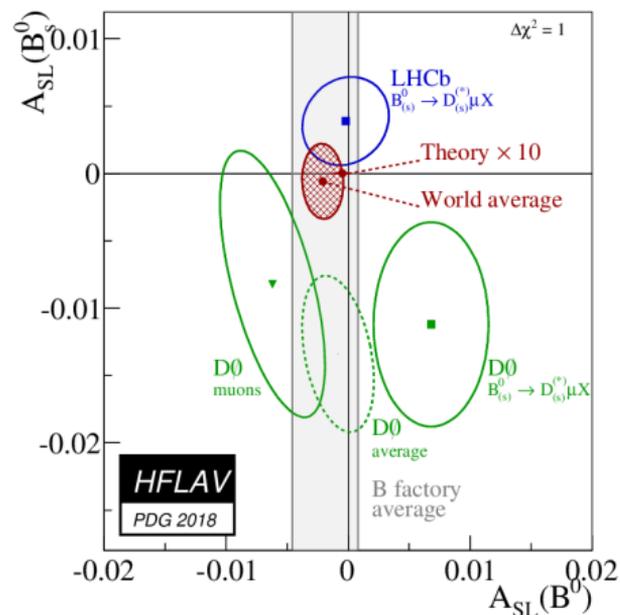
a_{sl}^d : Phys. Rev. Lett. 114 (2015), 041601

a_{sl}^s : Phys. Rev. Lett. 117 (2016), 061803

$$a_{\text{sl}}^s = (.39 \pm .26 (\text{stat.}) \pm .20 (\text{syst.})) \%$$

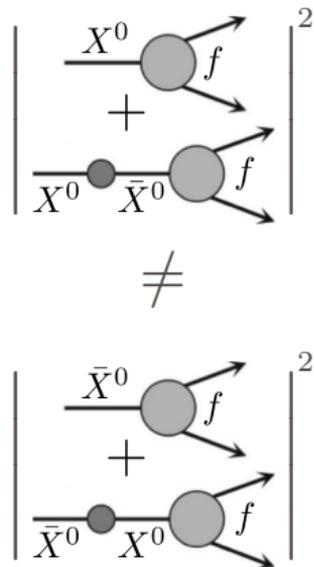
$$A_{\text{raw}} \equiv \frac{\Gamma[D_s^- \mu^+, t] - \Gamma[D_s^+ \mu^-, t]}{\Gamma[D_s^- \mu^+, t] + \Gamma[D_s^+ \mu^-, t]} \approx \frac{a_{\text{sl}}^s}{2}$$

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Golden Channel: $B^0 \rightarrow J/\psi K_s$

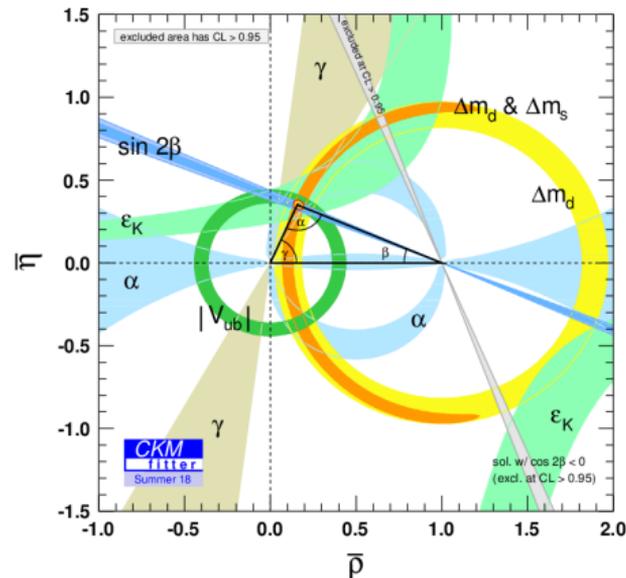
- B^0/\bar{B}^0 decays to CP eigenstate $J/\psi K_s$
- Theoretically clean mode
 - leading transitions via tree
 - penguin strongly suppressed
 - no CPV from mixing expected
- sensitive to $\sin 2\beta$ (*Golden Channel*)



CPV in Mixing and Decay

- Reconstruct: $B^0 \rightarrow [c\bar{c}]K_s$
 - $[c\bar{c}] = J/\psi \rightarrow \mu\mu^*$
 - $[c\bar{c}] = J/\psi \rightarrow ee, \psi(2S) \rightarrow \mu\mu^{**}$
- Measure:
 - $\Gamma(t) \equiv \Gamma(B^0(t=0) \rightarrow [c\bar{c}]K_s|t)$
 - $\bar{\Gamma}(t) \equiv \Gamma(\bar{B}^0(t=0) \rightarrow [c\bar{c}]K_s|t)$
- CPV observable: $(\bar{\Gamma} - \Gamma)/(\bar{\Gamma} + \Gamma)(t)$

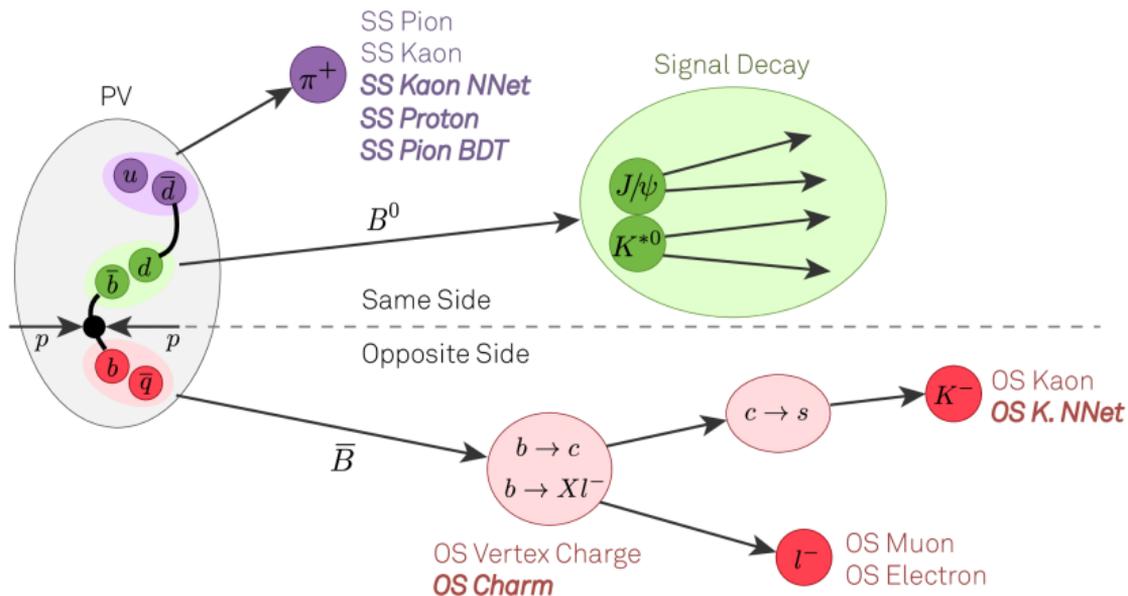
$$\approx -C \cos(\Delta mt) + S \sin(\Delta mt)$$
 - assuming negligible CPV in mixing
 - $C \approx 0$ (\sim (no) direct CPV)
 - $S = \sin 2\beta$



* Phys. Rev. Lett. 115 (2015), 031601

** JHEP11 (2017), 170

Tagging at LHCb

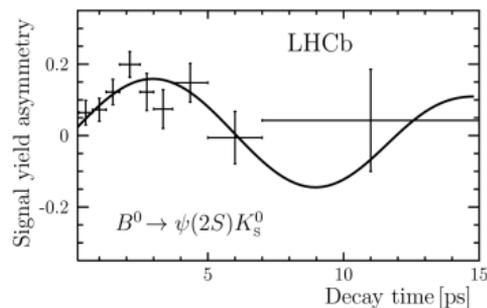
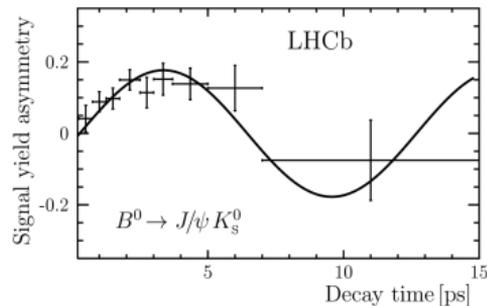
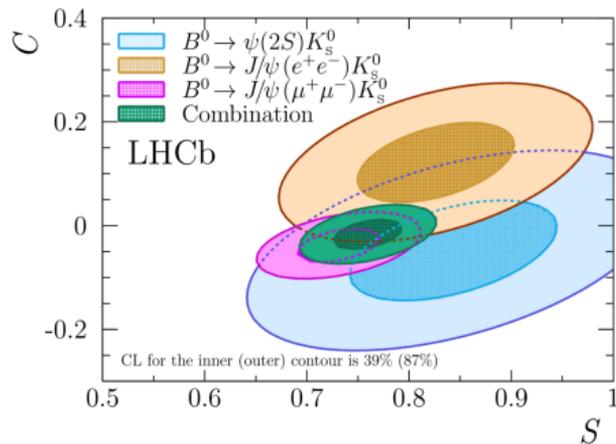


Time-dependent CP measurements need **flavor** of B^0 / \bar{B}^0 -system **at production time!**

CPV in Mixing and Decay

Combined results:

- Direct CPV: $C(B^0 \rightarrow [c\bar{c}]K_s) = -.014(30)$
- $\sin 2\beta = S(B^0 \rightarrow [c\bar{c}]K_s) = -.75(4)$

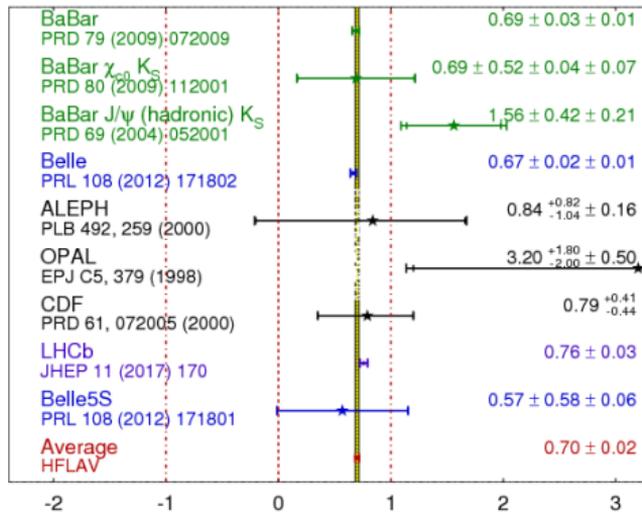


$\sin 2\beta$ has a two-fold ambiguity: $2\beta \leftrightarrow \pi - 2\beta$; cf. arXiv:1804.06152 (2018): combined dataset of Babar and Belle excludes $\pi - 2\beta$ (@ 7.3σ)

CPV in Mixing and Decay

$$\sin(2\beta) \equiv \sin(2\phi_1)$$

HFLAV
Moriond 2018
PRELIMINARY

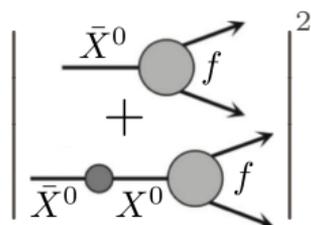
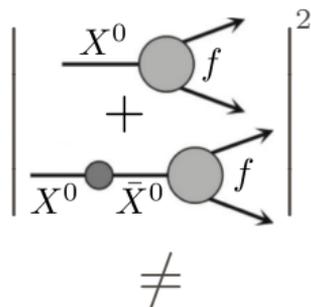


LHCb measurements are already competitive with B -factories!

Another Channel: $B_{(s)}^0 / \bar{B}_{(s)}^0 \rightarrow h^+ h^-$

- $B_{(s)}^0 / \bar{B}_{(s)}^0$ decays to CP eigenstates $\pi^+ \pi^-$ or $K^+ K^-$
- Theoretically nuisance: **tree** and **penguin** contributions
- CPV observable: $(\bar{\Gamma} - \Gamma) / (\bar{\Gamma} + \Gamma)(t)$

$$= \frac{-C_f \cos(\Delta m_{d,s} t) + S_f \sin(\Delta m_{d,s} t)}{\cosh(\Delta \Gamma_{d,s} t / 2) + A_f^{\Delta \Gamma} \sinh(\Delta \Gamma_{d,s} t / 2)}$$



\neq

CPV in interference of mixing and decay

$$A_{\text{CP}}(t) = \frac{-C_f \cos(\Delta m_{d,s} t) + S_f \sin(\Delta m_{d,s} t)}{\cosh(\Delta \Gamma_{d,s} t/2) + A_f^{\Delta \Gamma} \sinh(\Delta \Gamma_{d,s} t/2)}$$

→ Parameters C_f , S_f and $A_f^{\Delta \Gamma}$ are completely defined by $\lambda_f \in \mathbb{C}$

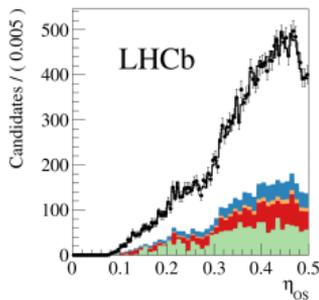
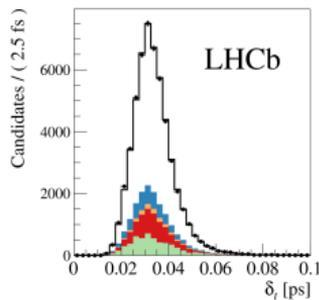
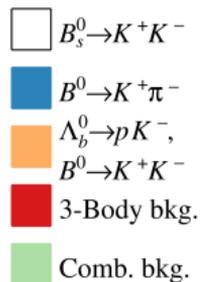
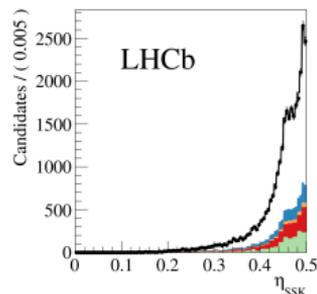
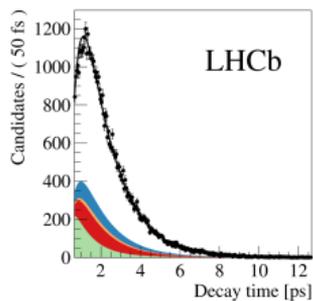
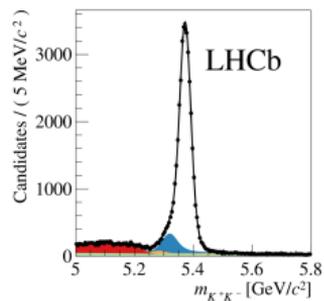
$$\lambda_f \equiv \underbrace{\frac{\langle \bar{X}^0 | X_L \rangle}{\langle X^0 | X_L \rangle}}_{q/p} \underbrace{\frac{\mathcal{A}(\bar{X}^0 \rightarrow f)}{\mathcal{A}(X^0 \rightarrow f)}}_{\mathcal{A}_f / \bar{\mathcal{A}}_f}$$

- direct CPV if $|\mathcal{A}_f / \bar{\mathcal{A}}_f| \neq 1$
- CPV in oscillation if $|q/p| \neq 1$
- CPV from interference of mixing and decay if $\text{Im}(\lambda_f) \neq 0$

$$\text{CPV} \Leftrightarrow \lambda_f \neq 1$$

CPV in Mixing and Decay

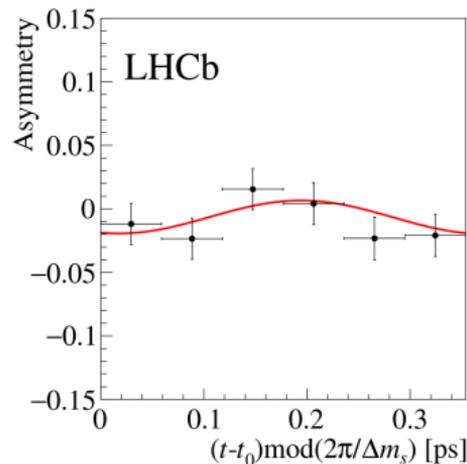
Fit simultaneously to unfold time-dependent CP asymmetries



Phys. Rev. D98 (2018), 032004

CPV in interference of mixing and decay

- Obtain $C_{KK}, S_{KK}, A_{KK}^{\Delta\Gamma}$ (and $A_{CP}(B_{(s)}^0)$) from simultaneous fit to:
 - invariant mass
 - decay time + uncertainty
 - tagging decision + associated mistag probabilities
- unfold time-dependent asymmetries
- Check: $\lambda_{KK} = 1$
 - $\Leftrightarrow (C_{KK}, S_{KK}, A_{KK}^{\Delta\Gamma}) = (0, 0, -1)$
 - Excluded by $4\sigma!$ (for B_s^0)

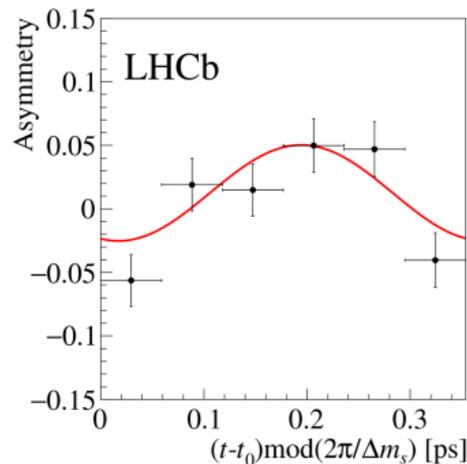


Same side tag for $B_s^0 \rightarrow K^+ K^-$
(Phys. Rev. D98 (2018), 032004)

Strongest evidence for time-dependent CPV in B_s^0 to date!

CPV in interference of mixing and decay

- Obtain $C_{KK}, S_{KK}, A_{KK}^{\Delta\Gamma}$ (and $A_{CP}(B_{(s)}^0)$) from simultaneous fit to:
 - invariant mass
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- Check: $\lambda_{KK} = 1$
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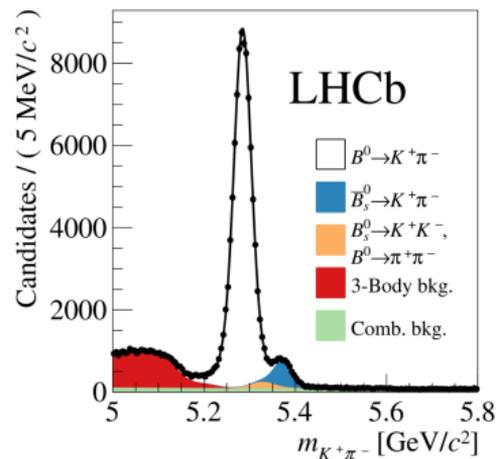


Opposite side tag for $B_s^0 \rightarrow K^+ K^-$
(Phys. Rev. D98 (2018), 032004)

Strongest evidence for time-dependent CPV in B_s^0 to date!

CPV in interference of mixing and decay

- Same machinery for $B^0 \rightarrow \pi^+ \pi^-$
 - most precise measurements of $C_{\pi\pi}$ and $S_{\pi\pi}$ from a single experiment to date
 - combination of both measurements allow stringent constraints on γ and β_s
- Measurement of CP asymmetry $A_{CP}^{B^0}$ and $A_{CP}^{B_s^0}$
 - most precise measurements from a single experiment to date
 - improve constraints on BSM that contributes to loop diagrams



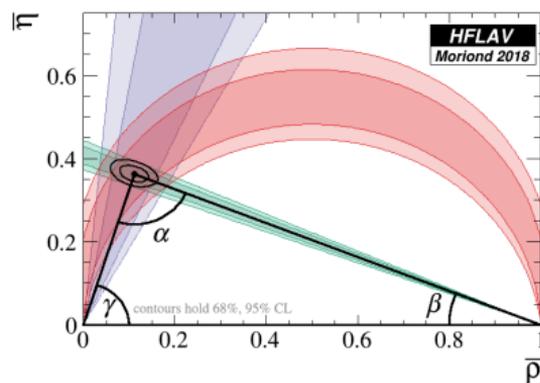
Phys. Rev. D98 (2018), 032004

CKM Angle γ

Why γ ?

- γ (had been) the CKM angle with largest uncertainties amongst α, β, γ
- γ does **not depend on any top quark coupling!**
 - measurements are dominated by **tree** contributions!
 - smaller uncertainties in physical observables!
 - but: **direct CPV**, i.e. strong phases and suppression by lowest amplitude

$$\gamma \equiv \arg \left(-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} \right)$$



LHCb Combination – LHCb-CONF-2018-002

→ Time-integrated amplitude analyses

(GLW, ADS, GLW+ADS, GGSZ, ...)

$$\rightarrow B^+ \rightarrow D^{(*)} K^+, DK^{(*)+}$$

$$\rightarrow B^0 \rightarrow DK^{*0}$$

$$\rightarrow B^0 \rightarrow DK^+ \pi^-$$

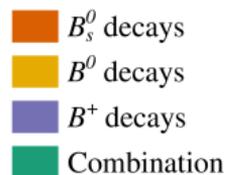
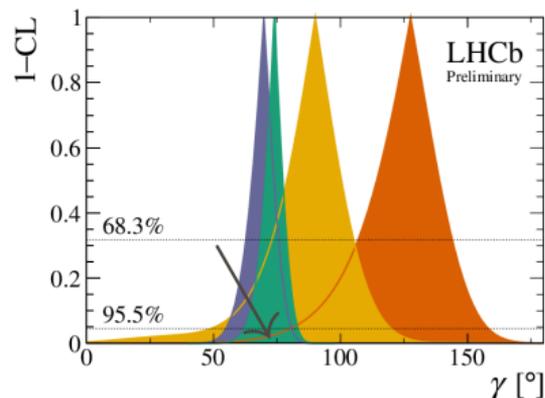
$$\rightarrow B^+ \rightarrow DK^+ \pi^+ \pi^-$$

→ Time-dependent

$$\rightarrow B_s^0 \rightarrow D_s^\mp K^\pm$$

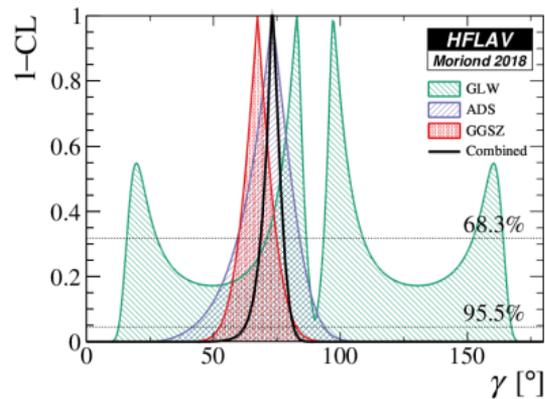
$$\rightarrow B^0 \rightarrow D^\mp \pi^\pm$$

$$\text{Best Fit value: } \gamma_{\text{LHCb}} = (74.0_{-5.8}^{+5.0})^\circ$$



Best Fit value: $\gamma_{\text{LHCb}} = (74.0^{+5.0}_{-5.8})^\circ$

- Most precise measurement from a single experiment (to date)
- cf. full BaBar dataset: $\gamma_{\text{BaBar}} = (69^{+17}_{-16})^\circ$
(Phys. Rev. D87 (2013), 052015)

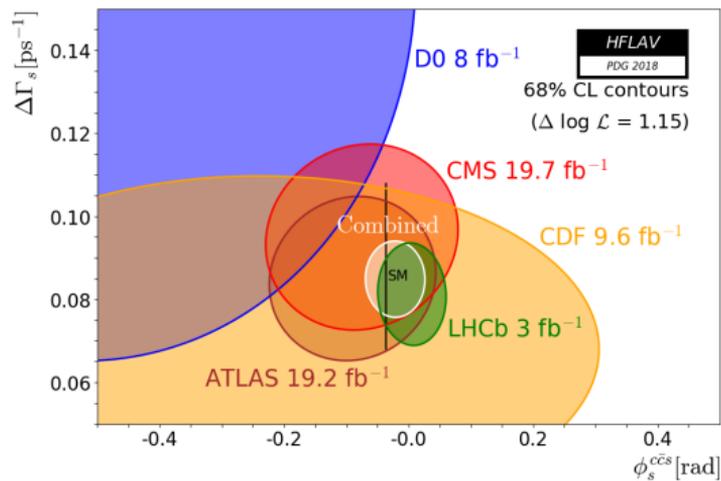


$$\gamma_{\text{HFLAV}} = (73.5^{+4.2}_{-5.1})^\circ$$

There is more!

There is more at LHCb

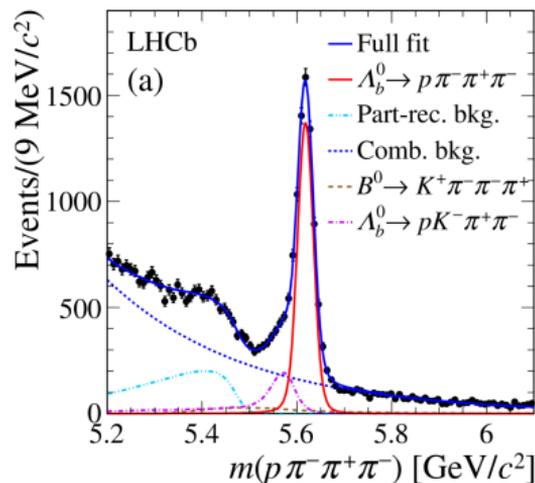
- There are more angles than β and γ and LHCb can measure these!
- LHCb also probes the baryon sector for CPV, e.g.
 - $\Lambda_b \rightarrow p\pi^-\pi^+\pi^-$, $p\pi^-K^+K^-$
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- Besides CPV: probing *lepton universality*, e.g.
 - $B^+ \rightarrow K^+\ell^+\ell^-$
 - $B^0 \rightarrow K^{*0}\ell^+\ell^-$
- ...



Including four LHCb measurements (i.a. first evidence for $\Delta\Gamma_s > 0$!)

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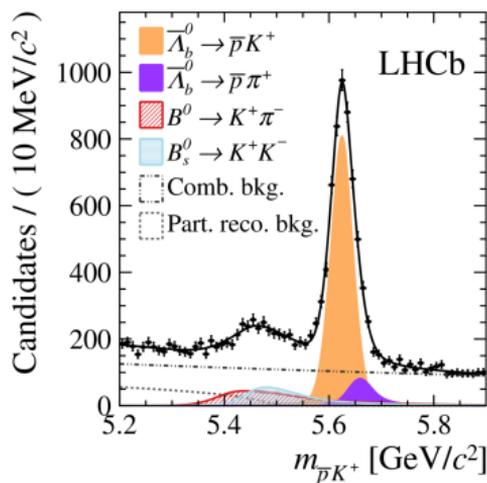
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First evidence (3.3σ) for CPV in baryon sector
(Nature Physics 13 (2017), 391)

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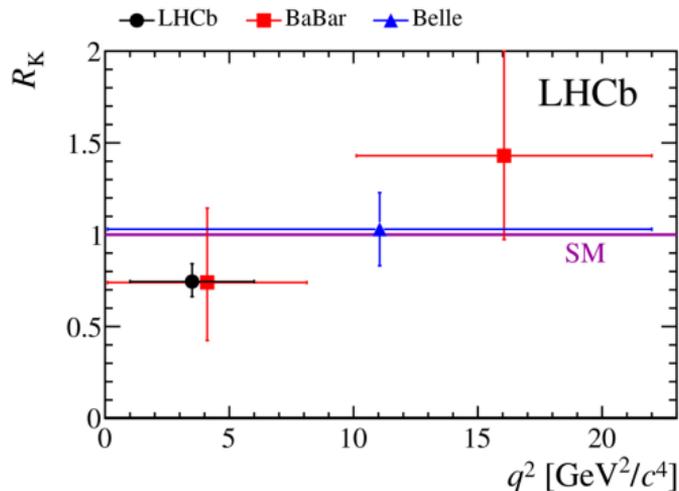
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Most precise measurement to date, compatible with no direct CP violation in $\Lambda_b \rightarrow ph^-$ (arXiv:1807.06544, 2018)

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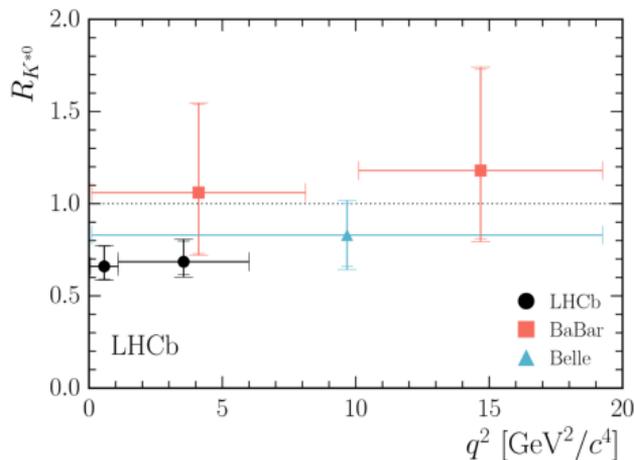
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Most precise measurements of R_{K^0} to date (lepton universality), compatible with SM at 2.6σ (Phys. Rev. Lett. 113 (2014), 151601)

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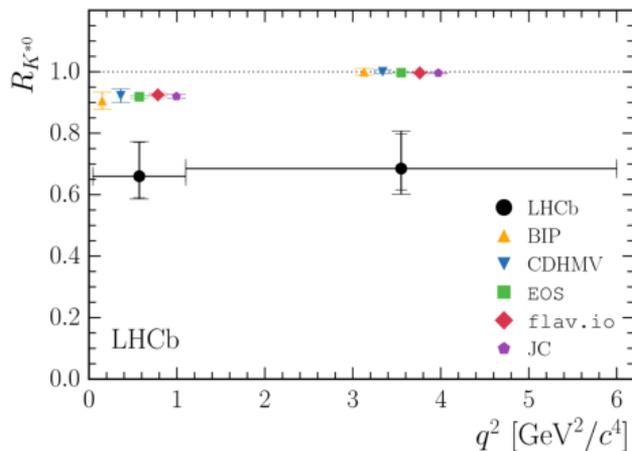
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Most precise measurements of $R_{K^{*0}}$ to date (lepton universality), compatible with SM at $2.1 \dots 2.3\sigma$ and $2.4 \dots 2.5\sigma$ (JHEP 08 (2017), 055)

There is more at LHCb

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Bundesministerium
für Bildung
und Forschung



Thank you for your attention!

Backup

Direct CP Violation

$$\left| \begin{array}{c} X^{0/\pm} \\ \text{---} \\ \bullet \\ \begin{array}{l} \nearrow \\ \searrow \end{array} \\ f \end{array} \right|^2 \neq \left| \begin{array}{c} \bar{X}^{0/\pm} \\ \text{---} \\ \bullet \\ \begin{array}{l} \nearrow \\ \searrow \end{array} \\ \bar{f} \end{array} \right|^2$$

Example: Direct CP Violation

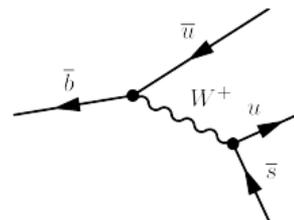
Decay $X \rightarrow Y$ via **two different channels** (weak phases $\phi_{1,2}$, strong phases $\delta_{1,2}$)

$$\mathcal{A}(X \rightarrow Y) = |\mathcal{A}_1|e^{+i\phi_1+i\delta_1} + |\mathcal{A}_2|e^{+i\phi_2+i\delta_2}$$

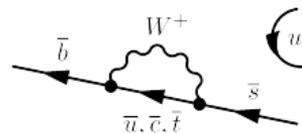
$$\bar{\mathcal{A}}(\bar{X} \rightarrow \bar{Y}) = |\mathcal{A}_1|e^{-i\phi_1+i\delta_1} + |\mathcal{A}_2|e^{-i\phi_2+i\delta_2}$$

$$\frac{|\bar{\mathcal{A}}|^2 - |\mathcal{A}|^2}{|\bar{\mathcal{A}}|^2 + |\mathcal{A}|^2} =$$

$$\frac{2|\mathcal{A}_1||\mathcal{A}_2| \sin(\phi_1 - \phi_2) \sin(\delta_1 - \delta_2)}{|\mathcal{A}_1|^2 + |\mathcal{A}_2|^2 + 2|\mathcal{A}_1||\mathcal{A}_2| \cos(\phi_1 - \phi_2) \cos(\delta_1 - \delta_2)}$$



tree diagram: $\bar{b} \rightarrow \bar{u} u \bar{s}$



penguin diagram: $\bar{b} \rightarrow \bar{u} u \bar{s}$

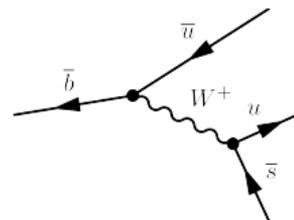
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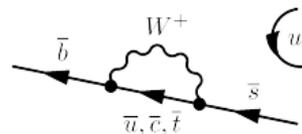
$$\frac{\Gamma(\bar{X} \rightarrow \bar{Y}) - \Gamma(X \rightarrow Y)}{\Gamma(\bar{X} \rightarrow \bar{Y}) + \Gamma(X \rightarrow Y)} = \frac{2|\mathcal{A}_1||\mathcal{A}_2| \sin(\phi_1 - \phi_2) \sin(\delta_1 - \delta_2)}{|\mathcal{A}_1|^2 + |\mathcal{A}_2|^2 + 2|\mathcal{A}_1||\mathcal{A}_2| \cos(\phi_1 - \phi_2) \cos(\delta_1 - \delta_2)}$$

→ Measurable numerical value of CPV depends on

- difference between the weak phases: $|\phi_1 - \phi_2|$ (👍)
- difference between the strong phases: $|\delta_1 - \delta_2|$ (👎)
- product of amplitudes: $|\mathcal{A}_1 \mathcal{A}_2| / |\mathcal{A}_1 + \mathcal{A}_2|^2$ (👎)



tree diagram: $\bar{b} \rightarrow \bar{u} u \bar{s}$



penguin diagram: $\bar{b} \rightarrow \bar{u} u \bar{s}$

The Thing about Direct CPV

→ **Key channels** (non-comprehensive and IMHO)

→ $\alpha: B^0 \rightarrow \rho\rho / \pi\pi$ (CPV in **mixing + decay**, i.e. time-dependent)

→ $\beta: B^0 \rightarrow J/\psi K_s$ (CPV in **mixing + decay**, i.e. time-dependent)

→ $\beta_s: B_s^0 \rightarrow J/\psi \phi$ (CPV in **mixing + decay**, i.e. time-dependent)

→ $\gamma: B\text{'ish} \rightarrow D\text{'ish} K\text{'ish}$ (*mostly* **direct** CPV, i.e. time-integrated)

→ Why direct CPV for γ ?

→ no need to tag initial flavor

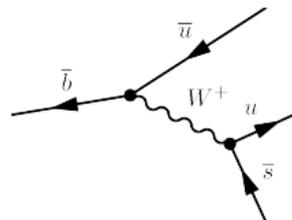
→ bring amplitudes on same level by nifty composing suppressed and favored modes

→ experimental extraction of strong phases in similar decays

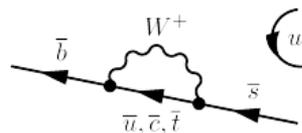
Example: $B^0 \rightarrow K^+ \pi^-$

- Interference between tree and penguin
- CPV observable: **time-integrated** asymmetry

$$\frac{\Gamma(B^0 \rightarrow K^+ \pi^-) - \Gamma(\bar{B}^0 \rightarrow K^- \pi^+)}{\Gamma(B^0 \rightarrow K^+ \pi^-) + \Gamma(\bar{B}^0 \rightarrow K^- \pi^+)}$$



tree diagram: $\bar{b} \rightarrow \bar{u}u\bar{s}$



penguin diagram: $\bar{b} \rightarrow \bar{u}u\bar{s}$

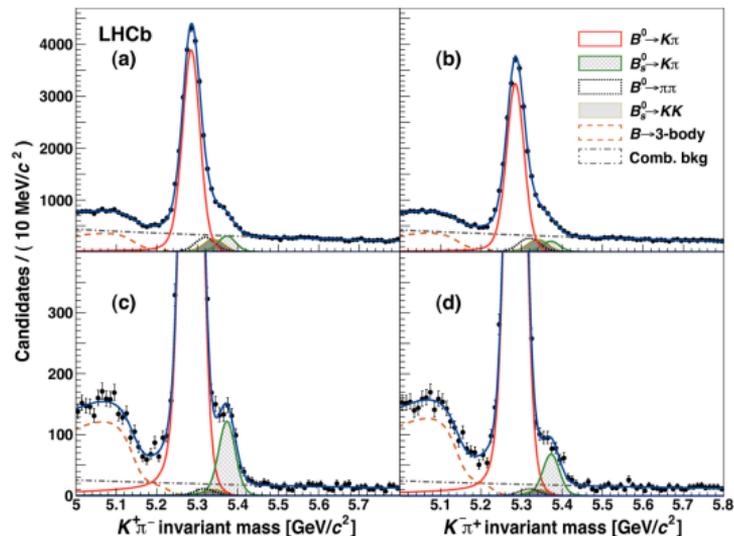
Early results in 2011

- improved determination:

$$A_{CP}(B^0 \rightarrow K^+\pi^-) \\ = -.080 \pm .007 \text{ (stat)} \pm .003 \text{ (syst)}$$

- first measurement:

$$A_{CP}(B_s^0 \rightarrow K^-\pi^+) \\ = +.27 \pm .04 \text{ (stat)} \pm .01 \text{ (syst)}$$



Phys. Rev. Lett. 110 (2013), 221601

Improved results in 2018

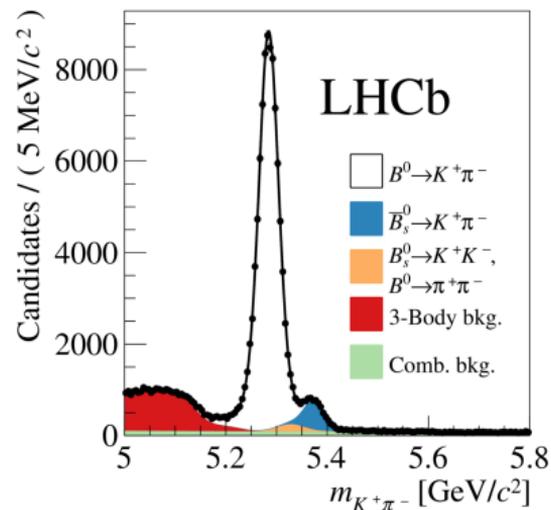
$$\rightarrow A_{\text{CP}}(B^0 \rightarrow K^+ \pi^-) = -0.084 \pm 0.004 (\text{stat}) \pm 0.003 (\text{syst})$$

$$\rightarrow A_{\text{CP}}(B_s^0 \rightarrow K^- \pi^+) = +0.213 \pm 0.015 (\text{stat}) \pm 0.007 (\text{syst})$$

Remember: direct CPV $\sim \sin \Delta\phi_{\text{weak}} \times \sin \Delta\delta_{\text{strong}}$

\rightarrow strong phase difference $\Delta\delta_{\text{strong}}$ not yet accessible from theory (with acceptable uncertainties)

\rightarrow just two (boring) manifestations of CPV?



Phys. Rev. D98 (2018), 032004

→ SM + U-Spin symmetry (assumption) – Phys. Lett. B492 (2000), 297

$$\begin{aligned} |\mathcal{A}(B_s^0 \rightarrow K^- \pi^+)|^2 - |\mathcal{A}(\bar{B}_s^0 \rightarrow K^+ \pi^-)|^2 \\ = |\mathcal{A}(\bar{B}^0 \rightarrow K^- \pi^+)|^2 - |\mathcal{A}(B^0 \rightarrow K^+ \pi^-)|^2 \end{aligned}$$

→ or:

$$\Delta = \frac{A_{\text{CP}}^{B^0}}{A_{\text{CP}}^{B_s^0}} + \frac{\mathcal{B}(B_s^0 \rightarrow K^- \pi^+) \tau(B^0)}{\mathcal{B}(B^0 \rightarrow K^+ \pi^-) \tau(B_s^0)} = 0$$

→ ...the equality follows from a “miracle” which occurs in the standard model and is not expected in common new physics models – Phys. Lett. B621 (2005), 126-132

Direct CP Violation

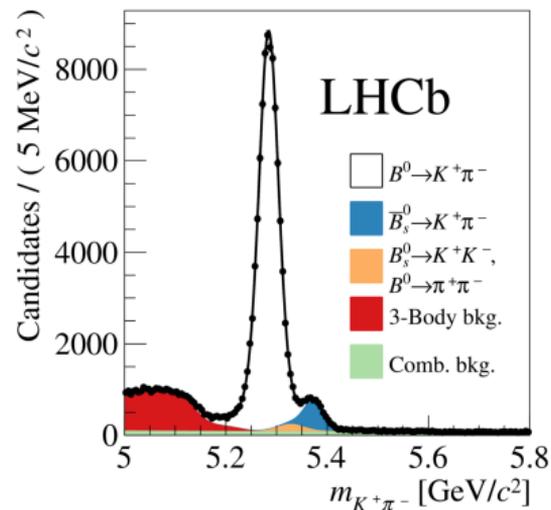
→ Theory (assuming U-Spin symmetry):

$$\Delta = \frac{A_{\text{CP}}^{B^0}}{A_{\text{CP}}^{B_s^0}} + \underbrace{\frac{\mathcal{B}(B_s^0 \rightarrow K^- \pi^+) \tau(B^0)}{\mathcal{B}(B^0 \rightarrow K^+ \pi^-) \tau(B_s^0)}}_{\text{from literature}^*} = 0$$

→ LHCb measurement:

$$\Delta = -0.11 \pm 0.04 (\text{meas.}) \pm 0.03 (\text{lit.}^*)$$

*HFLAV avg. arXiv:1612.07233 & JHEP 04 (2013), 001

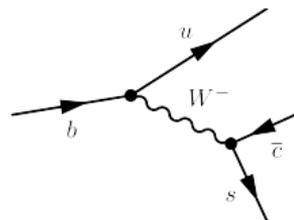
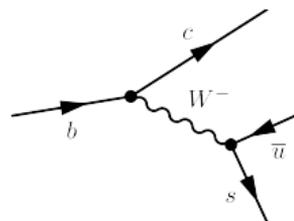


Phys. Rev. D98 (2018), 032004

Example: GLW* method in $B^+ \rightarrow D^0 K^+$

*Phys. Lett. B253 (1991), 483 & Phys. Lett. B265 (1991), 172

- Reconstruct D^0 in CP-even states, such as $K^+ K^-$
 - $B^+ \rightarrow D_{\text{CP-even}} K^+$
 - $B^- \rightarrow D_{\text{CP-even}} K^-$
- Reconstruct D^0 self-tagging: $D^0 \rightarrow K^- \ell^+ \nu_\ell$, or $K^- \pi^+$
 - $B^+ \rightarrow D^0 K^+$ and c.c.
 - $B^+ \rightarrow \bar{D}^0 K^+$ and c.c.
- Combination of all 6 (4) modes allow clean extraction of γ (no theory input for strong phase necessary)
- Note: CPV is not in D^0 , but needs entire decay chain!

tree: $b \rightarrow s [\bar{c}u]_{D^0}$ tree: $b \rightarrow s [c\bar{u}]_{D^0}$

$$\cos 2\beta > 0$$

Combined dataset Babar & Belle

arXiv:1804.06152

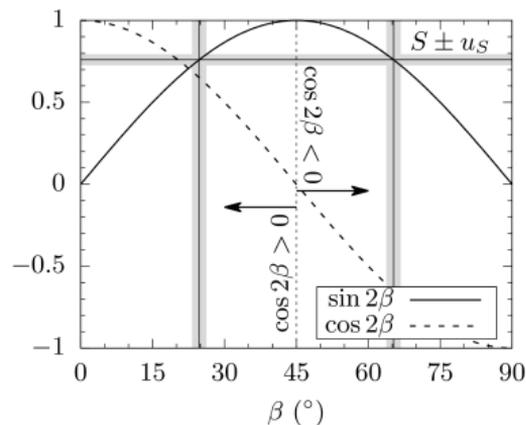
→ Time-dependent Dalitz analysis for

$$B^0 \rightarrow D^{(*)} h^0, D \rightarrow K_s \pi^+ \pi^-$$

→ $\Gamma(B^0(t=0) \rightarrow [K_s \pi^+ \pi^-]_{D^{(*)}} h^0 | t) = \dots$

$$\text{Im}(\mathcal{A}_{D^0} \mathcal{A}_{\bar{D}^0}^*) \cos 2\beta - \text{Re}(\mathcal{A}_{D^0} \mathcal{A}_{\bar{D}^0}^*) \sin 2\beta$$

→ Clean extraction of β (not just $\sin 2\beta$)



$S \pm u_S$ from JHEP11 (2017), 170